

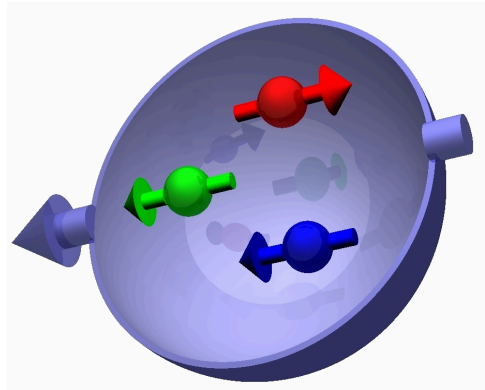
Spin Structure of the Proton and Neutron studied with Electron DIS

Recent HERMES Results on the Quark Structure
of the Nucleon

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Evolution of our Nucleon Picture

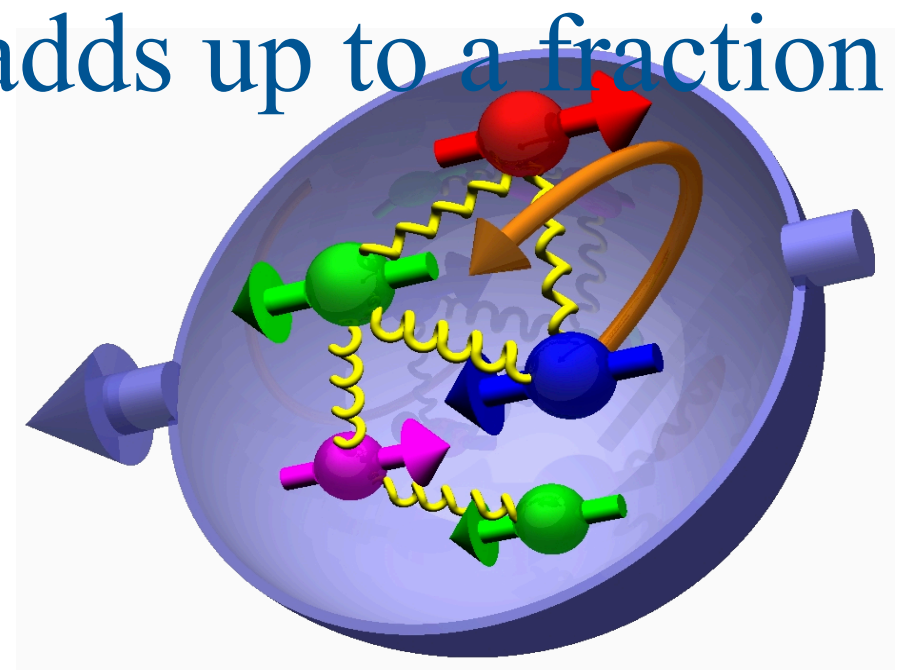


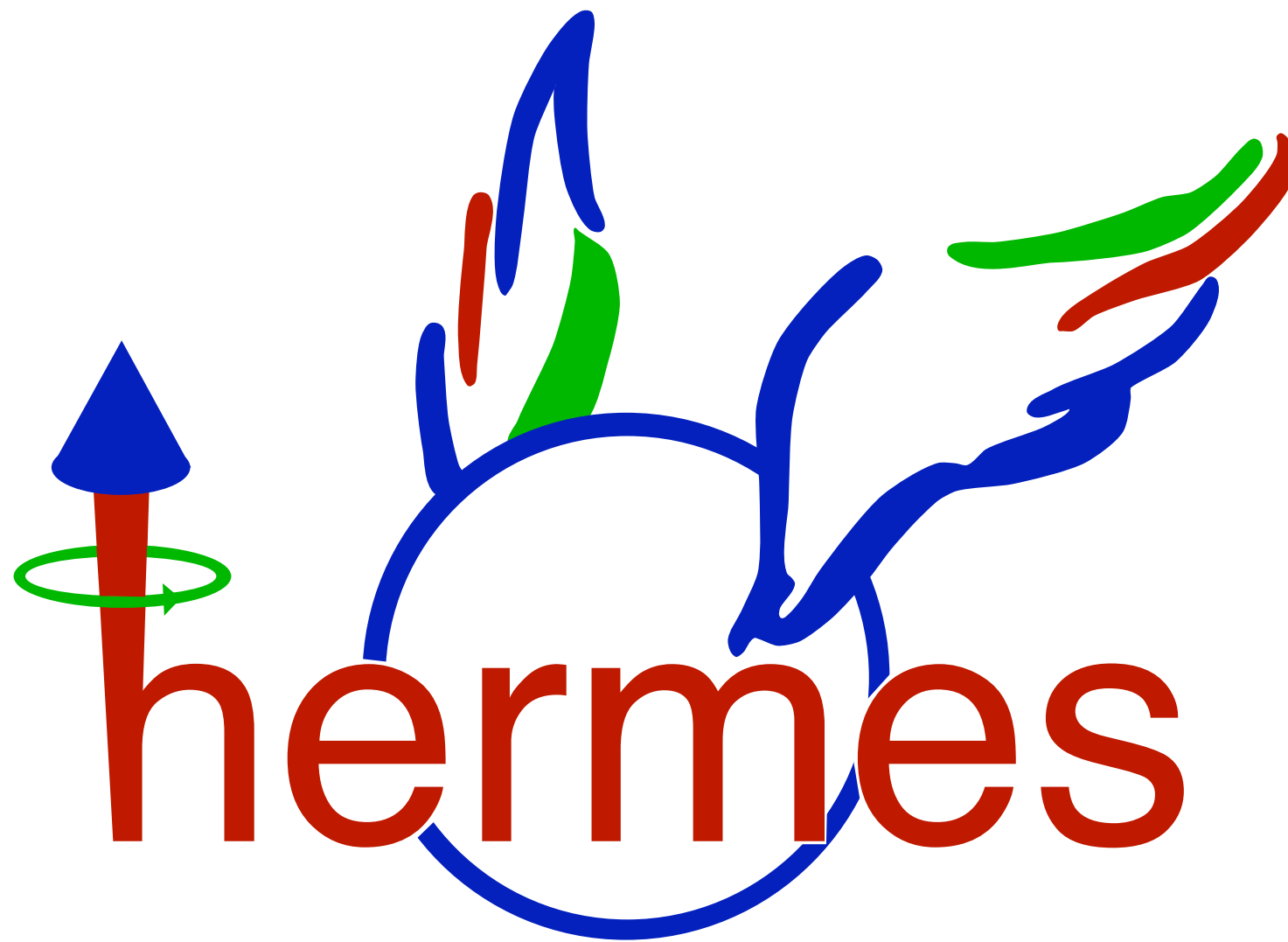
1960/70's : Nucleon consists of three (up and down) quarks whose spins add up to $1/2$)

1970/80's : ☹️ only ~50% of proton mass is coming from quarks as well!

☹️ quark spin only adds up to a fraction of $1/2$

1990/2000's : Don't forget orbital angular momentum!!





- proposed in 1988 to solve the “spin crisis”
- commissioned in 1995 at HERA
- data taking ended in 2007 because of closing of HERA
- however, many more exciting results to be expected from data on discs

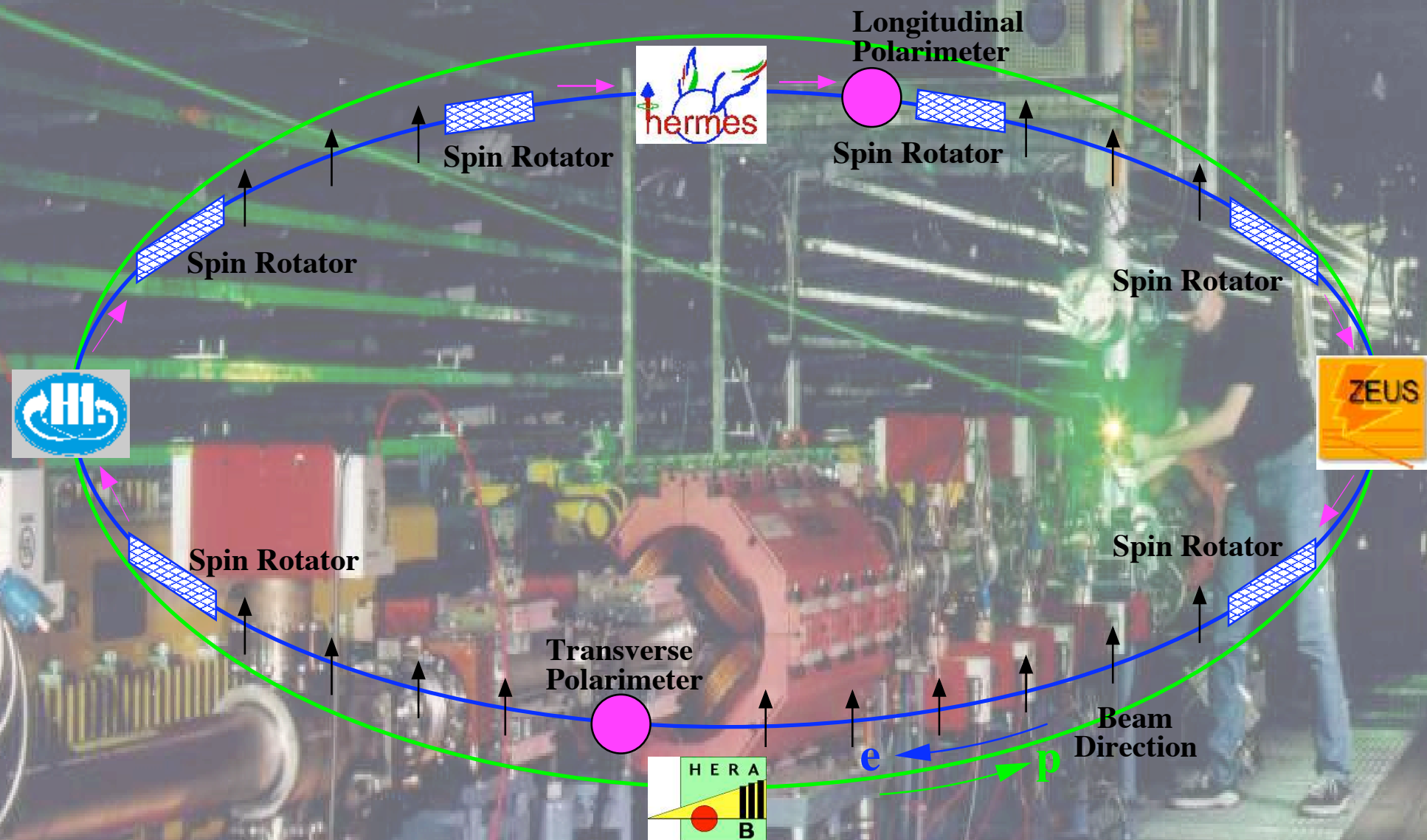
The HERMES Experiment

27.5 GeV e^+/e^- beam of HERA



- forward-acceptance spectrometer
⇒ $40\text{mrad} < \theta < 220\text{mrad}$
- high lepton ID efficiency and purity
- excellent hadron ID thanks to dual-radiator RICH

Beam Polarization at HERA

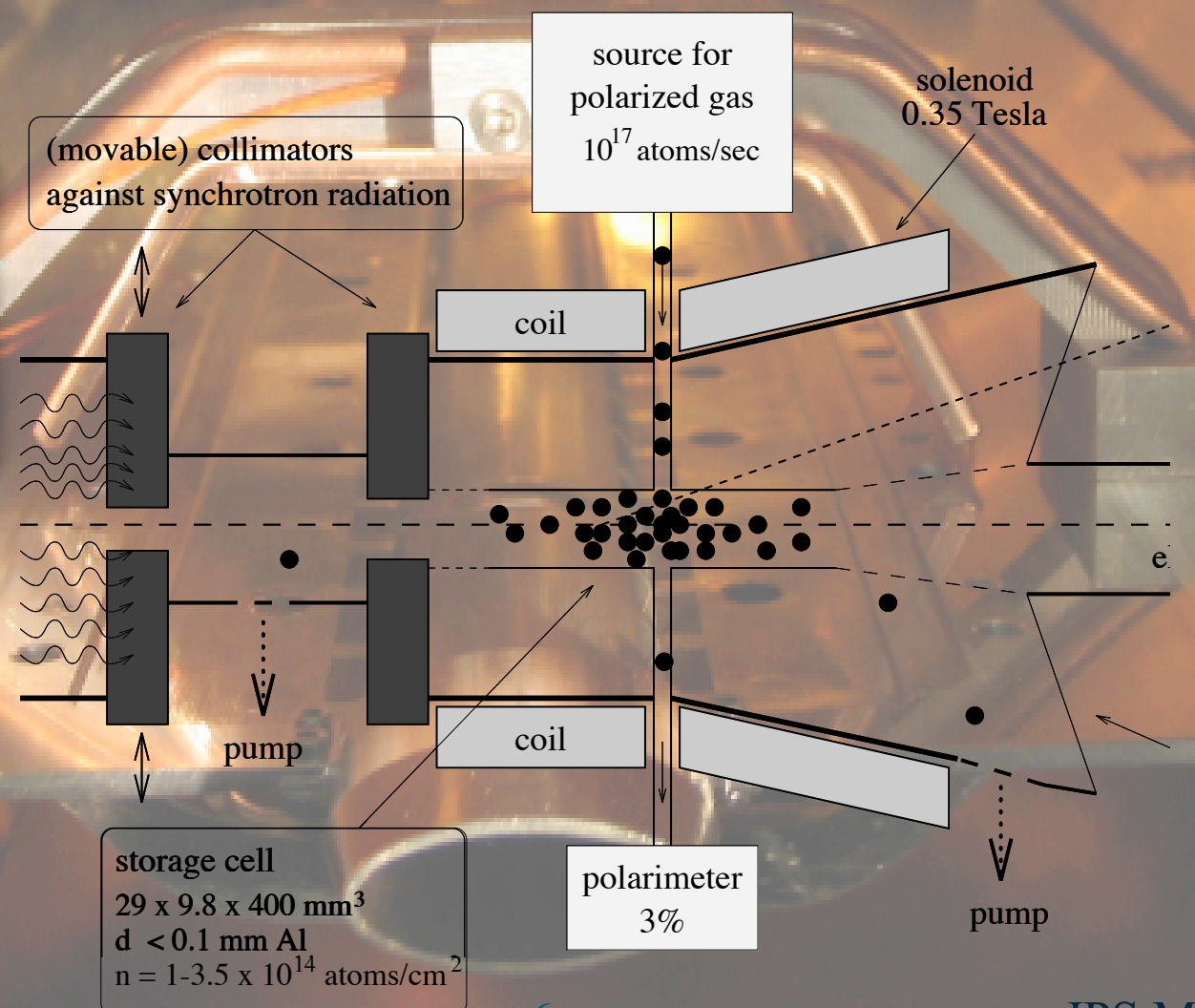


- 27.5 GeV e⁺/e⁻ beam
- Self-polarizing through Sokolov-Ternov-Effect
- Average beam polarization of about 55%

HERMES Polarized Target

- Storage cell with atomic beam source
- Pure target (NO dilution)
- Polarized or unpolarized targets possible
- Different gas targets available (H, D, He, N, Kr ...)

Polarization:
longitudinal: ~85%
transversal: ~75%

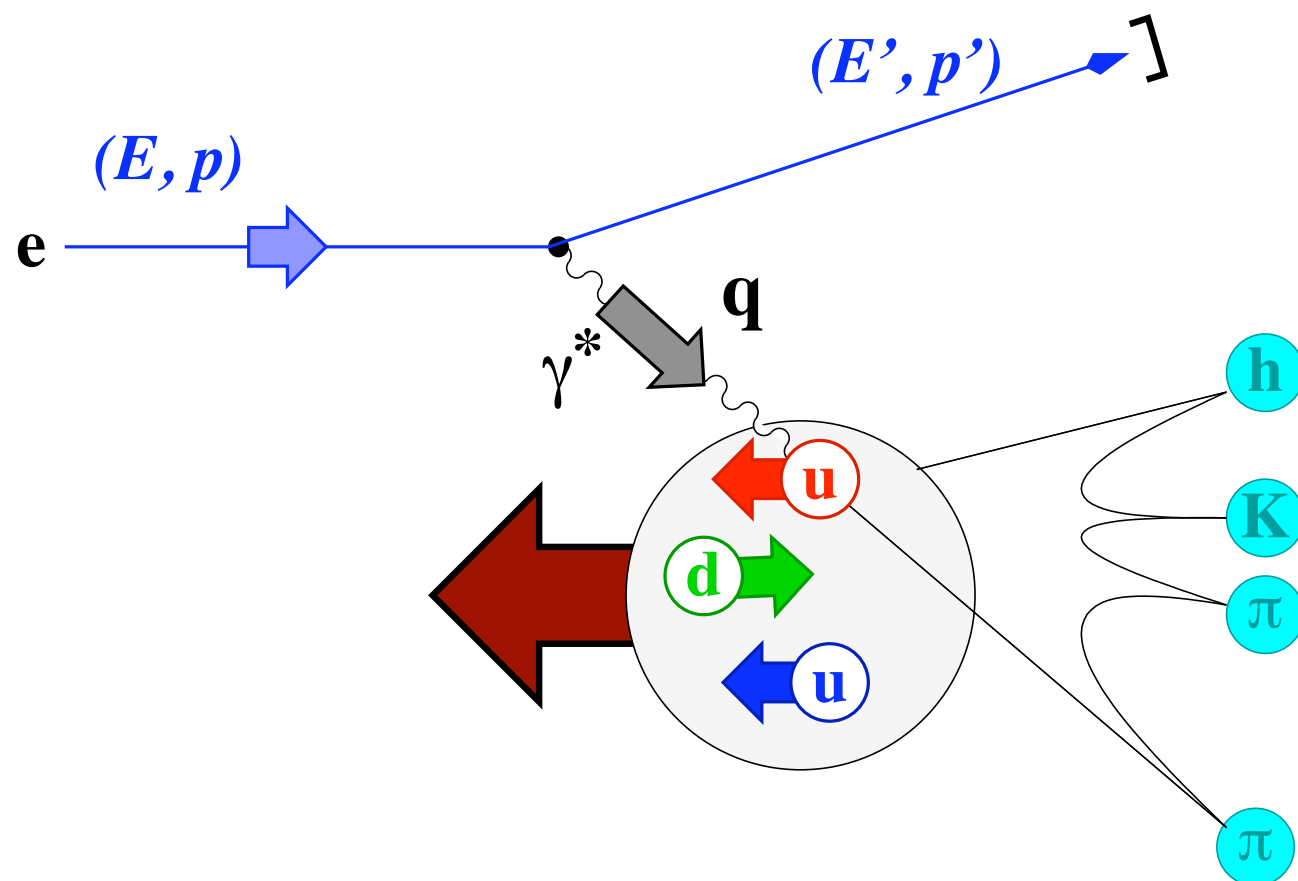


Deep-Inelastic Scattering

Probing the Structure of the Nucleon

Deep-Inelastic Scattering

use well-known probe to study hadronic structure:



Q^2	$\stackrel{\text{lab}}{=}$	$4EE' \sin^2(\frac{\Theta}{2})$
ν	$\stackrel{\text{lab}}{=}$	$E - E'$
W^2	$\stackrel{\text{lab}}{=}$	$M^2 + 2M\nu - Q^2$
y	$\stackrel{\text{lab}}{=}$	$\frac{\nu}{E}$
x	$\stackrel{\text{lab}}{=}$	$\frac{Q^2}{2M\nu}$

inclusive DIS: detect scattered lepton
 Factorization $\Rightarrow \sigma^{ep \rightarrow e h X} = \sum_q B_F^q \otimes \sigma^{q \rightarrow e \nu} \otimes F_F^{q \rightarrow h}$

semi-inclusive DIS: detect scattered lepton and some fragments
 exploit strong correlation between flavor structure of leading hadron and struck quark

