

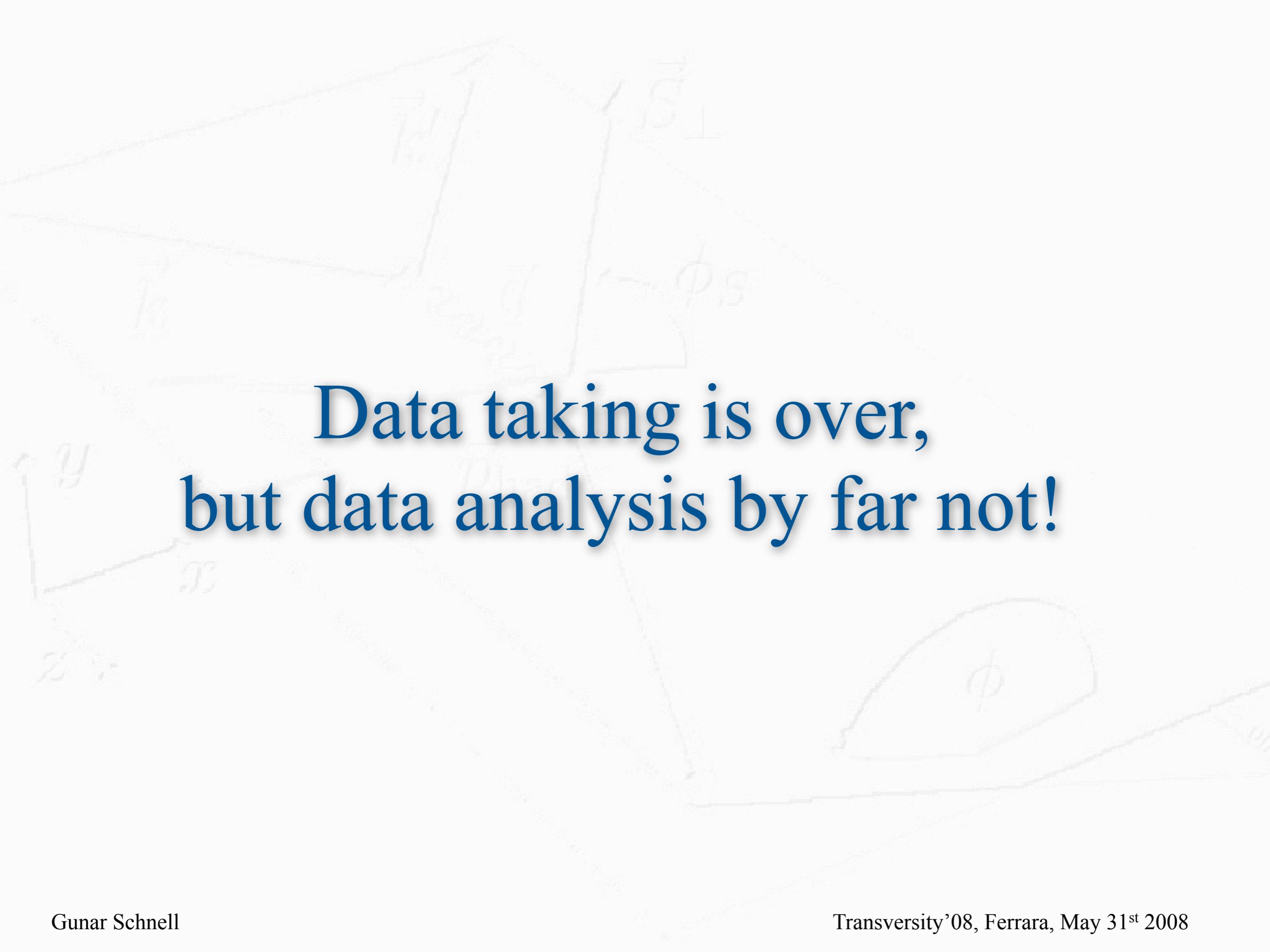
Transversity and Beyond at HERMES

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Transversity 2008
Ferrara, May 31st 2008

June 30th, 2007 (around midnight)





Data taking is over,
but data analysis by far not!

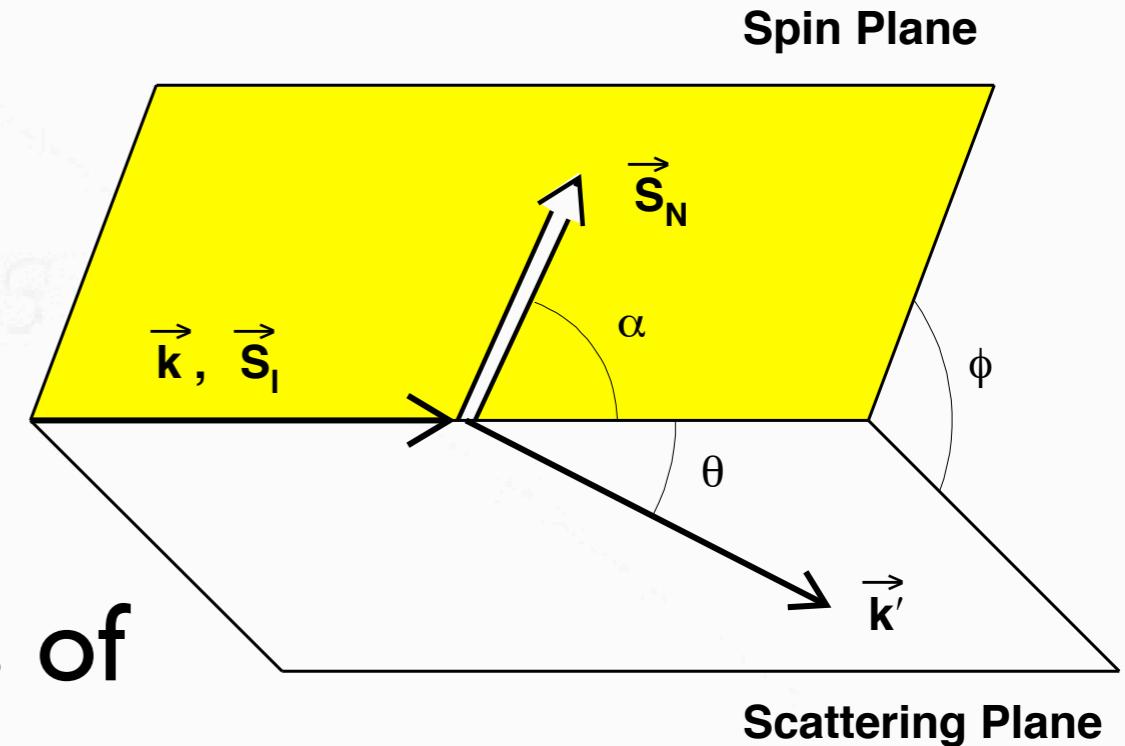
Inclusive DIS

$$\frac{d^2\sigma(s, S)}{dx \, dQ^2} = \frac{2\pi\alpha^2 y^2}{Q^6} \mathbf{L}_{\mu\nu}(s) \mathbf{W}^{\mu\nu}(S)$$

Lepton Tensor Hadron Tensor

parametrized in terms of

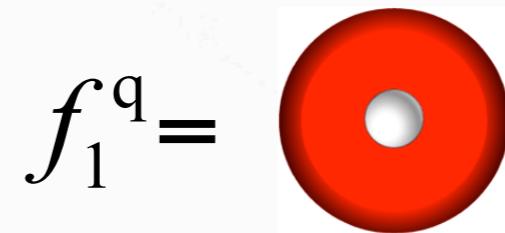
Structure Functions



$$\begin{aligned}
 \frac{d^3\sigma}{dxdy d\phi} \propto & \frac{y}{2} F_1(x, Q^2) + \frac{1 - y - \gamma^2 y^2 / 4}{2xy} F_2(x, Q^2) \\
 & - P_l P_T \cos \alpha \left[\left(1 - \frac{y}{2} - \frac{\gamma^2 y^2}{4} \right) g_1(x, Q^2) - \frac{\gamma^2 y}{2} g_2(x, Q^2) \right] \\
 & + P_l P_T \sin \alpha \cos \phi \gamma \sqrt{1 - y - \frac{\gamma^2 y^2}{4}} \left(\frac{y}{2} g_1(x, Q^2) + g_2(x, Q^2) \right)
 \end{aligned}$$

Parton-Model Interpretation of Structure Functions

$$F_1(x) = \frac{1}{2} \sum_q e_q^2 f_1^q(x)$$



$$F_2(x) = x \sum_q e_q^2 f_1^q(x)$$

$$g_1(x) = \frac{1}{2} \sum_q e_q^2 g_1^q(x)$$



$$g_2(x) = 0$$



quark-spin contribution to nucleon helicity

Closer Look at g_2

- higher-twist, thus
 - no probabilistic interpretation
 - probes parton correlations
- at LO: $g_2 = \frac{1}{2} \sum e_q^2 g_T(x)$

