

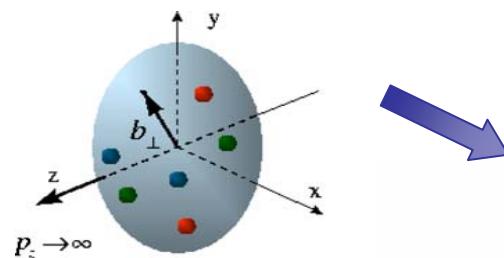
towards a 3D imaging of hadrons: GPDs

- a brief introduction (an experimentalists point of view)
- a personal selection of recent results
- models & data
- conclusion & perspectives

nucleon studied for decades:

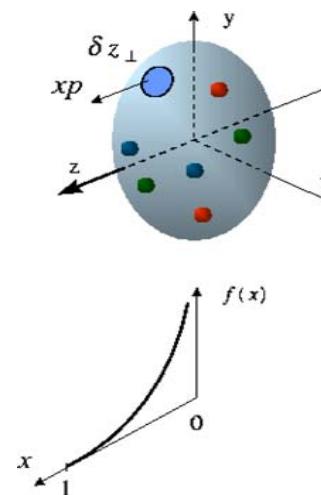
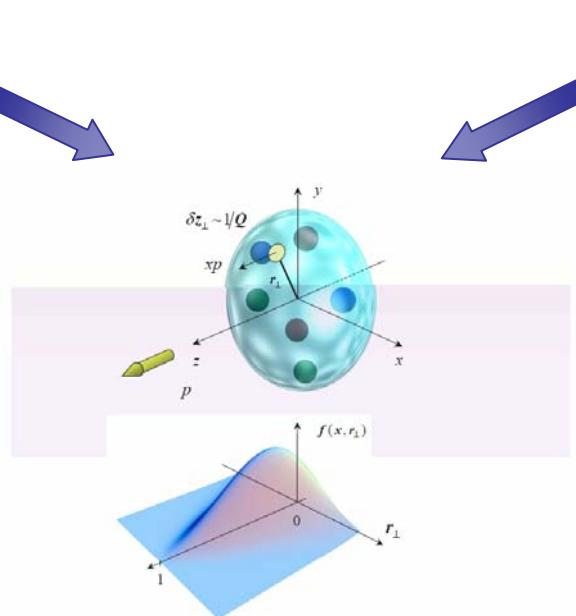
form factors

location of partons in nucleon



parton distributions

longitudinal momentum fraction x



generalised parton distributions (GPDs)

longitudinal momentum fraction x at transverse location b_{\perp}

only known framework to gain information on 3D picture of hadrons

why GPDs ?

→ 3D structure of hadrons : nucleon tomography

nucleon tomography

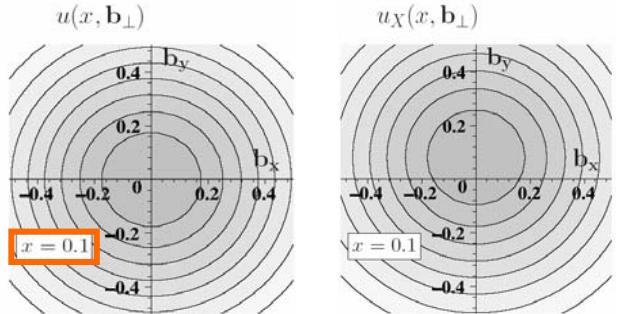
[M. Burkardt, M. Diehl 2002]

$\mathcal{FT}(\text{GPD})$: momentum space \rightarrow impact parameter space:

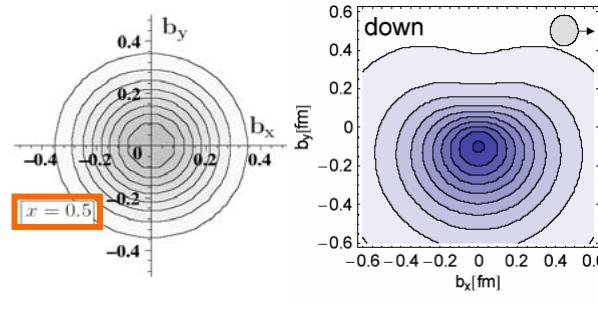
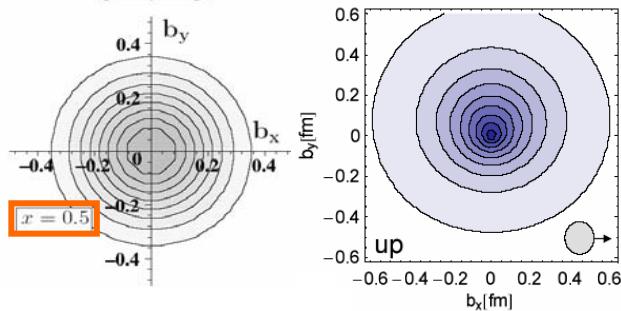
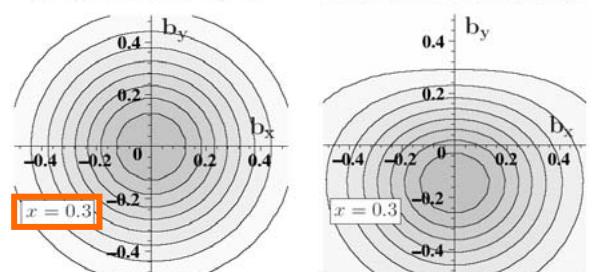
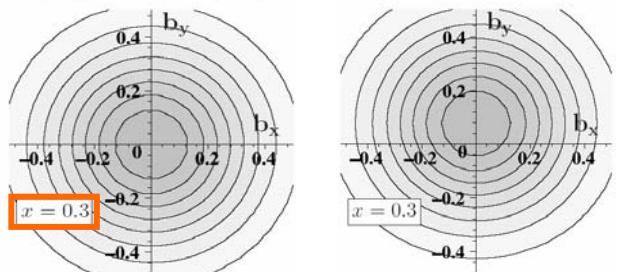
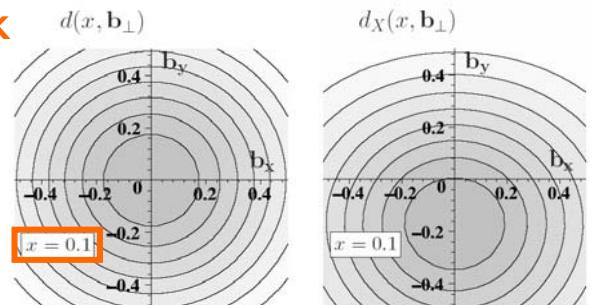
probing partons with specified long. momentum @transverse position \mathbf{b}_\perp

polarised nucleon:

u-quark



d-quark

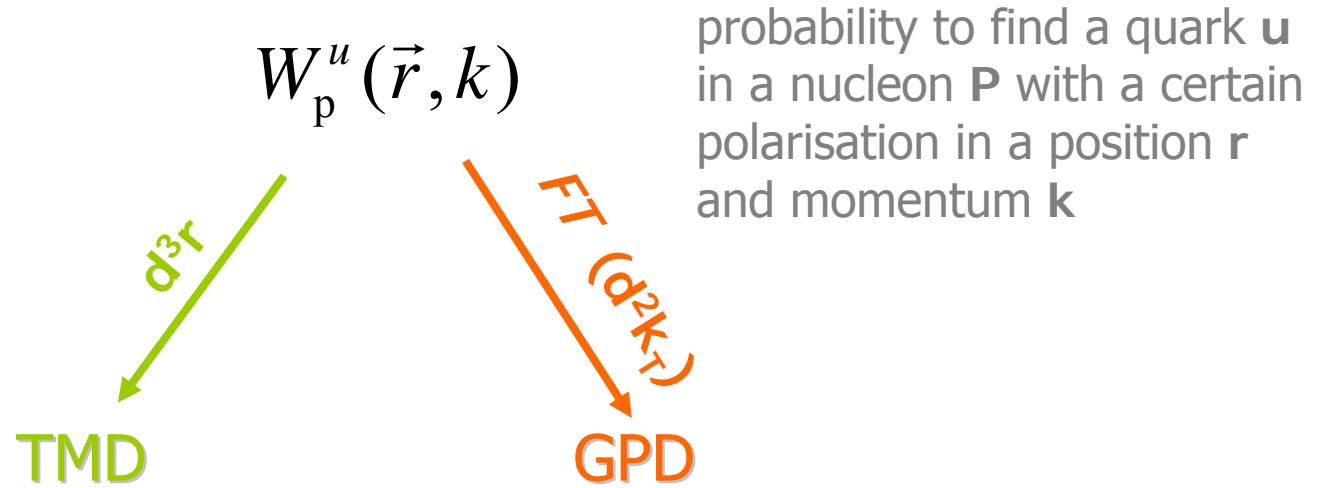


from
lattice

why GPDs ?

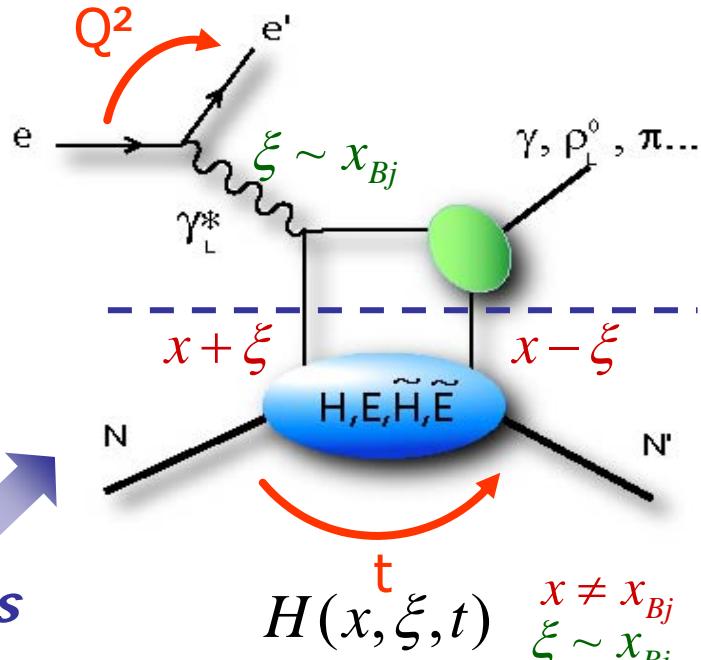
- 3D structure of hadrons : nucleon tomography
- complementary to TMDs :

Wigner distribution: ("mother" function)



- phenomena of single-spin asymmetries

what do we know about GPDs ?



form factors

$$\sum_q e_q \int dx H^q(x, \xi, t) = F_1(t)$$

:

$Q^2 \gg, t \ll$

appear in factorisation theorem for *hard exclusive processes*

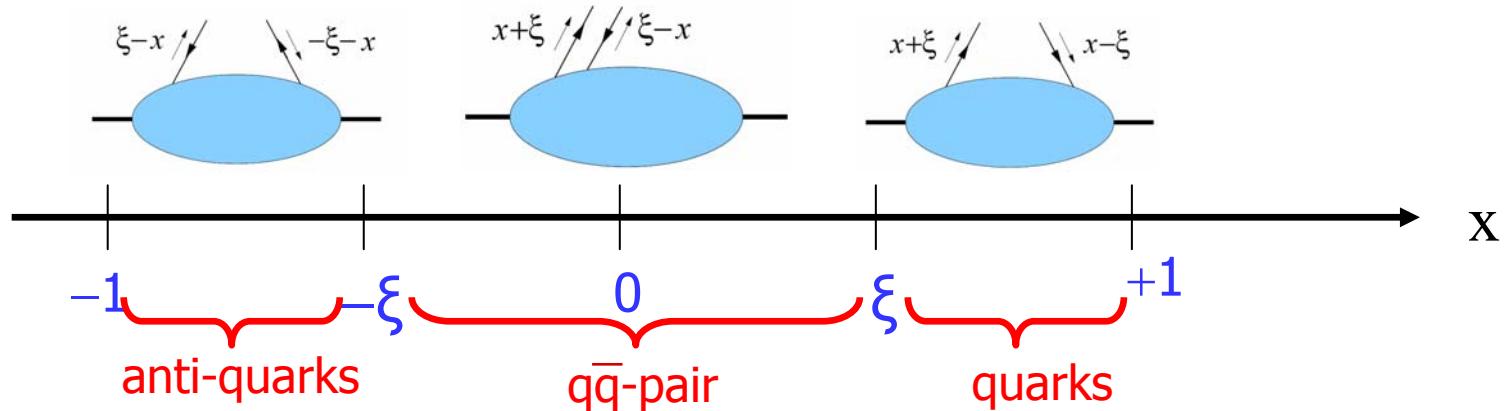
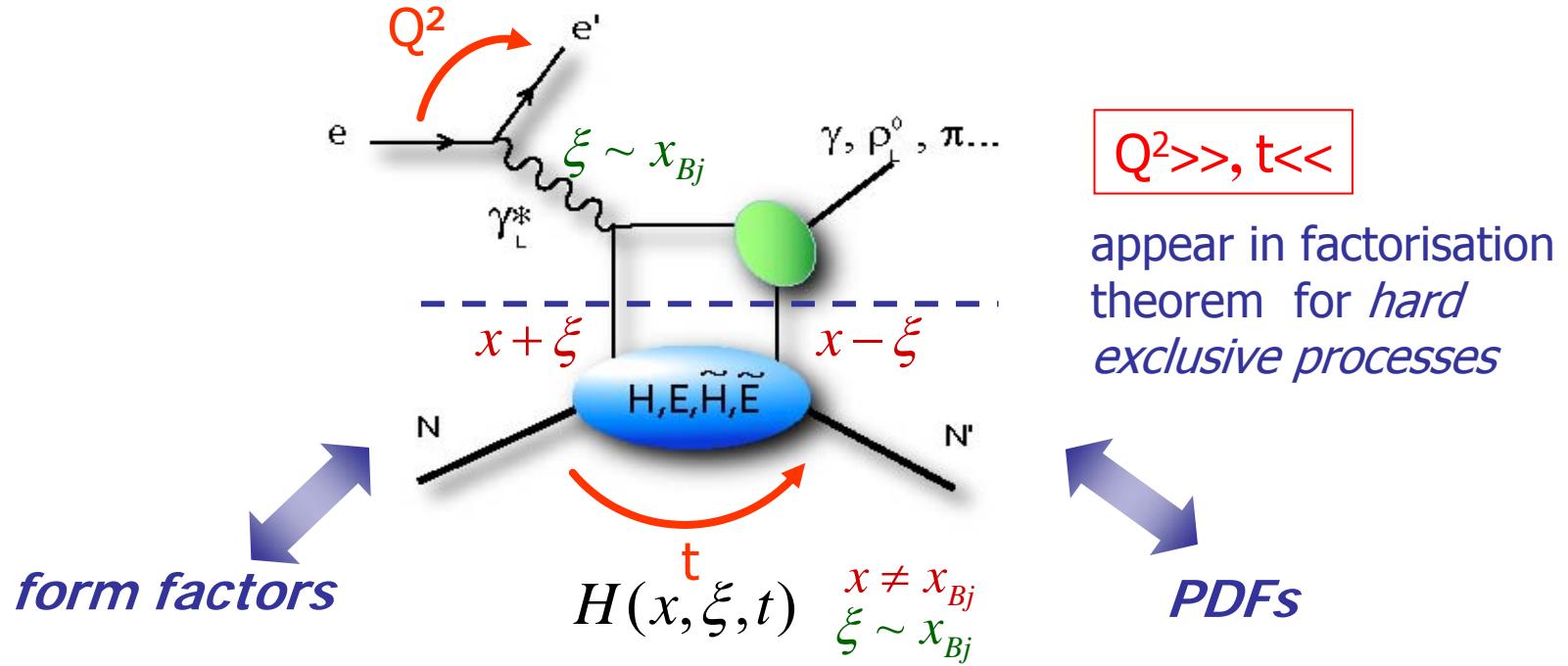
PDFs

$$H^{q,g}(x, 0, 0) = q(x)$$

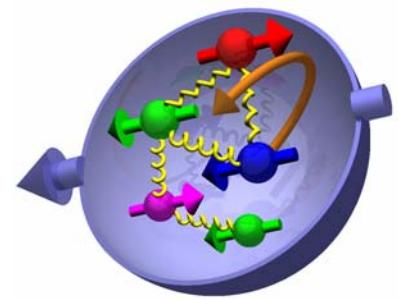
$$\tilde{H}^{q,g}(x, 0, 0) = \Delta q(x)$$

E, \tilde{E} : *nucleon helicity flip* \rightarrow don't appear in DIS
 \rightarrow new information !

what do we know about GPDs ?



GPDs and the spin puzzle



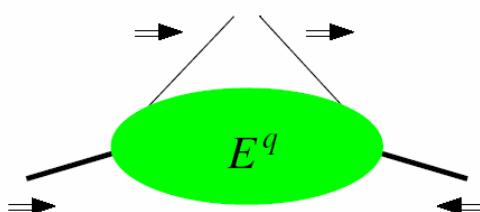
nucleon spin:

$$s_z^n = \frac{1}{2} = \frac{1}{2} \sum_q \Delta q + L_z^q + \Delta G + L_z^g = J_q + J_g$$

\uparrow \uparrow
 $\approx 30\%$ $\approx \text{zero}$

[X. Ji, 1997]

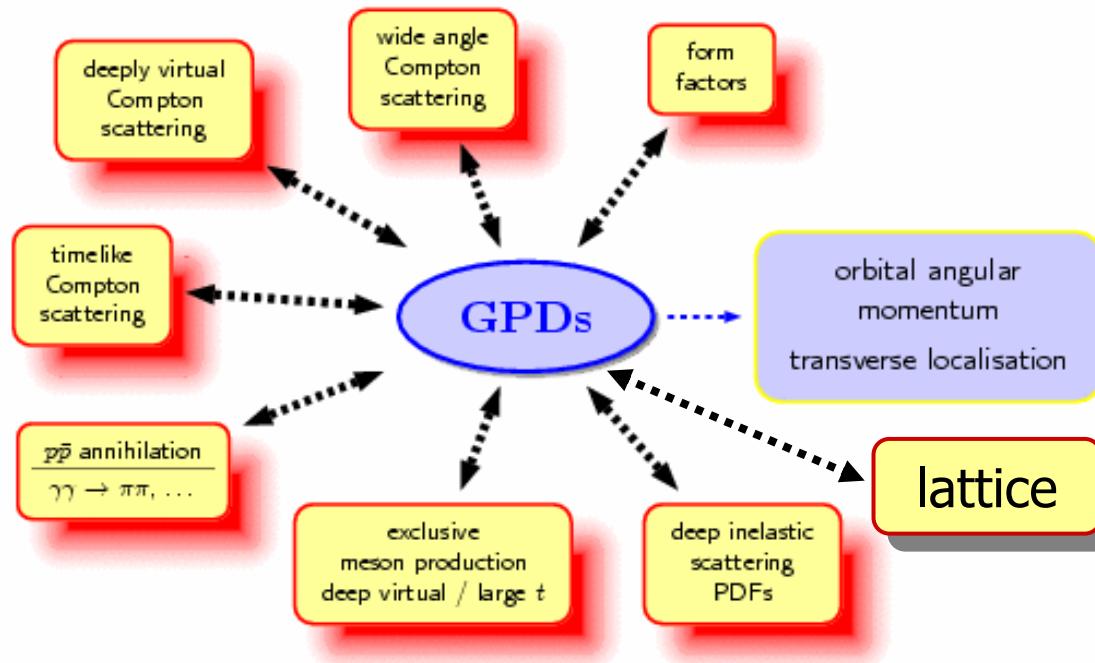
$$J_{q,g} = \lim_{t \rightarrow 0} \frac{1}{2} \int_{-1}^1 x dx \left[H^{q,g}(x, \xi, t) + E^{q,g}(x, \xi, t) \right]$$



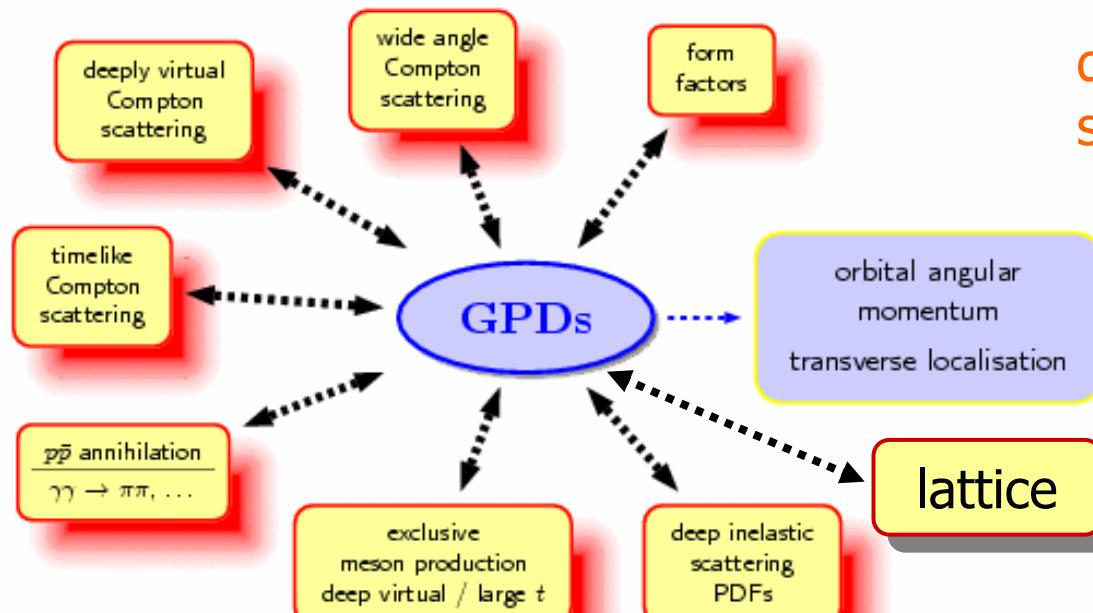
$E^q \neq 0$ requires orbital angular momentum

proton helicity flipped but quark helicity conserved

how to access GPDs ?



