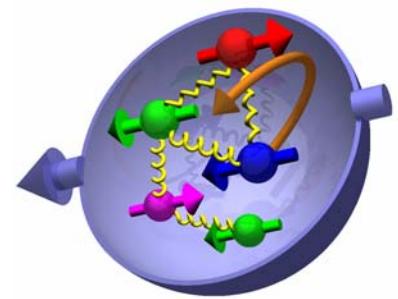


hunting the OAM @ ...



- a brief introduction
- GPDs & OAM
- observables: A_{UT} in DVCS & exclusive ρ^0
- conclusion & perspectives

GPDs and the spin puzzle



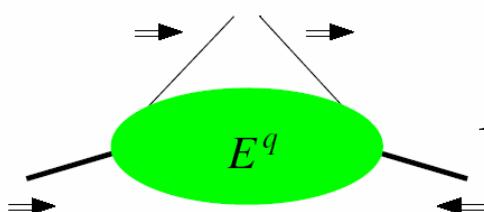
nucleon spin:

$$S_z^n = \frac{1}{2} = \frac{1}{2} \sum_q \Delta q + L_z^q + \Delta G + L_z^g = J_q + J_g$$

\uparrow \uparrow
 $\approx 30\%$ $\approx \text{zero}$

[X. Ji, 1997]

$$J_{q,g} = \lim_{t \rightarrow 0} \frac{1}{2} \int_{-1}^1 x dx \left[H^{q,g}(x, \xi, t) + E^{q,g}(x, \xi, t) \right]$$



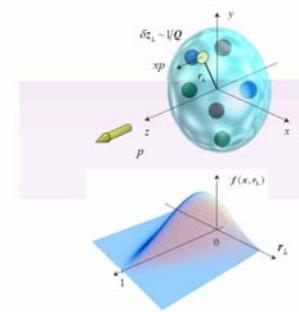
$E^q \neq 0$ requires orbital angular momentum

proton helicity flipped but quark helicity conserved

GPDs: nucleon tomography

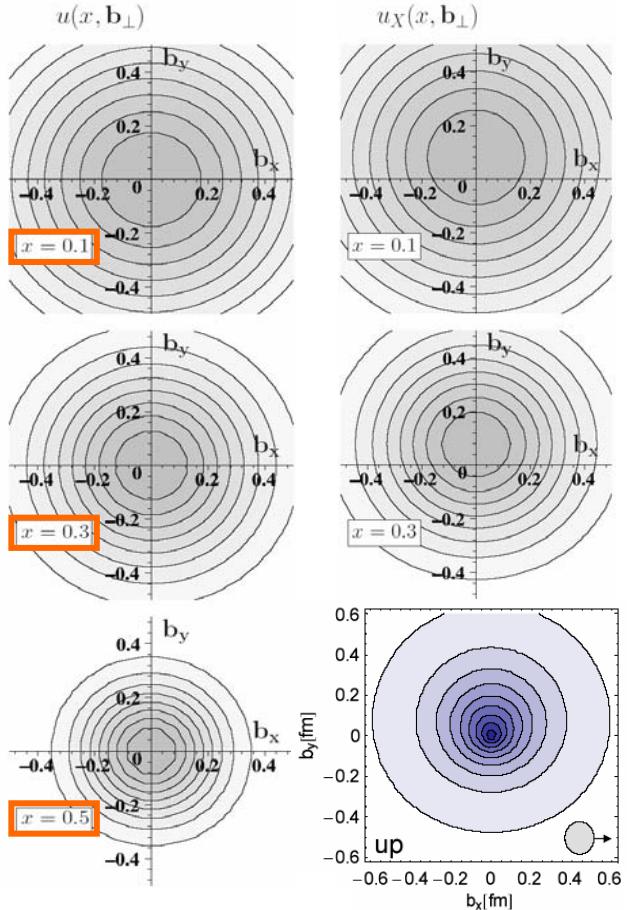
[M. Burkardt, M. Diehl 2002]

$\mathcal{FT}(\text{GPD})$: momentum space \rightarrow impact parameter space:
distribution of partons in plane transverse to longitudinal momentum x

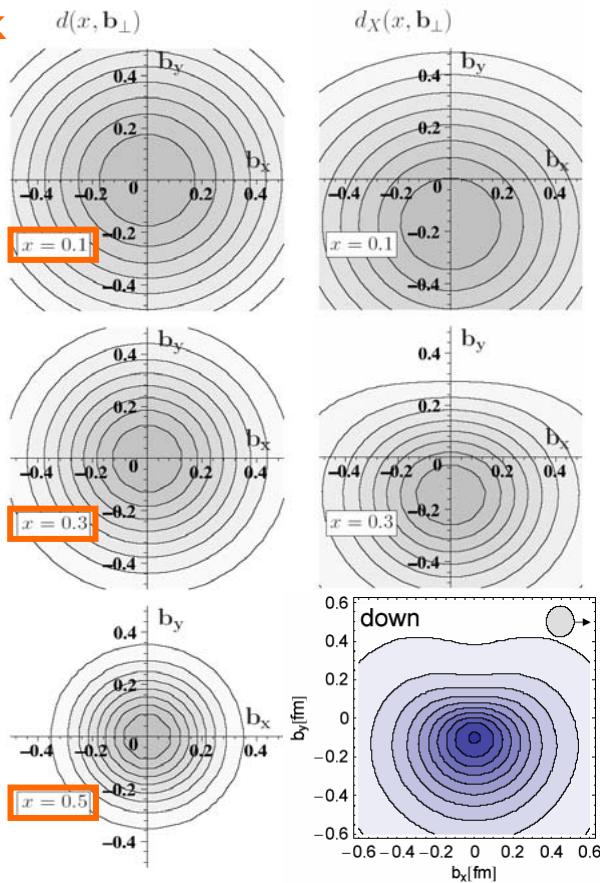


polarised nucleon: *spin-orbit correlations (TMDs)*

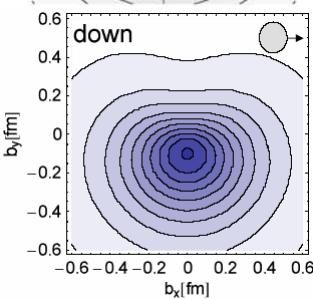
u-quark



d-quark



from
lattice
[QCDSF]

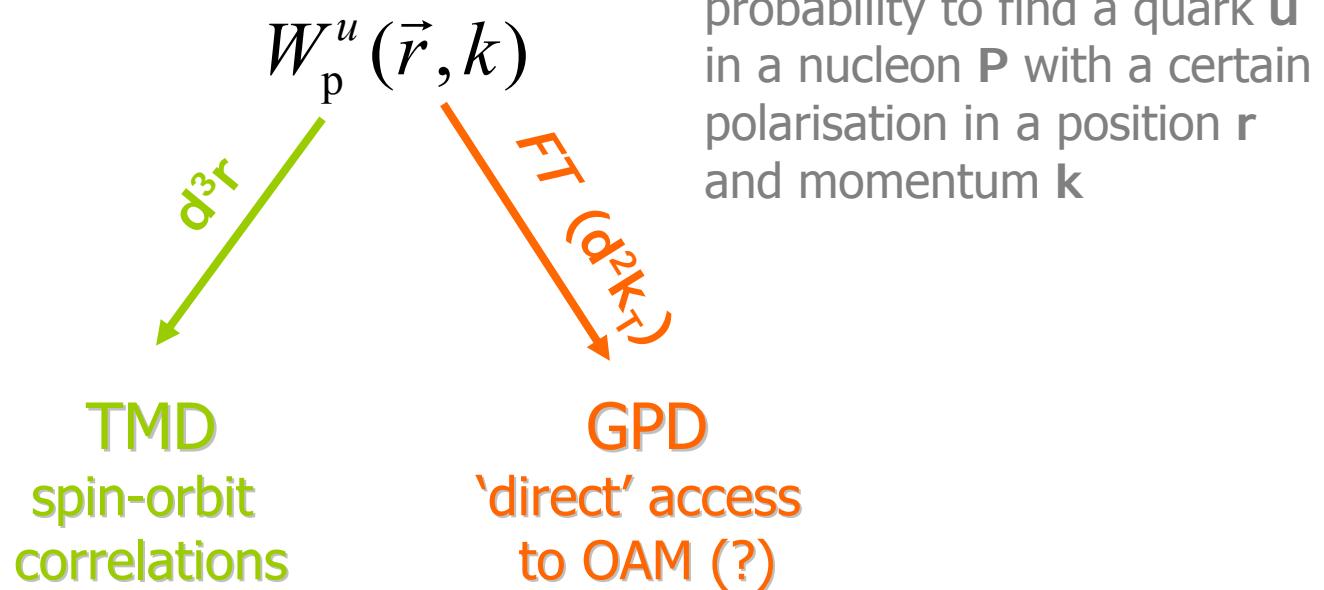


TMDs \longleftrightarrow GPDs

3D structure of hadrons : nucleon tomography

→ complementary:

Wigner distribution: ("mother" function)



