

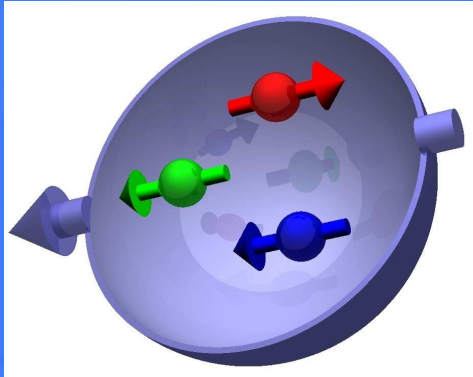
Where is the Spin of the Nucleon hidden?

E.C. Aschenauer

DESY-ZEUTHEN



The Spin Structure of the Nucleon



Naive Parton Model:

$$\Delta u_v + \Delta d_v = 1$$
$$\Rightarrow \Delta u_v = \frac{4}{3}, \Delta d_v = \frac{-1}{3}$$

BUT

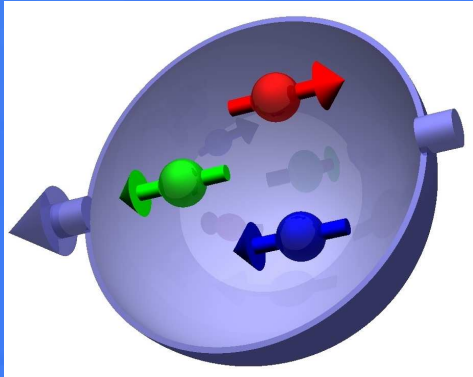
1988 EMC measured:

$$\Sigma = 0.123 \pm 0.013 \pm 0.019$$

$\Rightarrow$ Spin Puzzle

$$\frac{1}{2} = \frac{1}{2}(\Delta u_v + \Delta d_v)$$

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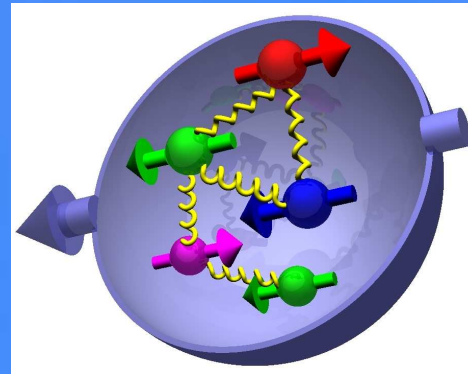
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F_2 from HERA tells:

Gluons are important !

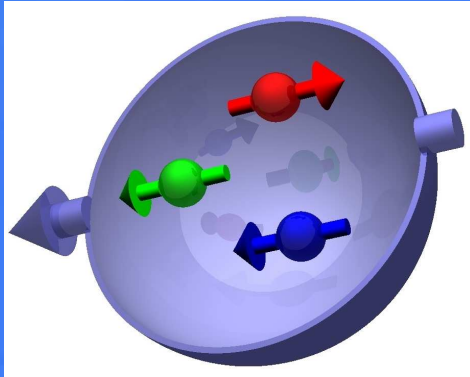
$\Rightarrow$ sea quarks Δq_s

$\Rightarrow \Delta G$

$$\frac{1}{2} = \frac{1}{2} (\Delta u_v + \Delta d_v + \underbrace{\Delta q_s}_{\Delta u_s, \Delta d_s, \Delta \bar{u}, \Delta \bar{d}, \Delta s, \Delta \bar{s}}) + \Delta G$$

$$\Delta u_s, \Delta d_s, \Delta \bar{u}, \Delta \bar{d}, \Delta s, \Delta \bar{s}$$

The Spin Structure of the Nucleon



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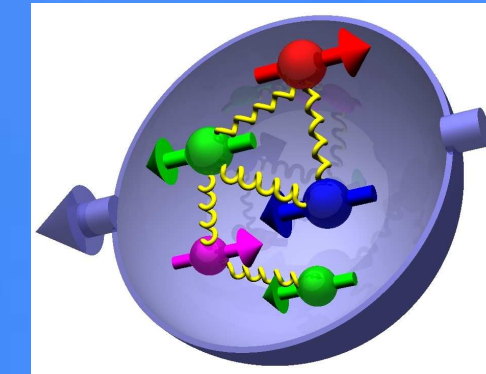
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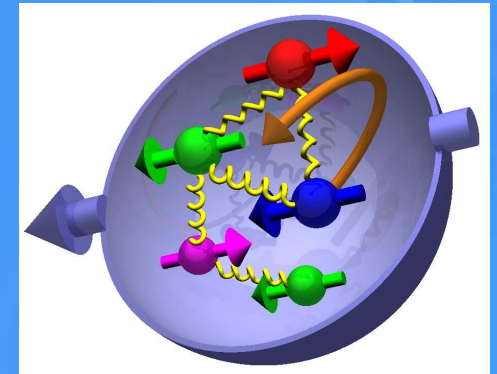


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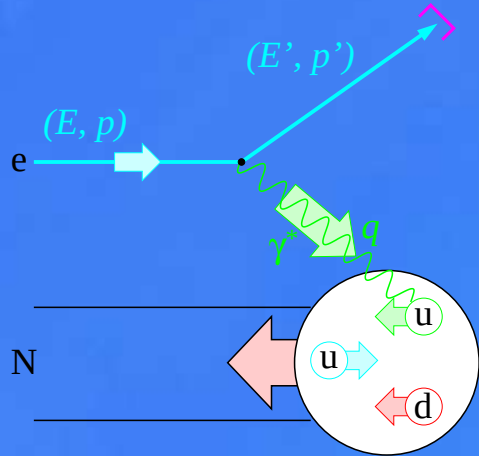
$\Rightarrow \Delta G$



Full description of J_q & J_g
needs
orbital angular momentum

$$\frac{1}{2} = \frac{1}{2} \underbrace{(\Delta u_v + \Delta d_v + \Delta q_s)}_{\Delta \Sigma} + \Delta G + L_q + L_g$$

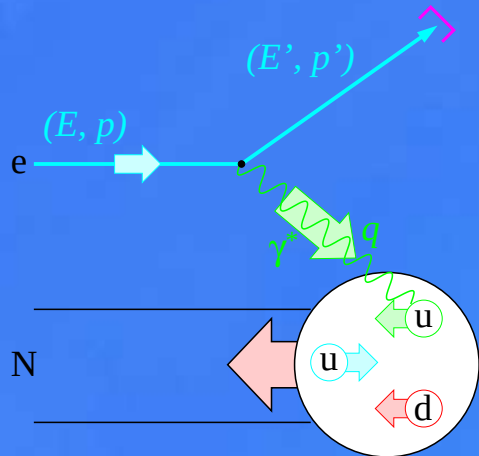
Inclusive Scattering:



detect scattered lepton

$$\begin{aligned}
 Q^2 &\stackrel{lab}{=} 4EE' \sin^2\left(\frac{\theta}{2}\right) & \nu &\stackrel{lab}{=} E - E' \\
 W^2 &\stackrel{lab}{=} M^2 + 2M\nu - Q^2 \\
 x &\stackrel{lab}{=} \frac{Q^2}{2M\nu} & y &\stackrel{lab}{=} \frac{\nu}{E} = \frac{p \cdot q}{p \cdot k}
 \end{aligned}$$

Inclusive Scattering:



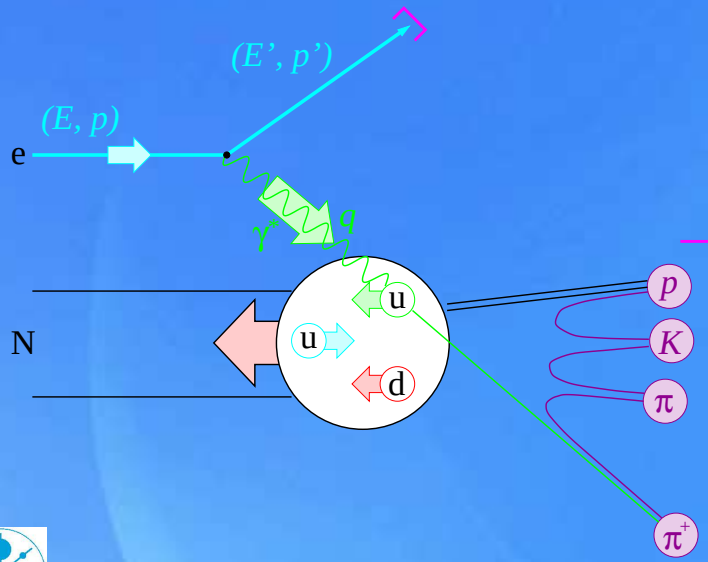
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$$Q^2 \stackrel{lab}{=} 4EE' \sin^2\left(\frac{\theta}{2}\right) \quad \nu \stackrel{lab}{=} E - E'$$

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$$x \stackrel{lab}{=} \frac{Q^2}{2M\nu} \quad y \stackrel{lab}{=} \frac{\nu}{E} = \frac{p \cdot q}{p \cdot k}$$

Semi-Inclusive Scattering:



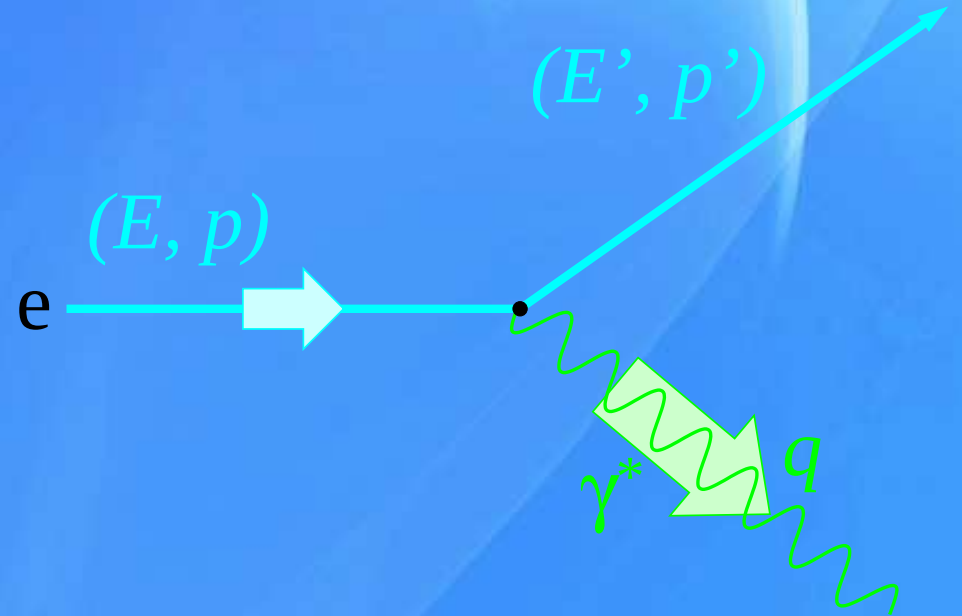
detect scattered lepton and produced hadrons

$$\eta = \frac{\log(E_{cm}^h - p_{||}^h)}{2 * (E_{cm}^h + p_{||}^h)}$$

$$z \stackrel{lab}{=} \frac{E_h}{\nu}$$

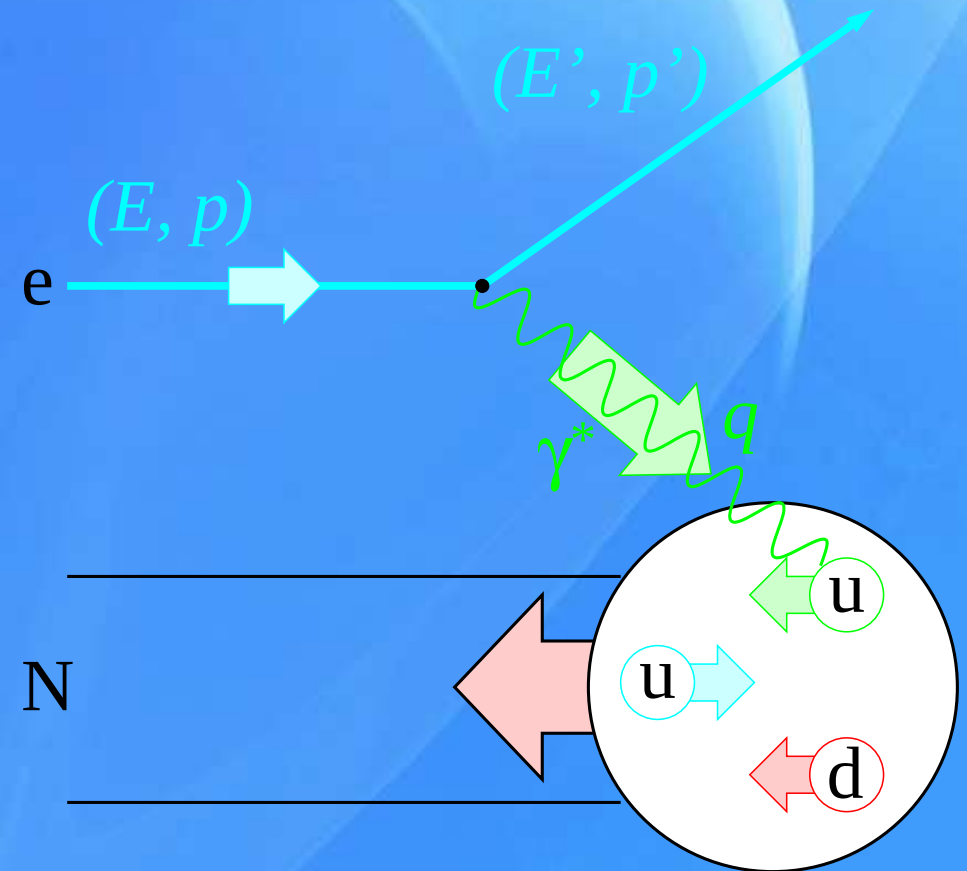
Experimental Prerequisites

- Longitudinally polarized lepton beam



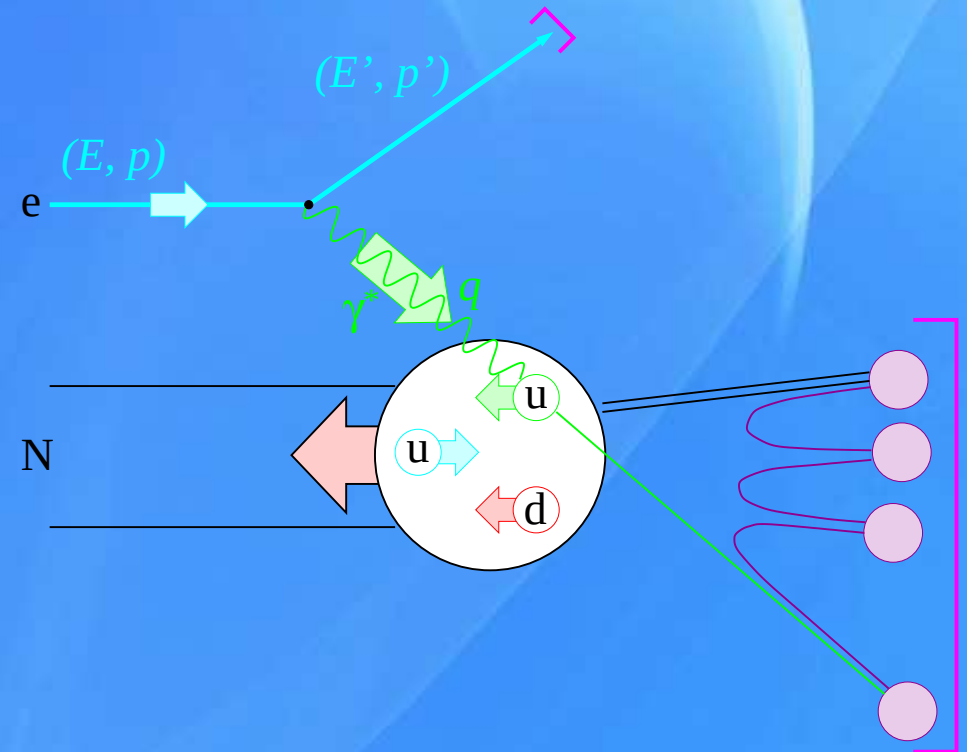
Experimental Prerequisites

- Longitudinally polarized lepton beam
- Longitudinally polarized nuclear target



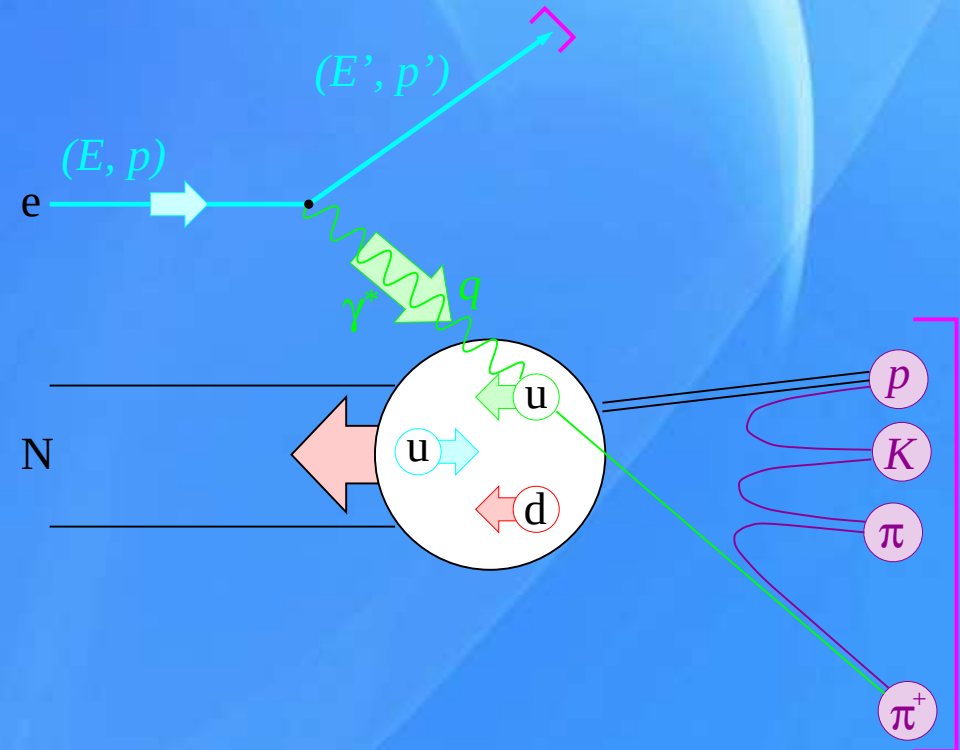
Experimental Prerequisites

- Longitudinally polarized lepton beam
- Longitudinally polarized nuclear target
- Large geometrical acceptance
 - small angles:
limited by synchrotron radiation
 - big angles:
limited by money



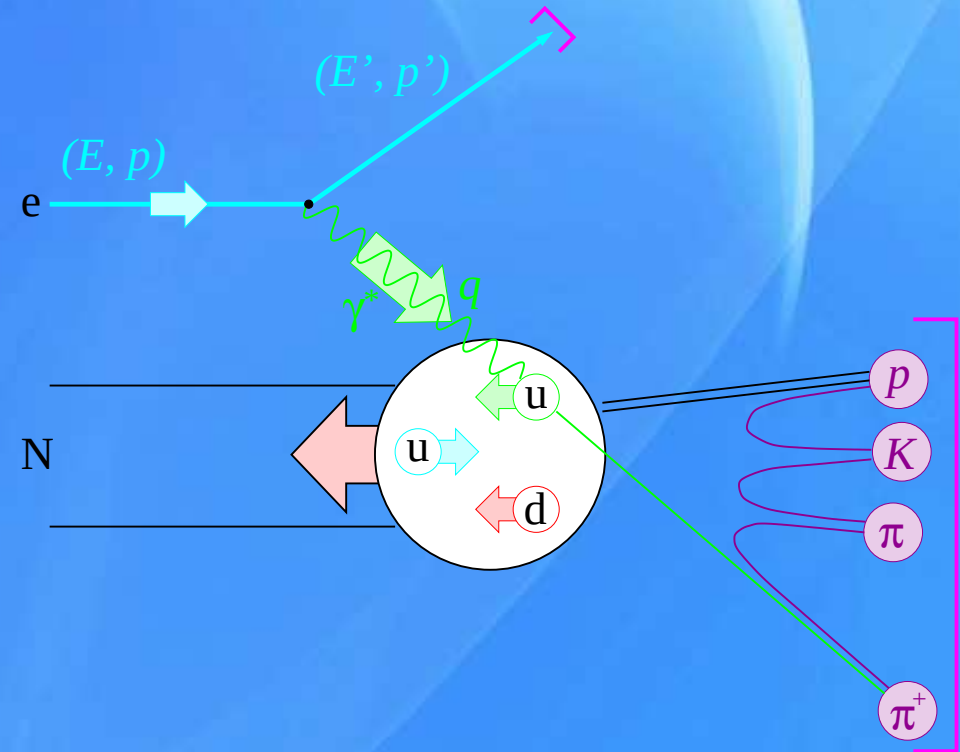
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- Good particle identification



Experimental Prerequisites

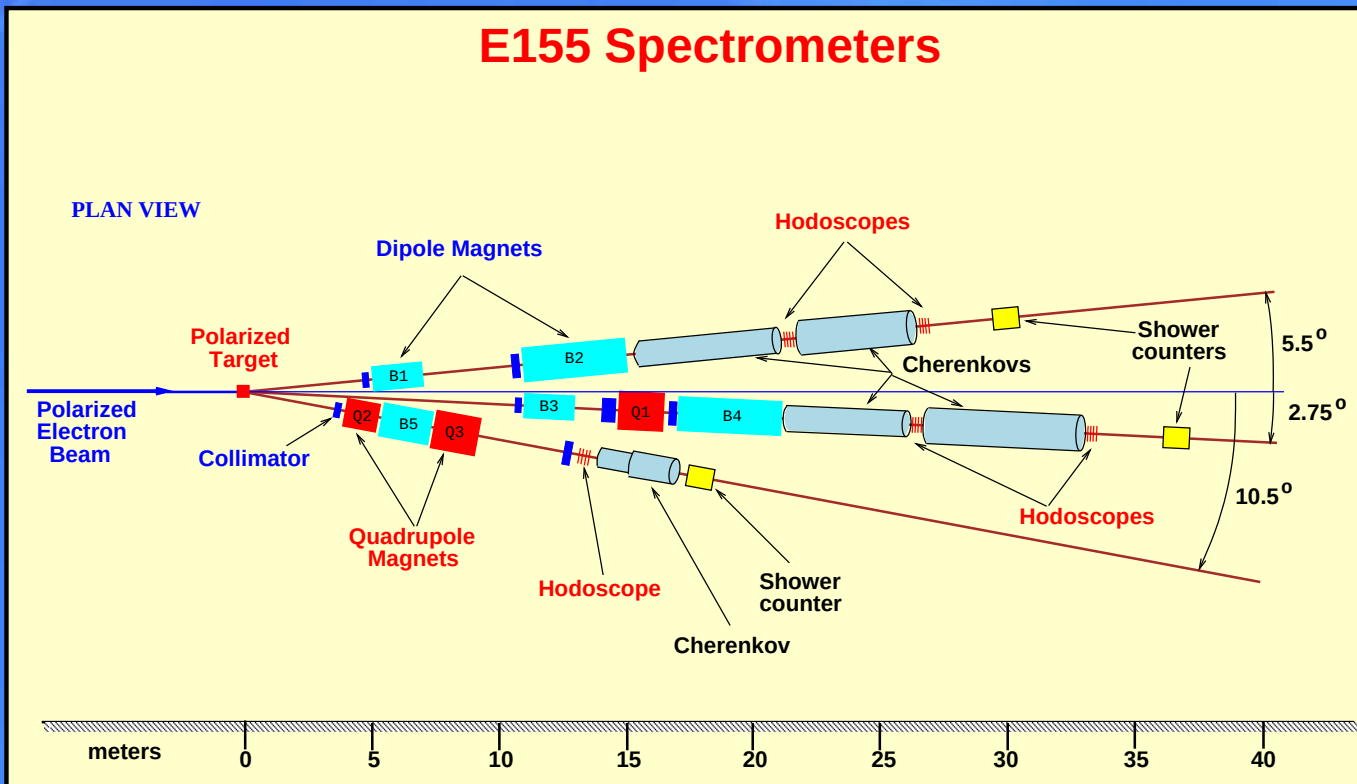
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so far only "fixed target" experiments at:



E155 Spectrometers



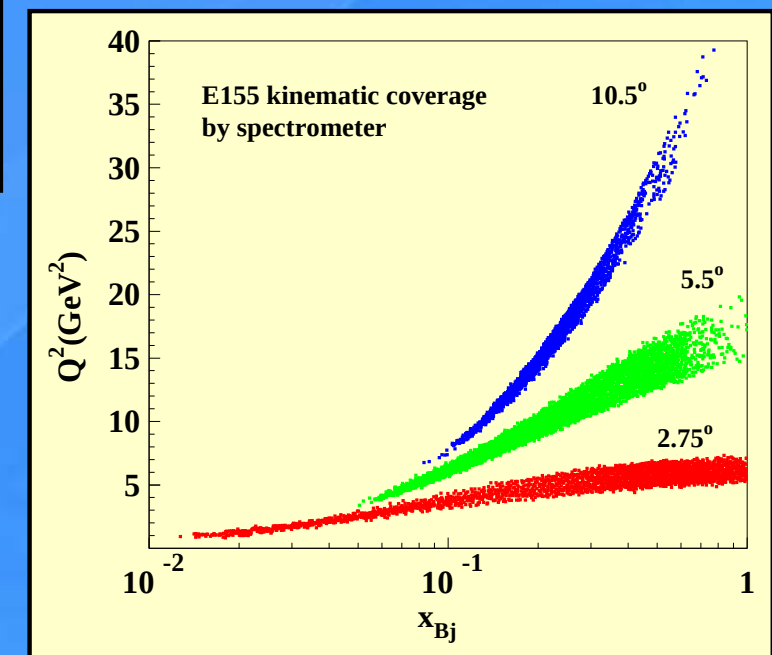
Inclusive
only

Scattering
only

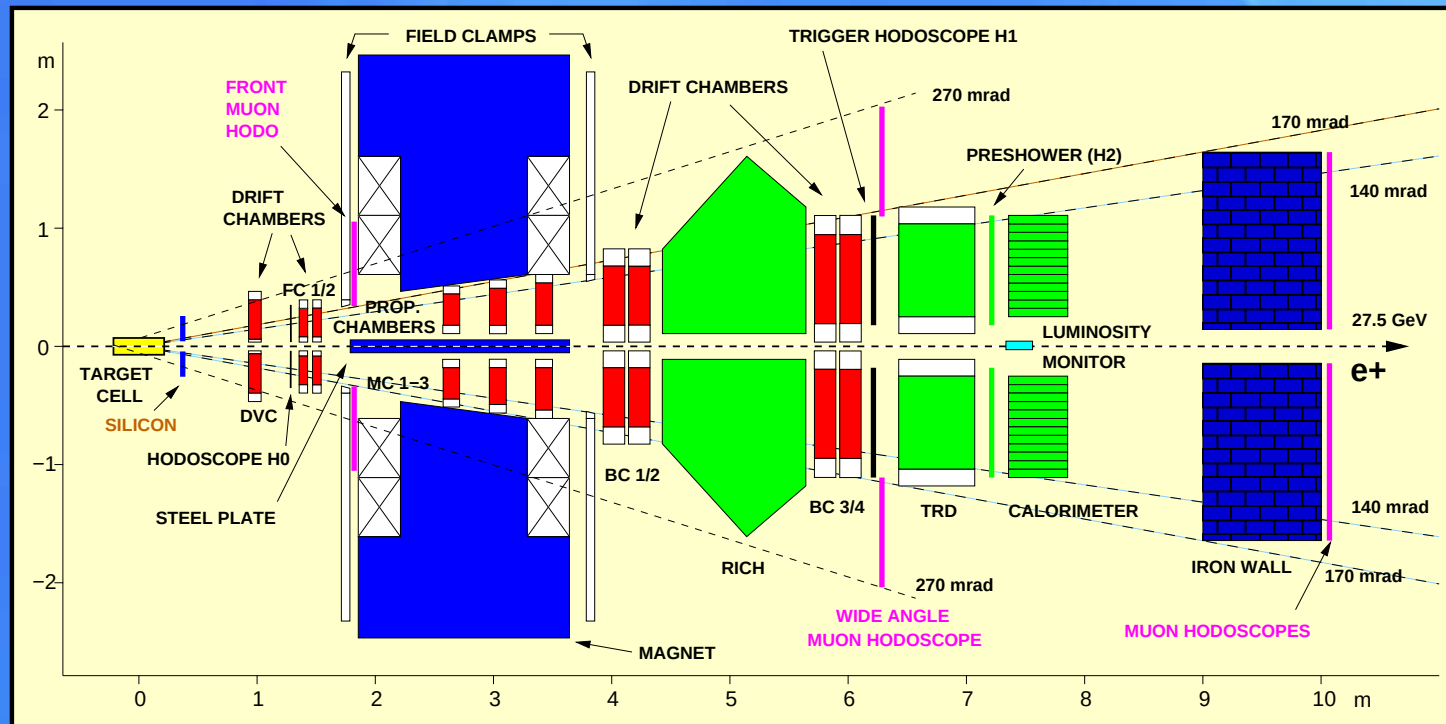
Reconstruction: very small angular coverage

