

HERMES Status Report

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for the  Collaboration

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Outline

- Introduction

- Physics Highlights from HERMES

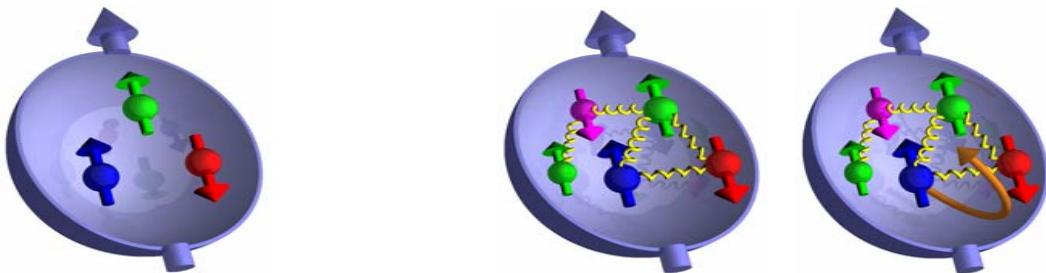
- Isoscalar extraction of ΔS
- Model-dependent constraint on J_u, J_d

- Recoil Detector Status

- Detector calibration
- Noise studies
- Tracking and alignment
- Particle identification

- Summary and Outlook

Spin Structure of the Nucleon

$$\frac{1}{2} = \frac{1}{2} \underbrace{(\Delta u_v + \Delta d_v + \Delta q_{sea})}_{\text{Quarks}} + \underbrace{\Delta G}_{\text{Gluons}} + \underbrace{L_q + L_g}_{\text{Orbital angular momentum}}$$


- Quark spin $\Delta\Sigma = \Delta u + \Delta d + \Delta \bar{u} + \Delta \bar{d} + \Delta s + \Delta \bar{s}$
 $\Delta\Sigma$ is small $\sim 1/3$ (HERMES, ...)

Phys. Rev. D 75, 012007 (2007)

ΔS

- Gluon spin ΔG

- Orbital angular momentum $L_q + L_g$

Physics Highlights

- Isoscalar extraction of ΔS

arXiv:0803.2993, DESY 07-223, submitted to PRL,

A. Airapetian et al., Measurement of Parton Distributions of Strange Quarks in the Nucleon from Charged Kaon Production in Deep-Inelastic Scattering on the Deuteron

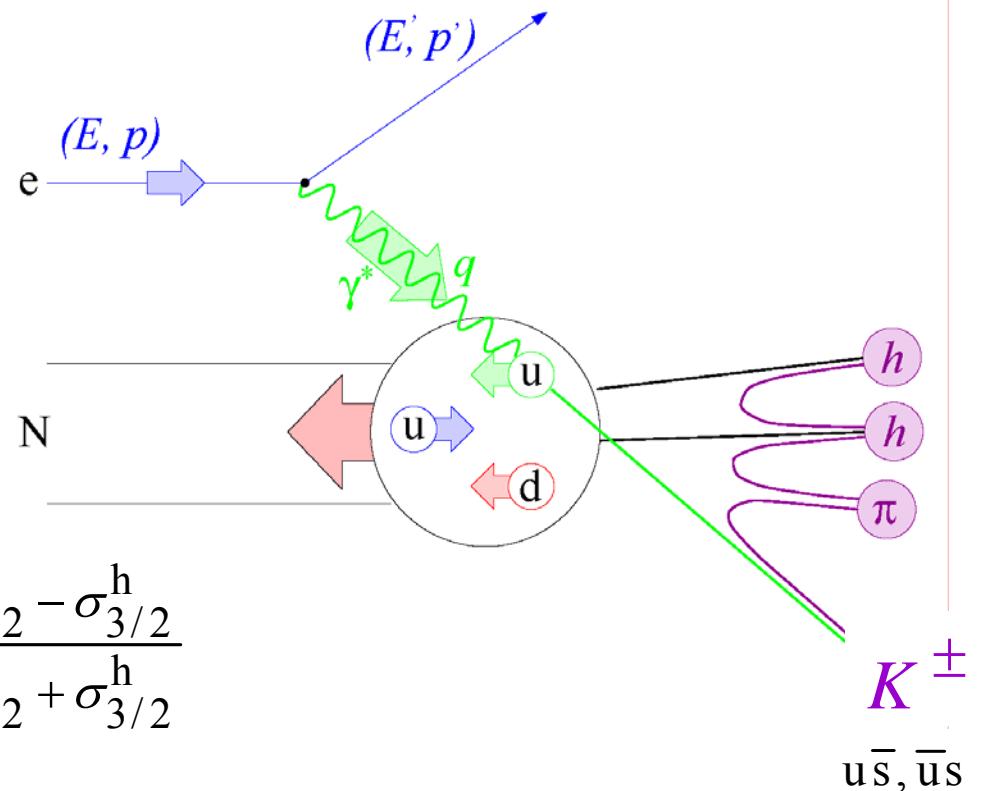
- Model-Dependent Constraint on J_u, J_d

arXiv:0802.2499, DESY 07-225, submitted to JHEP,

A. Airapetian et al., Measurement of Azimuthal Asymmetries with Respect to Both Beam Charge and Transverse Target Polarization in Exclusive Electroproduction of Real Photons

Semi-Inclusive DIS

- Semi-inclusive DIS, observation of a coincident hadron, allows probing of flavor in more detail because of the correlation between flavor structure of the hadron and struck quark



- Measure asymmetries $A_1^h = \frac{\sigma_{1/2}^h - \sigma_{3/2}^h}{\sigma_{1/2}^h + \sigma_{3/2}^h}$

$$A_1^h(x, Q^2, z)^{\text{LO}} = \sum_q P_q^h(x, Q^2, z) \frac{\Delta q(x, Q^2)}{q(x, Q^2)}, \quad z = \frac{E_h}{v}$$

$$P_q^h(x, Q^2, z)$$

'Purity' is a conditional probability that a hadron of type h observed in the final state originated from a struck quark of flavor q.
'Purity' depends on fragmentation functions and unpolarized PDFs

Isoscalar Extraction of ΔS

- Longitudinally polarized deuterium target
 - Strange quark sea in proton and neutron identical
 - Analysis of purities simplifies
- All needed information can be extracted from HERMES data alone:
 - Inclusive $A_{1,d}(x, Q^2)$ and Kaon $A_{1,d}^{K^+ + K^-}(x, Q^2, z)$ double spin asymmetries
 - Kaon multiplicities
- Only assumptions used
 - Isospin symmetry between proton and neutron
 - Charge-conjugation invariance in fragmentation

Isoscalar Extraction of ΔS

$$\begin{pmatrix} A_{l,d}(x) \\ A_{l,d}^{K^+ + K^-}(x) \end{pmatrix} = \begin{pmatrix} P_Q & P_S \\ P_Q^{K^+ + K^-} & P_S^{K^+ + K^-} \end{pmatrix} \begin{pmatrix} \Delta Q(x)/Q(x) \\ \Delta S(x)/S(x) \end{pmatrix}$$