$$\int_0^1 dx x g_2(x, Q^2) = \frac{1}{3} (-a_2(Q^2) + d_2(Q^2))$$

second moment of g_1

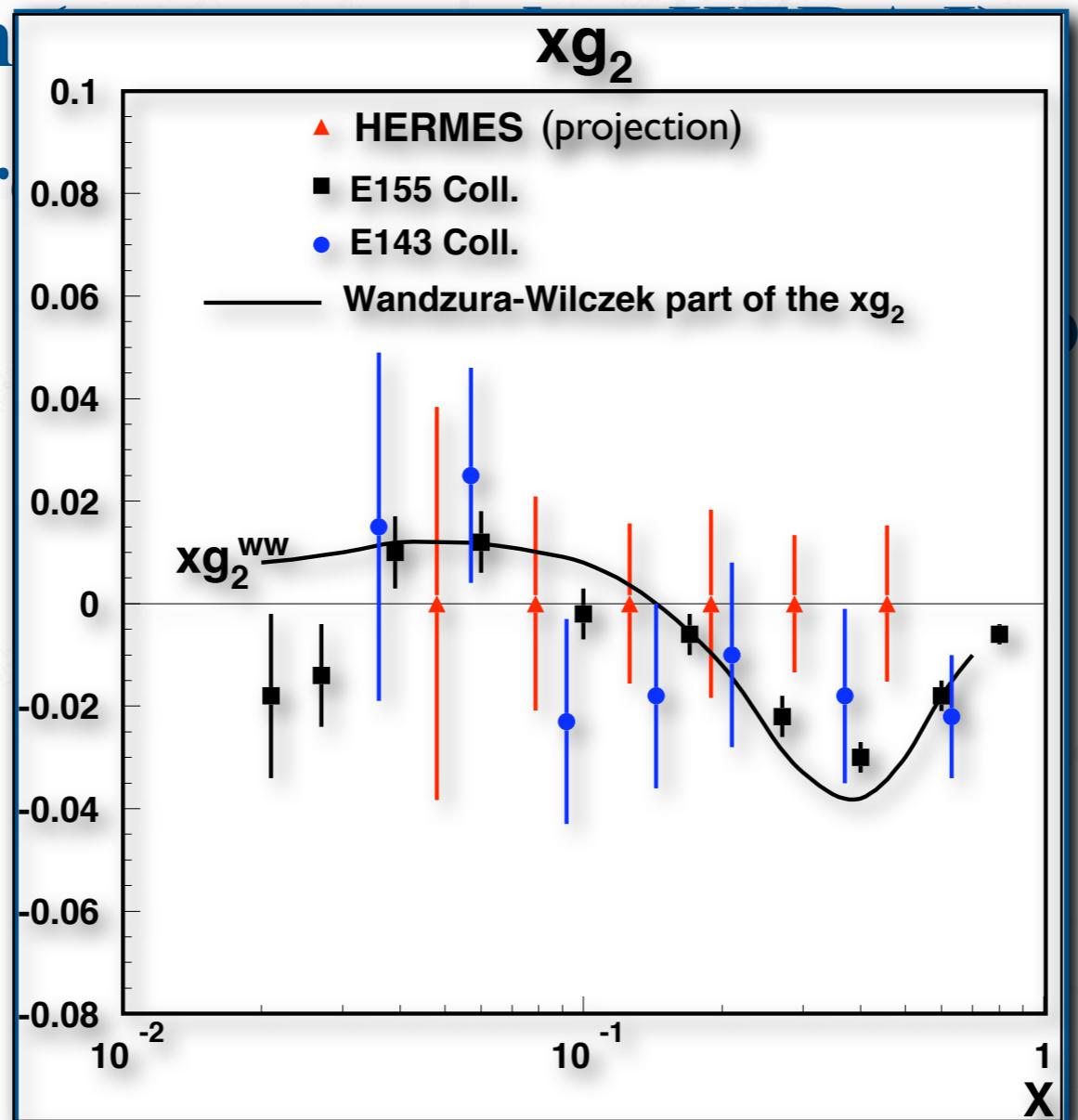
sensitive to a quark *and* a gluon amplitude

What HERMES can do for g_2

- unfortunately, HERA II with **no or only low beam polarization (as compared to HERA I)**
 - ➡ **low figure of merit**
 - ➡ **expected precision not comparable to E155**

What HERMES can do for g_2

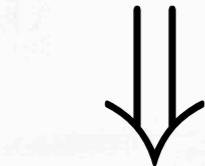
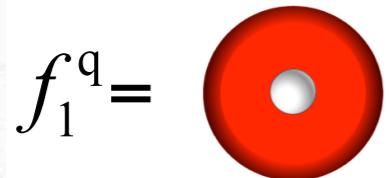
- unfortunately, HERA II with no or only low beam polarization
 - ➡ low figure
 - ➡ expected



E155

... back to parton distributions ...

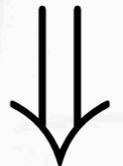
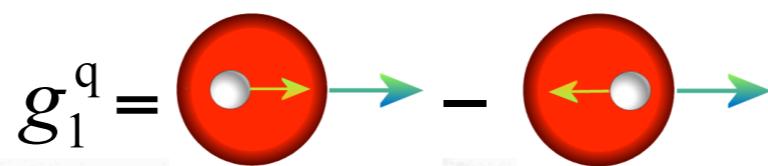
Quark Structure of the Nucleon



Unpolarized quarks
and nucleons

$f_1^q(x)$: spin averaged
(well known)

⇒ Vector Charge



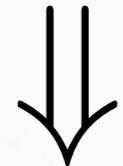
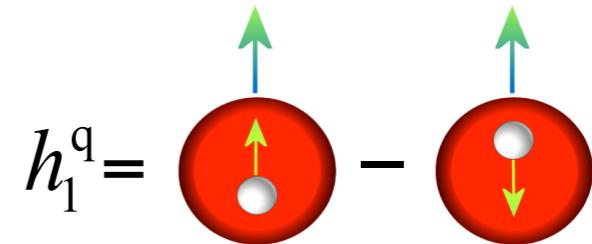
Longitudinally
polarized quarks
and nucleons

$g_1^q(x)$: helicity
difference (known)

Alexei & Co., THANKS!!

$$\langle PS | \bar{\Psi} \gamma^\mu \Psi | PS \rangle = \int dx (f_1^q(x) - f_1^{\bar{q}}(x))$$

$$\langle PS | \bar{\Psi} \gamma^\mu \gamma_5 \Psi | PS \rangle = \int dx (g_1^q(x) + g_1^{\bar{q}}(x))$$



Transversely
polarized quarks
and nucleons

$h_1^q(x)$: transversity
~~(hardly known!)~~

SSAs in One-Hadron Production

$$\begin{aligned}
 A_{UT}(\phi, \phi_S) &= \frac{1}{\langle |S_\perp| \rangle} \frac{N_h^\uparrow(\phi, \phi_S) - N_h^\downarrow(\phi, \phi_S)}{N_h^\uparrow(\phi, \phi_S) + N_h^\downarrow(\phi, \phi_S)} \\
 &\sim \sin(\phi + \phi_S) \sum_q e_q^2 \mathcal{I} \left[\frac{k_T \hat{P}_h}{M} f_{1T}^{\perp, q}(x, p_T^2) H_1^{\perp, q}(z, k_T^2) \right] \\
 &\quad + \sin(\phi - \phi_S) \sum_q e_q^2 \mathcal{I} \left[\frac{k_T \hat{P}_{h\perp}}{M} f_{1T}^{\perp, q}(x, p_T^2) D_1^q(z, k_T^2) \right]
 \end{aligned}$$

cf. L. Pappalardo's talk for results

$\mathcal{I}[\dots]$: convolution integral over initial (p_T) and final (k_T) quark transverse momenta

\Rightarrow 2D Max.Likelihd. fit to get Collins and Sivers amplitudes:

$$PDF(2\langle \sin(\phi \pm \phi_S) \rangle_{UT}, \dots, \phi, \phi_S) = \frac{1}{2} \{ 1 + P_T(2\langle \sin(\phi \pm \phi_S) \rangle_{UT} \sin(\phi \pm \phi_s) + \dots) \}$$

