

Deeply Virtual Compton Scattering

on the Deuteron and Heavier Nuclei at HERMES

Caroline Riedl



19th International Spin Physics Symposium
Jülich, Germany, October 1, 2010

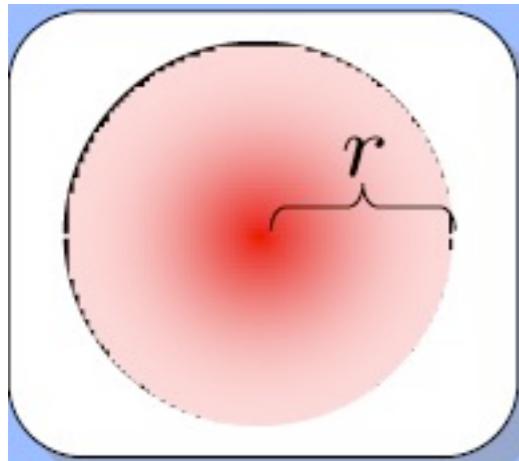


Outline: DVCS on nuclei

- Theoretical motivation
- HERMES experiment at HERA/DESY
- Azimuthal asymmetries in DVCS and GPDs
- Results on the deuteron and heavier nuclei
- Summary

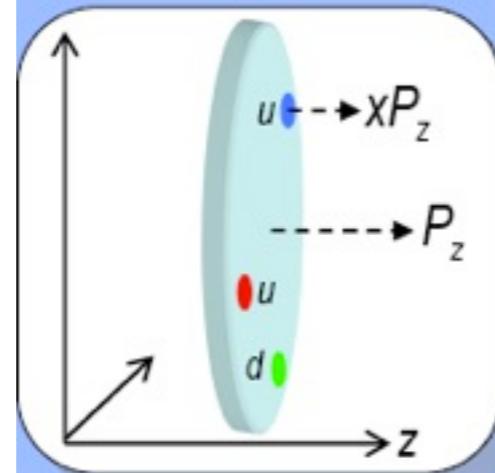
Generalized Parton Distributions

Elastic Form Factors



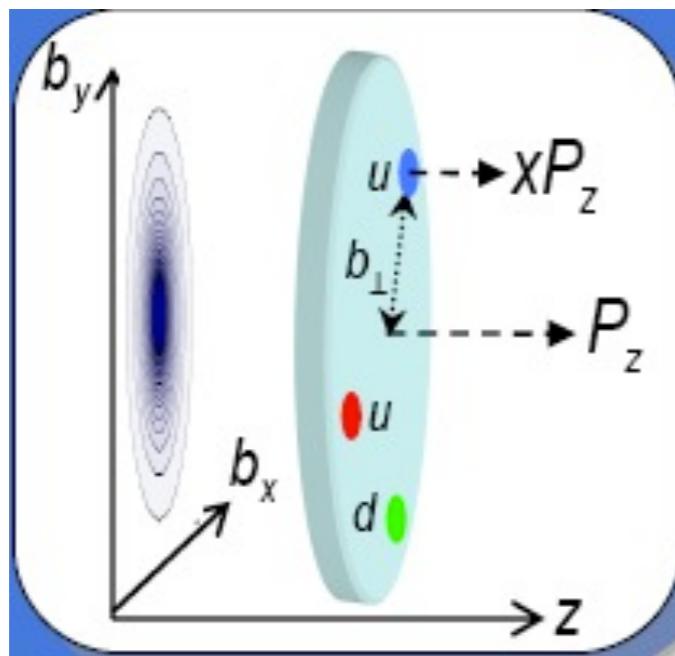
transverse position
of partons

Parton Distribution Functions (PDFs)



longitudinal momentum
of partons

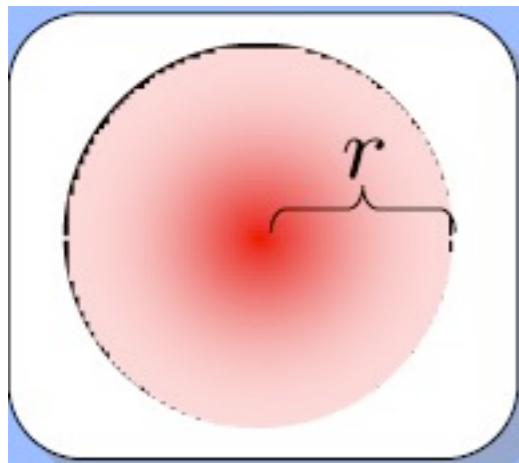
Generalized Parton Distributions



Nucleon
Tomography

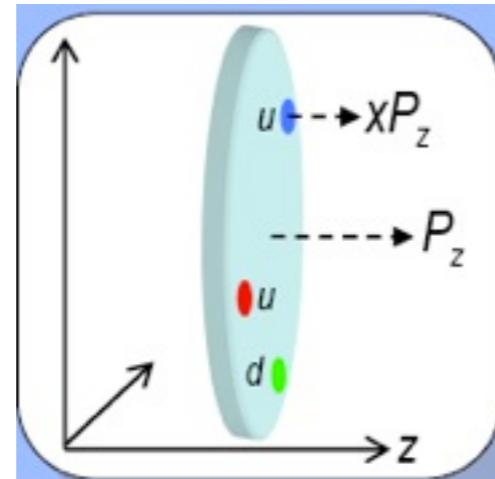
correlation between longitudinal momentum and transverse position

Elastic Form Factors



transverse position
of partons

Parton Distribution Functions (PDFs)

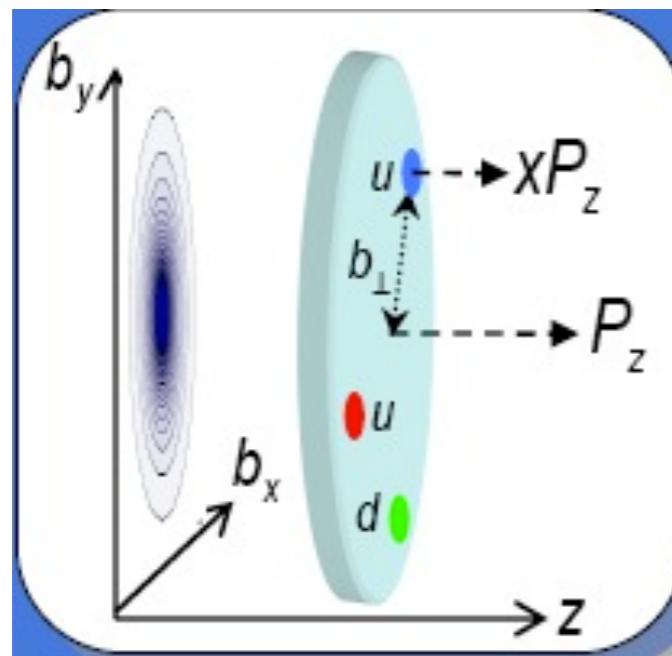


longitudinal momentum
of partons

Generalized Parton Distributions

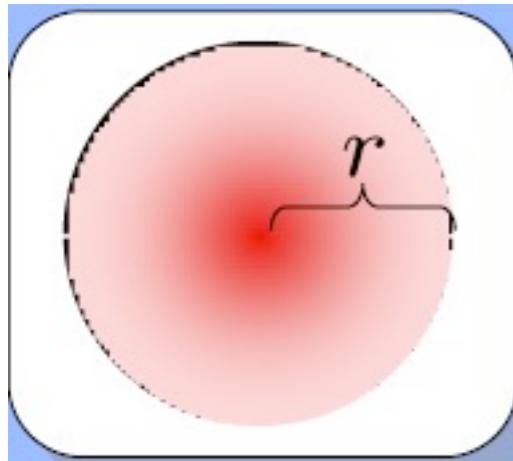
Vector Mesons: see talk by E. Avetisyan

Nucleon Tomography



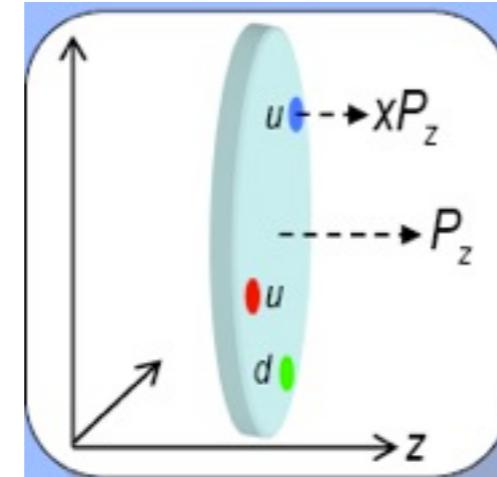
correlation between longitudinal momentum and transverse position

Elastic Form Factors



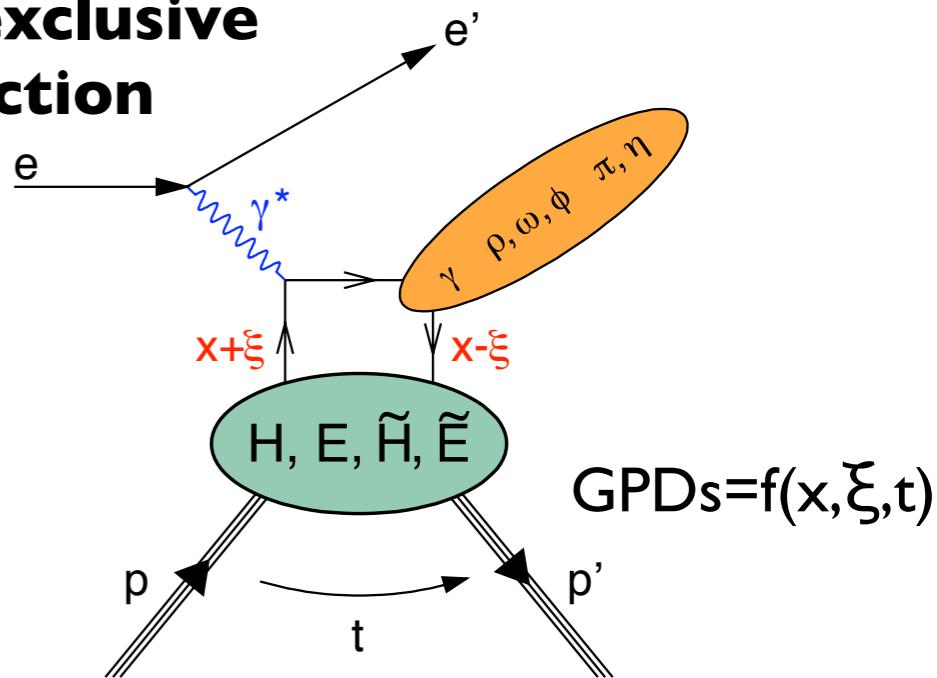
transverse position of partons

Parton Distribution Functions (PDFs)



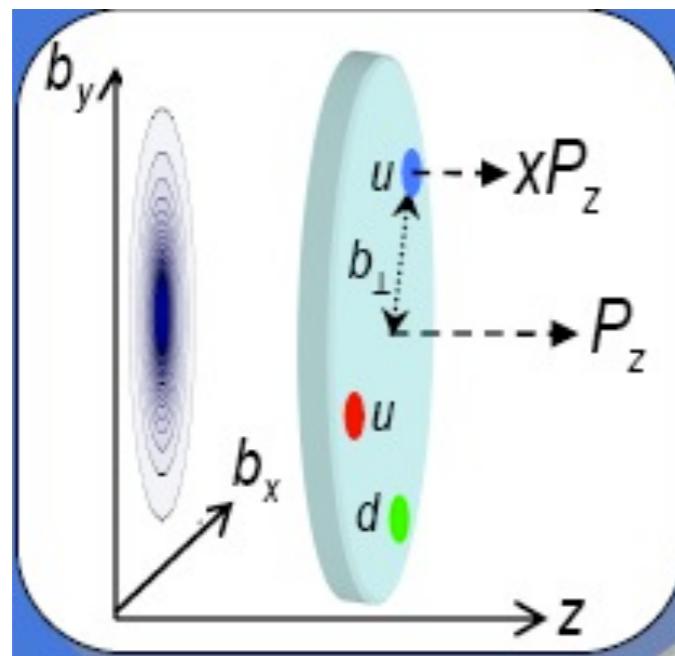
longitudinal momentum of partons

Hard exclusive reaction



Generalized Parton Distributions

Vector Mesons: see talk by E. Avetisyan

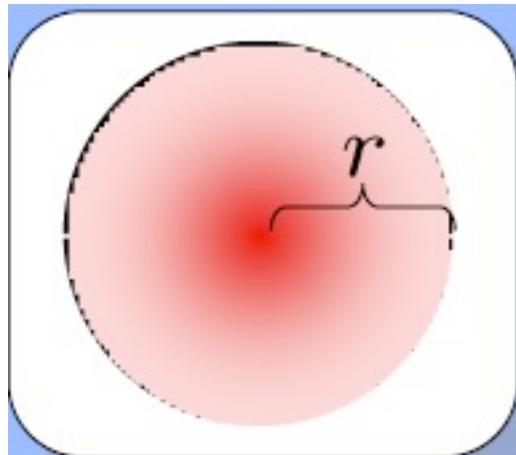


Nucleon Tomography

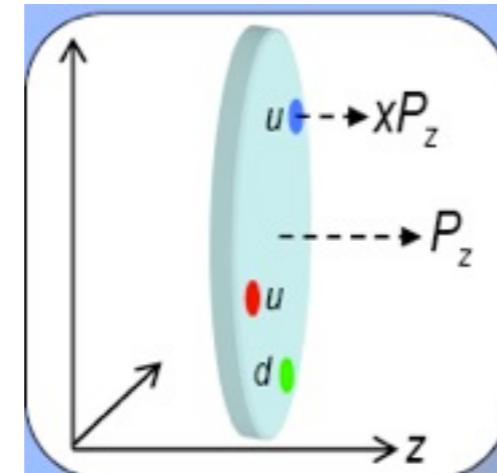
correlation between longitudinal momentum and transverse position

Elastic Form Factors

Parton Distribution Functions (PDFs)

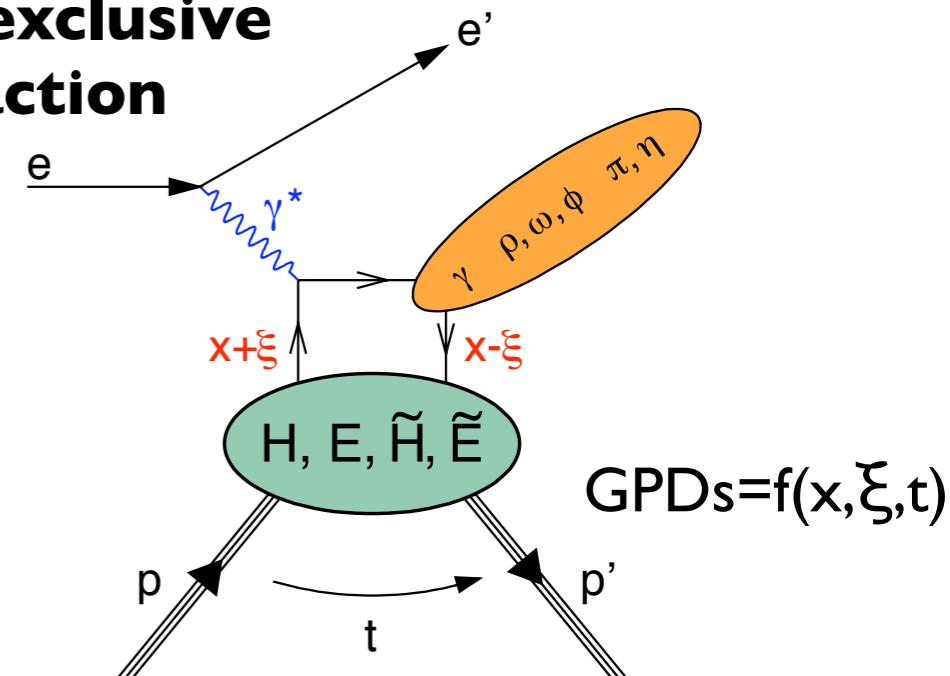


transverse position of partons



longitudinal momentum of partons

Hard exclusive reaction



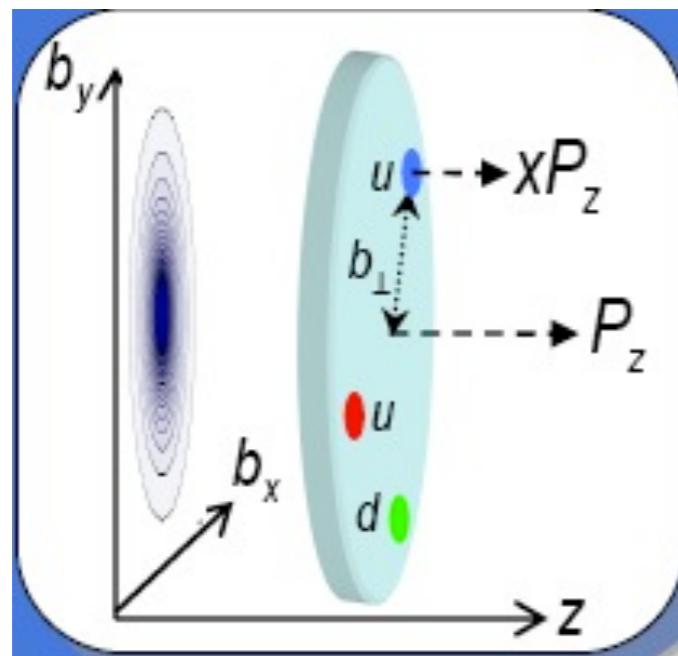
Ji sum rule for the nucleon

-Ji, PRL 78 (1997) 610-

$$J_q = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H^q(x, \xi, t) + E^q(x, \xi, t)]$$

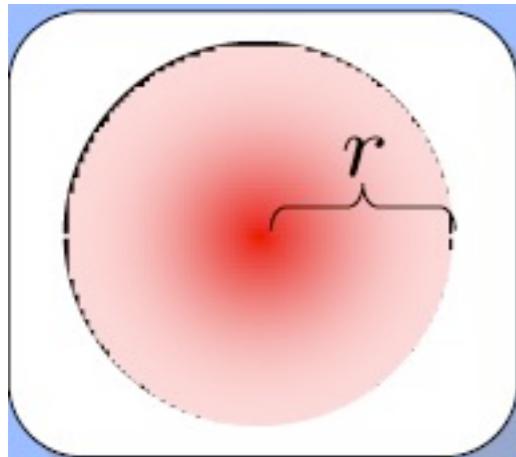
$$\frac{1}{2} = \frac{1}{2} \Delta \Sigma + L_q + J_g$$

Generalized Parton Distributions



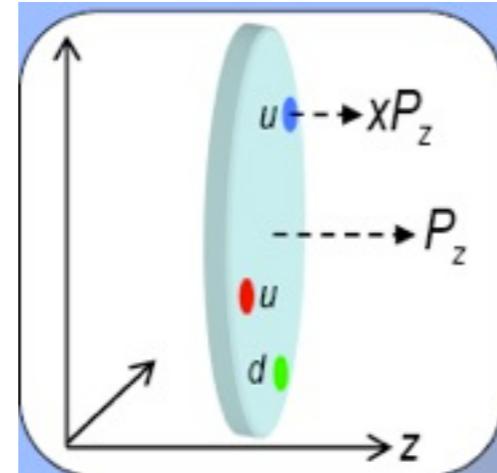
Nucleon Tomography

Elastic Form Factors



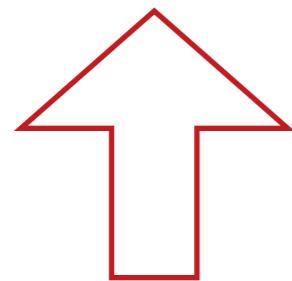
transverse position
of partons

Parton Distribution Functions (PDFs)

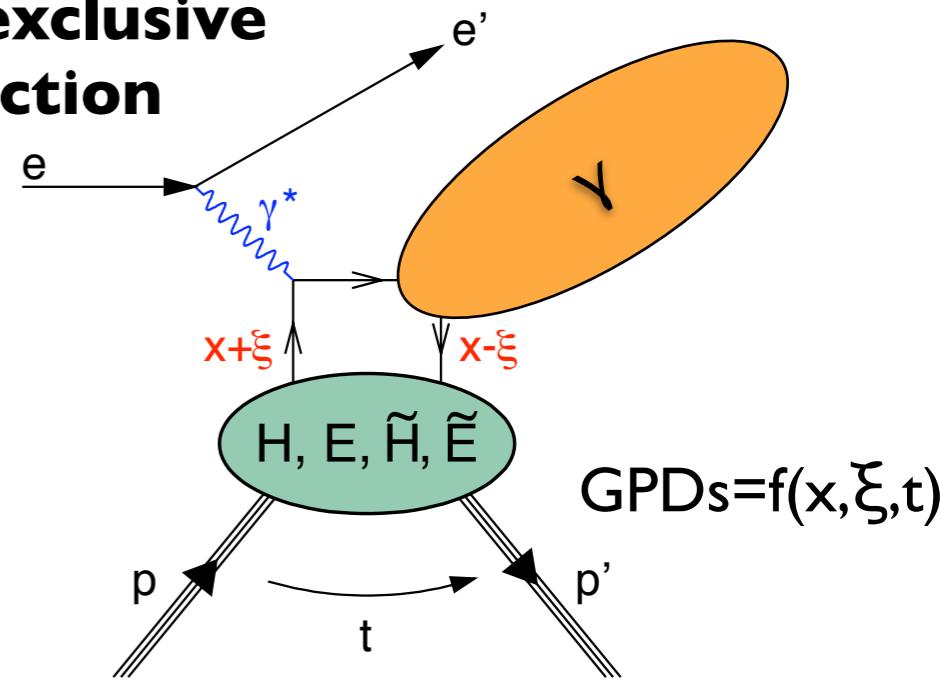


longitudinal momentum
of partons

correlation between longitudinal momentum and transverse position



Hard exclusive reaction



Ji sum rule for the nucleon

-Ji, PRL 78 (1997) 610-

$$J_q = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H^q(x, \xi, t) + E^q(x, \xi, t)]$$

$$\frac{1}{2} = \frac{1}{2} \Delta \Sigma + L_q + J_g$$

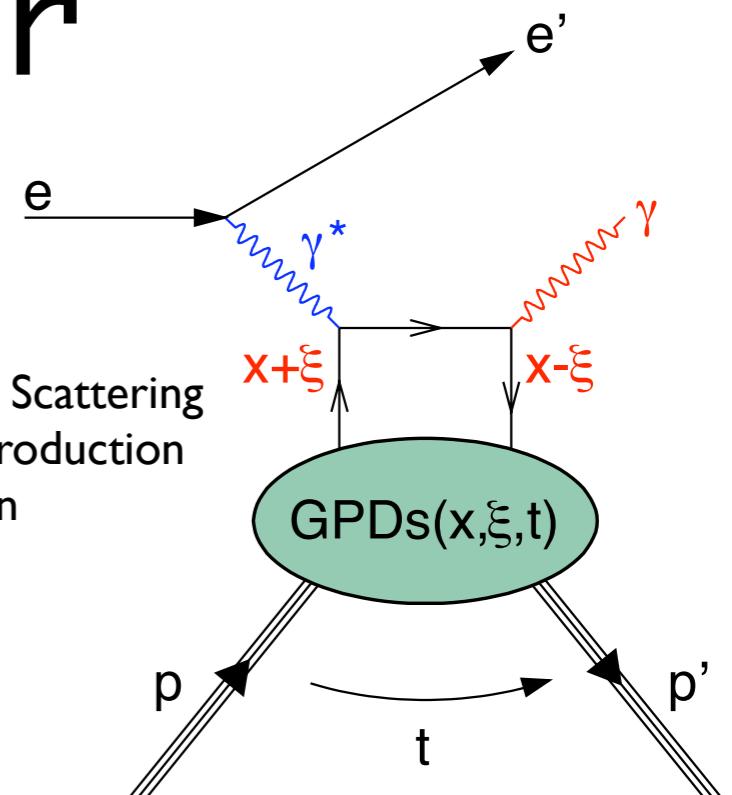
Vector Mesons: see talk by E. Avetisyan

DVCS on hadrons other than the nucleon

I.

Spin-1/2	flips nucleon helicity	conserves nucleon helicity
does not depend on quark helicity	E	H
depends on quark helicity	\tilde{E}	\tilde{H}

Deeply Virtual Compton Scattering
= hard exclusive lepto-production
of a real photon



4 chiral-even quark GPDs at leading twist

(4 more GPDs for chiral-odd case
related to transversity)

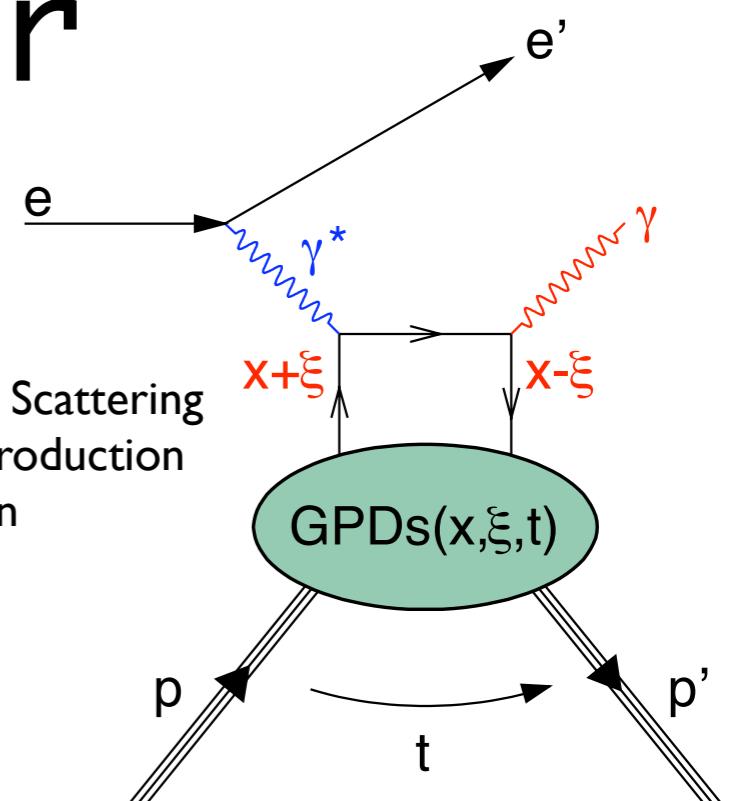
DVCS on hadrons other than the nucleon

I.

Spin-1/2	flips nucleon helicity	conserves nucleon helicity
does not depend on quark helicity	E	H
depends on quark helicity	\tilde{E}	\tilde{H}

$q(x)$ $\Delta q(x)$
 forward limit
 $\xi \rightarrow 0, t \rightarrow 0$

Deeply Virtual Compton Scattering
= hard exclusive lepto-production
of a real photon



4 chiral-even quark GPDs at leading twist

(4 more GPDs for chiral-odd case
related to transversity)

DVCS on hadrons other than the nucleon

I.