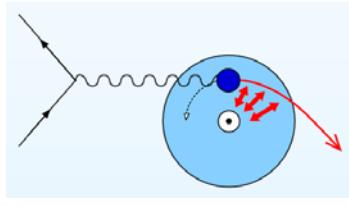
→ relations between TMDs and GPDs (?)

see talk by L.Gamberg

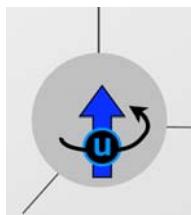
spin-orbit correlations @

Sivers fct., Boer-mulders fct., pretzelosity

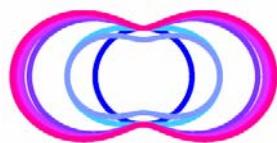
see talk by N. Makins



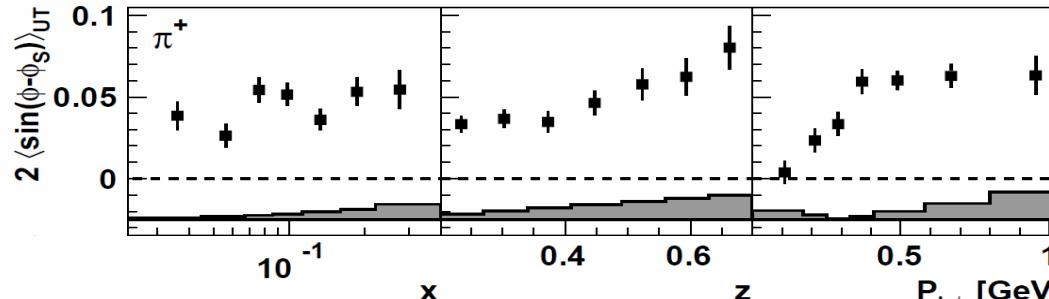
$$\Delta L = 1$$



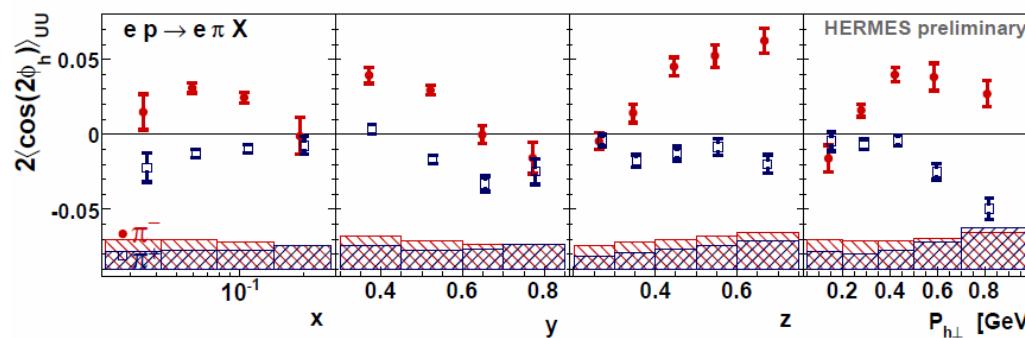
$$\Delta L = 1$$



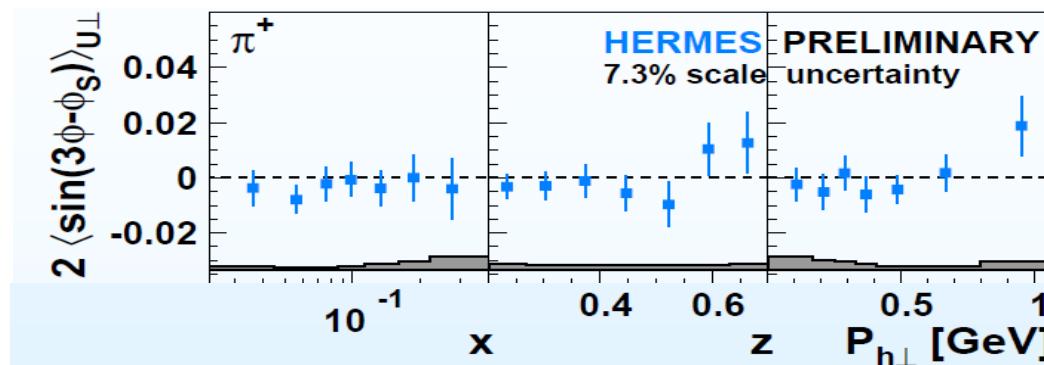
$$\Delta L = 2$$



Sivers

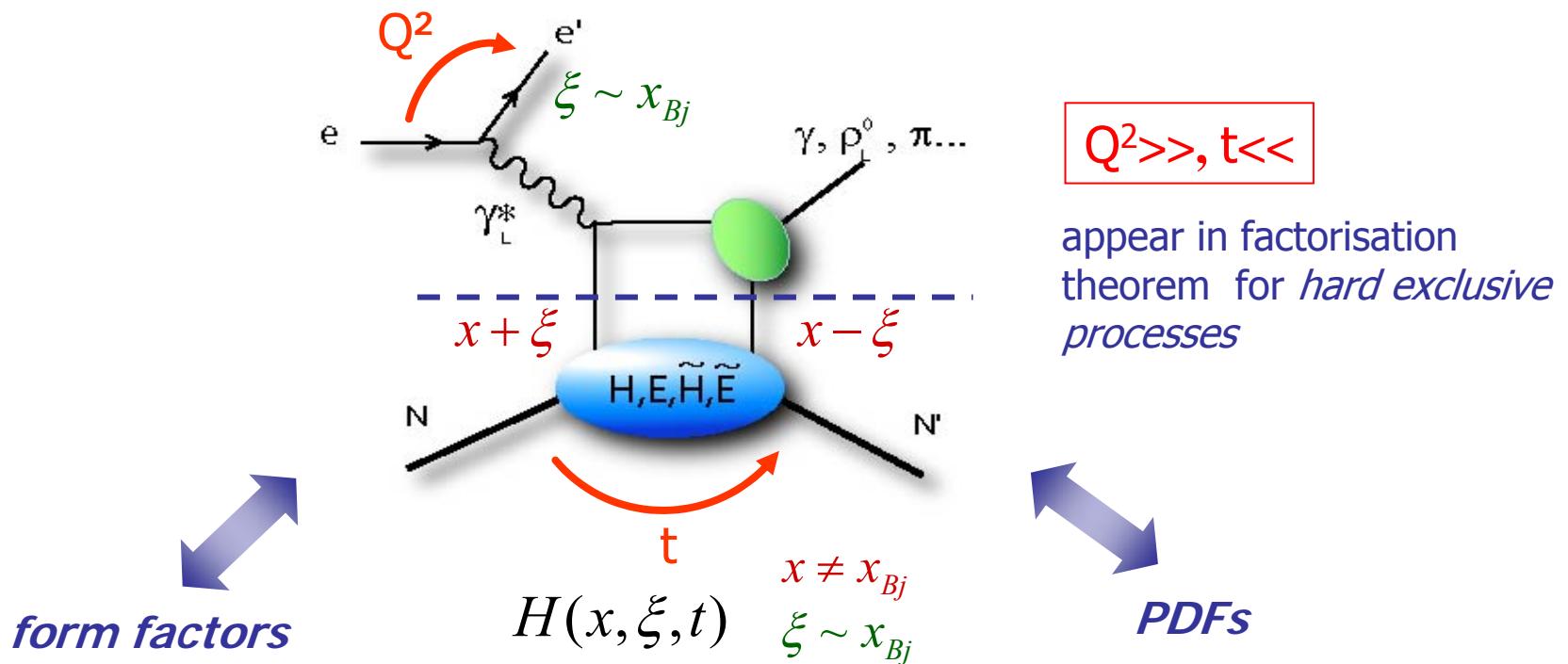


Boer-Mulders



pretzelosity

what do we know about GPDs ?



$$\sum_q e_q \int dx H^q(x, \xi, t) = F_1(t)$$

⋮

$$H^{q,g}(x, 0, 0) = q(x)$$

$$\tilde{H}^{q,g}(x, 0, 0) = \Delta q(x)$$

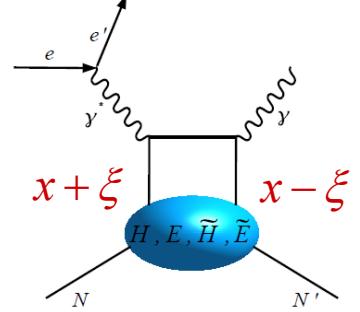
E, \tilde{E} : *nucleon helicity flip* → don't appear in DIS
→ new information !