Resolving the Convolution Integral

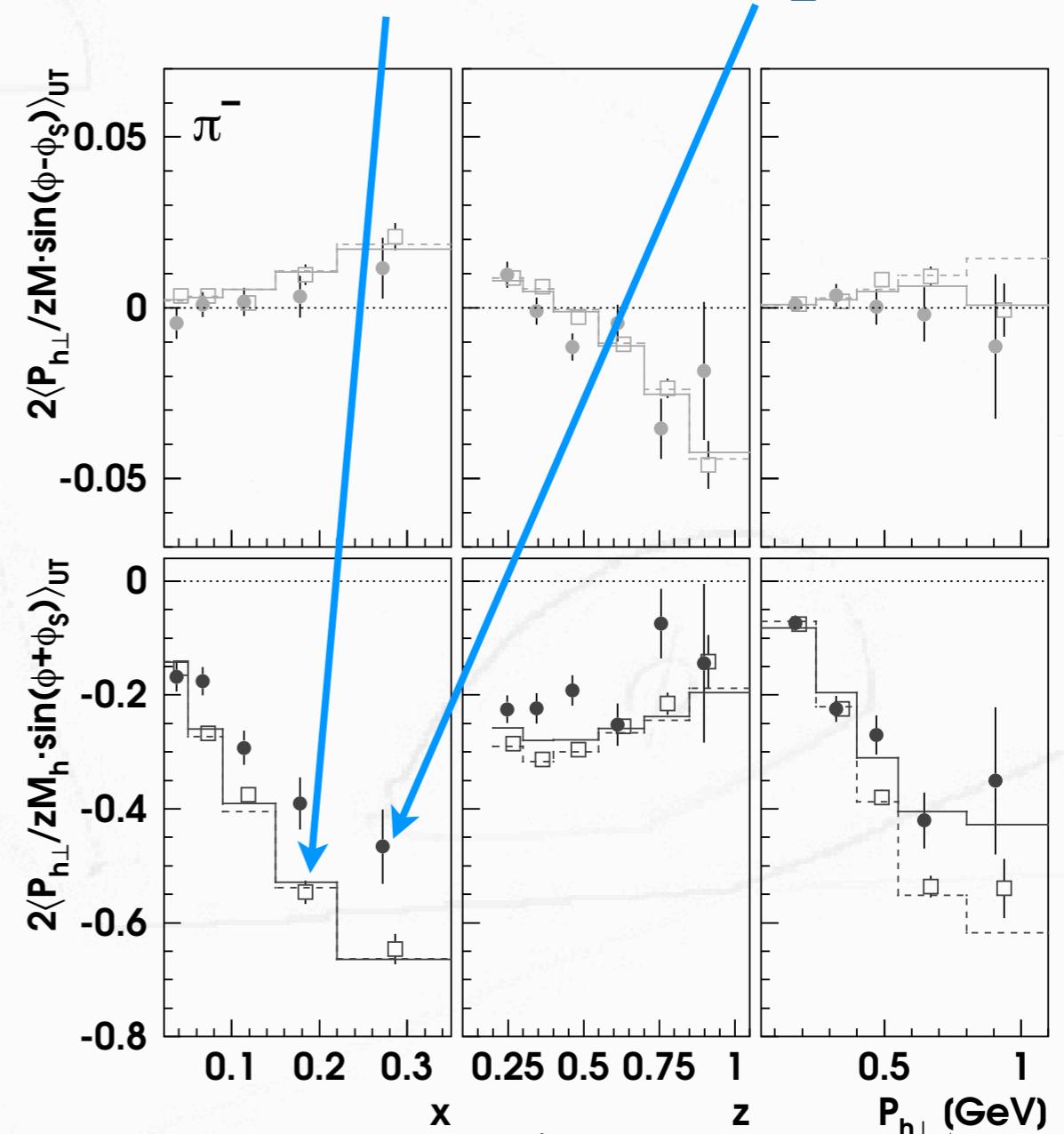
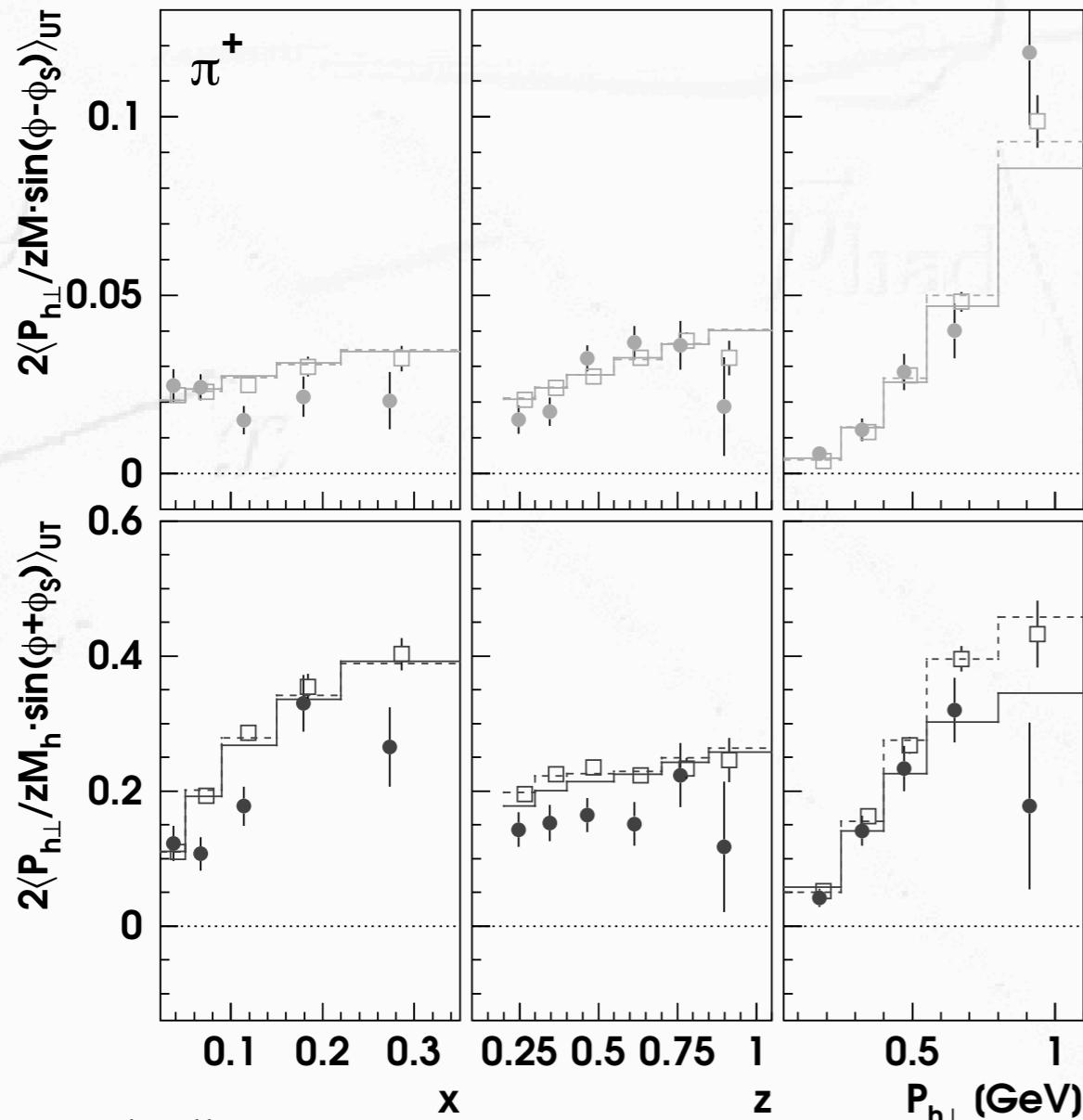
Weight with transverse hadron momentum $P_{h\perp}$ to resolve convolution:

$$\begin{aligned}\tilde{A}_{UT}(\phi, \phi_S) &= \frac{1}{\langle S_\perp \rangle} \frac{\sum_{i=1}^{N^+} P_{h\perp,i} - \sum_{i=1}^{N^-} P_{h\perp,i}}{N^+ + N^-} \\ &\sim \sin(\phi + \phi_S) \cdot \sum_q e_q^2 \, h_1^q(x) \, z \, H_1^{\perp(1),q}(z) \quad (1): \text{ } p_T^2/k_T^2\text{-moment of} \\ &\quad \text{distribution / fragmentation} \\ &- \sin(\phi - \phi_S) \cdot \sum_q e_q^2 \, f_{1T}^{\perp(1),q}(x) \, z \, D_1^q(z) \quad \text{function} \\ &+ \dots\end{aligned}$$

- idea goes back to Kotzinian & Mulders [Phys. Rev. D 54 (1996) 1229]
- factorized expressions
- Q^2 evolution under control
- model-independent analysis of DFs and FFs possible

What about Integration over Transverse Momentum?

Large acceptance effects observed when comparing reconstructed *weighted* amplitudes in 4π vs. acceptance



Extracting full kinematic dependence

1. The **full kinematic dependence** of the Collins and Sivers moments on $\bar{x} \equiv (x, Q^2, z, P_{h\perp})$ is **evaluated from the real data** through a fit of the full set of SIDIS events based on a Taylor expansion on \bar{x} :

$$f(\bar{x}, P_t; c) = 1 + P_t \cdot [A_{Collins}(\bar{x}; c_i) \cdot \sin(\phi + \phi_S) + A_{Sivers}(\bar{x}; c_i) \cdot \sin(\phi - \phi_S)]$$

e.g.: $A_{Collins}(\bar{x}, c) = c_0 + c_1 \cdot x + c_2 \cdot z + c_3 \cdot Q^2 + c_4 \cdot P_{h\perp} + c_5 \cdot x^2 + \dots + c_{22} \cdot x^2 \cdot z \cdot P_{h\perp}$

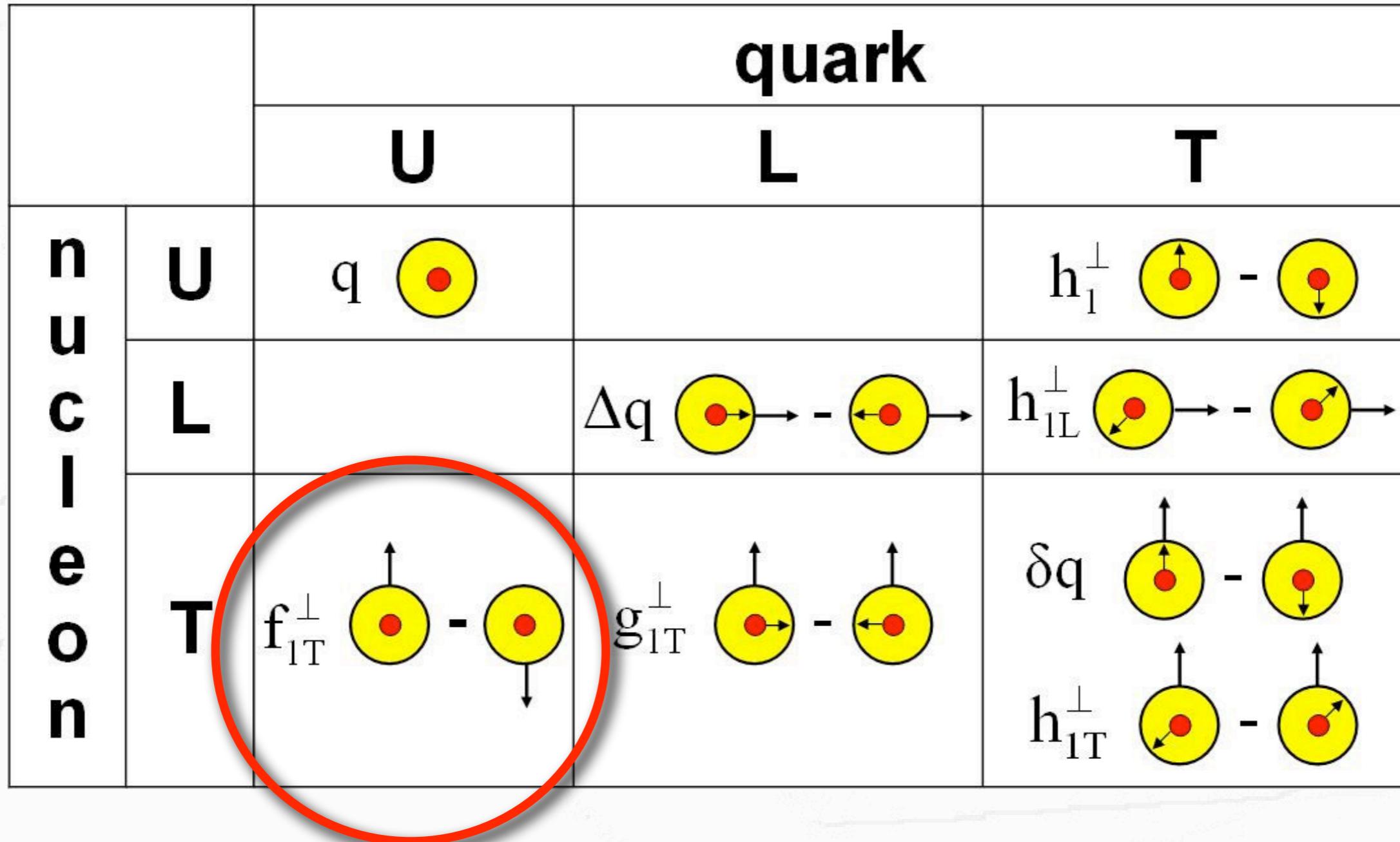
2. The extracted azimuthal moments $A_{Collins}(\bar{x}; c_i)$ and $A_{Sivers}(\bar{x}; c_i)$ are folded with the spin-independent cross section (known!) in 4π ($\sigma_{UU}^{4\pi}$) and within the HERMES acceptance ($\sigma_{UU}^{acc.}$):

$$\left\langle \frac{P_{h\perp}}{zM} \sin(\phi \pm \phi_S) \right\rangle_{UT}^{acc, 4\pi}(x) = \frac{\int P_{h\perp}/(zM) \sigma_{UU}^{acc, 4\pi}(\bar{x}) A_{Collins, Sivers}(\bar{x}; c_i)}{\int \sigma_{UU}^{acc, 4\pi}(\bar{x})}$$