Spin-1/2	flips nucleon helicity	conserves nucleon helicity
does not depend on quark helicity	E	H
depends on quark helicity	\tilde{E}	\tilde{H}

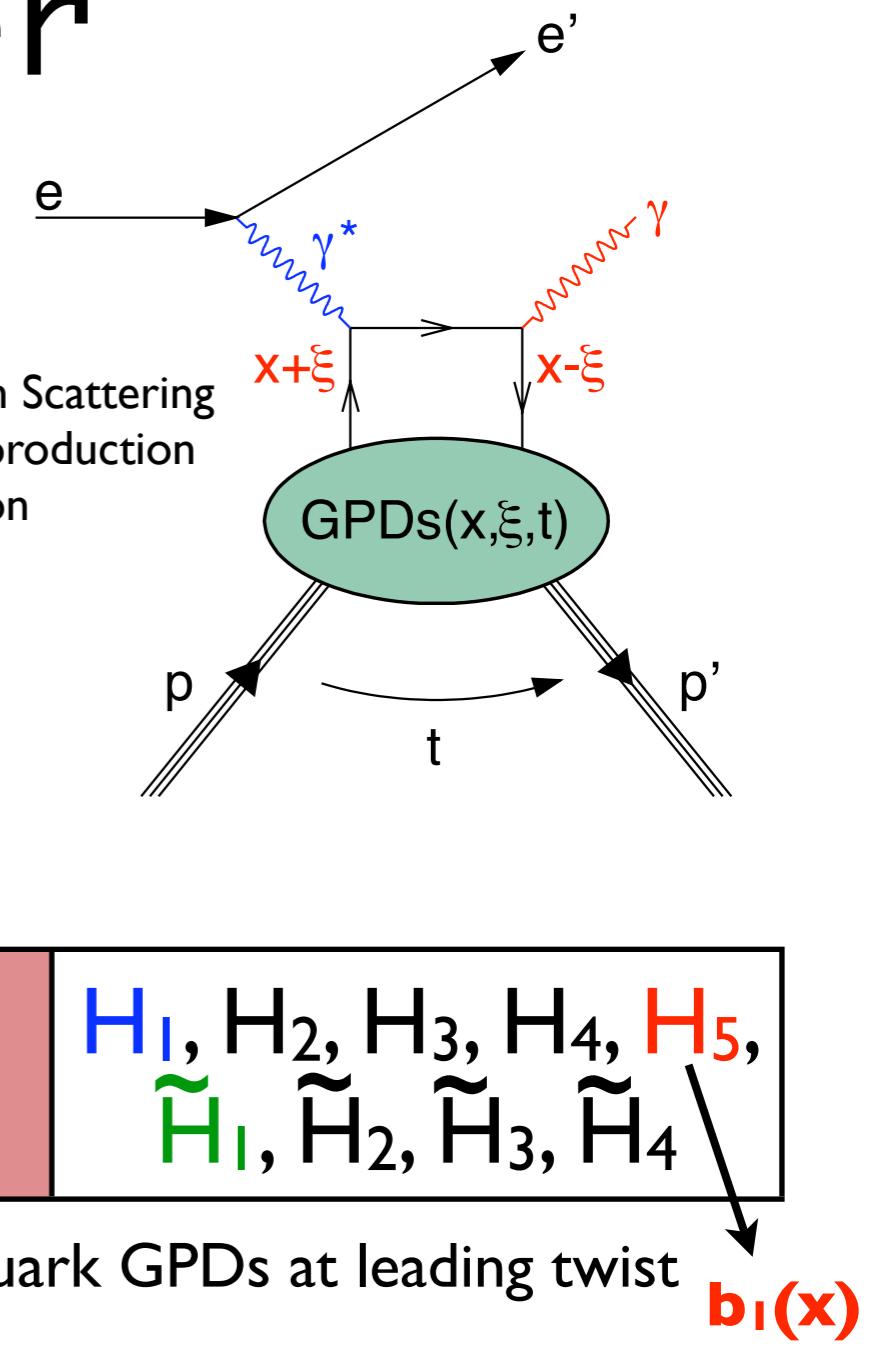
$\xrightarrow{\text{forward limit}} \xi \rightarrow 0, t \rightarrow 0$

$\mathbf{q}(x) \quad \Delta \mathbf{q}(x)$

4 chiral-even quark GPDs at leading twist

(4 more GPDs for chiral-odd case
related to transversity)

Deeply Virtual Compton Scattering
= hard exclusive lepto-production
of a real photon



Spin-1	$H_1, H_2, H_3, H_4, H_5,$ $\tilde{H}_1, \tilde{H}_2, \tilde{H}_3, \tilde{H}_4$
--------	--

9 chiral-even quark GPDs at leading twist
tensor structure function

H_3, H_5 associated with 5% D-wave component of deuteron wave function

DVCS on hadrons other than the nucleon

I.

Spin-1/2	flips nucleon helicity	conserves nucleon helicity
does not depend on quark helicity	E	H
depends on quark helicity	\tilde{E}	\tilde{H}

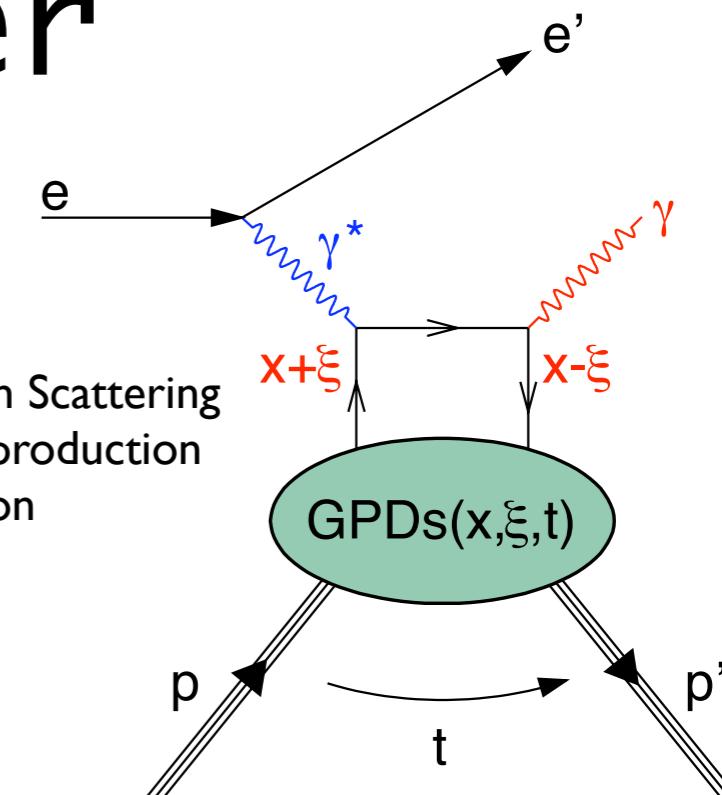
4 chiral-even quark GPDs at leading twist

(4 more GPDs for chiral-odd case related to transversity)

2.

How does the nuclear environment modify the DVCS amplitude?

Deeply Virtual Compton Scattering
= hard exclusive lepto-production of a real photon

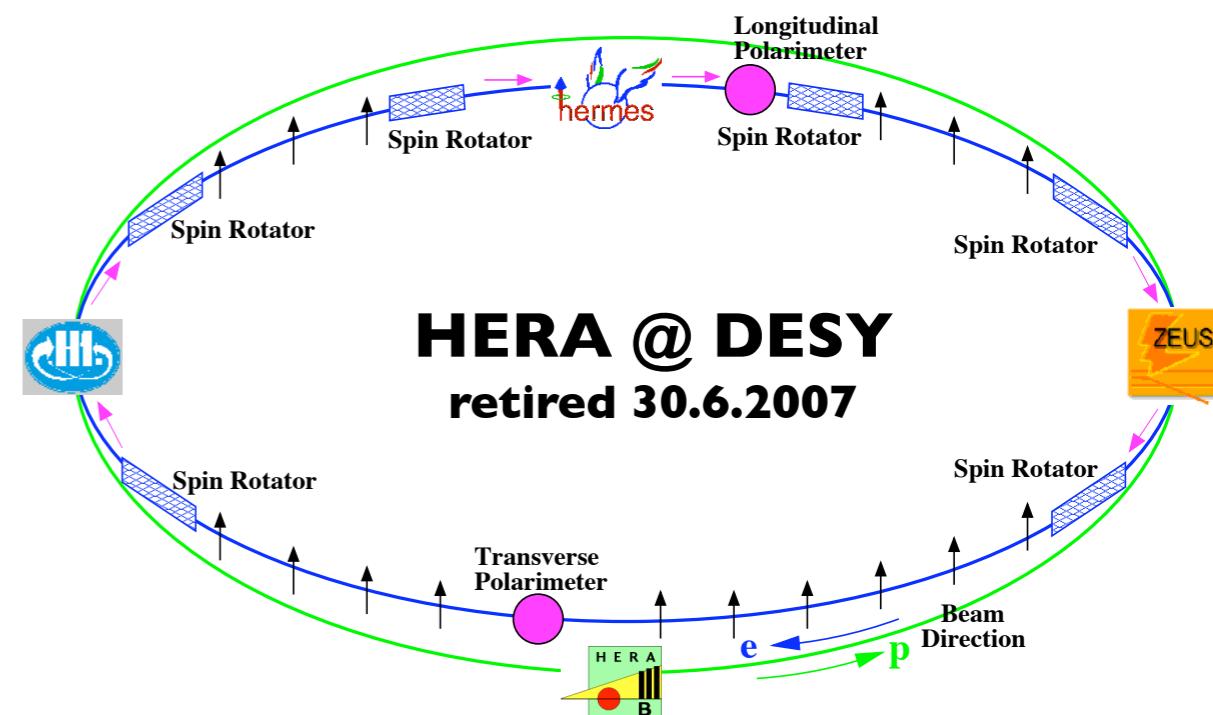


Spin-1	$H_1, H_2, H_3, H_4, H_5,$ $\tilde{H}_1, \tilde{H}_2, \tilde{H}_3, \tilde{H}_4$
	$b_1(x)$

9 chiral-even quark GPDs at leading twist
tensor structure function

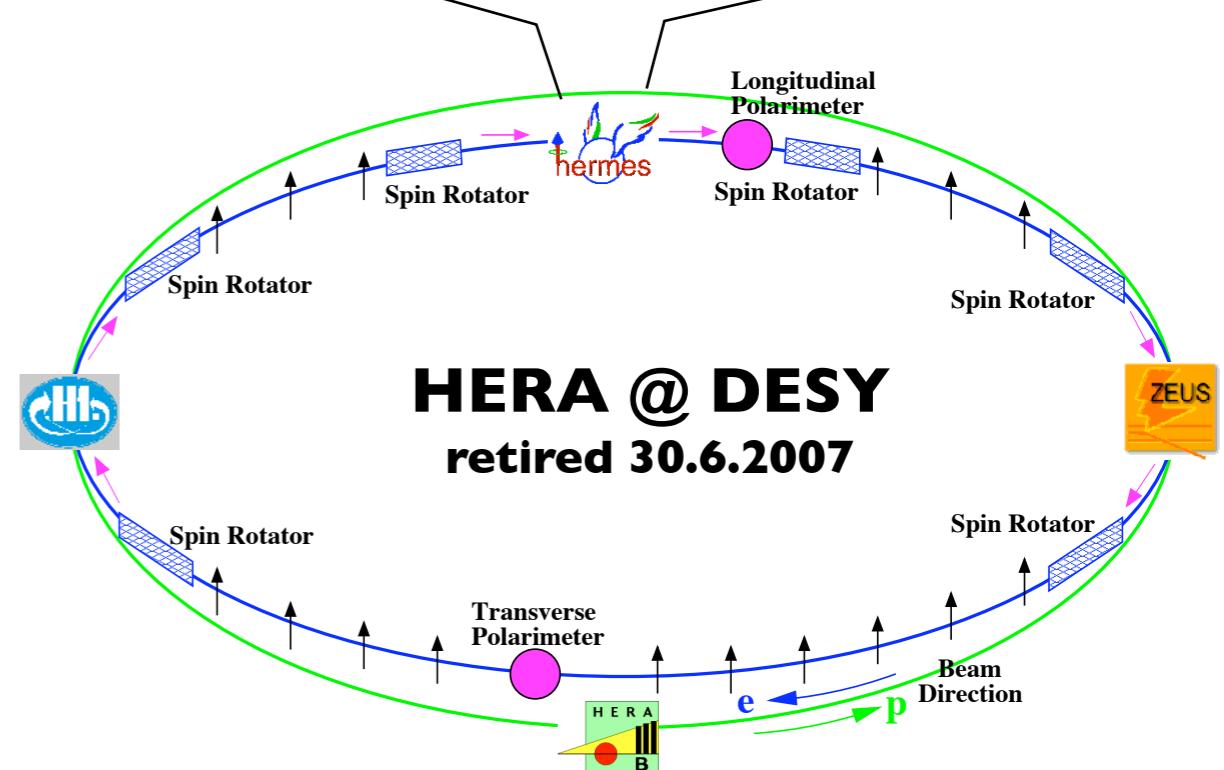
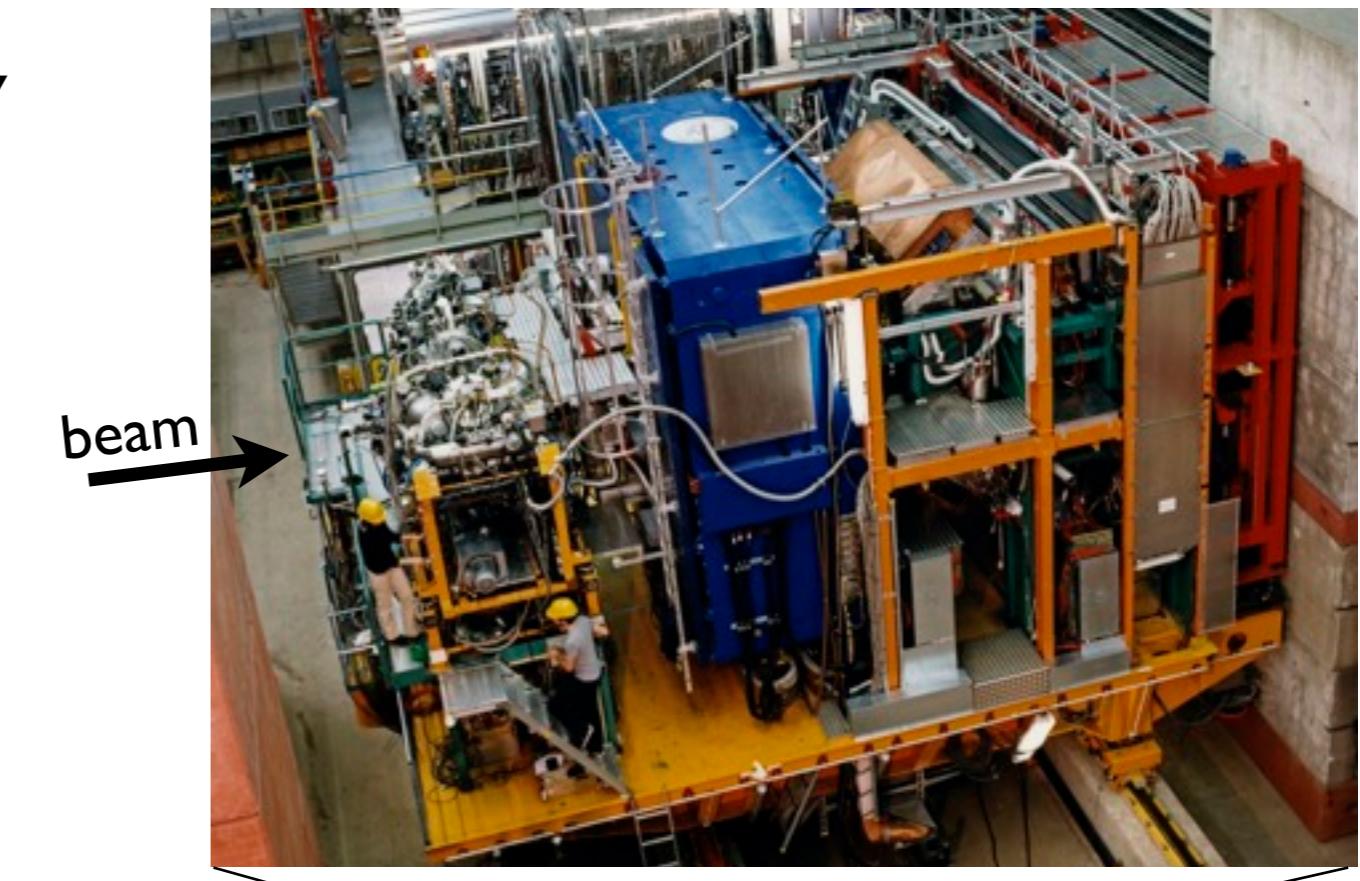
H_3, H_5 associated with 5% D-wave component of deuteron wave function

HERMES at DESY



self-polarized e+ and e- beams
helicity switched every few months

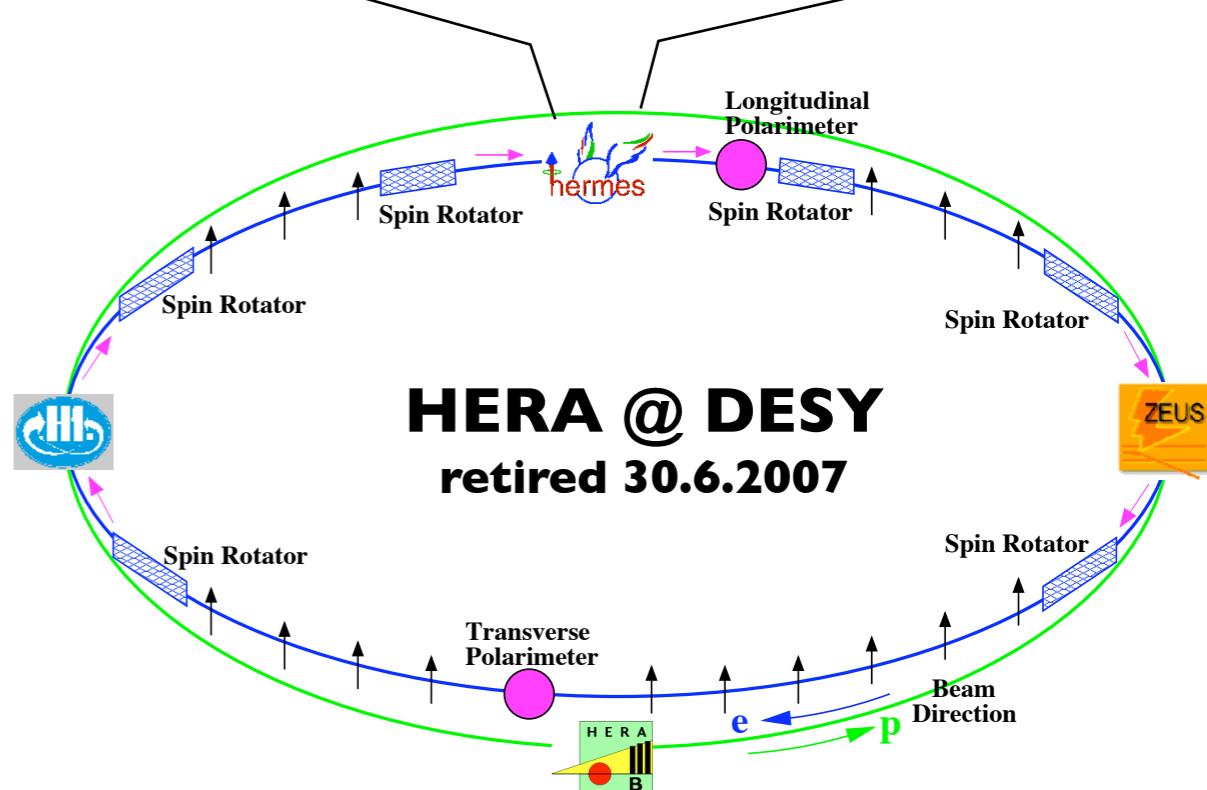
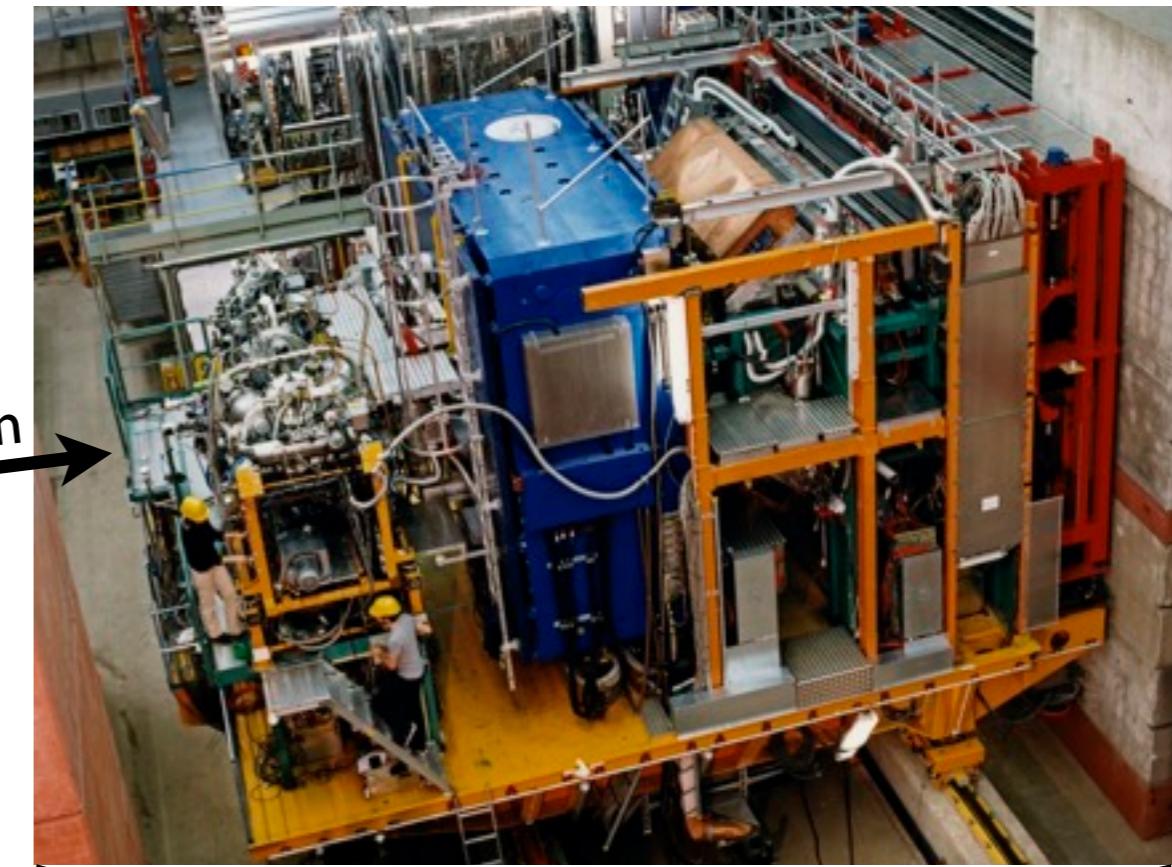
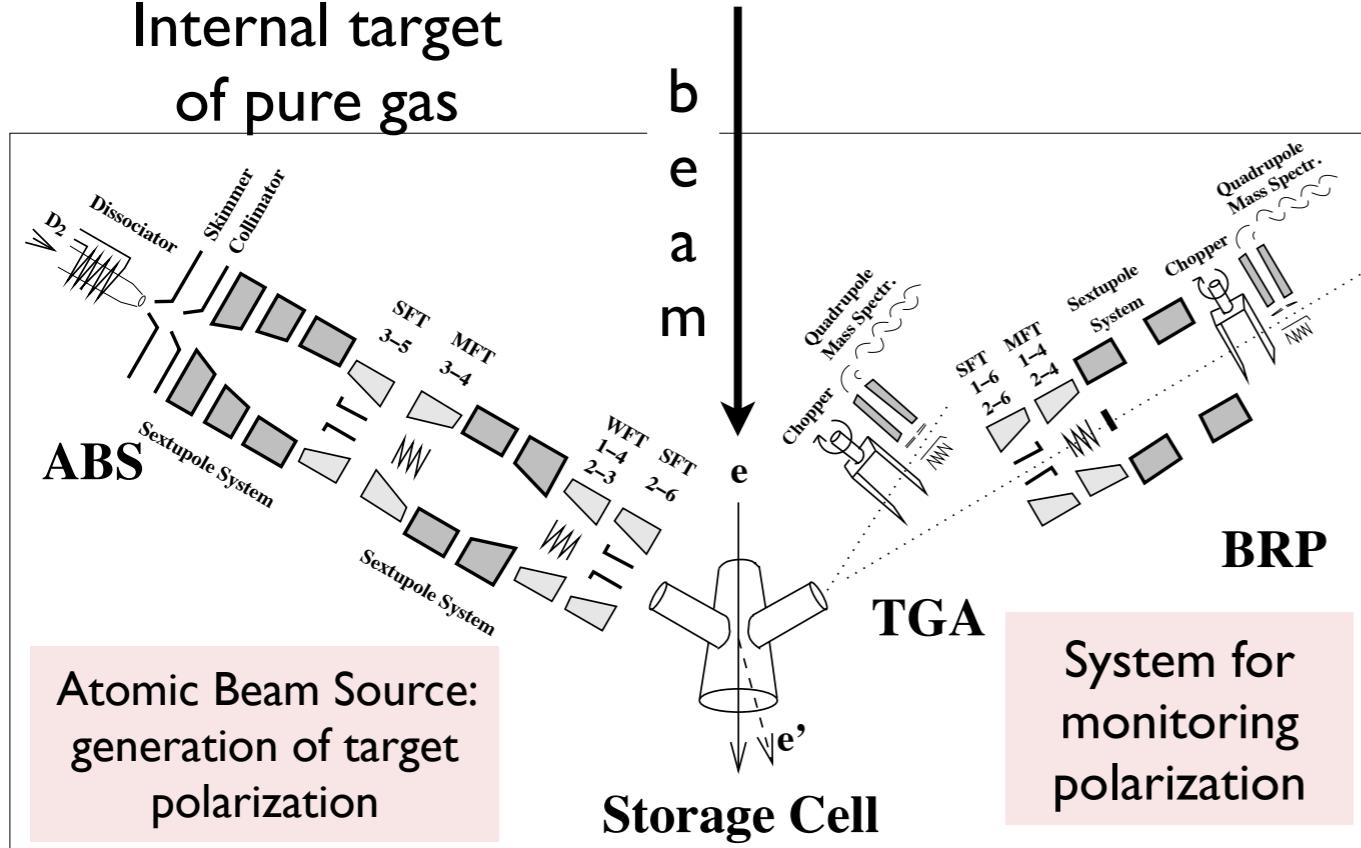
HERMES at DESY



