

Transverse Target-Spin Asymmetry Associated with DVCS on the Proton and a Resulting Model-Dependent Constraint on J_u vs J_d

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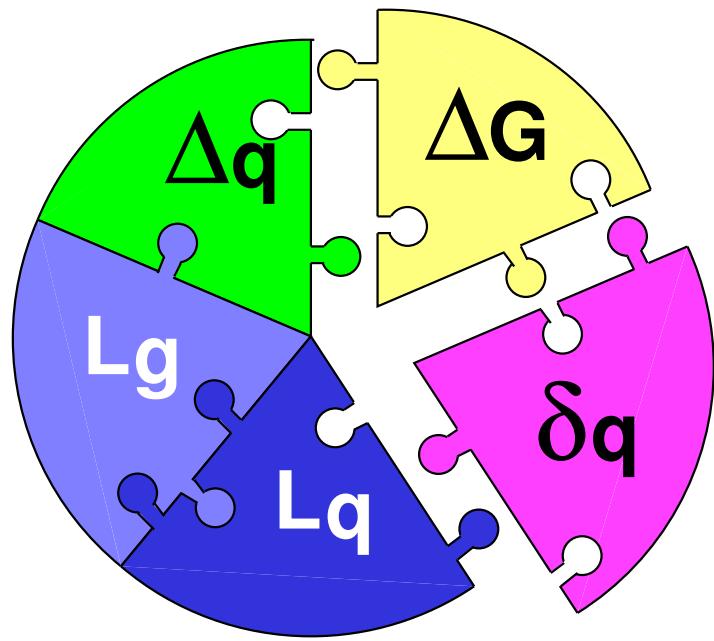
for the HERMES Collaboration

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- Motivation: Spin Structure of the Nucleon from Generalized Parton Distributions
 - Deeply Virtual Compton Scattering and Transverse Target-Spin Asymmetry
 - The HERMES Experiment and the Preliminary Result on the TTSA
 - A Model-Dependent Constraint on the Quark Total Angular Momenta in the Nucleon
 - Summary and Outlook



Motivation: Spin Structure of the Nucleon



Nucleon Spin

$$\frac{1}{2} = \underbrace{\frac{1}{2}(\Delta u + \Delta d + \Delta s)}_{J_q} + \underbrace{L_q}_{\textcolor{red}{J_g}} + \underbrace{\Delta G + L_g}_{\textcolor{red}{J_g}}$$

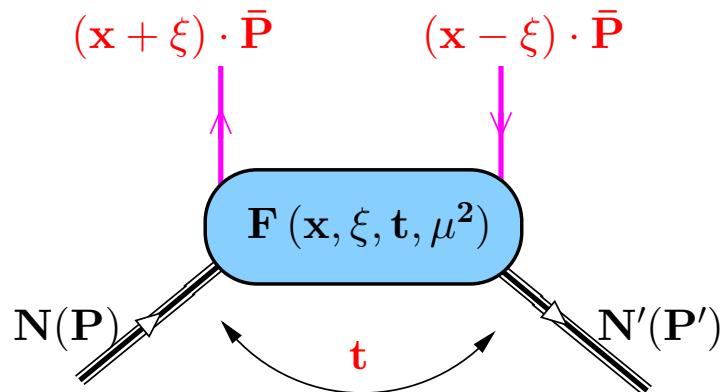
$\Delta \Sigma \sim 20 - 35\%$: Measured in DIS

ΔG : First measurements

L_q, L_g : Unknown!

Generalized Parton Distributions $\Rightarrow J_q, J_g (L_q, L_g)$

Generalized Parton Distributions



- F : GPDs, defined through ME $\langle P' | \mathcal{O}_{q/g} | P \rangle$
- $x \pm \xi$: Parton longitudinal momentum fractions
- t : Invariant momentum transfer to the target
- μ^2 : Renormalization scale

- For a $S = \frac{1}{2}$ hadron, there are 4 twist-2 parton-helicity non-flip GPDs, H , E , \tilde{H} , and \tilde{E} :

	unpolarized	polarized
nucleon-helicity non-flip	H	\tilde{H}
nucleon-helicity flip	E	\tilde{E}

- GPDs provide an access to J_q (Ji 1997):

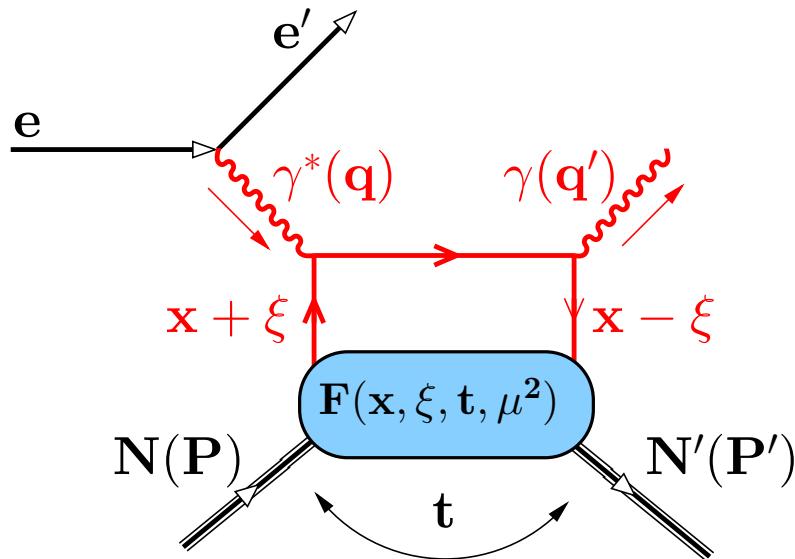
$$J_q(\mu^2) = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H_q(x, \xi, t, \mu^2) + E_q(x, \xi, t, \mu^2)].$$