how to access GPDs ?



quantum number of final state selects different GPDs:

- VM (ρ, ω, ϕ): H E
- PS mesons (π, η): $\tilde{H} \tilde{E}$
- DVCS (γ): H, E, \tilde{H} , \tilde{E}

π^0	$2\Delta u + \Delta d$
η	$2\Delta u - \Delta d$
ρ^0	$2u+d, 9g/4$
ω	$2u-d, 3g/4$
ϕ	s, g
ρ^+	$u-d$
J/ψ	g

→ DVCS most clean process for gaining information on GPDs

→ meson provide info on quark flavours
 VM: quark and gluon GPDs appear at same order α_s

accessing GPDs: caveats

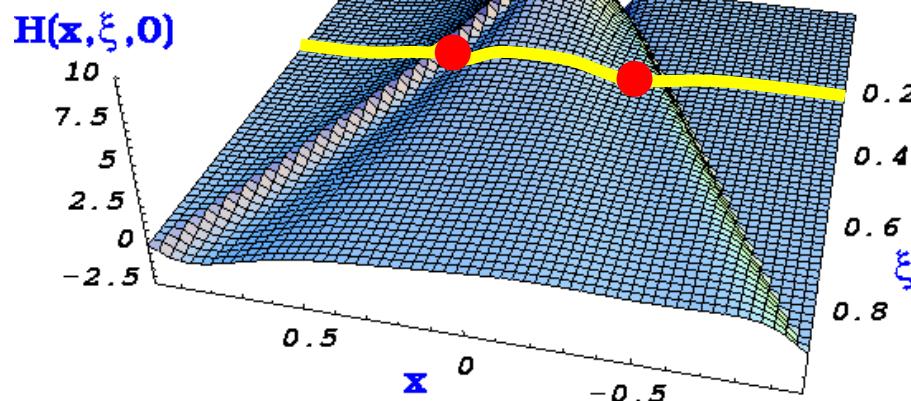
- $H(x, \xi, t)$ but only ξ and t accessible experimentally
- x is mute variable (integrated over):
 - apart from cross-over trajectory ($\xi=x$) GPDs not directly accessible: deconvolution needed ! (model dependent)
 - GPD moments cannot be directly revealed, extrapolations $t \rightarrow 0$ are model dependent

e.g.

$$T^{DVCS} \sim \int_{-1}^{+1} \frac{H(x, \xi, t)}{x \pm \xi + i\epsilon} dx + \dots$$

$$\sim P \int_{-1}^{+1} \frac{H(x, \xi, t)}{x \pm \xi} dx - i\pi H(\pm \xi, \xi, t) + \dots$$

$$\quad$$



cross sections & beam-charge asymmetry $\sim \text{Re}(T^{DVCS})$

beam or target-spin asymmetries $\sim \text{Im}(T^{DVCS})$

the ideal experiment for measuring hard exclusive processes

the ideal experiment for measuring hard exclusive processes

- high+variable beam energy
 - hard regime
 - wide kinematic range
 - high luminosity
 - small cross sections
 - measure in 3 kinematic variables simultaneously
 - complete event reconstruction
 - ensure exclusivity
- ... doesn't exist (yet)...

the menu



data from exclusive VM over wide kinematic range

JLab → HERMES → COMPASS → HERA-collider

→ role of quarks and gluons

→ NLO corrections



exclusive PS mesons production

→ role of power corrections



DVCS: from first signals → detailed measurements

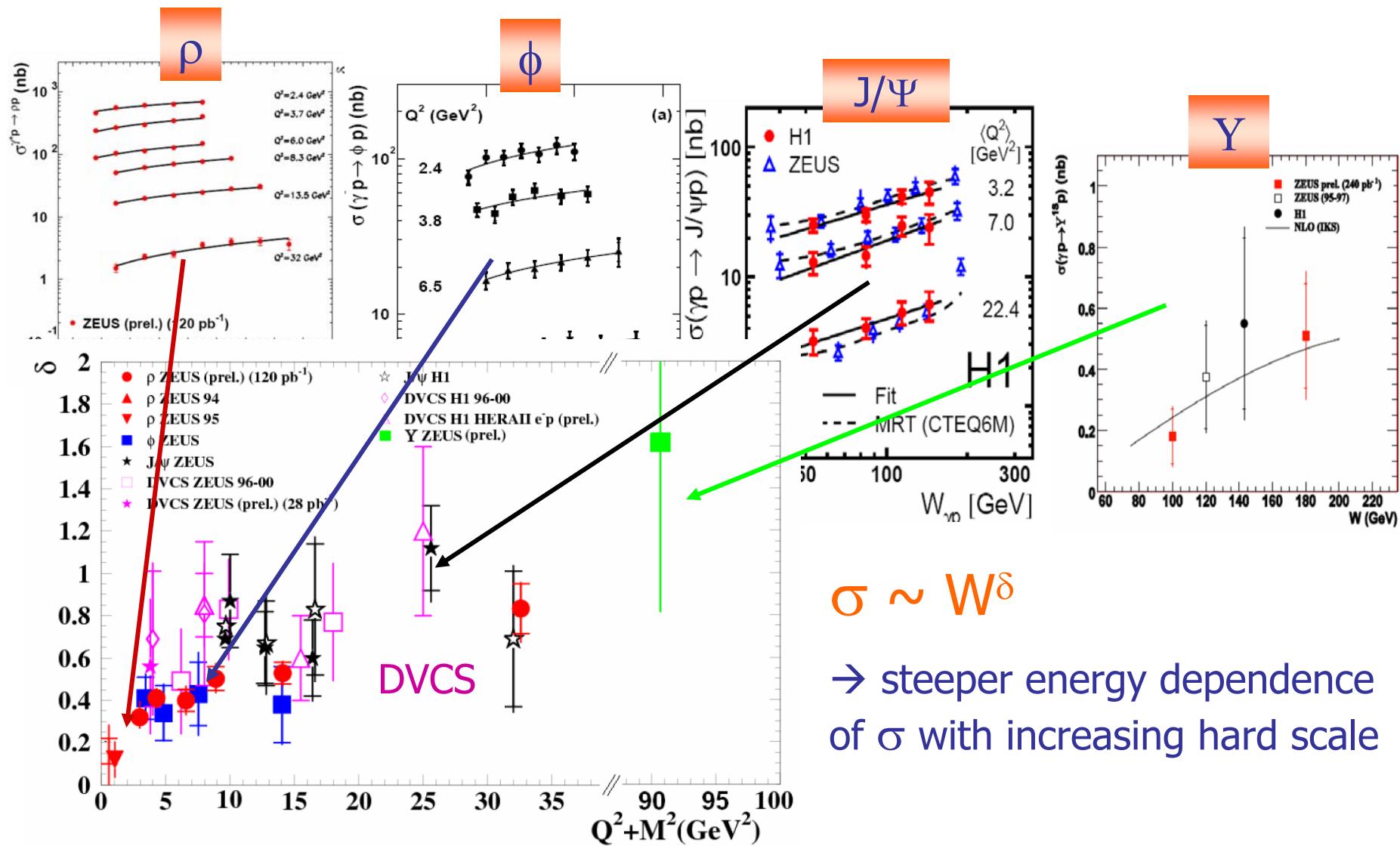
reminder: for meson production factorisation only for σ_L (σ_T suppressed by $1/Q^2$)