self-polarized e+ and e- beams
helicity switched every few months

HERMES at DESY

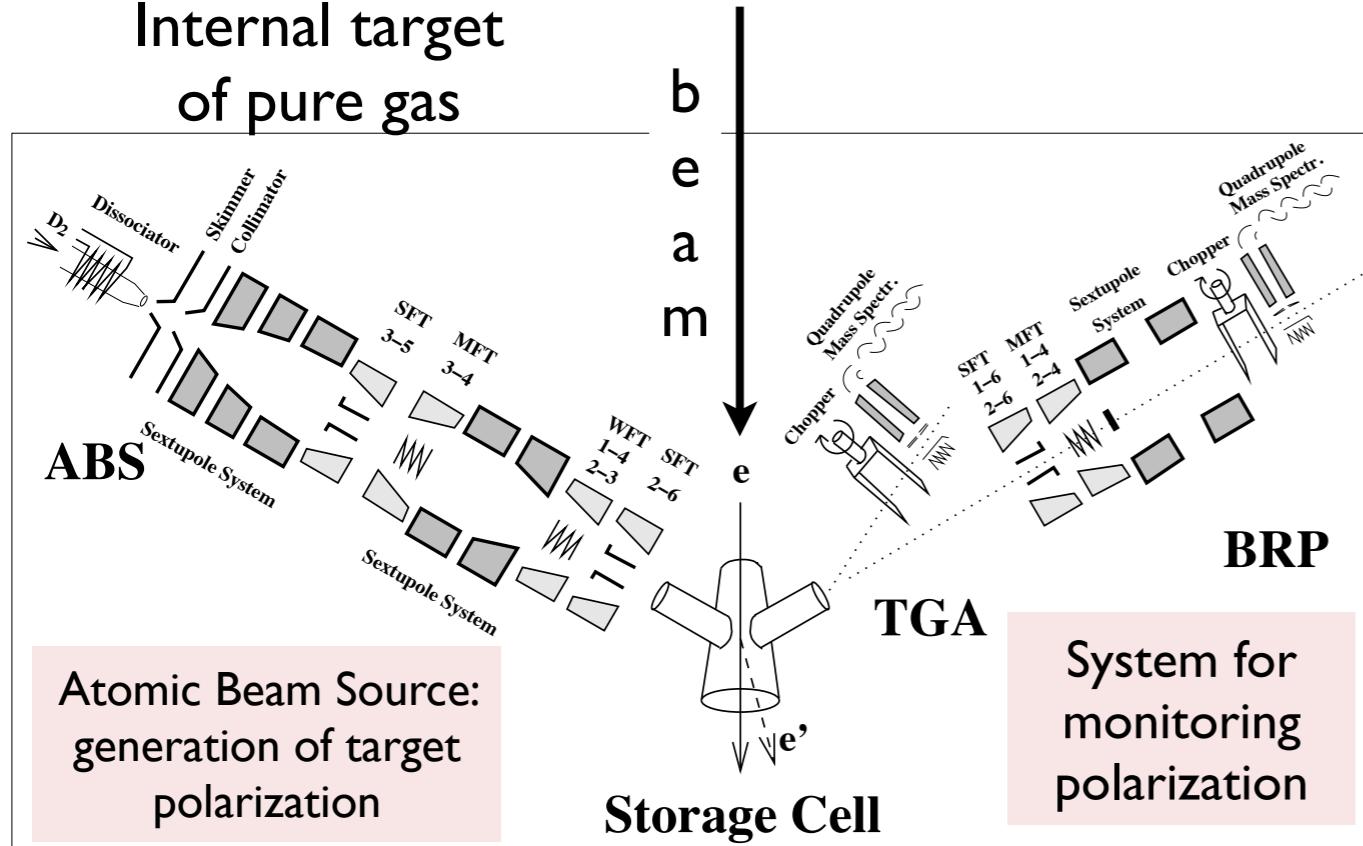
Internal target
of pure gas



self-polarized e+ and e- beams
helicity switched every few months

HERMES at DESY

Internal target
of pure gas



Atomic Beam Source:
generation of target polarization

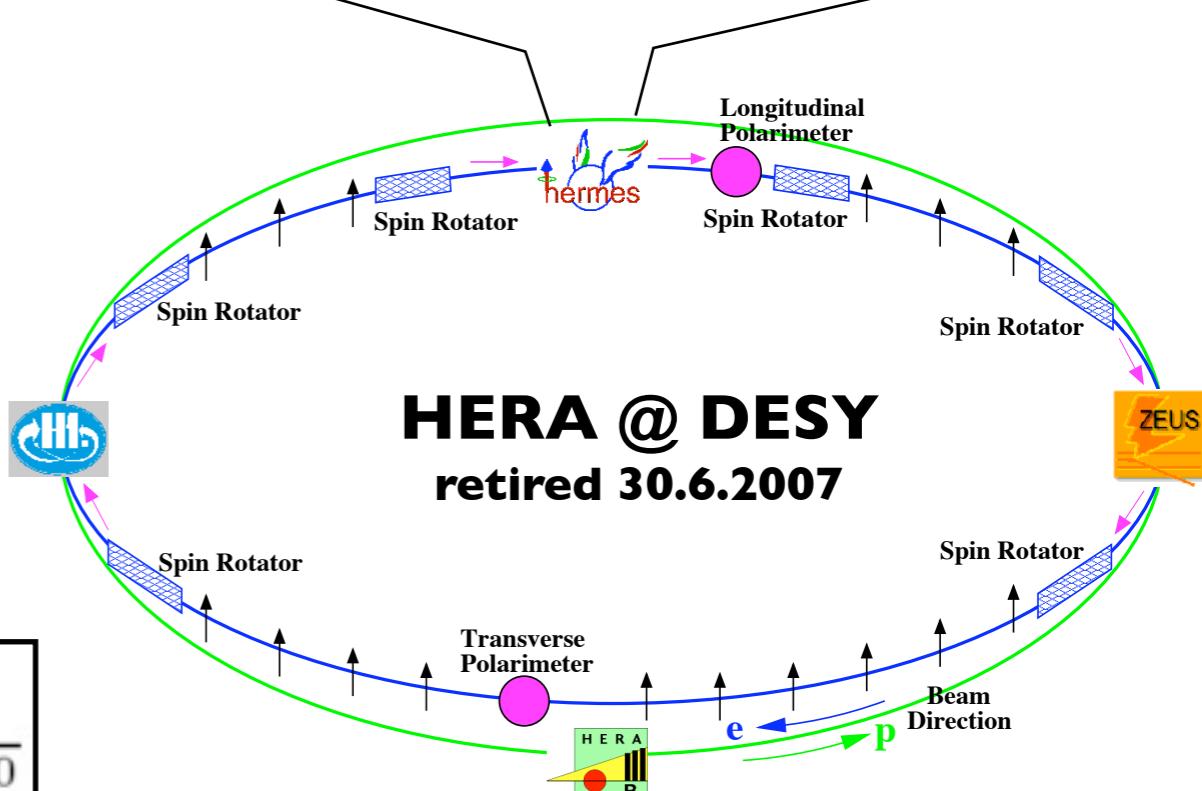
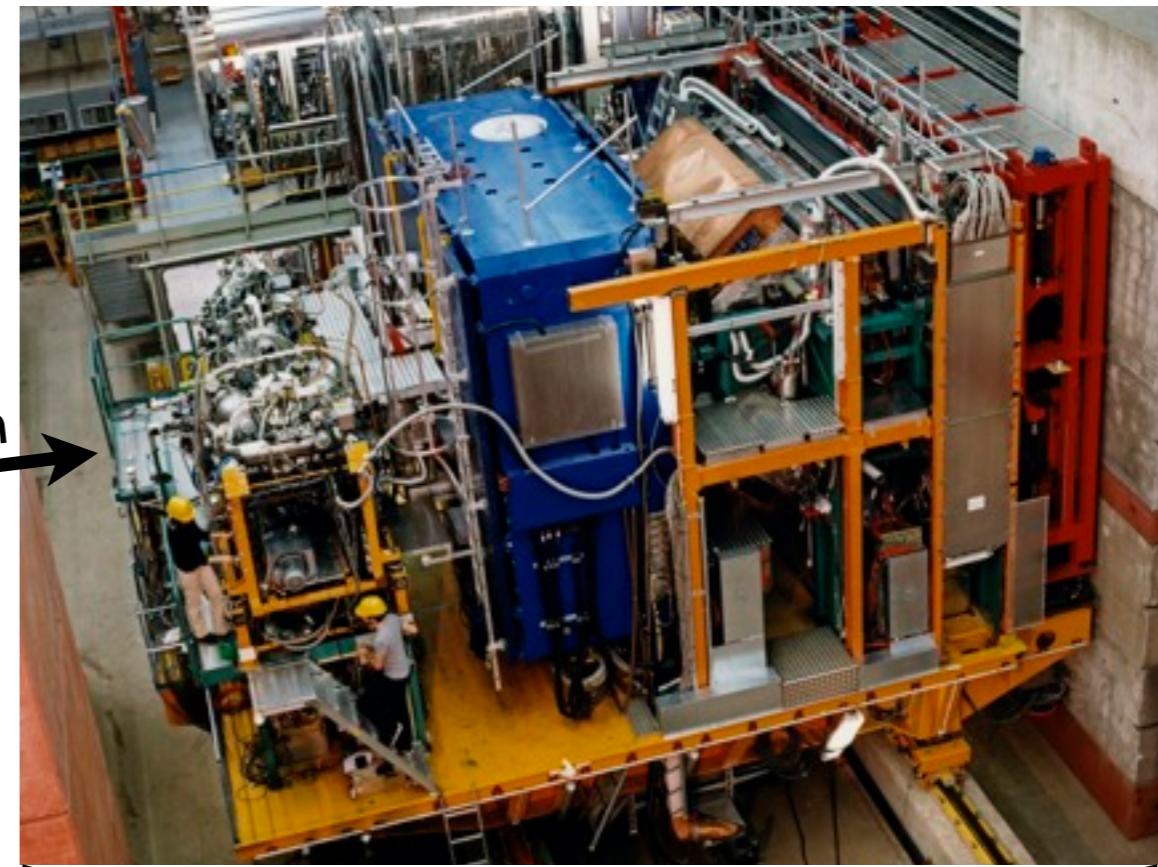
Longitudinally polarized deuterium
(1998-2000)

- » Vector polarization $P_z \approx 0.85$
- » **Tensor polarization** $P_{zz} \approx 0.83$
- » Dedicated data set with with $P_{zz} = -1.656 \text{ & } P_z \approx 0$

Spin-1 particle
with $\Lambda = -1, 0, +1$

$$P_z = \frac{n^+ - n^-}{n^+ + n^- + n^0}$$

$$P_{zz} = \frac{n^+ + n^- - 2n^0}{n^+ + n^- + n^0}$$

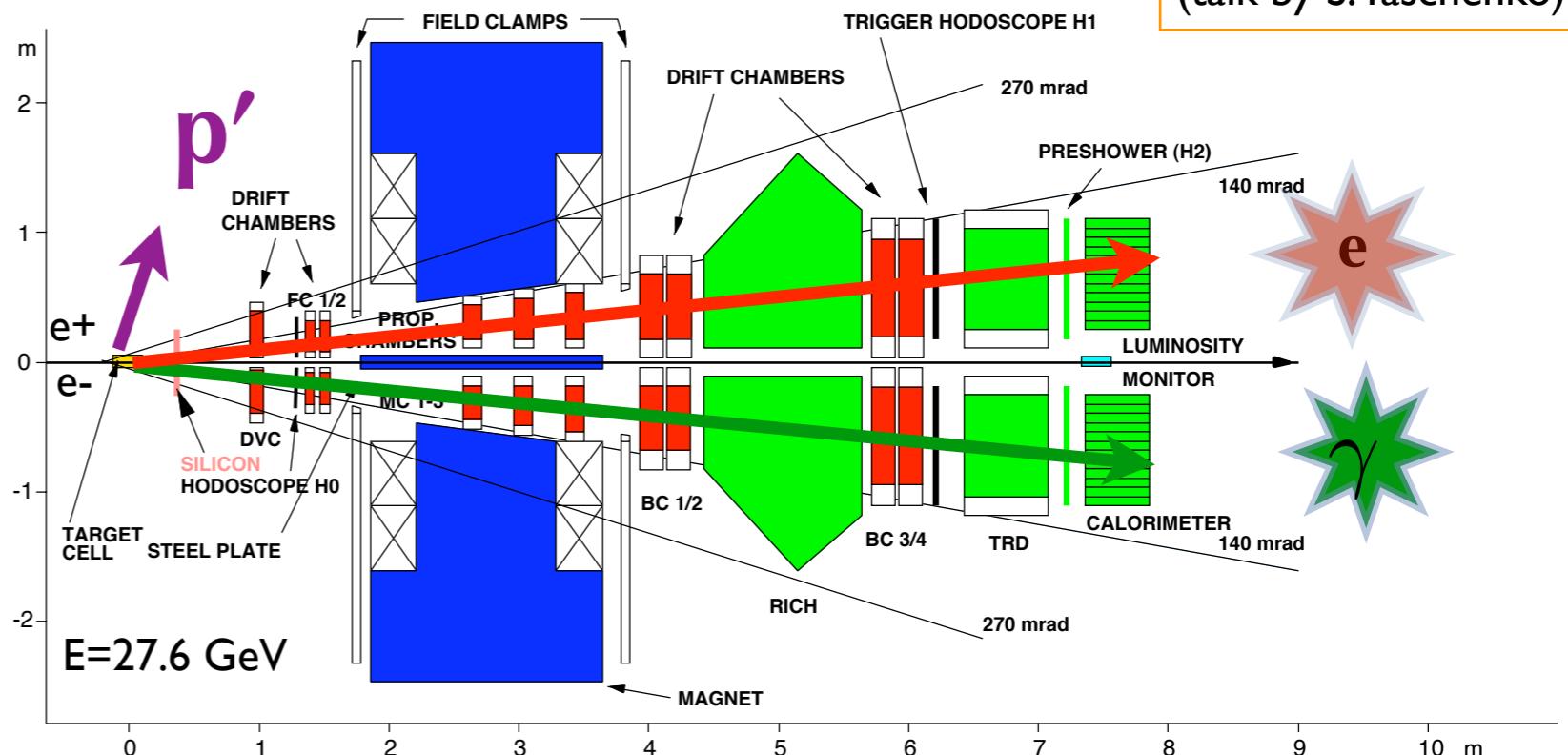


self-polarized e^+ and e^- beams
helicity switched every few months

DVCS at HERMES 1996-2005

2006/2007: recoil
proton detected
(talk by S.Yaschenko)

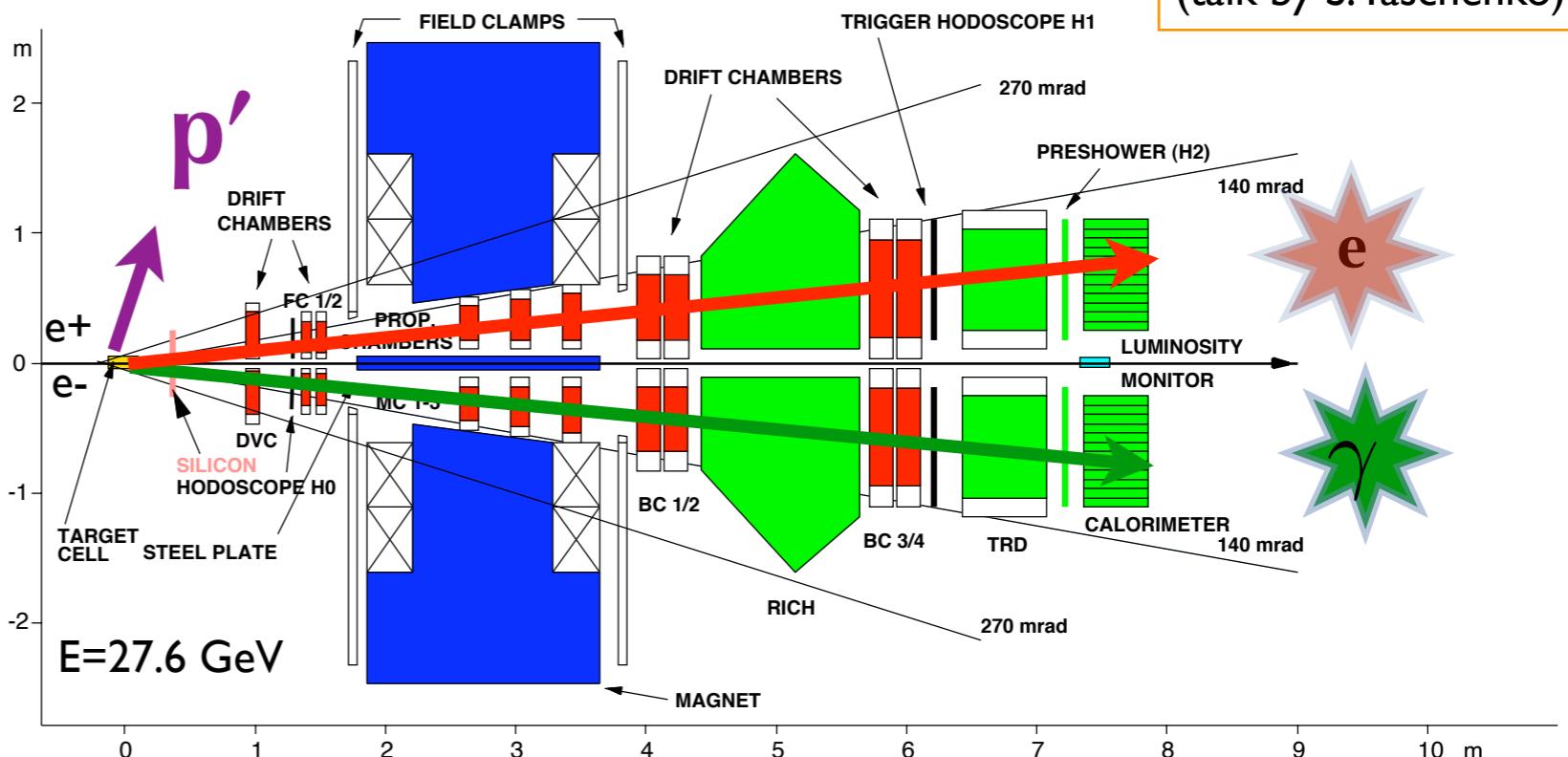
Hydrogen target:
400 pb⁻¹
unpolarized Deuterium:
300 pb⁻¹
L-polarized Deuterium:
200 pb⁻¹
Heavier Nuclear targets:
He, N, Ne, Kr, Xe
300pb⁻¹



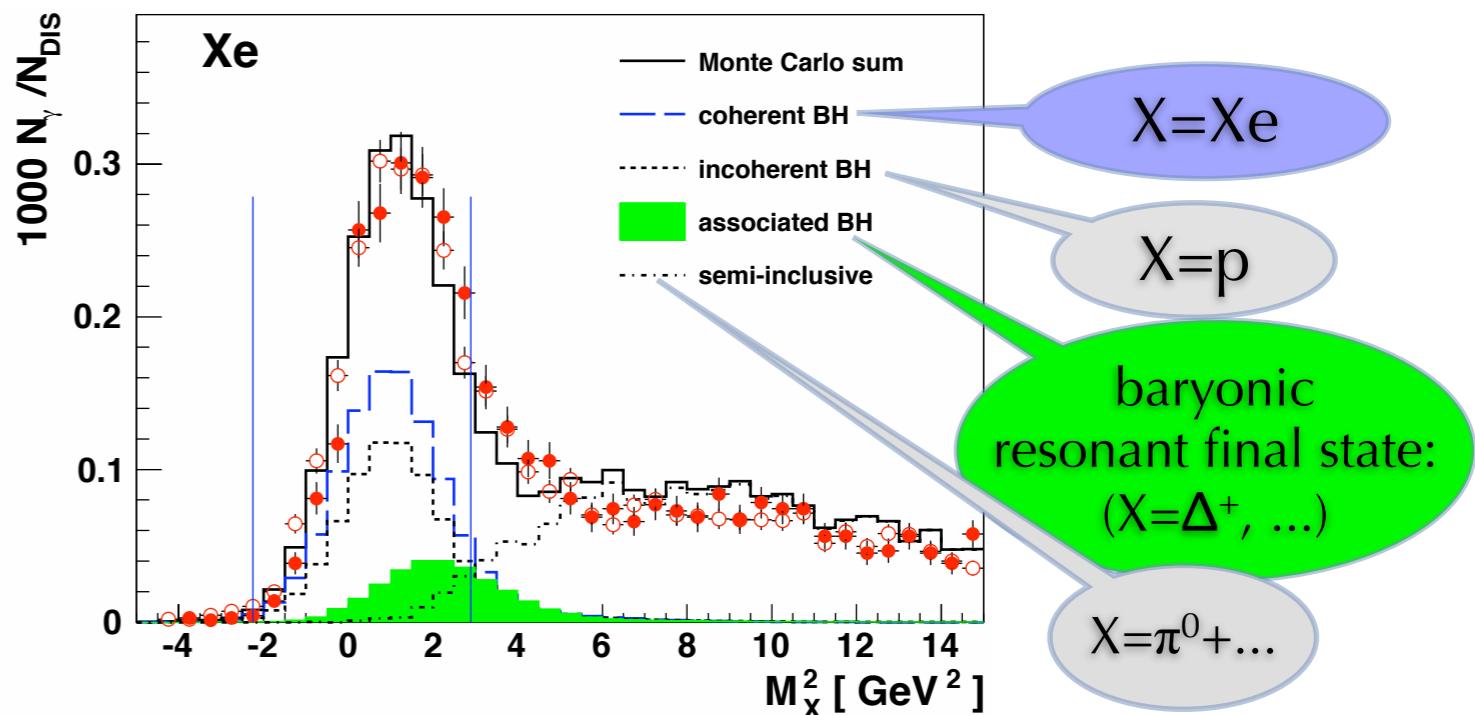
DVCS at HERMES 1996-2005

2006/2007: recoil
proton detected
(talk by S.Yaschenko)

Hydrogen target:
400 pb⁻¹
unpolarized Deuterium:
300 pb⁻¹
L-polarized Deuterium:
200 pb⁻¹
Heavier Nuclear targets:
He, N, Ne, Kr, Xe
300pb⁻¹



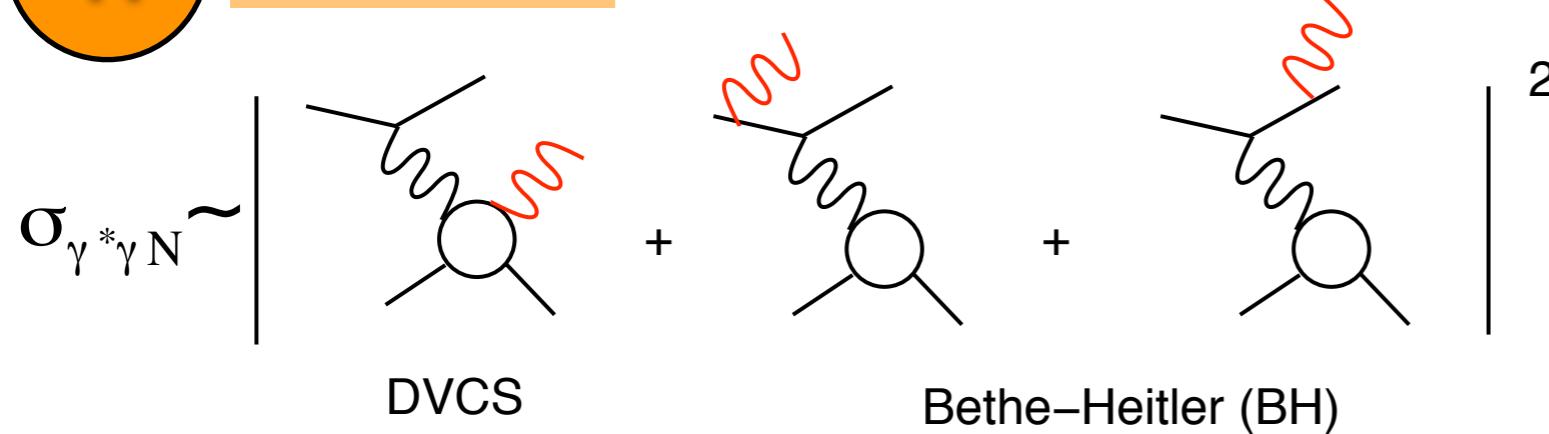
Missing mass technique for $e p \rightarrow e X \gamma$
 $M_X^2 = (p + q - p_\gamma)^2$



Deeply Virtual Compton Scattering

I.

Cross-section



Small at HERMES

Exactly calculable in QED
given the nucleon elastic
form factors F_1 and F_2

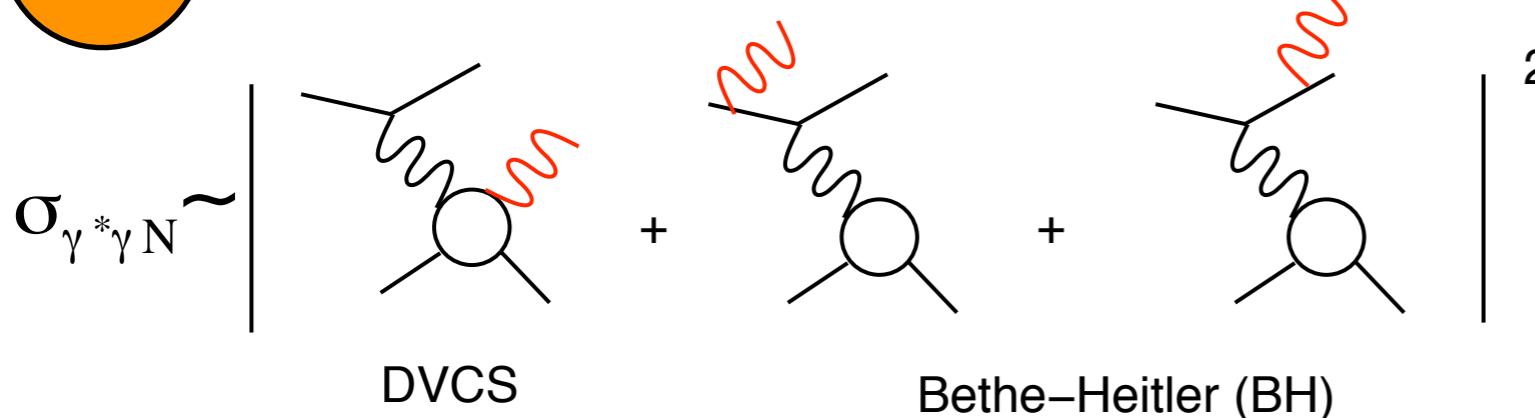
$$= |\tau_{\text{DVCS}}|^2 + |\tau_{\text{BH}}|^2 + (\tau_{\text{DVCS}} \tau_{\text{BH}}^* + \tau_{\text{DVCS}}^* \tau_{\text{BH}})$$

DVCS-BH interference term

Deeply Virtual Compton Scattering

I.

Cross-section



Small at HERMES

Exactly calculable in QED
given the nucleon elastic
form factors F_1 and F_2

$$= |\tau_{\text{DVCS}}|^2 + |\tau_{\text{BH}}|^2 +$$

$$(\tau_{\text{DVCS}} \tau_{\text{BH}}^* + \tau_{\text{DVCS}}^* \tau_{\text{BH}})$$

DVCS-BH interference term

2.