How to Study Generalized Parton Distributions

- GPDs are related to known quantities (parton densities, nucleon FFs):

in the forward limit:	$H_q(x, 0, 0, \mu^2) = q(x, \mu^2)$	E_q not measurable through DIS
nucleon FFs:	$\int_{-1}^1 dx H_q(x, \xi, t, \mu^2) = F_1^q(t)$	$\int_{-1}^1 dx E_q(x, \xi, t, \mu^2) = F_2^q(t)$

- GPDs enter in hard exclusive reactions, e.g., DVCS:



Kinematics: $Q^2 = -q^2$,
 $t = (P - P')^2$,
 $x_B = \frac{Q^2}{2Pq} \simeq \frac{2\xi}{1 + \xi}$.

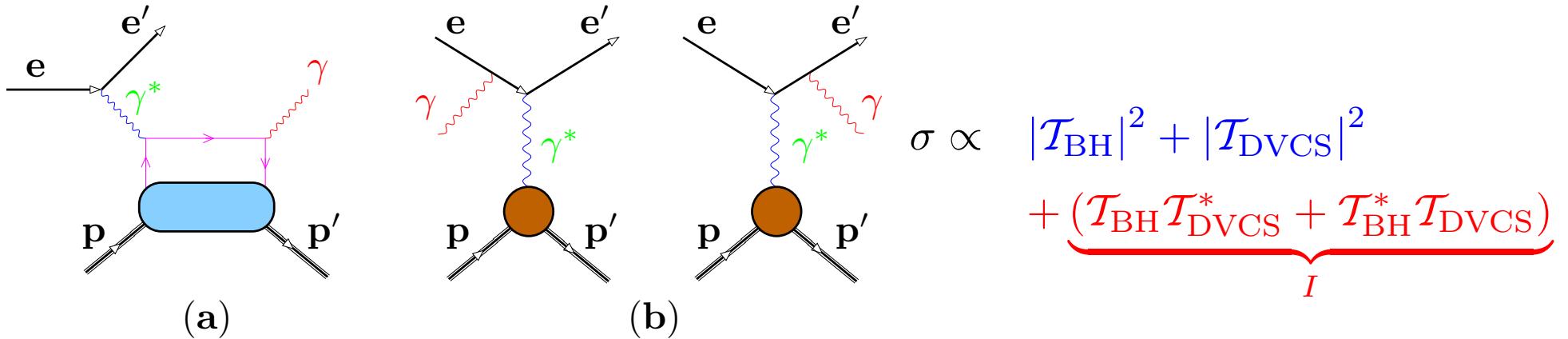
At large Q^2 , fixed x_B and t ,

$$A_{LO}^{\gamma^* p \rightarrow \gamma p} = \sum_q e_q^2 \int_{-1}^1 dx K_{\text{pert.}}(x, \xi, Q^2) F_q(x, \xi, t, Q^2).$$

- The Mellin moments in x of GPDs can be calculated in Lattice QCD.

