# Highlights from the HERMES experiment

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# HERMES: HERA MEasurement of Spin



### The HERMES experiment



#### <u>Beam</u>

longitudinally pol.  $e^+$  &  $e^-$ E=27.6 GeV <u>Gaseous internal target</u> transversely pol. H longitudinally pol. H, D, He unpol. H, D, He, Ne, Kr, Xe

- lepton-hadron PID:
- high efficiency (>98%)
- low contamination (<1%)
- hadron PID: RICH 2-15 GeV



- longitudinally polarised proton, deuteron, ...
- longitudinally polarised  $e^{\pm}$ ,  $\mu^{\pm}$  beam



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flip spin



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flip spin



• longitudinally polarised proton, deuteron, ...







flip spin

$$\frac{\stackrel{\leftarrow}{\sigma}}{\stackrel{\rightarrow}{\sigma}} \stackrel{\stackrel{\rightarrow}{\sigma}}{\stackrel{\rightarrow}{\sigma}} \propto g_1^p(x) = \frac{1}{2} \sum_q e_q^2 \Delta q(x) = \frac{1}{2} \sum_q e_q^2 \left( \stackrel{\leftarrow}{\stackrel{\leftarrow}{q}} (x) - \stackrel{\leftarrow}{\stackrel{\rightarrow}{q}} (x) \right)$$
 parton fractional longitudinal momentum

$$\frac{\stackrel{\leftarrow}{\sigma}}{\stackrel{\rightarrow}{\sigma}} \stackrel{\stackrel{\rightarrow}{\sigma}}{\stackrel{\rightarrow}{\sigma}} \propto g_1^p(x) = \frac{1}{2} \sum_q e_q^2 \Delta q(x) = \frac{1}{2} \sum_q e_q^2 \left( \stackrel{\leftarrow}{\stackrel{\leftarrow}{q}} (x) - \stackrel{\leftarrow}{\stackrel{\rightarrow}{q}} (x) \right)$$
 parton fractional longitudinal moments

longitudinal momentum



$$\begin{split} \overleftarrow{\sigma} & \overrightarrow{\sigma} \\ \overleftarrow{\sigma} & \overrightarrow{\sigma} \\ \overleftarrow{\sigma} & \overrightarrow{\sigma} \\ \overleftarrow{\sigma} & + & \overrightarrow{\sigma} \\ \hline \sigma & - & \overrightarrow{\sigma} \\ \hline \sigma & + & \overrightarrow{\sigma} \\ \hline \sigma & + & \overrightarrow{\sigma} \\ \hline \sigma & - & \overrightarrow{\sigma} \\ \hline \sigma & + & \overrightarrow{\sigma} \\ \hline \sigma & - & \overrightarrow{\sigma} \hline \sigma & - & \overrightarrow{\sigma} \\ \hline \sigma & - & \overrightarrow{\sigma} \hline \sigma & - & \overrightarrow$$

$$\begin{split} \stackrel{\overleftarrow{\alpha}}{\overleftarrow{\sigma}} & \stackrel{\overrightarrow{\alpha}}{\overrightarrow{\sigma}} \\ \stackrel{\overleftarrow{\alpha}}{\overleftarrow{\sigma}} & \stackrel{\overrightarrow{\alpha}}{\overrightarrow{\sigma}} \\ \stackrel{\overleftarrow{\alpha}}{\overleftarrow{\sigma}} & \stackrel{\overrightarrow{\alpha}}{\overrightarrow{\sigma}} \\ \stackrel{\overleftarrow{\alpha}}{\overleftarrow{\sigma}} & \stackrel{\overrightarrow{\alpha}}{\overrightarrow{\sigma}} \\ \stackrel{\overrightarrow{\alpha}}{\overleftarrow{\sigma}} & \stackrel{\overrightarrow{\alpha}}{\overrightarrow{\sigma}} \\ \stackrel{\overrightarrow{\alpha}}{\overrightarrow{\sigma}} & \stackrel{\overrightarrow{\alpha}}{\overrightarrow{\sigma}} \\ \stackrel{\overrightarrow{\alpha}}{\overrightarrow{\sigma}} \rightarrow \overrightarrow{\sigma} \overrightarrow{\sigma} \overrightarrow{\sigma}$$