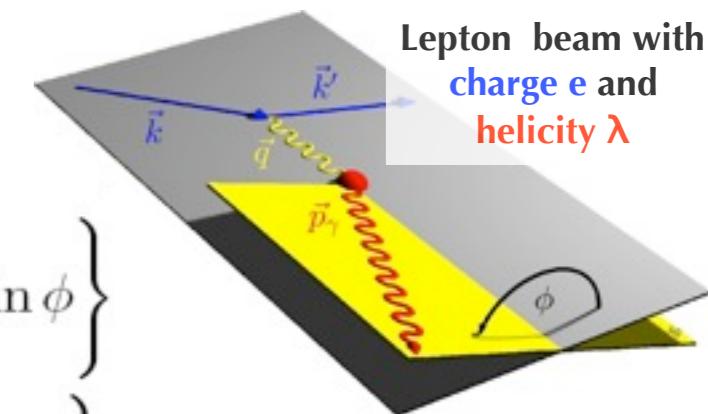
Harmonic expansion

Unpolarized target
**Case of polarized
target is more
complicated!**

$$|\tau_{\text{BH}}|^2 = \frac{K_{\text{BH}}}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^2 c_n^{\text{BH}} \cos(n\phi) \right\}$$

$$|\tau_{\text{DVCS}}|^2 = \frac{1}{Q^2} \left\{ \sum_{n=0}^2 c_n^{\text{DVCS}} \cos(n\phi) + \lambda s_1^{\text{DVCS}} \sin \phi \right\}$$

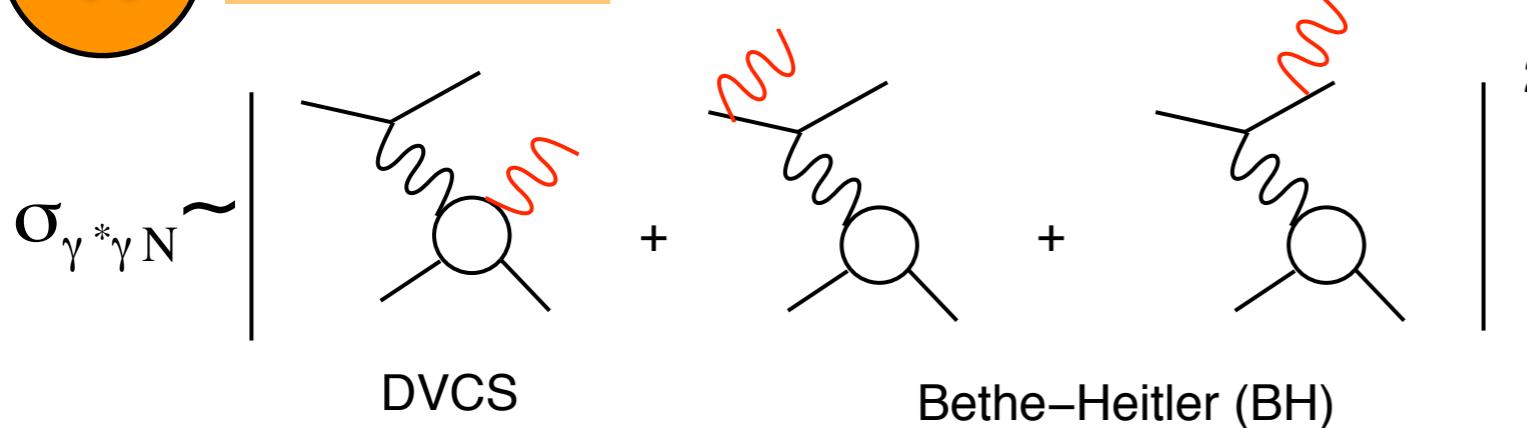
$$I = \frac{-e_\ell K_I}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^I \cos(n\phi) + \sum_{n=1}^2 \lambda s_n^I \sin(n\phi) \right\}$$



Deeply Virtual Compton Scattering

I.

Cross-section



Small at HERMES

Exactly calculable in QED
given the nucleon elastic
form factors F_1 and F_2

$$= |\tau_{\text{DVCS}}|^2 + |\tau_{\text{BH}}|^2 +$$

DVCS-BH interference term

2.

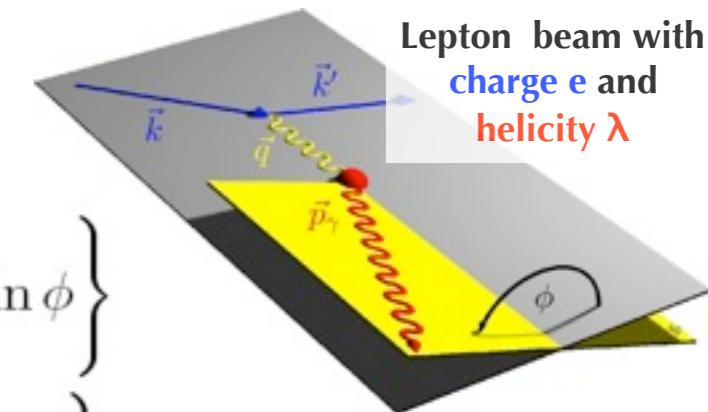
Harmonic expansion

Unpolarized target
**Case of polarized
target is more
complicated!**

$$|\tau_{\text{BH}}|^2 = \frac{K_{\text{BH}}}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^2 c_n^{\text{BH}} \cos(n\phi) \right\}$$

$$|\tau_{\text{DVCS}}|^2 = \frac{1}{Q^2} \left\{ \sum_{n=0}^2 c_n^{\text{DVCS}} \cos(n\phi) + \lambda s_1^{\text{DVCS}} \sin \phi \right\}$$

$$I = \frac{-e_\ell K_I}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^I \cos(n\phi) + \sum_{n=1}^2 \lambda s_n^I \sin(n\phi) \right\}$$



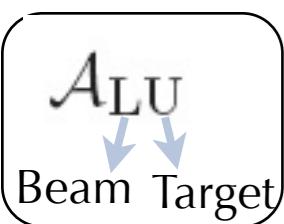
3.

Express cross-section in terms of azimuthal asymmetries

$$\sigma(\phi; P_\ell, e_\ell) = \sigma_{\text{UU}}(\phi) \times [1 + P_\ell \mathcal{A}_{\text{LU}}^{\text{DVCS}}(\phi) + e_\ell P_\ell \mathcal{A}_{\text{LU}}^I(\phi) + e_\ell \mathcal{A}_C(\phi)]$$

\mathcal{A}_{LU}
Beam Target

Azimuthal Asymmetries and GPDs



$$I = \frac{-e_\ell K_I}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^I \cos(n\phi) + \sum_{n=1}^2 \lambda s_n^I \sin(n\phi) \right\}$$

Fourier coefficients = azimuthal asymmetry amplitudes are related to certain linear or bi-linear combinations of CFFs.

Compton Form Factors (CFFs)

$$\mathcal{F}(\xi, t) = \sum_q \int_{-1}^1 dx \, C_q^\mp(\xi, x) F^q(x, \xi, t)$$

twist-2 GPD

Azimuthal Asymmetries and GPDs

$$\mathcal{A}_{LU}$$

Beam Target

**Single-charge
beam-helicity
asymmetry**

$$\mathcal{A}_{LU}(\phi) \equiv \frac{d\sigma^{\rightarrow} - d\sigma^{\leftarrow}}{d\sigma^{\rightarrow} + d\sigma^{\leftarrow}}$$

no separate access
to s_I^I and s_I^{DVCS}

$$I = \frac{-e_\ell K_I}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^I \cos(n\phi) + \sum_{n=1}^2 \lambda s_n^I \sin(n\phi) \right\}$$

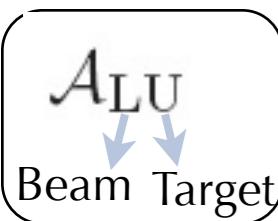
Fourier coefficients = azimuthal asymmetry
amplitudes are related to certain linear or
bi-linear combinations of CFFs.

Compton Form Factors (CFFs)

$$\mathcal{F}(\xi, t) = \sum_q \int_{-1}^1 dx C_q^{\mp}(\xi, x) F^q(x, \xi, t)$$

twist-2 GPD

Azimuthal Asymmetries and GPDs



**Single-charge
beam-helicity
asymmetry**

$$\mathcal{A}_{LU}(\phi) \equiv \frac{d\sigma^{\rightarrow} - d\sigma^{\leftarrow}}{d\sigma^{\rightarrow} + d\sigma^{\leftarrow}}$$

no separate access
to s_I^l and s_I^{DVCS}

**Beam-helicity asymmetries
with 2 beam charges**

Charge-average

$$\mathcal{A}_{LU}$$

Charge-difference

$$\mathcal{A}_{LU}$$

s_I^{DVCS} and s_I^l can be disentangled

$$I = \frac{-e_\ell K_I}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^I \cos(n\phi) + \sum_{n=1}^2 \lambda s_n^I \sin(n\phi) \right\}$$

Fourier coefficients = azimuthal asymmetry
amplitudes are related to certain linear or
bi-linear combinations of CFFs.

Compton Form Factors (CFFs)

$$\mathcal{F}(\xi, t) = \sum_q \int_{-1}^1 dx C_q^{\mp}(\xi, x) F^q(x, \xi, t)$$

twist-2 GPD

Azimuthal Asymmetries and GPDs

$$\mathcal{A}_{LU}$$

**Single-charge
beam-helicity
asymmetry**

$$\mathcal{A}_{LU}(\phi) \equiv \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-}$$

no separate access
to s_I^+ and s_I^-

**Beam-helicity asymmetries
with 2 beam charges**

Charge-average

$$\mathcal{A}_{LU}$$

Charge-difference

$$\mathcal{A}_{LU}$$

s_I^- ^{DVCS} and s_I^+ can be disentangled

**Beam-charge
asymmetry**

$$\mathcal{A}_C(\phi) \equiv \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-}$$

$$I = \frac{-e_\ell K_I}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^I \cos(n\phi) + \sum_{n=1}^2 \lambda s_n^I \sin(n\phi) \right\}$$

Fourier coefficients = azimuthal asymmetry
amplitudes are related to certain linear or
bi-linear combinations of CFFs.

Compton Form Factors (CFFs)

$$\mathcal{F}(\xi, t) = \sum_q \int_{-1}^1 dx C_q^\mp(\xi, x) F^q(x, \xi, t)$$

twist-2 GPD

Azimuthal Asymmetries and GPDs

$$\mathcal{A}_{LU}$$

**Single-charge
beam-helicity
asymmetry**

$$\mathcal{A}_{LU}(\phi) \equiv \frac{d\sigma^{\rightarrow} - d\sigma^{\leftarrow}}{d\sigma^{\rightarrow} + d\sigma^{\leftarrow}}$$

no separate access
to s_l^I and s_l^{DVCS}

**Beam-helicity asymmetries
with 2 beam charges**

Charge-average

$$\mathcal{A}_{LU}$$

Charge-difference

$$\mathcal{A}_{LU}$$

s_l^{DVCS} and s_l^I can be disentangled

**Beam-charge
asymmetry**

$$\mathcal{A}_C(\phi) \equiv \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-}$$

$$I = \frac{-e_\ell K_I}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^I \cos(n\phi) + \sum_{n=1}^2 \lambda s_n^I \sin(n\phi) \right\}$$

Fourier coefficients = azimuthal asymmetry
amplitudes are related to certain linear or
bi-linear combinations of CFFs.

Compton Form Factors (CFFs)

$$\mathcal{F}(\xi, t) = \sum_q \int_{-1}^1 dx C_q^\mp(\xi, x) F^q(x, \xi, t)$$

twist-2 GPD

$$\mathcal{A}_{UL}(\phi, e_\ell) \equiv$$

Target-spin asymmetry

$$\frac{[\sigma^{\leftarrow\rightarrow}(\phi, e_\ell) + \sigma^{\rightarrow\rightarrow}(\phi, e_\ell)] - [\sigma^{\leftarrow\leftarrow}(\phi, e_\ell) + \sigma^{\rightarrow\leftarrow}(\phi, e_\ell)]}{[\sigma^{\leftarrow\rightarrow}(\phi, e_\ell) + \sigma^{\rightarrow\rightarrow}(\phi, e_\ell)] + [\sigma^{\leftarrow\leftarrow}(\phi, e_\ell) + \sigma^{\rightarrow\leftarrow}(\phi, e_\ell)]}$$

Azimuthal Asymmetries and GPDs

$$\mathcal{A}_{LU}$$

**Single-charge
beam-helicity
asymmetry**

$$\mathcal{A}_{LU}(\phi) \equiv \frac{d\sigma^{\rightarrow} - d\sigma^{\leftarrow}}{d\sigma^{\rightarrow} + d\sigma^{\leftarrow}}$$

no separate access
to s_l^I and s_l^{DVCS}

$$I = \frac{-e_\ell K_I}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^I \cos(n\phi) + \sum_{n=1}^2 \lambda s_n^I \sin(n\phi) \right\}$$

Fourier coefficients = azimuthal asymmetry
amplitudes are related to certain linear or
bi-linear combinations of CFFs.

Compton Form Factors (CFFs)

$$\mathcal{F}(\xi, t) = \sum_q \int_{-1}^1 dx C_q^{\mp}(\xi, x) F^q(x, \xi, t)$$

twist-2 GPD

**Beam-helicity asymmetries
with 2 beam charges**

Charge-average

$$\mathcal{A}_{LU}$$

Charge-difference

$$\mathcal{A}_{LU}$$

s_l^{DVCS} and s_l^I can be disentangled

**Beam-charge
asymmetry**

$$\mathcal{A}_C(\phi) \equiv \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-}$$

$$\mathcal{A}_{UL}(\phi, e_\ell) \equiv$$

Target-spin asymmetry

$$\frac{[\sigma^{\leftarrow\rightarrow}(\phi, e_\ell) + \sigma^{\rightarrow\rightarrow}(\phi, e_\ell)] - [\sigma^{\leftarrow\leftarrow}(\phi, e_\ell) + \sigma^{\rightarrow\leftarrow}(\phi, e_\ell)]}{[\sigma^{\leftarrow\rightarrow}(\phi, e_\ell) + \sigma^{\rightarrow\rightarrow}(\phi, e_\ell)] + [\sigma^{\leftarrow\leftarrow}(\phi, e_\ell) + \sigma^{\rightarrow\leftarrow}(\phi, e_\ell)]}$$

$$\mathcal{A}_{LL}(\phi, e_\ell) \equiv$$

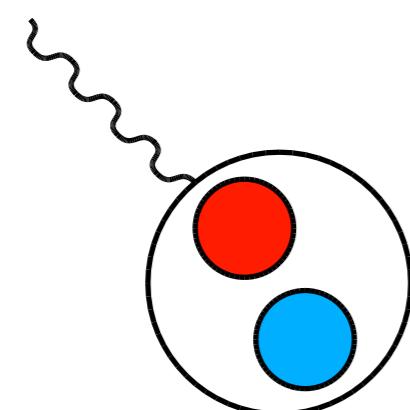
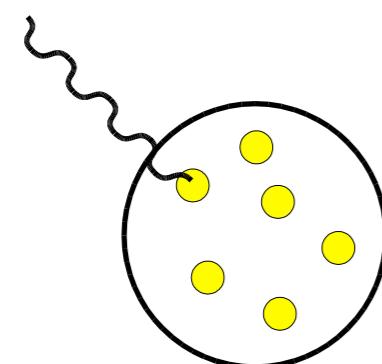
Double-spin asymmetry

$$\frac{[\sigma^{\rightarrow\rightarrow}(\phi, e_\ell) + \sigma^{\leftarrow\leftarrow}(\phi, e_\ell)] - [\sigma^{\leftarrow\rightarrow}(\phi, e_\ell) + \sigma^{\rightarrow\leftarrow}(\phi, e_\ell)]}{[\sigma^{\rightarrow\rightarrow}(\phi, e_\ell) + \sigma^{\leftarrow\leftarrow}(\phi, e_\ell)] + [\sigma^{\leftarrow\rightarrow}(\phi, e_\ell) + \sigma^{\rightarrow\leftarrow}(\phi, e_\ell)]}$$

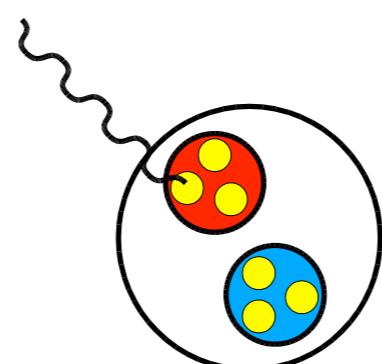
Coherent vs. incoherent

$$eA \rightarrow eA\gamma \quad eA \rightarrow e(A-1)N\gamma$$

DVCS

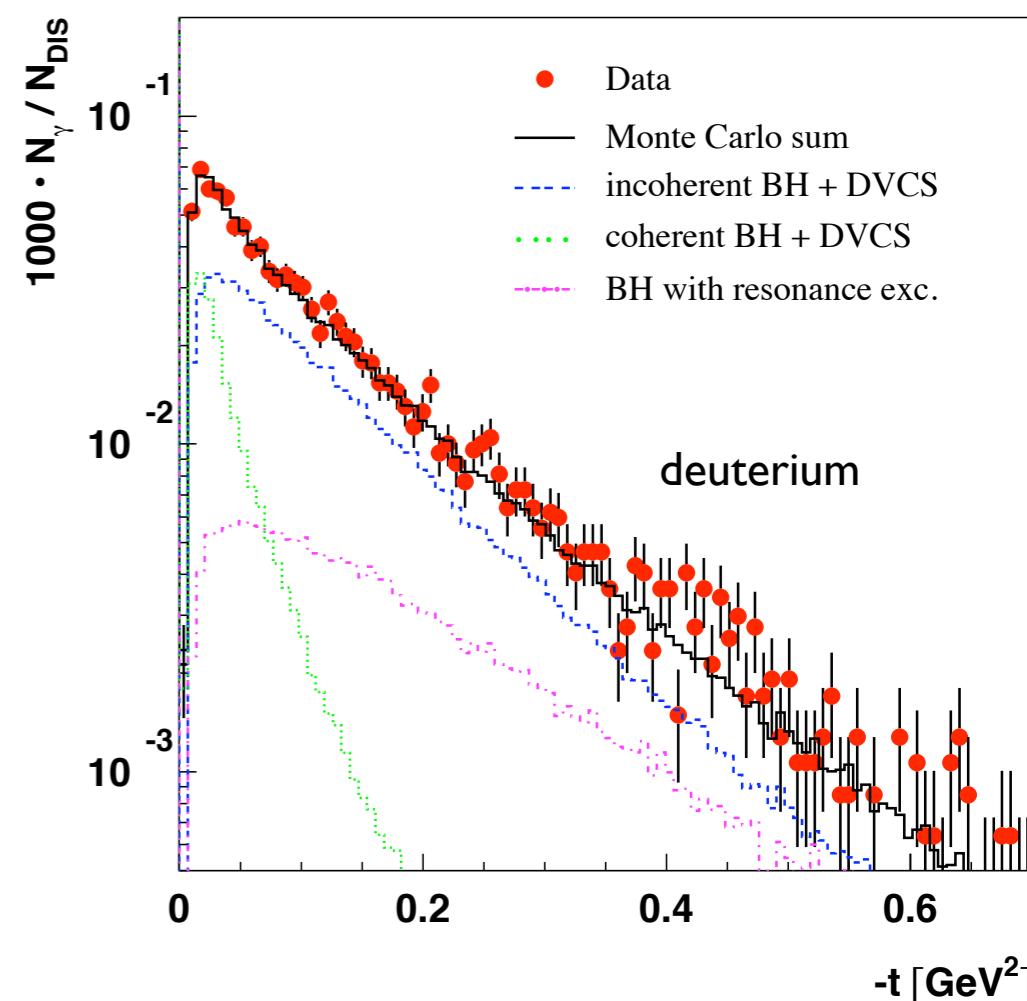


Deuteron: probe
spin-1 object



Nucleon: probe
spin-1/2 object

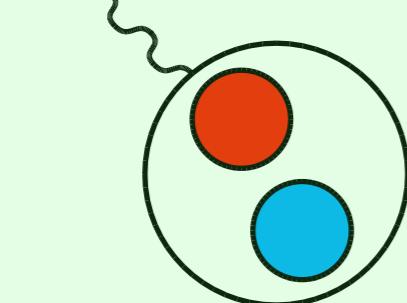
BH (proton) \gg BH (neutron)
due to elastic electric form factor



Coherent vs. incoherent

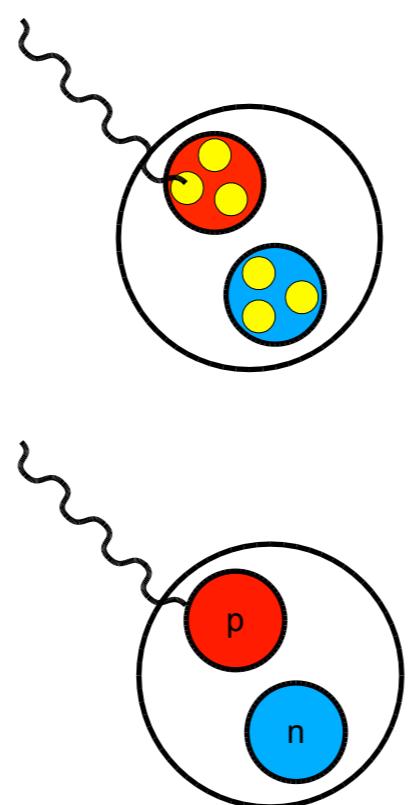
$$eA \rightarrow eA\gamma$$

DVCS



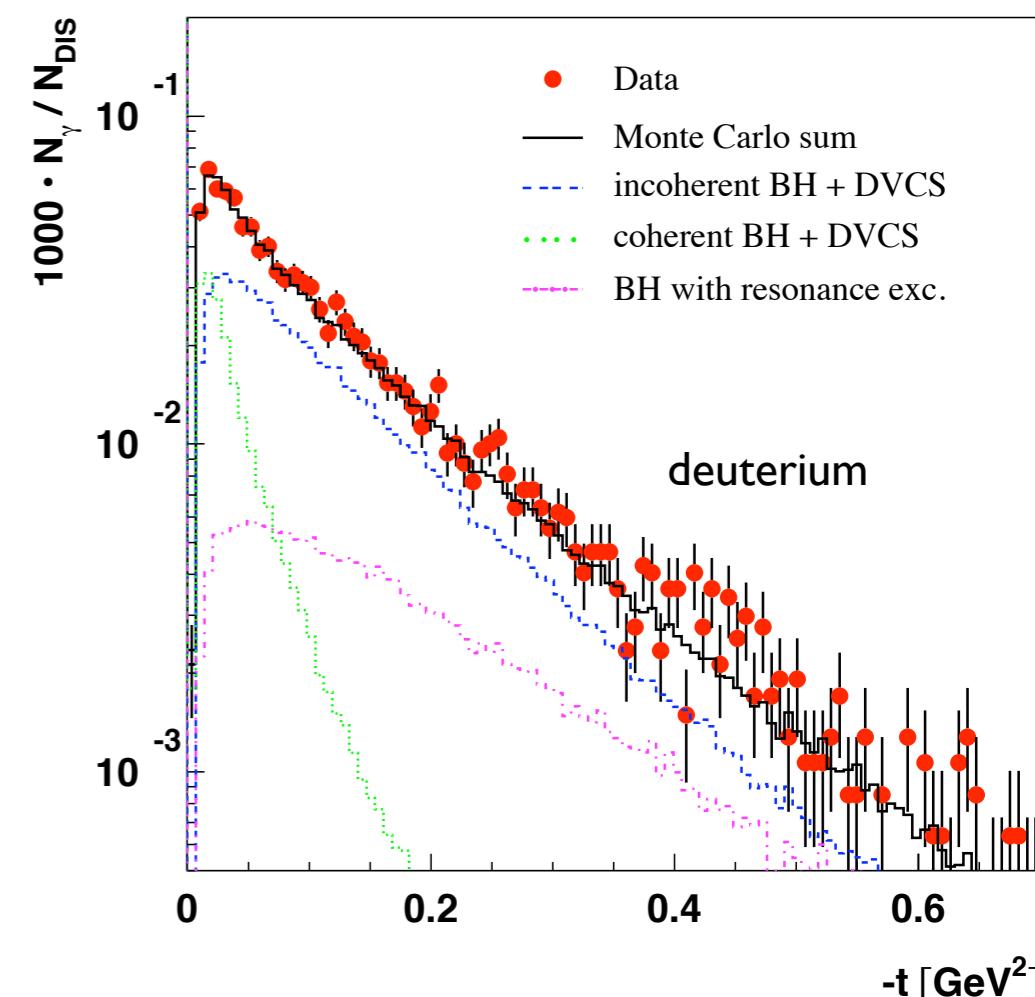
Deuteron: probe
spin-1 object

$$eA \rightarrow e(A-1)N\gamma$$



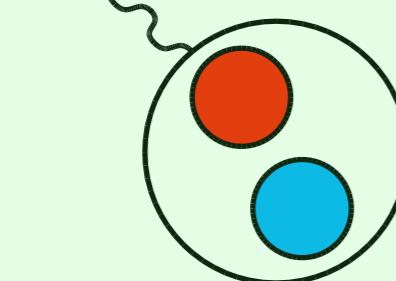
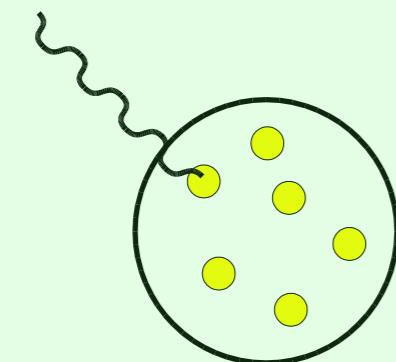
Nucleon: probe
spin-1/2 object

BH (proton) \gg BH (neutron)
due to elastic electric form factor



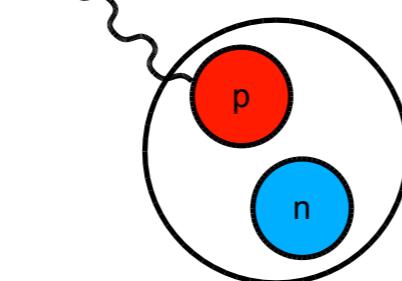
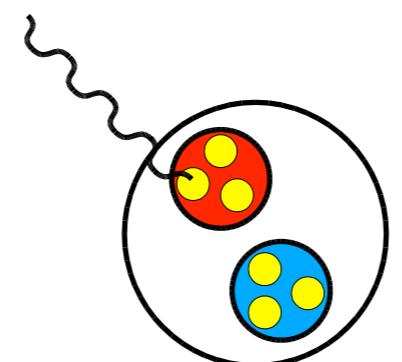
Coherent vs. incoherent

