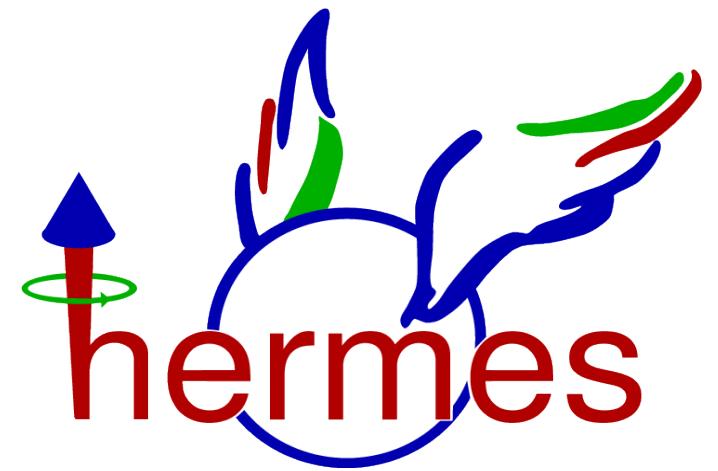


# Overview of recent HERMES results

Charlotte Van Hulse, University of Ghent  
on behalf of the HERMES collaboration

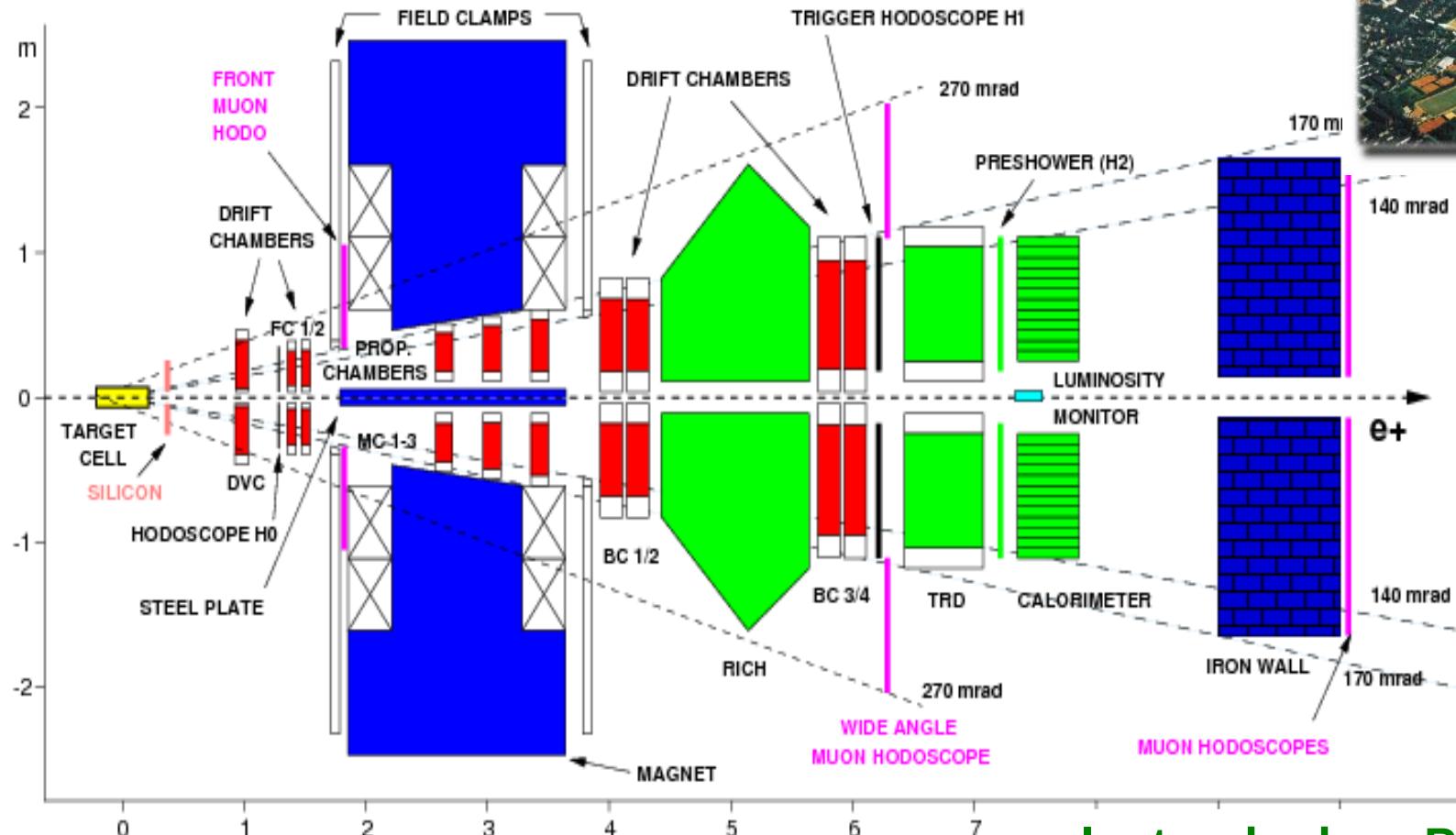
ENPC09



# Overview

- strange quark distributions
- transverse structure of the nucleon
  - transverse momentum dependent distributions
  - transverse position:  
generalized parton distributions and  
deeply virtual Compton scattering