Inclusive DIS

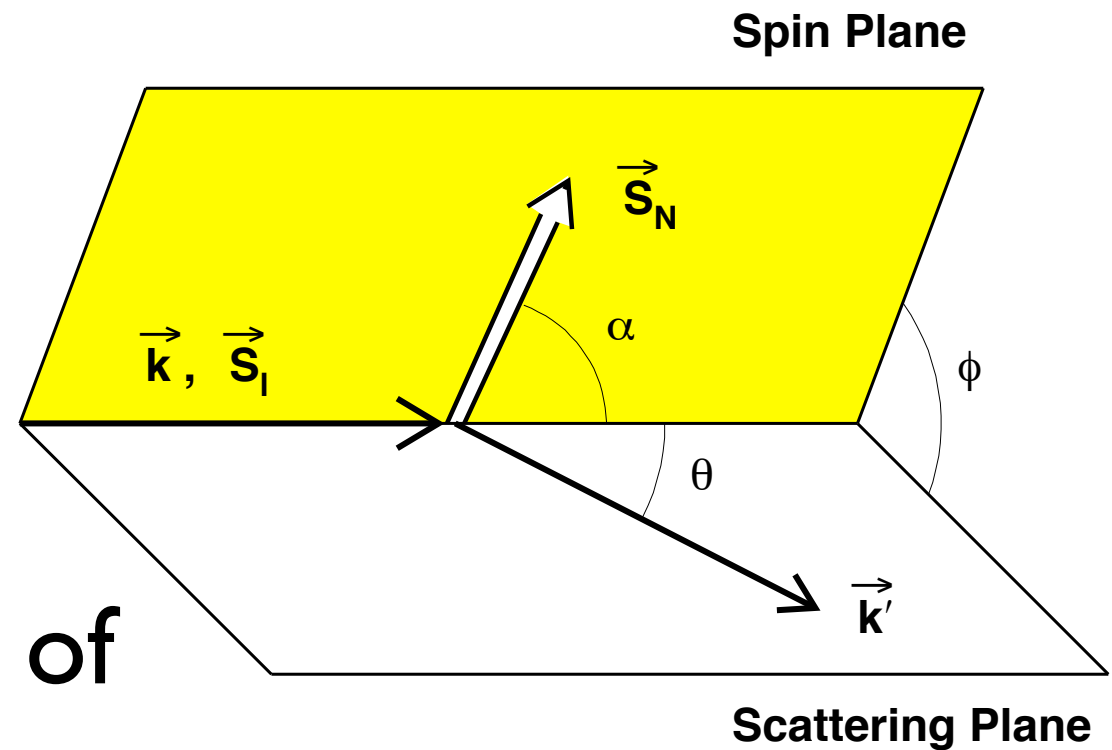
$$\frac{d^2\sigma(s, S)}{dx dQ^2} = \frac{2\pi\alpha^2 y^2}{Q^6} \mathbf{L}_{\mu\nu}(s) \mathbf{W}^{\mu\nu}(S)$$

Lepton Tensor

Hadron Tensor

parametrized in terms of

Structure Functions



$$\begin{aligned} \frac{d^3\sigma}{dxdy d\phi} \propto & \frac{y}{2} F_1(x, Q^2) + \frac{1 - y - \gamma^2 y^2 / 4}{2xy} F_2(x, Q^2) \\ & - P_l P_T \cos \alpha \left[\left(1 - \frac{y}{2} - \frac{\gamma^2 y^2}{4} g_1(x, Q^2) - \frac{\gamma^2 y}{2} g_2(x, Q^2) \right) \right] \\ & + P_l P_T \sin \alpha \cos \phi \gamma \sqrt{1 - y - \frac{\gamma^2 y^2}{4}} \left(\frac{y}{2} g_1(x, Q^2) + g_2(x, Q^2) \right) \end{aligned}$$

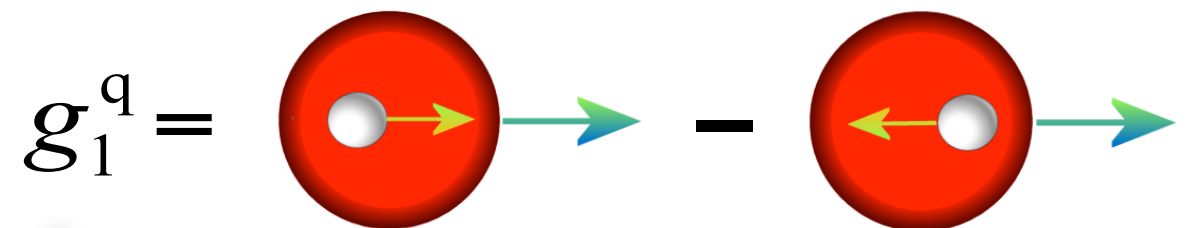
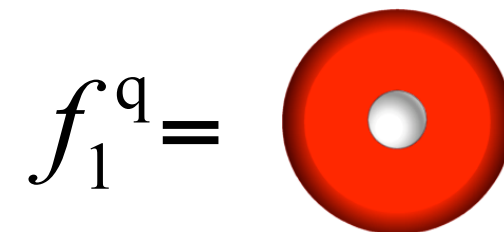
Parton-Model Interpretation of Structure Functions

$$F_1(x) = \frac{1}{2} \sum_q e_q^2 f_1^q(x)$$

$$F_2(x) = x \sum_q e_q^2 f_1^q(x)$$

$$g_1(x) = \frac{1}{2} \sum_q e_q^2 g_1^q(x)$$

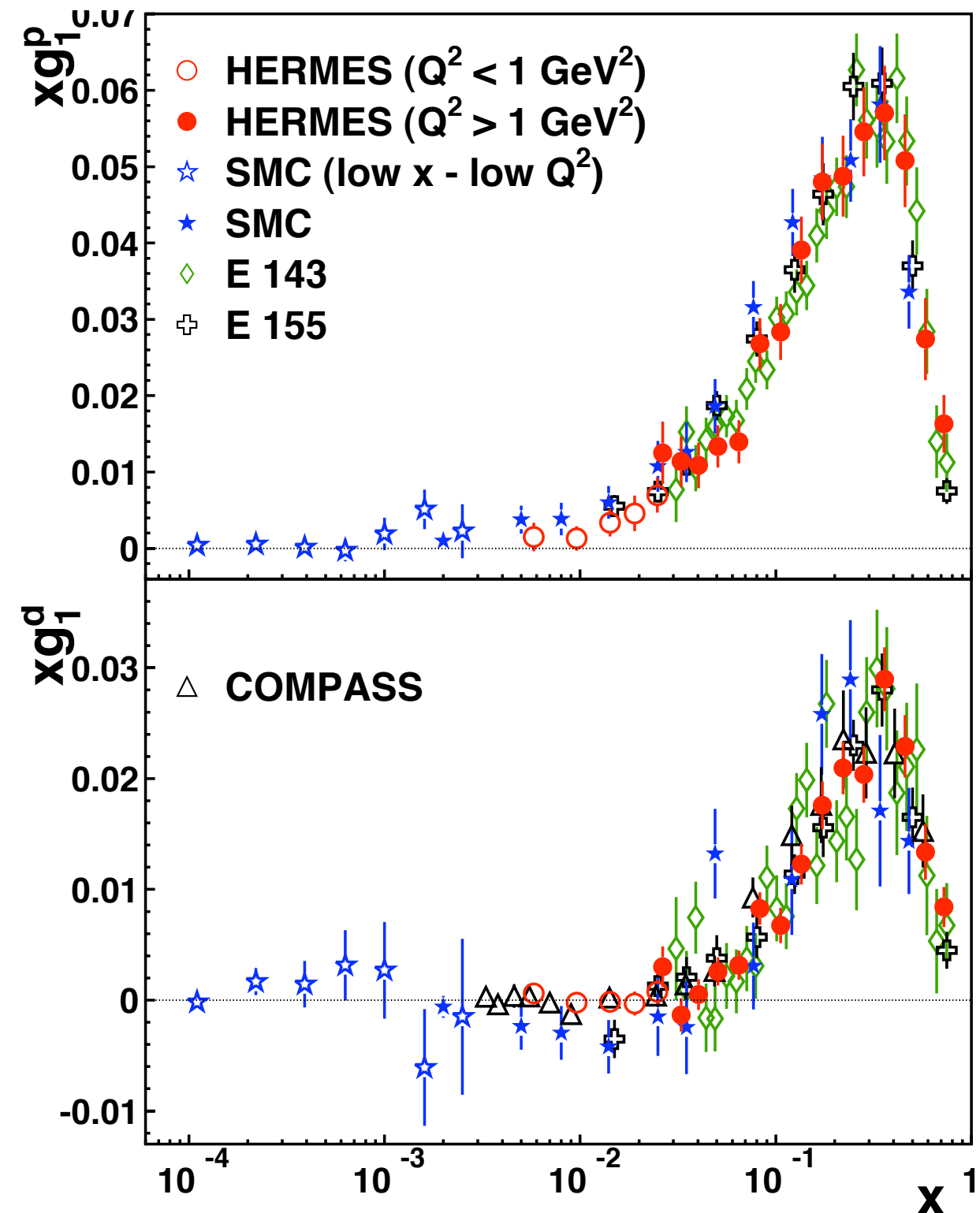
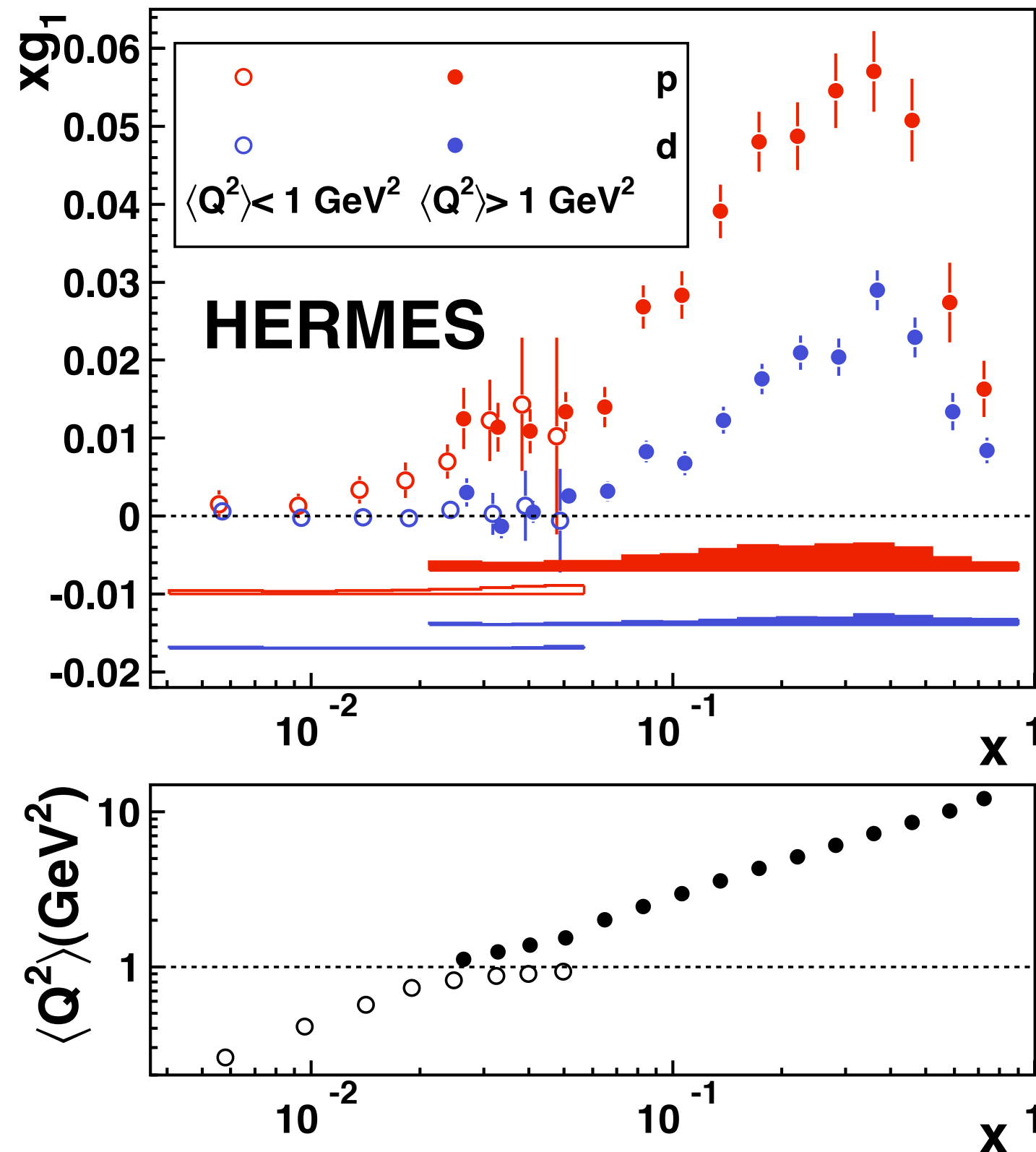
$$g_2(x) = 0$$



quark-spin contribution to nucleon helicity

Polarized Structure Function g_1

A. Airapetian et al., PRD 75 (2007)



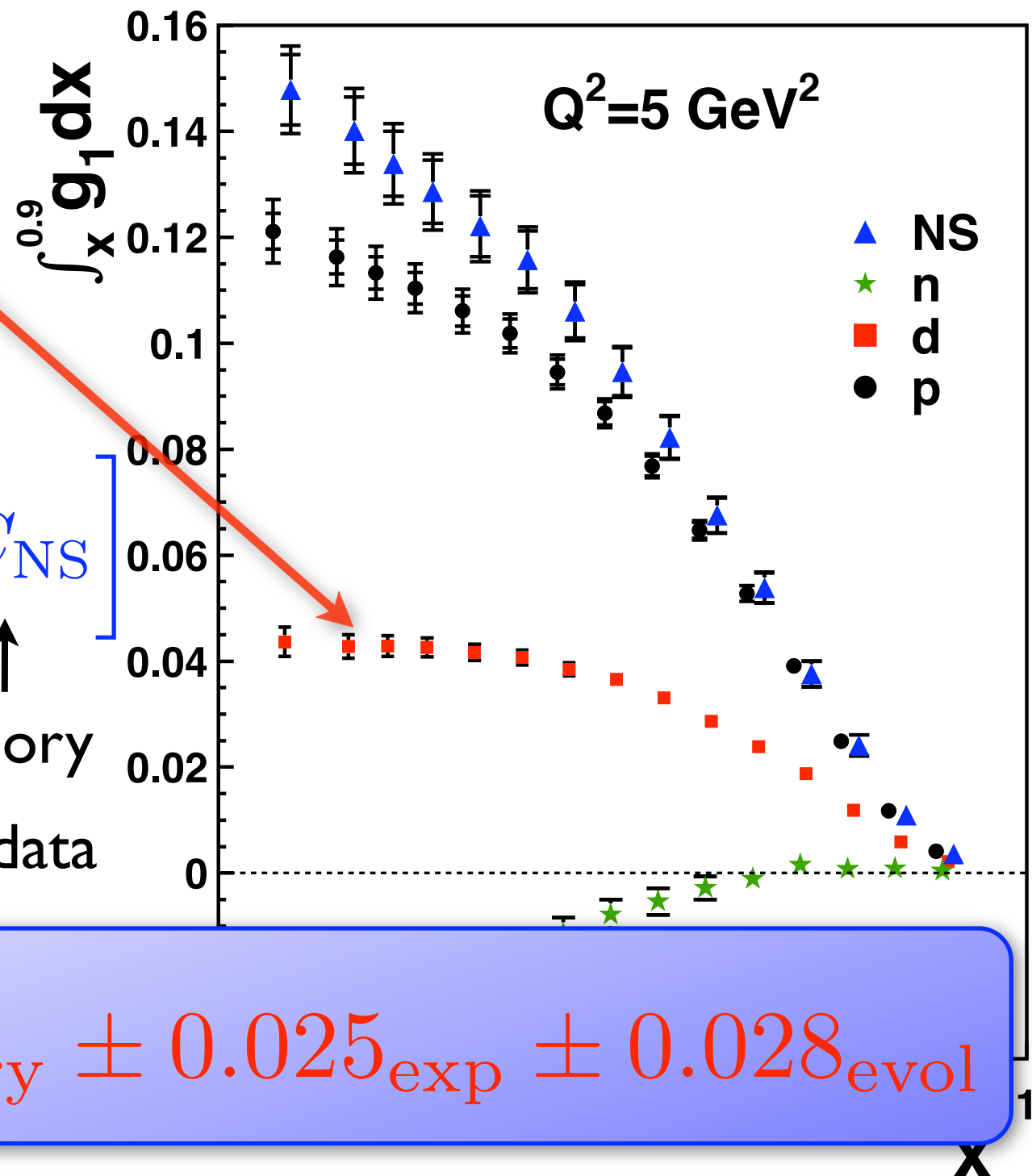
Integral of $g_1(x)$

Saturation

⇒ **close to full integral?**

$$\Delta\Sigma \stackrel{\overline{\text{MS}}}{=} \frac{1}{\Delta C_S} \left[\frac{9\Gamma_1^d}{1 - \frac{3}{2}\omega_D} - \frac{1}{4}a_8\Delta C_{\text{NS}} \right]$$

$\uparrow$ theory $\uparrow$ 0.05 ± 0.05 $\uparrow$ theory
 hyperon-decay data



$$\Delta\Sigma \stackrel{\overline{\text{MS}}}{=} 0.330 \pm 0.011_{\text{theory}} \pm 0.025_{\text{exp}} \pm 0.028_{\text{evol}}$$

most precise result; only 1/3 of nucleon spin from quarks

Flavor Separation using SIDIS

- Remember:

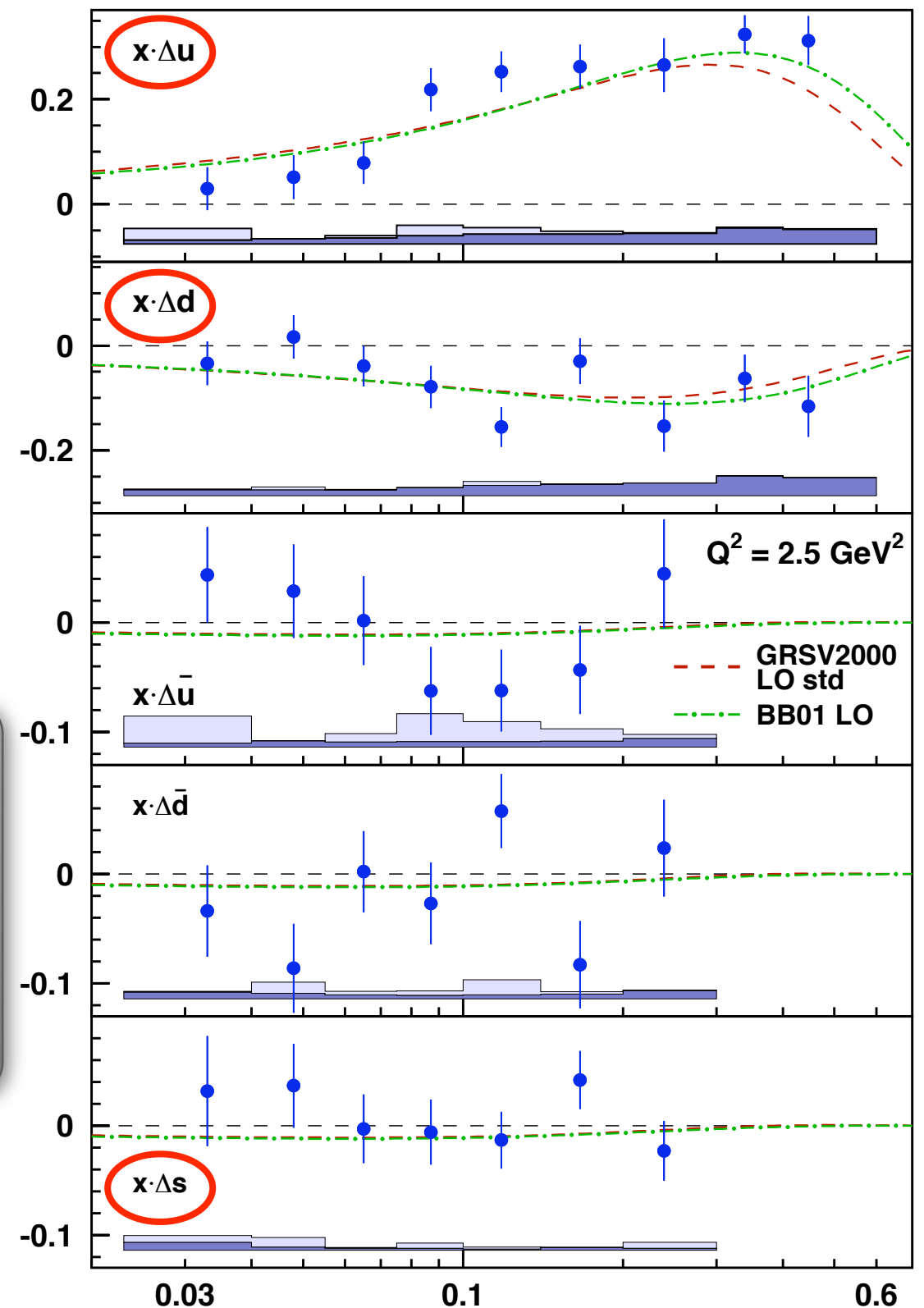
$$\Delta\Sigma = \sum_q \int dx g_1^q(x) \approx 1/3$$