$eA \rightarrow eA\gamma$



Deuteron: probe
spin-1 object

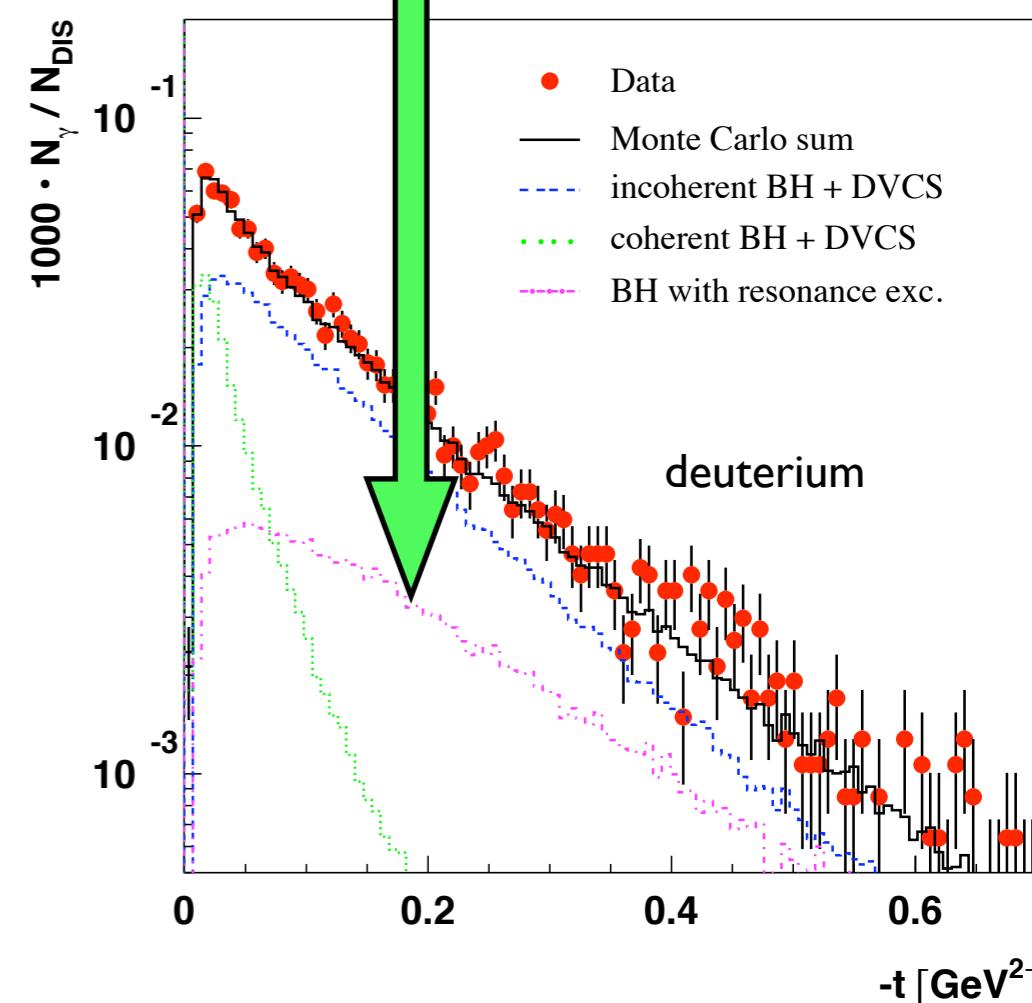
$eA \rightarrow e(A-1)N\gamma$



Nucleon: probe
spin-1/2 object

BH (proton) \gg BH (neutron)
due to elastic electric form factor

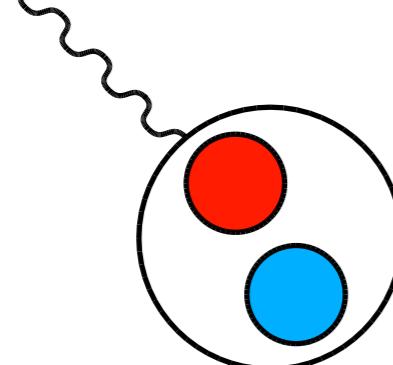
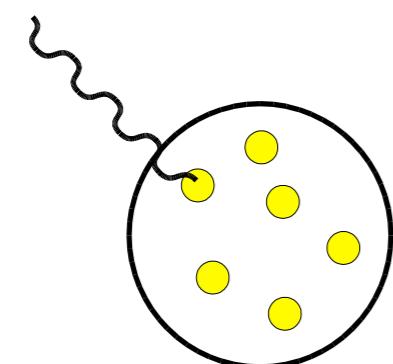
Coherent contribution
rapidly decreasing with $-t$



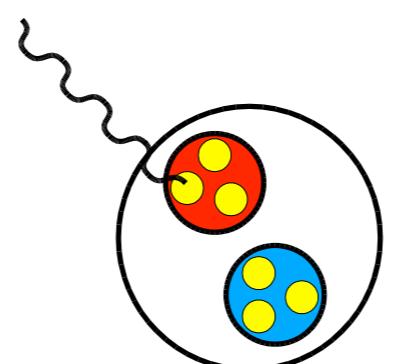
Coherent vs. incoherent

$$eA \rightarrow eA\gamma \quad eA \rightarrow e(A-1)N\gamma$$

DVCS



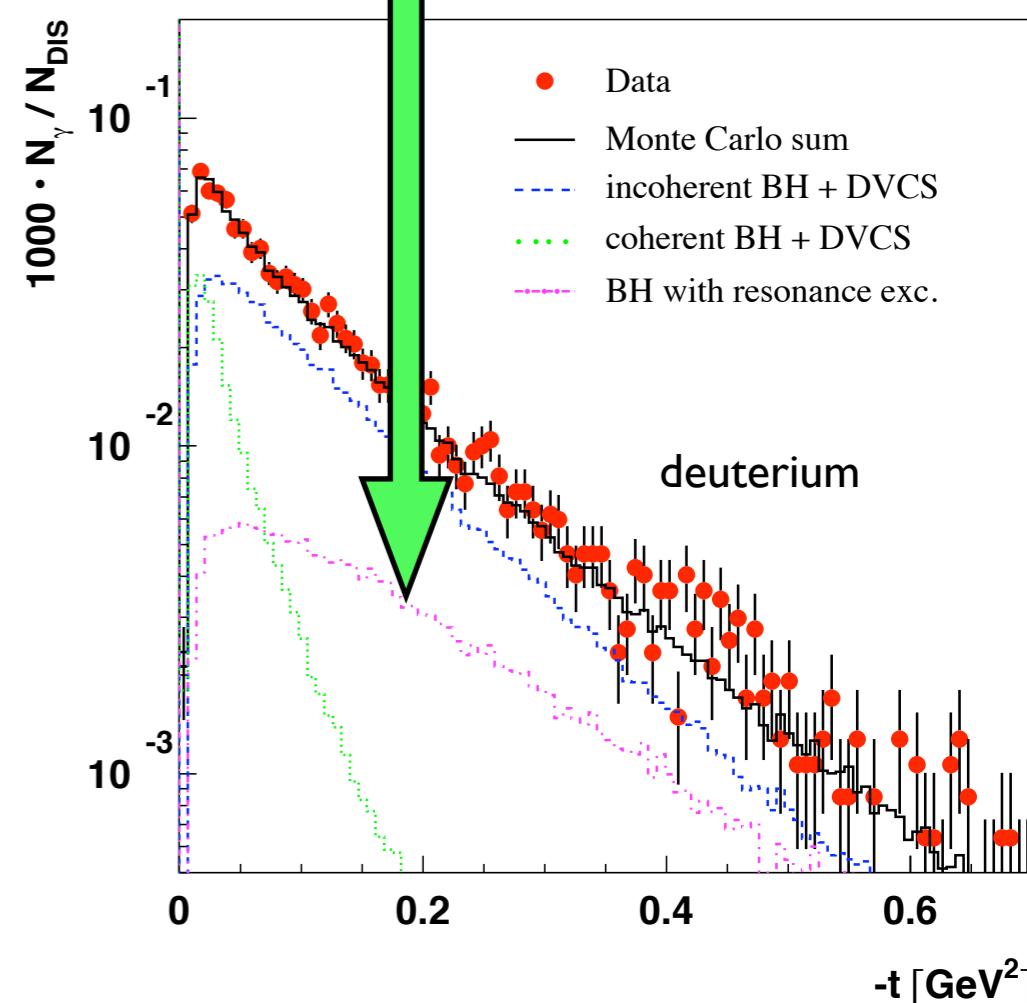
Deuteron: probe
spin-1 object



Nucleon: probe
spin-1/2 object

BH (proton) \gg BH (neutron)
due to elastic electric form factor

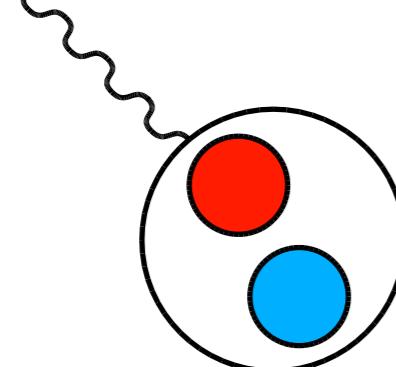
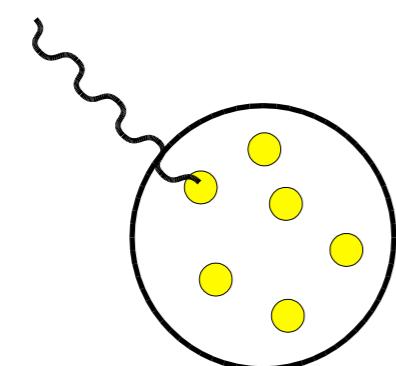
Coherent contribution
rapidly decreasing with $-t$



Coherent vs. incoherent

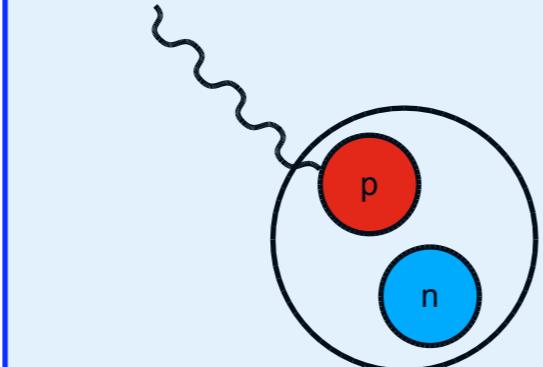
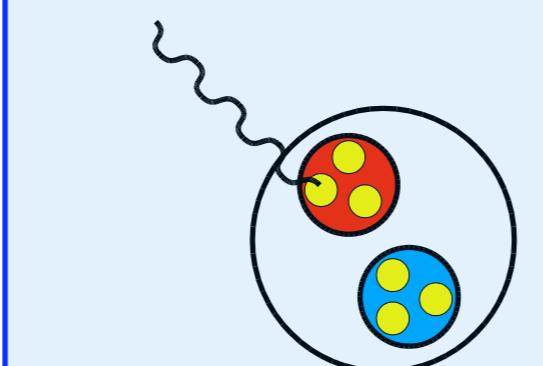
$$eA \rightarrow eA\gamma$$

DVCS



Deuteron: probe
spin-1 object

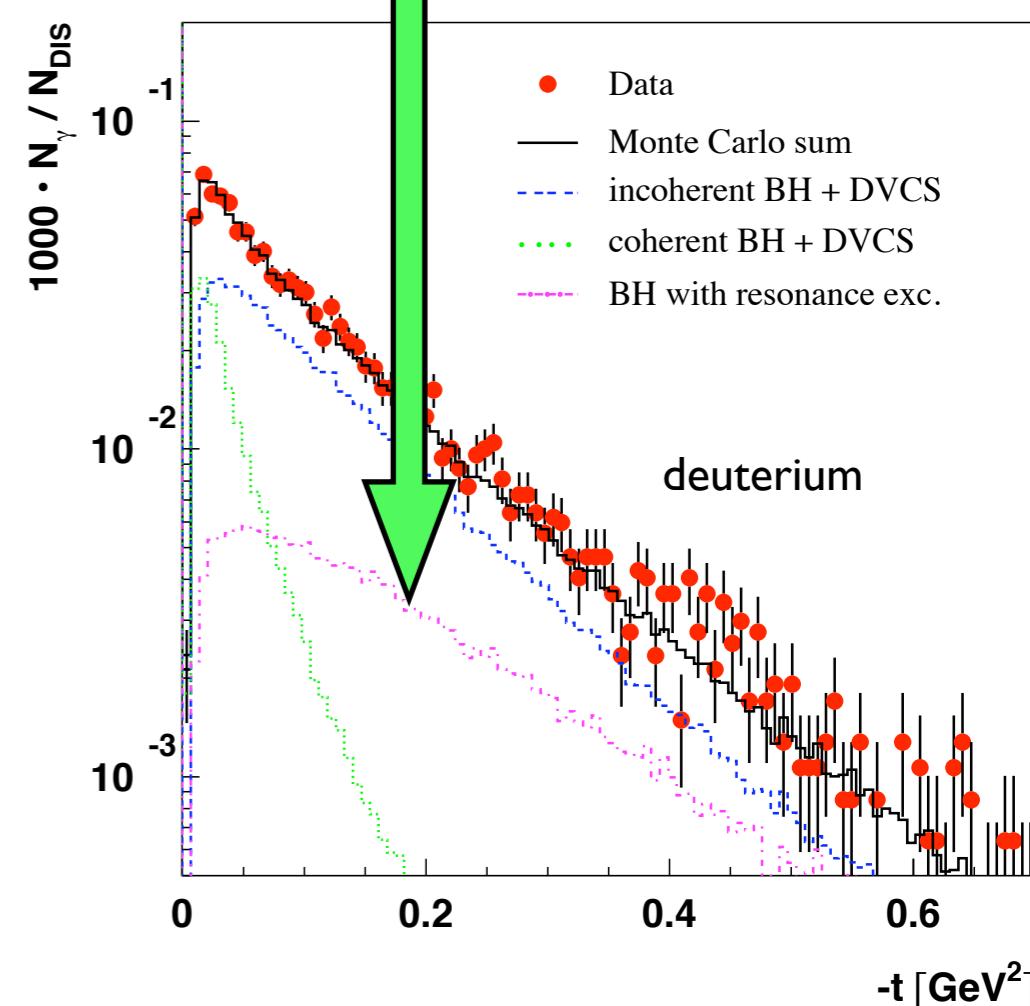
$$eA \rightarrow e(A-1)N\gamma$$



Nucleon: probe
spin-1/2 object

BH (proton) \gg BH (neutron)
due to elastic electric form factor

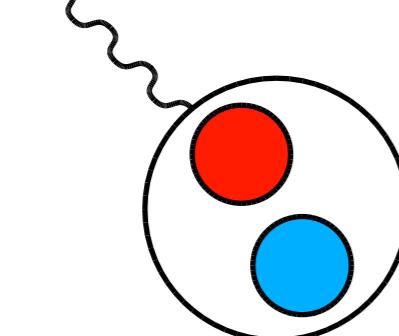
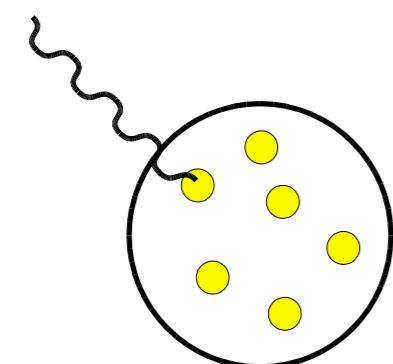
Coherent contribution
rapidly decreasing with $-t$



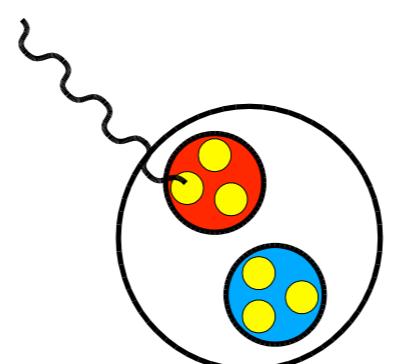
Coherent vs. incoherent

$$eA \rightarrow eA\gamma \quad eA \rightarrow e(A-1)N\gamma$$

DVCS



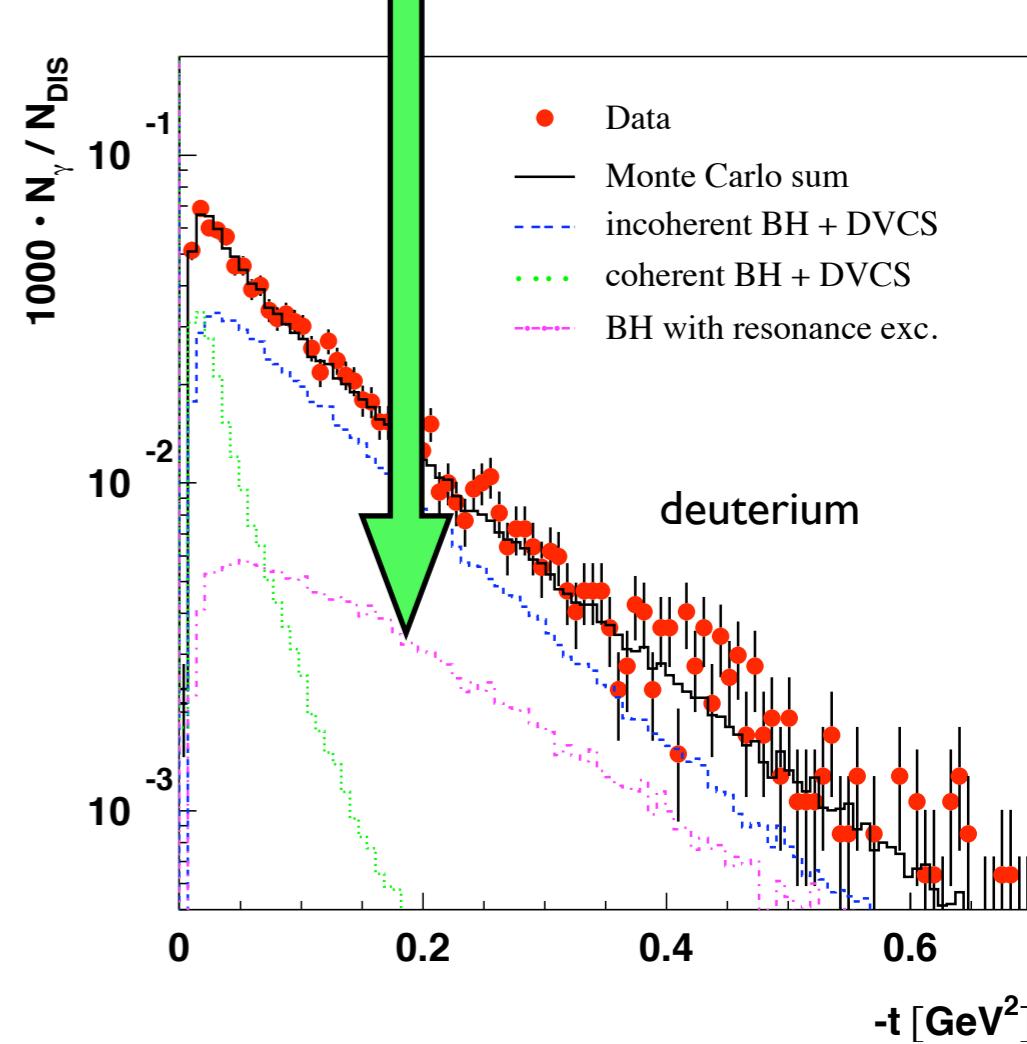
Deuteron: probe
spin-1 object



Nucleon: probe
spin-1/2 object

BH (proton) \gg BH (neutron)
due to elastic electric form factor

Coherent contribution
rapidly decreasing with $-t$



Asymmetries on polarized deuterons

	Lepton charge		Target population (deuterons)			Beam helicity		
	+1	-1	\Rightarrow	\Leftarrow	0	$\lambda = +1$	$\lambda = -1$	Coherent sensitivity
Single-charge	$\mathcal{A}_{L\uparrow\downarrow}$	■		■ + ■		■ - ■		$\Im m(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{UL}	■		■ - ■		■ + ■		$\Im m(\tilde{\mathcal{H}}_1)$
	\mathcal{A}_{LL}	■		■ - ■		■ - ■		(BH)
	\mathcal{A}_{Lzz}	■		■ + ■ - ■		■ - ■		$\Im m(\mathcal{H}_5)$
Single-helicity	$\mathcal{A}_{C\uparrow\downarrow}$	■ - ■		■ + ■			■	$\Im m/\Re e(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{o_L}	■ + ■		■ - ■			■	(BH)
	\mathcal{A}_{C_L}	■ - ■		■ - ■			■	$\Im m/\Re e(\tilde{\mathcal{H}}_1)$



HERMES data set available

Not all combinations of beam-charge and beam-helicity available!

Asymmetries on polarized deuterons

	Lepton charge		Target population (deuterons)			Beam helicity		
	+1	-1	$\Lambda = +1$	$\Lambda = -1$	$\Lambda = 0$	$\lambda = +1$	$\lambda = -1$	Coherent sensitivity
Single-charge	$\mathcal{A}_{L\uparrow\downarrow}$		■	+	■	■	-	■
	\mathcal{A}_{UL}		■	-	■	■	+	■
	\mathcal{A}_{LL}		■	-	■	■	-	■
	\mathcal{A}_{Lzz}		■	+	■	-	■	■
Single-helicity	$\mathcal{A}_{C\uparrow\downarrow}$		■	-	■			■
	\mathcal{A}_{o_L}		■	+	■			■
	\mathcal{A}_{C_L}		■	-	■			■

■ HERMES data set available

Not all combinations of beam-charge and beam-helicity available!

Asymmetries on polarized deuterons

	Lepton charge		Target population (deuterons)			Beam helicity		
	+1	-1	\Rightarrow	\Leftarrow	0	$\lambda = +1$	$\lambda = -1$	Coherent sensitivity
Single-charge	$\mathcal{A}_{L\uparrow\downarrow}$	■		■ + ■		■ - ■		$\Im m(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{UL}	■		■ - ■		■ + ■		$\Im m(\tilde{\mathcal{H}}_1)$
	\mathcal{A}_{LL}	■		■ - ■		■ - ■		(BH)
	\mathcal{A}_{Lzz}	■		■ + ■ - ■		■ - ■		$\Im m(\mathcal{H}_5)$
Single-helicity	$\mathcal{A}_{C\uparrow\downarrow}$	■ - ■		■ + ■			■	$\Im m/\Re e(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{o_L}	■ + ■		■ - ■			■	(BH)
	\mathcal{A}_{C_L}	■ - ■		■ - ■			■	$\Im m/\Re e(\tilde{\mathcal{H}}_1)$



HERMES data set available

Not all combinations of beam-charge and beam-helicity available!

Asymmetries on polarized deuterons

	Lepton charge		Target population (deuterons)			Beam helicity		
	+1	-1	\Rightarrow	\Leftarrow	0	$\lambda = +1$	$\lambda = -1$	Coherent sensitivity
Single-charge	$\mathcal{A}_{L\Leftarrow}$	■	■	+	■	■	—	■
	\mathcal{A}_{UL}	■	■	—	■	■	+	■
	\mathcal{A}_{LL}	■	■	—	■	■	—	■
	\mathcal{A}_{Lzz}	■	■	+	■	—	■	■
Single-helicity	$\mathcal{A}_{C\Updownarrow}$	■	—	■	■	■	—	■
	\mathcal{A}_{o_L}	■	+	■	■	■	—	■
	\mathcal{A}_{C_L}	■	—	■	■	■	—	■



HERMES data set available

Not all combinations of beam-charge and beam-helicity available!

Asymmetries on polarized deuterons

	Lepton charge		Target population (deuterons)			Beam helicity		
	+1	-1	\Rightarrow	\Leftarrow	0	$\lambda = +1$	$\lambda = -1$	Coherent sensitivity
Single-charge	$\mathcal{A}_{L\uparrow\downarrow}$	■		■ + ■		■ - ■		$\Im m(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{UL}	■		■ - ■		■ + ■		$\Im m(\tilde{\mathcal{H}}_1)$
	\mathcal{A}_{LL}	■		■ - ■		■ - ■		(BH)
	\mathcal{A}_{Lzz}	■		■ + ■ - ■		■ - ■		$\Im m(\mathcal{H}_5)$
Single-helicity	$\mathcal{A}_{C\uparrow\downarrow}$	■ - ■		■ + ■			■	$\Im m/\Re e(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{o_L}	■ + ■		■ - ■			■	(BH)
	\mathcal{A}_{C_L}	■ - ■		■ - ■			■	$\Im m/\Re e(\tilde{\mathcal{H}}_1)$

■ HERMES data set available

Not all combinations of beam-charge and beam-helicity available!

Asymmetries on polarized deuterons

	Lepton charge		Target population (deuterons)			Beam helicity		
	+1	-1	\Rightarrow	\Leftarrow	0	$\lambda = +1$	$\lambda = -1$	Coherent sensitivity
Single-charge	$\mathcal{A}_{L\uparrow\downarrow}$	■		■ + ■		■ - ■	■ - ■	$\Im m(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{UL}	■		■ - ■		■ + ■	■ + ■	$\Im m(\tilde{\mathcal{H}}_1)$
	\mathcal{A}_{LL}	■		■ - ■		■ - ■	■ - ■	(BH)
	\mathcal{A}_{Lzz}	■		■ + ■ - ■		■ - ■	■ - ■	$\Im m(\mathcal{H}_5)$
Single-helicity	$\mathcal{A}_{C\uparrow\downarrow}$	■ - ■		■ + ■		■	■	$\Im m/\Re e(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{o_L}	■ + ■		■ - ■		■	■	(BH)
	\mathcal{A}_{C_L}	■ - ■		■ - ■		■	■	$\Im m/\Re e(\tilde{\mathcal{H}}_1)$



HERMES data set available

Not all combinations of beam-charge and beam-helicity available!

Asymmetries on polarized deuterons

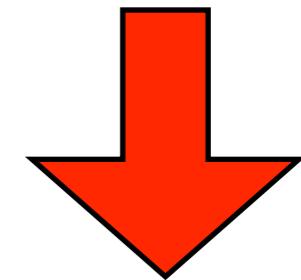
	Lepton charge		Target population (deuterons)			Beam helicity		
	+1	-1	\Rightarrow	\Leftarrow	0	$\lambda = +1$	$\lambda = -1$	Coherent sensitivity
Single-charge	$\mathcal{A}_{L\uparrow\downarrow}$	■		■ + ■		■ - ■		$\Im m(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{UL}	■		■ - ■		■ + ■		$\Im m(\tilde{\mathcal{H}}_1)$
	\mathcal{A}_{LL}	■		■ - ■		■ - ■		(BH)
	\mathcal{A}_{Lzz}	■		■ + ■ - ■		■ - ■		$\Im m(\mathcal{H}_5)$
Single-helicity	$\mathcal{A}_{C\uparrow\downarrow}$	■ - ■		■ + ■			■	$\Im m/\Re e(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{o_L}	■ + ■		■ - ■			■	(BH)
	\mathcal{A}_{C_L}	■ - ■		■ - ■			■	$\Im m/\Re e(\tilde{\mathcal{H}}_1)$



HERMES data set available

Not all combinations of beam-charge and beam-helicity available!

Asymmetries on polarized deuterons



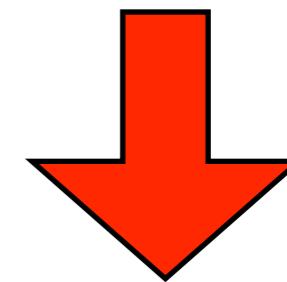
	Lepton charge		Target population (deuterons)			Beam helicity		
	+1	-1	\Rightarrow	\Leftarrow	0	$\lambda = +1$	$\lambda = -1$	Coherent sensitivity
Single-charge	$\mathcal{A}_{L\Leftarrow}$	■	■	+	■	■	—	■ $\Im m(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{UL}	■	■	—	■	■	+	■ $\Im m(\tilde{\mathcal{H}}_1)$
	\mathcal{A}_{LL}	■	■	—	■	■	—	■ (BH)
	\mathcal{A}_{Lzz}	■	■	+	■	—	■	■ $\Im m(\mathcal{H}_5)$
Single-helicity	$\mathcal{A}_{C\Leftarrow}$	■	—	■	■	+	■	■ $\Im m/\Re e(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{o_L}	■	+	■	■	—	■	■ (BH)
	\mathcal{A}_{C_L}	■	—	■	■	—	■	■ $\Im m/\Re e(\tilde{\mathcal{H}}_1)$



HERMES data set available

Not all combinations of beam-charge and beam-helicity available!