Target: polarized solid state targets (NH_3 / LiD)
Big dilution factor

Particle ID: Lepton-Hadron separation



The HERMES Detector



Kinematic Range: $0.02 \leq x \leq 0.8$ at $Q^2 \geq 1\text{GeV}^2$ and $W \geq 2\text{GeV}$
 $\Theta_x \leq 175 \text{ mrad}$, $40 \text{ mrad} \leq \Theta_y \leq 140 \text{ mrad}$

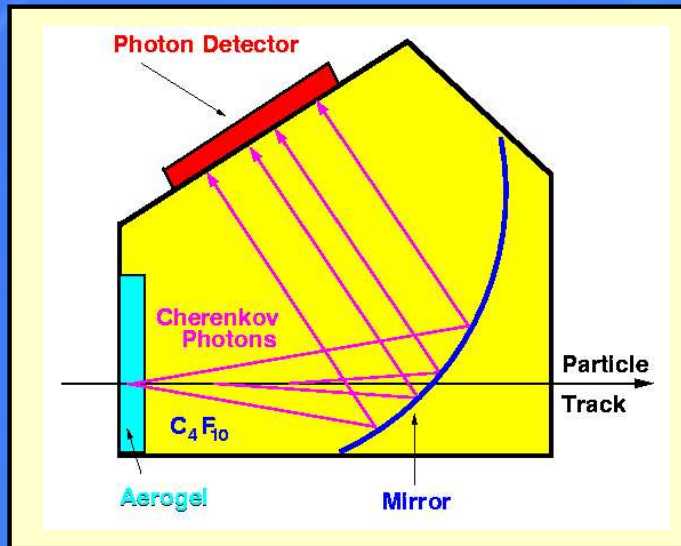
Reconstruction: $\delta p/p$ 1.0 - 2.0%, $\delta\Theta \leq 0.6\text{mrad}$

Internal Gas Target: $\vec{H}_e, \vec{D}, \vec{H}, \textcolor{red}{H}^\uparrow$ unpol: $\text{H}_2, \text{D}_2, \text{He}, \text{N}_2, \text{Ne}, \text{Ar}, \text{Kr}$

Particle ID: TRD, Preshower, Calorimeter

$\Rightarrow$ 1997: Čerenkov 1998 $\Rightarrow$: **RICH** + **Muon-ID**

The HERMES RICH

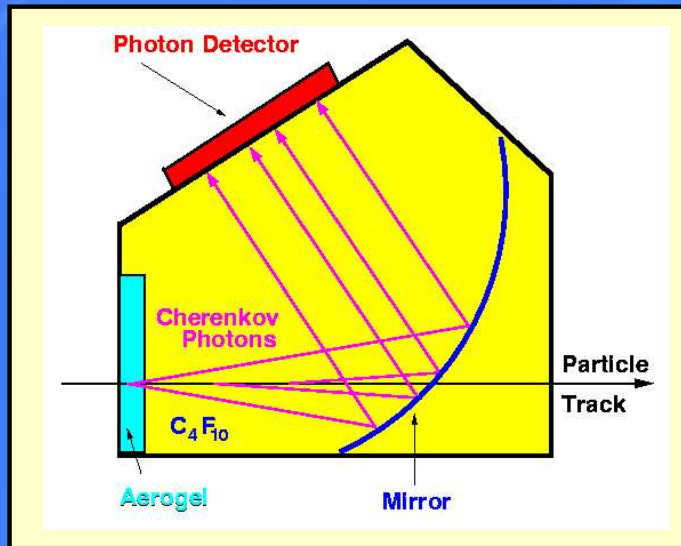


Dual Radiator RICH

Aerogel: $n = 1.03$

C₄F₁₀: $n = 1.0014$

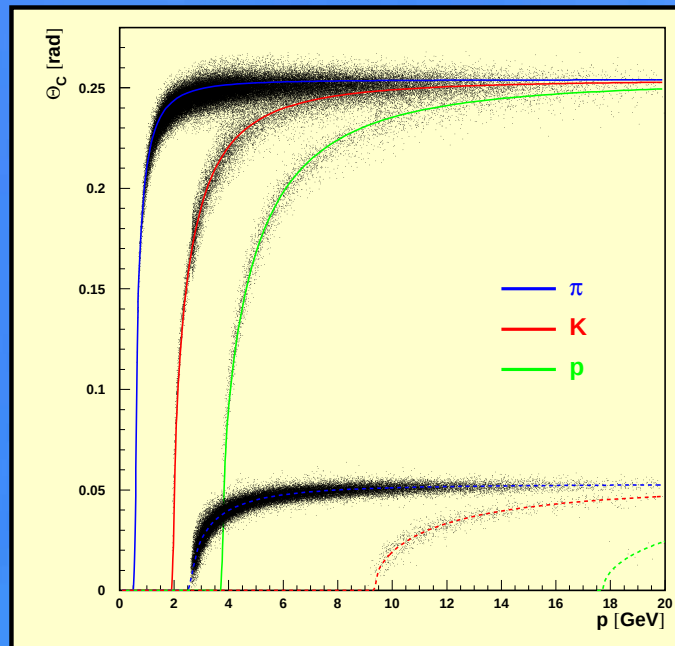
The HERMES RICH



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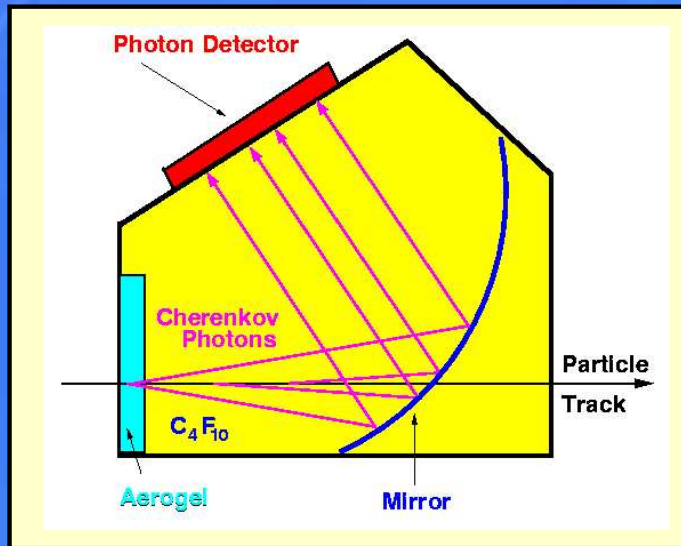
Particle identification:

$$\cos \Theta_c = \frac{1}{\beta n}$$

threshold behaviour:

$$p = \frac{m \beta c}{\sqrt{1 - \beta^2}}$$

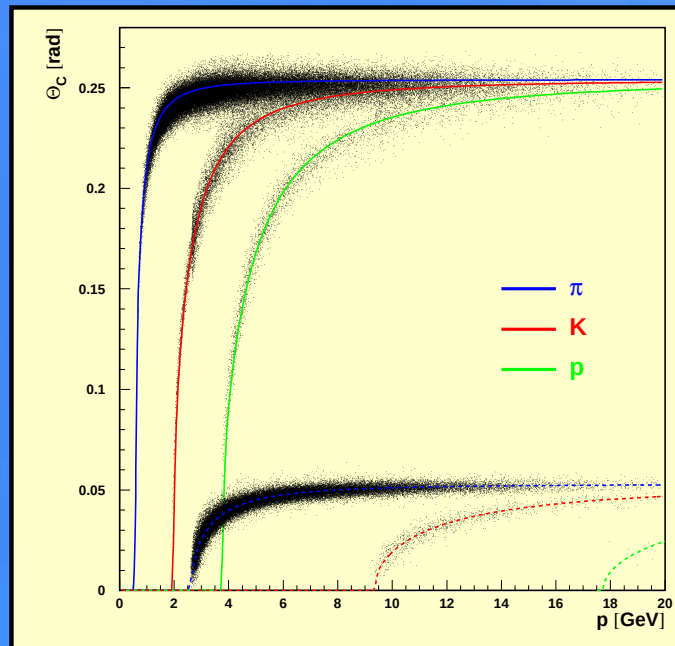
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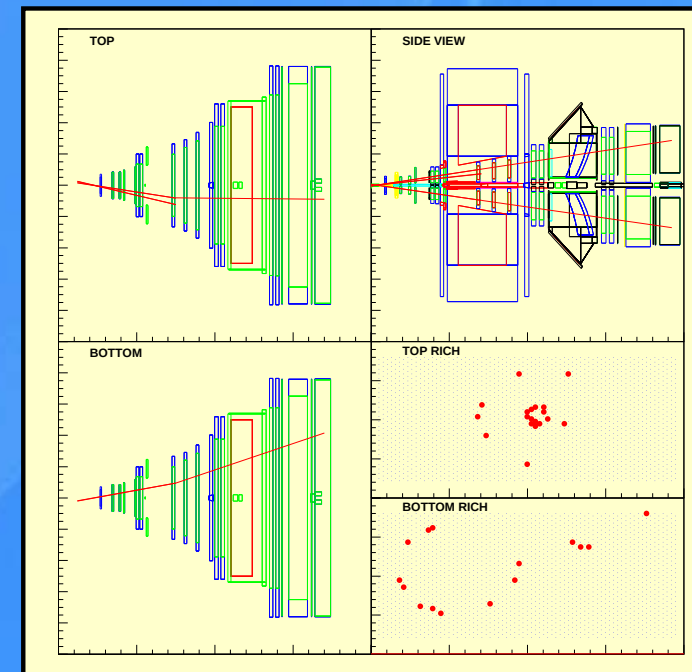


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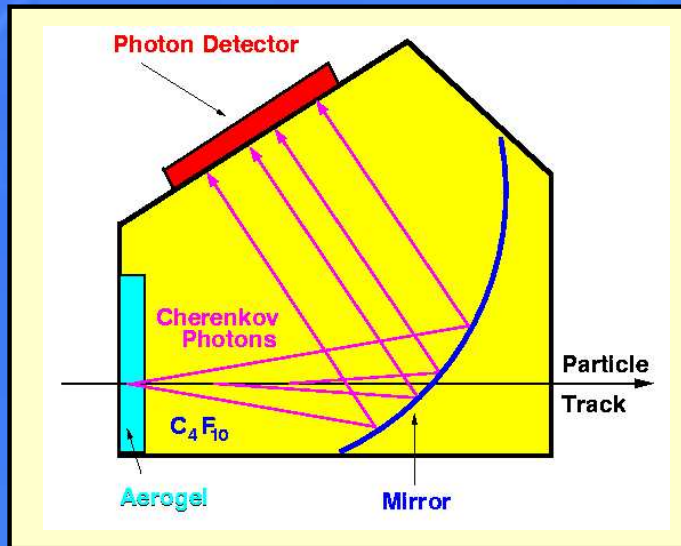
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Real πK event

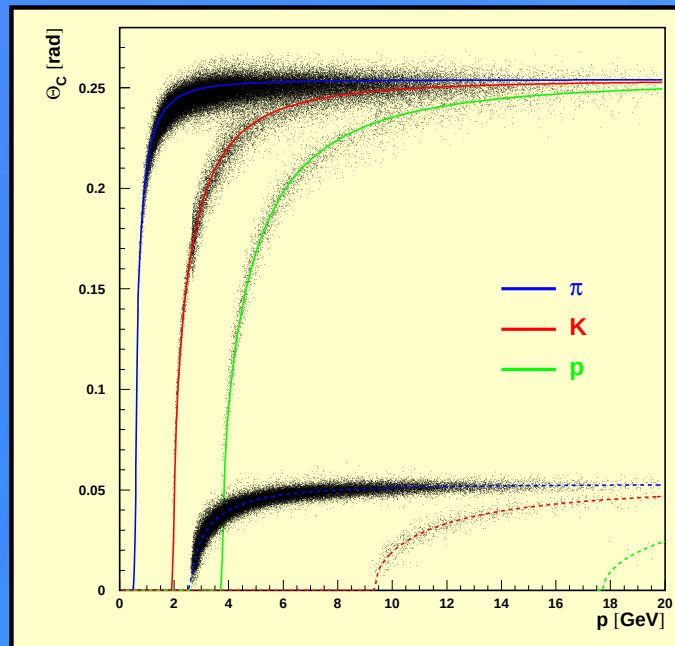
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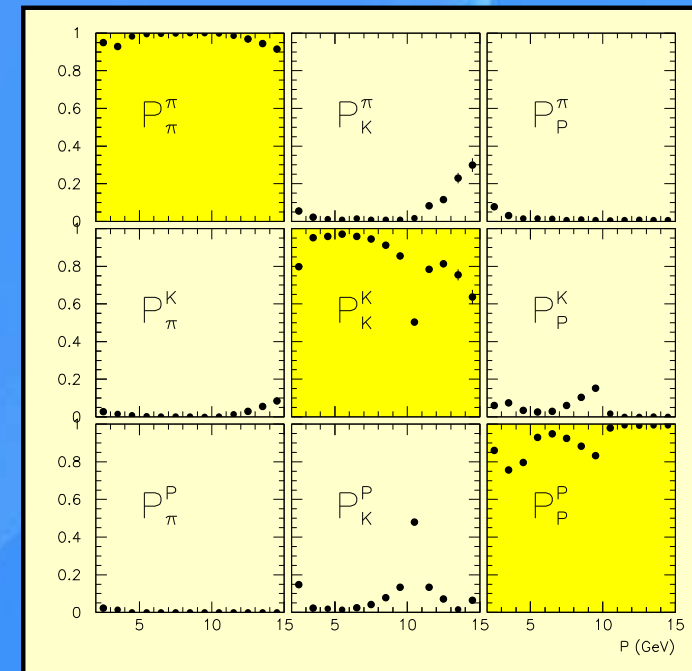


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detection efficiencies
misidentifications

Lab	Experiment	Year	Beam			Tar get			Measure
			Energy	Type	P_B	Type	P_T	f	
SLAC	E80	75	10-16 GeV	e^-	0.85	H-butanol	0.50	0.13	A_1^p
	E130	80	16-23 GeV	e^-	0.81	H-butanol	0.58	0.15	A_1^p
	E142	92	19-26 GeV	e^-	0.39	^3He	0.35	0.35	g_1^n
	E143	93	10-29 GeV	e^-	0.85	NH ₃	0.70	0.15	g_1^p
	"	"	"	"	"	ND ₃	0.25	0.24	g_1^d
	E154	95	48 GeV	e^-	0.83	^3He	0.38	0.55	g_1^n
	E155	97	48 GeV	e^-	0.81	NH ₃	0.80	0.15	g_1^p
	"	"	"	"	"	LiD	0.22	0.36	g_1^d
	E155X	99	29/32 GeV	e^-	0.81	NH ₃	0.70	0.16	g_2^p
	"	"	"	"	"	LiD	0.22	0.36	g_2^d
CERN	EMC	85	200 GeV	μ^+	0.79	NH ₃	0.78	0.16	g_1^p
	SMC	92	100 GeV	μ^+	0.81	D-butanol	0.40	0.19	$g_1^{d,n}$
	"	93	190 GeV	μ^+	0.80	H-butanol	0.86	0.12	g_1^p, g_2^p
	"	94/95	"	μ^+	0.80	D-butanol	0.50	0.20	g_2^d
	"	96	"	μ^+	0.80	NH ₃	0.89	0.16	g_1^p
DESY	HERMES	95	28 GeV	e^+	0.55	^3He	0.46	1.00	g_1^n
	"	96/97	"	"	"	H	0.88	1.00	g_1^p
	"	98-00	"	$e^-/+$	"	D	0.85	1.00	g_1^d, b_1^d
	"	≥ 01	"	$e^\pm$	"	H	0.85	1.00	
CERN	COMPASS	≥ 01	160 GeV	μ^+	0.80	NH ₃	0.90	0.16	
		"	"	"	0.80	LiD	0.50	0.50	
BNL	RHIC	≥ 01	200 GeV	p	0.70	200 GeV p	0.70	1.0	

Polarized Deep Inelastic Scattering

Cross Section:
$$\frac{d^2\sigma}{d\Omega dE^2} = \frac{\alpha^2 E'}{Q^2 E} \underbrace{L_{\mu\nu}(k, q, s)}_{\text{leptonic}} \underbrace{W^{\mu\nu}(P, q, S)}_{\text{hadronic}}$$

$L_{\mu\nu}$: purely electromagnetic $\Rightarrow$ calculable in QED

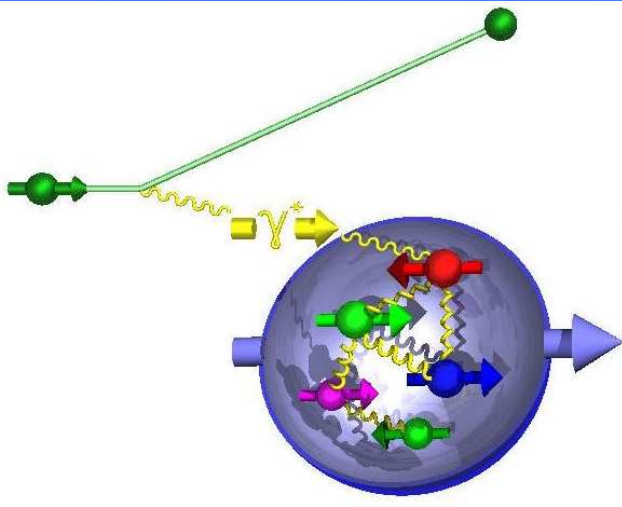
$W^{\mu\nu} = -g^{\mu\nu} F_1(x, Q^2) + \frac{p^\mu p^\nu}{\nu} F_2(x, Q^2) + i\epsilon^{\mu\nu\lambda\sigma} \frac{q_\lambda}{\nu} (S_\sigma g_1(x, Q^2) + \frac{1}{\nu} (p \cdot q S_\sigma - S \cdot q p_\sigma) g_2(x, Q^2))$

(for spin 1 Target) $-b_1(x, Q^2) r_{\mu\nu} + \frac{1}{6} b_2(x, Q^2) (s_{\mu\nu} + t_{\mu\nu} + u_{\mu\nu}) + \frac{1}{2} b_3(x, Q^2) (s_{\mu\nu} - u_{\mu\nu}) + \frac{1}{2} b_4(x, Q^2) (s_{\mu\nu} - t_{\mu\nu})$

F_1, F_2 : **Unpolarized** Structure Functions
 $\Rightarrow$ describe momentum distribution of quarks

g_1, g_2 : **Polarized** Structure Functions
 $\Rightarrow$ describe spin distribution of quarks

Virtual Photon Asymmetry



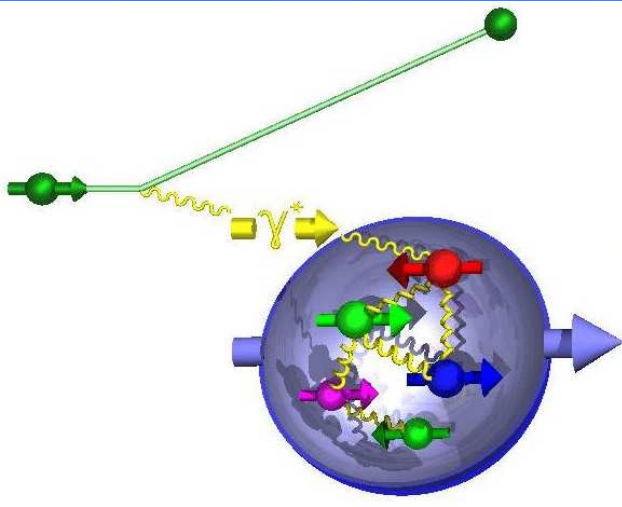
$$\sigma_{3/2} \sim \mathbf{q}^-(\mathbf{x})$$

$$\vec{S}_\gamma + \vec{S}_N = 3/2$$

$$\vec{S}_N = -\vec{S}_q$$

- Virtual photon γ^* can only couple to quarks of opposite helicity

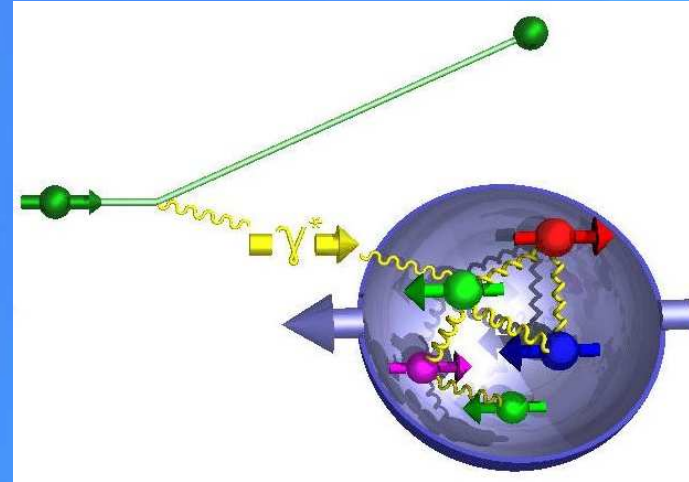
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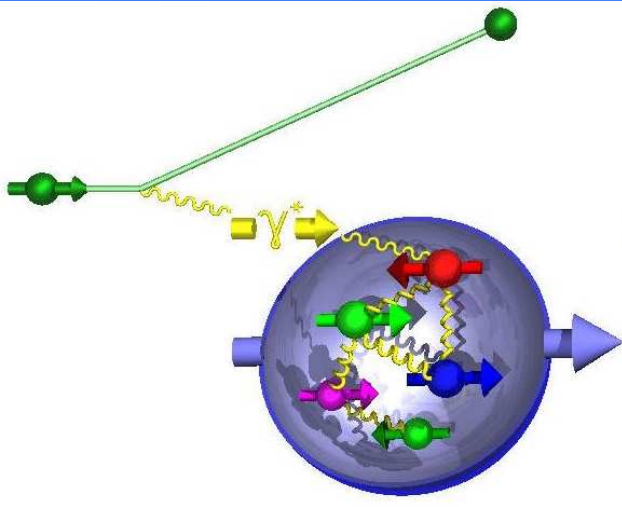
$$\sigma_{1/2} \sim \mathbf{q}^+(\mathbf{x})$$

$$\vec{S}_\gamma + \vec{S}_N = 1/2$$

$$\vec{S}_N = \vec{S}_q$$

- Virtual photon γ^* can only couple to quarks of opposite helicity
- Select $\mathbf{q}^+(\mathbf{x})$ or $\mathbf{q}^-(\mathbf{x})$ by changing the orientation of target nucleon spin or helicity of incident lepton beam

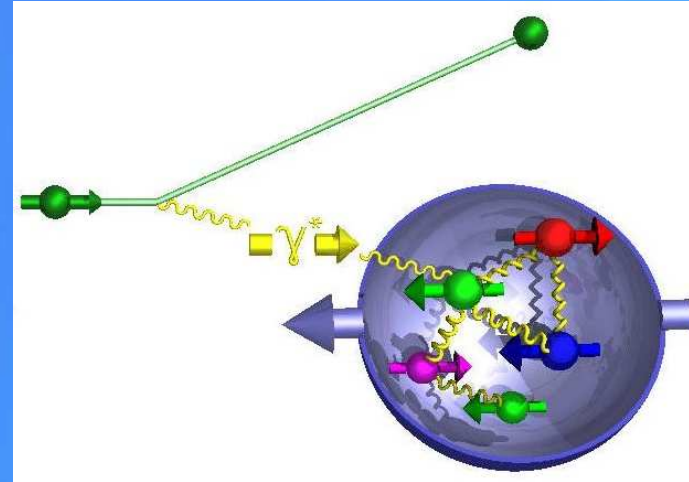
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Quark Helicity Distributions:

$$\Delta q_f(x) := q_f^+(x) - q_f^-(x) \quad (f : u, d, s, \bar{u}, \bar{d}, \bar{s})$$

$$F_1(x) = \frac{1}{2} \sum_i e_i^2 (q_i^+(x) + q_i^-(x)) = \frac{1}{2} \sum_i e_i^2 q_i(x)$$

2xF₁ measures the momentum distribution of quarks

$$g_1(x) = \frac{1}{2} \sum_i e_i^2 (q_i^+(x) - q_i^-(x)) = \frac{1}{2} \sum_i e_i^2 \Delta q_i(x)$$

g₁ measures the spin distribution of quarks

Virtual Photon Asymmetries:

$$A_1 = \frac{\sigma_{\frac{1}{2}} - \sigma_{\frac{3}{2}}}{\sigma_{\frac{1}{2}} + \sigma_{\frac{3}{2}}} = \frac{g_1 - \gamma^2 g_2}{F_1} \quad A_2 = \frac{\sigma_{TL}}{\sigma_T} = \frac{\gamma(g_1 + g_2)}{F_1}$$