- use different hadron flavors in final state to *tag* quark flavor

$$\begin{aligned} g_1^u(x) &\equiv \Delta u > 0 \text{ and large} \\ g_1^d(x) &\equiv \Delta d < 0 \text{ and smaller} \\ g_1^s(x) &\equiv \Delta s \approx 0 \end{aligned}$$

A.Airapetian et al., PRL 92 (2004)

A.Airapetian et al., PRD 71 (2005)

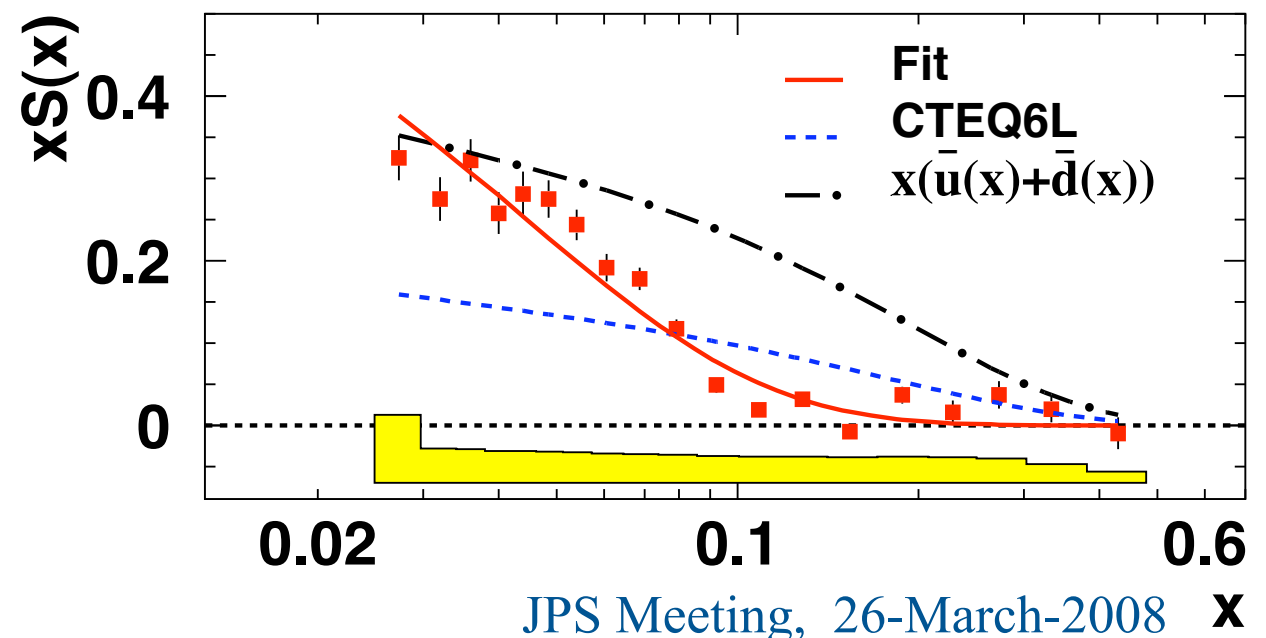
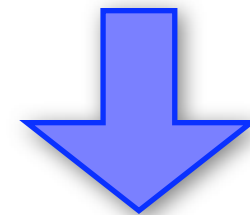
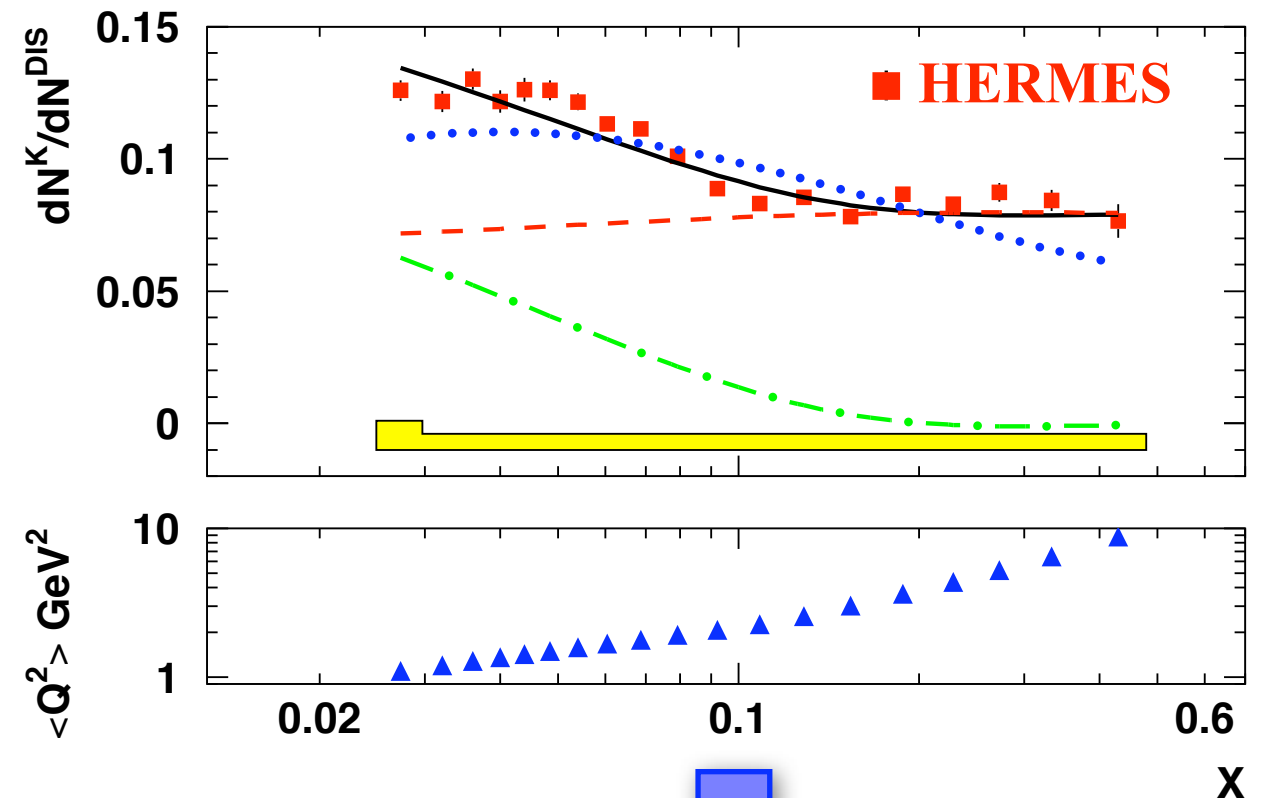


Distribution of Strange Quark

- strange quarks carry no isospin, thus the same in proton and neutron
- use isoscalar probe and target to extract strange-quark distributions
- only need inclusive asymmetries and **$K^+ + K^-$ asymmetries**, i.e., $A_{1,d}(x, Q^2)$ and $A_{1,d}^{K^+ + K^-}(x, z, Q^2)$, as well as **$K^+ + K^-$ multiplicities**
- strange-quark fragmentation function either directly from data or from parametrizations

Unpolarized Strange Quarks

- $S(x)$ non-zero for $x < 0.1$
- vanishes for $x > 0.1$
- apparent discrepancy with CTEQ6L
- $S(x)$ not an average of an isoscalar non-strange sea
 $\Rightarrow$ violation of SU(3) symmetry
- HERMES data provides extremely valuable input to extraction of unpolarized PDFs



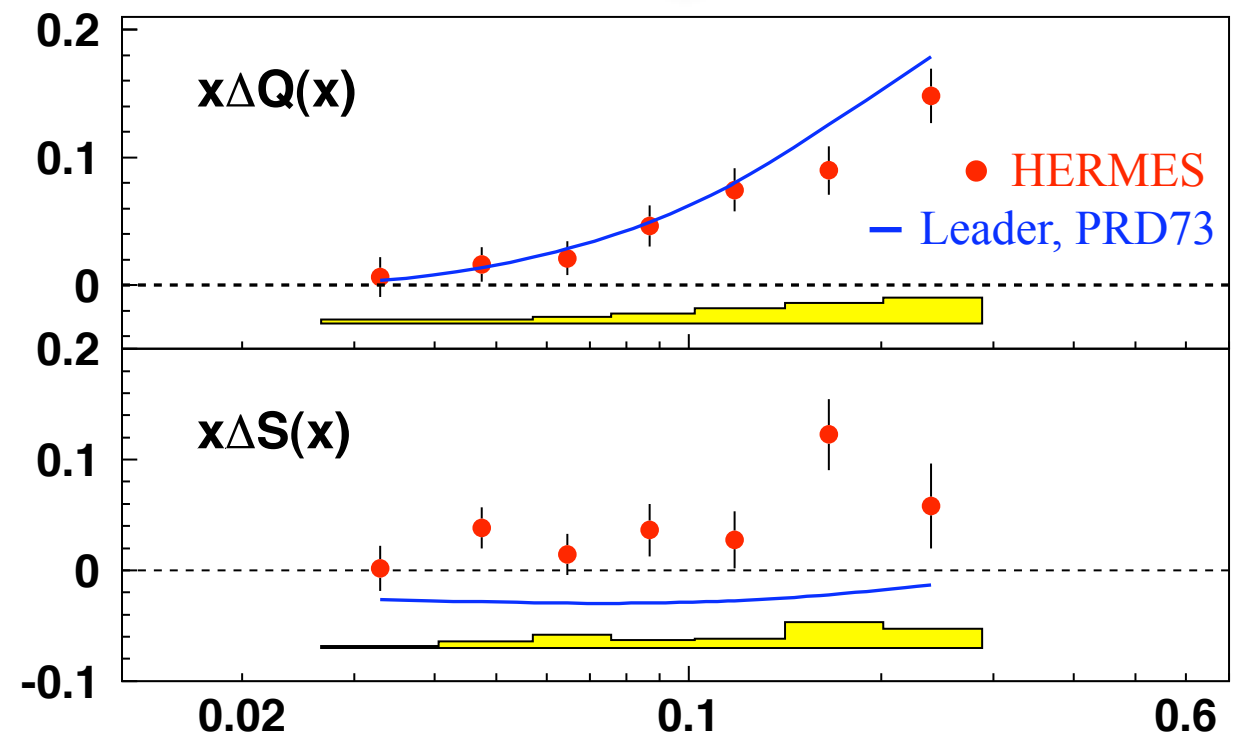
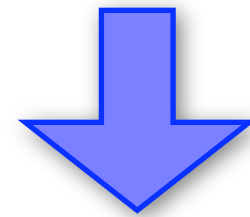
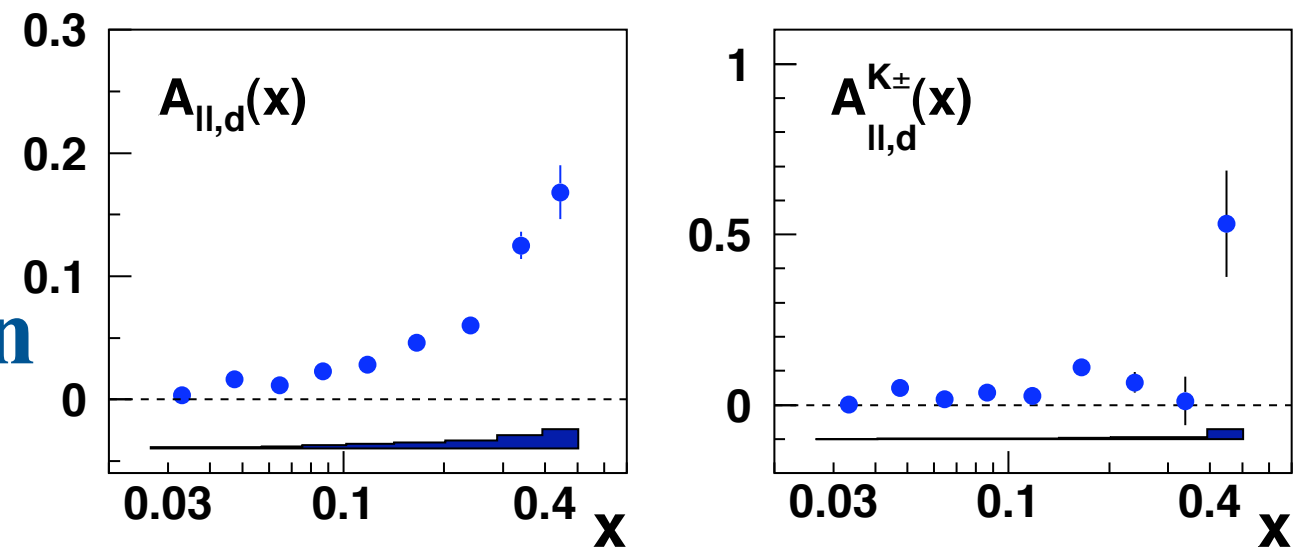
A. Airapetian et al., arXiv:0803.2993

Polarized Strange Quarks

- results consistent with previous flavor decomposition
- no sizeable negatively polarized strange sea as expected from inclusive DIS results
- ➡ sign of violation of SU(3) symmetry or of substantial contributions from low- x region!

A. Airapetian et al., arXiv:0803.2993

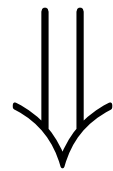
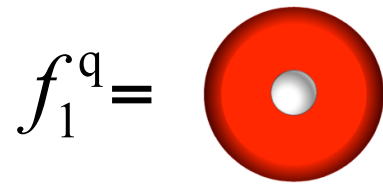
Gunar Schnell



JPS Meeting, 26-March-2008 X

Transverse-Spin Effects

Quark Structure of the Nucleon

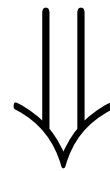
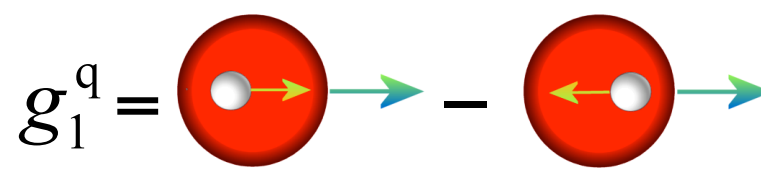


Unpolarized quarks
and nucleons

$f_1^q(x)$: spin averaged
(well known)

⇒ Vector Charge

$$\langle PS | \bar{\Psi} \gamma^\mu \Psi | PS \rangle = \int dx (f_1^q(x) - f_1^{\bar{q}}(x))$$

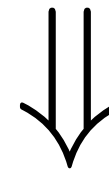
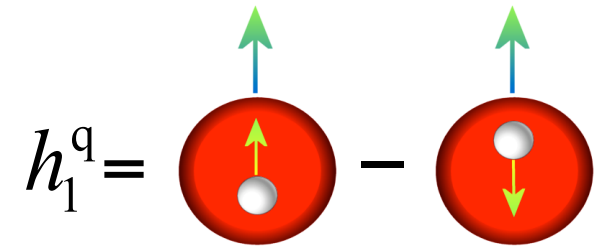


Longitudinally
polarized quarks
and nucleons

$g_1^q(x)$: helicity
difference (known)

⇒ Axial Charge

$$\langle PS | \bar{\Psi} \gamma^\mu \gamma_5 \Psi | PS \rangle = \int dx (g_1^q(x) + g_1^{\bar{q}}(x))$$



Transversely
polarized quarks
and nucleons

$h_1^q(x)$: transversity
(hardly known!)

⇒ Tensor Charge

$$\langle PS | \bar{\Psi} \sigma^{\mu\nu} \gamma_5 \Psi | PS \rangle = \int dx (h_1^q(x) - h_1^{\bar{q}}(x))$$

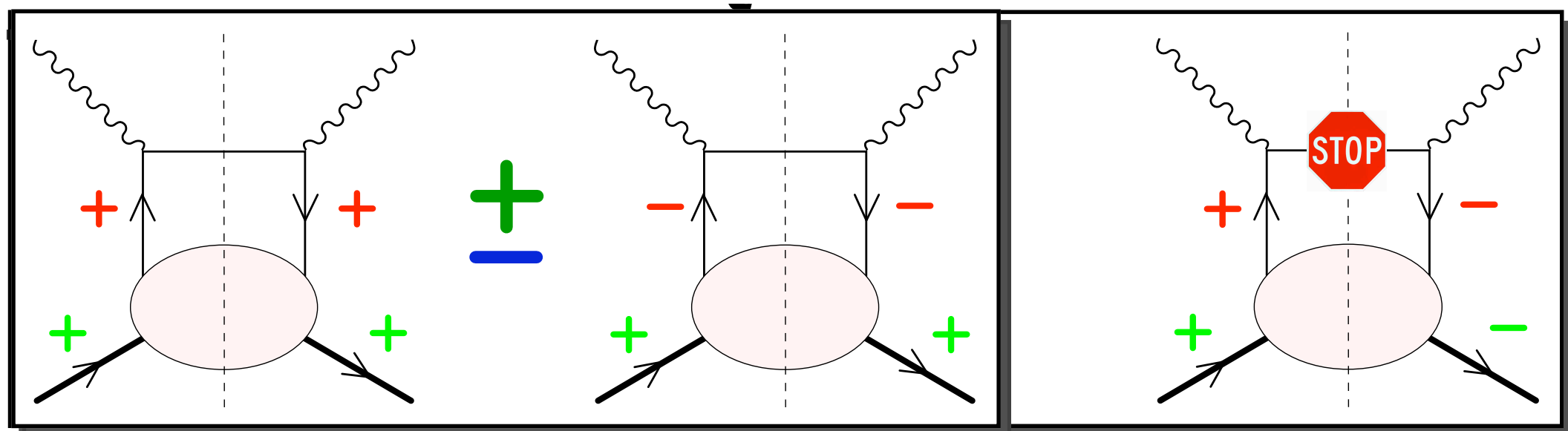
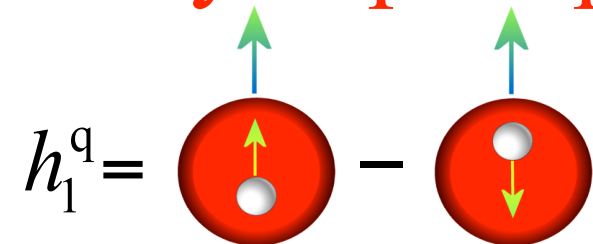
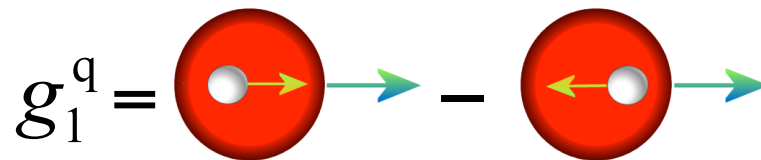
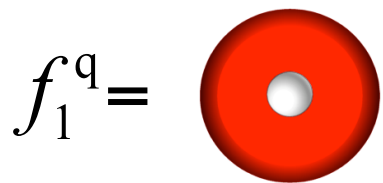
The “Trouble” with Transversity

Transverse-spin states written in terms of helicity states:

$$|\uparrow\rangle = \frac{1}{\sqrt{2}} [|+\rangle + i|-\rangle]$$

$$|\downarrow\rangle = \frac{1}{\sqrt{2}} [|+\rangle - i|-\rangle]$$

Transverse-spin asymmetries involve helicity-flip amplitudes



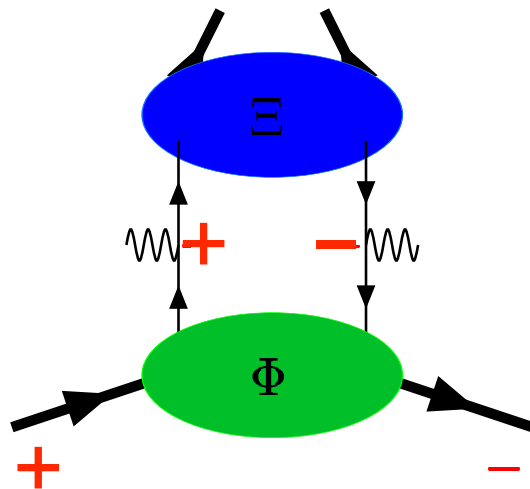
Transversity Measurement

How can one measure transversity?