From deuterium data:  $\Delta \Sigma (Q^2 = 5 \text{ GeV}^2) = 0.330 \pm 0.011 \text{(theo.)} \pm 0.025 \text{(exp.)} \pm 0.028 \text{(evol.)}$ 



### Disentangling quark flavours



 $\sigma^{ep \to eh} = \sum \Delta q(x, Q^2) \ \sigma^{eq \to eq} \ D_q^h(z, Q^2)$  $\boldsymbol{q}$ 



### Disentangling quark flavours



# Disentangling quark flavours



# Transverse momentum dependent semi-inclusive DIS



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#### Transverse-momentum-dependent (TMD) PDFs



### Transverse-momentum-dependent (TMD) PDFs



$$\Delta q = g_{1L} = \begin{array}{c} & & & \\$$

$$g_{1T}^{\perp} = - +$$

 $h_{1T}^{\perp} = 2$  -2







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# Semi-inclusive DIS single-hadron production



$$\sigma^{h}(\phi,\phi_{S}) = \sigma^{h}_{UU} \left\{ 1 + 2\langle \cos(\phi) \rangle^{h}_{UU} \cos(\phi) + 2\langle \cos(2\phi) \rangle^{h}_{UU} \cos(2\phi) \right\}$$

+  $\lambda_l 2 \langle \sin(\phi) \rangle_{LU}^h \sin(\phi)$ 

+ 
$$S_L \left[ 2 \langle \sin(\phi) \rangle_{UL}^h \sin(\phi) + 2 \langle \sin(2\phi) \rangle_{UL}^h \sin(2\phi) \right]$$

- +  $\lambda_l \left( 2 \langle \cos(0\phi) \rangle_{LL}^h \cos(0\phi) + 2 \langle \cos(\phi) \rangle_{LL}^h \cos(\phi) \right) \right]$
- +  $S_T \left[ 2 \langle \sin(\phi \phi_S) \rangle_{UT}^h \sin(\phi \phi_S) + 2 \langle \sin(\phi + \phi_S) \rangle_{UT}^h \sin(\phi + \phi_S) \rangle_{UT}^h \sin(\phi + \phi_S) \right]$
- +  $2\langle \sin(3\phi \phi_S) \rangle_{UT}^h \sin(3\phi \phi_S) + 2\langle \sin(\phi_S) \rangle_{UT}^h \sin(\phi_S)$

+ 
$$2\langle \sin(2\phi - \phi_S) \rangle_{UT}^h \sin(2\phi - \phi_S)$$

+ 
$$\lambda_l \left( 2 \langle \cos(\phi - \phi_S) \rangle_{LT}^h \cos(\phi - \phi_S) \right)$$

+ 
$$2\langle\cos(\phi_S)\rangle_{LT}^h\cos(\phi_S) + 2\langle\cos(2\phi - \phi_S)\rangle_{LT}^h\cos(2\phi - \phi_S)\rangle\right]$$











Azimuthal amplitudes related to structure functions  $F_{XY}$ :

 $2\langle \sin(\phi + \phi_S) \rangle_{UT}^h = \epsilon F_{UT}^{\sin(\phi + \phi_S)}$ 

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P

TMD

FF

 $\sigma$ 

TMD

PDF

k

е

 $F_{XY} \propto \mathcal{C} [\text{TMD PDF}(x, k_{\perp}) \times \text{TMD FF}(z, p_{\perp})]$ 

_	quark polarization				
DOIALIZAUOL		U	L	т	
	U	$f_1$		$h_1^\perp$	
nucleon	L		$g_{1L}$	$h_{1L}^{\perp}$	
	т	$f_{1T}^{\perp}$	$g_{1T}^{\perp}$	$h_{1T} h_{1T}^{\perp}$	