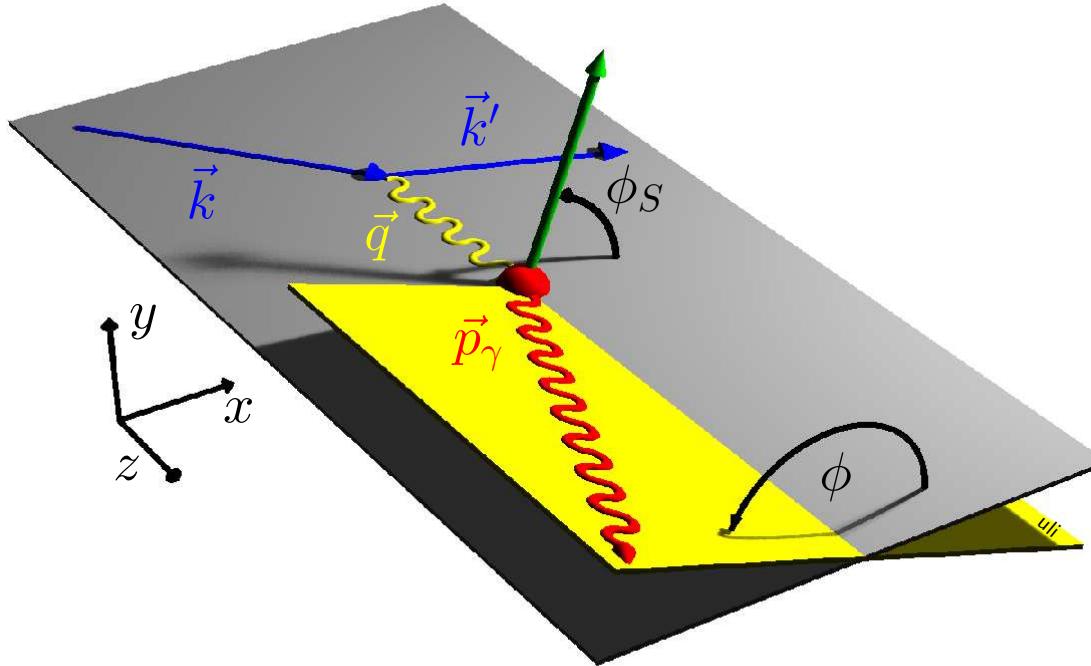
Deeply Virtual Compton Scattering

- The same final state in DVCS (a) and Bethe-Heitler (b) \Rightarrow interference:



- \mathcal{T}_{BH} is parameterized in terms of nucleon FFs F_1 and F_2 , calculable in QED.
- $\mathcal{T}_{\text{DVCS}}$ is parameterized in terms of Compton FFs \mathcal{H} , \mathcal{E} , $\tilde{\mathcal{H}}$, and $\tilde{\mathcal{E}}$, which are convolutions of the respective GPDs with the hard-scattering kernels.
- At HERMES, $\mathcal{T}_{\text{BH}} \gg \mathcal{T}_{\text{DVCS}}$, $\mathcal{T}_{\text{DVCS}}$ can be accessed through \mathcal{I} : both its amplitude and phase!

Transverse Target-Spin Asymmetry on the Proton



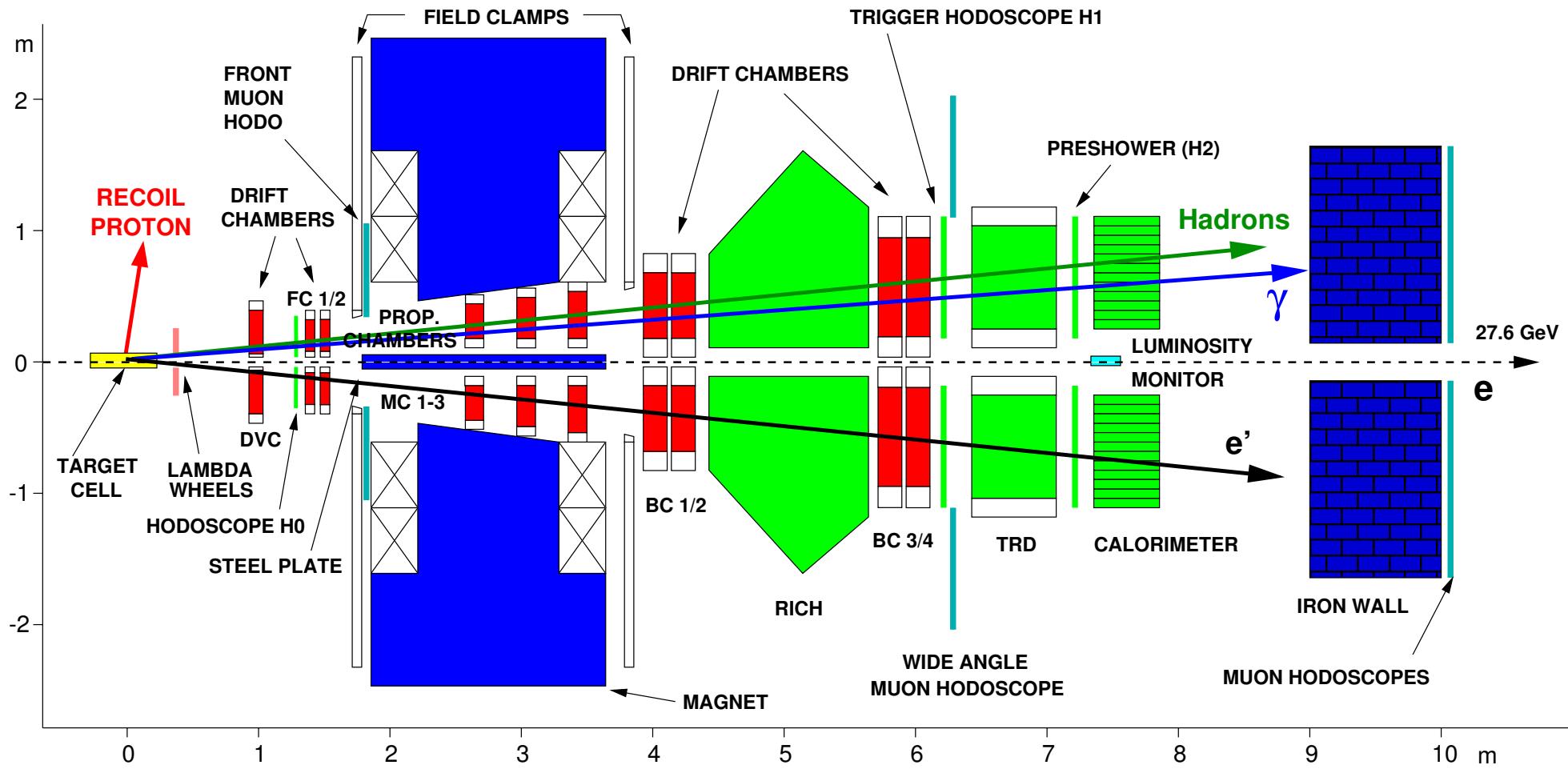
Transverse target-spin asymmetry (Ellighaus, Nowak, Vinnikov, Ye, hep-ph/0506012)

