

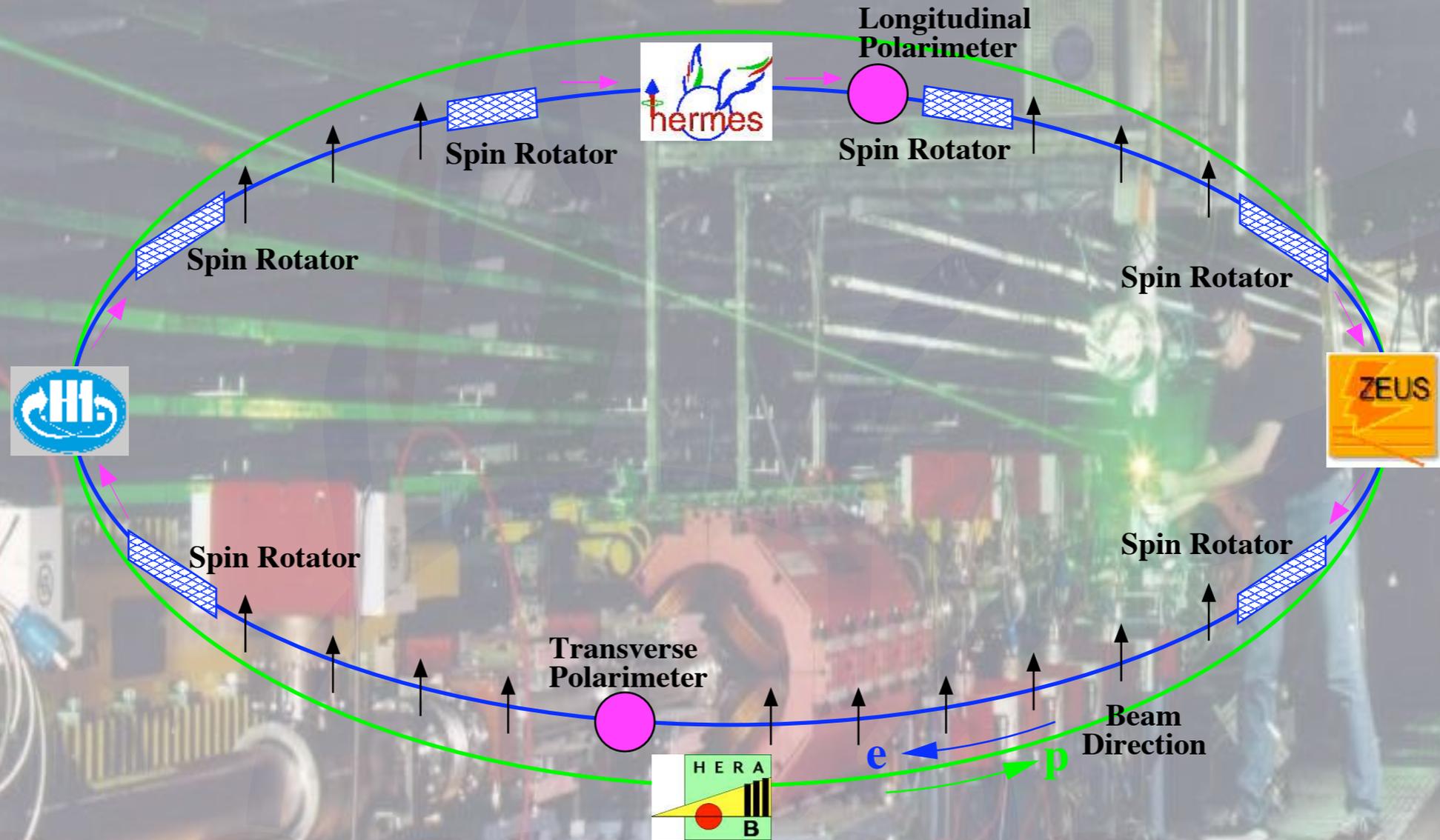
POETIC 2013

*Physics Opportunities
at an ElecTron Ion Collider*

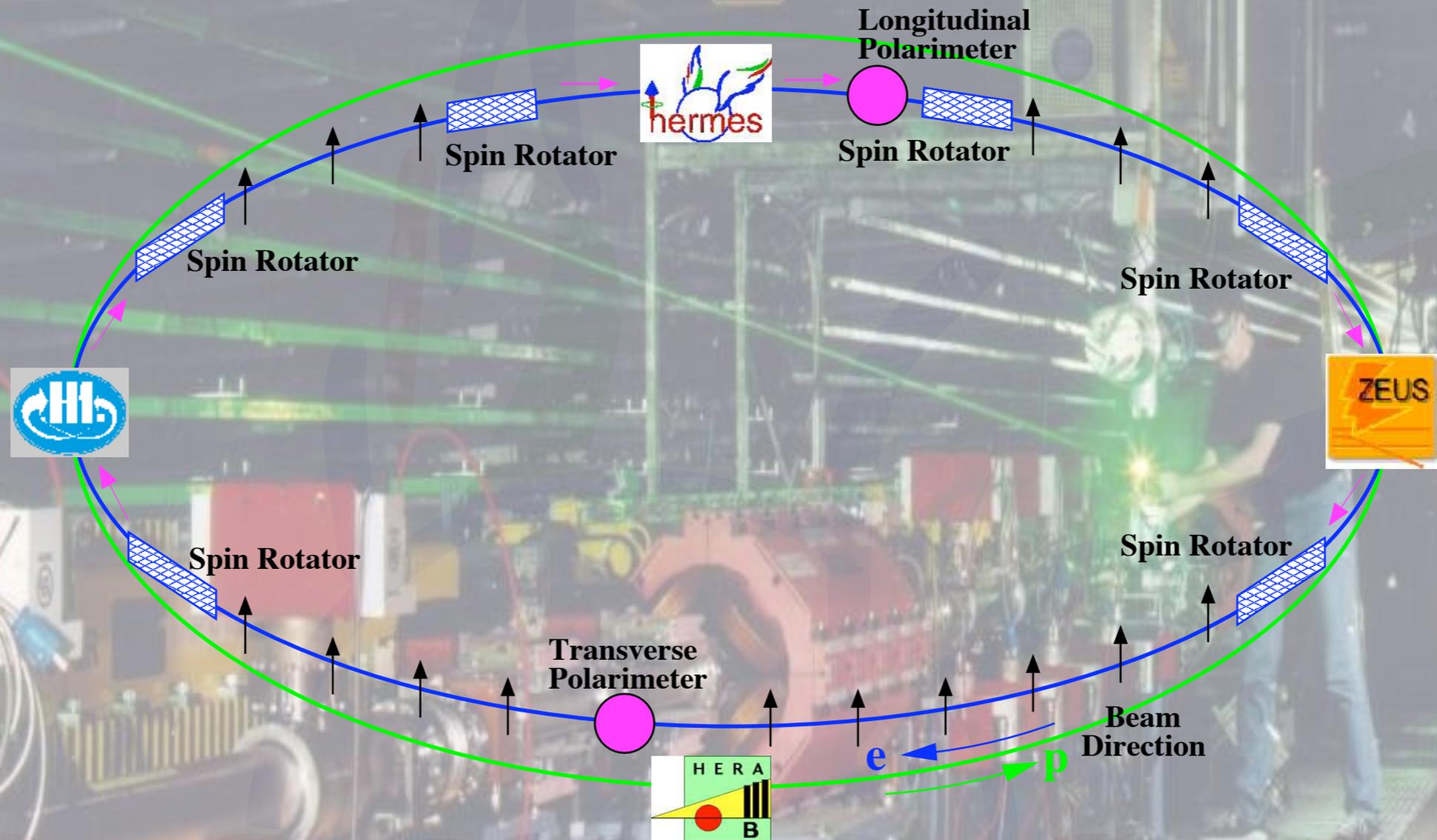
highlights from the  hermes collaboration

- the HERMES experiment
- inclusive DIS
- semi-inclusive DIS
- 3D structure via TMDs and GPDs

Beam polarization at HERA

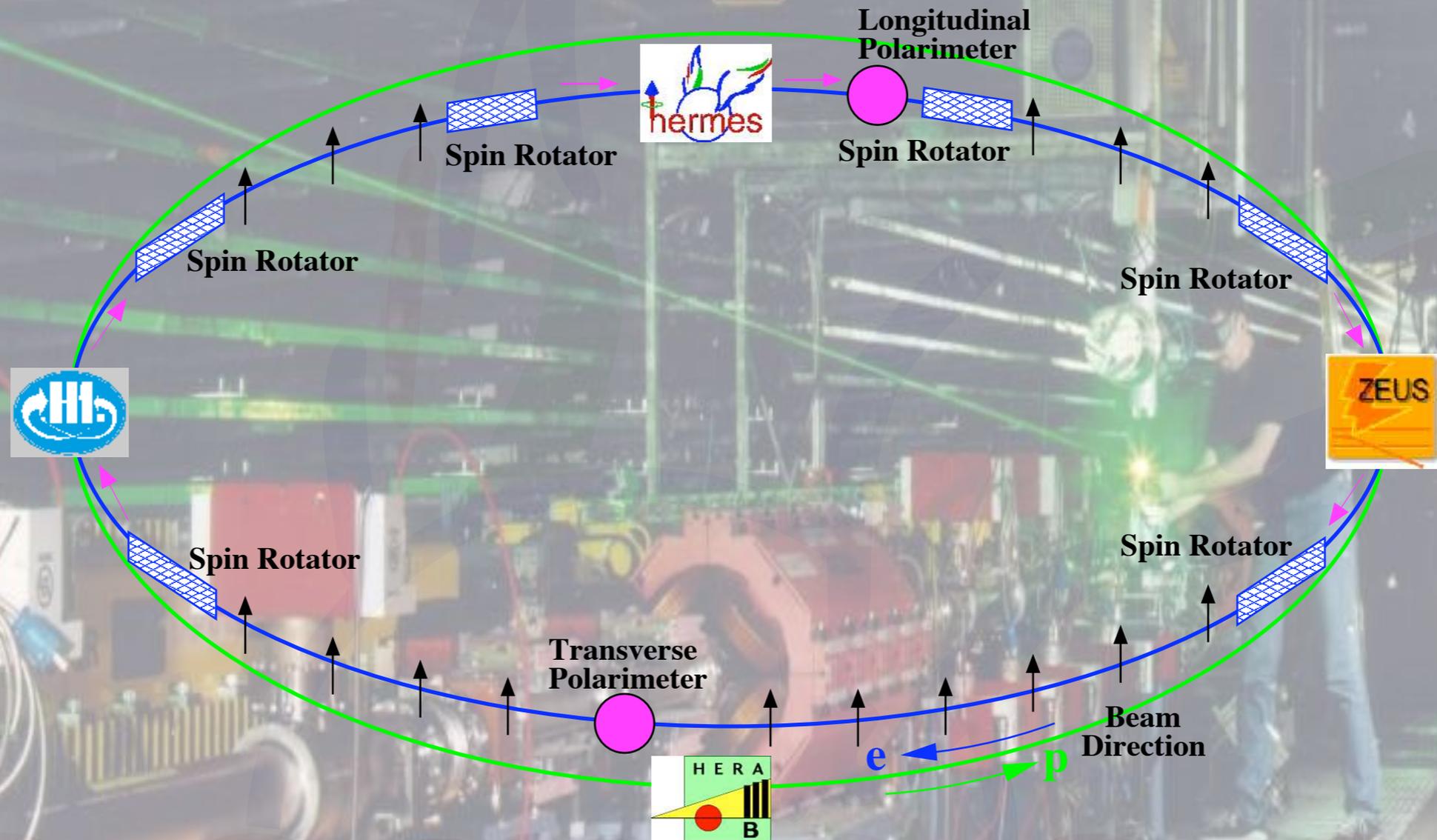


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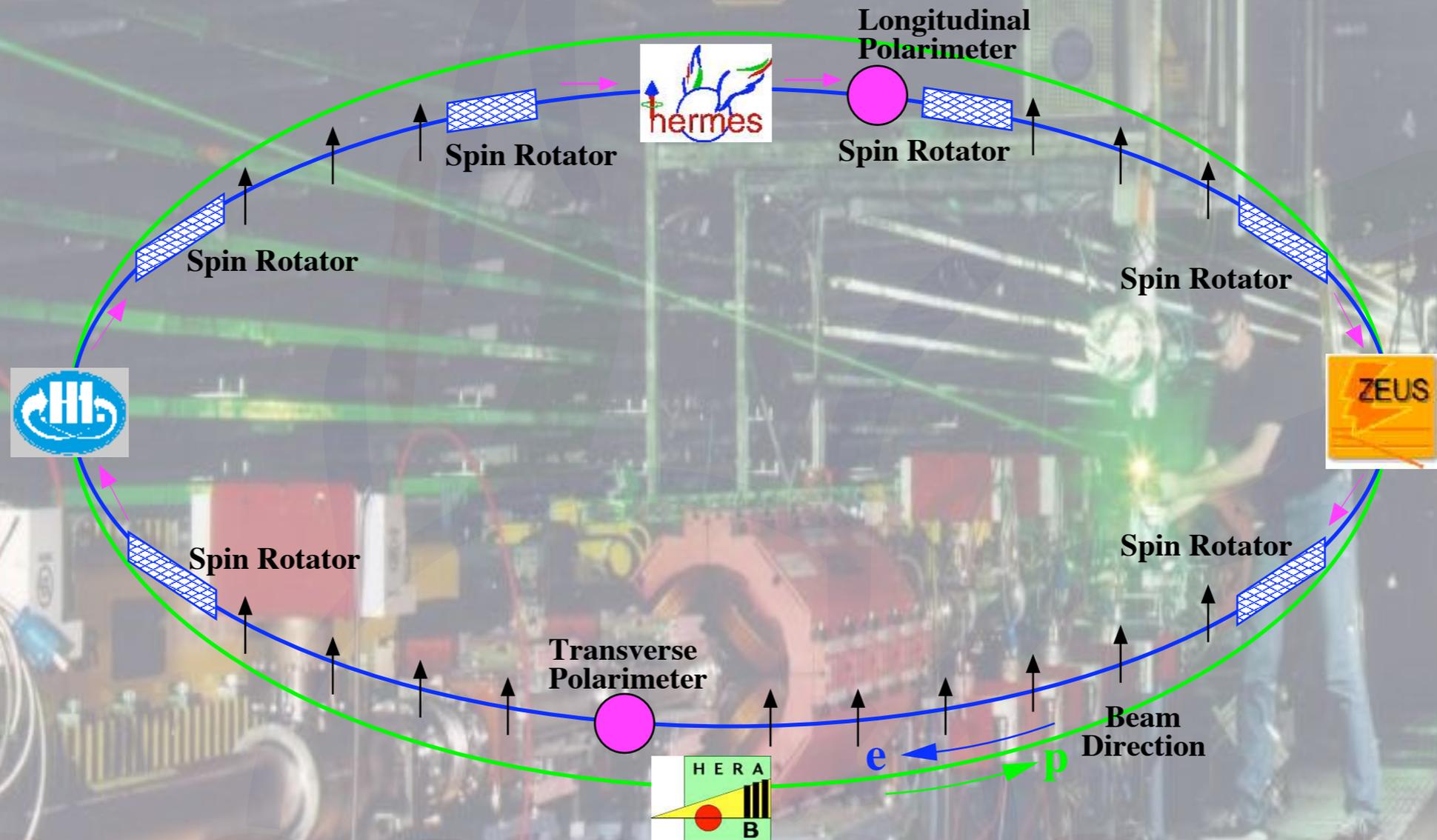
- 27.5 GeV electron/positron beam

Beam polarization at HERA



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- transversely polarized through Sokolov-Ternov effect

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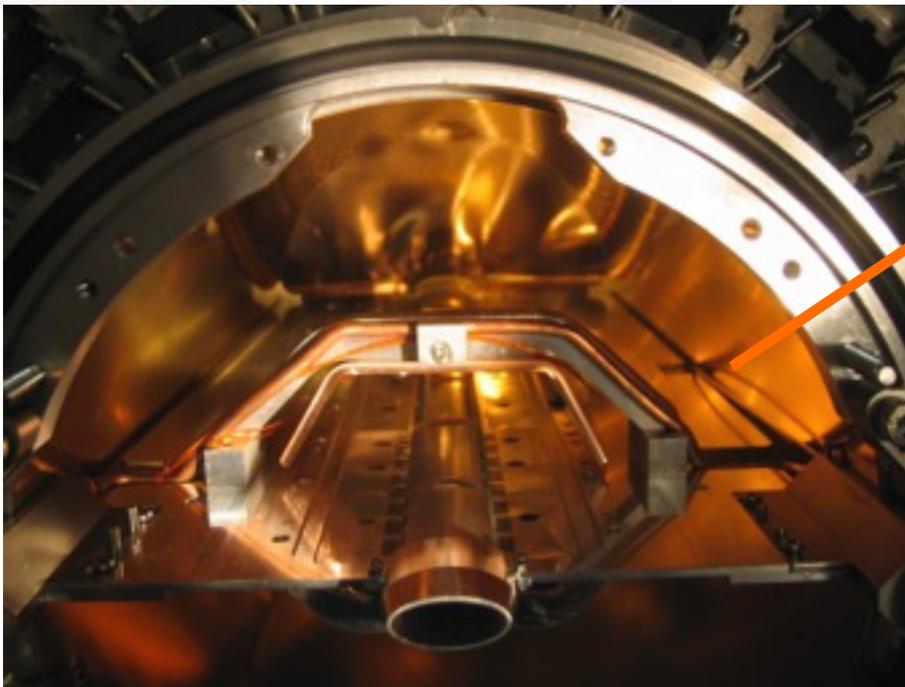


- 27.5 GeV electron/positron beam
- transversely polarized through Sokolov-Ternov effect
- average beam polarization up to 55%

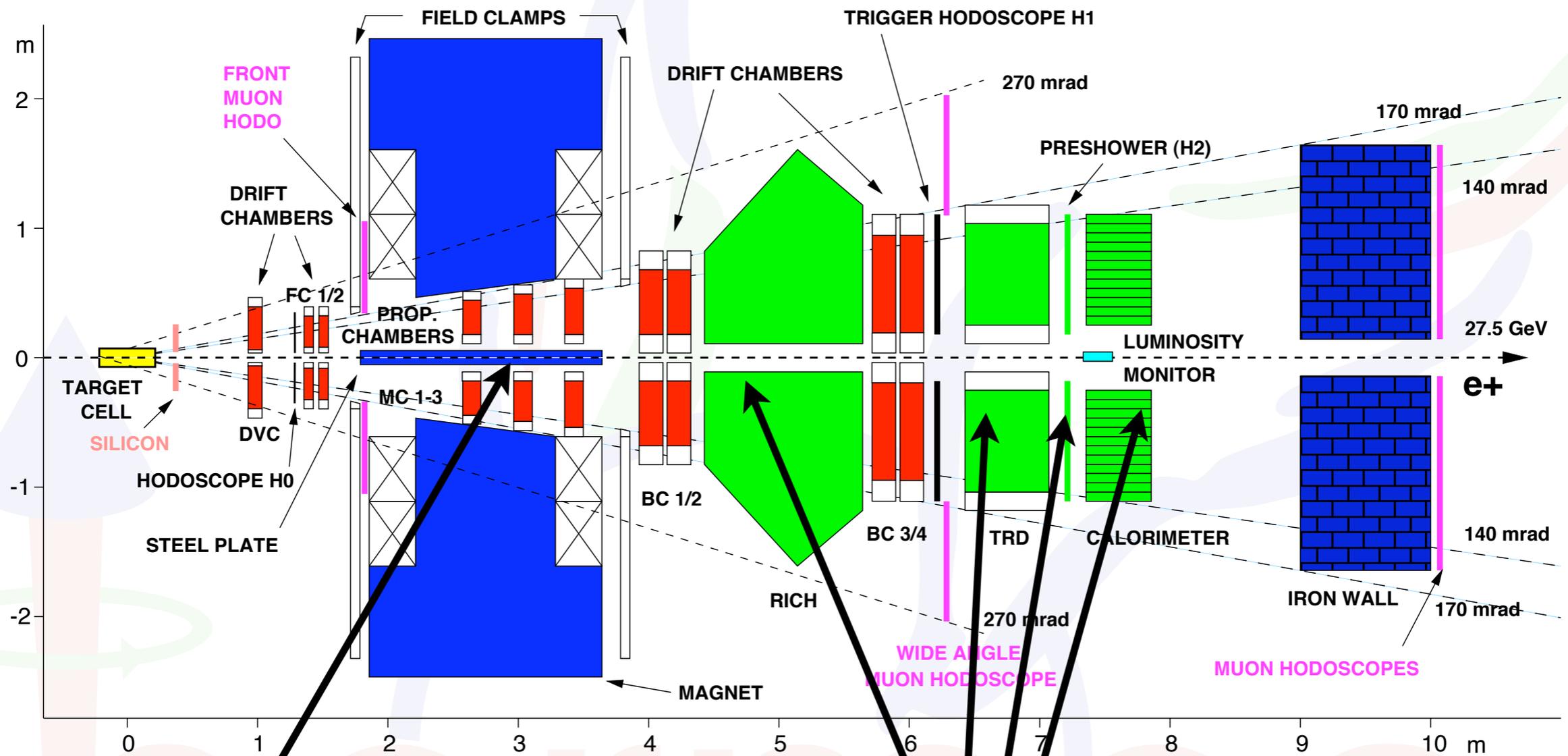
The HERMES experiment (1995-2007)

novel pure gas target:

- internal to HERA lepton ring
- unpolarized (^1H ... Xe)
- longitudinally polarized: ^1H , ^2H , ^3He
- transversely polarized: ^1H



HERMES schematically

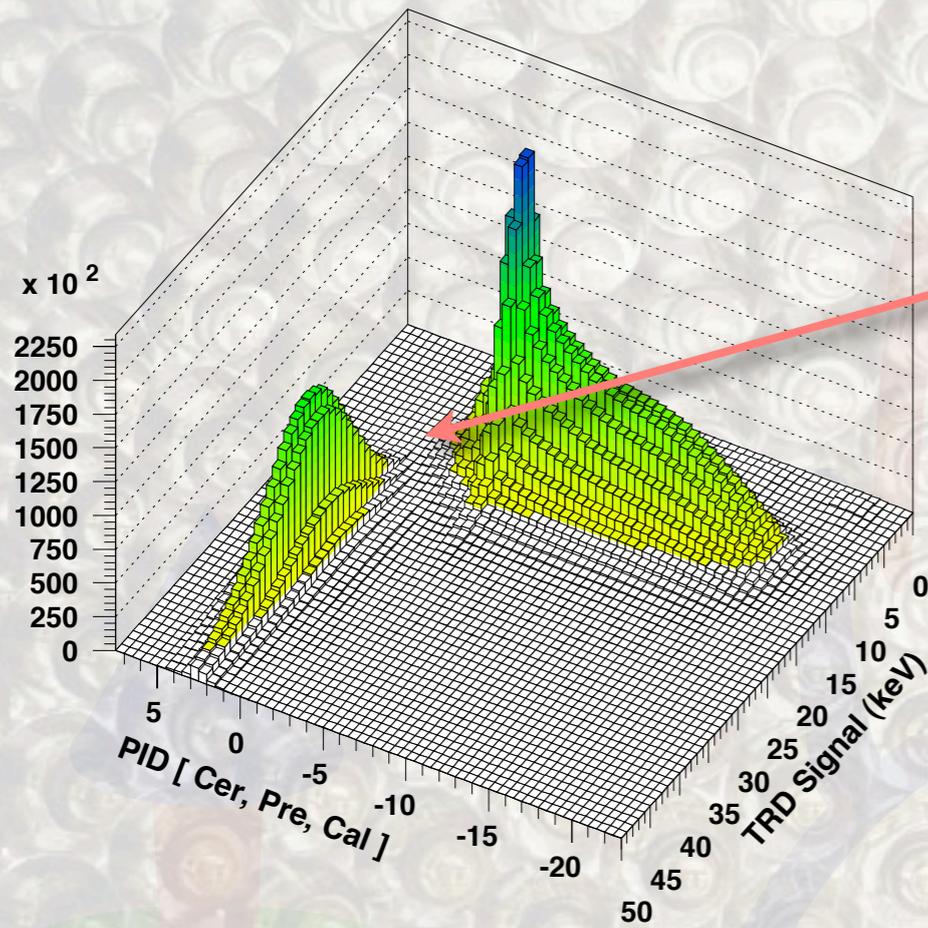


two (mirror-symmetric) halves
 -> no homogenous azimuthal coverage

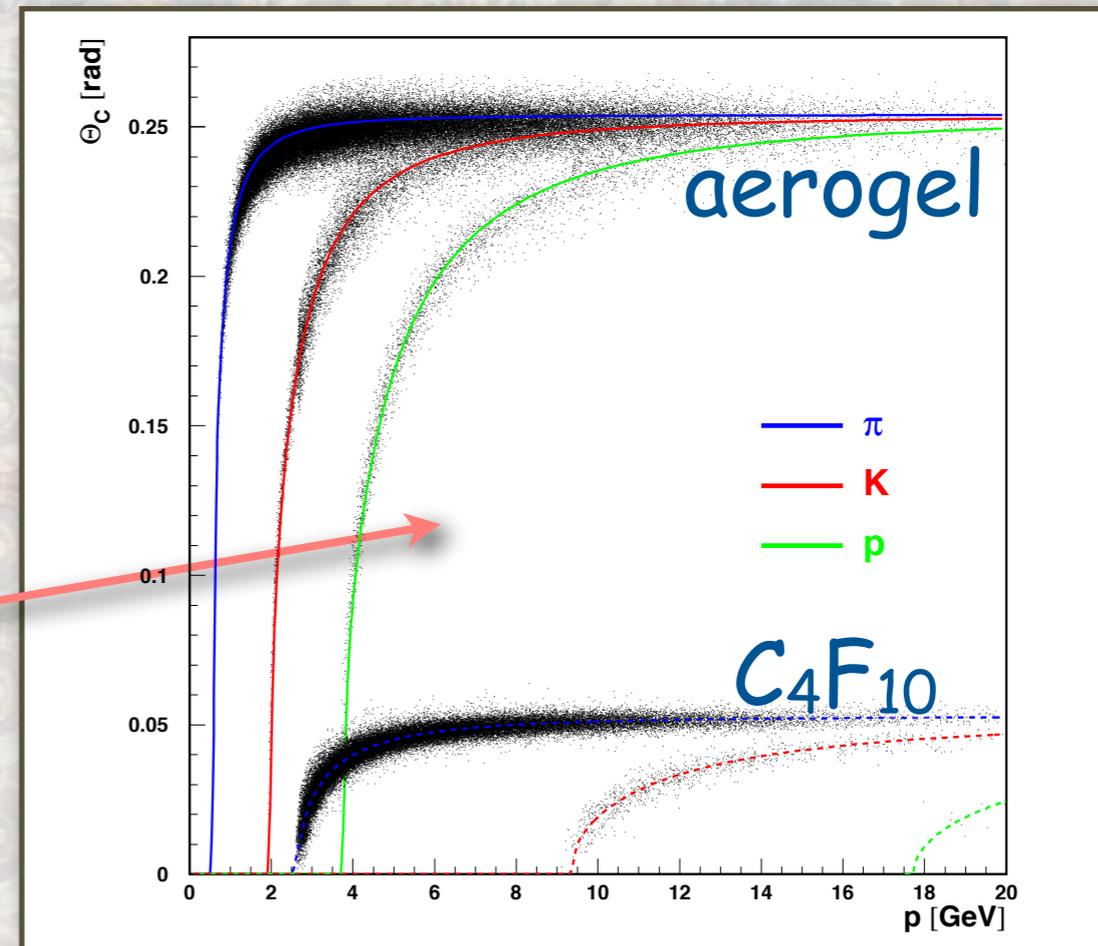
Particle ID detectors allow for
 - lepton/hadron separation
 - RICH: pion/kaon/proton discrimination $2\text{GeV} < p < 15\text{GeV}$

Particle identification

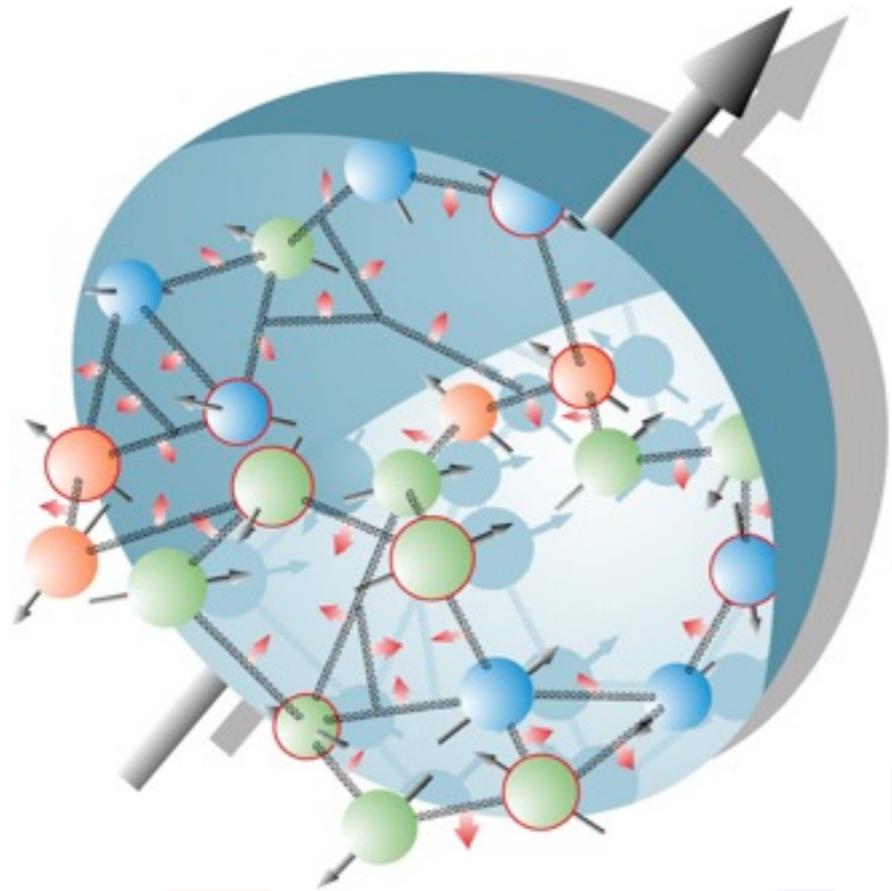
excellent lepton/hadron
separation



Dual-Radiator RICH
hadron ID for
momenta 2-15 GeV



The HERA-I (1995-2000) harvest



$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma$$

← quark spin

$$+ \Delta G$$

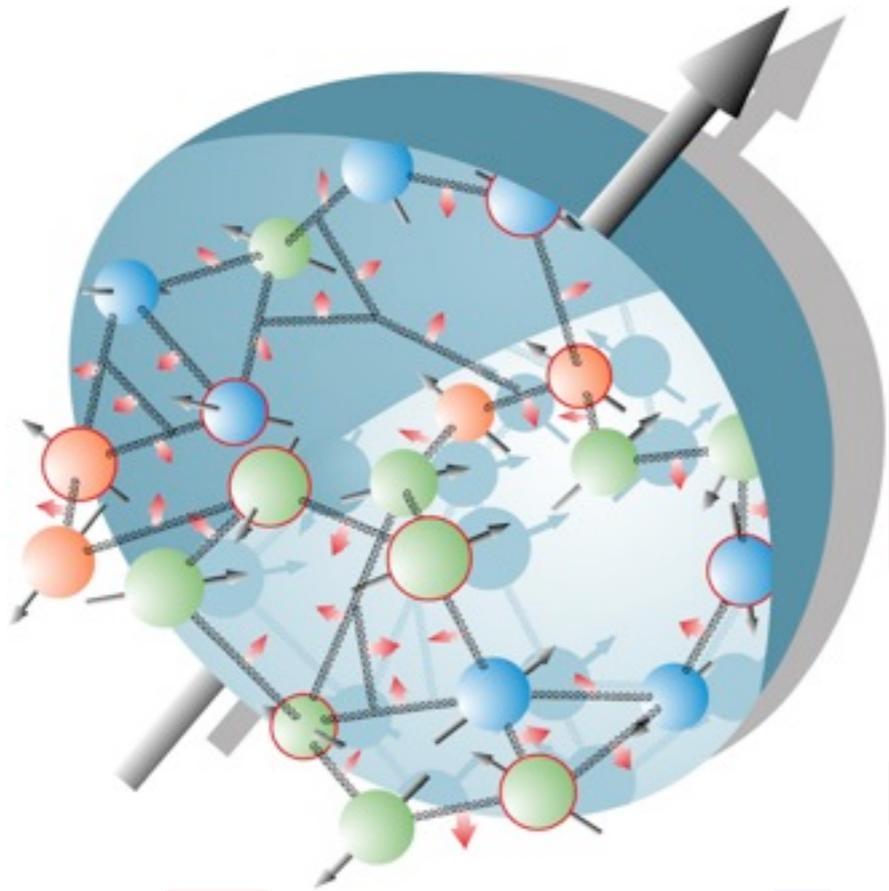
← gluon spin

$$+ L_q + L_g$$

← orbital angular momentum

hermes

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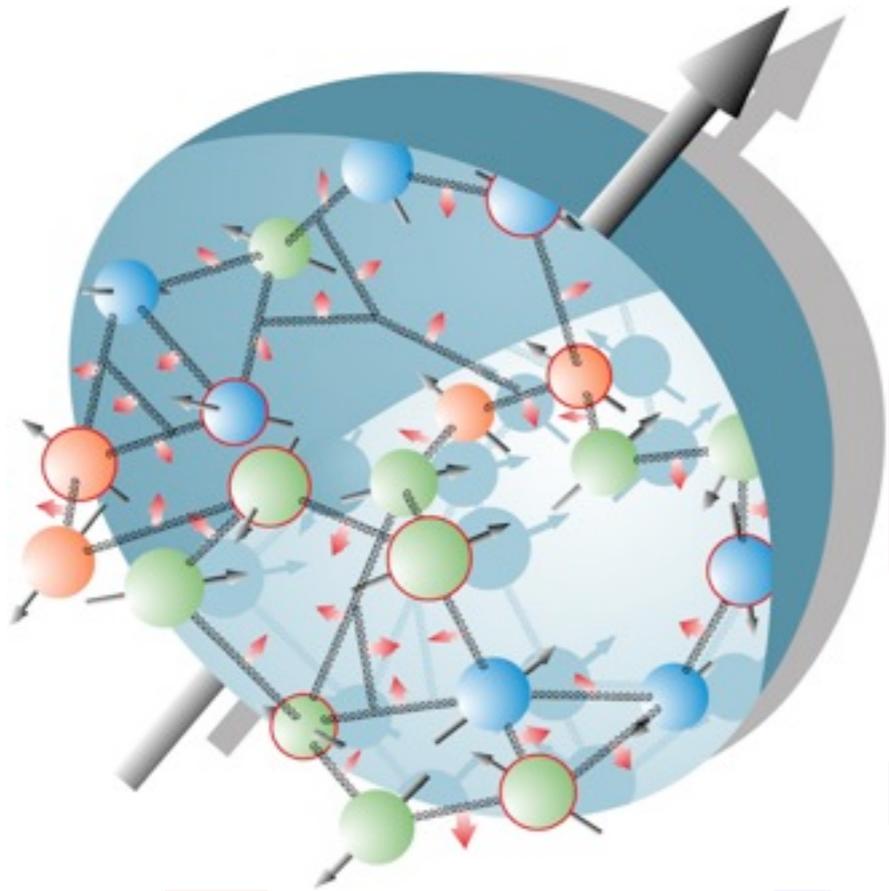
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- Inclusive **DIS** from longitudinally polarized deuterium target:

$$\Delta\Sigma = 0.330 \pm 0.025 \text{ (exp.)} \pm 0.011 \text{ (theory)} \pm 0.028 \text{ (evol.)}$$

PRD 75 (2007) 012007

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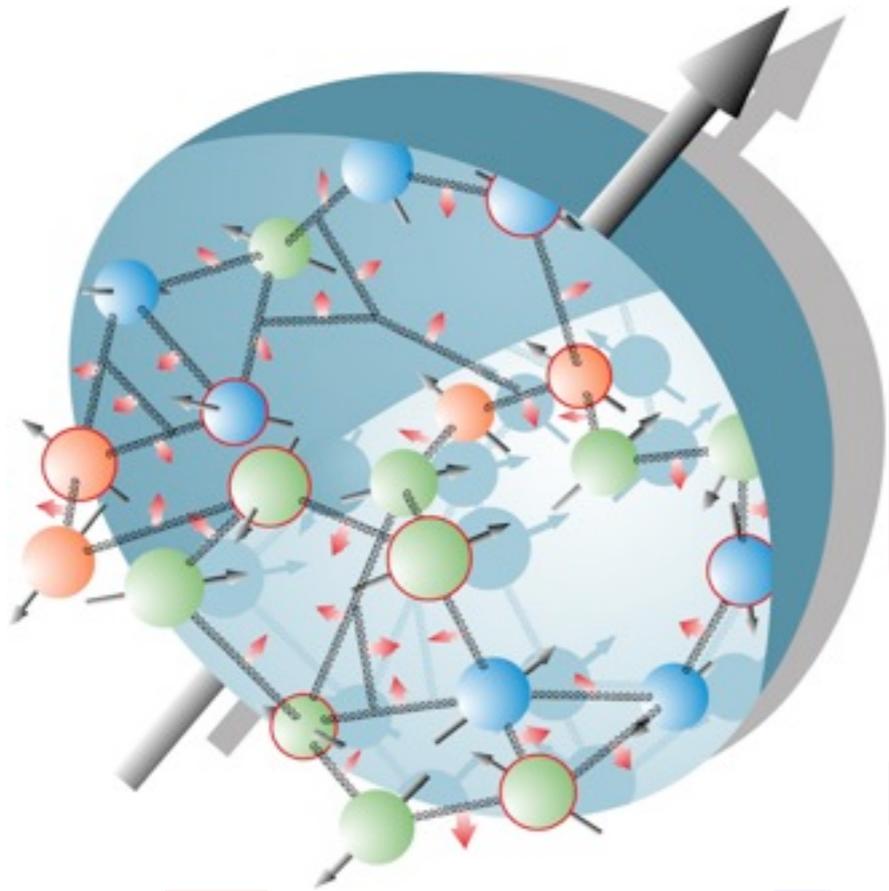
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PRD 75 (2007) 012007

- High- p_T hadrons at HERMES:

$$\Delta G/G = 0.071 \pm 0.034 \text{ (stat)} \pm 0.010 \text{ (sys-exp)} \begin{matrix} +0.127 \\ -0.105 \end{matrix} \begin{matrix} \text{JHEP 1008 (2010) 130} \\ \text{(sys-model)} \end{matrix}$$

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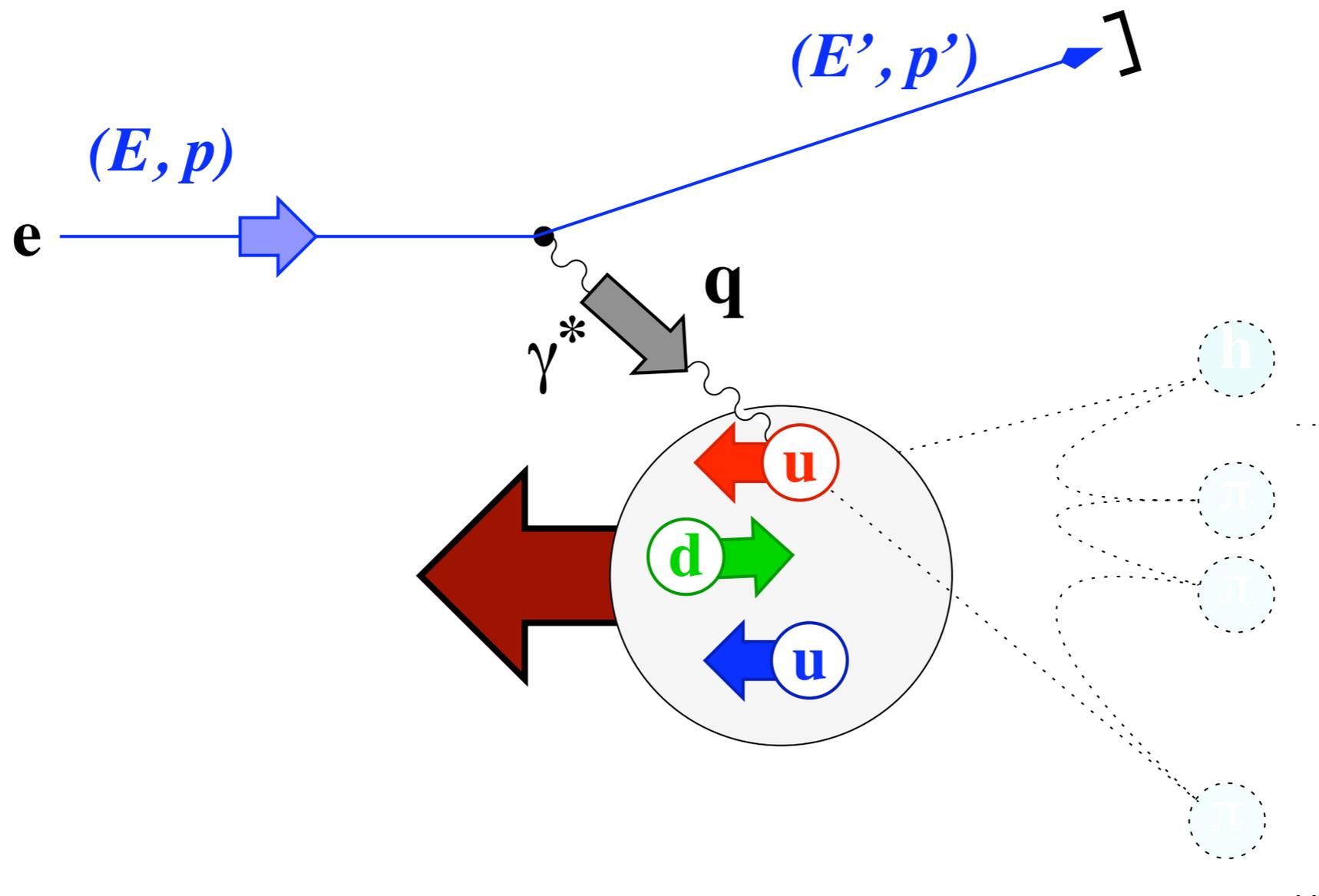
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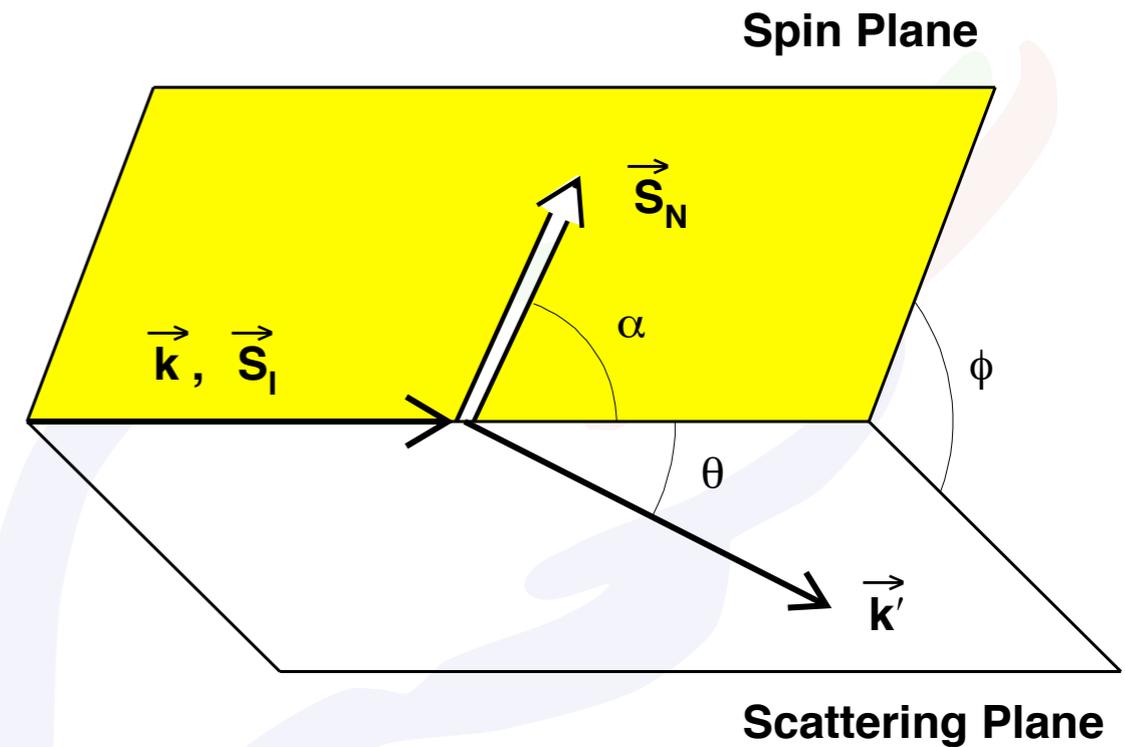
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Inclusive DIS



Inclusive DIS (one-photon exchange)

$$\frac{d^2\sigma(s, S)}{dx dQ^2} = \frac{2\pi\alpha^2 y^2}{Q^6} \mathbf{L}_{\mu\nu}(s) \mathbf{W}^{\mu\nu}(S)$$

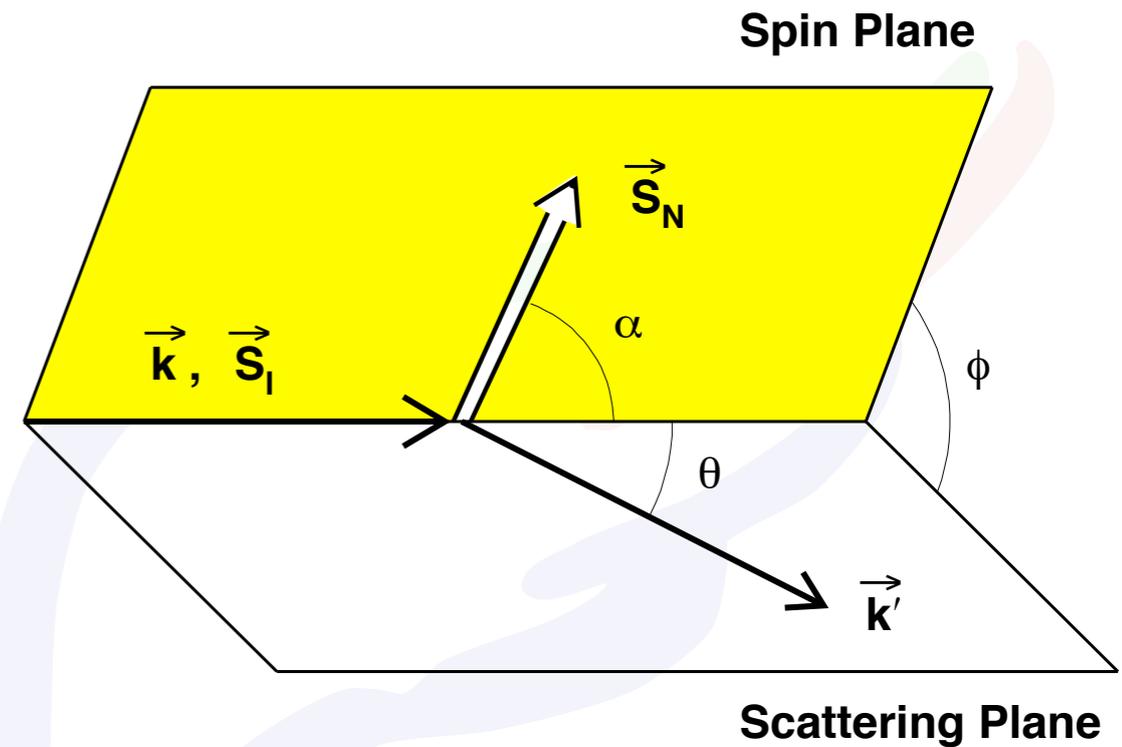


hermes

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Lepton Tensor



Inclusive DIS (one-photon exchange)

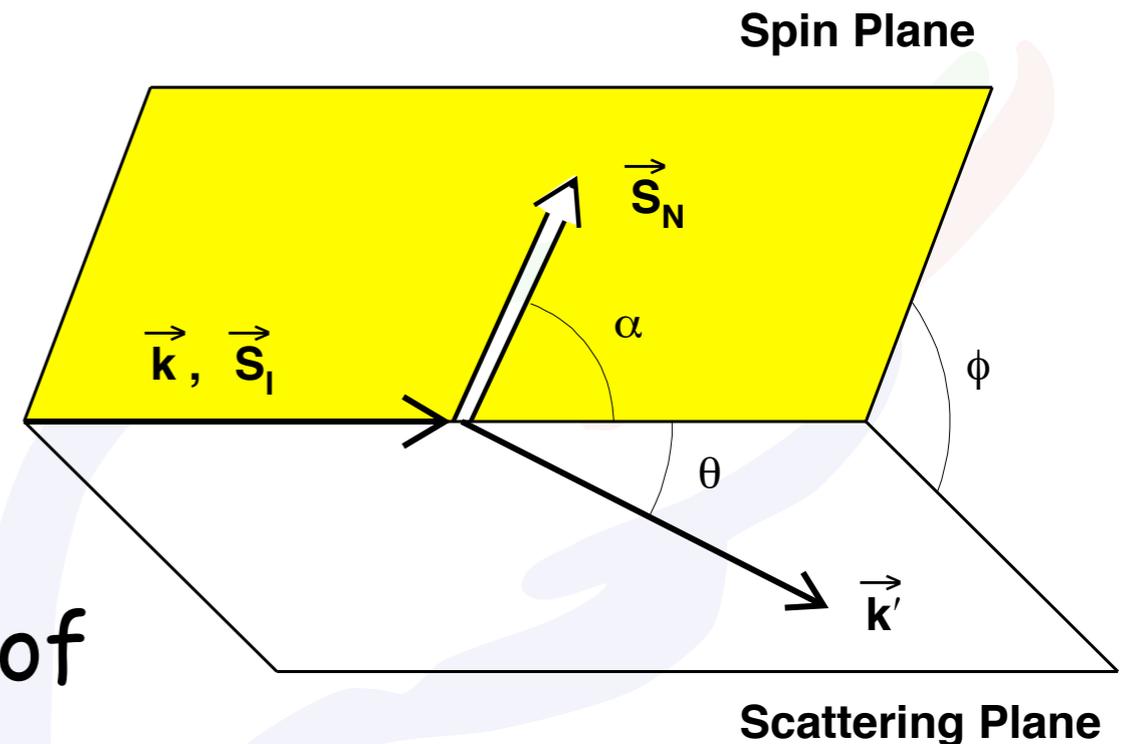
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Lepton Tensor

Hadron Tensor

parametrized in terms of

Structure Functions



Inclusive DIS (one-photon exchange)

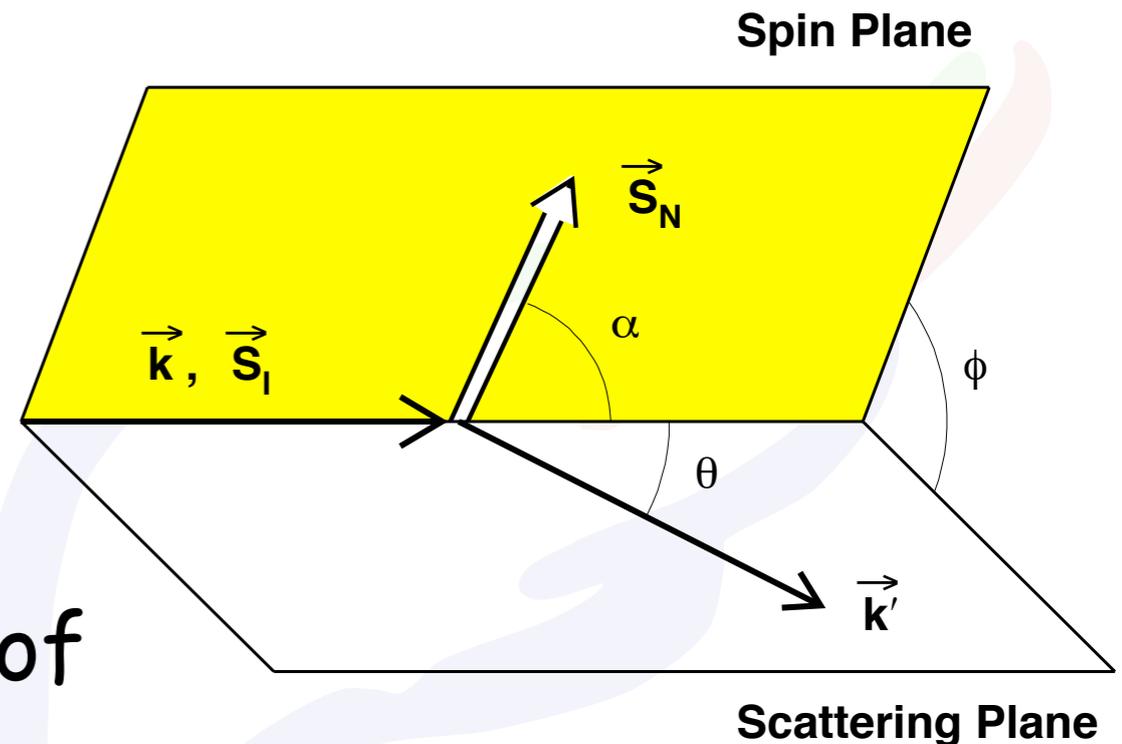
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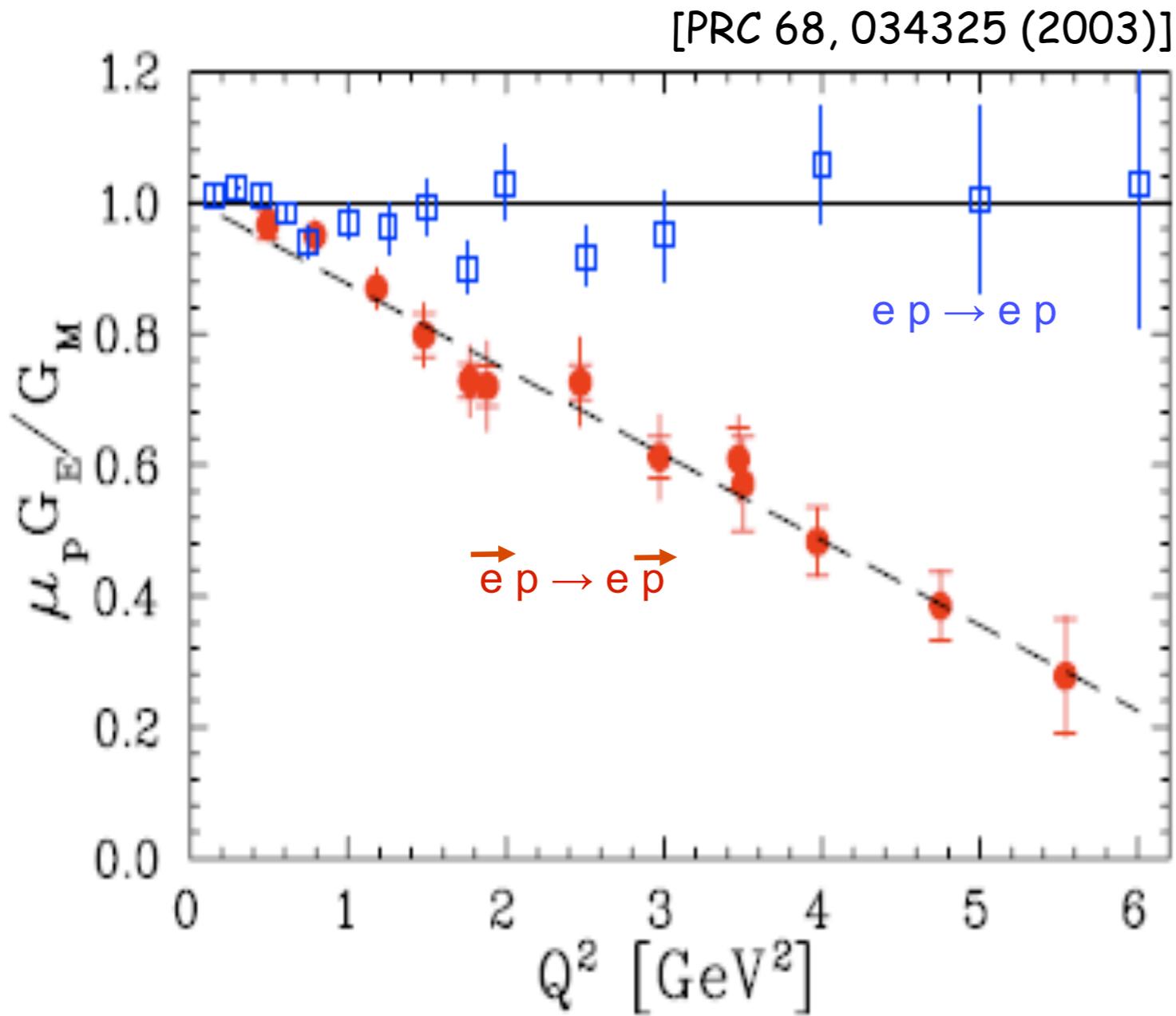


$$\frac{d^3\sigma}{dx dy d\phi} \propto \frac{y}{2} F_1(x, Q^2) + \frac{1 - y - \gamma^2 y^2 / 4}{2xy} F_2(x, Q^2) - S_l S_N \cos \alpha \left[\left(1 - \frac{y}{2} - \frac{\gamma^2 y^2}{4} \right) g_1(x, Q^2) - \frac{\gamma^2 y}{2} g_2(x, Q^2) \right] + S_l S_N \sin \alpha \cos \phi \gamma \sqrt{1 - y - \frac{\gamma^2 y^2}{4}} \left(\frac{y}{2} g_1(x, Q^2) + g_2(x, Q^2) \right)$$

Check the details!



Check the details!