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zatior		U	L	т	
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leon	L		$g_{1L}$	$h_{1L}^{\perp}$	
nuc	т	$f_{1T}^{\perp}$	$g_{1T}^{\perp}$	$h_{1T} h_{1T}^{\perp}$	

quark polarization

hadron polarization



P

k

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### Presented amplitudes

$$\begin{aligned} \sigma^{h}(\phi,\phi_{S}) &= \sigma_{UU}^{h} \left\{ 1 + 2\langle\cos(\phi)\rangle_{UU}^{h} \cos(\phi) + 2\langle\cos(2\phi)\rangle_{UU}^{h} \cos(2\phi) \\ &+ \lambda_{l} 2\langle\sin(\phi)\rangle_{LU}^{h} \sin(\phi) \\ &+ S_{L} \left[ 2\langle\sin(\phi)\rangle_{UL}^{h} \sin(\phi) + 2\langle\sin(2\phi)\rangle_{UL}^{h} \sin(2\phi) \\ &+ \lambda_{l} \left( 2\langle\cos(0\phi)\rangle_{LL}^{h} \cos(0\phi) + 2\langle\cos(\phi)\rangle_{LL}^{h} \cos(\phi) \right) \right] \\ &+ S_{T} \left[ 2\langle\sin(\phi - \phi_{S})\rangle_{UT}^{h} \sin(\phi - \phi_{S}) + 2\langle\sin(\phi + \phi_{S})\rangle_{UT}^{h} \sin(\phi + \phi_{S}) \\ &+ 2\langle\sin(3\phi - \phi_{S})\rangle_{UT}^{h} \sin(3\phi - \phi_{S}) + 2\langle\sin(\phi_{S})\rangle_{UT}^{h} \sin(\phi_{S}) \\ &+ 2\langle\sin(2\phi - \phi_{S})\rangle_{UT}^{h} \sin(2\phi - \phi_{S}) \\ &+ \lambda_{l} \left( 2\langle\cos(\phi - \phi_{S})\rangle_{LT}^{h} \cos(\phi - \phi_{S}) \\ &+ 2\langle\cos(\phi_{S})\rangle_{LT}^{h} \cos(\phi_{S}) + 2\langle\cos(2\phi - \phi_{S})\rangle_{LT}^{h} \cos(2\phi - \phi_{S}) \right) \right] \end{aligned}$$

Presented here

$$Q^2 > 1 \text{ GeV}^2$$
  
 $W^2 > 10 \text{ GeV}^2$   
 $0.023 < x < 0.6$ 

### Presented amplitudes

- Unpolarized and longitudinally polarized e+/e- beam
- Transversely polarized H target: fit all amplitudes simultaneously
  - → Results for charged pions, kaons, (anti-)protons, neutral pions

$$+ S_{T} \left[ 2 \langle \sin(\phi - \phi_{S}) \rangle_{UT}^{h} \sin(\phi - \phi_{S}) + 2 \langle \sin(\phi + \phi_{S}) \rangle_{UT}^{h} \sin(\phi + \phi_{S}) \right]$$

$$+ 2 \langle \sin(3\phi - \phi_{S}) \rangle_{UT}^{h} \sin(3\phi - \phi_{S}) + 2 \langle \sin(\phi_{S}) \rangle_{UT}^{h} \sin(\phi_{S})$$

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# Collins amplitudes

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$$\mathcal{C}[h_{1T}^q \times H_1^{\perp,q}]$$



• Positive amplitudes for  $\pi^+$ , negative amplitudes for  $\pi^-$ :

$$H_1^{\perp, u \to \pi^+} \approx -H_1^{\perp, u \to \pi^-}$$
#### Artru model

X. Artru et al., Z. Phys. C73 (1997) 527

polarisation component in lepton scattering plane reversed by photoabsorption:



string break, quark-antiquark pair with vacuum numbers:





orbital angular momentum creates transverse momentum:



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 $\mathcal{C}[h_{1T}^q \times H_1^{\perp,q}]$ 