GPDs: caveats

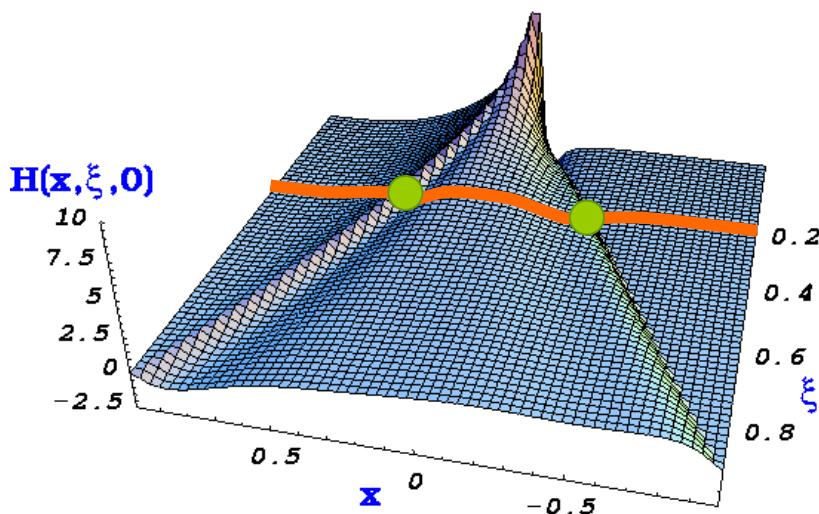
$$T_{\mu\nu} = [\mathcal{H}, \mathcal{E}, \tilde{\mathcal{H}}, \tilde{\mathcal{E}}](\xi, t, Q^2),$$

$$\boxed{\mathcal{F}(\xi, t, Q^2)} = \int_{-1}^1 dx C^-(\xi, x) \boxed{F(x, \xi, t, Q^2)}.$$

CFF hard scatt. part *GPD*



- x is mute variable (integrated over):
→ apart from cross-over trajectory ($\xi=x$) GPDs not directly accessible
- extrapolation $t \rightarrow 0$ is model dependent



cross sections & beam-charge asymmetry $\sim \text{Re}(T^{DVCS})$

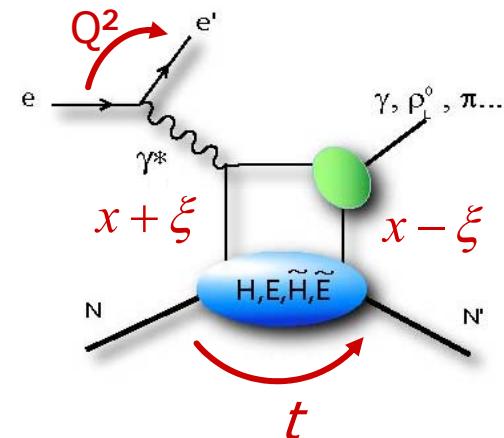
beam or target-spin asymmetries $\sim \text{Im}(T^{DVCS})$

→ double DVCS: $|x| < \xi$

attempts to constrain J_q

$$J_q = \lim_{t \rightarrow 0} \frac{1}{2} \int_{-1}^1 x dx \left[H^q(x, \xi, t) + E^q(x, \xi, t) \right]$$

GPD models: J_q free parameter in ansatz for E



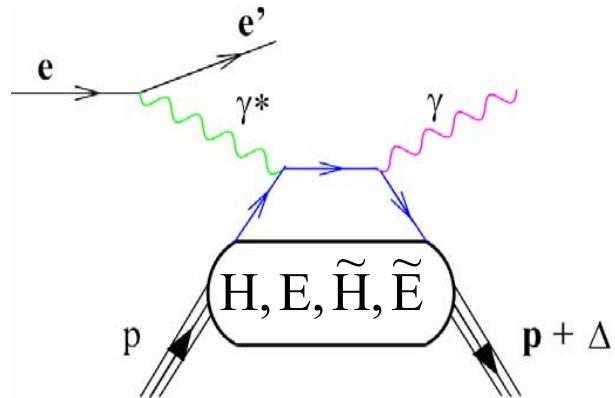
observables sensitive to E :

- DVCS A_{UT} : HERMES
- nDVCS A_{LU} : Hall A
- excl. $\rho^0 A_{UT}$: HERMES, COMPASS

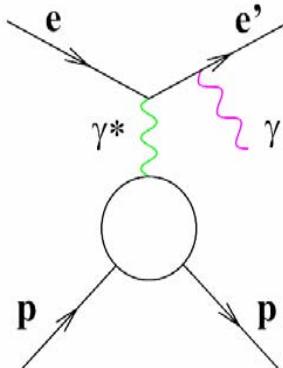
A_{UT}
 ↙ ↘
 beam target
 U, L U, T
 Unpolarised
 Longitudinally
 Transversely polarised

deeply virtual compton scattering

DVCS



Bethe-Heitler



@HERMES, JLab:

DVCS << Bethe-Heitler

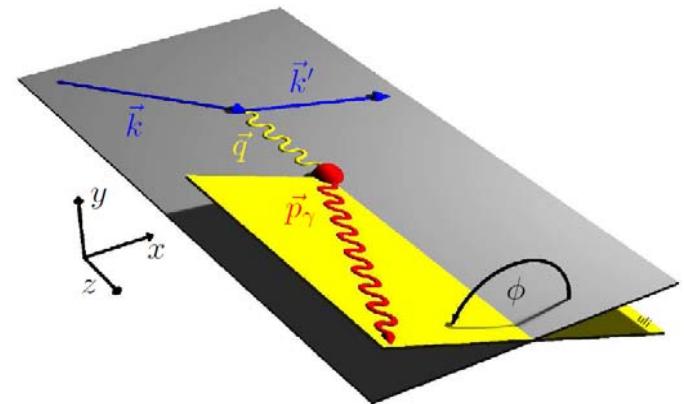
$$d\sigma \propto |\tau_{BH}|^2 + |\tau_{DVCS}|^2 + (\tau_{BH}^* \tau_{DVCS} + \tau_{DVCS}^* \tau_{BH})$$

DVCS interference term

$$d\sigma \propto |\tau_{\text{BH}}|^2 + |\tau_{\text{DVCS}}|^2 + (\tau_{\text{BH}}^* \tau_{\text{DVCS}} + \tau_{\text{DVCS}}^* \tau_{\text{BH}})$$

→ different charges: $e^+ e^-$

$$d\sigma(e^+, \phi) - d\sigma(e^-, \phi) \propto \text{Re}[F_1 \mathcal{H}] \cdot \cos \phi$$



DVCS interference term

$$d\sigma \propto |\tau_{\text{BH}}|^2 + |\tau_{\text{DVCS}}|^2 + (\tau_{\text{BH}}^* \tau_{\text{DVCS}} + \tau_{\text{DVCS}}^* \tau_{\text{BH}})$$

→ different charges: $e^+ e^-$

$$d\sigma(e^+, \phi) - d\sigma(e^-, \phi) \propto \text{Re}[F_1 \mathcal{H}] \cdot \cos \phi$$

→ polarisation observables:

- beam spin asymmetry A_{LU} :

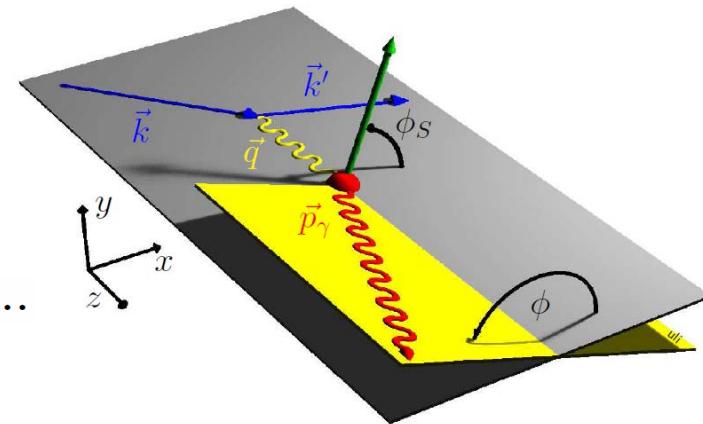
$$d\sigma(\vec{e}, \phi) - d\sigma(\overleftarrow{e}, \phi) \propto \text{Im}[F_1 \mathcal{H}] \cdot \sin \phi + \dots$$

- longitudinal target spin asymmetry A_{UL} :

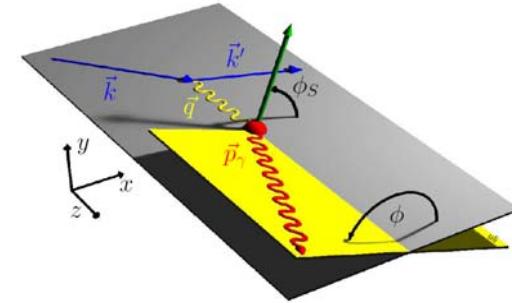
$$d\sigma(\overleftarrow{P}, \phi) - d\sigma(\overrightarrow{P}, \phi) \propto \text{Im}[F_1 \tilde{\mathcal{H}}] \cdot \sin \phi + \dots$$

- transverse target spin asymmetry A_{UT} :

$$\begin{aligned} d\sigma(\phi, \phi_S) - d\sigma(\phi, \phi_S + \pi) &\propto \text{Im}[F_2 \mathcal{H} - F_1 \mathcal{E}] \cdot \sin(\phi - \phi_S) \cos \phi \\ &+ \text{Im}[F_2 \tilde{\mathcal{H}} - F_1 \tilde{\mathcal{E}}] \cdot \sin(\phi - \phi_S) + \dots \end{aligned}$$