Measurable Asymmetries:

$$A_{\parallel} = \frac{\sigma^{\rightarrow\leftarrow} - \sigma^{\rightarrow\rightarrow}}{\sigma^{\rightarrow\leftarrow} + \sigma^{\rightarrow\rightarrow}} \quad A_{\perp} = \frac{\sigma^{\uparrow\rightarrow} - \sigma^{\uparrow\leftarrow}}{\sigma^{\uparrow\rightarrow} + \sigma^{\uparrow\leftarrow}}$$

$$A_{\parallel} = D(A_1 + \eta A_2) \quad A_{\perp} = d(A_2 + \xi A_1)$$

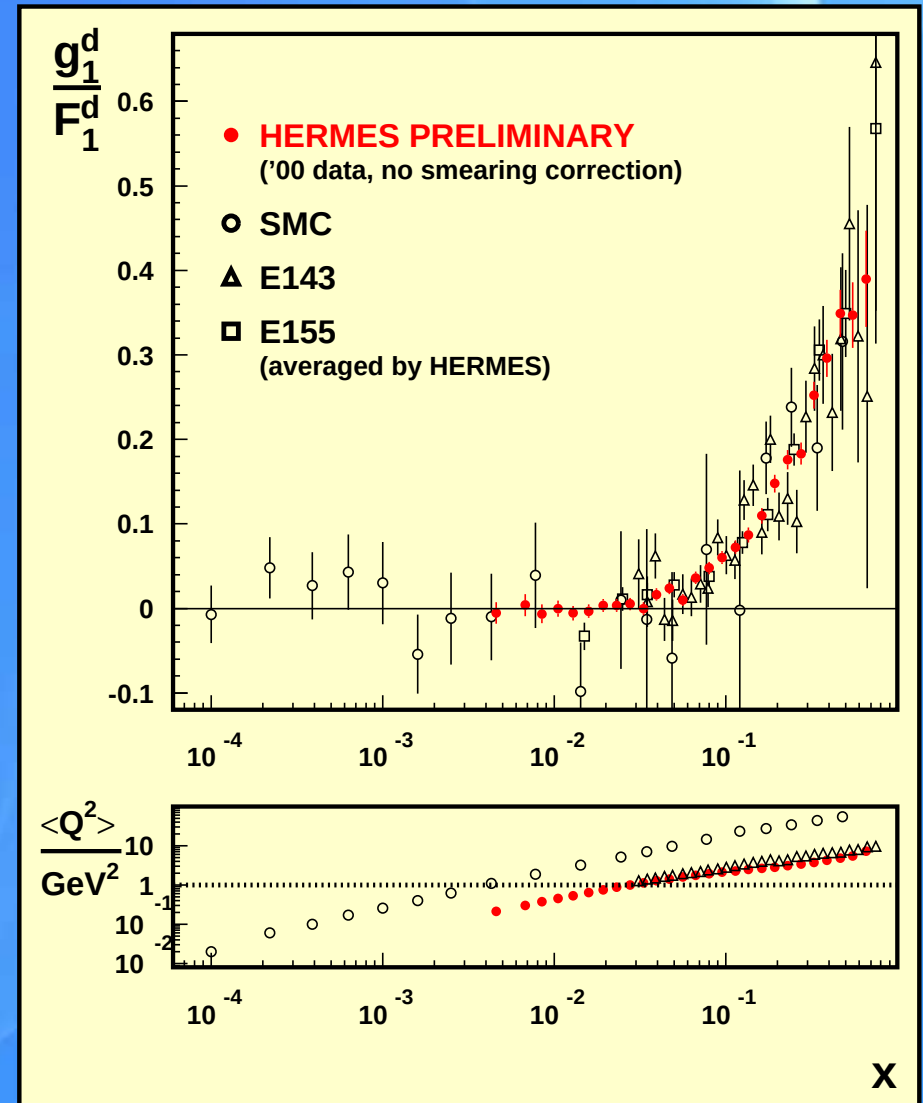
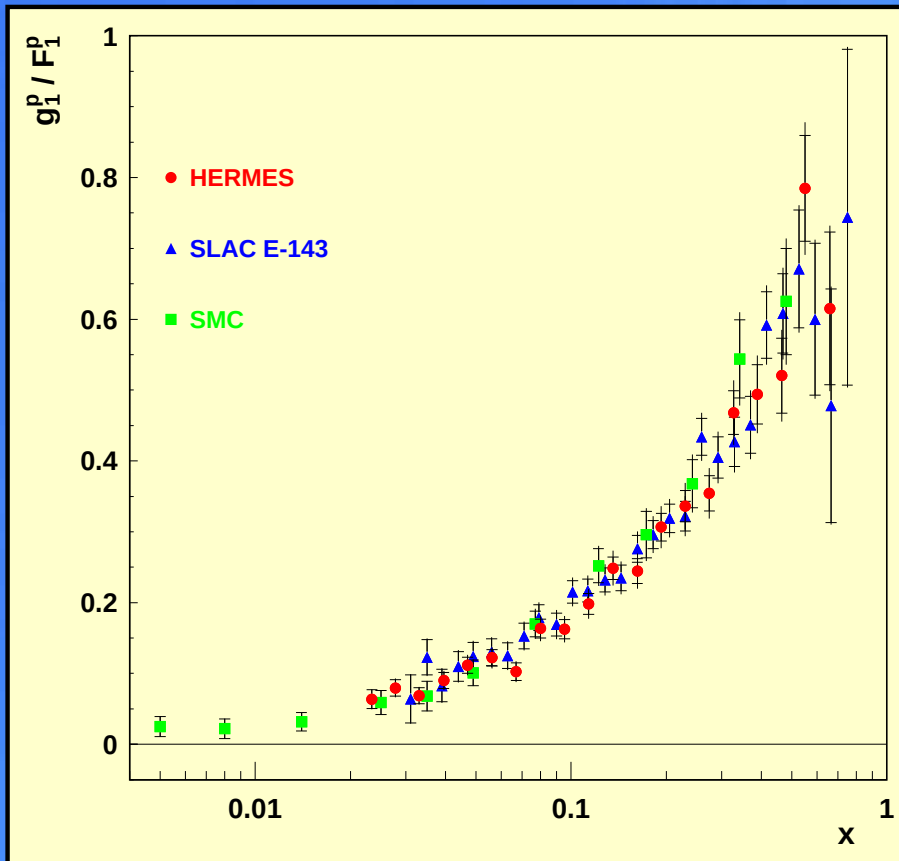
With: $D, d, R, \epsilon, \gamma, \xi, \eta$ being kinematic factors

World data on $g_1(x, Q^2)/F_1(x, Q^2)$

Proton

>

Deuterium



Data given at measured $\langle Q^2 \rangle$: 0.02 - 58 GeV²

$$A_{||} = \frac{1}{P_b P_t} \frac{N^{\rightleftharpoons} L^{\rightleftharpoons} - N^{\rightarrow} L^{\rightarrow}}{N^{\rightleftharpoons} L^{\rightleftharpoons} + N^{\rightarrow} L^{\rightarrow}}$$

$$\frac{g_1}{F_1} = \frac{1}{1 + \gamma^2} \left[\frac{A_{||}}{D} + (\gamma - \eta) A_2 \right]$$

World data on $g_1(x, Q^2)$

$$g_1^p > g_1^D > g_1^n$$

Neglecting sea quarks

$$p : 2 \cdot \frac{4}{9} \Delta u_p + \frac{1}{9} \Delta d_p$$

$$d : p + n$$

$$n : 2 \cdot \frac{1}{9} \Delta d_n + \frac{4}{9} \Delta u_n$$

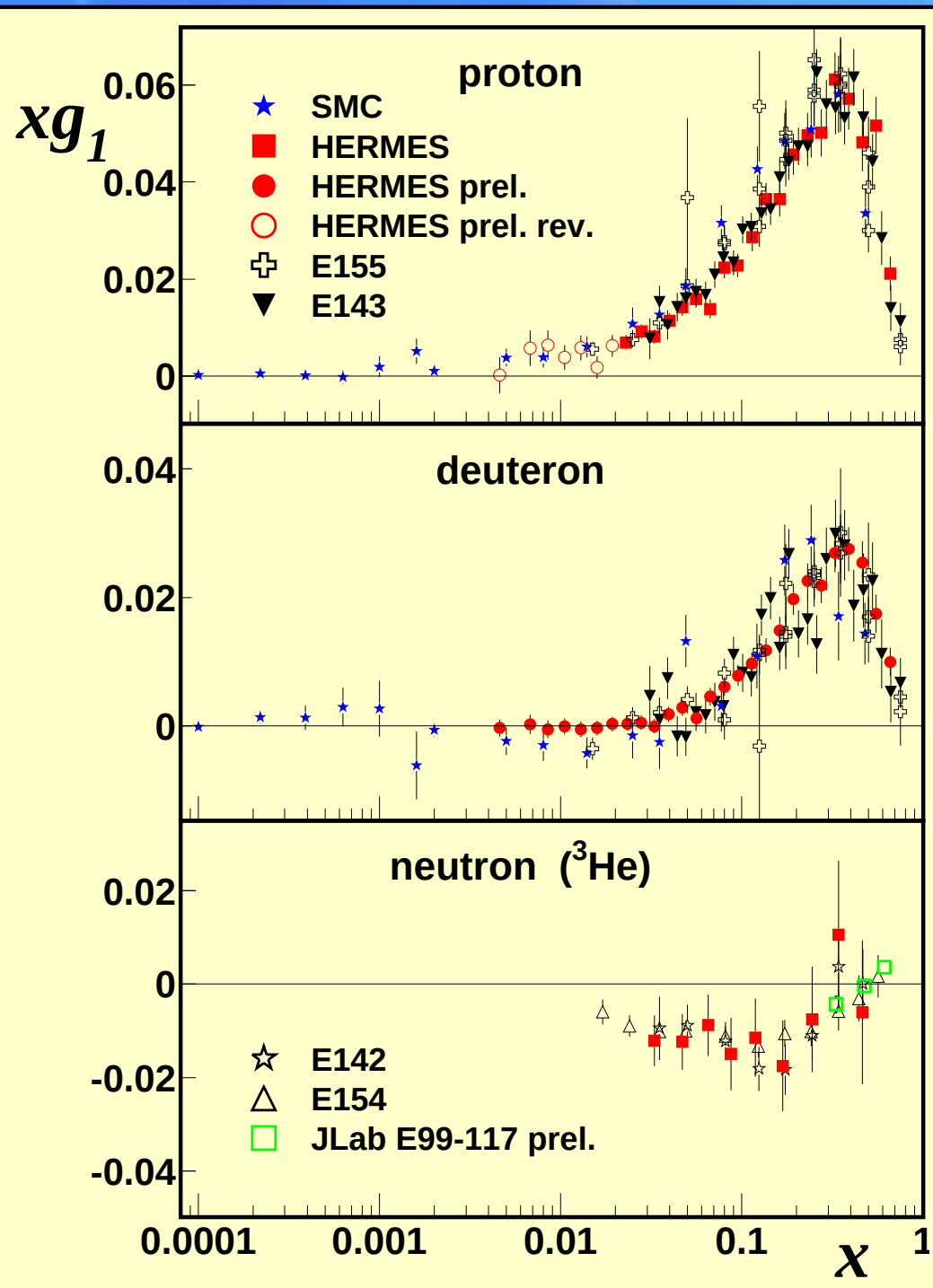
Isospin rotation

$$n : 2 \cdot \frac{1}{9} \Delta u_p + \frac{4}{9} \Delta d_p$$

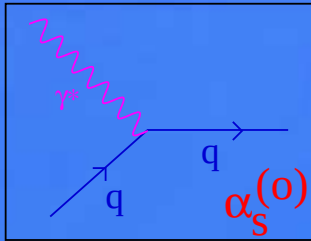
$$\Delta u_p > 0$$

$$\Delta d_p < 0$$

What does a sophisticated model say?

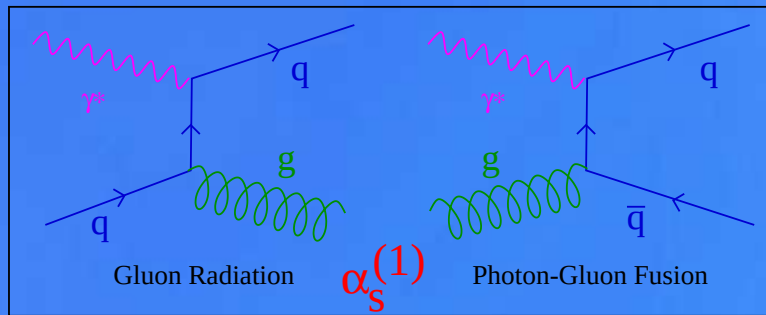


Beyond the Naive Parton Model



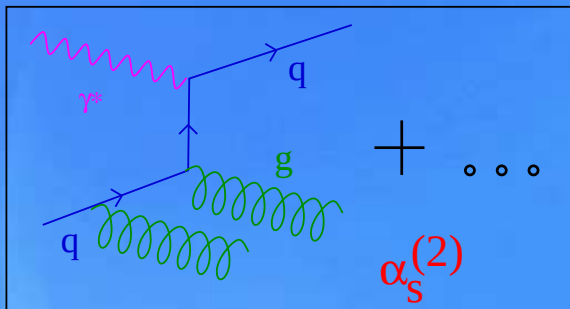
⇒ No gluons

$$g_1^0(x) = \frac{1}{2} \sum e_q^2 \Delta q(x)$$



⇒ Quarks are re-defined with the inclusion of ΔG (weak dependence)

$$g_1^{LO}(x, Q^2) = \frac{1}{2} \sum e_q^2 \Delta q(x, Q^2)$$



⇒ g_1 becomes explicitly ΔG dependent

$$g_1^{NLO}(x, Q^2) = g_1^{LO}(x, Q^2) + \frac{\alpha_s}{2\pi} \frac{1}{2} \sum e_q^2 [\Delta q(x, Q^2) \otimes C_q + \Delta G(x, Q^2) \otimes C_G]$$

- In NLO there are two independent NS distributions and $\Delta\Sigma$ and ΔG

$$\begin{aligned}\Delta q_{NS}^p &= \frac{1}{2} (2\Delta u - \Delta d - \Delta s) & \Delta q_{NS}^n &= \frac{1}{2} (2\Delta d - \Delta u - \Delta s) \\ \Delta\Sigma &= \Delta u + \Delta d + \Delta s\end{aligned}$$

- In NLO there are two independent NS distributions and $\Delta\Sigma$ and ΔG

$$\Delta q_{NS}^p = \frac{1}{2} (2\Delta u - \Delta d - \Delta s) \quad \Delta q_{NS}^n = \frac{1}{2} (2\Delta d - \Delta u - \Delta s)$$

$$\Delta\Sigma = \Delta u + \Delta d + \Delta s$$

- Their Q^2 dependence is regulated by the evolution equations

$$\frac{d}{d \ln Q^2} \Delta q_{NS} = \frac{\alpha_s}{2\pi} P_{qq}^{NS} \otimes \Delta q_{NS}$$

$$\frac{d}{d \ln Q^2} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix} = \frac{\alpha_s}{2\pi} \begin{pmatrix} P_{qq}^S & 2n_f P_{qG} \\ P_{Gq}^S & P_{GG} \end{pmatrix} \otimes \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix}$$

- In NLO there are two independent NS distributions and $\Delta\Sigma$ and ΔG

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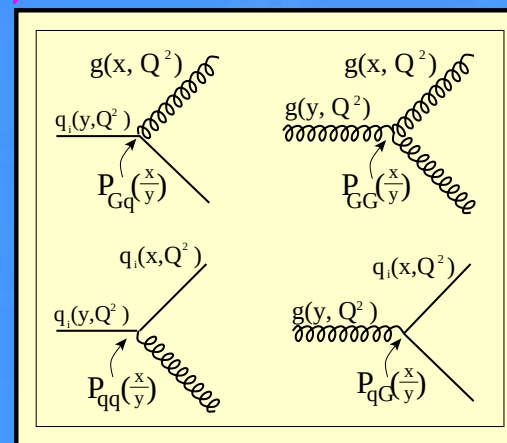
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$$\frac{d}{d \ln Q^2} \Delta q_{NS} = \frac{\alpha_s}{2\pi} P_{qq}^{NS} \otimes \Delta q_{NS}$$

$$\frac{d}{d \ln Q^2} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix} = \frac{\alpha_s}{2\pi} \begin{pmatrix} P_{qq}^S & 2n_f P_{qG} \\ P_{Gq}^S & P_{GG} \end{pmatrix} \otimes \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix}$$

Splitting Functions



- In NLO there are two independent NS distributions and $\Delta\Sigma$ and ΔG

$$\Delta q_{NS}^p = \frac{1}{2} (2\Delta u - \Delta d - \Delta s) \quad \Delta q_{NS}^n = \frac{1}{2} (2\Delta d - \Delta u - \Delta s)$$

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- Each distribution is parameterized at a starting Q_0^2 :
medium x

$$x \Delta q_i(x, Q_0^2) = \eta_i A_i \underbrace{x^{a_i}}_{\text{low } x} \overbrace{(1 + \gamma_i x + \rho_i x^{\frac{1}{2}})}^{\text{medium } x} \underbrace{(1 - x)^{b_i}}_{\text{high } x}$$

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- They are evolved to Q_{meas}^2 using the evolution equations to obtain g_1^{calc}
- The χ^2 is minimized $\chi^2 = \sum_{\text{data}} (g_1^{\text{meas}} - g_1^{\text{calc}})^2 / \sigma_{\text{stat}}^2$

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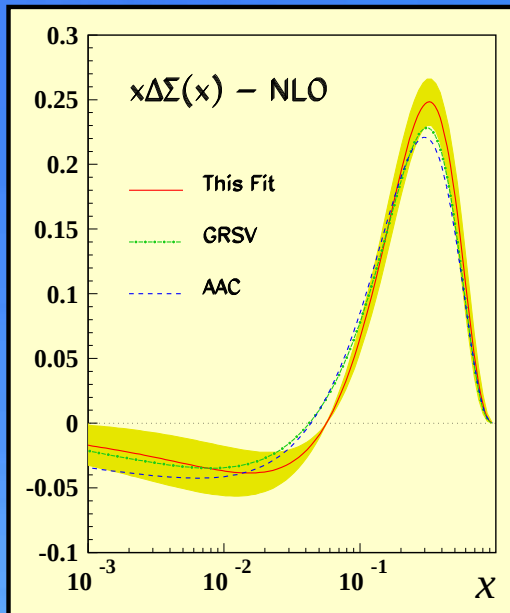
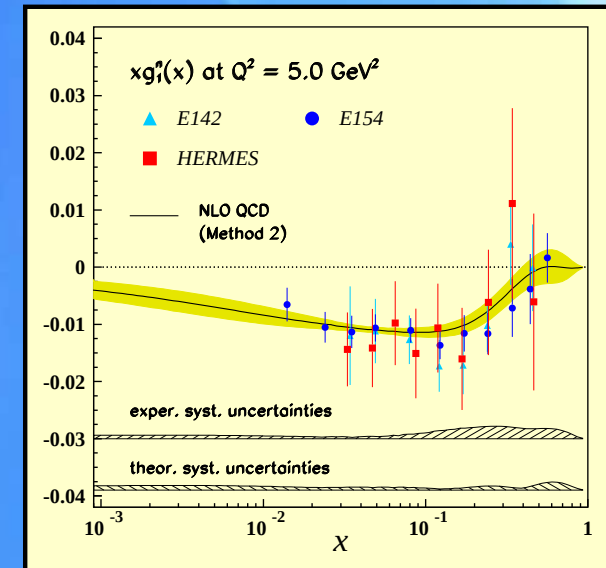
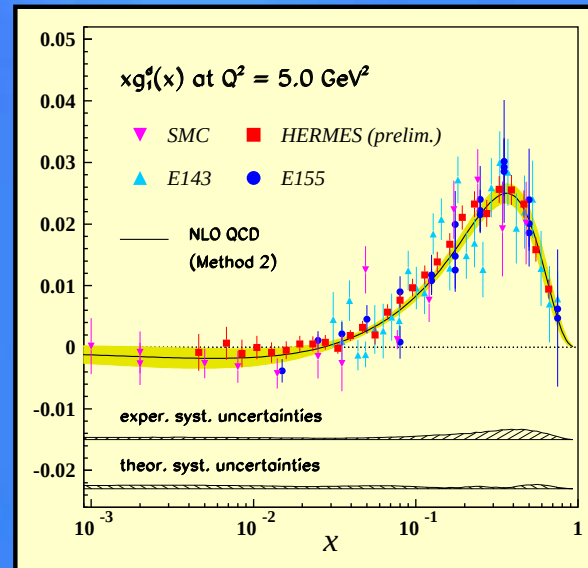
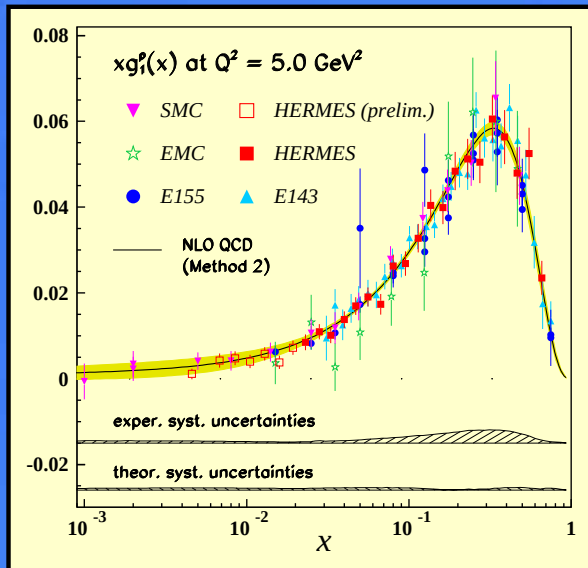
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- The χ^2 is minimized $\chi^2 = \sum_{\text{data}} (g_1^{\text{meas}} - g_1^{\text{calc}})^2 / \sigma_{\text{stat}}^2$
- The parameters $a, b, \gamma, \dots$ are evaluated

NLO-QCD Fit Results for $g_1(x, Q^2)$



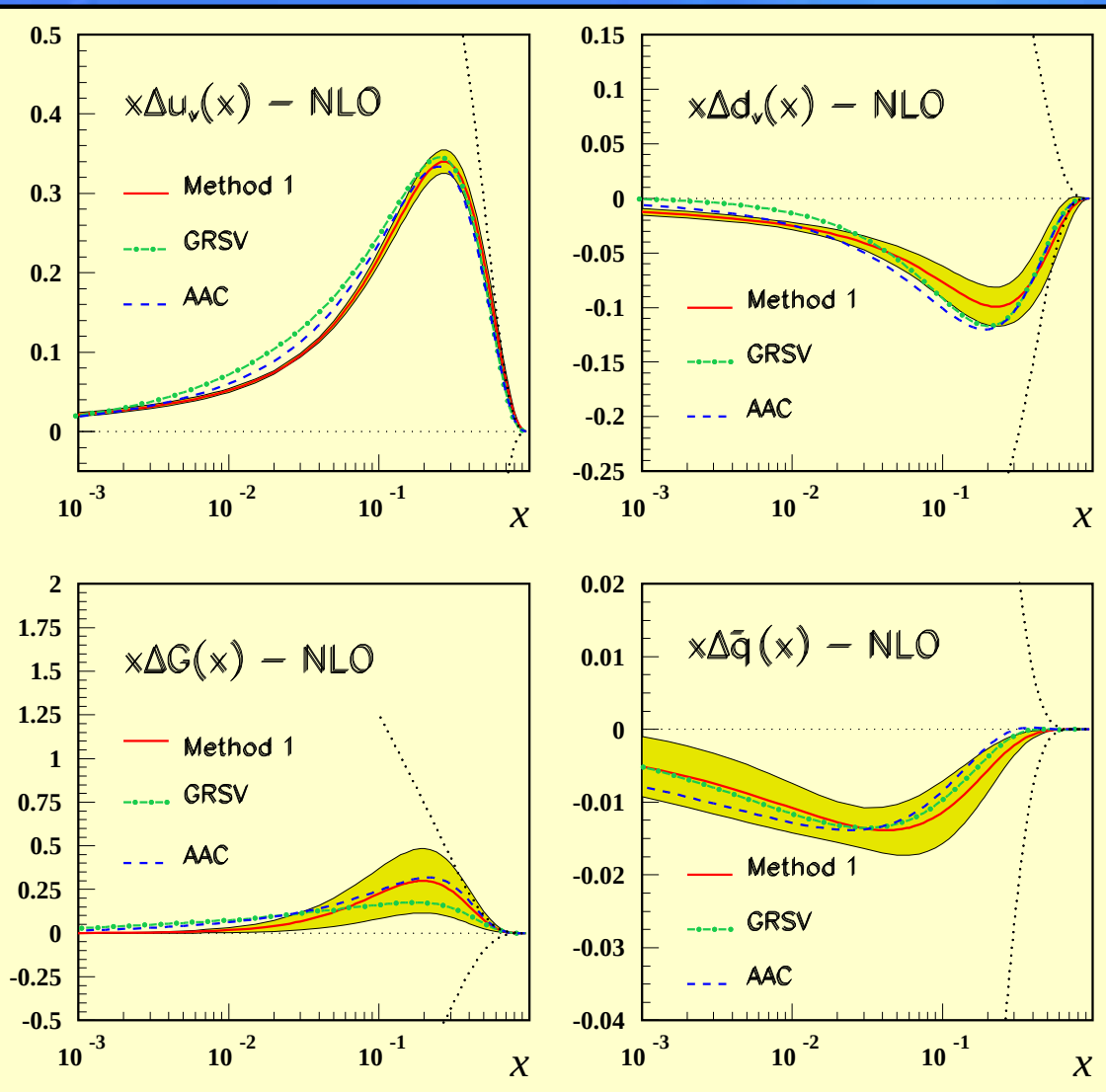
Good Description of $g_1^{p,d,n}(x, Q^2)$

BUT:

$$\Delta\Sigma = 0.201 \pm 0.103 \neq 1$$

Who carries the spin ?