L. Pappalardo's talk at Transversity'08

From SSA amplitudes to TMDs

Leading-Twist TMDs



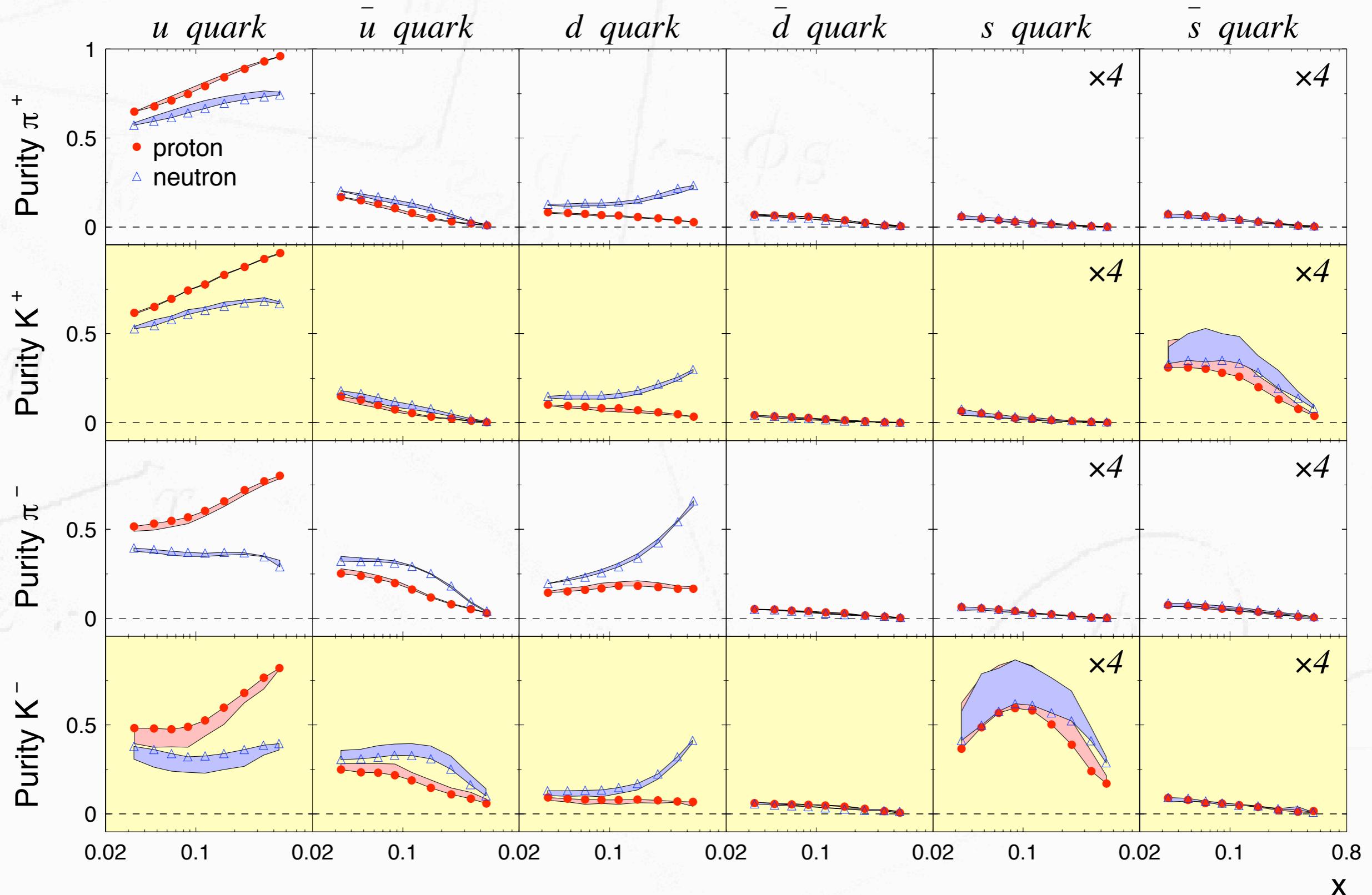
Sivers function

Using *Purities* to extract PDFs

$$\begin{aligned}
 \tilde{A}_{UT}^{\sin(\phi - \phi_S), h}(x) &= C \cdot \frac{\sum_q e_q^2 f_{1T}^{\perp(1), q}(x) \int dz D_1^{q, h}(z) \mathcal{A}(x, z)}{\sum_{q'} e_{q'}^2 f_1^{q'}(x) \int dz D_1^{q', h}(z) \mathcal{A}(x, z)} \\
 &= C \cdot \sum_q \frac{e_q^2 f_1^q(x) \mathcal{D}_1^{q, h}(x)}{\sum_{q'} e_{q'}^2 f_1^{q'}(x) \mathcal{D}_1^{q', h}(x)} \cdot \frac{f_{1T}^{\perp(1), q}}{f_1^q}(x) \\
 &= C \cdot \sum_q \mathcal{P}_q^h(x) \cdot \frac{f_{1T}^{\perp(1), q}}{f_1^q}(x)
 \end{aligned}$$

- when using weighted asymmetries, model-independent extraction of PDFs (e.g., Sivers) possible
- used successfully in HERMES helicity-DF analysis
- purities are completely unpolarized objects → present MC can be used
- Sivers case: easy as the only FF that appears is D_1

Purities at HERMES



Using *Purities* to extract PDFs

$$\begin{aligned}\tilde{A}_{UT}^{\sin(\phi+\phi_S),h}(x) &= C \cdot \frac{\sum_q e_q^2 h_1^q(x) \int dz H_1^{\perp(1),q,h}(z) \mathcal{A}(x,z)}{\sum_{q'} e_{q'}^2 f_1^{q'}(x) \int dz D_1^{q',h}(z) \mathcal{A}(x,z)} \\ &= C \cdot \sum_q \frac{e_q^2 f_1^q(x) \mathcal{H}_1^{\perp(1),q,h}(x)}{\sum_{q'} e_{q'}^2 f_1^{q'}(x) \mathcal{D}_1^{q',h}(x)} \cdot \frac{h_1^q}{f_1^q}(x) \\ &= C \cdot \sum_q \mathcal{P}_q^h(x) \cdot \frac{h_1^q}{f_1^q}(x)\end{aligned}$$

- when using weighted asymmetries, model-independent extraction of PDFs (e.g., Sivers) possible
- used successfully in HERMES helicity-DF analysis
- purities are completely unpolarized objects → present MC can be used
- Collins case: need to consider Collins FF

Valence-Quark Sivers DF

- look at difference in charged-pion yields: $\Delta N = N^{\pi^+} - N^{\pi^-}$:

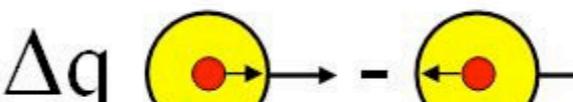
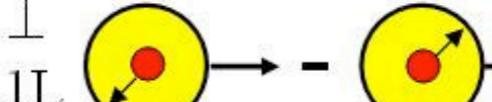
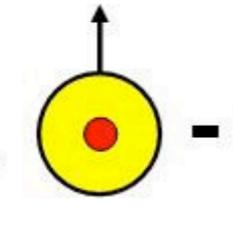
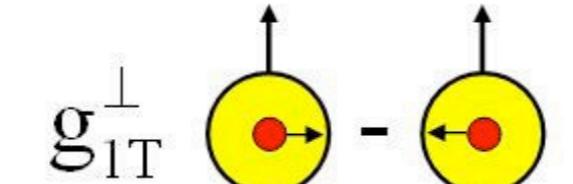
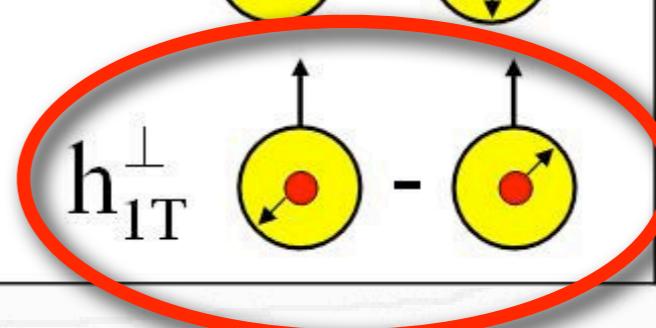
$$A_{UT}^{\pi^+ - \pi^-}(\phi, \phi_S) \equiv \frac{1}{S_T} \frac{\Delta N^\uparrow(\phi, \phi_S) - \Delta N^\downarrow(\phi, \phi_S)}{\Delta N^\uparrow(\phi, \phi_S) + \Delta N^\downarrow(\phi, \phi_S)}$$