VM production @small x



W & t dependences: probe transition from soft → hard regime

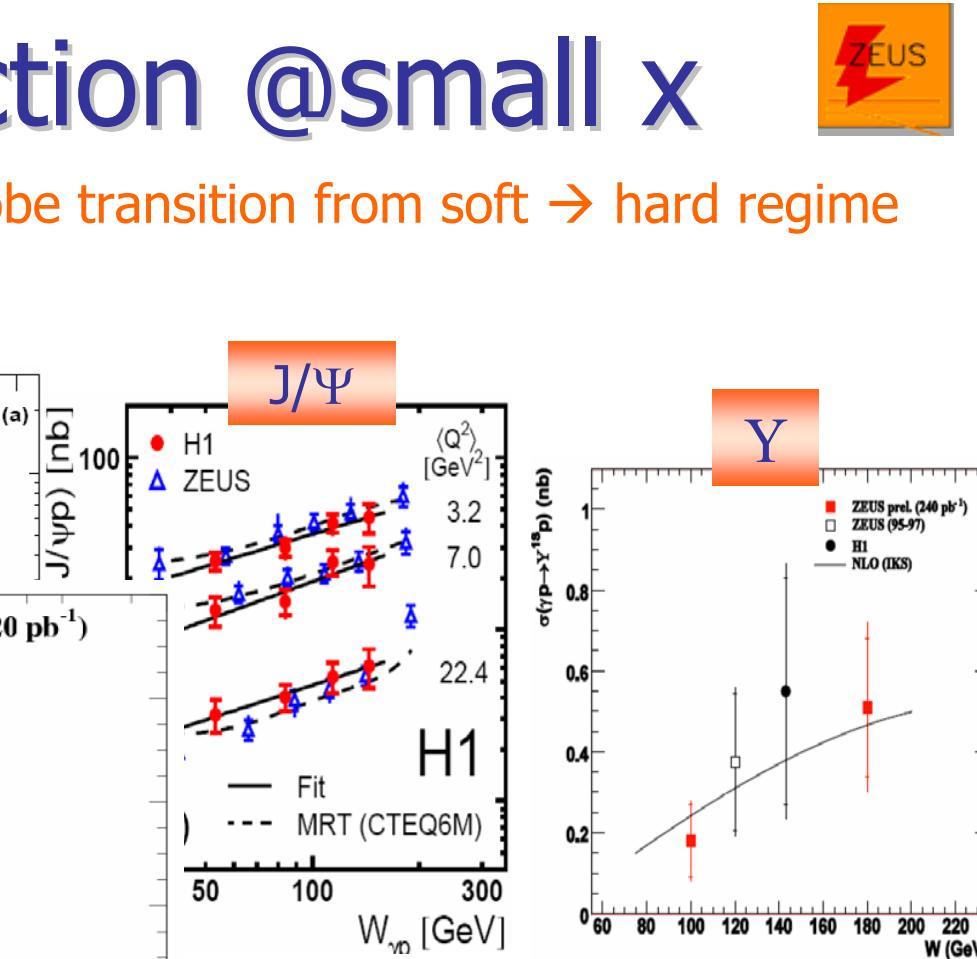
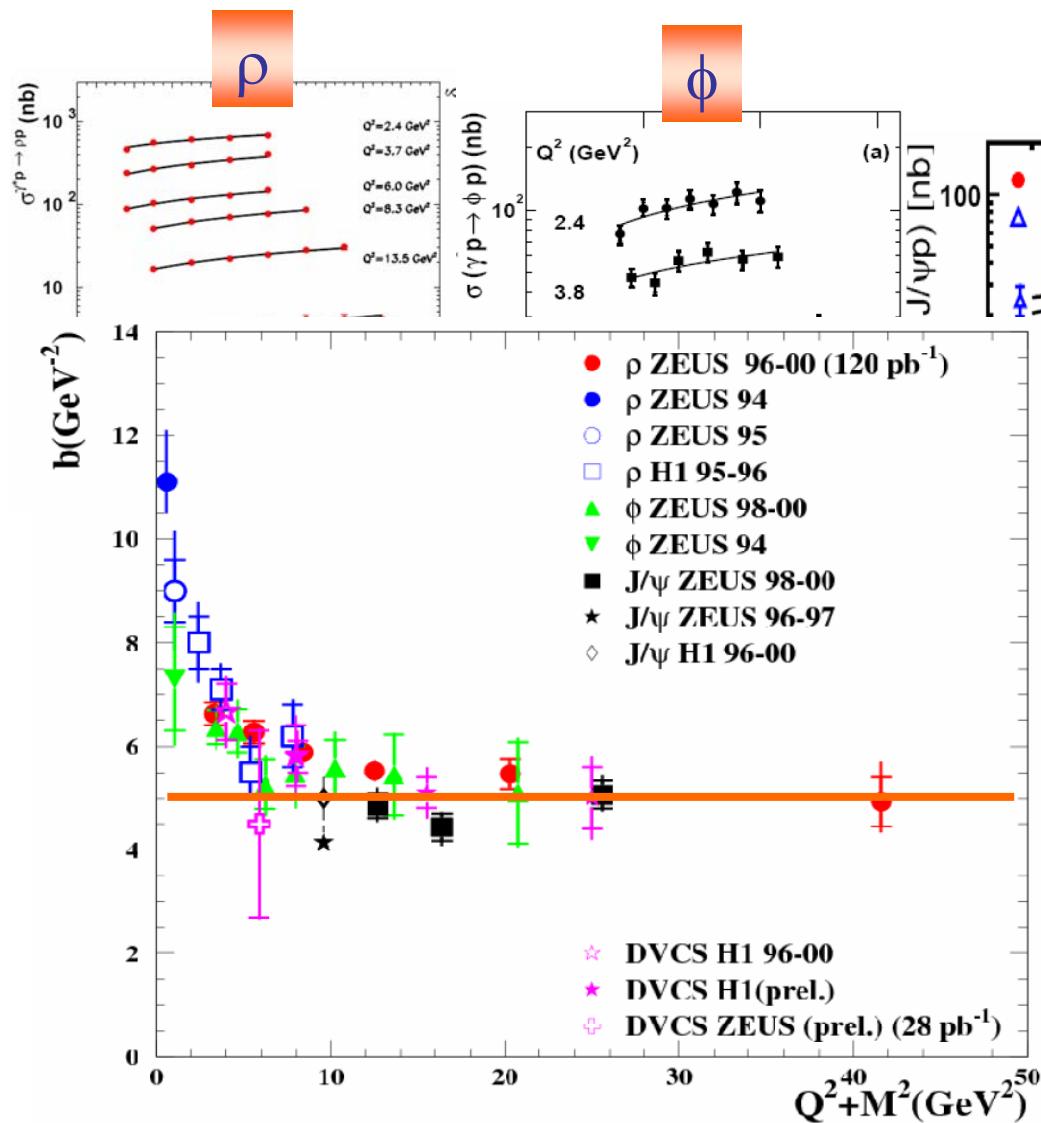




VM production @small x



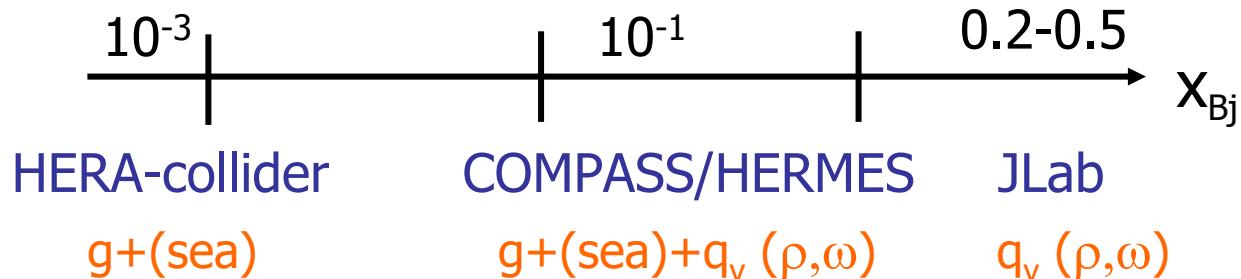
W & t dependences: probe transition from soft → hard regime



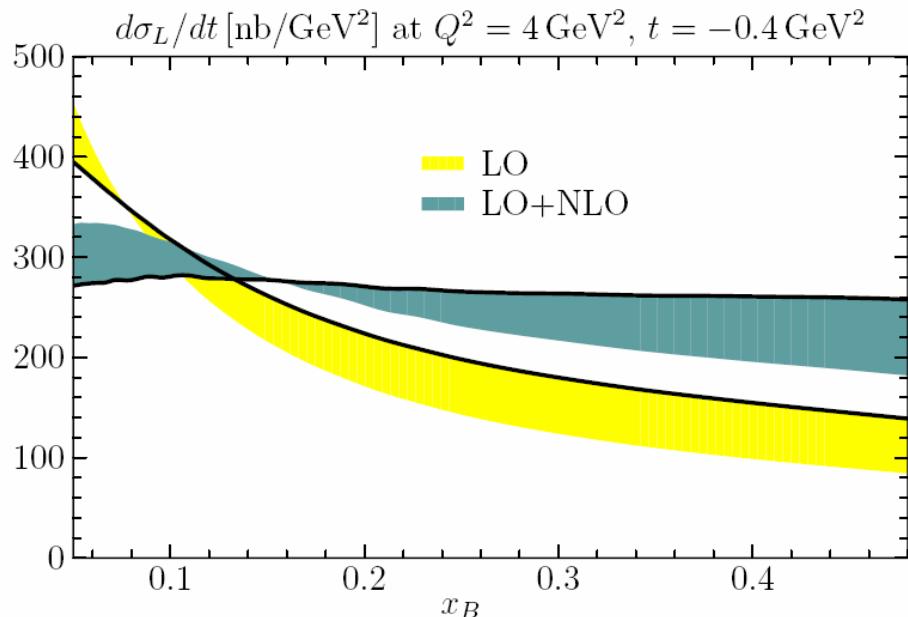
$$\sigma \sim e^{-b|t|}$$

→ universality of b-slope parameter: point-like configurations dominate

VM production: small \rightarrow high x

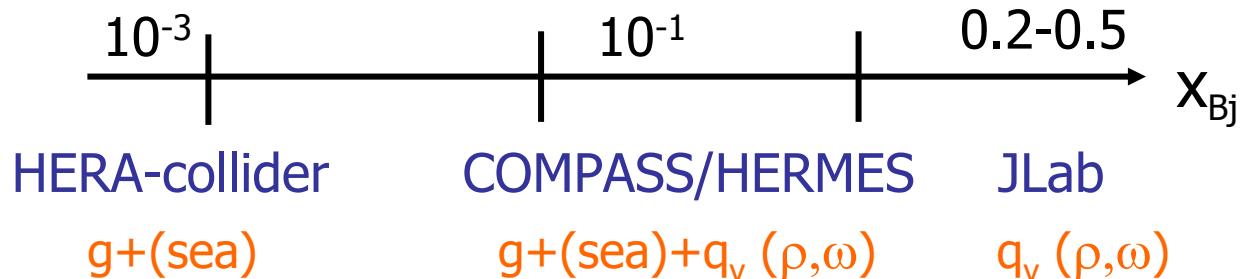


- NLO corrections to VM production are large: [M.Diehl, W.Kugler arXiv0708.1121]