DVCS A_{UT}



sensitivity to E (J_q) from both interference and DVCS² term:

$$\sigma(\phi, P_\ell, S_T) = \sigma_{UU}(\phi) \times [1 + S_T \mathcal{A}_{UT}^{\text{DVCS}}(\phi, \phi_s) + S_T e_\ell \mathcal{A}_{UT}^I(\phi, \phi_s) + e_\ell \mathcal{A}_C(\phi)]$$

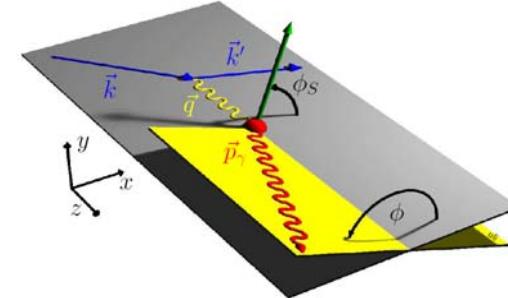
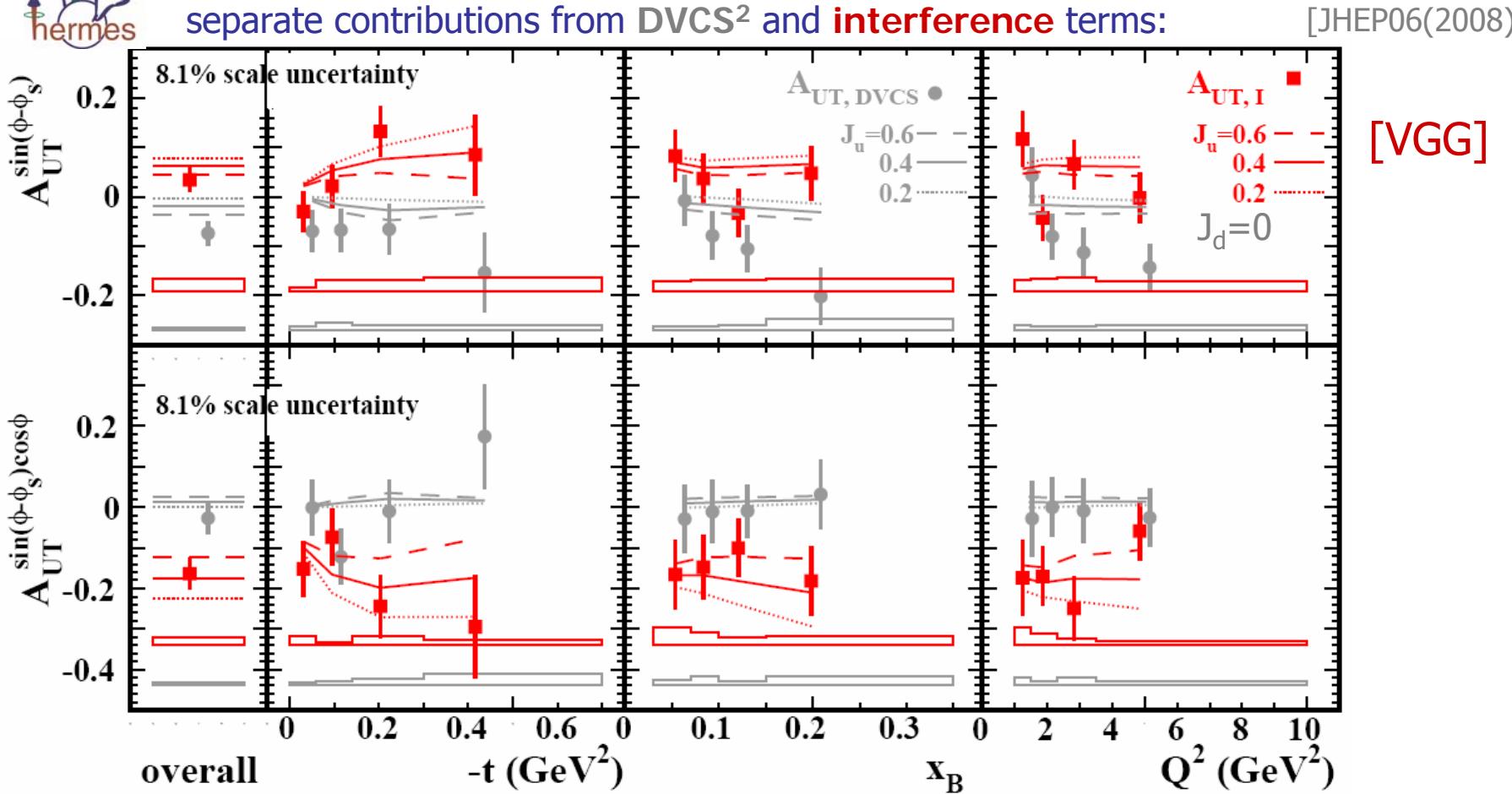
$$\begin{aligned} \mathcal{A}_{UT}^I(\phi, \phi_s) &= \sum_{n=0}^2 A_{UT, I}^{\sin(\phi - \phi_s) \cos(n\phi)} \sin(\phi - \phi_s) \cos(n\phi) \\ &+ \sum_{n=1}^2 A_{UT, I}^{\cos(\phi - \phi_s) \sin(n\phi)} \cos(\phi - \phi_s) \sin(n\phi) \end{aligned}$$

analogous modulations for DVCS² term

$n = 0, 1$ terms found to be most sensitive to values of $J_u \rightarrow$

attempts to constrain J_q

J_q free parameter in ansatz for E



attempts to constrain J_q

Hall-A

J_q free parameter in ansatz for E

difference of polarised cross sections on LH_2 & $LD_2 \rightarrow nDVCS$: [PRL99(2007)]

$$[\mathcal{C}_n^I] = F_1 \mathcal{H} + \xi(F_1 + F_2) \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E}$$

p target: $\mathcal{C}_n^I \sim \mathcal{H}$

n target: $\mathcal{C}_n^I \sim \mathcal{E} \quad \rightarrow F_1$ small

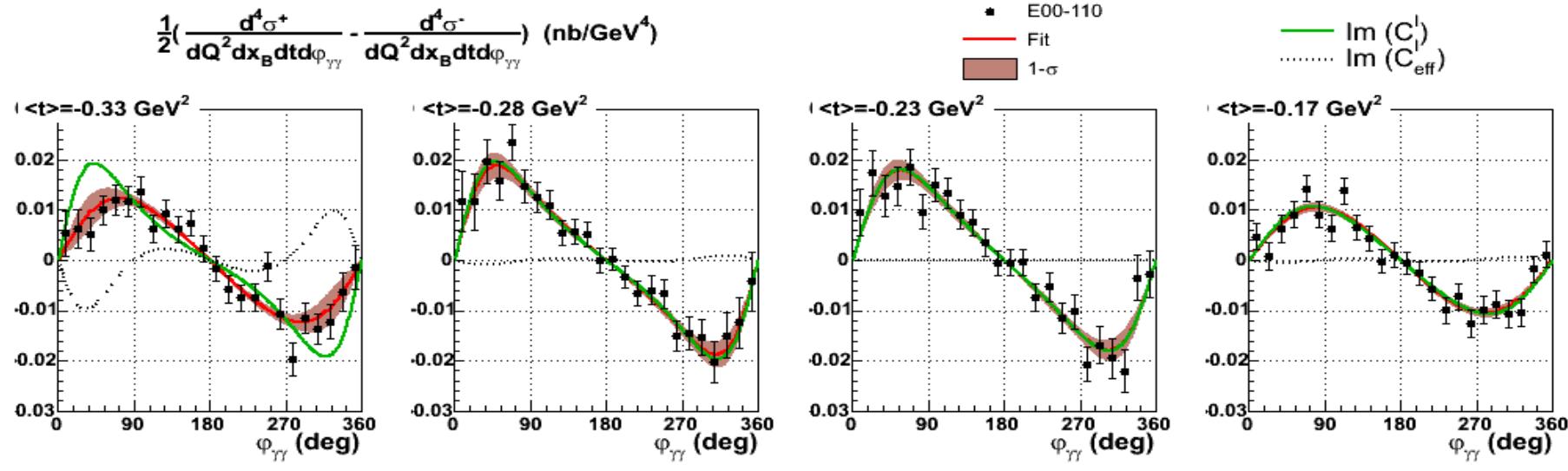
\rightarrow cancellation between u and d quark pol. pdfs in $\tilde{\mathcal{H}}$

attempts to constrain J_q

Hall-A

J_q free parameter in ansatz for E

difference of polarised cross sections on LH_2 & $\text{LD}_2 \rightarrow \text{nDVCS}$: [PRL99(2007)]