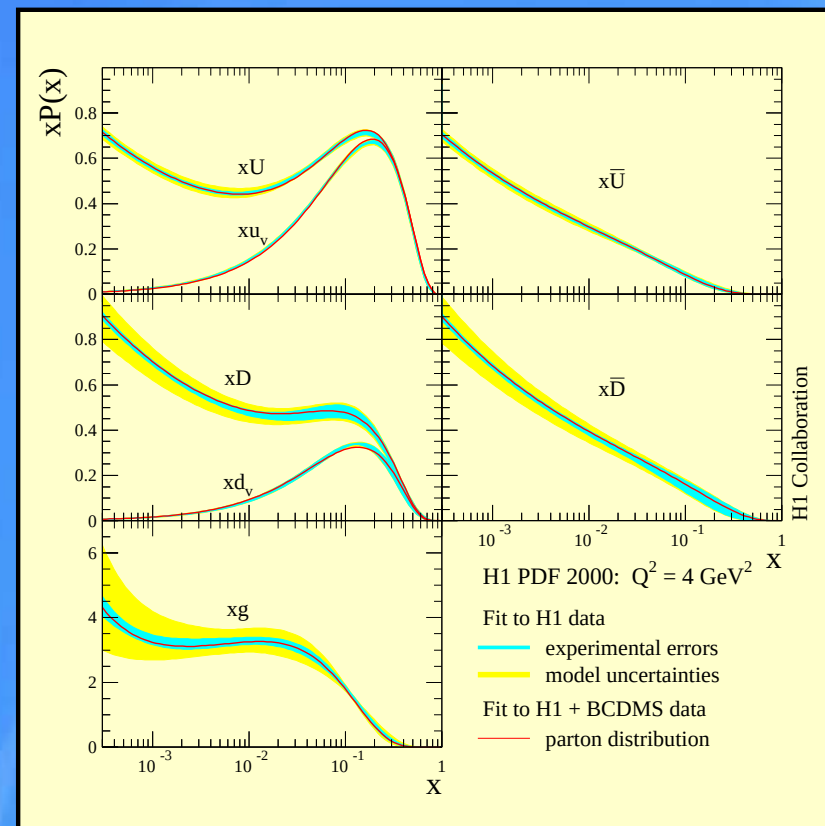
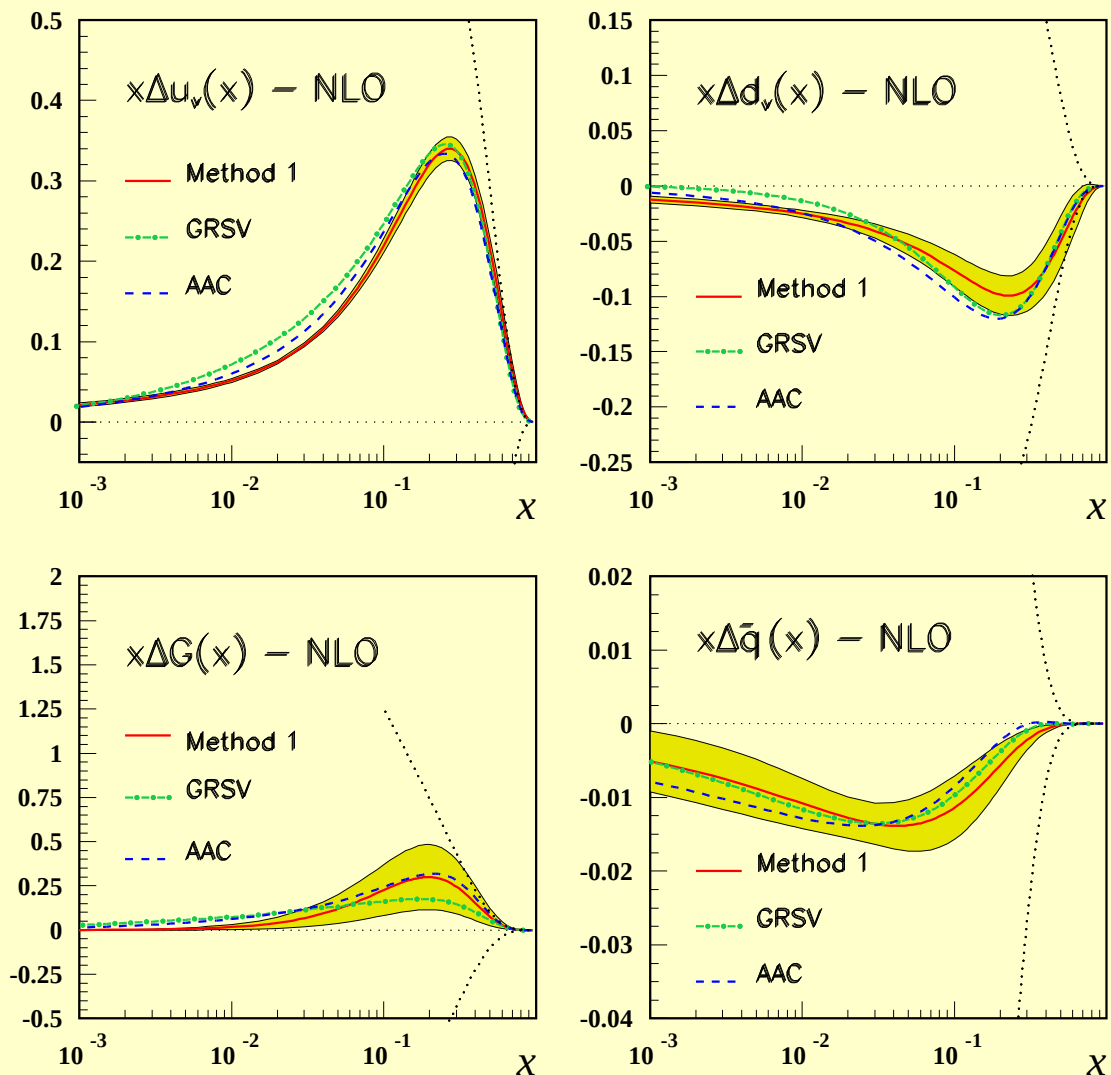
NLO-QCD Fit Results for Different Flavors



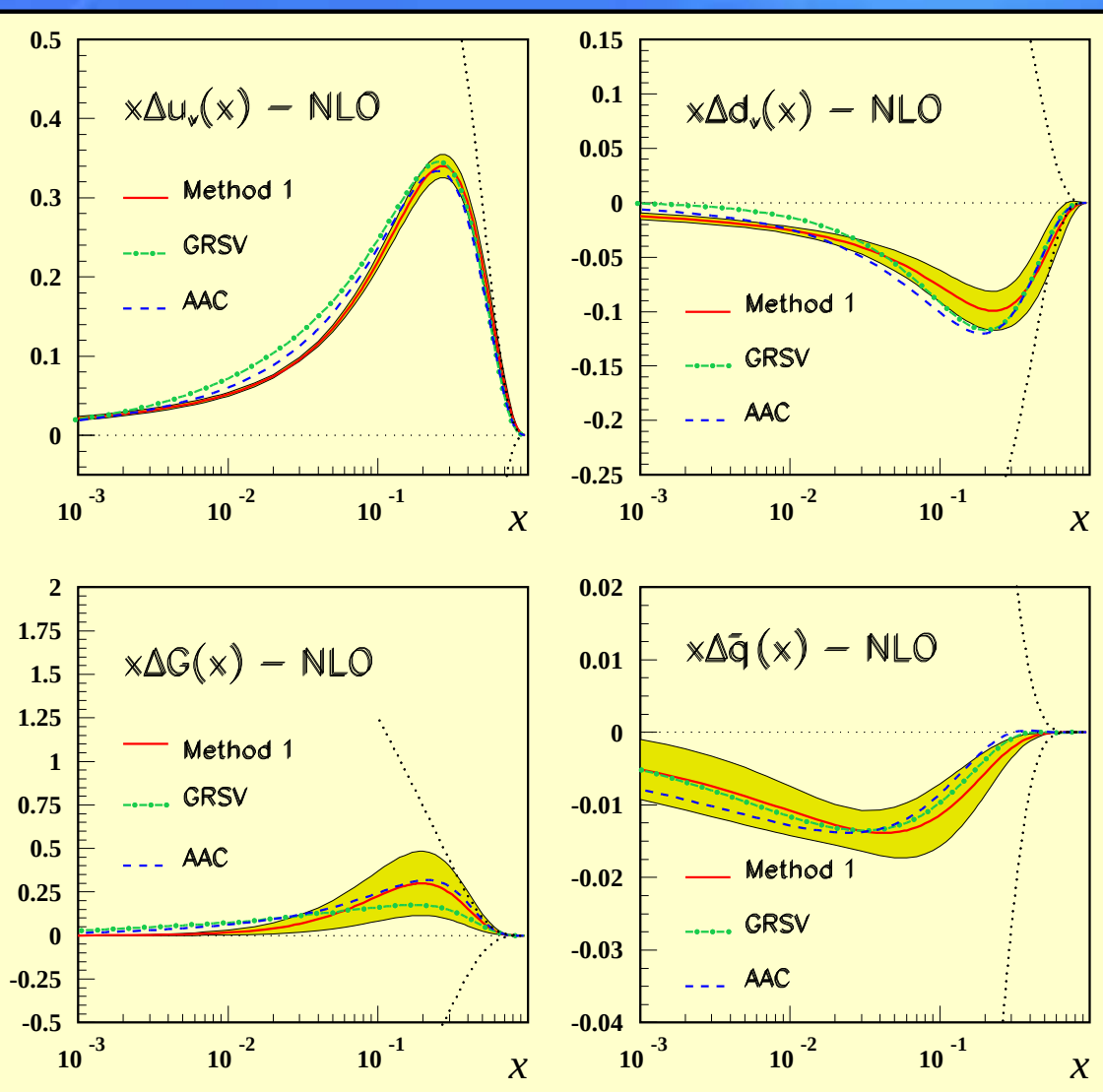
NLO-QCD Fit Results for Different Flavors

Δu_v and Δd_v well determined

Comparable accuracy as in unpolarised case



NLO-QCD Fit Results for Different Flavors



Δu_v and Δd_v well determined

$$\Delta s(\bar{q}) = -0.070 \pm 0.028 \text{ (stat.)} < 0$$

BUT

$$\Delta u_s = \Delta \bar{u} = \Delta d_s = \Delta \bar{d} = \Delta s = \Delta \bar{s}$$

is assumed

unpolarized PDFs:

SU(2) flavor symmetry breaking

$$\Delta G = 0.616 \pm 0.388 \text{ (stat.)} > 0$$

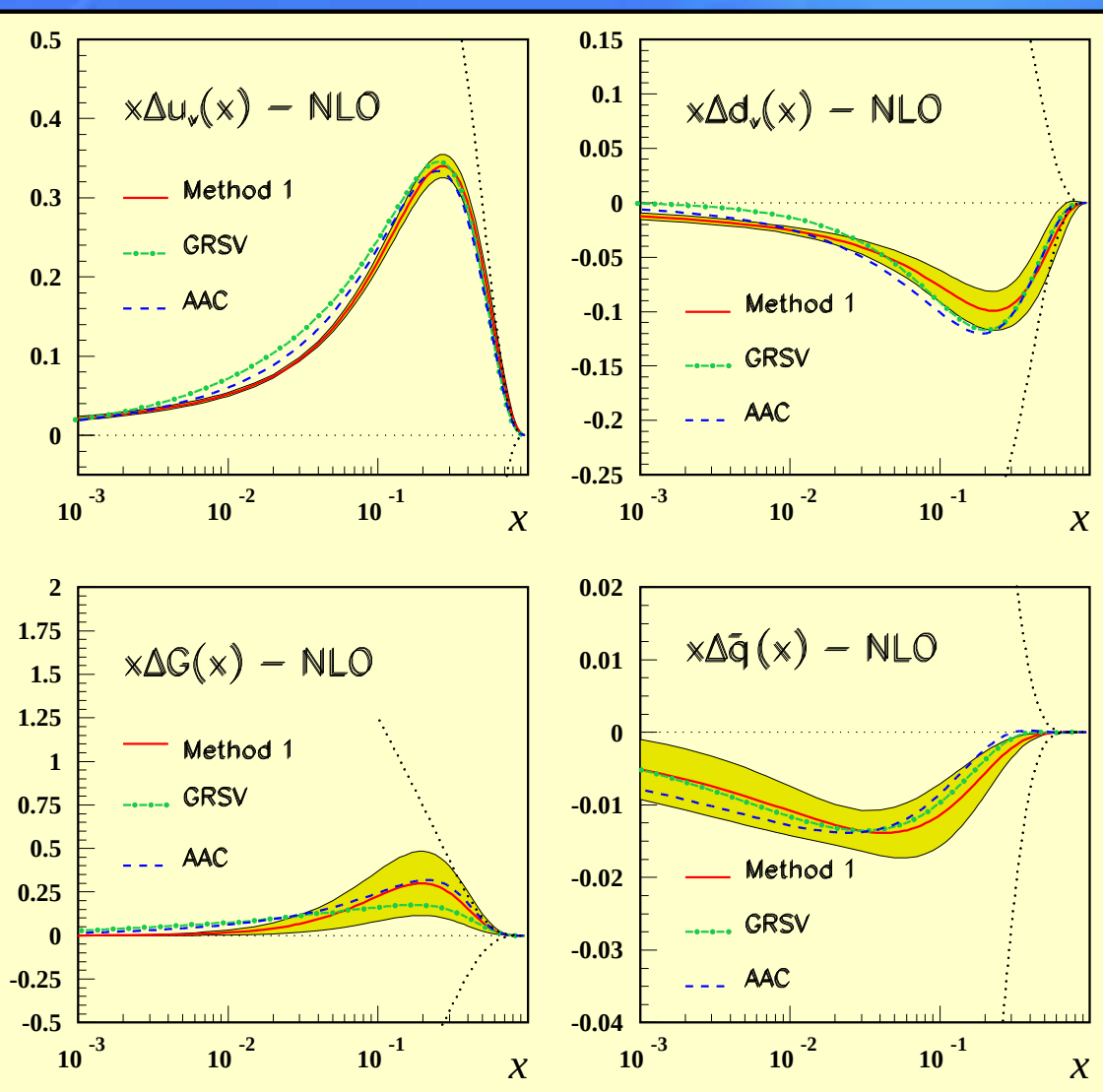
BUT Big error $\Rightarrow$ small Q^2 lever arm

GRSV: Glück et al. hep/ph 0011215

AAC: Goto et al. hep/ph 0001046

at $Q^2 = 4 \text{ GeV}^2$

NLO-QCD Fit Results for Different Flavors



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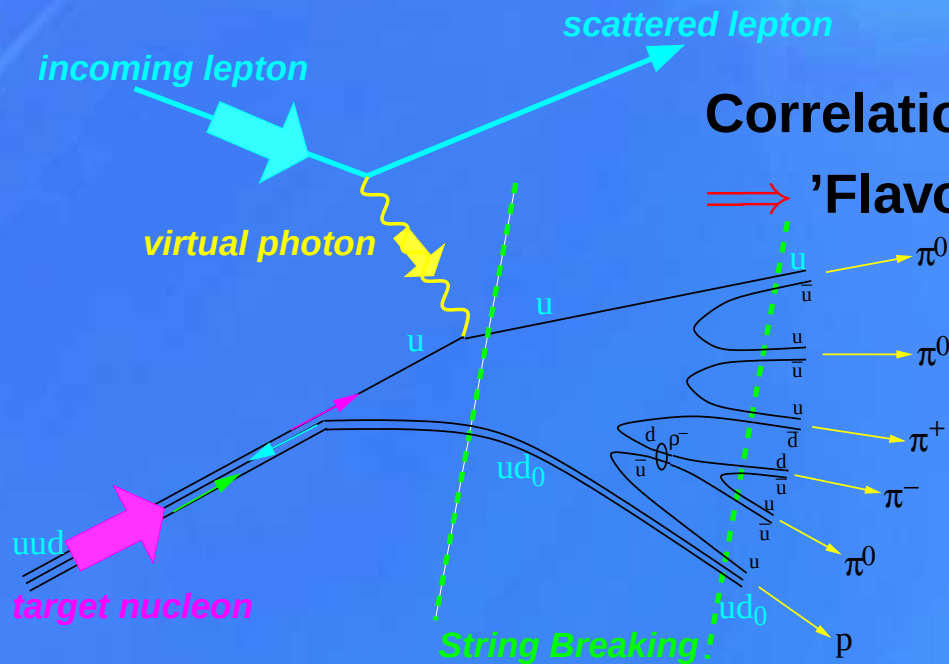
AAC: Goto et al. hep/ph 0001046

at $Q^2 = 4 \text{ GeV}^2$

more direct access to ΔG and Δq_s desirable



Semi-inclusive DIS



Correlation between detected hadron and struck q_f
 $\Rightarrow$ 'Flavor - Separation'

Inclusive DIS:

$$\Delta\Sigma = (\Delta u + \Delta\bar{u} + \Delta d + \Delta\bar{d} + \Delta s + \Delta\bar{s})$$

Semi-inclusive DIS:

$$\Delta u, \Delta\bar{u}, \Delta d, \Delta\bar{d}, \Delta s, \Delta\bar{s}$$

In LO-QCD:

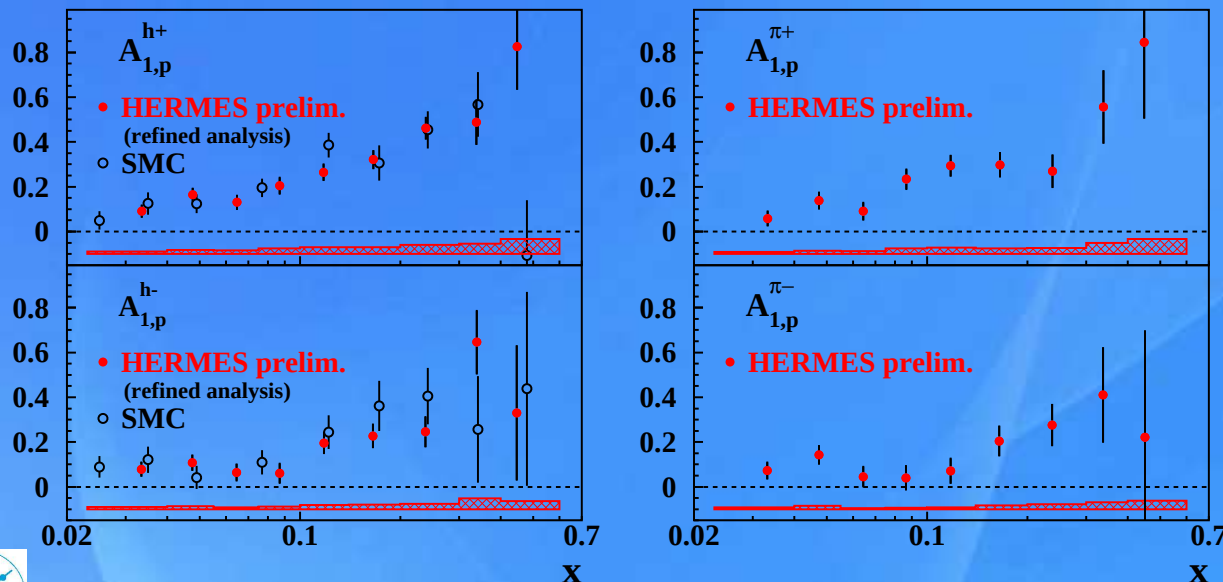
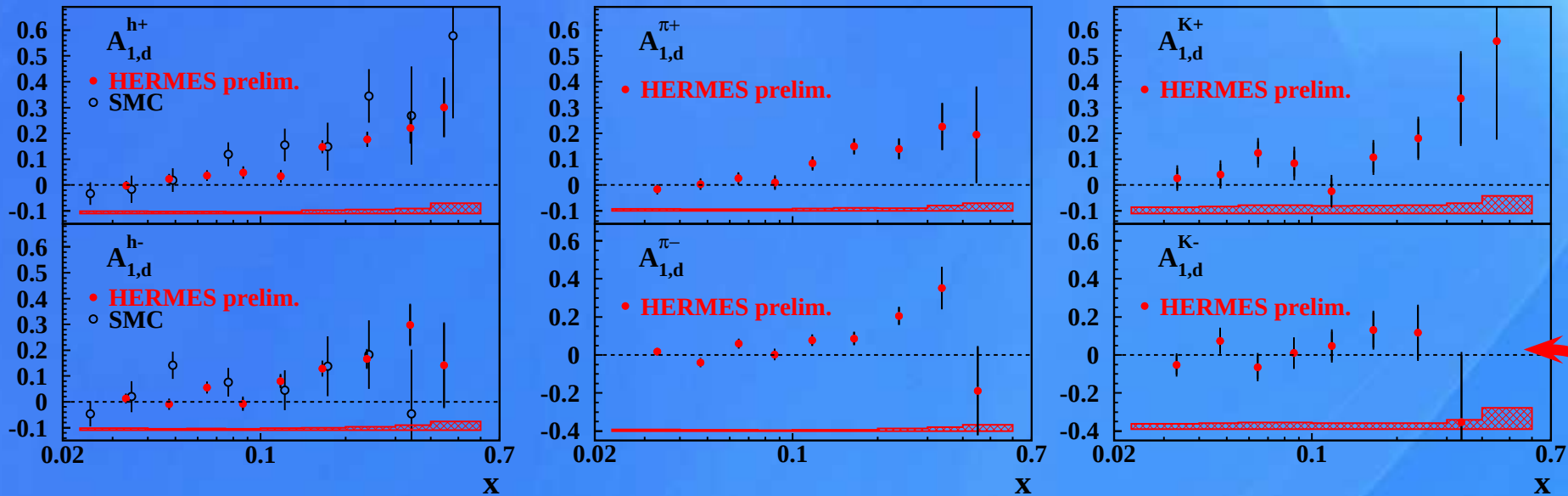
$$A_1^h(x, Q^2) = \frac{\sigma_{1/2}^h - \sigma_{3/2}^h}{\sigma_{1/2}^h + \sigma_{3/2}^h} = \frac{1 + R(x, Q^2)}{1 + \gamma^2} \cdot \frac{\sum_f e_f^2 \Delta q_f(x, Q^2) \int dz D_f^h(z, Q^2)}{\sum_f e_f^2 q_f(x, Q^2) \int dz D_f^h(z, Q^2)}$$

$(\Delta q_f), q_f$
 $D_f^h(z)$

(Polarized) quark distributions

fragmentation functions giving the probability that a (struck) **quark** of flavor f fragments into a hadron of type h

Measured Hadron Asymmetries



- $A_1^{K^-}(x) \approx 0$!!
- $K^- = (\bar{u}s)$ is an all-sea object
- Covered range:
 $0.2 \leq z \leq 0.8$
 $W^2 > 10 \text{ GeV}^2$

- Rewrite Photon-Nucleon Asymmetry

$$A_1^h(\mathbf{x}) \stackrel{g_2=0}{\simeq} C \cdot \sum_{\mathbf{q}} \underbrace{\frac{e_q^2 q(x) \int dz D_q^h(z)}{\sum_{q'} e_{q'}^2 q'(x) \int dz D_{q'}^h(z)}}_{P_q^h(x, z)} \frac{\Delta \mathbf{q}(\mathbf{x})}{\mathbf{q}(\mathbf{x})}$$

- Purity $P_q^h(x, z)$: probability hadron **h** originates from an event with struck **quark f**; completely **unpolarized** quantity
- Extract $\Delta \mathbf{q}$ by solving:

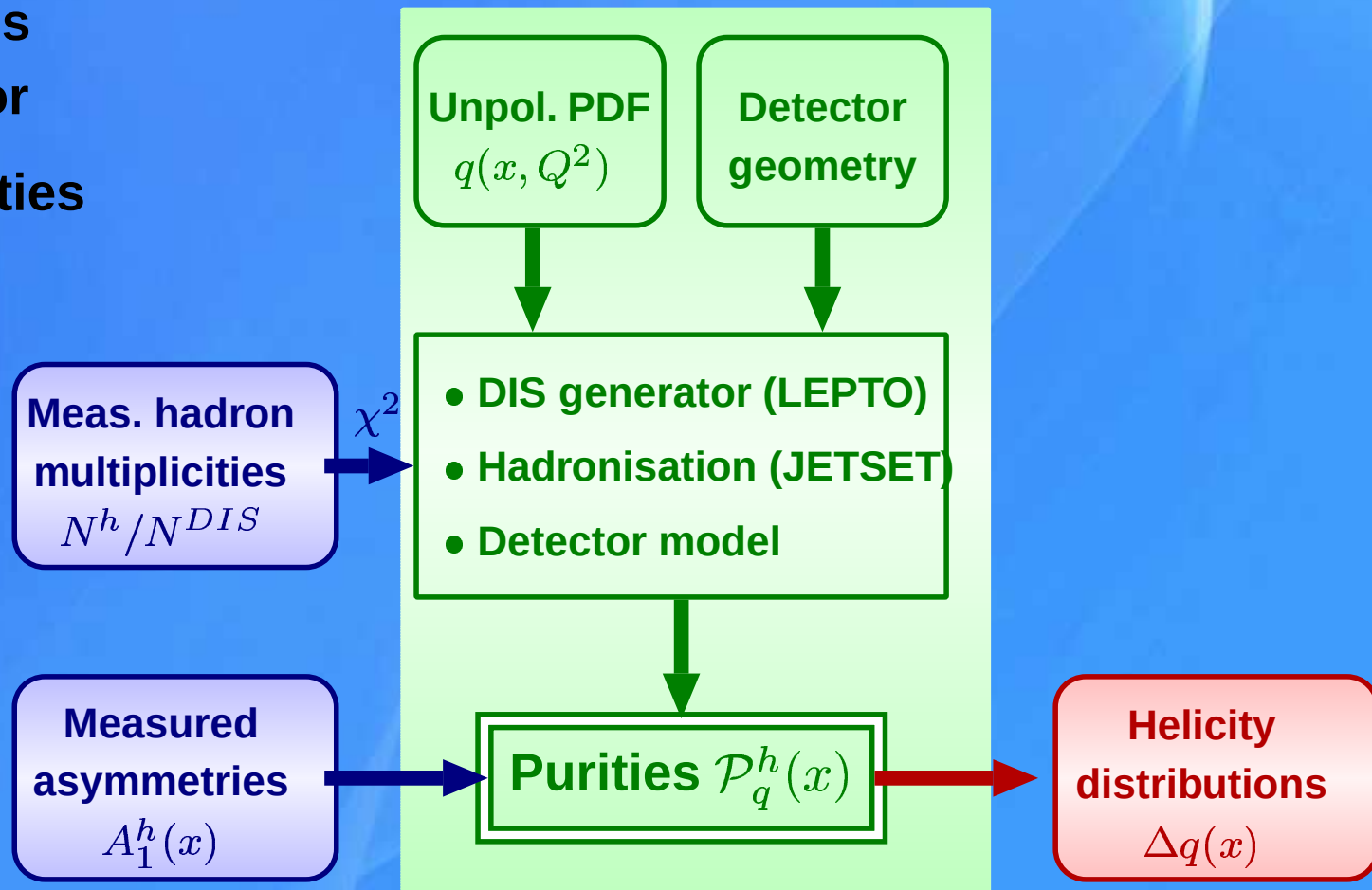
$$\vec{A} = \mathcal{P} \vec{Q}$$

$$\vec{A} = (A_{1,p}(\mathbf{x}), A_{1,d}(\mathbf{x}), A_{1,p}^{\pi^\pm}(\mathbf{x}), A_{1,d}^{\pi^\pm}(\mathbf{x}), A_{1,d}^{K^\pm}(\mathbf{x}))$$

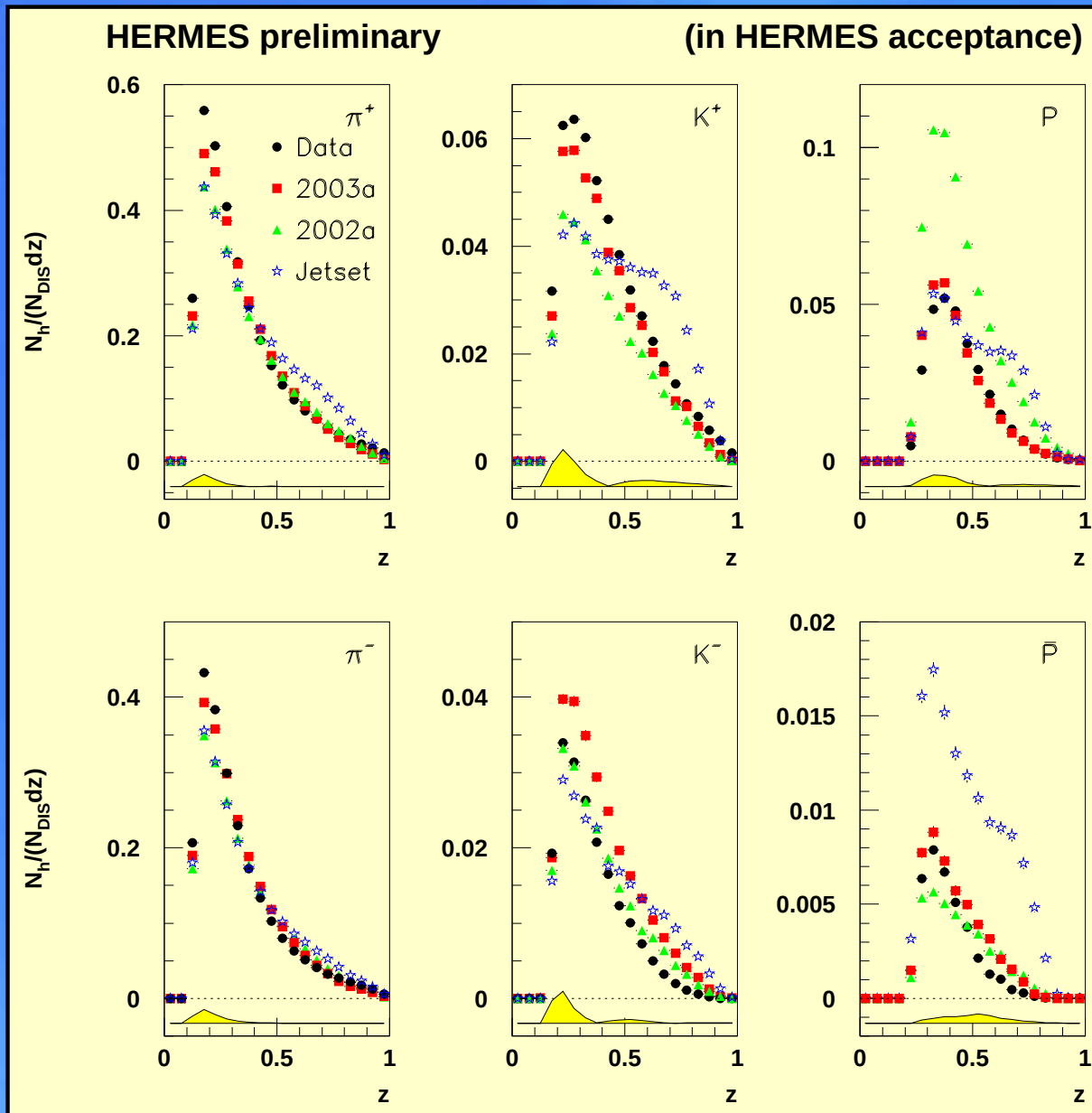
$$\vec{Q} = \left(\frac{\Delta u}{u}, \frac{\Delta d}{d}, \frac{\Delta \bar{u}}{\bar{u}}, \frac{\Delta \bar{d}}{\bar{d}}, \frac{\Delta s + \Delta \bar{s}}{s + \bar{s}} \right)$$