- Extract

$$\Delta S(x) \equiv \Delta s(x) + \Delta \bar{s}(x)$$

$$\Delta Q(x) \equiv \Delta u(x) + \Delta \bar{u}(x) + \Delta d(x) + \Delta \bar{d}(x)$$

- Inclusive purities are simple combinations of unpolarized PDFs

$$P_Q(x) = \frac{5Q(x)}{5Q(x) + 2S(x)}, P_S(x) = \frac{2S(x)}{5Q(x) + 2S(x)}$$

- Kaon purities can be computed from the unpolarized Kaon multiplicity assuming only charge symmetry in fragmentation

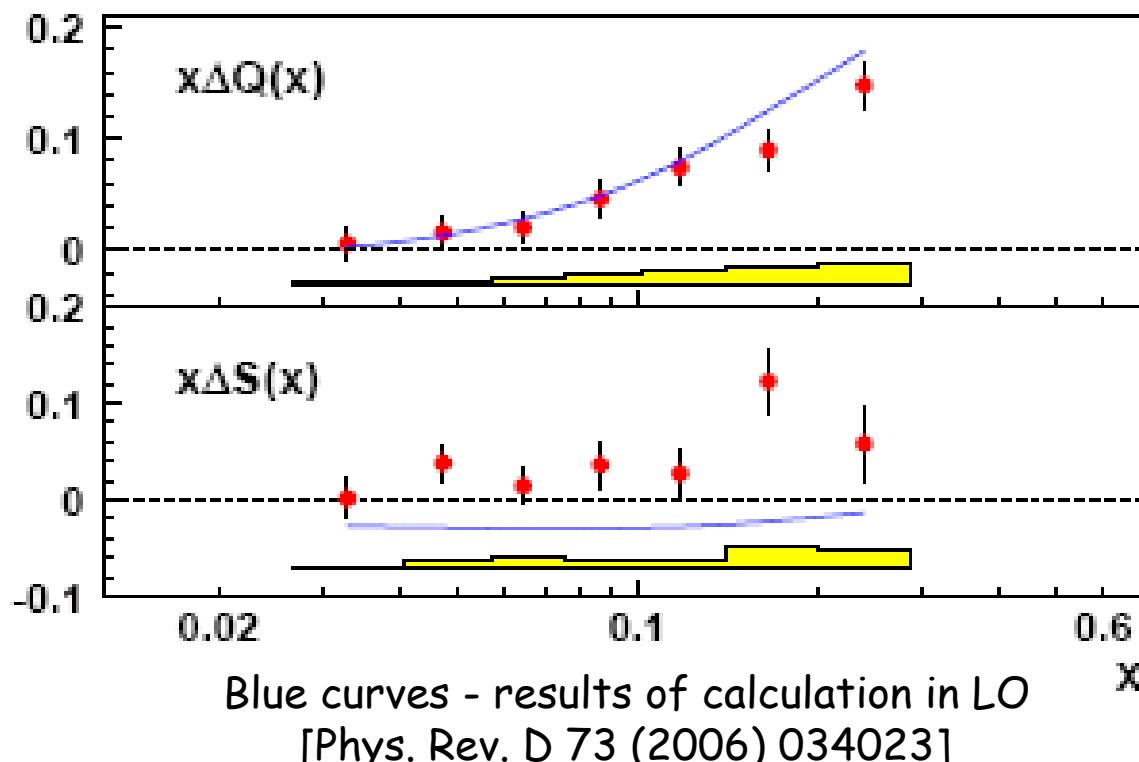
$$D_q^{K^+ + K^-}(z) = D_{\bar{q}}^{K^+ + K^-}(z)$$

Results on Isoscalar Extraction of ΔS

- No dependence on fragmentation model
- No hint of negative strange sea in measured range

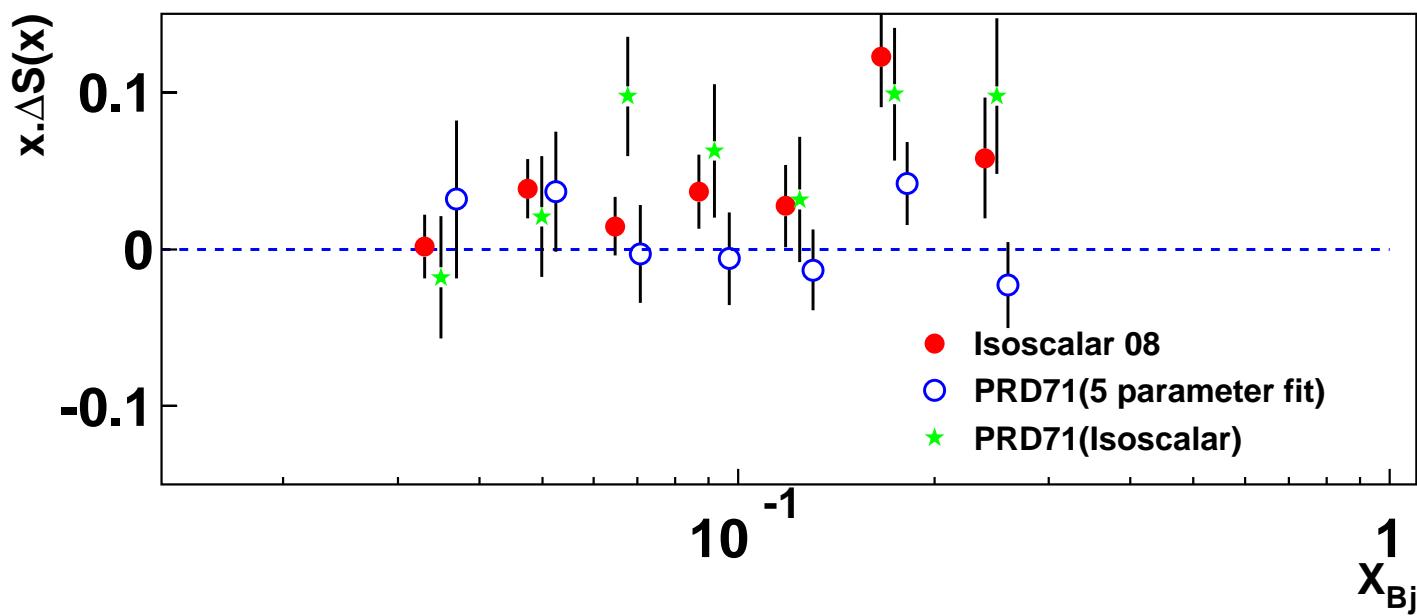
$$\Delta Q = 0.359 \pm 0.026 \pm 0.018$$

$$\Delta S = 0.037 \pm 0.019 \pm 0.027$$



Results on Isoscalar Extraction of ΔS

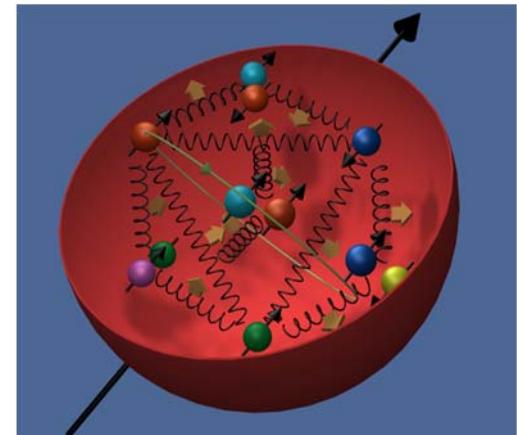
- No dependence on fragmentation model
- No hint of negative strange sea in measured range
- Result is consistent with full 5 flavor extraction at HERMES