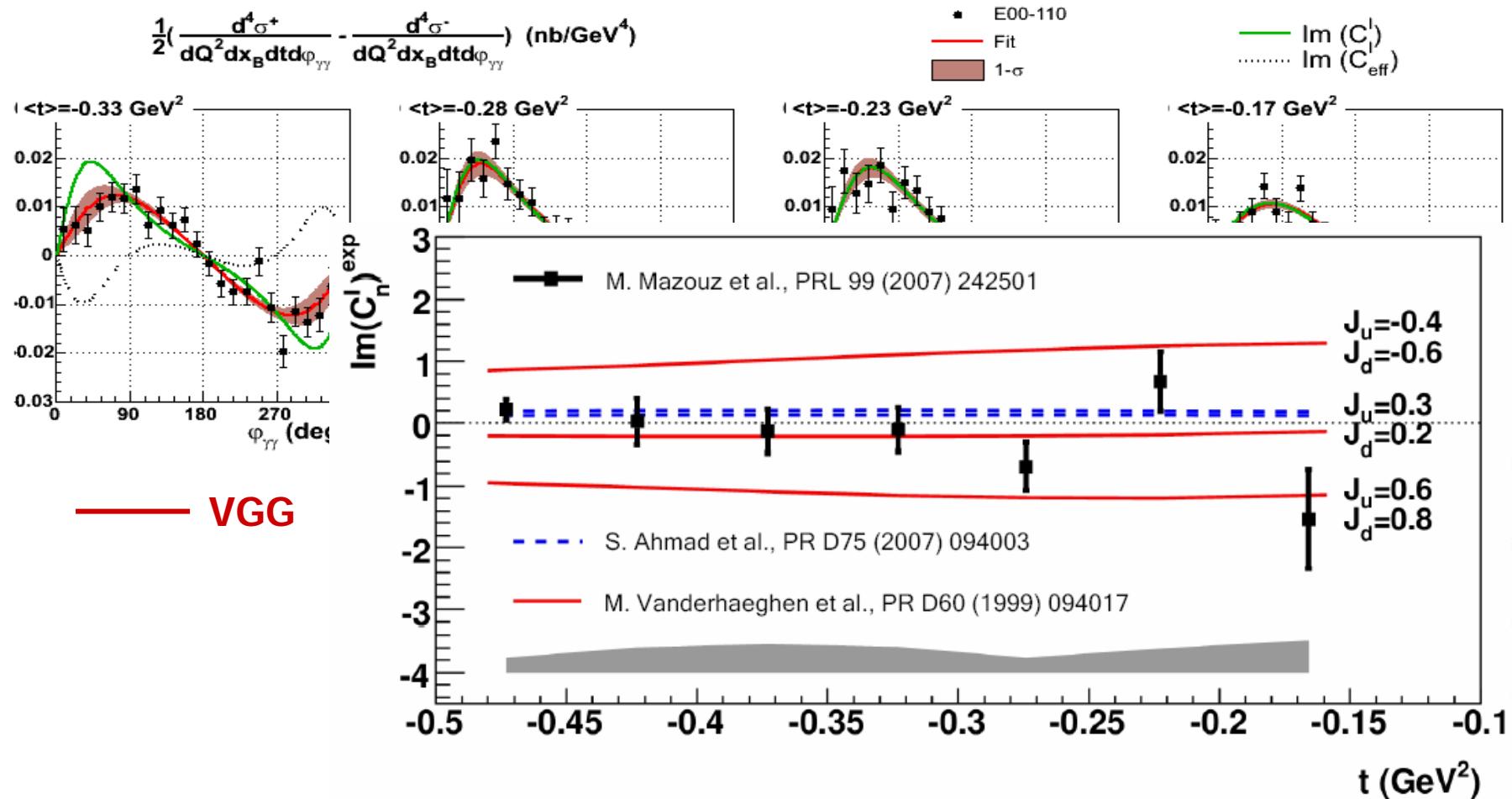
— VGG : tw-2 CFF

attempts to constrain J_q

Hall-A

J_q free parameter in ansatz for E

difference of polarised cross sections on LH_2 & $\text{LD}_2 \rightarrow \text{nDVCS}$: [PRL99(2007)]



a word about 'user friendly' GPD models

VGG: [Vanderhaegen, Guichon, Guidal 1999]

- double distributions [Radyshkin]; factorised or regge-inspired t-dependence
- D-term to restore full polynomiality
- skweness depending on free parameters b_{val} & b_{sea}
- includes tw-3 (WW approx)

Dual: [Guzey, Teckentrup 2006, 2009]

- GPDs based on infinite sum of t channel resonances (minimal: truncated $k=[0,2]$)
- factorised or regge-inspired t-dependence
- tw-2 only

→ more models & new approaches [... an incomplete list]

- polynomials [Belitsky et al.(00), Liuti et al.(07), Moutarde(09)]
- analytical [Belitsky, Muller, Kirchner(01)]
- dispersion integral fits & flexible GPD modelling [Kumericki, Muller(08,09)]

a word about 'user friendly' GPD models

VGG: [Vanderhaegen, Guichon, Guidal 1999]

- double distributions [Radys]
 - D-term to restore full poly
 - skweness depending on fr
 - includes tw-3 (WW approx)
- describes well beam charge & target spin asymmetries
→ fails for beam spin asymm. & cross sections
→ charge asymm. favours 'no D-term' ← contradicts
 χ QSM & lattice results

Dual: [Guzey, Teckentrup 2006, 2009]

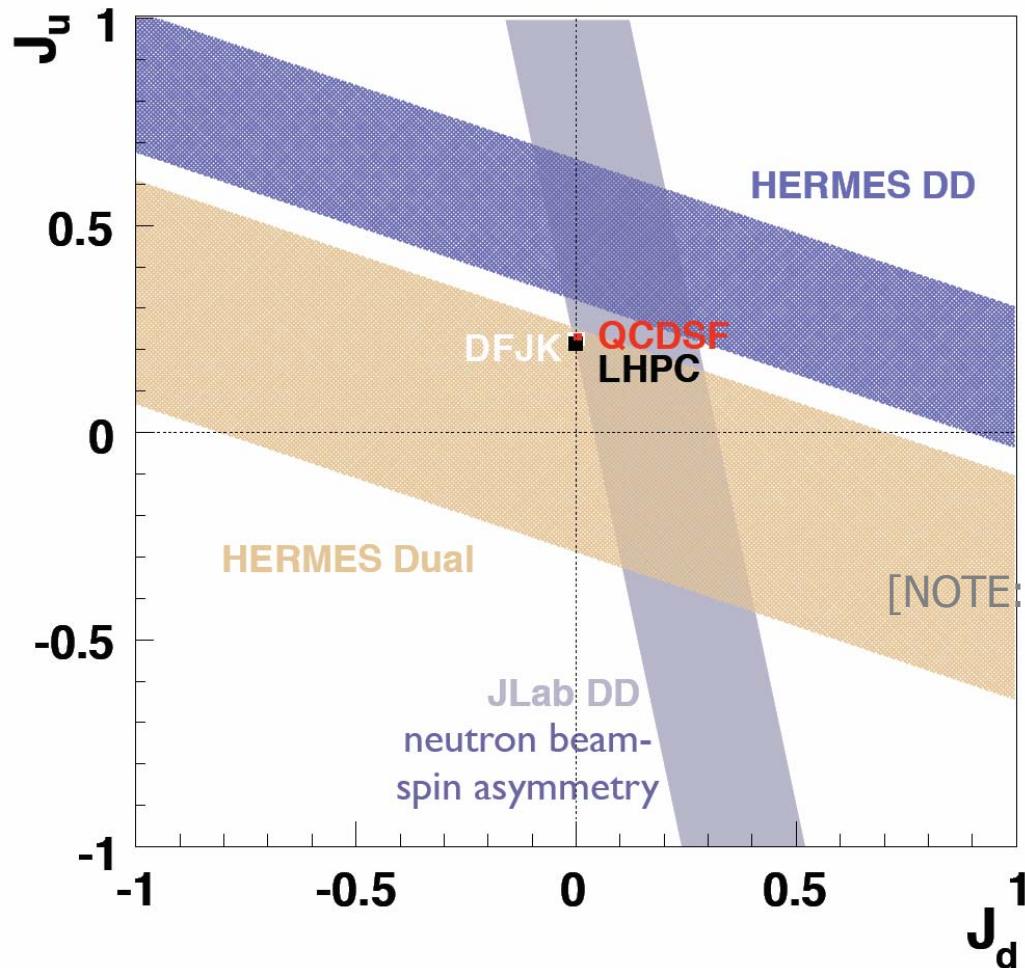
- GPDs based on infinite sum
 - factorised or regge-inspire
 - tw-2 only
- describes well kinematic dependencies of beam charge & beam spin asymmetries
→ after correction in calculations: magnitude off by factor 2-4

→ more models & new ap

- polynomials [Belitsky et al.(00), Liuti et al.(07), Moutarde(09)]
- analytical [Belitsky, Muller, Kirchner(01)]
- dispersion integral fits & flexible GPD modelling [Kumericki, Muller(08,09)]

...nevertheless: constraining J_q

J_q free parameter in ansatz for E



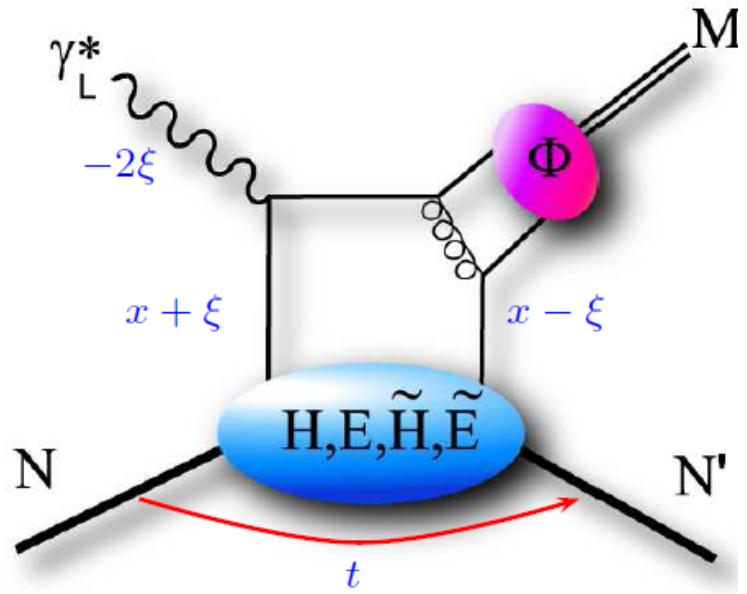
→ highly model
dependent extraction !!!

