



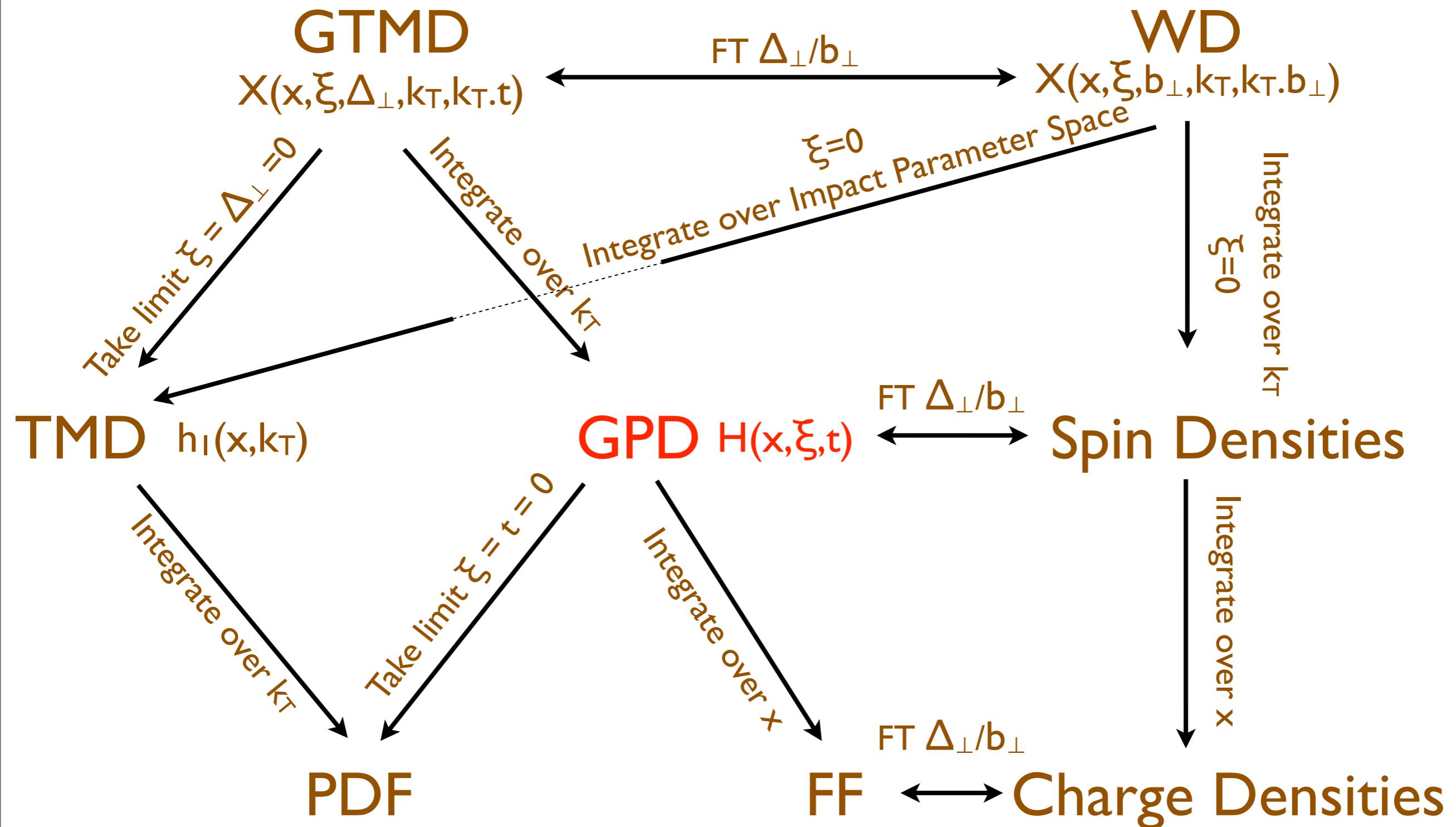
University  
ofGlasgow

# DVCS @ HERMES

M. MURRAY, UNIVERSITY OF GLASGOW  
SPIN 2012



# Distribution Graph



# GPD Physics

GPDs describe only the soft part of the interaction

Accessed via cross-sections and asymmetries:  
requires convolution with a hard scattering kernel

$$H \rightarrow \mathcal{H} \quad \tilde{H} \rightarrow \tilde{\mathcal{H}} \quad E \rightarrow \mathcal{E} \quad \tilde{E} \rightarrow \tilde{\mathcal{E}}$$

Results in “Compton Form Factors” accessible through DVCS, which have real and imaginary parts

# GPD Physics

GPDs describe only the soft part of the interaction

Accessed via cross-sections and asymmetries:  
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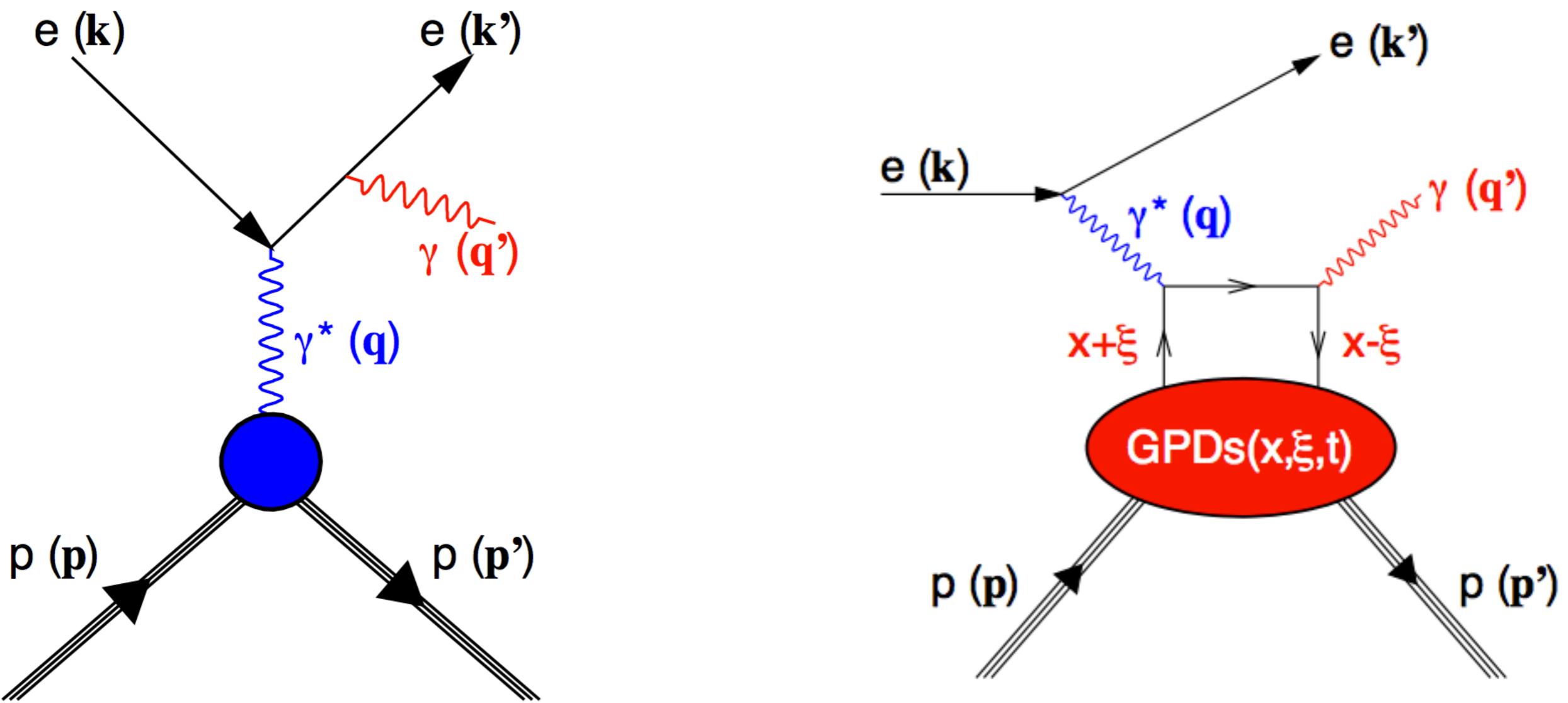
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$$\Im m \mathcal{F}(\xi, t) = F(\xi, \xi, t) \pm F(-\xi, \xi, t),$$

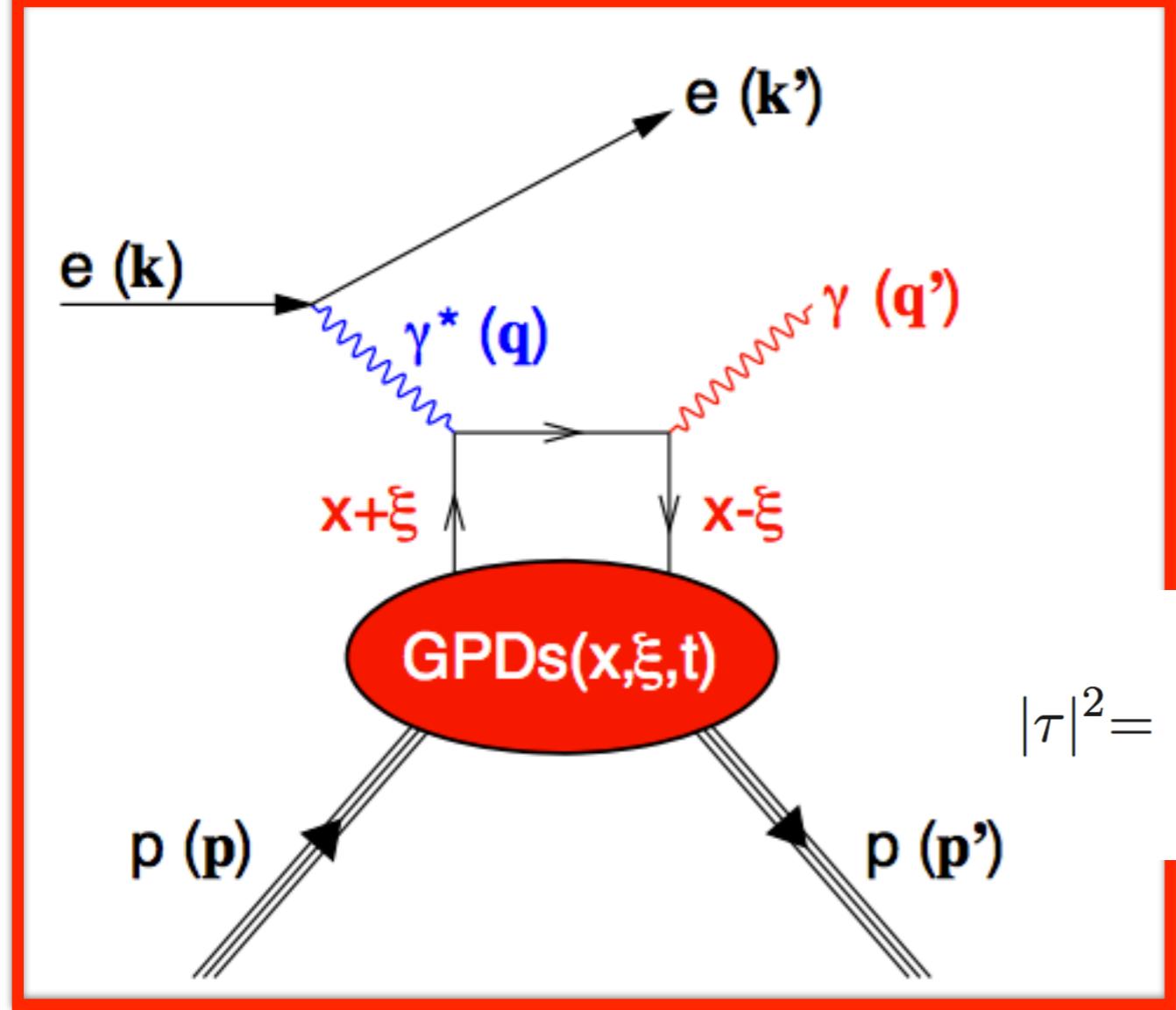
$$\Re e \mathcal{F}(\xi, t) = \mathcal{P}_C \int_{-1}^1 \frac{F(x, \xi, t)}{x - \xi} \pm \frac{F(x, \xi, t)}{x + \xi} dx$$

# Deeply Virtual Compton Scattering

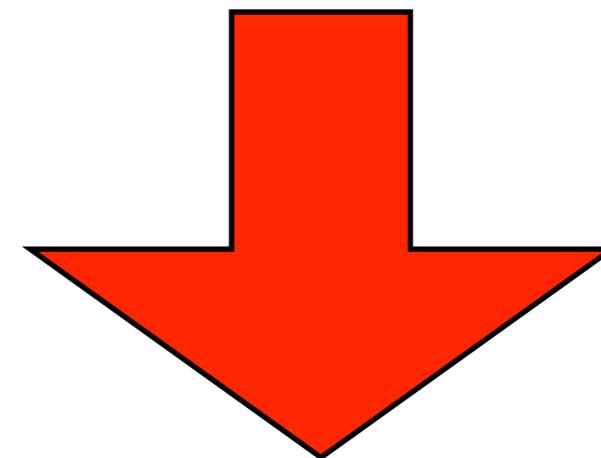
$e p \rightarrow e p \gamma$



# Deeply Virtual Compton Scattering



$$\frac{d\sigma}{dx_B dQ^2 d|t| d\phi} = \frac{x_B e^6 |\tau|^2}{32(2\pi)^4 Q^4 \sqrt{1 + \epsilon^2}}$$



$$|\tau|^2 = |\tau_{BH}|^2 + \overbrace{|\tau_{DVCS}|^2 + \tau_{BH}\tau_{DVCS}^* + \tau_{BH}^*\tau_{DVCS}}^{\mathcal{I}}$$

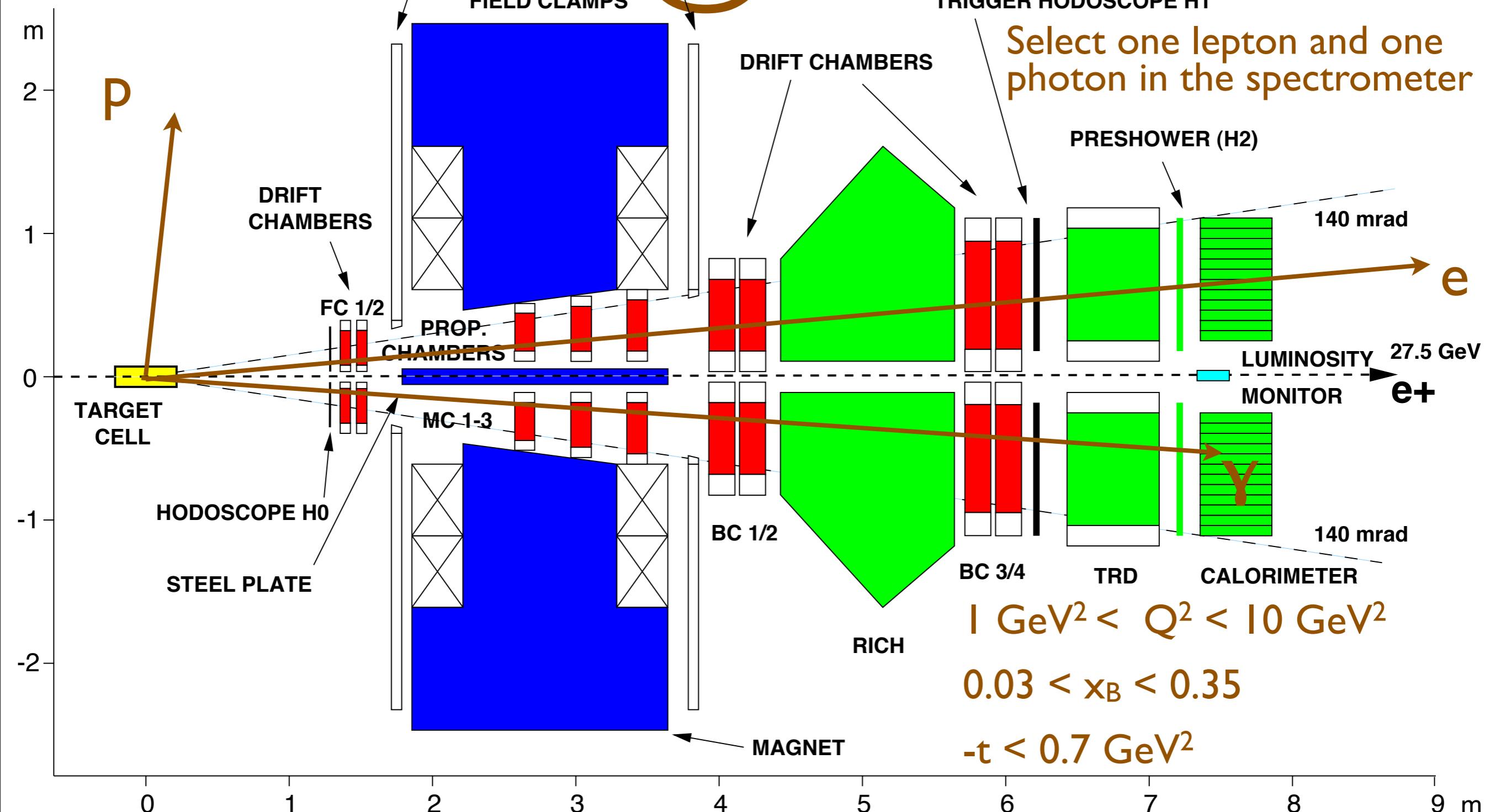
# DVCS @ HERMES

$\mathcal{A}_C(\phi) \equiv \frac{d\sigma^+(\phi) - d\sigma^-(\phi)}{d\sigma^+(\phi) + d\sigma^-(\phi)}$	$\approx$	$\text{Re}(\mathcal{H})$
$\mathcal{A}_{\text{LU}}^{\text{I}}(\phi) \equiv \frac{(d\sigma(\phi)^{+\rightarrow} - d\sigma(\phi)^{+\leftarrow}) - (d\sigma(\phi)^{-\rightarrow} - d\sigma(\phi)^{-\leftarrow})}{(d\sigma(\phi)^{+\rightarrow} + d\sigma(\phi)^{+\leftarrow}) + (d\sigma(\phi)^{-\rightarrow} + d\sigma(\phi)^{-\leftarrow})}$	$\approx$	$\text{Im}(\mathcal{H})$
$\mathcal{A}_{\text{LU}}^{\text{DVCS}}(\phi) \equiv \frac{(d\sigma(\phi)^{+\rightarrow} + d\sigma(\phi)^{-\rightarrow}) - (d\sigma(\phi)^{+\leftarrow} + d\sigma(\phi)^{-\leftarrow})}{(d\sigma(\phi)^{+\rightarrow} + d\sigma(\phi)^{-\rightarrow}) + (d\sigma(\phi)^{+\leftarrow} + d\sigma(\phi)^{-\leftarrow})}$	$\approx$	$\text{Im}[\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*]$
$\mathcal{A}_{\text{UT}}^{\text{I}}(\phi, \phi_S) \equiv \frac{d\sigma^+(\phi, \phi_S) - d\sigma^+(\phi, \phi_S + \pi) - d\sigma^-(\phi, \phi_S) + d\sigma^-(\phi, \phi_S + \pi)}{d\sigma^+(\phi, \phi_S) + d\sigma^+(\phi, \phi_S + \pi) + d\sigma^-(\phi, \phi_S) + d\sigma^-(\phi, \phi_S + \pi)}$	$\approx$	$\text{Im}(E)$
$\mathcal{A}_{\text{UT}}^{\text{DVCS}}(\phi, \phi_S) \equiv \frac{d\sigma^+(\phi, \phi_S) - d\sigma^+(\phi, \phi_S + \pi) + d\sigma^-(\phi, \phi_S) - d\sigma^-(\phi, \phi_S + \pi)}{d\sigma^+(\phi, \phi_S) + d\sigma^+(\phi, \phi_S + \pi) + d\sigma^-(\phi, \phi_S) + d\sigma^-(\phi, \phi_S + \pi)}$	$\approx$	$\text{Im}(E)$
$\mathcal{A}_{\text{LT}}^{\text{BH+DVCS}}(\phi, \phi_S) \equiv \frac{1}{8d\sigma_{\text{UU}}} [(d\vec{\sigma}^{+\uparrow} - d\vec{\sigma}^{+\downarrow} - d\vec{\sigma}^{-\uparrow} + d\vec{\sigma}^{-\downarrow}) + (d\vec{\sigma}^{-\uparrow} - d\vec{\sigma}^{-\downarrow} - d\vec{\sigma}^{+\uparrow} + d\vec{\sigma}^{+\downarrow})]$	$\approx$	$\text{Re}(\mathcal{H} + E)$
$\mathcal{A}_{\text{LT}}^{\text{I}}(\phi, \phi_S) \equiv \frac{1}{8d\sigma_{\text{UU}}} [(d\vec{\sigma}^{+\uparrow} - d\vec{\sigma}^{+\downarrow} - d\vec{\sigma}^{-\uparrow} + d\vec{\sigma}^{-\downarrow}) - (d\vec{\sigma}^{-\uparrow} - d\vec{\sigma}^{-\downarrow} - d\vec{\sigma}^{+\uparrow} + d\vec{\sigma}^{+\downarrow})]$	$\approx$	$\text{Re}(\mathcal{H})$
$\mathcal{A}_{\text{UL}}(\phi) \equiv \frac{[\sigma^{\leftarrow\rightarrow}(\phi) + \sigma^{\rightarrow\Rightarrow}(\phi)] - [\sigma^{\leftarrow\leftarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]}{[\sigma^{\leftarrow\rightarrow}(\phi) + \sigma^{\rightarrow\Rightarrow}(\phi)] + [\sigma^{\leftarrow\leftarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]}$	$\approx$	$\text{Im}(\tilde{\mathcal{H}})$
$\mathcal{A}_{\text{LL}}(\phi) \equiv \frac{[\sigma^{\rightarrow\Rightarrow}(\phi) + \sigma^{\leftarrow\leftarrow}(\phi)] - [\sigma^{\leftarrow\Rightarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]}{[\sigma^{\rightarrow\Rightarrow}(\phi) + \sigma^{\leftarrow\leftarrow}(\phi)] + [\sigma^{\leftarrow\Rightarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]}$	$\approx$	$\text{Re}(\tilde{\mathcal{H}})$