Generation of Purities (HERMES)

- Use Monte Carlo model of DIS process (LEPTO), fragmentation process (JETSET) and detector
- Systematic uncertainties from
 - Variation of fragmentation parameters
 - Use of alternative PDF set GRV98LO vs. CTEQ5L



JETSET-Fragmentation Function



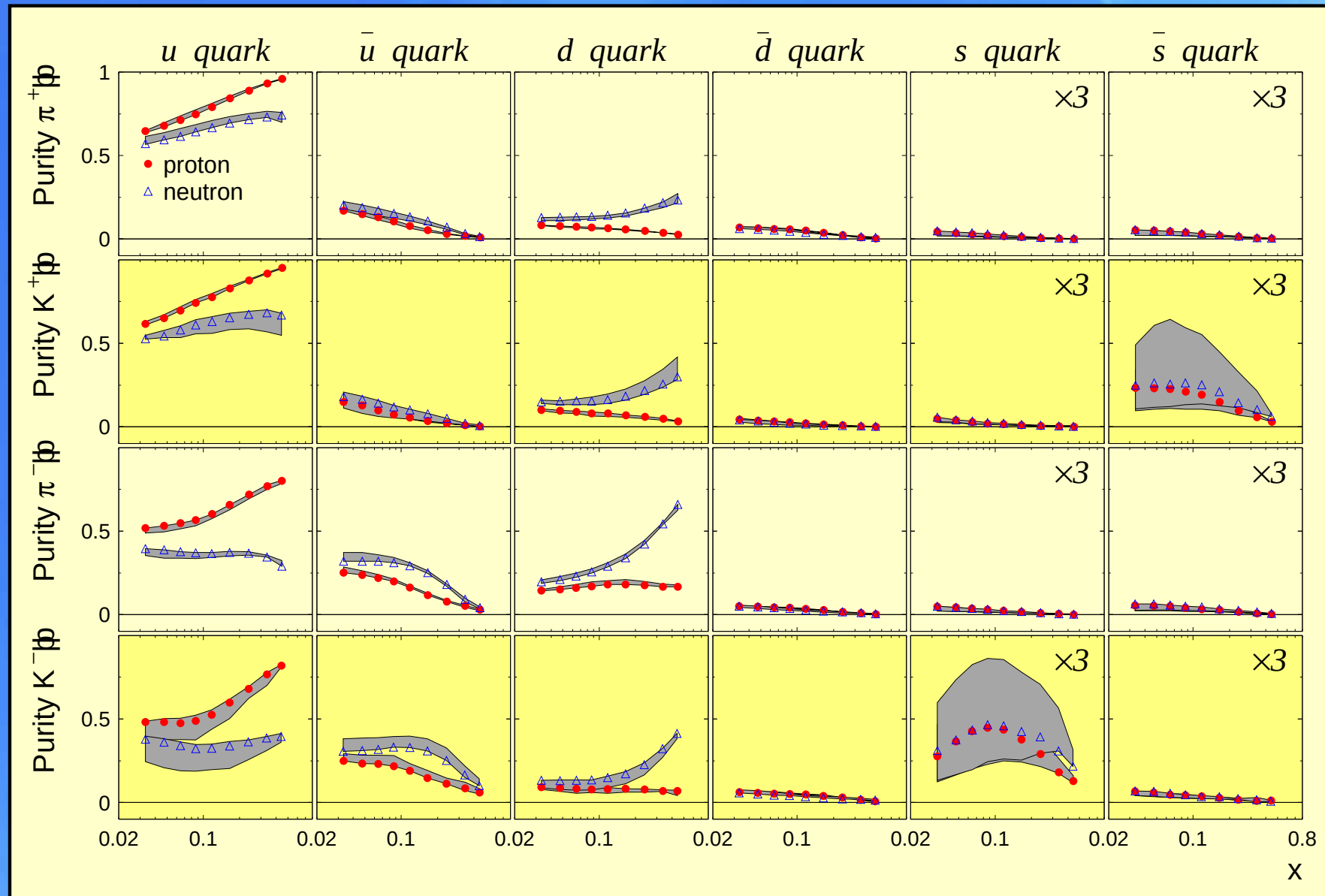
- JETSET default does not describe HERMES data
- Use measured hadron multiplicities $N^h(z)/N_{DIS}$ to tune the JETSET fragmentation model

⇒ FF

$D_u^{\pi^+}, D_d^{\pi^+}, D_u^{\pi^-}, D_d^{\pi^-},$

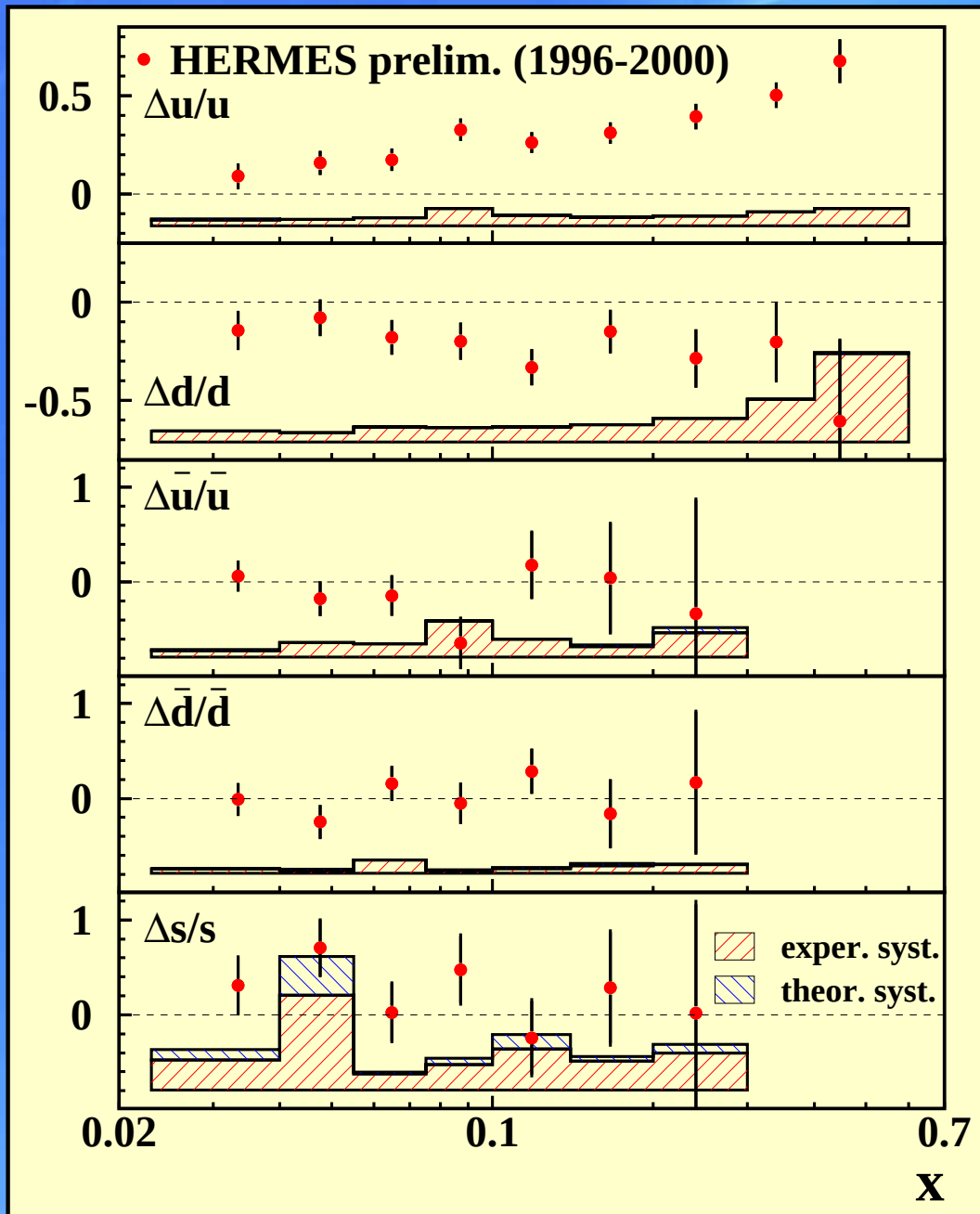
$D_u^{K^+}, D_s^{K^+}, \dots$

Purities (HERMES)



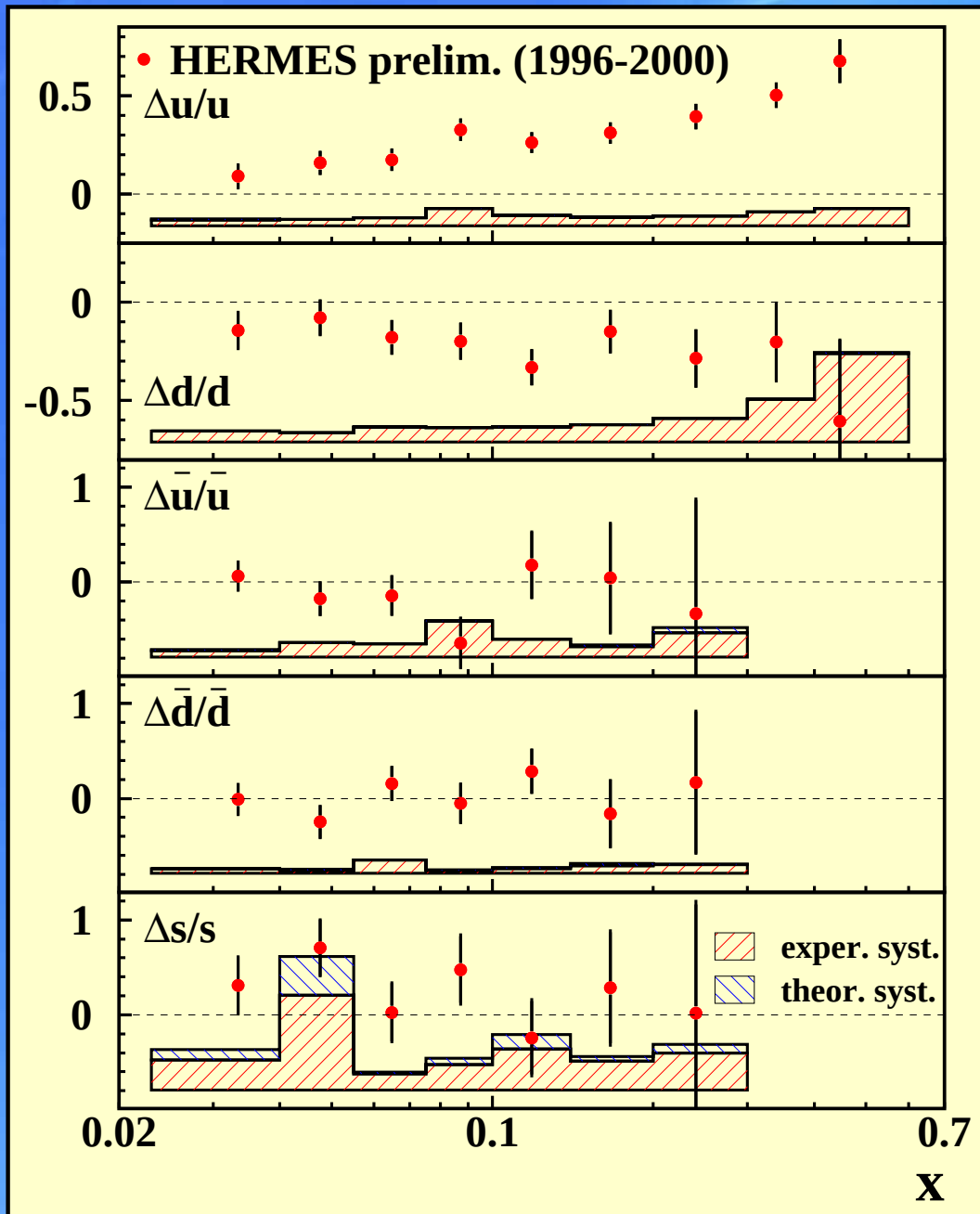
Syst. uncertainties from PDF sets and LUND parameters

Quark Polarizations



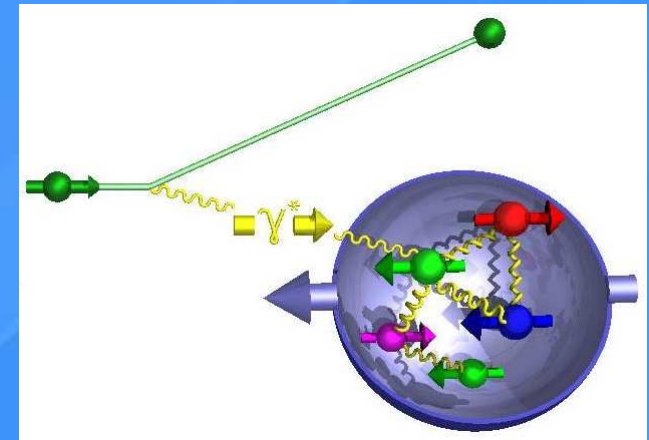
$$\frac{\Delta q_f(x)}{q_f(x)} := \frac{q_f^+(x)}{q_f(x)} - \frac{q_f^-(x)}{q_f(x)}$$

Quark Polarizations

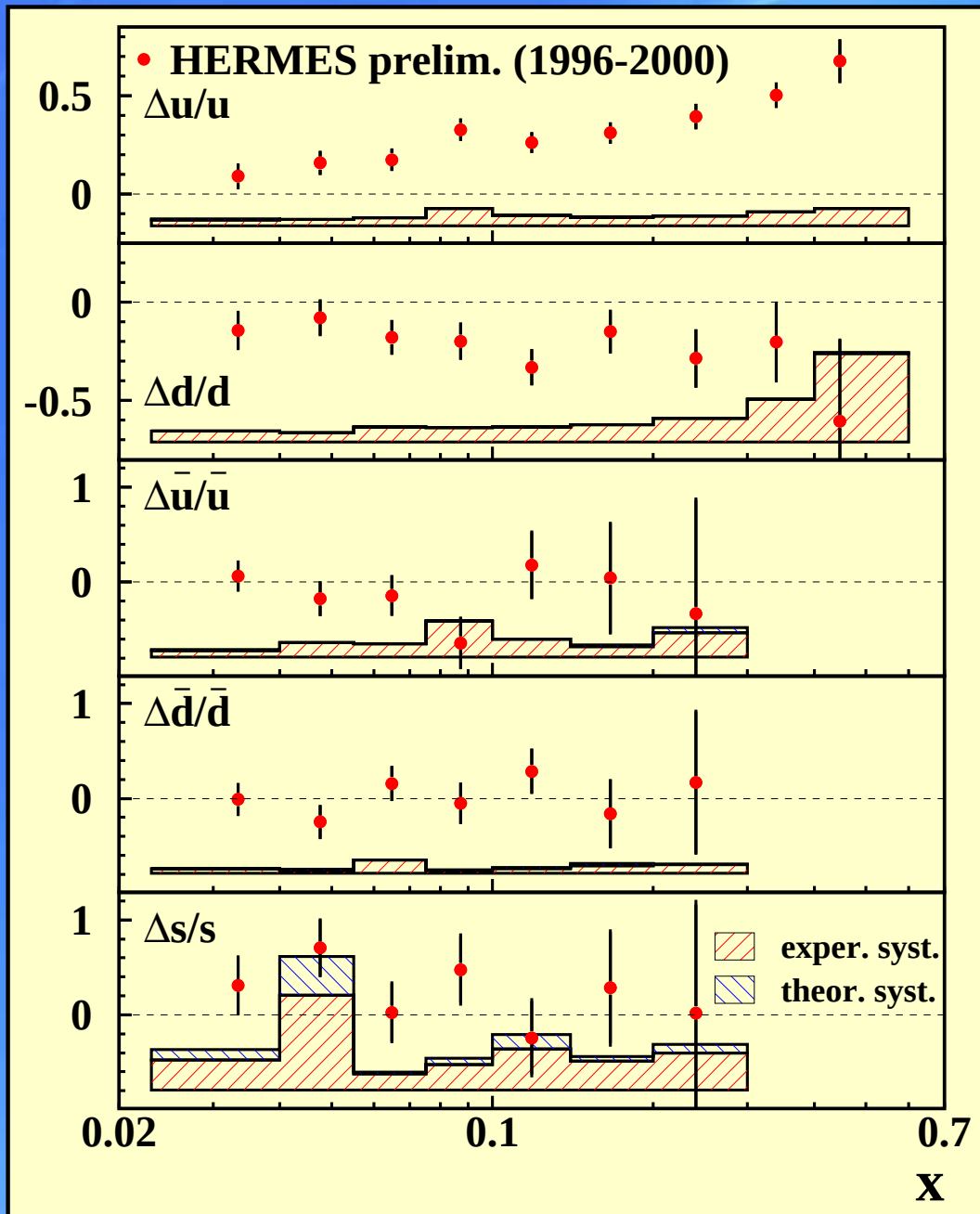


$$\frac{\Delta q_f(x)}{q_f(x)} := \frac{q_f^+(x)}{q_f(x)} - \frac{q_f^-(x)}{q_f(x)}$$

- $\Delta u(x)/u(x) > 0$
 $\Rightarrow$ polarized parallel to the proton spin

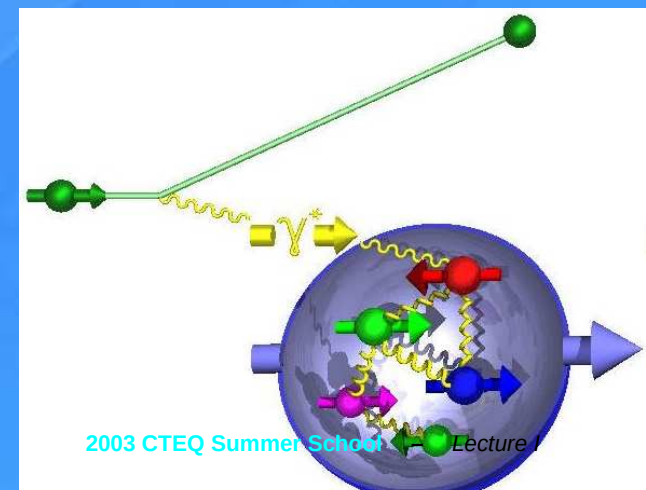


Quark Polarizations

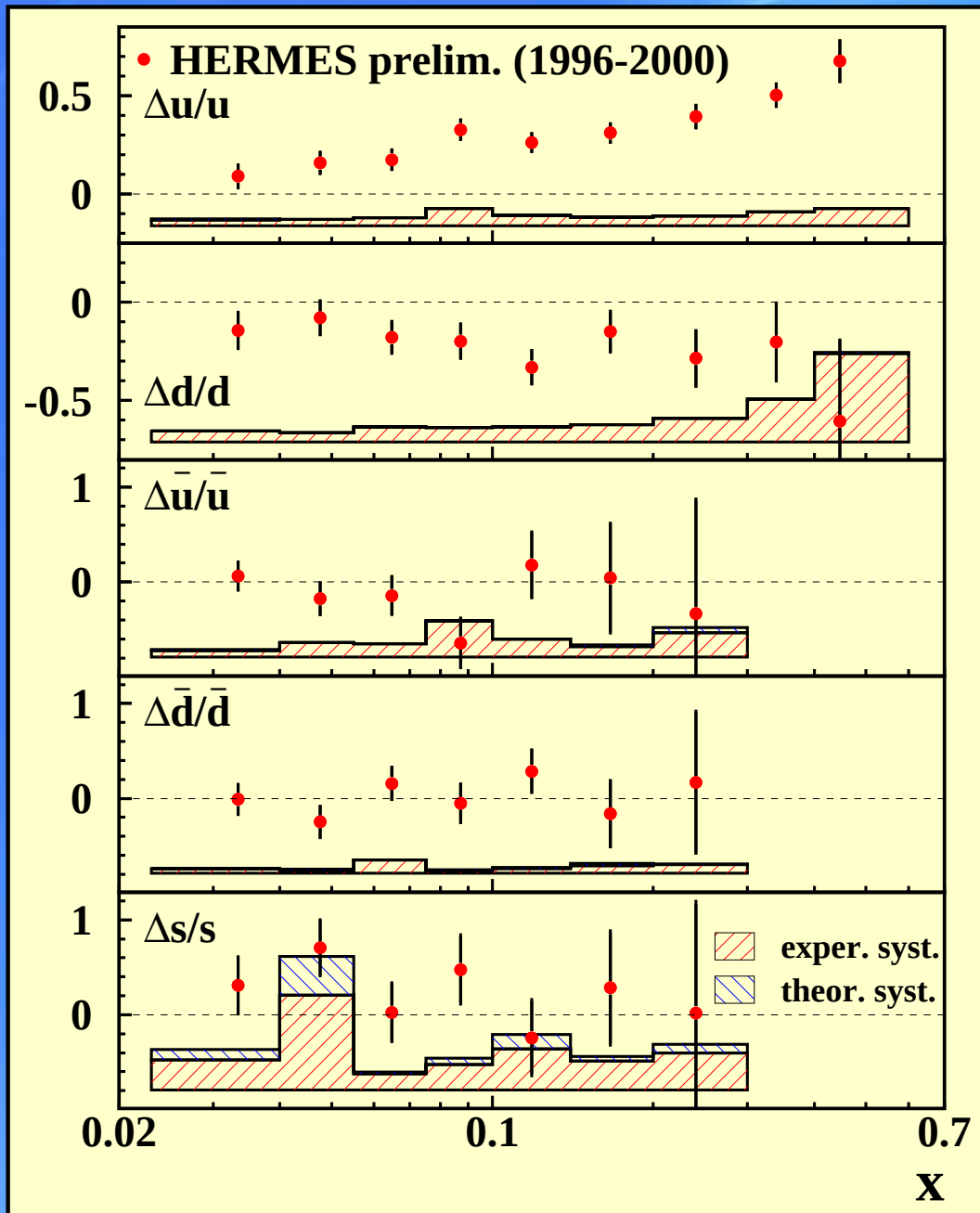


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Quark Polarizations

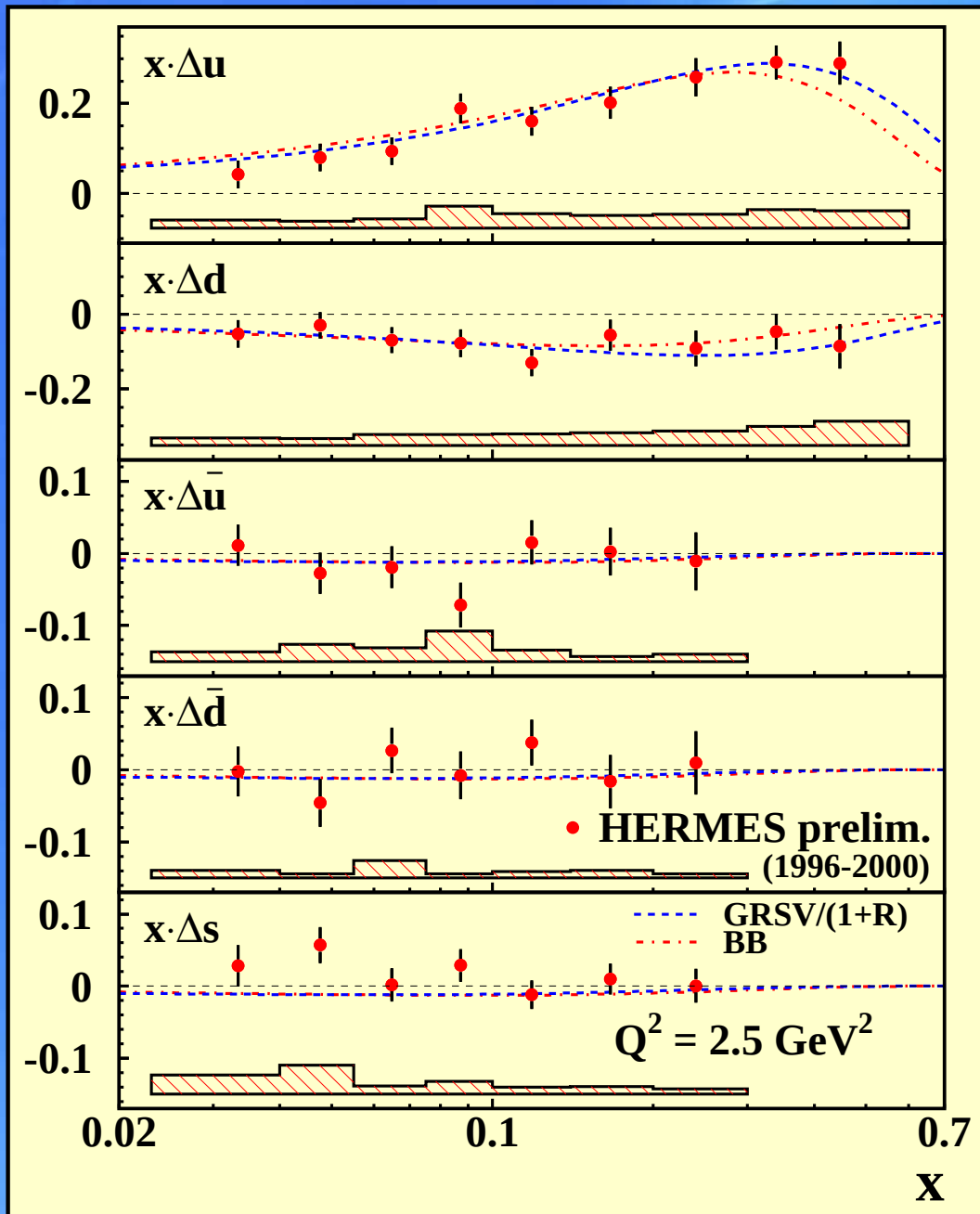


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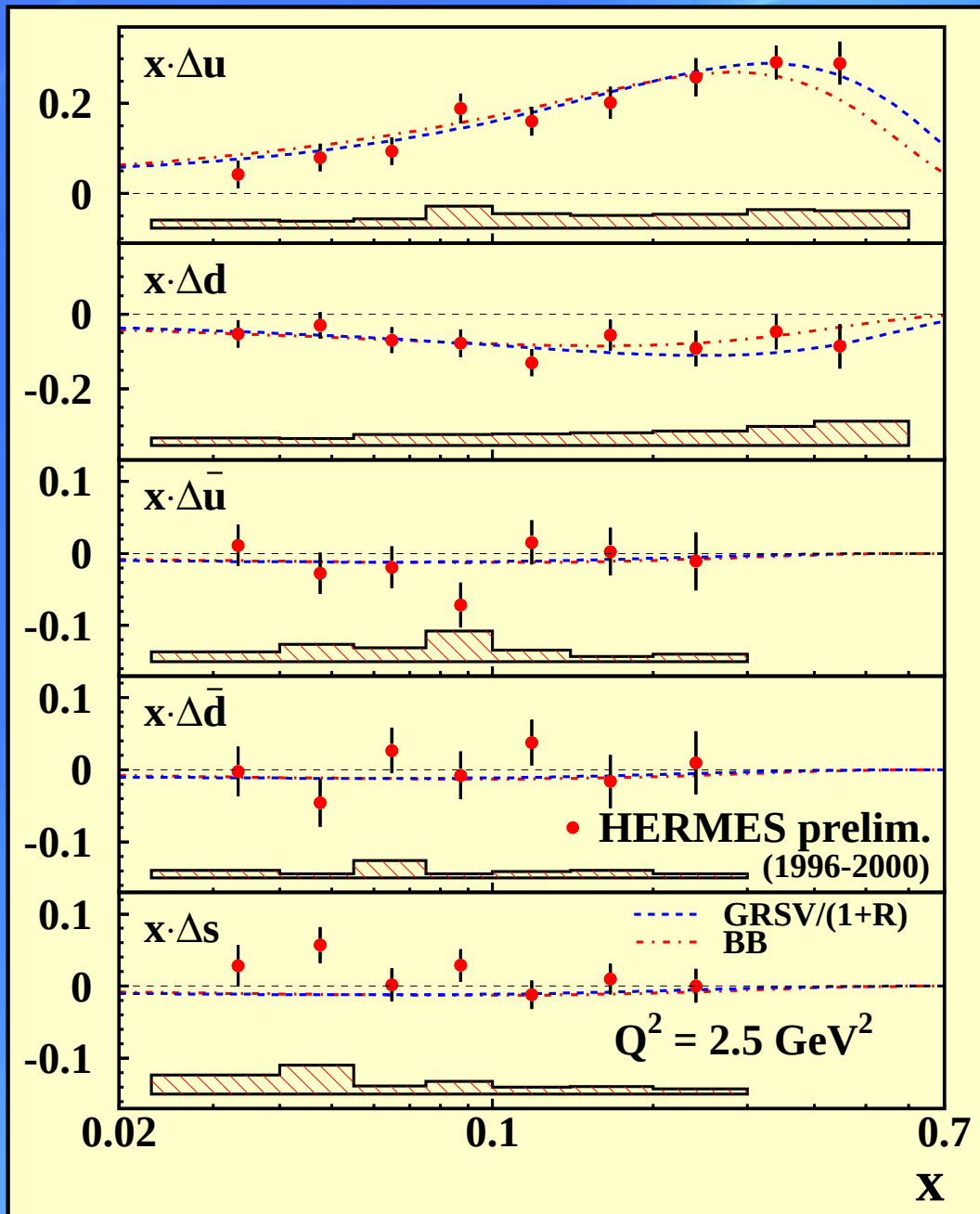
$$\bullet \quad \frac{\Delta \bar{u}}{\bar{u}} \sim \frac{\Delta \bar{d}}{\bar{d}} \sim \frac{\Delta s + \Delta \bar{s}}{s + \bar{s}} \sim 0$$