# HERMES: HERA MEASUREMENT of SPIN



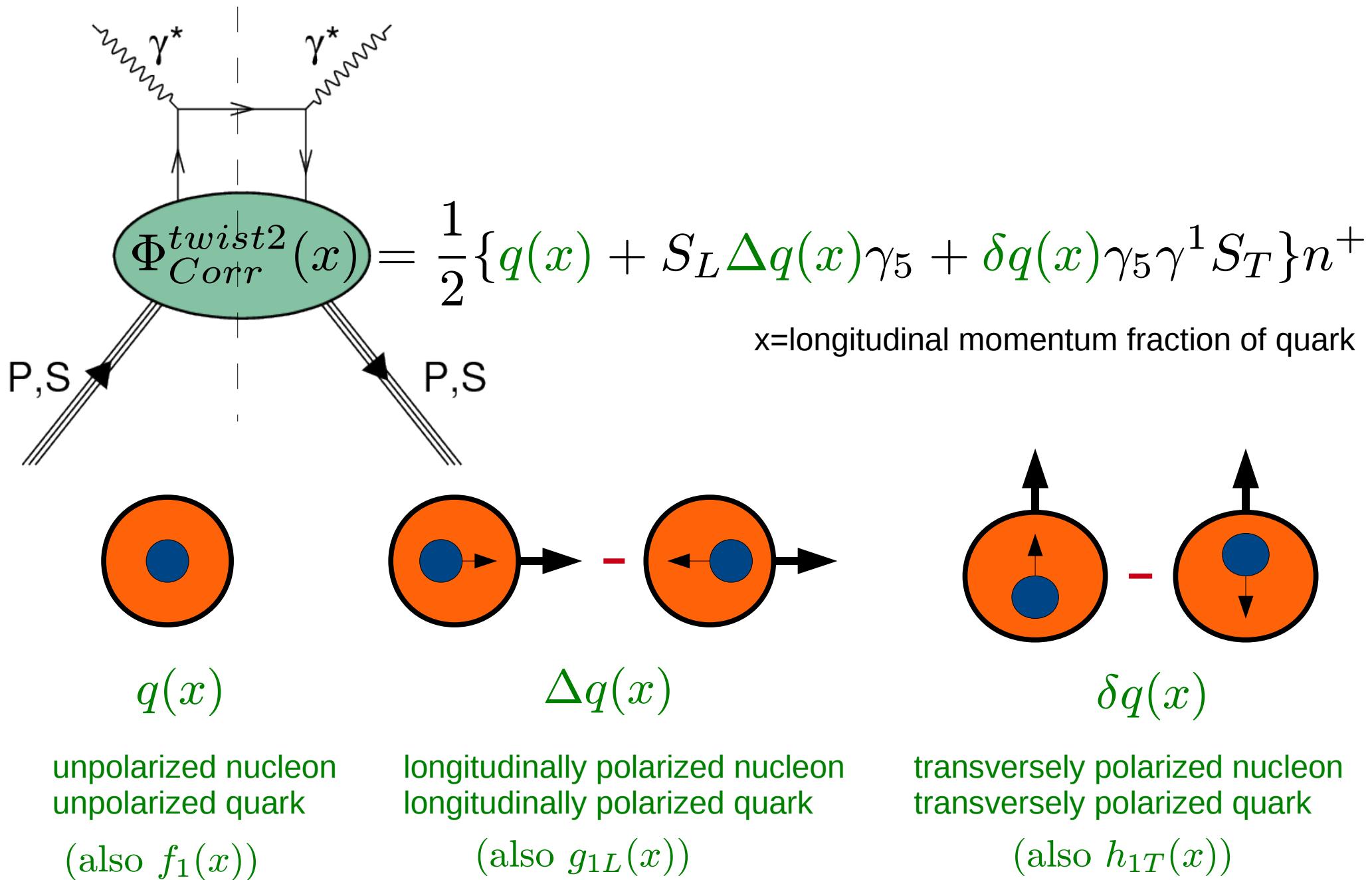
data taking from  
1995  
until  
June, 30 2007

Beam  
longitudinally pol.  
 $e^+$  &  $e^-$   
 $E = 27.6 \text{ GeV}$

Gaseous internal target  
longitudinally pol. H,D,He  
transversely pol. H  
unpol. H,D,Ne,Kr,..

lepton-hadron PID efficiency:~98%  
hadron PID: RICH 2-15 GeV/c  
 $\delta E\gamma/E\gamma \approx 5\%$   
 $\delta P/P < 2\%$   
 $\Delta\theta < 1\text{mrad}$

# Probing the nucleon in DIS



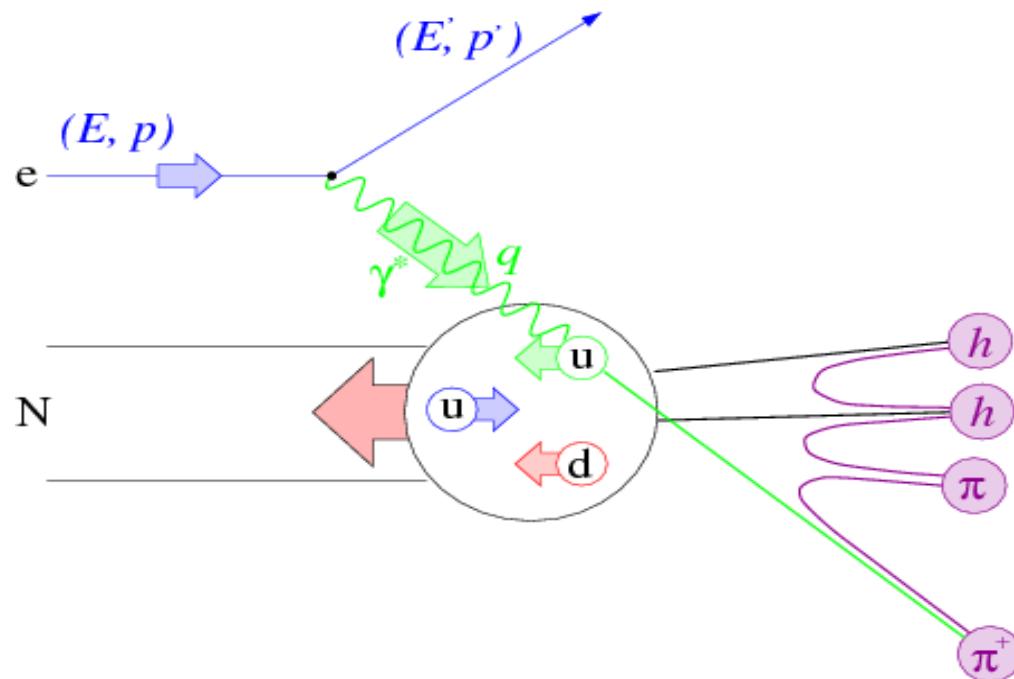
unpolarized nucleon  
unpolarized quark  
(also  $f_1(x)$ )

longitudinally polarized nucleon  
longitudinally polarized quark  
(also  $g_{1L}(x)$ )

transversely polarized nucleon  
transversely polarized quark  
(also  $h_{1T}(x)$ )

# How to access flavored quark distributions

## Semi-inclusive deep-inelastic scattering (SIDIS)



$$\begin{aligned}Q^2 &= -q^2 \\ \nu &\stackrel{lab}{=} E - E' \\ y &\stackrel{lab}{=} \frac{\nu}{E} \\ x &\stackrel{lab}{=} \frac{Q^2}{2M\nu} \\ z &\stackrel{lab}{=} \frac{E_h}{\nu}\end{aligned}$$

$$\sigma^{ep \rightarrow eh} = \sum_q DF^{p \rightarrow q}(x) \otimes \sigma^{eq \rightarrow eq} \otimes FF^{q \rightarrow h}(z)$$