Access to the Orbital Angular Momentum

Nucleon spin:

$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + L_q + J_g$$



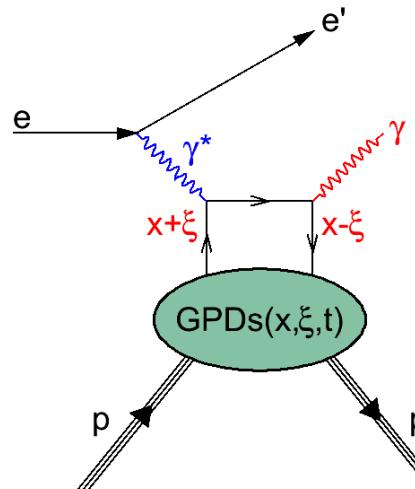
- Quark spin: $\Delta\Sigma$ measured $\sim 1/3$ (HERMES, ...)
- Quark orbital momentum: L_q unknown

Generalized parton distributions (GPDs)

- Incorporate our knowledge about PDFs and FF to a new theoretical framework
- Access the orbital angular momentum of quarks in nucleon

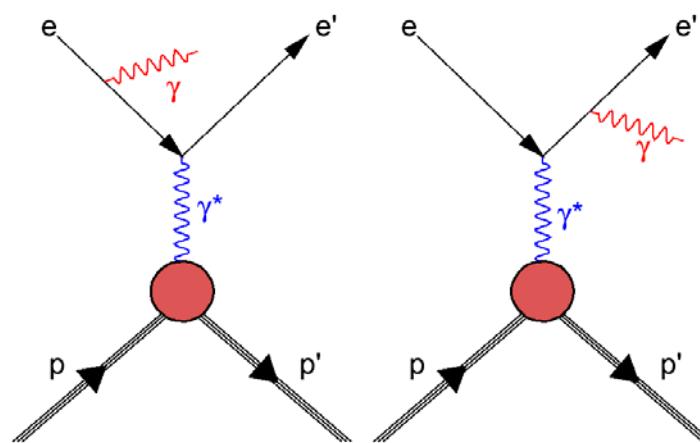
Access to GPD via Deeply Virtual Compton Scattering (DVCS)

Cleanest way to access GPDs: DVCS



$$d\sigma(eN \rightarrow eN\gamma) \propto |T_{BH}|^2 + |T_{DVCS}|^2 + \underbrace{T_{BH} T_{DVCS}^* + T_{BH}^* T_{DVCS}}_{\text{interference term}}$$

Bethe-Heitler



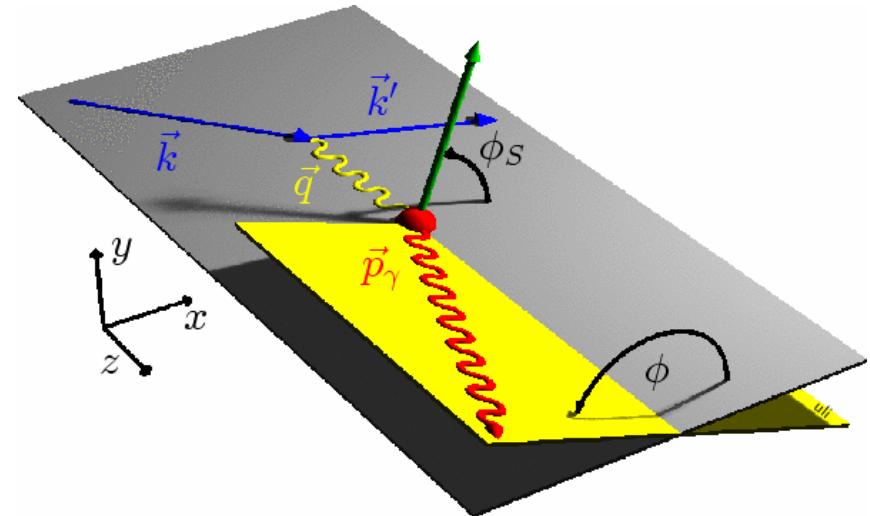
- DVCS and Bethe-Heitler: the same initial and final state
- Bethe-Heitler dominates at HERMES kinematics
- GPDs accessible through cross-section differences and azimuthal asymmetries via interference term

Transverse Target Spin Asymmetry (TTSA)

- Ji relation → access to the total angular momentum of quarks

$$J_q = \lim_{t \rightarrow 0} \int_{-1}^1 dx x \left[\underbrace{H_q(x, \xi, t) + E_q(x, \xi, t)}_{\text{GPDs}} \right]$$

- Bethe-Heitler - DVCS interference term in $ep \rightarrow e p \gamma$ induces azimuthal asymmetries → access to GPDs H and E