# DVCS @ HERMES

FIELD CLAMPS

TRIGGER HODOSCOPE H1

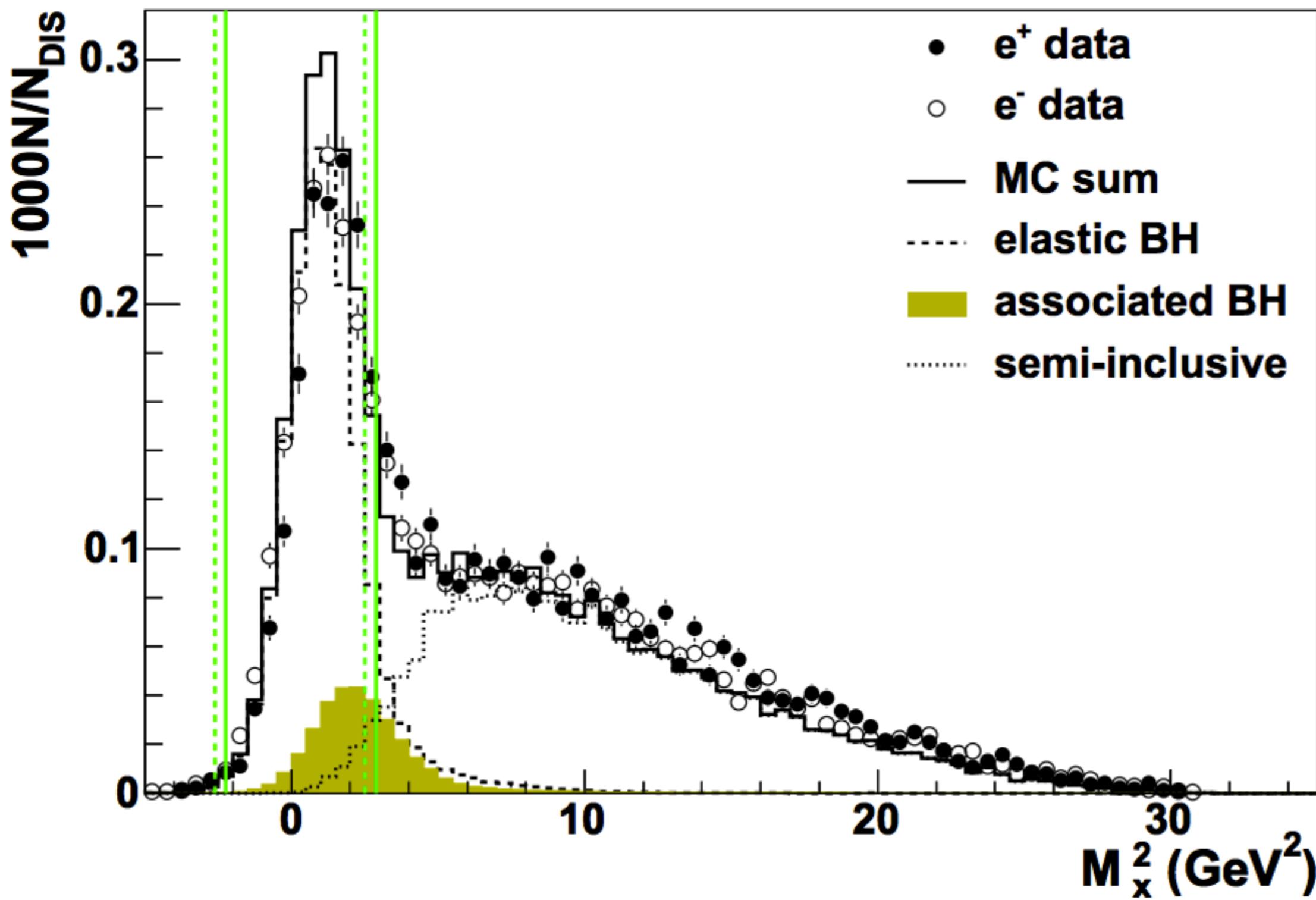


$$\langle Q^2 \rangle \approx 2.4 \text{ GeV}^2$$

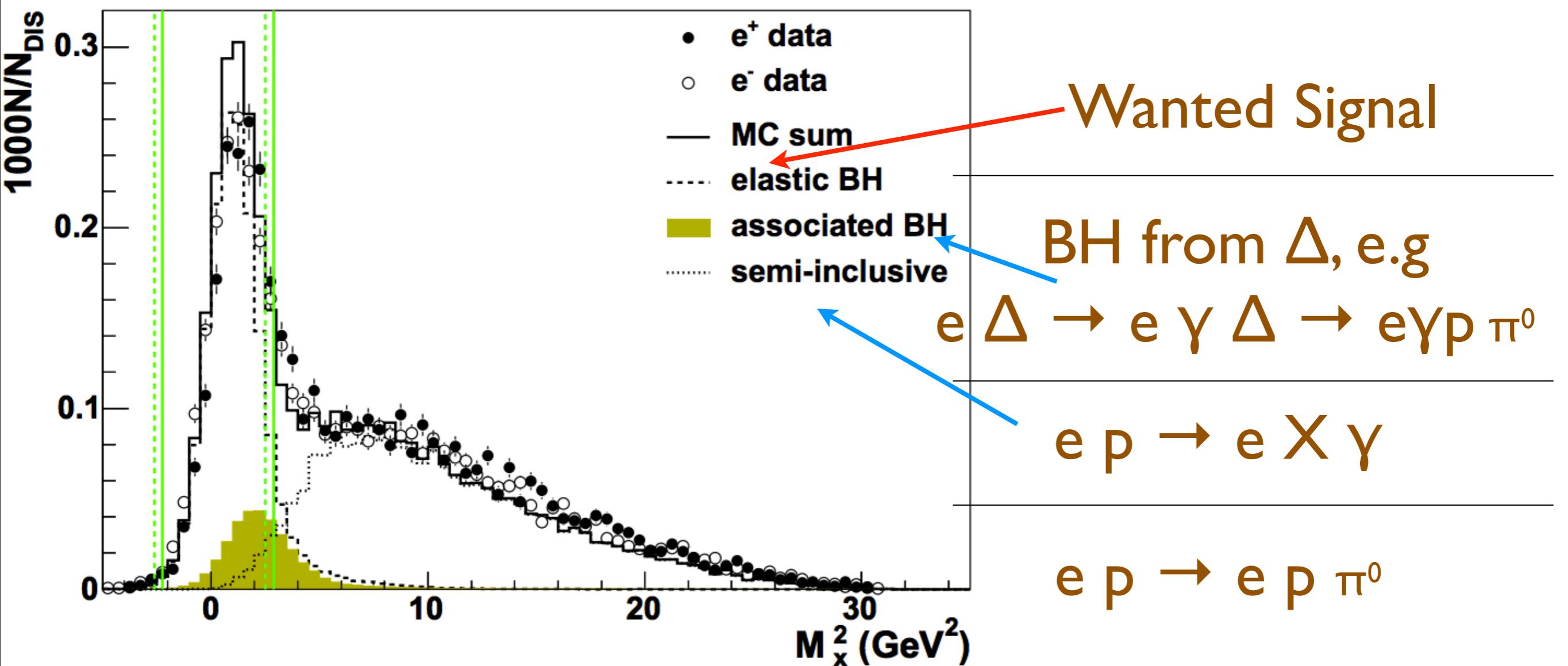
$$\langle x_B \rangle \approx 0.1$$

$$\langle -t \rangle \approx 0.1 \text{ GeV}^2$$

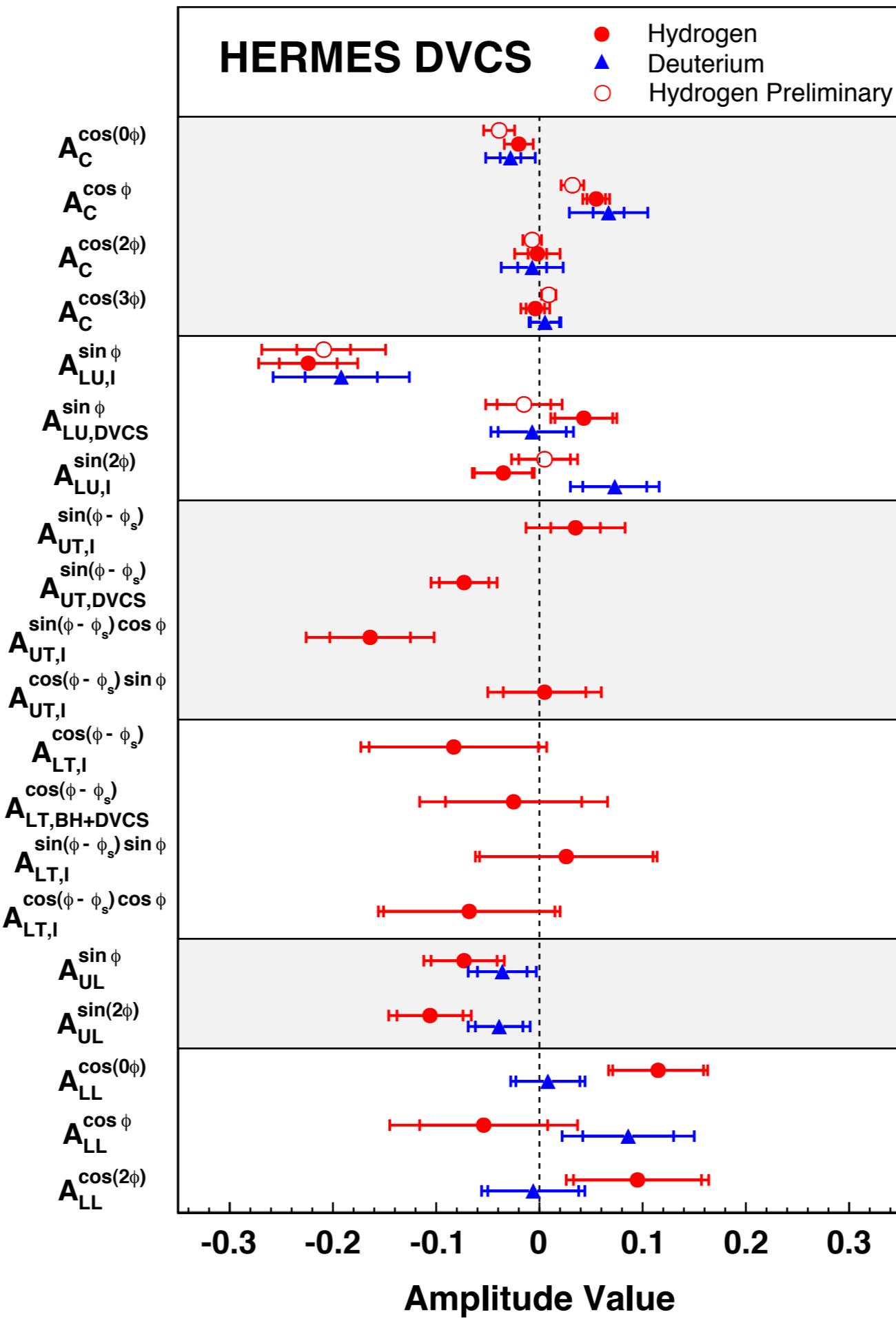
# DVCS @ HERMES



# DVCS @ HERMES



D  
V  
C  
S  
@

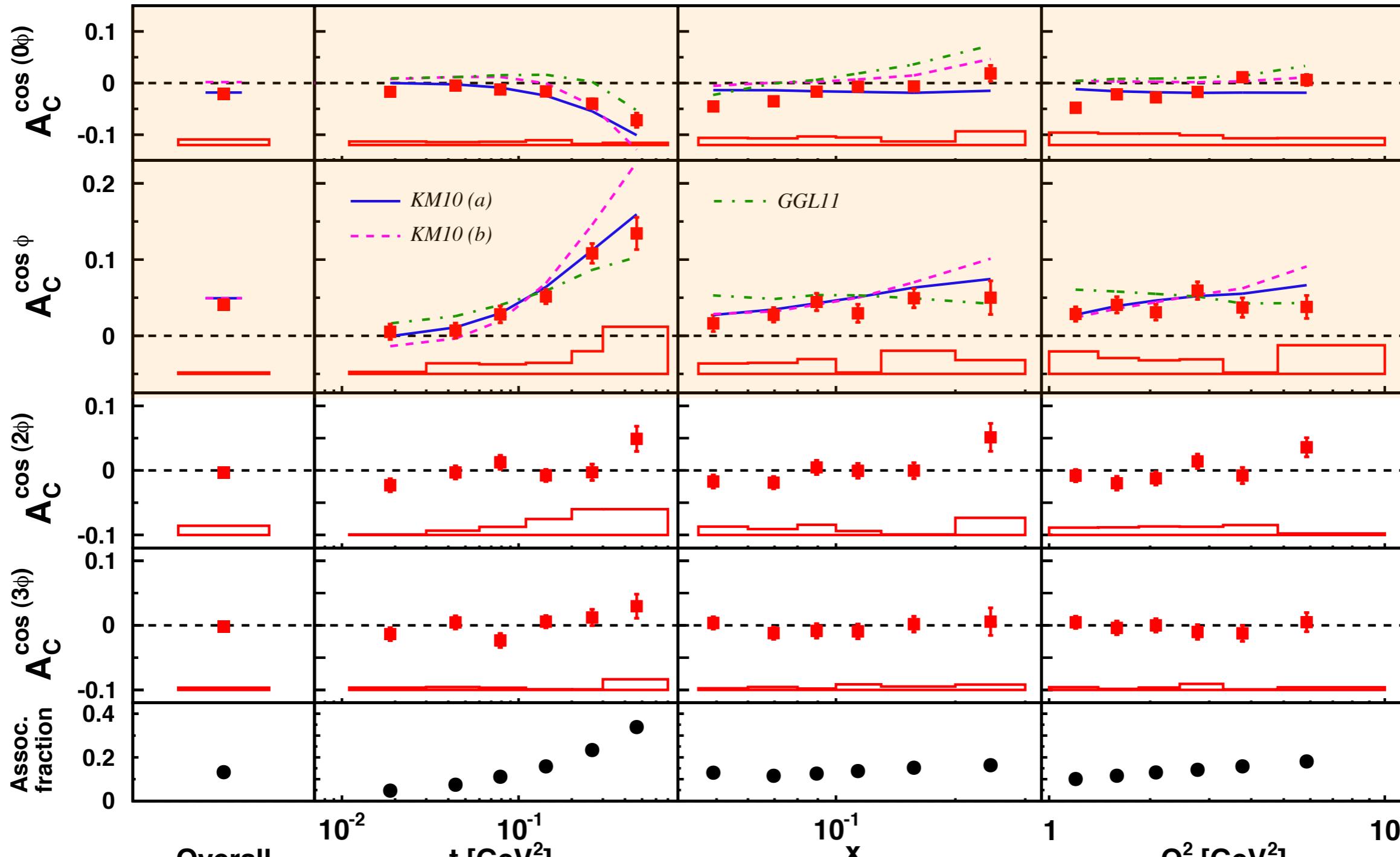


H  
E  
R  
M  
E  
S

# Beam-Charge Asymmetries

A. Airapetian *et al*, JHEP 07 (2012) 032

<http://arxiv.org/abs/1203.6287>



Beam Charge Asymmetries access  $\text{Re}(\mathcal{H})$

G. Goldstein, J. Hernandez and S. Liuti, *Phys. Rev. D84* (2011)

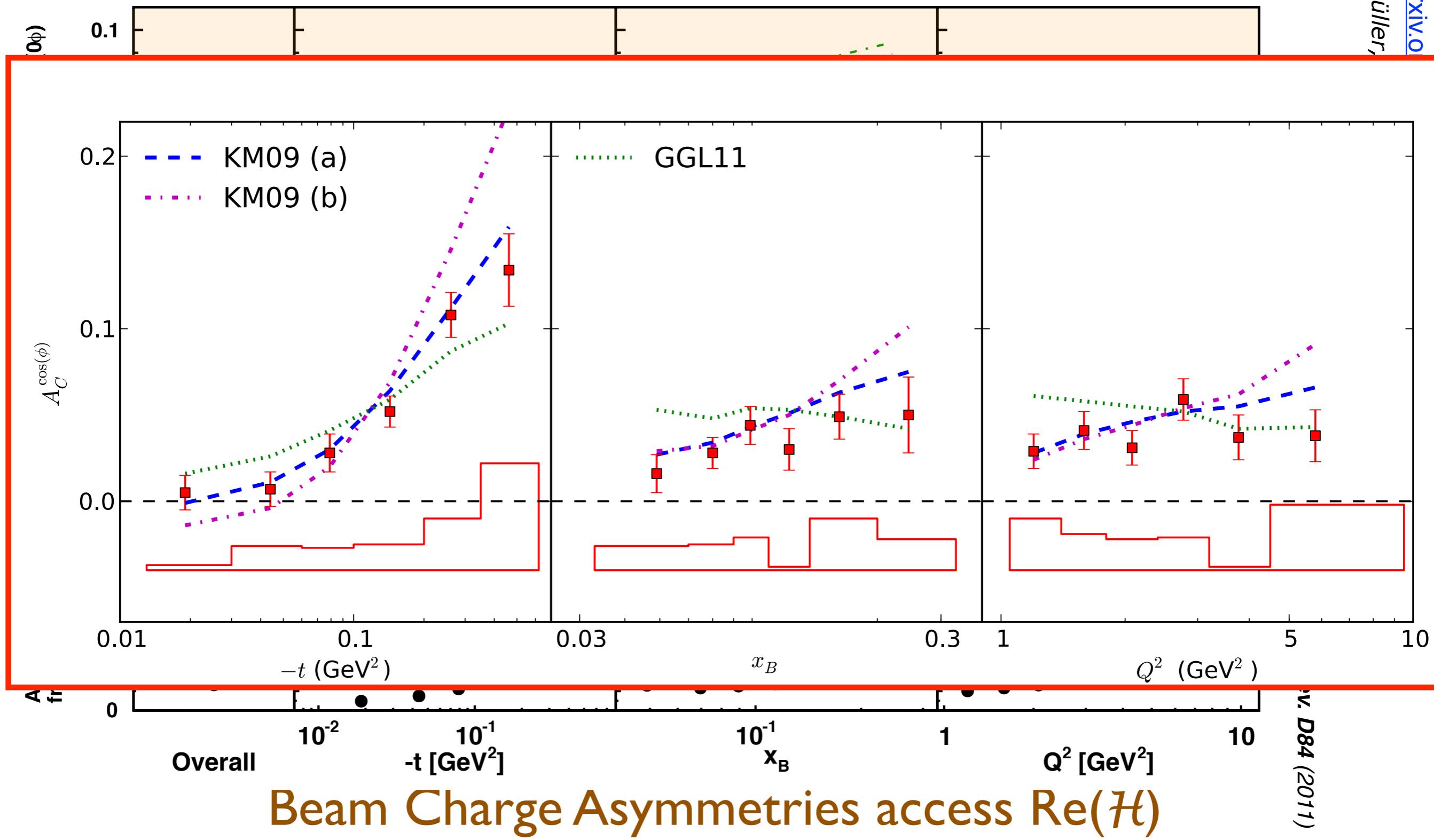
<http://arxiv.org/abs/1012.3776>

Kumerički and Müller, *Nucl. Phys. B841* (2010)