**Distribution Function (DF):** distribution of quarks in nucleon

**Fragmentation Function (FF):** fragmentation of struck quark into final-state hadron

# Strange quark distributions

# Extraction of the strange quark distribution

- leading order extraction in  $\alpha_s$
- charged kaon production on (longitudinally pol.) deuteron:
$$e + d \rightarrow K^\pm + X$$
- strange quarks carry no isospin  $\rightarrow s_p(x) = s_n(x)$ 
$$\Delta s_p(x) = \Delta s_n(x)$$
- isoscalar target &  $K^+ + K^- \rightarrow$  FF without isospin dependence
- hypothesis: isospin symmetry between p and n  
charge-conjugation invariance of FF

A. Airapetian et al., Phys. Lett. B **666**, 446-450 (2008)

# Kaon multiplicities

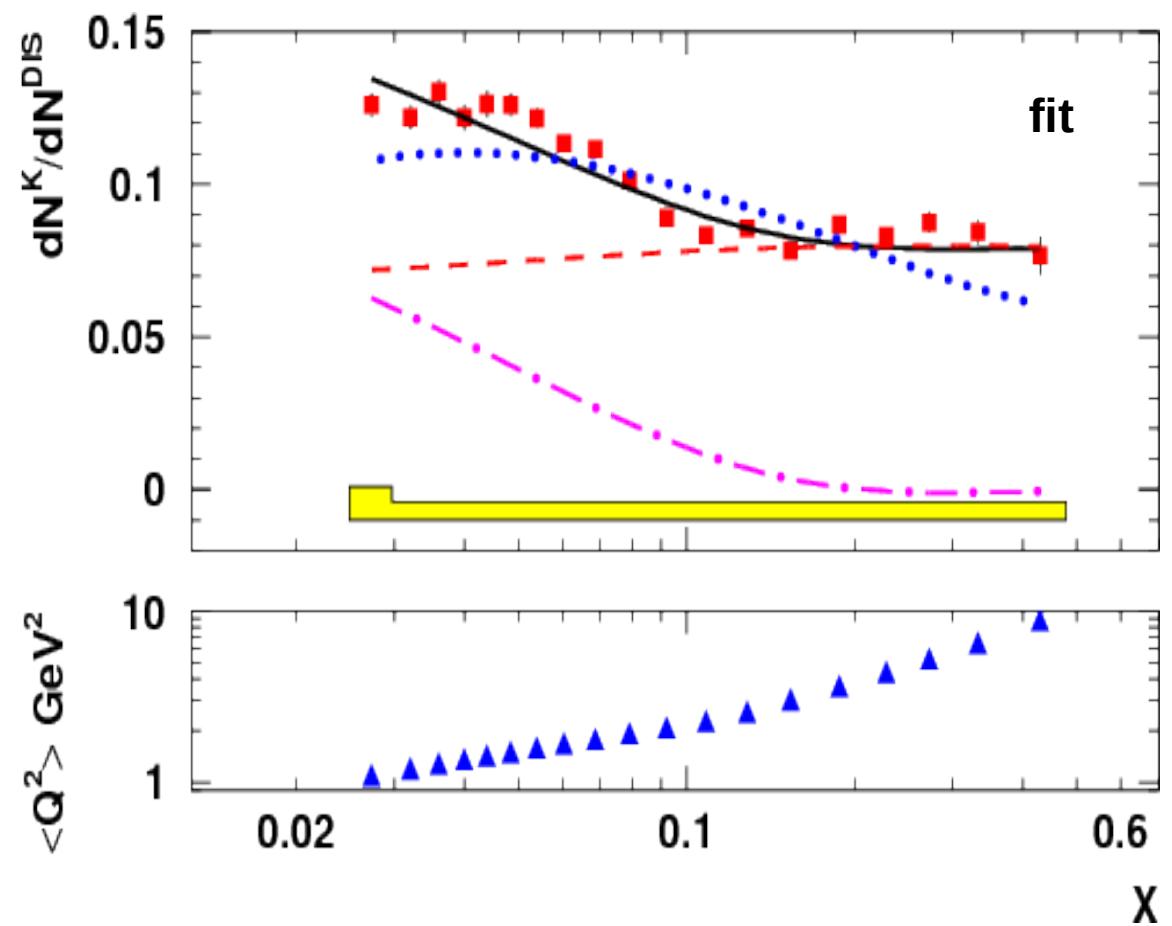
$$\frac{d^2 N^K(x)}{d^2 N^{DIS}(x)} = \frac{Q(x) \int_{0.2}^{0.8} \mathcal{D}_Q^K(z) dz + S(x) \int_{0.2}^{0.8} \mathcal{D}_S^K(z) dz}{5Q(x) + 2S(x)}$$