[VGG]

[NOTE: uncorrected Dual !]

→ data are free to be re-
used at any time with
new models ☺

exclusive ρ^0 production



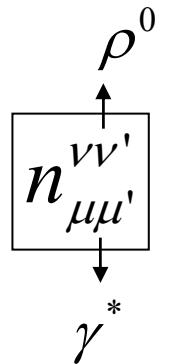
$$A_{UT}^{\gamma_L^*}(\phi, \phi_s) \propto \frac{\text{Im}(\mathcal{E}_\rho^* \mathcal{H}_\rho)}{|\mathcal{H}_\rho|^2} \propto \left| \frac{\mathcal{E}_\rho}{\mathcal{H}_\rho} \right|$$

exclusive ρ^0 production

after the full glory of transverse SDME extraction [formalism: M. Diehl (2007)] :

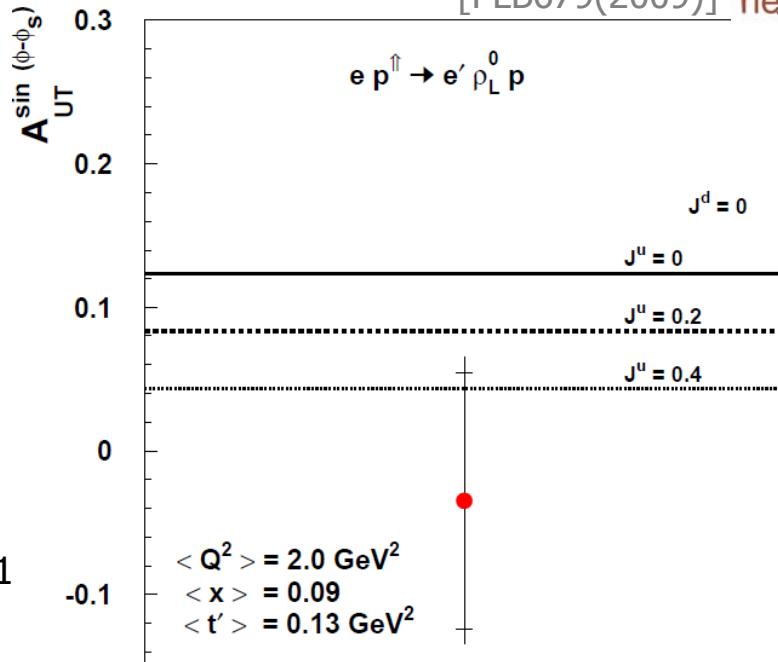
$(\gamma_L^* \rightarrow \rho_L^0)$:

$$A_{UT}^{\gamma^*}(\phi, \phi_s) = \frac{\text{Im } n_{00}^{00}}{u_{00}^{00}}$$



$$\mu, \nu = 0, \pm 1$$

long.pol: 0
transv.pol: ± 1

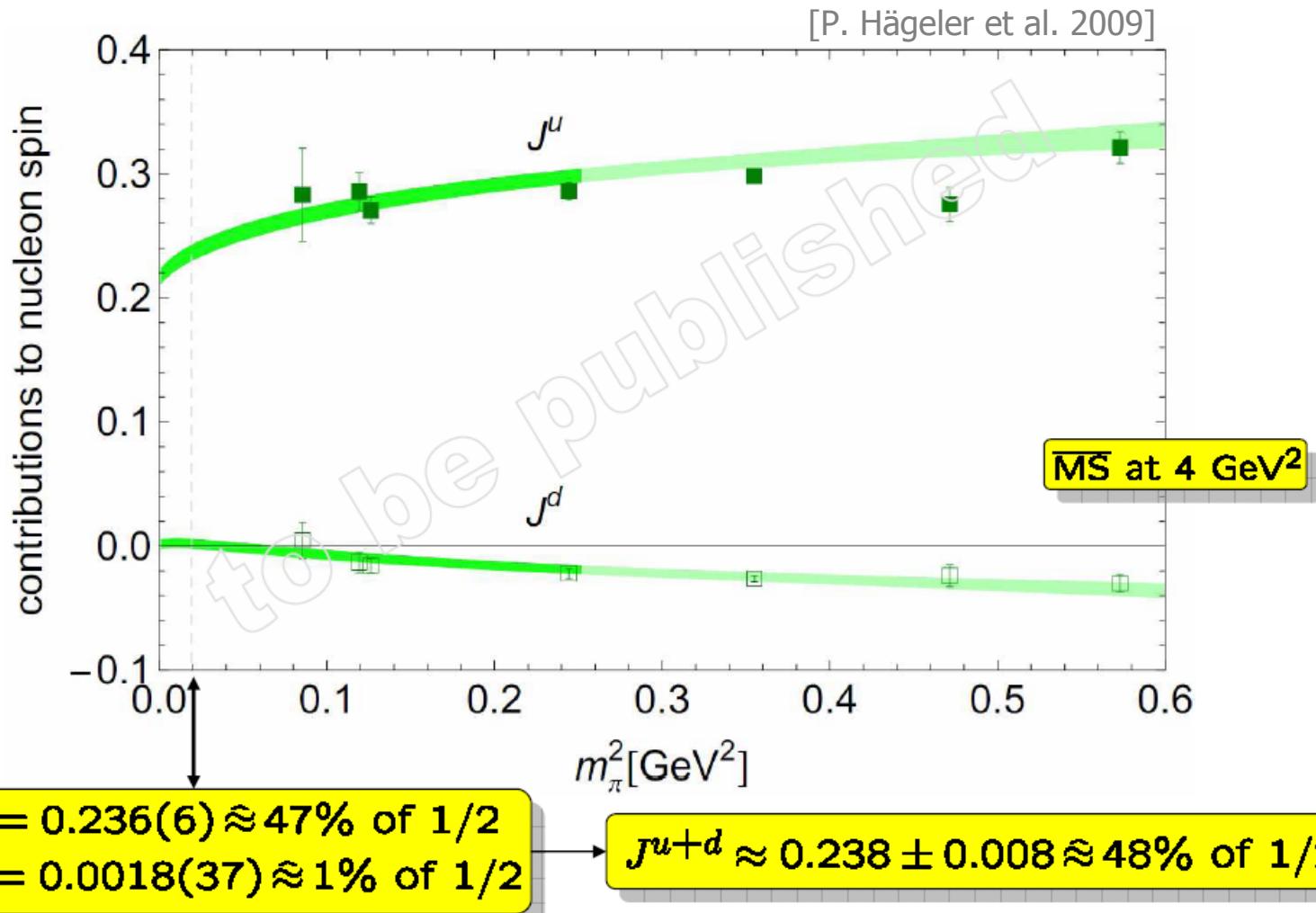


[GPD model:
Ellighaus et al. (2004)]