<http://arxiv.org/abs/0904.0458>

# Beam-Charge Asymmetries

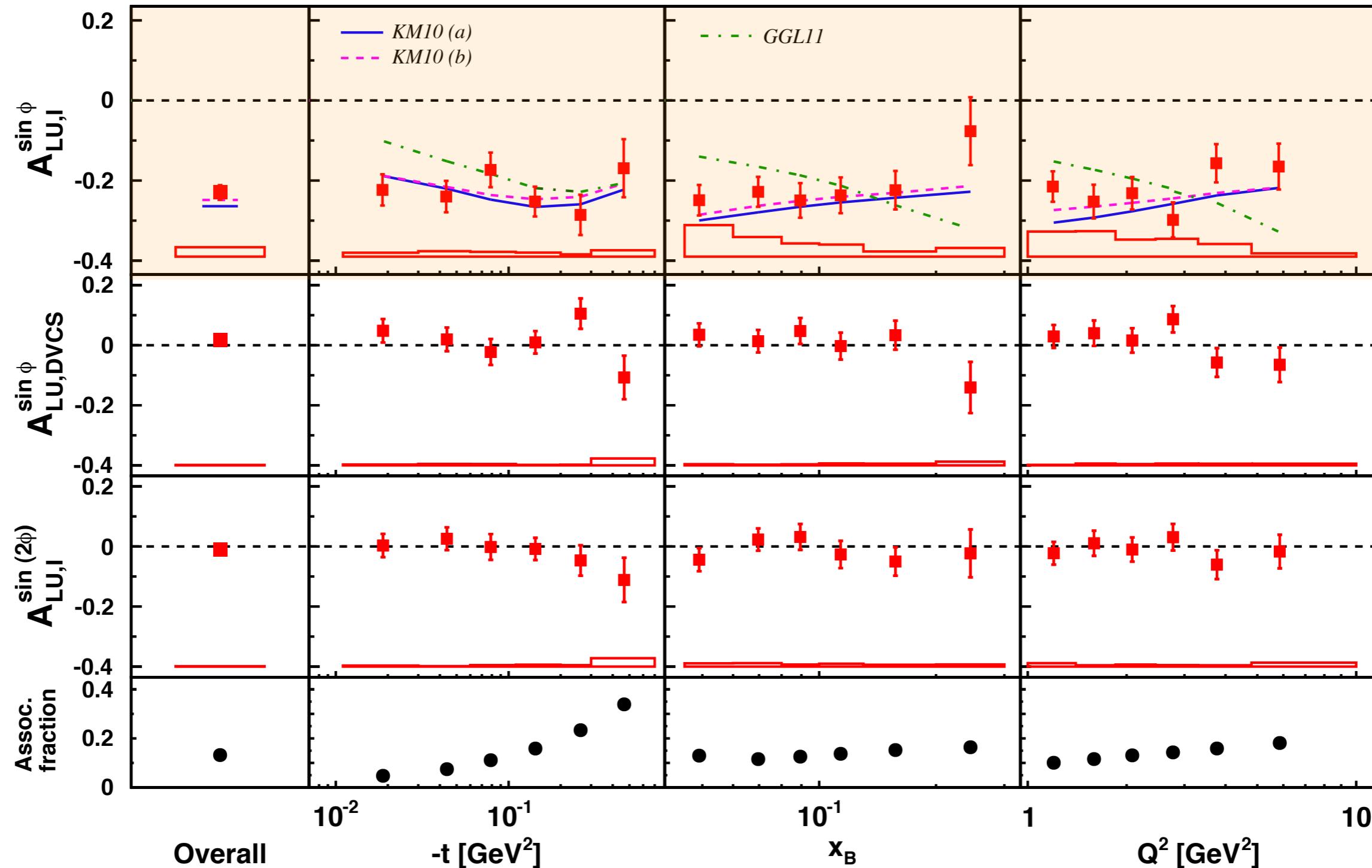
<http://arxiv.org/abs/1203.6287>



# Beam-Spin Asymmetries

A. Airapetian *et al*, JHEP 07 (2012) 032

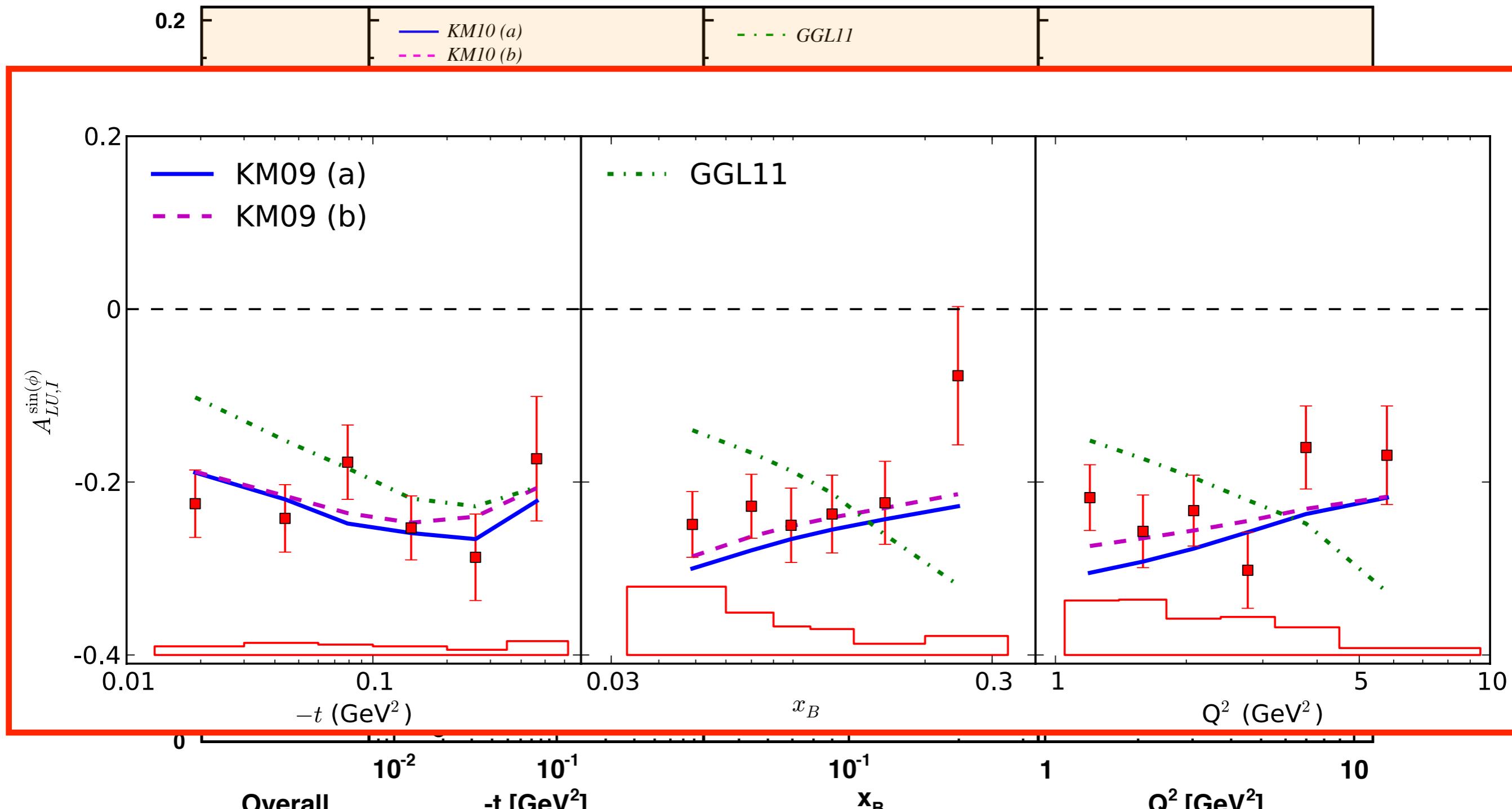
<http://arxiv.org/abs/1203.6287>



Beam Helicity Asymmetries access  $\text{Im}(\mathcal{H})$

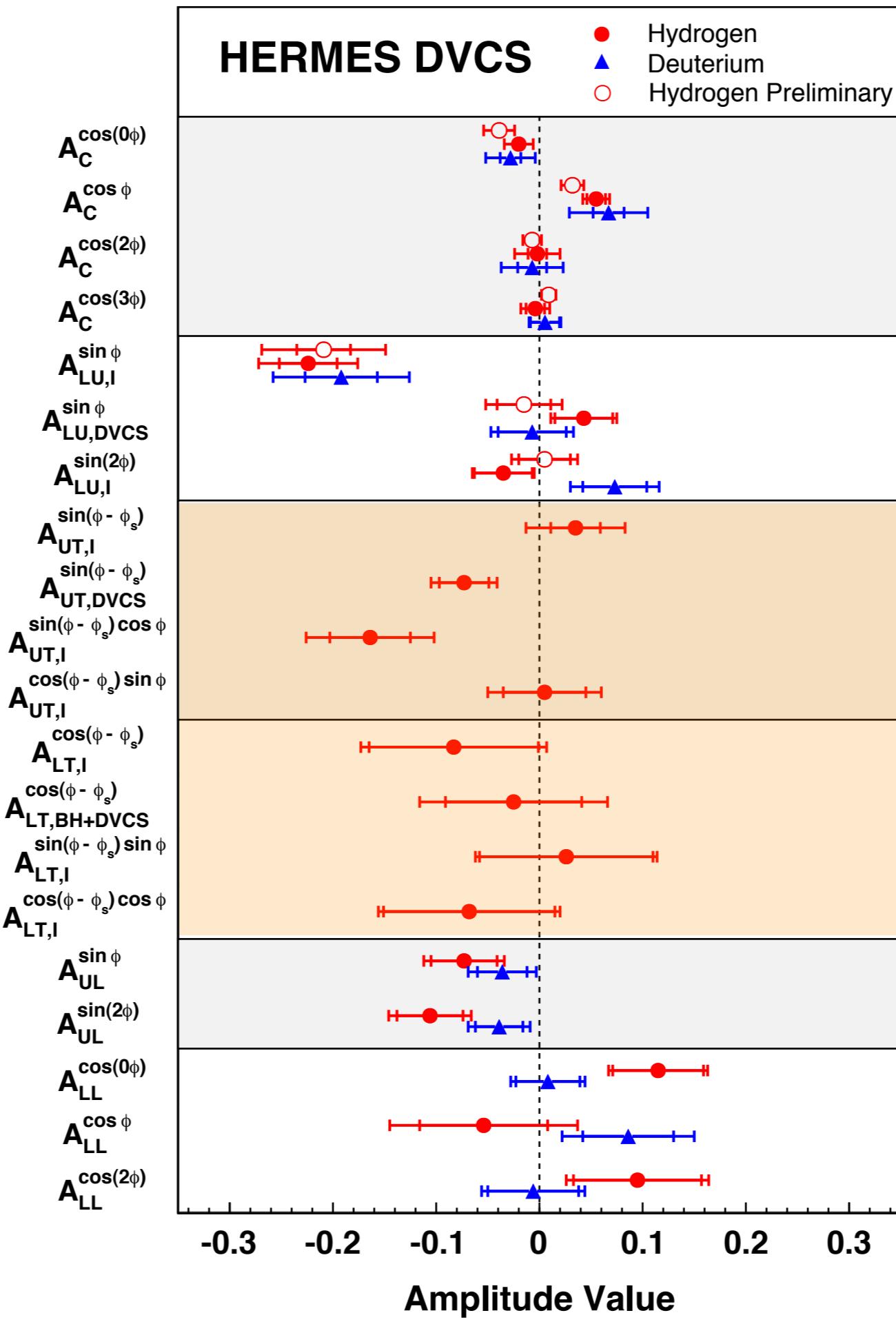
# Beam-Spin Asymmetries

<http://arxiv.org/abs/1203.6287>



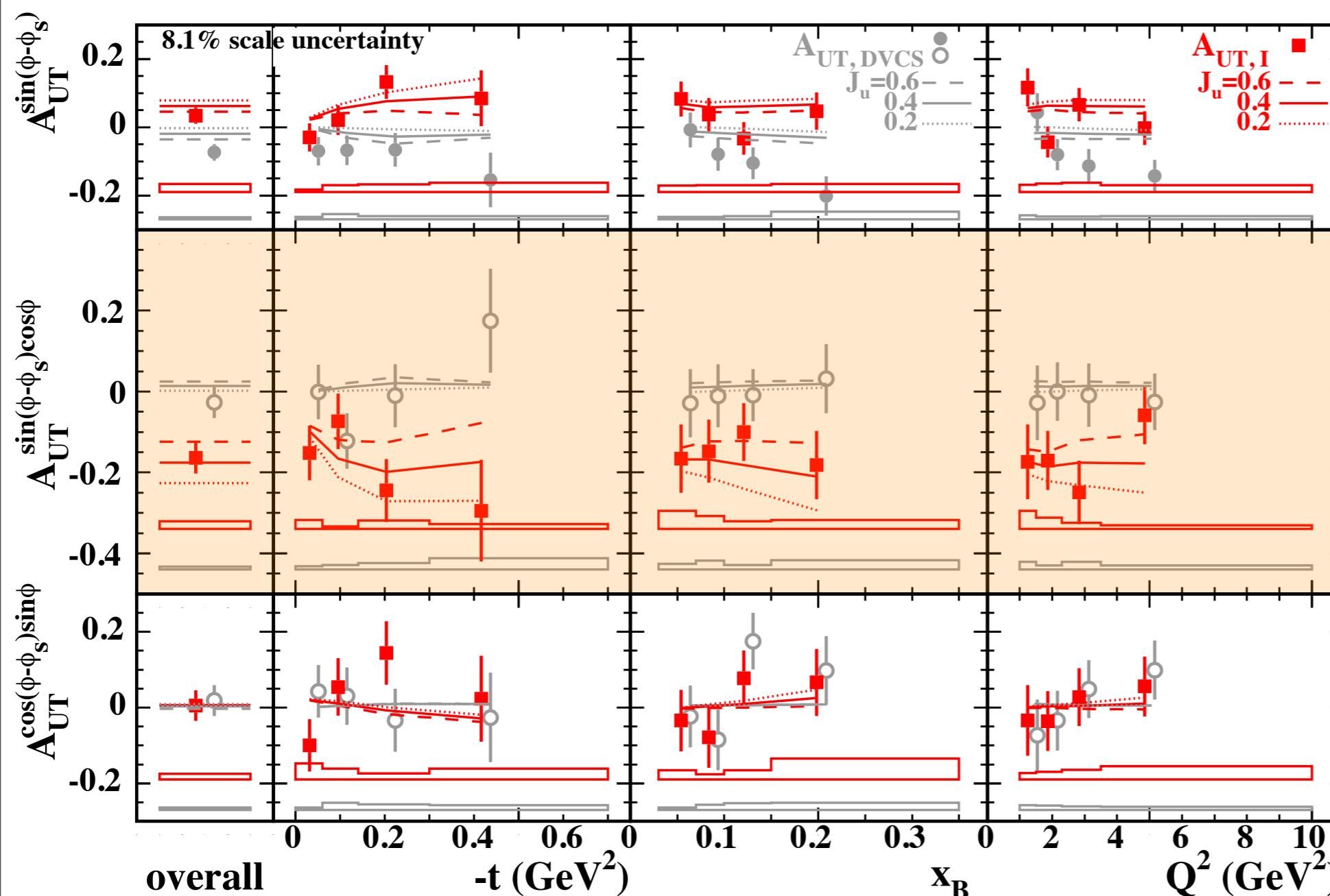
Overall Beam Helicity Asymmetries access  $\text{Im}(\mathcal{H})$

# D V C S @



H  
E  
R  
M  
E  
S

# Transverse-Target Asymmetries



VGG Model:

<http://arxiv.org/abs/hep-ph/9905372>

*Phys.Rev. D60 (1999) 094017*

HERMES Data:

<http://arxiv.org/abs/0802.2499>

A. Airapetian *et al*, JHEP 06 (2008) 066, 24pp

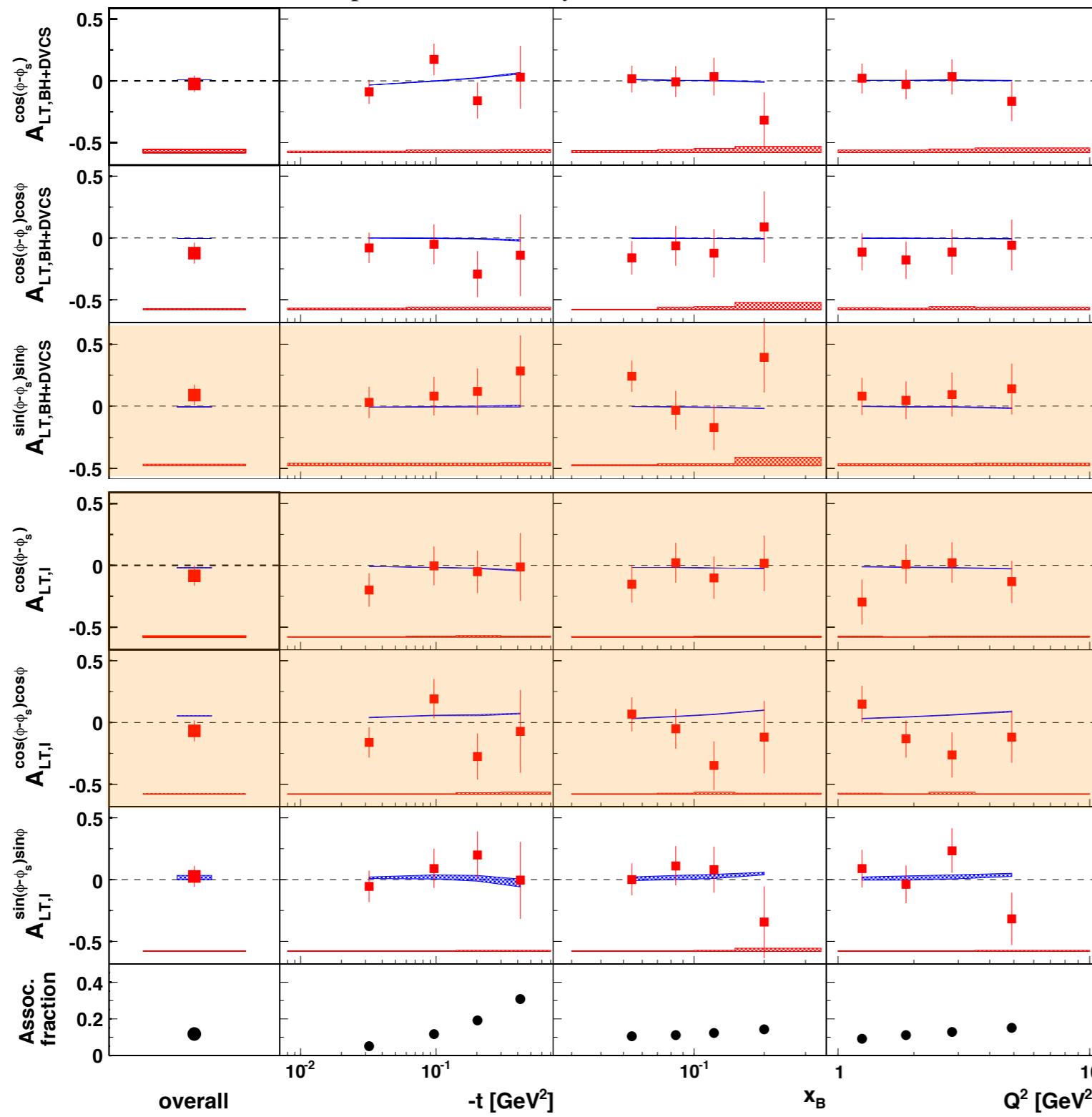
DVCS amplitude  
involves transversity  
GPDs

Transverse Target  
Asymmetries can  
access E?

Pioneering  
measurement to be  
repeated at CLAS12  
and the EIC

# Double-Spin Asymmetries

A. Airapetian et al, Phys. Lett. B 704 (2011) 15-23



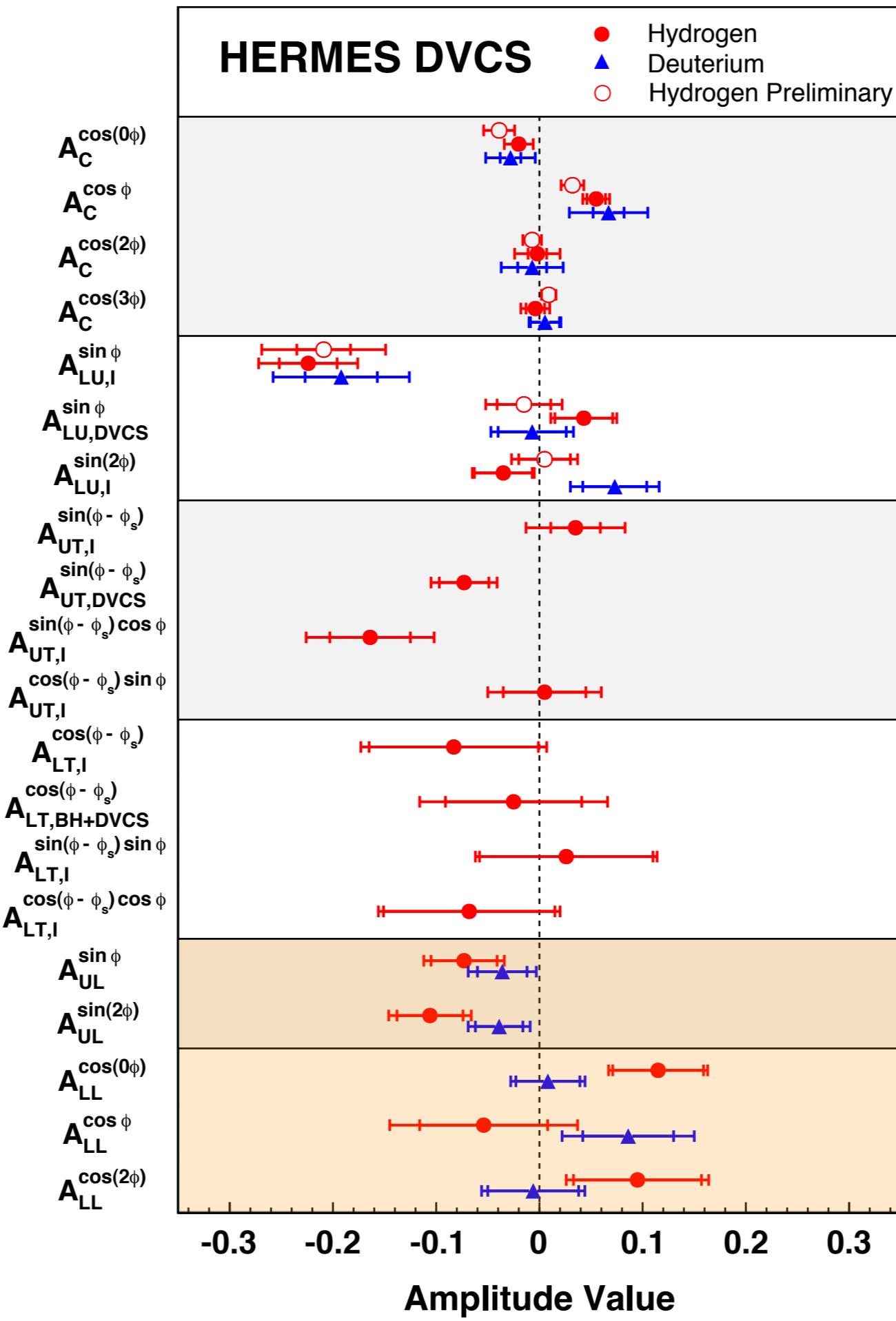
Tran. Pol. target /  
Long. Pol. Beam

Real parts of  $\mathcal{H}$   
and  $\mathcal{E}$

Extracted to be 0;  
compatible with  
VGG predictions.