$$Q(x) \equiv u(x) + \bar{u}(x) + d(x) + \bar{d}(x)$$

$$S(x) \equiv s(x) + \bar{s}(x)$$

$$\mathcal{D}_Q^K(z) \equiv 4D_u^K(z) + D_d^K(z)$$

$$\mathcal{D}_S^K(z) \equiv 2D_s^K(z)$$



# Kaon multiplicities

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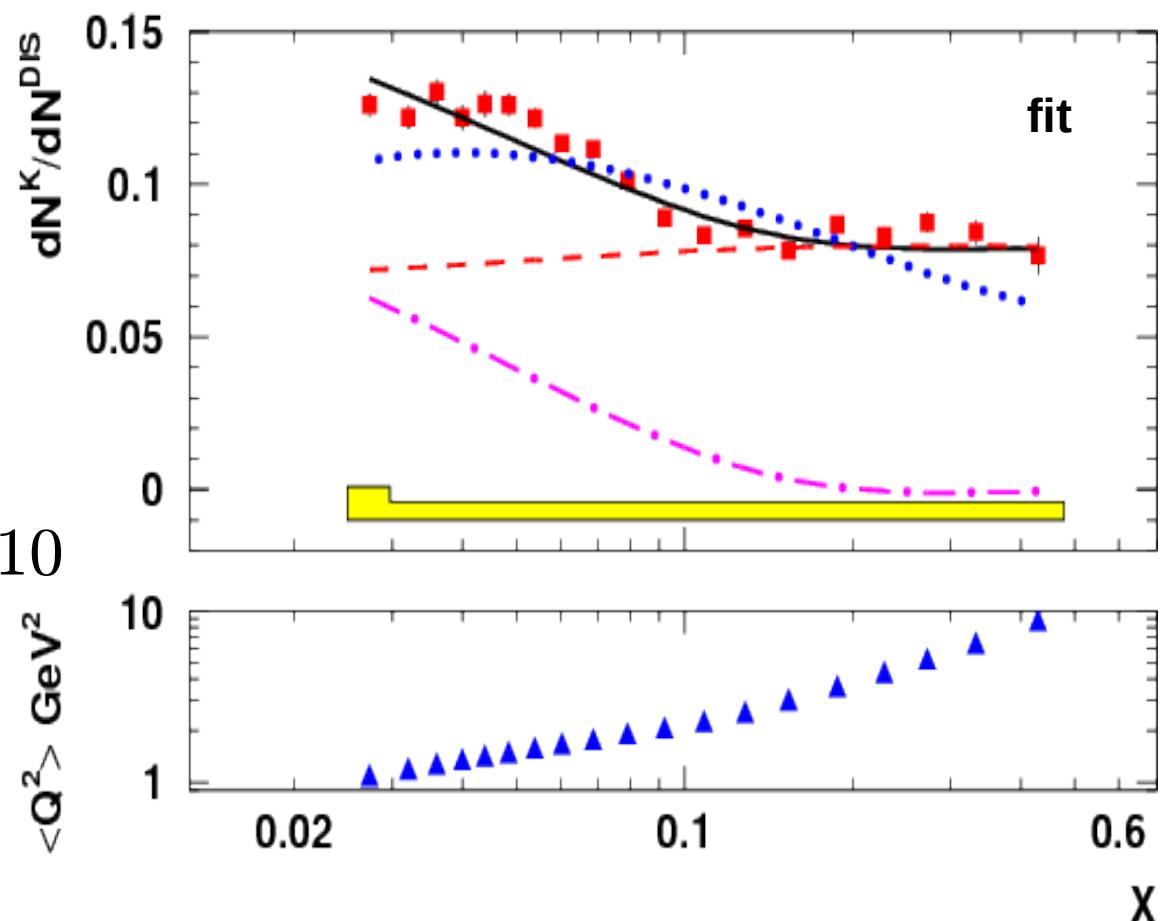
$$\mathcal{D}_Q^K(z) \equiv 4D_u^K(z) + D_d^K(z)$$

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$S(x)=0$  for  $x>0.15$

$Q(x)$  from CTEQ6L

$$\int_{0.2}^{0.8} \mathcal{D}_Q^K(z) dz = 0.398 \pm 0.010$$



# Kaon multiplicities

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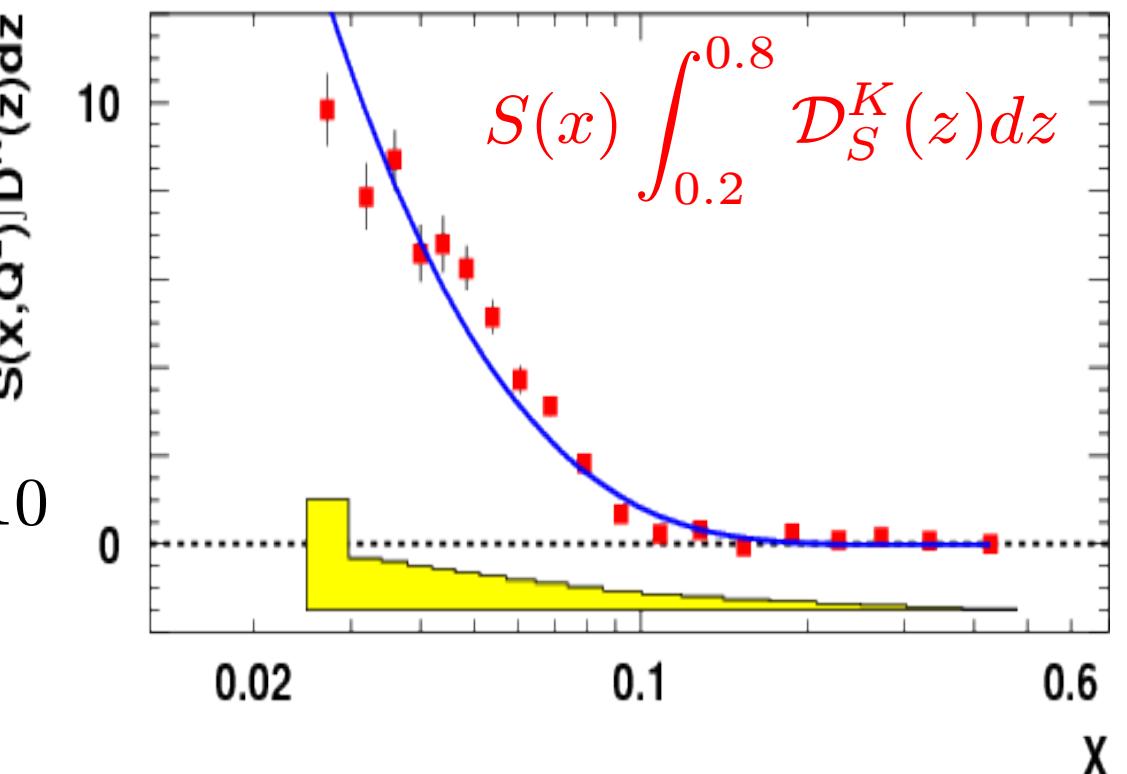
$$S(x)=0 \text{ for } x>0.15$$