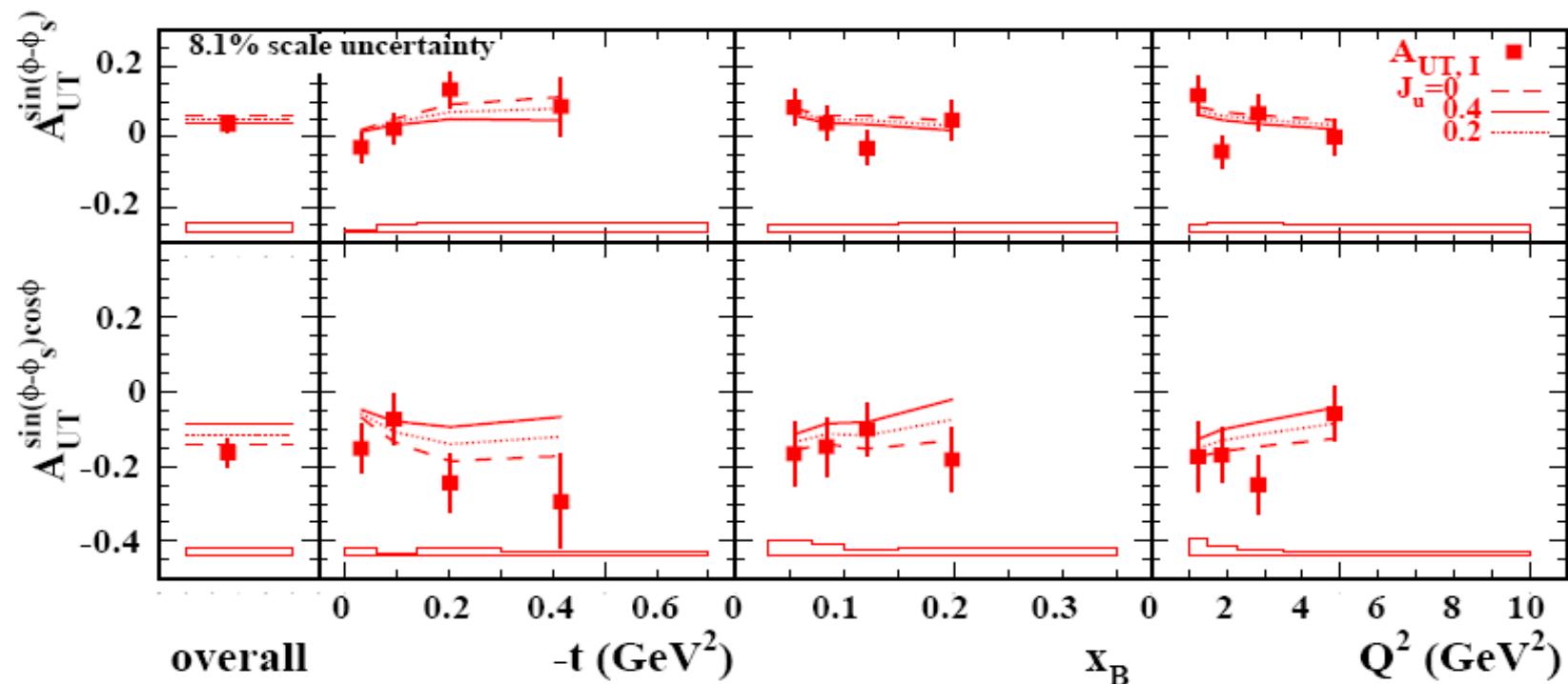


- Measure Transverse Target Spin Asymmetry (TTSA)

$$A_{UT}^I(\phi, \phi_S) \propto [d\sigma^+(\phi, \phi_S) - d\sigma^-(\phi, \phi_S)] - [d\sigma^+(\phi, \phi_S + \pi) - d\sigma^-(\phi, \phi_S + \pi)]$$

- TTSA provides access to GPD E

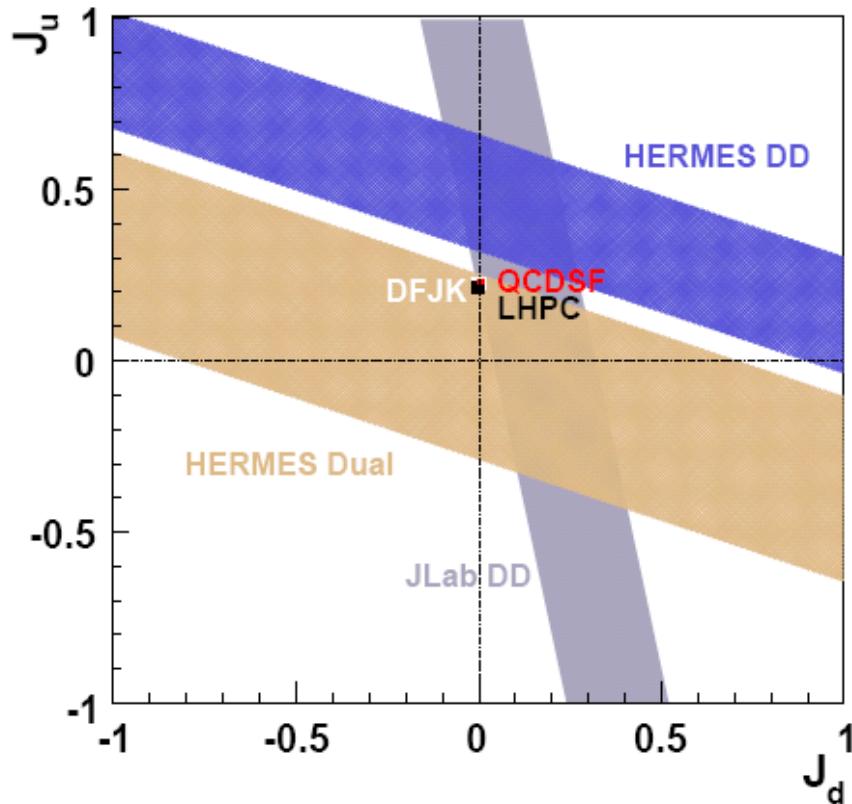
Results on TTSA Asymmetry Amplitudes



Sensitivity of GPD model predictions to J_u at fixed $J_d=0$
[Phys. Rev. D74(2006) 054027]