Polarized Quark Densities



$$\Delta q_f(x) := q_f^+(x) - q_f^-(x)$$

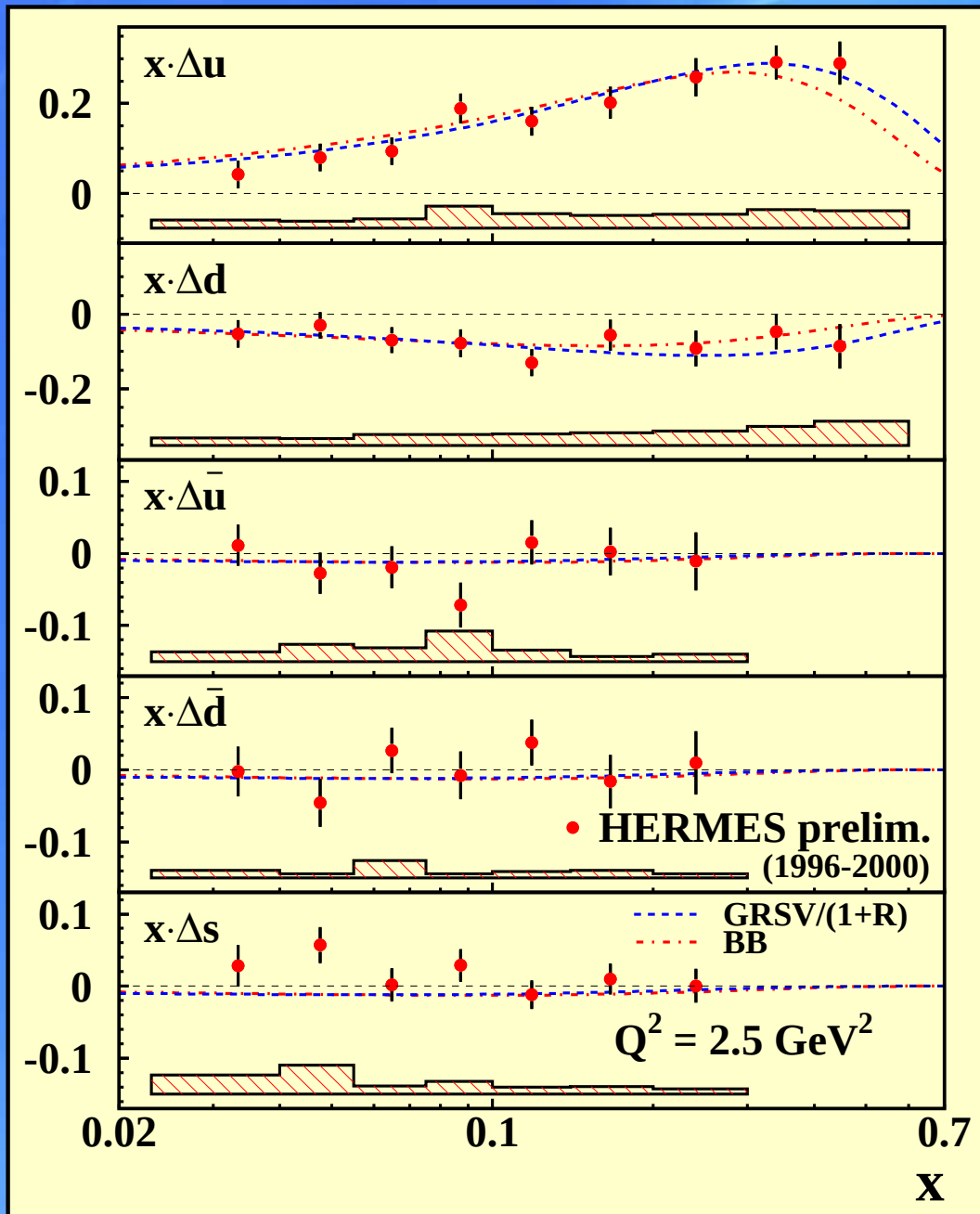
Polarized Quark Densities



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good agreement with NLO-QCD fit

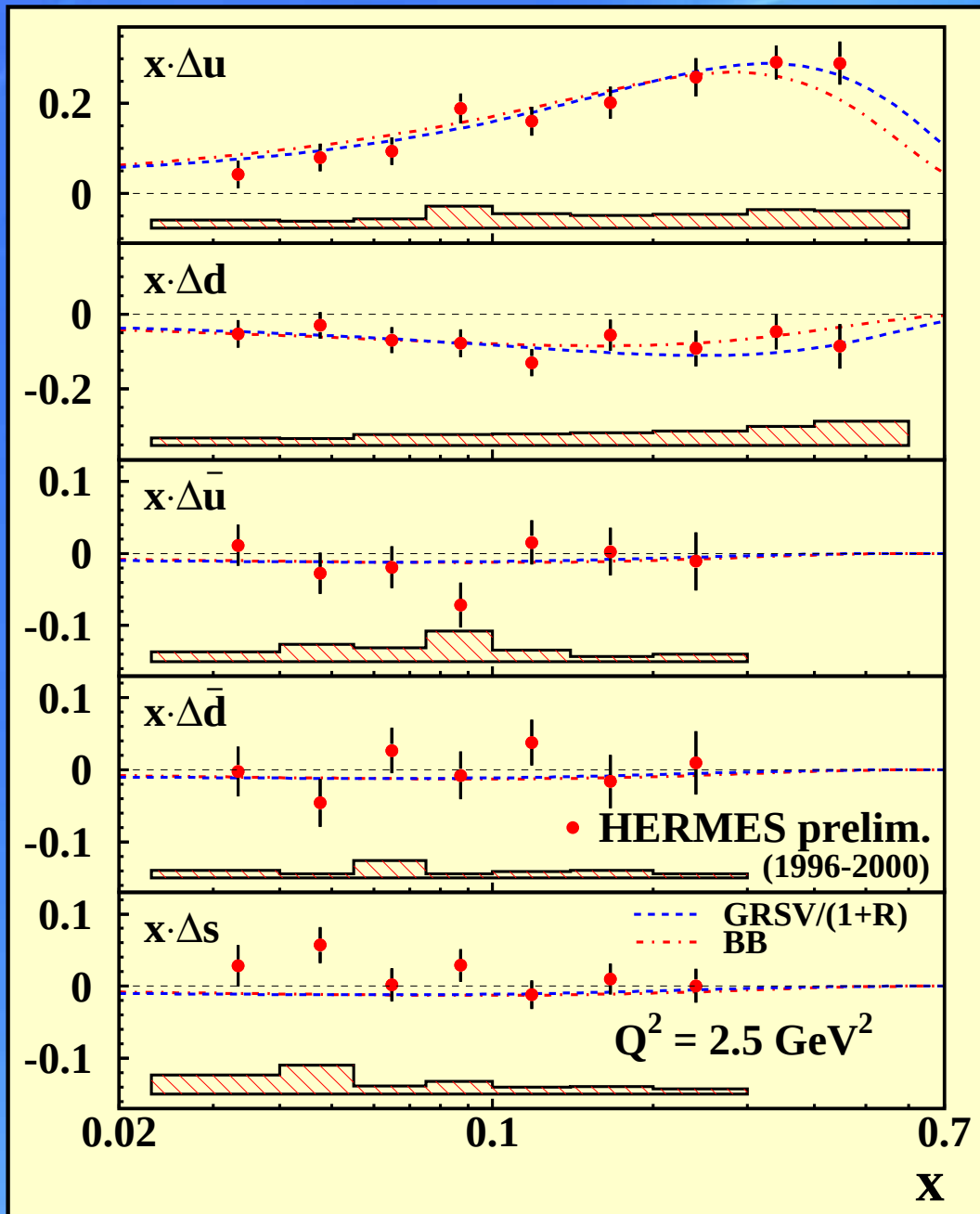
Polarized Quark Densities



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Polarized Quark Densities

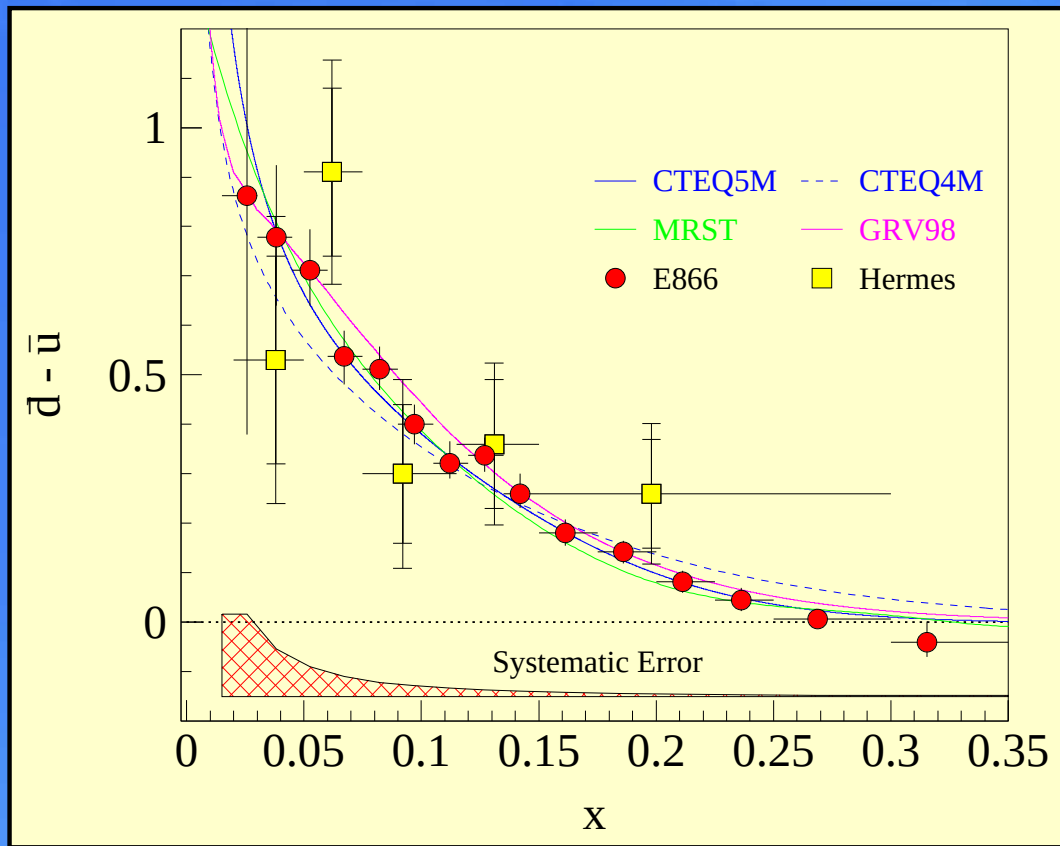


$$\Delta q_f(x) := q_f^+(x) - q_f^-(x)$$

- $\Delta u(x)$ and $\Delta d(x)$
good agreement with NLO-QCD fit
- $\Delta \bar{u}(x), \Delta \bar{d}(x) \sim 0$
- No indication for $\Delta s(x) < 0$

SU(2)-Flavor Symmetry Breaking(?)

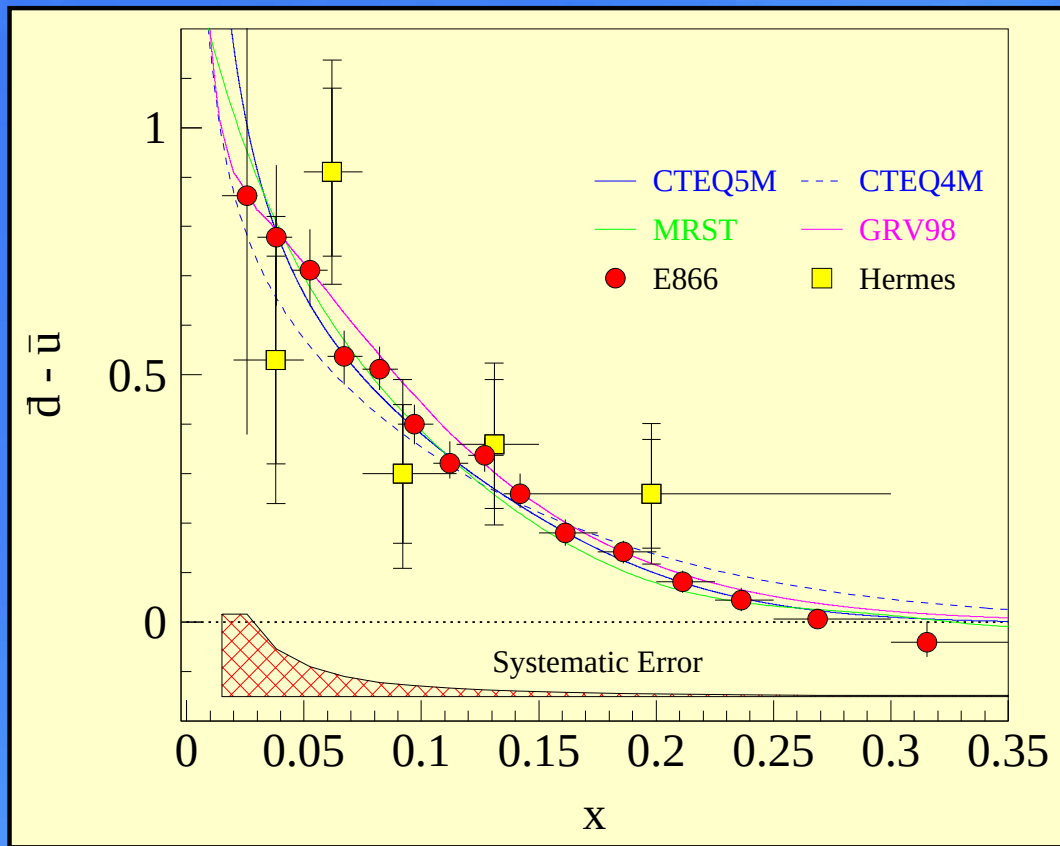
Unpolarised



Strong breaking of SU(2)-Flavor Symmetry

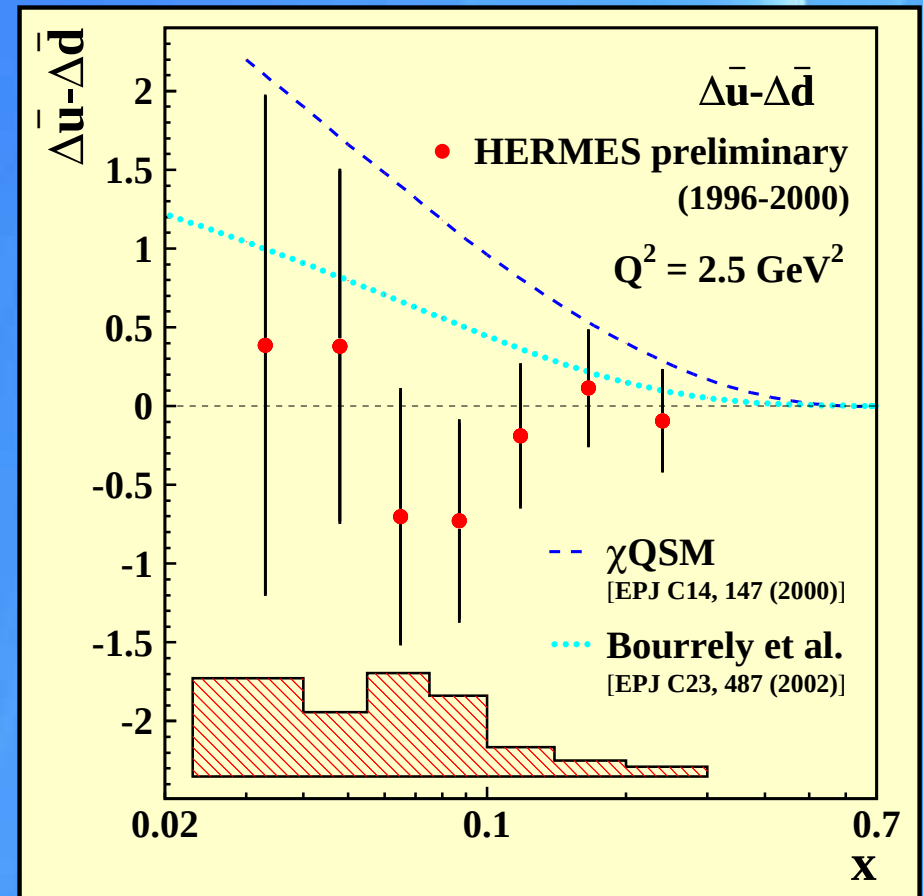
SU(2)-Flavor Symmetry Breaking(?)

Unpolarised



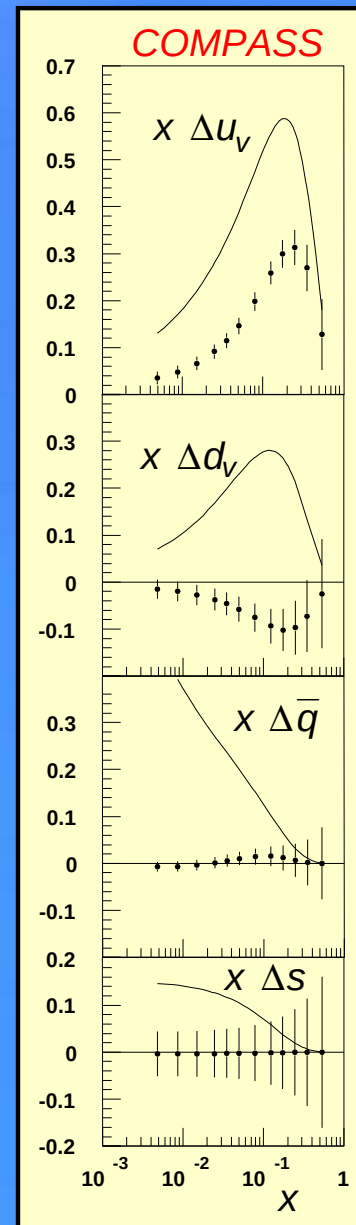
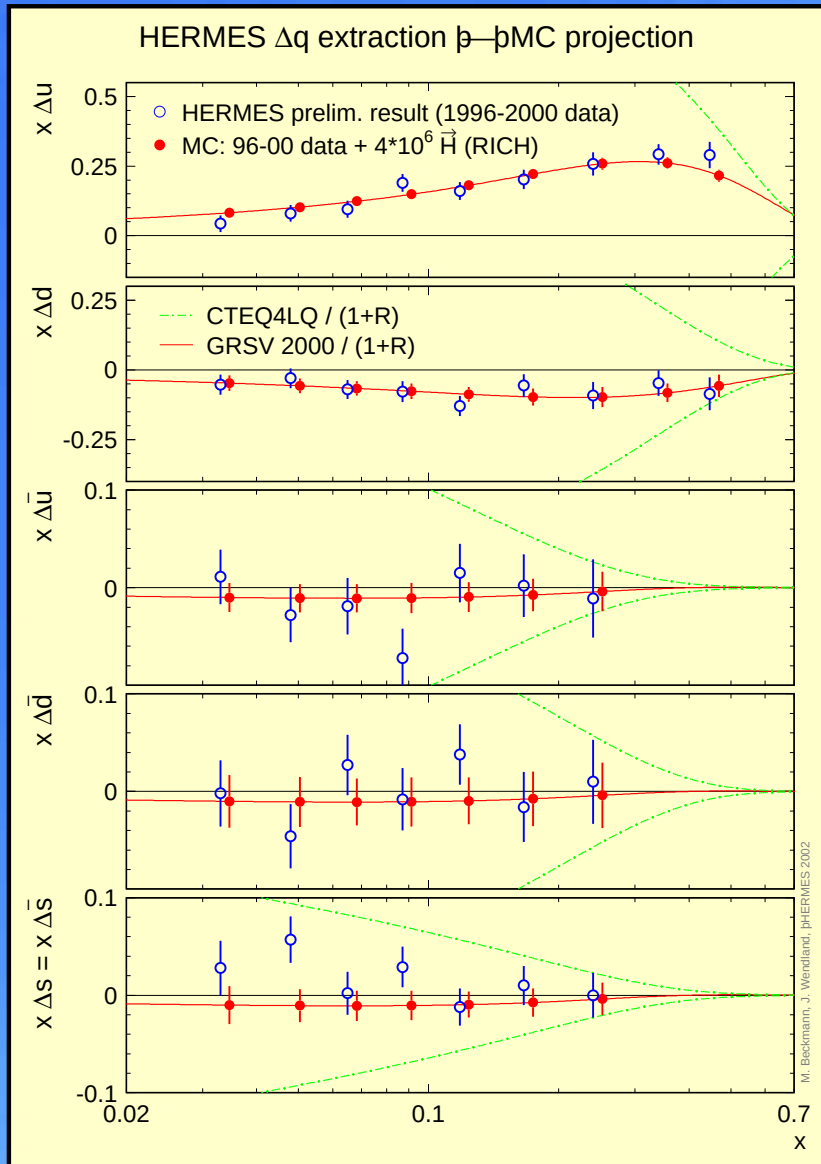
Strong breaking of SU(2)-Flavor Symmetry

Polarised



No significant breaking of SU(2)-Flavor Symmetry $\Delta \bar{u} \sim \Delta \bar{d}$
More Data needed

Future of Polarized Quark Densities



In DIS:

HERMES

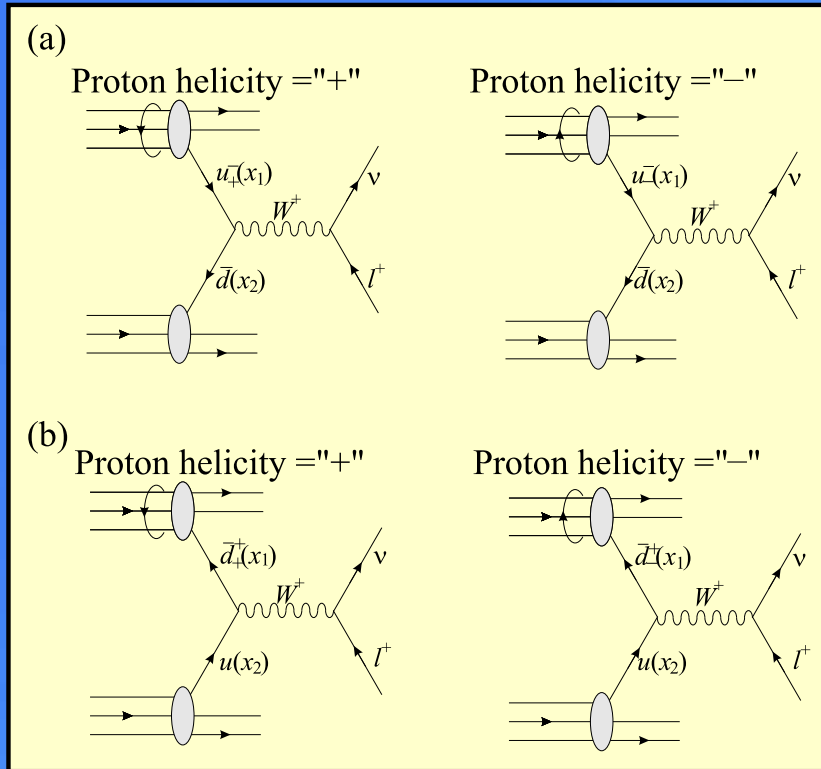
additional 4 Million DIS
with polarized $\vec{H}$ and the RICH