COMPASS, HERMES, Jefferson Lab Hall A

Collins amplitudes

 $\mathcal{C}[h_{1T}^q \times H_1^{\perp,q}]$ 





 $\mathcal{C}[h_{1T}^q \times H_1^{\perp,q}]$ 





 $\mathcal{C}[h_{1T}^q \times H_1^{\perp,q}]$ 





# Pretzelosity amplitudes $C[h_{1T}^{\perp,q} \times H_1^{\perp,q}]$

- Pretzelosity
  - requires non-zero orbital angular momentum
  - models:

$$h_{1T}^{\perp(1),q}(x) = g_{1L}^q(x) - h_{1T}^q(x)$$

- $\twoheadrightarrow$  measure for relativistic effects
- suppressed as  $P_{h\perp}^2$  compared to Collins amplitude



#### Pretzelosity amplitudes



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data from Jefferson Lab Hall A preliminary data from COMPASS, HERMES

Twist-3:  $\langle \sin(\phi_S) \rangle_{UT}$ 

$$\begin{split} &\langle \sin(\phi_S) \rangle_{UT} \\ &\propto \mathcal{C}[f_T^q \times D_1^q, h_{1T}^q \times \tilde{H}^q, h_T^q \times H_1^{\perp,q}, g_{1T}^{\perp,q} \times \tilde{G}^{\perp,q}, h_T^{\perp,q} \times H_1^{\perp,q}, f_{1T}^{\perp,q} \times \tilde{D}^{\perp,q}] \\ & \text{twist-3} \end{split}$$

Twist-3: 
$$\langle \sin(\phi_S) \rangle_{UT}$$

 $\langle \sin(\phi_S) \rangle_{UT} \\ \propto \mathcal{C}[f_T^q \times D_1^q, h_{1T}^q \times \tilde{H}^q, h_T^q \times H_1^{\perp,q}, g_{1T}^{\perp,q} \times \tilde{G}^{\perp,q}, h_T^{\perp,q} \times H_1^{\perp,q}, f_{1T}^{\perp,q} \times \tilde{D}^{\perp,q} ]$ twist-3 integrate over hadron transverse momentum  $P_{h\perp}$ 

$$\langle \sin(\phi_S) \rangle_{UT} = -x \frac{2M_h}{Q} \sum_q e_q^2 h_{1T}^q \frac{H^q}{z}$$

no convolution

Twist-3:  $\langle \sin(\phi_S) \rangle_{UT}$ 



• Significant non-zero signal for  $\pi^-$ , increasing with x, z

Twist-3:  $\langle \sin(\phi_S) \rangle_{UT}$ 







Wigner distributions  $W(x, \vec{k}_T, \vec{b}_\perp)$   $\int d^2 \vec{b}_\perp$ transverse-momentum dependent PDFs (TMDs)  $\gamma^*$ 























- x=average longitudinal momentum fraction
- 2ξ=average longitudinal momentum transfer
- t=squared momentum transfer to nucleon



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Four quark helicity-conserving twist-2 GPDs

$H(x,\xi,t)$	$E(x,\xi,t)$	spin independent
$ ilde{H}(x,\xi,t)$	$ ilde{E}(x,\xi,t)$	spin dependent
proton helicity non flip	proton helicity flip	



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Four quark helicity-flip twist-2 GPDs

$H_T(x,\xi,t)$	$E_T(x,\xi,t)$
$\tilde{H}_T(x,\xi,t)$	$\tilde{E}_T(x,\xi,t)$



- x=average longitudinal momentum fraction
- 2ξ=average longitudinal momentum transfer
- t=squared momentum transfer to nucleon

Four quark helicity-conserving twist-2 GPDs

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Four quark helicity-flip twist-2 GPDs

$H_T(x,\xi,t)$	$E_T(x,\xi,t)$
$\tilde{H}_T(x,\xi,t)$	$\tilde{E}_T(x,\xi,t)$

$$J = \lim_{t \to 0} \frac{1}{2} \int_{-1}^{1} dx \, x \left[ H(x,\xi,t) + E(x,\xi,t) \right]$$
 X. Ji, Phys. Rev. Lett. 78 (1997) 610

