

# First results on two-hadron interference fragmentation on a transversely polarized hydrogen target

Paul van der Nat  
(on behalf of the HERMES collaboration)

## Layout:

- Introduction
- Results on longitudinally polarized target
- Results on transversely polarized target
- Status of the analysis
- Conclusions & Outlook



# Introduction - 1p fragmentation



$h_1$  couples to  $H_1^\perp(z, z^2 \mathbf{k}_T^2)$

$$\curvearrowleft \mathcal{I}[\dots] \equiv \int d^2 p_T d^2 k_T \mathbf{d}(p_T - k_T - \frac{\mathbf{P}_{h\perp}}{z}) [\dots]$$

$$d\mathbf{S}_{UT}^{Collins} \propto \sum_q e_q^2 \sin(\mathbf{f}_h + \mathbf{f}_s) \mathcal{I} \left[ \frac{\mathbf{k}_T \cdot \hat{\mathbf{P}}_{h\perp}}{M_h} h_1^q H_1^{\perp q} \right]$$

## Difficulties:

- extraction of  $h_1 H_1^\perp$  difficult, needs weighting with  $P_h^\perp$
- Sivers & Collins entangled:

$$d\mathbf{S}_{UT}^{Sivers} \propto \sum_q e_q^2 \sin(\mathbf{f}_h - \mathbf{f}_s) \mathcal{I} \left[ \frac{\mathbf{p}_T \cdot \hat{\mathbf{P}}_{h\perp}}{M_h} f_{1T}^{\perp q} D_1^q \right]$$



# Introduction - 2p fragmentation



$h_1$  couples to:

$$H_1^\perp(z, V, M_h^2, \mathbf{k}_T^2, \mathbf{k}_T \cdot \mathbf{R}_T) \& H_1^\times(z, V, M_h^2, \mathbf{k}_T^2, \mathbf{k}_T \cdot \mathbf{R}_T)$$

Integrate over  $P_h^\perp$ :

left with only  $H_1^\times(z, V, M_h^2)$



$$V \propto z_1 / (z_1 + z_2)$$

$$s_{UT} \propto \sum_q e_q^2 \sin(\mathbf{f}_{R\perp} + \mathbf{f}_S) h_1 H_1^\times$$

Advantages:

- cross section asymmetry directly proportional to  $h_1 H_1^\times$  (no weighting needed)
- No Collins/Sivers “entanglement”
- Completely independent from  $1\pi$  analysis

Disadvantages:

- less statistics
- $H_1^\times$  unknown (but can be measured at Belle & Babar)

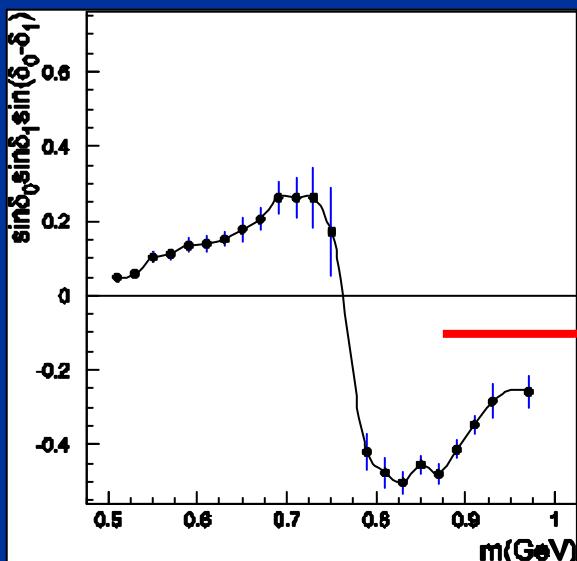
# Interference FF

$$A_{UT} \sim \sin(\mathbf{f}_{R\perp} + \mathbf{f}_S) \sin \mathbf{q} h_1 H_1^\prec$$

Expansion of  $H_1^\prec$  in Legendre moments:

$$H_1^\prec(z, \cos \mathbf{q}, M_{pp}^2) = H_1^{\prec, sp}(z, M_{pp}^2) + \cos \mathbf{q} H_1^{\prec, pp}(z, M_{pp}^2)$$

describe interference between 2 pion pairs  
coming from different production channels



Jaffe et al. [hep-ph/9709322]:

$$\begin{aligned} H_1^{\prec, sp}(z, M_{pp}^2) &= \sin \mathbf{d}_0 \sin \mathbf{d}_1 \sin(\mathbf{d}_0 - \mathbf{d}_1) H_1^{\prec, sp'}(z, M_{pp}^2) \\ &\quad (\mathbf{d}_0(\mathbf{d}_1) \rightarrow S(P)\text{-wave phase shift}) \\ &= \mathcal{P}(M_{pp}^2) H_1^{\prec, sp'}(z, M_{pp}^2) \end{aligned}$$