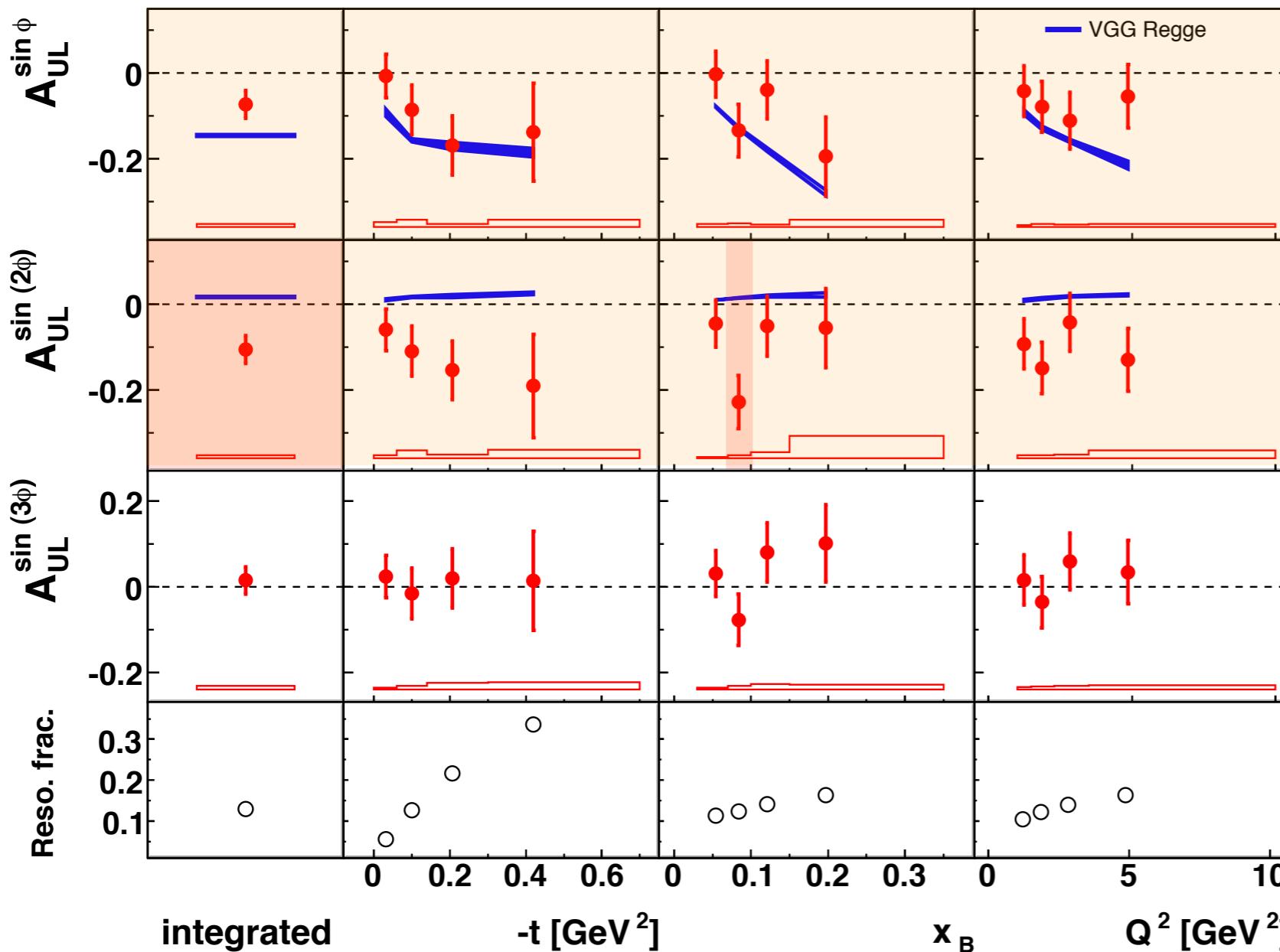
<http://arxiv.org/abs/1106.2990>

D  
V  
C  
S  
@



H  
E  
R  
M  
E  
S

# Longitudinal-Target Asymmetries



Long. Pol. target  
asymmetries  
access  $\text{Im}(\tilde{\mathcal{H}})$

<http://arxiv.org/abs/1004.0177>

*A. Airapetian et al, JHEP 06 (2010) 019*

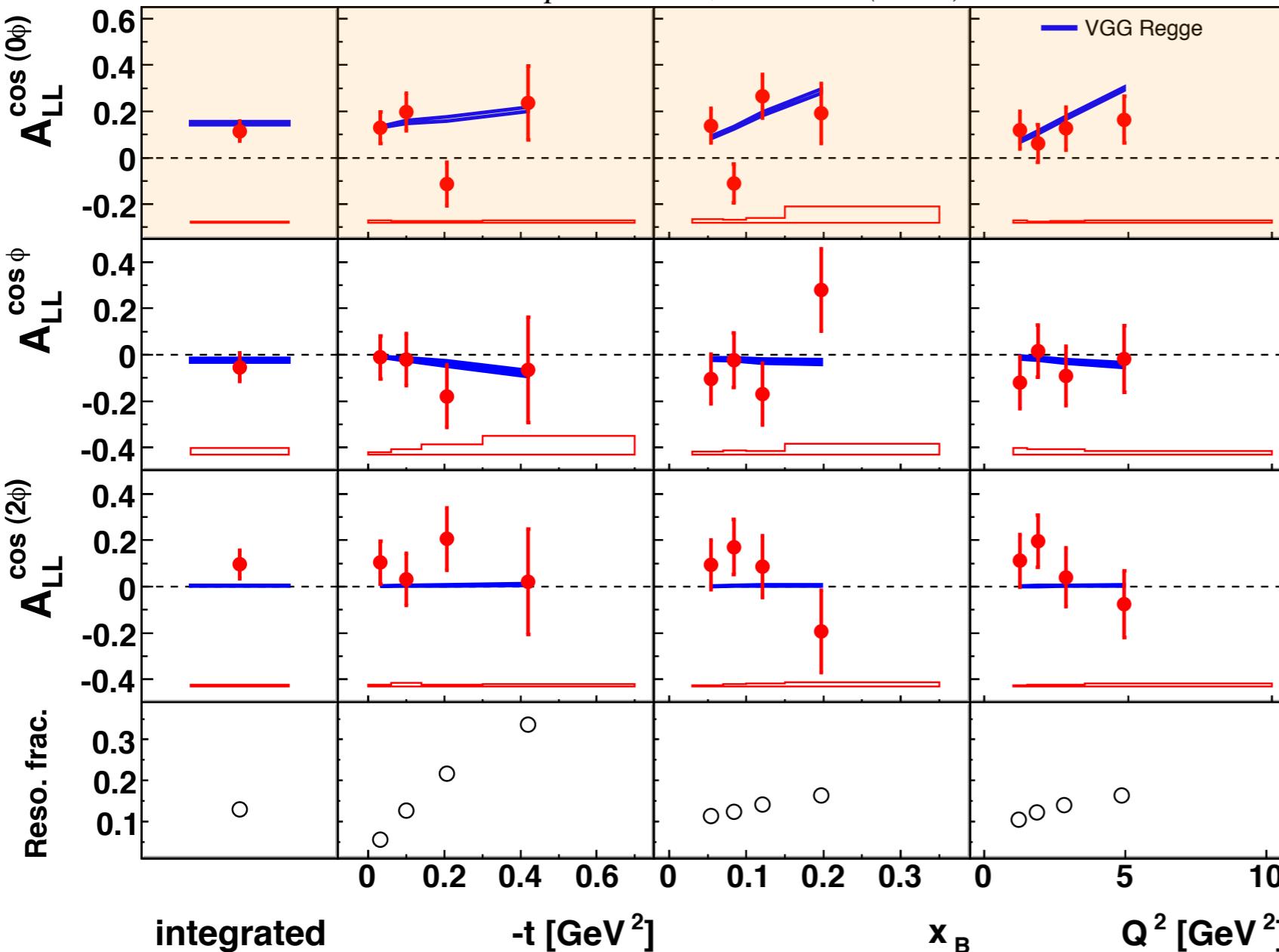
VGG Model:

<http://arxiv.org/abs/hep-ph/9905372>

*Phys.Rev. D60 (1999) 094017*

# Double-Spin Asymmetries

A. Airapetian et al, JHEP 06 (2010) 019

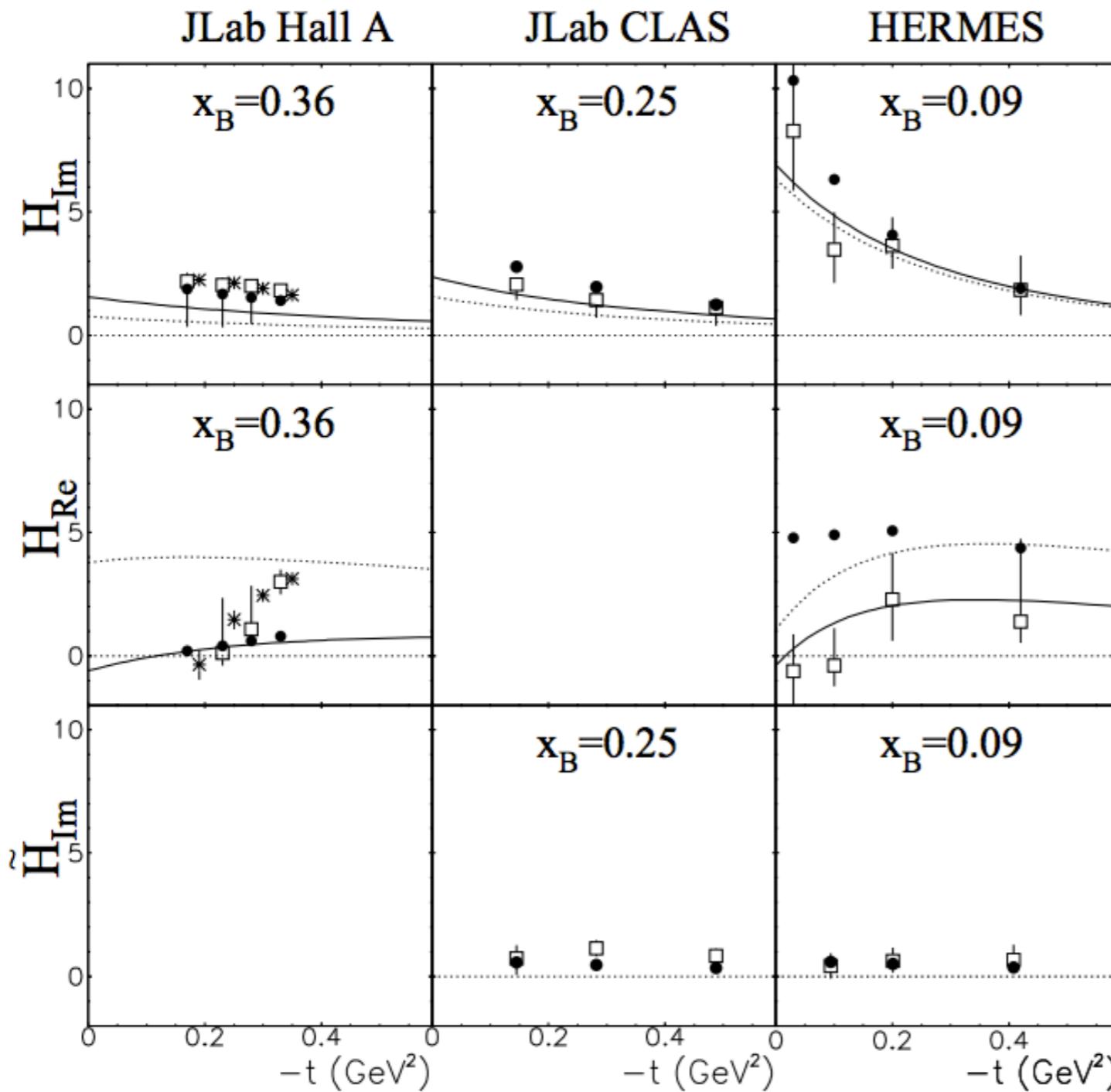


Long. Pol. target /  
Long. Pol. Beam  
access  $\text{Re}(\tilde{\mathcal{H}})$

Caveat! Relatively  
large BH  
contribution to  
these asymmetries!

<http://arxiv.org/abs/1004.0177>

# GPD Extraction



Even for  $H, VGG$  model  
GPDs are shown **not to**  
**be consistent with**  
**experimental**  
**measurements** when  
**CFFs are extracted from**  
**data.**