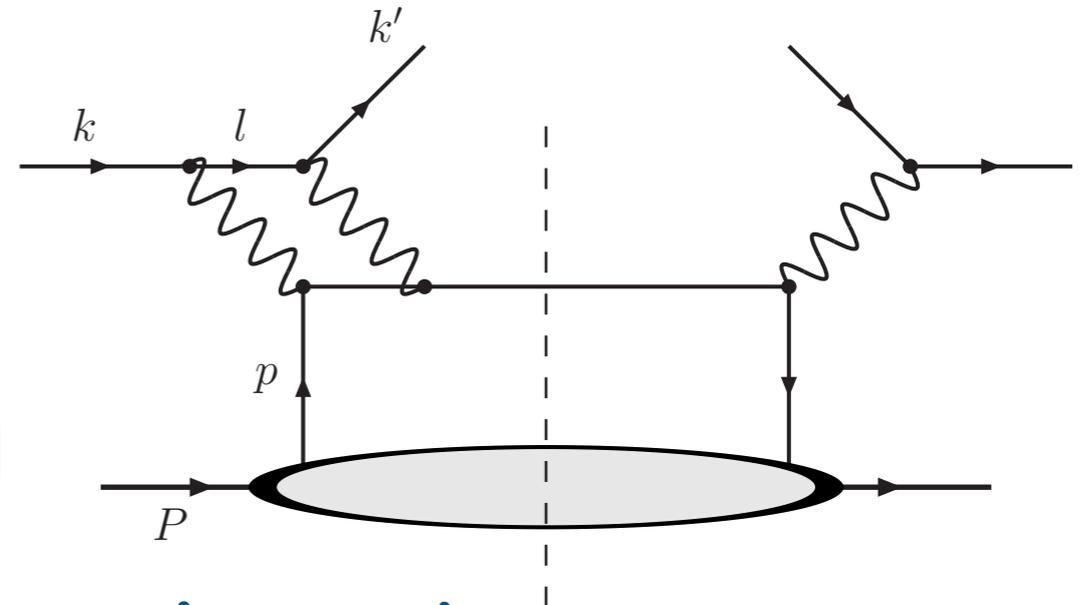
☞ two-photon exchange can be important!

x_F

nes

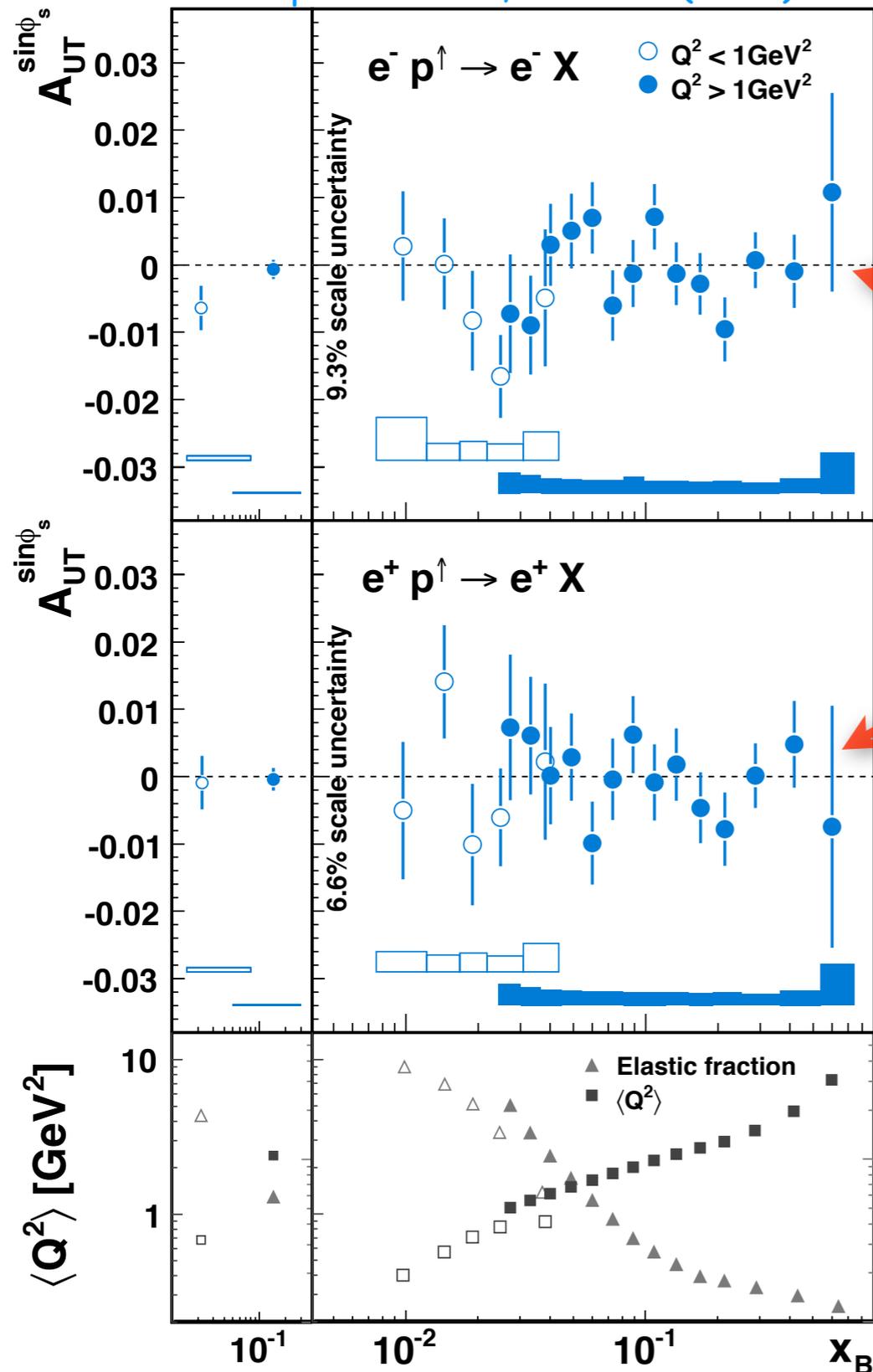
Two-photon exchange

- Candidate to explain discrepancy in form-factor measurements
- Interference between one- and two-photon exchange amplitudes leads to SSAs in inclusive DIS off transversely polarized targets
- cross section proportional to $S(kxk')$ - either measure left-right asymmetries or sine modulation
- sensitive to beam charge due to odd number of e.m. couplings to beam



No sign of two-photon exchange

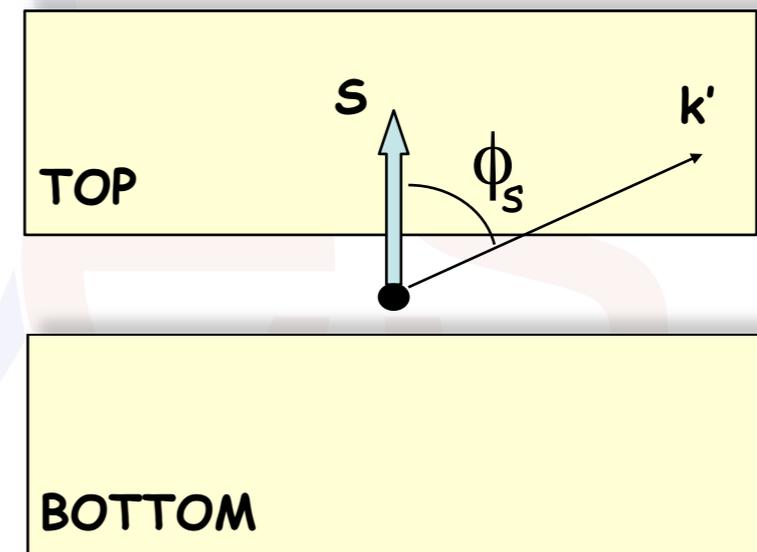
A. Airapetian et al., PLB 682 (2010) 350



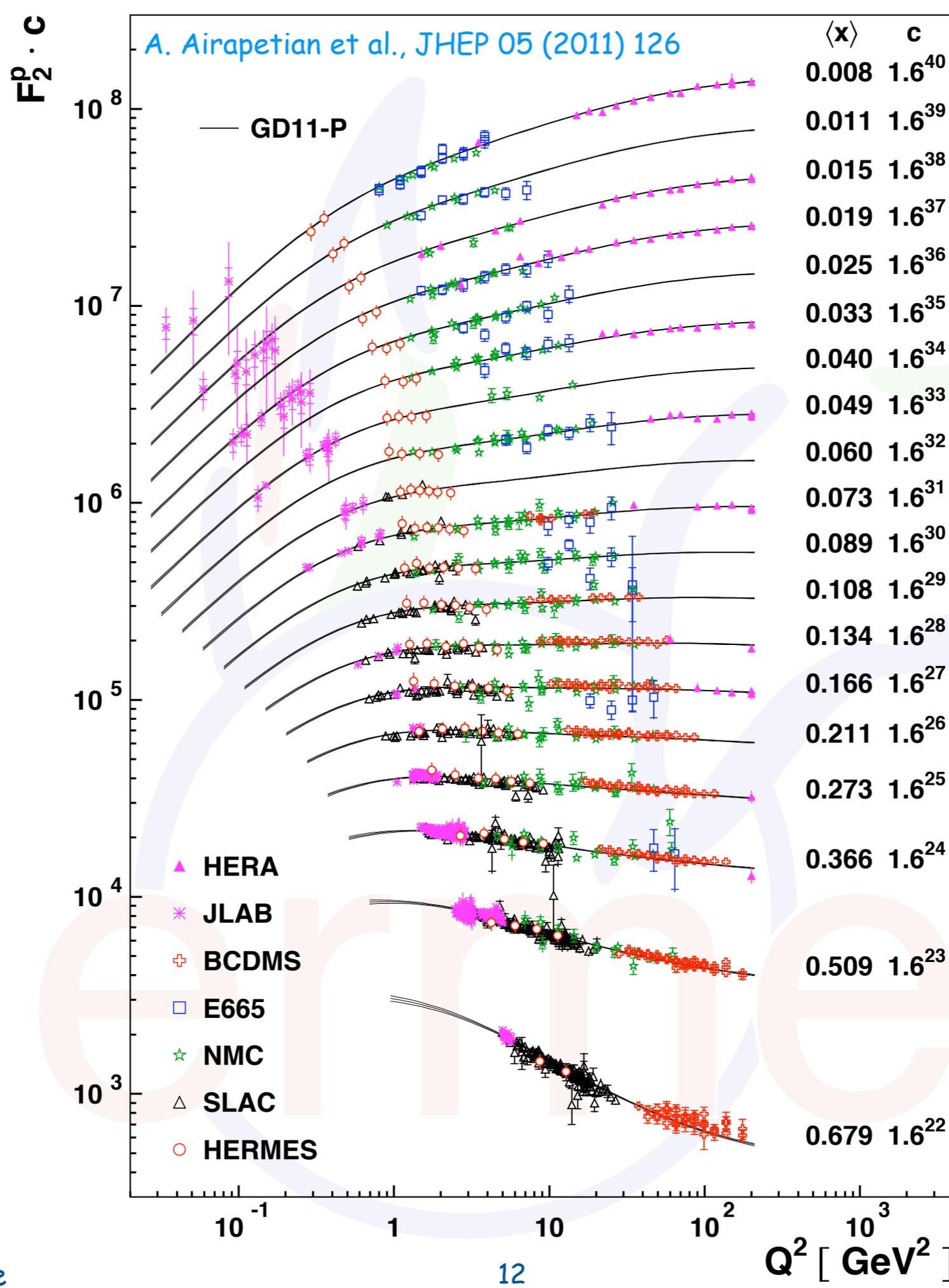
Beam	$A_{UT}^{\sin\phi_s} \times 10^{-3}$	$\delta A_{UT}^{\sin\phi_s} (\text{stat.}) \times 10^{-3}$	$\delta A_{UT}^{\sin\phi_s} (\text{syst.}) \times 10^{-3}$	$\langle x_B \rangle$	$\langle Q^2 \rangle [\text{GeV}^2]$
e^+	-0.61	3.97	0.63	0.02	0.68
e^-	-6.55	3.40	0.63		
e^+	-0.60	1.70	0.29	0.14	2.40
e^-	-0.85	1.50	0.29		

consistent with zero for both e^+/e^-

Front view of HERMES detector



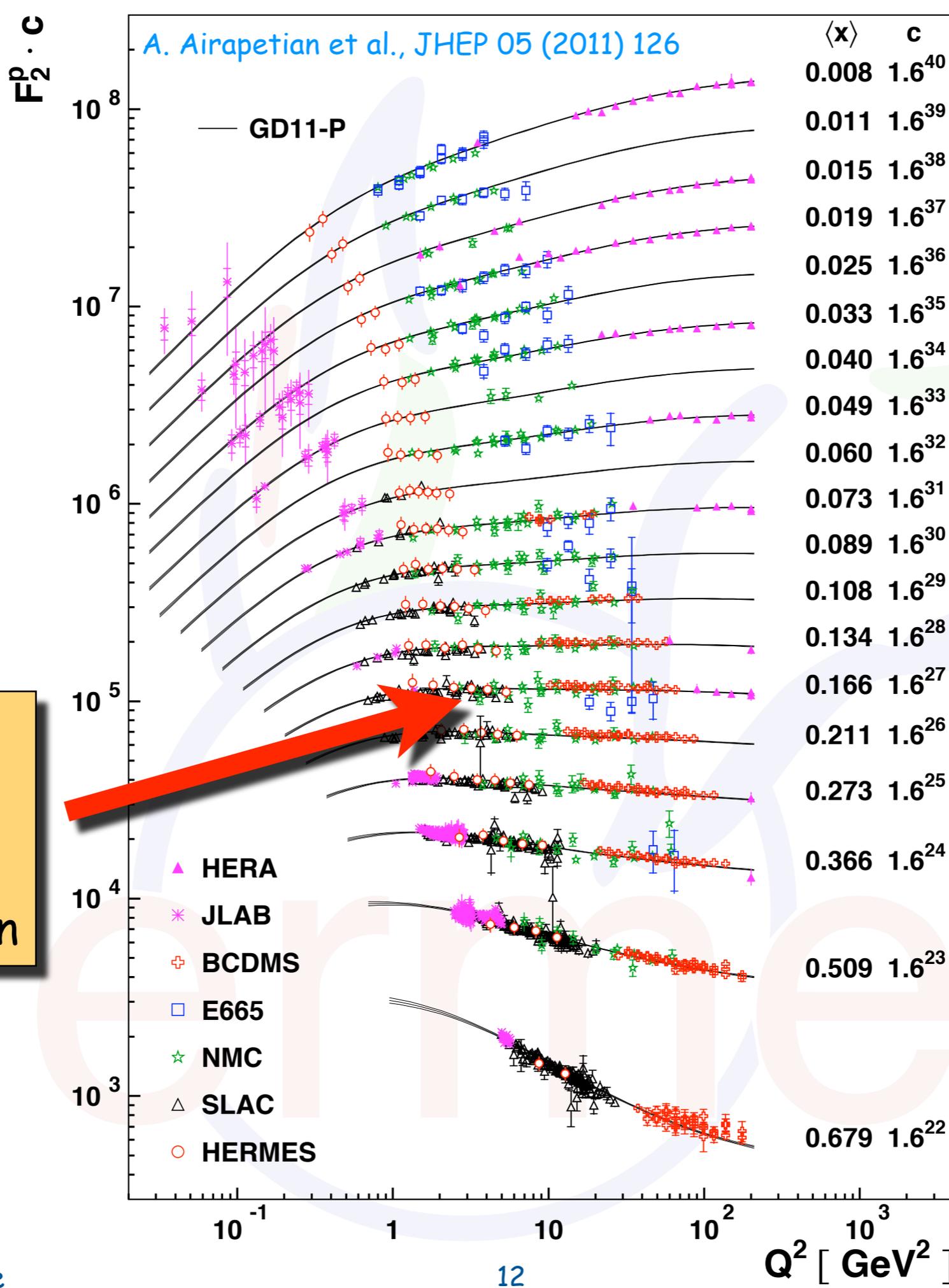
F_2 proton



GD11 - global fit

From global fit GD11:
 HERMES relative normalization
 ~2% for proton and deuteron
 ~0.5% for their ratio

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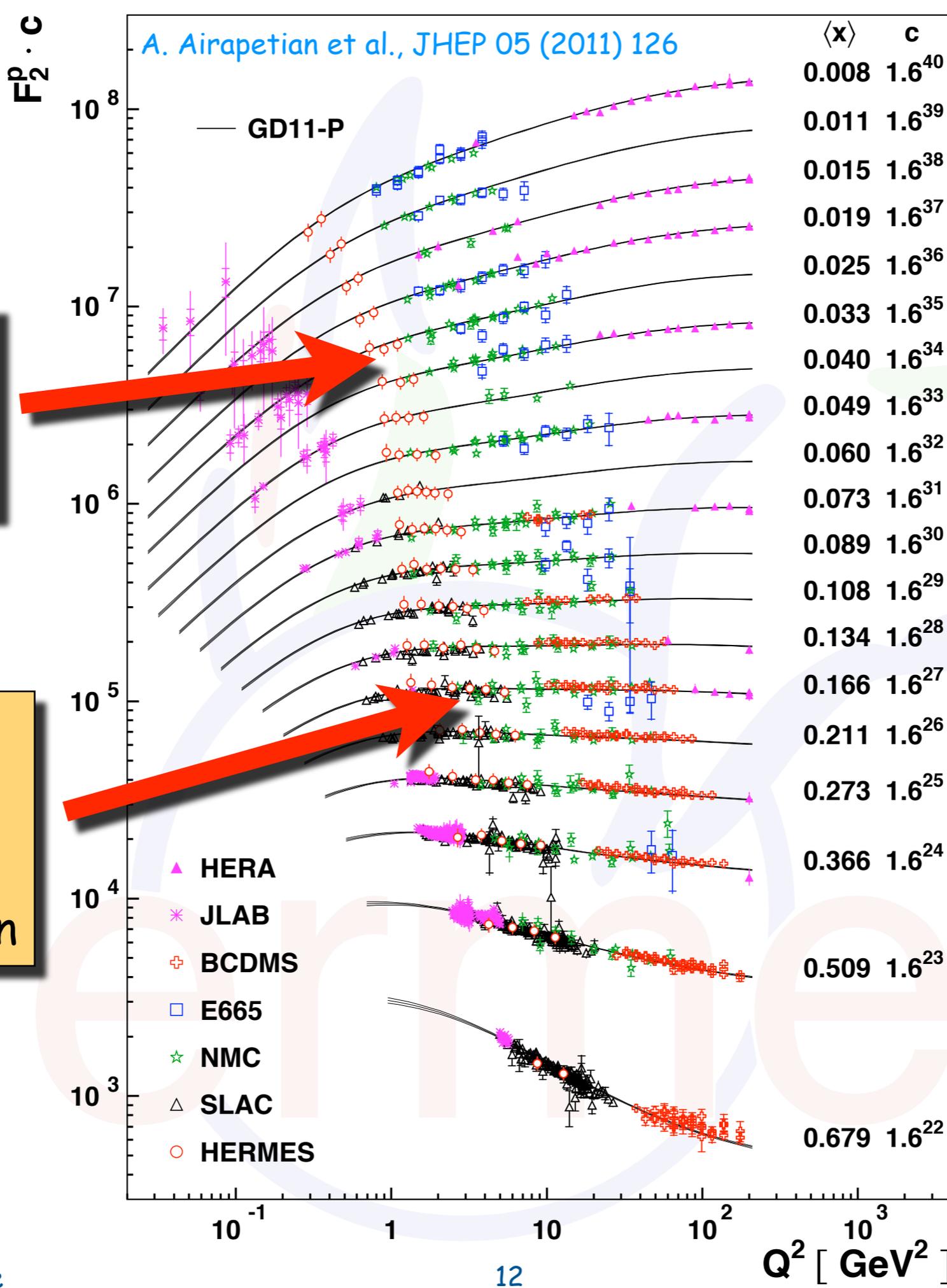
Agreement with world data in the overlap region

F_2 proton

A. Airapetian et al., JHEP 05 (2011) 126

New region covered by HERMES

Agreement with world data in the overlap region



$\langle x \rangle$	c
0.008	1.6^{40}
0.011	1.6^{39}
0.015	1.6^{38}
0.019	1.6^{37}
0.025	1.6^{36}
0.033	1.6^{35}
0.040	1.6^{34}
0.049	1.6^{33}
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0.509	1.6^{23}
0.679	1.6^{22}

GD11 - global fit

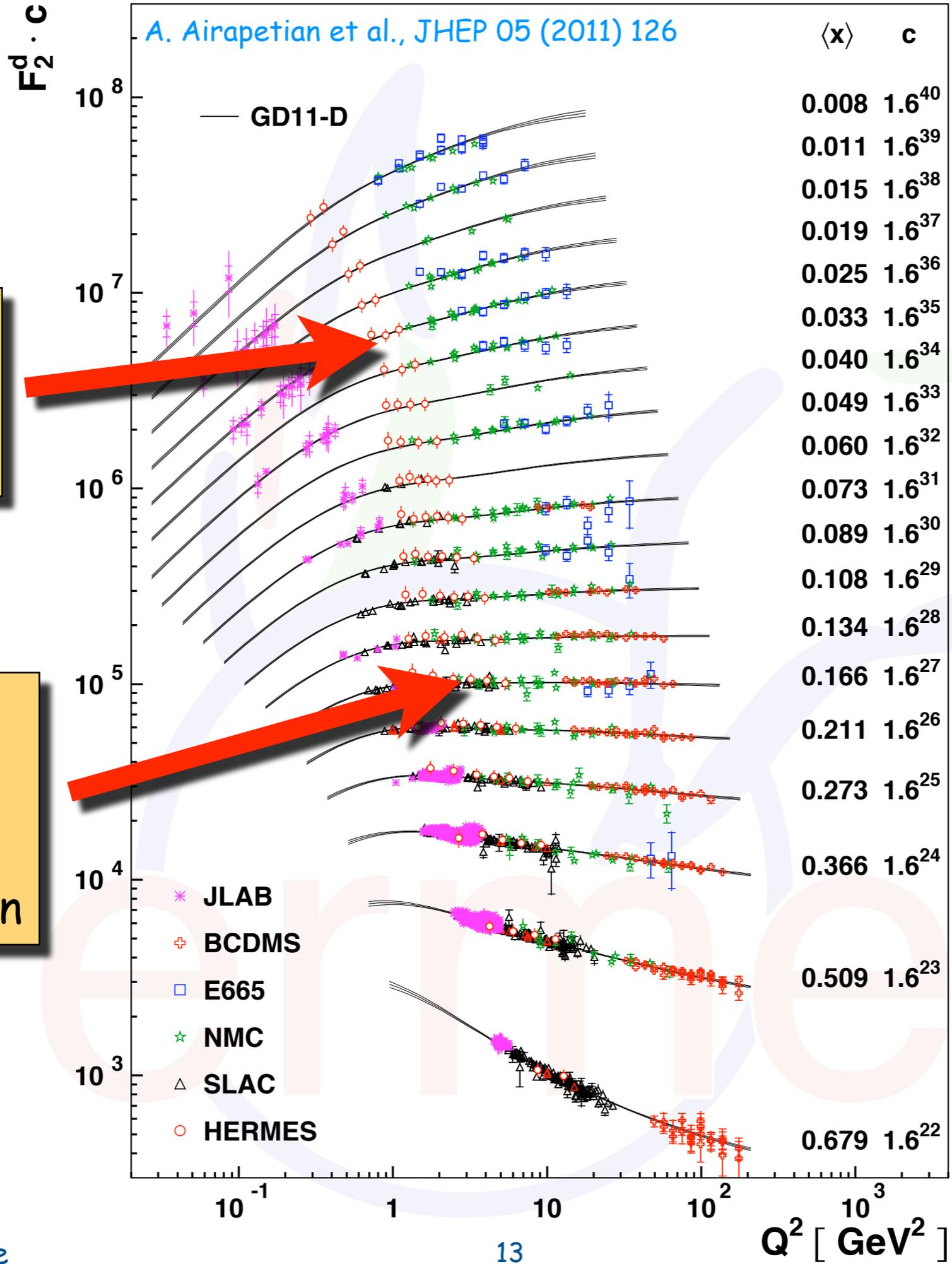
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F_2 deuteron

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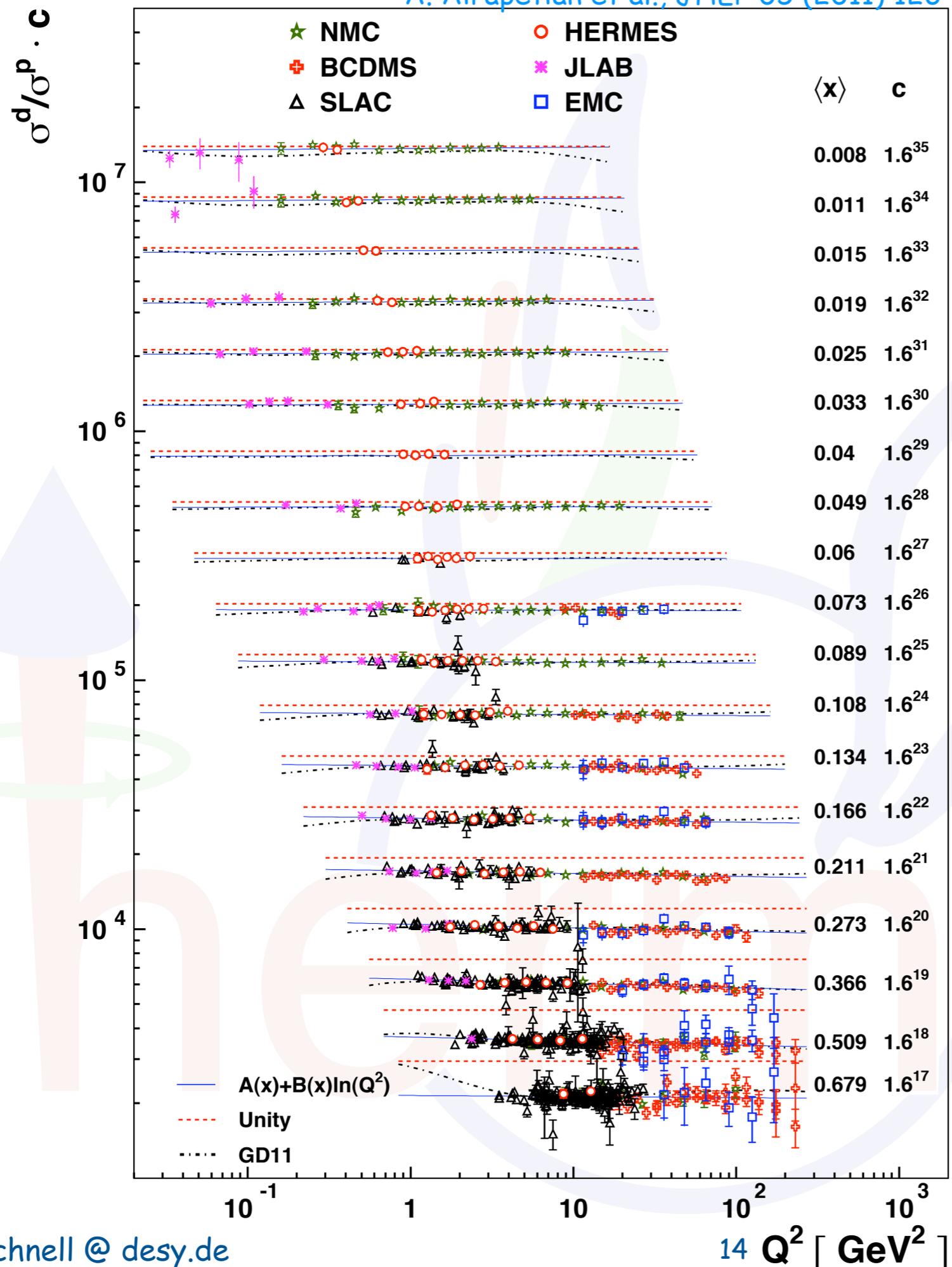


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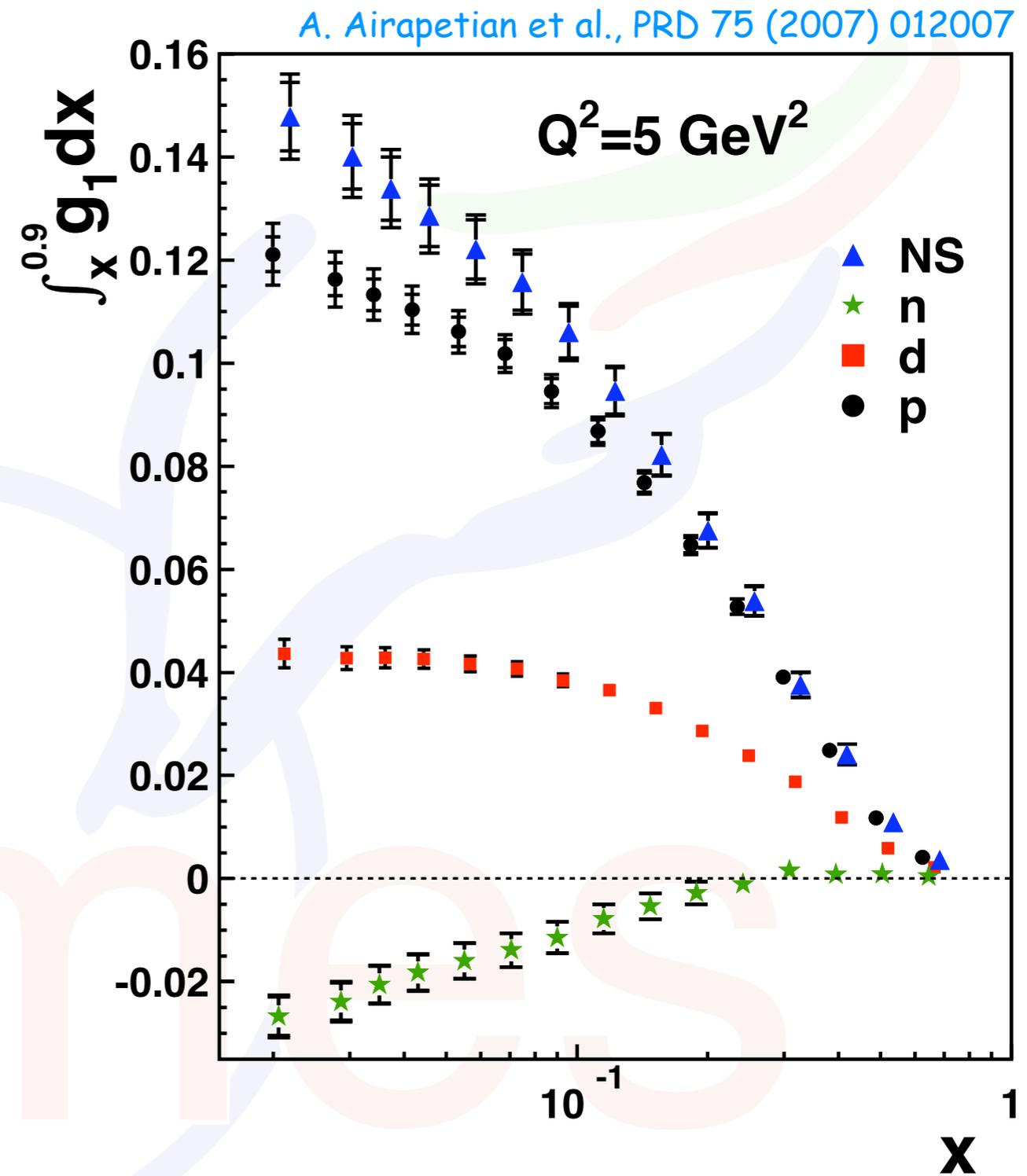
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$$\sigma^d / \sigma^p$$



Many systematic errors common to proton and deuteron cross sections cancel in ratio

Polarized structure function g_1



Polarized structure function g_1

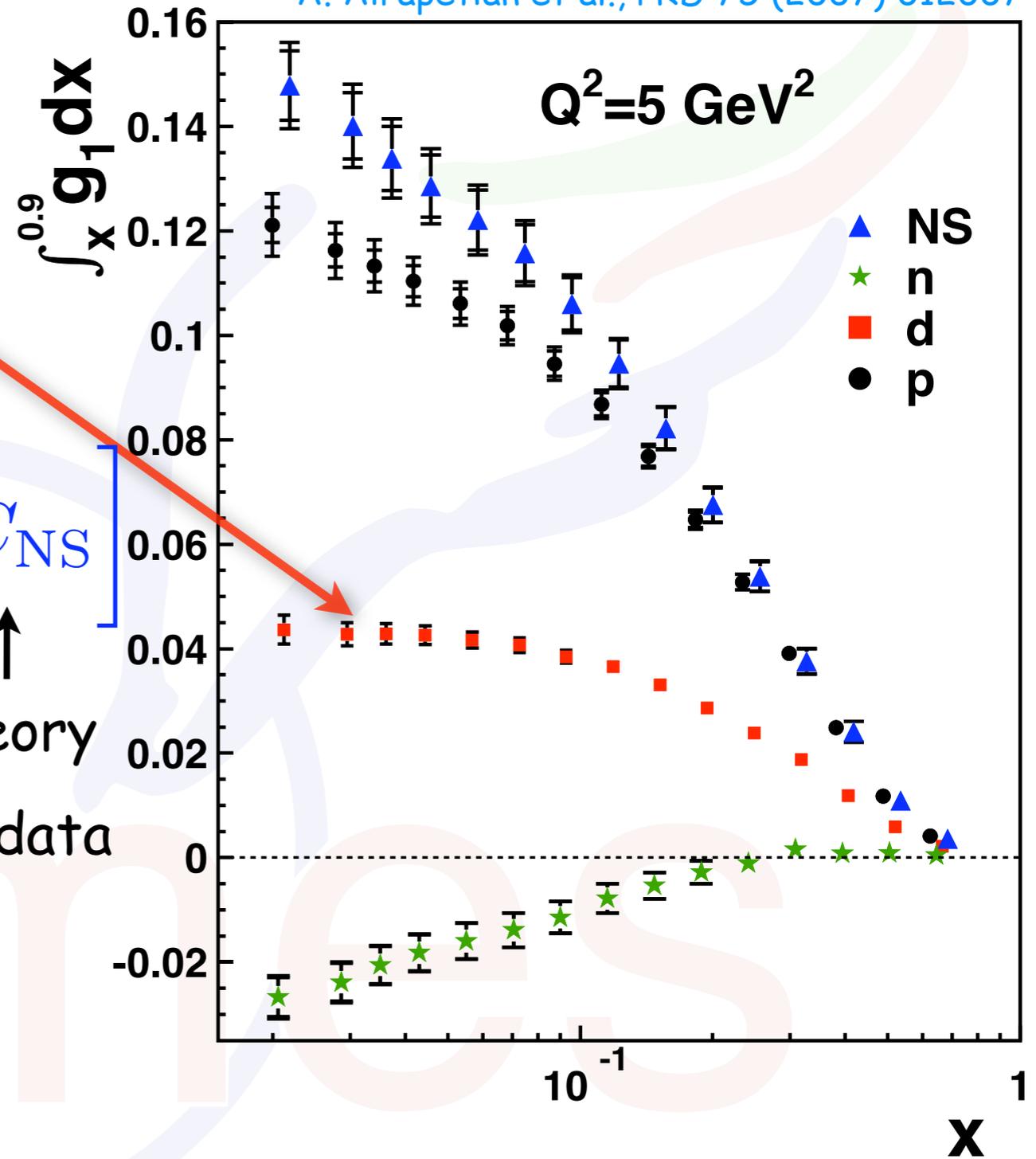
Saturation

→ **close to full integral?**

$$\Delta\Sigma \stackrel{\overline{\text{MS}}}{=} \frac{1}{\Delta C_S} \left[\frac{9\Gamma_1^d}{1 - \frac{3}{2}\omega_D} - \frac{1}{4} a_8 \Delta C_{NS} \right]$$

\uparrow theory \uparrow 0.05 \pm 0.05 \uparrow theory
 hyperon-decay data

A. Airapetian et al., PRD 75 (2007) 012007

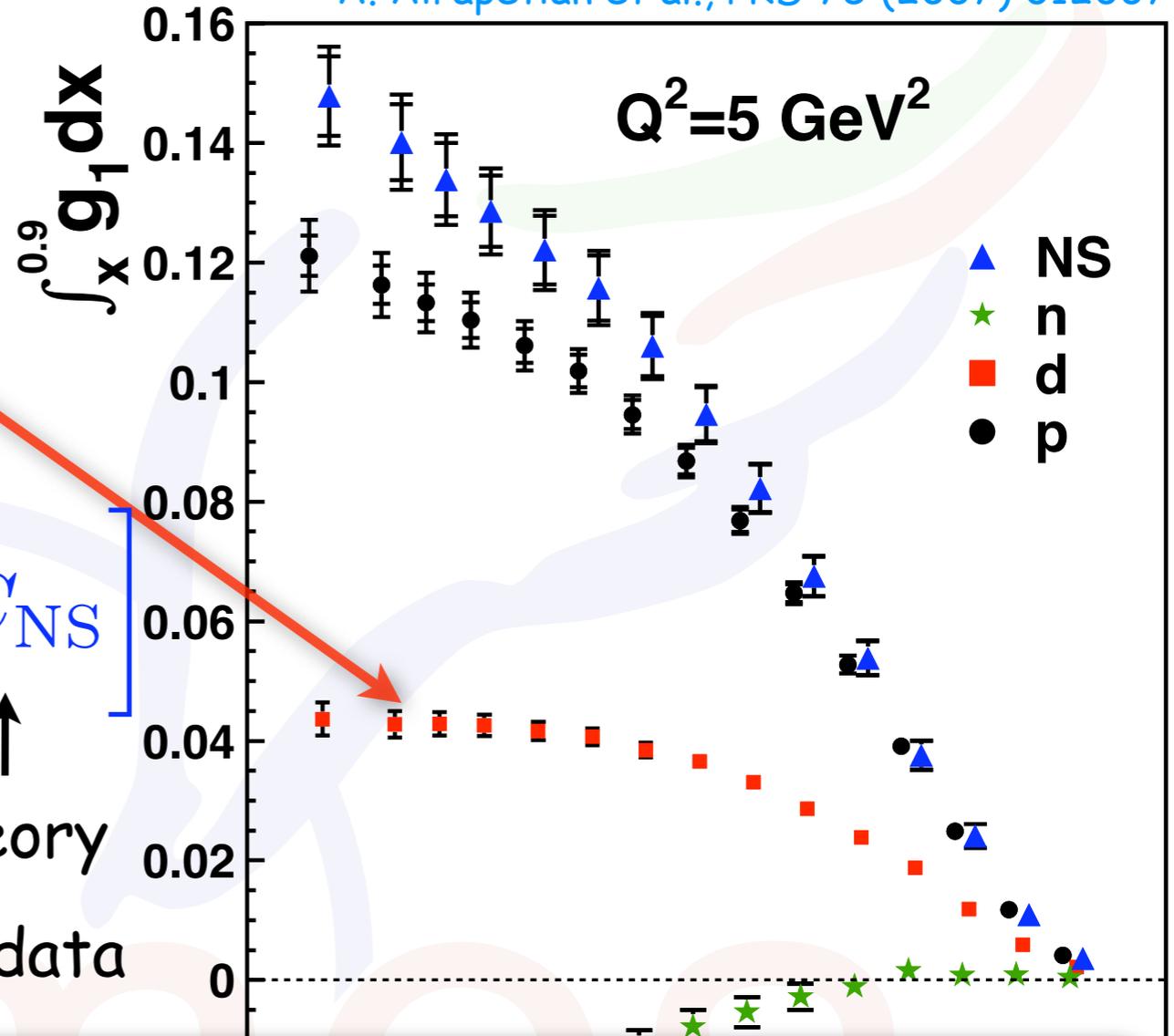


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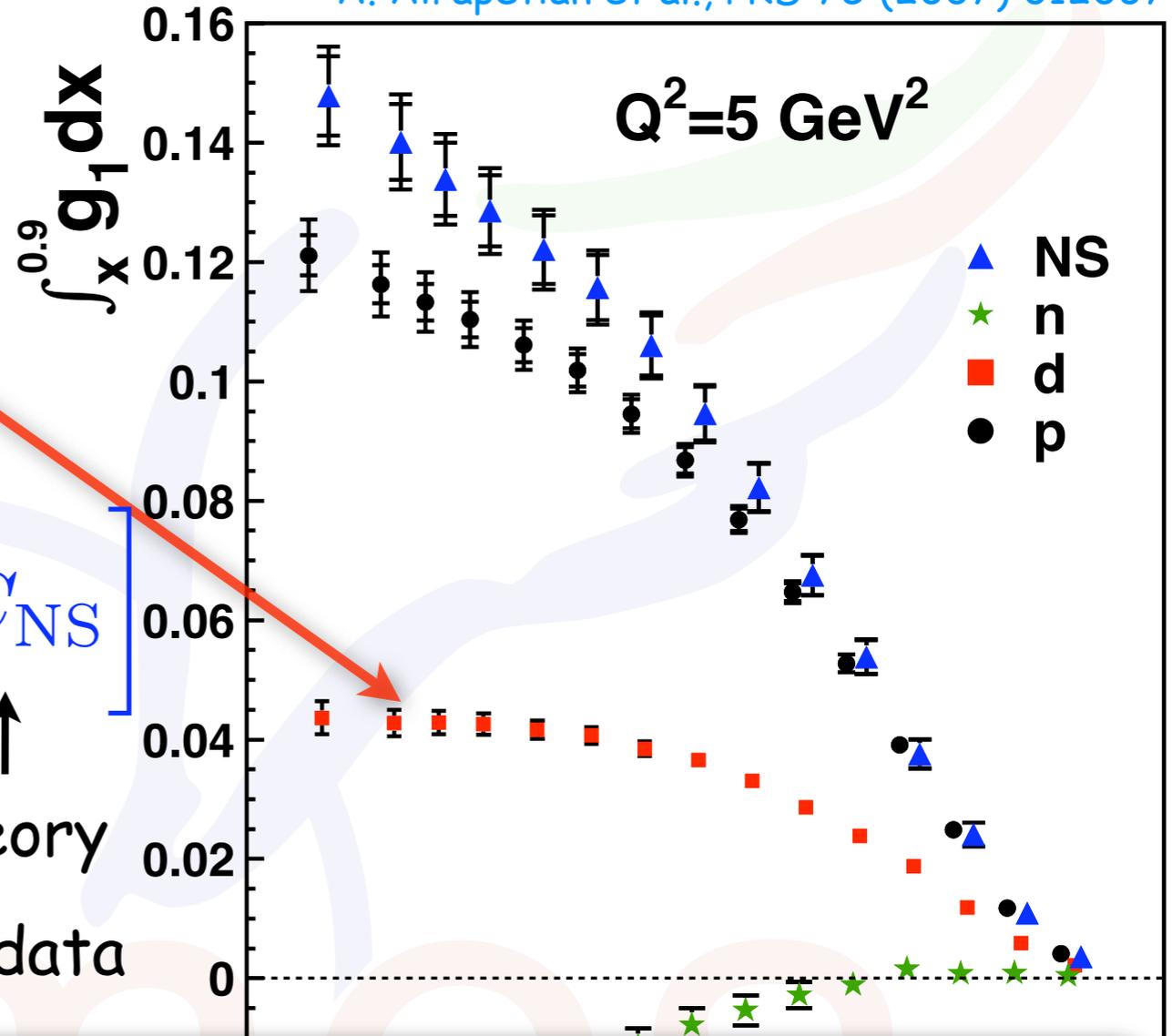
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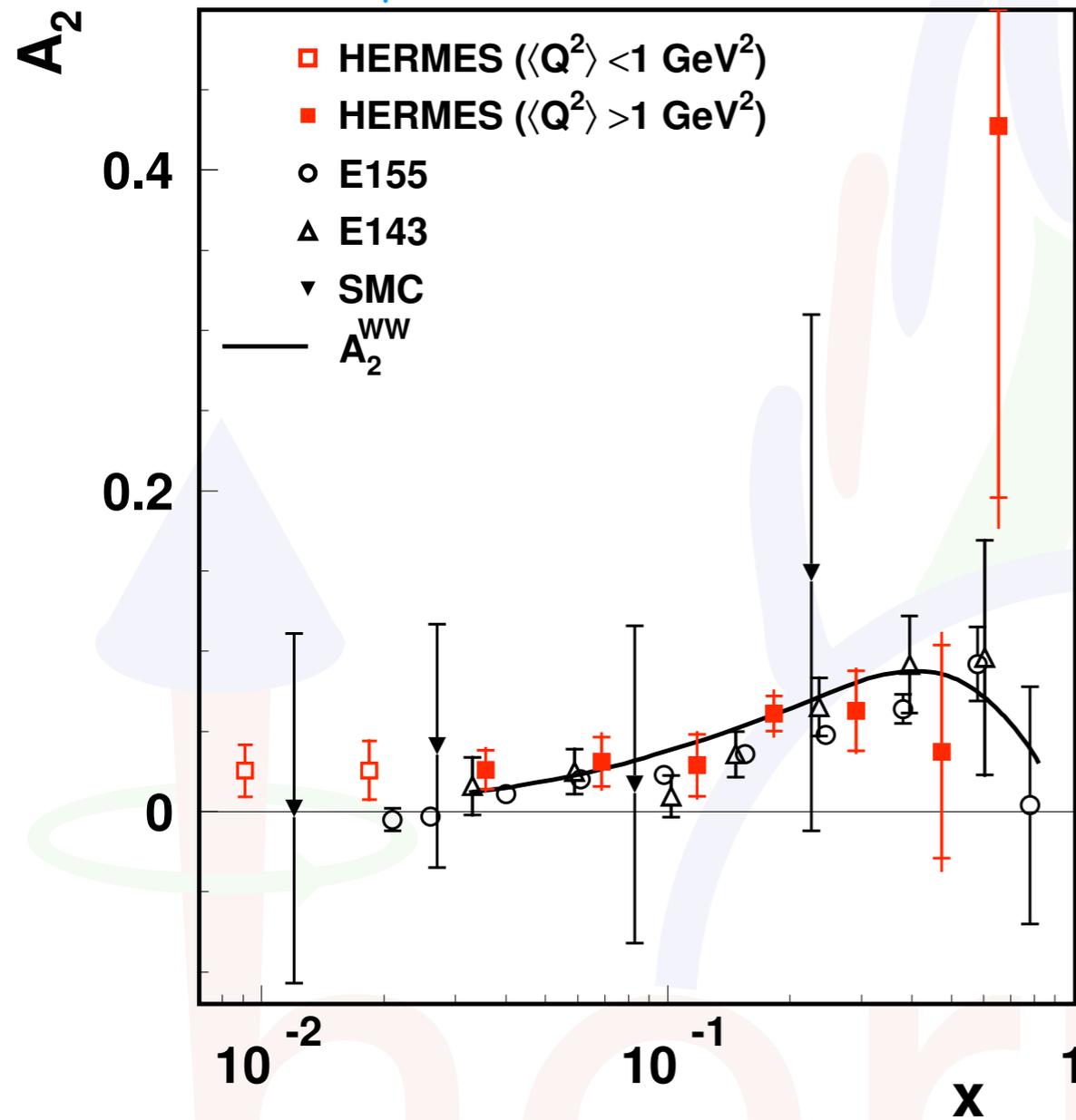
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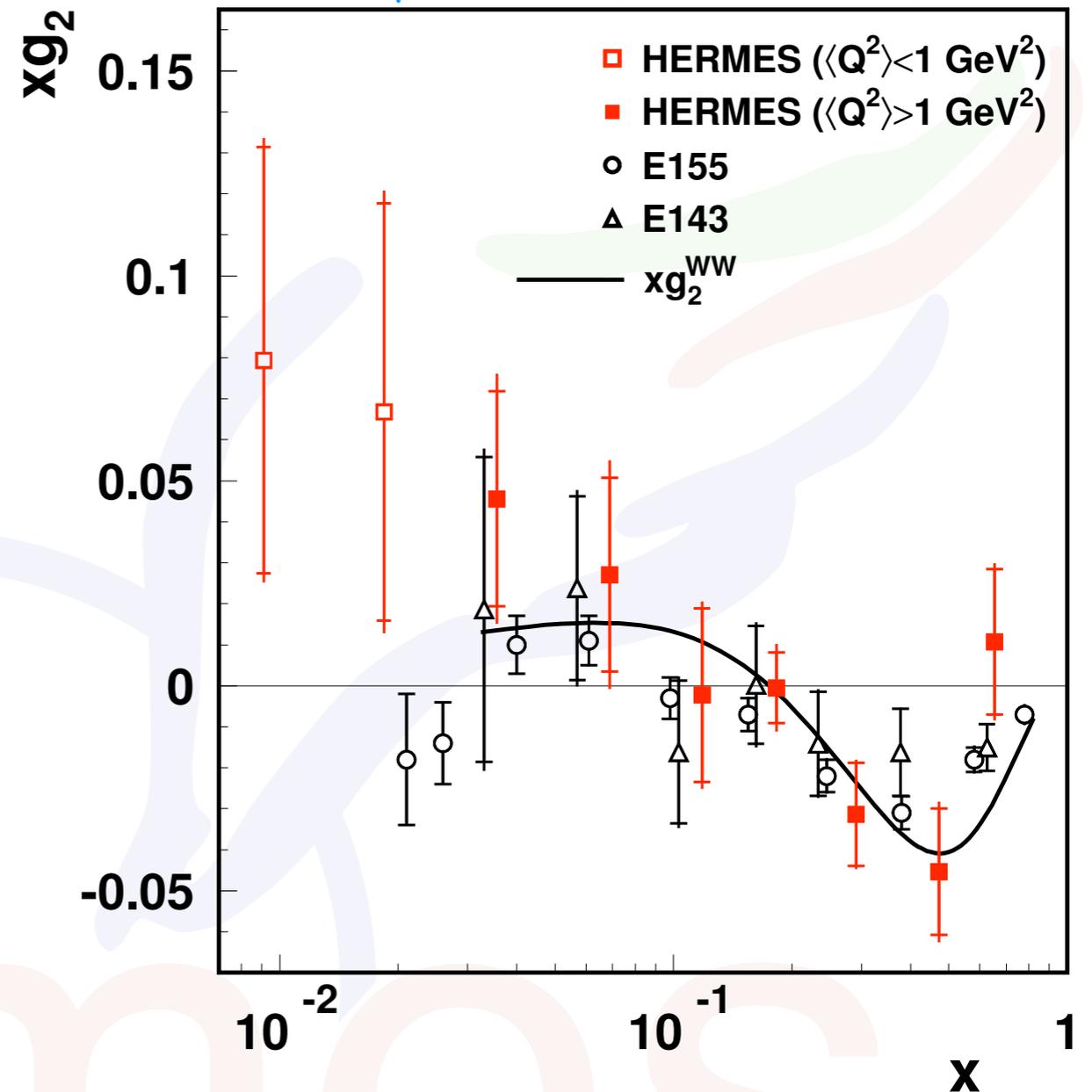
most precise result; **only 1/3** of nucleon spin from quarks

Results on A_2 and xg_2

A. Airapetian et al., EPJ C72 (2012) 1921



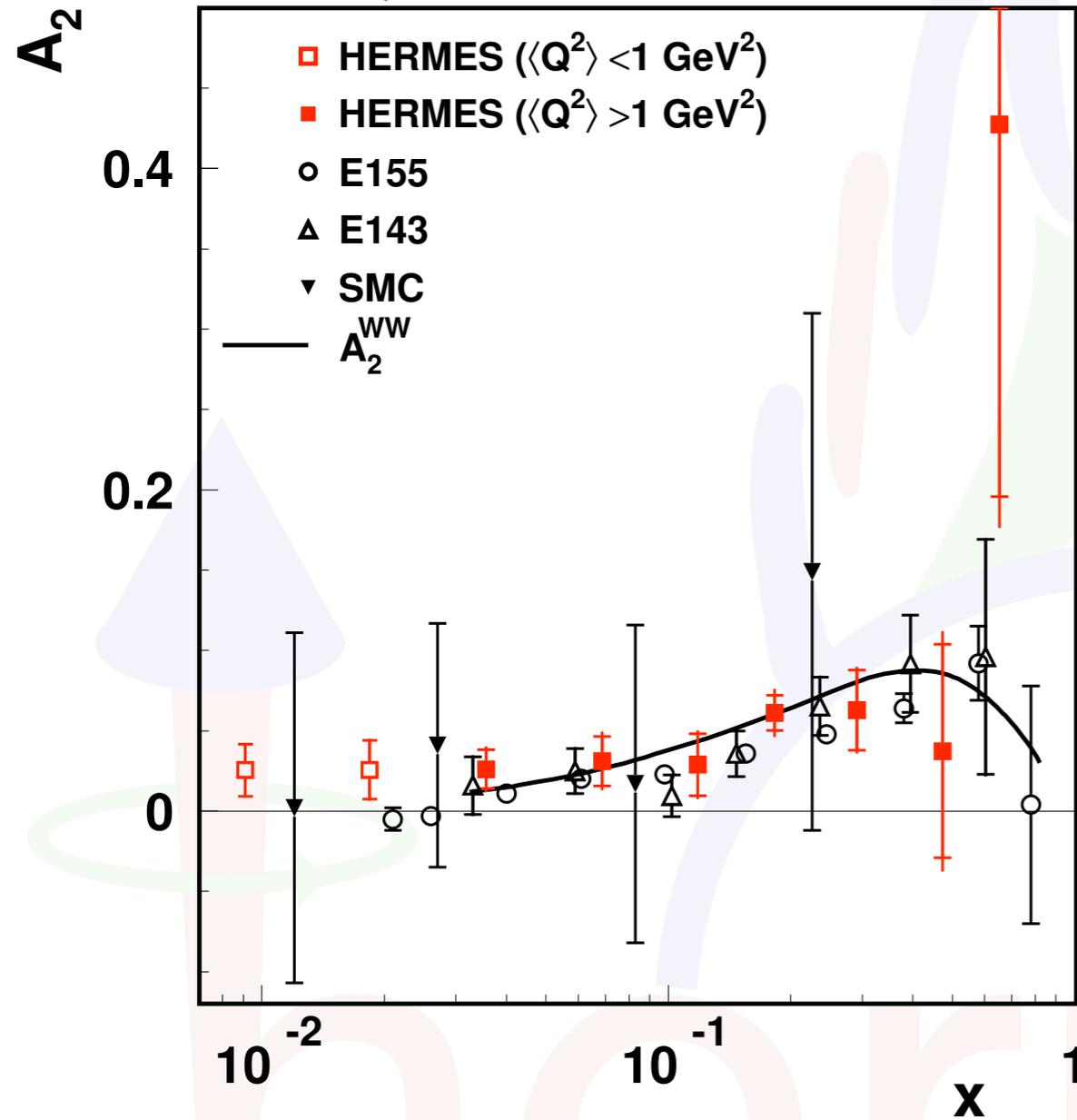
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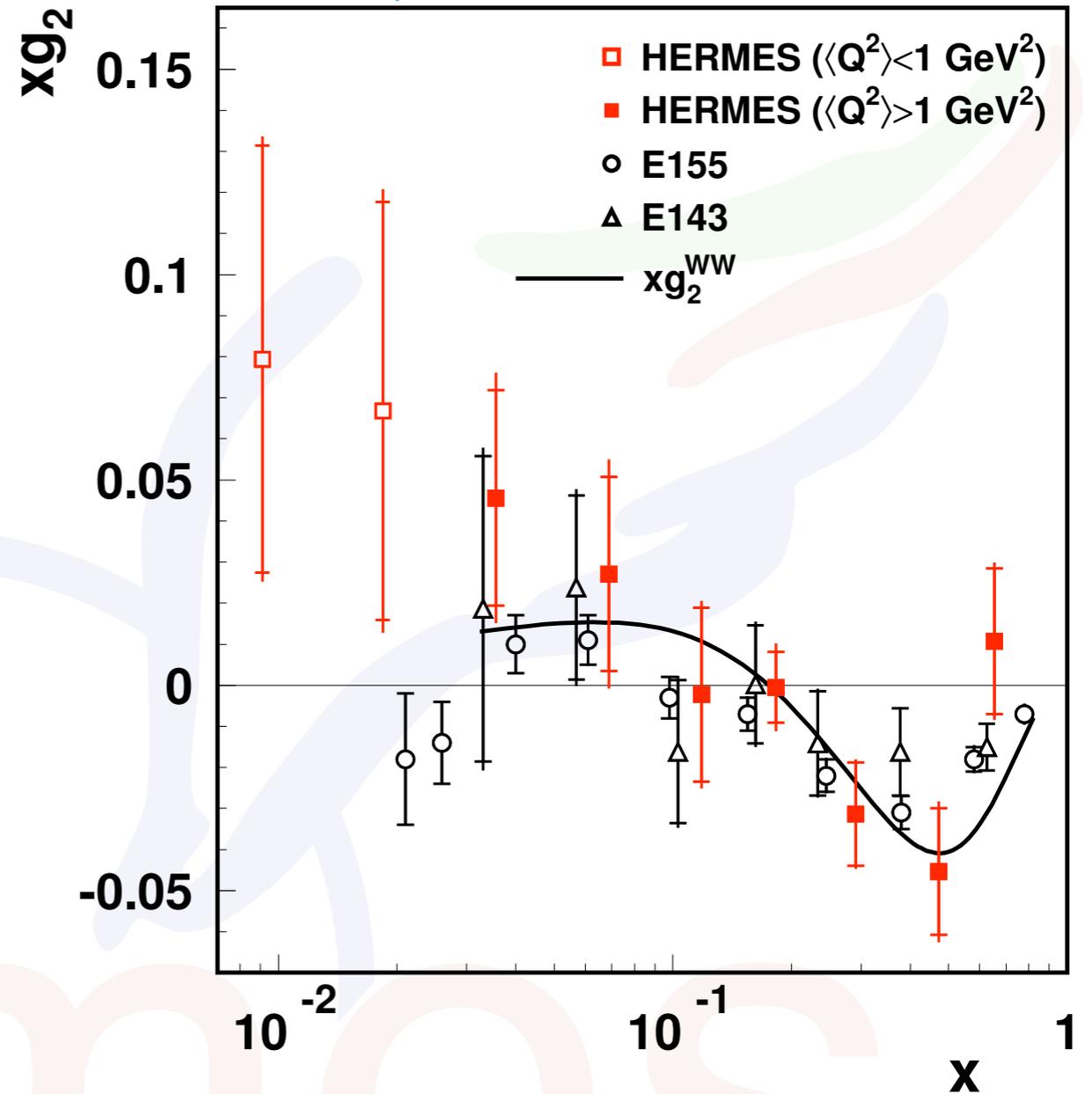
- consistent with (sparse) world data
- low beam polarization during HERA II \rightarrow small f.o.m.