- simple interpretation in terms of valence distributions:

$$\langle \sin(\phi - \phi_S) \rangle_{UT}^{\pi^+ - \pi^-}(\phi, \phi_S) = -\frac{4f_{1T}^{\perp, u_v} - f_{1T}^{\perp, d_v}}{4f_1^{u_v} - f_1^{d_v}}$$

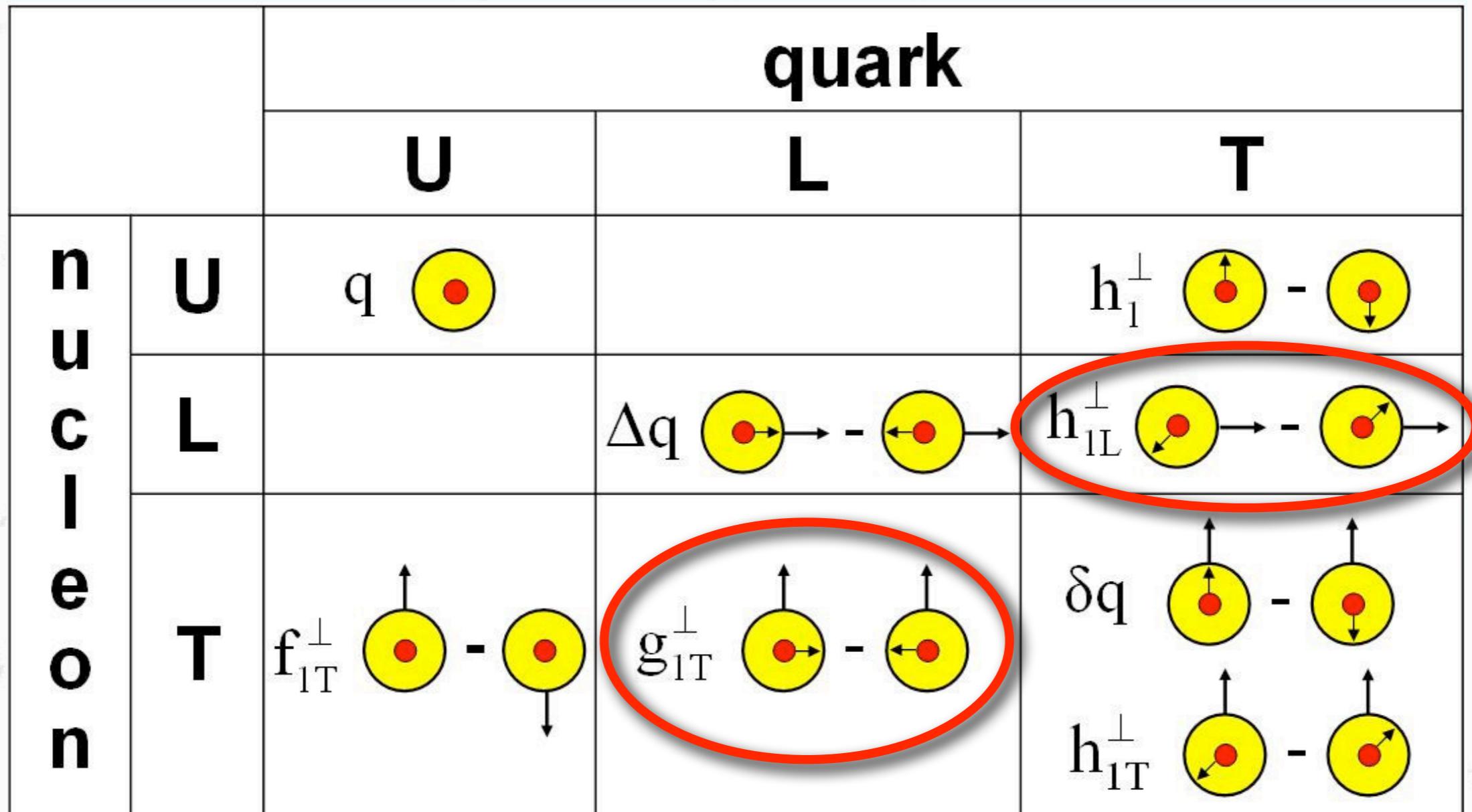
at least for weighted asymmetries, but also for unweighted case convolutions simplify

Leading-Twist TMDs

	quark			
	U	L	T	
nucleon	U	q 	Δq 	h_1^\perp 
	L		h_{1L}^\perp 	
	T	f_{1T}^\perp 	g_{1T}^\perp 	δq  

“Pretzelosity”

Leading-Twist TMDs



the “weird/boost” distributions

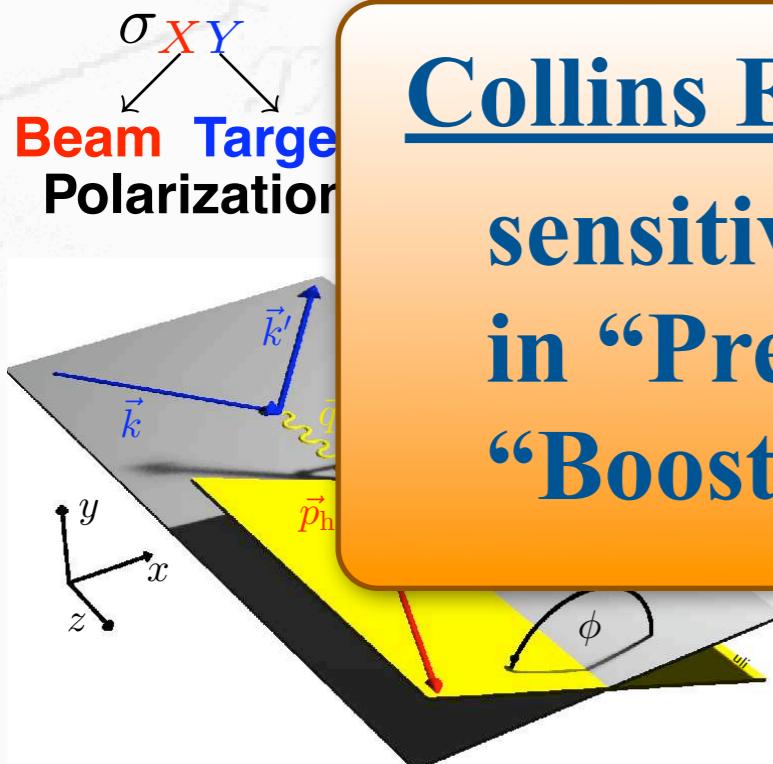
1-Hadron Production ($e p \rightarrow e h X$)

$$d\sigma = d\sigma_{UU}^0 + \cos 2\phi d\sigma_{UU}^1 + \frac{1}{Q} \cos \phi d\sigma_{UU}^2 + \lambda_e \frac{1}{Q} \sin \phi d\sigma_{LU}^3$$

$$+ S_L \left\{ \sin 2\phi d\sigma_{UL}^4 + \frac{1}{Q} \sin \phi d\sigma_{UL}^5 + \lambda_e \left[d\sigma_{LL}^6 + \frac{1}{Q} \cos \phi d\sigma_{LL}^7 \right] \right\}$$

$$+ S_T \left\{ \sin(\phi - \phi_S) d\sigma_{UT}^8 + \sin(\phi + \phi_S) d\sigma_{UT}^9 + \sin(3\phi - \phi_S) d\sigma_{UT}^{10} \right.$$

$$\left. + \frac{1}{Q} (\sin(2\phi - \phi_S) d\sigma_{UT}^{11} + \sin \phi_S d\sigma_{UT}^{12}) \right)$$



Collins Effect:

sensitive to quark transverse spin
in “Pretzelosity” distribution or
“Boostisity” h_{1L}^\perp

Bacchetta et al., JHEP 0702 (2007) 093

“Trento Conventions”, Phys. Rev. D 70 (2004) 117504

1-Hadron Production ($e p \rightarrow e h X$)

$$d\sigma = d\sigma_{UU}^0 + \cos 2\phi d\sigma_{UU}^1 + \frac{1}{Q} \cos \phi d\sigma_{UU}^2 + \lambda_e \frac{1}{Q} \sin \phi d\sigma_{LU}^3$$

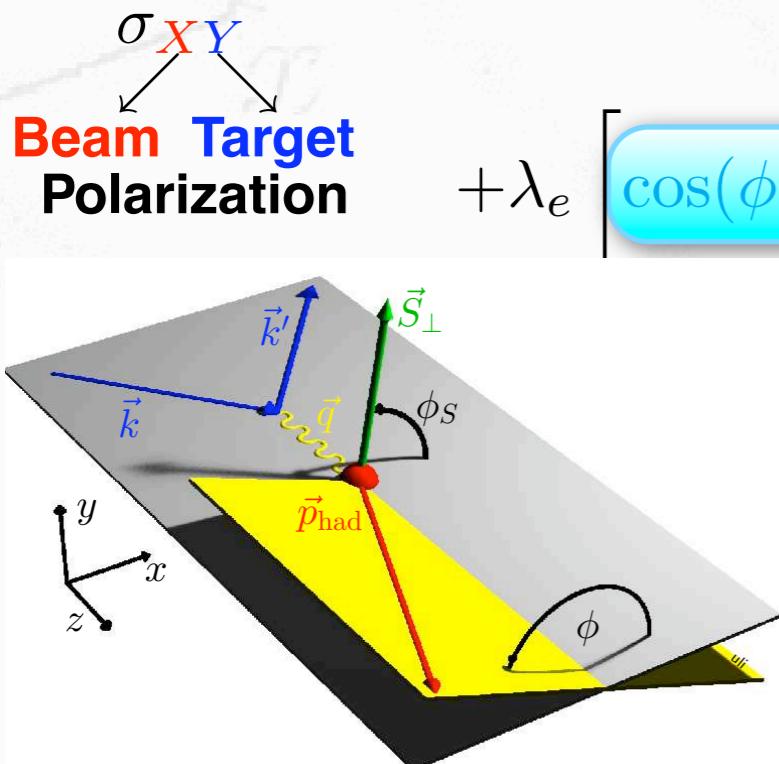
$$+ S_L \left\{ \sin 2\phi d\sigma_{UL}^4 + \dots \right\}$$

$$+ S_T \left\{ \sin(\phi - \phi_S) d\sigma_{UT}^8 - \dots \right\}$$

DSA involving spin-independent FF and “Boostisity” g_{1T}^\perp

$$+ \frac{1}{Q} (\sin(2\phi - \phi_S) d\sigma_{UT}^{\tilde{1}} + \sin \phi_S d\sigma_{UT}^{\tilde{2}})$$

$$+ \lambda_e \left[\cos(\phi - \phi_S) d\sigma_{LT}^{13} + \frac{1}{Q} (\cos \phi_S d\sigma_{LT}^{14} + \cos(2\phi - \phi_S) d\sigma_{LT}^{15}) \right]$$



Mulders and Tangermann, Nucl. Phys. B 461 (1996) 197

Boer and Mulders, Phys. Rev. D 57 (1998) 5780

Bacchetta et al., Phys. Lett. B 595 (2004) 309

Bacchetta et al., JHEP 0702 (2007) 093

“Trento Conventions”, Phys. Rev. D 70 (2004) 117504

... more twist-3 ...