<http://arxiv.org/abs/1011.4195>

*Guidal, ICHEP Procs. (2010)*

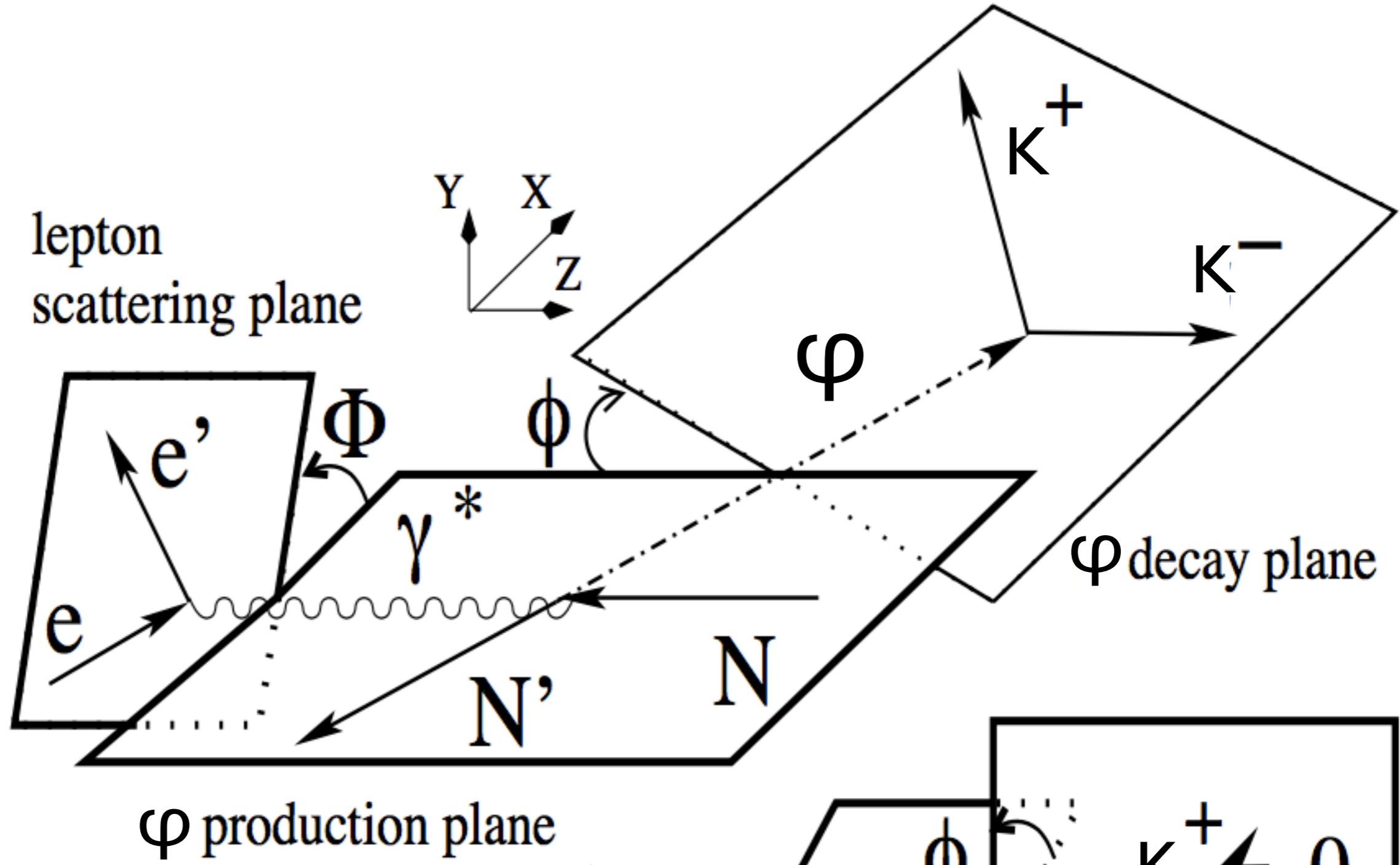
<http://arxiv.org/abs/0904.1648>

*H. Moutarde, Phys. Rev. D79 (2009)*

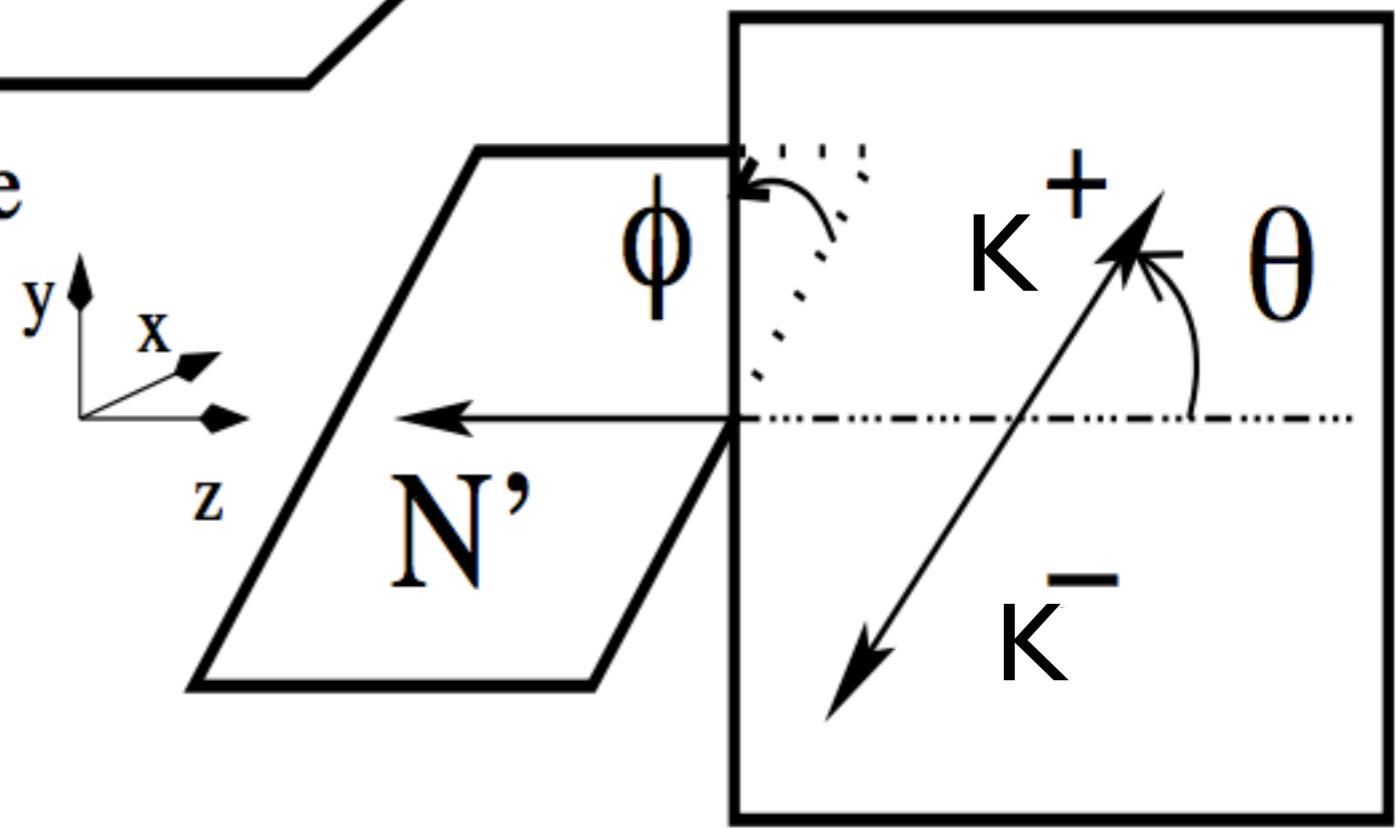
<http://arxiv.org/abs/0904.0458>

*Kumerički and Müller, Nucl. Phys. B841 (2010)*

# Other Data?



# Exclusive $\varphi$ Meson Production

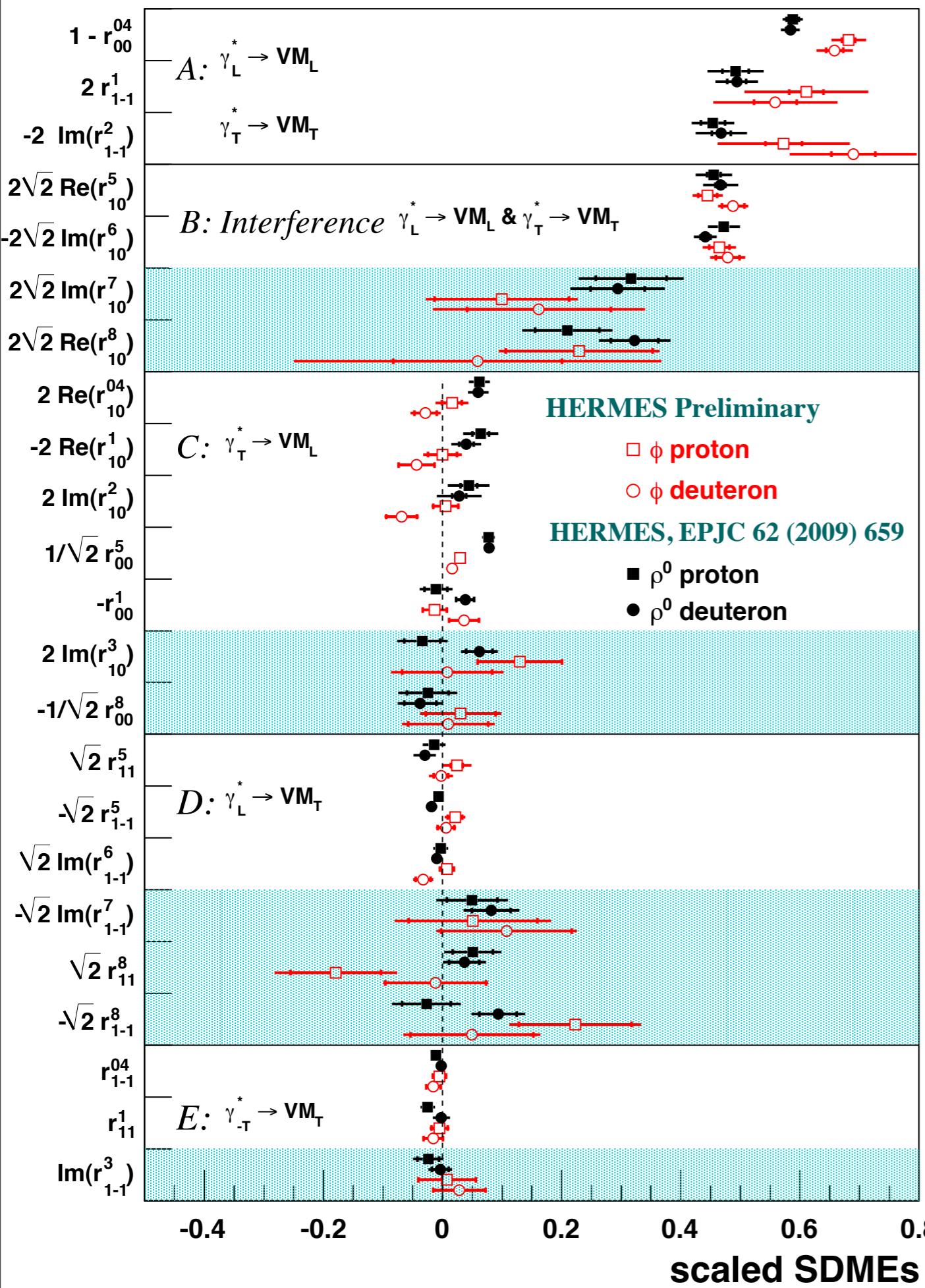


# Exclusive Meson Production

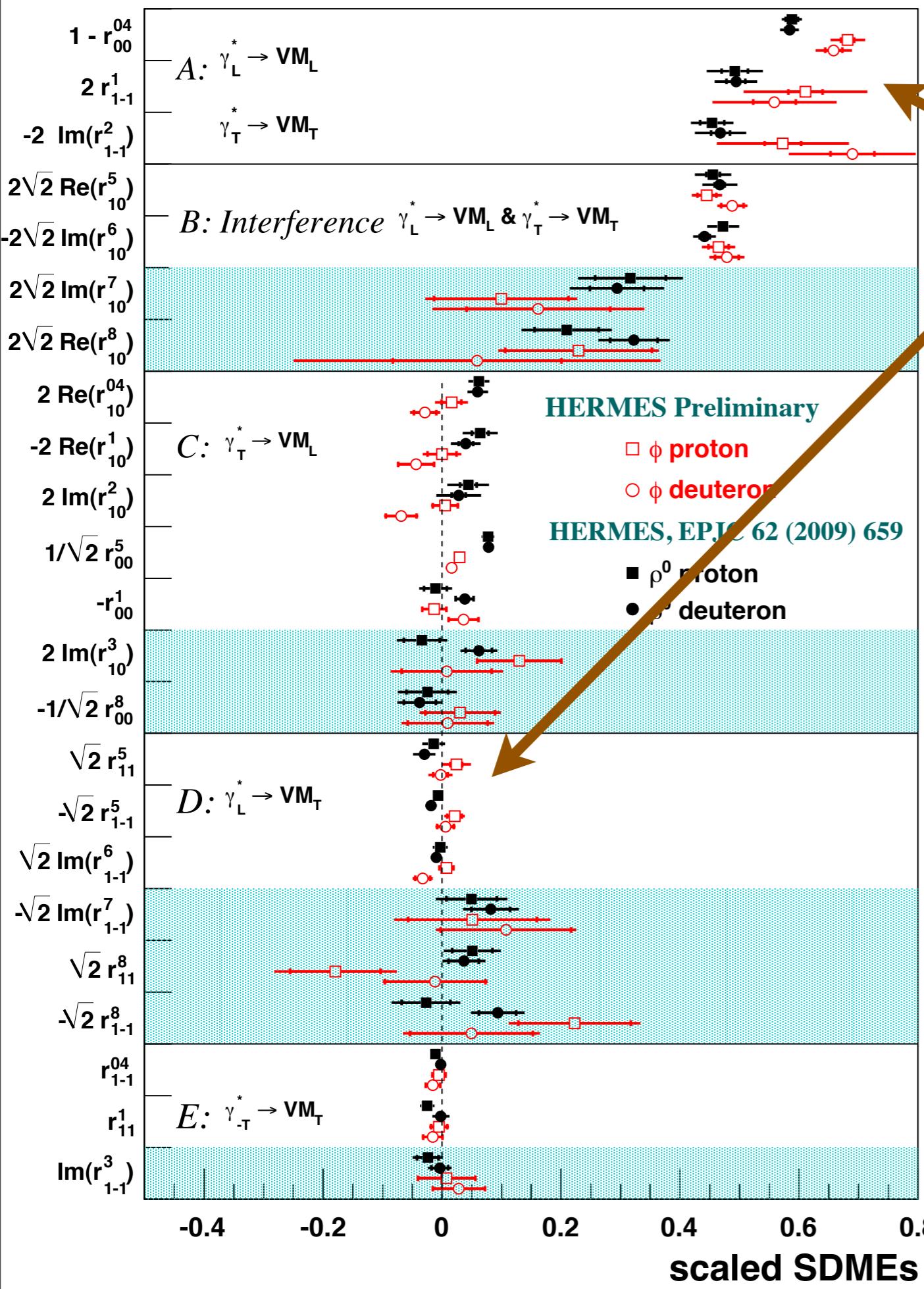
Results taken from measurement of  $ep \rightarrow eX\varphi$ .

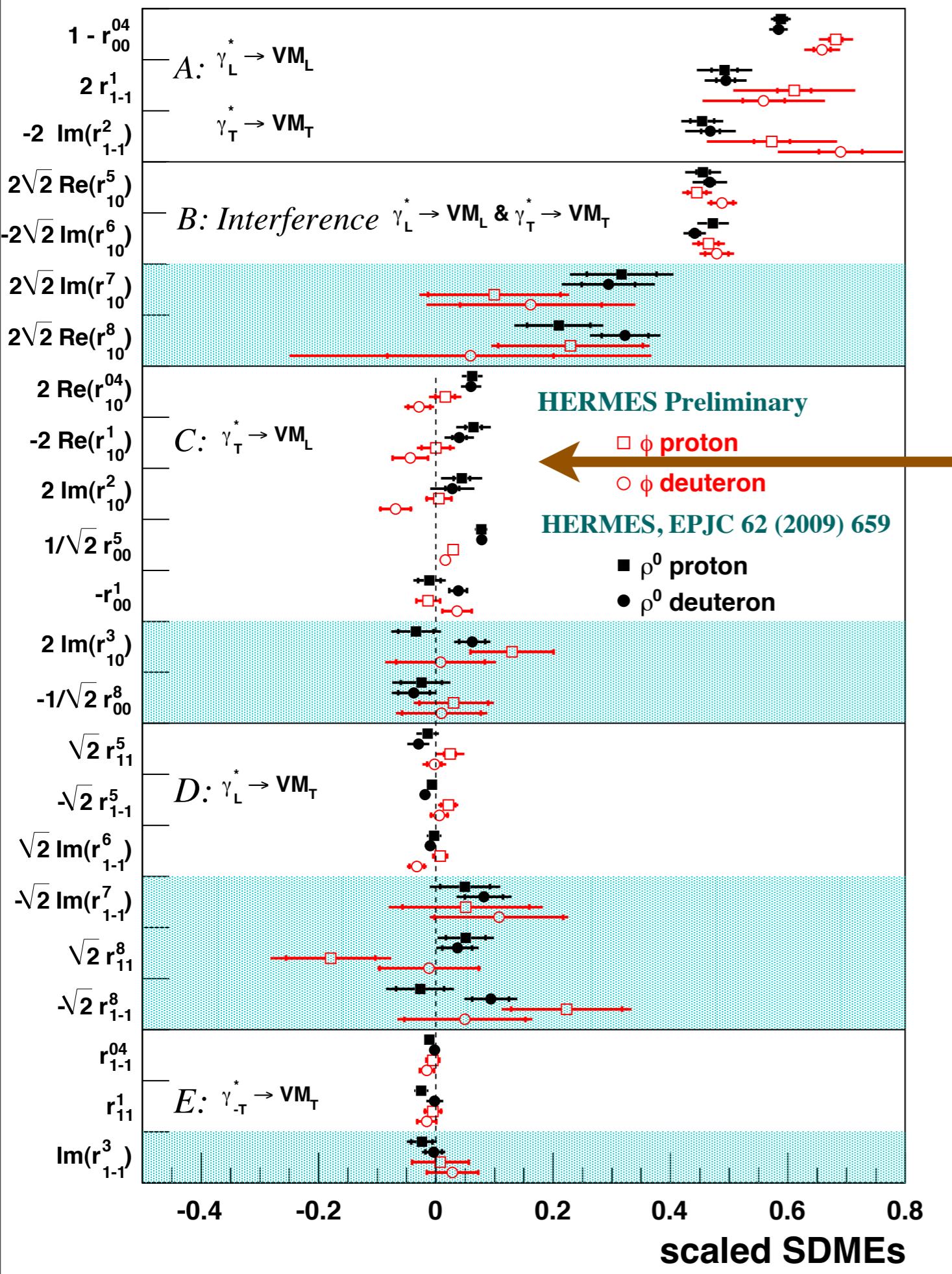
No measured distinction between proton and deuteron data.

Leading-twist transitions are typically larger than the  $\rho^0$ -equivalent.



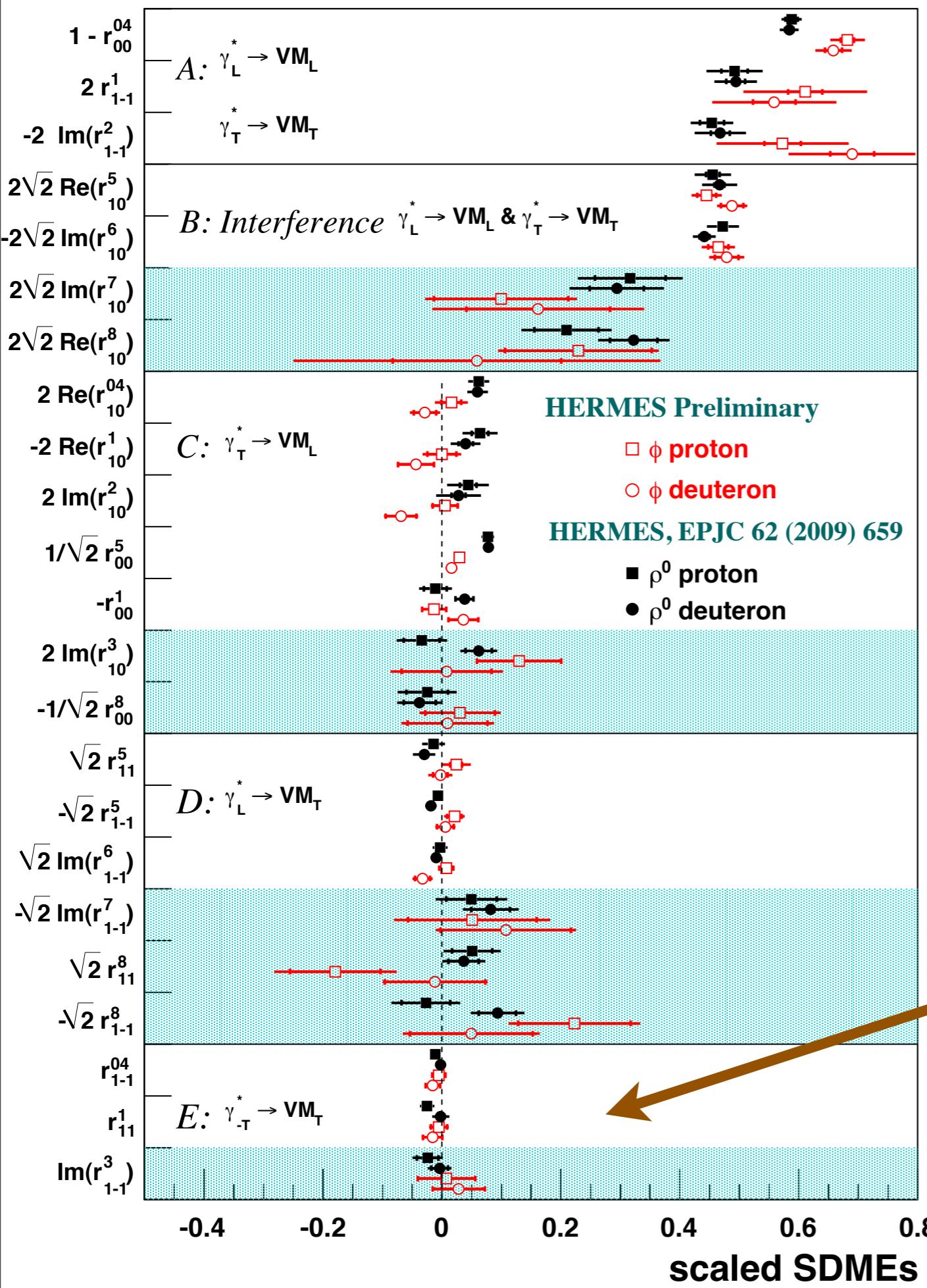
Longitudinal photons  
mostly produce  
longitudinal mesons





Longitudinal photons mostly produce longitudinal mesons

Some small indication that transverse photons can produce longitudinal mesons



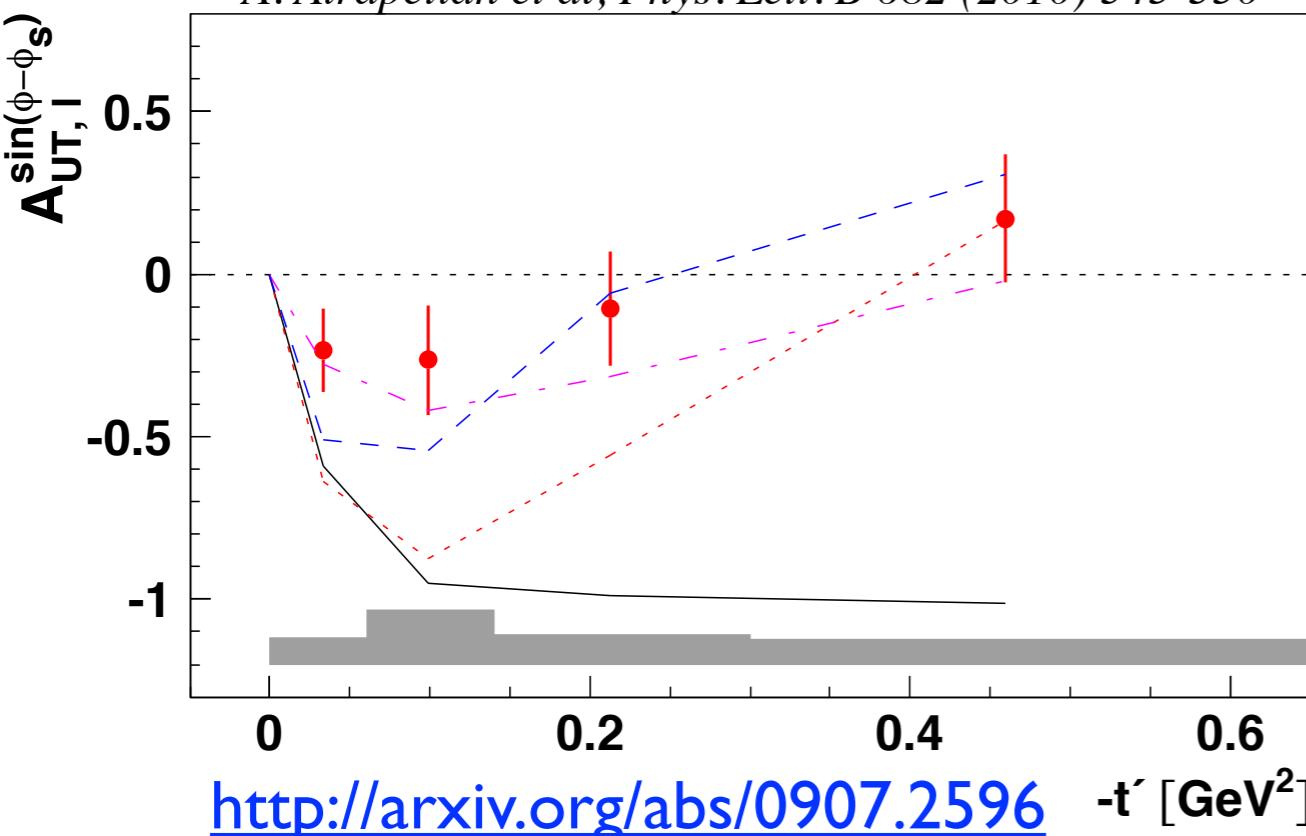
Longitudinal photons  
mostly produce  
longitudinal mesons

Some small indication  
that transverse photons  
can produce longitudinal  
mesons

Zero indication of two  
units of angular  
momentum change  
( $-T\gamma$  makes  $+T\varphi$ )

# Other Data?

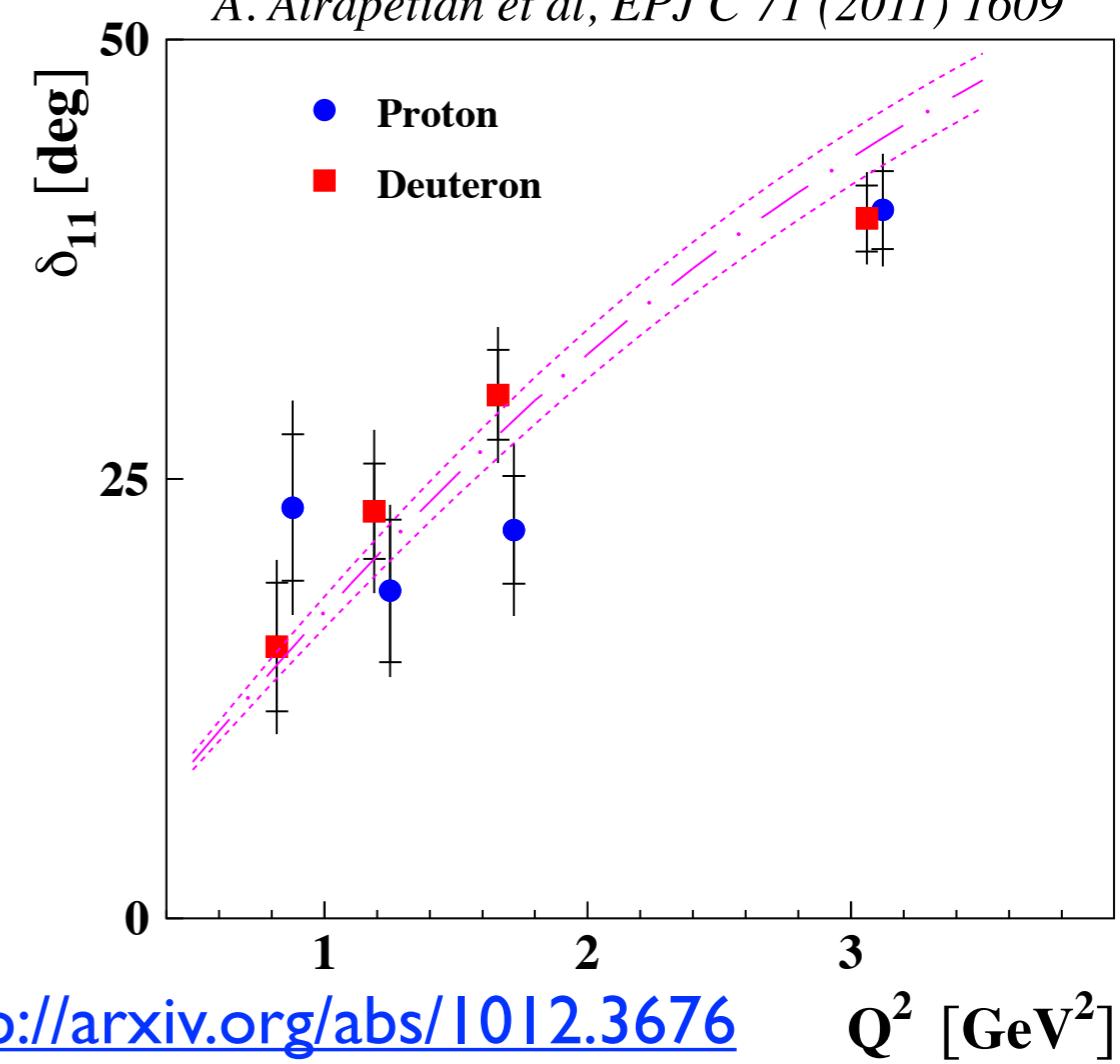
A. Airapetian et al, Phys. Lett. B 682 (2010) 345-350



Extraction of SDMES and Helicity Amplitude Ratios at HERMES for  $\rho$  mesons have shown that the handbag approximation is insufficient!

Meson data can also play a vital role in accessing GPDs - especially the “polarised” GPDs  $\tilde{H}$  and  $\tilde{E}$ !

A. Airapetian et al, EPJ C 71 (2011) 1609

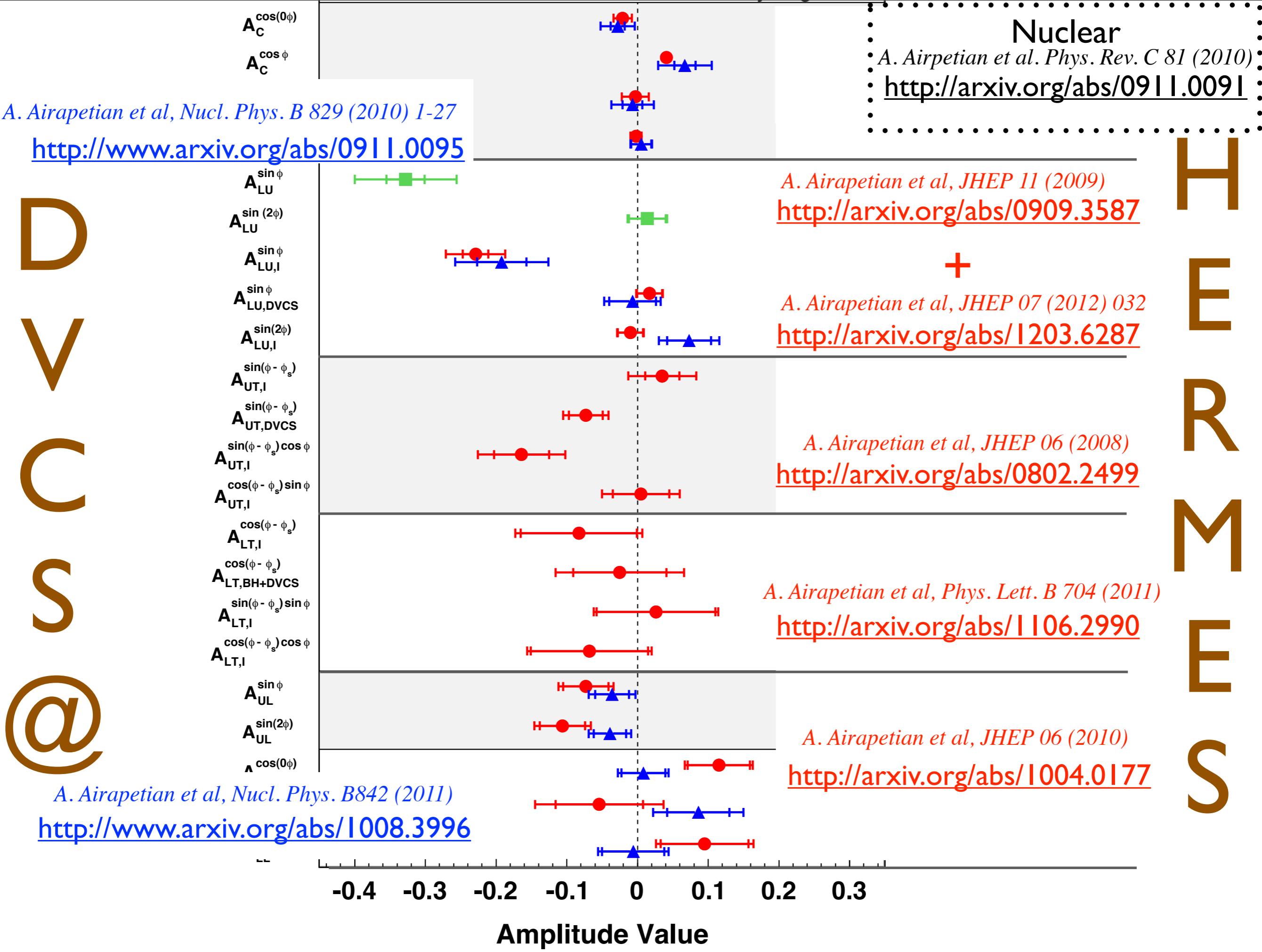


# Conclusions - What did we learn at HERMES?

- DVCS can be used to access information on Generalised Parton Distributions
- HERMES has the most diverse DVCS measurements of any experiment.
- Polarised target experiments are essential for the extraction of GPDs; should be seen as a fundamental experimental priority!