Q(x) from CTEQ6L

$$\int_{0.2}^{0.8} \mathcal{D}_Q^K(z) dz = 0.398 \pm 0.010$$

neglecting  $2S(x)$

$$S(x) \int_{0.2}^{0.8} \mathcal{D}_S^K(z) dz$$

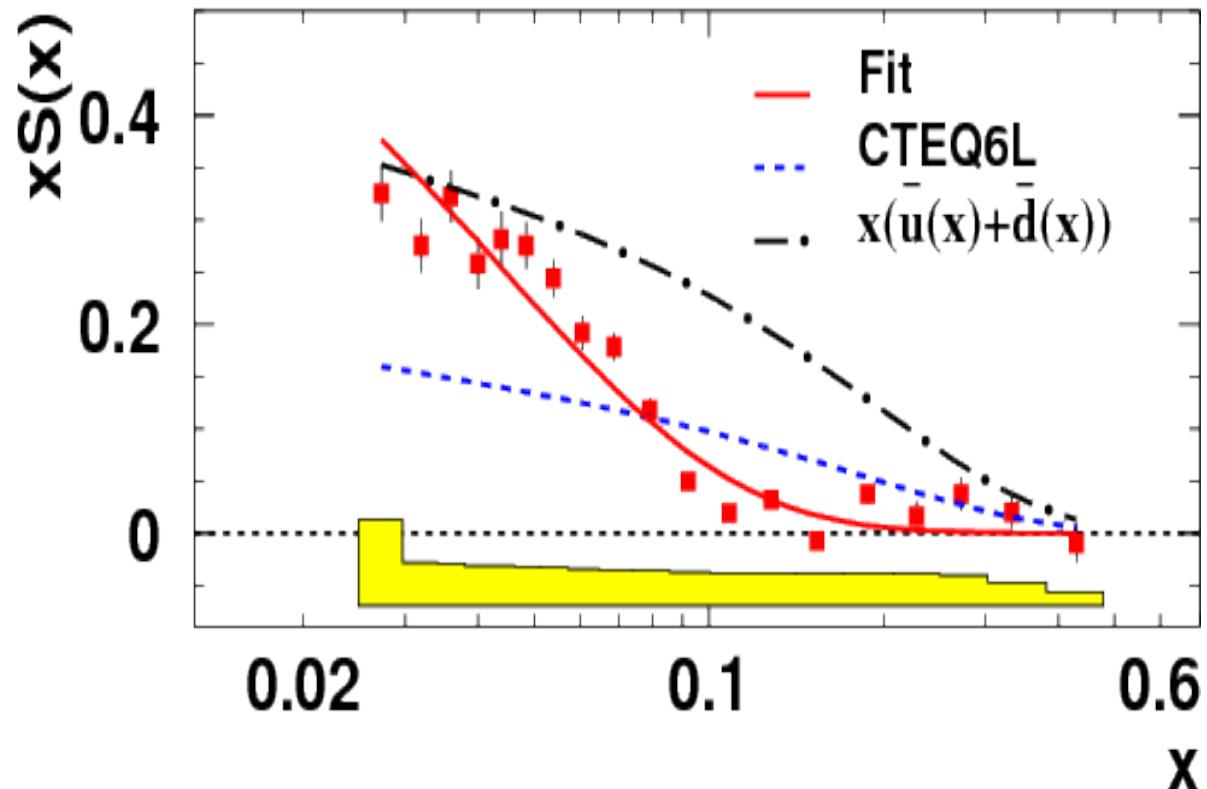
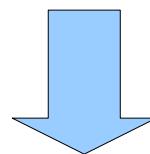


# Spin-independent strange quark distribution

- evolution of data to  $Q^2=2.5 \text{ GeV}^2$  (CTEQ6L)

- $\int_{0.2}^{0.8} \mathcal{D}_S^K(z) dz = 1.27 \pm 0.13$

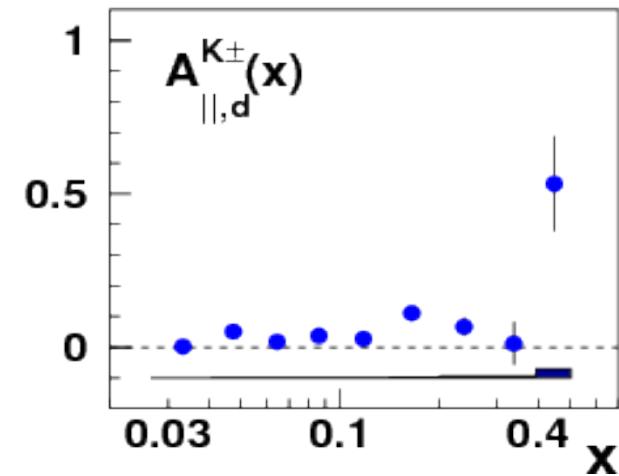
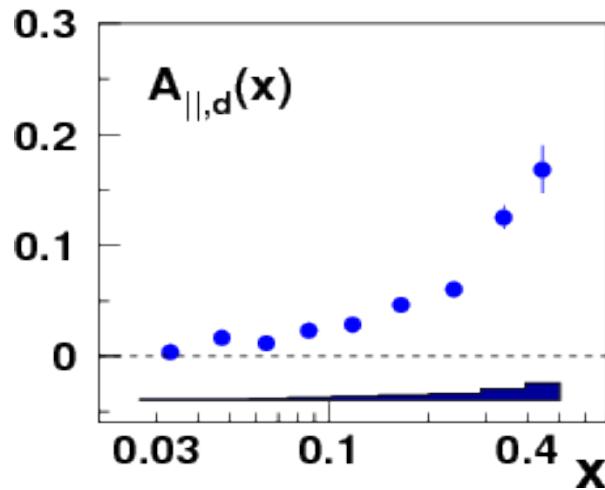
D. de Florian et al., PRD75, 114010



- discrepancy with CTEQ6L
- $S(x)$  not an average of isoscalar non-strange sea  $S(x)$  towards lower  $x$ -values

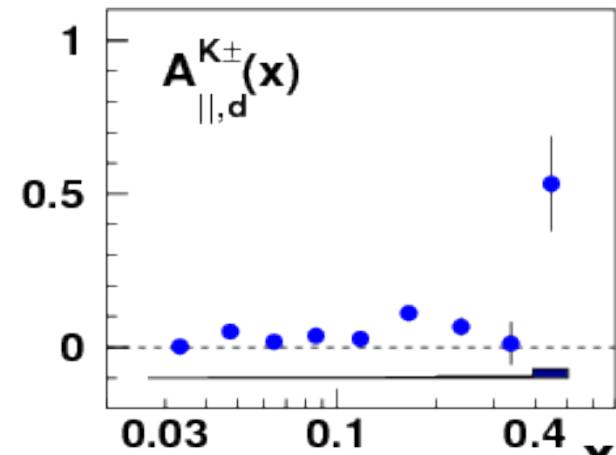
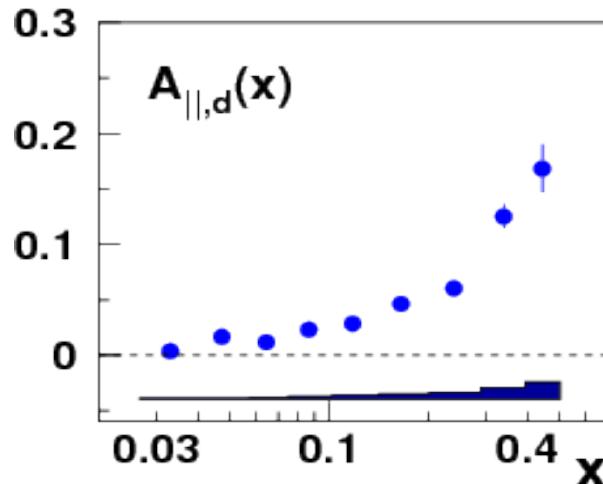
# Double spin asymmetry from longitudinally polarized deuterons

$$A_{\parallel}^{(h)} = \frac{\sigma^{\leftarrow,(h)} - \sigma^{\rightarrow,(h)}}{\sigma^{\leftarrow,(h)} + \sigma^{\rightarrow,(h)}}$$