Model-Dependent Constraint on J_u, J_d

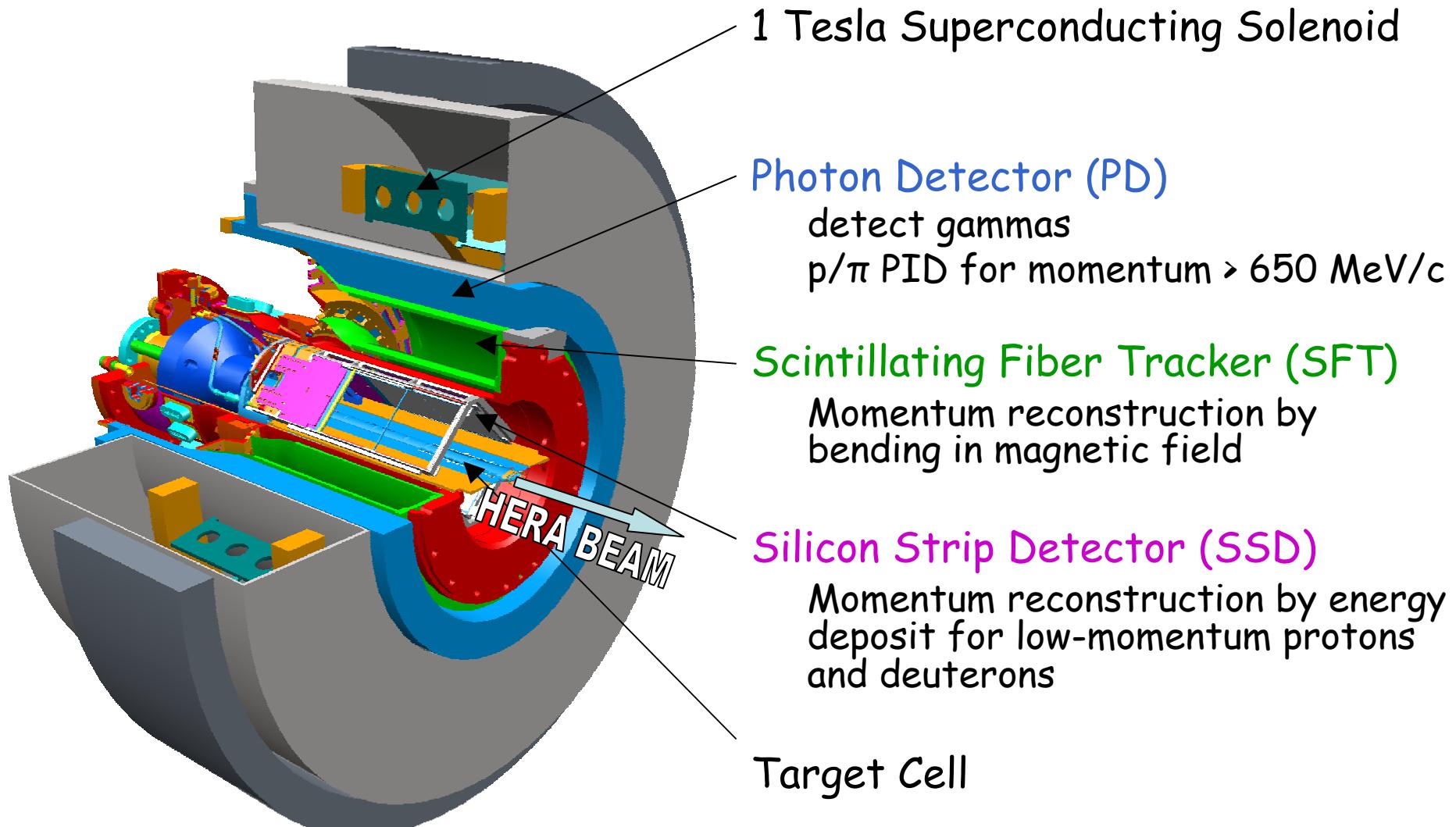
$$\chi^2(J_u, J_d) = \left(A_{\text{UT,I}}^{\sin(\varphi - \varphi_s) \cos n\varphi} |_{\text{exp}} - A_{\text{UT,I}}^{\sin(\varphi - \varphi_s) \cos n\varphi} |_{\text{theo}}(J_u, J_d) \right)^2 / (\delta A_{\text{stat}}^2 + \delta A_{\text{syst}}^2)$$



- J_u, J_d are free parameters in GPD models
- Double Distribution (DD)
[Phys. Rev. D 60 (1999) 094017,
Prog. Part. Nucl. Phys. 47(2001)401]
 $J_u + J_d / 2.8 = 0.48 \pm 0.17$
- Dual Parameterization (Dual)
[hep-ph/0207153,
Phys. Rev. D74(2006) 054027]
 $J_u + J_d / 2.8 = -0.02 \pm 0.27$
- Jlab DD (neutron cross section data)
[Phys. Rev. Lett. 99(2007)242501]
- Lattice calculations QCDSF, LHPC

Recoil Detector

Motivation: exclusivity, background suppression, t-resolution



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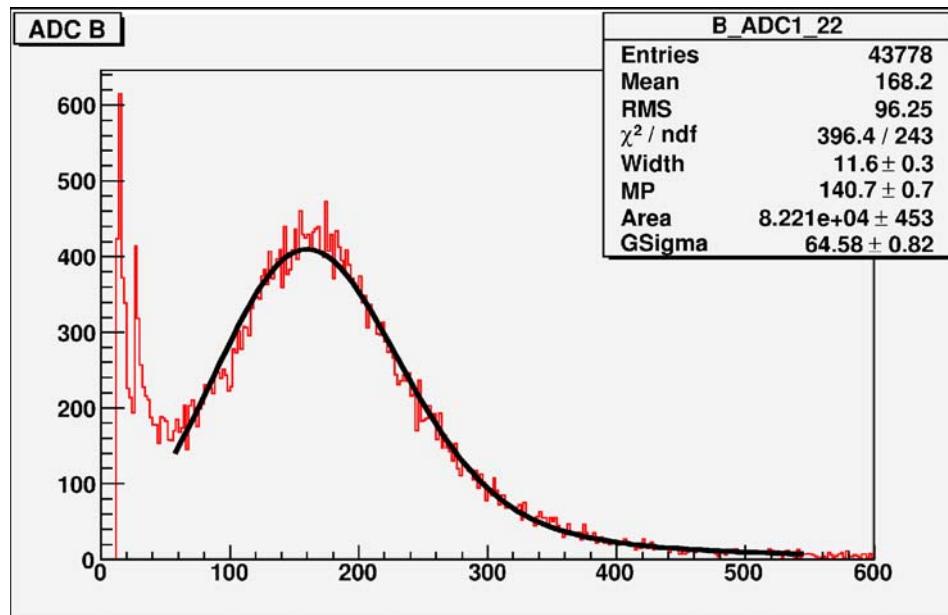
- Recoil Detector Status

- Detector calibration
 - Use MIPs for calibration by selecting π^+ and π^-
- Noise studies
- Tracking and alignment
- Particle identification