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A. Airapetian et al., EPJ C72 (2012) 1921



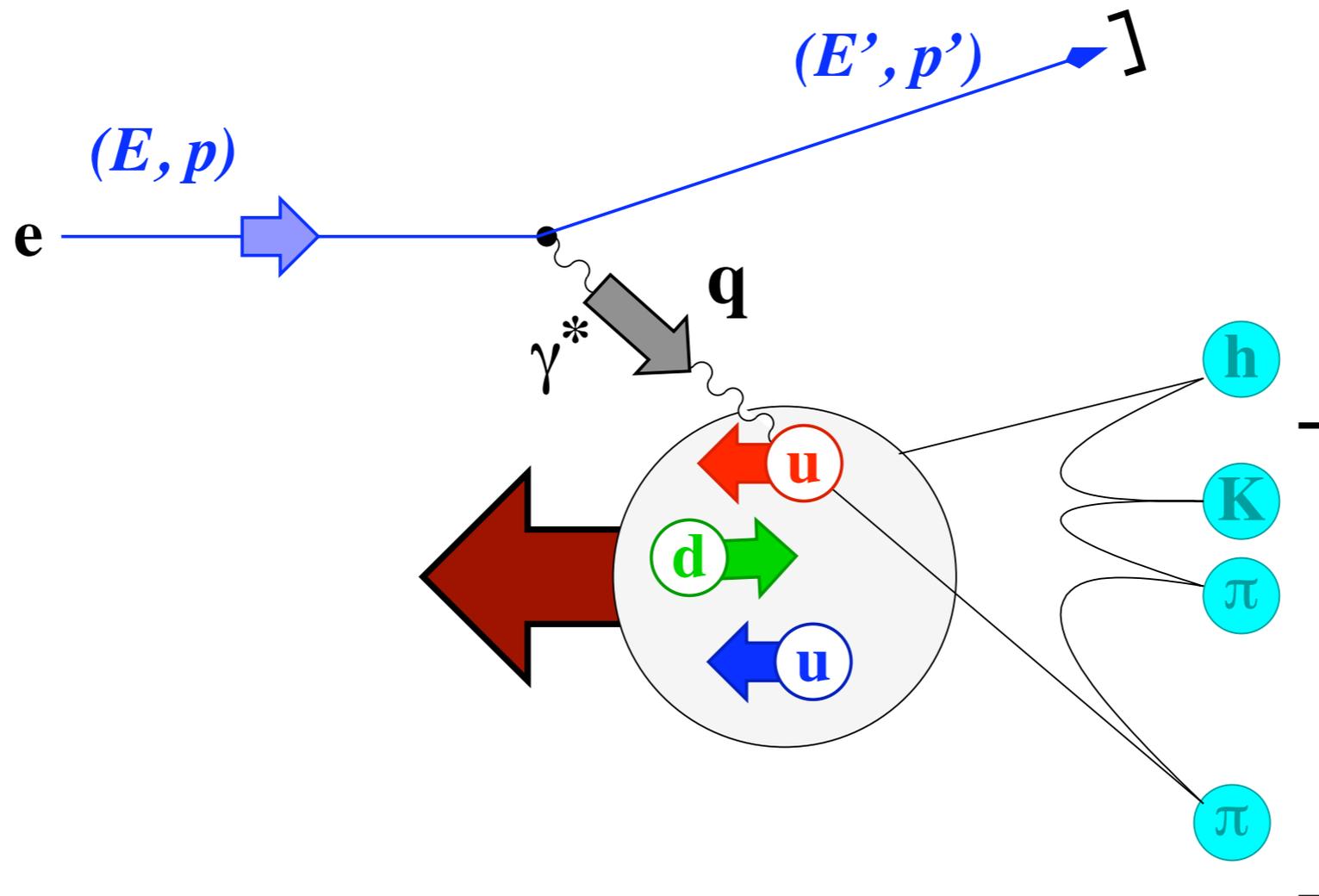
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$$\int_{0.023}^{0.9} g_2(x, Q^2) dx = 0.006 \pm 0.024_{\text{stat}} \pm 0.017_{\text{syst}}$$

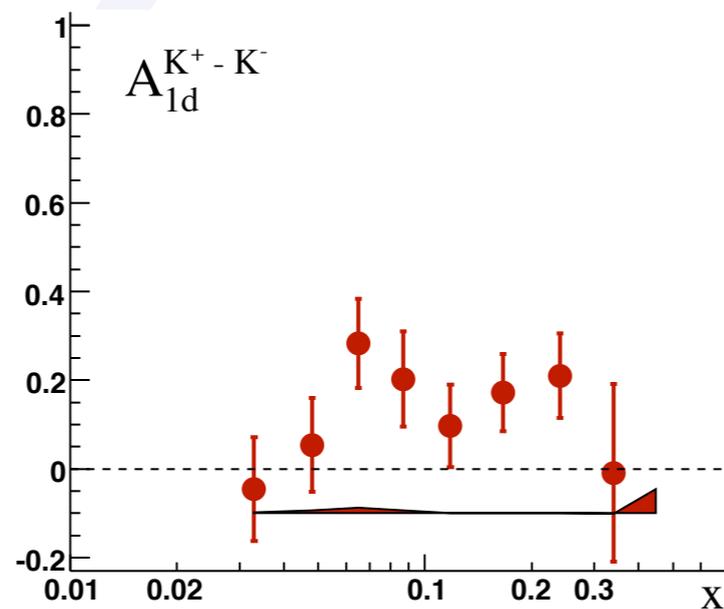
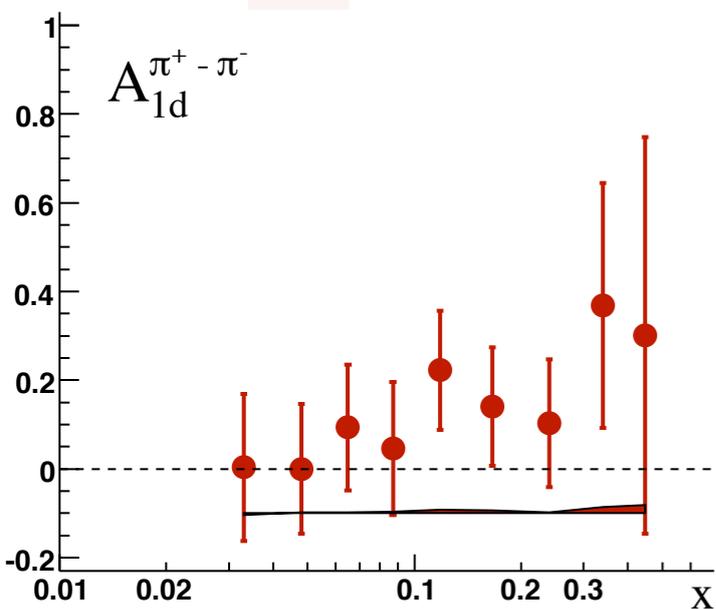
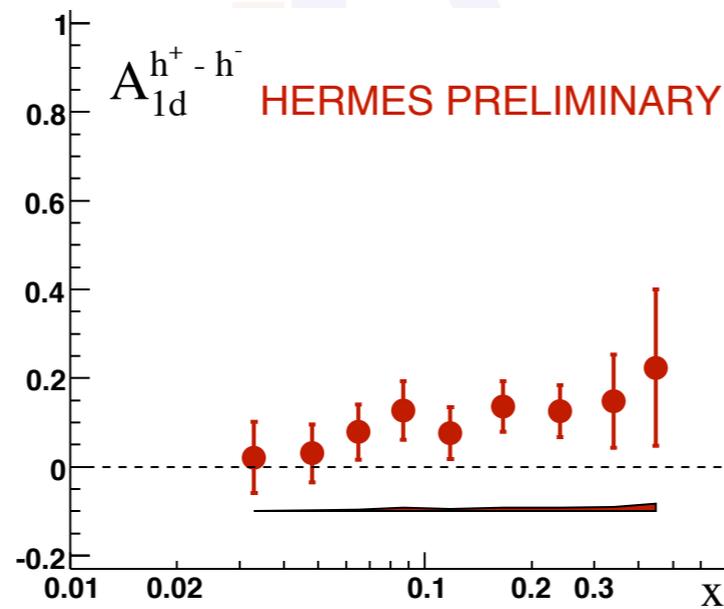
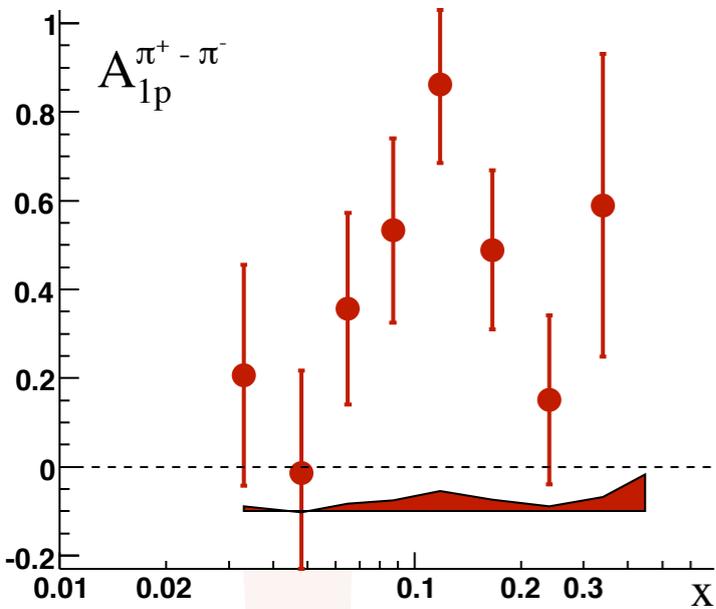
$$d_2(Q^2) \equiv 3 \int_0^1 x^2 \bar{g}_2(x, Q^2) dx = 0.0148 \pm 0.0096_{\text{stat}} \pm 0.0048_{\text{syst}}$$

Semi-Inclusive DIS



Helicity density - valence quarks

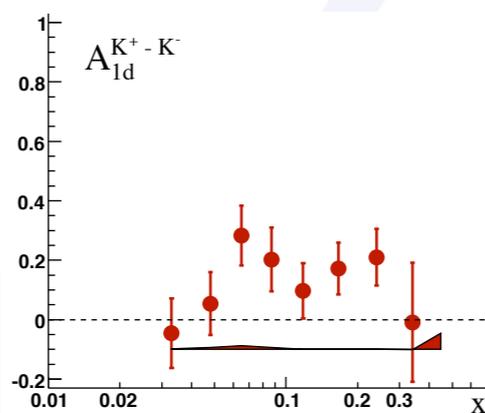
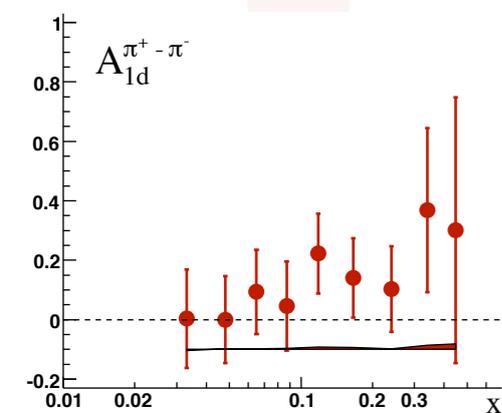
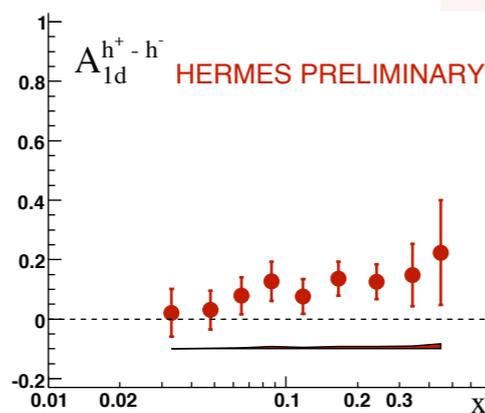
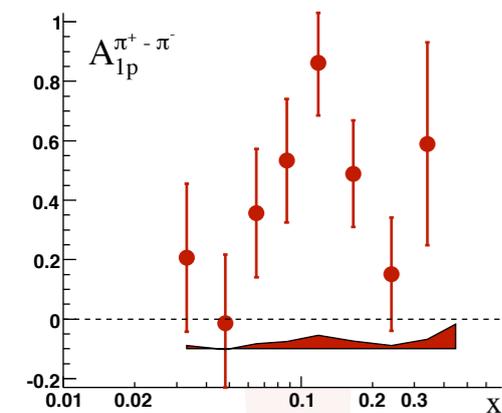
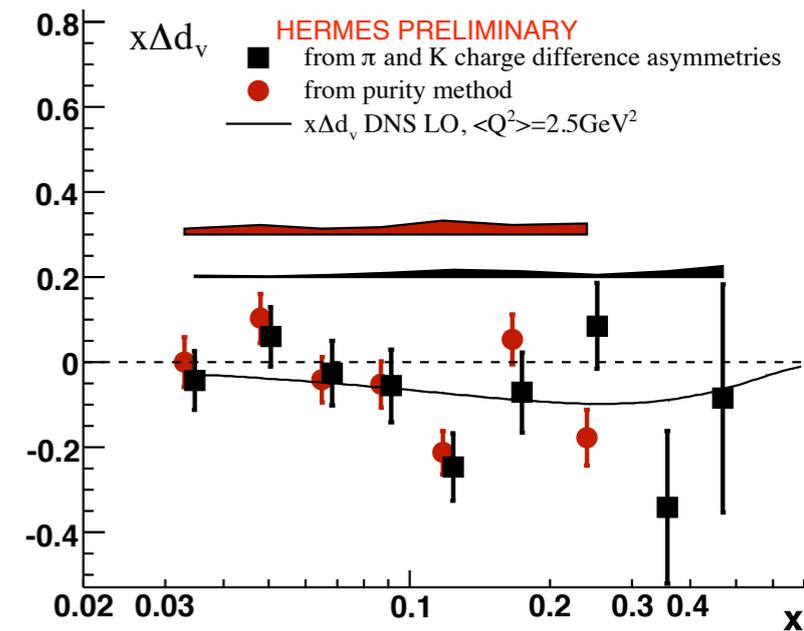
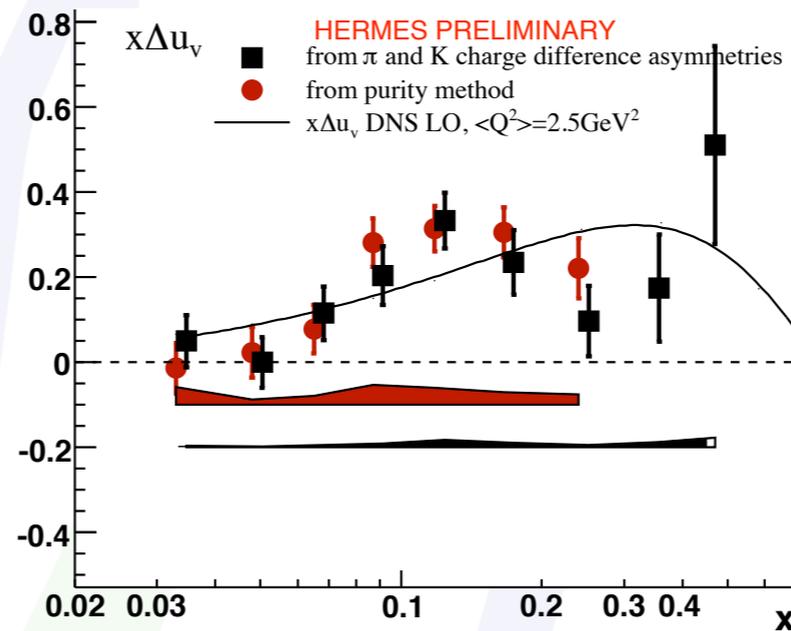
$$A_1^{h^+ - h^-} = \frac{(d\sigma_{h^+}^{\leftarrow} - d\sigma_{h^-}^{\leftarrow}) - (d\sigma_{h^+}^{\rightarrow} - d\sigma_{h^-}^{\rightarrow})}{(d\sigma_{h^+}^{\leftarrow} - d\sigma_{h^-}^{\leftarrow}) + (d\sigma_{h^+}^{\rightarrow} - d\sigma_{h^-}^{\rightarrow})}$$



● charge-difference double-spin asymmetries

Helicity density - valence quarks

$$A_1^{h^+ - h^-} = \frac{(d\sigma_{h^+}^{\leftarrow} - d\sigma_{h^-}^{\leftarrow}) - (d\sigma_{h^+}^{\rightarrow} - d\sigma_{h^-}^{\rightarrow})}{(d\sigma_{h^+}^{\leftarrow} - d\sigma_{h^-}^{\leftarrow}) + (d\sigma_{h^+}^{\rightarrow} - d\sigma_{h^-}^{\rightarrow})}$$



- charge-difference double-spin asymmetries
- use charge-conjugation symmetry to extract, at LO(!), valence distributions

$$A_{1p}^{h^+ - h^-} \cong \frac{4\Delta u_v - \Delta d_v}{4u_v - d_v}$$

$$A_{1d}^{h^+ - h^-} \cong \frac{\Delta u_v + \Delta d_v}{u_v + d_v}$$

... going 3D

Spin-momentum structure of the nucleon

$$\frac{1}{2}\text{Tr}\left[(\gamma^+ + \lambda\gamma^+\gamma_5)\Phi\right] = \frac{1}{2}\left[f_1 + S^i\epsilon^{ij}k^j\frac{1}{m}f_{1T}^\perp + \lambda\Lambda g_1 + \lambda S^i k^i\frac{1}{m}g_{1T}\right]$$

$$\frac{1}{2}\text{Tr}\left[(\gamma^+ - s^j i\sigma^{+j}\gamma_5)\Phi\right] = \frac{1}{2}\left[f_1 + S^i\epsilon^{ij}k^j\frac{1}{m}f_{1T}^\perp + s^i\epsilon^{ij}k^j\frac{1}{m}h_1^\perp + s^i S^i h_1\right. \\ \left.+ s^i(2k^i k^j - \mathbf{k}^2\delta^{ij})S^j\frac{1}{2m^2}h_{1T}^\perp + \Lambda s^i k^i\frac{1}{m}h_{1L}^\perp\right]$$

quark pol.

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

nucleon pol.

- each TMD describes a particular spin-momentum correlation
- functions in black survive integration over transverse momentum
- functions in green box are chirally odd
- functions in red are naive T-odd

Spin-momentum structure of the nucleon

$$\frac{1}{2} \text{Tr} \left[(\gamma^+ + \lambda \gamma^+ \gamma_5) \Phi \right] = \frac{1}{2} \left[f_1 + S^i \epsilon^{ij} k^j \frac{1}{m} f_{1T}^\perp + \lambda \Lambda g_1 + \lambda S^i k^i \frac{1}{m} g_{1T} \right]$$

$$\frac{1}{2} \text{Tr} \left[(\gamma^+ - s^j i \sigma^{+j} \gamma_5) \Phi \right] = \frac{1}{2} \left[f_1 + S^i \epsilon^{ij} k^j \frac{1}{m} f_{1T}^\perp + s^i \epsilon^{ij} k^j \frac{1}{m} h_1^\perp + s^i S^i h_1 \right. \\ \left. + s^i (2k^i k^j - k^2 \delta^{ij}) S^j \frac{1}{2m^2} h_{1T}^\perp + \Lambda s^i k^i \frac{1}{m} h_{1L}^\perp \right]$$

helicity

quark pol.

nucleon pol.

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

Boer-Mulders

- each TMD describes a particular spin-relation

- functions in black survive integration over transverse momentum

- functions in green box are chirally odd

Sivers

pretzelosity

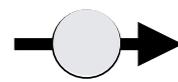
red are naive T-odd

transversity

worm-gear

TMDs - probabilistic interpretation

Proton goes out of the screen/ photon goes into the screen

  nucleon with transverse or longitudinal spin

  parton with transverse or longitudinal spin

 parton transverse momentum

$$f_1 = \text{[Diagram: circle with red dot]}$$

$$g_1 = \text{[Diagram: circle with black dot and red dot]} - \text{[Diagram: circle with black dot and red cross]}$$

$$h_1 = \text{[Diagram: circle with red dot and right arrow]} - \text{[Diagram: circle with red dot and left arrow]}$$

$$f_{1T}^\perp = \text{[Diagram: circle with blue down arrow and red dot]} - \text{[Diagram: circle with blue up arrow and red dot]}$$

$$h_1^\perp = \text{[Diagram: circle with blue down arrow, red dot, and red right arrow]} - \text{[Diagram: circle with blue up arrow, red dot, and red right arrow]}$$

$$g_{1T} = \text{[Diagram: circle with blue right arrow and red dot]} - \text{[Diagram: circle with blue left arrow and red dot]}$$

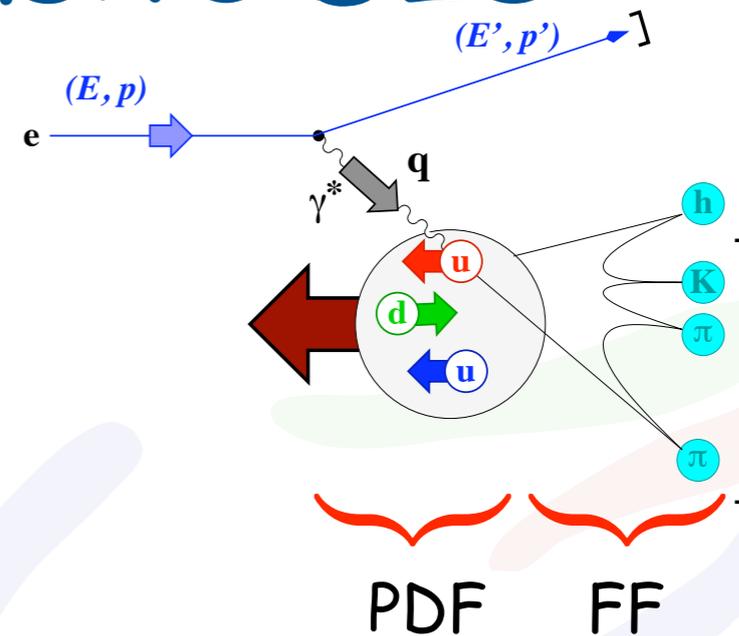
$$h_{1L}^\perp = \text{[Diagram: circle with black dot, red dot, and blue right arrow]} - \text{[Diagram: circle with black dot, red dot, and blue left arrow]}$$

$$h_{1T}^\perp = \text{[Diagram: circle with blue right arrow and red dot]} - \text{[Diagram: circle with blue left arrow and red dot]}$$

[courtesy of A. Bacchetta]

Probing TMDs in semi-inclusive DIS

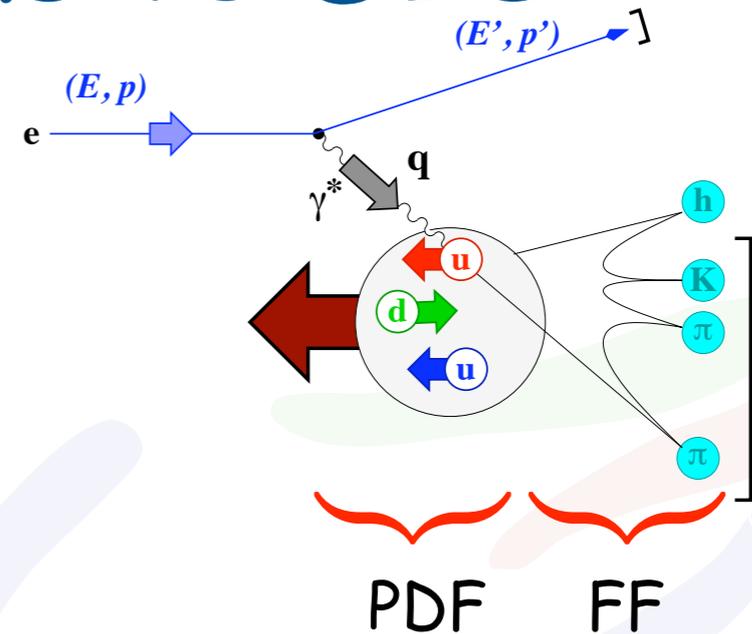
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nucleon pol.	U	f_1		h_1^\perp
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in SIDIS*) couple PDFs to:

*) semi-inclusive DIS with unpolarized final state

Probing TMDs in semi-inclusive DIS



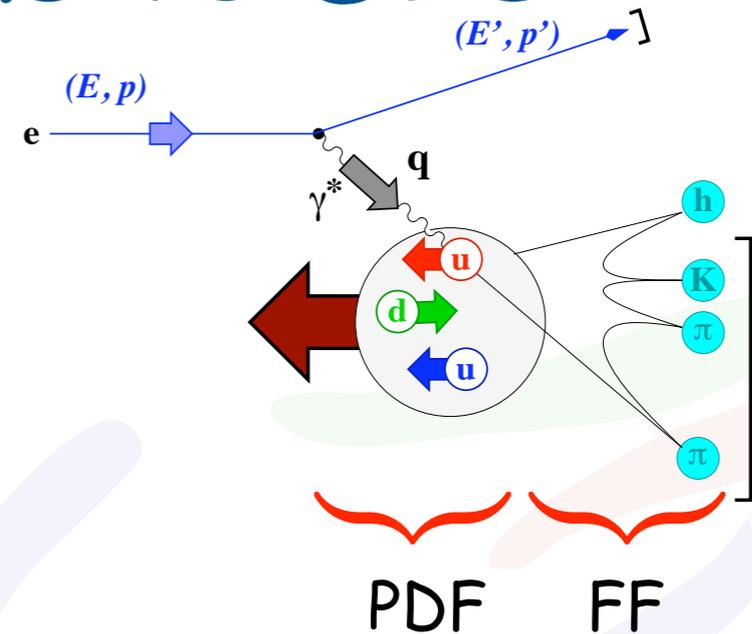
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Collins FF: $H_1^\perp, q \rightarrow h$

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Probing TMDs in semi-inclusive DIS



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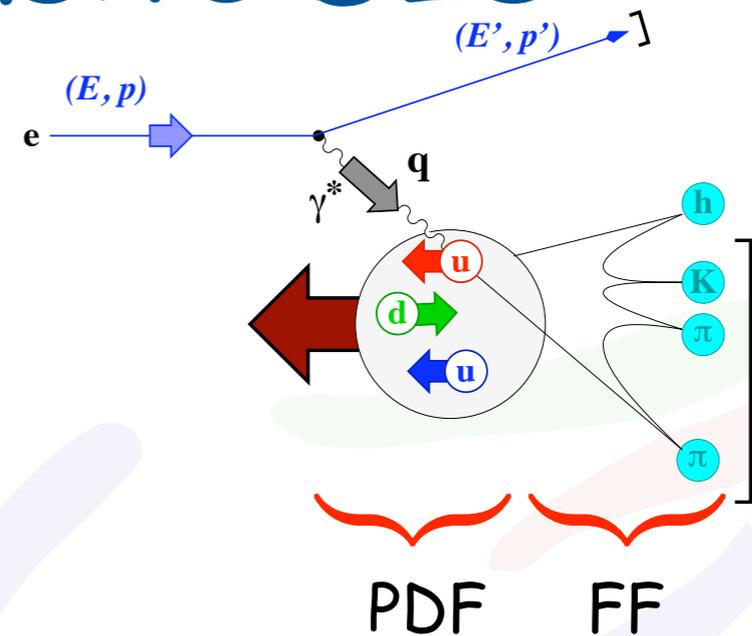
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ordinary FF: $D_1^{q \rightarrow h}$

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⇒ gives rise to characteristic azimuthal dependences

*) semi-inclusive DIS with unpolarized final state

1-Hadron production ($ep \rightarrow ehX$)

$$d\sigma = d\sigma_{UU}^0 + \cos 2\phi d\sigma_{UU}^1 + \frac{1}{Q} \cos \phi d\sigma_{UU}^2 + \lambda_e \frac{1}{Q} \sin \phi d\sigma_{LU}^3$$

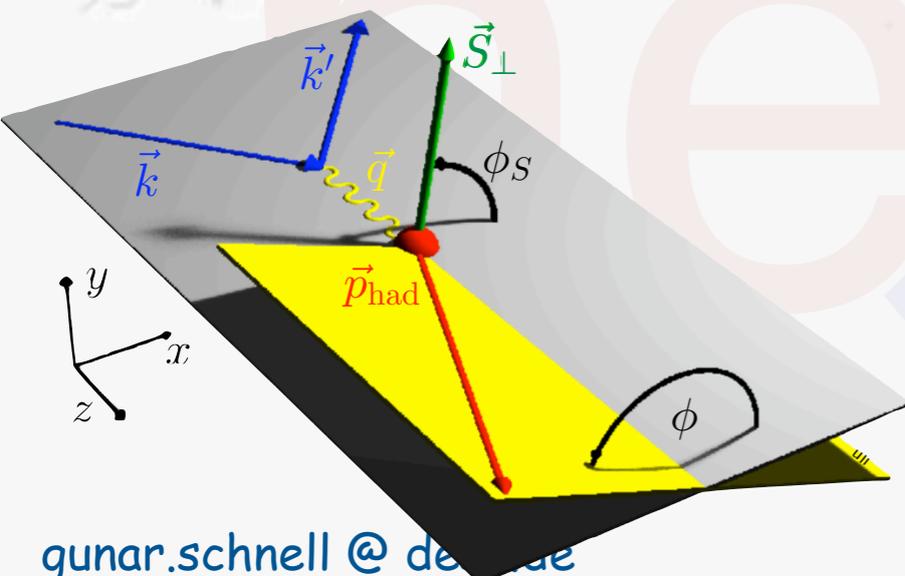
$$+ S_L \left\{ \sin 2\phi d\sigma_{UL}^4 + \frac{1}{Q} \sin \phi d\sigma_{UL}^5 + \lambda_e \left[d\sigma_{LL}^6 + \frac{1}{Q} \cos \phi d\sigma_{LL}^7 \right] \right\}$$

$$+ S_T \left\{ \sin(\phi - \phi_S) d\sigma_{UT}^8 + \sin(\phi + \phi_S) d\sigma_{UT}^9 + \sin(3\phi - \phi_S) d\sigma_{UT}^{10} \frac{1}{Q} \right.$$

$$\left. + \frac{1}{Q} (\sin(2\phi - \phi_S) d\sigma_{UT}^{11} + \sin \phi_S d\sigma_{UT}^{12}) \right.$$

$$\left. + \lambda_e \left[\cos(\phi - \phi_S) d\sigma_{LT}^{13} + \frac{1}{Q} (\cos \phi_S d\sigma_{LT}^{14} + \cos(2\phi - \phi_S) d\sigma_{LT}^{15}) \right] \right\}$$

σ_{XY}
 ↙ ↘
Beam Target
Polarization



Mulders and Tangermann, Nucl. Phys. B 461 (1996) 197

Boer and Mulders, Phys. Rev. D 57 (1998) 5780

Bacchetta et al., Phys. Lett. B 595 (2004) 309

Bacchetta et al., JHEP 0702 (2007) 093

"Trento Conventions", Phys. Rev. D 70 (2004) 117504

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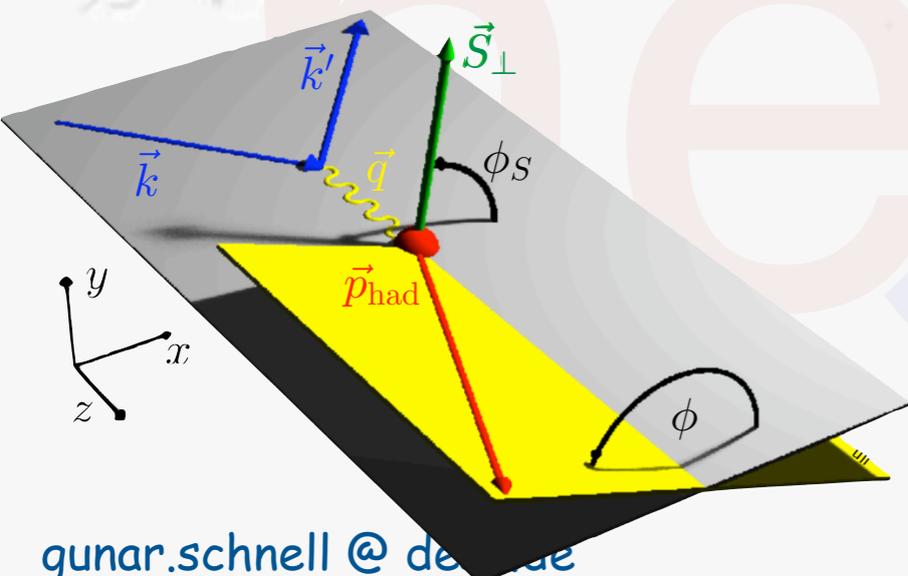
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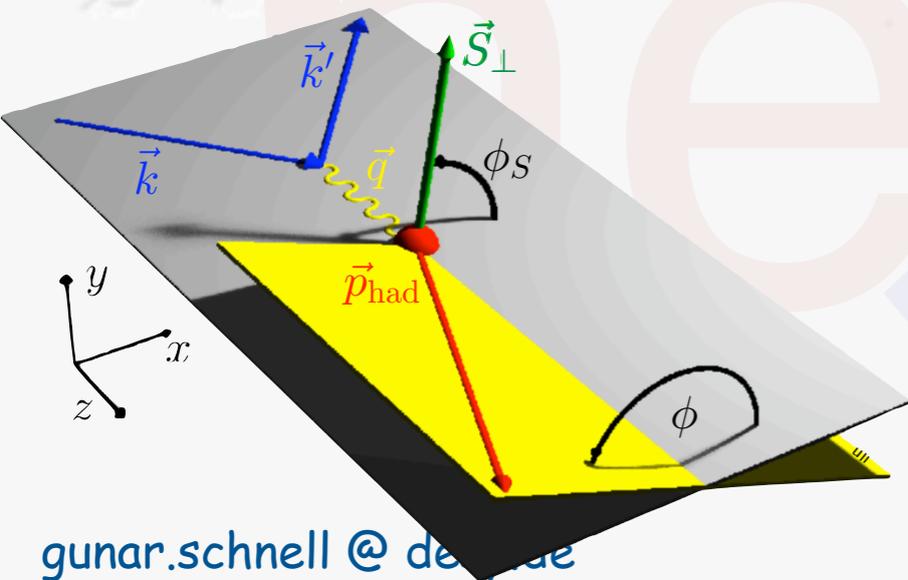
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Hadron multiplicities in DIS

$$\frac{d^5\sigma}{dx dy dz d\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \{F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1-\epsilon)} F_{UU}^{\cos\phi_h} \cos\phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h\}$$

$$F_{XY,Z} = F_{XY,Z}(x, y, z, P_{h\perp})$$

target polarization \downarrow
 \uparrow beam polarization \uparrow virtual-photon polarization

[see, e.g., Bacchetta et al., JHEP 0702 (2007) 093]

$$\gamma = \frac{2Mx}{Q}$$

$$\epsilon = \frac{1 - y - \frac{1}{4}\gamma^2 y^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2 y^2}$$

Hadron multiplicities in DIS

hadron multiplicity:
normalize to inclusive DIS
cross section

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Hadron multiplicities in DIS

hadron multiplicity:
normalize to inclusive DIS
cross section

$$\frac{d^2 \sigma^{\text{incl. DIS}}}{dxdy} \propto F_T + \epsilon F_L$$

$$\frac{d^4 \mathcal{M}^h(x, y, z, P_{h\perp}^2)}{dxdydzdP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \frac{F_{UU,T} + \epsilon F_{UU,L}}{F_T + \epsilon F_L}$$

$$\approx \frac{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^{q \rightarrow h}(z, K_T^2)}{\sum_q e_q^2 f_1^q(x)}$$

$$\frac{d^5 \sigma}{dxdydzd\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1-\epsilon)} F_{UU}^{\cos \phi_h} \cos \phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h \right\}$$

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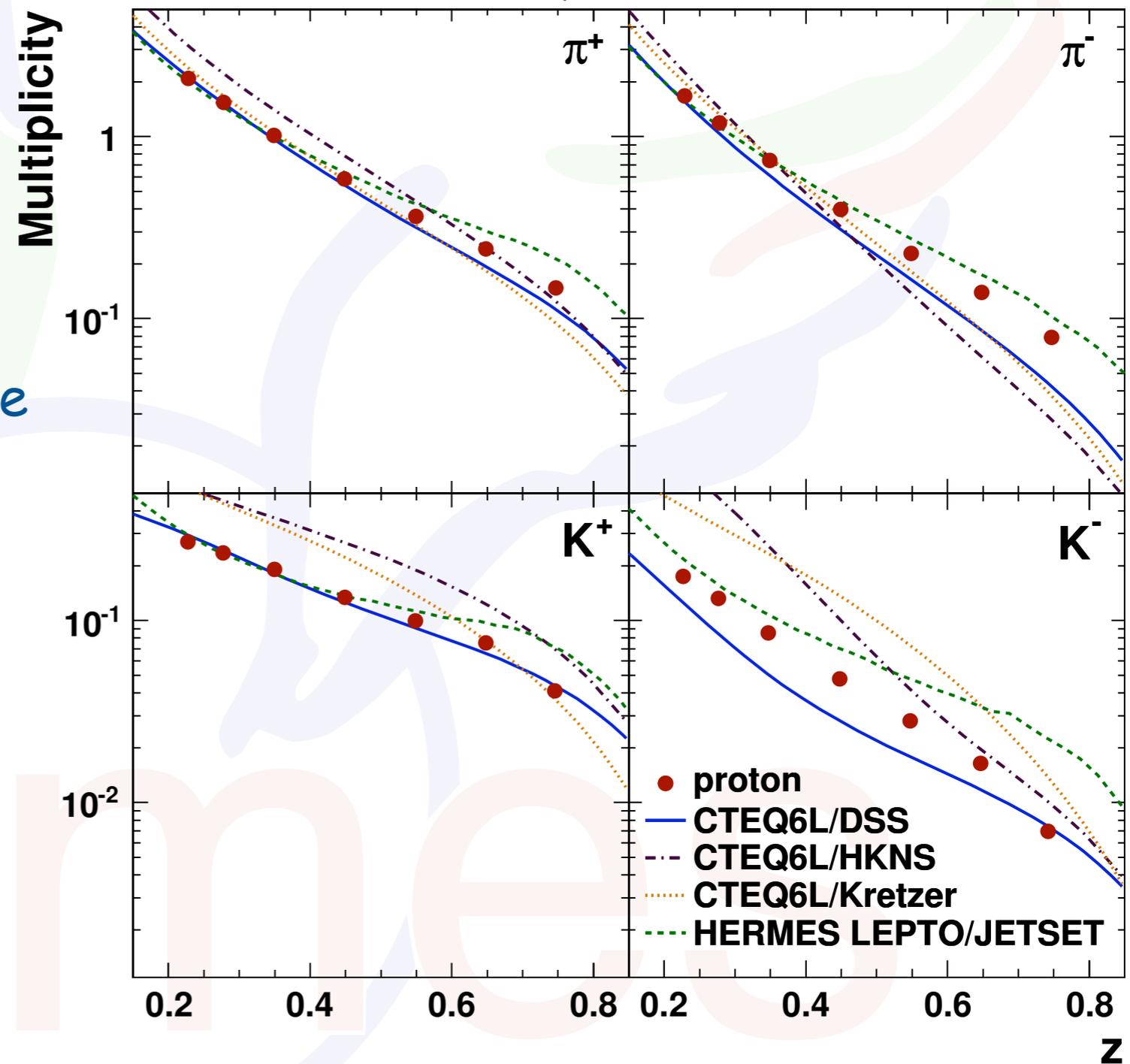
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Multiplicities @ HERMES

- extensive data set on pure proton and deuteron targets for identified charged mesons
- extracted in a multi-dimensional unfolding procedure
- fair agreement between DSS and positive mesons
- poor description of negative mesons
- p/d differences due to flavor dependence of fragmentation

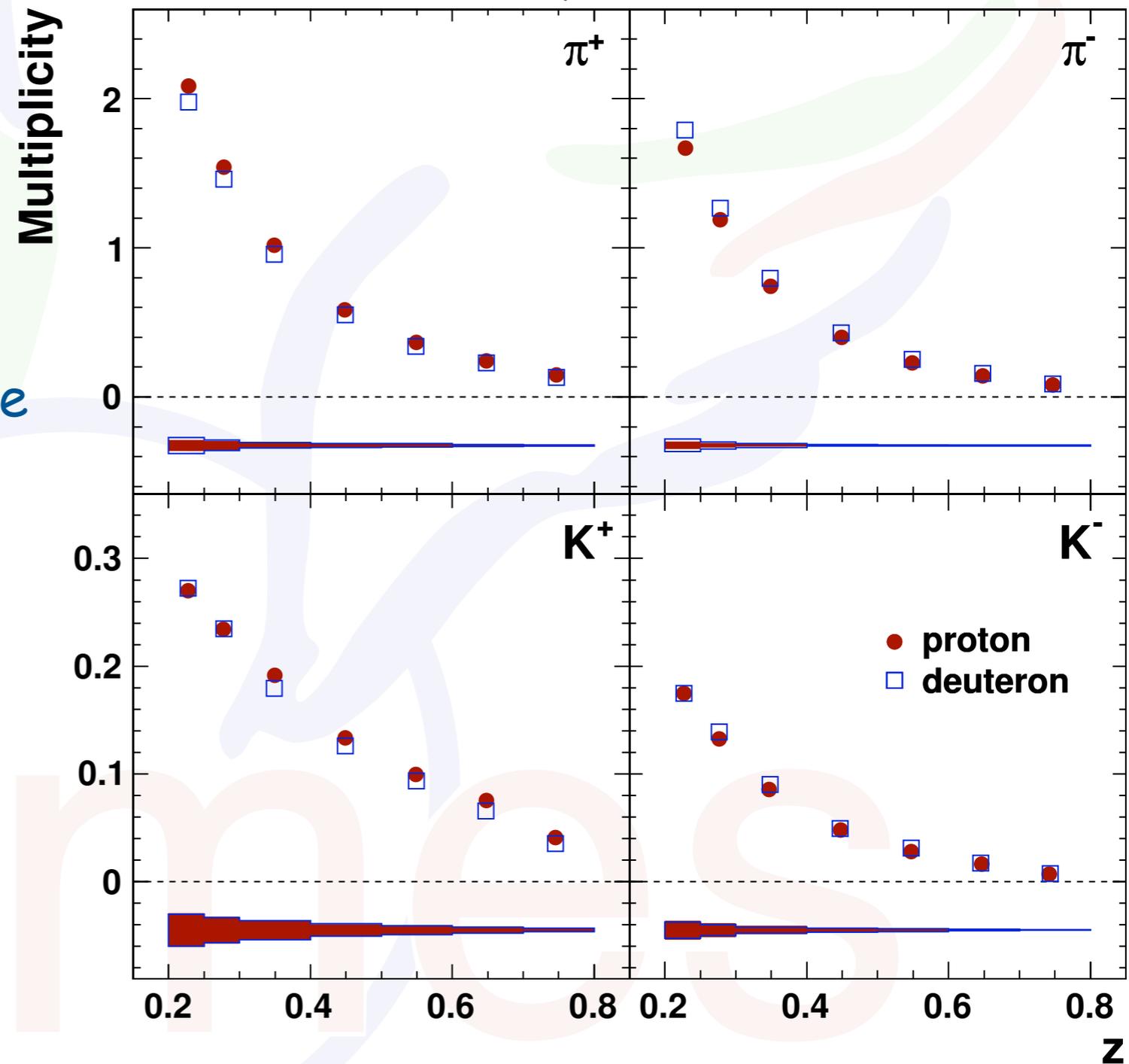
[Airapetian et al., arXiv:1212.5407]



Multiplicities @ HERMES

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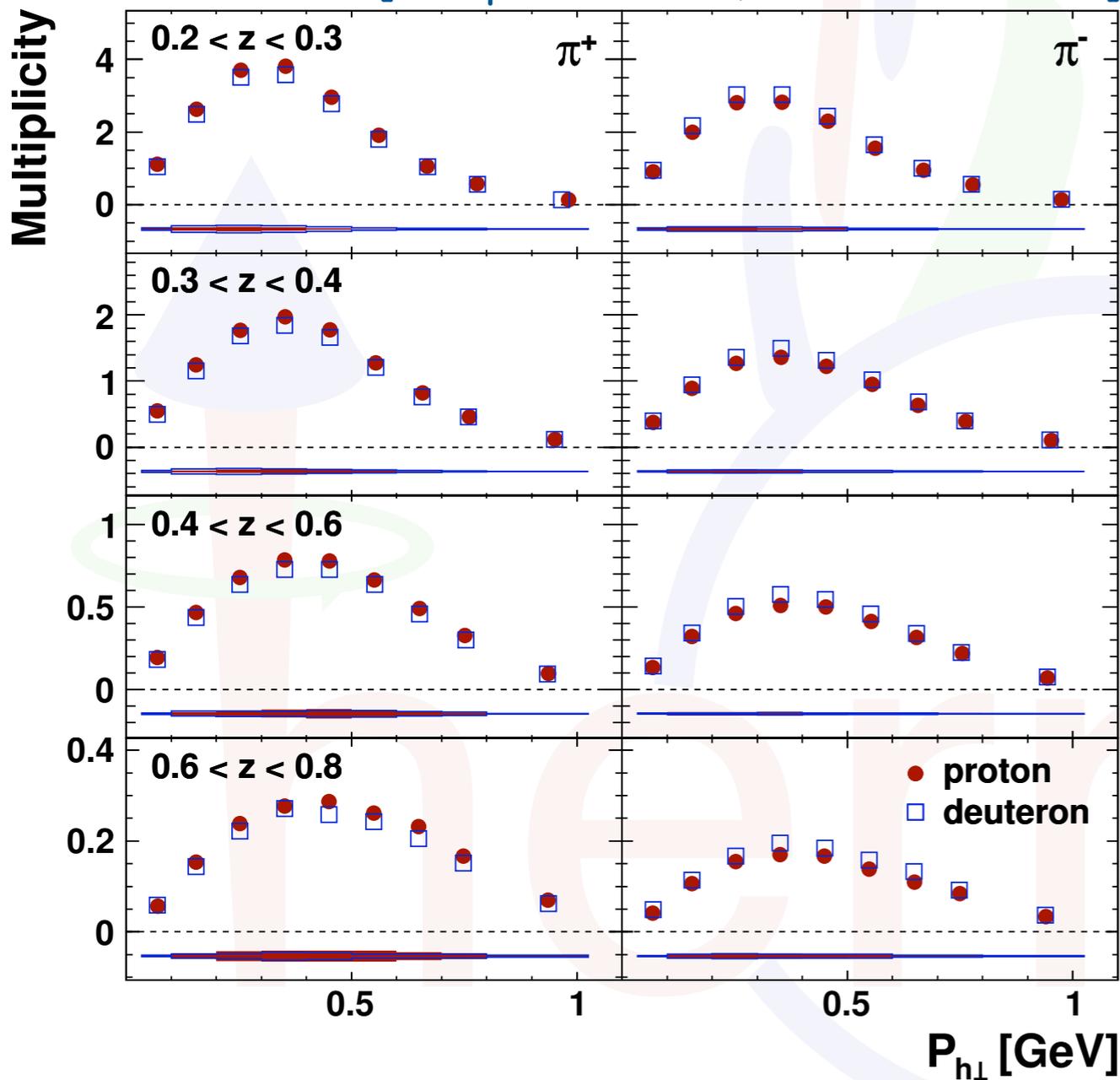
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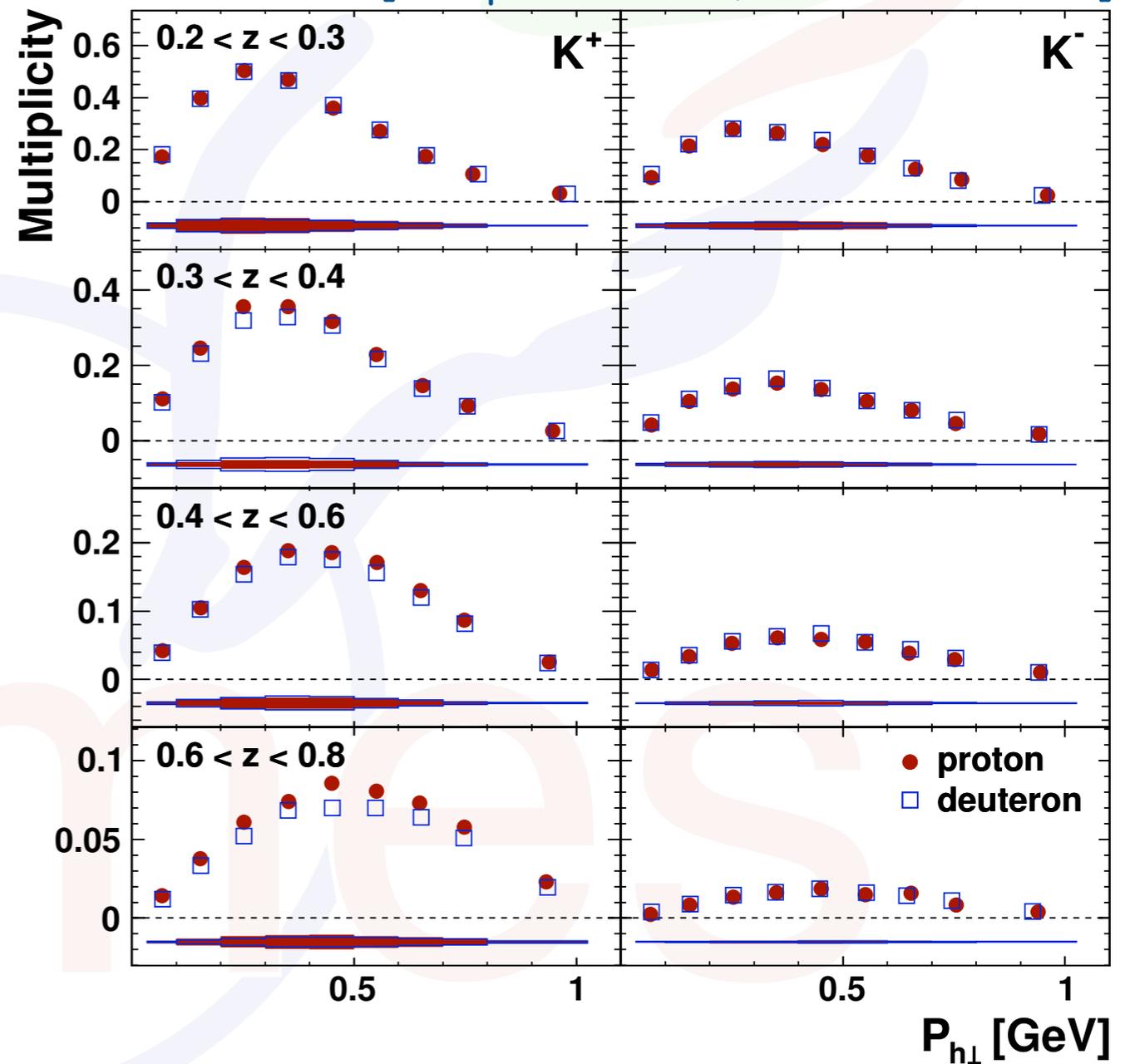
Multiplicities @ HERMES

multi-dimensional analysis allows exploration of new kinematic dependences

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Transversely polarized quarks?

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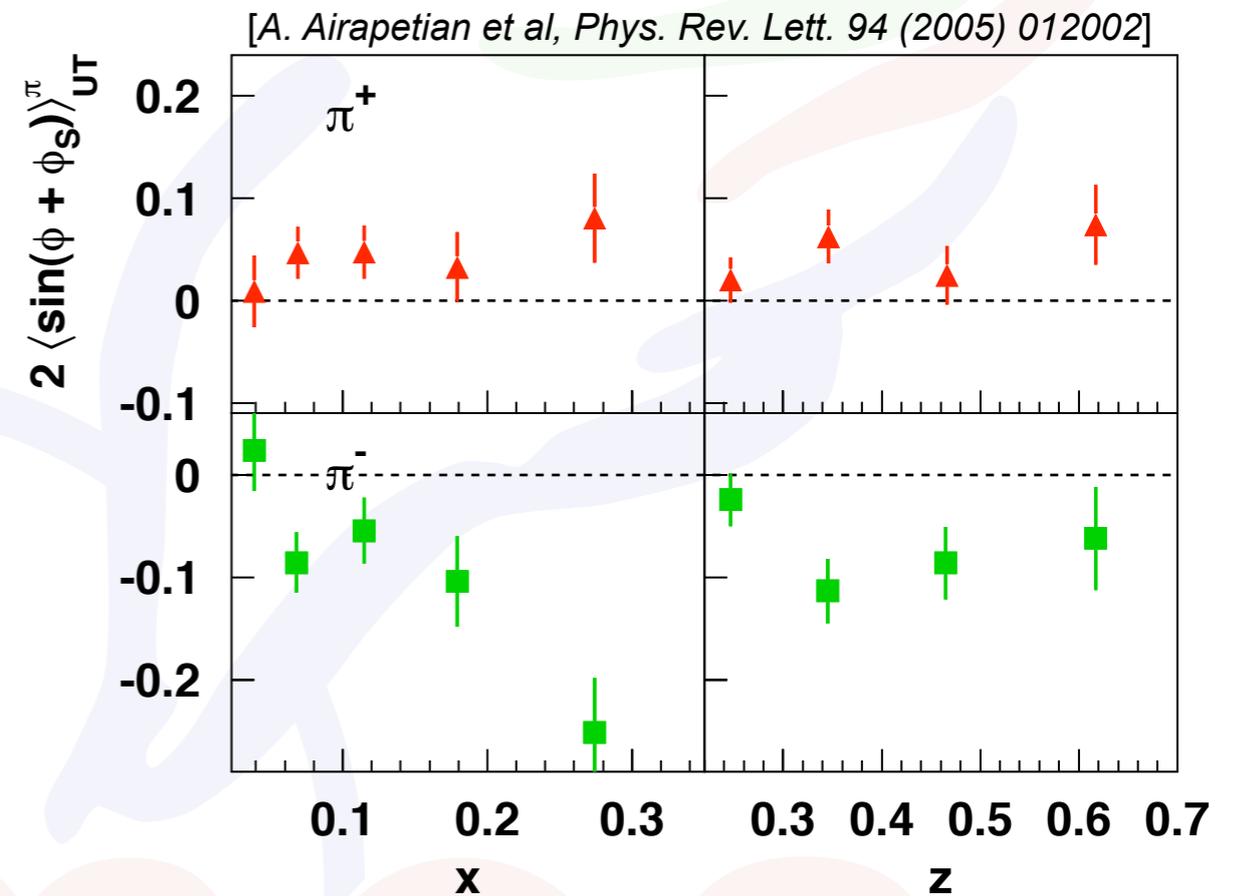


hermes

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- transverse polarization of quarks leads to large effects!