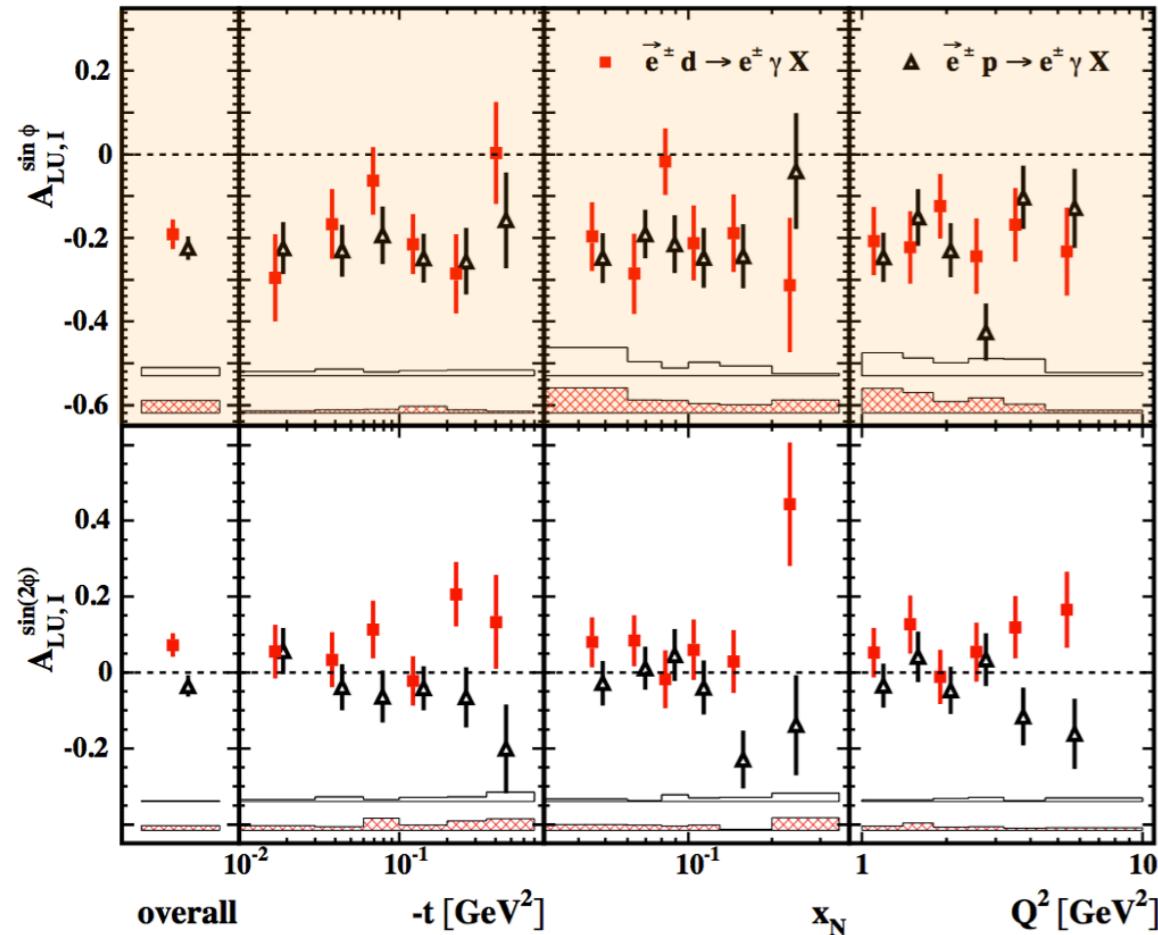
# Conclusions - What did we learn at HERMES?

- Lack of data means that **nuclear effects** on GPDs are not quantified! Incentive for new experiments at JLab, COMPASS and the EIC!
- Already, **GPDs can be constrained** - but there is much left to do!
- What effects do **chiral-odd GPDs** or **higher-twist distributions** have?



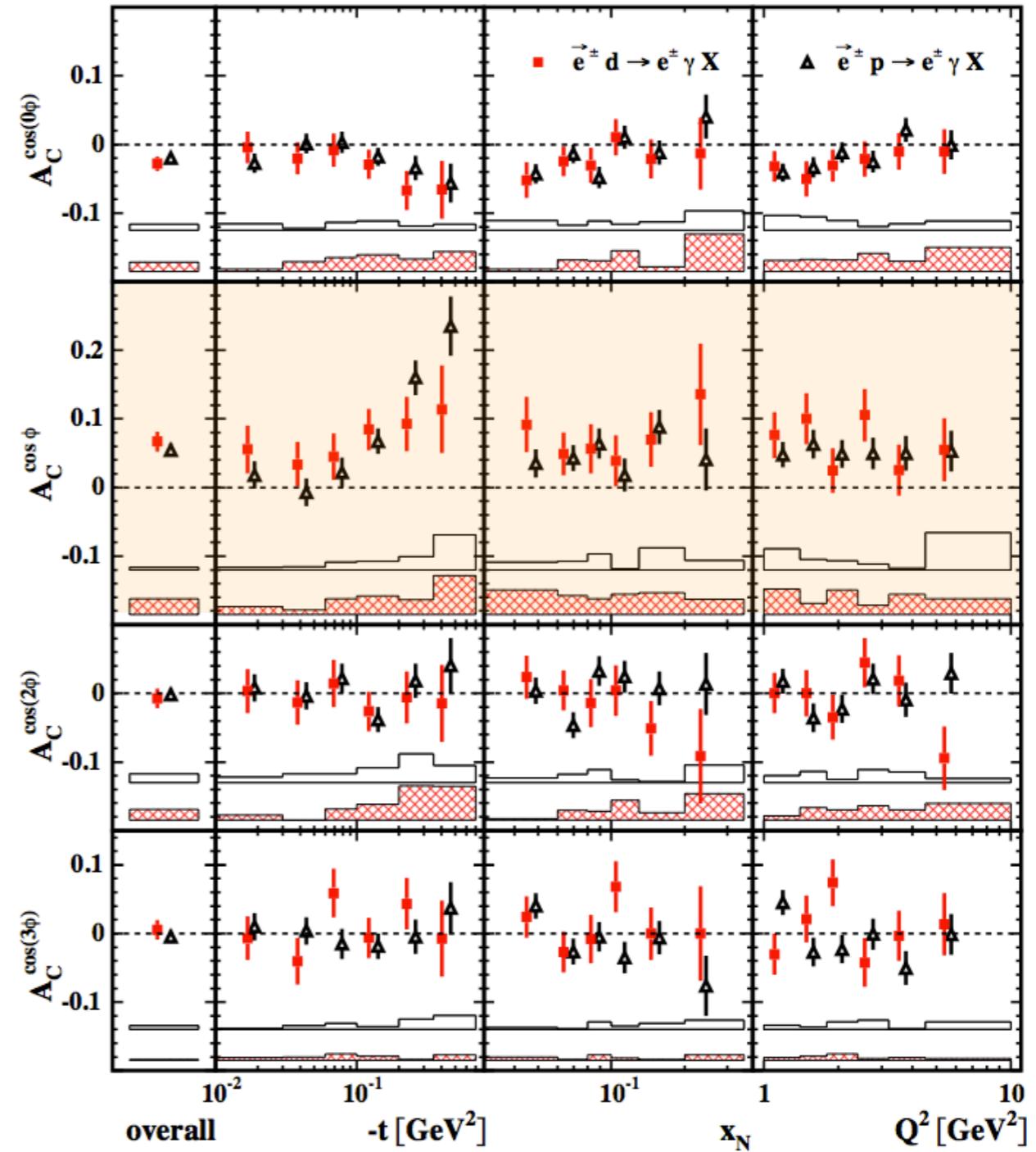
# Deuterium Beam-Asymmetries

A. Airapetian *et al*, Nucl. Phys. B 829 (2010) 1-27



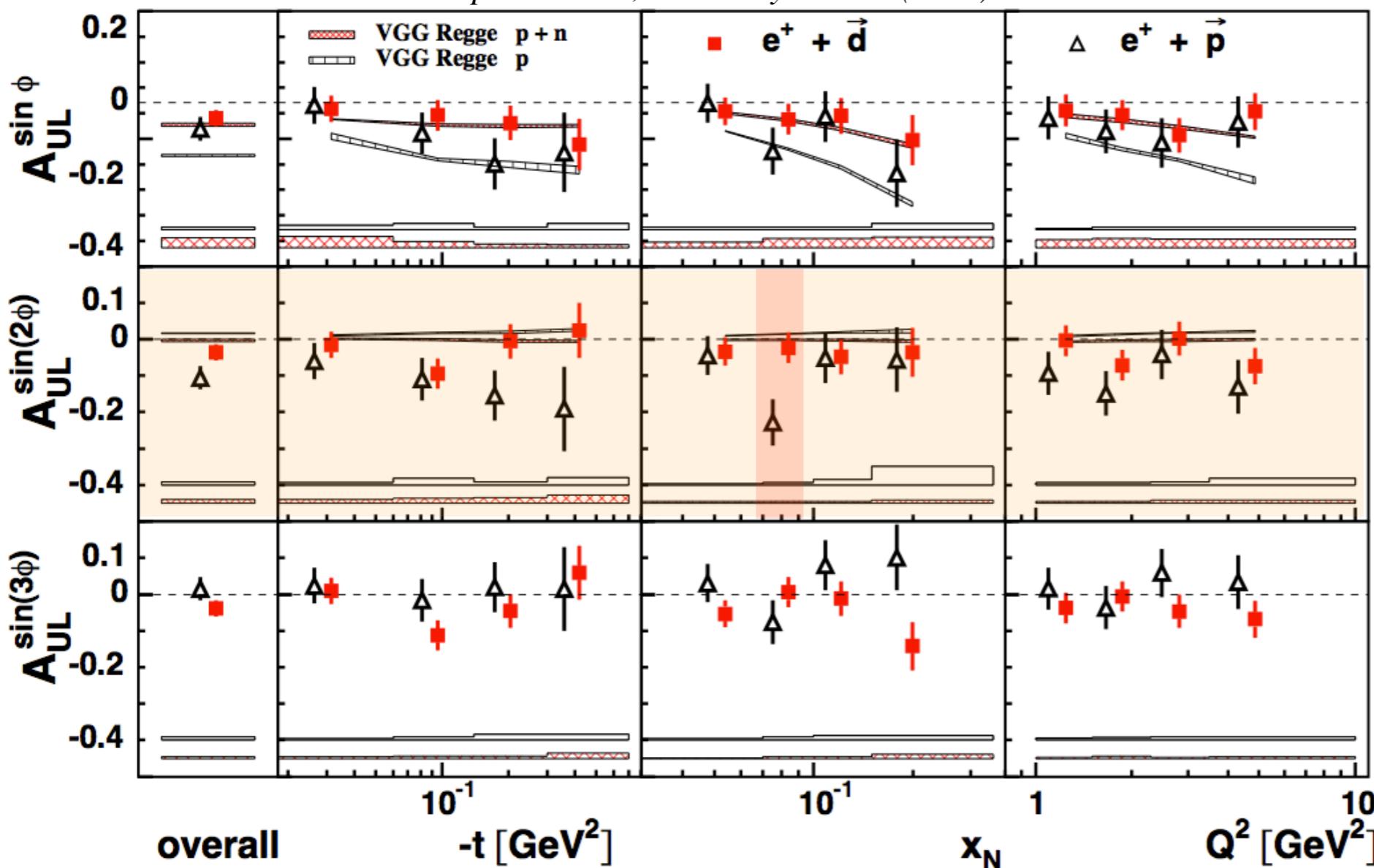
<http://arxiv.org/abs/0911.0095>

Deuterium is governed by different GPDs - but the asymmetry data is not so different even at low  $t$ !



# Deuterium-Target Asymmetries

A. Airapetian et al, Nucl. Phys. B842 (2011) 265-298

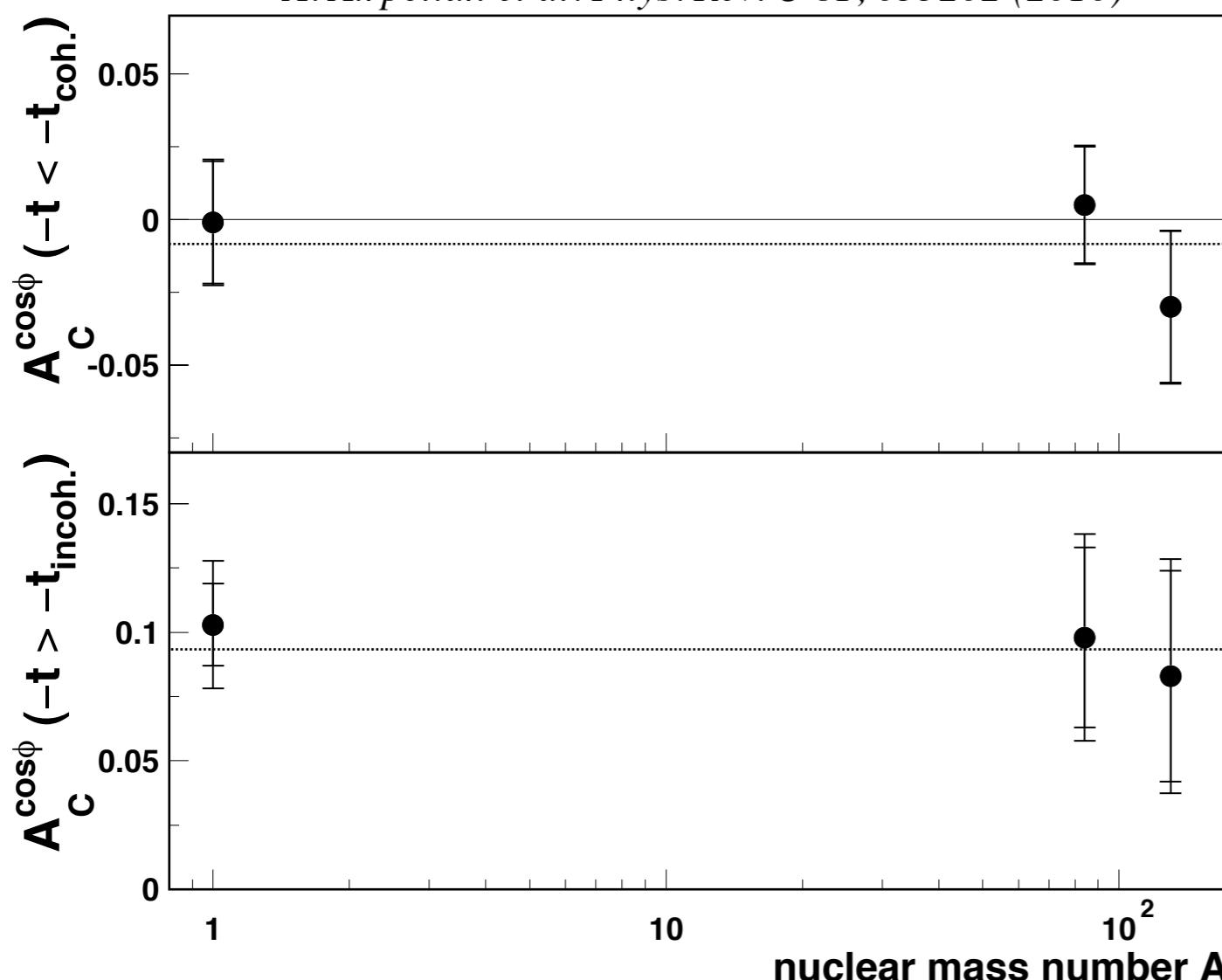


No good idea  
how to model  
long. pol.  
deuterium  
GPDs. Currently  
use a proton/  
neutron hybrid

<http://arxiv.org/abs/1008.3996>

# Nuclear Mass Dependence

A. Airpetian et al. Phys. Rev. C 81, 035202 (2010)

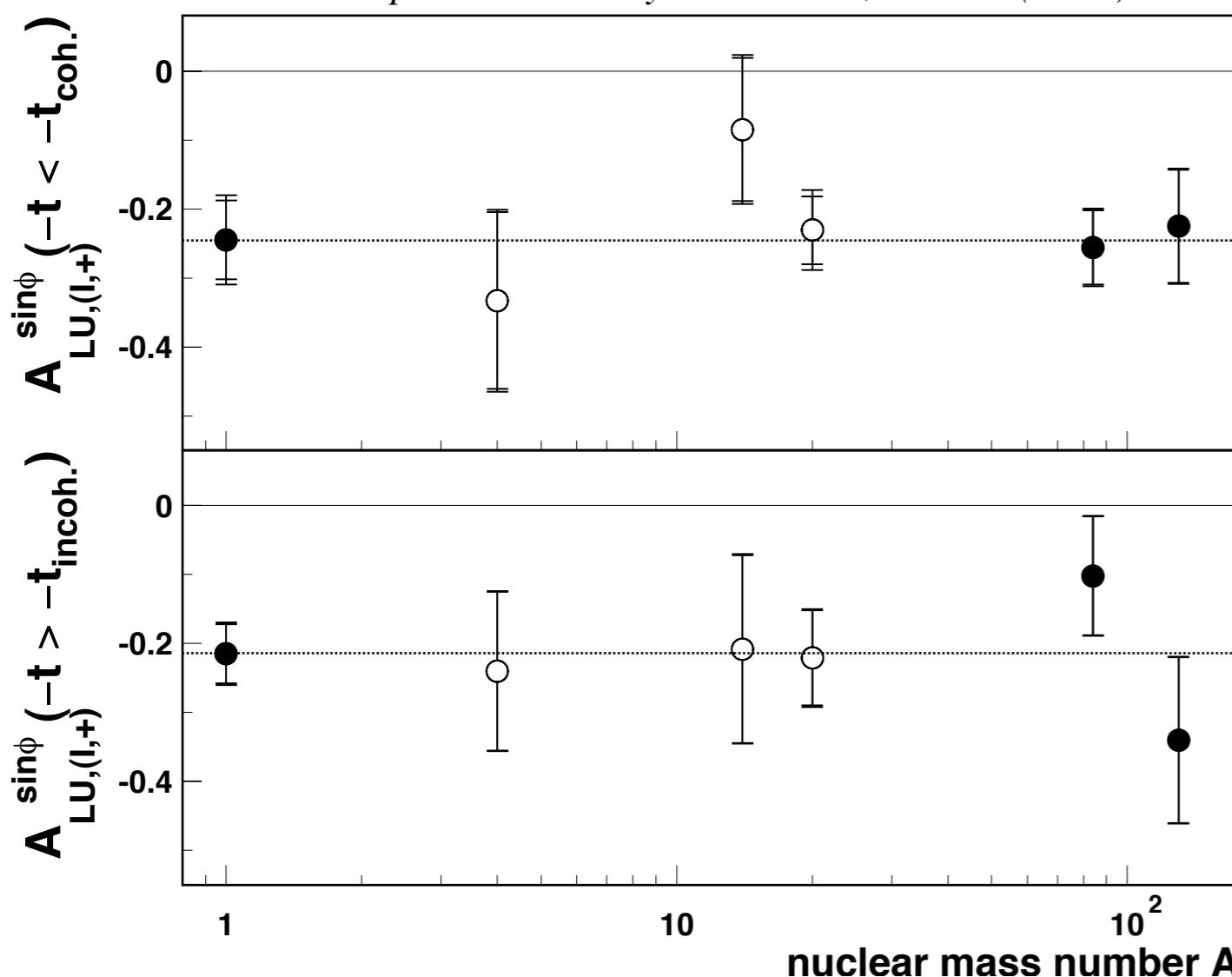


Nuclear-Binding models expected the DVCS asymmetry for nuclear targets to be 160-180% of the Hydrogen asymmetry.