1-Hadron Production (ep → ehX)

$$d\sigma = d\sigma_{UU}^0 + \cos 2\phi d\sigma_{UU}^1 + \frac{1}{Q} \cos \phi d\sigma_{UU}^2 + \lambda_e \frac{1}{Q} \sin \phi d\sigma_{LU}^3 \\ + S_L \left\{ \sin 2\phi d\sigma_{UL}^4 + \frac{1}{Q} \sin \phi d\sigma_{UL}^5 + \lambda_e \left[d\sigma_{LL}^6 + \frac{1}{Q} \cos \phi d\sigma_{LL}^7 \right] \right\}$$

$\int_S \sin(\phi - \phi_S) d\sigma^8 + \sin(\phi + \phi_S) d\sigma^9_{UT} + \sin(3\phi - \phi_S) d\sigma^{10}_{UT}$
**sensitivity to or needed
for transversity**
 $(d\sigma_{UT}^{11} + \sin \phi_S d\sigma_{UT}^{12})$
 $+ \lambda_e \left[\cos(\phi - \phi_S) d\sigma_{LT}^{13} + \frac{1}{Q} (\cos \phi_S d\sigma_{LT}^{14} + \cos(2\phi - \phi_S) d\sigma_{LT}^{15}) \right]$

Beam Target Polarization

Mulders and Tangermann, Nucl. Phys. B 461 (1996) 197
Boer and Mulders, Phys. Rev. D 57 (1998) 5780
Bacchetta et al., Phys. Lett. B 595 (2004) 309
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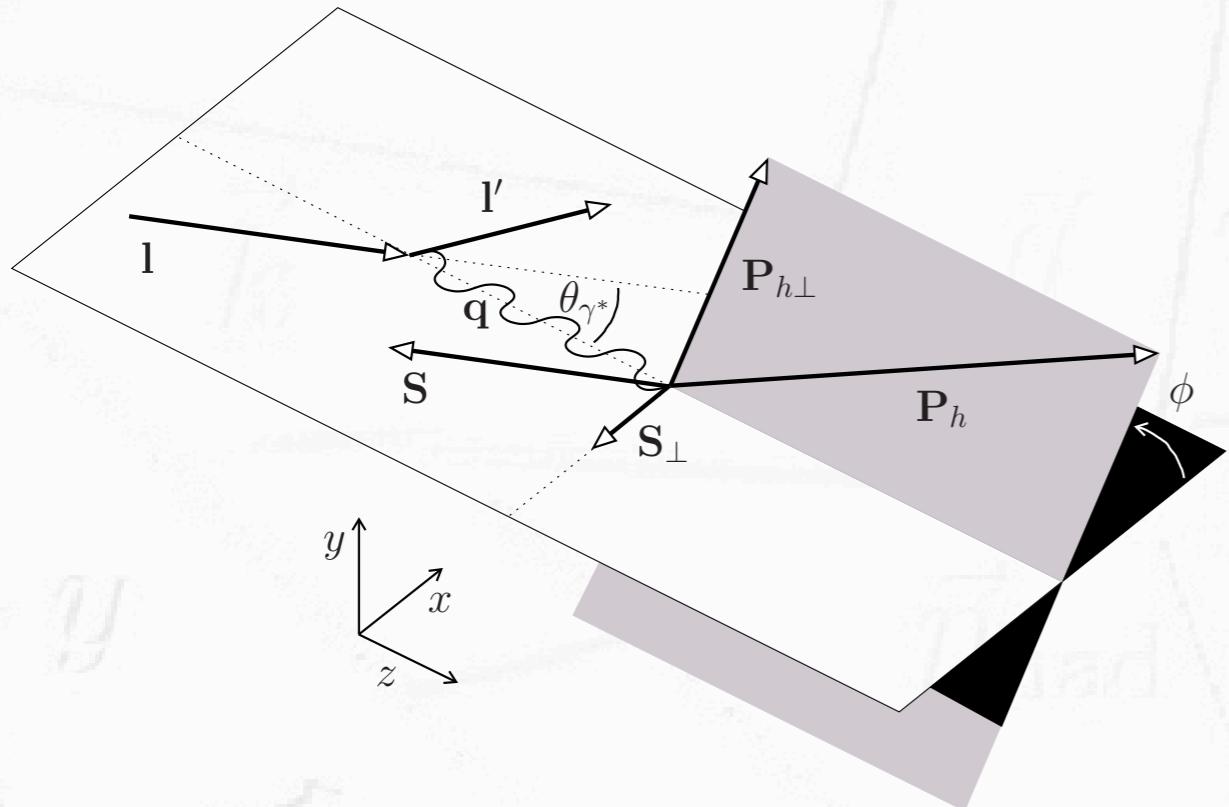
$\sin \varphi_s$ - term in AUT

$$-\mathcal{I} \left[\frac{\mathbf{k}_T \cdot \mathbf{p}_T}{2MM_h} \left(x h_T H_1^\perp - x h_T^\perp H_1^\perp - \frac{M_h}{M} g_{1T} \frac{\tilde{G}^\perp}{z} + \frac{M_h}{M} f_{1T}^\perp \frac{\tilde{D}^\perp}{z} \right) \right]$$

$-h_1 + \tilde{h}_T - \frac{p_T^2}{2M^2x} h_{1T}^\perp$
↑
 $h_1 - \tilde{h}_T^\perp - \frac{p_T^2}{2M^2x} h_{1T}^\perp$

\rightarrow
 $-2h_1$
 \rightarrow
 $-(\tilde{h}_T + \tilde{h}_T^\perp)$

AUL & Mixing of Azimuthal Moments



Experiment: Target Polarization w.r.t. Beam Direction (l)!

Theory: Polarization along virtual photon direction (q)

⇒ mixing of “experimental” and “theory” asymmetries via:

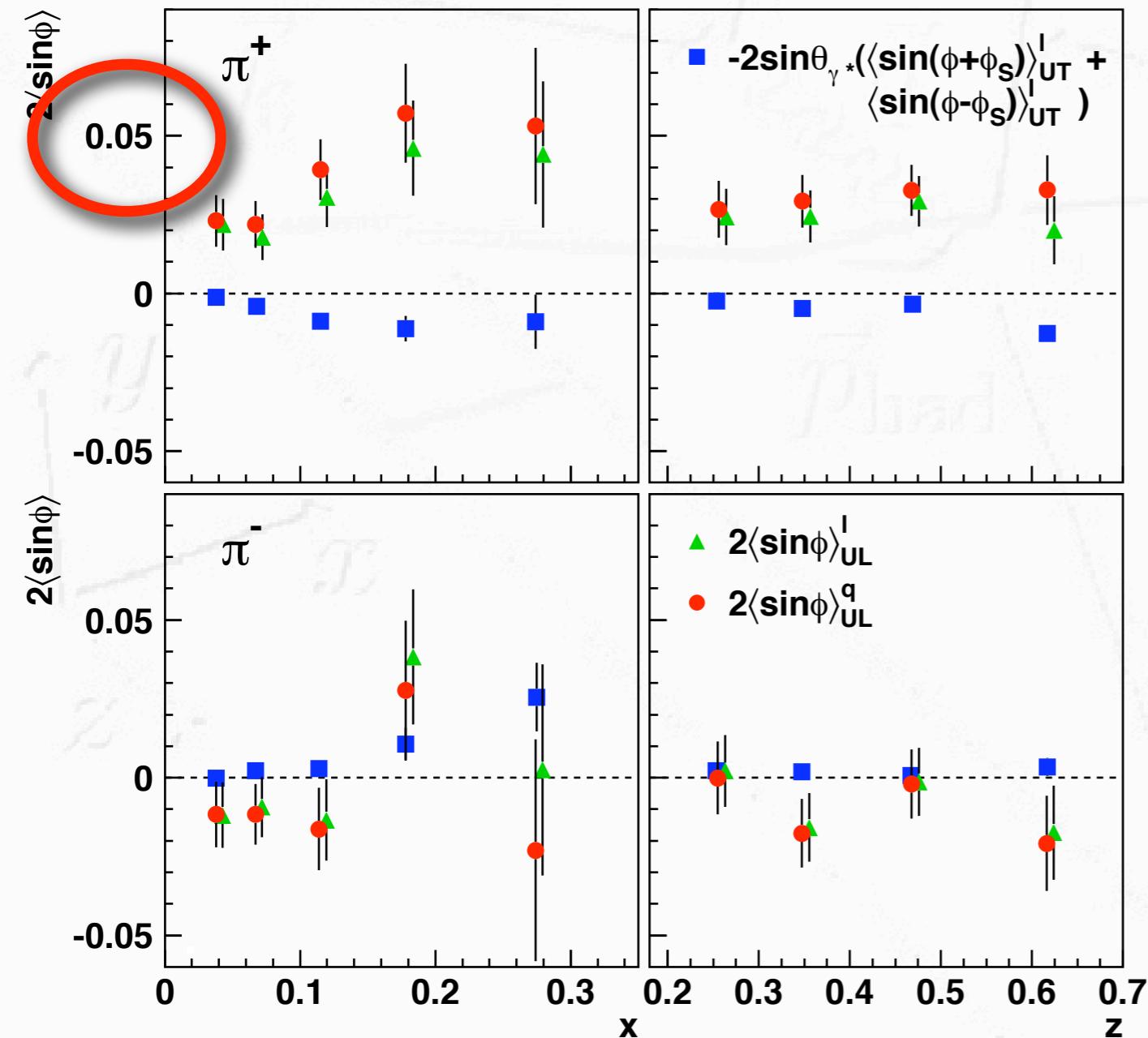
[Diehl and Sapeta, Eur. Phys. J. C41 (2005)]