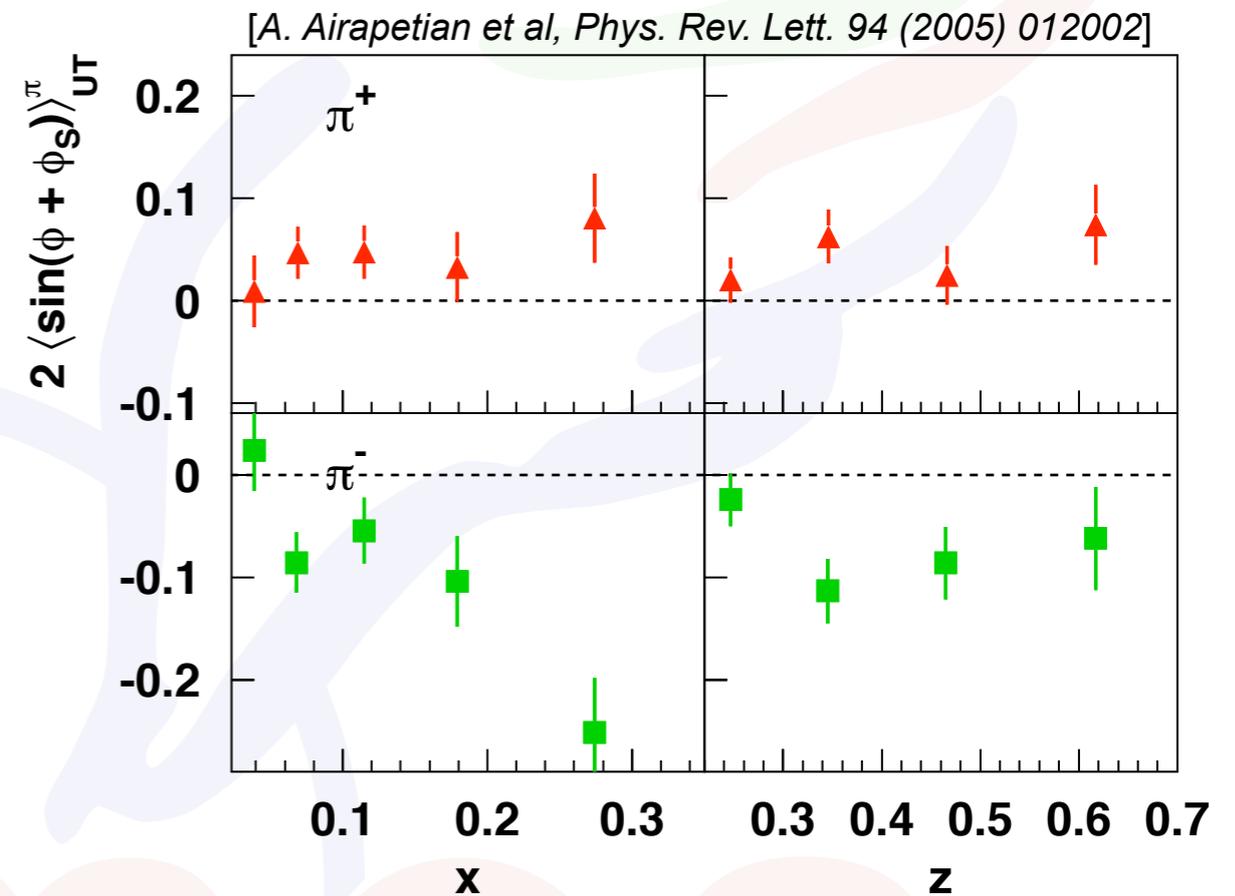
2005: First evidence from HERMES
SIDIS on proton

Non-zero transversity
Non-zero Collins function

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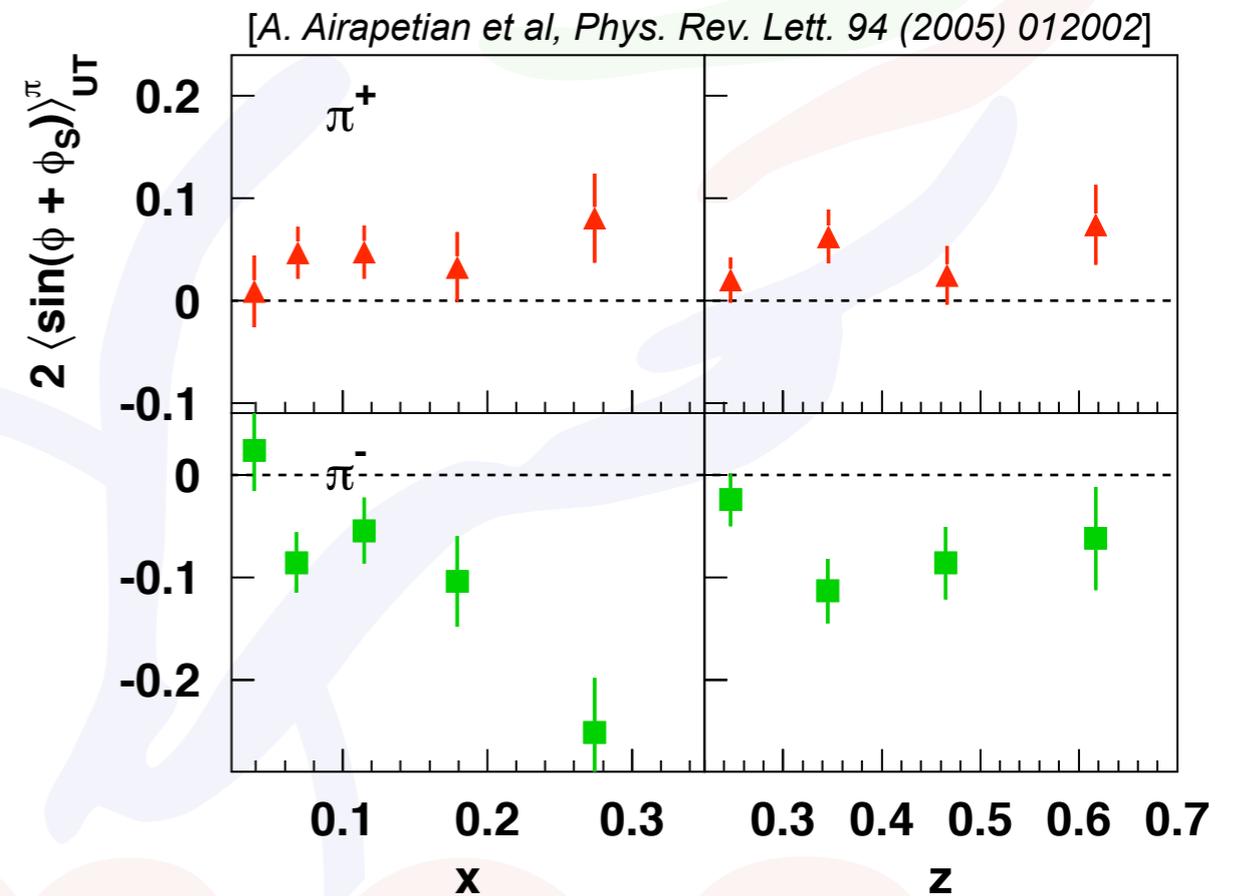
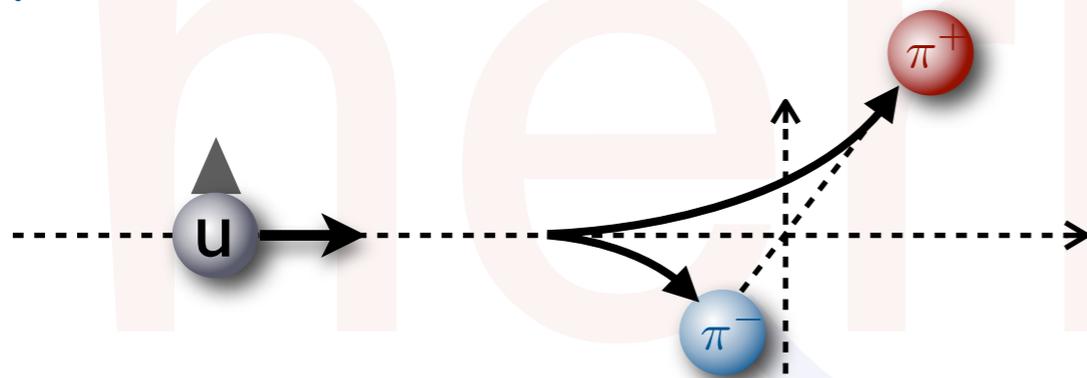
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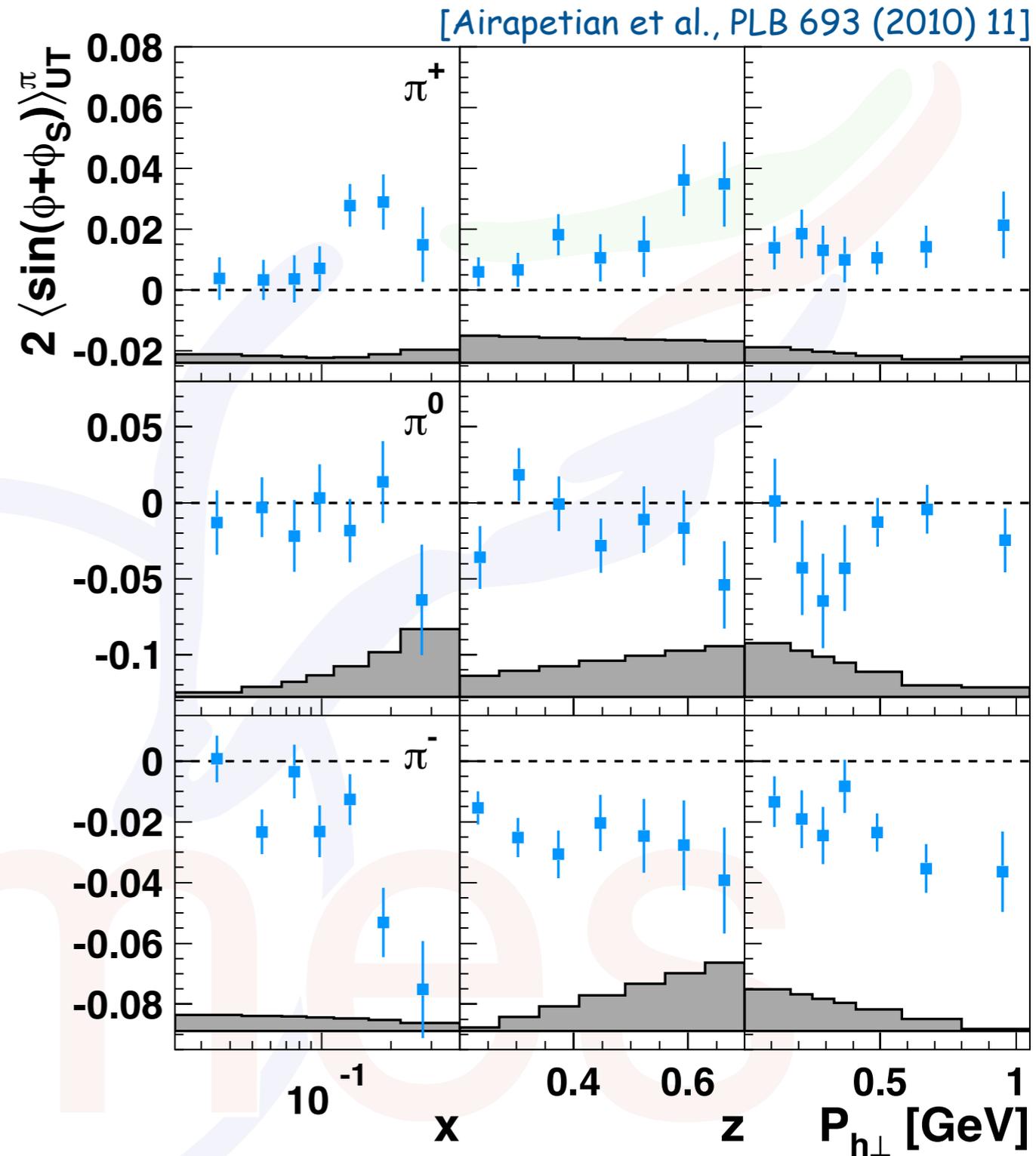
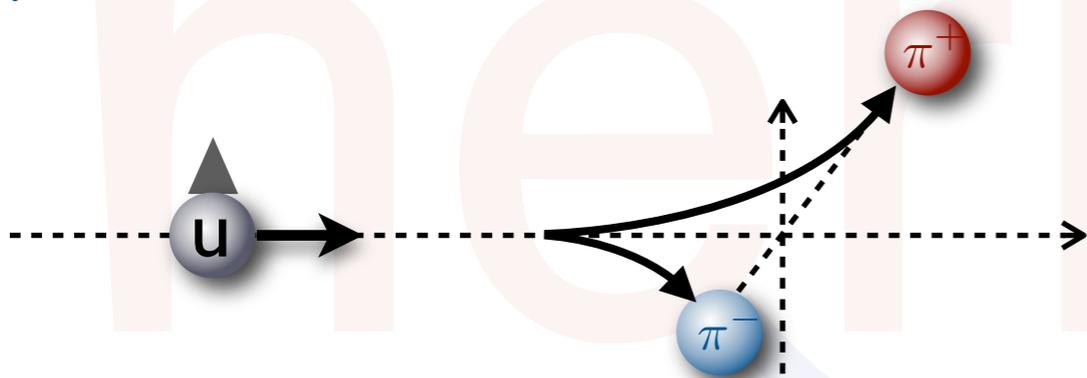
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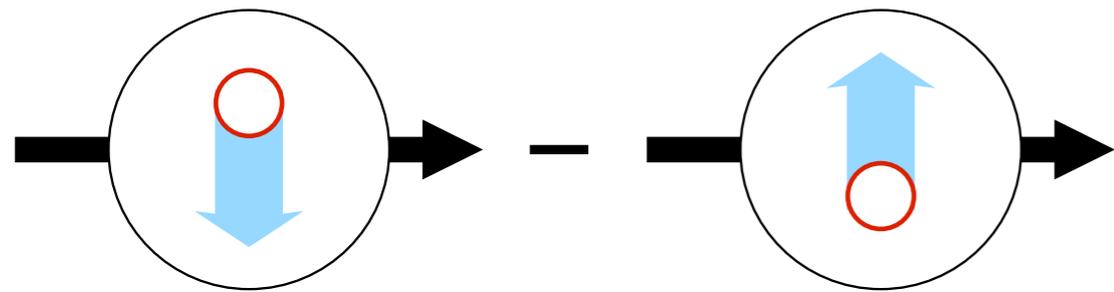
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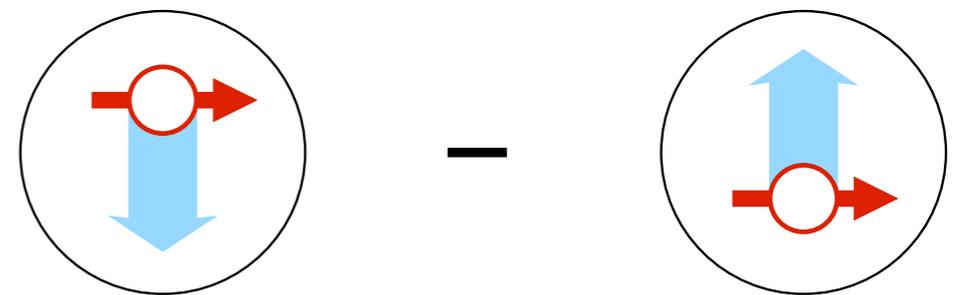




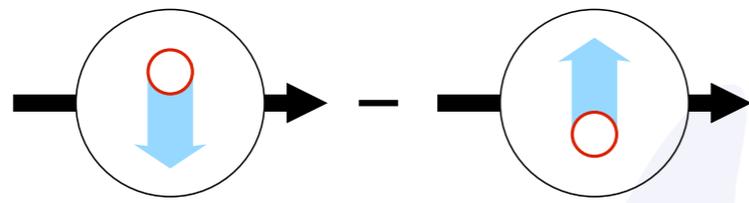
Sivers effect

naively T-odd distributions
 "Wilson-line physics"

Boer-Mulders effect



$f_{1T}^\perp =$



Sivers function



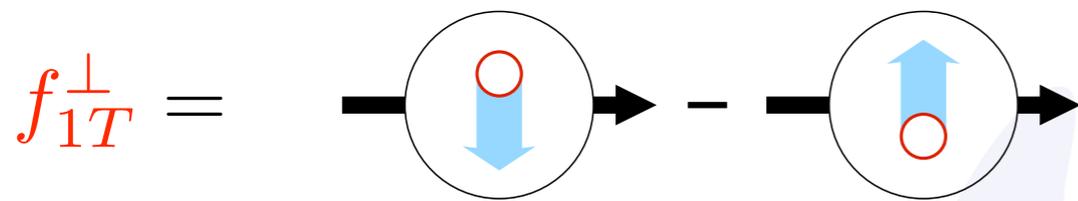
hermes

$$f_{1T}^\perp = \text{---} \left(\begin{array}{c} \circ \\ \downarrow \\ \circ \end{array} \right) \rightarrow \text{---} \left(\begin{array}{c} \circ \\ \uparrow \\ \circ \end{array} \right) \rightarrow$$

Sivers function

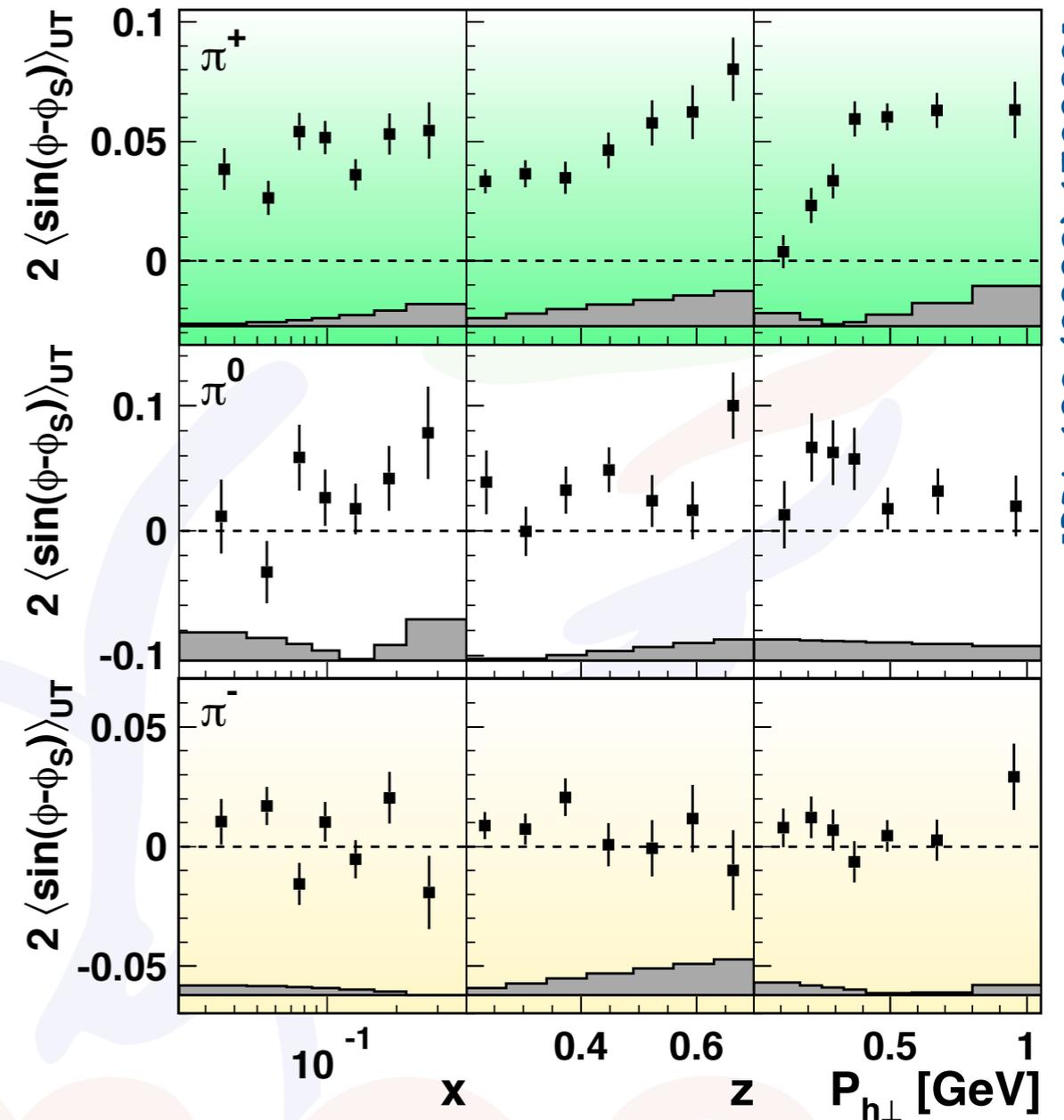
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- candidate for large (30-50%) asymmetries in $p^\uparrow p \rightarrow hX$



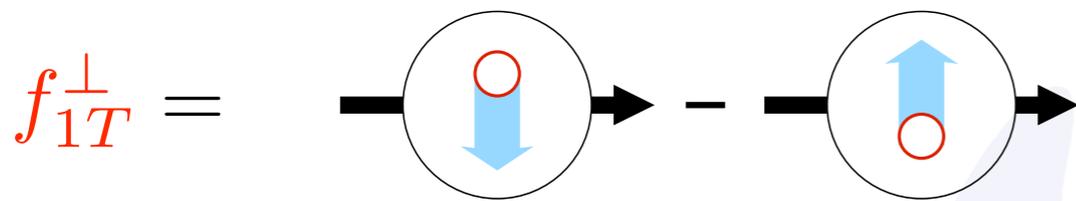


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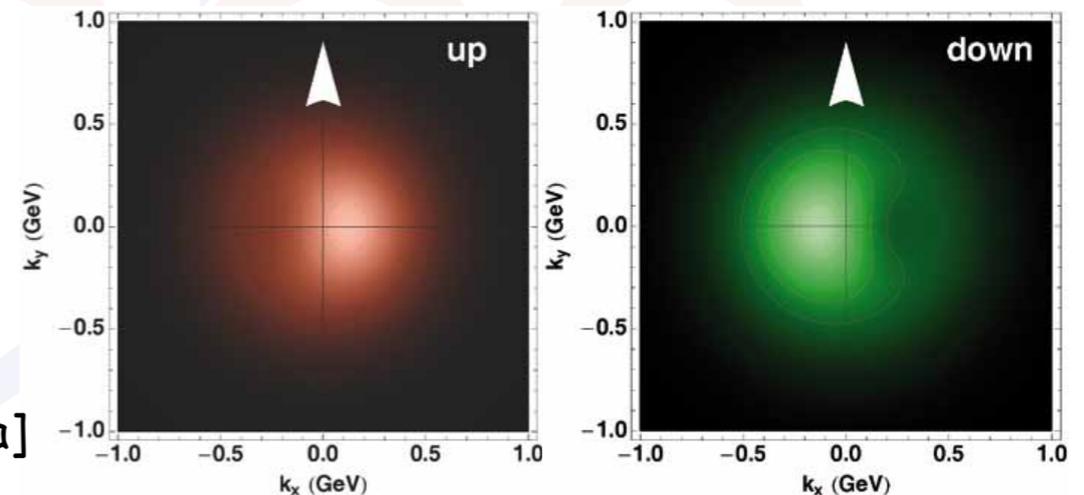
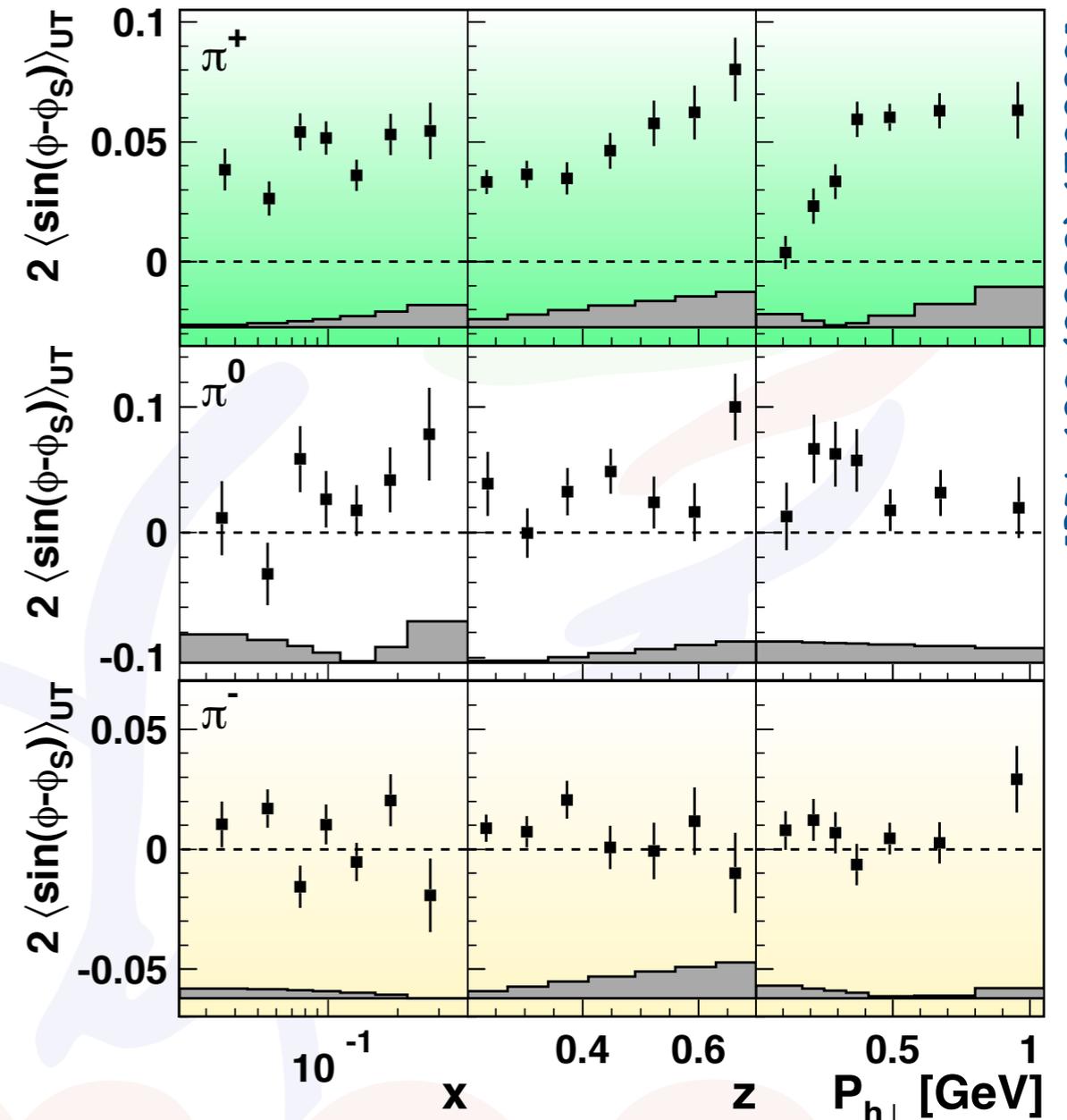


[PRL 103(2009) 152002]

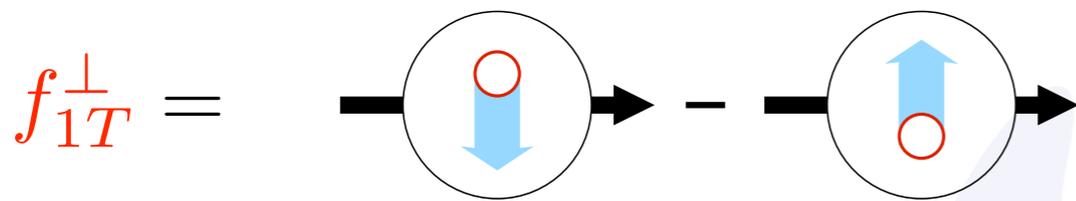


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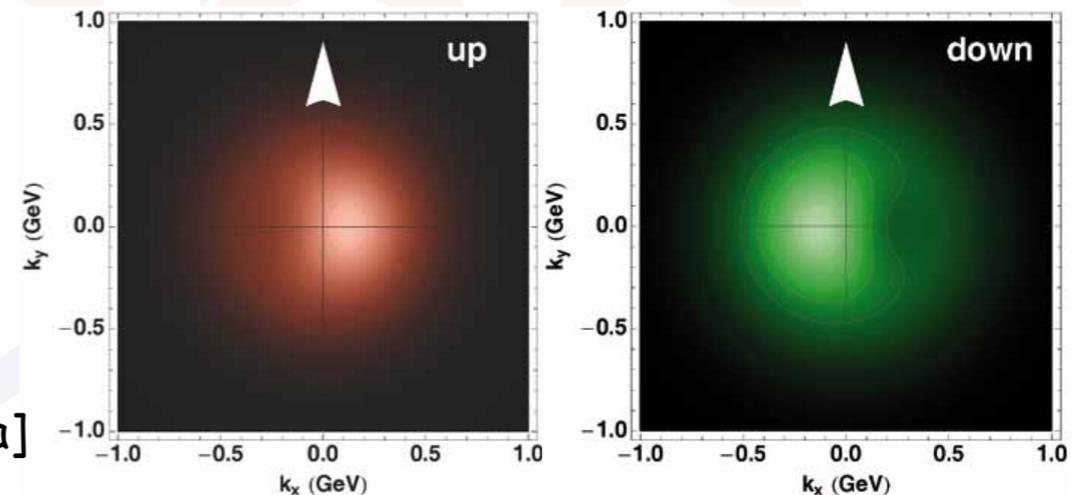
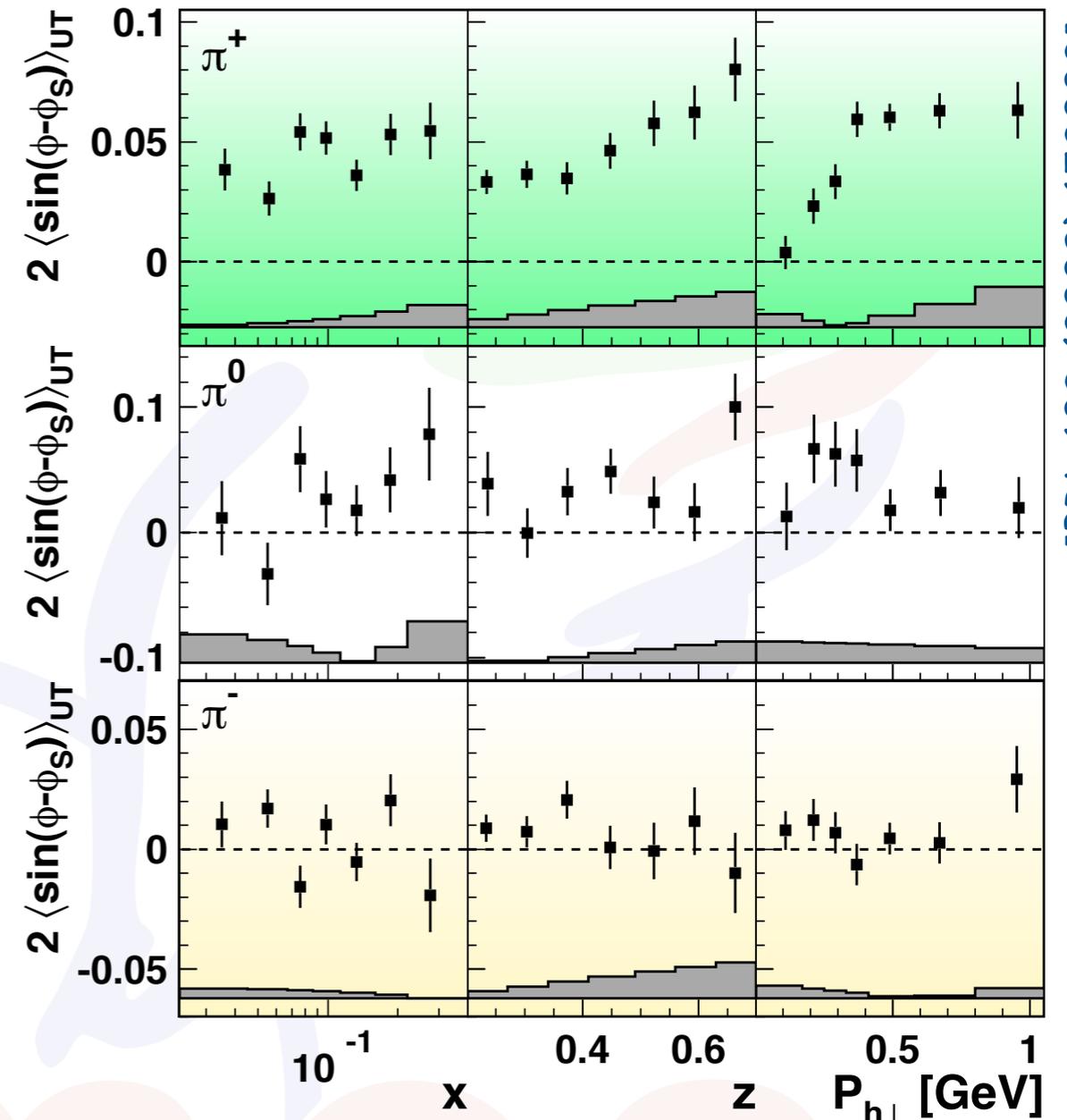


[courtesy of A. Bacchetta]



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- leads to peculiar calculable universality breaking (DIS vs. Drell-Yan)

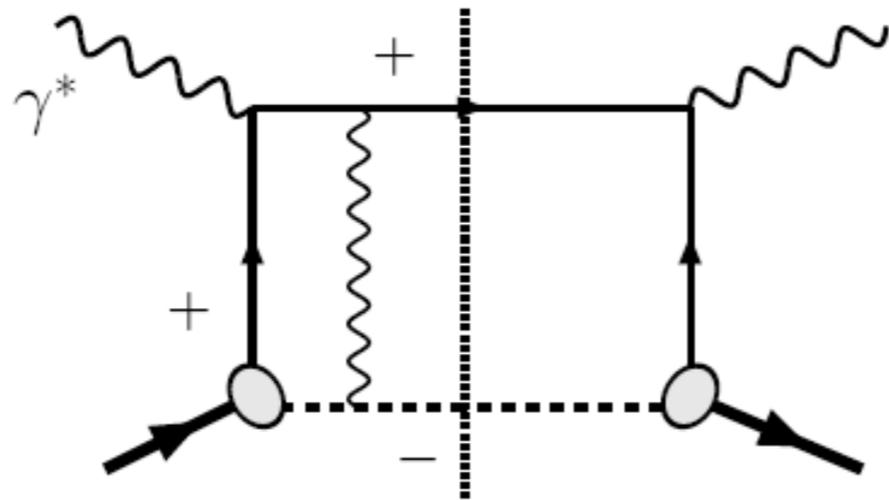
Sivers function



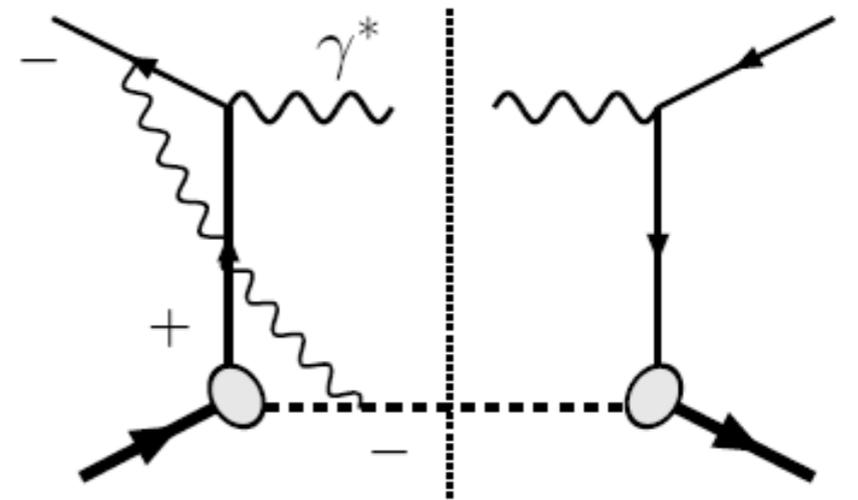
[courtesy of A. Bacchetta]

Process dependence

simple QED
example



DIS: attractive



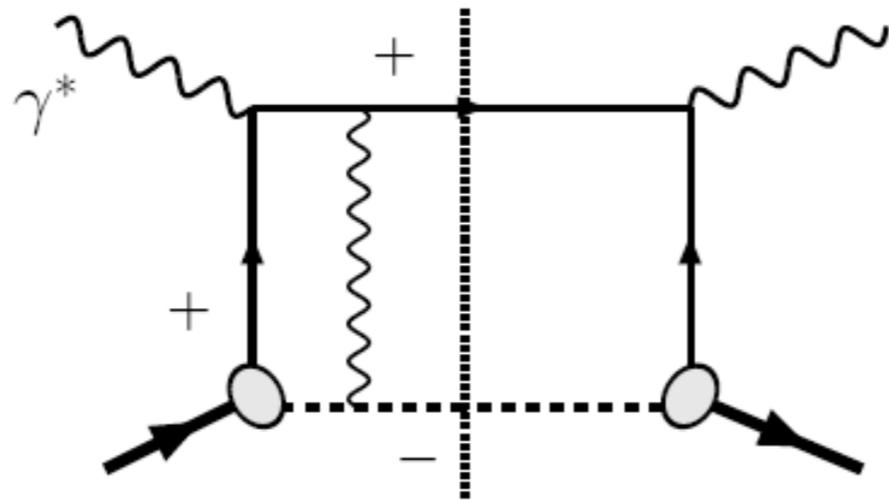
Drell-Yan: repulsive



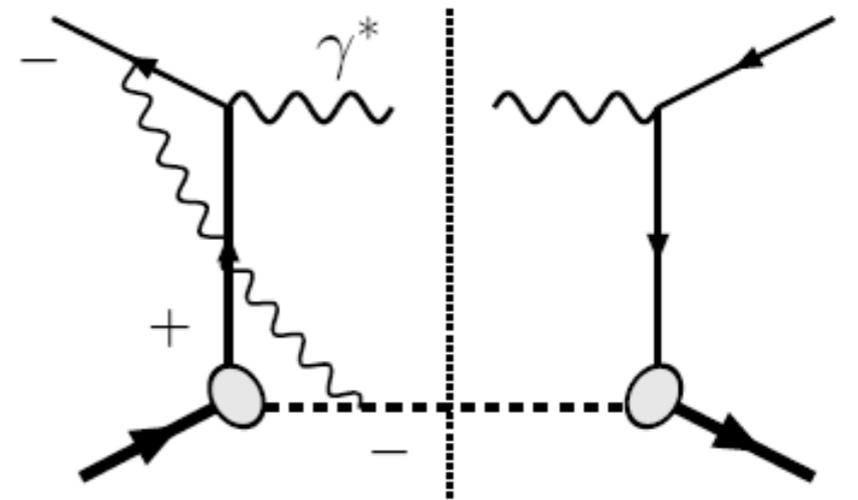
hermes

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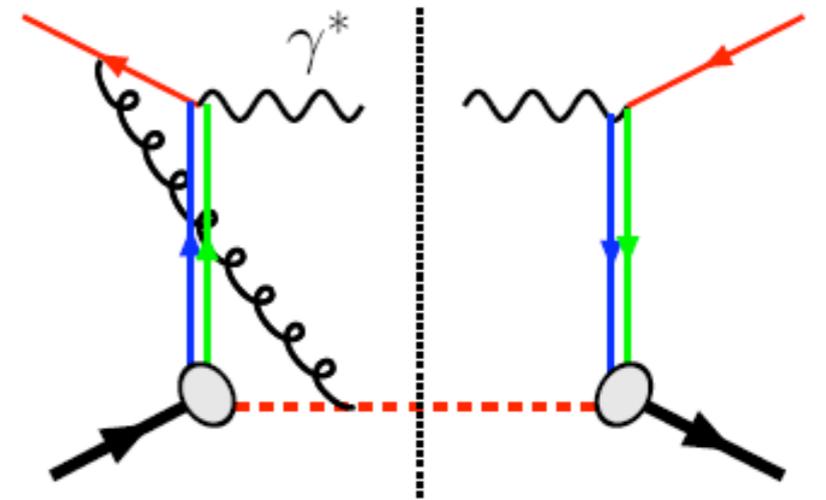
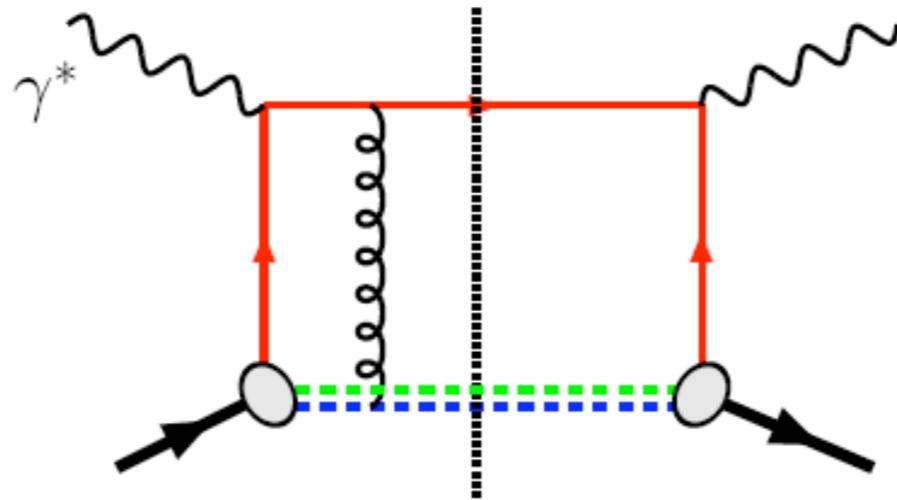


DIS: attractive



Drell-Yan: repulsive

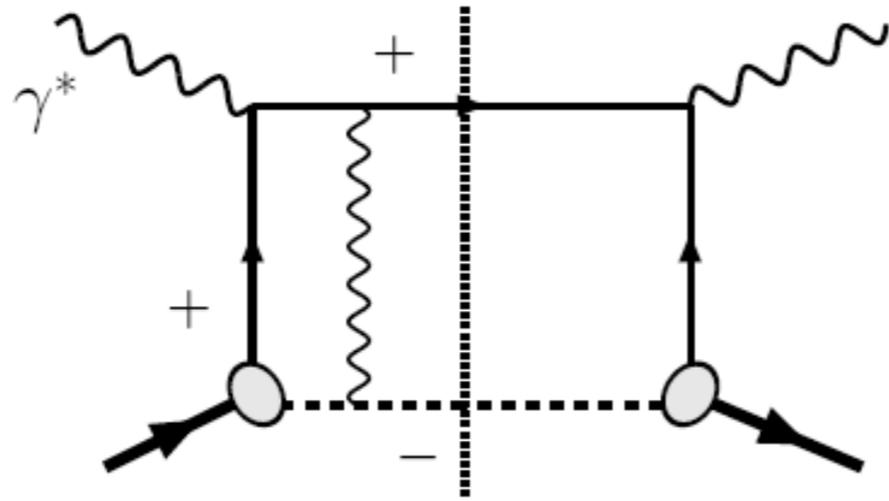
add color:
QCD



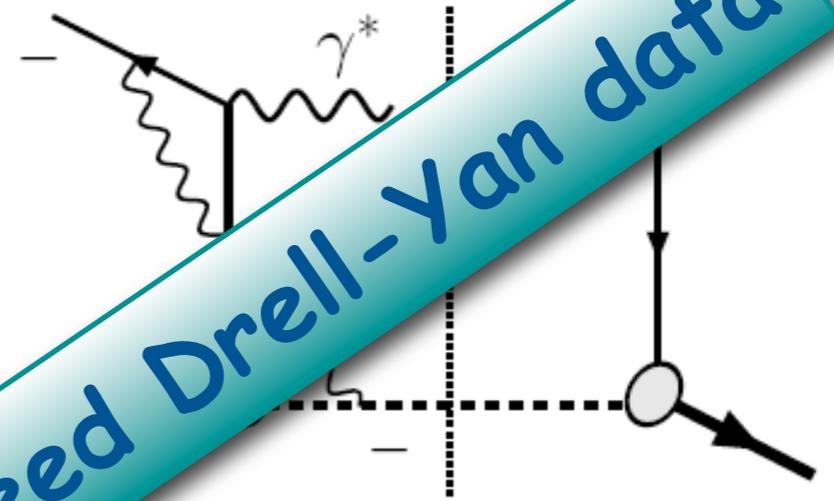
result: $Sivers|_{DIS} = - Sivers|_{DY}$

Process dependence

simple QED
example

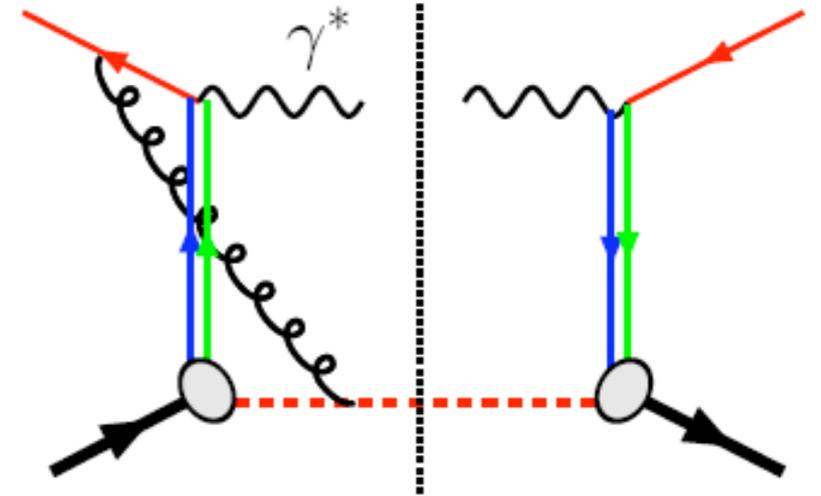
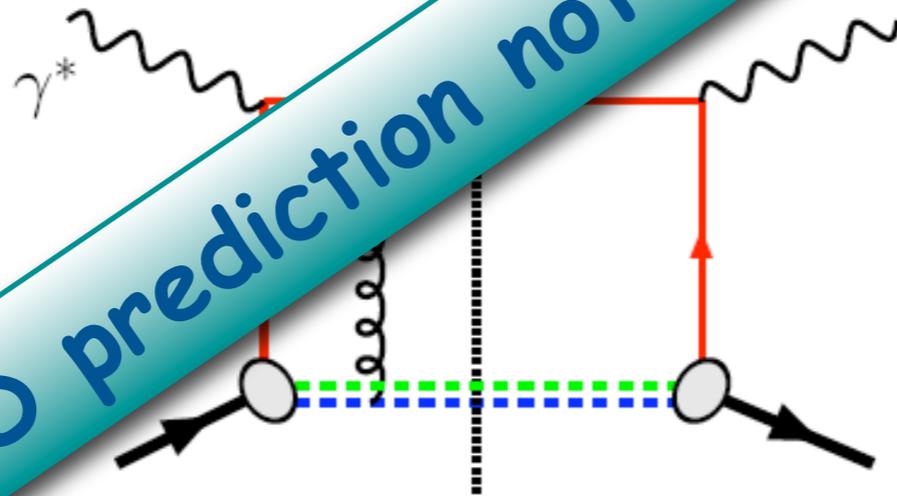


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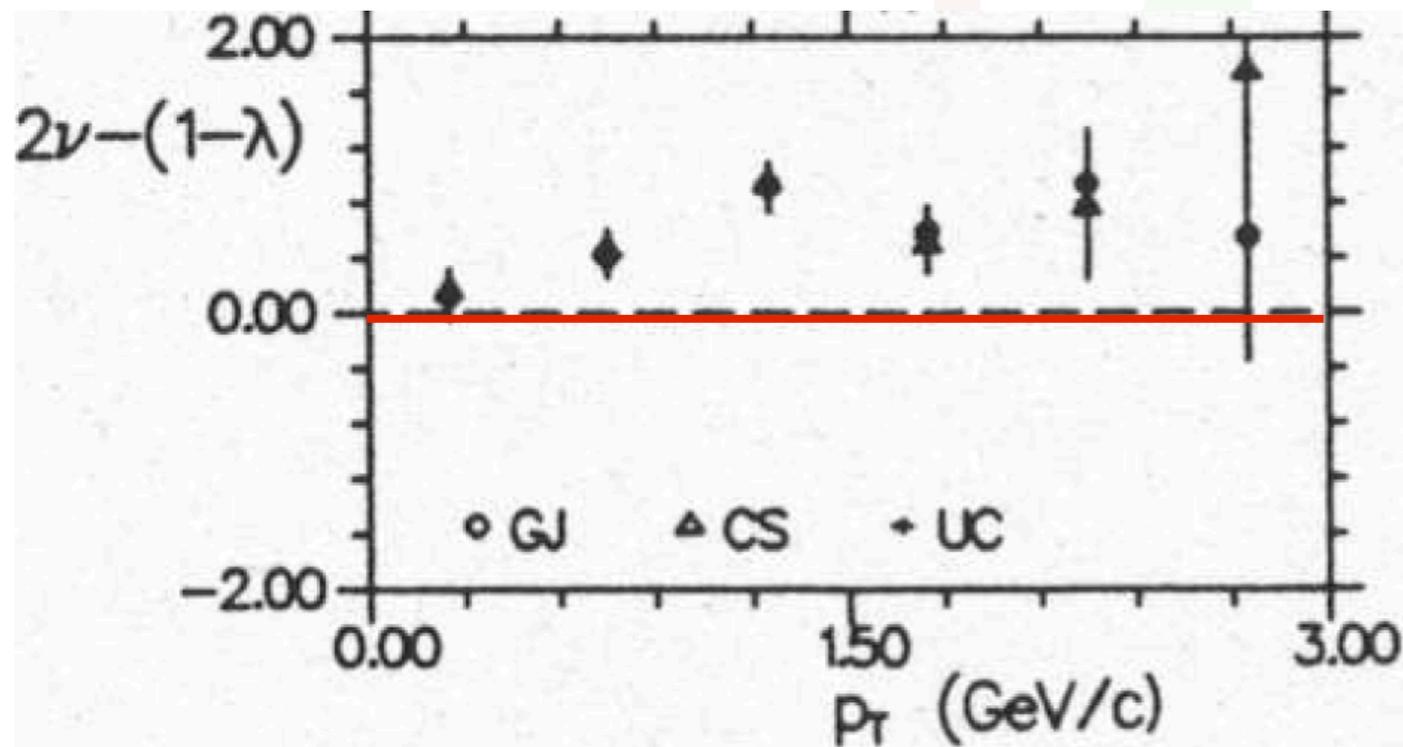
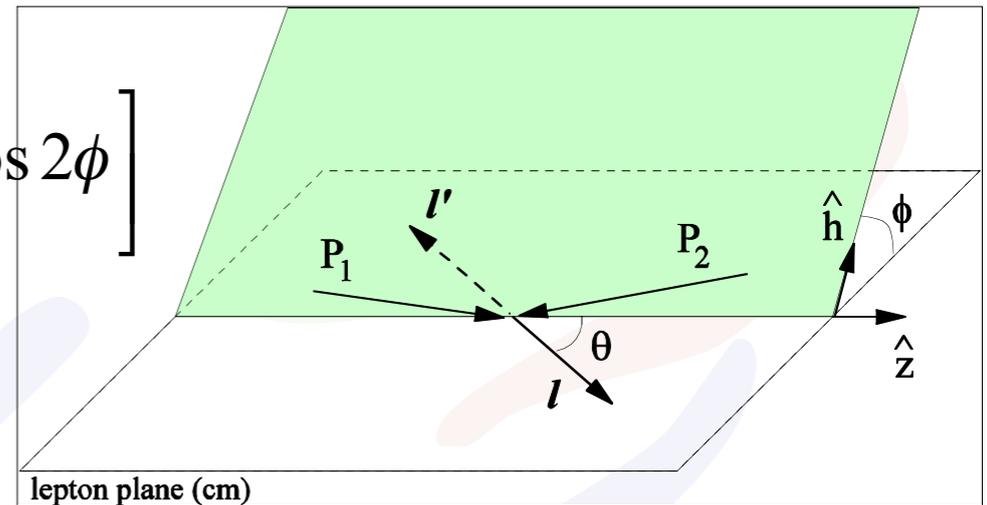
add color:
QCD



result: $Sivers|_{DIS} = -Sivers|_{DY}$

Unpolarized Drell-Yan

$$\left(\frac{1}{\sigma}\right)\left(\frac{d\sigma}{d\Omega}\right) = \left[\frac{3}{4\pi}\right] \left[1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi\right]$$

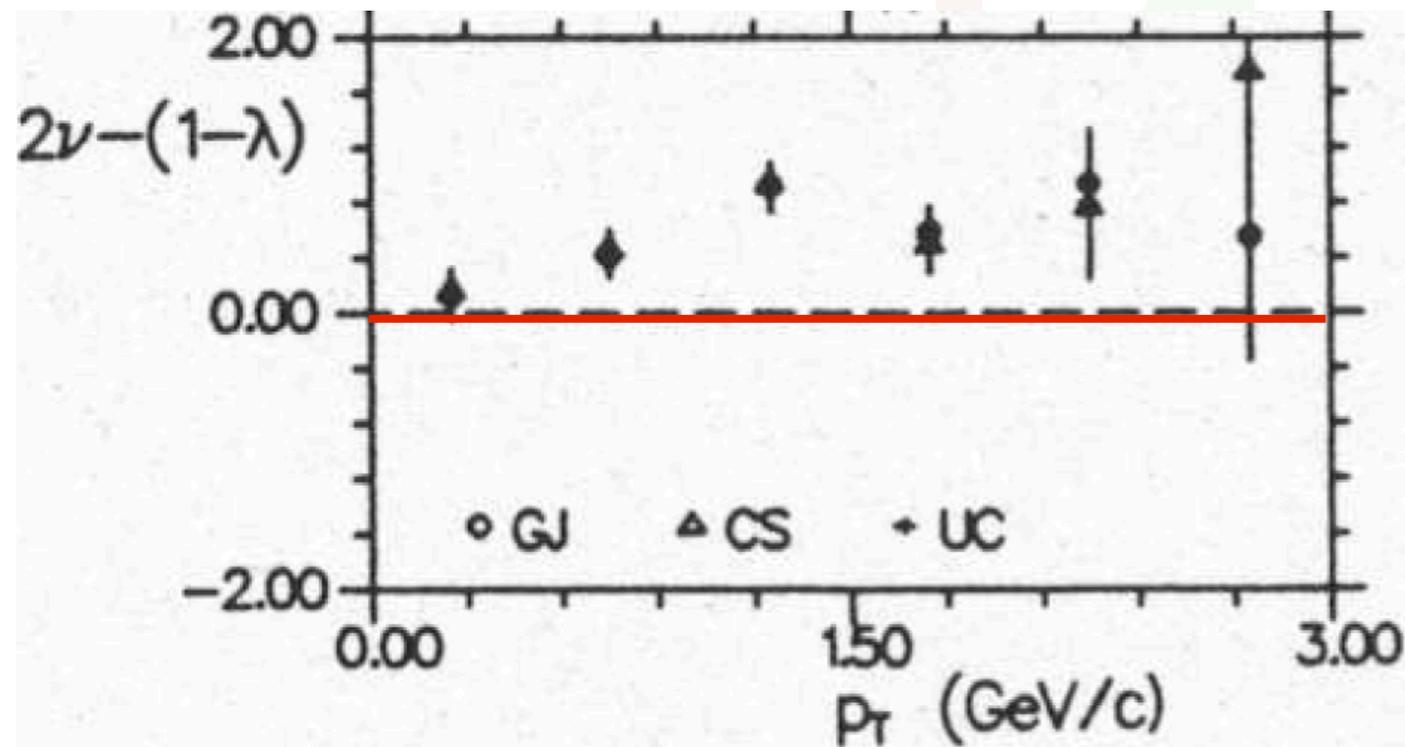
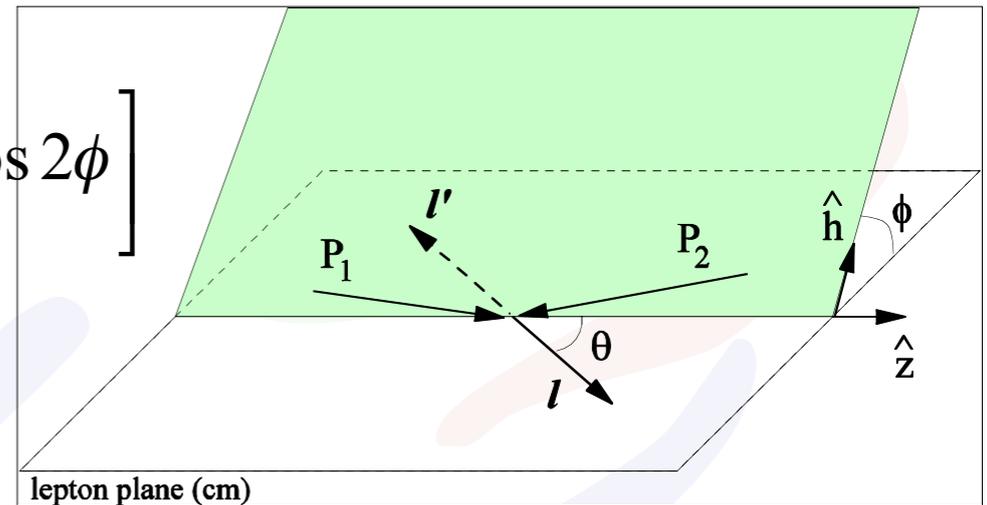


$$1 - \lambda - 2\nu = 0$$

Large deviations from Lam-Tung relation observed in DY
[NA10 ('86/'88) & E615 ('89)]

Unpolarized Drell-Yan

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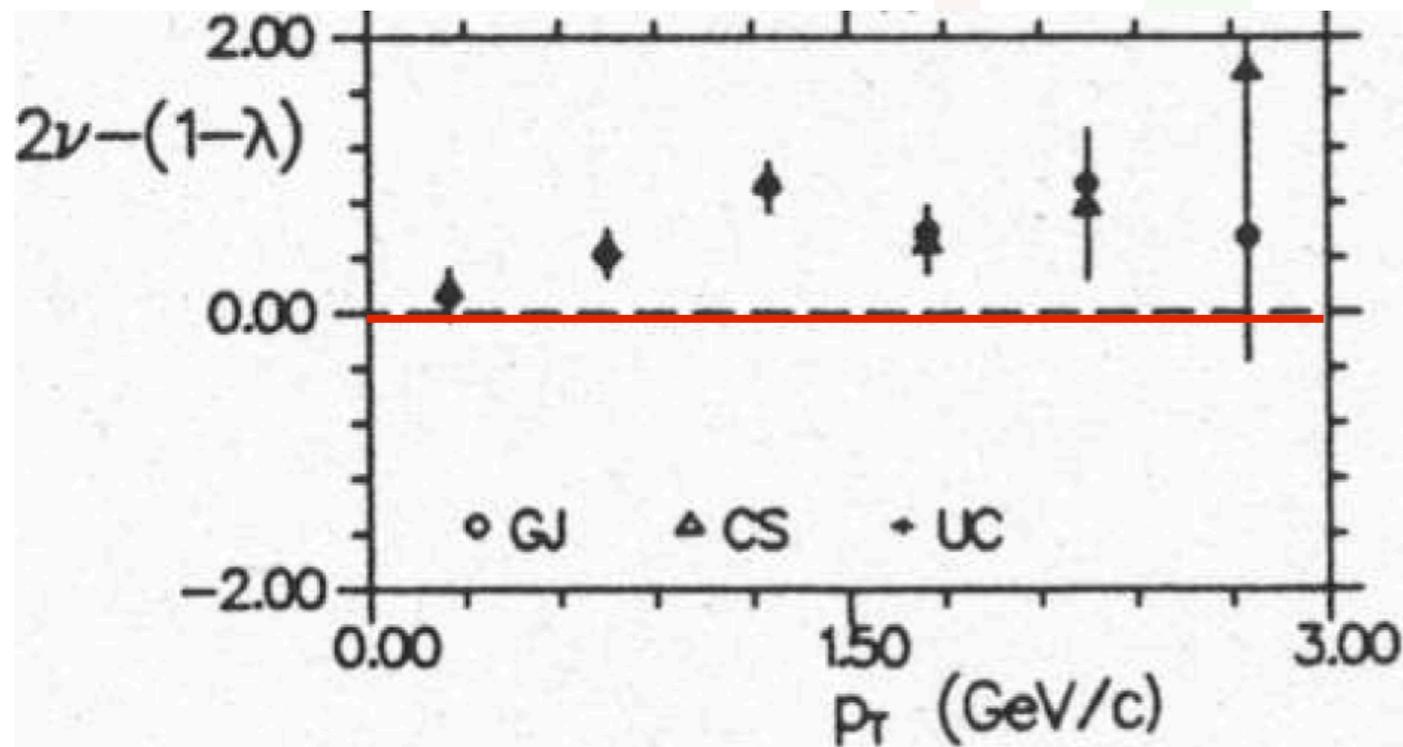
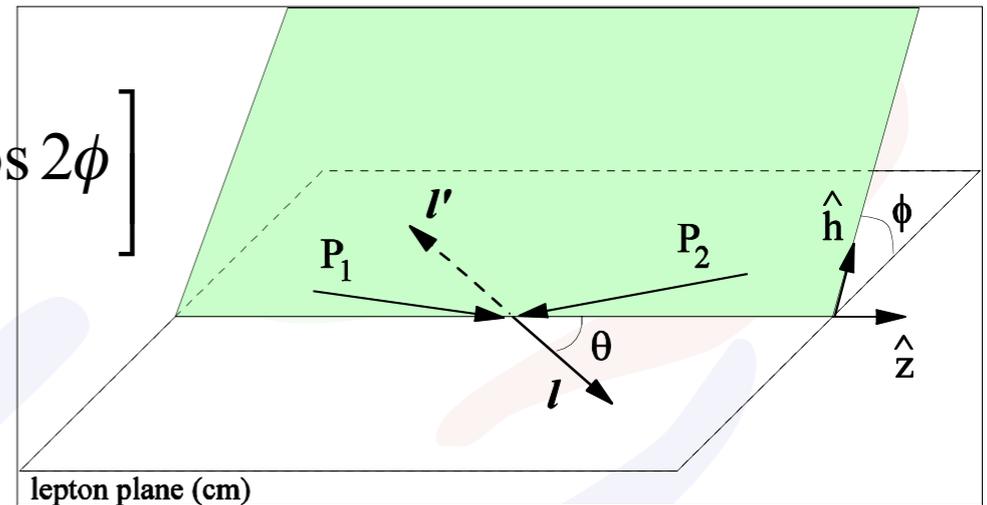
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- “failure” of collinear pQCD

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Large deviations from Lam-Tung relation observed in DY
[NA10 ('86/'88) & E615 ('89)]

- “failure” of collinear pQCD
- possible source: Boer-Mulders effect

Boer-Mulders effect

Sivers effect:

$$f_{1T}^{\perp} = \text{---} \left(\begin{array}{c} \circ \\ \downarrow \\ \circ \end{array} \right) \text{---} \text{---} \left(\begin{array}{c} \uparrow \\ \circ \end{array} \right) \text{---}$$