COMPASS

will extend to lower x
expected luminosity:
 $2 \text{ fb}^{-1} / \text{year}$

Future of Polarized Quark Densities at RHIC

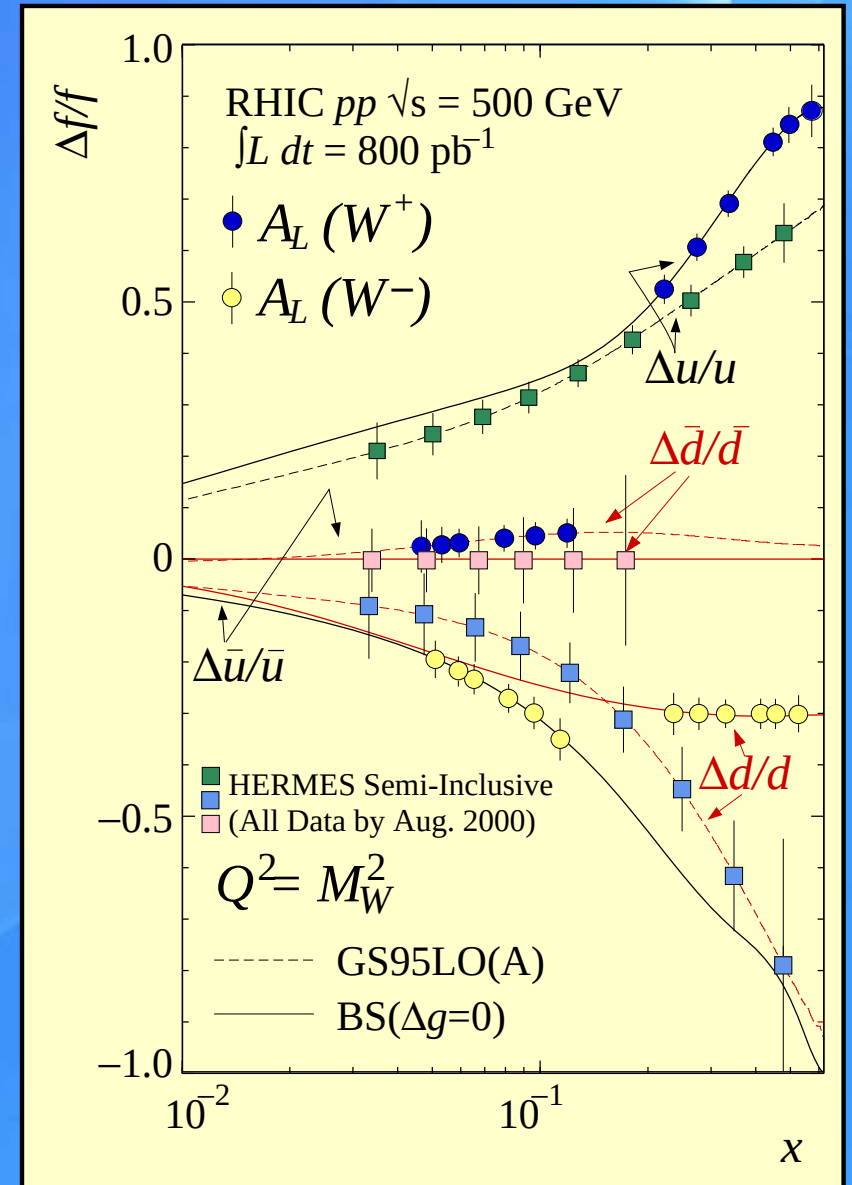
$\frac{\Delta q}{q}$ via parity violating asymmetry in $W^\pm$ production



$$A_L(W^+) = \frac{\Delta u(x_1)\bar{d}(x_2) - \Delta\bar{d}(x_1)u(x_2)}{u(x_1)\bar{d}(x_2) + \bar{d}(x_1)u(x_2)}$$

$$W^+ \Rightarrow u + \bar{d} \quad W^- \Rightarrow d + \bar{u}$$

$$x_1 \gg x_2 \Rightarrow \frac{\Delta u}{u} \quad x_2 \gg x_1 \Rightarrow \frac{\Delta\bar{d}}{\bar{d}}$$

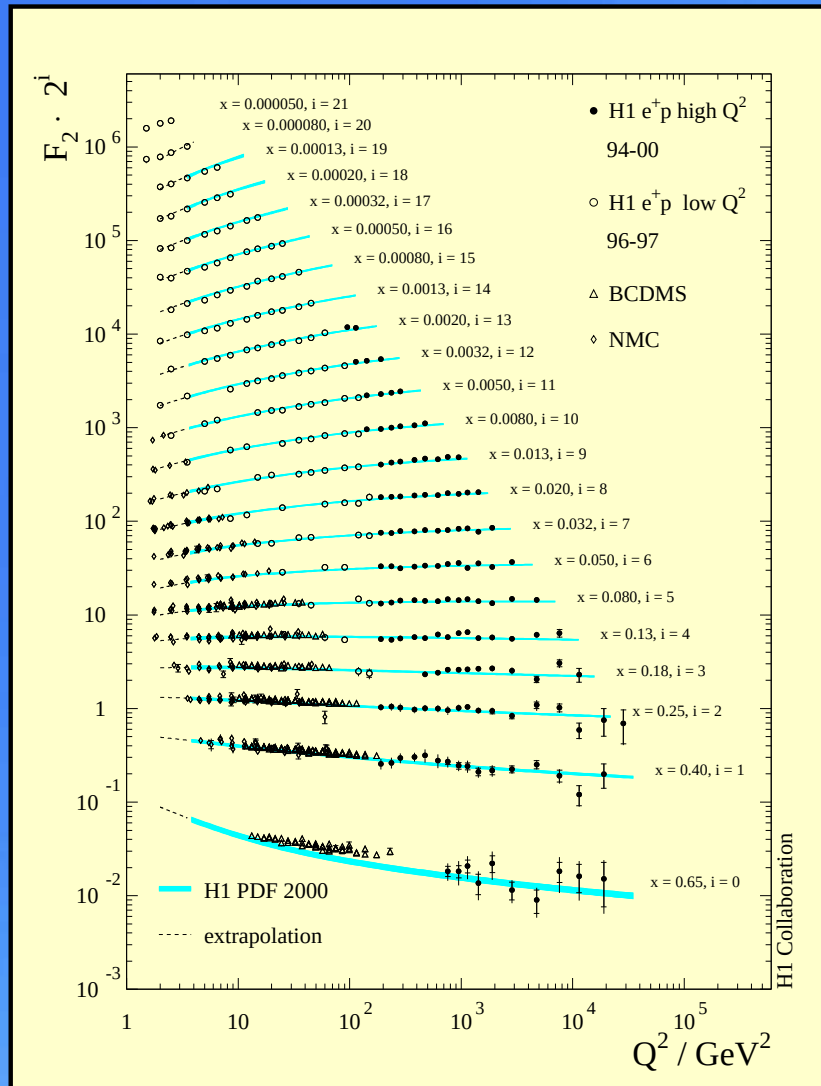


How to measure ΔG

- **"Indirect" from scaling violation**

How to measure ΔG

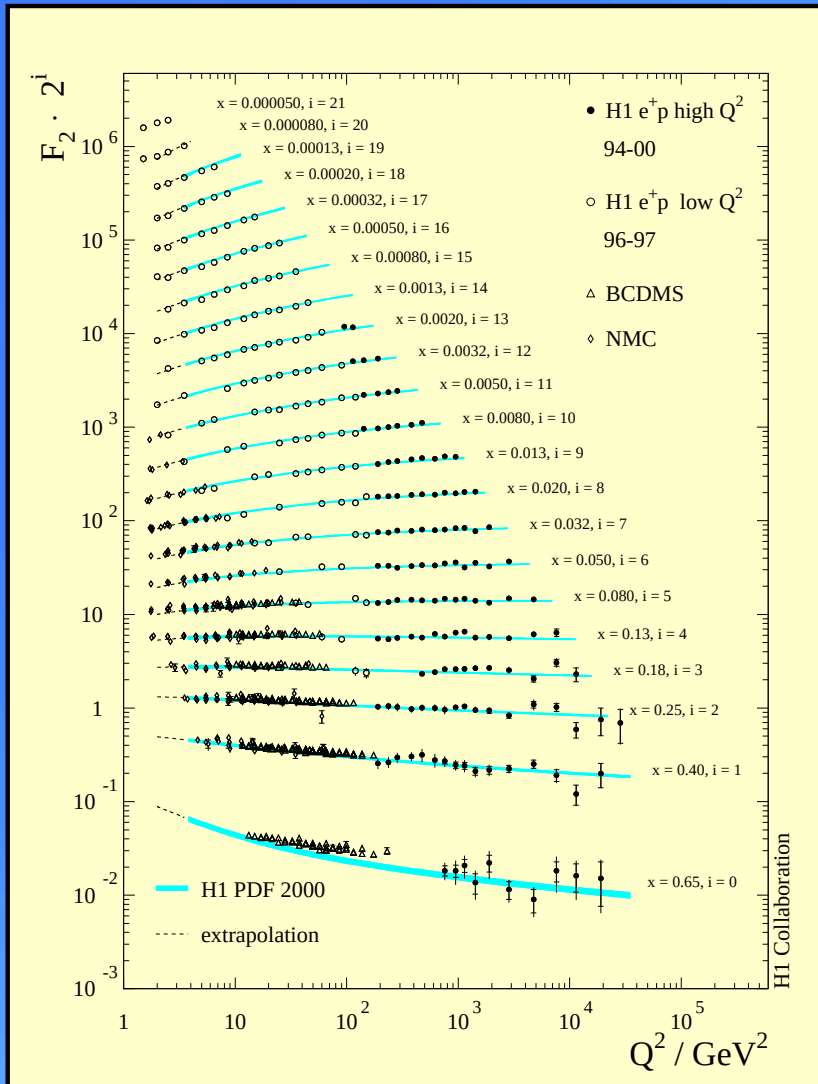
- "Indirect" from scaling violation
Remember unpolarised case:



How to measure ΔG

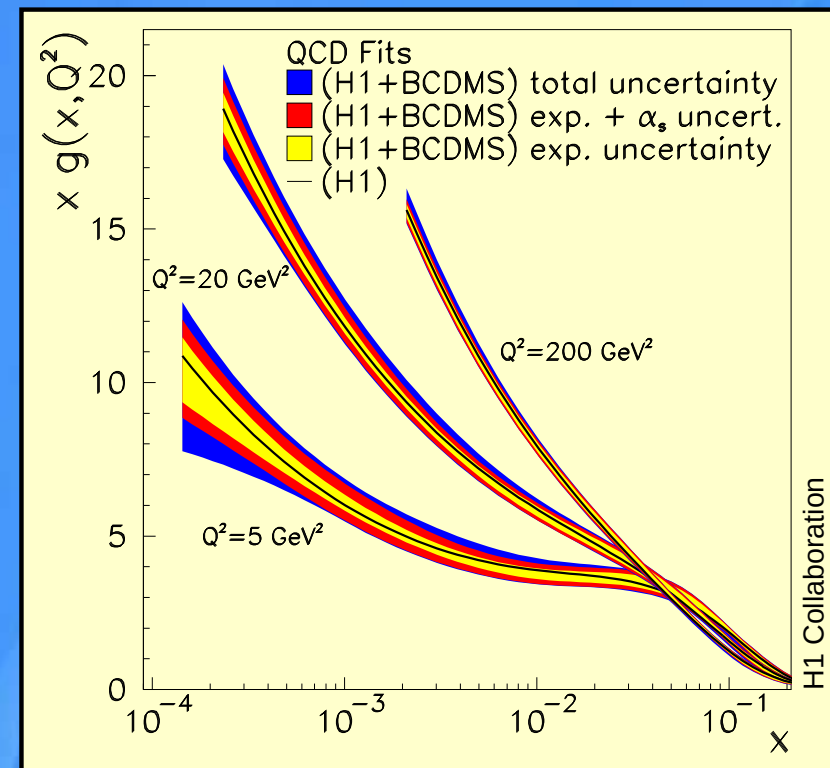
- “Indirect” from scaling violation

Remember **unpolarised case**:



⇒ big $Q^2 - x_{bj}$ lever arm

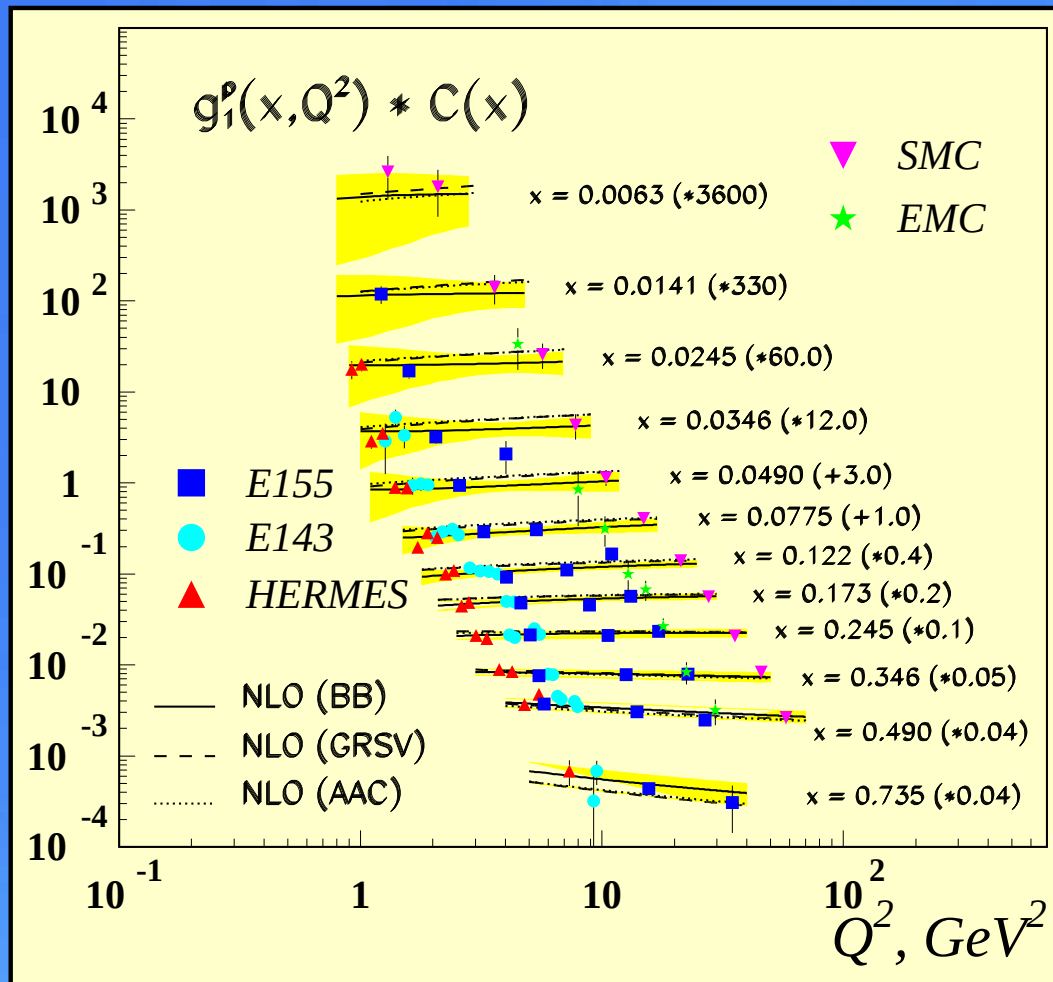
⇒ very accurate $G(x)$



How to measure ΔG

- "Indirect" from scaling violation

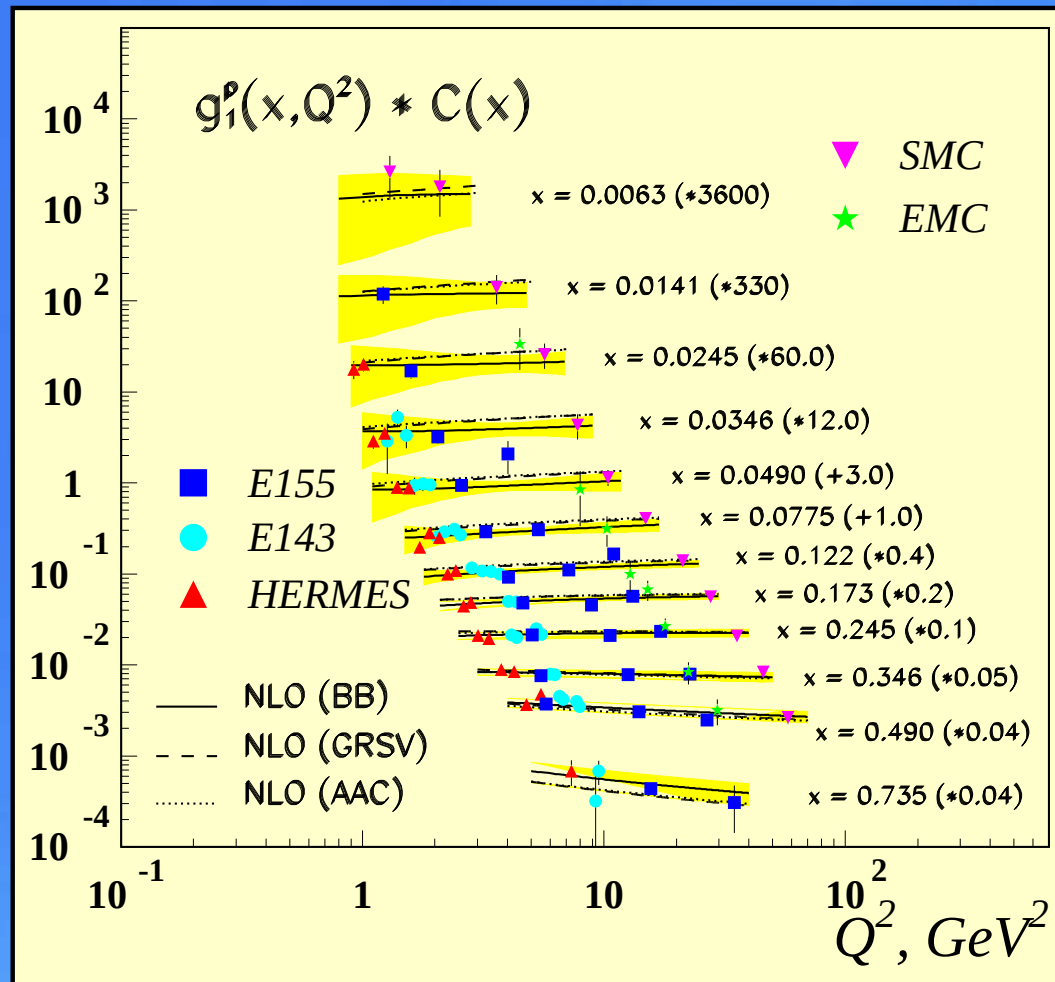
Polarised case:



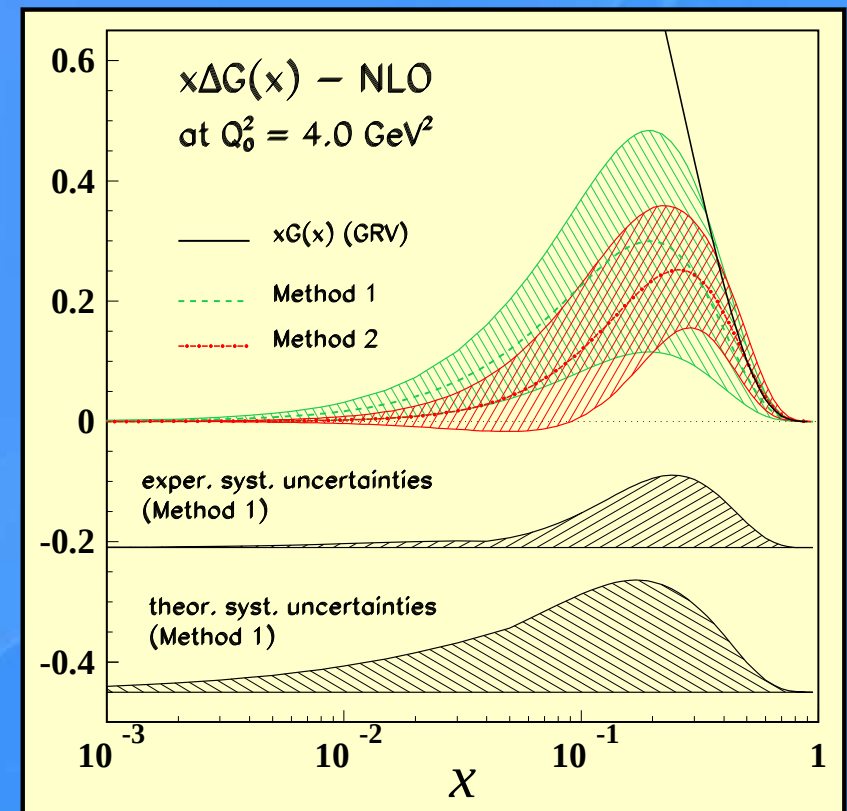
How to measure ΔG

- "Indirect" from scaling violation

Polarised case:



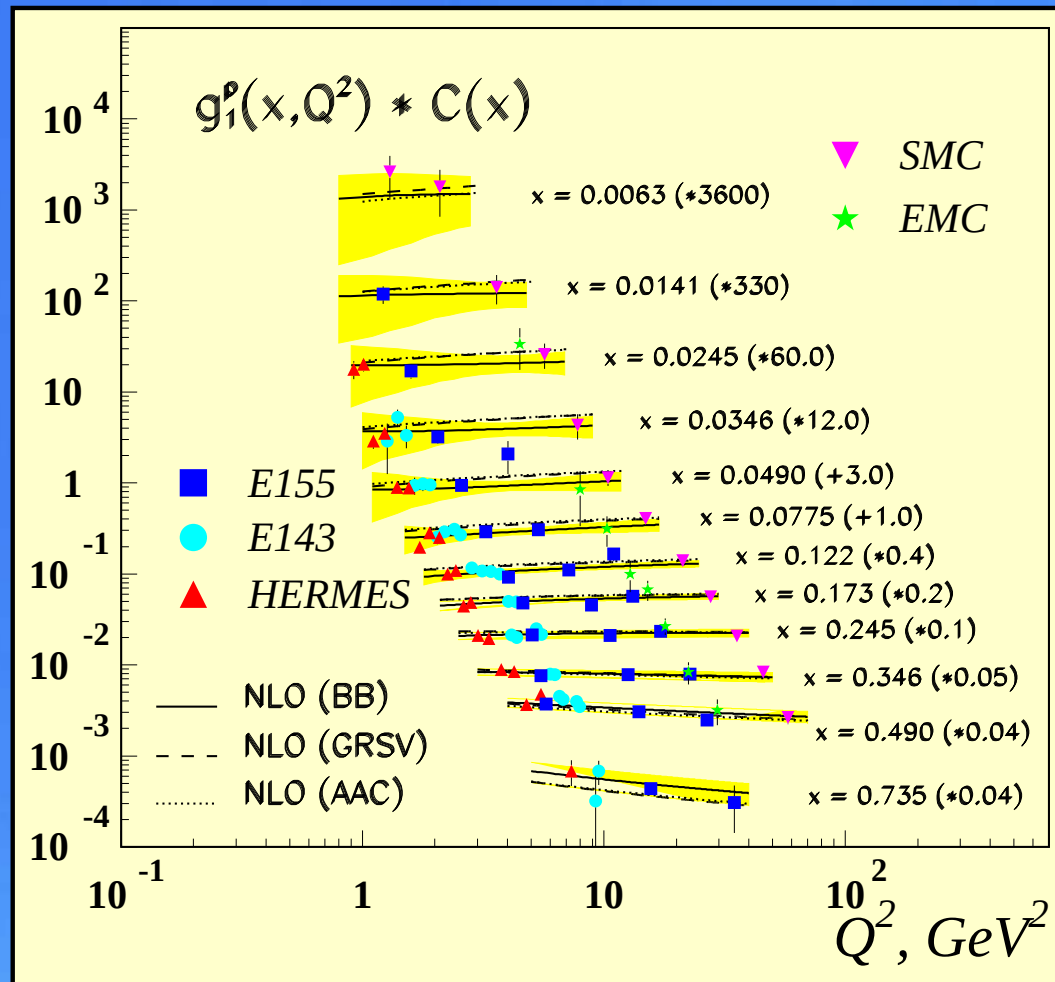
- fixed target experiments
 $\Rightarrow$ small $Q^2 - x_{bj}$ lever arm
- determines only sign of $\Delta G(x)$



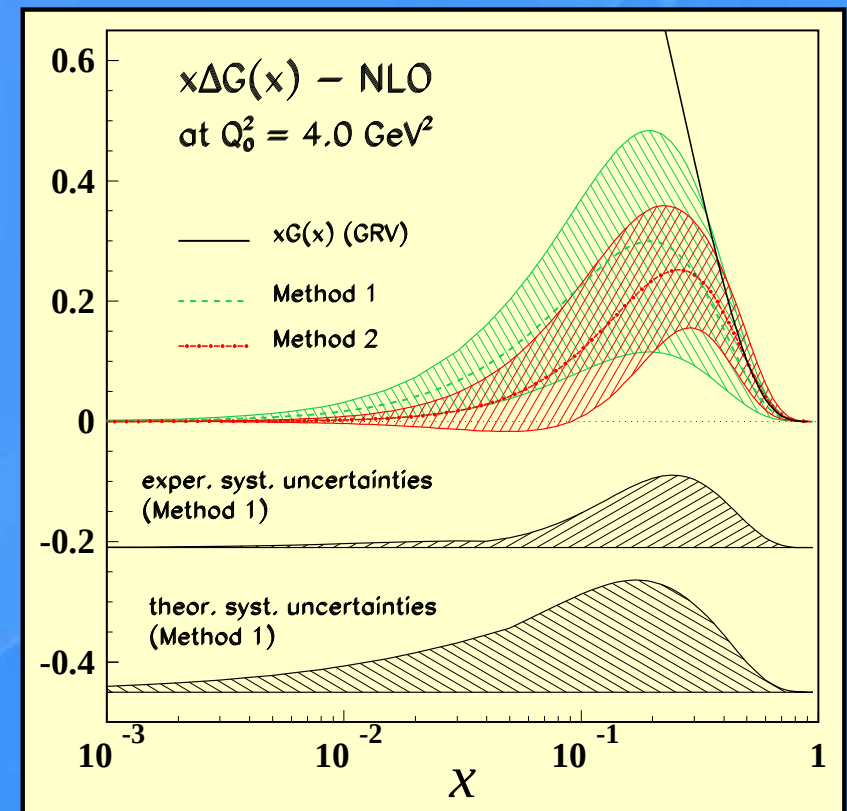
How to measure ΔG

- "Indirect" from scaling violation

Polarised case:



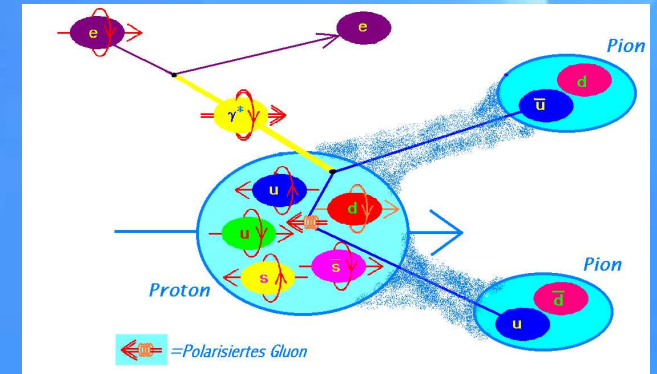
- fixed target experiments
 $\Rightarrow$ small $Q^2 - x_{bj}$ lever arm
- determines only sign of $\Delta G(x)$



$\Rightarrow$ Alternative extraction methods are needed

Direct Measurements of ΔG

Isolate the photon-gluon fusion process (PGF)