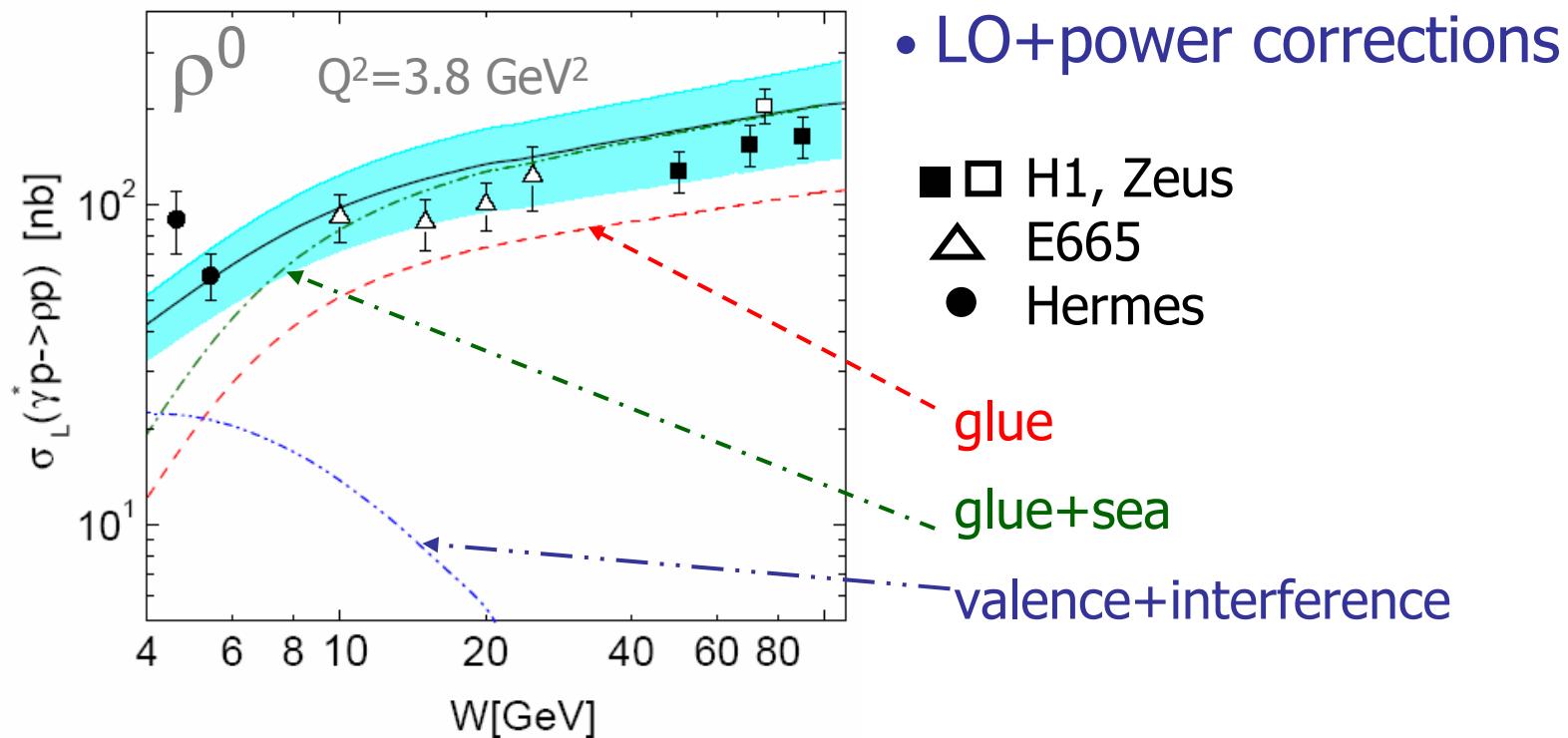


- ρ^0 cross section @typical kinematics of compass / hermes / jlab12

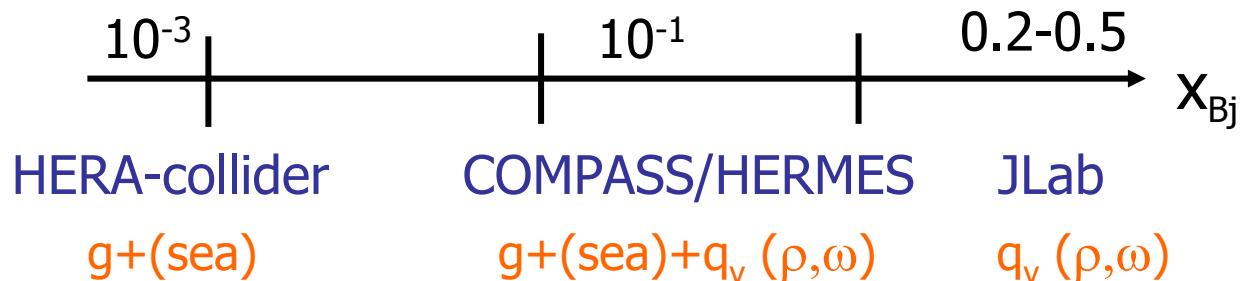
VM production: small \rightarrow high x



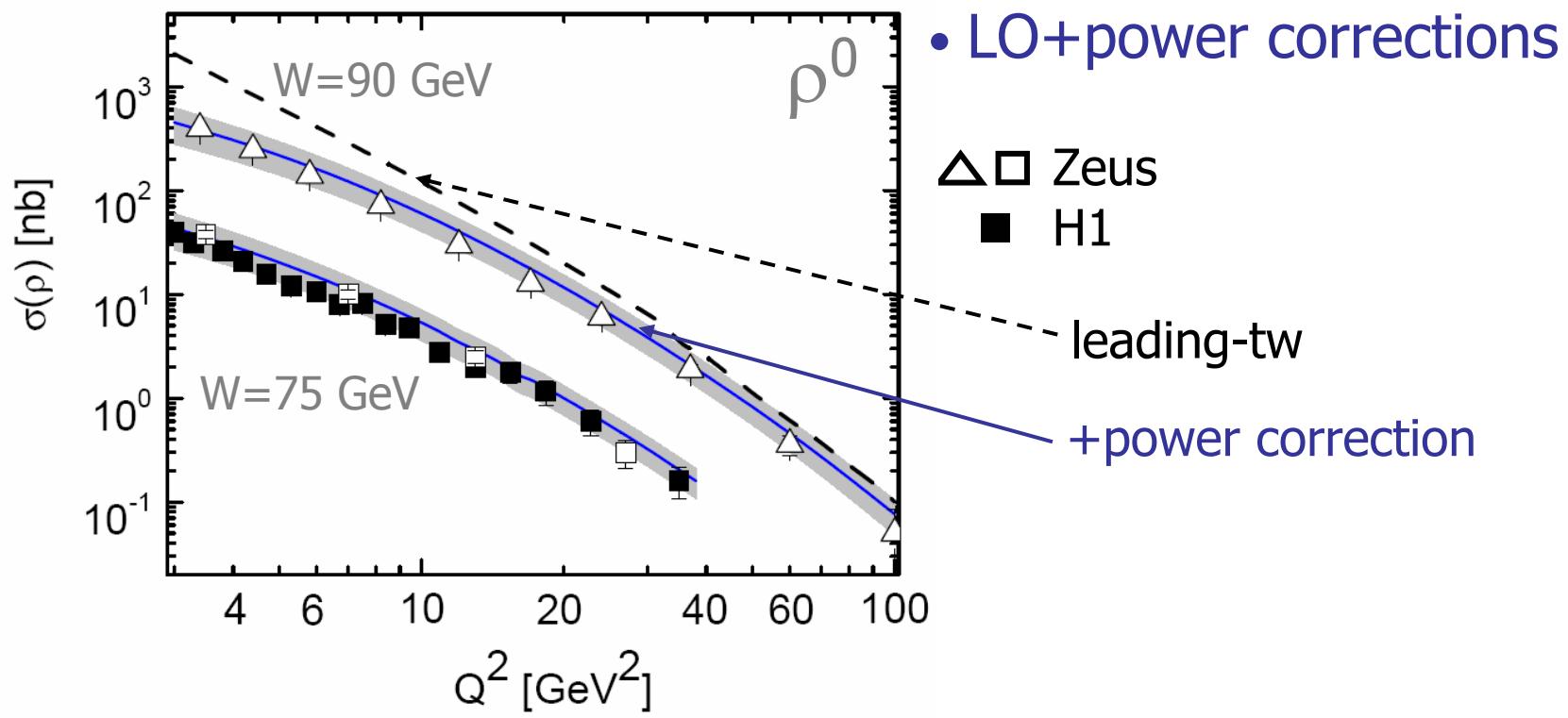
- ...despite: LO GPD model (handbag fact.) [S.Goloskokov, P.Kroll arXiv0711.4736]



VM production: small \rightarrow high x



- ...despite: LO GPD model (handbag fact.) [S.Goloskokov, P.Kroll arXiv0711.4736]

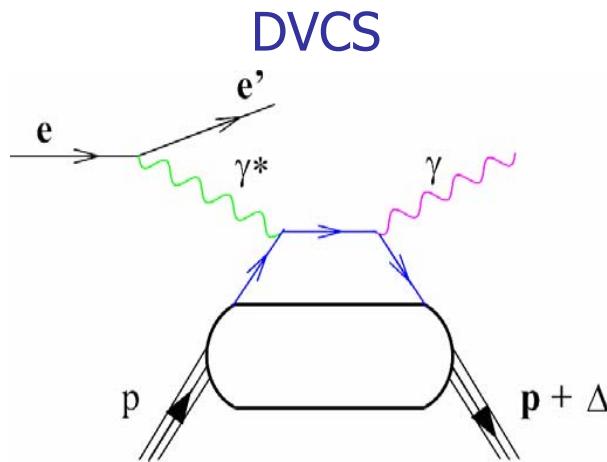


deeply virtual compton scattering

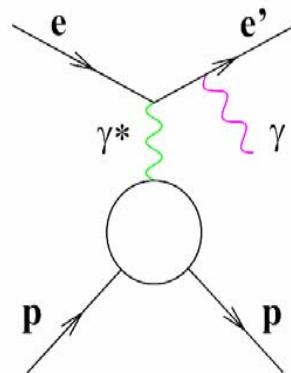
DVCS

→ $H, \tilde{H}, E, \tilde{E}$

most clean channel for interpretation in terms of GPDs (full factorisation proof)



Bethe-Heitler



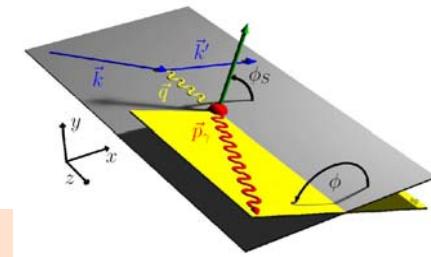
@HERMES/JLab:

DVCS << Bethe-Heitler

$$\frac{d^4\sigma}{dx_B dQ^2 d|t| d\phi} \propto |T_{\text{DVCS}} + T_{\text{BH}}|^2 = |T_{\text{DVCS}}|^2 + |T_{\text{BH}}|^2 + \underbrace{T_{\text{DVCS}} T_{\text{BH}}^* + T_{\text{DVCS}}^* T_{\text{BH}}}_{I}$$

→ leads to non-zero azimuthal asymmetries:

DVCS @amplitude level



$$d\sigma \propto |\tau_{\text{BH}}|^2 + |\tau_{\text{DVCS}}|^2 + (\tau_{\text{BH}}^* \tau_{\text{DVCS}} + \tau_{\text{DVCS}}^* \tau_{\text{BH}})$$

$$I \sim \Delta\sigma$$

→ different charges: $e^+ e^-$ (*only @HERA!*):

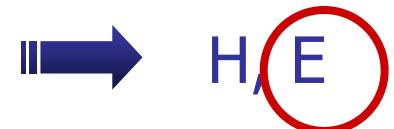
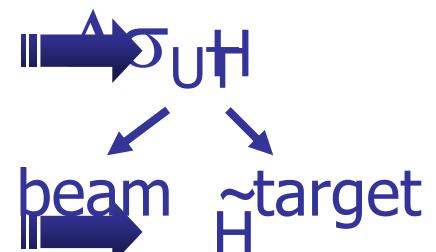
$$\Delta\sigma_C \sim \cos\phi \cdot \text{Re}\{ H + \xi \tilde{H} + \dots \} \quad \Rightarrow H$$