- Summary and Outlook

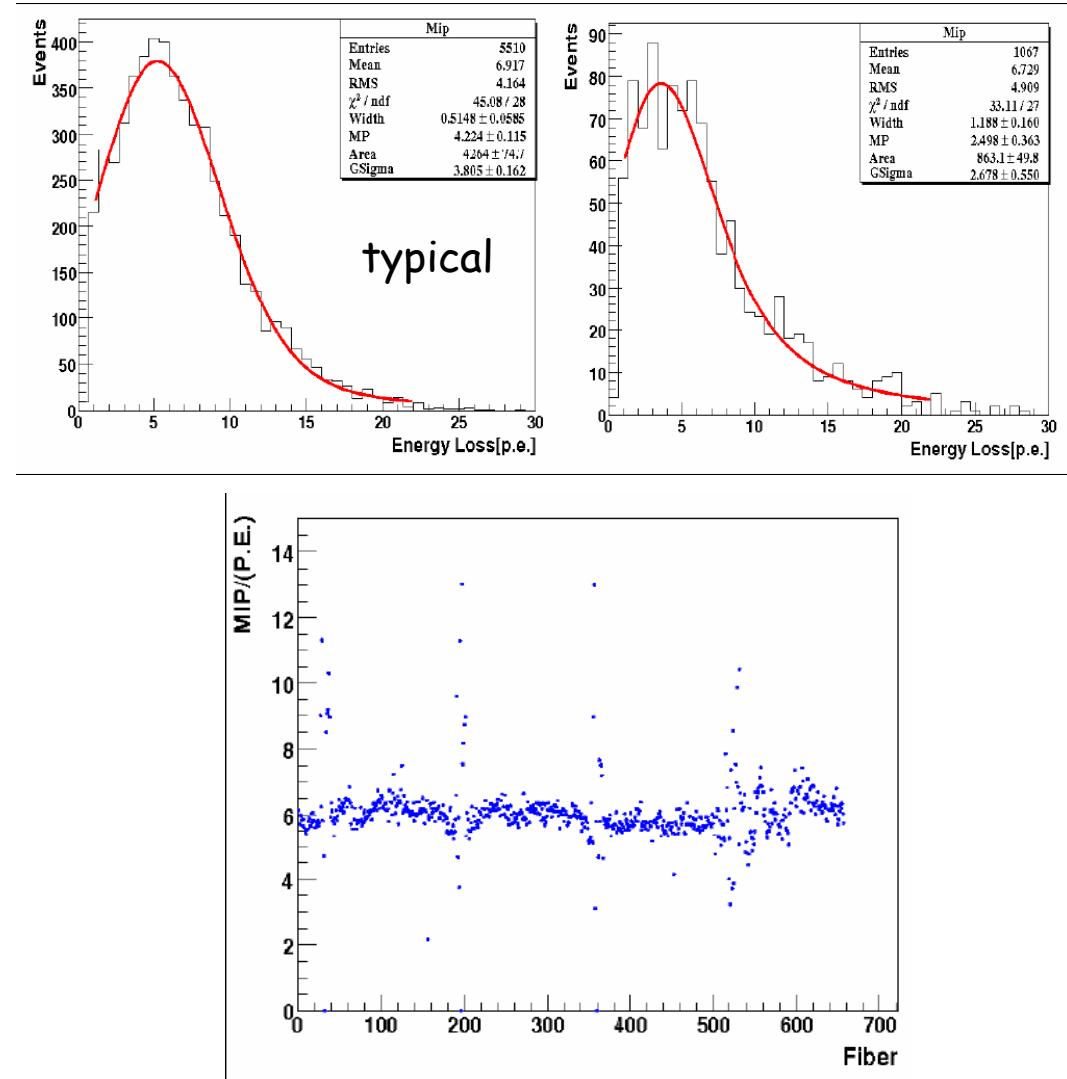
Detector Calibration: Photon detector

- MIP peaks seen for all strips in all layers
- Fit Landau-Gaussian convolution to all ADC spectra
- Need to understand large width of MIP peak
 - Implement correction for attenuation
 - Refine MC digitization
- Get absolute energy calibration
 - Final check with reconstructed π^0



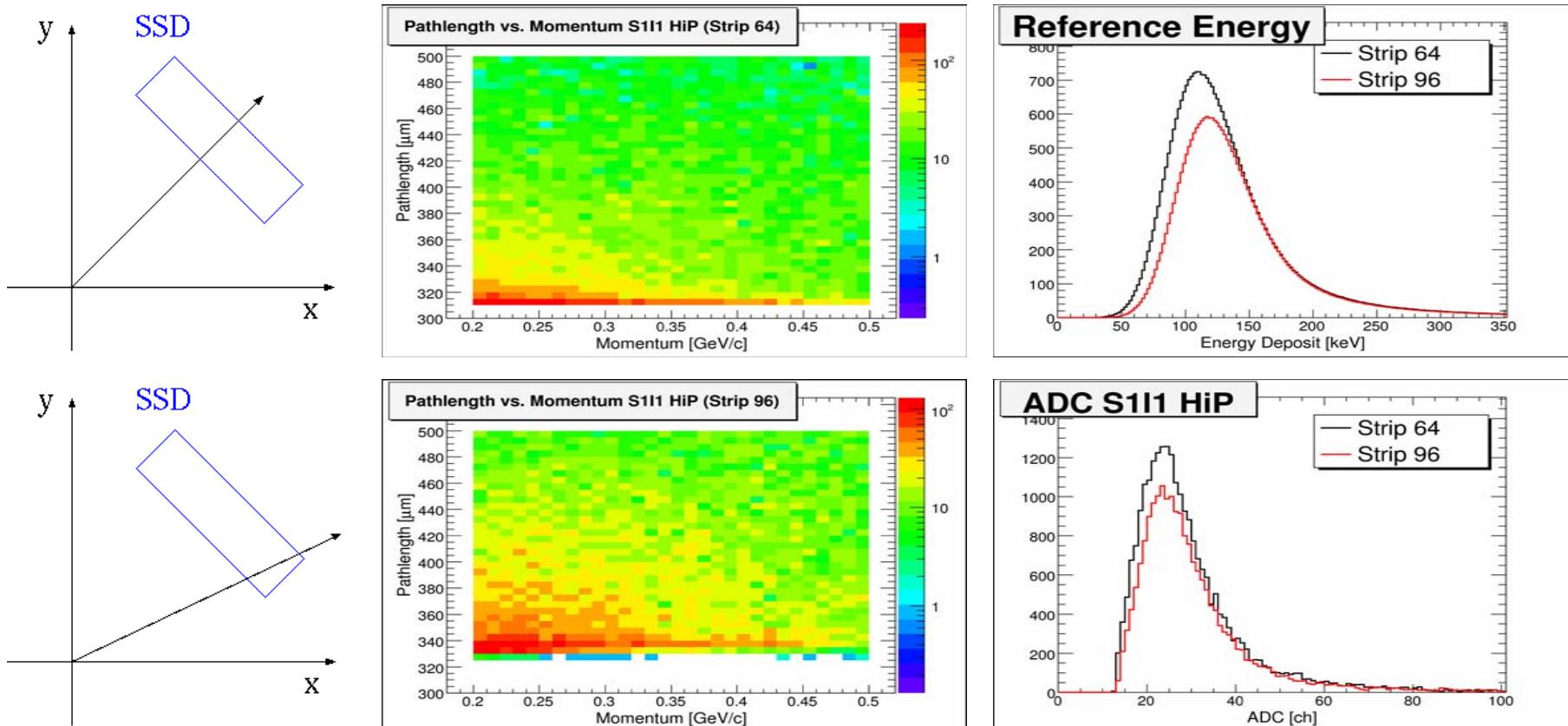
Detector Calibration: Scintillating Fiber Tracker

- MIP peak seen in all fibers
- Some fibers show higher noise
- Calibration currently in number of photoelectrons
- Use MC to get reference energy deposits and get the absolute calibration