Direct Measurements of ΔG

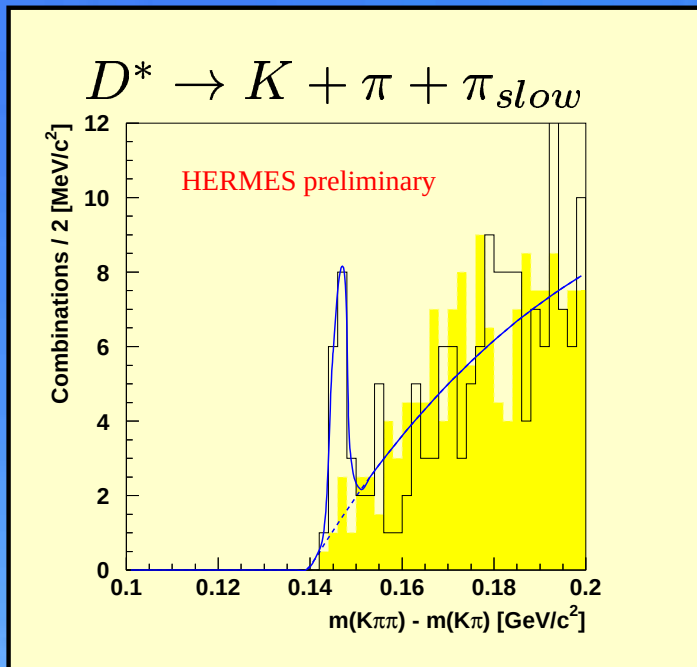
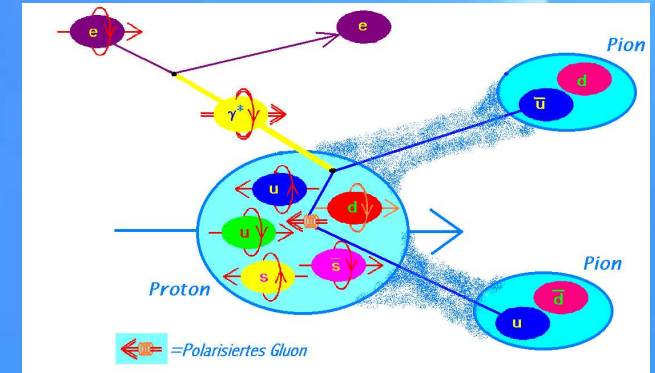
Isolate the photon-gluon fusion process (PGF)

Heavy Quark Production

$c\bar{c} \Rightarrow$ reconstruct D^* , D^0

$$A_{||} = \frac{N_{c\bar{c}}^{\rightarrow\rightarrow} - N_{c\bar{c}}^{\rightarrow\leftarrow}}{N_{c\bar{c}}^{\rightarrow\rightarrow} + N_{c\bar{c}}^{\rightarrow\leftarrow}}$$

$$A^{\gamma p \rightarrow c\bar{c}} \sim \Delta G / G$$



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HIGH- P_T

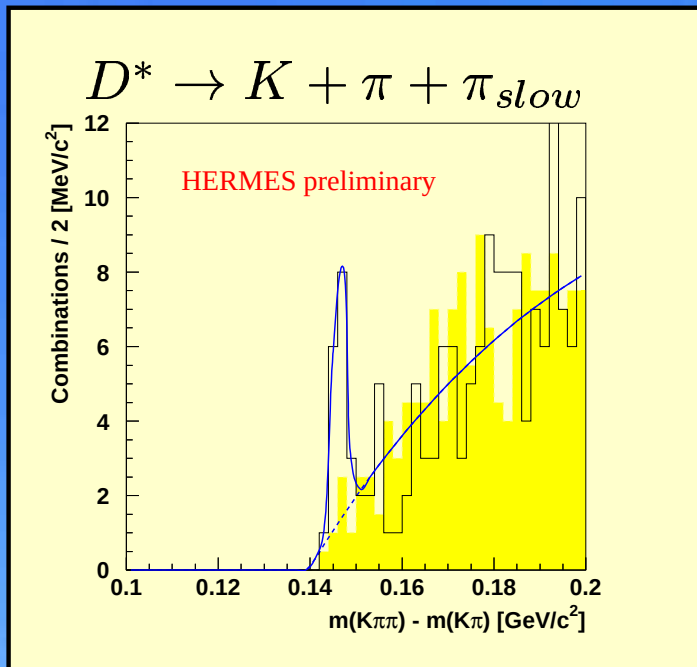
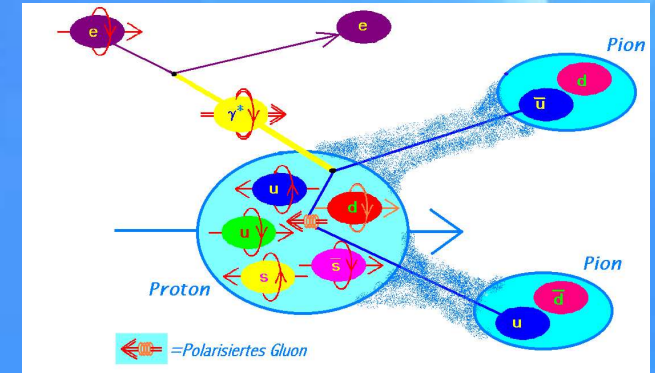
pairs of high- P_T hadrons

$$p_T(h_1^\pm, h_2^\mp) > 1 \text{ GeV}$$

$$A_{||} = \frac{N_{h_1^\pm h_2^\mp}^{\leftarrow} - N_{h_1^\pm h_2^\mp}^{\rightarrow}}{N_{h_1^\pm h_2^\mp}^{\leftarrow} + N_{h_1^\pm h_2^\mp}^{\rightarrow}}$$

$$A^{\gamma^* p \rightarrow h_1^\pm + h_2^\mp} \sim \Delta G/G$$

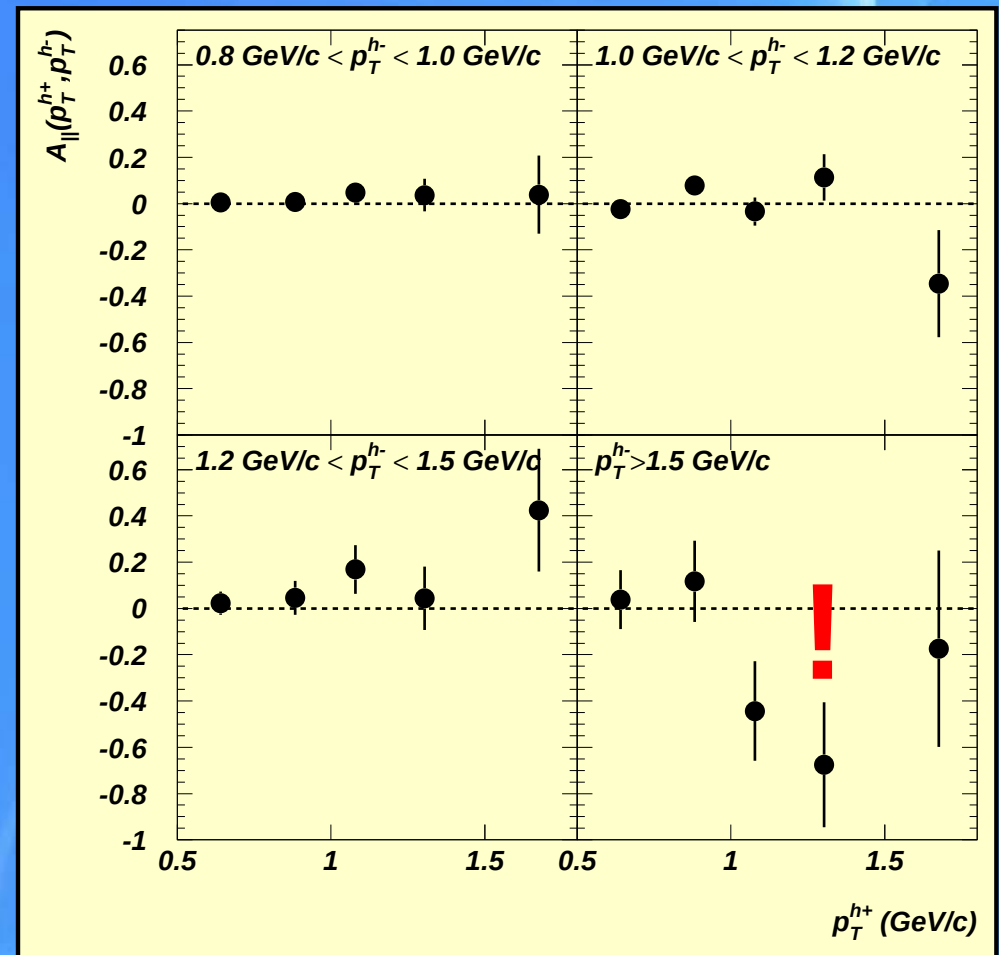
additionally:
use identified hadrons



The HERMES hunt for ΔG

- Reaction: $\gamma^* p \rightarrow h_1^\pm + h_2^\mp + X$
- select two oppositely charged hadrons
- low Q^2 range ($Q^2 > 0 \text{ GeV}^2$)
 - ⇒ increase statistics
 - ⇒ scale of hard sub-processes $\hat{p}_T^2$
- require $p_T(h_1, h_2) > 0.5 \text{ GeV}$,
 $M(2\pi) > 1.0 \text{ GeV}$
 - ⇒ removes resonances ρ and ϕ

- $A_{||} = \frac{N_{h_1^\pm h_2^\mp}^{\leftarrow} - N_{h_1^\pm h_2^\mp}^{\rightarrow}}{N_{h_1^\pm h_2^\mp}^{\leftarrow} + N_{h_1^\pm h_2^\mp}^{\rightarrow}} \text{ vs. } p_T(h_1^\pm, h_2^\mp)$



Negative asymmetry in DIS unexpected

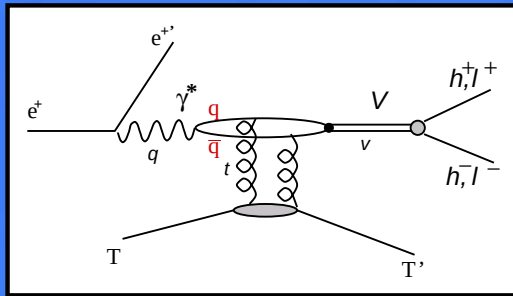
Interpretation ?!

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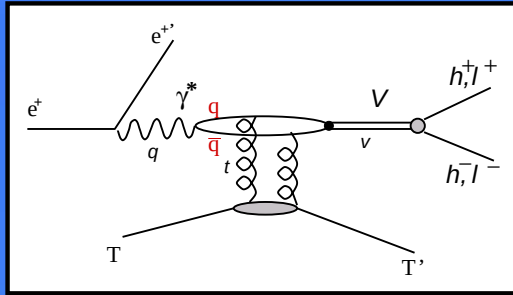


VMD

$$A_{\text{VMD}} = 0$$

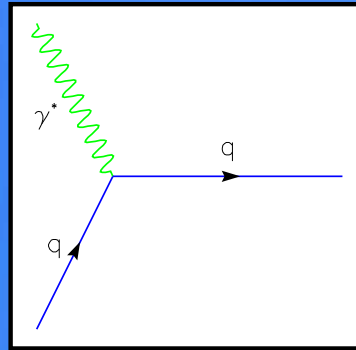
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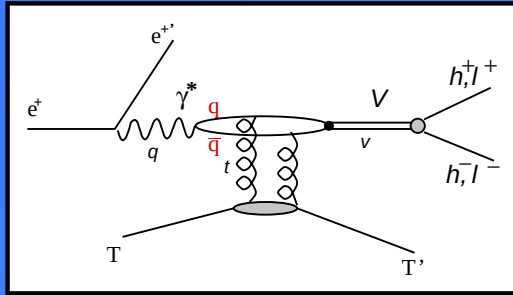


DIS

$$A_{\text{DIS}} \sim \frac{\Delta q}{q}$$

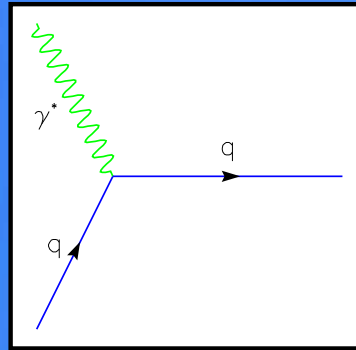
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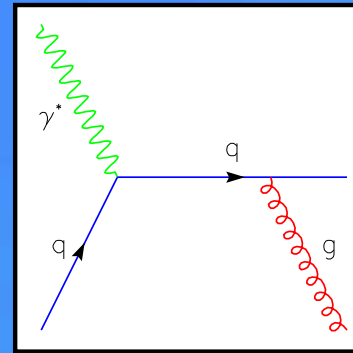
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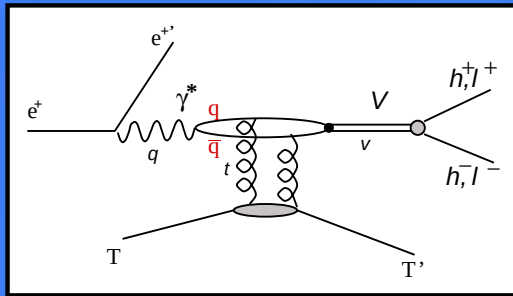


QCDC

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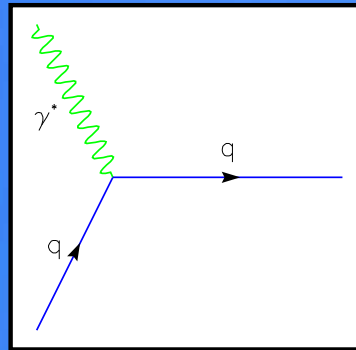
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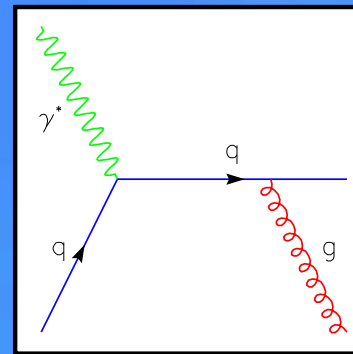
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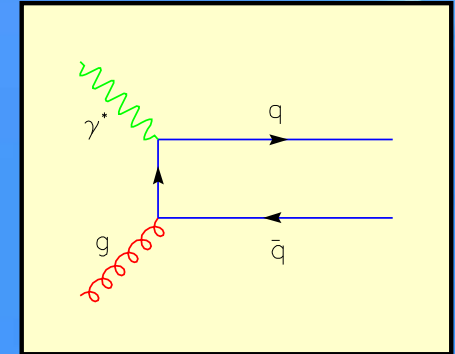
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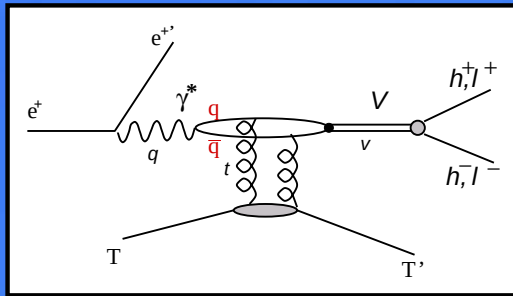


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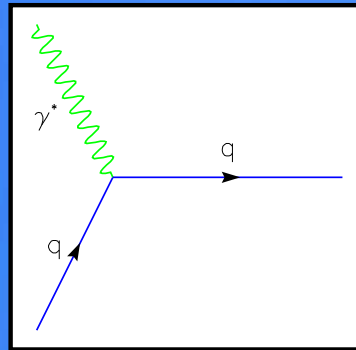
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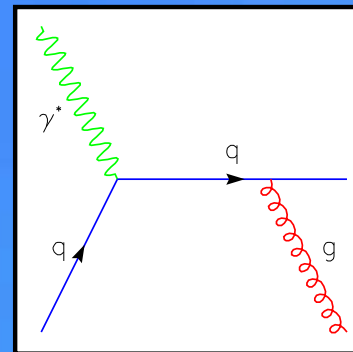
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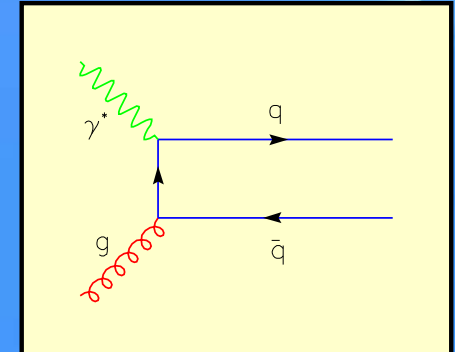
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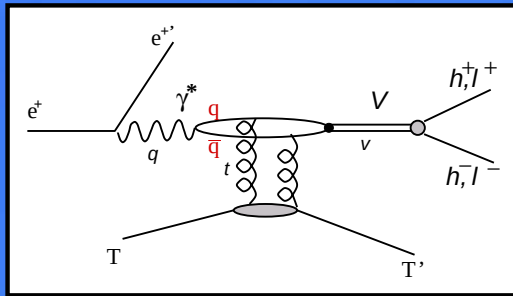
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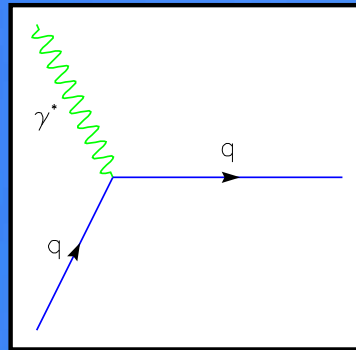
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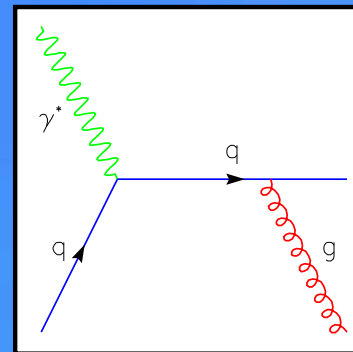
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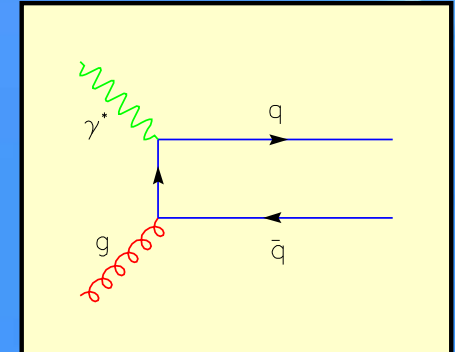
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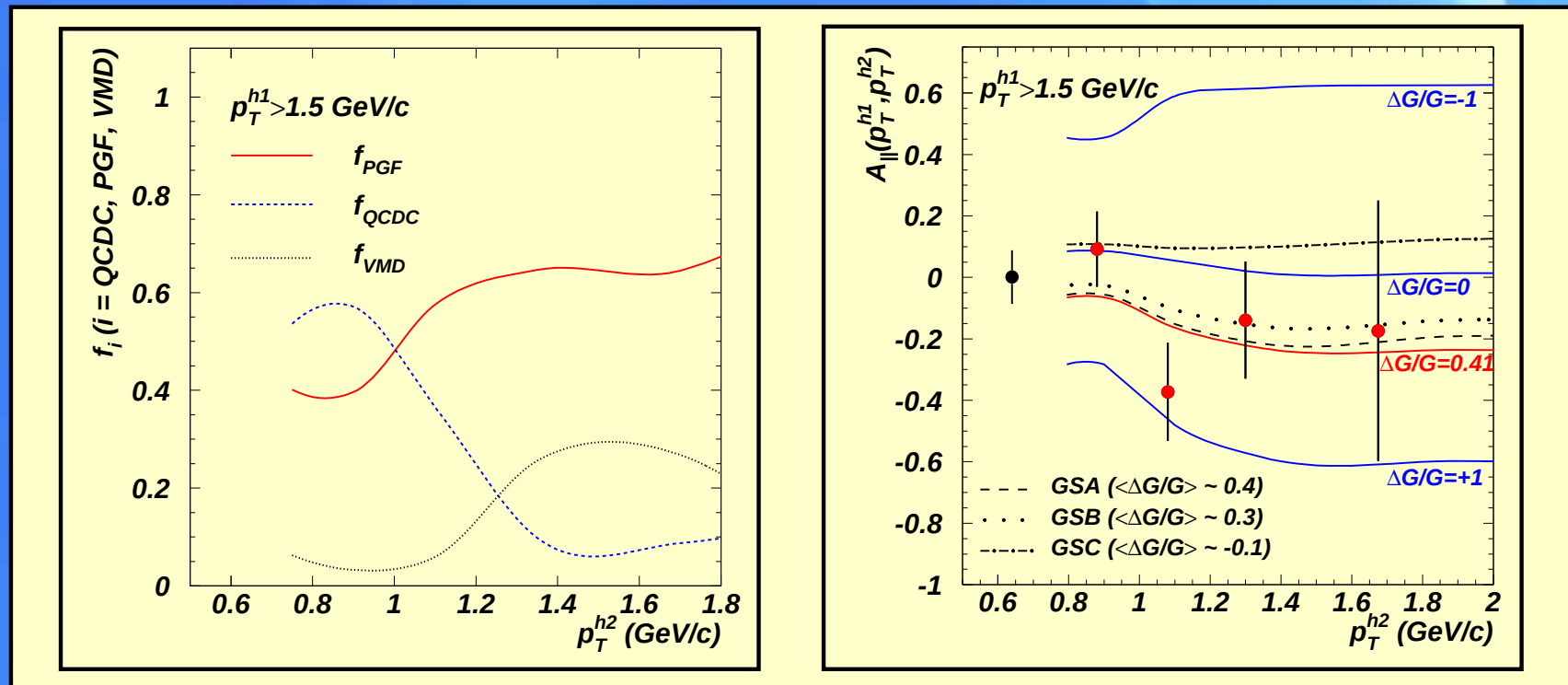
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.. estimate their
relative contributions f_i
using Monte Carlo
 $\Rightarrow \Delta G/G$

Pairs of high- P_T Hadrons



within LO pQCD and PYTHIA5 MC model

$$\Delta G/G = 0.41 \pm 0.18 \text{ (stat.)} \pm 0.03 \text{ (exp.syst.)}$$

$$\text{at } \langle x_G \rangle = 0.17 \text{ and } \langle \hat{p}_T^2 \rangle = 2.1 \text{ GeV}^2$$

Extraction strongly Model dependent

More data on polarised Deuterium available

⇒ Different Data needed to get $\Delta G/G$

COMPASS

- pairs of high p_T hadrons
 $\Rightarrow \Delta G/G : 0.04 < x_g < 0.2$
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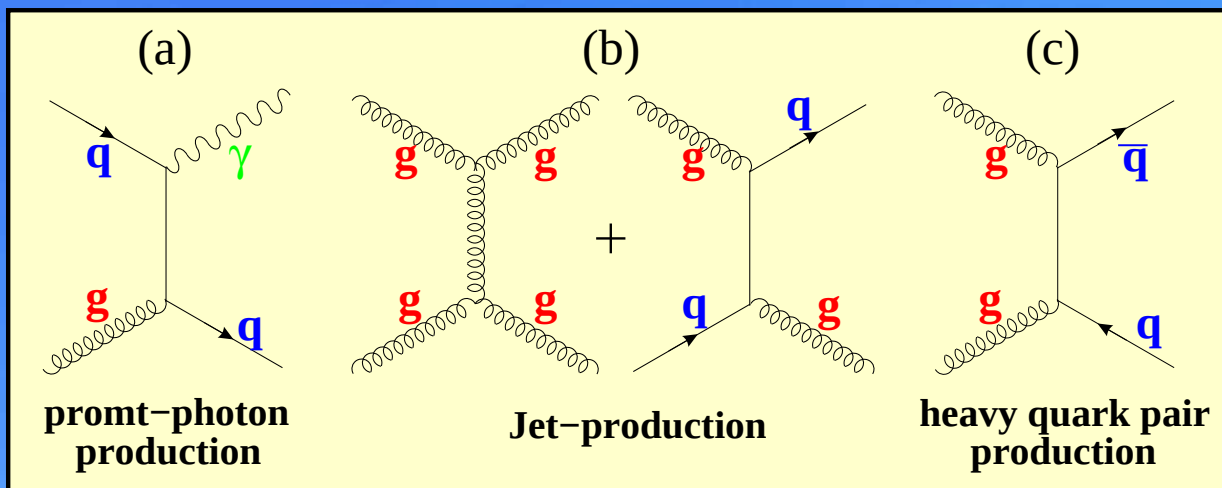
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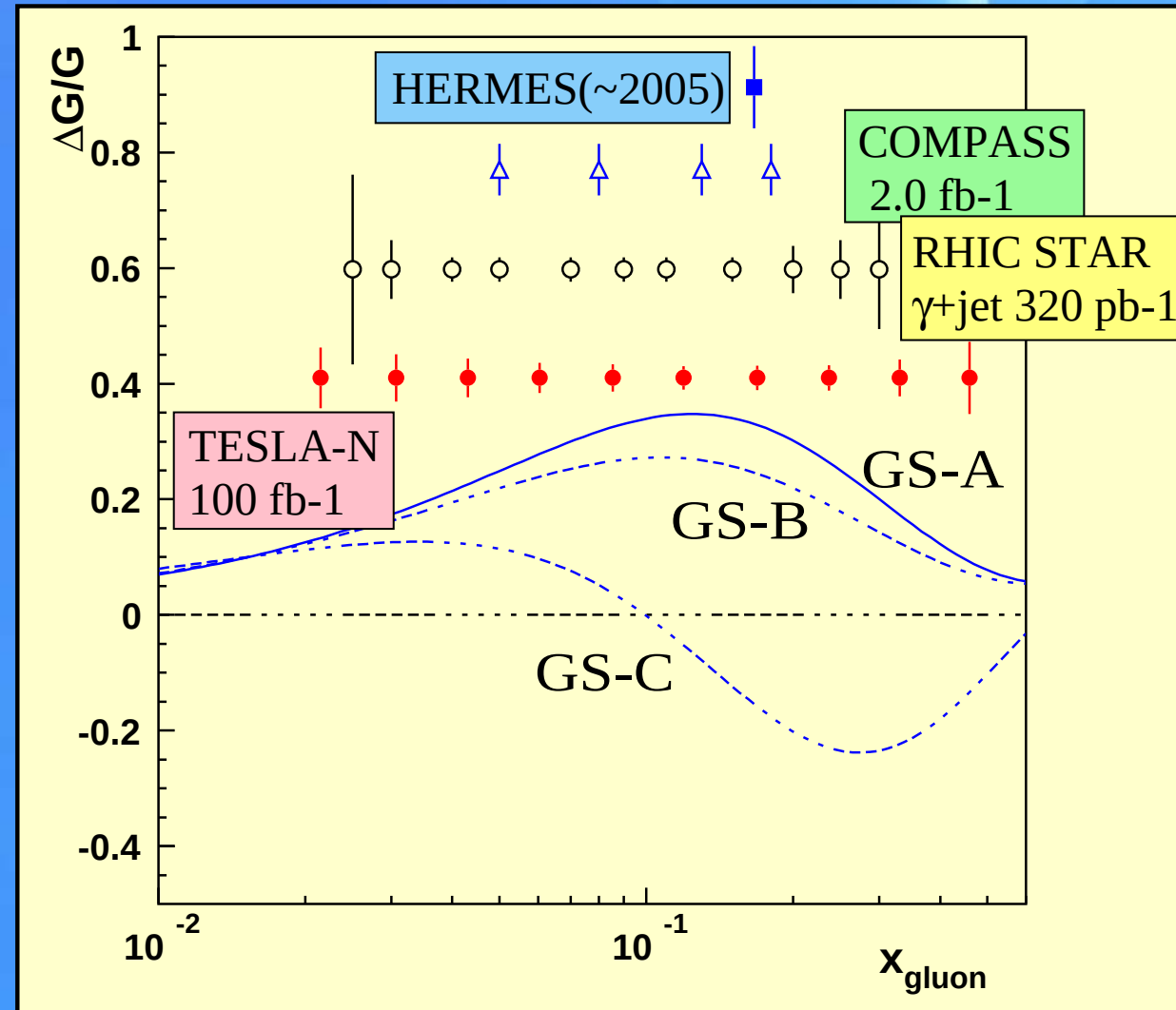
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