$$\begin{pmatrix} \langle \sin \phi \rangle_{UL}^l \\ \langle \sin(\phi - \phi_S) \rangle_{UT}^l \\ \langle \sin(\phi + \phi_S) \rangle_{UT}^l \end{pmatrix} = \begin{pmatrix} \cos \theta_{\gamma^*} & -\sin \theta_{\gamma^*} & -\sin \theta_{\gamma^*} \\ \frac{1}{2} \sin \theta_{\gamma^*} & \cos \theta_{\gamma^*} & 0 \\ \frac{1}{2} \sin \theta_{\gamma^*} & 0 & \cos \theta_{\gamma^*} \end{pmatrix} \begin{pmatrix} \langle \sin \phi \rangle_{UL}^q \\ \langle \sin(\phi - \phi_S) \rangle_{UT}^q \\ \langle \sin(\phi + \phi_S) \rangle_{UT}^q \end{pmatrix}$$

($\cos \theta_{\gamma^*} \simeq 1$, $\sin \theta_{\gamma^*}$ up to 15% at HERMES energies)

Twist-3 at HERMES

$$\langle \sin \phi \rangle_{UL}^q = \langle \sin \phi \rangle_{UL}^I + \sin \theta_{\gamma^*} \left(\langle \sin(\phi + \phi_S) \rangle_{UT}^I + \langle \sin(\phi - \phi_S) \rangle_{UT}^I \right)$$



- twist-3 dominates measured asymmetries on longitudinally polarized targets!
- significantly positive for π^+
- consistent with zero for π^-
- twist-3 not necessarily small

Airapetian et al., Phys. Lett. B 622 (2005) 14

Another Longitudinal SSA: ALU

longitudinally pol. beam & unpol. target \Rightarrow subleading-twist

$$\langle \sin \phi \rangle_{LU} \propto \lambda_e \frac{M}{Q} \mathcal{I} \left[x e(x) H_1^\perp(z) - \frac{M_h}{z M} h_1^\perp(x) E(z) \right]$$

\Rightarrow for long time candidate to access $e(x)$
 $(h_1^\perp(x) \text{ contribution either assumed to be zero (T-odd!) or small(??)})$

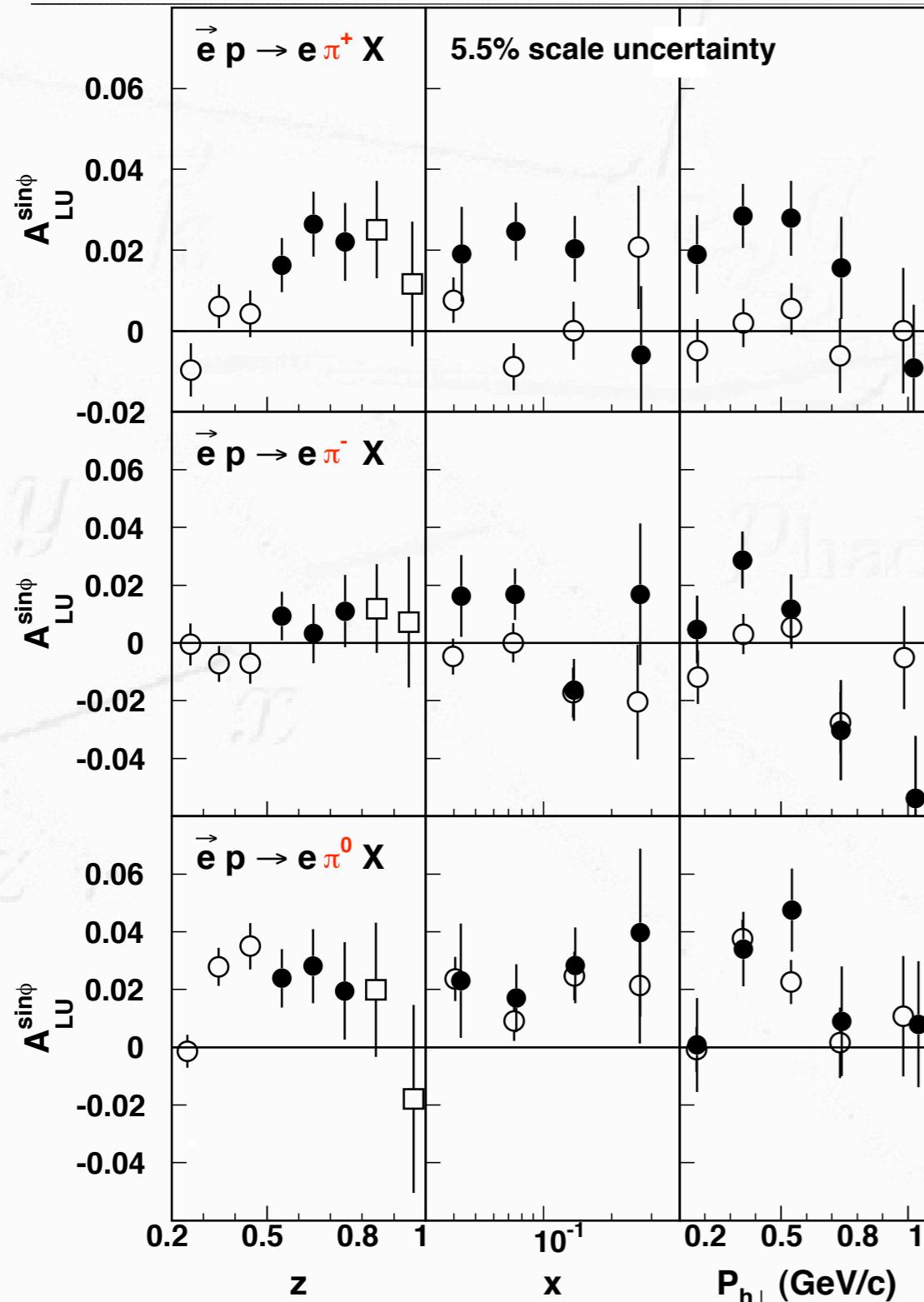
Another Longitudinal SSA: ALU

longitudinally pol. beam & unpol. target \Rightarrow subleading-twist

$$\begin{aligned} \langle \sin \phi \rangle_{LU} \propto & \lambda_e \frac{M}{Q} \mathcal{I} \left[xe(x) H_1^\perp(z) - \frac{M_h}{zM} h_1^\perp(x) E(z) \right. \\ & + \frac{M_h}{zM} f_1(x) G^\perp(z) - x g^\perp(x) D_1(z) \\ \text{quark-mass suppressed} \Rightarrow & \left. + \frac{m_q}{M} h_1^\perp(x) D_1(z) - \frac{m_q}{M} f_1(x) H_1^\perp(z) \right] \end{aligned}$$

Bacchetta et al., Phys. Lett. B 595 (2004) 309

Longit. Beam-Spin Asymmetries



- Significantly positive amplitudes for neutral and positive pions
- much more data on tape:
 - in total factor 7 for H
 - factor 3 for D
 - mostly with RICH detector, thus kaon amplitudes possible

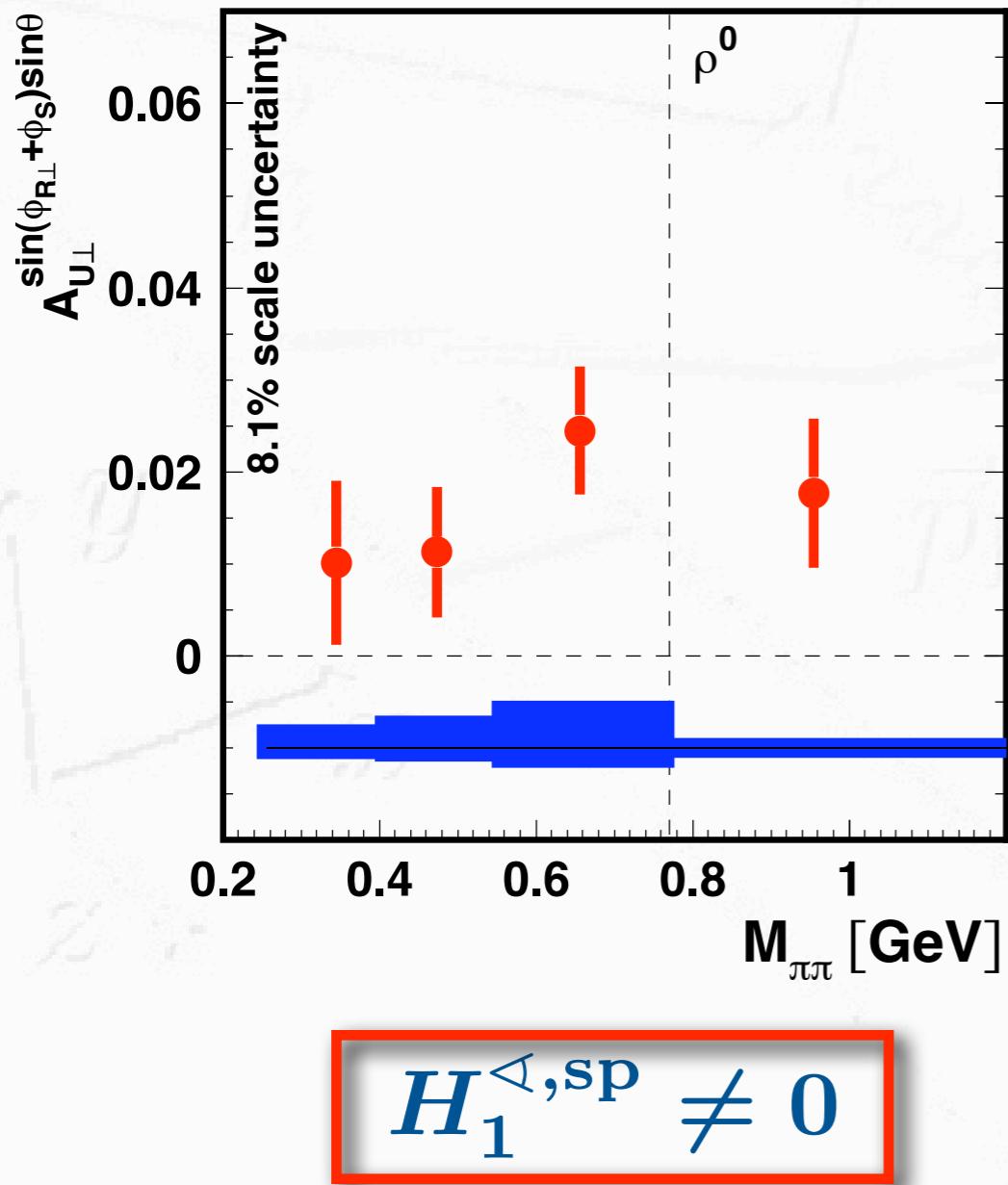
What can we learn from AUL?