$$\begin{aligned}
 A_{UT}(\phi, \phi_S) &= \frac{d\sigma(\phi, \phi_S) - d\sigma(\phi, \phi_S + \pi)}{d\sigma(\phi, \phi_S) + d\sigma(\phi, \phi_S + \pi)} \simeq \frac{\mathcal{T}^{\text{TP}}}{|\mathcal{T}_{\text{BH}}^{\text{unp}}|^2} \\
 &\propto \text{Im}[\mathcal{F}_2 \mathcal{H} - \mathcal{F}_1 \mathcal{E}] \cdot \sin(\phi - \phi_S) \cos \phi + \text{Im}[\mathcal{F}_2 \tilde{\mathcal{H}} - \mathcal{F}_1 \xi \tilde{\mathcal{E}}] \cdot \cos(\phi - \phi_S) \sin \phi
 \end{aligned}$$

$$\implies A_{\text{UT}}^{\sin(\phi - \phi_S) \cos \phi} \text{ sensitive to } J_q = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^1 dx x (H_q + E_q)$$

The HERMES Experiment

- Transversely polarized hydrogen target data taking in 2002-2005.
- Recoiling protons were not detected.

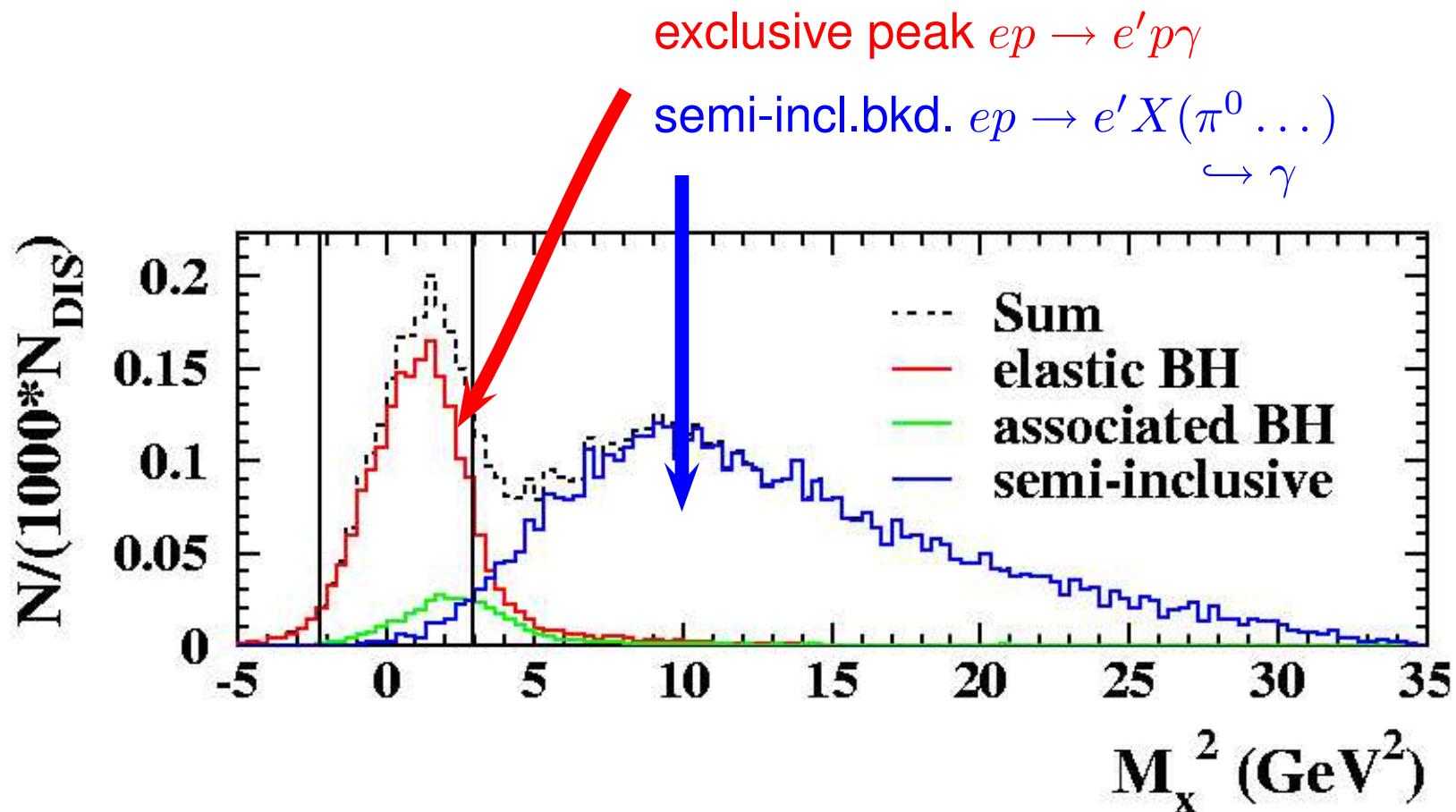


The HERMES Experiment

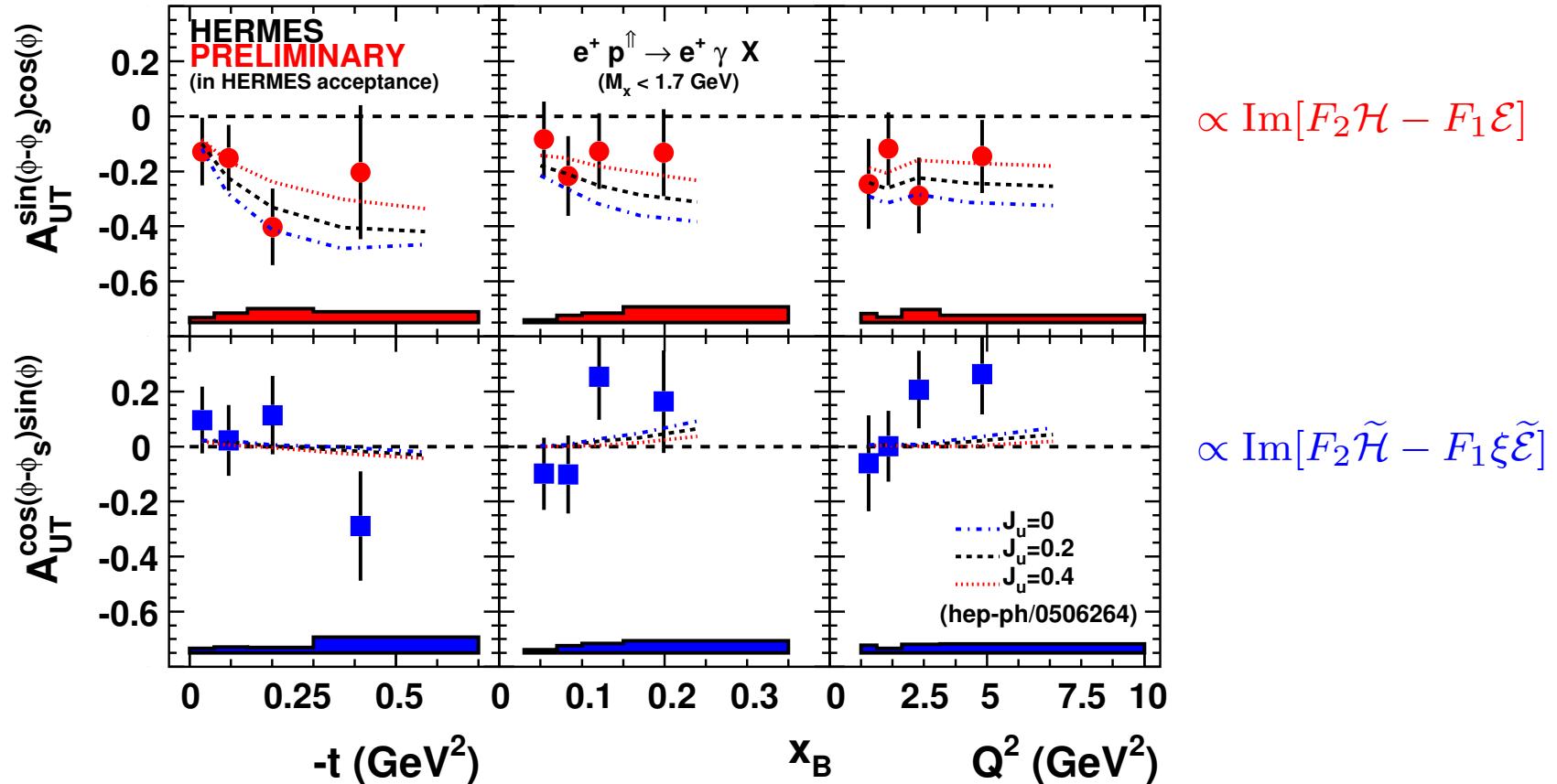
- Exclusivity of the measurement is maintained from the missing mass:

$$M_x^2 = (P_e + P_p - P_{e'} - P_\gamma)^2$$

- Background contribution $\sim 5\%$ is determined from MC and corrected.