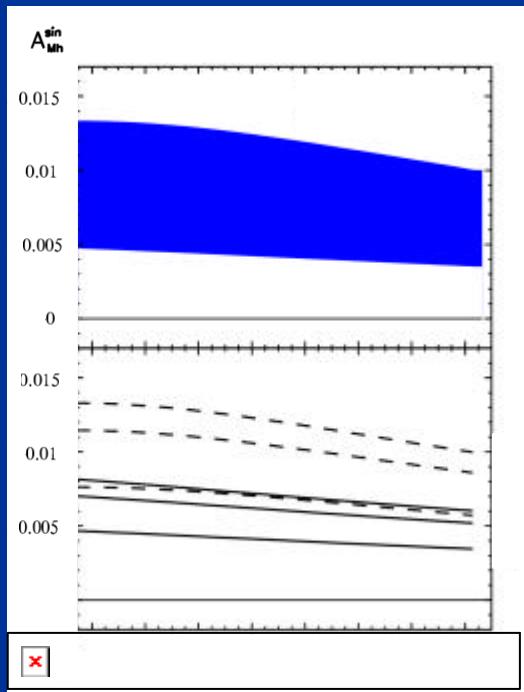
→  $A_{UT}^{\sin \mathbf{f}_{R\perp}}$  might depend strongly on  $M_{pp}$  !!

# Interference FF

$$A_{UT} \sim \sin(\mathbf{f}_{R\perp} + \mathbf{f}_S) \sin \mathbf{q} h_1 H_1^{\star}$$

Expansion of  $H_1^{\star}$  in Legendre moments:

$$H_1^{\star}(z, \cos \mathbf{q}, M_{pp}^2) = H_1^{\star, sp}(z, M_{pp}^2) + \cos \mathbf{q} H_1^{\star, pp}(z, M_{pp}^2)$$



Radici et al. [hep-ph/0110252]:

completely different model, not predicting  
a sign change of the asymmetry around  
the  $r_0$  (times  $-1 \rightarrow$  trento conventions)  
based on the spectator model



# The polarized cross sections



The  $A_{UL}$  &  $A_{UT}$  asymmetries are related to these polarized cross sections (subleading twist, Bacchetta et al.):

$$\mathbf{s}_{UL} \sim \sum_q e_q^2 \sin \mathbf{f}_{R\perp} \sin \mathbf{q} \left[ K_1 |S_{\parallel}| h_L - K_2 |S_{\perp}| h_l \right] (H_1^{\prec,sp} + H_1^{\prec,pp} \cos \mathbf{q})$$

$$\begin{aligned} \mathbf{s}_{UT} \sim & \sum_q e_q^2 |S_{\perp}| \sin(\mathbf{f}_{R\perp} + \mathbf{f}_s) \sin \mathbf{q} K_3 h_l (H_1^{\prec,sp} + H_1^{\prec,pp} \cos \mathbf{q}) \\ & + K_4 \sin \mathbf{f}_s (\dots) \end{aligned}$$