Need another chiral-odd object!

⇒ Semi-Inclusive DIS

$$\sigma^{ep \rightarrow ehX} = \sum_q h_1^q \otimes \sigma^{eq \rightarrow eq} \otimes F F^{q \rightarrow h}$$



$\Downarrow$
chiral-odd
DF
 $\Downarrow$
chiral-odd
FF
 CHIRAL EVEN

→ chiral-odd FF as a **polarimeter** of transv. quark polarization

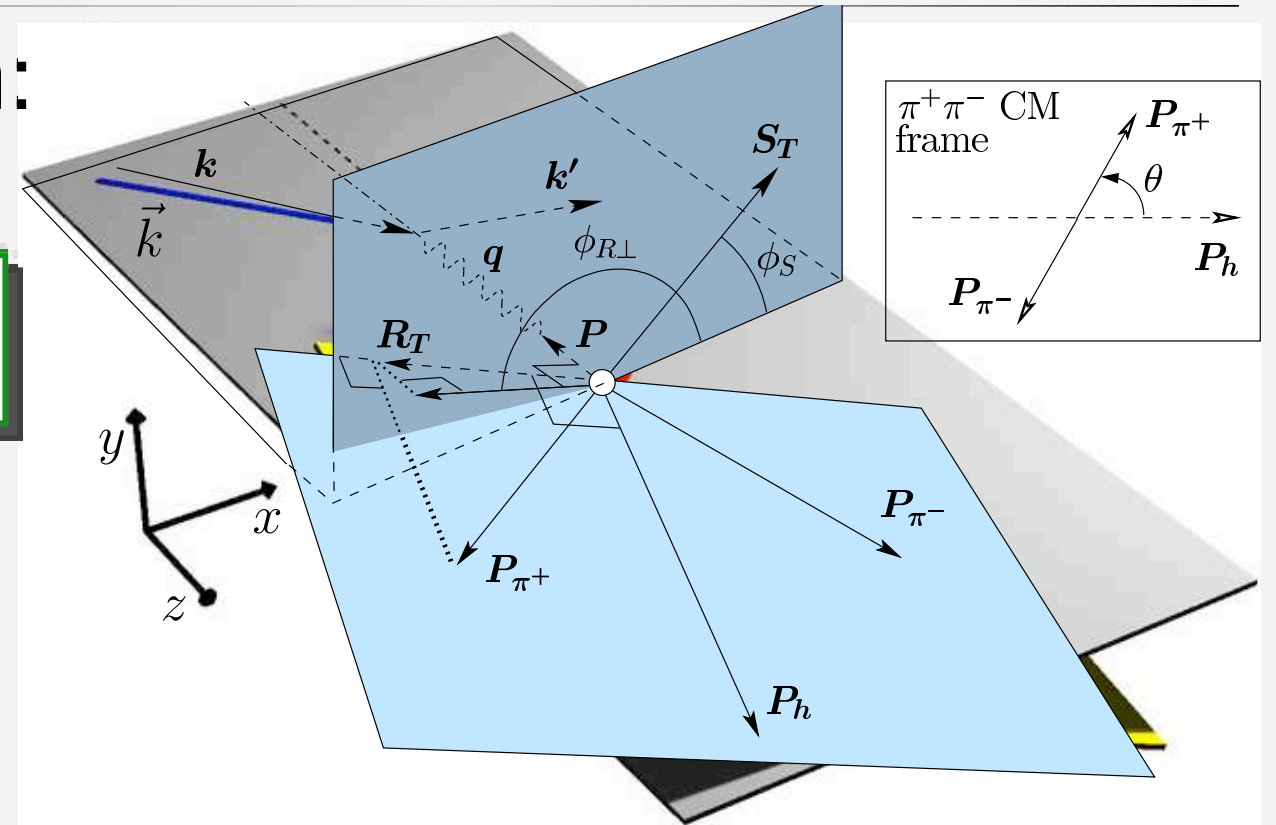
2-Hadron Fragmentation

polarized 2-hadron cross section:
(Unpolarized beam, Transversely pol. target)

$$\sigma_{UT} \sim \sin(\phi_{R\perp} + \phi_S) \sum e_q^2 h_1^q H_1^\Delta$$

$$H_1^\Delta = H_1^\Delta(z, \zeta, M_{\pi\pi}^2)$$

$$(\zeta \sim z_1/(z_1 + z_2))$$



- 😊 only relative momentum of hadron pair relevant
⇒ integration over transverse momentum of hadron pair simplifies factorization and Q^2 evolution
- 😞 however, cross section becomes quite complex
(differential in 9 variables)

Model for 2-Hadron Fragmentation

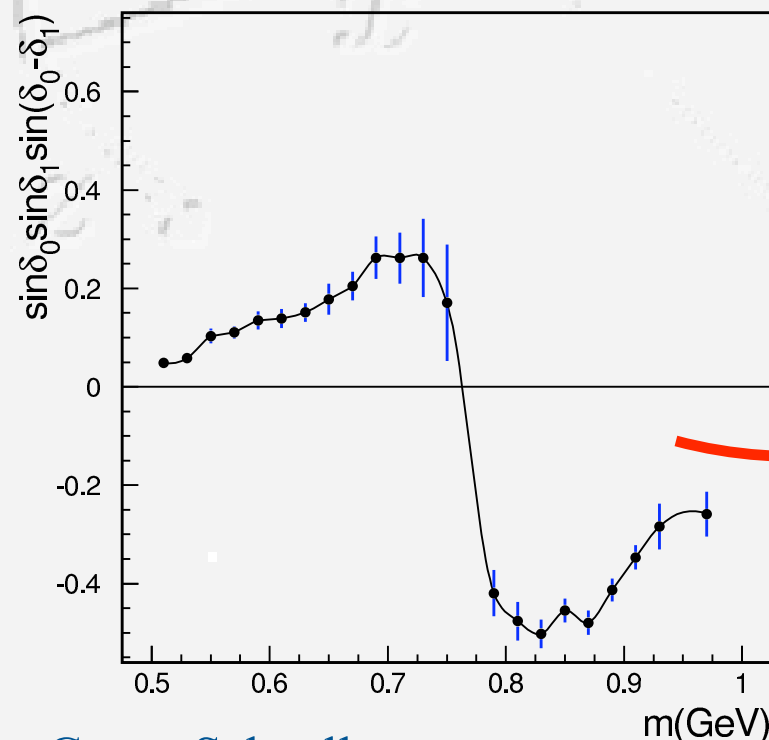
$$A_{UT} \sim \sin(\phi_{R\perp} + \phi_S) \sin \theta h_1 H_1^{\triangleleft}$$

Expansion of $H_1^{\triangleleft}$ in Legendre moments:

$$H_1^{\triangleleft}(z, \cos \theta, M_{\pi\pi}^2) = H_1^{\triangleleft,sp}(z, M_{\pi\pi}^2) + \cos \theta H_1^{\triangleleft,pp}(z, M_{\pi\pi}^2)$$

describe interference between 2 pion pairs coming from different production channels.

about $H_1^{\triangleleft,sp}$:



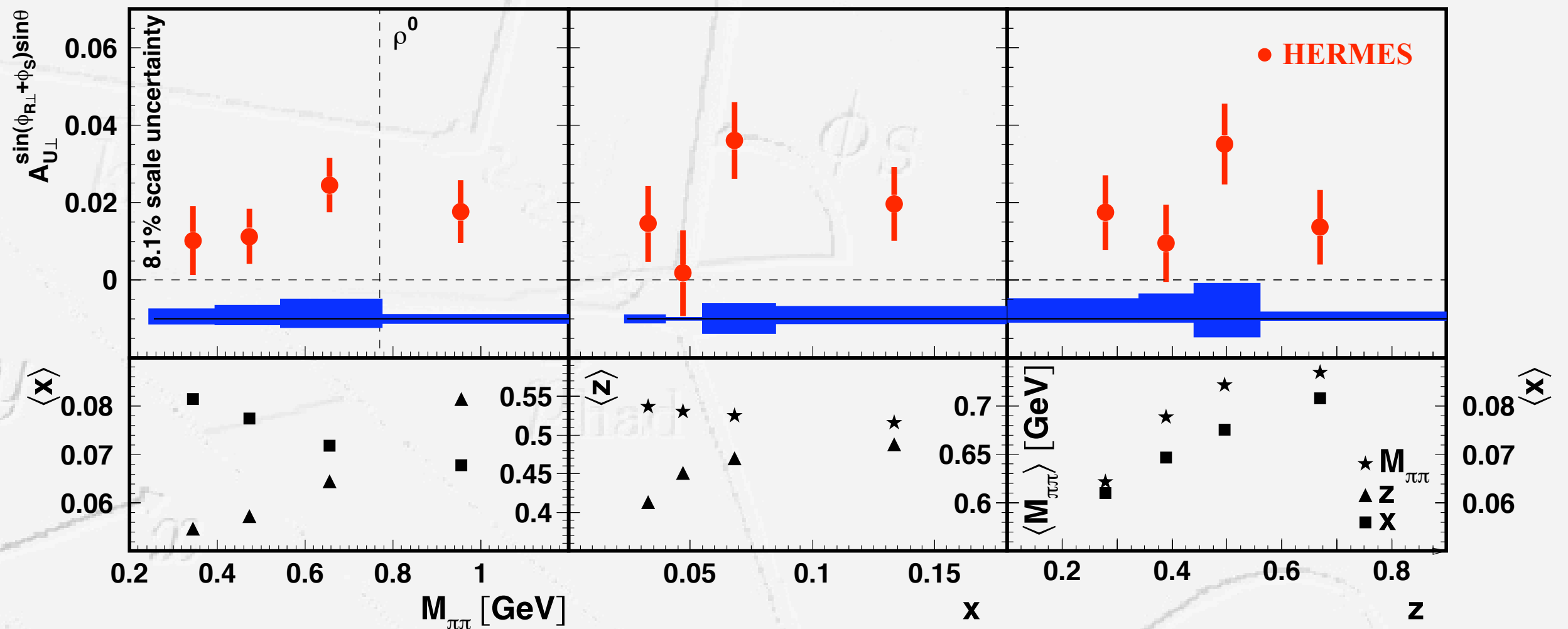
Jaffe et al. [[hep-ph/9709322](#)]:

$$H_1^{\triangleleft,sp}(z, M_{\pi\pi}^2) = \frac{\sin \delta_0 \sin \delta_1 \sin(\delta_0 - \delta_1) H_1^{\triangleleft,sp'}(z)}{\delta_0 (\delta_1) \rightarrow \text{S(P)-wave phase shifts}}$$

$$= \mathcal{P}(M_{\pi\pi}^2) H_1^{\triangleleft,sp'}(z)$$

$\Rightarrow A_{UT}$ might depend strongly on $M_{\pi\pi}$

HERMES Results (complete data set)



- ✓ first evidence for T-odd 2-hadron fragmentation function in SIDIS!
- ✓ invariant-mass dependence rules out Jaffe model

A. Airapetian et al., arXiv:0803.2367

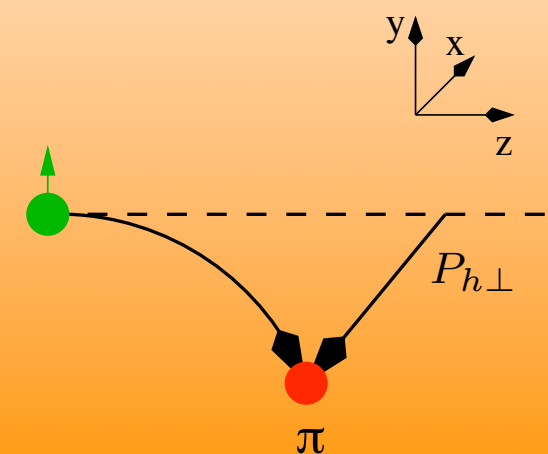
1-Hadron Production ($ep \rightarrow ehX$)

$$\begin{aligned}
 d\sigma = & d\sigma_{UU}^0 + \cos 2\phi d\sigma_{UU}^1 + \frac{1}{Q} \cos \phi d\sigma_{UU}^2 + \lambda_e \frac{1}{Q} \sin \phi d\sigma_{LU}^3 \\
 & + S_L \left\{ \sin 2\phi d\sigma_{UL}^4 + \frac{1}{Q} \sin \phi d\sigma_{UL}^5 + \lambda_e \left[d\sigma_{LL}^6 + \frac{1}{Q} \cos \phi d\sigma_{LL}^7 \right] \right\} \\
 & + S_T \left\{ \sin(\phi - \phi_S) d\sigma_{UT}^8 + \sin(\phi + \phi_S) d\sigma_{UT}^9 + \sin(3\phi - \phi_S) d\sigma_{UT}^{10} \right\}
 \end{aligned}$$

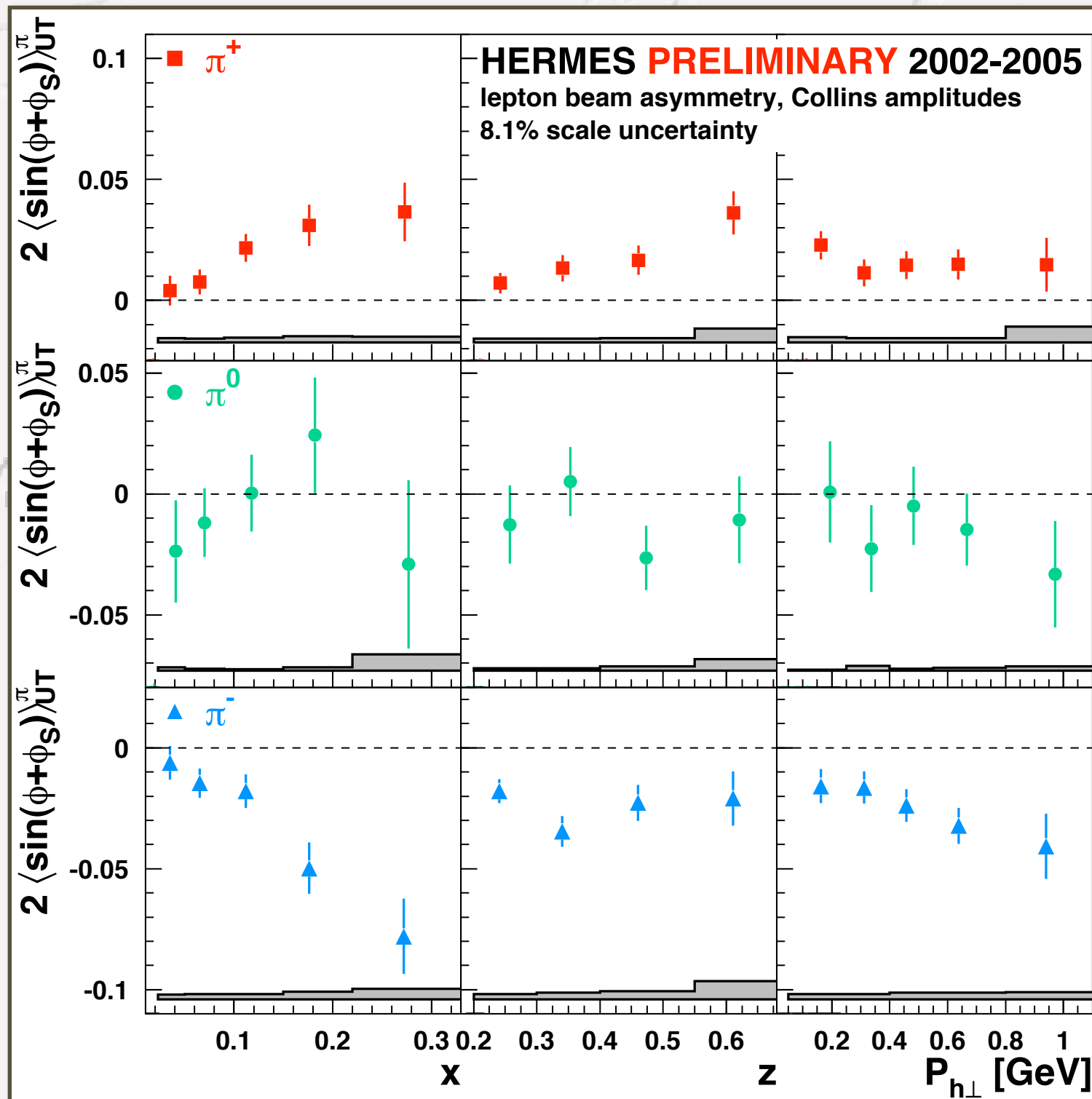
Sivers Effect:

- correlates hadron's transverse momentum with nucleon spin
- requires orbital angular momentum

hadron transverse spin



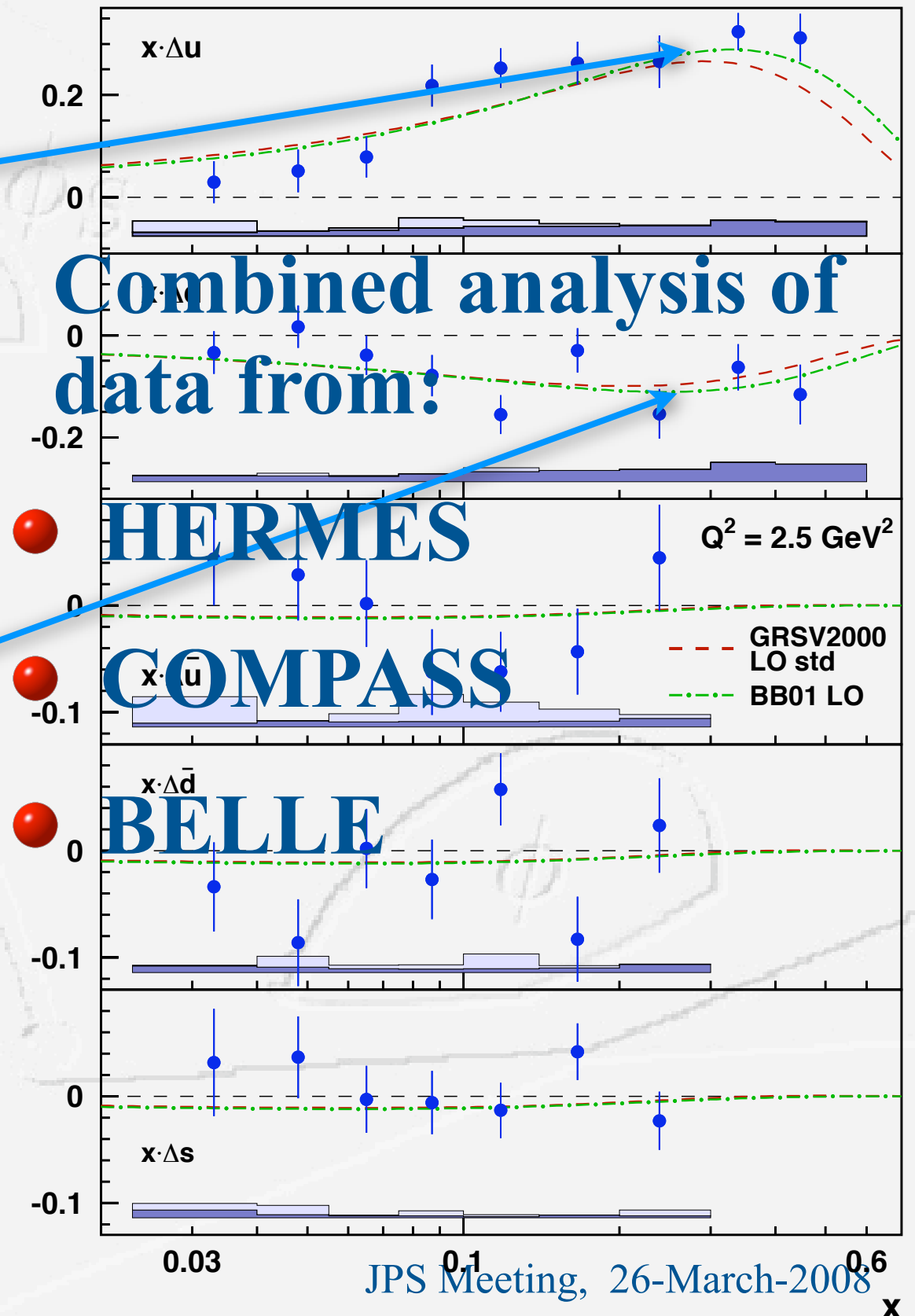
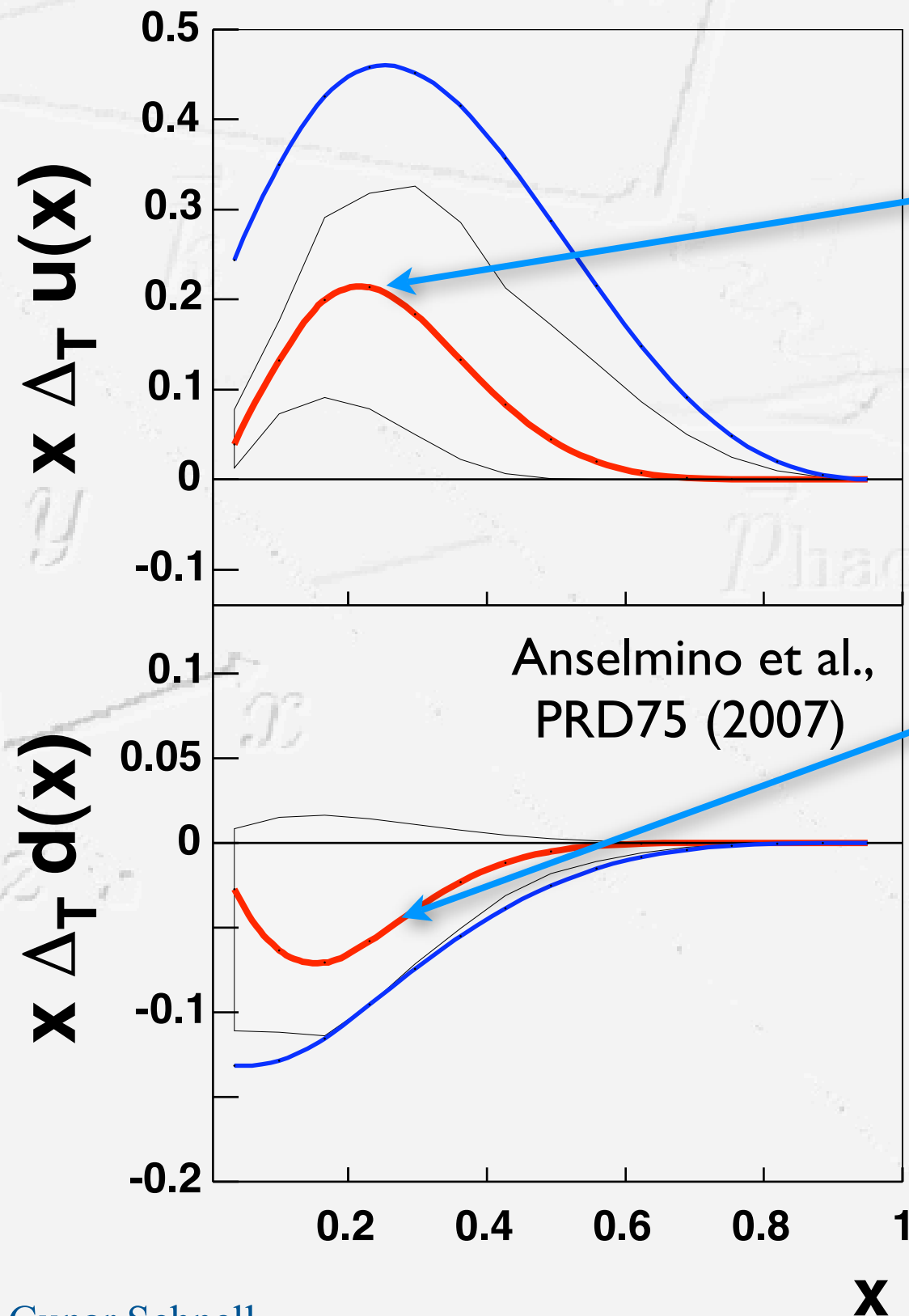
The HERMES Collins Results



- published[†] results **confirmed** with much higher statistical precision
- overall scale uncertainty of 8.1%
- positive for π^+ and negative for π^- as maybe expected ($\delta u \equiv h_1^u > 0$ and $\delta d \equiv h_1^d < 0$)
- ☒ **non-zero Collins effect observed!**
- unexpected **large π^- asymmetry**
⇒ role of **disfavored Collins FF**
- most likely: $H_1^{\perp, disf} \approx -H_1^{\perp, fav}$
- ☒ **both Collins FF and isospin symmetry among charged and neutral pions fulfilled**
- transversity sizeable**

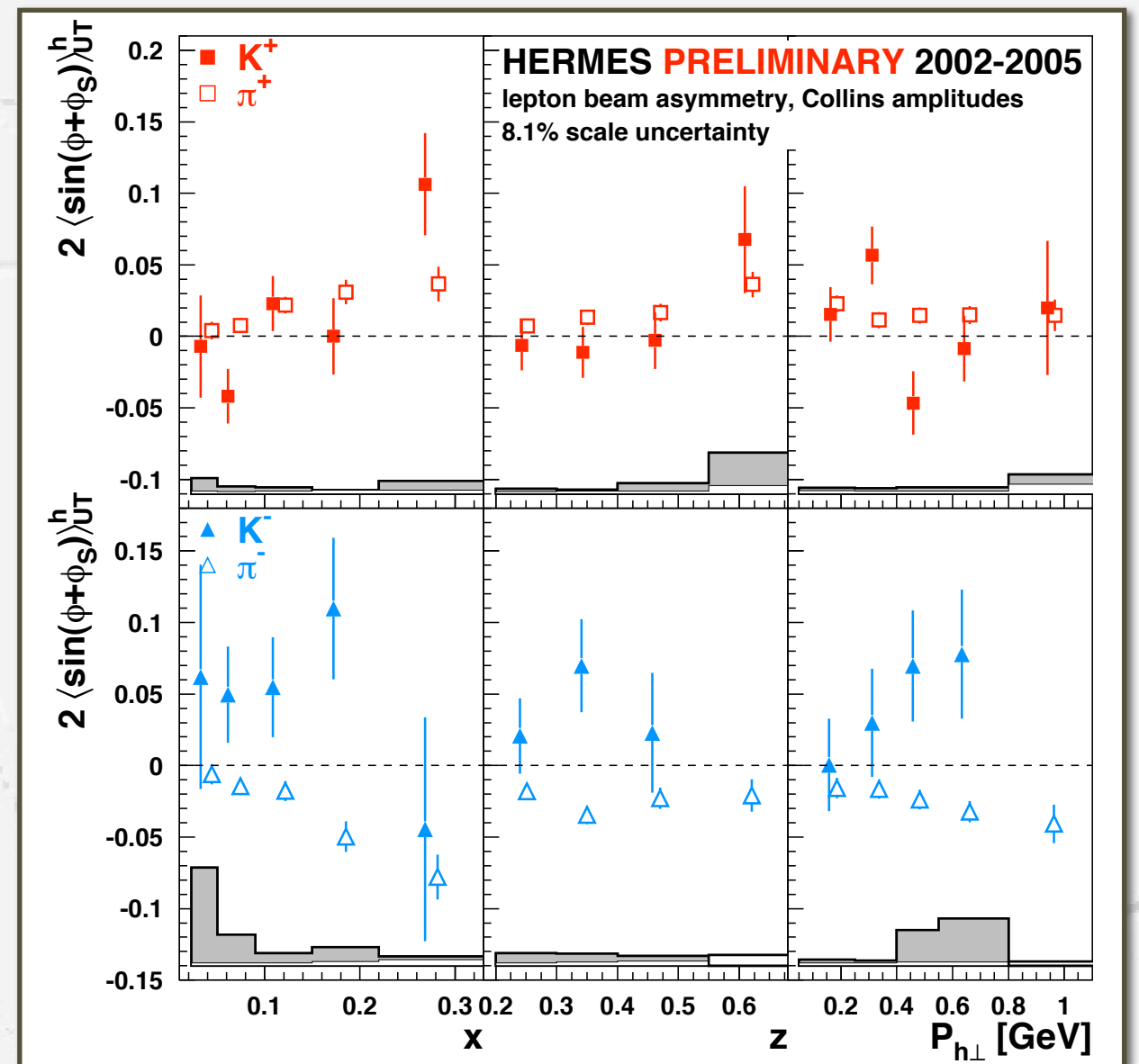
[†] [A. Airapetian et al, Phys. Rev. Lett. 94 (2005) 012002]

First Glimpse at Transversity

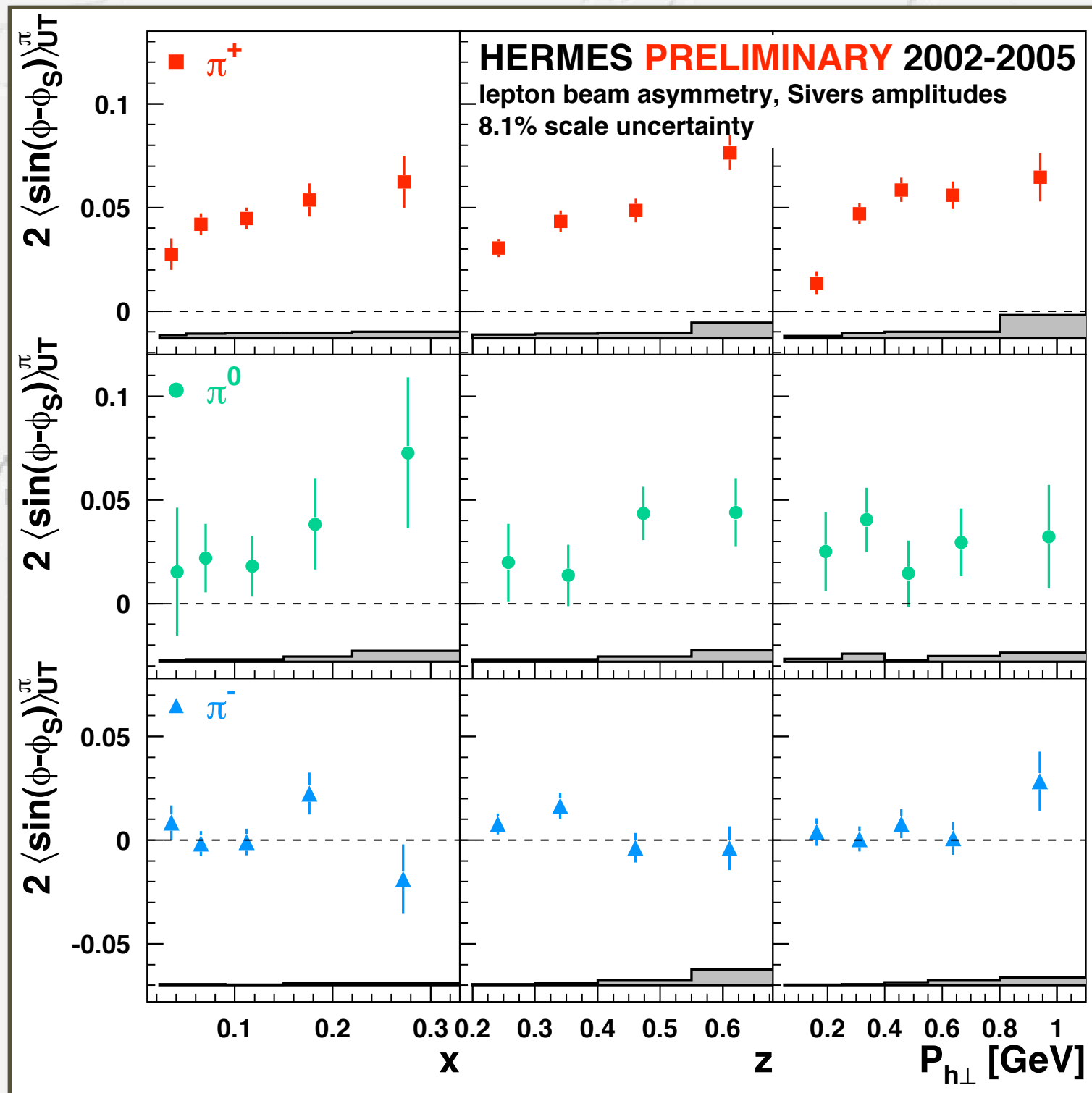


Collins Amplitudes for Kaons

- none of the kaon amplitudes significantly nonzero
- K^+ amplitudes not really different from π^+ amplitudes
- K^- amplitudes slightly positive, contrary to large negative π^- amplitudes
- K^- is pure “sea object”



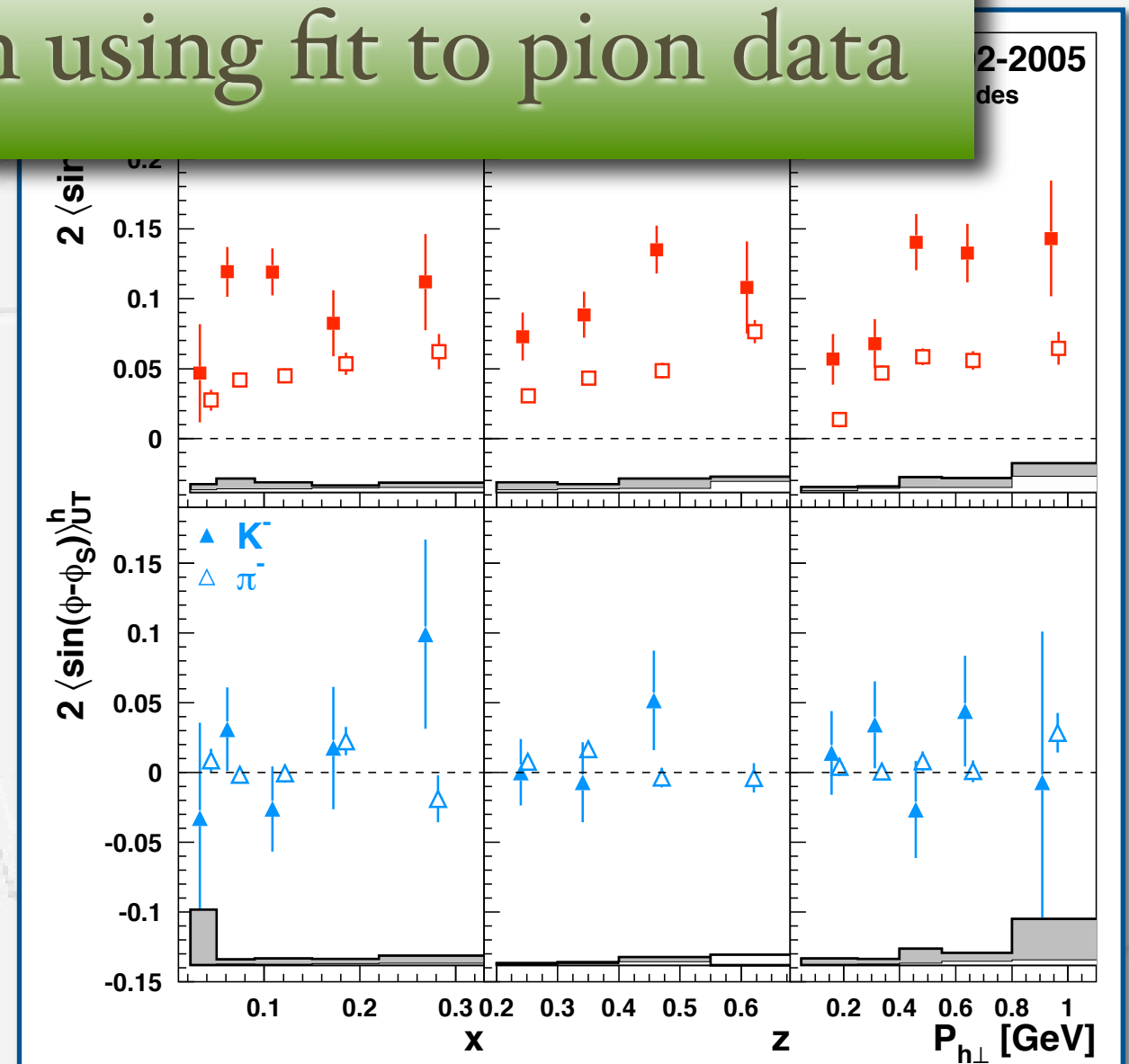
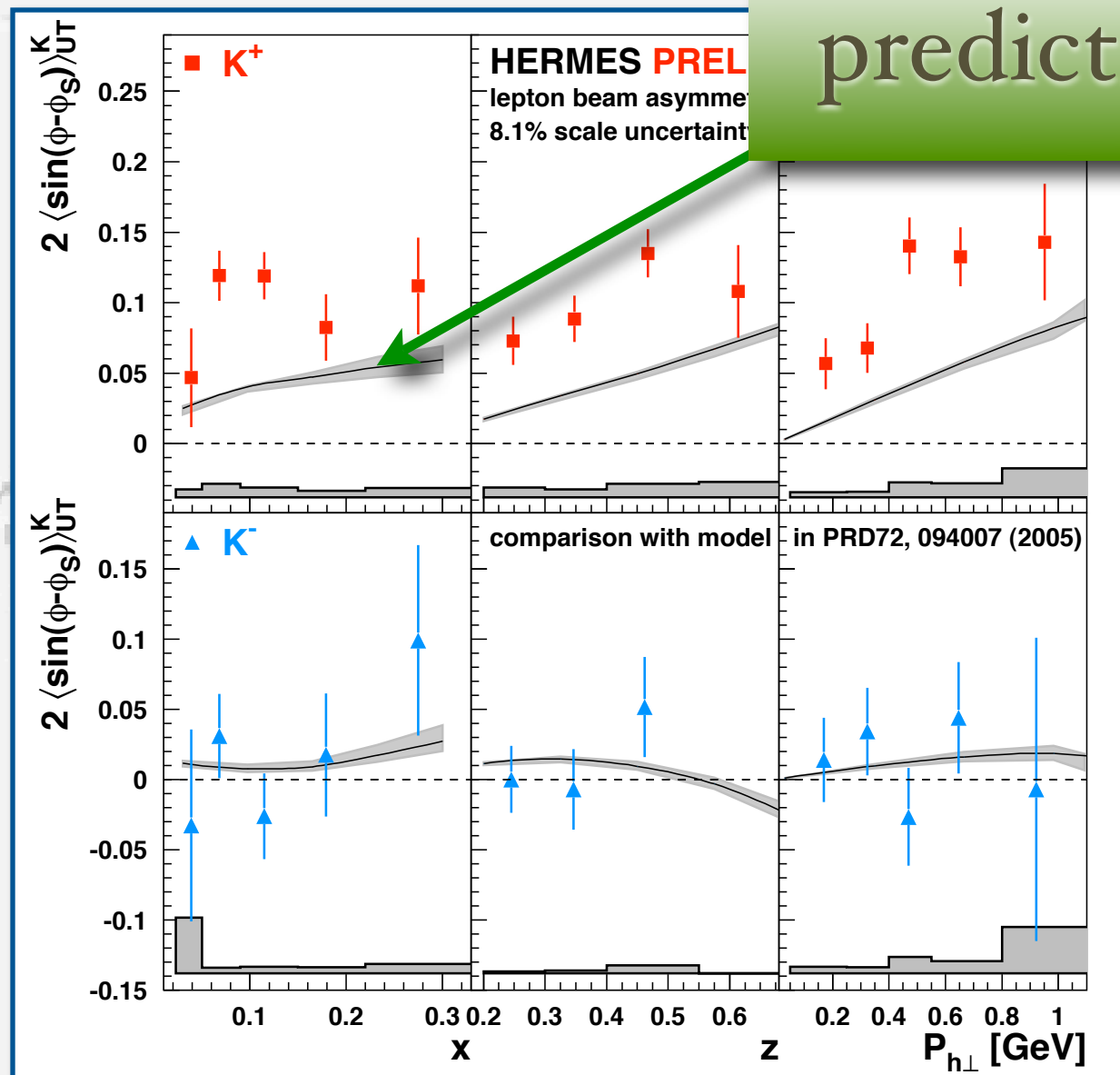
HERMES Sivvers Results



- ☑ first observation of T-odd Sivvers effect in SIDIS!
- ☑ u-quark dominance suggests sizeable u-quark orbital motion

The Intriguing Kaon Amplitudes

prediction using fit to pion data



non-trivial role of sea quarks!