#### Hard exclusive processes



- Deeply virtual Compton scattering (DVCS): theoretically cleanest probe
- Exclusive meson production
  - probe various types of GPDs with different sensitivity and different flavour combinations: also access to quark-helicity-flip GPDs
  - complementary to DVCS
- Target polarization state: access to different GPDs











**Bethe-Heitler** 



 $d\sigma \propto |\tau_{BH}|^2 + |\tau_{DVCS}|^2 + \tau_{DVCS}\tau_{BH}^* + \tau_{DVCS}^*\tau_{BH}$ 

 $d\sigma \propto |\tau_{BH}|^2 + |\tau_{DVCS}|^2 + \tau_{DVCS}\tau_{BH}^* + \tau_{DVCS}^*\tau_{BH}$ 

Unpolarized nucleon Longitudinally polarized lepton beam

 $d\sigma \propto |\tau_{BH}|^2 + |\tau_{DVCS}|^2 + \tau_{DVCS}\tau_{BH}^* + \tau_{DVCS}^*\tau_{BH}$ 

#### Unpolarized nucleon Longitudinally polarized lepton beam

$$|\tau_{BH}|^2 = \frac{K_{BH}}{\mathcal{P}_1(\phi) \,\mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^2 c_n^{BH} \cos(n\phi) \right\}$$

calculable with knowledge Pauli & Dirac form factors

$$|\tau_{DVCS}|^{2} = \frac{1}{Q^{2}} \left\{ \sum_{n=0}^{2} c_{n}^{DVCS} \cos(n\phi) + \lambda s_{1}^{DVCS} \sin(\phi) \right\}$$

coefficients: bilinear in GPDs

$$\mathcal{I} = \frac{-e_l K_{\mathcal{I}}}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^{\mathcal{I}} \cos(n\phi) + \lambda \sum_{n=1}^2 s_n^{\mathcal{I}} \sin(n\phi) \right\} \quad \text{coefficients:}$$

coefficients: linear in GPDs



 $d\sigma \propto |\tau_{BH}|^2 + |\tau_{DVCS}|^2 + \tau_{DVCS}\tau_{BH}^* + \tau_{DVCS}^*\tau_{BH}$ 

## Unpolarized nucleon Longitudinally polarized lepton beam $|\tau_{BH}|^2 = \frac{K_{BH}}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^{2} c_n^{BH} \cos(n\phi) \right\}$ calculable with knowledge Pauli & Dirac form factors $|\tau_{DVCS}|^{2} = \frac{1}{Q^{2}} \left\{ \sum_{n=0}^{2} c_{n}^{DVCS} \cos(n\phi) + \lambda s_{1}^{DVCS} \sin(\phi) \right\}$ coefficients: bilinear in GPDs $\mathcal{I} = \frac{-e_l K_{\mathcal{I}}}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^{3} c_n^{\mathcal{I}} \cos(n\phi) + \lambda \sum_{n=1}^{2} s_n^{\mathcal{I}} \sin(n\phi) \right\} \quad \text{coefficients: linear in GPDs}$ beam polarization 25

 $d\sigma \propto |\tau_{BH}|^2 + |\tau_{DVCS}|^2 + \tau_{DVCS}\tau_{BH}^* + \tau_{DVCS}^*\tau_{BH}$ 



$$\mathcal{I} = \frac{-e_l K_{\mathcal{I}}}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^{\mathcal{I}} \cos(n\phi) + \lambda \sum_{n=1}^2 s_n^{\mathcal{I}} \sin(n\phi) \right\}$$