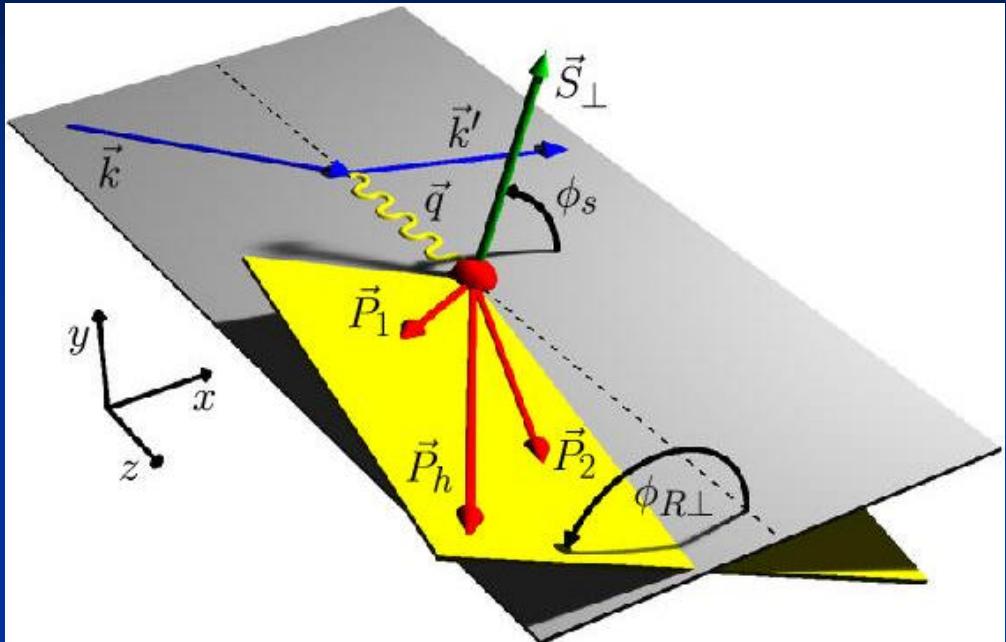
# before integration over $P_{h\perp}$



$$\begin{aligned}
 d^9 \mathbf{s}_{UT} = & \sum_a \frac{\mathbf{a}^2 e_a^2}{2 \mathbf{p} Q^2 y} \bar{S}_T |A(y) \left\{ \frac{\bar{R}_T}{M_h} \sin(\mathbf{j}_R - \mathbf{j}_s) \mathcal{I} \left[ \frac{\vec{p}_T \cdot \vec{k}_T}{2 M M_h} g_{1T} G_1^\perp \right] \right. \\
 & - \frac{|\bar{R}_T|}{M_h} \cos(\mathbf{j}_R - \mathbf{j}_s) \mathcal{I} \left[ \frac{(\vec{p}_T \cdot \vec{P}_{h\perp})(\vec{P}_{h\perp} \times \vec{k}_T) - (\vec{k}_T \cdot \vec{P}_{h\perp})(\vec{P}_{h\perp} \times \vec{p}_T)}{2 M M_h} g_{1T} G_1^\perp \right] \\
 & - \frac{|\bar{R}_T|}{M_h} \sin(2\mathbf{j}_h - \mathbf{j}_R - \mathbf{j}_s) \mathcal{I} \left[ \frac{2(\vec{p}_T \cdot \vec{P}_{h\perp})(\vec{k}_T \cdot \vec{P}_{h\perp}) - \vec{p}_T \cdot \vec{k}_T}{2 M M_h} g_{1T} G_1^\perp \right] \\
 & - \frac{|\bar{R}_T|}{M_h} \cos(2\mathbf{j}_h - \mathbf{j}_R - \mathbf{j}_s) \mathcal{I} \left[ \frac{(\vec{p}_T \cdot \vec{P}_{h\perp})(\vec{P}_{h\perp} \times \vec{k}_T) + (\vec{k}_T \cdot \vec{P}_{h\perp})(\vec{P}_{h\perp} \times \vec{p}_T)}{2 M M_h} g_{1T} G_1^\perp \right] \\
 & + \sin(\mathbf{j}_h - \mathbf{j}_s) \mathcal{I} \left[ \frac{(\vec{p}_T \cdot \vec{P}_{h\perp})}{M} f_{1T}^\perp D_1 \right] + \cos(\mathbf{j}_h - \mathbf{j}_s) \mathcal{I} \left[ \frac{(\vec{P}_{h\perp} \times \vec{p}_T)}{M} f_{1T}^\perp D_1 \right] \Big\} \\
 & + \sum_a \frac{\mathbf{a}^2 e_a^2}{2 \mathbf{p} Q^2 y} \bar{S}_T |B(y) \left\{ \sin(\mathbf{j}_h + \mathbf{j}_s) \mathcal{I} \left[ \frac{(\vec{k}_T \cdot \vec{P}_{h\perp})}{M_h} h_i H_1^\perp \right] \right. \\
 & + \cos(\mathbf{j}_h + \mathbf{j}_s) \mathcal{I} \left[ \frac{(\vec{P}_{h\perp} \times \vec{k}_T)}{M_h} h_i H_1^\perp \right] \left. + \frac{|\bar{R}_T|}{M_h} \sin(\mathbf{j}_R + \mathbf{j}_s) \mathcal{I} \left[ h_i \bar{H}_1^\perp \right] + 3 \sin(3\mathbf{f}_h - \mathbf{f}_s) \right. \\
 & \times \mathcal{I} \left[ \frac{4(\vec{p}_T \cdot \vec{P}_{h\perp})^2 (\vec{k}_T \cdot \vec{P}_{h\perp}) - 2(\vec{p}_T \cdot \vec{P}_{h\perp})(\vec{p}_T \cdot \vec{k}_T) - \vec{p}_T^2 (\vec{k}_T \cdot \vec{P}_{h\perp})}{2 M^2 M_h} h_{1T}^\perp H_1^\perp \right] \\
 & + \cos(3\mathbf{j}_h - \mathbf{j}_s) \mathcal{I} \left[ \left( \frac{2(\vec{p}_T \cdot \vec{P}_{h\perp})^2 (\vec{P}_{h\perp} \times \vec{k}_T) + 2(\vec{k}_T \cdot \vec{P}_{h\perp})(\vec{p}_T \cdot \vec{P}_{h\perp})(\vec{P}_{h\perp} \times \vec{p}_T)}{2 M^2 M_h} \right. \right. \\
 & \left. \left. - \frac{\vec{p}_T^2 (\vec{P}_{h\perp} \times \vec{k}_T)}{2 M^2 M_h} \right) h_{1T}^\perp H_1^\perp \right] + \frac{|\bar{R}_T|}{M_h} \sin(2\mathbf{j}_h + \mathbf{j}_R - \mathbf{j}_s) \mathcal{I} \left[ \frac{2(\vec{p}_T \cdot \vec{P}_{h\perp})^2 - \vec{p}_T^2}{2 M^2} h_{1T}^\perp \bar{H}_1^\perp \right] \\
 & \left. + \frac{|\bar{R}_T|}{M_h} \cos(2\mathbf{j}_h + \mathbf{j}_R - \mathbf{j}_s) \mathcal{I} \left[ \frac{(\vec{p}_T \cdot \vec{P}_{h\perp})(\vec{P}_{h\perp} \times \vec{p}_T)}{2 M^2} h_{1T}^\perp \bar{H}_1^\perp \right] \right\}
 \end{aligned}$$