Transverse Target-Spin Asymmetry from HERMES



- The presented result is based on the HERMES 2002-2004 data, $\int L dt \simeq 60 \text{ pb}^{-1}$: $\sim 4 \text{ k events in } |t| < 0.7 \text{ GeV}^2, 0.03 < x_B < 0.35, 1 < Q^2 < 10 \text{ GeV}^2$.
- Goeke et al., Prog.Part.Nucl.Phys.47 (2001) 401: The nucleon-helicity flip GPD E in the forward limit is modeled by $e(x) = A \cdot q_{val}(x) + B \cdot \delta(x)$, according to χ QSM model. The values A and B are related to J_q by: $\int dx x[q(x) + e(x)] = J_q$, $\int dx e(x) = F_2^q(0) = k^q$.
- hep-ph/0506264: $A_{UT}^{\sin(\phi-\phi_s)\cos(\phi)}$ sensitive to J_u and insensitive to the other parameters.

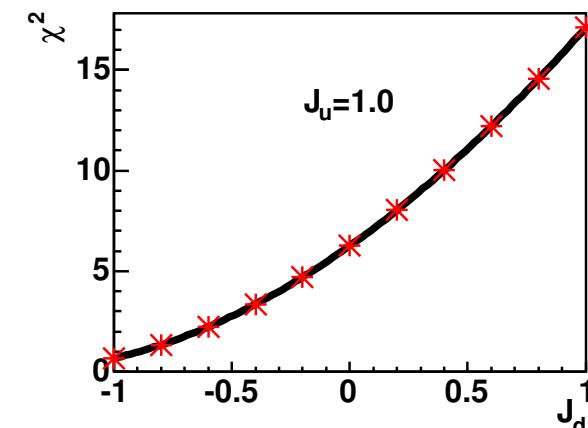
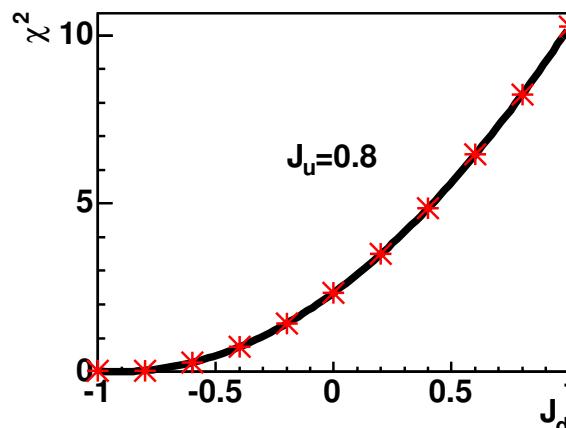
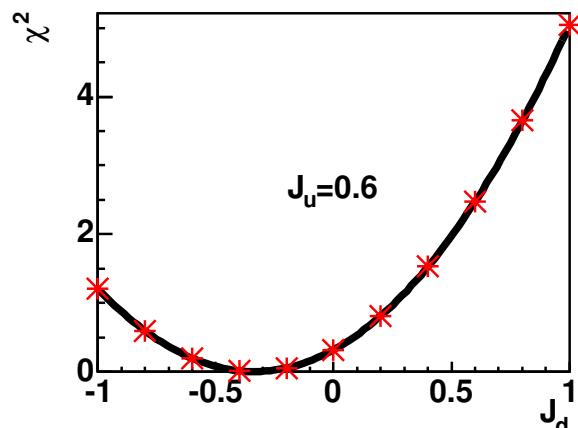
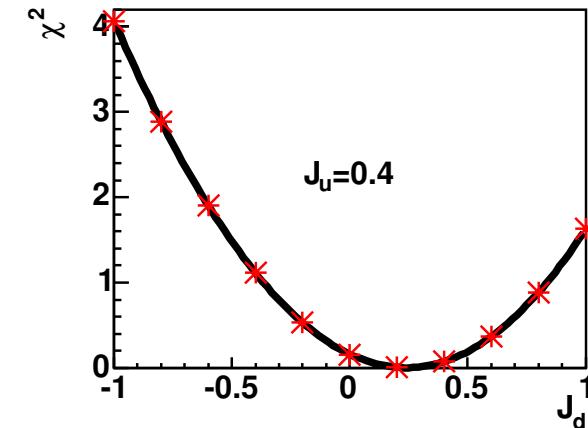
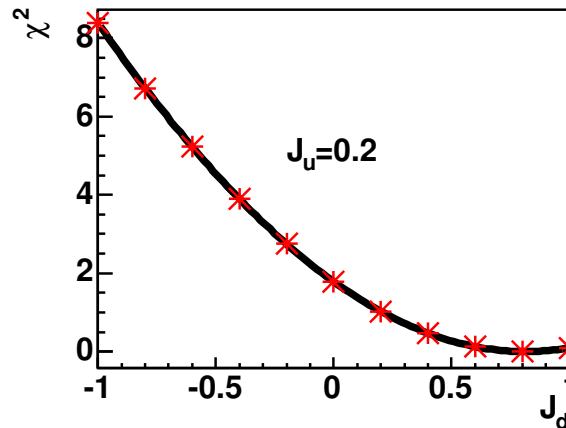
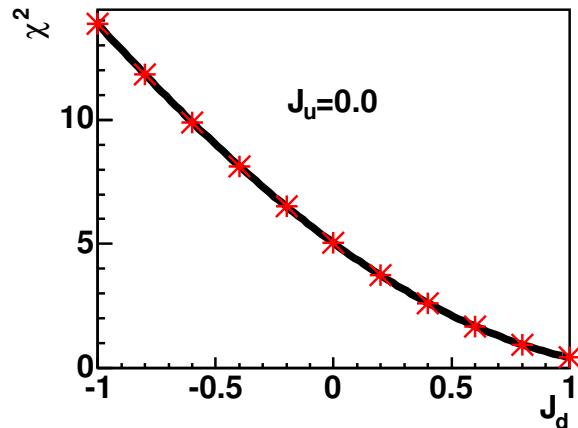
A Model-Dependent Constraint on J_u vs J_d