Asymmetries on polarized deuterons



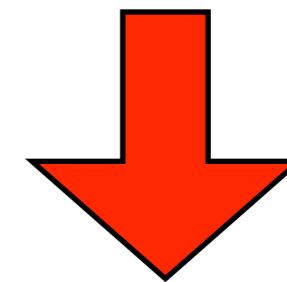
	Lepton charge		Target population (deuterons)			Beam helicity		
	$\Lambda = +1$	$\Lambda = -1$	\Rightarrow	\Leftarrow	0	$\lambda = +1$	$\lambda = -1$	Coherent sensitivity
Single-charge	$\mathcal{A}_{L\Leftarrow}$	■	■	+	■	■	—	■
	\mathcal{A}_{UL}	■	■	—	■	■	+	■
	\mathcal{A}_{LL}	■	■	—	■	■	—	■
	\mathcal{A}_{Lzz}	■	■	+	■	—	■	■
Single-helicity	$\mathcal{A}_{C\Leftarrow}$	■	—	■	■	■	—	■
	\mathcal{A}_{o_L}	■	+	■	■	■	—	■
	\mathcal{A}_{C_L}	■	—	■	■	■	—	■



HERMES data set available

Not all combinations of beam-charge and beam-helicity available!

Asymmetries on polarized deuterons



	Lepton charge		Target population (deuterons)			Beam helicity		
	$\Lambda = +1$	$\Lambda = -1$	\Rightarrow	\Leftarrow	0	$\lambda = +1$	$\lambda = -1$	Coherent sensitivity
Single-charge	$\mathcal{A}_{L\Leftarrow}$	■	■	+	■	■	—	■
	\mathcal{A}_{UL}	■	■	—	■	■	+	■
	\mathcal{A}_{LL}	■	■	—	■	■	—	■
	\mathcal{A}_{Lzz}	■	■	+	■	—	■	■
Single-helicity	$\mathcal{A}_{C\Updownarrow}$	■	—	■	■	■	■	$\Im m(\mathcal{H}_1, \mathcal{H}_5)$
	\mathcal{A}_{o_L}	■	+	■	■	■	■	$\Im m(\tilde{\mathcal{H}}_1)$
	\mathcal{A}_{C_L}	■	—	■	■	■	■	(BH)

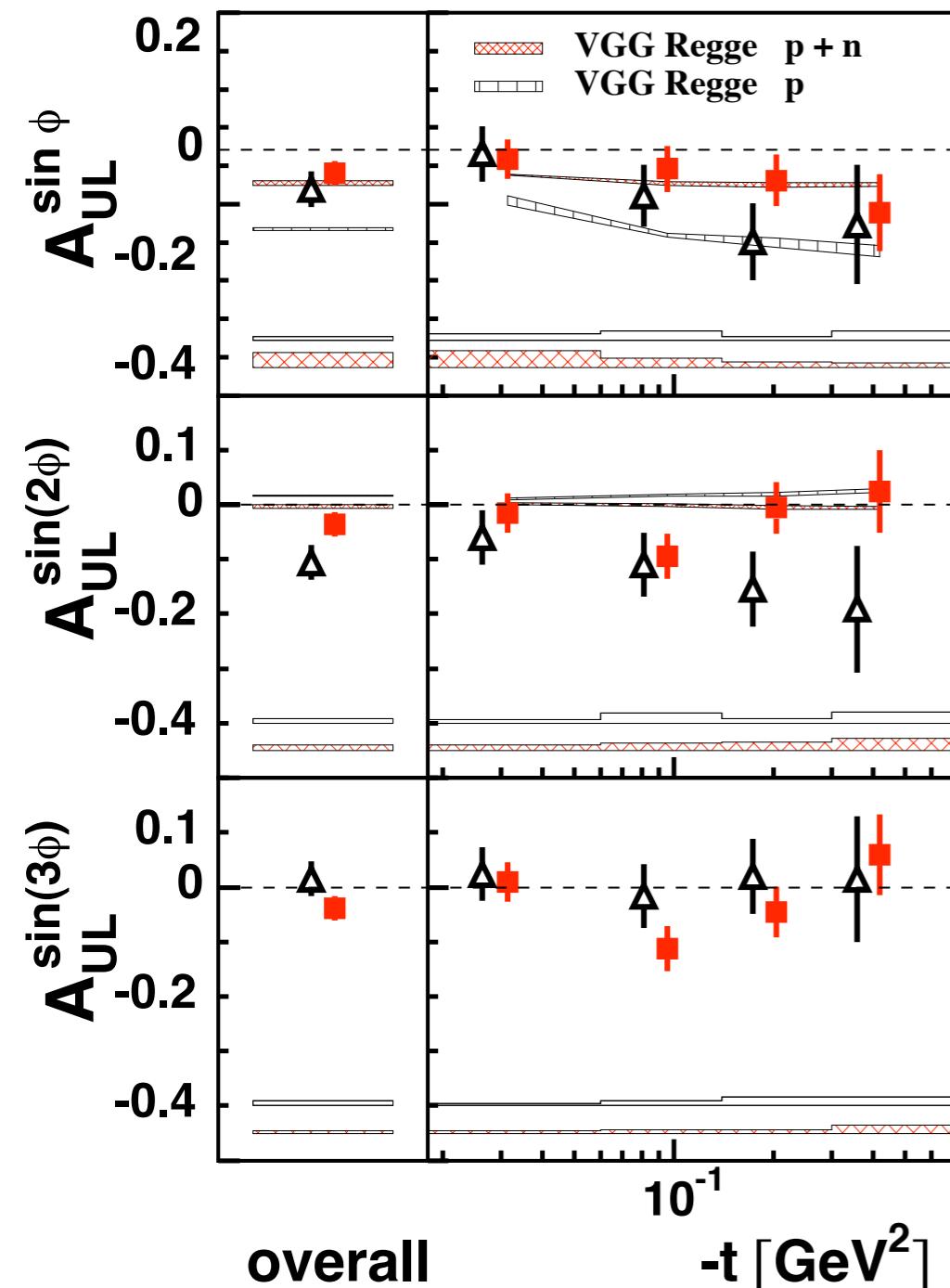


HERMES data set available

Not all combinations of beam-charge and beam-helicity available!

Target-spin asymmetries on
the deuteron (spin-1)
in comparison to
the proton (spin-1/2)

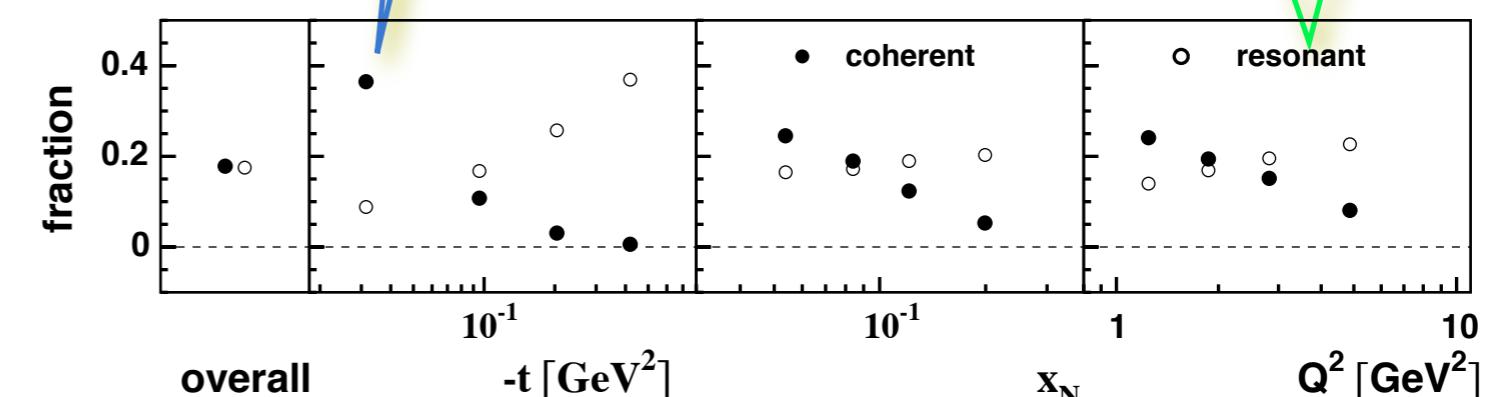
HERMES \mathcal{A}_{UL} on 1H and 2H



Coherent contribution
dominates at low values of the
squared momentum transfer

Search for
coherent
signature

Fraction of resonant
excitation stays part
of the signal



low t
scatter off deuteron
= spin-1 object

higher t
scatter off nucleon
= spin-1/2 object

$$\text{Im}(\tilde{\mathcal{H}}_1)$$

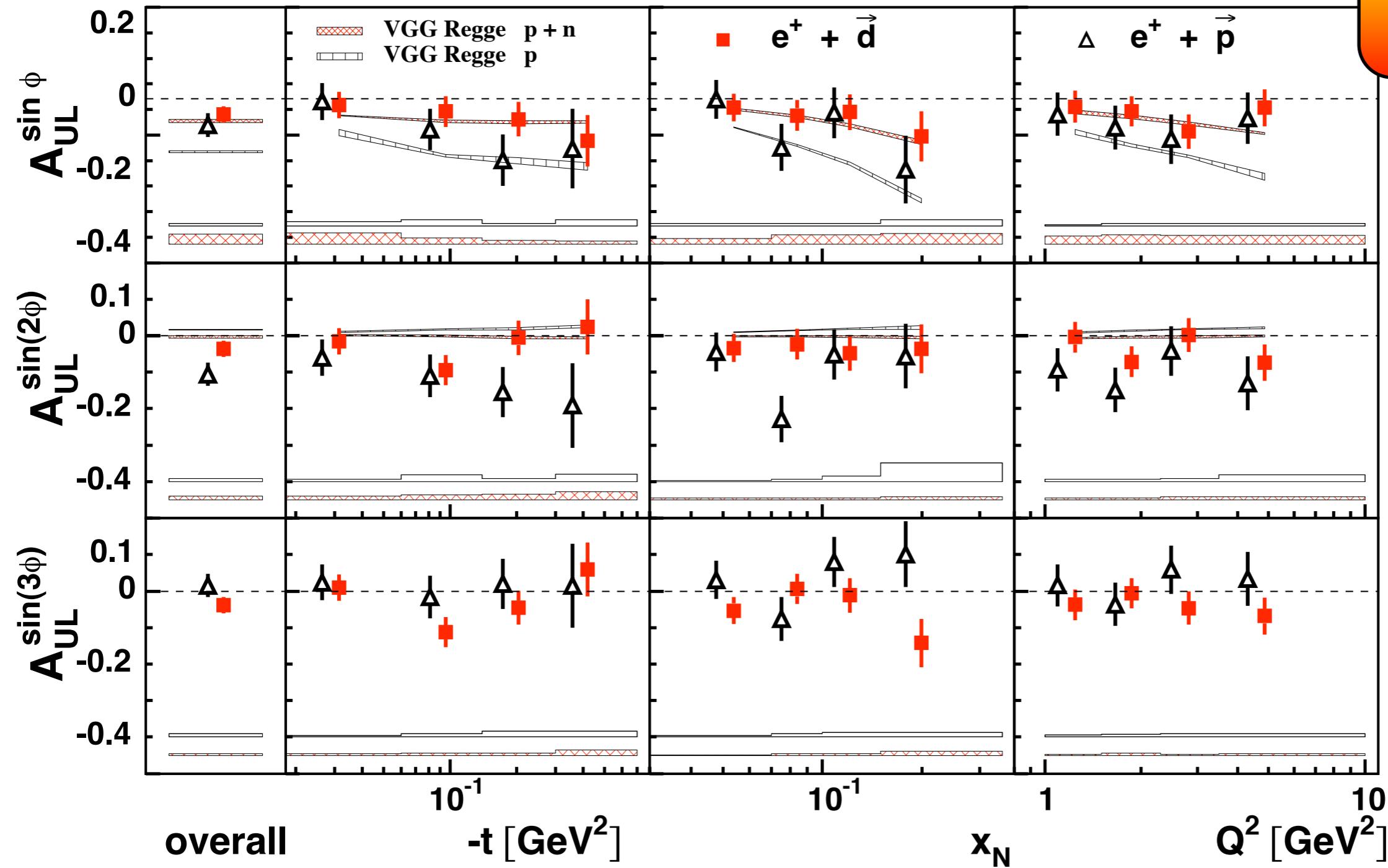
$$\text{Im}(\tilde{\mathcal{H}})$$

HERMES \mathcal{A}_{UL} on 1H and 2H

Accepted by
Nucl. Phys. B

arXiv:1008.3996 (hep-ex)

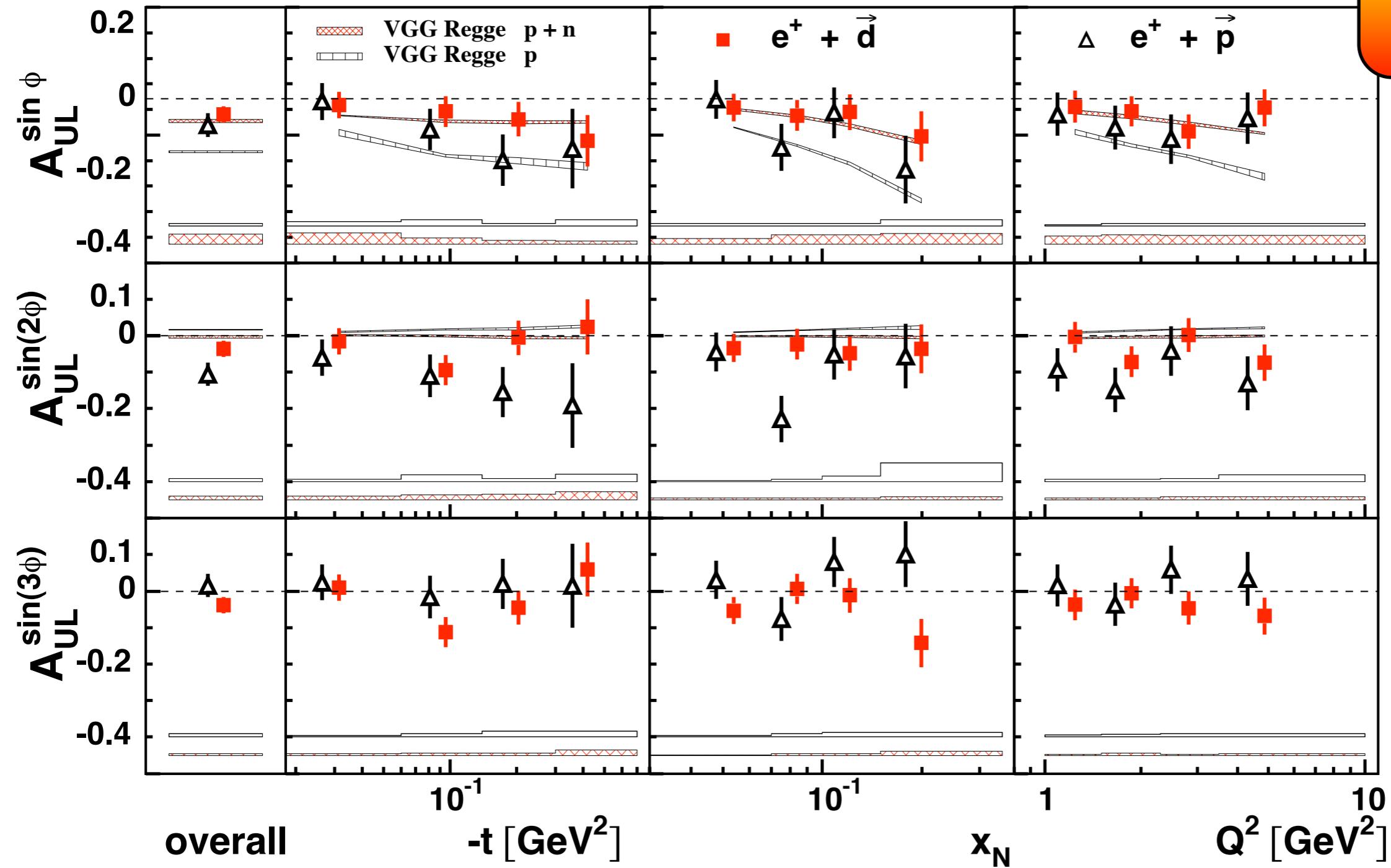
Search for
coherent
signature



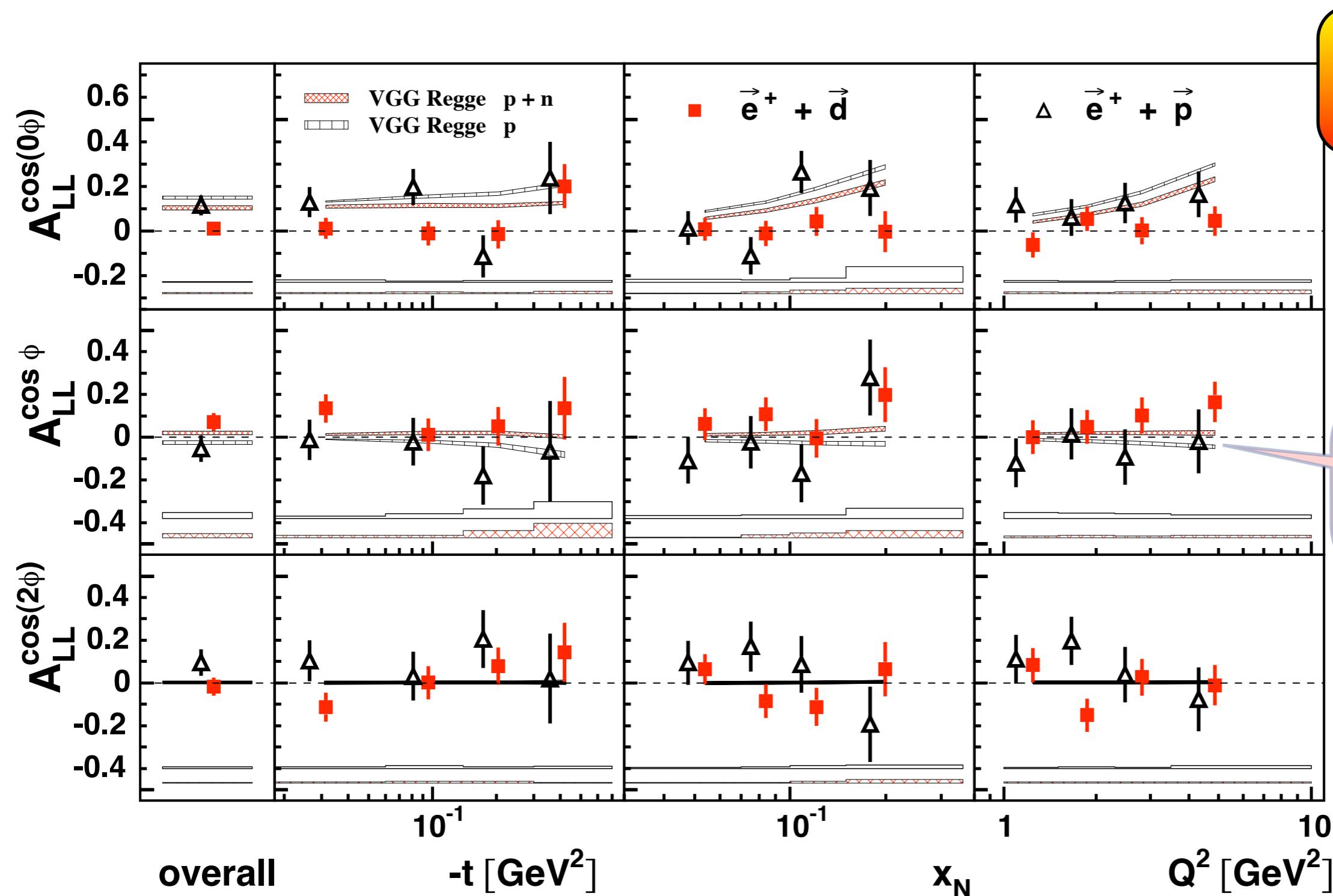
HERMES \mathcal{A}_{UL} on 1H and 2H

Accepted by
Nucl. Phys. B
arXiv:1008.3996 (hep-ex)

Search for
coherent
signature



HERMES \mathcal{A}_{LL} on ^1H and ^2H



@low t , scatter off deuteron
= spin 1 object: $\text{Re}(\tilde{\mathcal{H}}_1)$

@higher t , scatter off nucleon
= spin 1/2 object: $\text{Re}(\tilde{\mathcal{H}})$

**Search for
coherent
signature**

VGG:
Phys.Rev. **D60** (1999)
094017 and
Prog.Nucl.Phys. **47**
(2001) 401

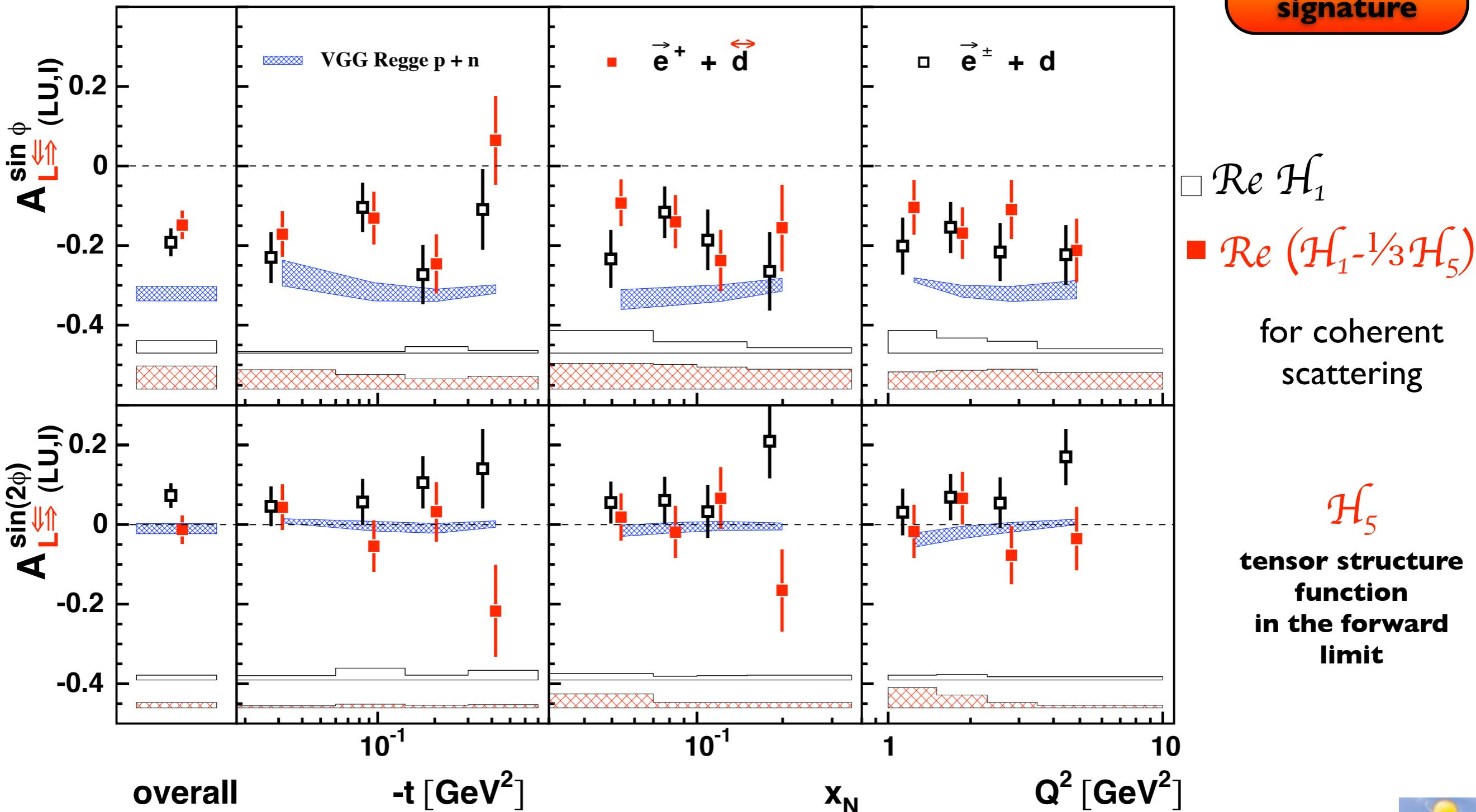
Beam-helicity and
beam-charge asymmetries
on tensor-polarized
deuterons with $P_{zz}=0.827$
in comparision to
unpolarized deuterons

HERMES \mathcal{A}_{LU} on (un)polarized ${}^2\text{H}$

Accepted by
Nucl. Phys. B

arXiv:1008.3996 (hep-ex)

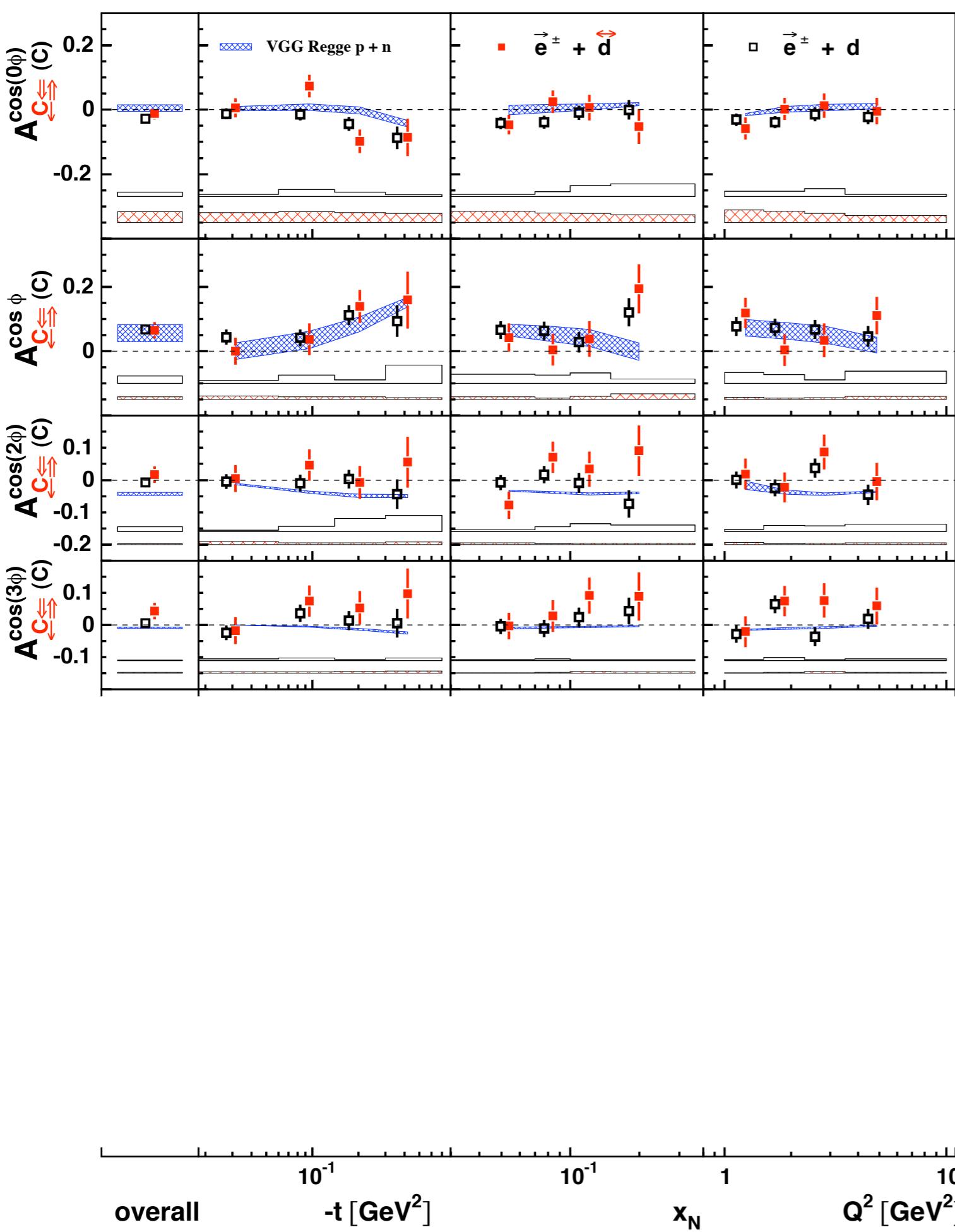
Search for
tensor
signature



HERMES \mathcal{A}_C on (un)polarized ^2H

**Search for
tensor
signature**

- $\mathcal{R}e \mathcal{H}_1$
 - $\mathcal{R}e (\mathcal{H}_1^{-1/3} \mathcal{H}_5)$
- for coherent
scattering