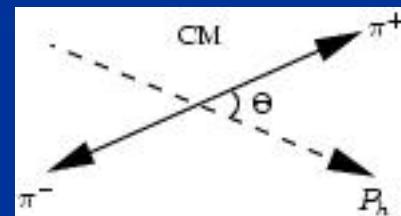
Bacchetta et al. [hep-ph/0212300]

# Single Spin Asymmetry



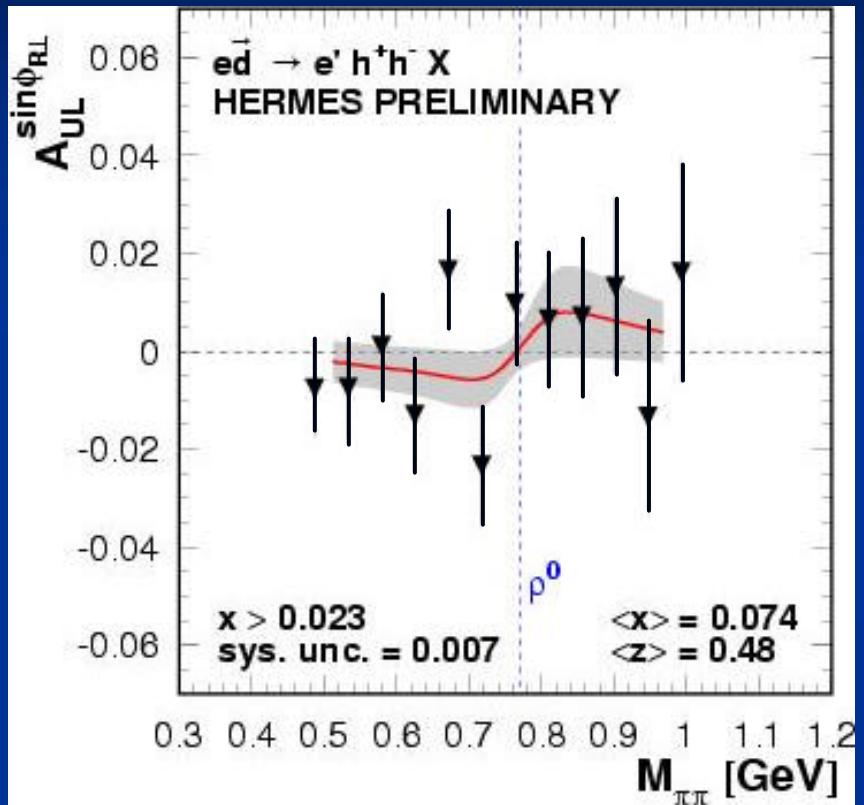
transversely polarized hydrogen target

$$\vec{P}_h \equiv \vec{P}_1 + \vec{P}_2$$



What is measured:

$$A_{UT}(\mathbf{f}_{R\perp}, \mathbf{f}_s, \mathbf{q}) = \frac{1}{S_T} \frac{N^\uparrow(\mathbf{f}_{R\perp}, \mathbf{f}_s, \mathbf{q}) / N_{DIS}^\uparrow - N^\downarrow(\mathbf{f}_{R\perp}, \mathbf{f}_s, \mathbf{q}) / N_{DIS}^\downarrow}{N^\uparrow(\mathbf{f}_{R\perp}, \mathbf{f}_s, \mathbf{q}) / N_{DIS}^\uparrow + N^\downarrow(\mathbf{f}_{R\perp}, \mathbf{f}_s, \mathbf{q}) / N_{DIS}^\downarrow}$$