→ polarisation observables:

$$\Delta\sigma_{LU} \sim \sin\phi \cdot \text{Im}\{ H + \xi \tilde{H} + kE \}$$

$$\Delta\sigma_{UL} \sim \sin\phi \cdot \text{Im}\{ \tilde{H} + \xi H + \dots \}$$

$$\Delta\sigma_{UT} \sim \sin(\phi - \phi_s) \cos\phi \cdot \text{Im}\{ k(H - E) + \dots \}$$



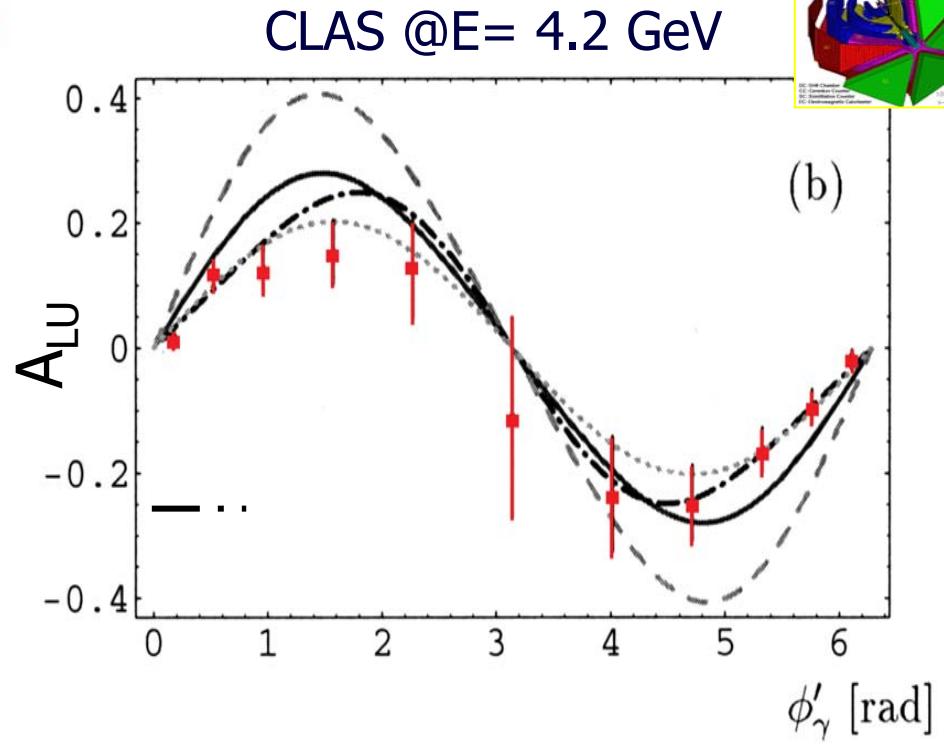
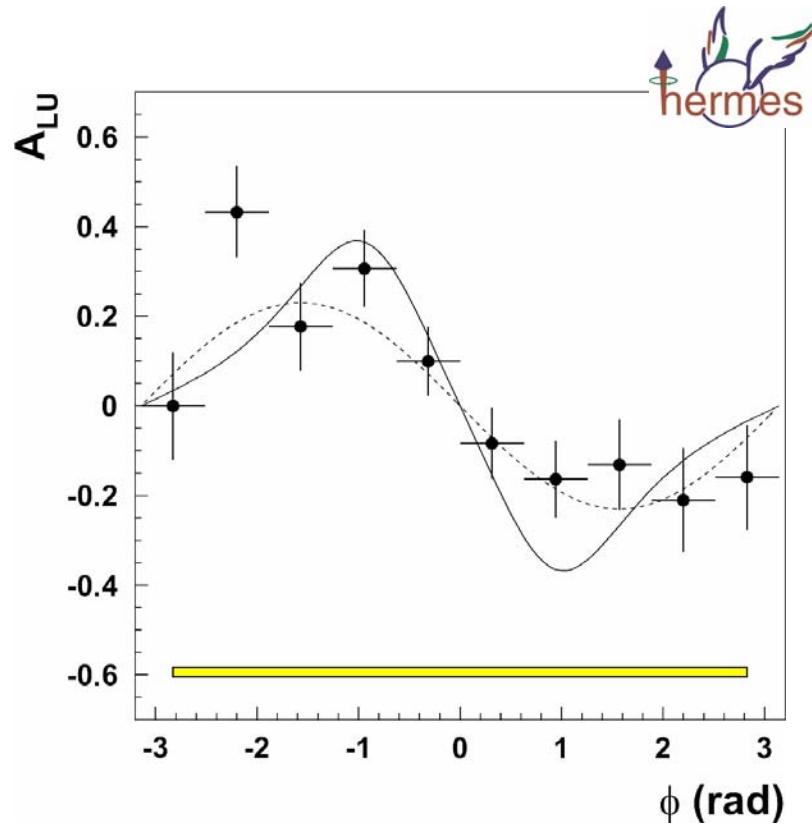
$$\xi = x_B/(2-x_B), k = t/4M^2$$

kinematically suppressed
@HERMES and JLab energies

first DVCS signals: A_{LU}

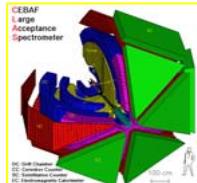
-- from interference term --

[PRL87(2001)]

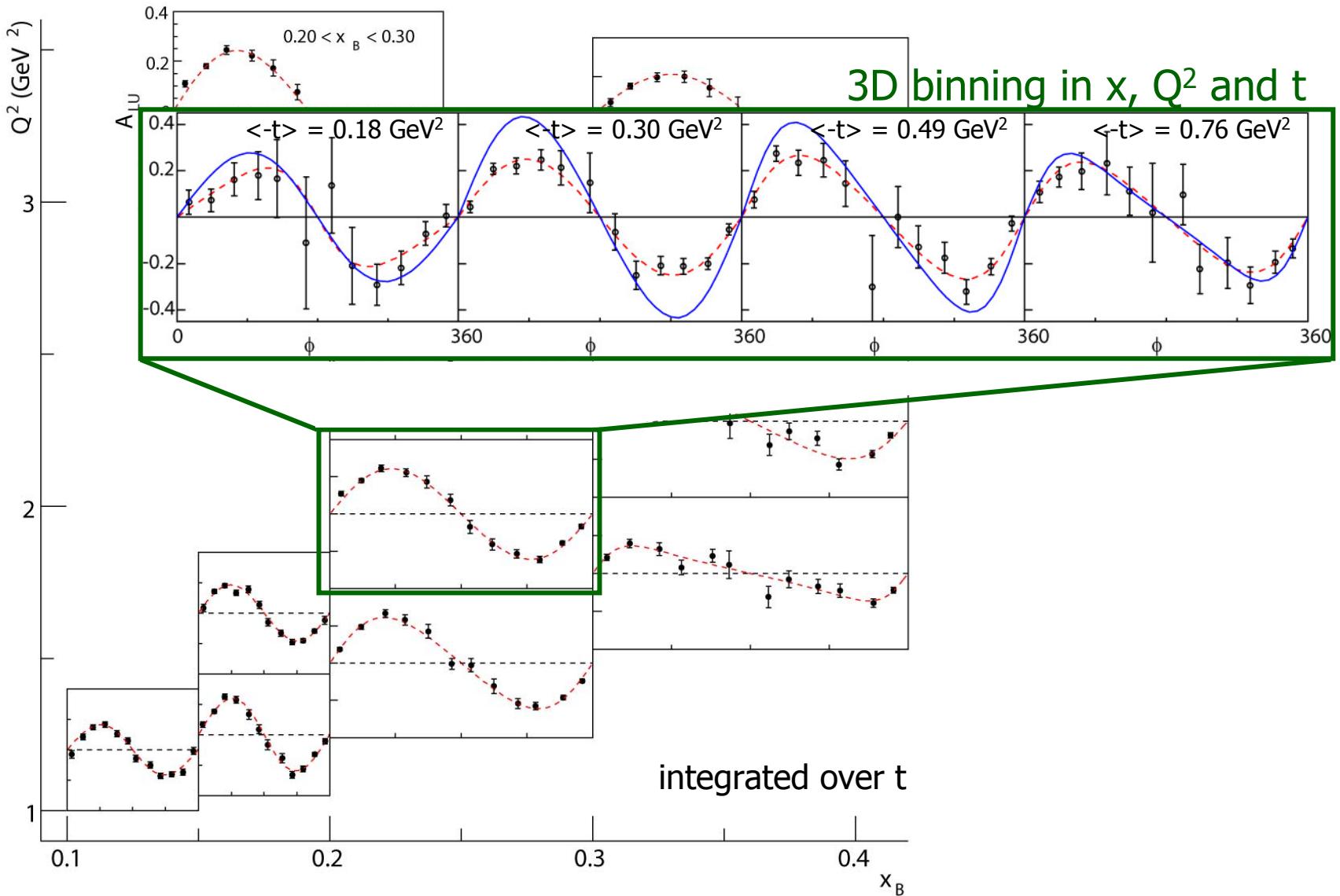


→ $\sin\phi$ dependence indicates dominance of handbag contribution

call for high statistics



JLab: E1-DVCS beam-spin asymmetry

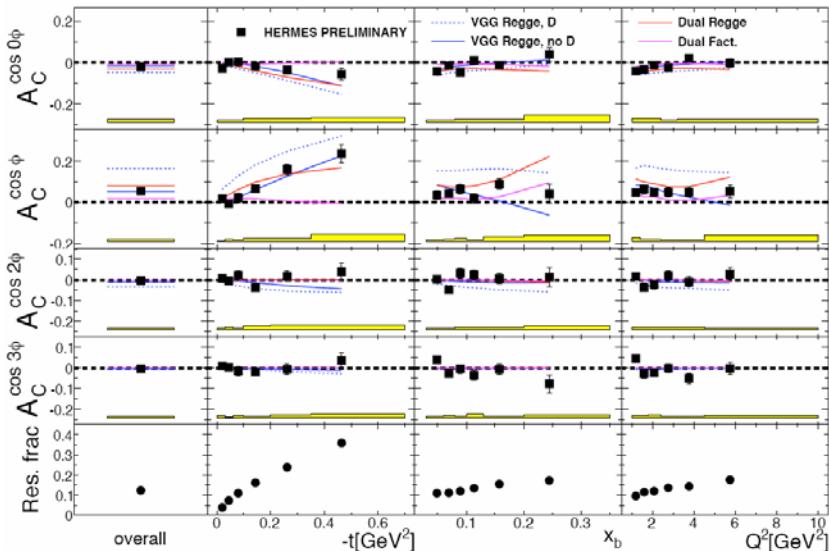


call for new analysis methods

HERMES: combined analysis of charge & polarisation dependent data
 → separation of interference term + DVCS²