<http://arxiv.org/abs/0911.0091>

# Nuclear Mass Dependence

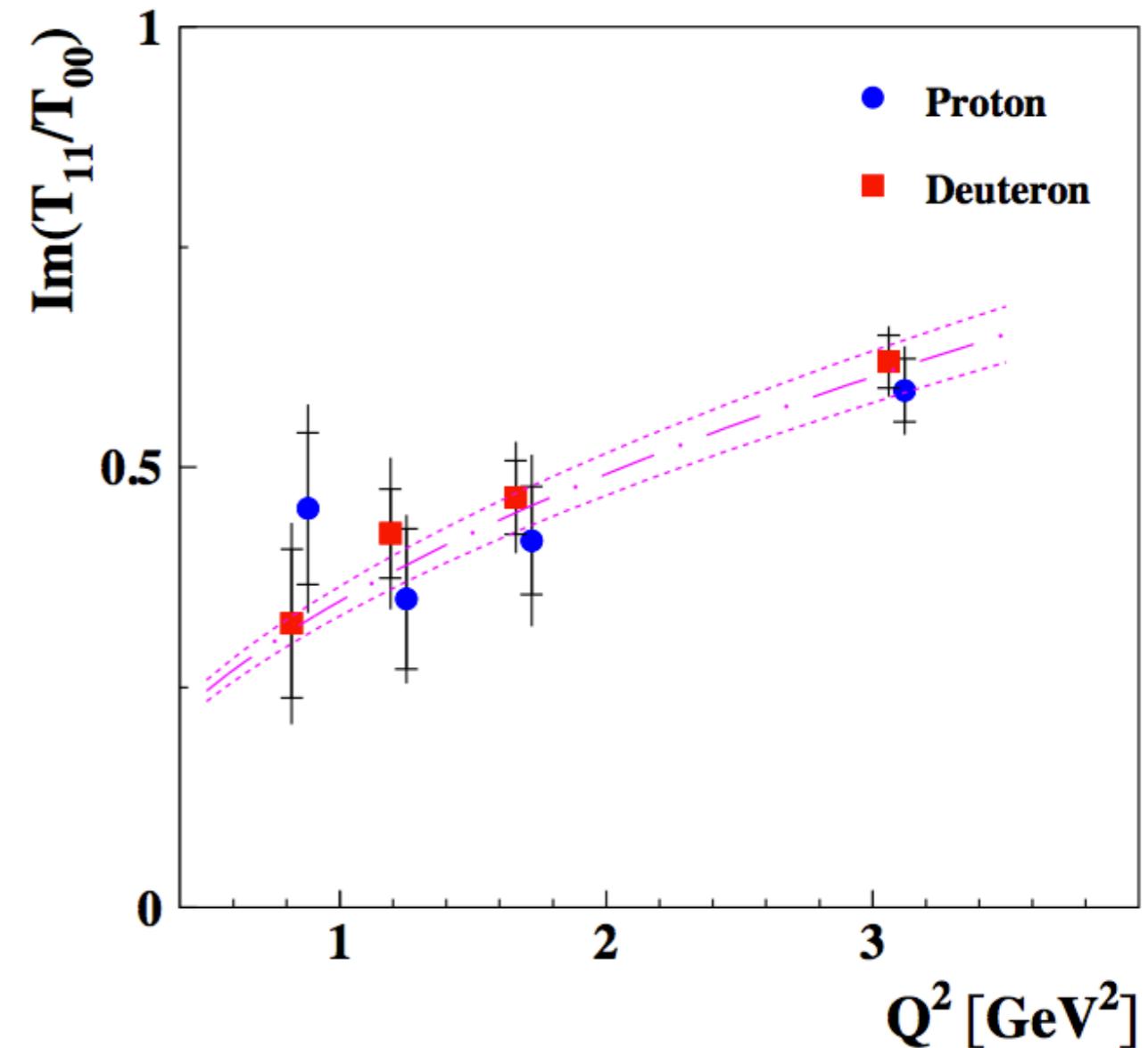
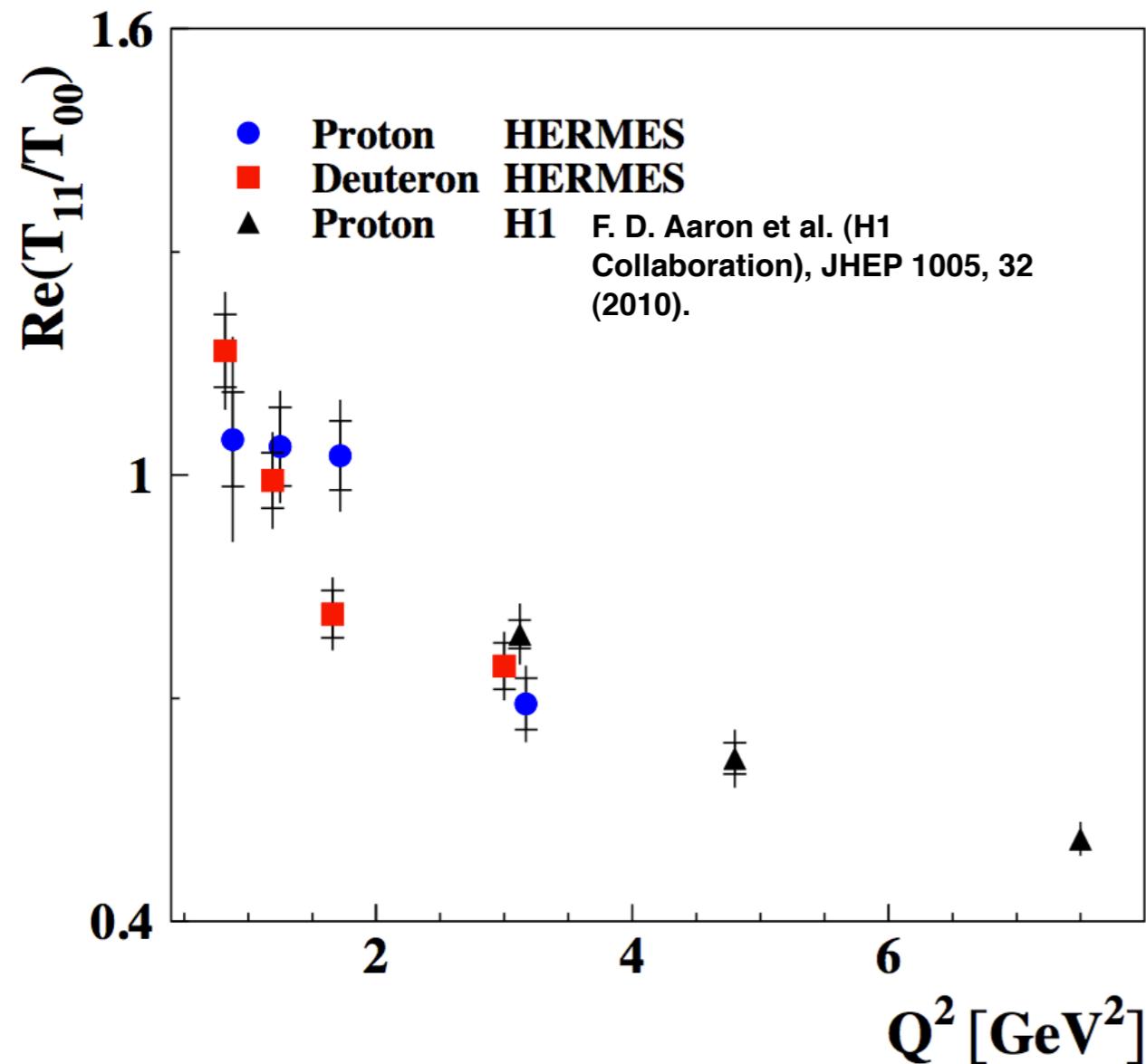
A. Airpetian et al. Phys. Rev. C 81, 035202 (2010)



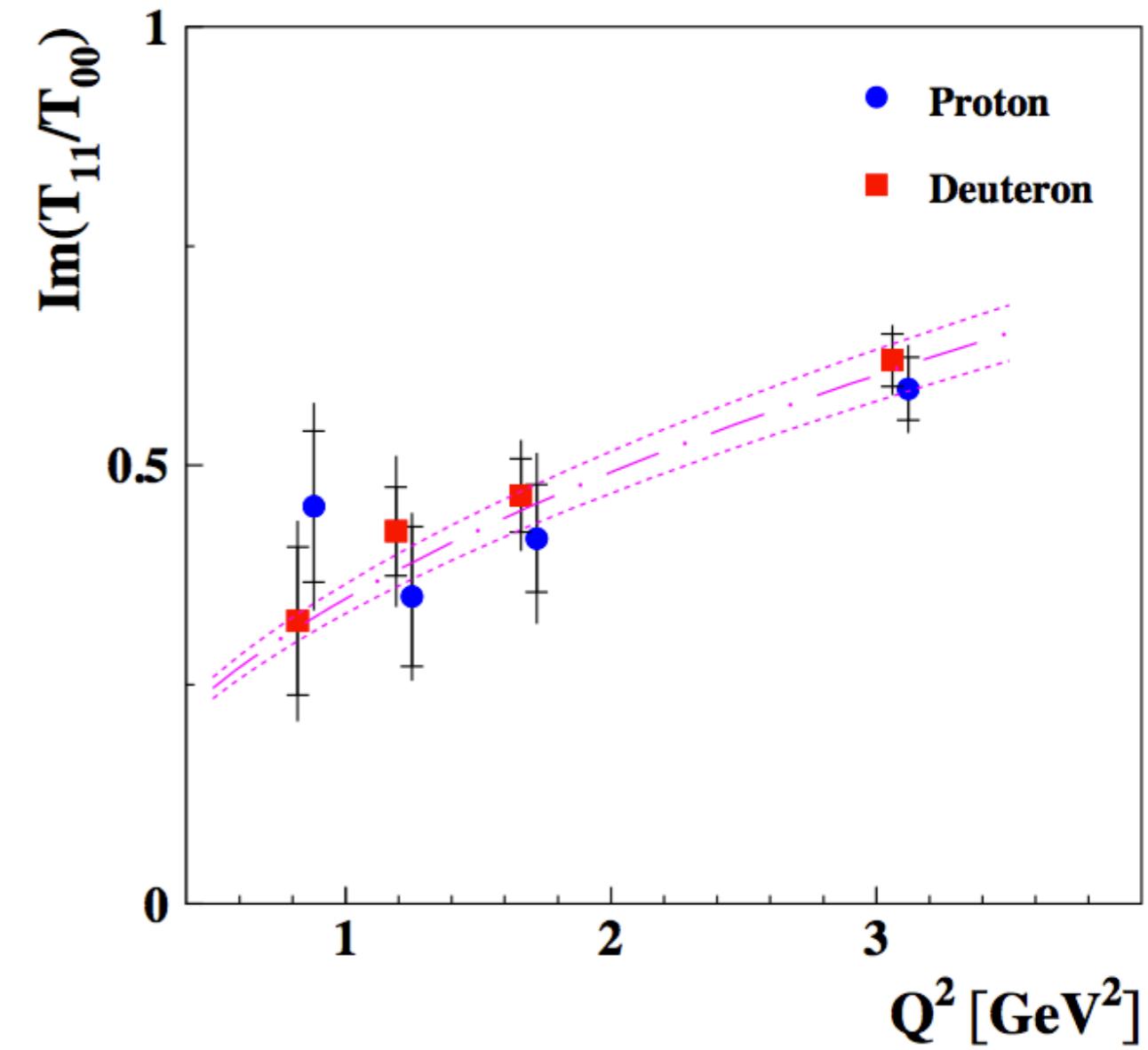
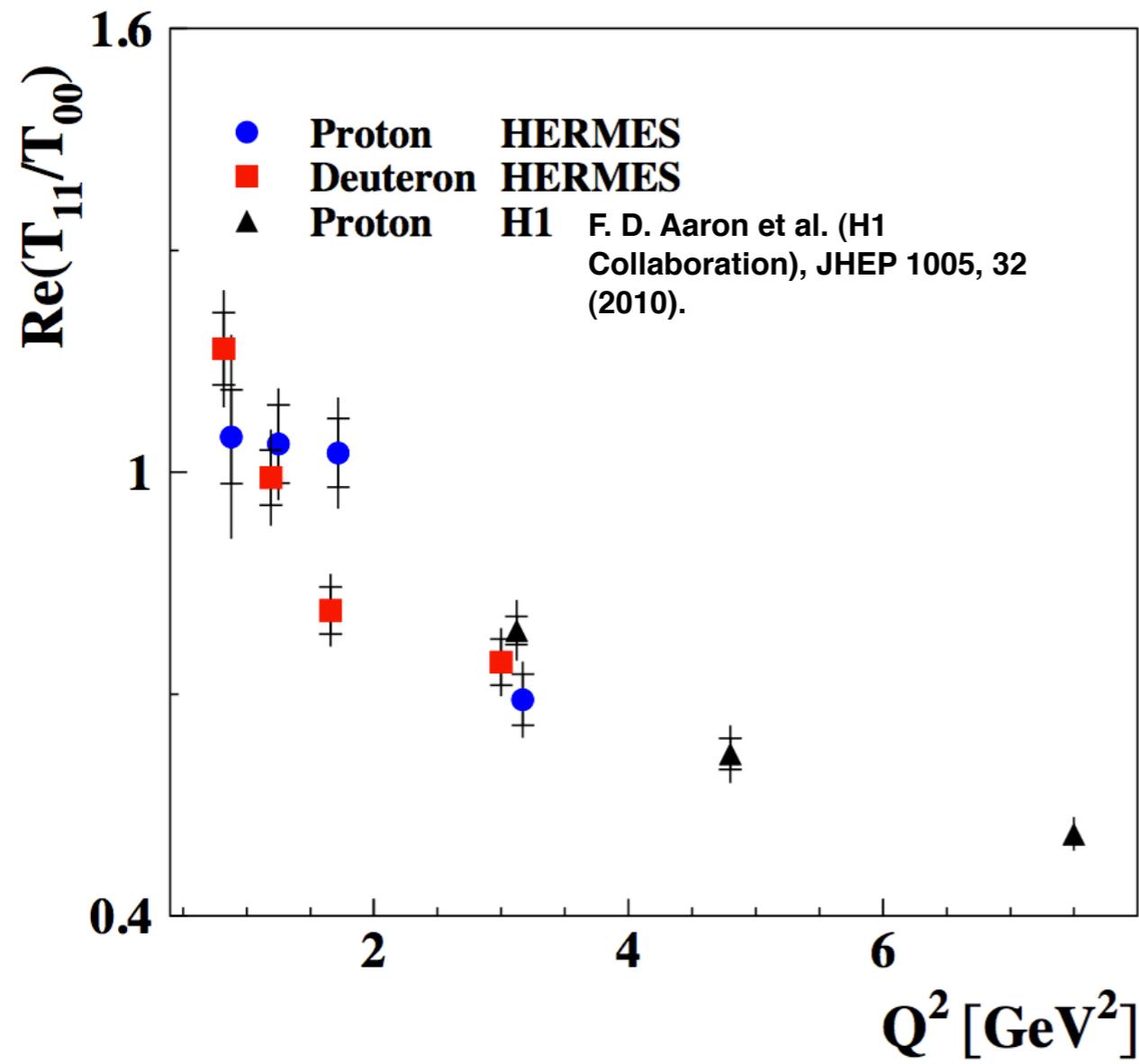
spin  $\frac{1}{2}$   
spin 1

The data shows  
**no significant difference**  
between coherent and  
incoherent DVCS  
processes

<http://arxiv.org/abs/0911.0091>



# Kinematic Dependence of $t_{11}$



Real Part follows  $a/Q$   
with  $a=1.11 \pm 0.03 \text{ GeV}$   
as expected!

Imaginary Part follows  $bQ$   
with  $b=0.34 \pm 0.02 \text{ GeV}^{-1}$   
(fit has no basis in theory)

# Phase Differences of HARs

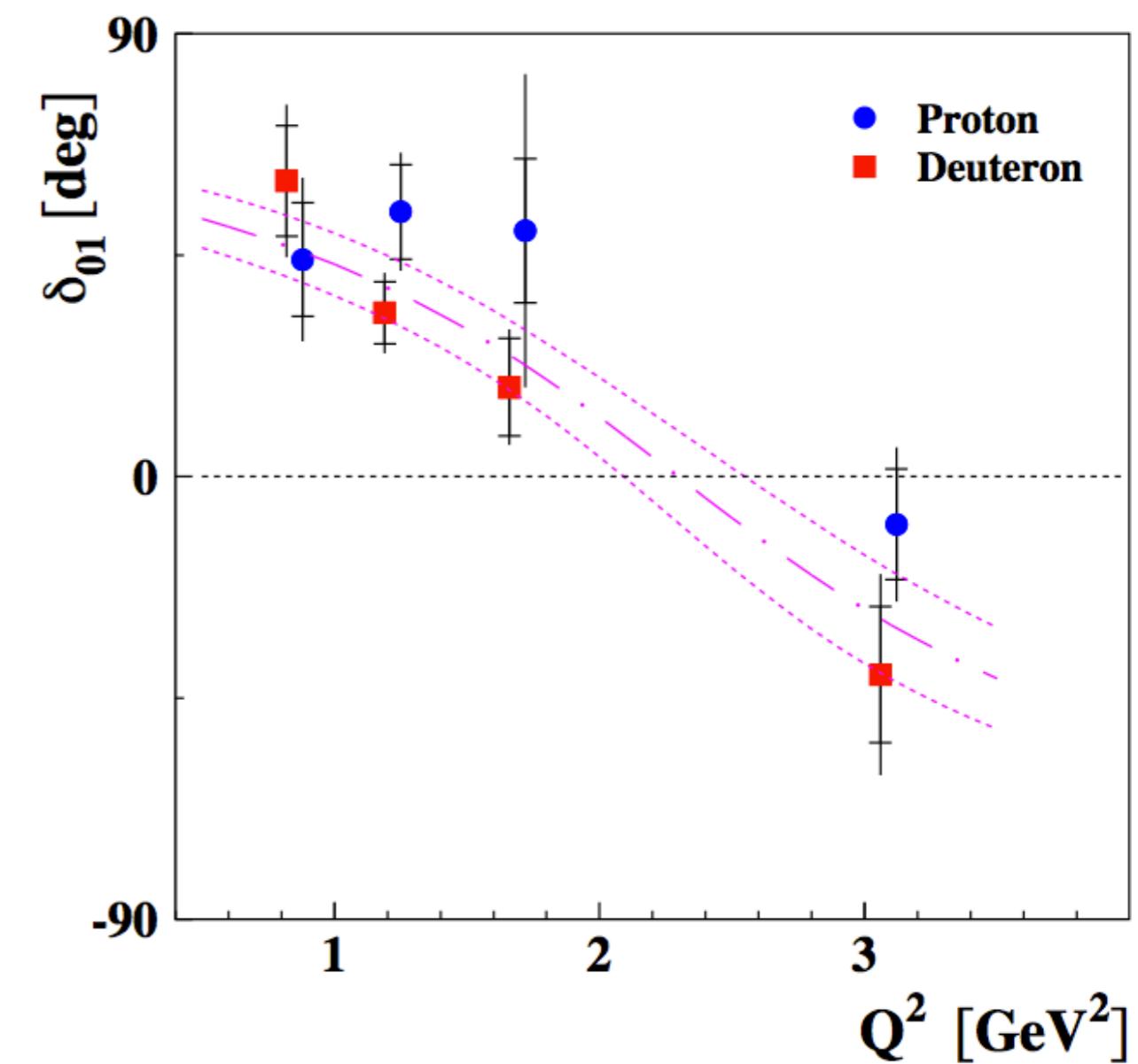
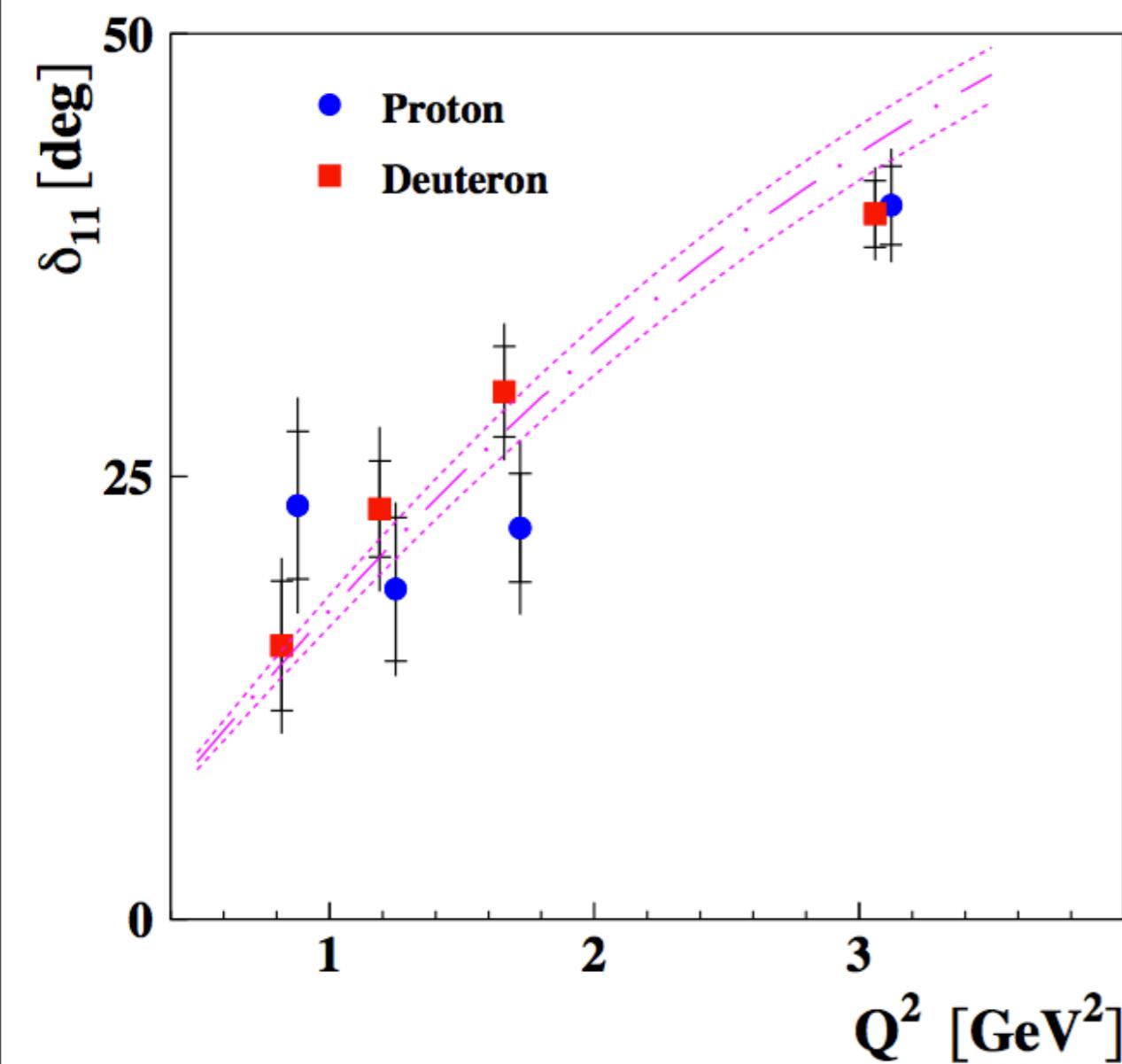
- GPD model predicts small phase difference for  $\tan(\delta_{11}) = \text{Im}(t_{11})/\text{Re}(t_{11})$

[S. V. Goloskokov and P. Kroll,  
Eur. Phys. J. C 53, 367 \(2008\)](#)

- $t_{01}$  is expected to be the largest SCHC-violating amplitude and  $\delta_{01}$  should be constant

[D. Yu. Ivanov and R. Kirschner,  
Phys. Rev. D 58, 114026 \(1998\)](#)