 $c_1^{\mathcal{I}} \propto \Re M^{1,1}$   $s_1^{\mathcal{I}} \propto \Im M^{1,1}$ 

$$\begin{split} M^{1,1} &= F_1(t) \,\mathcal{H}(\xi,t) + \frac{x_B}{2 - x_B} (F_1(t) + F_2(t)) \,\tilde{\mathcal{H}}(\xi,t) - \frac{t}{4M_p^2} \,F_2(t) \mathcal{E}(\xi,t) \\ & \quad \text{CFF} \,\mathcal{H}, \tilde{\mathcal{H}}, \mathcal{E} = \text{convolution GPD x hard scattering amplitude} \end{split}$$

At LO:  $\Im$  direct access to GPDs at  $x = \pm \xi$  $\Re$  convolution integral over x+ access to D-term



#### Beam-charge asymmetry

$$\mathcal{A}_{\rm C}(\phi) \equiv \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-} = \frac{-\frac{K_{\rm I}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \sum_{n=0}^3 c_n^I \cos(n\phi)}{\frac{K_{\rm BH}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \sum_{n=0}^2 c_n^{\rm BH} \cos(n\phi) + \frac{1}{Q^2} \sum_{n=0}^2 c_n^{\rm DVCS} \cos(n\phi)}$$

#### Beam-charge asymmetry


HERMES, JHEP 11 (2009) 083 0.1 VGG Regge, no D **Dual-GT Regge**  ${\sf A}_{\sf C}^{\cos{(0\varphi)}}$ 0 -0.1 0.3 0.2  ${f A}_{C}^{\cos\phi}$ 0.1 0 0.1  ${\sf A}_{\sf C}^{\sf cos\,(2\phi)}$ 0 -0.1 0.1  ${\sf A}_{\sf C}^{\cos{(3\psi)}}$ 0 -0.1 Assoc. fraction 0.3 0.2 0.1 10<sup>-2</sup> **10**<sup>-1</sup> **10**<sup>-1</sup> 10 1  $Q^2$  [GeV<sup>2</sup>] -t [GeV<sup>2</sup>] overall **X**B







Unpolarized nucleon Longitudinally polarized lepton beam

$$\begin{aligned} |\tau_{BH}|^2 &= \frac{K_{BH}}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^2 c_n^{BH} \cos(n\phi) \right\} & \text{Calculable with knowledge Pauli & Dirac form factors} \\ |\tau_{DVCS}|^2 &= \frac{1}{Q^2} \left\{ \sum_{n=0}^2 c_n^{DVCS} \cos(n\phi) + \lambda s_1^{DVCS} \sin(\phi) \right\} & \text{coefficients: bilinear in GPDs} \\ \mathcal{I} &= \frac{-e_l K_{\mathcal{I}}}{\mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^{\mathcal{I}} \cos(n\phi) + \lambda \sum_{n=1}^2 s_n^{\mathcal{I}} \sin(n\phi) \right\} & \text{coefficients: linear in GPDs} \\ & \text{beam} \\ & \text{charge} & \text{beam} \\ \end{aligned}$$

Charge-difference beam-helicity asymmetry

$$\mathcal{A}_{\mathrm{LU}}^{\mathrm{I}}(\phi) \equiv \frac{(d\sigma^{+\rightarrow} - d\sigma^{+\leftarrow}) - (d\sigma^{-\rightarrow} - d\sigma^{-\leftarrow})}{(d\sigma^{+\rightarrow} + d\sigma^{+\leftarrow}) + (d\sigma^{-\rightarrow} + d\sigma^{-\leftarrow})}$$
$$= \frac{-\frac{K_{\mathrm{I}}}{\mathcal{P}_{1}(\phi)\mathcal{P}_{2}(\phi)} \left[\sum_{n=1}^{2} s_{n}^{\mathrm{I}} \sin(n\phi)\right]}{\frac{K_{\mathrm{BH}}}{\mathcal{P}_{1}(\phi)\mathcal{P}_{2}(\phi)} \sum_{n=0}^{2} c_{n}^{\mathrm{BH}} \cos(n\phi) + \frac{1}{Q^{2}} \sum_{n=0}^{2} c_{n}^{\mathrm{DVCS}} \cos(n\phi)}$$