Detector Calibration: Silicon Strip Detector

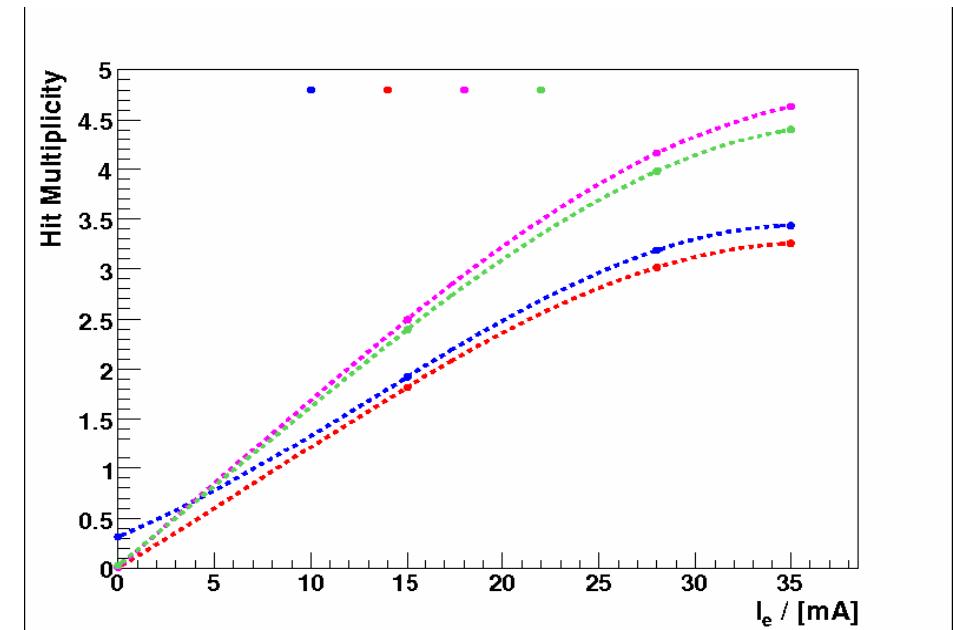
- Strong variations in momentum and path length seen by strips
 - Acceptance, geometry
- Introduce Path Length vs. Momentum histograms for each strip and lookup table to get nominal energy deposit histogram for each strip



- Stability has been studied, studies of nonlinearities in progress

Noise Studies

- Noise studies with random bunches
- Different data sets
 - Beginning of the fill (positron beam current 28mA)
 - End of the fill (15mA)
 - Empty target (35mA)
- Noise studies for SFT
 - Occupancy of the background is strongly correlated with the positron beam current
 - No correlation between hits in the SFT - random distributed noise is implemented in MC
- Similar studies are done for the Silicon and Photon detectors



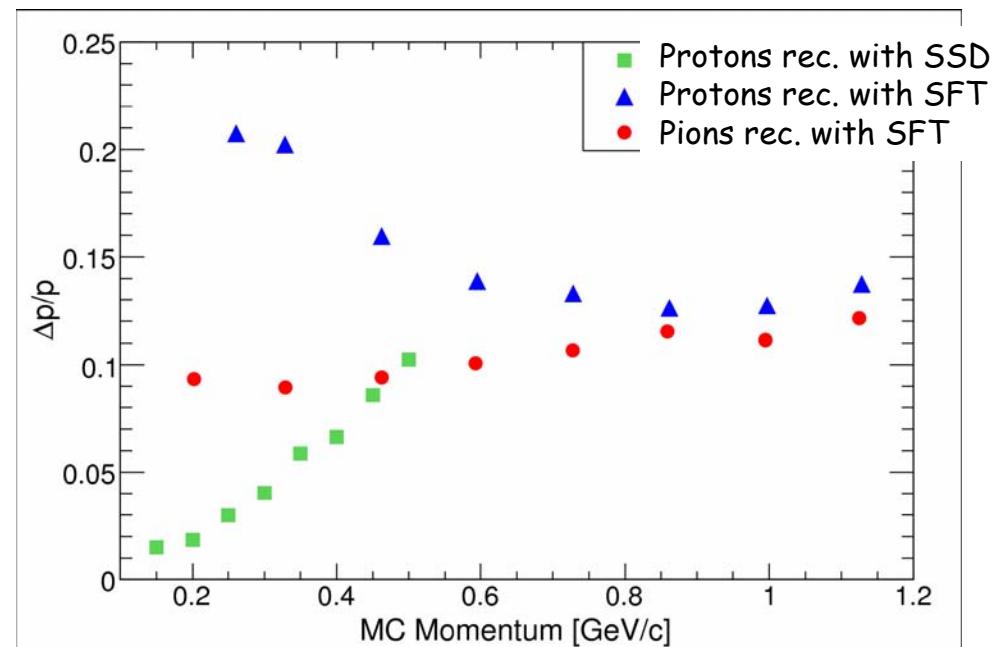
Tracking status

- Track finding works and is efficient, studies of ghost tracks in progress
- Detector alignment is done
- Track fitting takes energy losses and multiple scattering into account
- Track fitting done for proton and pion hypotheses
 - Different tracking reconstructed parameters for future PID
- Accuracy of momentum reconstruction by bending in magnetic field is in agreement with the technical design report
- Tracking is used for calibration
- Until detector calibration is done - time for improvements

Tracking status

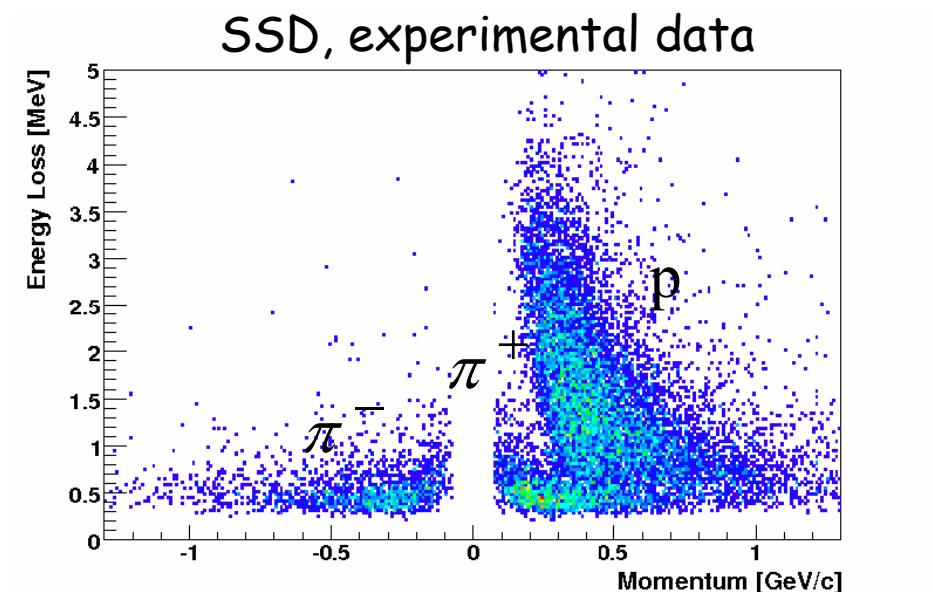
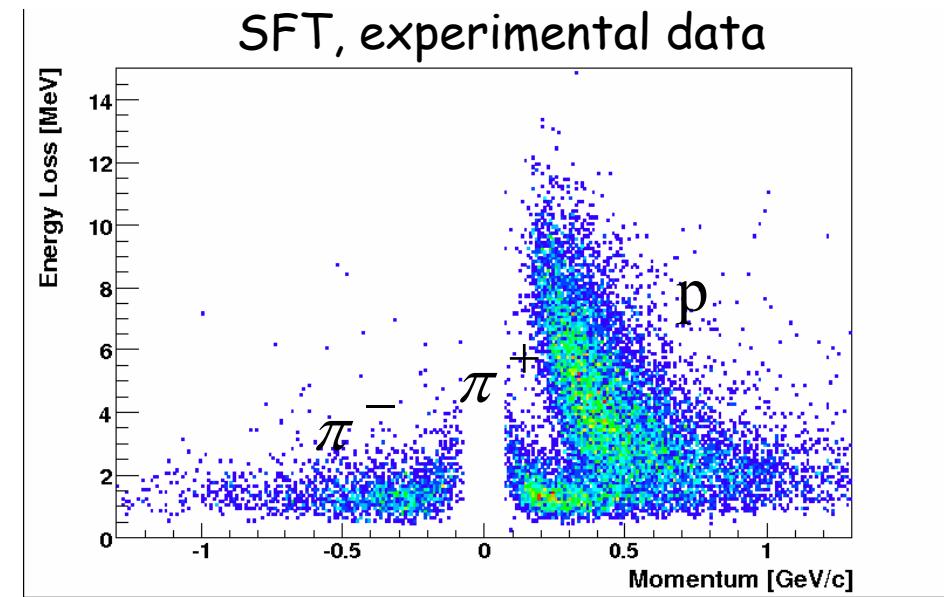
Momentum reconstruction by energy deposit in the Silicon Strip Detector