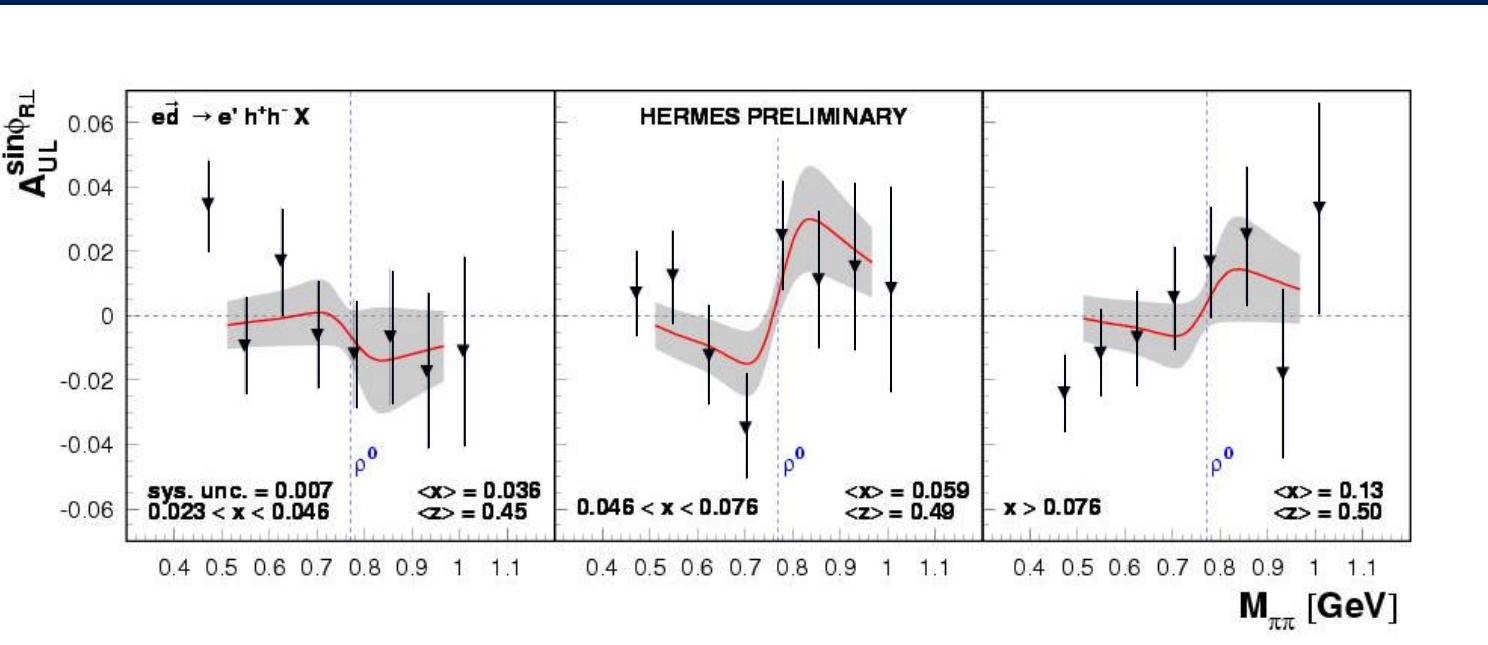


$$c_1 = -0.040 \pm 0.036$$

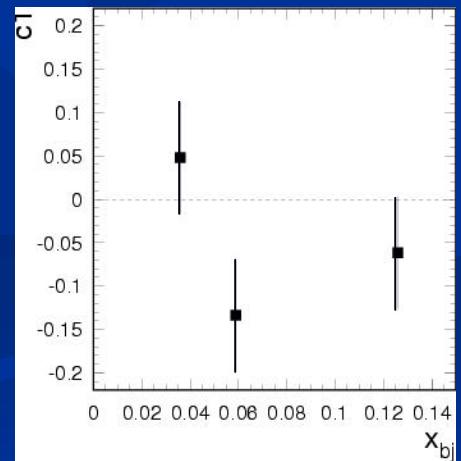
$$c_2 = -0.001 \pm 0.004$$

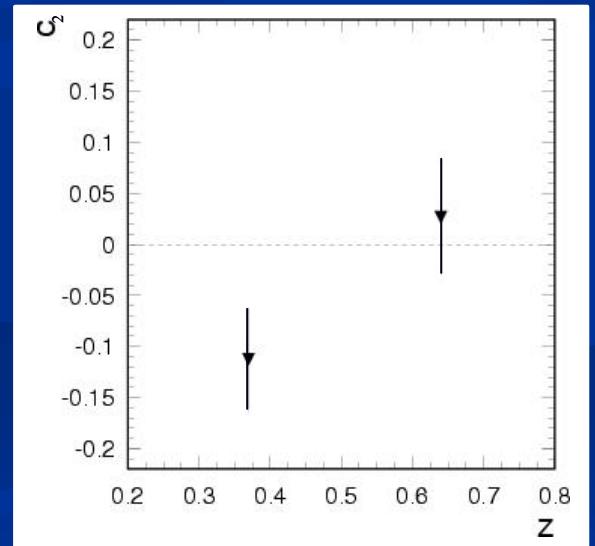