$$\langle \sin \phi \rangle_{LU} \propto \lambda_e \frac{M}{Q} \mathcal{I} \left[x e(x) H_1^\perp(z) - \frac{M_h}{zM} h_1^\perp(x) E(z) \right. \\ \left. - x g^\perp(x) D_1(z) + \frac{M_h}{zM} f_1(x) G^\perp(z) \right]$$

- any help from other observables to separate contributions?
- jet* SIDIS \Rightarrow only g^\perp -term survives
- 2-hadron production nonzero! (cf. R. Fabbri's talk)

$$\sigma_{LU} \propto \sin \phi_{R\perp} \left[x e(x) H_1^\triangleleft(z, \zeta, M_h^2) + \frac{1}{z} f_1(x) \tilde{G}^\triangleleft(z, \zeta, M_h^2) \right]$$

2-Hadron Fragmentation

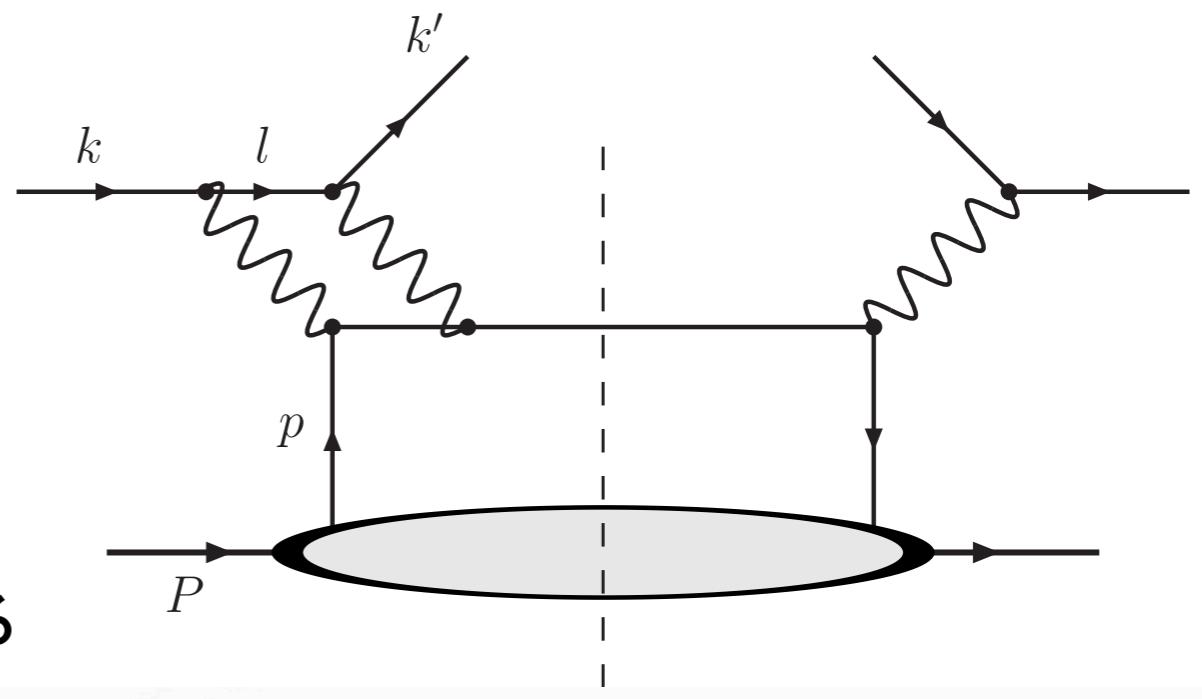


- so far: sp interference
→ look at pp interference
- $\pi K, KK$ pair production
- spin-1 fragmentation (ρ^0)
- ...
- for A_{UT} and A_{LU} (the latter also on D target)

now something completely different

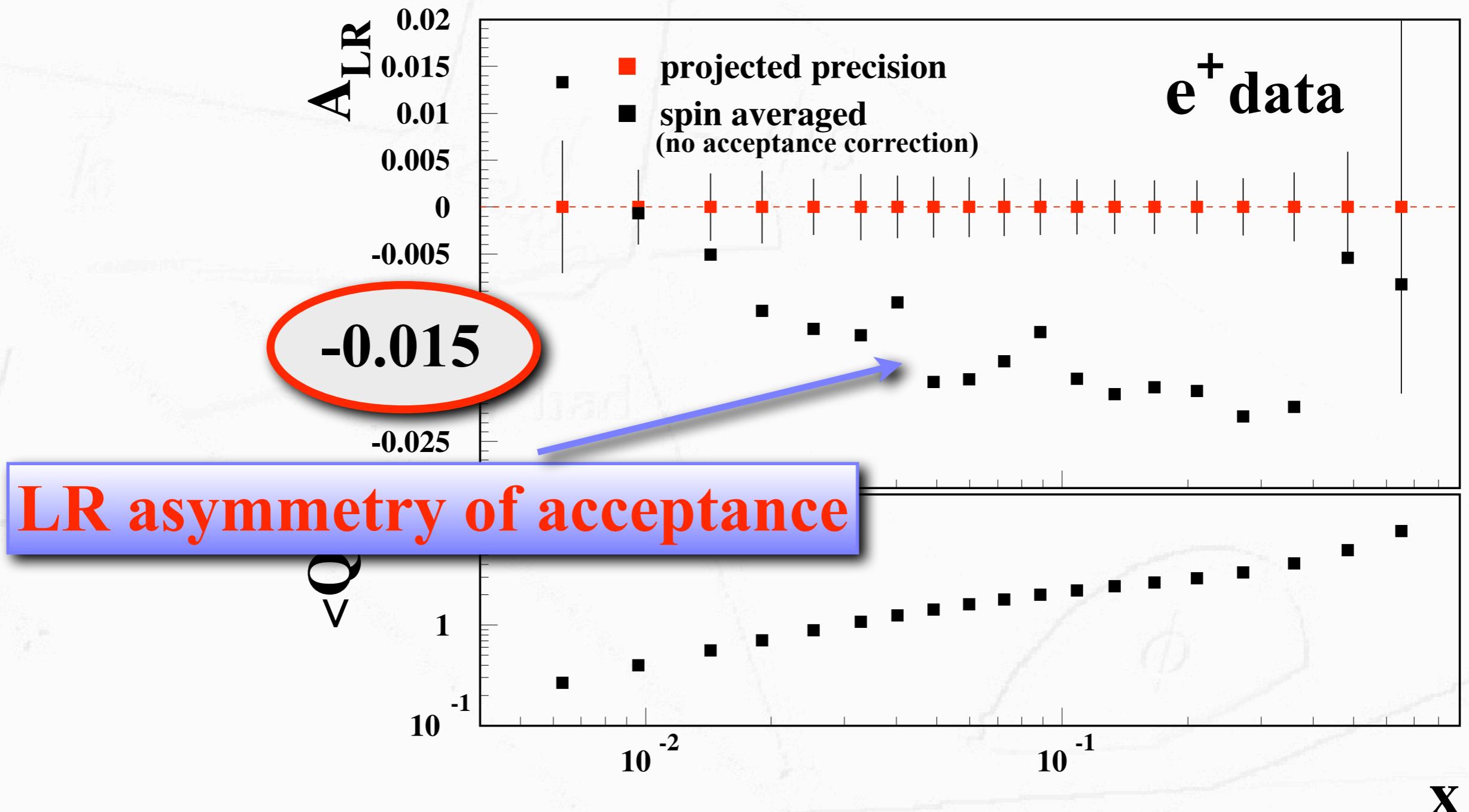
SSA in inclusive DIS

- transverse SSA require interference of amplitudes with different phases
- achievable via loop diagrams, e.g.
 - Sivers DF includes gauge link (soft gluon exchange)
- How about inclusive DIS?
- 2-photon exchange could provide such mechanism in inclusive DIS



A. Metz et al., Phys.Lett.B643:319-324,2006

2γ -exchange sensitivity @ HERMES



even more e^- data available!

The other inclusive SSA

- instead of inclusive DIS, look at inclusive hadron production (a la E704 etc., but photo-production)
- plenty of data available
- can they be related to Sivers effect in any way?
- or what is the physics of Left-Right asymmetries in photo-production?

Conclusions

- plenty of projects
 - inclusive DIS (g_2 , 2-photon exchange)
 - inclusive hadron left-right asymmetries
 - SSA and DSA in semi-inclusive DIS \Rightarrow access to various TMDs like Sivers, Pretzelosity and the “weird ones”
 - flavor decomposition of Sivers function via purity analysis and pion-yield difference asymmetries
 - 2-hadron fragmentation: sp-&pp-interference etc., in ALU and AUT
- “only the sky*) is the limit”

*) sky = manpower