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$$A_{\parallel}^{(h)} = \frac{\sigma^{\leftarrow,(h)} - \sigma^{\rightarrow,(h)}}{\sigma^{\leftarrow,(h)} + \sigma^{\rightarrow,(h)}}$$



$$A_{\parallel,d}(x) \frac{d^2 N^{DIS}(x)}{dx dQ^2} = \mathcal{K}_{LL}(x, Q^2) [5\Delta Q(x) + 2\Delta S(x)]$$

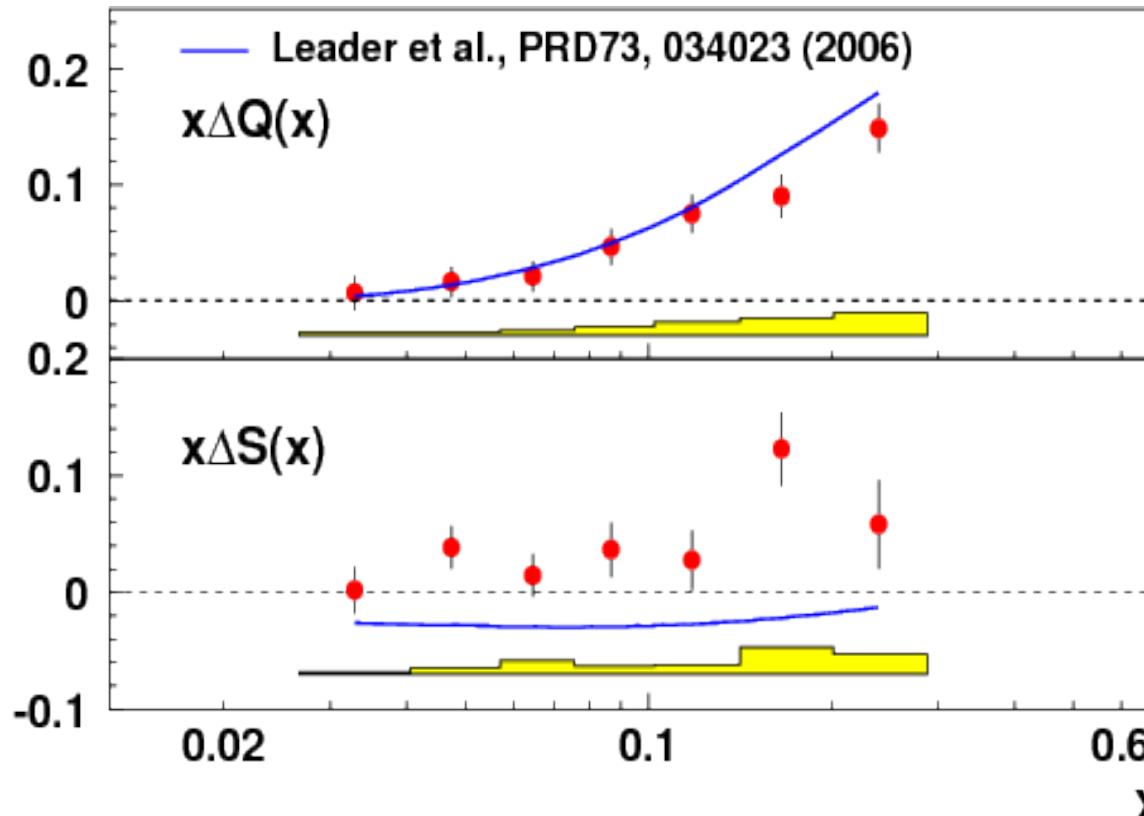
$$A_{\parallel,d}^{K\pm}(x) \frac{d^2 N^K(x)}{dx dQ^2} = \mathcal{K}_{LL}(x, Q^2) \left[ \Delta Q(x) \int_{0.2}^{0.8} \mathcal{D}_Q^K(z) dz + \Delta S(x) \int_{0.2}^{0.8} \mathcal{D}_S^K(z) dz \right]$$

$$\Delta Q(x) \equiv \Delta u(x) + \Delta \bar{u}(x) + \Delta d(x) + \Delta \bar{d}(x)$$

$$\Delta S(x) \equiv \Delta s(x) + \Delta \bar{s}(x)$$

$$\int_{0.2}^{0.8} \mathcal{D}_Q^K(z) dz, \int_{0.2}^{0.8} \mathcal{D}_S^K(z) dz \quad \text{from } S(x) \text{ extraction}$$

# Helicity distribution @ $Q^2=2.5 \text{ GeV}^2$



- $\Delta Q$  and  $\Delta S$  in agreement with previous HERMES results
- $\Delta q_8 = \int [\Delta Q(x) - 2\Delta S(x)] dx$  for  $x: 0.02 \rightarrow 0.6$   
 $= 0.285 \pm 0.046(\text{stat.}) \pm 0.057(\text{sys.})$   
≠  
 $\Delta q_8$  from hyperon decay, assuming SU(3) symmetry ( $x: 0 \rightarrow 1$ )  
 $= 0.586 \pm 0.031$

SU(3) symmetry violation  $\longleftrightarrow$  large contribution from  $x < 0.02$  region

# Transverse structure of the nucleon

# Transverse structure of the nucleon

transverse momentum  
dependent distributions

# Transverse structure of the nucleon

transverse momentum  
dependent distributions

transverse position:  
generalized parton  
distributions

# Transverse momentum dependent distributions (TMDs)

$$\sigma^{ep \rightarrow eh} = \sum_q \mathcal{I}[DF^{p \rightarrow q}(x, k_T^2) \otimes \sigma^{eq \rightarrow eq} \otimes FF^{q \rightarrow h}(z, p_T^2)]$$

$k_T/p_T$  = transverse momentum of struck/fragmenting quark

$\mathcal{I}[\dots]$  = convolution integral over  $k_T$  and  $p_T$

# Transverse momentum dependent distributions (TMDs)

$$\sigma^{ep \rightarrow eh} = \sum \mathcal{I}[DF^{p \rightarrow q}(x, k_T^2) \otimes \sigma^{eq \rightarrow eq} \otimes FF^{q \rightarrow h}(z, p_T^2)]$$

Distribution functions

$$\begin{aligned} f_1 &= \text{Diagram } 1 \\ g_{1L} &= \text{Diagram } 2 - \text{Diagram } 3 \\ h_{1T} &= \text{Diagram } 4 - \text{Diagram } 5 \\ f_{1T}^\perp &= \text{Diagram } 6 - \text{Diagram } 7 \\ h_1^\perp &= \text{Diagram } 8 - \text{Diagram } 9 \\ h_{1L}^\perp &= \text{Diagram } 10 - \text{Diagram } 11 \end{aligned}$$