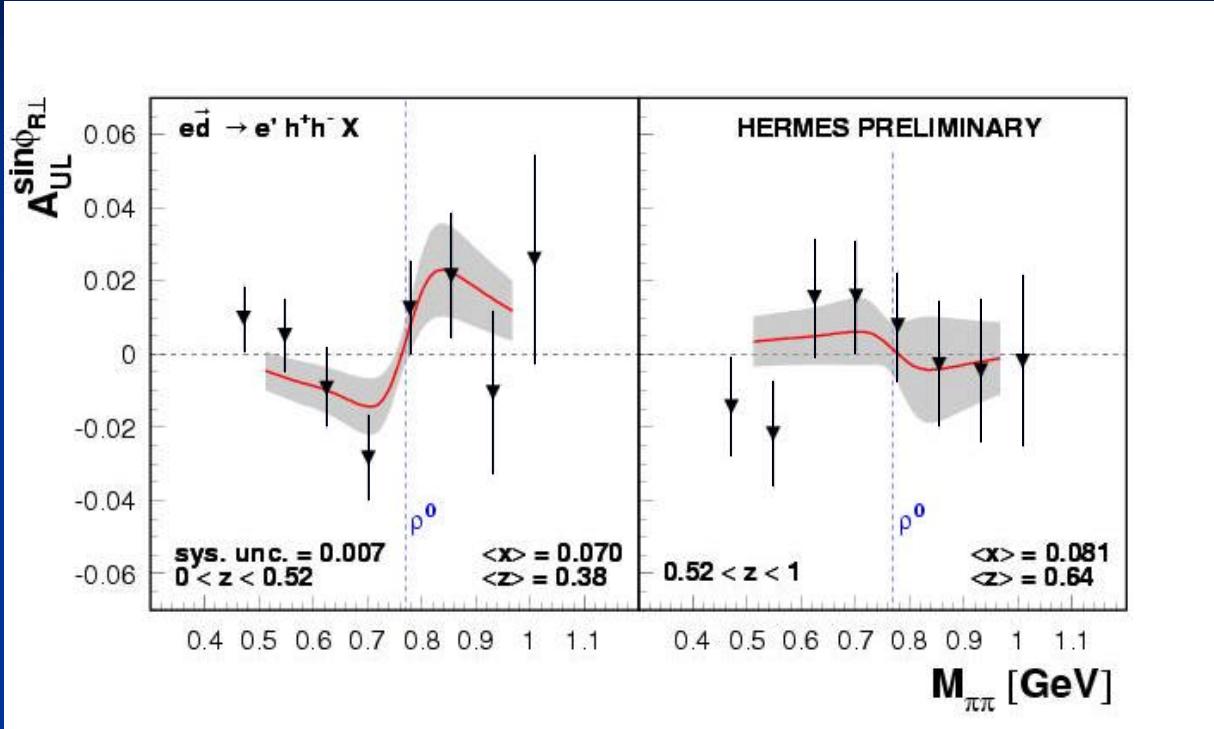
$$g(M_{pp}^2) \simeq c_1 \mathcal{P}(M_{pp}^2) + c_2$$

Hint of a sign change  
at the  $\rho^0$  mass!