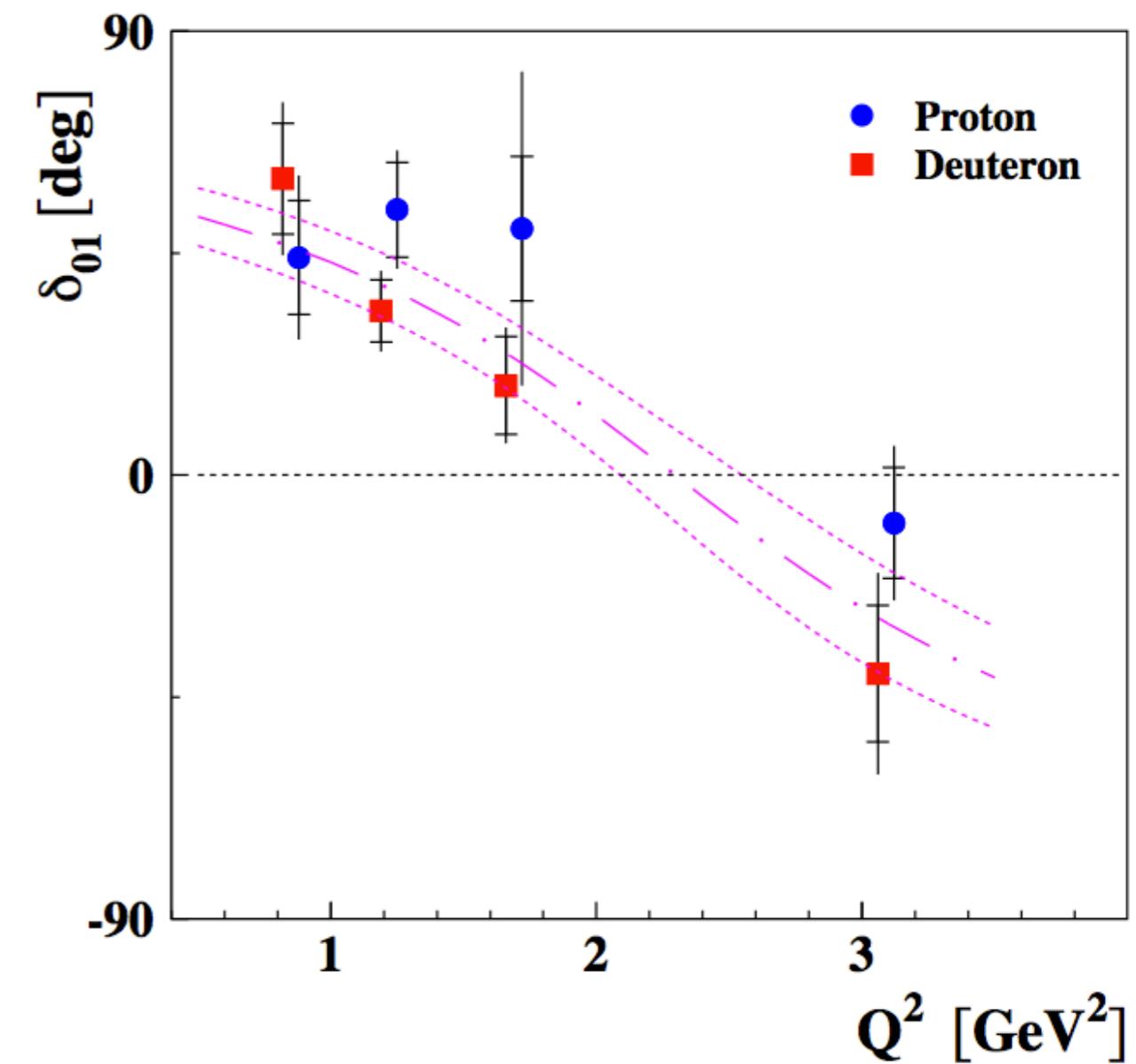
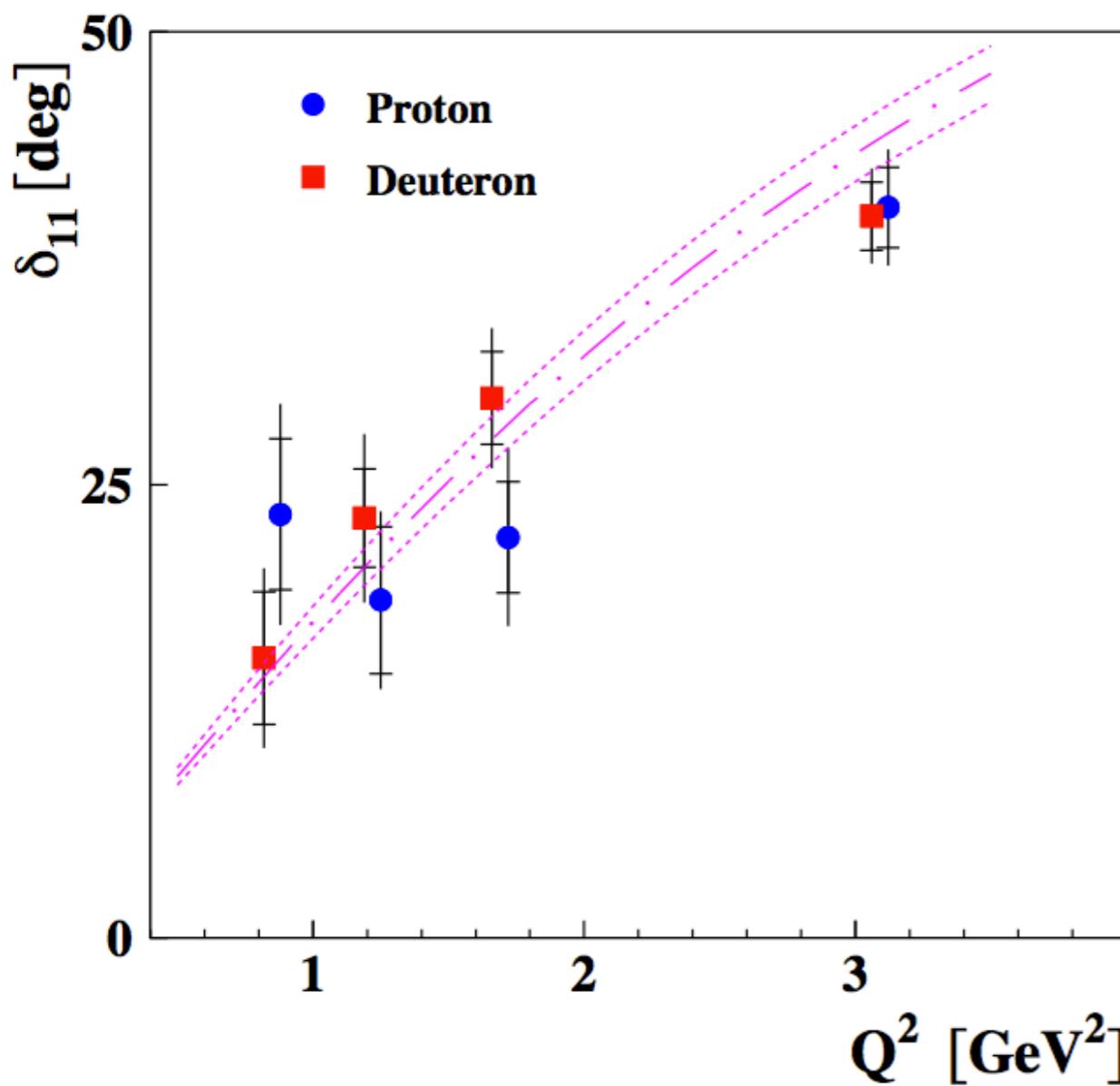
# Phase Differences of HARs



Large value **contradicts**  
GPD-based models

Should be a **constant**

(Neither  $\text{Re}(t_{01})$  nor  $\text{Im}(t_{01})$  follow theoretical dependence predictions!!!)



N.B: Fits have no basis in theory

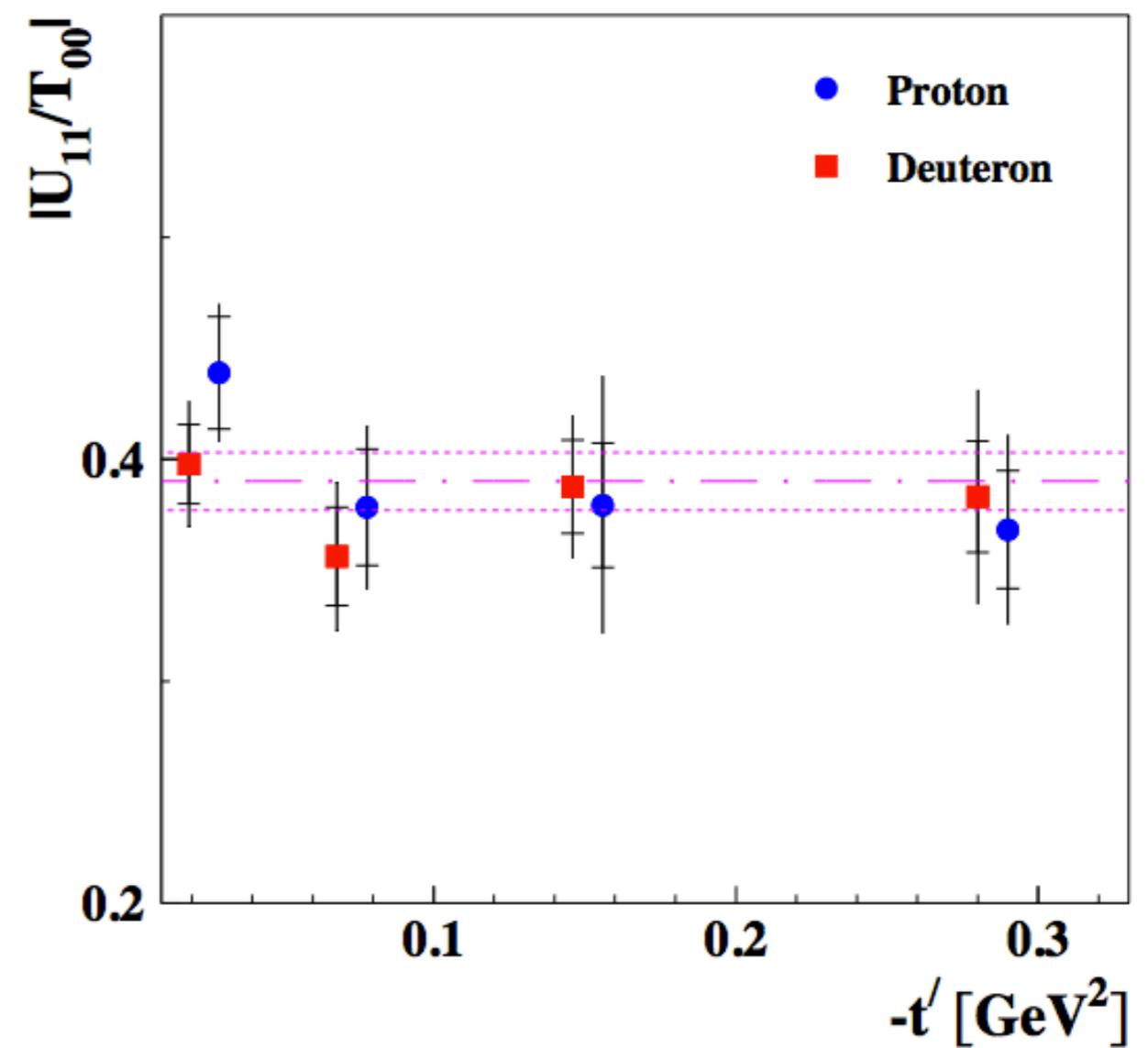
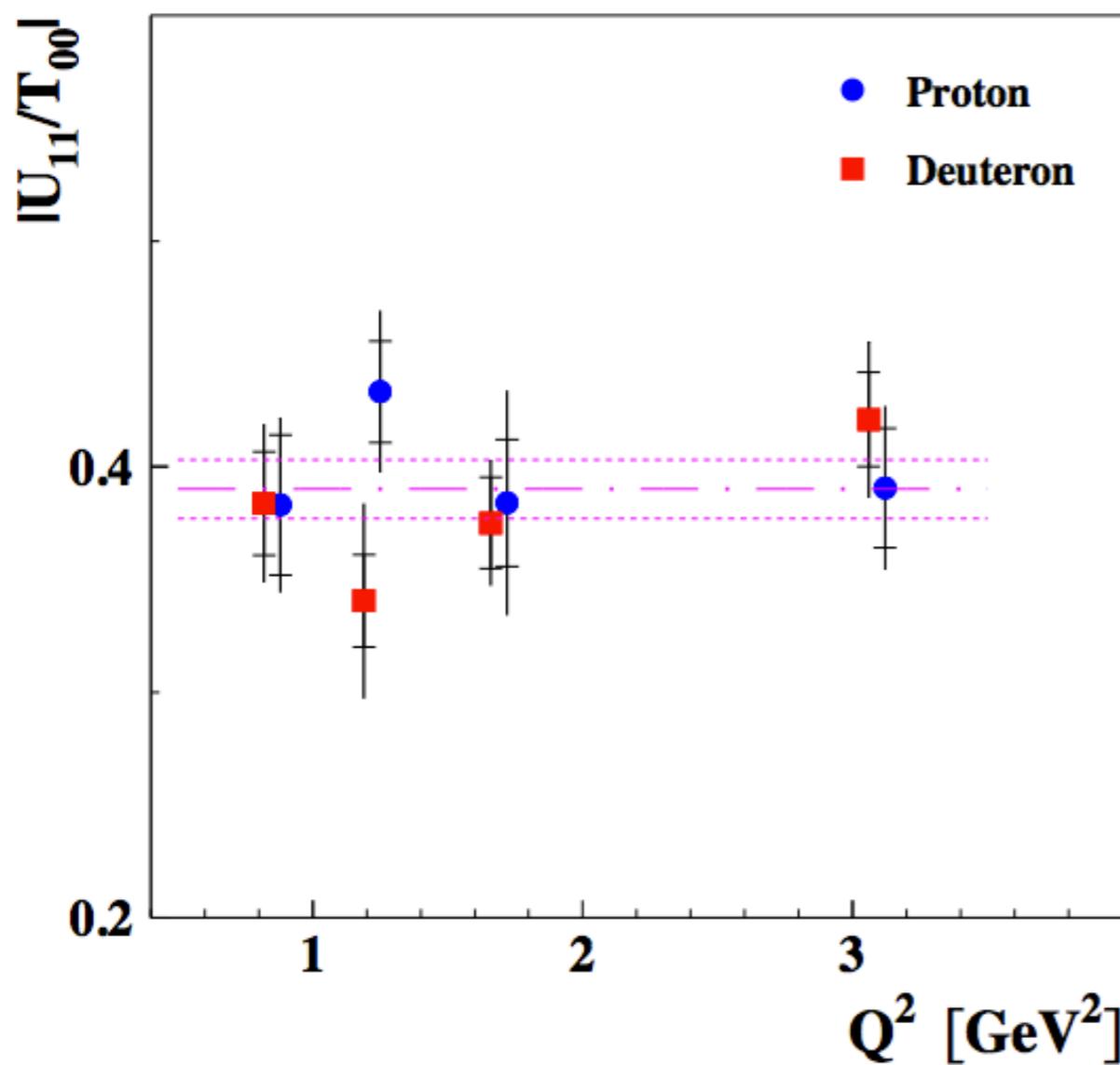
# Helicity Amplitude Hierarchy

## Behaviour of UPE

$$|T_{00}|^2 \approx |T_{\perp\perp}|^2 \gg |U_{\perp\perp}|^2 > |T_{0\perp}|^2 \gg |T_{\perp 0}|^2 \dots$$

- $u_{\perp\perp} = |U_{\perp\perp}|/|T_{00}|$  should be small ( $u_{\perp\perp} \approx 0.2$ ) but **visible** (only) for  $\rho^0$  at HERMES!
- May naively expect a  $I/Q$  dependence in  $u_{\perp\perp}$
- UPE is one-pion exchange => may also see some influence of the pion-pole at small  $t$ ?

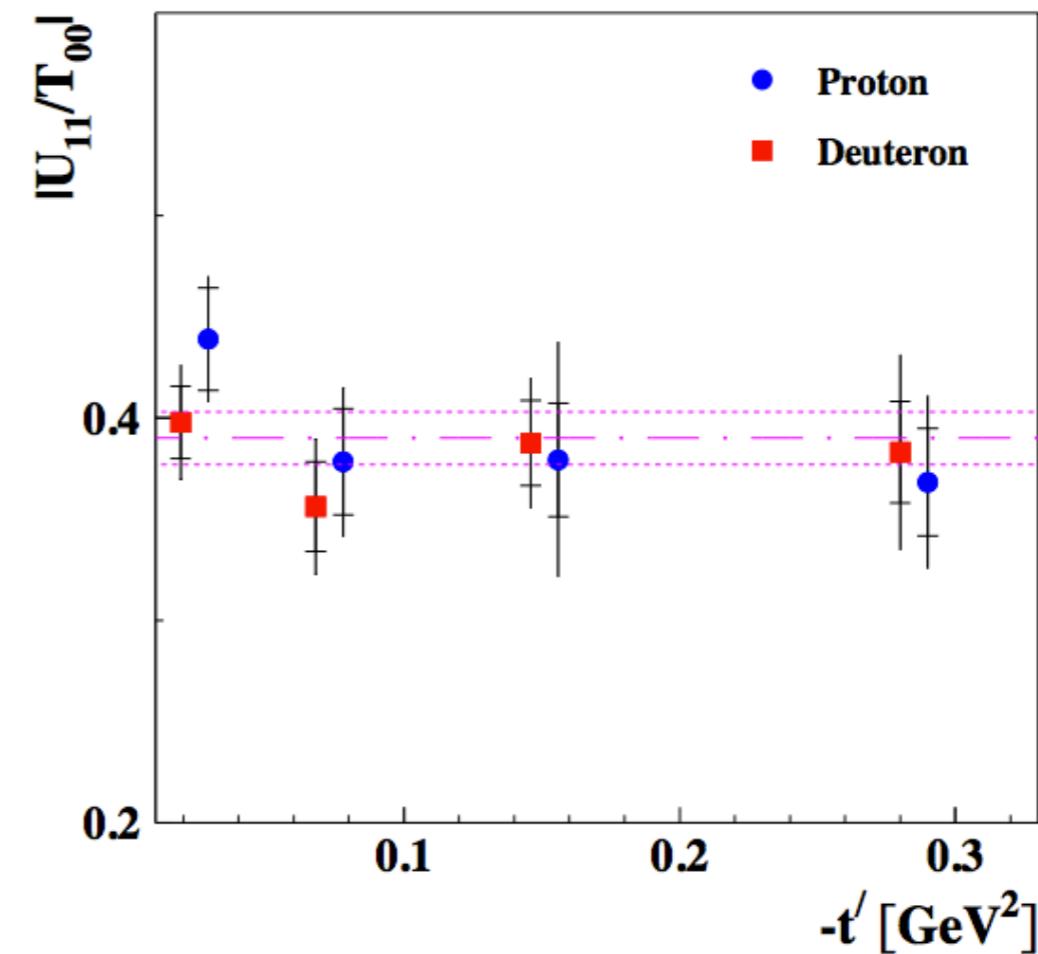
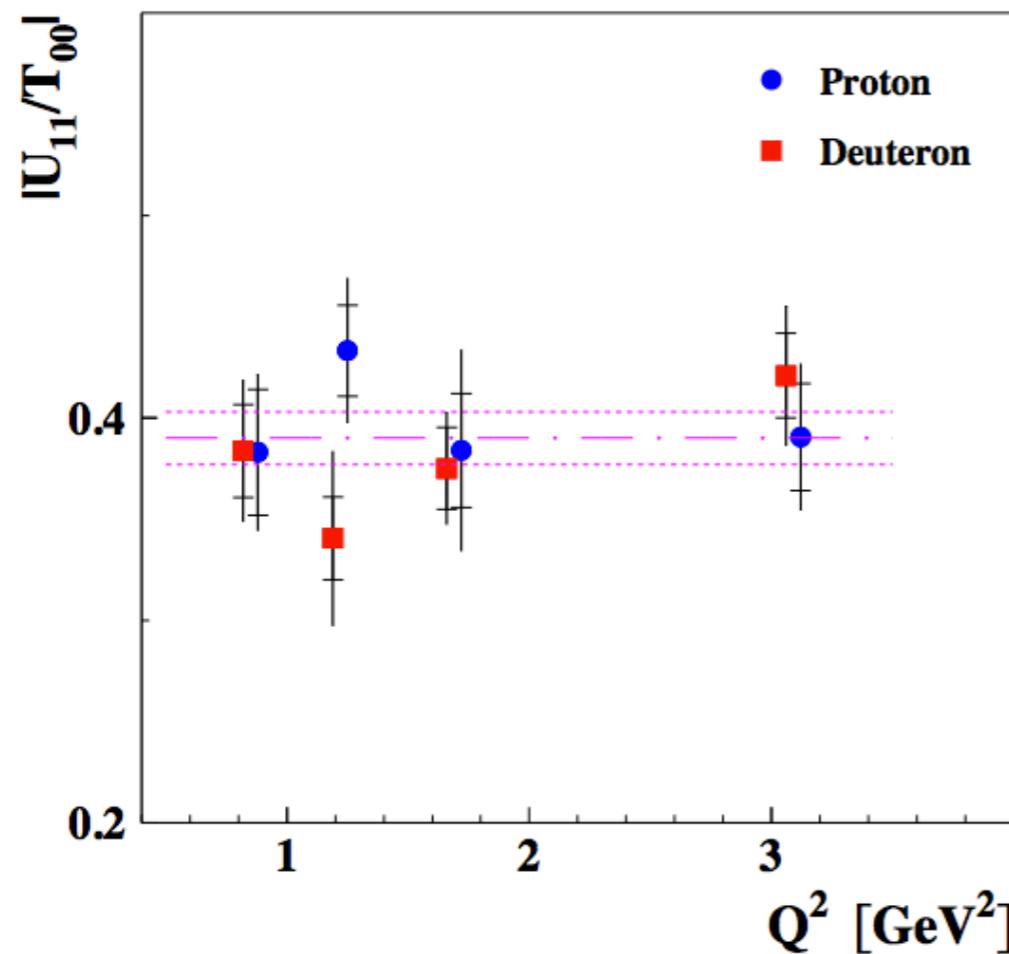
# Unnatural Parity Exchange



# Unnatural Parity Exchange

No dependence on  $Q^2$ !

No dependence on  $t'$ !



Existence established to  $20\sigma$  (integrated extraction)  
Magnitude of  $U_{11}$  is 2.5x smaller than  $T_{00}$

# Unnatural Parity Exchange

- No dependence on  $Q^2$  may be because  
**HERMES is far from the asymptotic region ?**
- No dependence on  $t'$ 
  - ➡ **Too far from pion-pole ?**
  - ➡  **$U_{11}$  not dominated by one-pion exchange ?**
  - ➡ **An underlying dependence of  $T_{00}$  on  $t'$  ?**