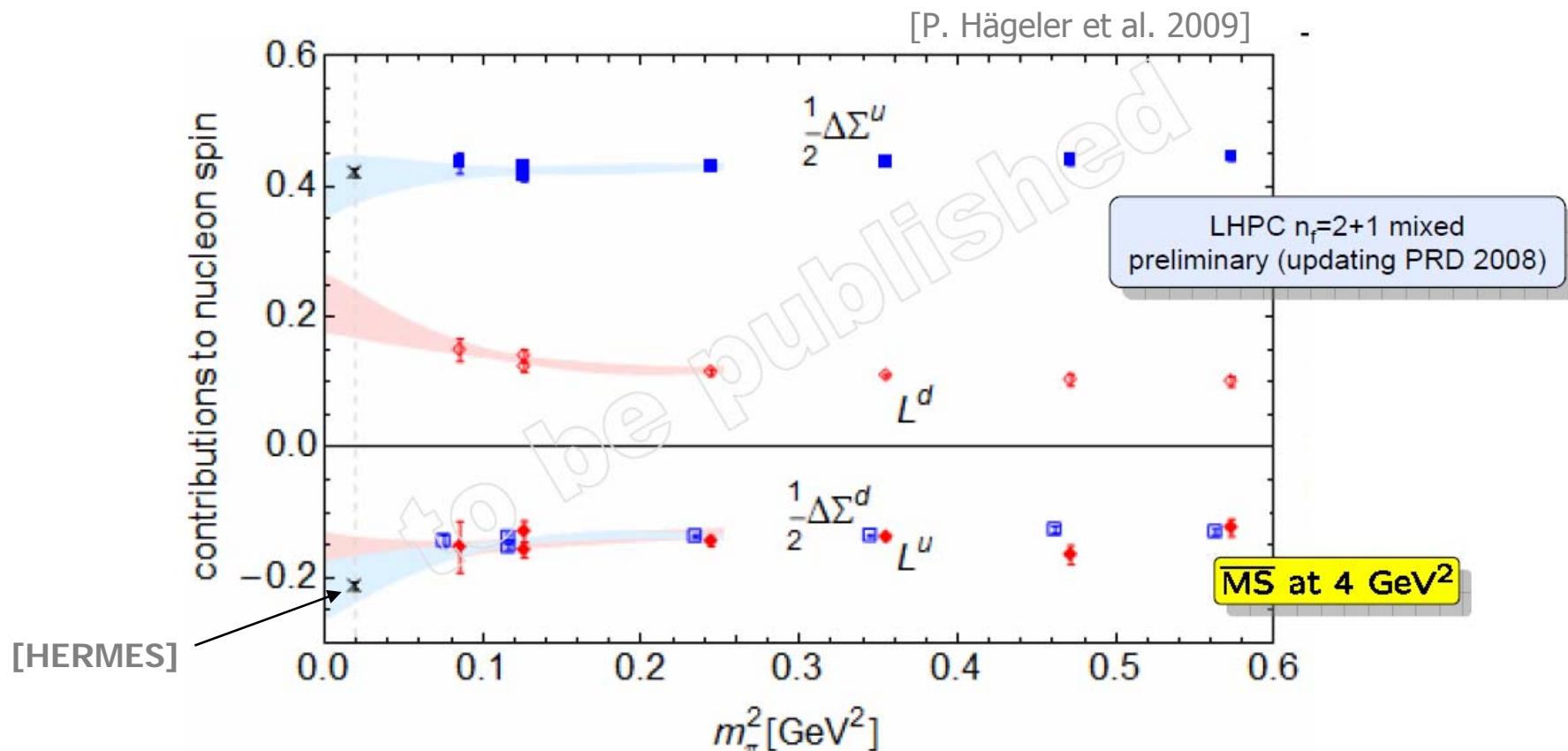
overall

- more data coming: COMPASS , CLAS12 with transverse target
- more models: Goloskokov, Kroll (09)

lattice's opinion about \not{q}

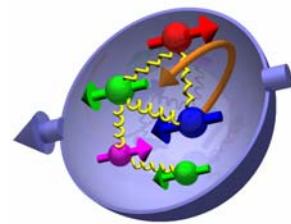


Lattice's opinion about $J^q \rightarrow L^q$



$L^d \approx -L^u \approx 0.185 \pm 0.06 \approx 36\% \text{ of } 1/2$
 $L^{u+d} \approx 0.030 \pm 0.012 \approx 6\% \text{ of } 1/2$

conclusion



presence of OAM w/o debate → how to measure it ?

■ GPDs: only known frame work to *quantify* OAM [Ji - SR]

we got an idea how to measure it but still a long way to go:

→ more data needed over a much wider kinematic range

→ call for more sophisticated GPD models & new approaches

■ prior to any interpretation of data: what is OAM ? [M. Burkardt, ...]

■ complementary information from TMDs :

role of transverse momenta & *spin-orbit correlations*

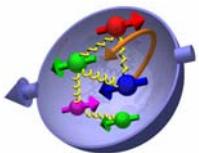
GPDs – a much wider concept: nucleon *tomography*

→ correlated information on longitudinal momentum fraction of quarks and their spatial distribution in the transverse plane → multi-D picture

relations GPDs ↔ TMDs ?

perspectives

hunting the OAM



contribution to nucleon spin:

- determination of $\Delta\Sigma$ and ΔG → missing piece attributed to OAM

quest for

ΔG :

- from scaling violation of g_1
- charm production & high pT hadrons over wide x_B range



EIC @highest possible energies



'direct' measurement via Ji-SR (GPDs)

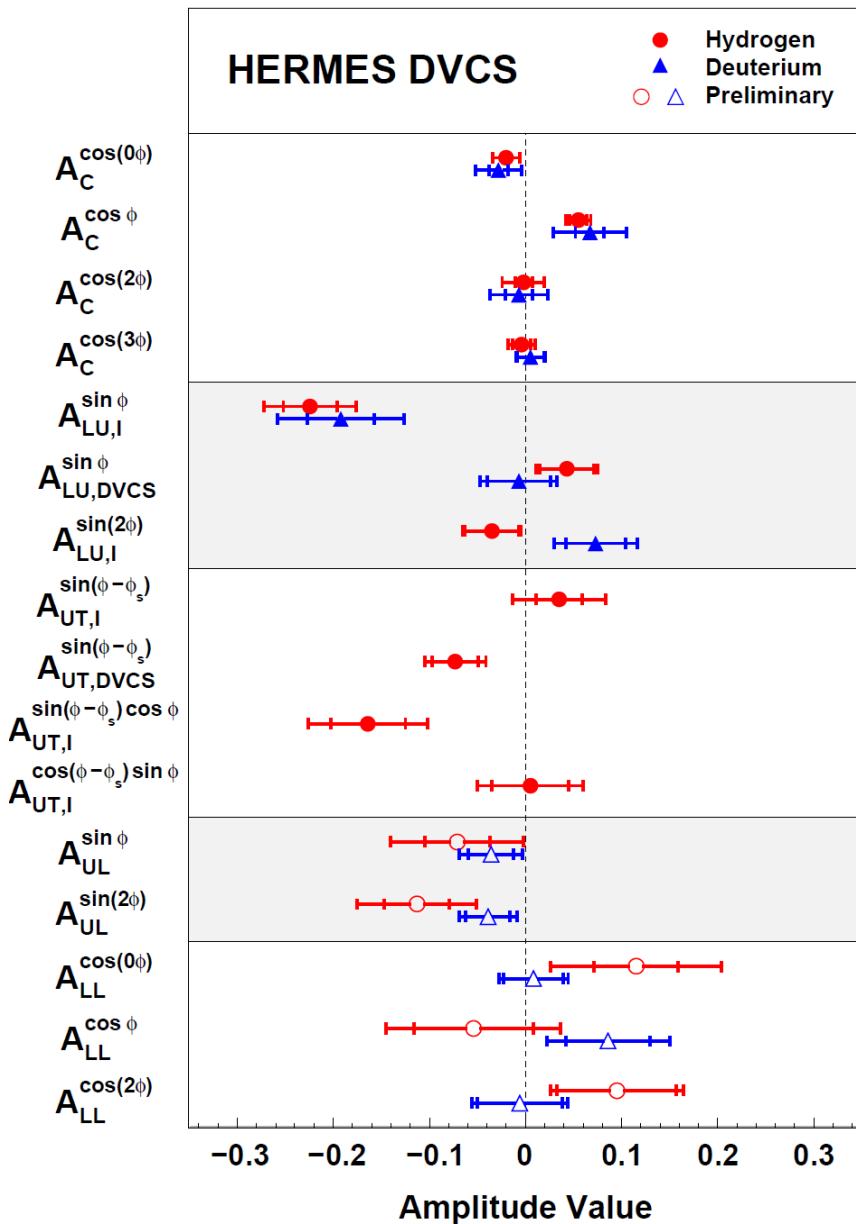


spin-orbit correlations from TMDs

EIC @modest energies, high lumi (symmetric beams)

JLab12

backup



→ beam charge asymmetry

$$\text{Re}\mathcal{H}$$

→ beam spin asymmetry

$$\text{Im}\mathcal{H}$$

→ transverse target spin asymm.