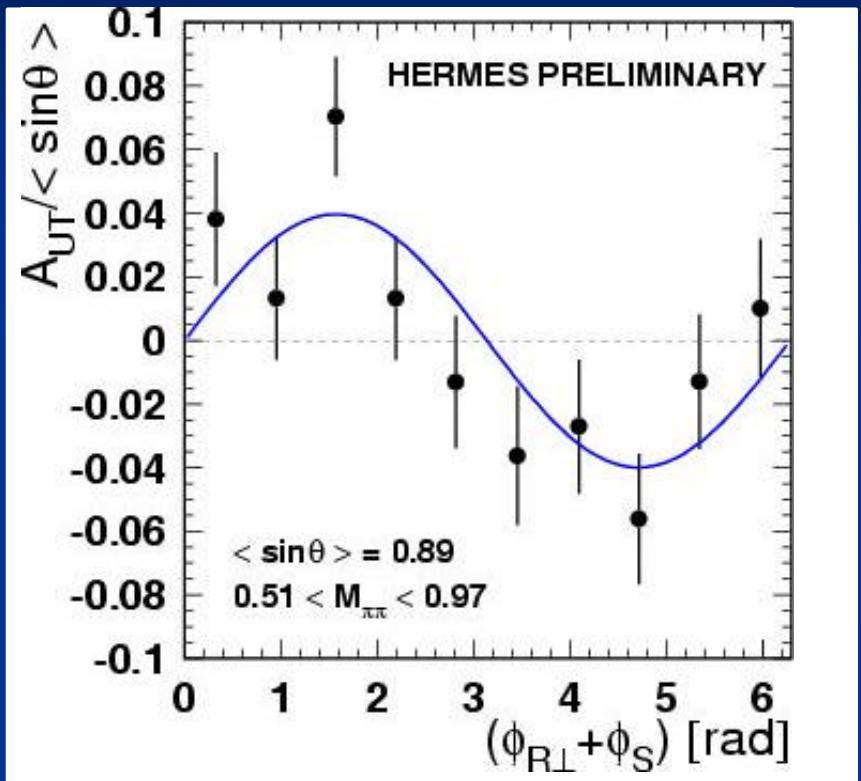
- higher  $x$ : hint of sign change at  $\rho^0$  mass according to Jaffe's model
- $c_1(x) \propto h_l(x)$  ?





- sign change at  $\rho^0$  according to Jaffe's model for low  $z$
- $c_2(z) \propto H_1^{\alpha, sp}(z, M_{pp})$

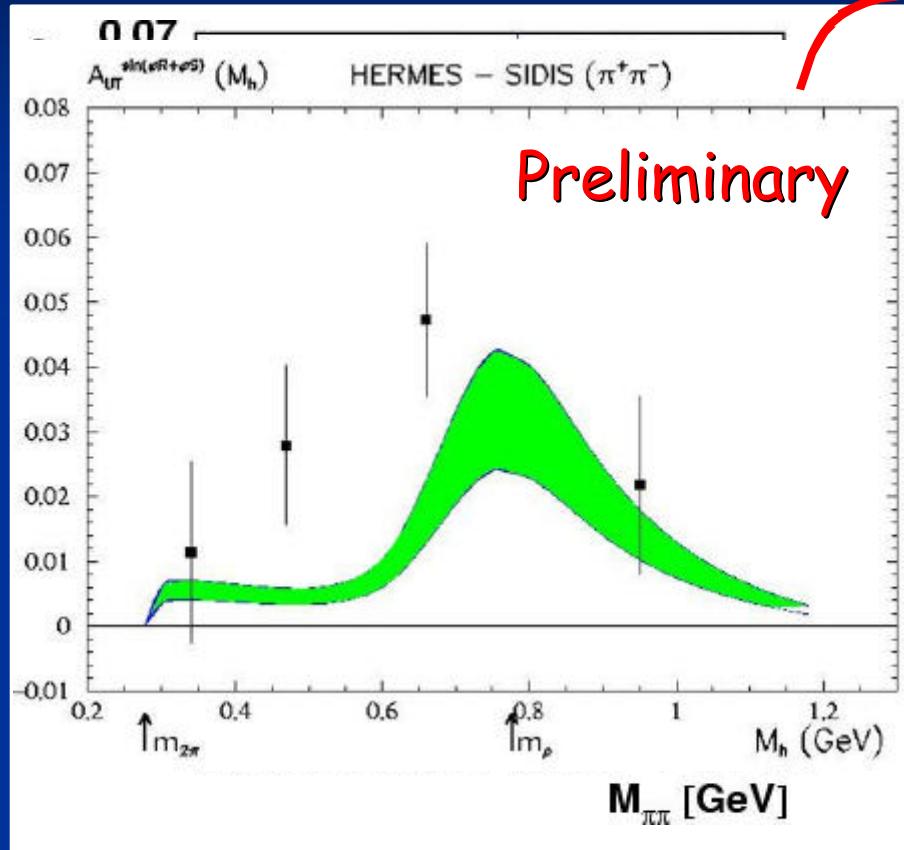
# The $A_{UT}$ asymmetry



Significant  $\sin(f_R + f_S)$  behavior!