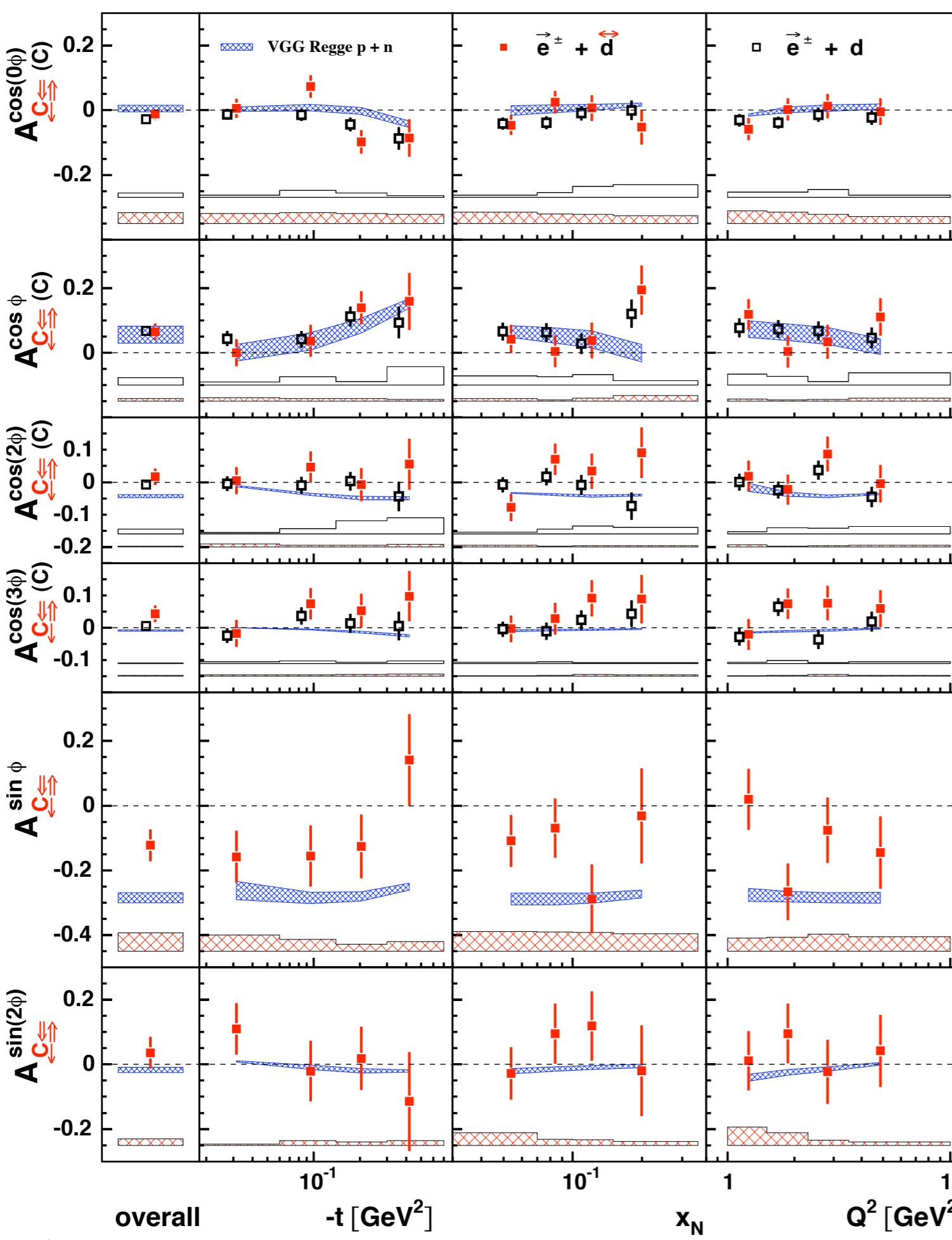


Accepted by
Nucl. Phys. B

arXiv:1008.3996 (hep-ex)



HERMES \mathcal{A}_C on (un)polarized ^2H



Search for
tensor
signature

$\square \mathcal{R}e \mathcal{H}_1$
 $\blacksquare \mathcal{R}e (\mathcal{H}_1^{-1/3} \mathcal{H}_5)$

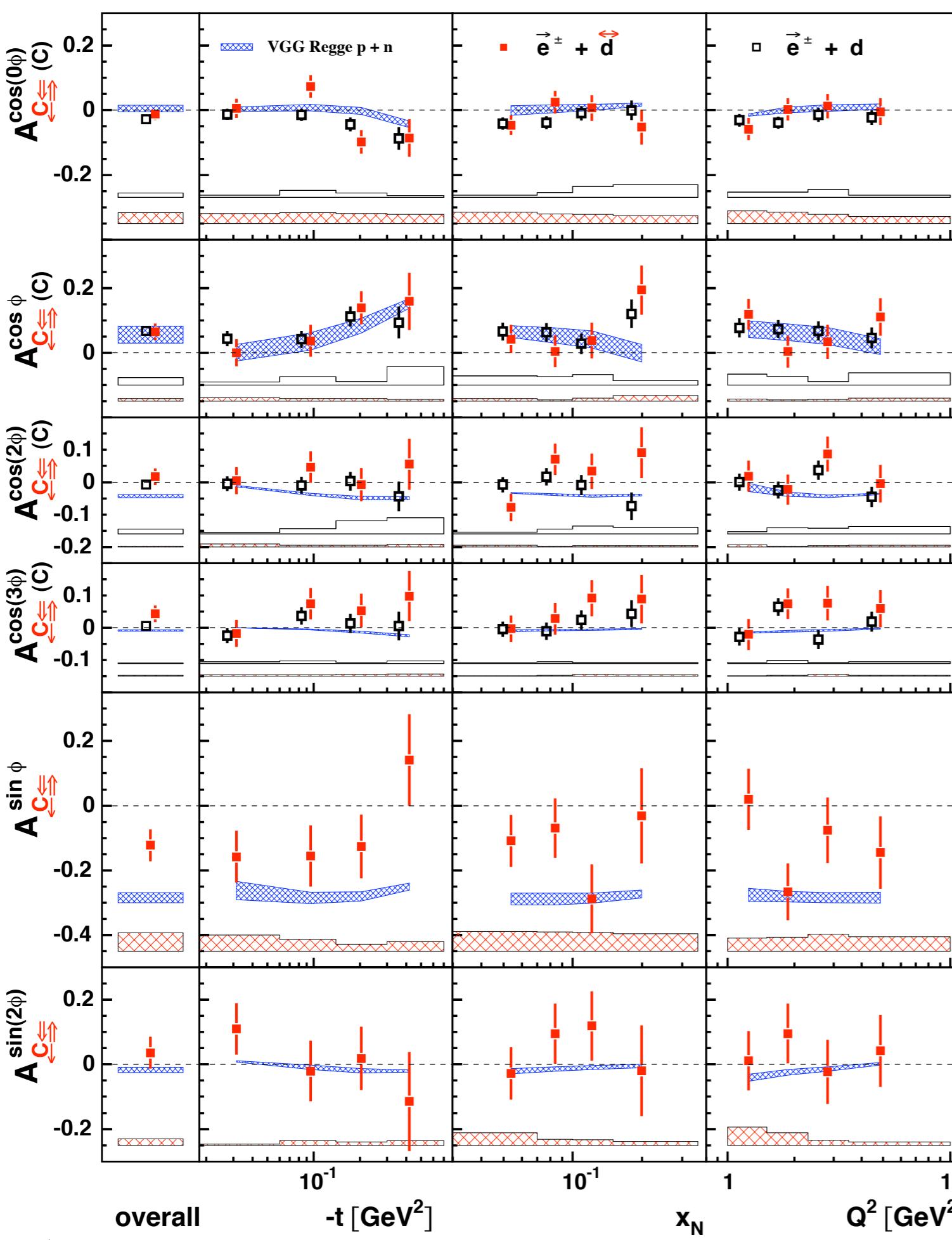
for coherent
scattering

$\blacksquare \mathcal{I}m (\mathcal{H}_1^{-1/3} \mathcal{H}_5)$

Accepted by
Nucl. Phys. B

arXiv:1008.3996 (hep-ex)

HERMES \mathcal{A}_C on (un)polarized ^2H



Search for tensor signature

□ $\Re \mathcal{H}_1$
■ $\Re (\mathcal{H}_1 - \frac{1}{3}\mathcal{H}_5)$

for coherent scattering

■ $\Im (\mathcal{H}_1 - \frac{1}{3}\mathcal{H}_5)$

DVCS $\mathcal{A}_{LZZ} \sin\phi$ amplitude:
 $0.074 \pm 0.196 \pm 0.022$
($-t < 0.06 \text{ GeV}^2$, 40% coherent)

**Accepted by
Nucl. Phys. B**

arXiv:1008.3996 (hep-ex)



Beam-spin and beam-charge asymmetries on heavier nuclei

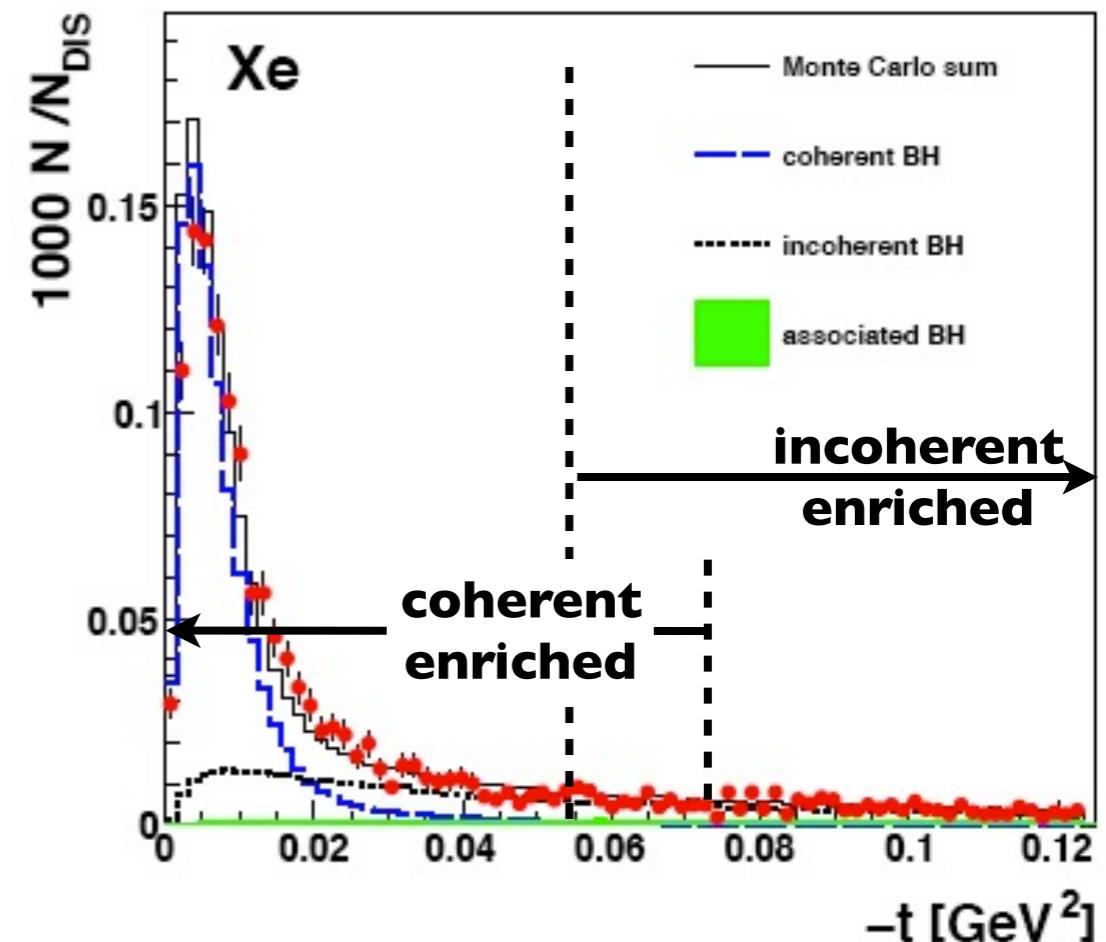
Nuclear data sets

HERMES
Phys. Rev. C 81 (2010)
035202
arXiv:0911.0091

Target	Spin	L (pb^{-1})
^1H	1/2	227
He	0	32
N	1	51
Ne	0	86
Kr	0	77
Xe	0, 1/2, 3/2	47

Heavy target data taken at
the end of each HERA fill
("high density runs")

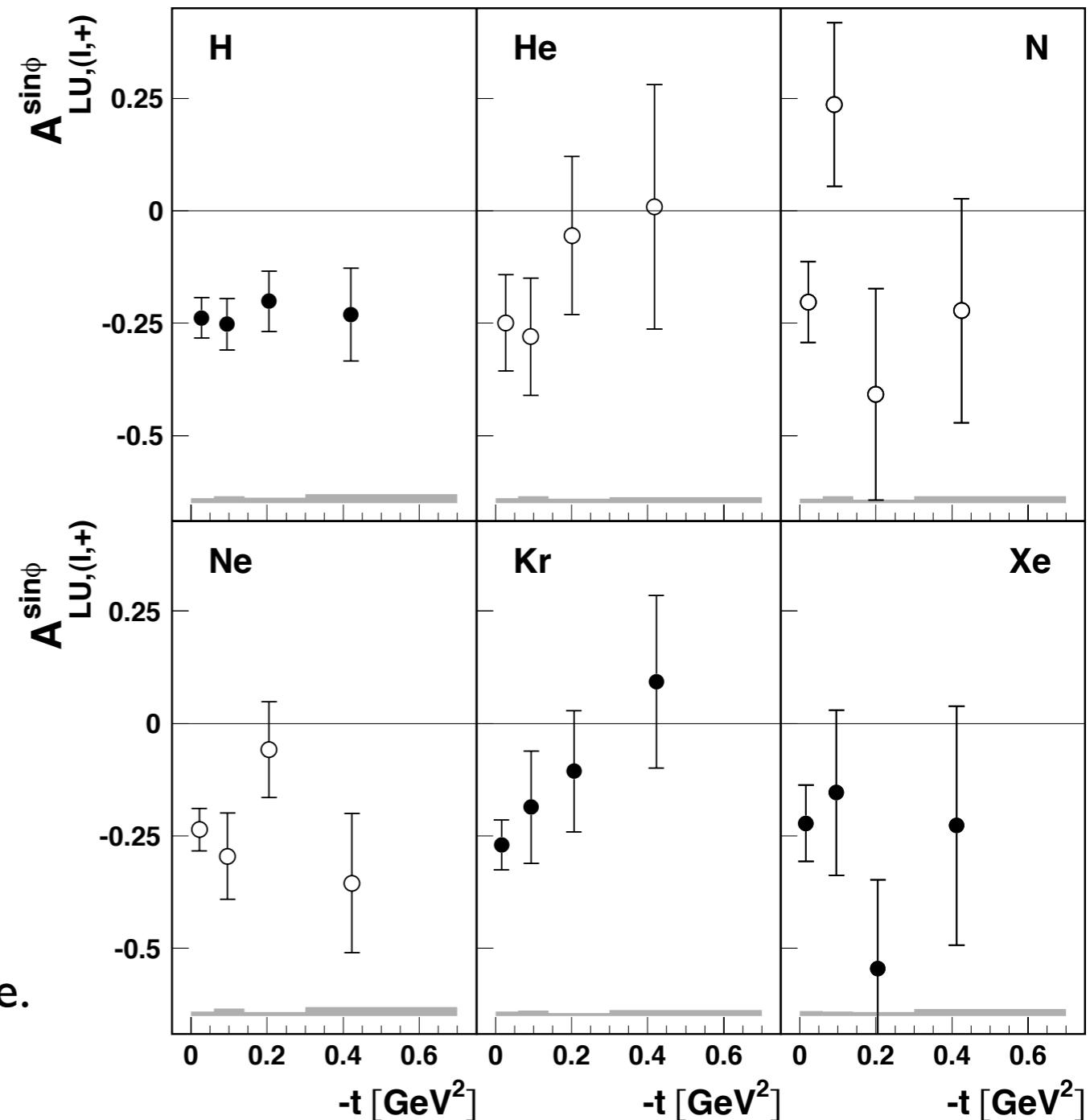
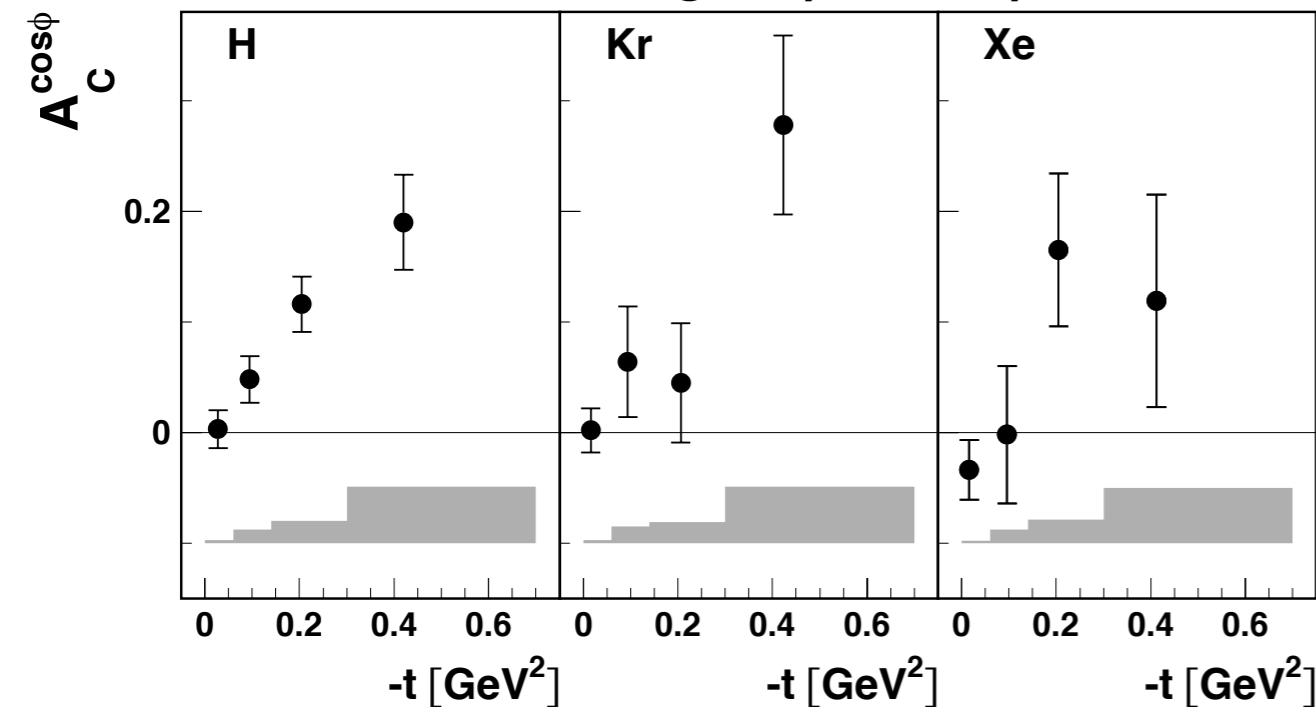
- Separation of coherent-enriched and incoherent-enriched data samples by t-cutoff such that \approx same average kinematics for each target.
- Coherent enriched samples: $\approx 65\%$ coherent fraction
- Incoherent enriched samples: $\approx 60\%$ incoherent fraction



DVCS asymmetries on Nuclei

Leading amplitude of
beam-helicity asymmetry

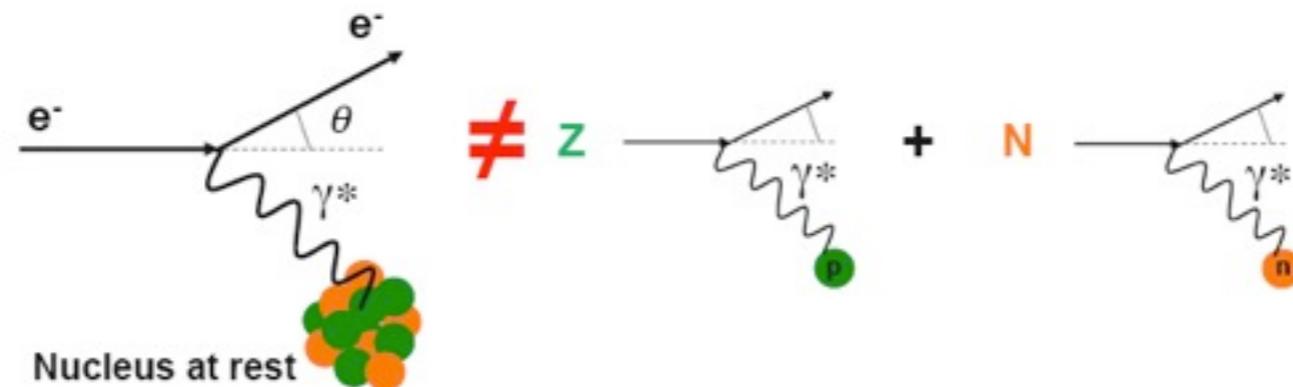
Leading amplitude of
beam-charge asymmetry



- Targets with 2 beam charges available.
 A_C and charge-difference A_{LU} sensitive to DVCS-BH interference term
- Targets with only one beam charge available.
No A_C and single-charge A_{LU} with entangled s_1 coefficients

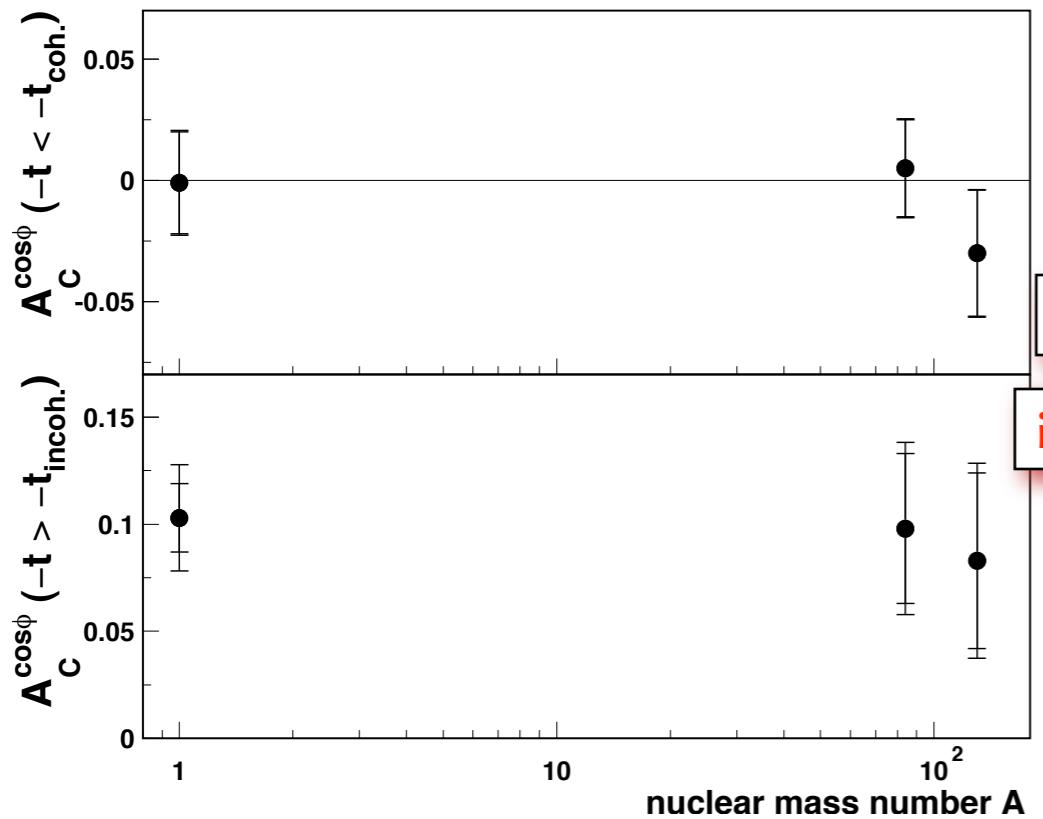
A-dependence of DVCS asymmetries

- How does the nuclear medium modify parton-parton correlations?
- How do the nucleon properties change in the nuclear medium?
- Is there an enhanced ‘generalized EMC effect’, which could be revealed through the rise if τ_{DVCS} with A ?