Exclusive Reactions

Accessing Generalized Parton Distributions

Angular Momentum and GPDs

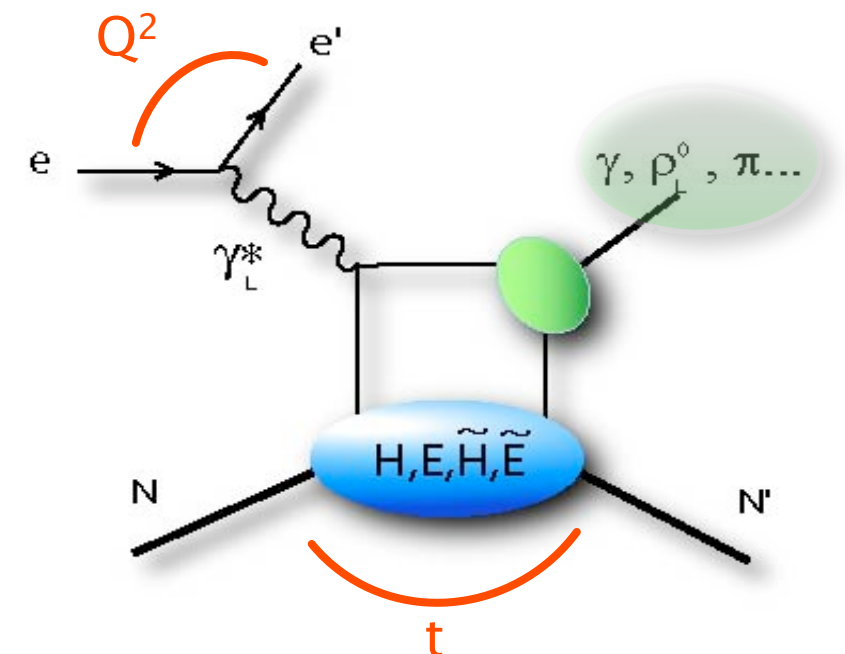
1997: Ji Relation for Nucleon Spin

$$\frac{1}{2}\hbar = \lim_{t \rightarrow 0} \underbrace{\sum_q \int dx [H^q(x, \xi, t) + E^q(x, \xi, t)]}_{J^q} + \lim_{t \rightarrow 0} \underbrace{\int dx [H^g(x, \xi, t) + E^g(x, \xi, t)]}_{J^g}$$

- “Ji’s Recipe” provides way to measure angular momenta
- involves moment over new class of PDFs: Generalized PDFs
- at leading twist there are 8 GPDs: $E, H, \tilde{E}, \tilde{H}, E_T, H_T, \tilde{E}_T, \tilde{H}_T$
- but only 2 of them needed for Ji’s sum-rule recipe: E, H
- GPDs provide info about transverse position and long. mom.

GPDs in Exclusive Reactions

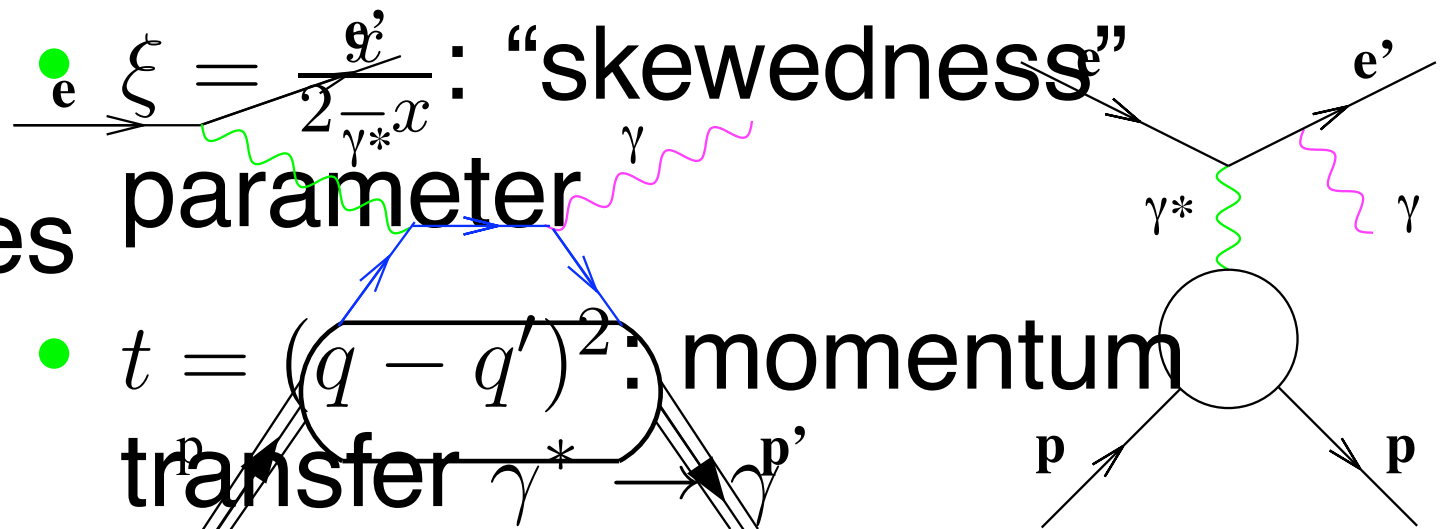
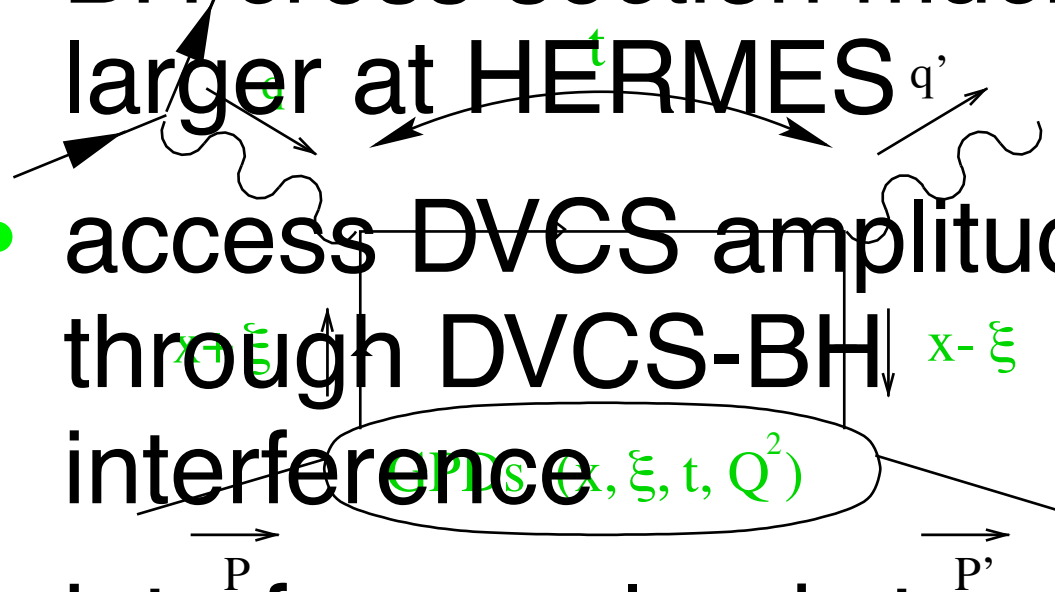
- GPDs involve off-forward matrix elements
- moments give Form Factors, e.g., $\int dx H^q(x, \xi, t) = F_1^q(t)$
- Forward limit give ordinary PDFs, e.g., $H^q(x, 0, 0) = f_1^q(x)$
- at HERMES accessed in exclusive reactions:
 - Deeply Virtual Compton Scattering (DVCS)
 - Exclusive Vector-Meson Production
 - Exclusive Pseudoscalar-Meson Production



Deeply Virtual Compton Scattering

$e + p \rightarrow e' + p' + \gamma$ clean & simple exclusive process:

- DVCS and Bethe-Heitler (BH) indistinguishable
- BH cross section much larger at HERMES
- access DVCS amplitudes through DVCS-BH interference
- interference leads to characteristic azimuthal asymmetries
- x : long. momentum fraction
- $\xi = \frac{q' - q}{2\gamma^* x}$: “skewedness” parameter
- $t = (q - q')^2$: momentum transfer
- $Q^2 = -q^2$



Azimuthal Asymmetries in DVCS

Interference DVCS & BH cause azimuthal asymmetries in cross-section:

- **Beam-charge asymmetry**

$$d\sigma(e^+, \phi) - d\sigma(e^-, \phi) \propto \text{Re}[F_1 \mathcal{H}] \cdot \cos \phi$$

- **Beam-spin asymmetry** $A_{\text{BSA}}(\phi)$:

$$d\sigma(\vec{e}, \phi) - d\sigma(\vec{e}, \phi) \propto \text{Re}[F_1 \mathcal{H}] \cdot \sin \phi$$

- **Long. target-spin asymmetry** $A_{UL}(\phi)$:

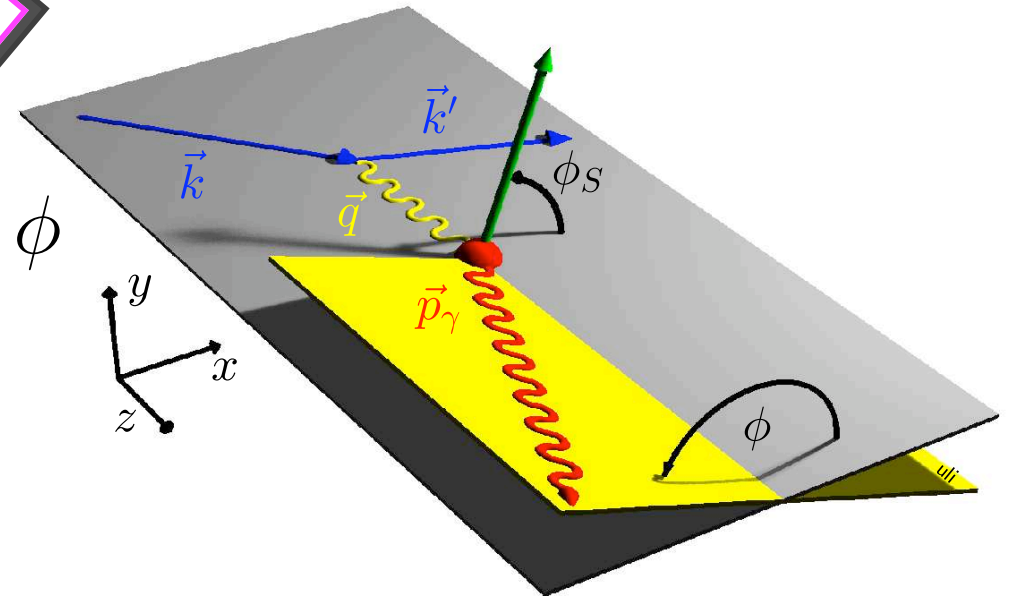
$$d\sigma(\vec{P}, \phi) - d\sigma(\vec{P}, \phi) \propto \text{Im}[F_1 \tilde{\mathcal{H}}] \cdot \sin \phi$$

- **Transversely polarized target-spin asymmetry** $A_{UT}(\phi, \phi_s)$ [TTSA]:

$$d\sigma(\phi, \phi_s) - d\sigma(\phi, \phi_s + \pi) \propto \text{Im}[F_2 \mathcal{H} - F_1 \mathcal{E}] \cdot \sin(\phi - \phi_s) \cos \phi \\ + \text{Im}[F_2 \tilde{\mathcal{H}} - F_1 \xi \tilde{\mathcal{E}}] \cdot \cos(\phi - \phi_s) \sin \phi$$

(F_1, F_2 are the Dirac and Pauli form factors, calculable in QED)

($\tilde{\mathcal{H}}, \tilde{\mathcal{E}}, \dots$ Compton form factors involving GPDs $H, E, \dots$)



Only TTSA sensitive to E!

Constraining E - Transverse TSA

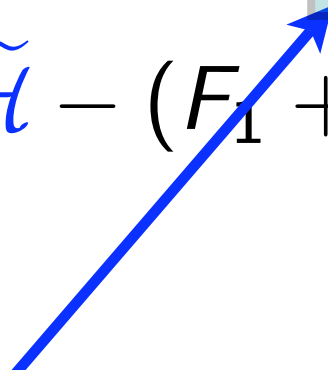
$$e^{\pm} p^{\uparrow} \longrightarrow e^{\pm} p \gamma$$

$$A_{\text{UT}}^{\mathcal{I}}(\phi, \phi_s) \propto [\text{d}\sigma^+(\phi, \phi_s) - \text{d}\sigma^-(\phi, \phi_s)] - [\text{d}\sigma^+(\phi, \phi_s + \pi) - \text{d}\sigma^-(\phi, \phi_s + \pi)]$$

$$A_{\text{UT}}^{\mathcal{I}}(\phi, \phi_s) \propto \text{Im}(F_2 \mathcal{H} - F_1 \mathcal{E}) \sin(\phi - \phi_s) \cos \phi$$

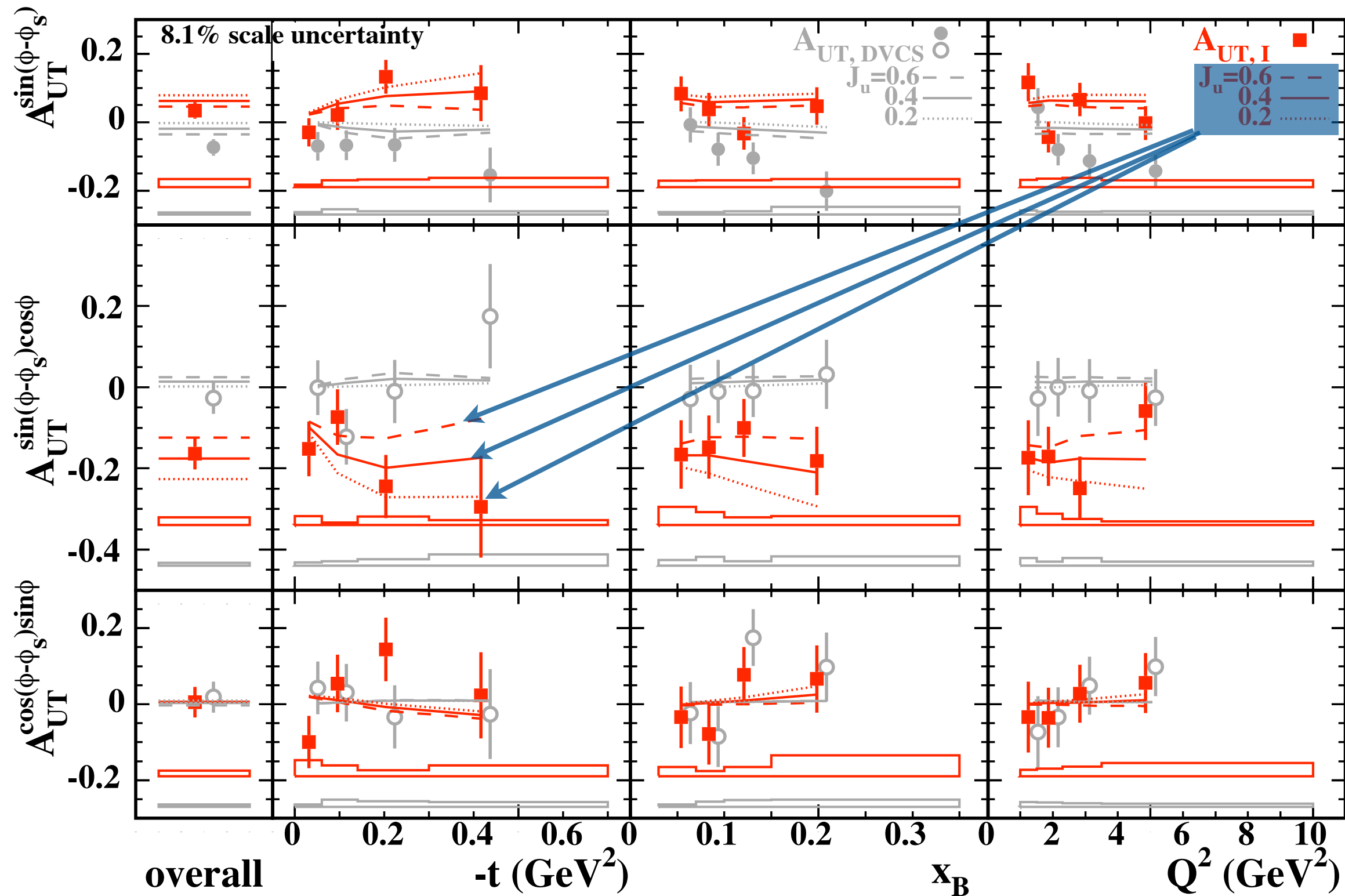
$$+ \text{Im}\left(F_2 \tilde{\mathcal{H}} - (F_1 + \xi F_2) \tilde{\mathcal{E}}\right) \cos(\phi - \phi_s) \sin \phi$$