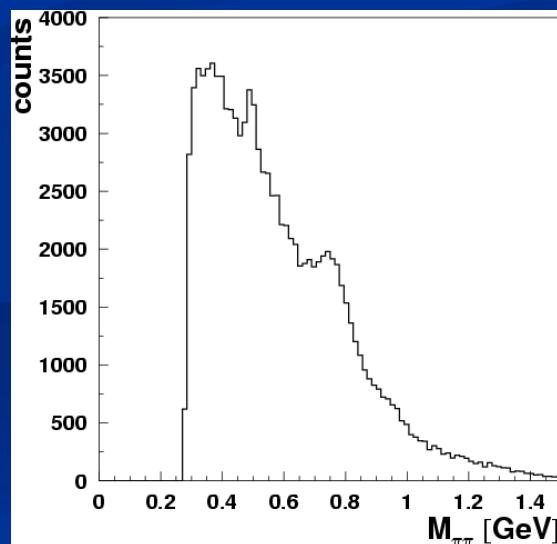
$$A_{UT}^{\sin(f_R + f_S) \sin q} = 0.040 \pm 0.009 \text{ (stat)} \pm 0.003 \text{ (syst)}$$

# Invariant mass dependence



see f.i. M. Radici, hep-ph/0510165

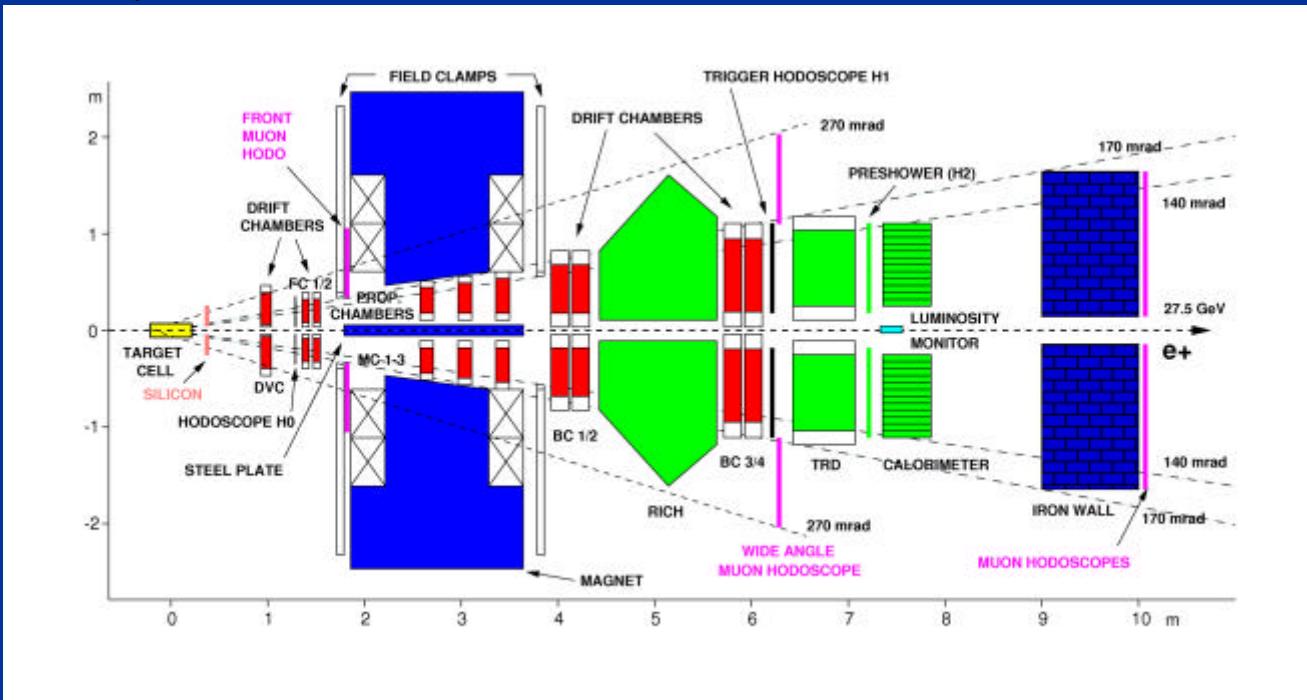
- positive asymmetry moments for all invariant mass bins
- result rules out predicted sign change at the  $\rho^0$  mass (Jaffe et al.)



# Status of the analysis

Dealing with low statistics:

- forced to integrate the asymmetry over many variables combined with the HERMES acceptance: watch out for acceptance effects!



tion of the  
fit



# Conclusions & Outlook



## Conclusions:

- A (significantly) non-zero asymmetry-moment has been measured providing evidence for a non-zero interference fragmentation function.
- This also implies that interference fragmentation can be used to study transversity!
- The new results using a transversely polarized hydrogen target rule out the invariant mass behavior as predicted by R. Jaffe.

## Outlook:

- Increase the statistics using the 2005 data
- Extract asymmetry moment relating to  $H_1^{\star,pp}$
- Extract quantitative information on  $h_1 H_1^\star$