$S_N \cdot (\mathbf{p}_{\perp} \times \mathbf{P}_N)$

Boer-Mulders effect:

$$h_1^{\perp} = \left(\begin{array}{c} \circ \rightarrow \\ \downarrow \\ \circ \rightarrow \end{array} \right) \text{---} \left(\begin{array}{c} \uparrow \\ \circ \end{array} \right)$$

$S_q \cdot (\mathbf{p}_{\perp} \times \mathbf{P}_N)$



hermes

Boer-Mulders effect

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- spin-effect in unpolarized reactions



hermes

Boer-Mulders effect

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$$f_{1T}^\perp = \text{---} \left(\begin{array}{c} \circ \\ \downarrow \\ \circ \end{array} \right) \text{---} \text{---} \left(\begin{array}{c} \uparrow \\ \circ \end{array} \right) \text{---}$$

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- "QCD Sokolov-Ternov effect" - transverse polarization of "orbiting" quarks

Boer-Mulders effect

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- "QCD Sokolov-Ternov effect" - transverse polarization of "orbiting" quarks
- QCD: sign change for DIS vs. Drell-Yan

Boer-Mulders effect

Sivers effect:

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$S_N \cdot (\mathbf{p}_\perp \times \mathbf{P}_N)$

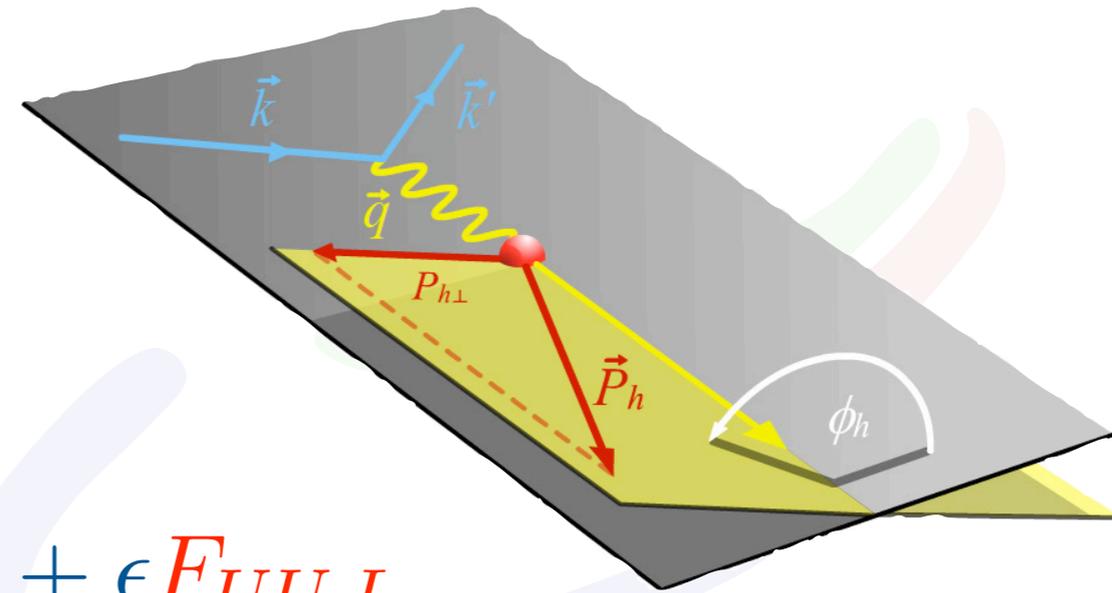
Boer-Mulders effect:

$$h_1^\perp = \left(\begin{array}{c} \circ \rightarrow \\ \downarrow \\ \circ \rightarrow \end{array} \right) \text{---} \left(\begin{array}{c} \uparrow \\ \circ \rightarrow \end{array} \right)$$

$S_q \cdot (\mathbf{p}_\perp \times \mathbf{P}_N)$

- spin-effect in unpolarized reactions
 - "QCD Sokolov-Ternov effect" - transverse polarization of "orbiting" quarks
 - QCD: sign change for DIS vs. Drell-Yan
 - up to now little data from DIS
- ➔ HERMES with most comprehensive data set

Cross section without polarization



$$\frac{d^5 \sigma}{dx dy dz d\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \{F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1-\epsilon)} F_{UU}^{\cos \phi_h} \cos \phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h\}$$

$$F_{XY,Z} = F_{XY,Z}(x, y, z, P_{h\perp})$$

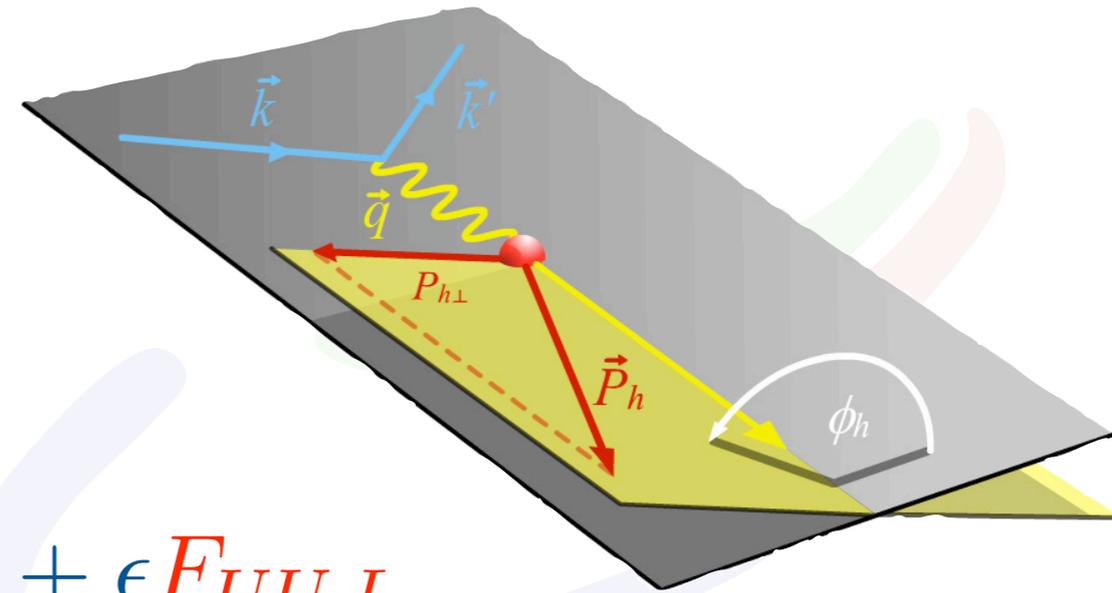
target polarization \downarrow
 \uparrow beam polarization \uparrow virtual-photon polarization

$$\gamma = \frac{2Mx}{Q}$$

$$\epsilon = \frac{1 - y - \frac{1}{4}\gamma^2 y^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2 y^2}$$

[see, e.g., Bacchetta et al., JHEP 0702 (2007) 093]

Cross section without polarization



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leading twist
 $F_{UU}^{\cos 2\phi_h} \propto C \left[\frac{2(\hat{P}_{h\perp} \cdot \vec{k}_T)(\hat{P}_{h\perp} \cdot \vec{p}_T) - \vec{k}_T \cdot \vec{p}_T}{MM_h} h_1^\perp H_1^\perp \right]$

next to leading twist
 $F_{UU}^{\cos \phi_h} \propto \frac{2M}{Q} C \left[\frac{\hat{P}_{h\perp} \cdot \vec{p}_T}{M_h} x h_1^\perp H_1^\perp - \frac{\hat{P}_{h\perp} \cdot \vec{k}_T}{M} x f_1 D_1 + \dots \right]$

Annotations:
 - Blue arrow: BOER-MULDERS EFFECT (points to the leading twist term)
 - Red arrow: CAHN EFFECT (points to the next to leading twist term)
 - Black arrow: Interaction dependent terms neglected (points to the ellipsis in the next to leading twist term)

$$\gamma = \frac{2Mx}{Q}$$

$$\epsilon = \frac{1 - y - \frac{1}{4}\gamma^2 y^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2 y^2}$$

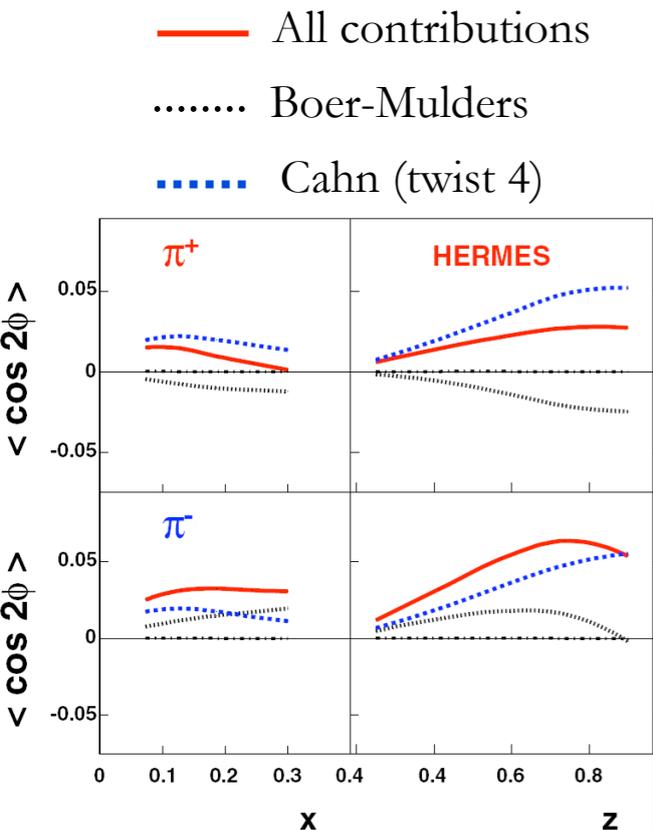
[see, e.g., Bacchetta et al., JHEP 0702 (2007) 093]

(Implicit sum over quark flavours)

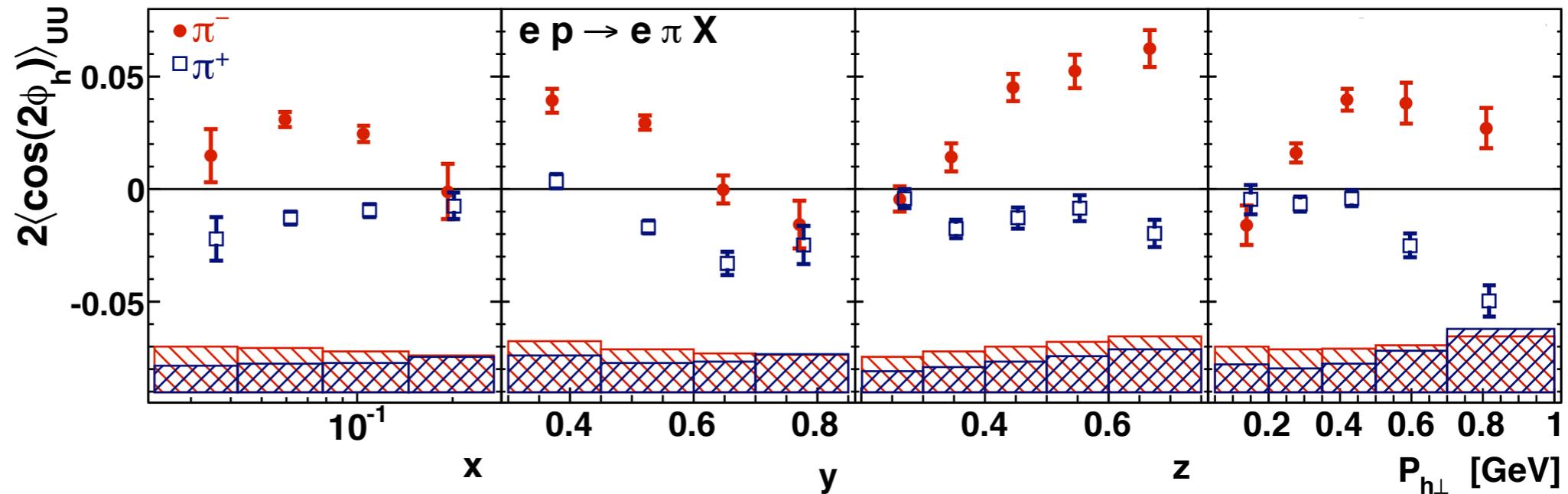
Signs of Boer-Mulders

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

[Airapetian et al., PRD 87 (2013) 012010]



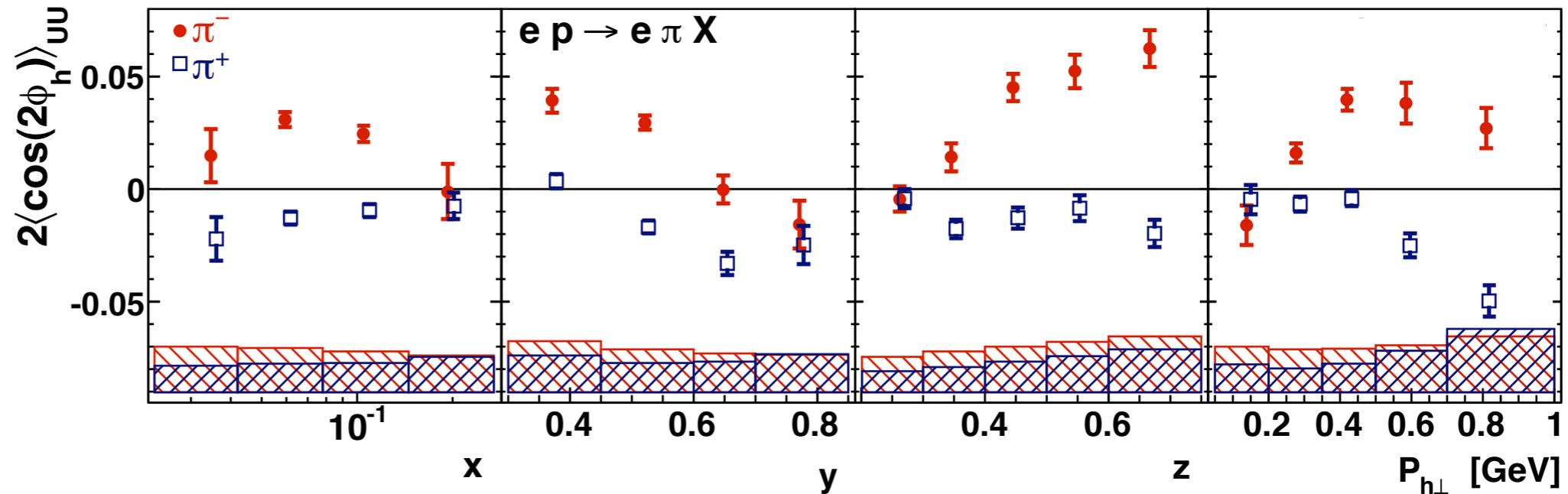
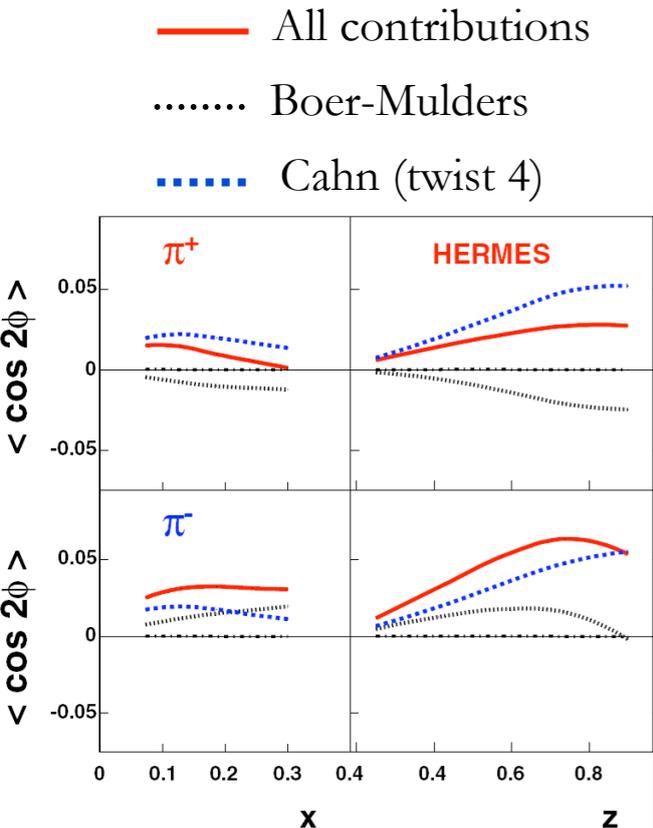
[V. Barone et al., Phys. Rev.D78 (2008) 045022]



Signs of Boer-Mulders

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

[Airapetian et al., PRD 87 (2013) 012010]



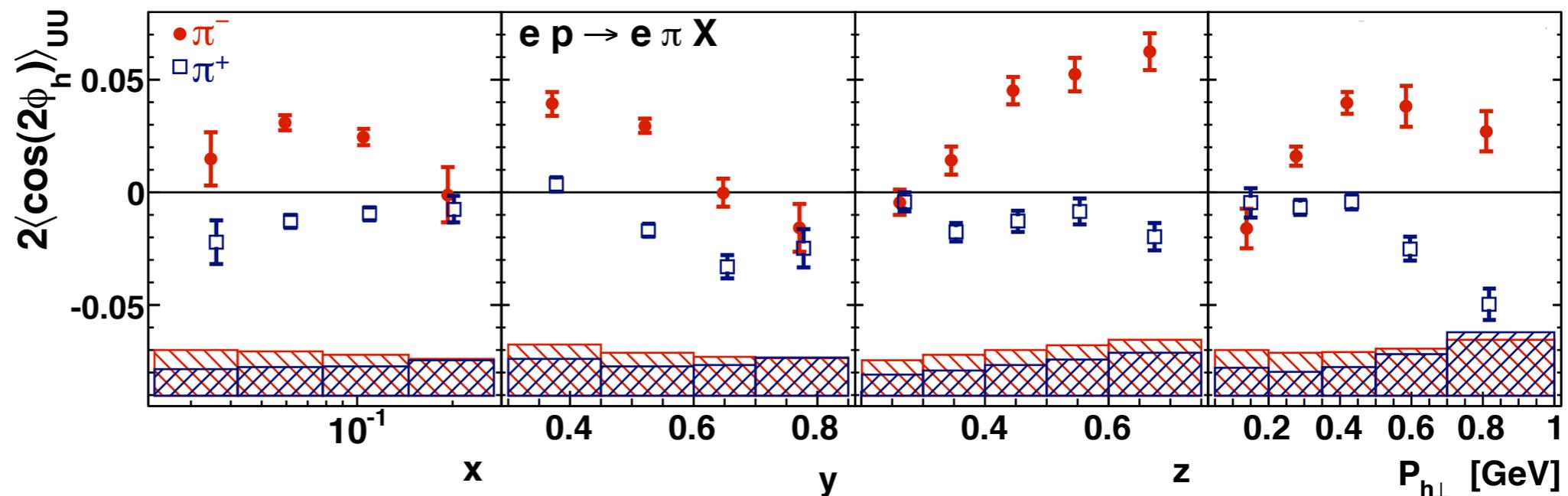
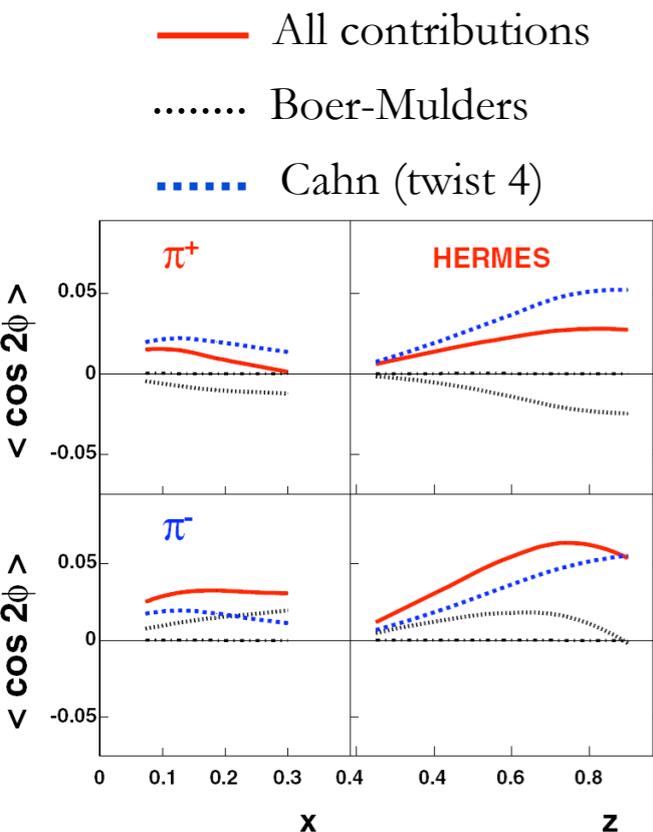
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● Cahn effect only does not describe data

Signs of Boer-Mulders

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

[Airapetian et al., PRD 87 (2013) 012010]



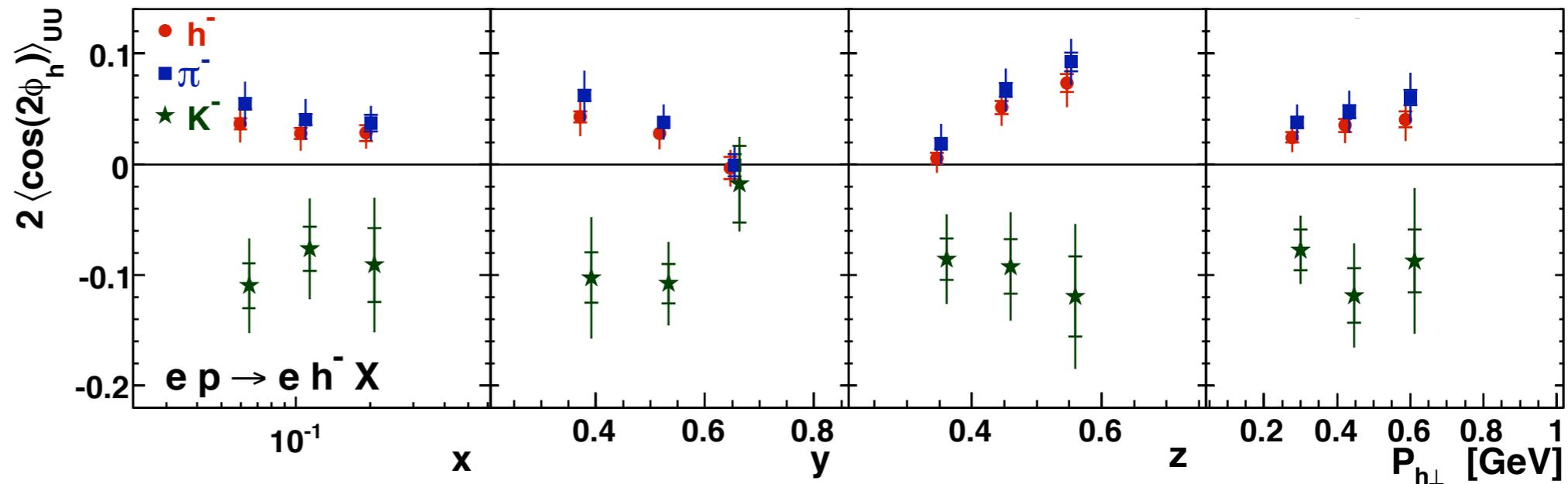
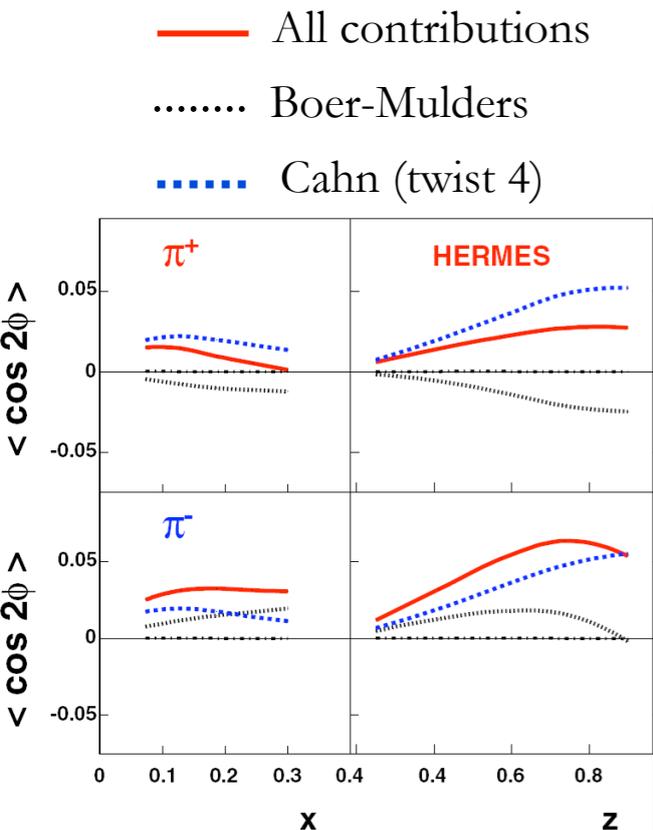
[V. Barone et al., Phys. Rev.D78 (2008) 045022]

- Cahn effect only does not describe data
- opposite sign for charged pions with larger magnitude for π^- (as expected)
 -> same-sign BM-function for valence quarks

Signs of Boer-Mulders

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

[Airapetian et al., PRD 87 (2013) 012010]



[V. Barone et al., Phys. Rev.D78 (2008) 045022]

- Cahn effect only does not describe data
- opposite sign for charged pions with larger magnitude for π^- (as expected)
 -> same-sign BM-function for valence quarks
- intriguing behavior for kaons

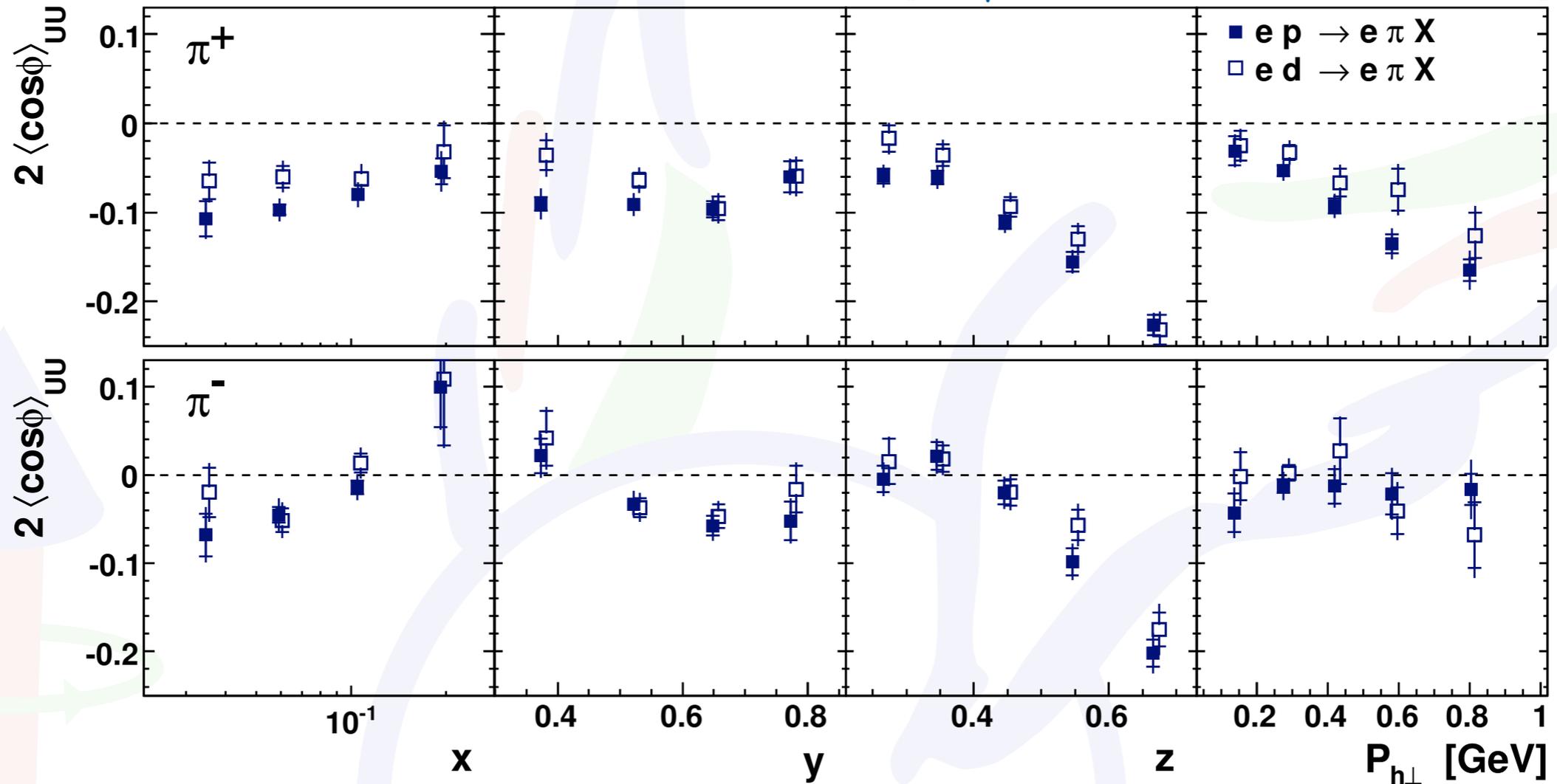
Cahn effect?

next to leading twist
 $F_{UU}^{\cos\phi_h} \propto \frac{2M}{Q} C$

$$C \left[\underbrace{-\frac{\hat{P}_{h\perp} \cdot \vec{p}_T}{M_h} x h_1^\perp H_1^\perp}_{\text{BOER-MULDERS EFFECT}} - \underbrace{\frac{\hat{P}_{h\perp} \cdot \vec{k}_T}{M} x f_1 D_1}_{\text{CAHN EFFECT}} + \dots \right]$$

Interaction dependent terms neglected

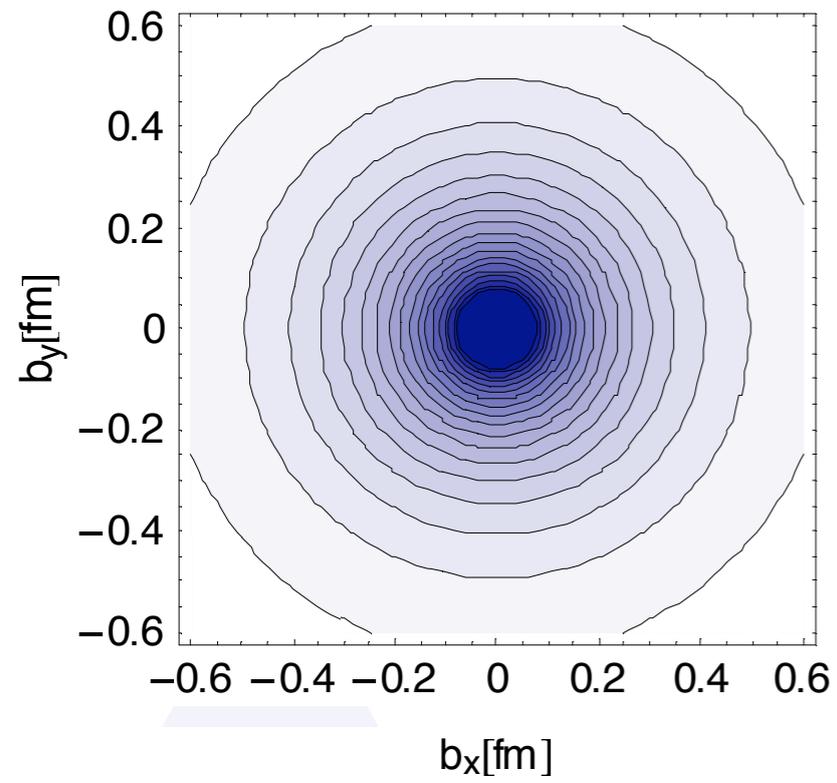
[Airapetian et al., PRD 87 (2013) 012010]



- no dependence on hadron charge expected for Cahn effect
- ➔ flavor dependence of transverse momentum
- ➔ sign of Boer-Mulders in $\cos\phi$ modulation (indeed, overall pattern resembles B-M modulations)
- ➔ additional "genuine" twist-3?

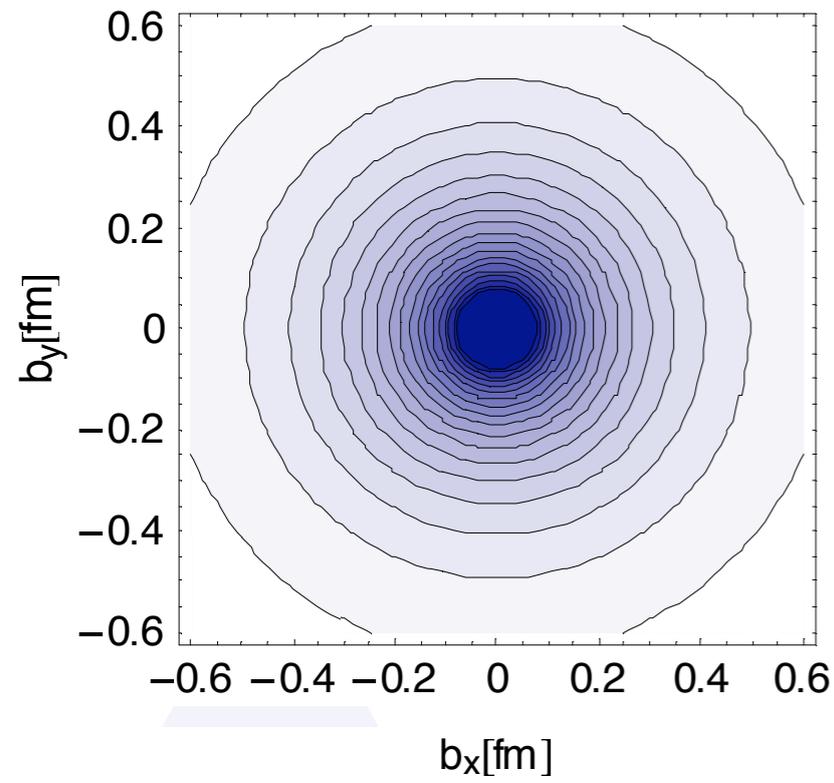
Exclusive reactions

A complementary 3D picture of the nucleon

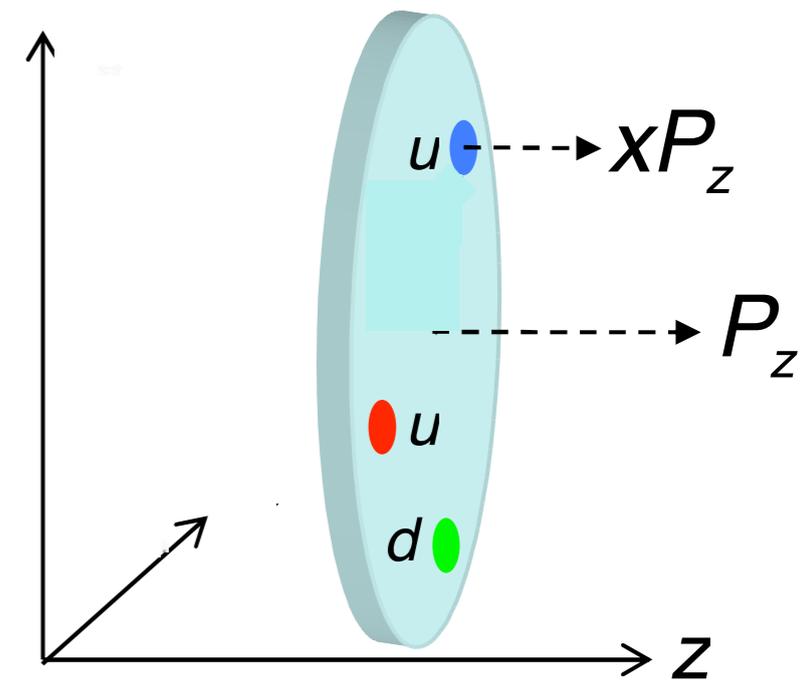


Form factors:
transverse distribution
of partons

A complementary 3D picture of the nucleon

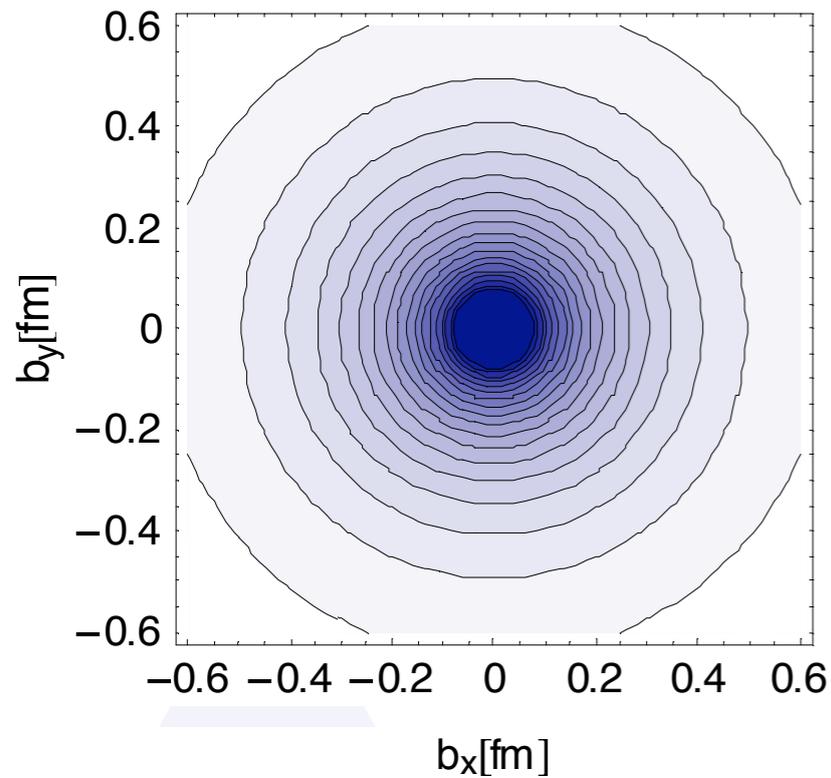


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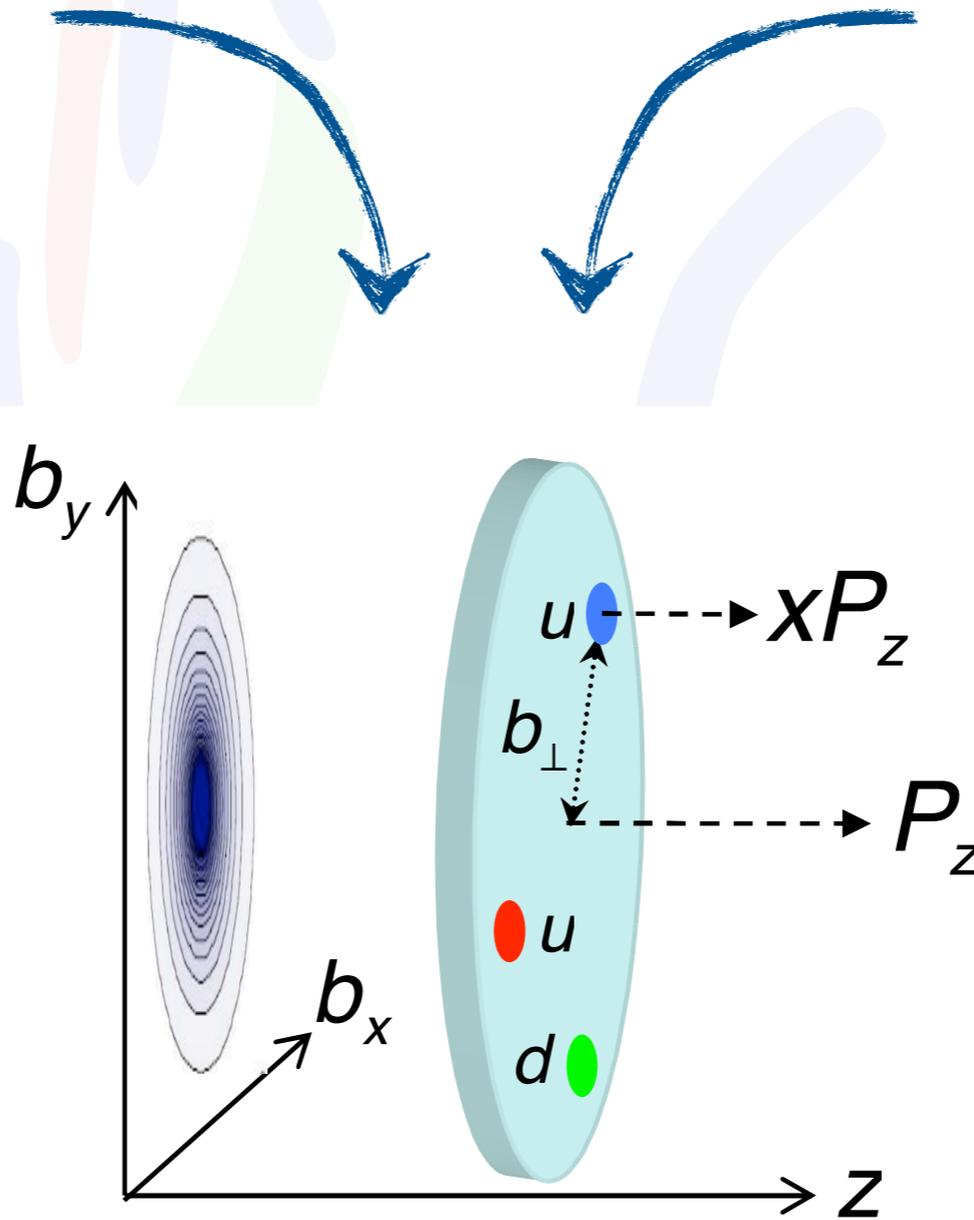


Parton distributions:
longitudinal momentum
of partons

A complementary 3D picture of the nucleon

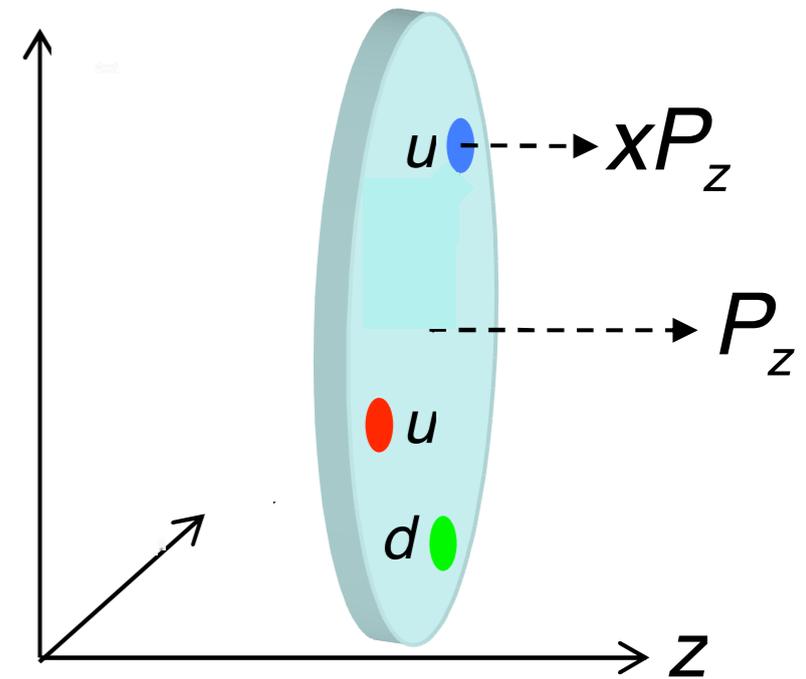


Form factors:
transverse distribution
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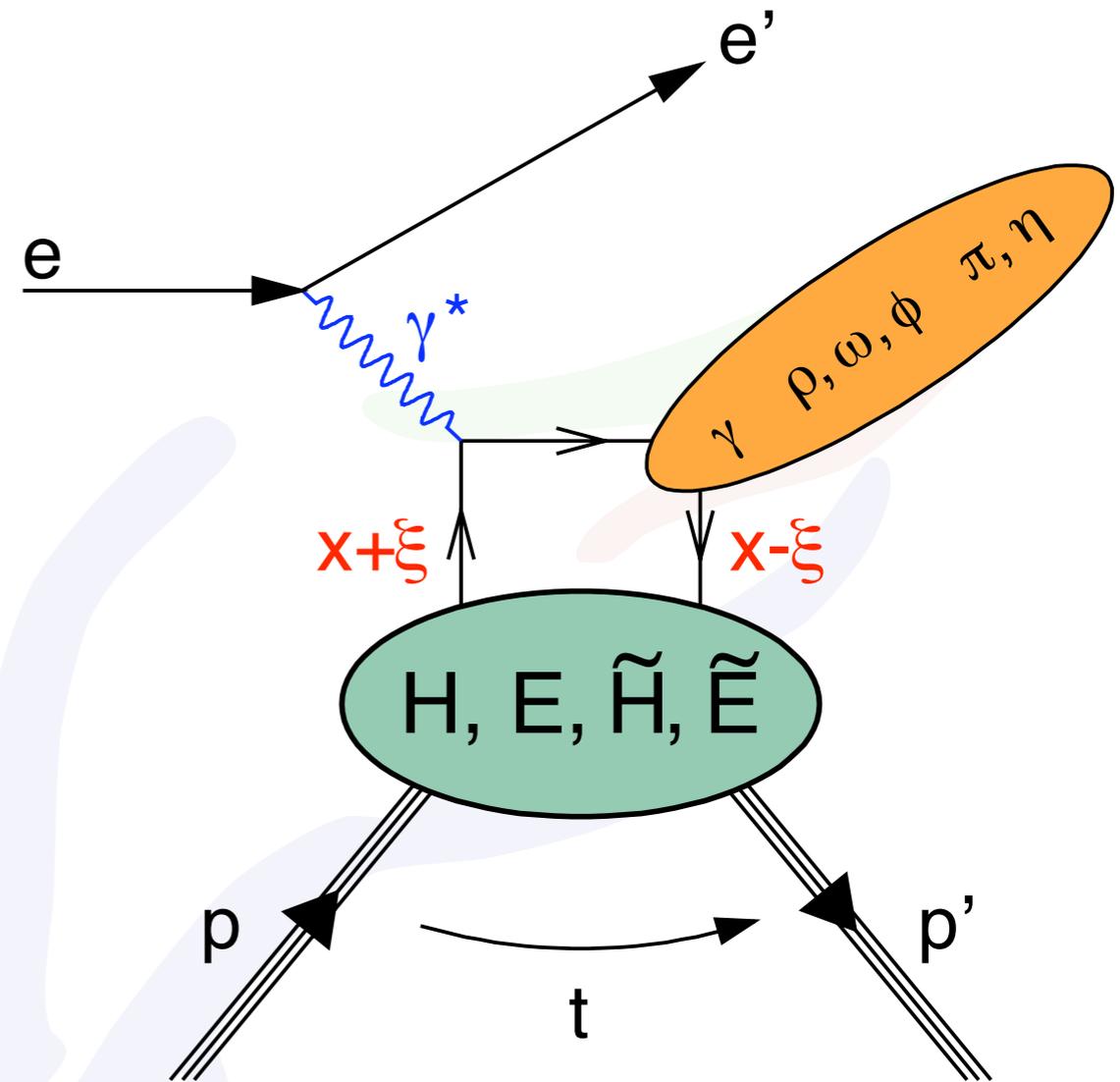
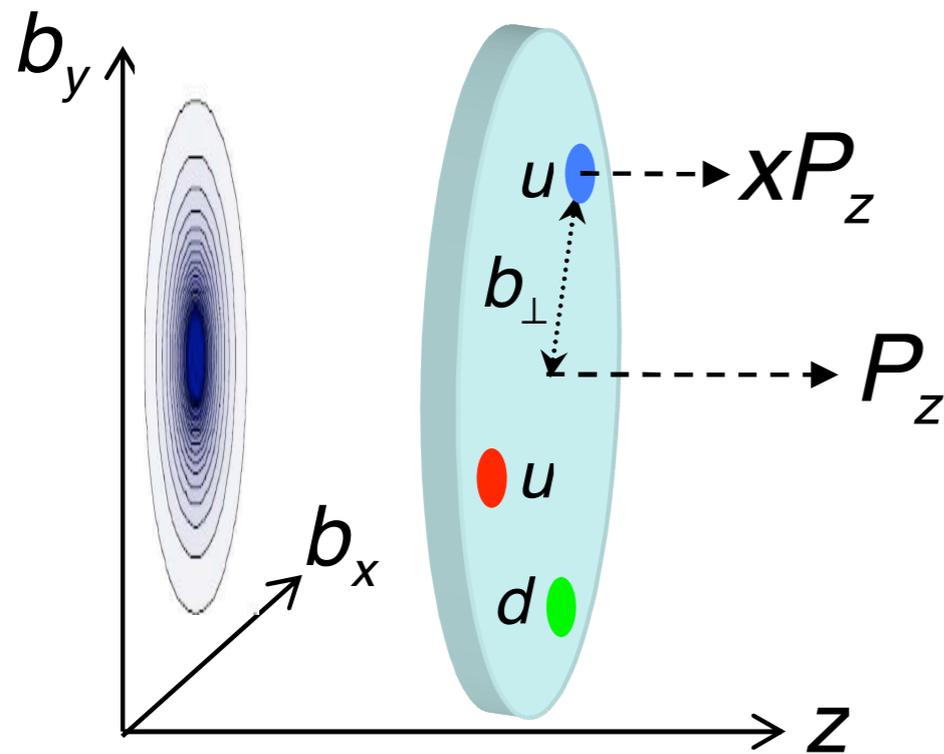
Nucleon Tomography

correlated info on transverse position and longitudinal momentum



Parton distributions:
longitudinal momentum
of partons

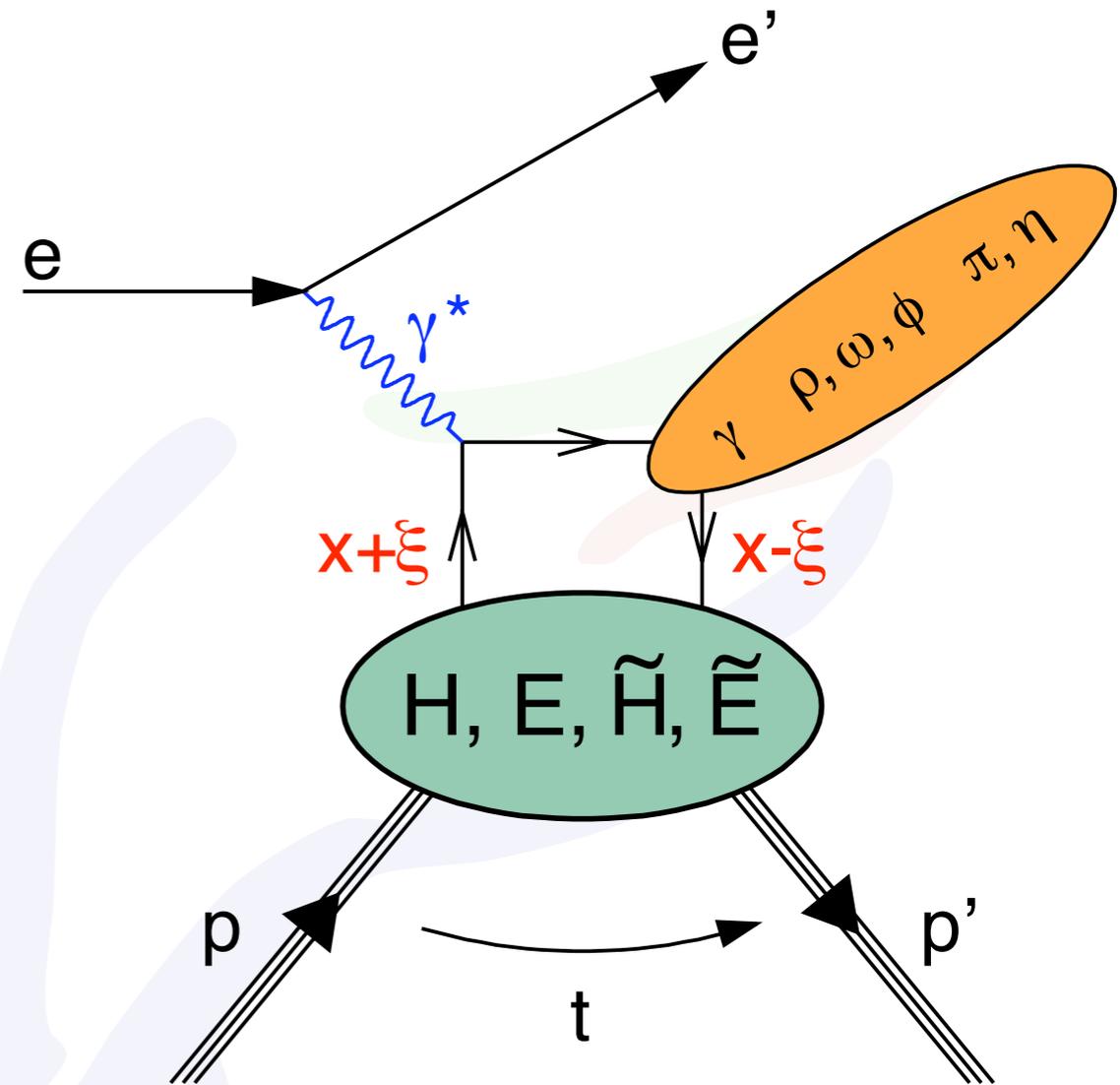
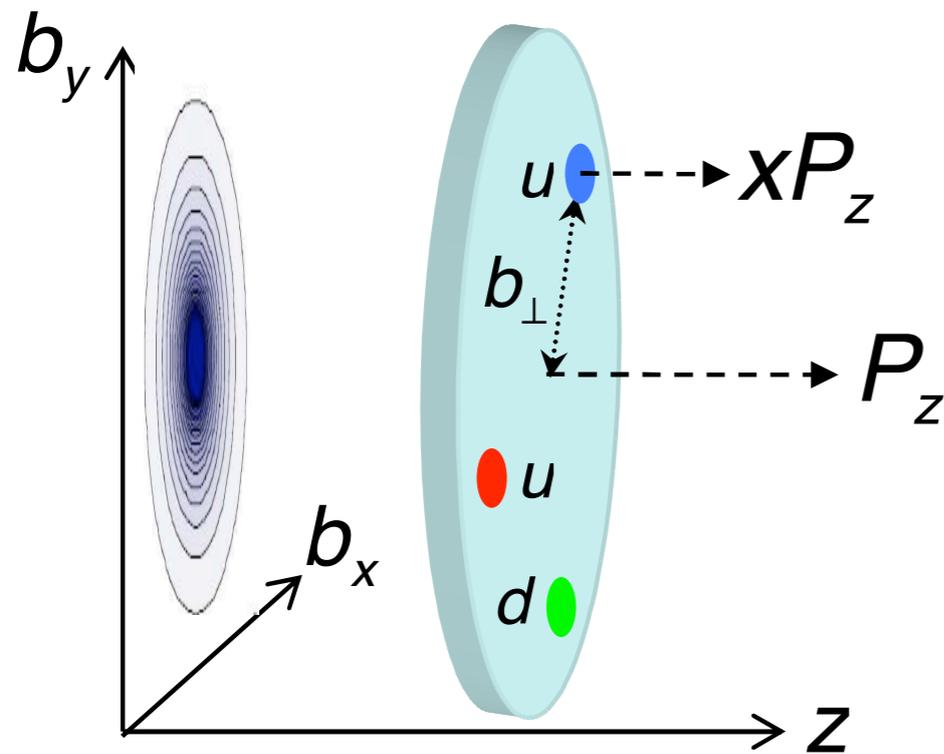
Probing GPDs in Exclusive Reactions



x : average longitudinal momentum fraction of active quark (usually not observed & $x \neq x_B$)