<sup>q</sup> leading twist

Fragmentation functions

$$\begin{aligned} D_1 &= \text{Diagram } 12 \\ G_{1L} &= \text{Diagram } 13 - \text{Diagram } 14 \\ H_{1T} &= \text{Diagram } 15 - \text{Diagram } 16 \\ D_{1T}^\perp &= \text{Diagram } 17 - \text{Diagram } 18 \\ H_1^\perp &= \text{Diagram } 19 - \text{Diagram } 20 \\ H_{1L}^\perp &= \text{Diagram } 21 - \text{Diagram } 22 \\ H_{1T}^\perp &= \text{Diagram } 23 - \text{Diagram } 24 \end{aligned}$$

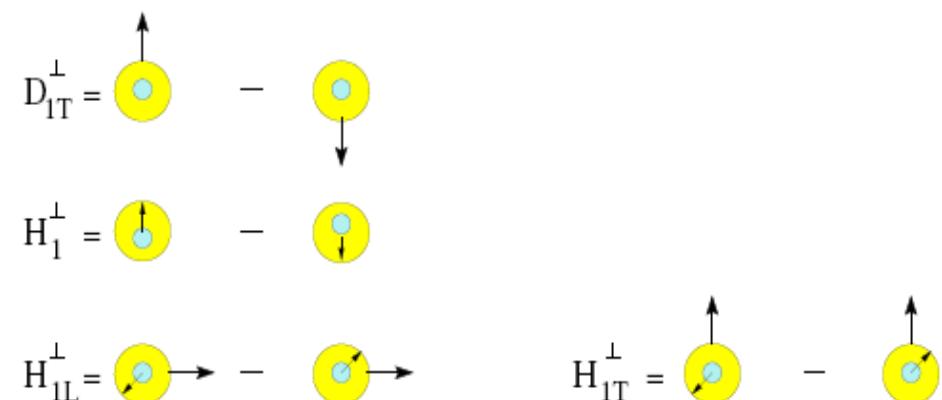
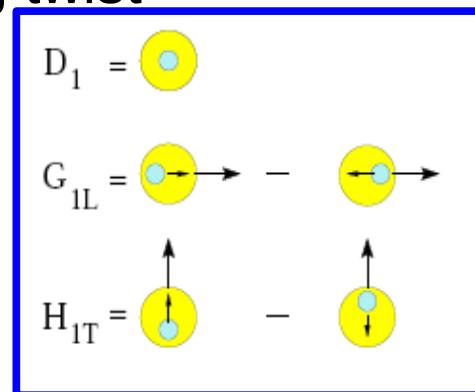
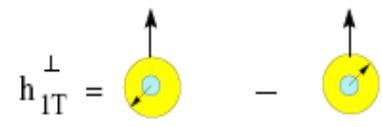
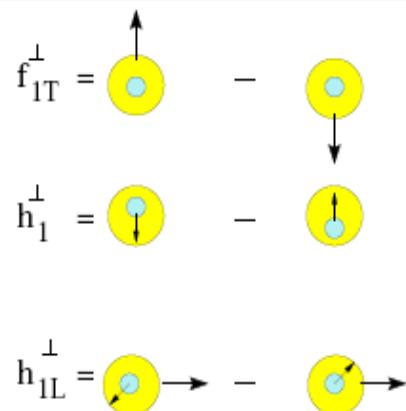
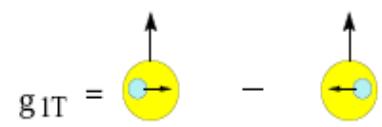
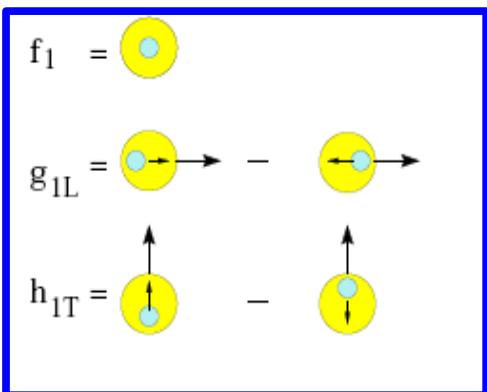
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Distribution functions

${}^q$  leading twist

Fragmentation functions



only distributions that survive integration over transverse momentum

# Transverse momentum dependent distributions (TMDs)

$$\sigma^{ep \rightarrow eh} = \sum \mathcal{I}[DF^{p \rightarrow q}(x, k_T^2) \otimes \sigma^{eq \rightarrow eq} \otimes FF^{q \rightarrow h}(z, p_T^2)]$$