sensitive to E



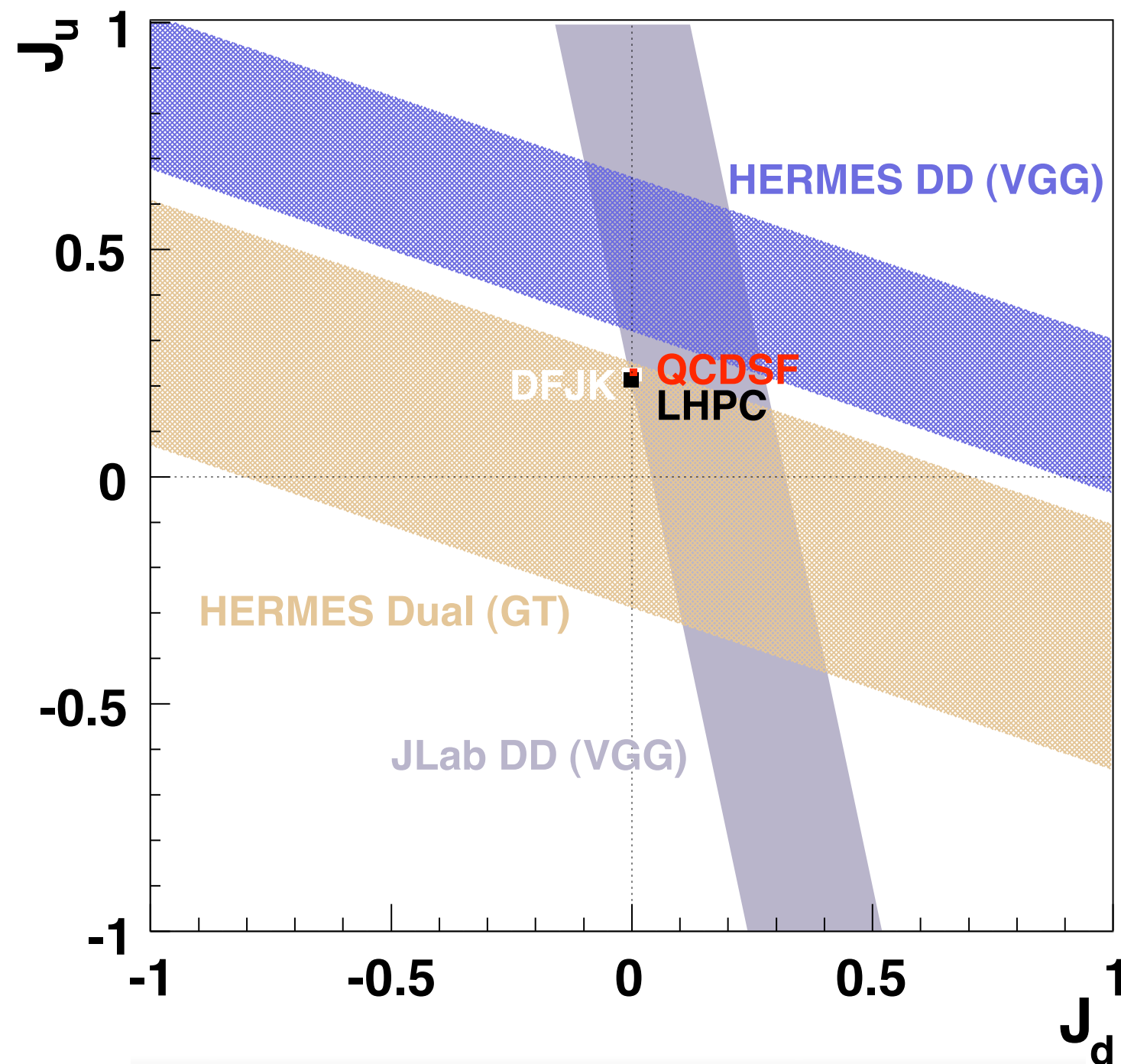
Amplitudes of TTSA

A. Airapetian et al., arXiv:0802.2499



Model-Dependent Extraction of J_u vs. J_d

$$\chi^2(J_u, J_d) = \left(A_{UT,I}^{\sin(\phi-\phi_s) \cos n\phi} |_{\text{exp}} - A_{UT,I}^{\sin(\phi-\phi_s) \cos n\phi} |_{\text{theo}}(J_u, J_d) \right)^2 / (\delta A_{\text{stat}}^2 + \delta A_{\text{sys}}^2)$$



- J_u and J_d are free param's in GPD models:

- **Double Distribution (DD):**
(Phys.Rev. D60 (1999) 094017)

$$J_u + J_d/2.8 = 0.49 \pm 0.17$$

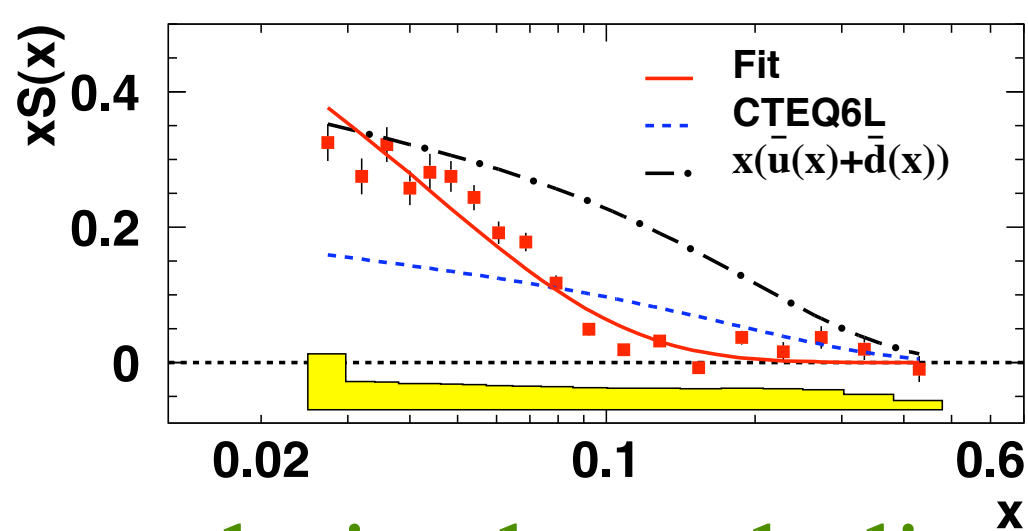
- **Dual Parameterization:**
(hep-ph/0207153 and Phys.Rev. D74 (2006) 054027)

$$J_u + J_d/2.8 = -0.02 \pm 0.27$$

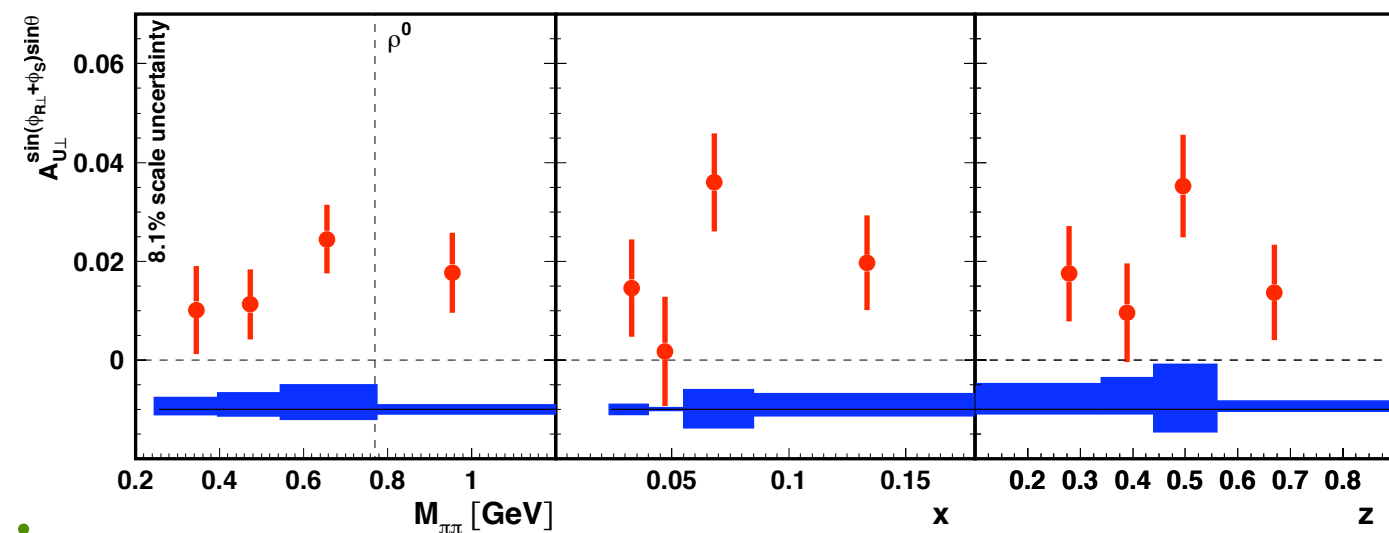
- **Bands are 1σ constraint on J_u vs. J_d**

A. Airapetian et al., JHEP 0806:066,2008

Conclusions

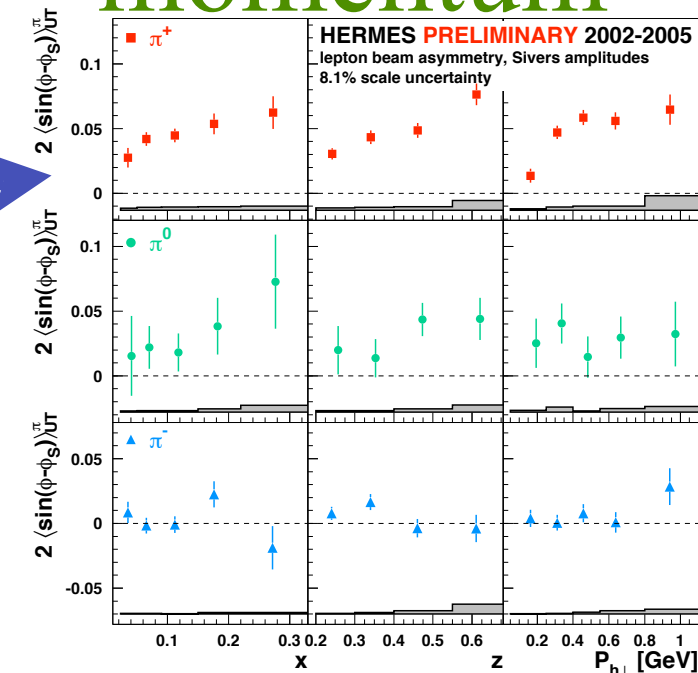


unpolarized quark distributions

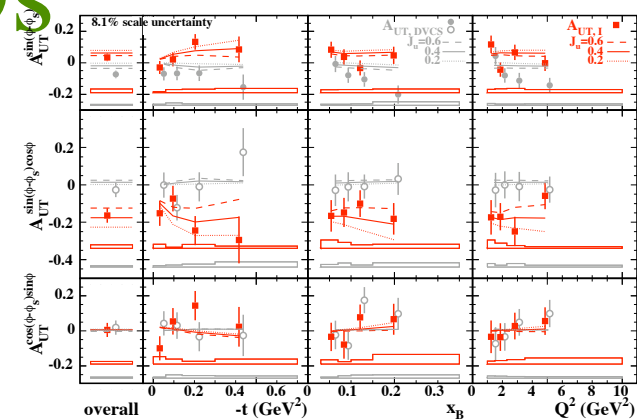


transversity

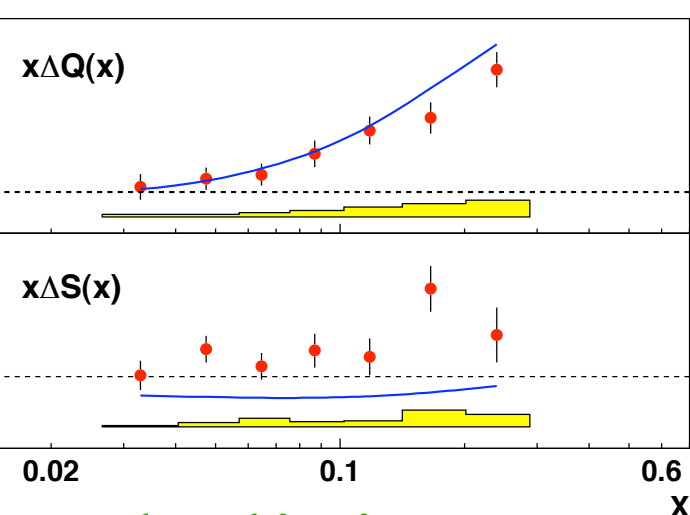
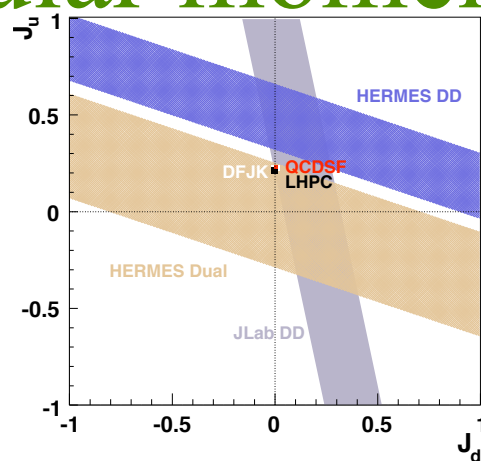
orbital angular momentum



GPDs



angular momentum

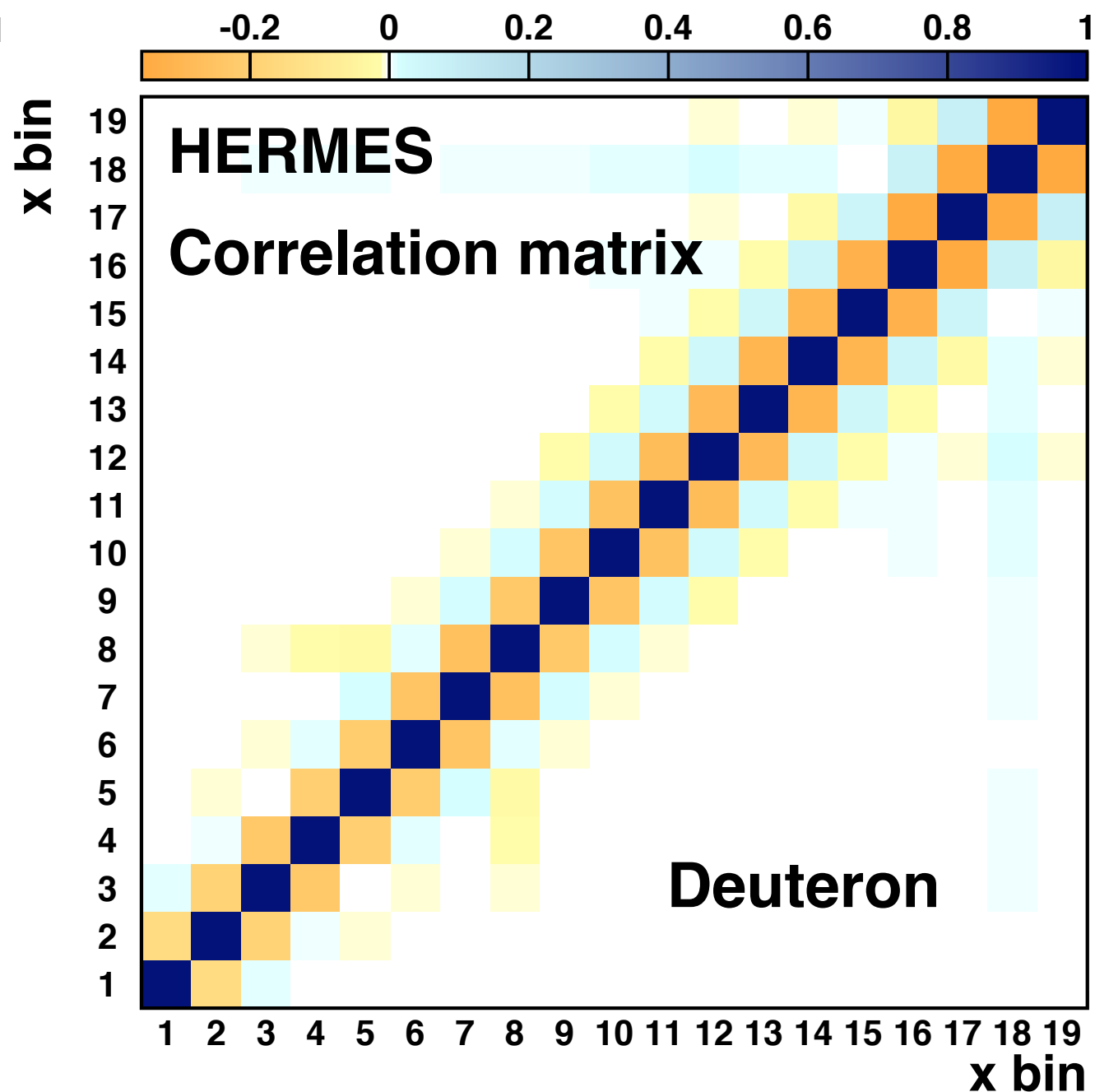
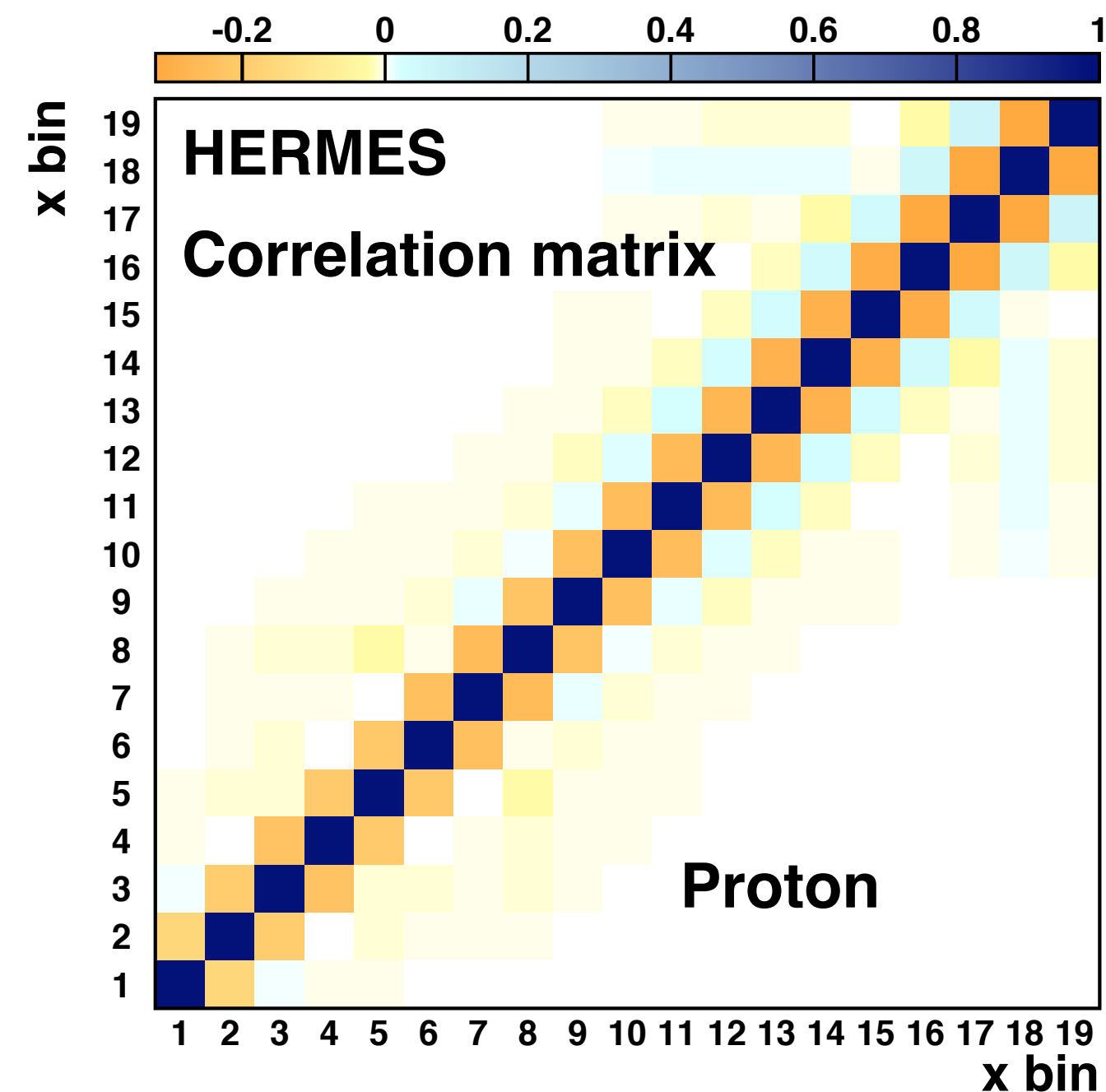


helicity distributions

hermes

Backup Slides

Statistical Correlations



Unknown systematic correlations replaced by known
statistical correlations through *unfolding*

Partial Moments of Polarized Quark Distributions

TABLE I: First moments of various helicity distributions in the Bjorken x range 0.02–0.6 at a scale of $Q_0^2 = 2.5 \text{ GeV}^2$.