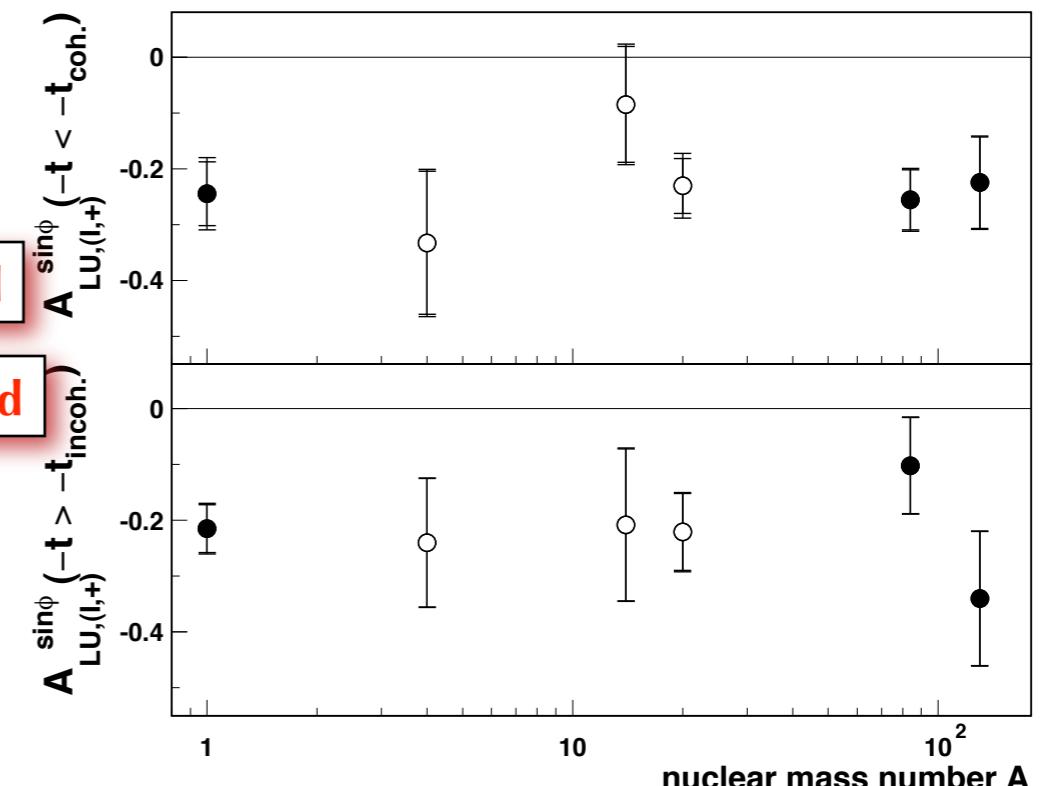
DVCS Nuclear Mass Dependence

$A_C^{\cos\phi}$ vs. A



Beam-charge asymmetry

$A_{LU}^{\sin\phi}$ vs. A



Beam-helicity asymmetry

HERMES
Phys. Rev. C 81
(2010) 035202
arXiv:0911.0091

Normalization to
hydrogen ${}^1\text{H} \Rightarrow$

Average A_{LU}^A / A_{LU}^H :
coherent enriched: 0.91 ± 0.19
incoherent enriched: 0.93 ± 0.23

HERMES results for spin-1/2 GPDs

See talk by M. Düren

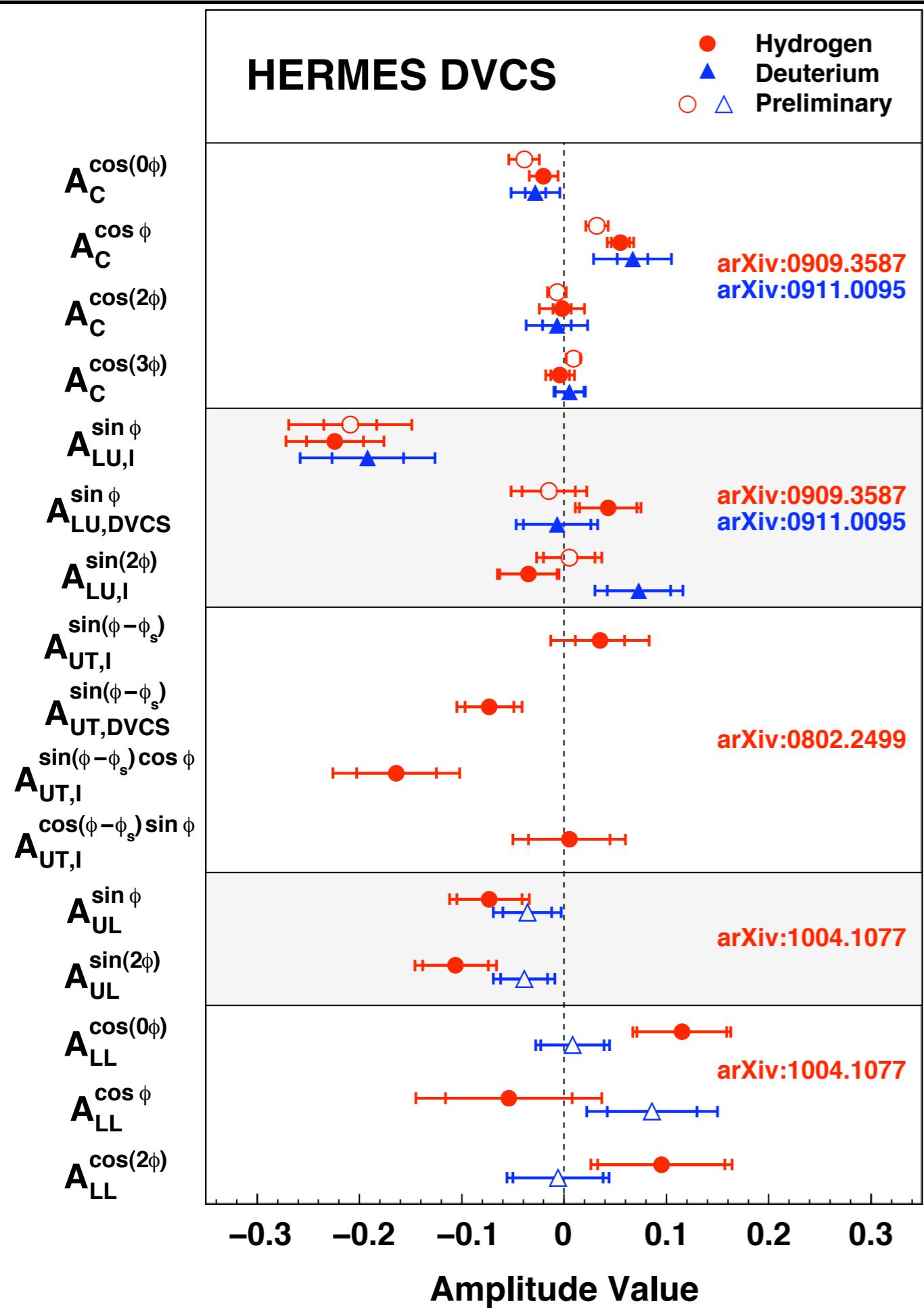
(A) Beam charge asymmetry:
GPD H

(B) Beam helicity asymmetry:
GPD H

(C) Transverse target spin asymmetry:
GPD E from proton target

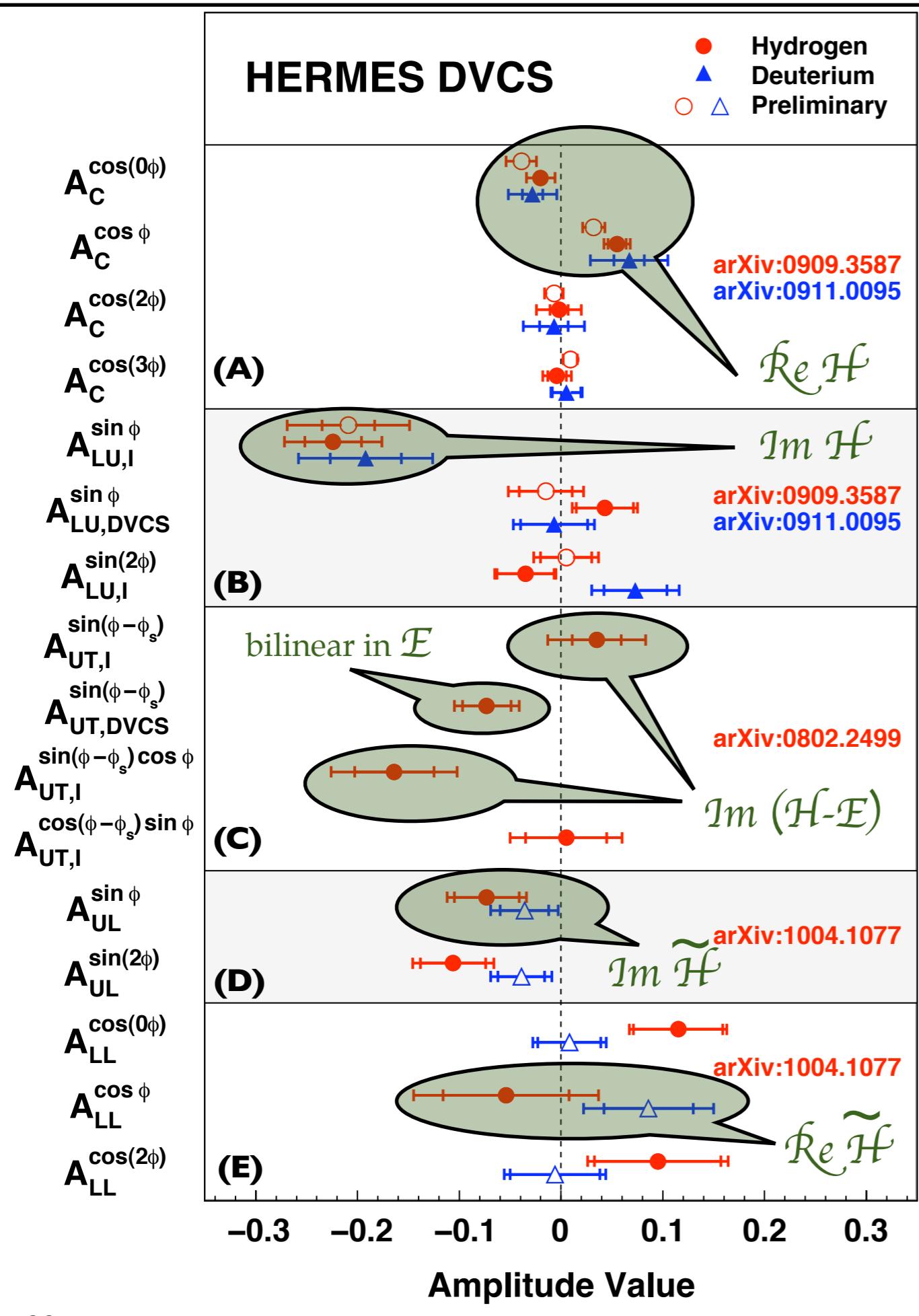
(D) Longitudinal target spin asymmetry:
GPD H

(E) Double-spin asymmetry:
GPD H



HERMES results for spin-1/2 GPDs

See talk by M. Düren



Summary: DVCS on nuclear targets at HERMES

- HERMES has analyzed its rich set of DVCS data on nuclear targets (Deuterium, Helium, Neon, Nitrogen, Krypton, Xenon):
➤beam-helicity, beam-charge, target-spin and double-spin asymmetries.
Final results. **Unique data set!**
- These measurements allow for the search for
 - the coherent signature expected for a deuteron target at low t .
 - the tensor signature expected for tensor-polarized deuterons.No obvious signatures revealed within the experimental uncertainties.
- No A-dependence of the DVCS asymmetries was found within the experimental uncertainties.
- Measure these observables with greater precision at an EIC

Backup slides

Systematics and corrections

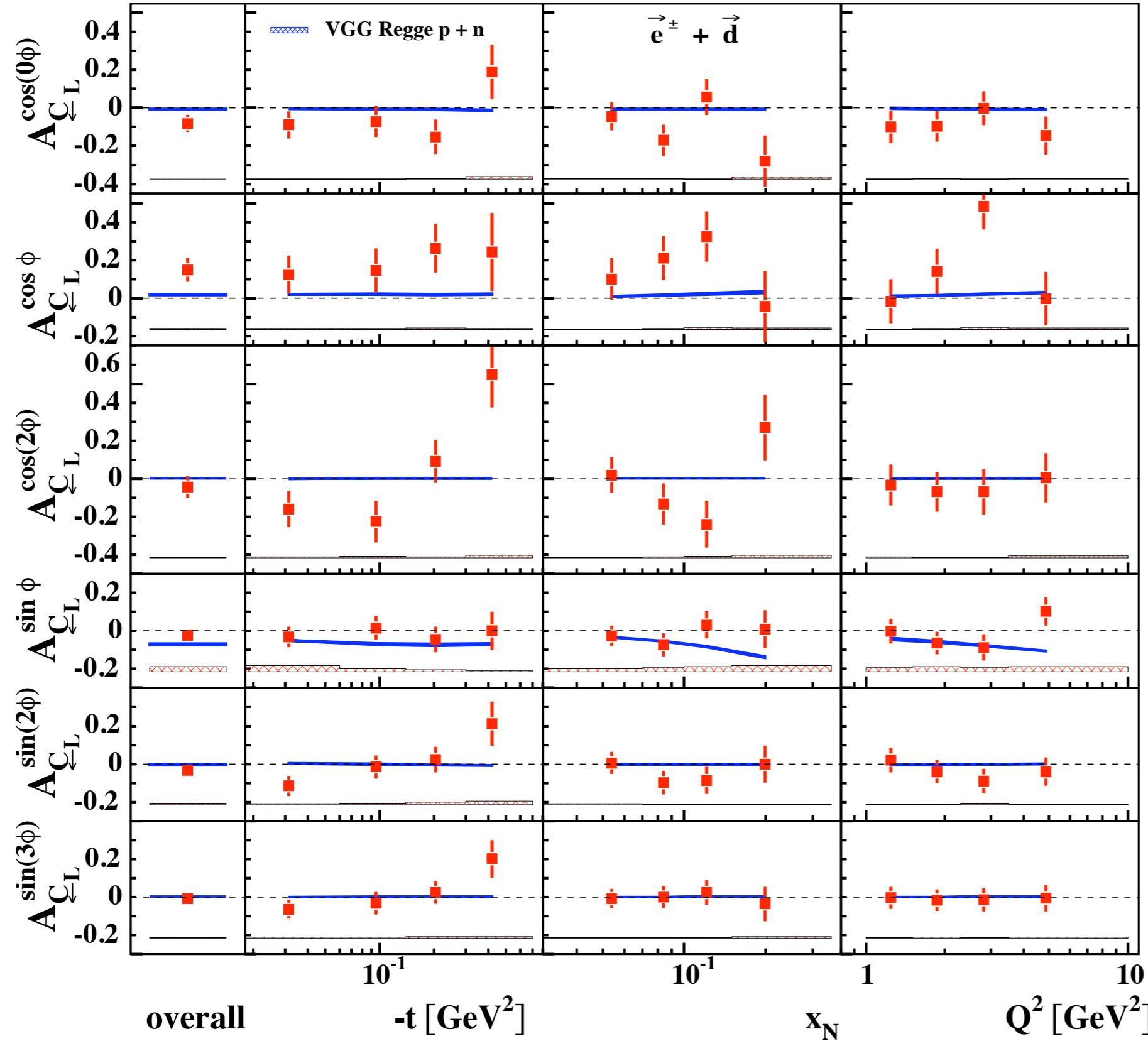
- Results corrected for background from decays of neutral mesons (semi-inclusive: 4.6%, exclusive: <0.7%)
- Dominant contribution to systematics: limited spectrometer acceptance and finite bin widths.
 - Combined contribution from acceptance, finite bin width, detector misalignment and kinematic smearing determined from MC simulation.
- Scale uncertainties due to polarization measurement:
 - deuterium target polarization: 4.0%
 - beam polarization, unpolarized deuterium: 2.4%, polarized deuterium: 1.9%

HERMES \mathcal{A}_{CL}
 (single helicity)
 on polarized
 2H

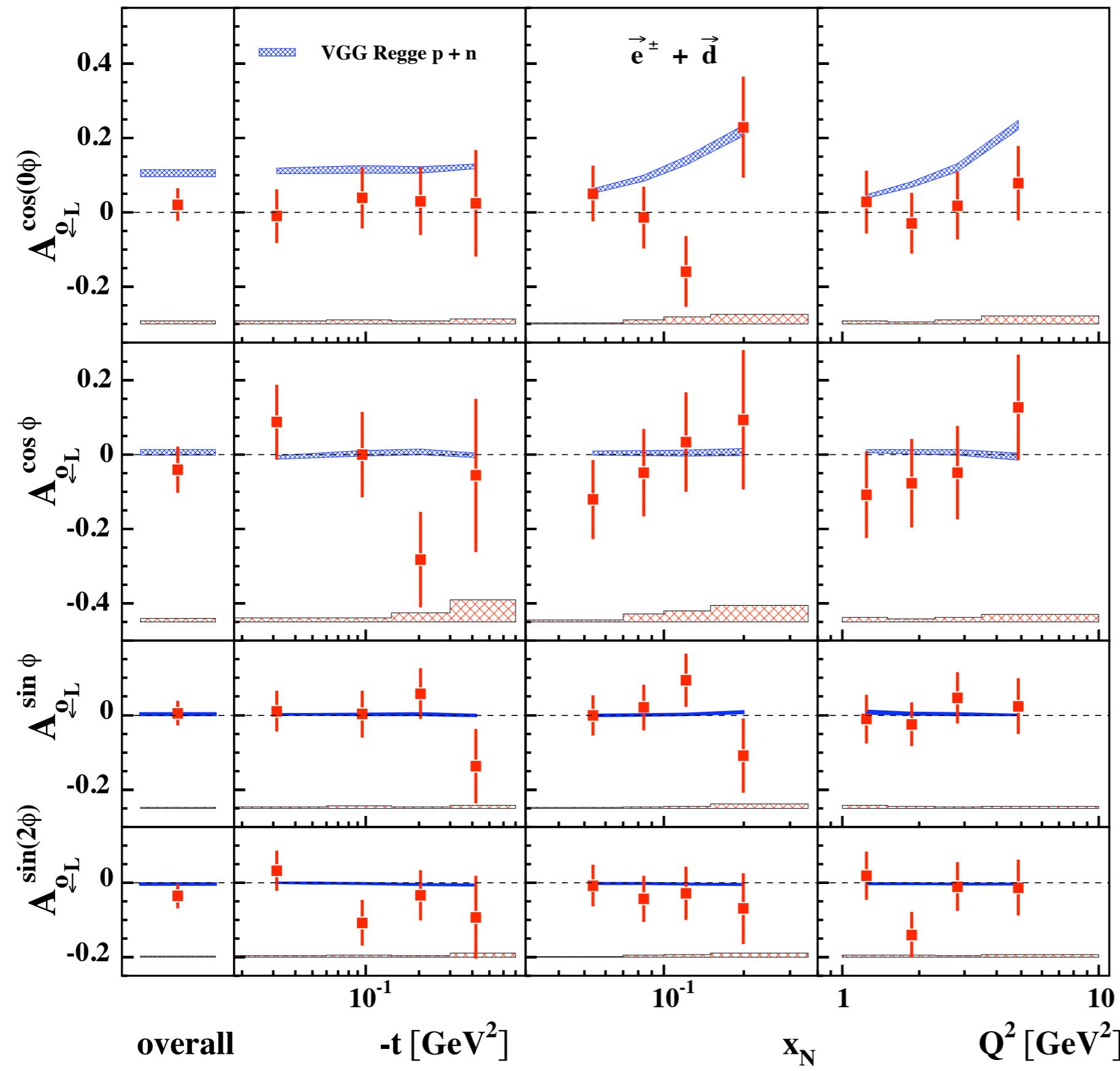
Re/Im($\tilde{\mathcal{H}}_1$)

Accepted by
 Nucl. Phys. B

arXiv:1008.3996 (hep-ex)



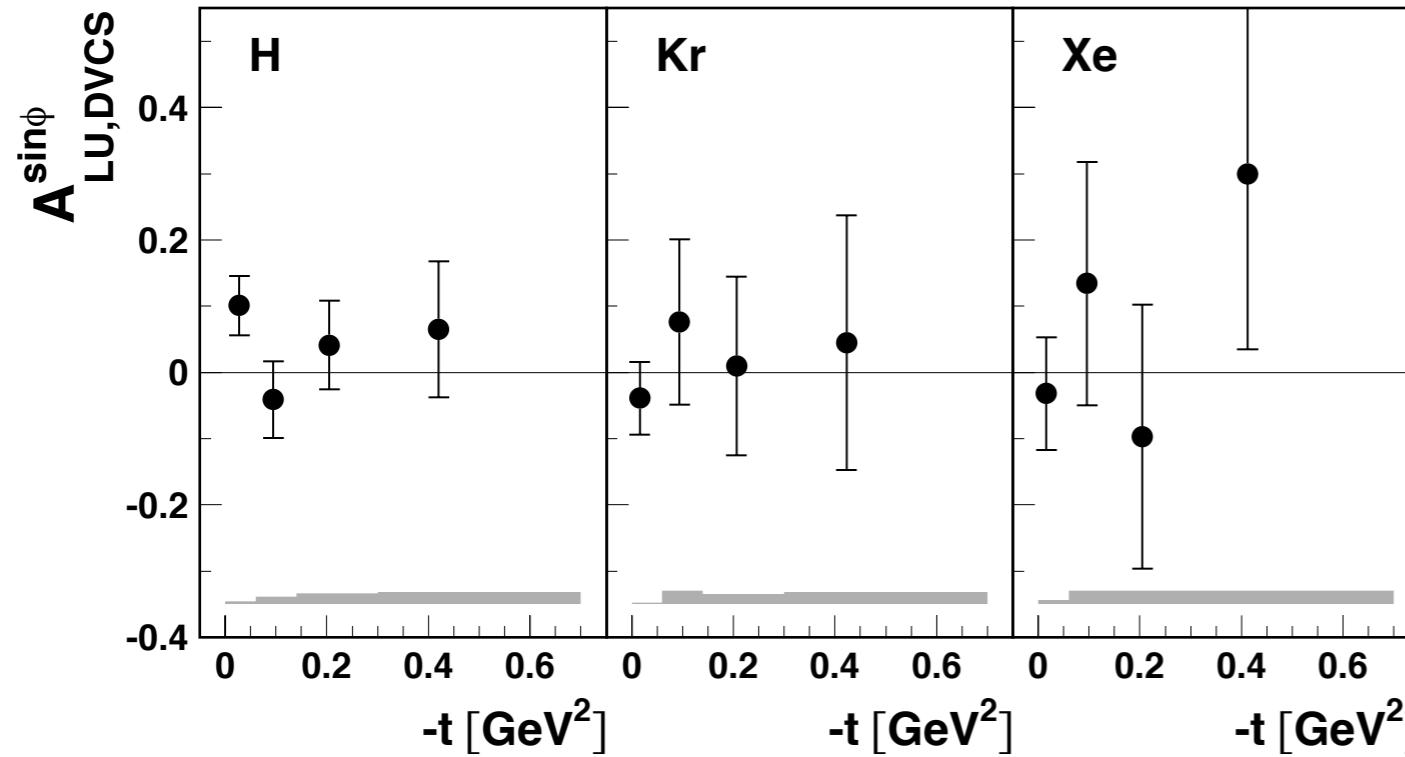
HERMES \mathcal{A}_{UL}
 (single helicity)
 on polarized
 ^2H



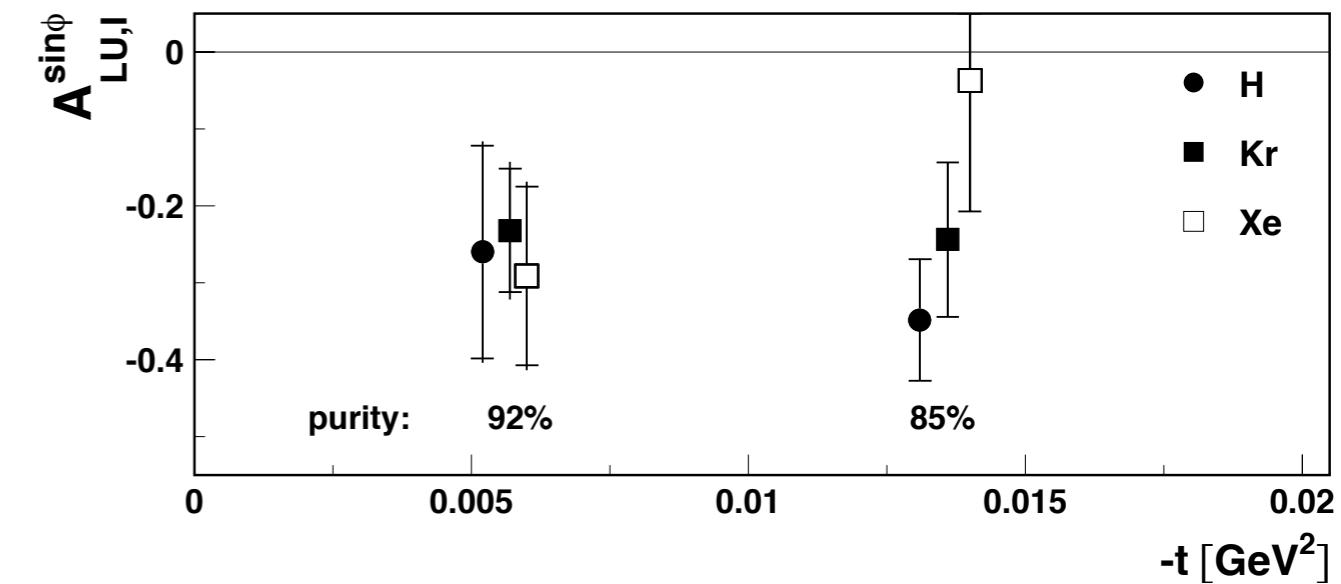
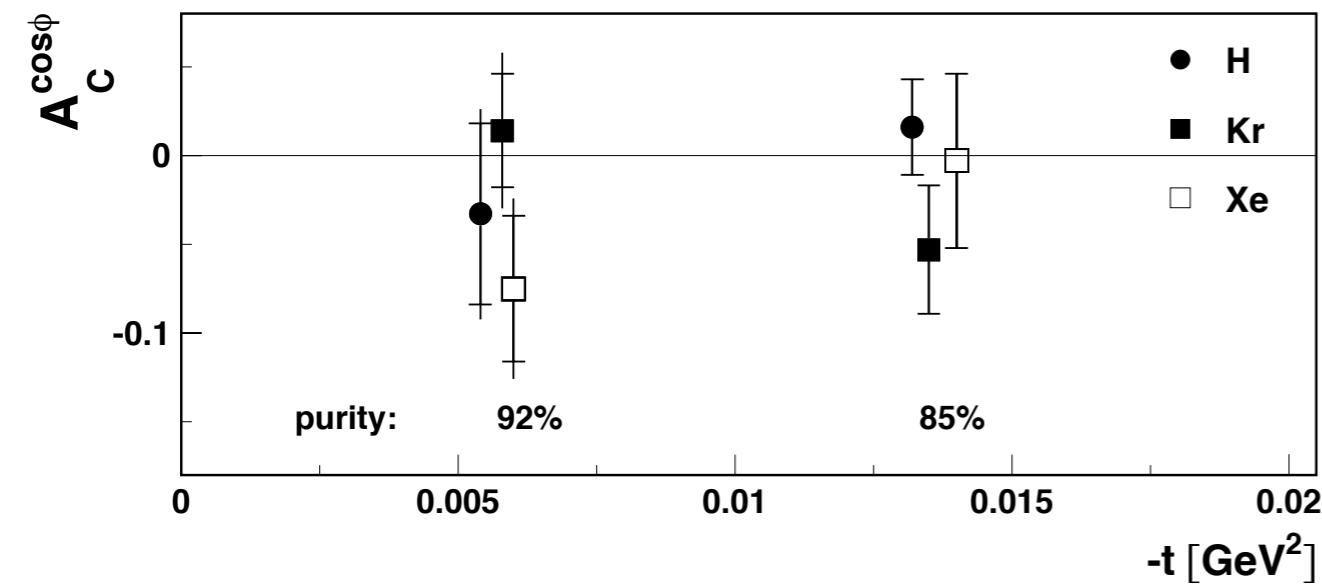
Accepted by
Nucl. Phys. B

arXiv:1008.3996 (hep-ex)

Nuclear DVCS backup



\mathcal{A}_{LU} amplitude sensitive to squared DVCS-term for nuclear targets with 2 beam charges



2 coherent-enriched t-bins