$$\text{Im}(\mathcal{HE})$$

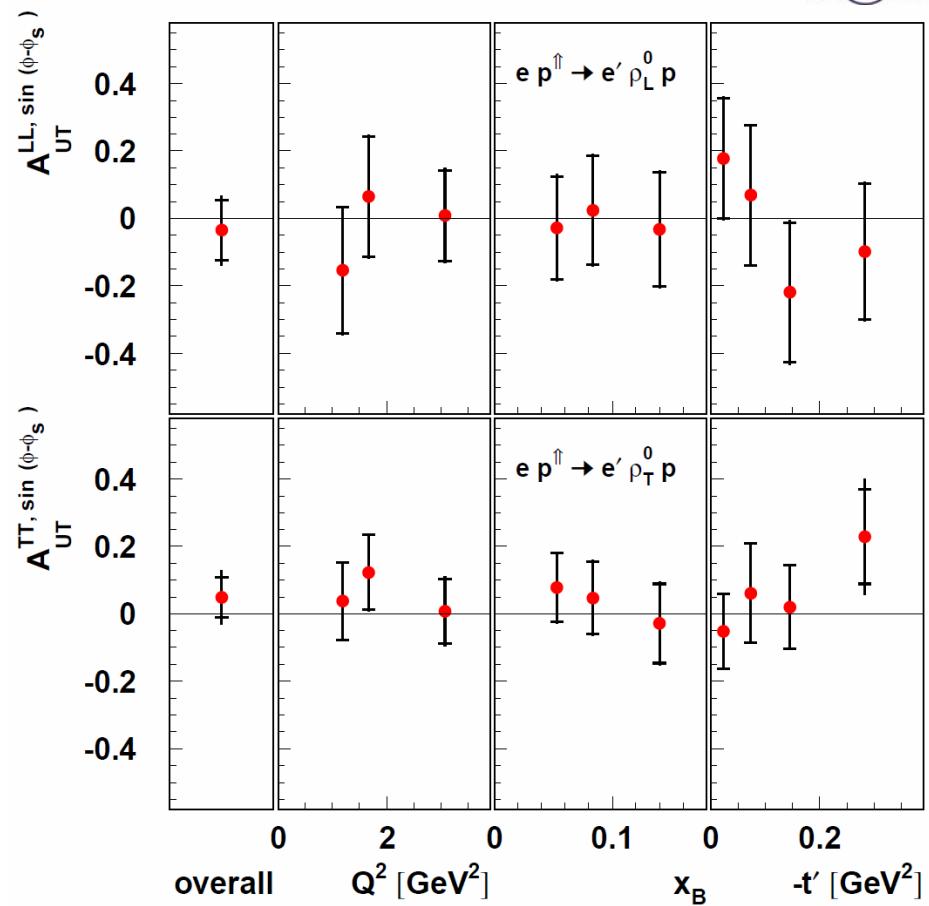
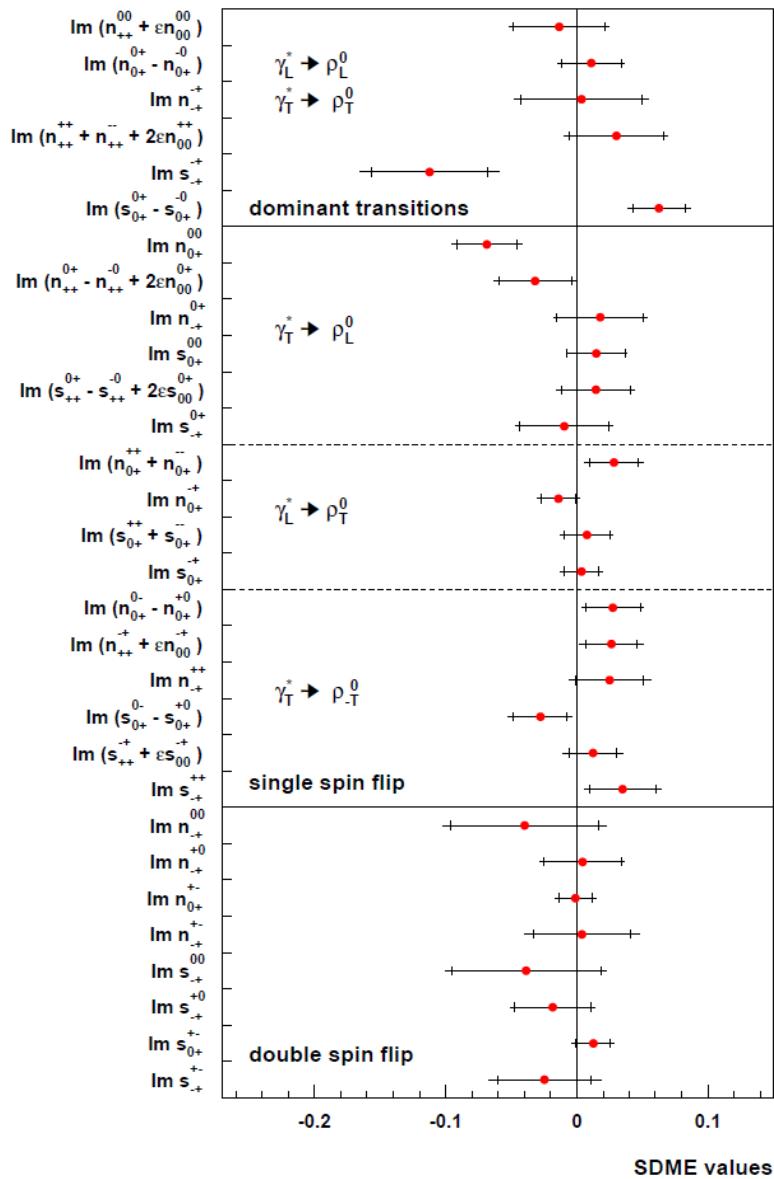
→ longitudinal target spin asymm.

$$\text{Im}\tilde{\mathcal{H}}$$

→ double spin asymmetry

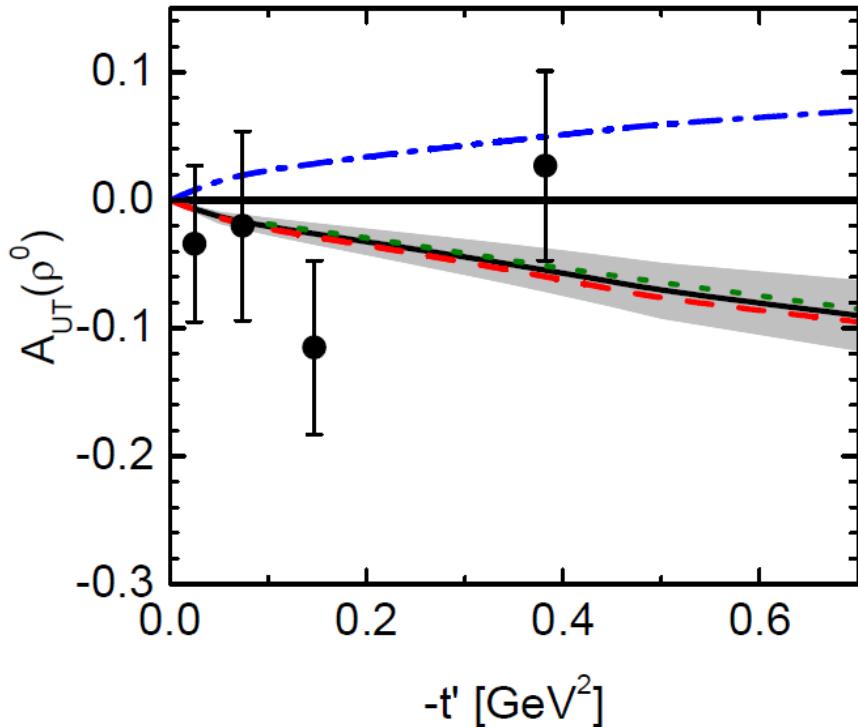
$$\text{Re}\tilde{\mathcal{H}}$$

exclusive ρ^0 production



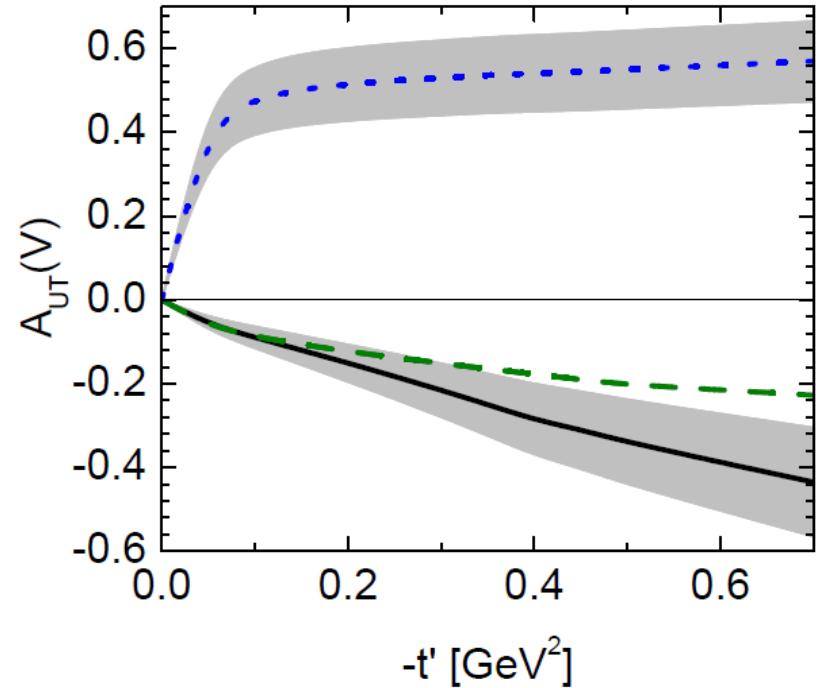
exclusive meson production

GPD model: Goloskokov, Kroll



$$W = 5 \text{ GeV} \quad Q^2 = 3 \text{ GeV}^2$$

variant 1, 2, 3, 4



variant 1 for ω , ρ^+ , K^{*0}