	Moments in measured range
ΔQ	$0.359 \pm 0.026(\text{stat.}) \pm 0.018(\text{sys.})$
ΔS	$0.037 \pm 0.019(\text{stat.}) \pm 0.027(\text{sys.})$
Δq_8	$0.285 \pm 0.046(\text{stat.}) \pm 0.057(\text{sys.})$

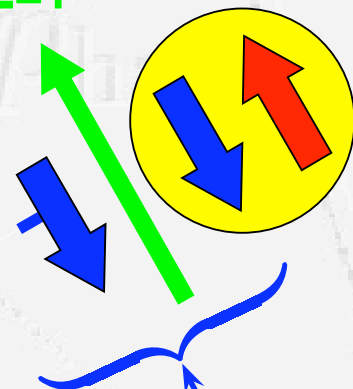
Artru Model for Collins Effect

transverse spin
of struck quark

(polarization component in lepton scattering
plane reversed by photoabsorption)

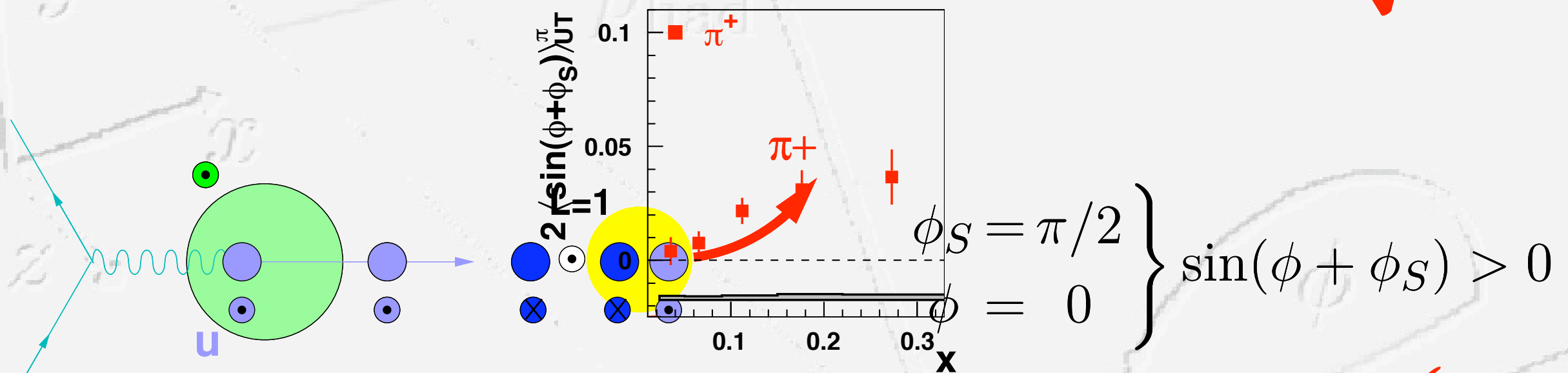
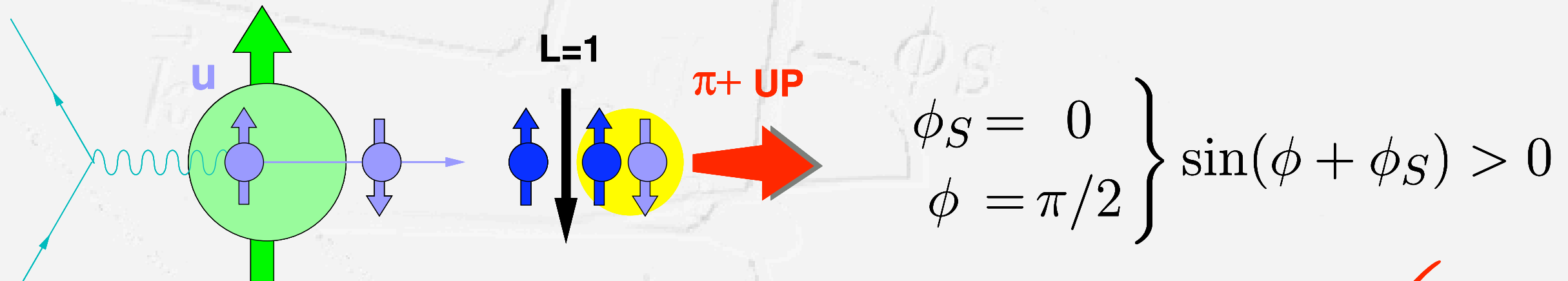
outgoing pion
deflected into page
(positive Collins FF)

$L=1$



$q\bar{q}$ -pair with vacuum
quantum numbers (3P_0 -state)

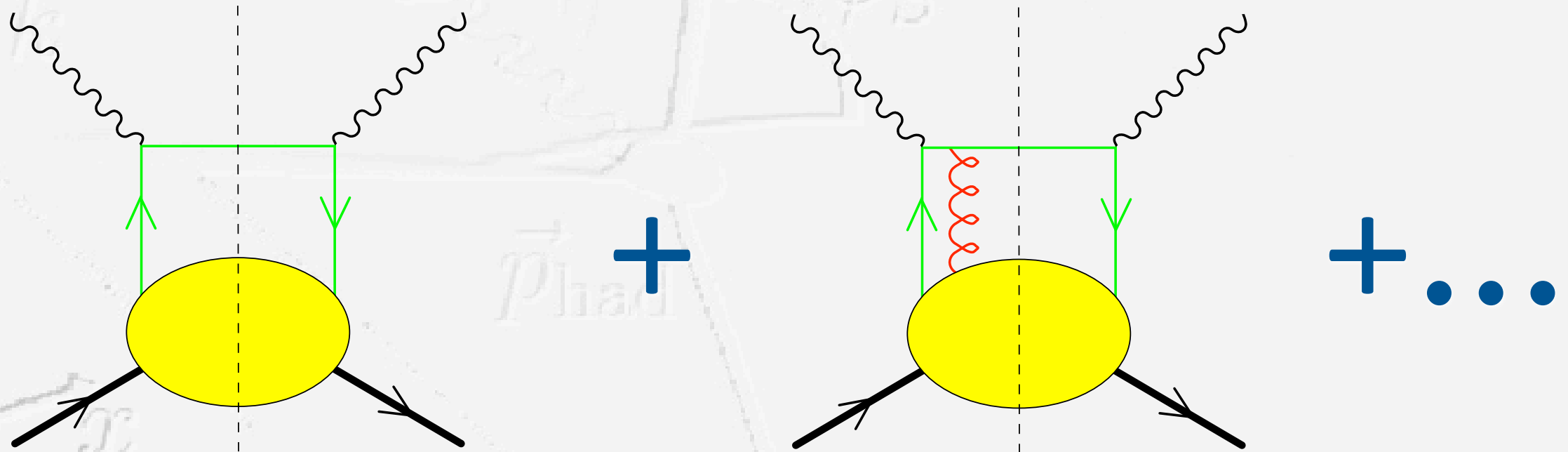
Artru Model vs. HERMES



Artru model and HERMES results in agreement!

The Sivvers Effect - A QCD Challenge?

- Distribution Function (DF) in handbag representation:

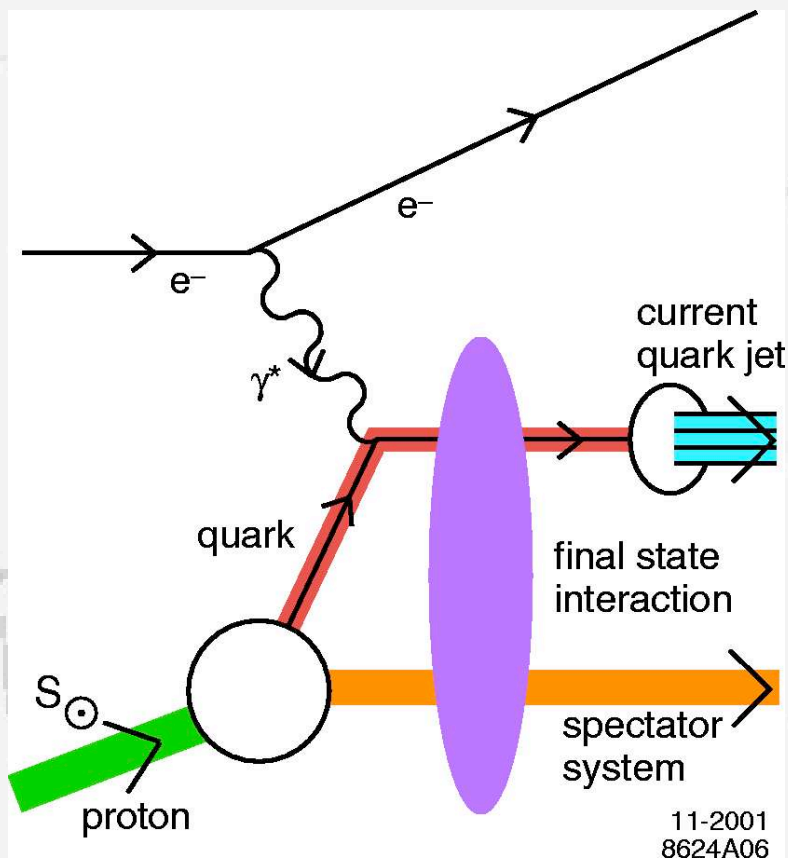


- Soft gluons needed for color gauge invariance
- Interference of amplitudes with different numbers of soft gluon exchanges possible
- represent color field of remnant seen by outgoing quark

The Sivers Effect

Thanks to Brodsky, Hwang & Schmidt:

- quark rescattering via soft gluon exchange
- correlates transverse spin of nucleon with direction of struck quark
- requires orbital angular momentum!
- leading twist!



Thanks to Collins, Ji, Yuan, Belitzky ...:

- ➡ Soft gluon is model for gauge link needed for gauge invariance
- ➡ Gauge links provide necessary complex phase for interference
- ➡ T-Symmetry of QCD requires **opposite sign of Sivers function in DIS and DY**
- ➡ slightly different approach by Burkardt using impact parameter dependent PDF's ("chromodynamic lensing")

“Chromodynamic Lensing”

approach by M. Burkardt:

[hep-ph/0309269]

spatial distortion of q-distribution

(obtained using anom. magn. moments
& impact parameter dependent PDFs)

+ attractive QCD potential

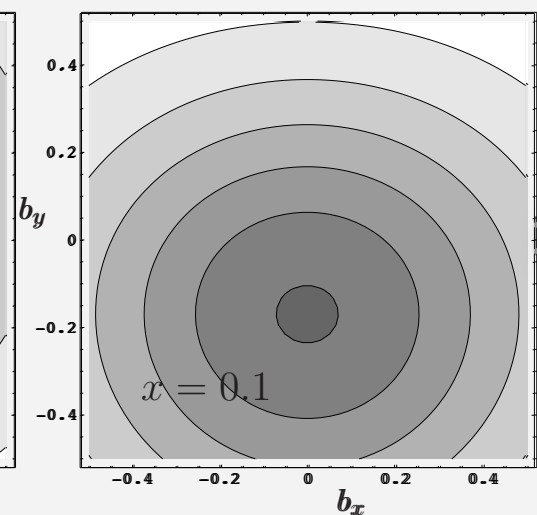
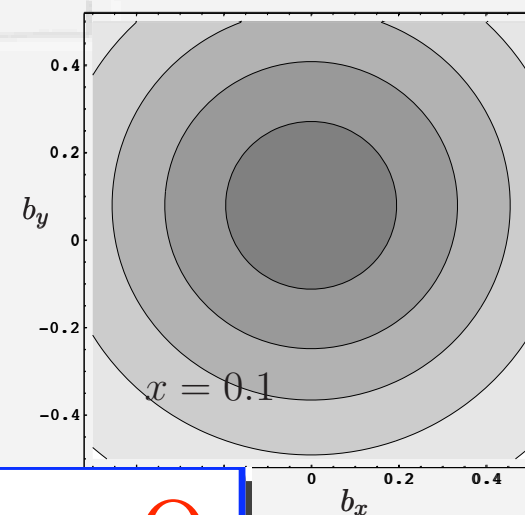
(gluon exchange)

⇒ transverse asymmetry

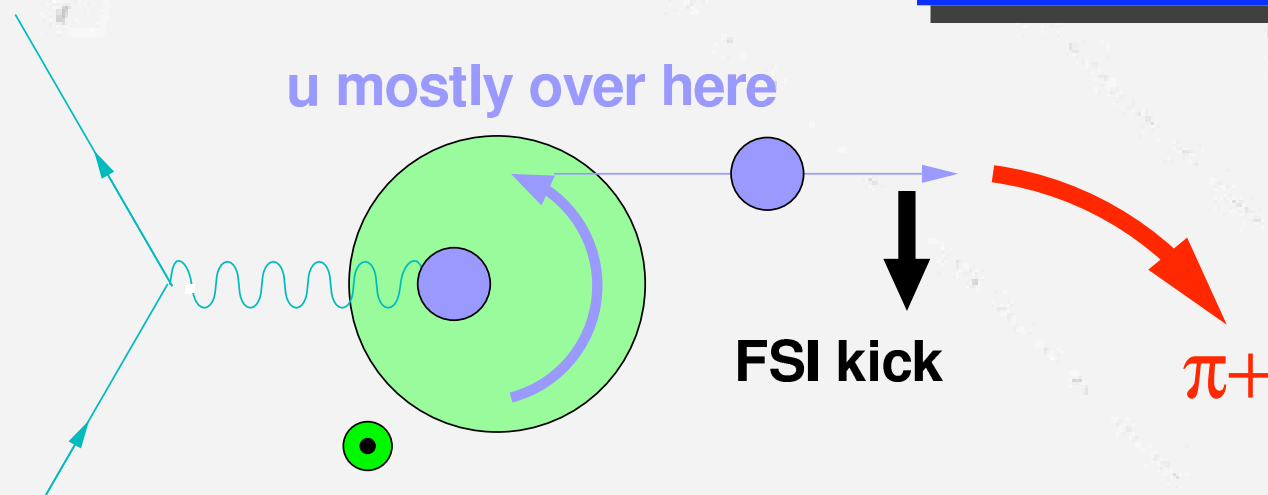
$$L_z^u > 0$$

$u_X(x, \mathbf{b}_\perp)$

$d_X(x, \mathbf{b}_\perp)$



u mostly over here

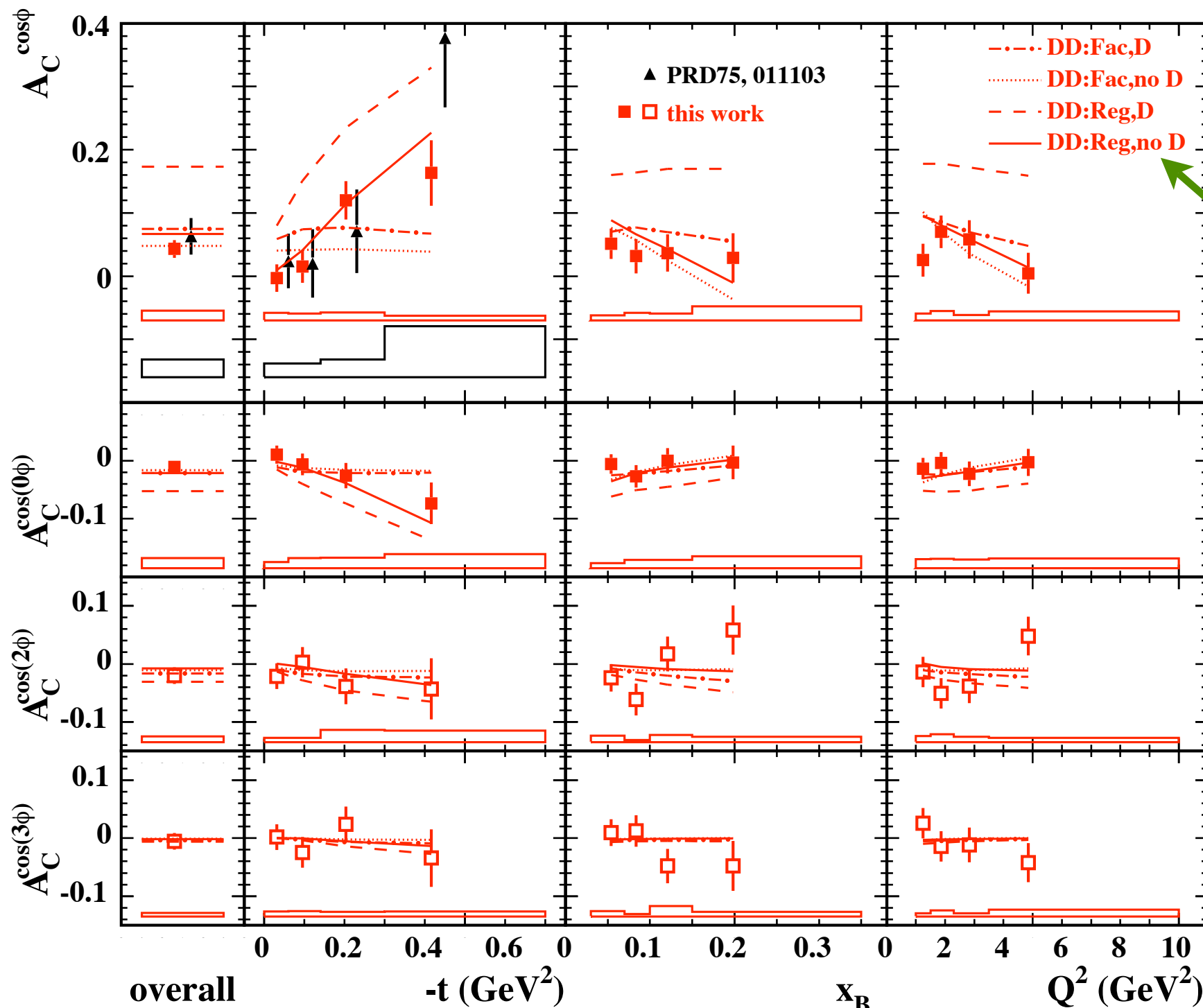


$$\left. \begin{array}{l} \phi_S = \pi/2 \\ \phi = \pi \end{array} \right\} \sin(\phi - \phi_S) > 0$$



Constraining GPD H

Beam-charge asymmetry $A_C(\phi) \propto \text{Re}[F_1 \mathcal{H}] \cdot \cos \phi$
(extracted from same data set as used for TTSA)



A. Airapetian et al.,
arXiv:0802.2499

clearly favored:

- double distribution
- factorized (Regge-inspired) Ansatz for t-independent part
- no “D-term”

(cf. PRD60 (1997))

JPS Meeting, 26-March-2008