- First method with simplifications in use
- Two methods under development
 - Lookup tables for momentum as a function of energy deposit and path length, direct reconstruction of momentum
 - Fitting method - will treat energy deposits the same way as measured coordinates during track fitting



Particle Identification Plan

- Particle identification possible in all 3 subdetectors of the Recoil
- Proton-pion separation in momentum range below 700 MeV/c possible with the SSD and SFT
- Photon detector will be used for proton-pion separation at energies above 650 MeV/c



Summary and Outlook

- Analysis of HERMES data is ongoing, new physics results coming soon
- 10 HERMES talks accepted for DIS08
- Recoil Detector calibration and tracking nearly finished, tracking improvements under implementation
- First results from the Recoil Detector are expected this fall

Backup Slides

Method of 5-Flavor Δq Extraction

$$\vec{A}_1 = P \cdot \vec{Q}$$

$$\vec{A}_1 = \begin{pmatrix} A_{1,p} \\ A_{1,p}^+ \\ A_{1,p}^- \\ A_{1,d} \\ \vdots \\ \vdots \\ A_{1,d}^- \end{pmatrix}$$

Measured

$$P_q^h(x) = \frac{e_q^2 q(x) \int dz D_q^h(z)}{\sum_{q'} e_{q'}^2 q'(x) \int dz D_{q'}^h(z)}$$

Simulated

Purity is a conditional probability that a hadron of type h observed in the final state originated from a struck quark of flavor q .
 Purities extracted from LUND string model MC simulation (LEPTO+JETSET) tuned to HERMES unpolarized data

$$\vec{Q} = \begin{pmatrix} \Delta u/u \\ \Delta d/d \\ \Delta \bar{u}/\bar{u} \\ \Delta \bar{d}/\bar{d} \\ \Delta s/s \end{pmatrix}$$

Extracted

System of asymmetries and purities solved through χ^2 minimization of polarizations

Generalized Parton Distributions

- GPDs → PDF

$$H_q(x, 0, 0) = q(x)$$

$$\tilde{H}_q(x, 0, 0) = \Delta q(x)$$

- GPD → FF

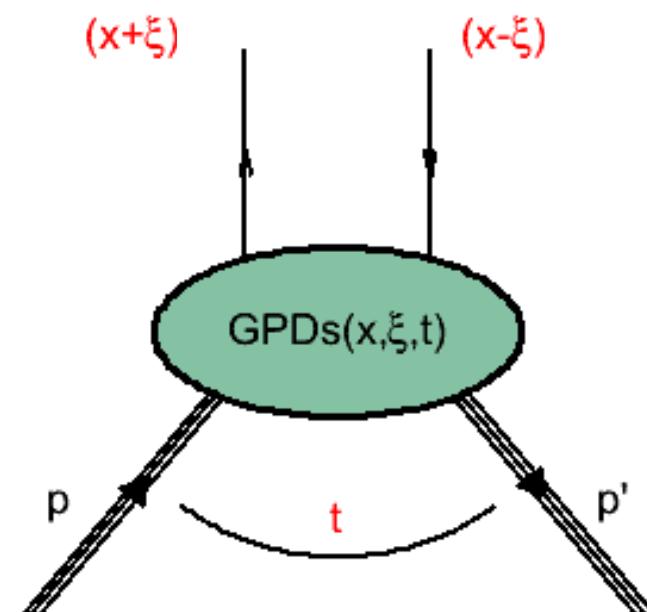
$$\int_{-1}^1 dx H_q(x, \xi, t) = F_1^q(t)$$

$$\int_{-1}^1 dx E_q(x, \xi, t) = F_2^q(t)$$

- H_q, \tilde{H}_q conserve nucleon helicity

- E_q, \tilde{E}_q flip nucleon helicity

$x \pm \xi$ parton longitudinal momentum fractions
 ξ fraction of the momentum transfer
 t invariant momentum transfer to the nucleon



Asymmetries Measurable at HERMES

➤ Beam-Spin Asymmetry (BSA)

$$A_{LU} = \frac{d\sigma(\mathbf{e}^{\rightarrow}, \phi) - d\sigma(\mathbf{e}^{\leftarrow}, \phi)}{d\sigma(\mathbf{e}^{\rightarrow}, \phi) + d\sigma(\mathbf{e}^{\leftarrow}, \phi)} \propto \Im m(\mathcal{H}) \sin(\phi)$$

➤ Beam-Charge Asymmetry (BCA)

$$A_C = \frac{d\sigma(\mathbf{e}^+, \phi) - d\sigma(\mathbf{e}^-, \phi)}{d\sigma(\mathbf{e}^+, \phi) + d\sigma(\mathbf{e}^-, \phi)} \propto \Re e(\mathcal{H}) \cos(\phi)$$

➤ Longitudinal Target Spin Asymmetry (LTSA)

$$A_{UL} = \frac{d\sigma(\mathbf{p}^{\rightarrow}, \phi) - d\sigma(\mathbf{p}^{\leftarrow}, \phi)}{d\sigma(\mathbf{p}^{\rightarrow}, \phi) + d\sigma(\mathbf{p}^{\leftarrow}, \phi)} \propto \Im m(\tilde{\mathcal{H}}) \sin(\phi)$$

➤ Transverse Target Spin Asymmetry (TTSA)

$$A_{UT} = \frac{d\sigma(\mathbf{p}^{\uparrow}, \phi) - d\sigma(\mathbf{p}^{\downarrow}, \phi)}{d\sigma(\mathbf{p}^{\uparrow}, \phi) + d\sigma(\mathbf{p}^{\downarrow}, \phi)} \propto f(\mathcal{H}, \mathcal{E}, \tilde{\mathcal{H}}, \tilde{\mathcal{E}}, \phi, \phi_S)$$