ξ : half the longitudinal momentum change $\approx x_B/(2-x_B)$

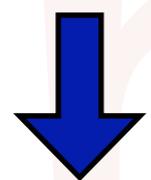
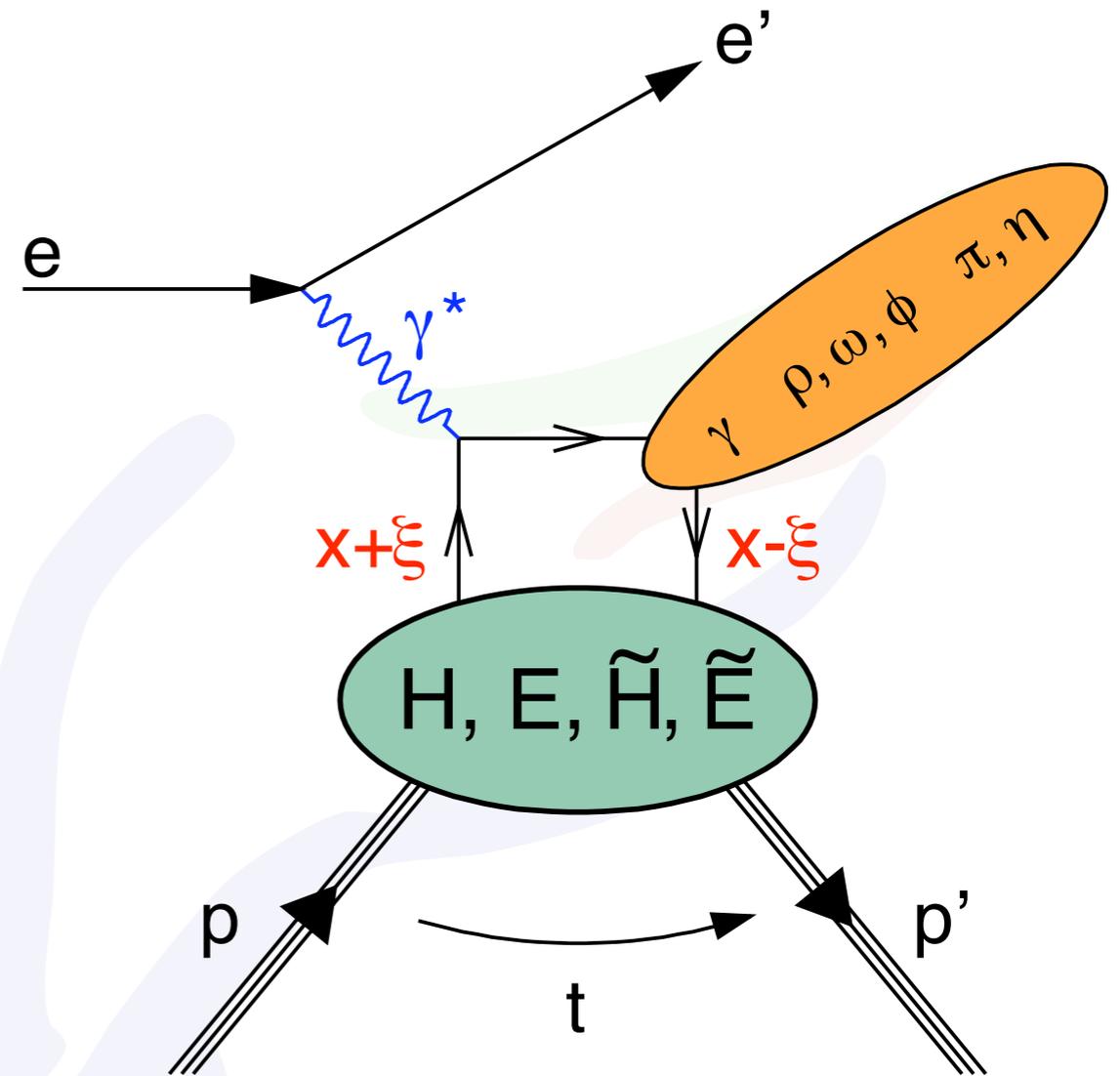
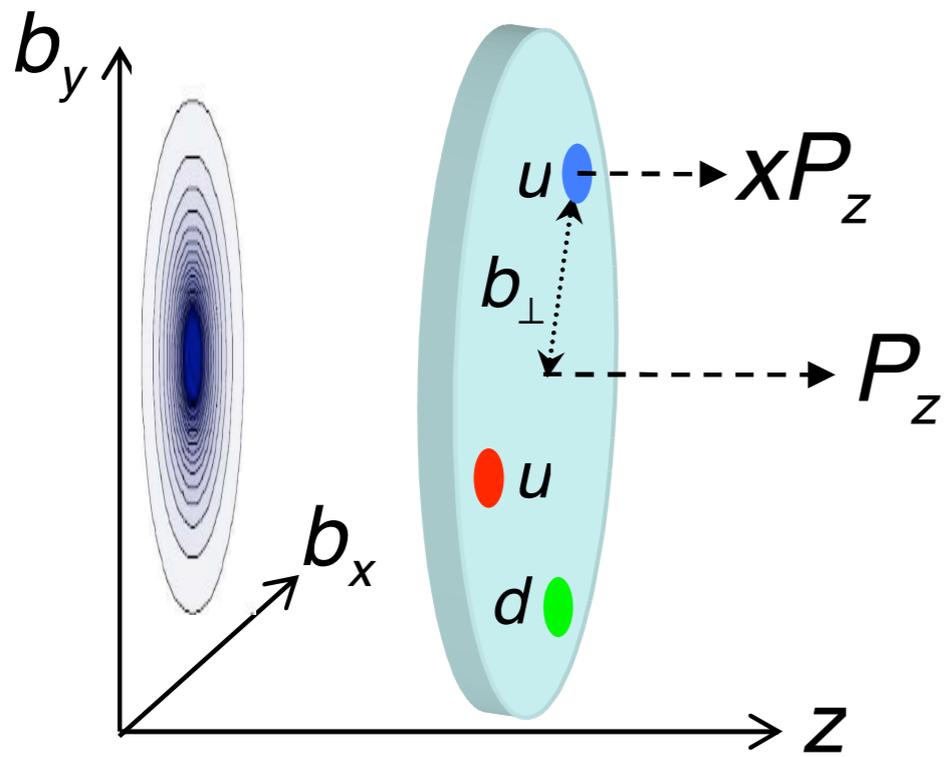
Probing GPDs in Exclusive Reactions



	no quark helicity flip	quark helicity flip
no nucleon helicity flip	H	\tilde{H}
nucleon helicity flip	E	\tilde{E}

(+ 4 more chiral-odd functions)

Probing GPDs in Exclusive Reactions



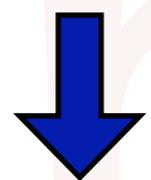
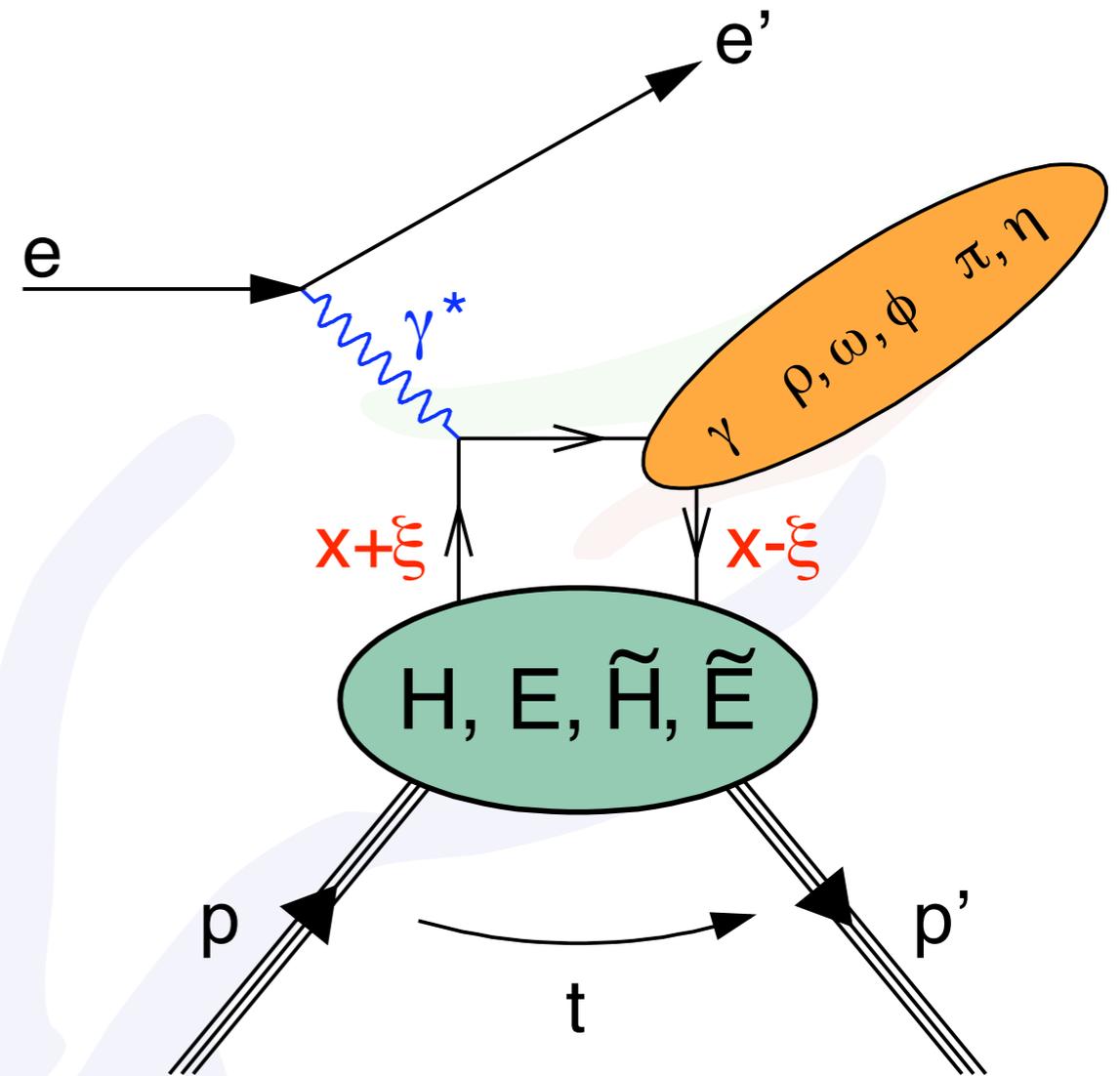
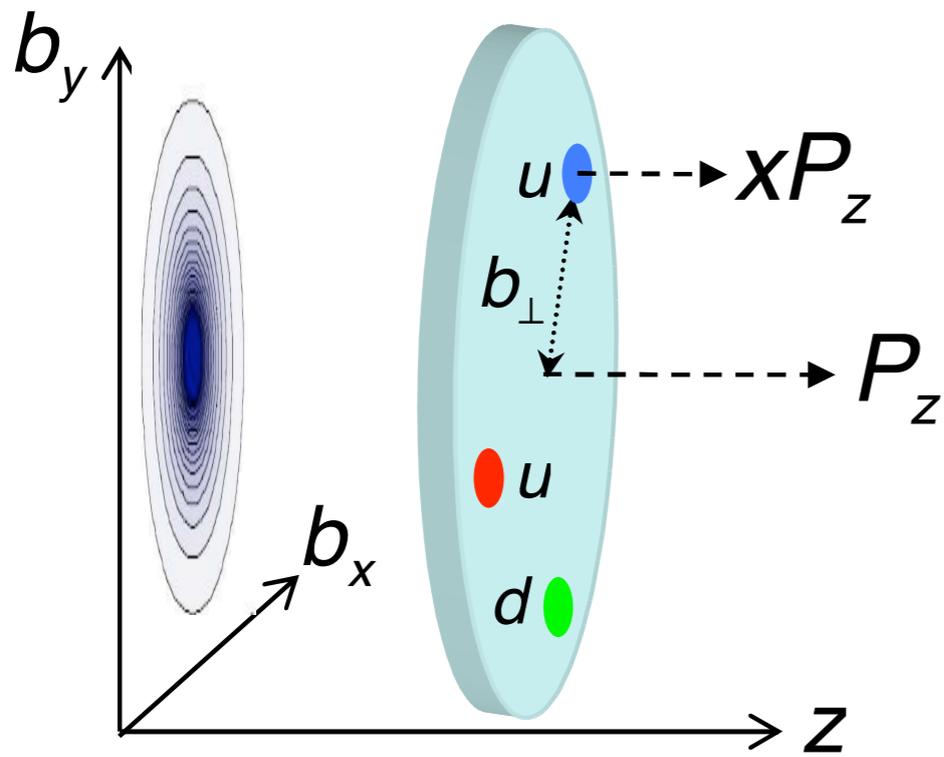
$$\int dx H^q(x, \xi, t) = F_1^q(t)$$

$$\int dx E^q(x, \xi, t) = F_2^q(t)$$

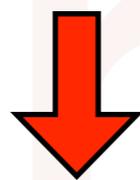
	no quark helicity flip	quark helicity flip
no nucleon helicity flip	H	\tilde{H}
nucleon helicity flip	E	\tilde{E}

(+ 4 more chiral-odd functions)

Probing GPDs in Exclusive Reactions



$$\int dx H^q(x, \xi, t) = F_1^q(t)$$



$$H^q(x, \xi = 0, t = 0) = q(x)$$

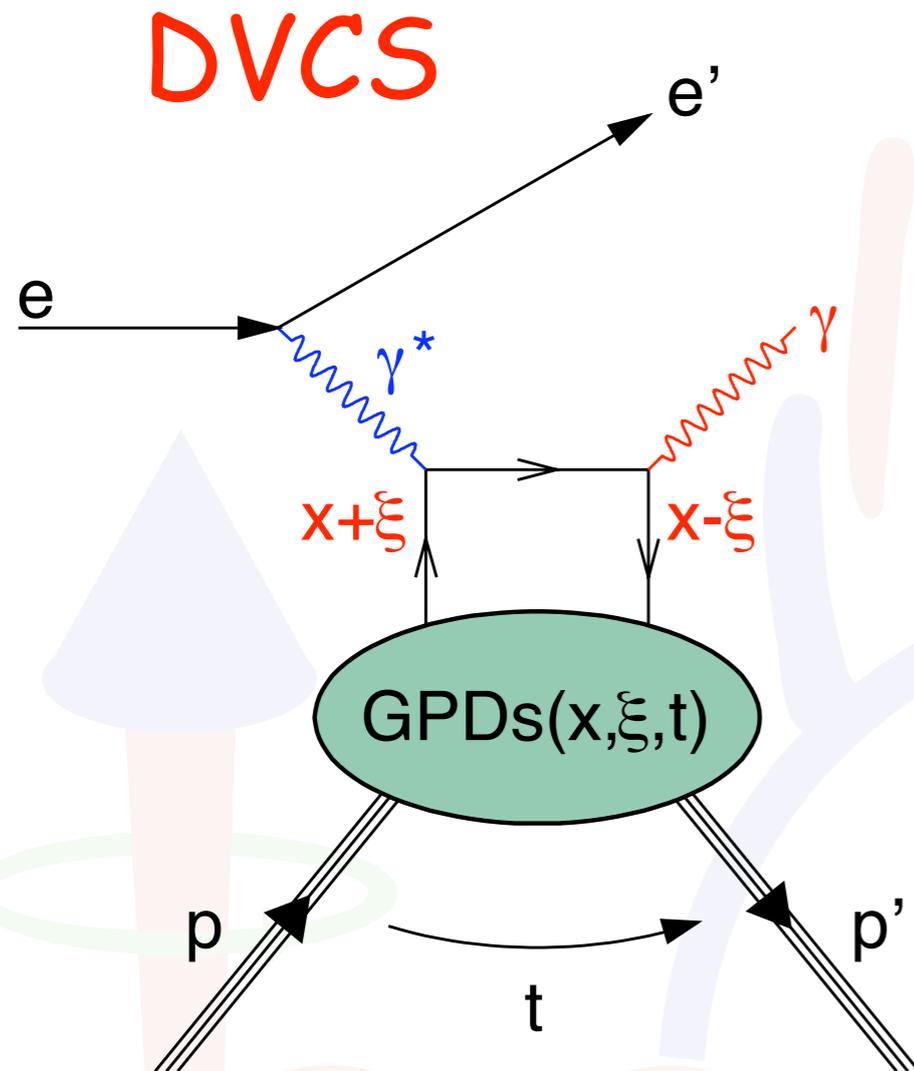
$$\int dx E^q(x, \xi, t) = F_2^q(t)$$

$$\tilde{H}^q(x, \xi = 0, t = 0) = \Delta q(x)$$

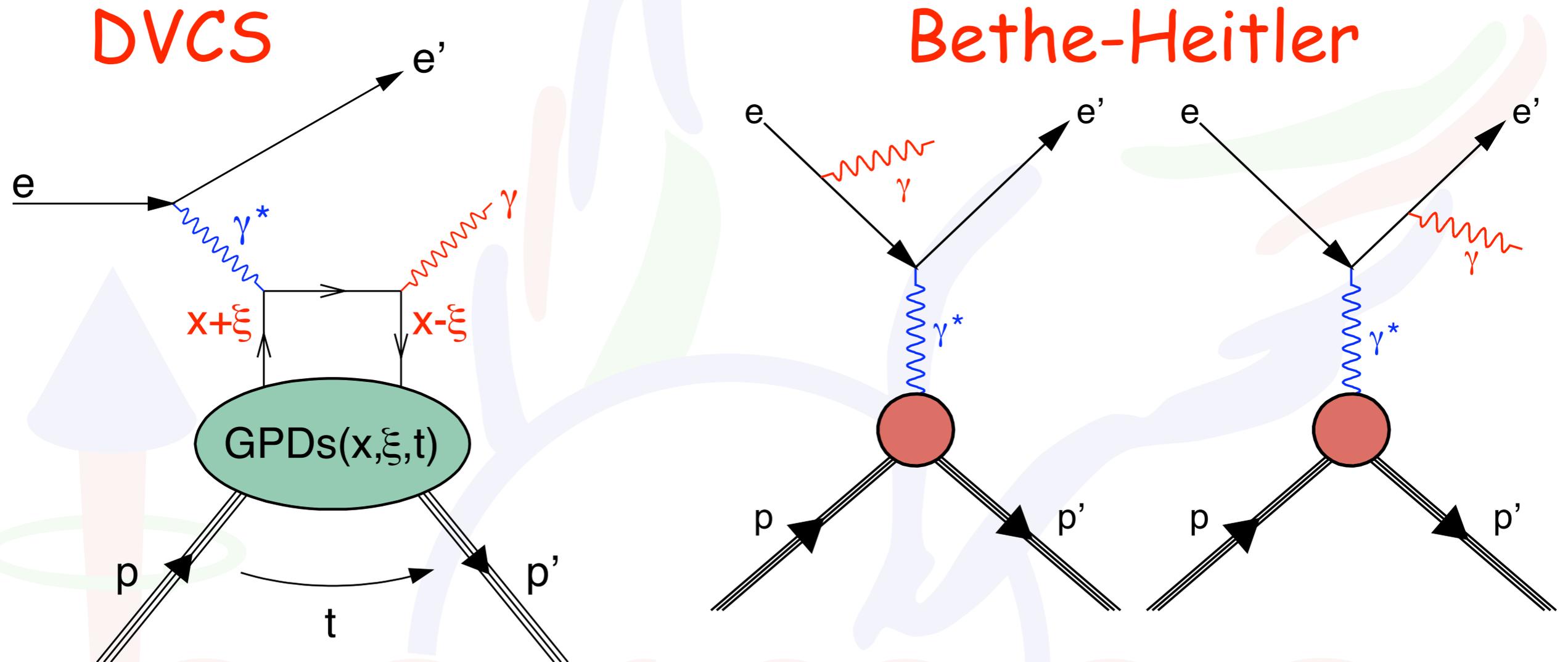
	no quark helicity flip	quark helicity flip
no nucleon helicity flip	H	\tilde{H}
nucleon helicity flip	E	\tilde{E}

(+ 4 more chiral-odd functions)

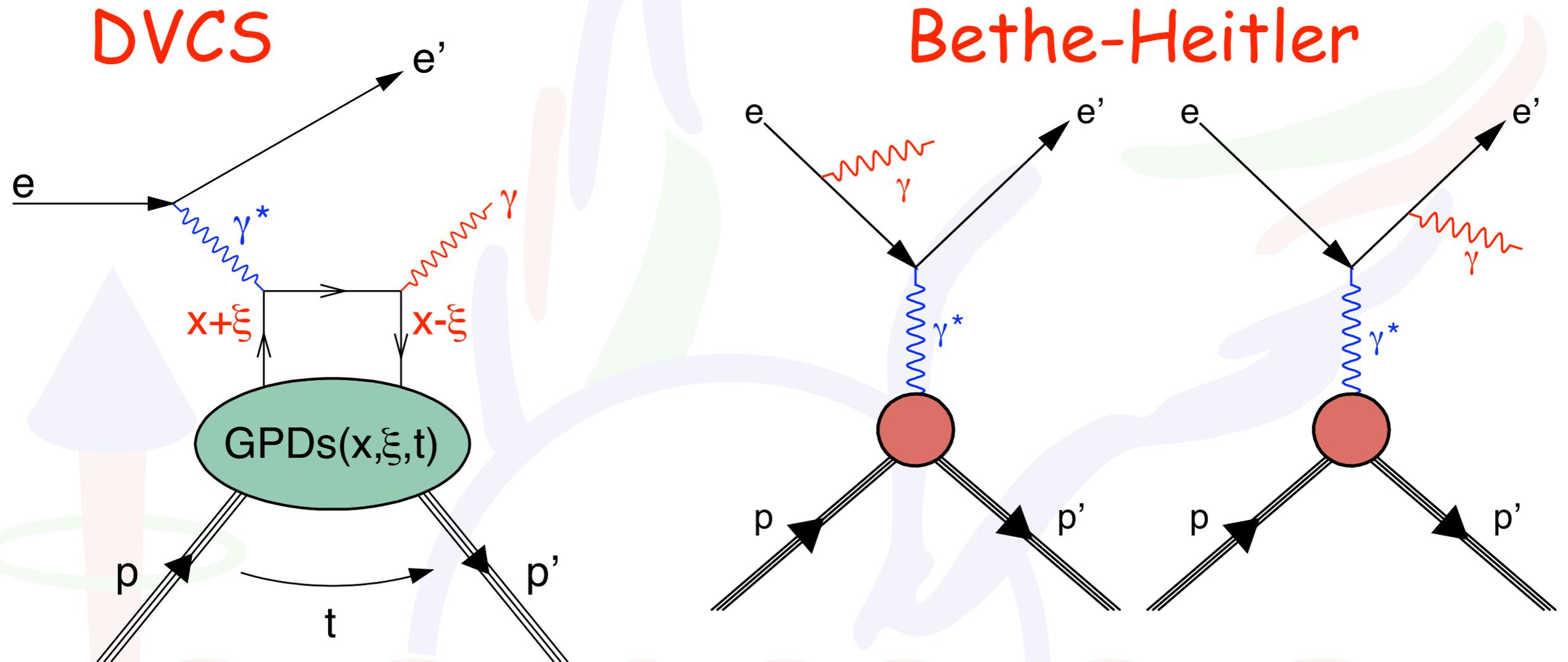
Real-photon production



Real-photon production



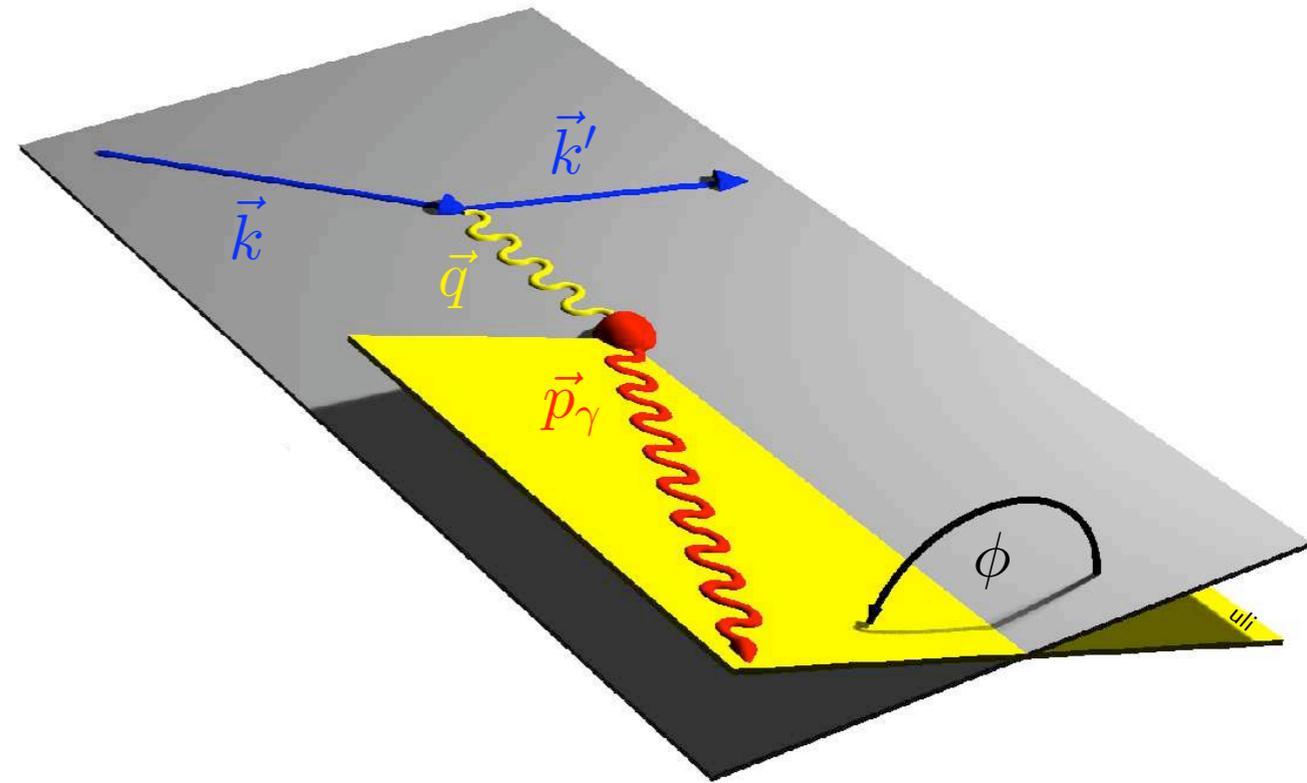
Real-photon production



$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi} = \frac{y^2}{32(2\pi)^4 \sqrt{1 + \frac{4M^2 x_B^2}{Q^2}}} (|\mathcal{T}_{\text{DVCS}}|^2 + |\mathcal{T}_{\text{BH}}|^2 + \mathcal{I})$$

Azimuthal dependences in DVCS/BH

- beam polarization P_B
- beam charge C_B
- here: unpolarized target



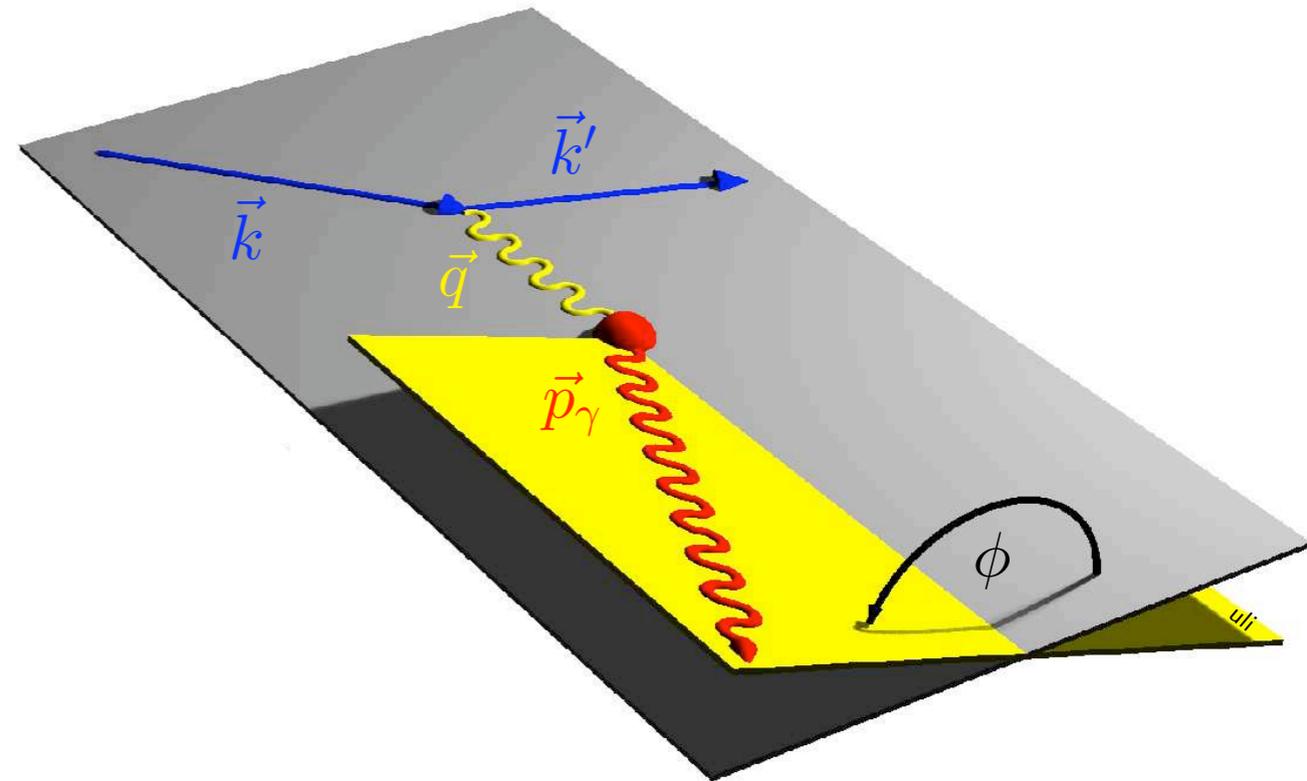
Fourier expansion for ϕ :

$$|\mathcal{T}_{\text{BH}}|^2 = \frac{K_{\text{BH}}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \sum_{n=0}^2 c_n^{\text{BH}} \cos(n\phi)$$

calculable in QED
(using FF measurements)

Azimuthal dependences in DVCS/BH

- beam polarization P_B
- beam charge C_B
- here: unpolarized target



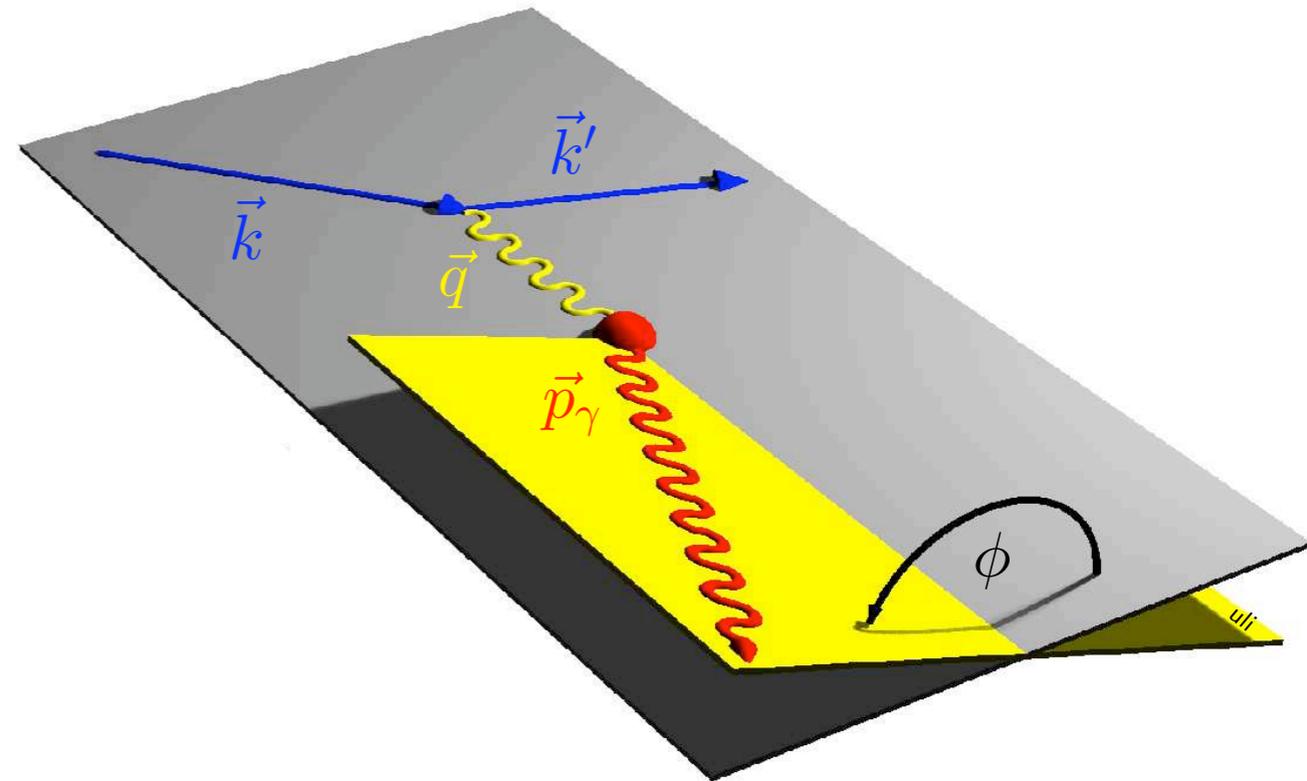
Fourier expansion for ϕ :

$$|\mathcal{T}_{BH}|^2 = \frac{K_{BH}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \sum_{n=0}^2 c_n^{BH} \cos(n\phi)$$

$$|\mathcal{T}_{DVCS}|^2 = K_{DVCS} \left[\sum_{n=0}^2 c_n^{DVCS} \cos(n\phi) + P_B \sum_{n=1}^1 s_n^{DVCS} \sin(n\phi) \right]$$

Azimuthal dependences in DVCS/BH

- beam polarization P_B
- beam charge C_B
- here: unpolarized target



Fourier expansion for ϕ :

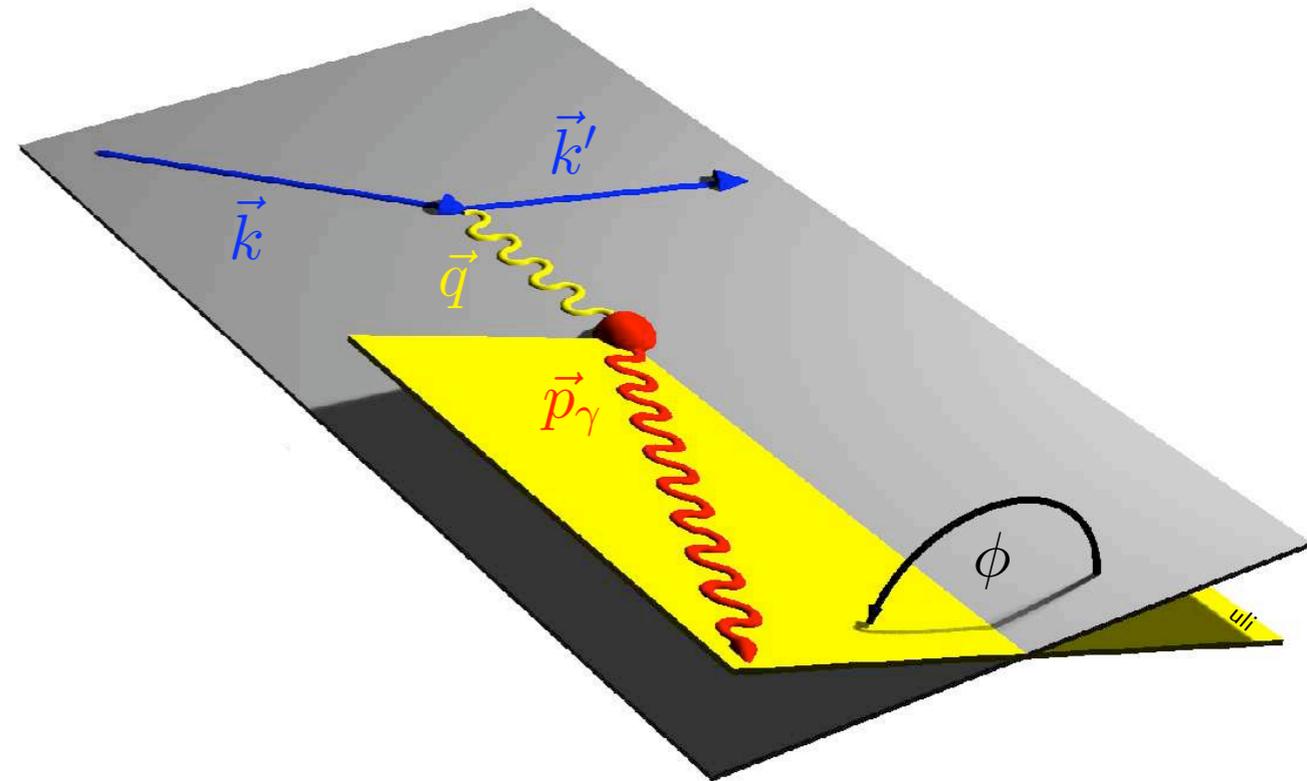
$$|\mathcal{T}_{\text{BH}}|^2 = \frac{K_{\text{BH}}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \sum_{n=0}^2 c_n^{\text{BH}} \cos(n\phi)$$

$$|\mathcal{T}_{\text{DVCS}}|^2 = K_{\text{DVCS}} \left[\sum_{n=0}^2 c_n^{\text{DVCS}} \cos(n\phi) + P_B \sum_{n=1}^1 s_n^{\text{DVCS}} \sin(n\phi) \right]$$

$$\mathcal{I} = \frac{C_B K_{\mathcal{I}}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \left[\sum_{n=0}^3 c_n^{\mathcal{I}} \cos(n\phi) + P_B \sum_{n=1}^2 s_n^{\mathcal{I}} \sin(n\phi) \right]$$

Azimuthal dependences in DVCS/BH

- beam polarization P_B
- beam charge C_B
- here: unpolarized target



Fourier expansion for ϕ :

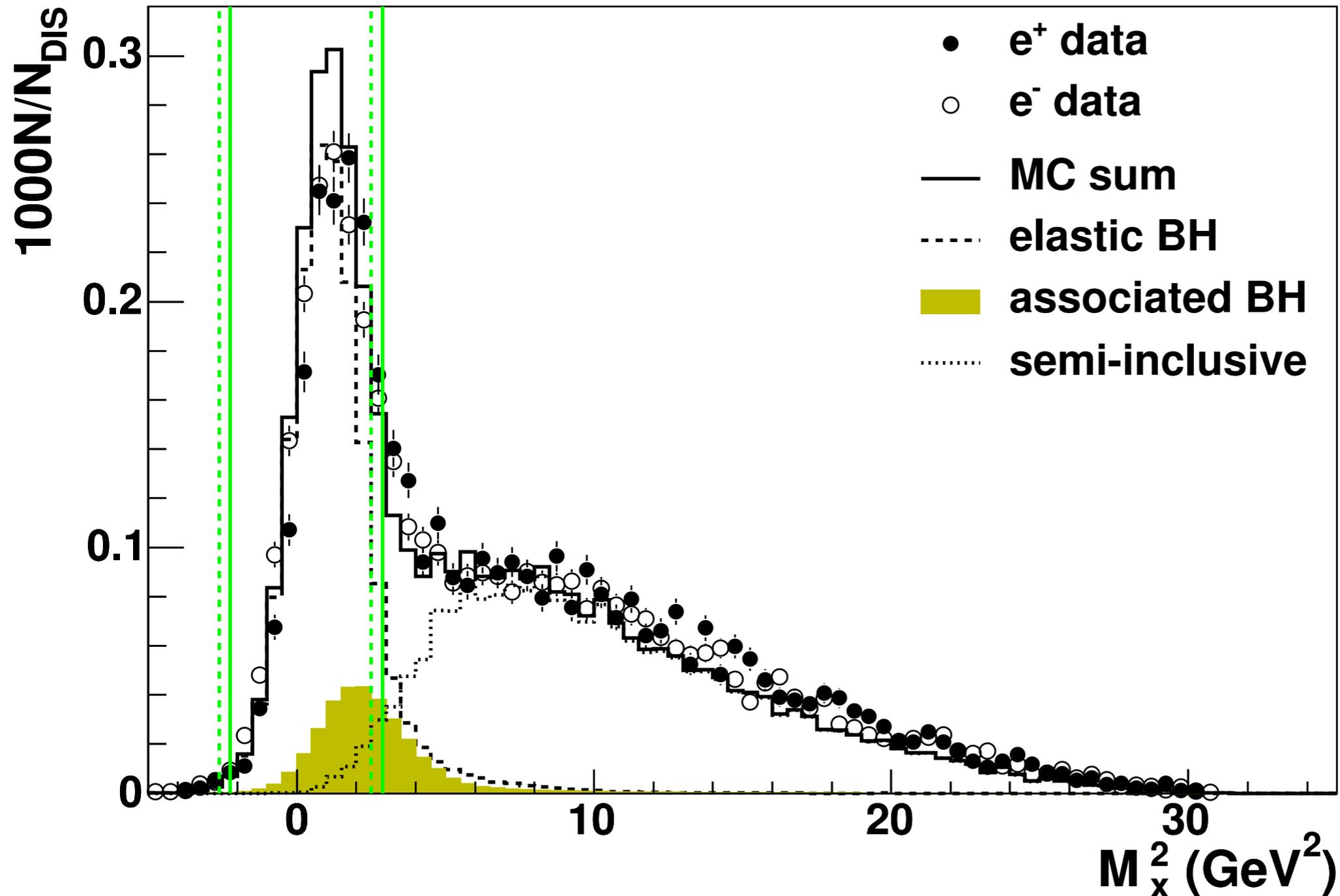
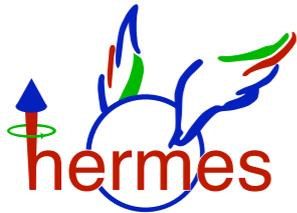
$$|\mathcal{T}_{\text{BH}}|^2 = \frac{K_{\text{BH}}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \sum_{n=0}^2 c_n^{\text{BH}} \cos(n\phi)$$

$$|\mathcal{T}_{\text{DVCS}}|^2 = K_{\text{DVCS}} \left[\sum_{n=0}^2 c_n^{\text{DVCS}} \cos(n\phi) + P_B \sum_{n=1}^1 s_n^{\text{DVCS}} \sin(n\phi) \right]$$

$$\mathcal{I} = \frac{C_B K_{\mathcal{I}}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \left[\sum_{n=0}^3 c_n^{\mathcal{I}} \cos(n\phi) + P_B \sum_{n=1}^2 s_n^{\mathcal{I}} \sin(n\phi) \right]$$

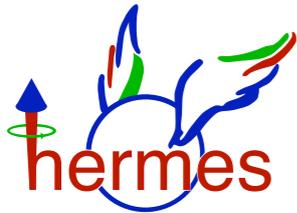
bilinear ("DVCS") or linear in GPDs

Exclusivity: missing-mass technique

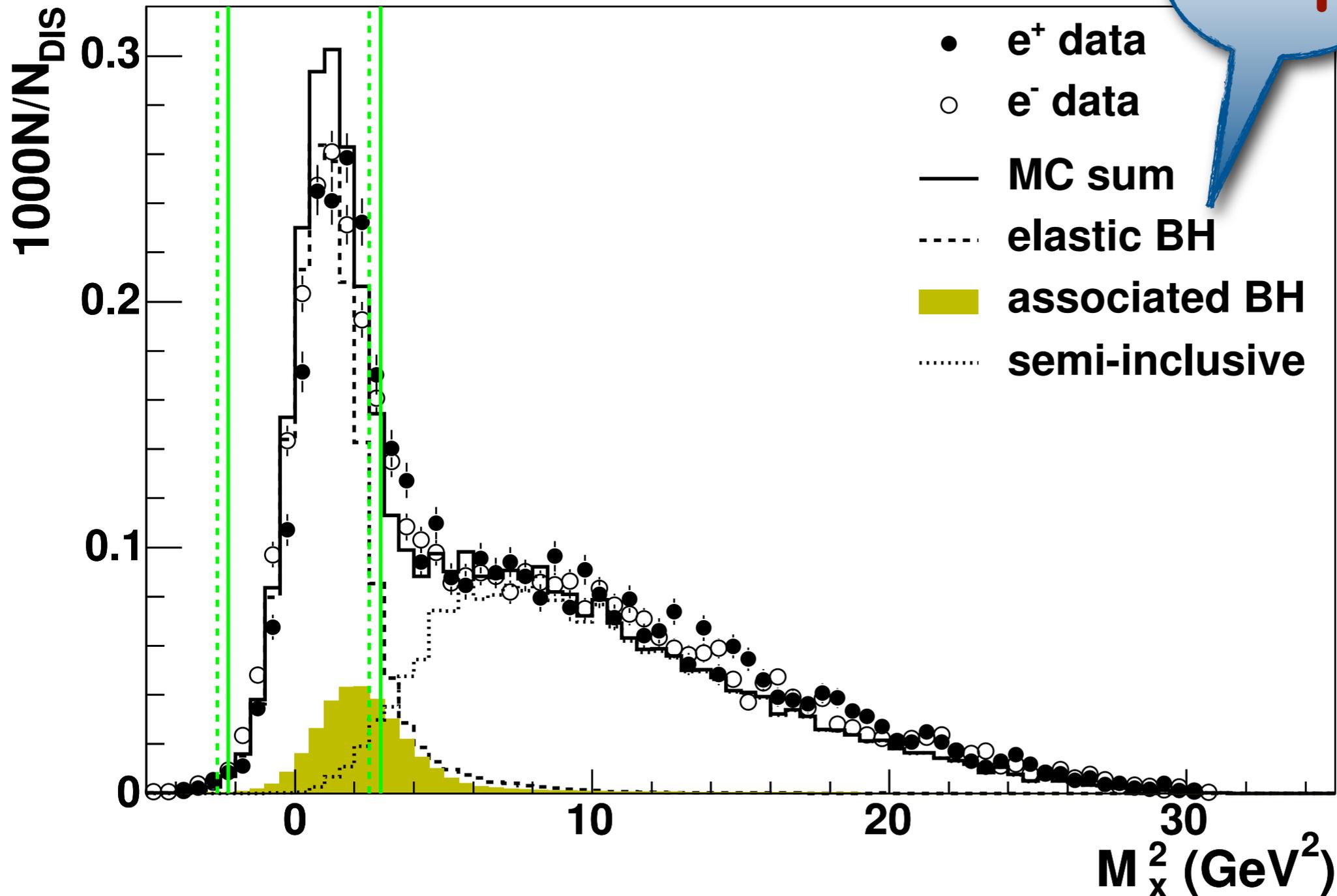


S

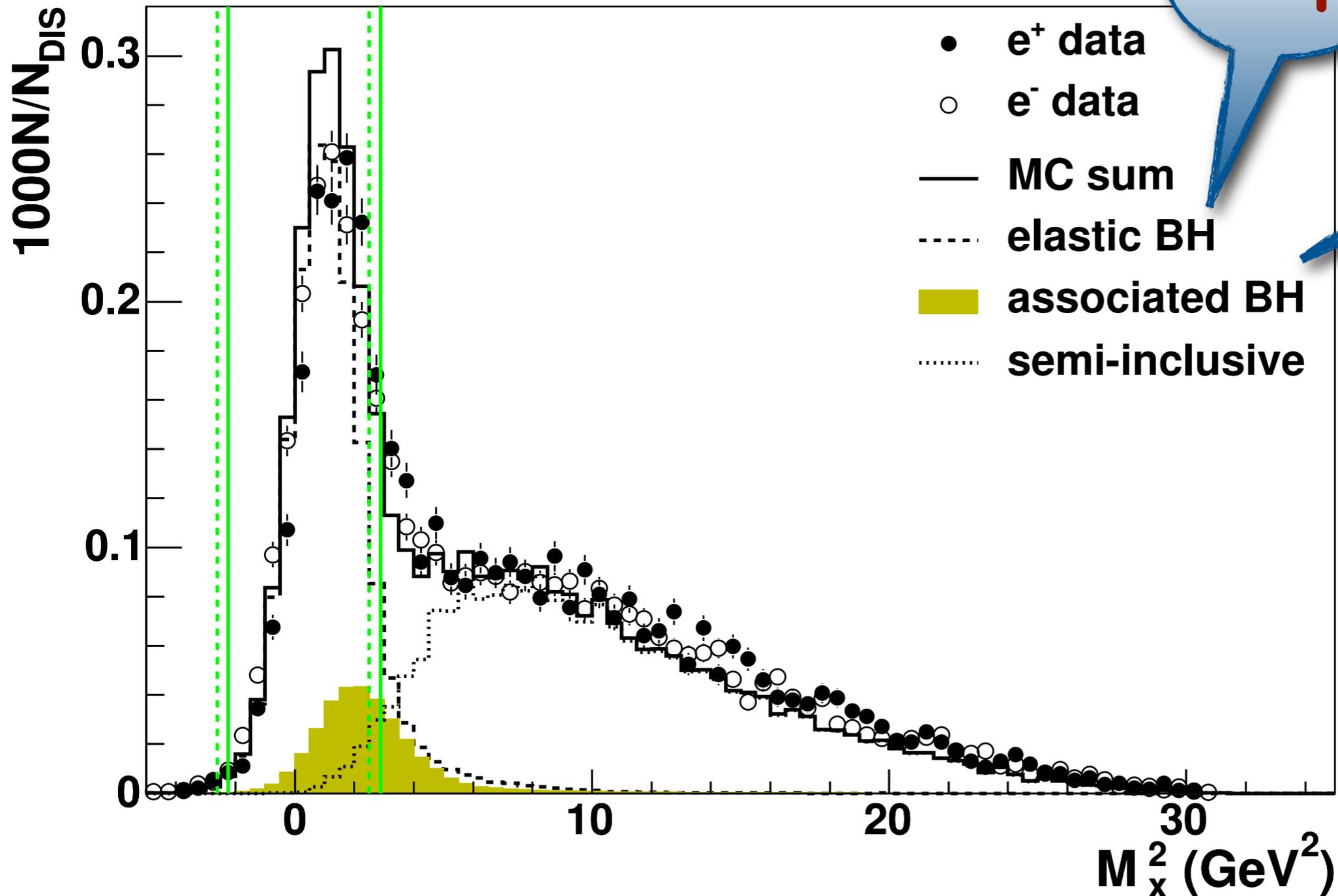
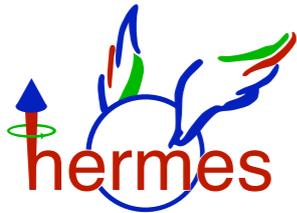
Exclusivity: missing-mass technique



$X=p$



Exclusivity: missing-mass technique

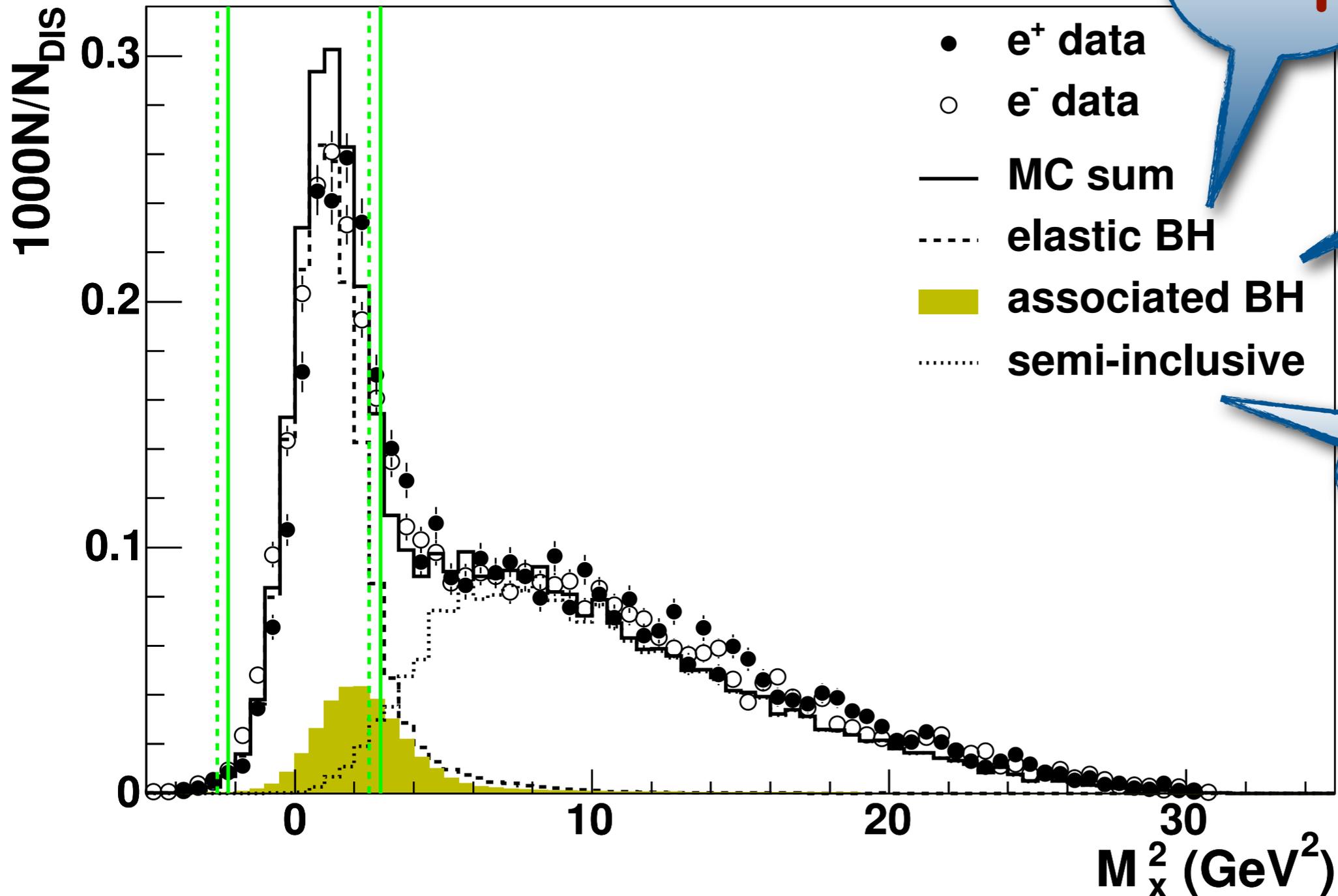
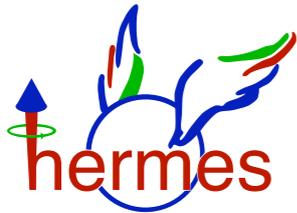


$X=p$

$X=\Delta^+$

S

Exclusivity: missing-mass technique

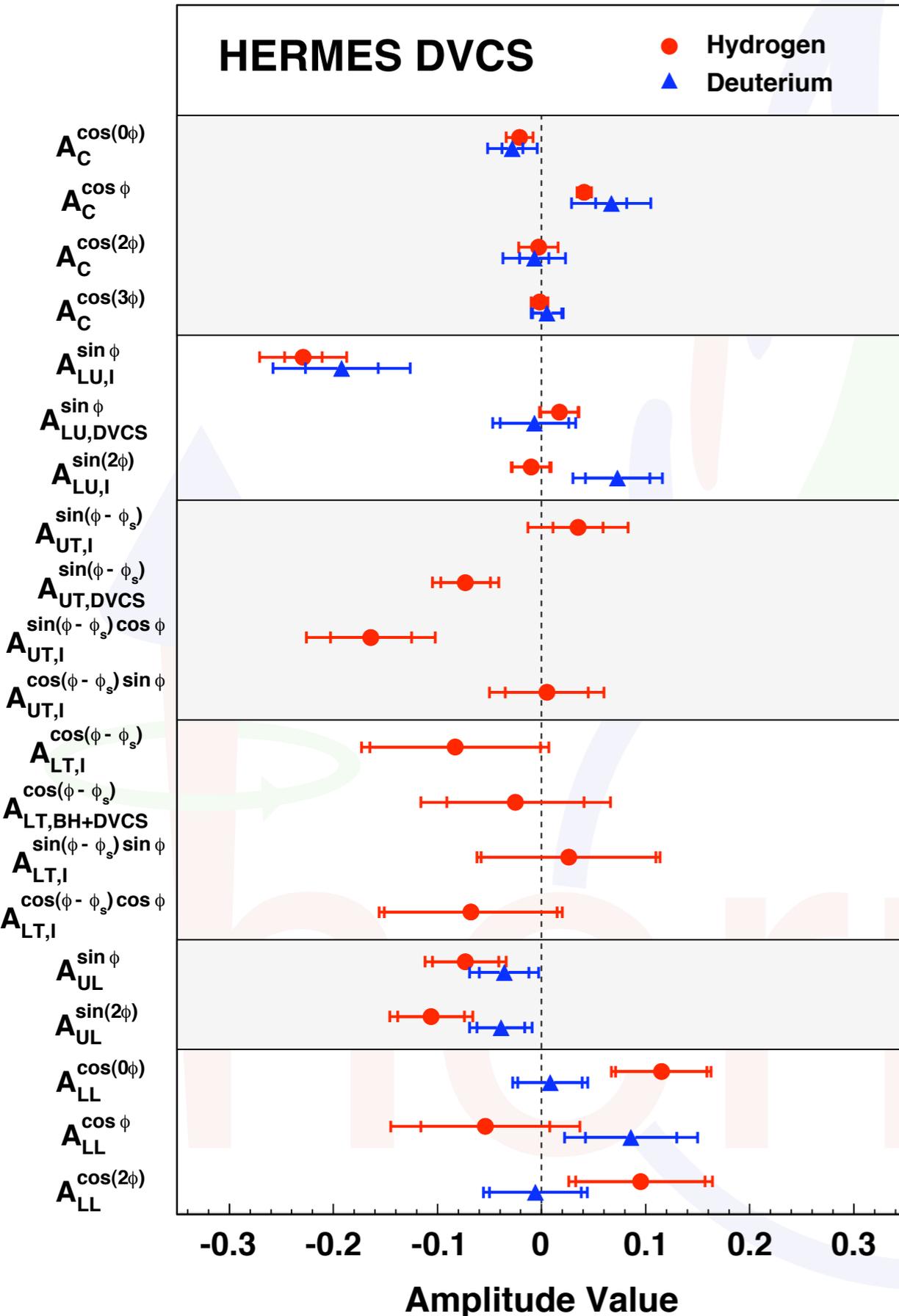


$X=p$

$X=\Delta^+$

$X=\pi^0 + \dots$

A wealth of azimuthal amplitudes



Beam-charge asymmetry:

GPD H

PRD 75 (2007) 011103

NPB 829 (2010) 1

JHEP 11 (2009) 083

PRC 81 (2010) 035202

Beam-helicity asymmetry:

GPD H

PRL 87 (2001) 182001

JHEP 07 (2012) 032

Transverse target spin asymmetries:

GPD E from proton target

JHEP 06 (2008) 066

PLB 704 (2011) 15

Longitudinal target spin asymmetry:

GPD \tilde{H}

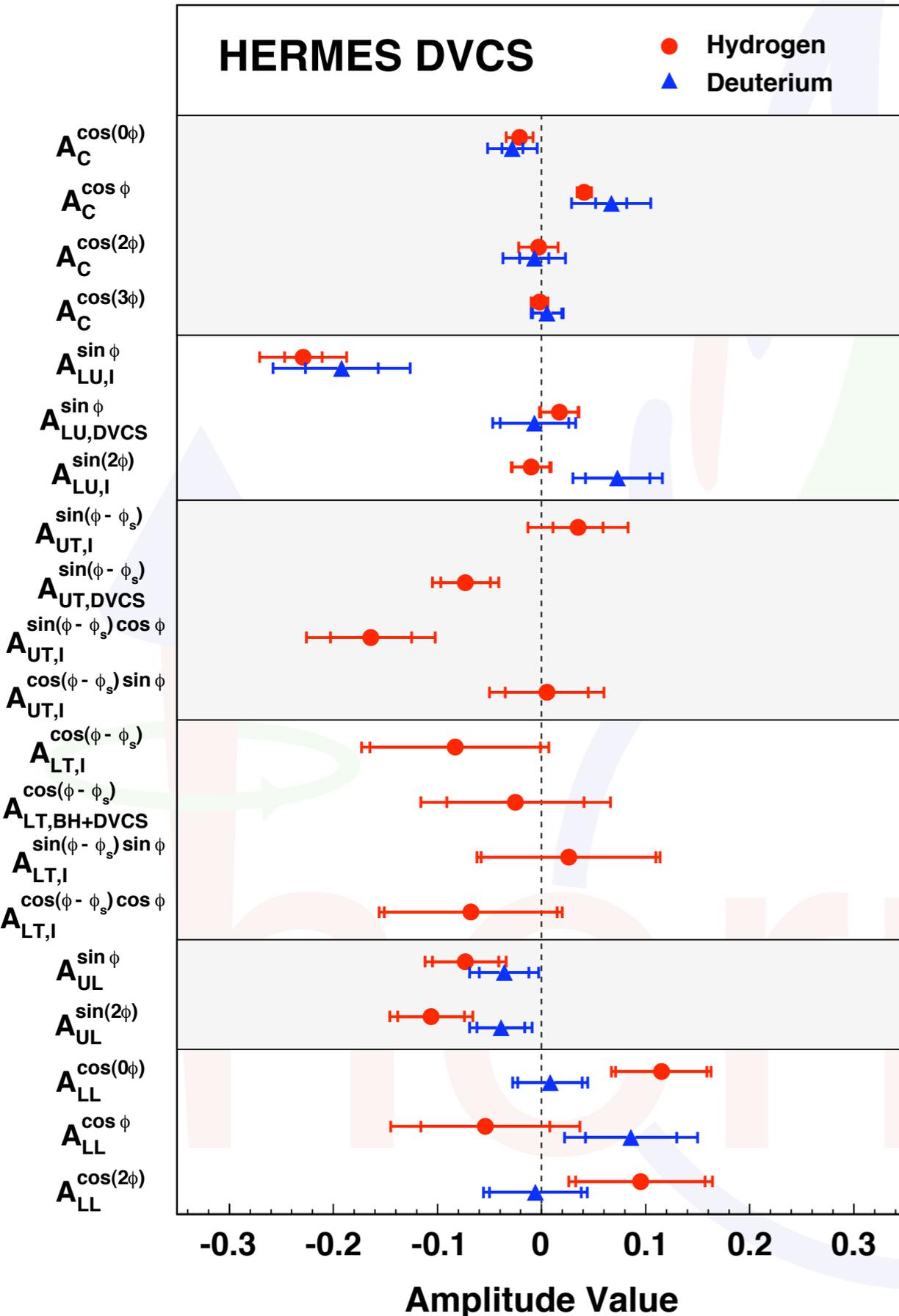
JHEP 06 (2010) 019

Double-spin asymmetry:

GPD \tilde{H}

NPB 842 (2011) 265

A wealth of azimuthal amplitudes



Beam-charge asymmetry:

GPD H

PRD 75 (2007) 011103

NPB 829 (2010) 1

JHEP 11 (2009) 083

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Longitudinal target spin asymmetry:

GPD \tilde{H}

JHEP 06 (2010) 019

Double-spin asymmetry:

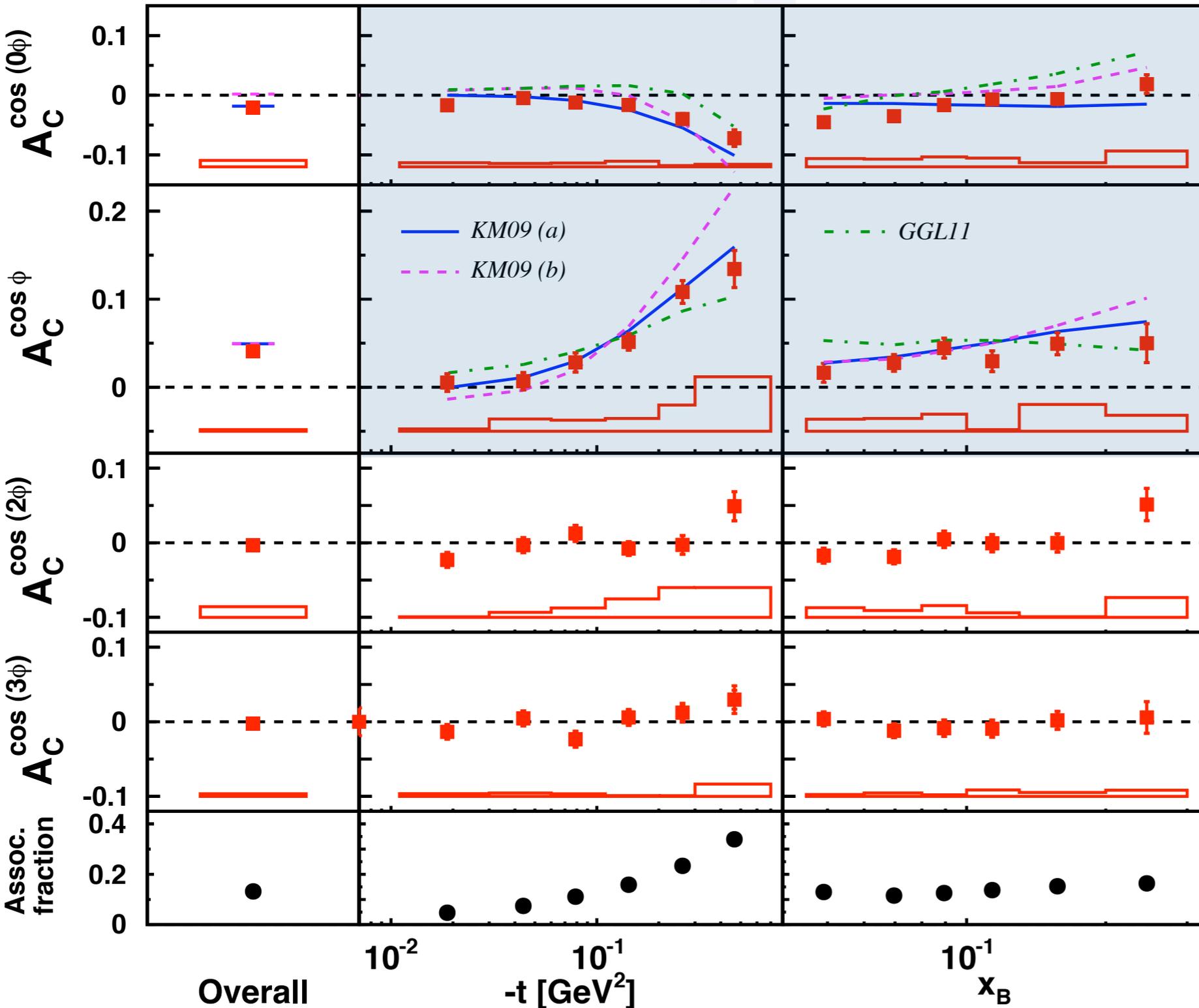
GPD \tilde{H}

NPB 842 (2011) 265

complete data set!