Charge-difference beam-helicity asymmetry

linear access to GPDs

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Charge-difference beam-helicity asymmetry

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Charge-averaged beam-helicity asymmetry

$$\mathcal{A}_{\mathrm{LU}}^{\mathrm{DVCS}}(\phi) \equiv \frac{(d\sigma^{+\to} - d\sigma^{+\leftarrow}) + (d\sigma^{-\to} - d\sigma^{-\leftarrow})}{(d\sigma^{+\to} + d\sigma^{+\leftarrow}) + (d\sigma^{-\to} + d\sigma^{-\leftarrow})}$$
$$= \frac{\frac{1}{Q^2} s_1^{\mathrm{DVCS}} \sin \phi}{\frac{K_{\mathrm{BH}}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \sum_{n=0}^2 c_n^{\mathrm{BH}} \cos(n\phi) + \frac{1}{Q^2} \sum_{n=0}^2 c_n^{\mathrm{DVCS}} \cos(n\phi)}$$

Charge-difference beam-helicity asymmetry

linear access to GPDs

$$\begin{aligned} \mathcal{A}_{\mathrm{LU}}^{\mathrm{I}}(\phi) &\equiv \frac{(d\sigma^{+\rightarrow} - d\sigma^{+\leftarrow}) - (d\sigma^{-\rightarrow} - d\sigma^{-\leftarrow})}{(d\sigma^{+\rightarrow} + d\sigma^{+\leftarrow}) + (d\sigma^{-\rightarrow} + d\sigma^{-\leftarrow})} \\ &= \frac{-\frac{K_{\mathrm{I}}}{\mathcal{P}_{1}(\phi)\mathcal{P}_{2}(\phi)} \left[\sum_{n=1}^{2} s_{n}^{\mathrm{I}} \sin(n\phi)\right]}{\frac{K_{\mathrm{BH}}}{\mathcal{P}_{1}(\phi)\mathcal{P}_{2}(\phi)} \sum_{n=0}^{2} c_{n}^{\mathrm{BH}} \cos(n\phi) + \frac{1}{Q^{2}} \sum_{n=0}^{2} c_{n}^{\mathrm{DVCS}} \cos(n\phi)} \end{aligned}$$

bilinear access to GPDs

Charge-averaged beam-helicity asymmetry

$$\mathcal{A}_{\mathrm{LU}}^{\mathrm{DVCS}}(\phi) \equiv \frac{(d\sigma^{+\rightarrow} - d\sigma^{+\leftarrow}) + (d\sigma^{-\rightarrow} - d\sigma^{-\leftarrow})}{(d\sigma^{+\rightarrow} + d\sigma^{+\leftarrow}) + (d\sigma^{-\rightarrow} + d\sigma^{-\leftarrow})}$$
$$= \frac{\frac{1}{Q^2} s_1^{\mathrm{DVCS}} \sin \phi}{\frac{K_{\mathrm{BH}}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \sum_{n=0}^2 c_n^{\mathrm{BH}} \cos(n\phi) + \frac{1}{Q^2} \sum_{n=0}^2 c_n^{\mathrm{DVCS}} \cos(n\phi)}$$

## Charge-difference and charge-average beam-helicity asymmetry



# Charge-difference and charge-average beam-helicity asymmetry



#### DVCS at HERMES



beam-charge asymmetry

JHEP 07 (2012) 32 Nucl. Phys. B 829 (2010) 1

beam-helicity asymmetry

JHEP 07 (2012) 32 Nucl. Phys. B 829 (2010) 1 JHEP10(2012)042

transverse target-spin asymmetry

JHEP 06 (2008) 066

double spin (LT) asymmetry

Phys. Lett. B 704 (2011) 15

longitudinal target-spin asymmetry

JHEP 06 (2010) 019 Nucl. Phys. B 842 (2011) 265

double spin (LL) asymmetry

JHEP 06 (2010) 019 Nucl. Phys. B 842 (2011) 265

## Summary







 $xP_z$ 

 $P_z$ 







**SDME values** 