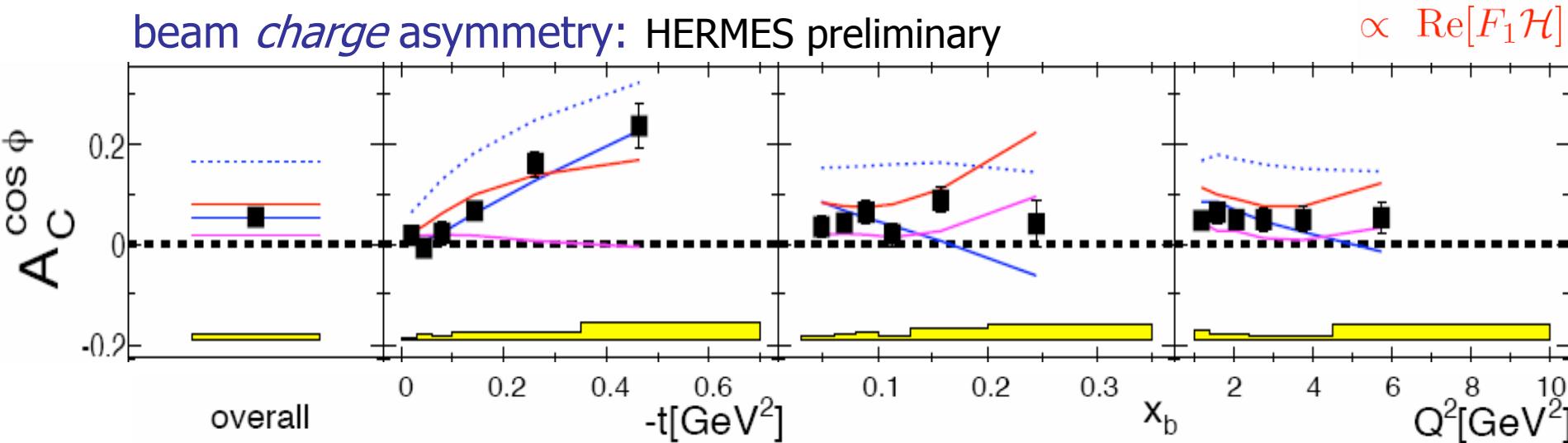
$$\sigma_{LU}(\phi; P_l, e_l) = \sigma_{UU}(\phi) \cdot \{ 1 + P_l A_{LU}^{\text{DVCS}}(\phi) + e_l P_l A_{LU}^{\mathcal{I}}(\phi) + e_l A_C(\phi) \}$$

$\sum_{n=1}^2 s_n^I \sin(n\phi)$ $\sum_{n=0}^3 c_n^I \cos(n\phi)$



call for new analysis methods

HERMES: combined analysis of charge & polarisation dependent data
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GPD models: VGG  w/o D-term  regge-ansatz for t-depend.

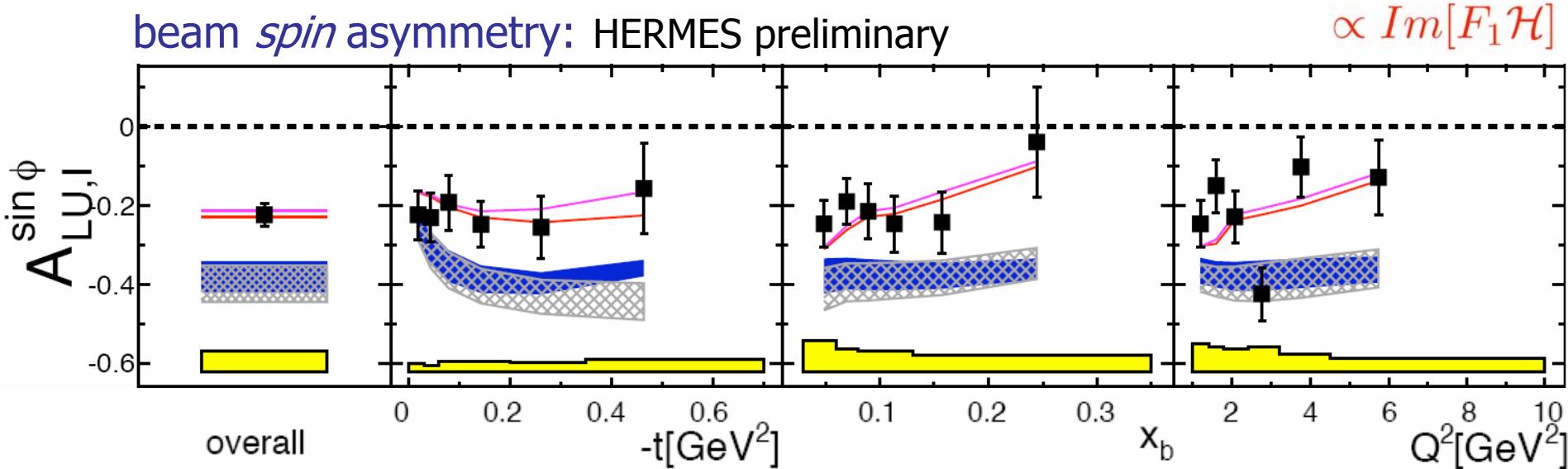
dual  regge-ansatz for t-dependence
 [Guzey, Teckentrup 2006]  factorised t-dependence

call for new analysis methods



HERMES: combined analysis of charge & polarisation dependent data

→ separation of interference term + DVCS²



GPD models: VGG

regge-ansatz for t-dependence

factorised t-dependence

dual

1

regge-ansatz for t-dependence

factorised t-dependence

a word about GPD models

VGG: [Vanderhaegen, Guichon, Guidal 1999]

- double distributions ; factorised or regge-inspired t-dependence
- D-term to restore full polynomiality
- skweness depending on free parameters b_{val} & b_{sea}
- includes tw-3 (WW approx)

dual: [Guzey, Teckentrup 2006]

- GPDs based on infinite sum of t channel resonances
- factorised or regge-inspired t-dependence
- tw-2 only

a word about GPD models

VGG: [Vanderhaegen, Guichon, Guidal 1999]

- double distributions ; fact → describes well A_C and A_{UT} data
- D-term to restore full pol → fails for A_{LU}
- skweness depending on → A_c favour 'no D-term' ← contradicts
- includes tw-3 (WW approc) χ QSM & lattice results

dual: [Guzey, Teckentrup 2006]

- GPDs based on infinite series
 - factorised or regge-inspired
 - tw-2 only
- describes well spin-asymmetries
→ fails for unpol. cross sections (HallA)

→ call for new, more sophisticated parametrisations of GPDs

... more models on the way: e.g. generalisation of Mellin transform technique

...nevertheless: first attempts to constrain J_q

observables sensitive to E:
(J_q input parameter in ansatz for E)

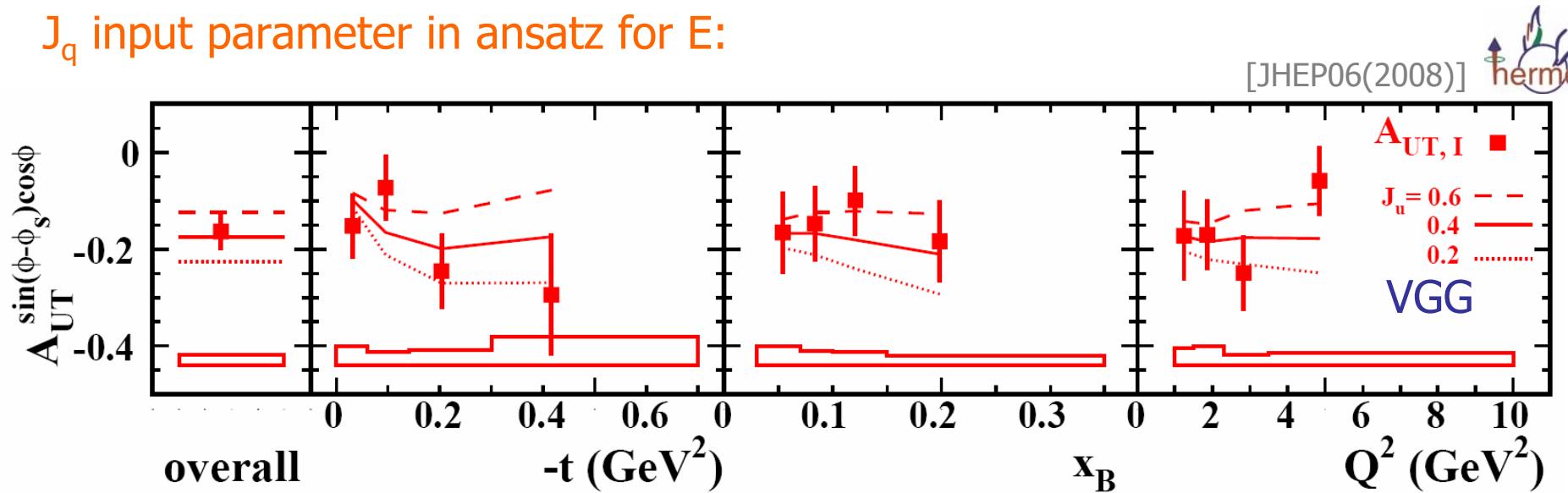
$$J_q = \frac{1}{2} \int_{-1}^1 x dx \left(H^q + E^q \right)$$