- In order to compare the theoretical predictions with the experimental results, calculate

$$\chi^2_{exp}(J_u, J_d) = \frac{\left[A_{UT}^{\sin(\phi - \phi_S) \cos \phi} |_{exp} - A_{UT}^{\sin(\phi - \phi_S) \cos \phi} |_{VGG}(J_u, J_d) \right]^2}{\delta A_{stat}^2 + \delta A_{syst}^2}$$

in a step of 0.2 in J_u and J_d , and interpolate inbetween by a 5th order polynomial.

- The $1-\sigma$ constraint on J_u vs J_d is determined by $\chi^2(J_u, J_d) \leq \chi^2_{min} + 1$.



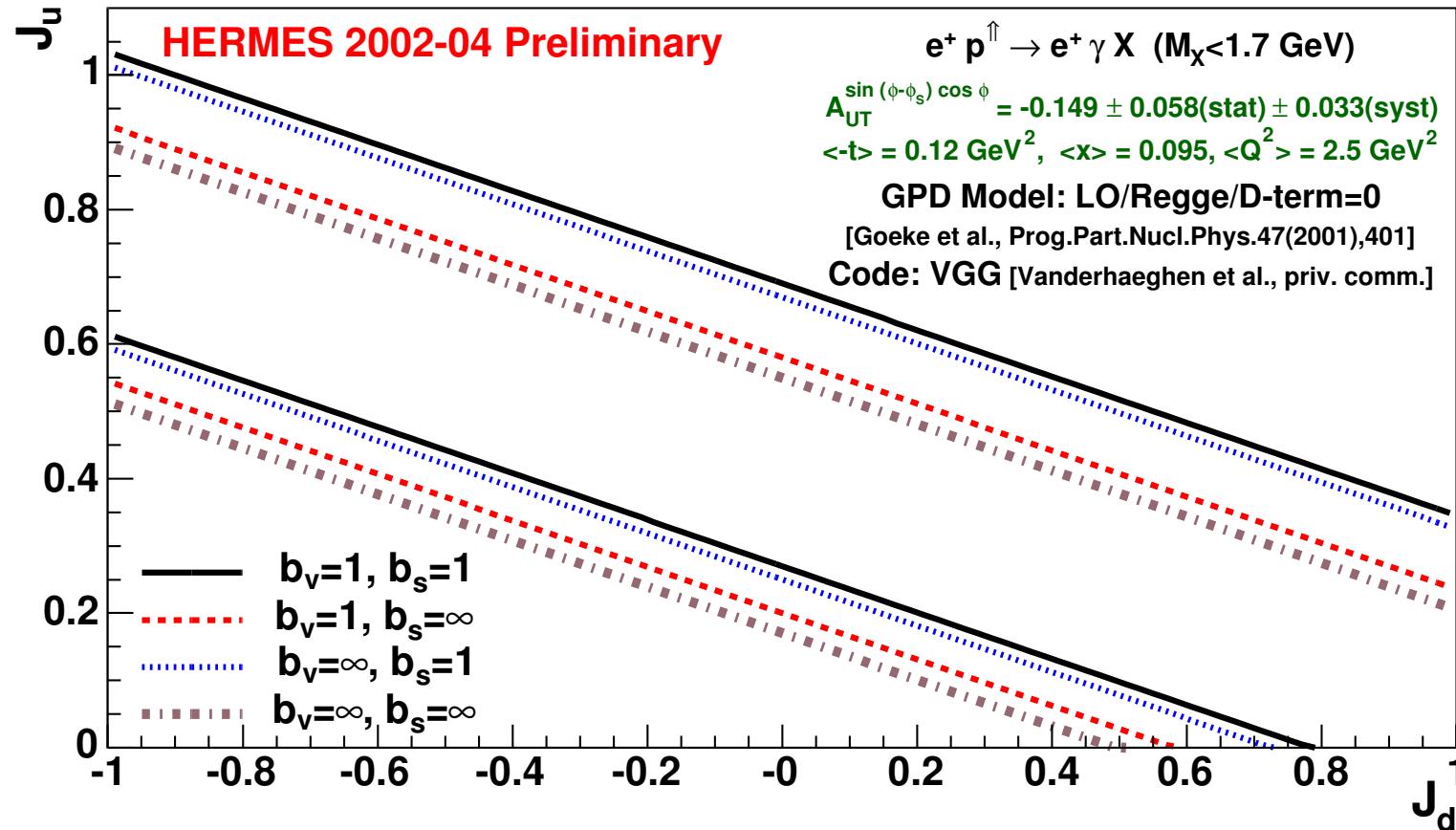
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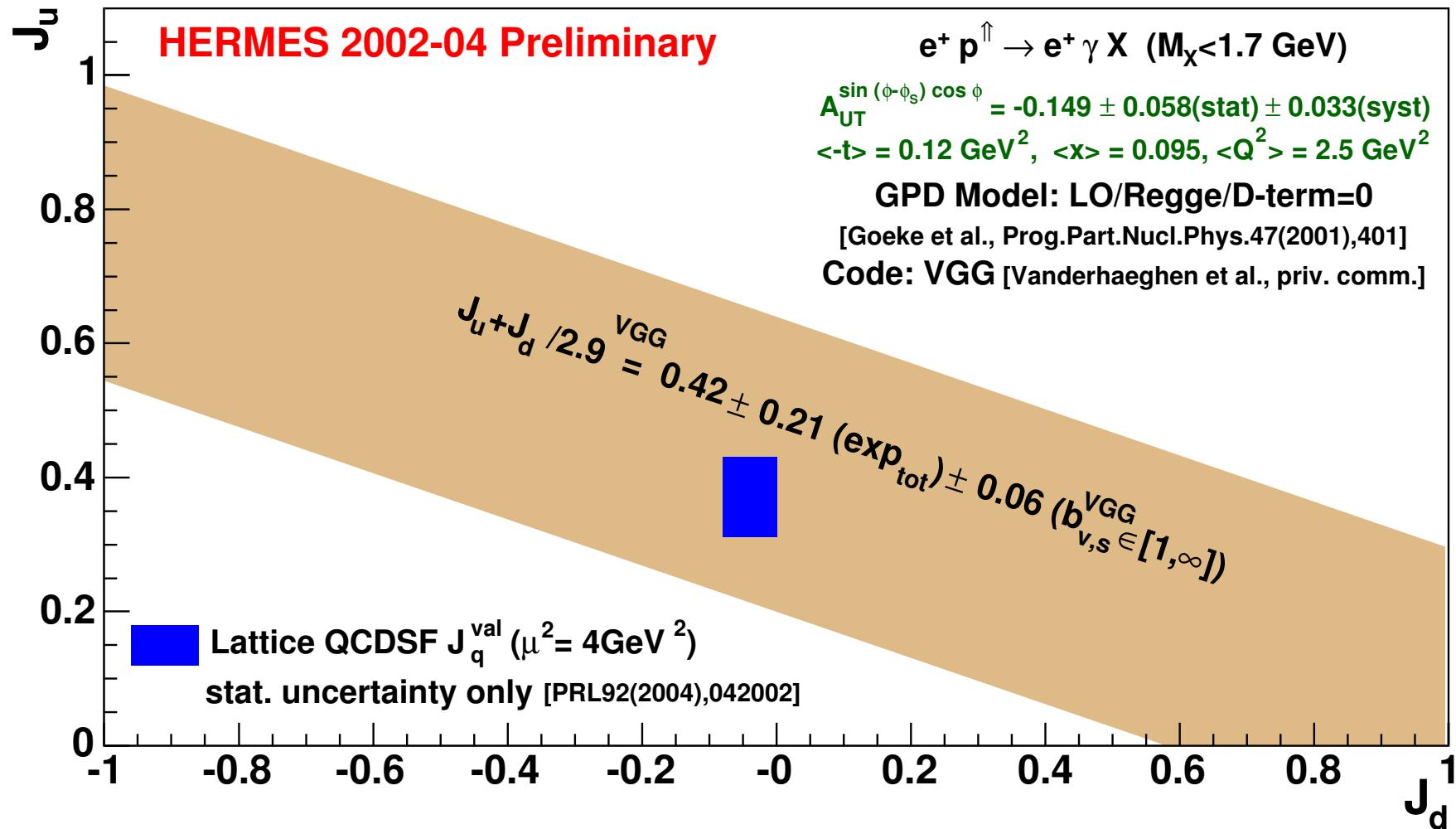
$$\chi^2_{exp}(J_u, J_d) = \frac{\left[A_{UT}^{\sin(\phi - \phi_S) \cos \phi} |_{exp} - A_{UT}^{\sin(\phi - \phi_S) \cos \phi} |_{VGG(J_u, J_d)} \right]^2}{\delta A_{stat}^2 + \delta A_{syst}^2}$$

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A Model-Dependent Constraint on J_u vs J_d



- The quenched Lattice calculation was done with the pion masses 1070, 870, and 640 MeV, and extrapolated linearly in m_π^2 to the physical value.

Summary and Outlook

Summary

- The TTSA associated with DVCS on the proton has been firstly measured at HERMES. This asymmetry is sensitive to the GPD E and to the quark total angular momentum J_q .
- A model-dependent constraint on J_u vs J_d is obtained by comparing the HERMES result on the TTSA and the theoretical predication based on a GPD model.

Outlook

- At present, the uncertainty is dominated by the statistical one. The situation will be improved after including the 2005 data: the statistics will be doubled.

HERMES is aiming at providing a more complete picture of nucleon spin.