Beam-charge asymmetry

[Airapetian et al., JHEP 07 (2012) 032]



constant term:

$$\propto -A_C^{\cos\phi}$$

$$\propto \text{Re}[F_1 \mathcal{H}]$$

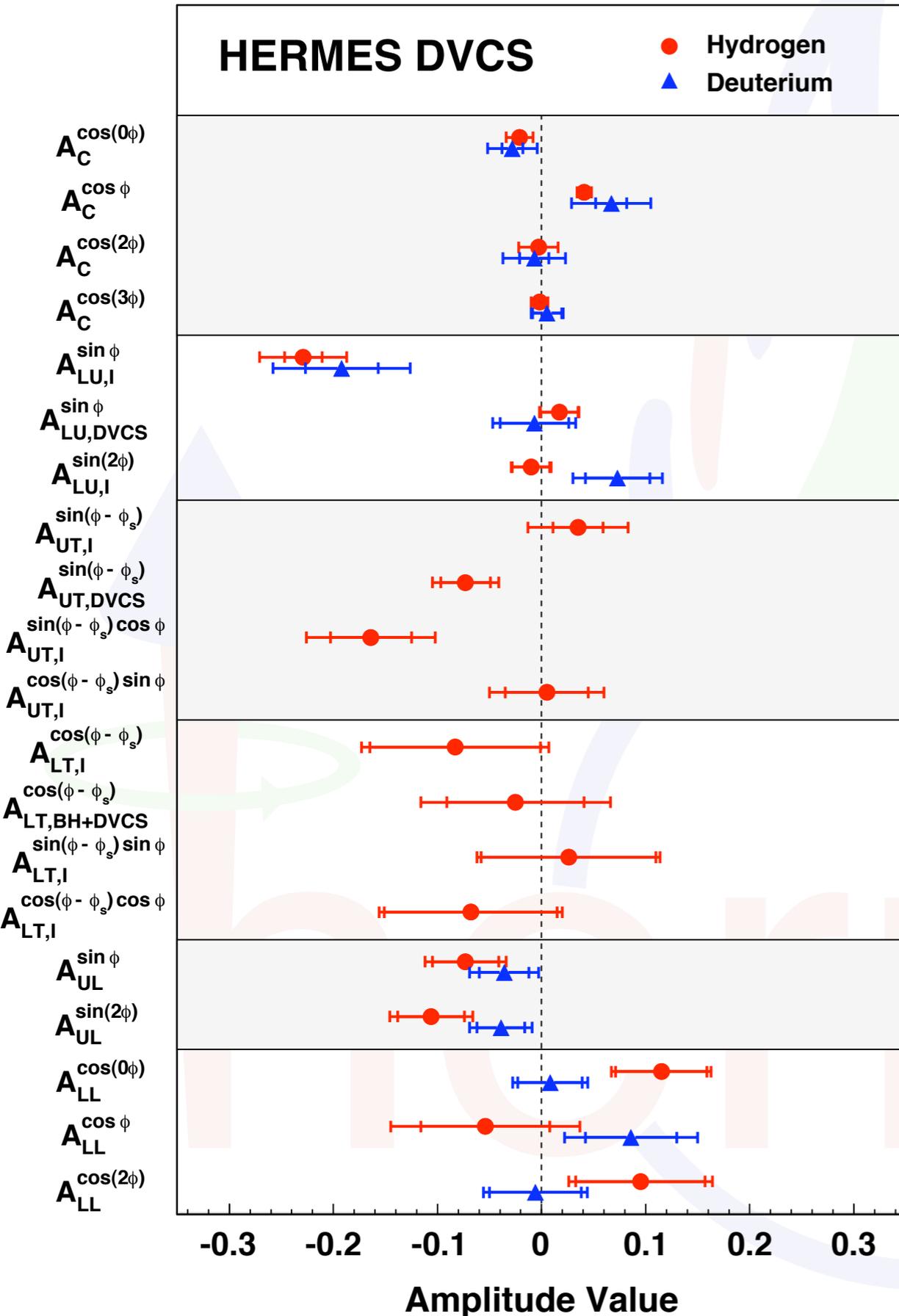
[higher twist]

[gluon leading twist]

Resonant fraction:



A wealth of azimuthal amplitudes



Beam-charge asymmetry:

GPD H

PRD 75 (2007) 011103

NPB 829 (2010) 1

JHEP 11 (2009) 083

PRC 81 (2010) 035202

Beam-helicity asymmetry:

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PLB 704 (2011) 15

Longitudinal target spin asymmetry:

GPD \tilde{H}

JHEP 06 (2010) 019

Double-spin asymmetry:

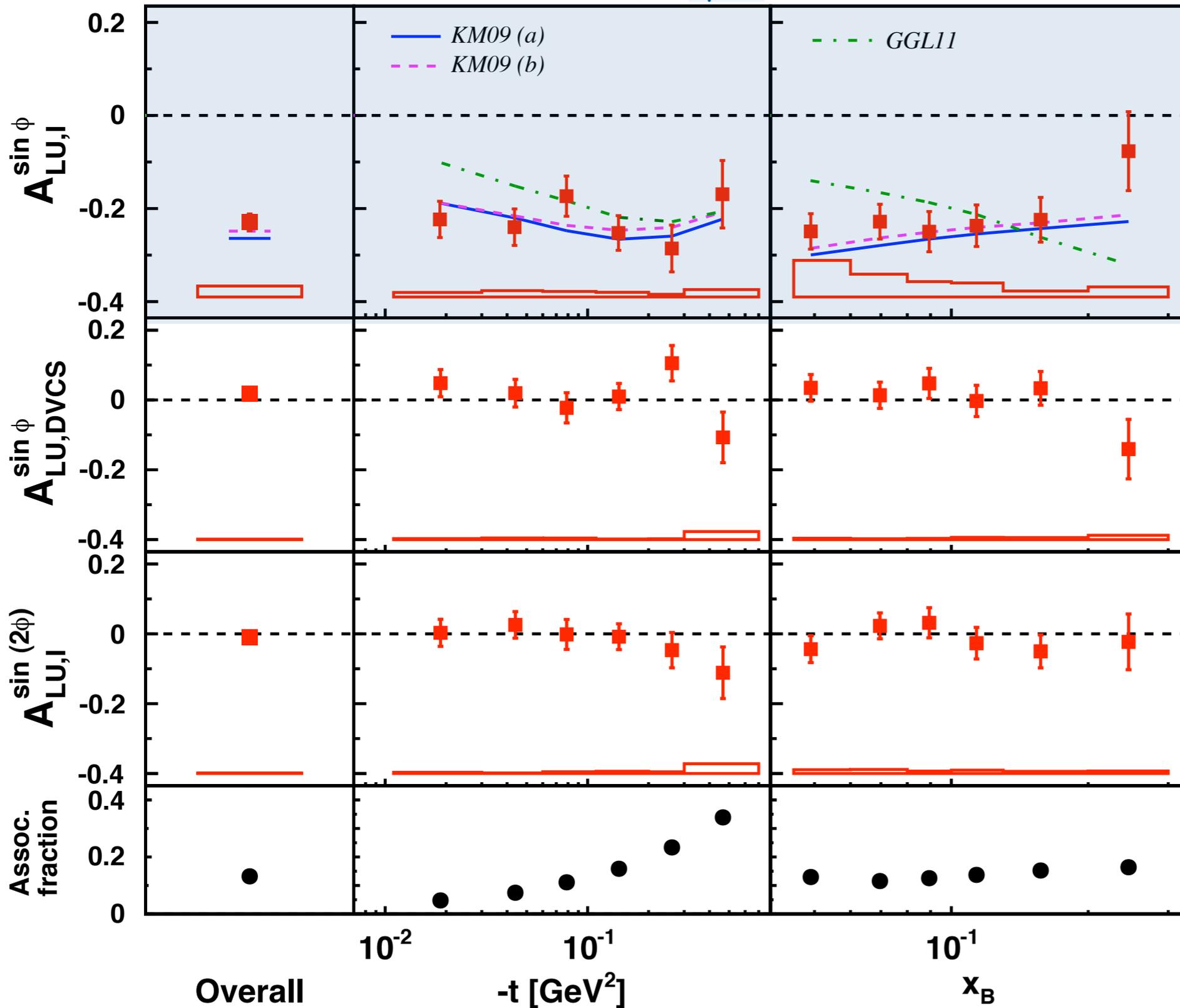
GPD \tilde{H}

NPB 842 (2011) 265

complete data set!

Beam-spin asymmetry

[Airapetian et al., JHEP 07 (2012) 032]



$$\propto \text{Im}[F_1 \mathcal{H}]$$

[higher twist]

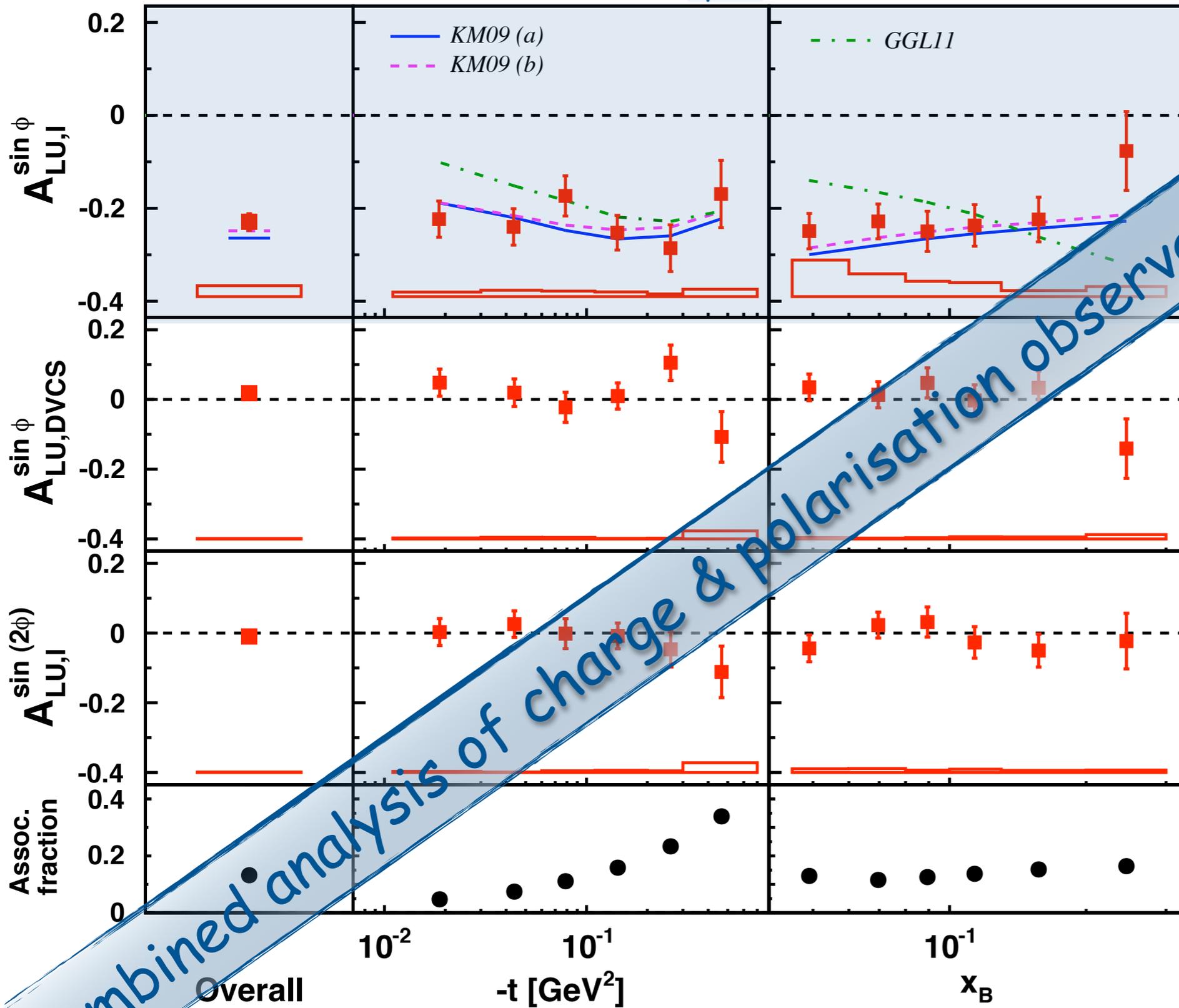
Resonant fraction:



complete data set!

Beam-spin asymmetry

[Airapetian et al., JHEP 07 (2012) 032]



combined analysis of charge & polarisation observables unique to HERA!

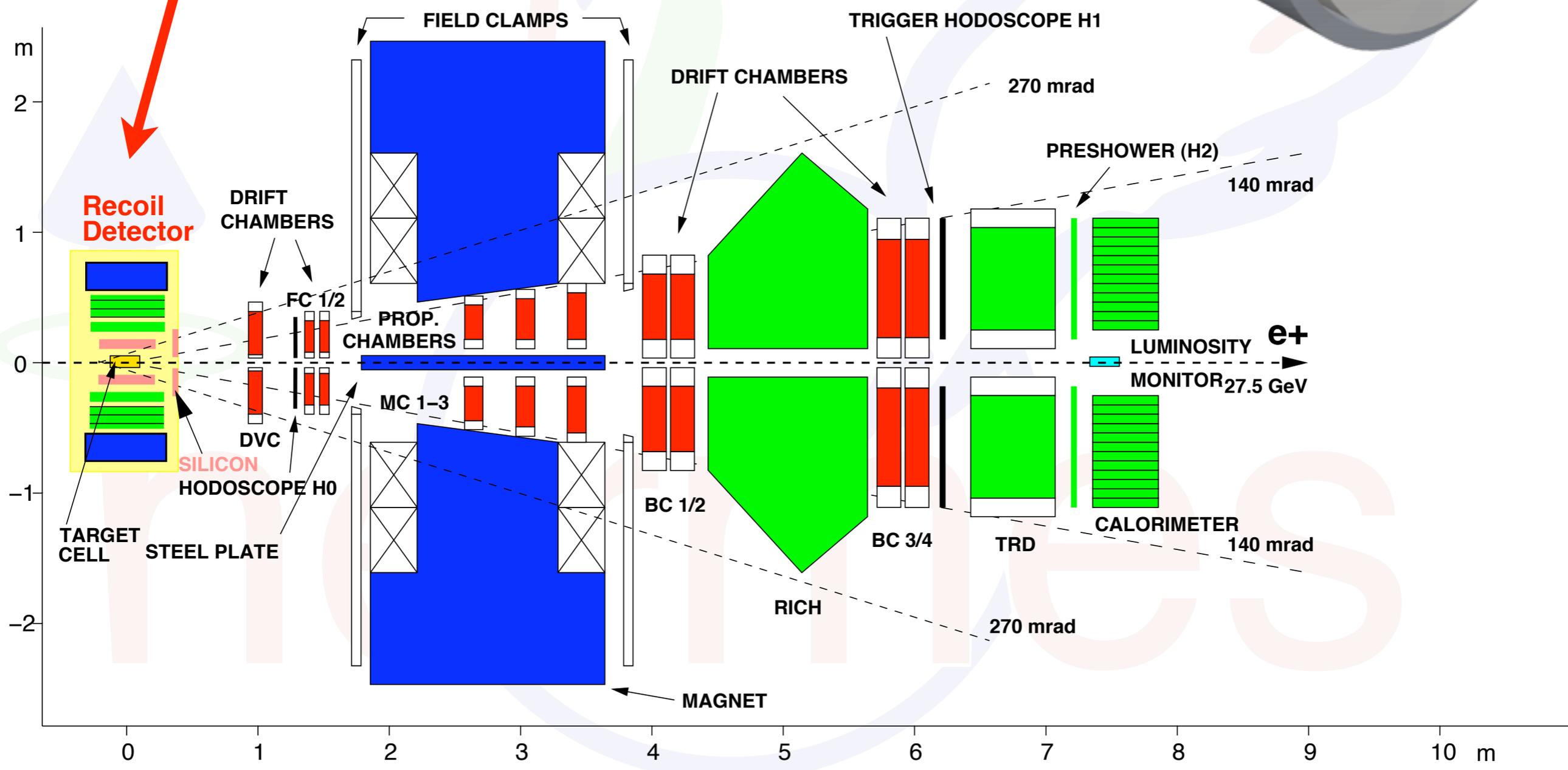
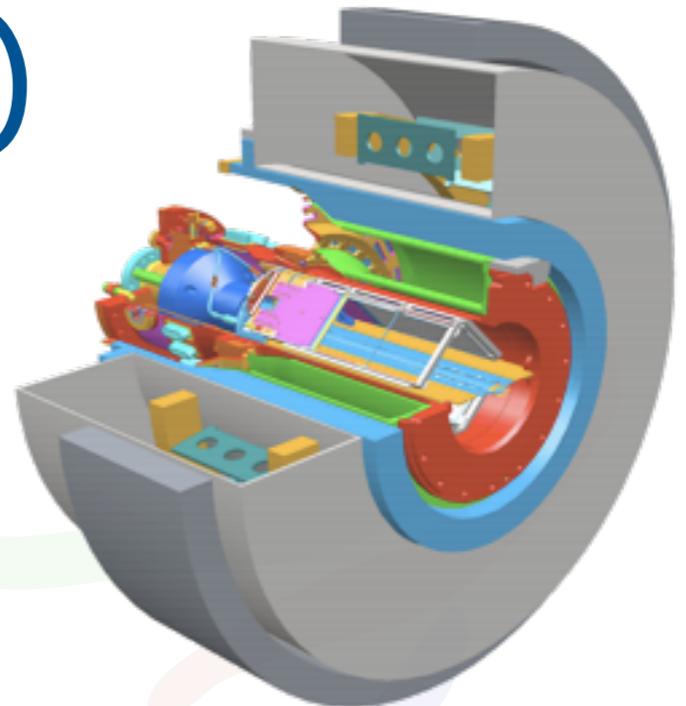
[higher twist]

Resonant fraction:

$$ep \rightarrow e\Delta^+\gamma$$

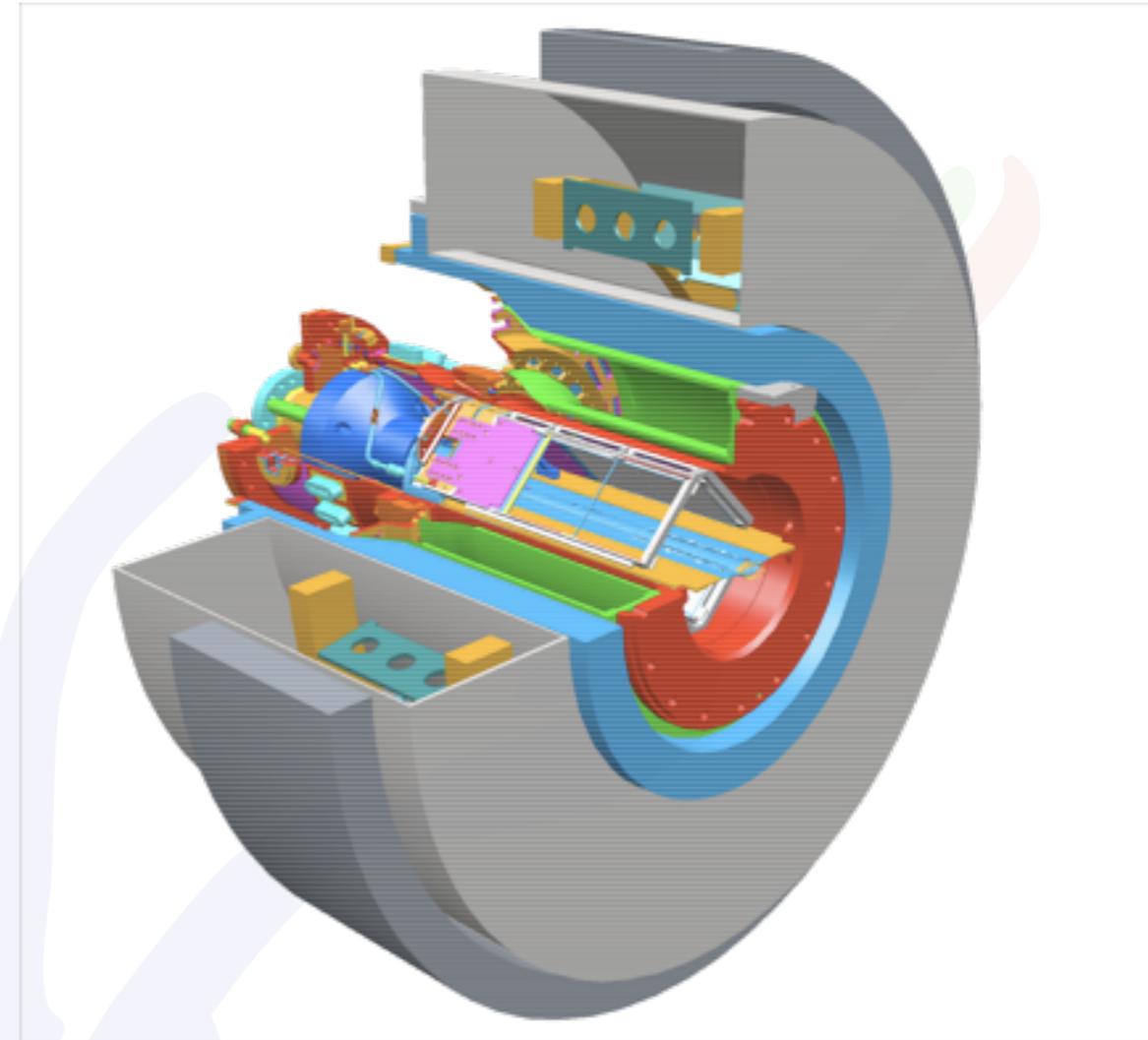
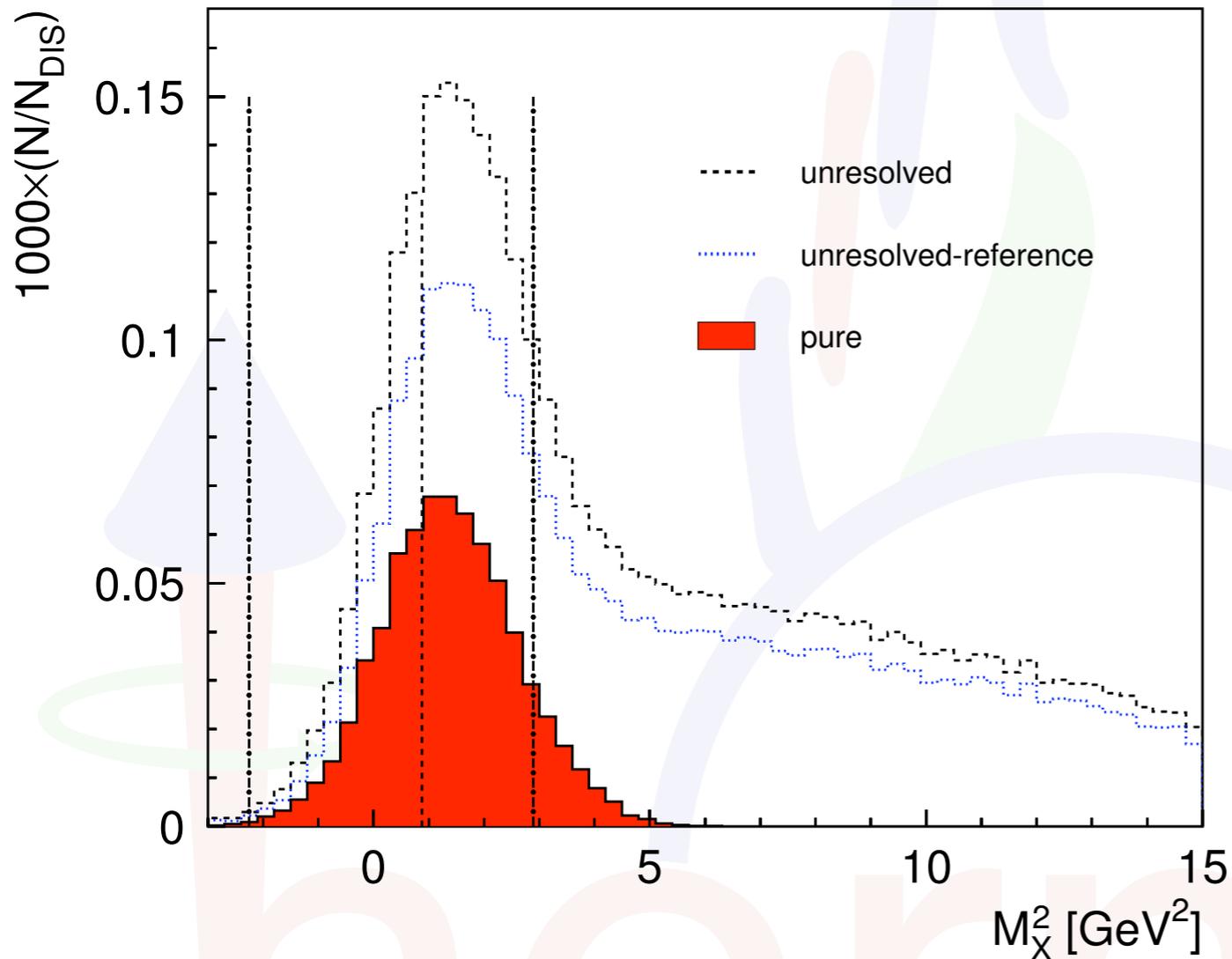
HERMES detector (2006/07)

detection of recoiling proton



HERMES detector (2006/07)

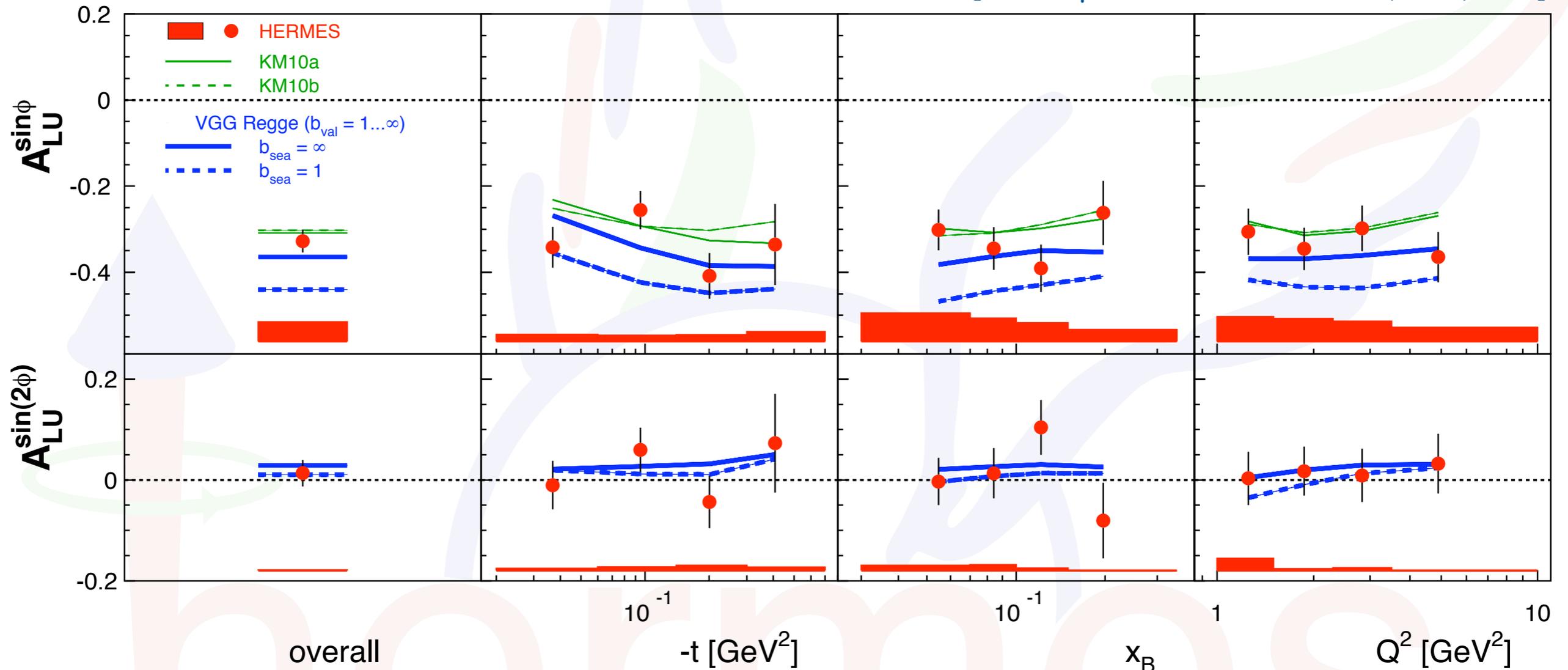
kinematic fitting



- All particles in final state detected → 4 constraints from energy-momentum conservation
- Selection of **pure BH/DVCS** ($ep \rightarrow ep \gamma$) with high efficiency (**~83%**)
- Allows to suppress background from associated and semi-inclusive processes to a negligible level (**<0.2%**)

DVCS with recoil detector

[A. Airapetian et al., JHEP 10 (2012) 042]

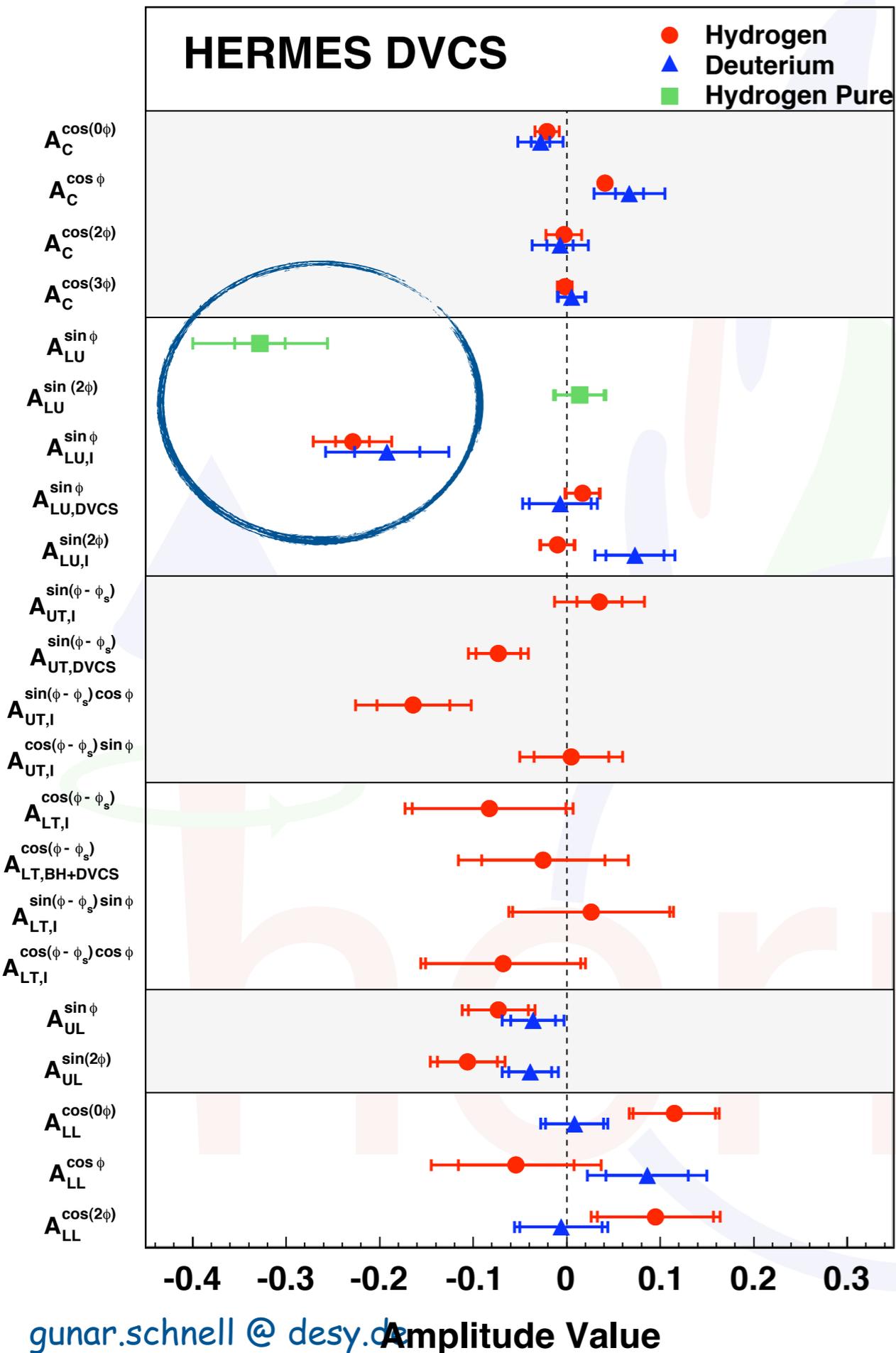


good agreement with models

KM10 - K. Kumericki and D. Müller, Nucl. Phys. B 841 (2010) 1

VGG - M. Vanderhaeghen et al., Phys. Rev. D 60 (1999) 094017

DVCS at HERMES

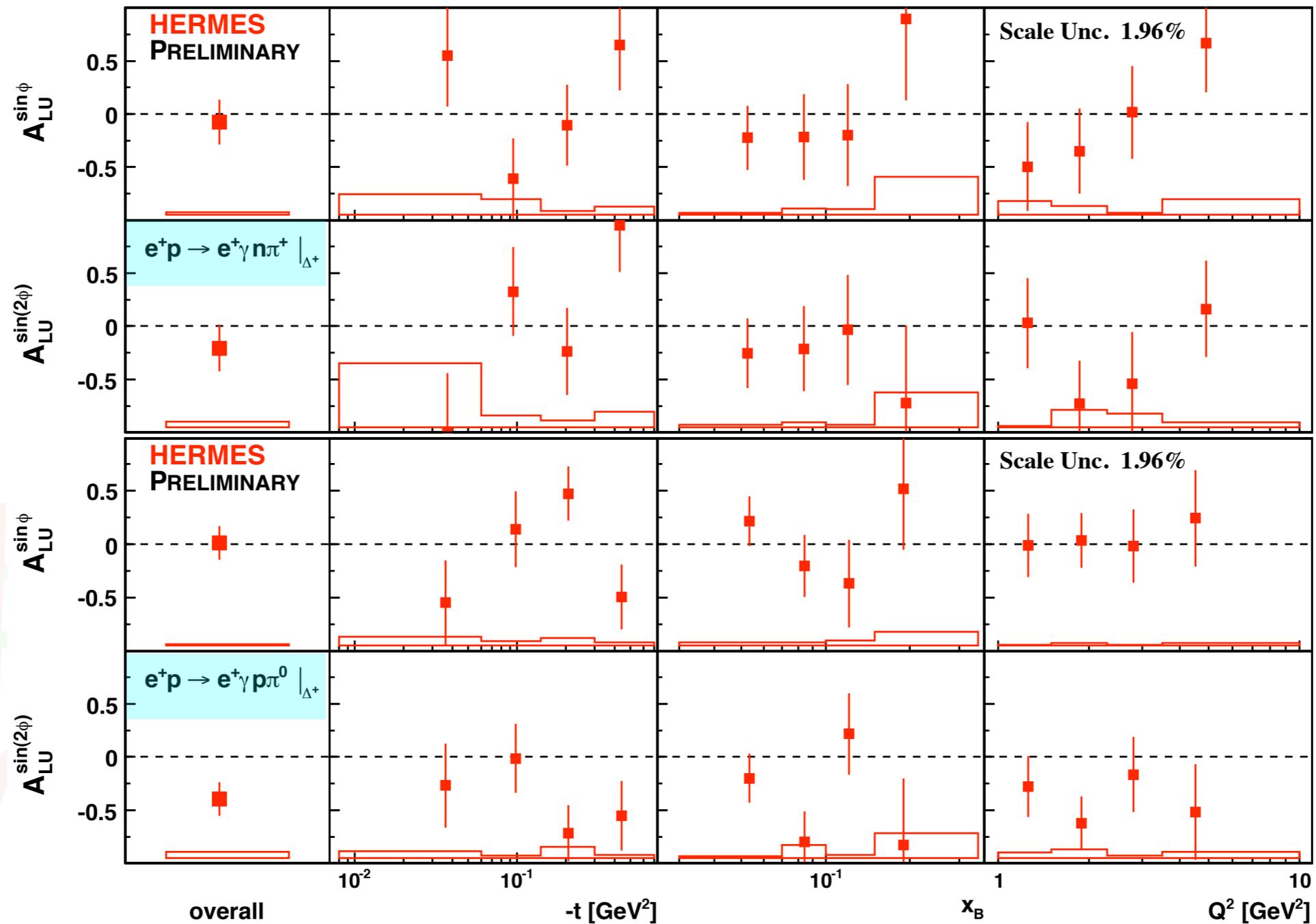


indication of larger amplitudes
for pure sample

(-> assoc. DVCS in "traditional"
analysis mainly dilution)

basically no contamination
-> clean interpretation

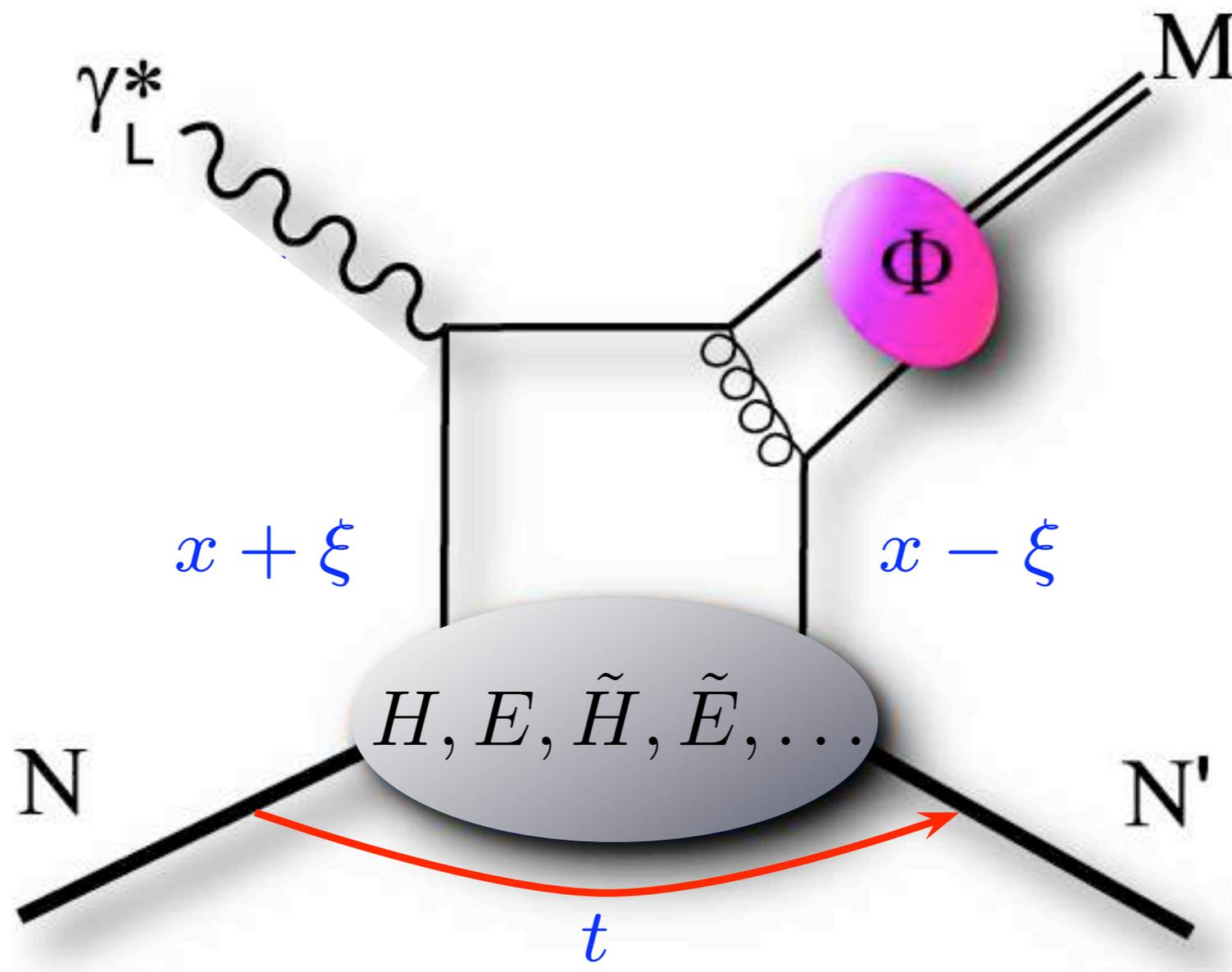
Associated DVCS with recoil detector



- asymmetry amplitudes consistent with zero

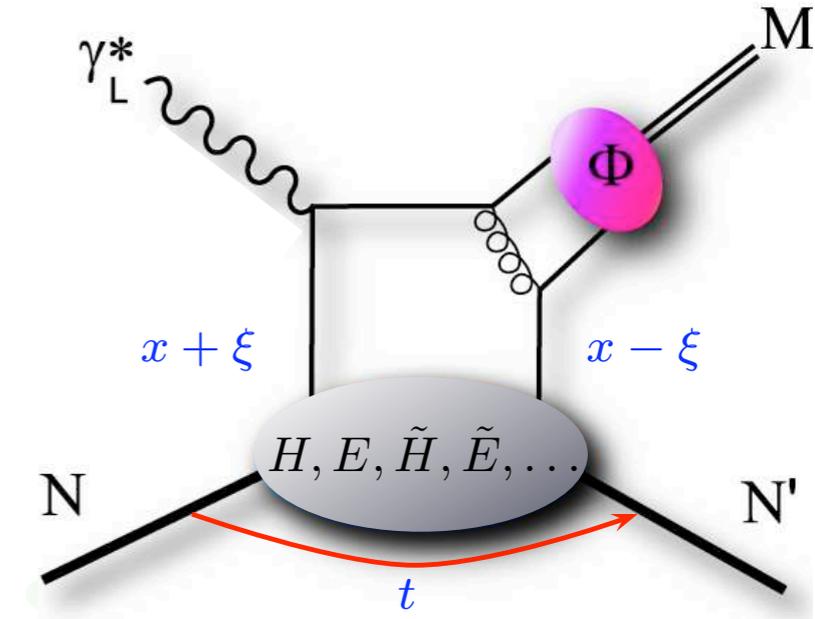
- consistent with pure DVCS results (e.g., dilution in traditional analysis)

Exclusive meson production



Exclusive meson production

- GPDs convoluted with meson amplitude



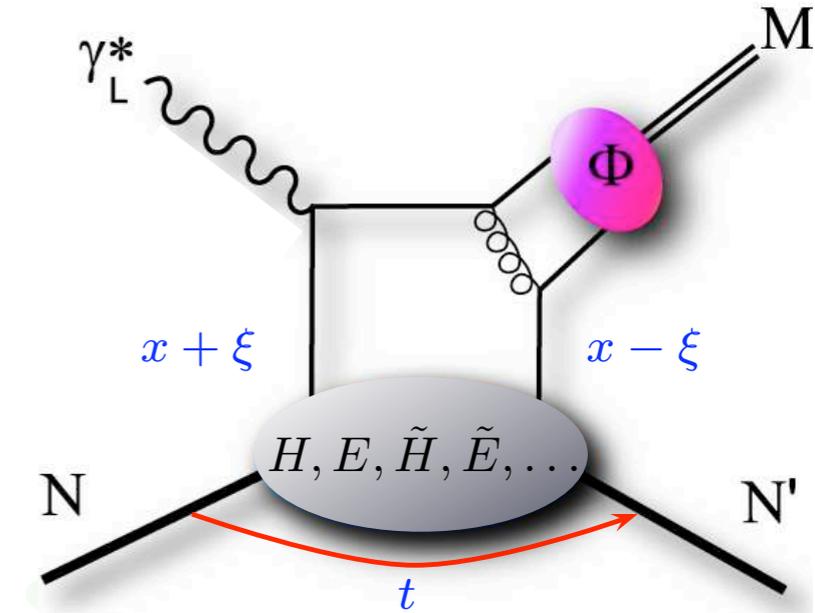
π^0	$2\Delta u + \Delta d$
η	$2\Delta u - \Delta d$
ρ^0	$2u + d, 9g/4$
ω	$2u - d, 3g/4$
ϕ	s, g
ρ^+	$u - d$
J/ψ	g



HERMES

Exclusive meson production

- GPDs convoluted with meson amplitude
- access to various quark-flavor combinations



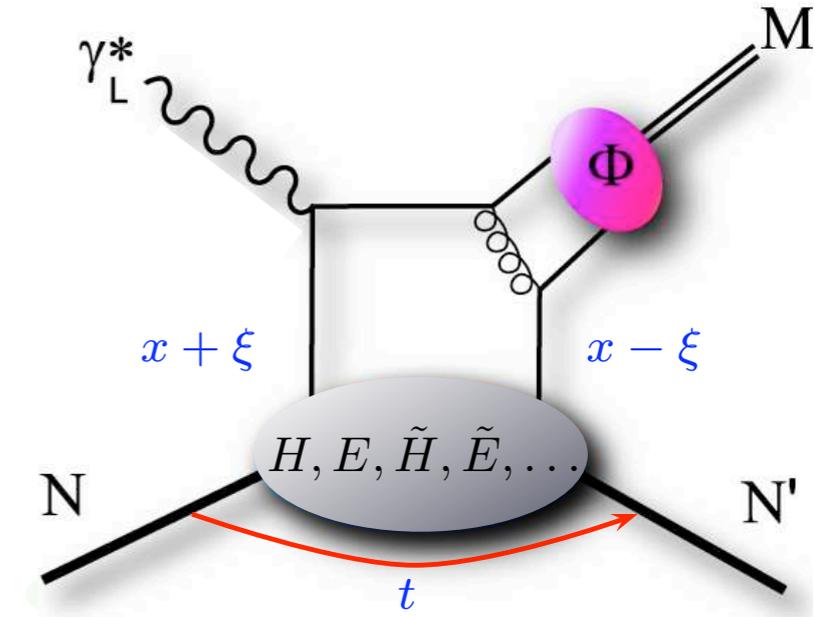
π^0	$2\Delta u + \Delta d$
η	$2\Delta u - \Delta d$
ρ^0	$2u + d, 9g/4$
ω	$2u - d, 3g/4$
ϕ	s, g
ρ^+	$u - d$
J/ψ	g



HERMES

Exclusive meson production

- GPDs convoluted with meson amplitude
- access to various quark-flavor combinations
- factorization proven for longitudinal photons

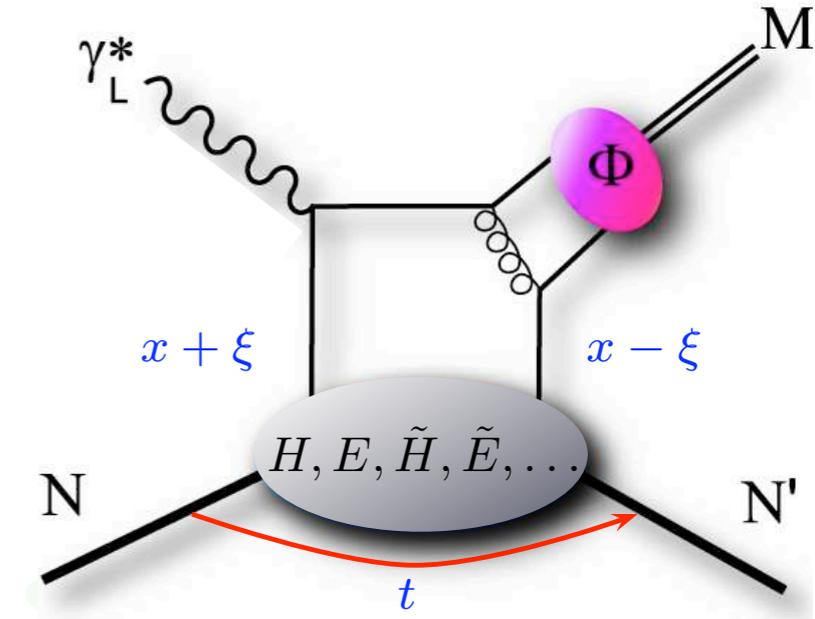


π^0	$2\Delta u + \Delta d$
η	$2\Delta u - \Delta d$
ρ^0	$2u + d, 9g/4$
ω	$2u - d, 3g/4$
ϕ	s, g
ρ^+	$u - d$
J/ψ	g



Exclusive meson production

- GPDs convoluted with meson amplitude
- access to various quark-flavor combinations
- factorization proven for longitudinal photons
- vector-meson cross section:



π^0	$2\Delta u + \Delta d$
η	$2\Delta u - \Delta d$
ρ^0	$2u + d, 9g/4$
ω	$2u - d, 3g/4$
ϕ	s, g
ρ^+	$u - d$
J/ψ	g

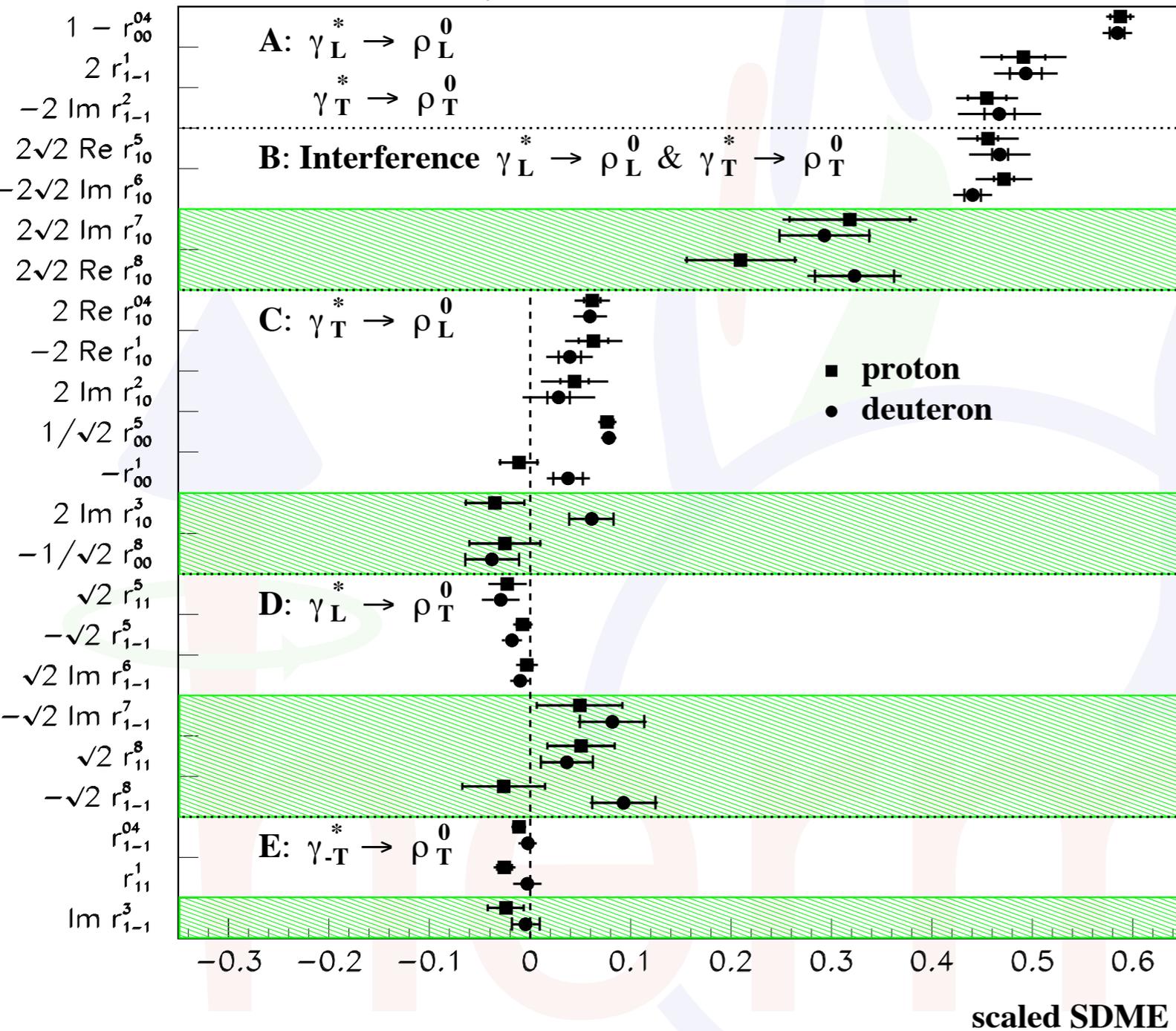
$$\frac{d\sigma}{dx_B dQ^2 dt d\phi_S d\phi d\cos\theta d\varphi} = \frac{d\sigma}{dx_B dQ^2 dt} W(x_B, Q^2, t, \phi_S, \phi, \cos\theta, \varphi)$$

$$W = W_{UU} + P_B W_{LU} + S_L W_{UL} + P_B S_L W_{LL} + S_T W_{UT} + P_B S_T W_{LT}$$

look at various angular (decay) distributions to study helicity transitions ("spin-density matrix elements")