- DVCS A_{UT} : HERMES
- nDVCS A_{LU} : Hall A
- ρ^0 A_{UT} : HERMES

...nevertheless: first attempts to constrain J_q

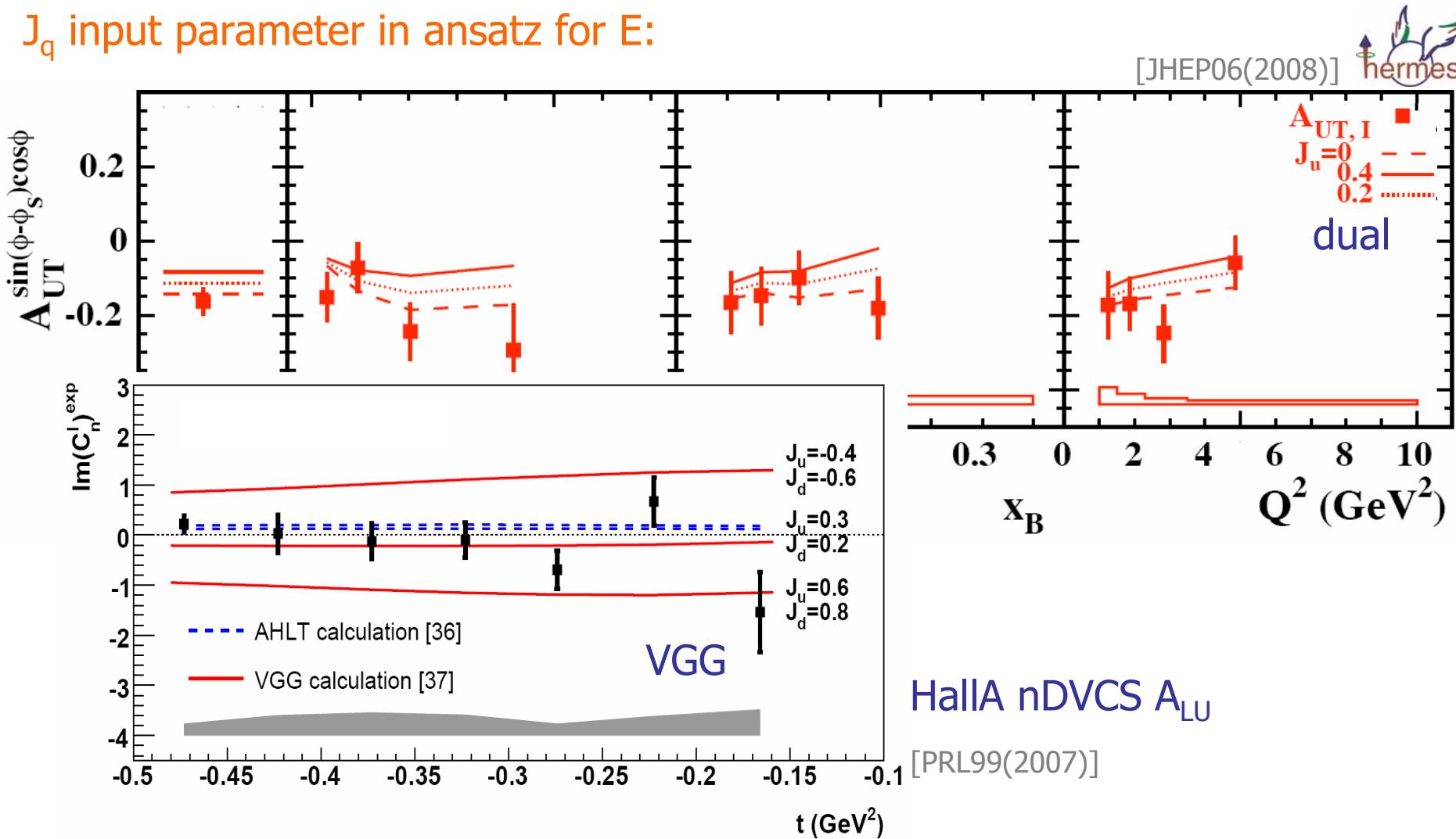
J_q input parameter in ansatz for E:

[JHEP06(2008)]



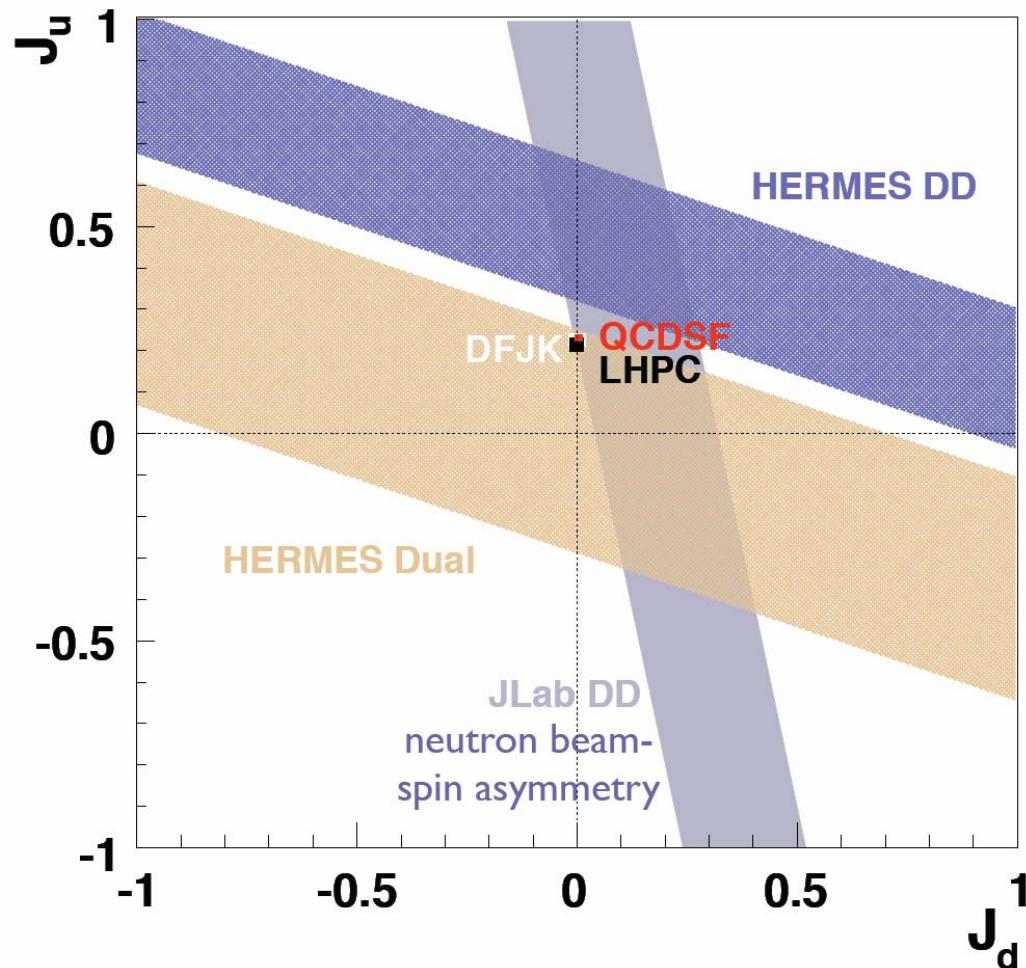
...nevertheless: first attempts to constrain J_q

J_q input parameter in ansatz for E:



...nevertheless: first attempts to constrain J_q

J_q input parameter in ansatz for E:



→ demonstrates model dependence of these analyses

→ data are free to be re-used at any time with new models ☺

conclusions

GPDs

contain a wealth of new information on hadron structure at parton level
→ only known framework allowing a 3D imaging of hadrons ←

... BUT they are intricate functions...

complementary to TMDs : relations GPDs \leftrightarrow TMDs [M. Burkardt, M. Schlegel]



GPDs offer a way to measure transversity!

→ see talk by
G. Goldstein

- increasing amount and precision of experimental data
- large “flow” of new data expected soon (JLab, HERMES, COMPASS)
- ‘standard’ models/parametrisations of GPDs too simple
 - models should describe large variety of different observables over wide kinematic range



prior to any conclusion about GPDs from data: call for new,
more sophisticated parametrisations

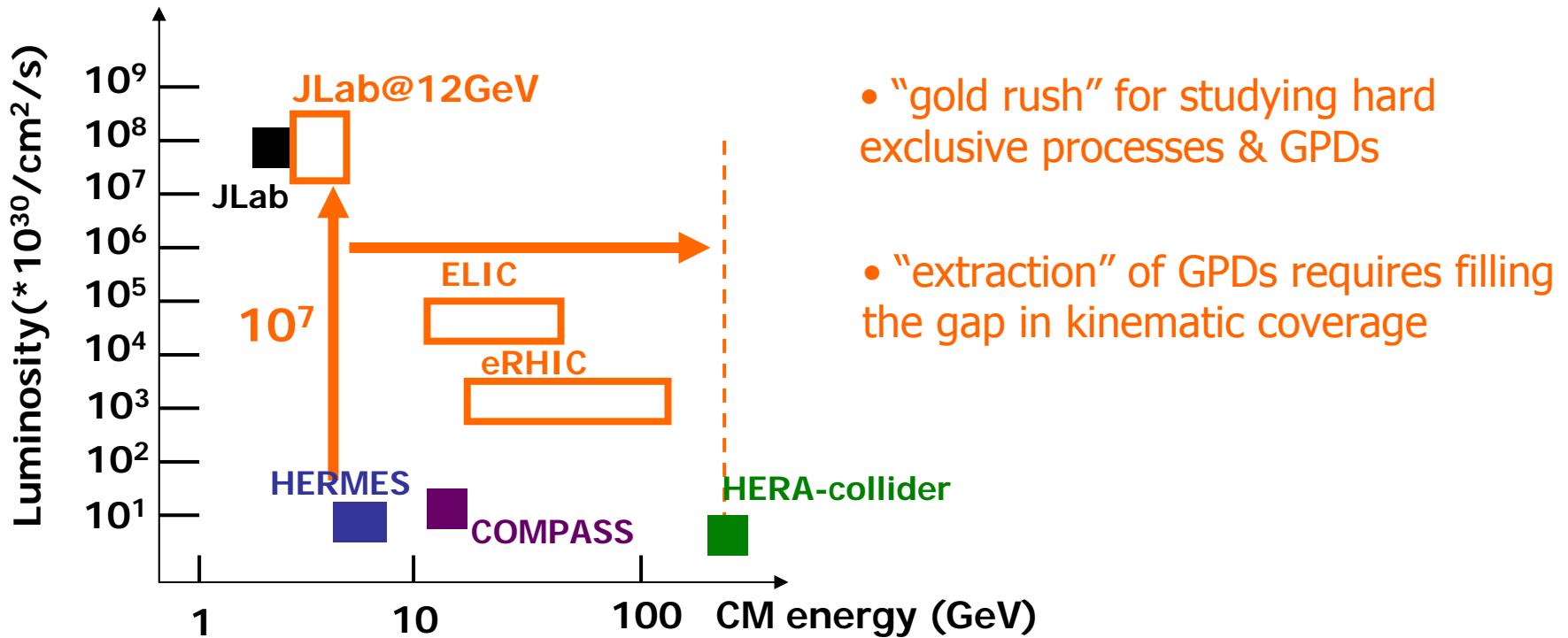
perspectives for GPDs

@ new facilities:

- high beam energy (hard regime, wide kinematic range)
- very high luminosity (small xsections, multi-D analyses)
- complete event reconstruction (ensure exclusivity)

→ exploration of new channels: WACS, time like DVCS, ...

→ ideas for accessing GPDs @LHC, @GSI, ...



perspectives for GPDs & TMDs

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