ρ^0 SDMEs from HERMES

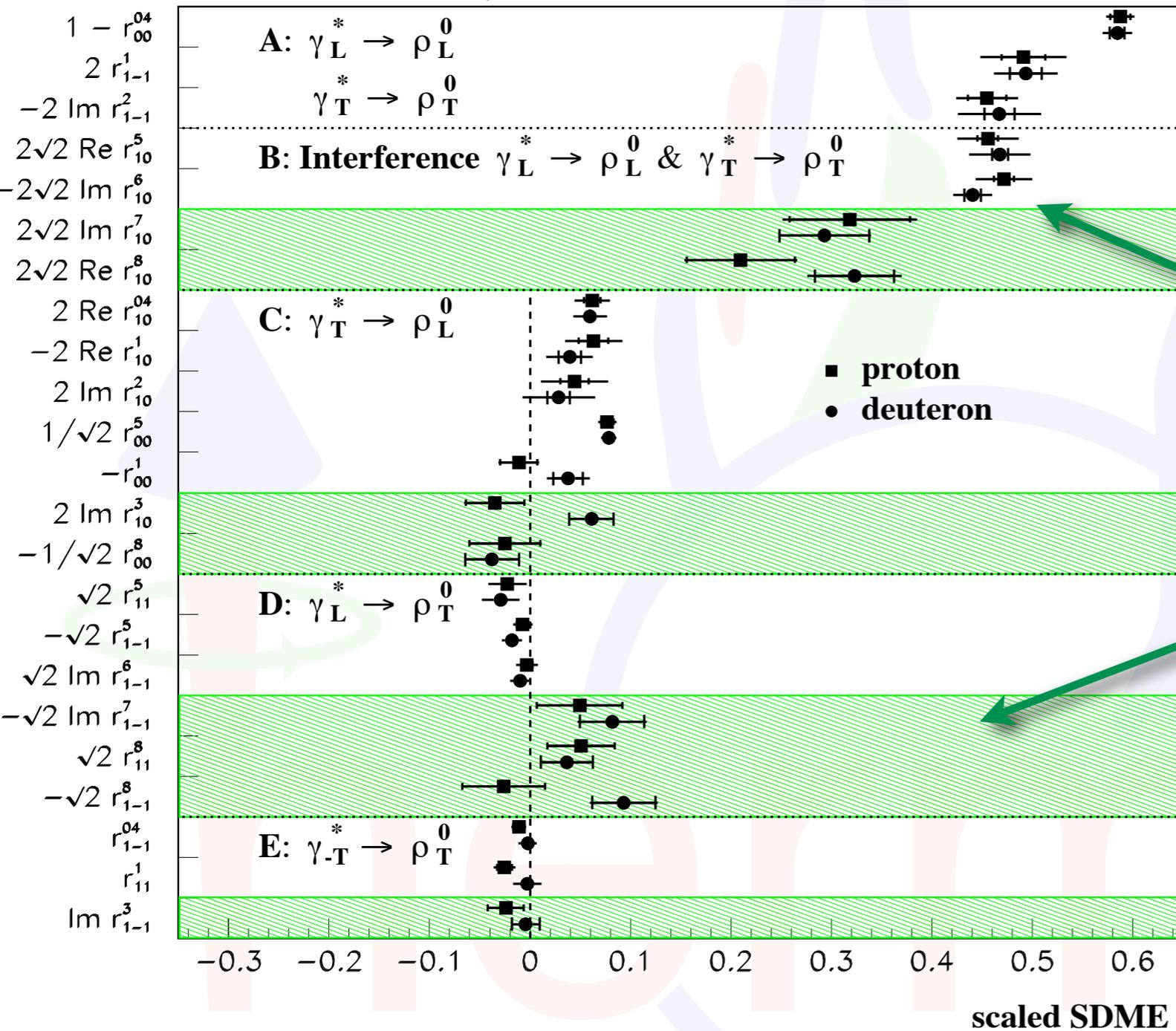
[A. Airapetian et al., EPJ C62 (2009) 659]



target-polarization independent SDMEs

ρ^0 SDMEs from HERMES

[A. Airapetian et al., EPJ C62 (2009) 659]

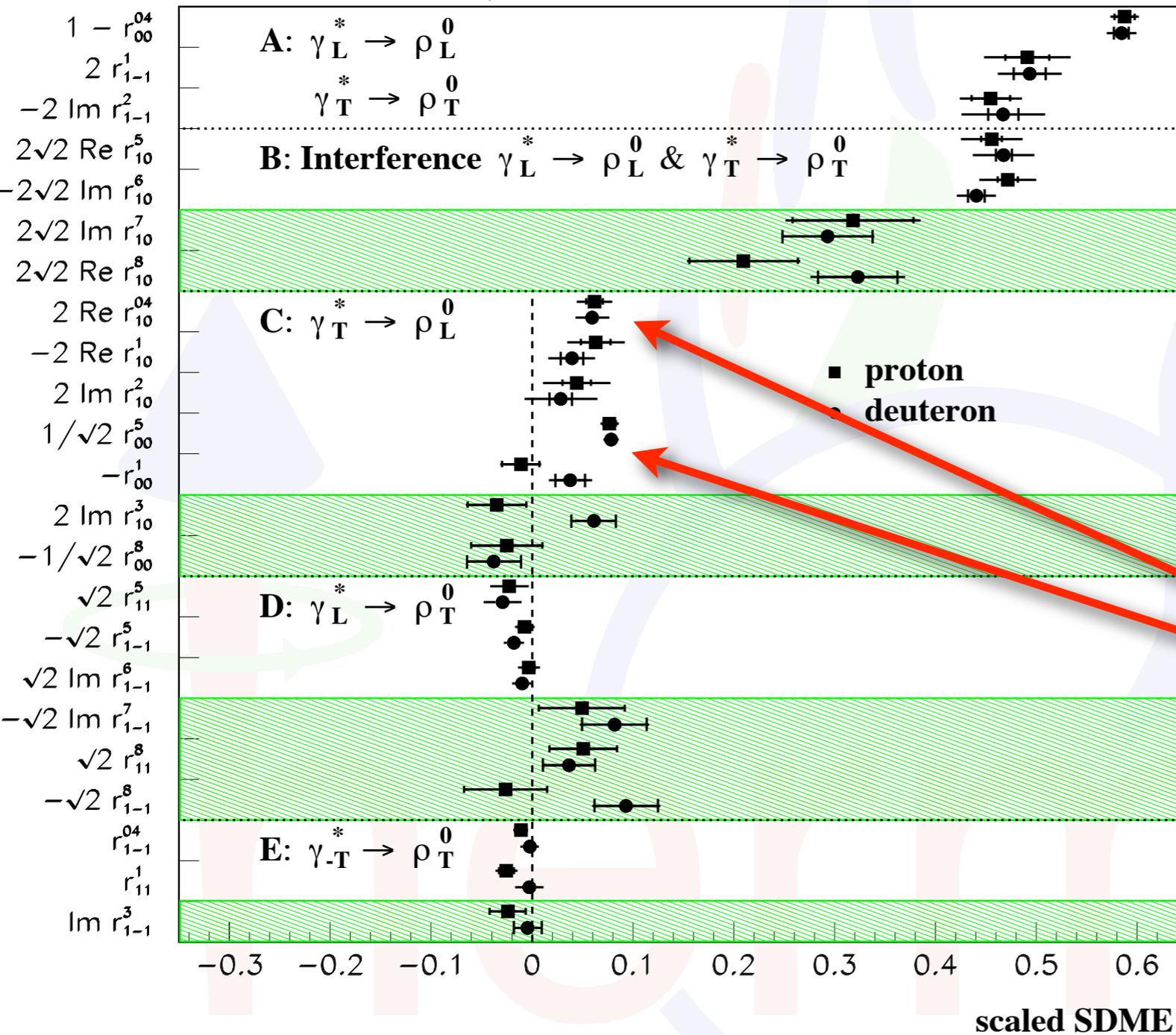


helicity non-flip much larger than helicity-flip and double helicity-flip

target-polarization independent SDMEs

ρ^0 SDMEs from HERMES

[A. Airapetian et al., EPJ C62 (2009) 659]



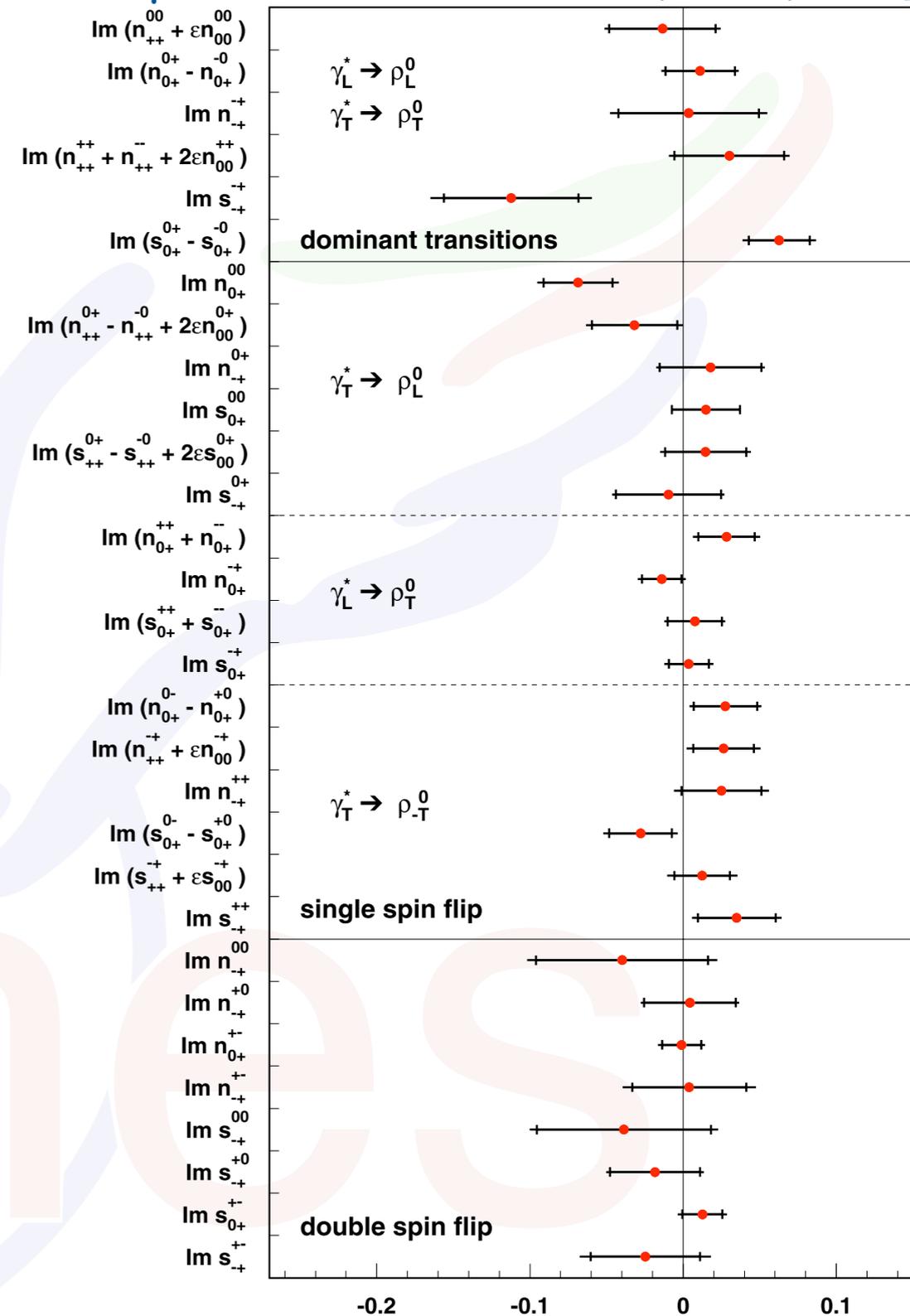
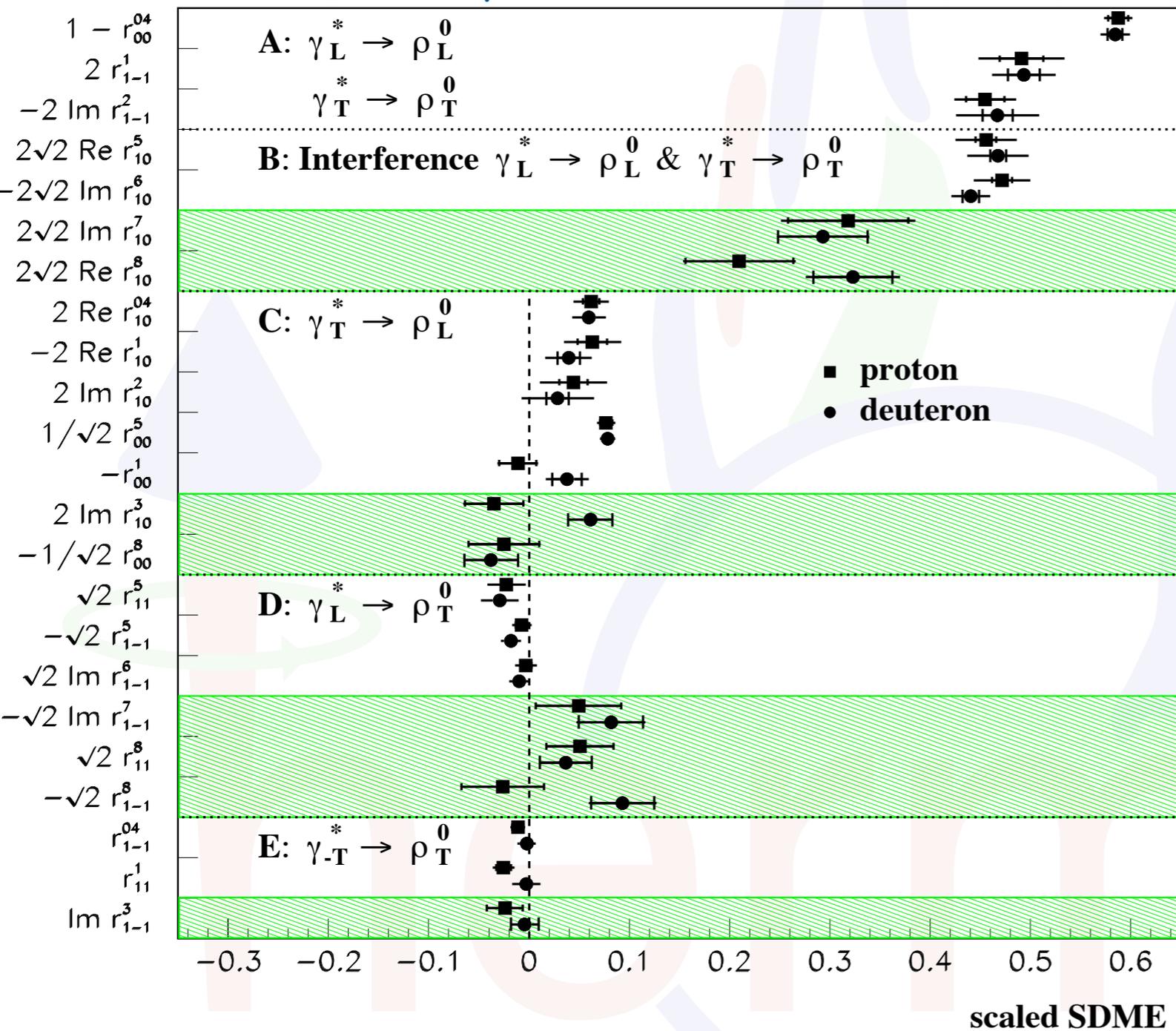
clear breaking of s-channel
helicity conservation

target-polarization independent SDMEs

ρ^0 SDMEs from HERMES

[A. Airapetian et al., PLB 679 (2009) 100]

[A. Airapetian et al., EPJ C62 (2009) 659]



target-polarization independent SDMEs

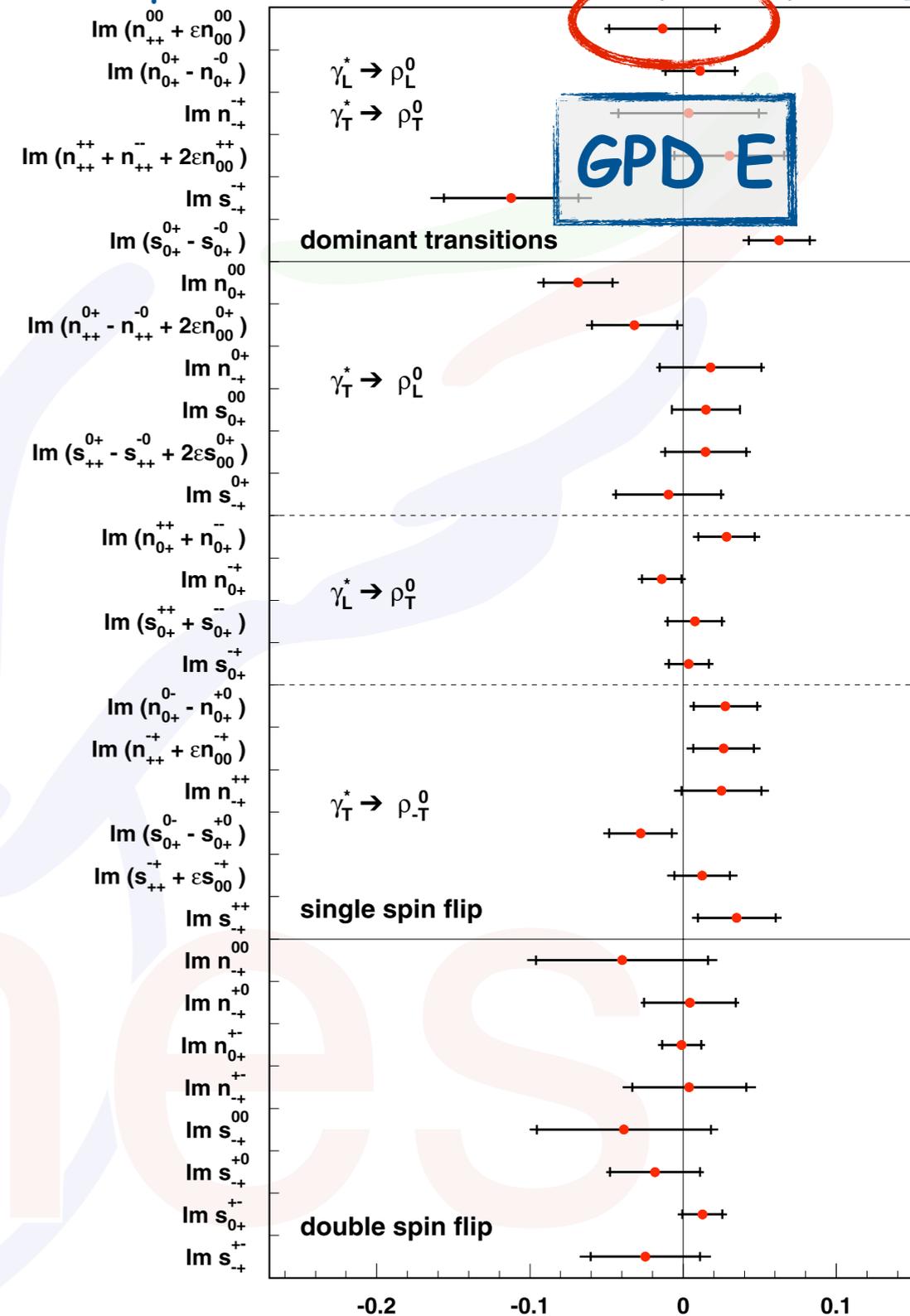
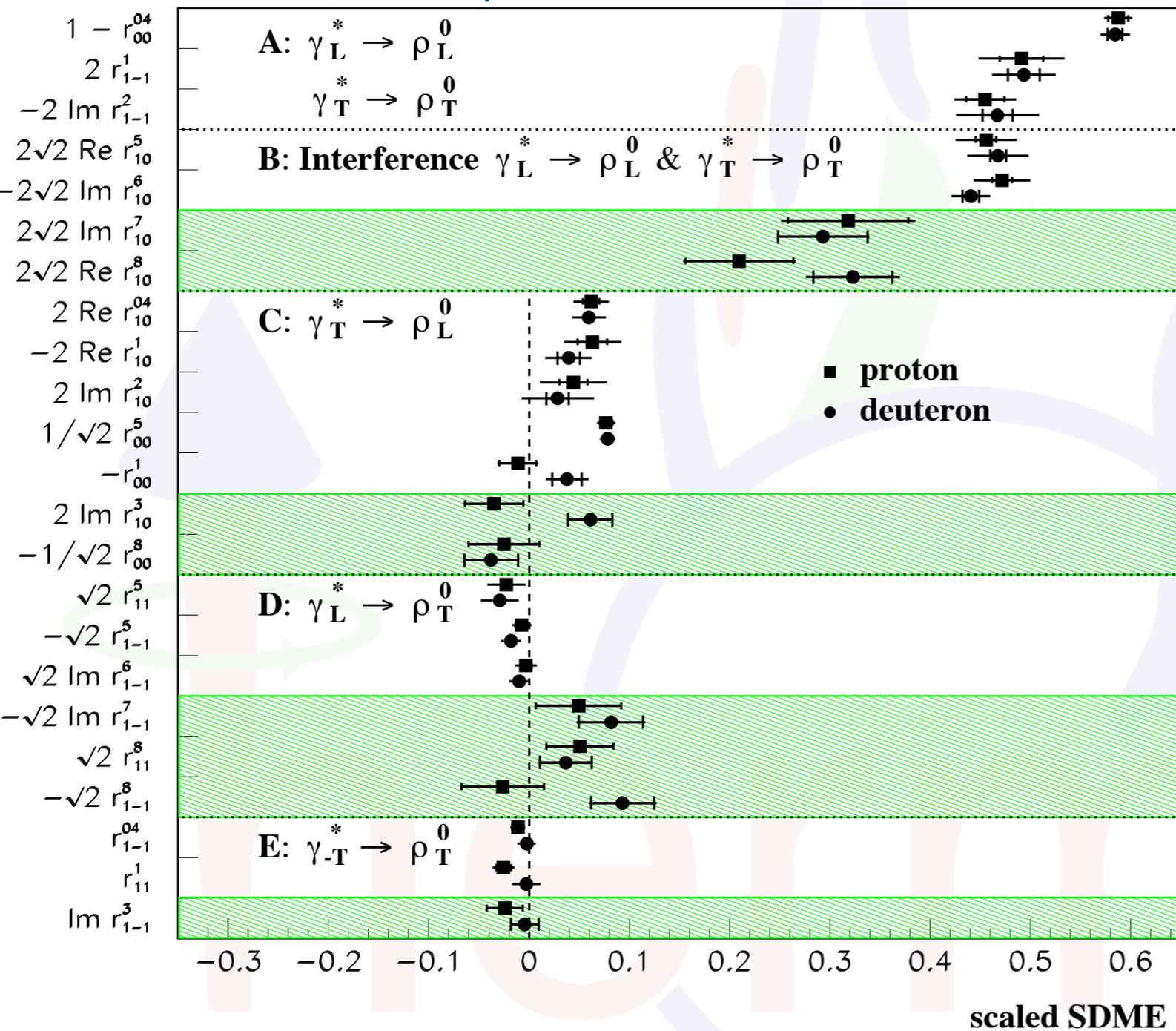
"transverse" SDMEs

POETIC 2013 - March 5th, 2013

ρ^0 SDMEs from HERMES

[A. Airapetian et al., PLB 679 (2009) 100]

[A. Airapetian et al., EPJ C62 (2009) 659]



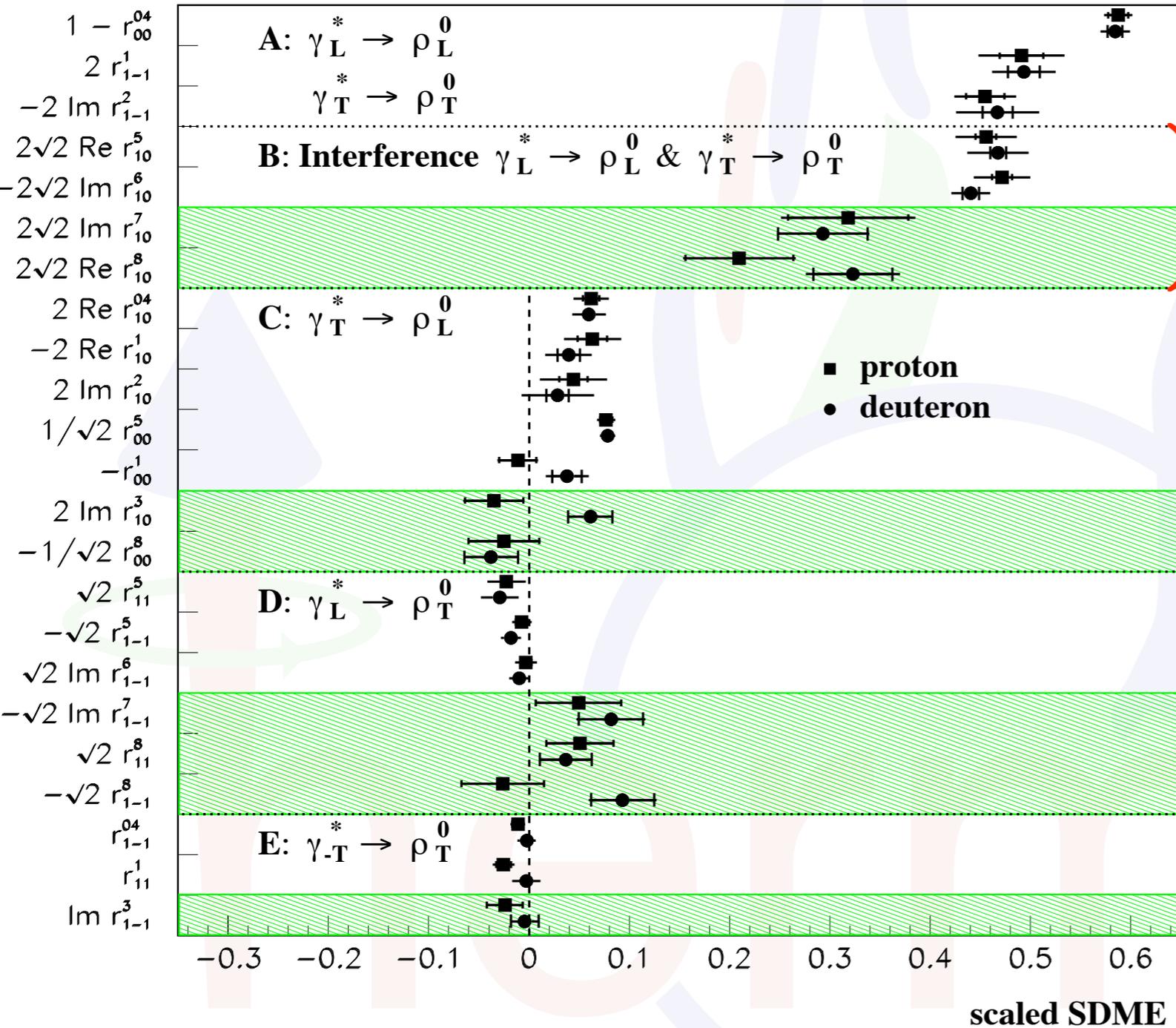
target-polarization independent SDMEs

"transverse" SDMEs

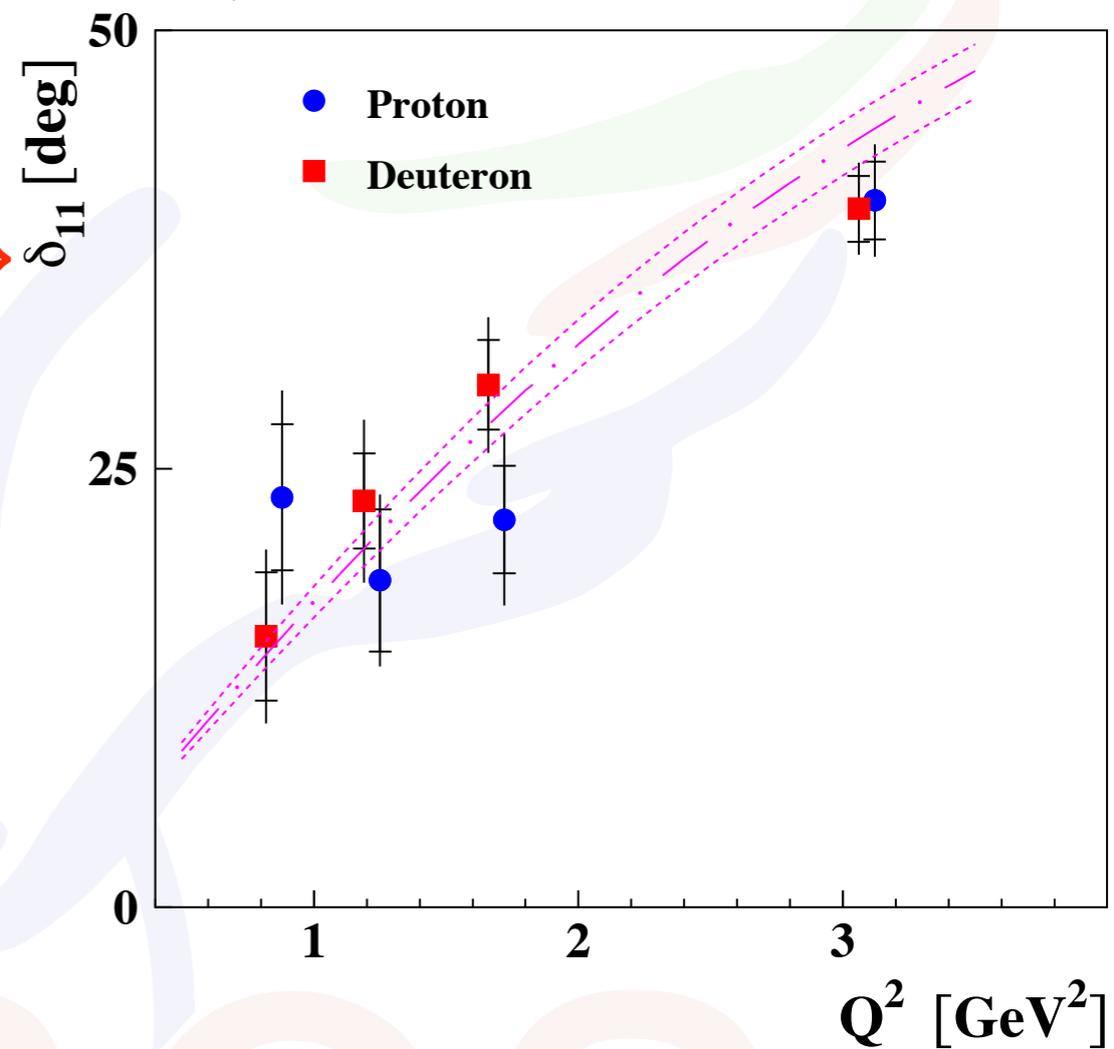
POETIC 2013 - March 5th, 2013

ρ^0 SDMEs from HERMES: challenges

[A. Airapetian et al., EPJ C62 (2009) 659]

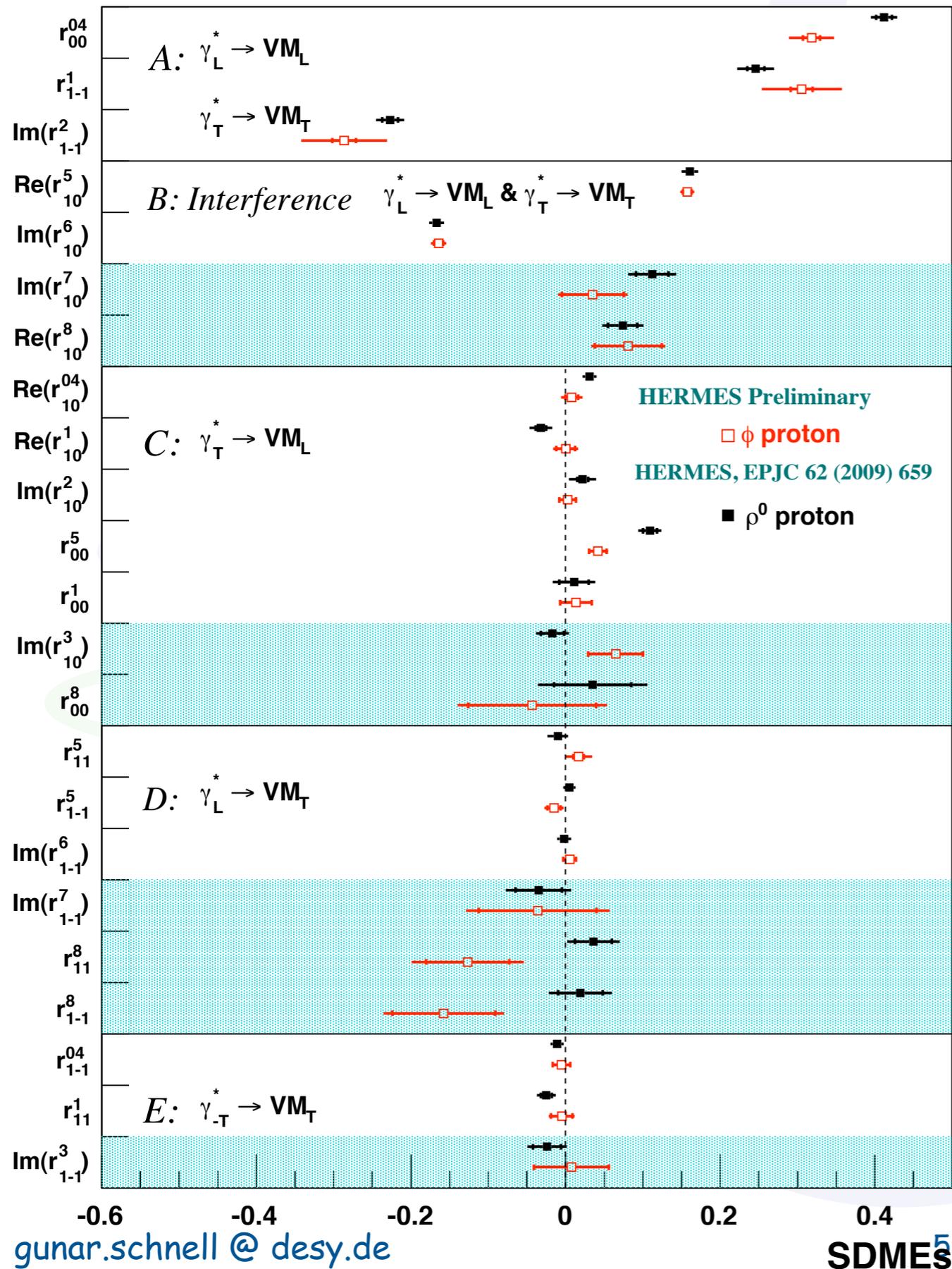


[A. Airapetian et al, EPJ C71 (2011) 1609]



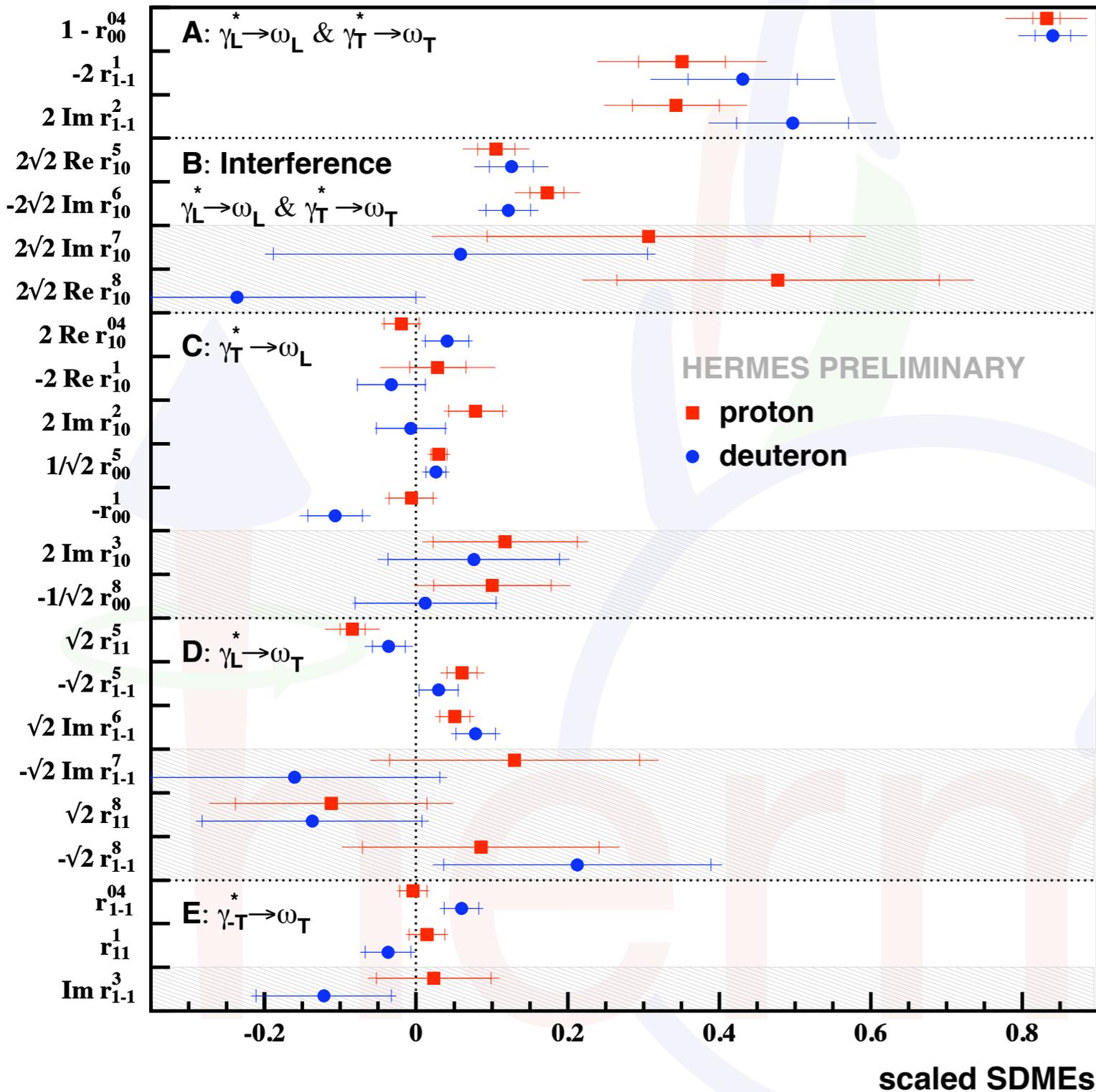
Extraction of SDMEs and helicity amplitude ratios at HERMES for ρ mesons challenges GPD-based calculations (giving small values)

ρ^0 and ϕ production



- similar for both phi and rho production: helicity-conserving SDMEs dominate
- hardly any violation of SCHC observed for phi while some indication for rho production
- also here: large interference effects

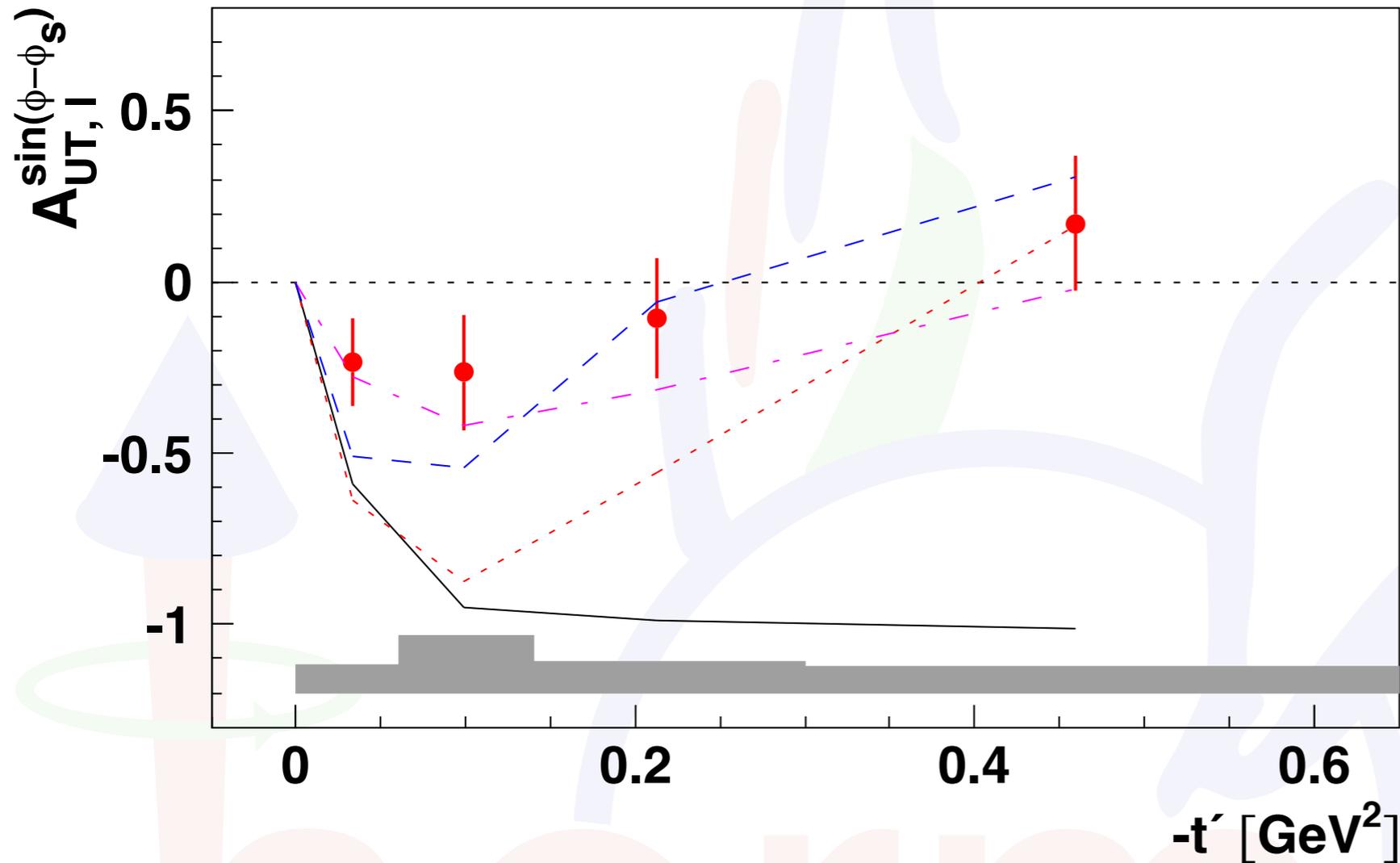
... ω production



- helicity-conserving SDMEs dominate
- hardly any violation of SCHC
- interference smaller than for phi and rho

Exclusive π^+ production

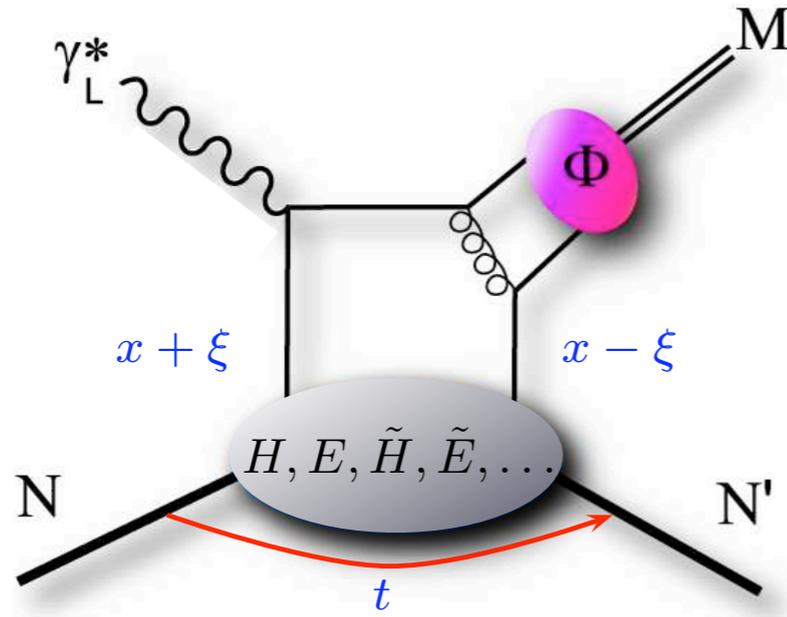
[A. Airapetian et al, PLB 682 (2010) 345]



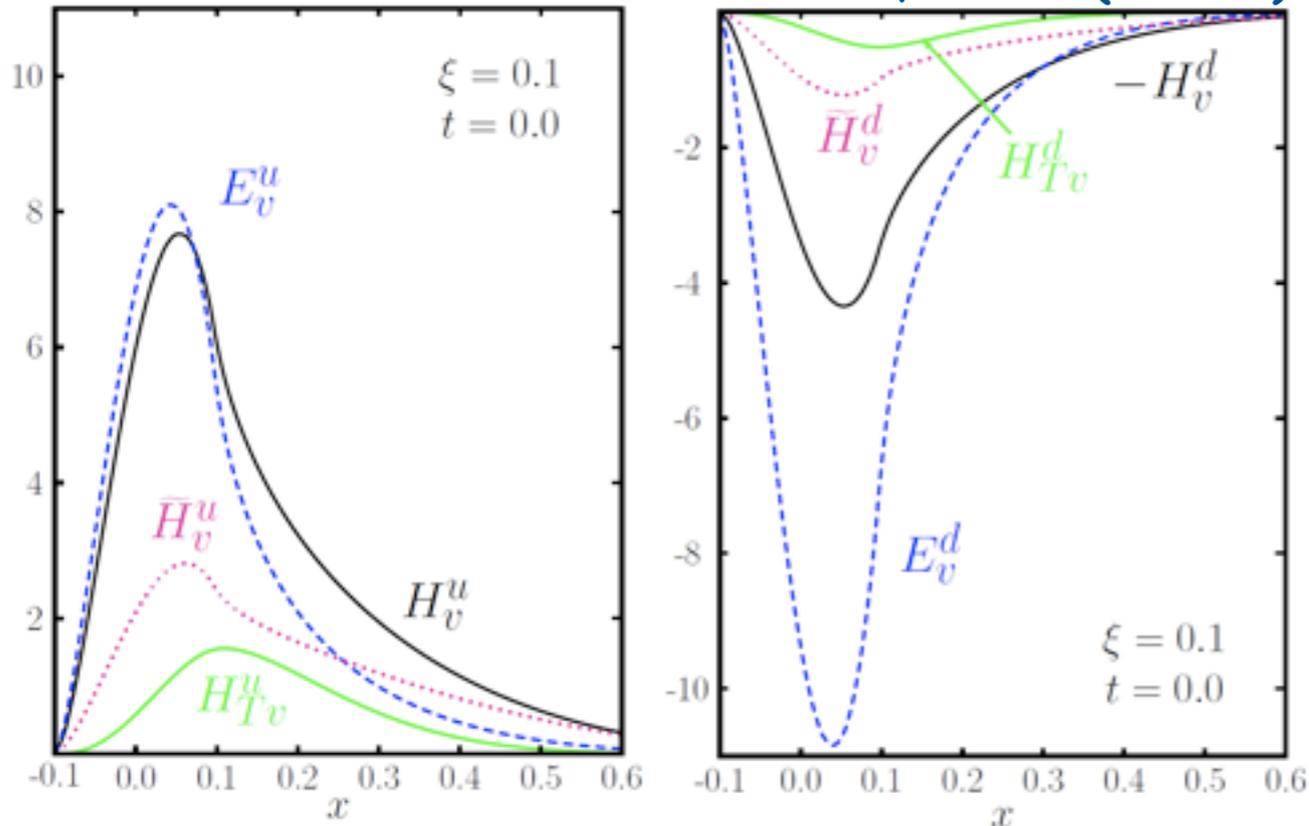
transverse target spin
asymmetry

meson data can also play a vital
role in accessing the "polarised"
GPDs, e.g., \tilde{H} and \tilde{E}

GPDs - a nice success story!

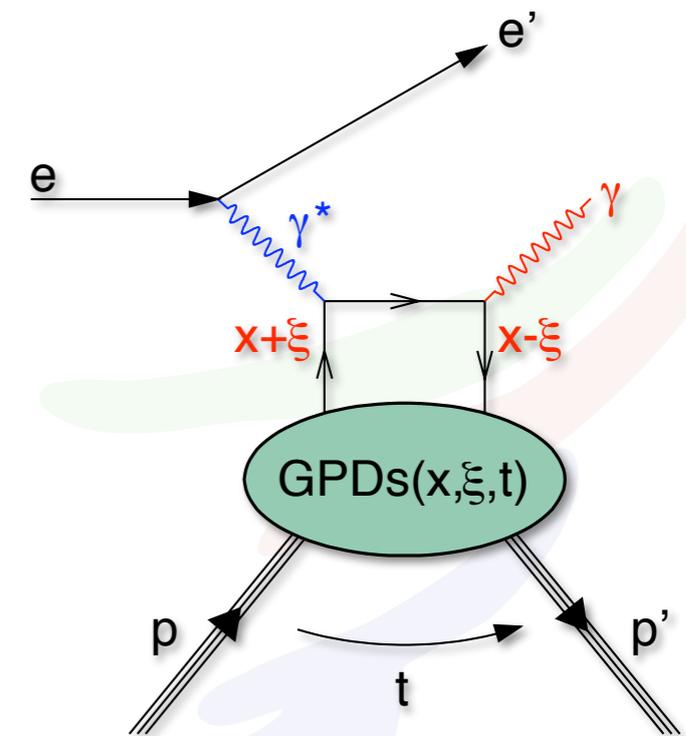
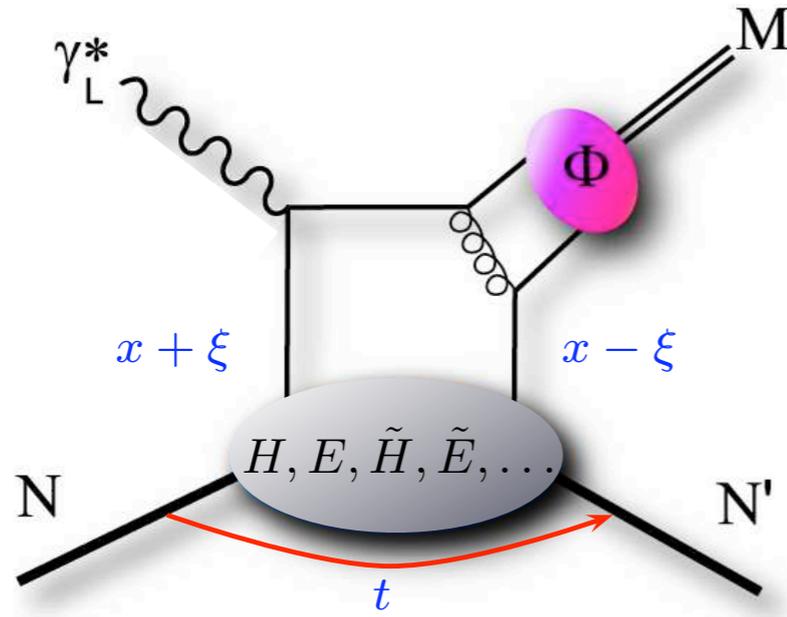


Goloskokov, Kroll (2007)

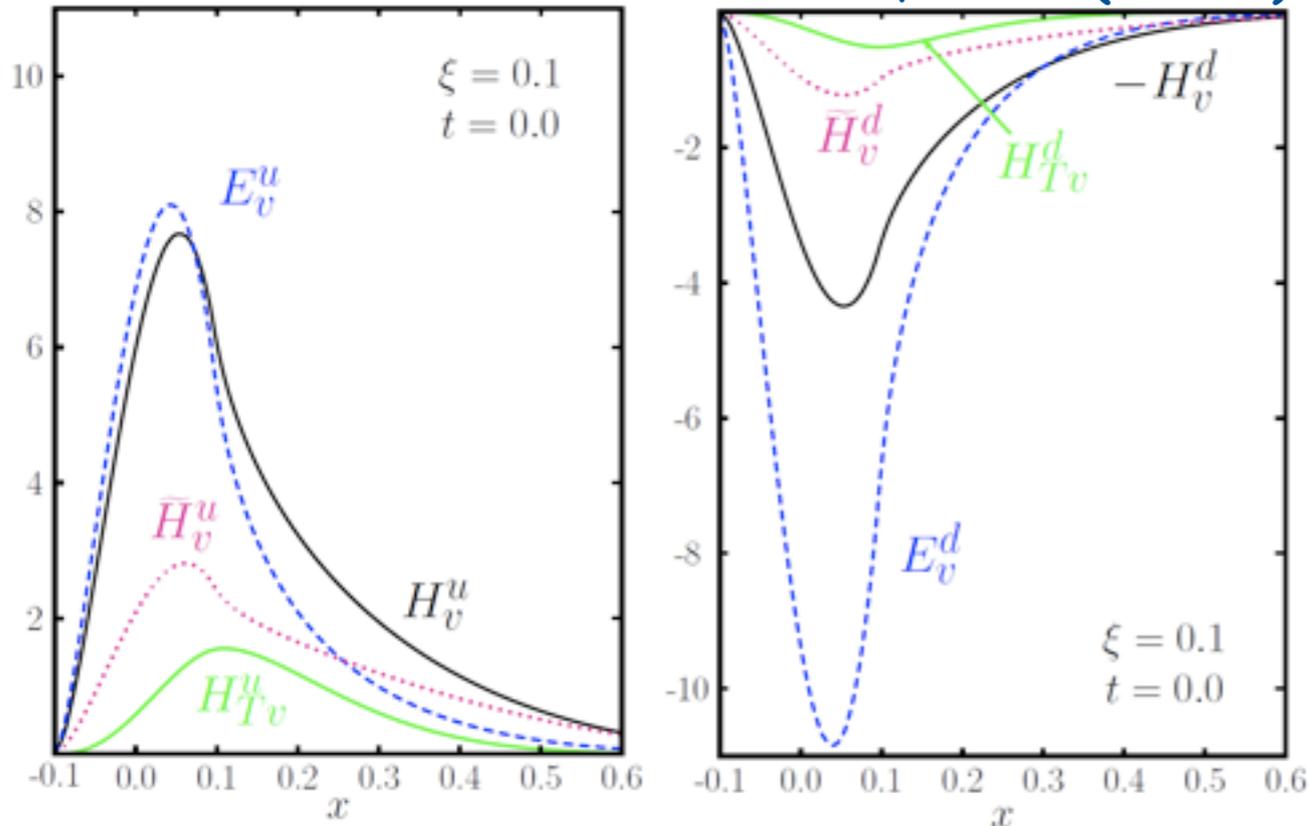


mes

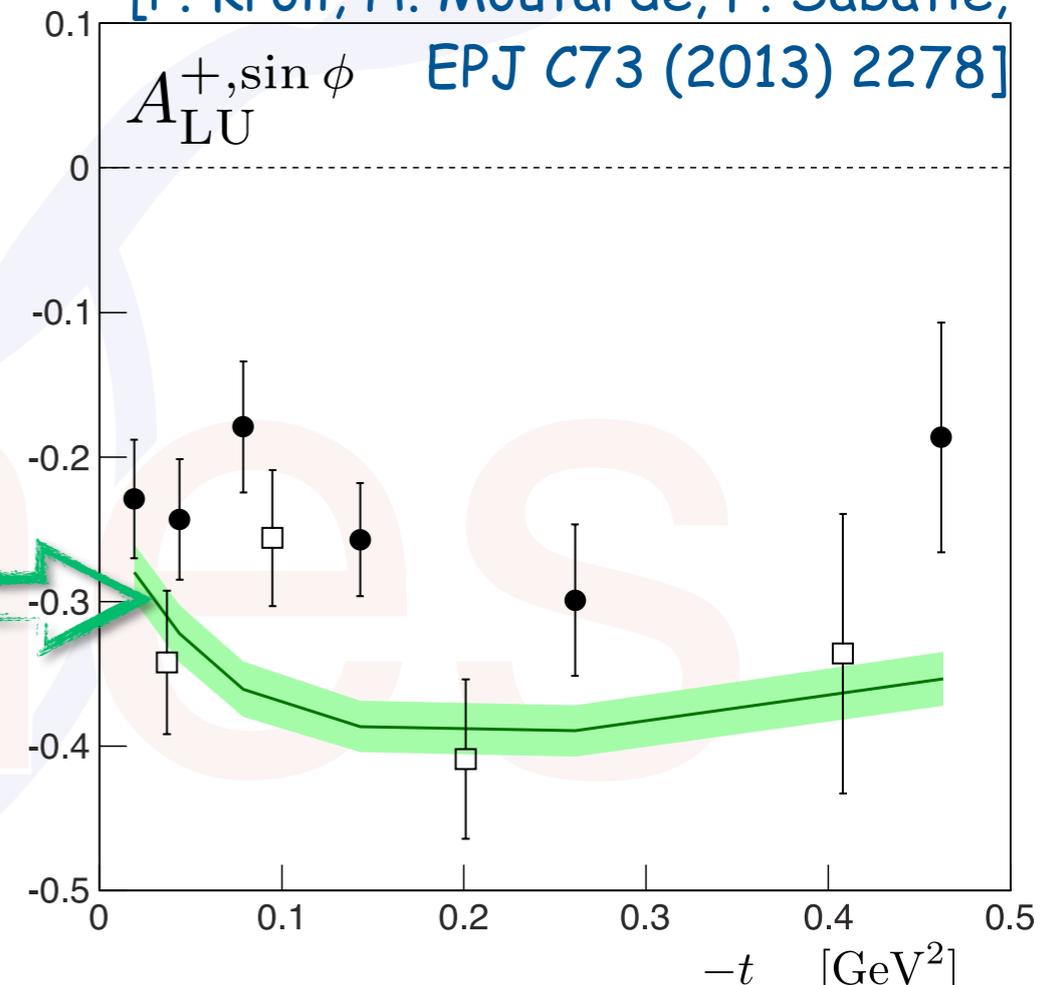
GPDs - a nice success story!



Goloskokov, Kroll (2007)



[P. Kroll, H. Moutarde, F. Sabatie, EPJ C73 (2013) 2278]



Conclusions

If you want to understand the proton, look at it deeply from various perspectives!