Distribution functions

<sup>q</sup> leading twist

Fragmentation functions

$$f_1 =$$

$$g_{1L} =$$

$$h_{1T} =$$

$$f_{1T}^\perp =$$

$$h_1^\perp =$$

$$h_{1L}^\perp =$$

$$g_{1T} =$$

$$D_1 =$$

$$G_{1L} =$$

$$H_{1T} =$$

$$G_{1T} =$$

$$D_{1T}^\perp =$$

$$H_1^\perp =$$

$$H_{1L}^\perp =$$

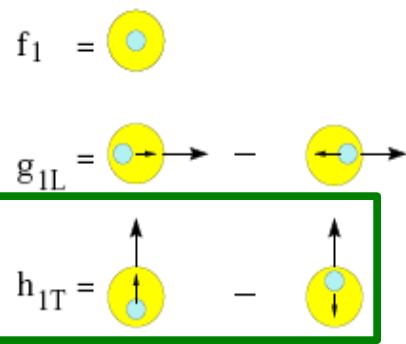
$$H_{1T}^\perp =$$

Chiral odd: involve helicity flip of quark  
appear in pairs in cross section

# Transverse momentum dependent distributions (TMDs)

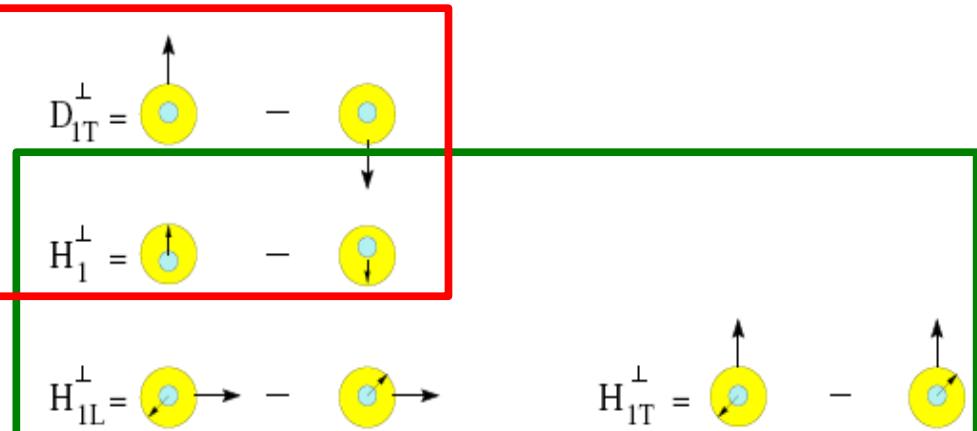
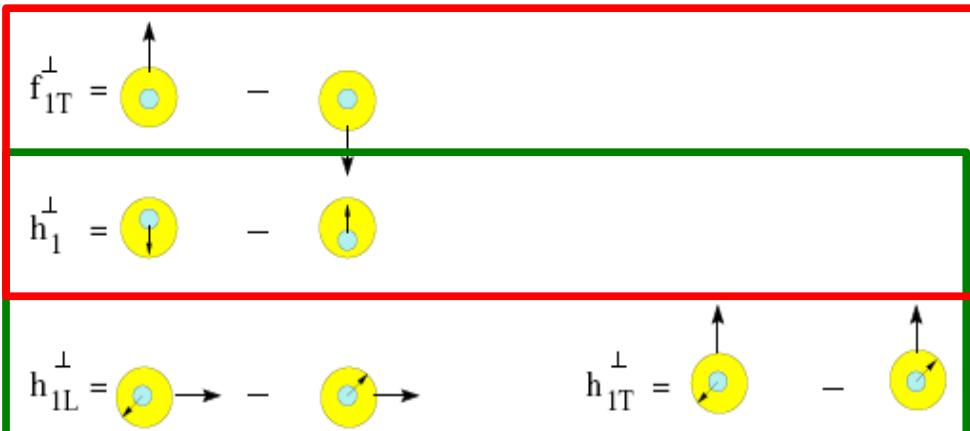
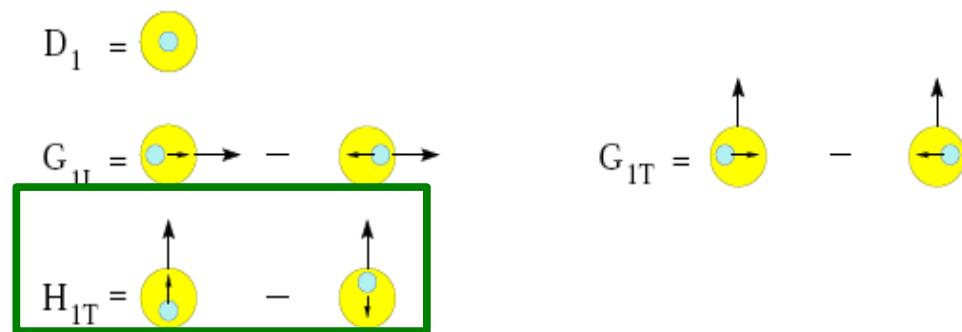
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Distribution functions



<sup>q</sup> leading twist

Fragmentation functions



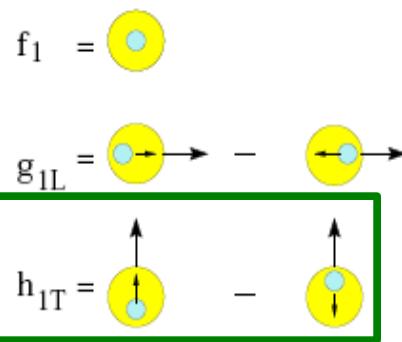
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single in Single Spin Asymmetries (SSAs)

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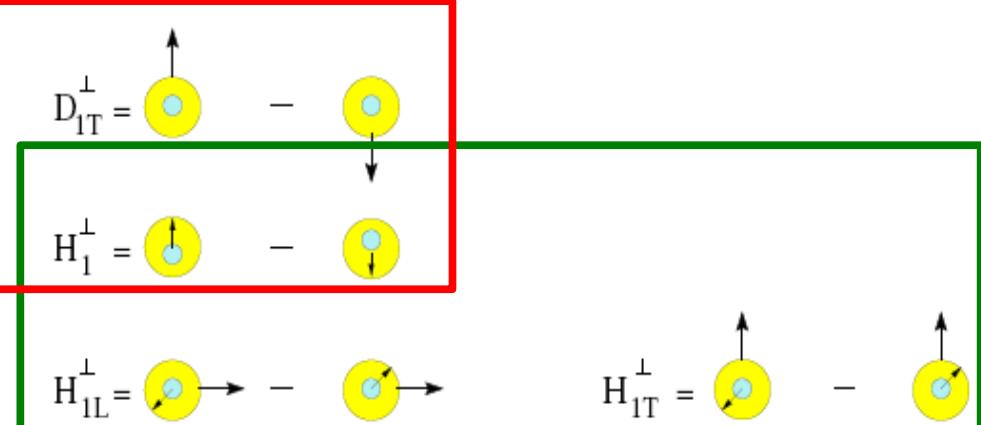
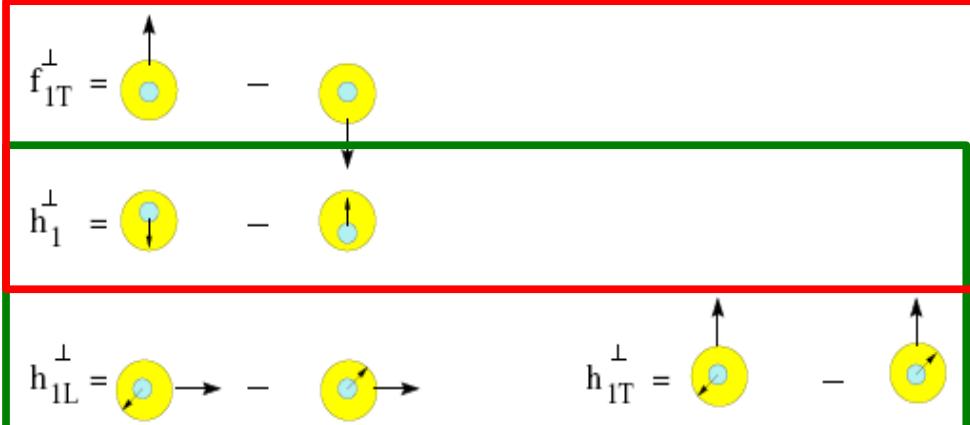
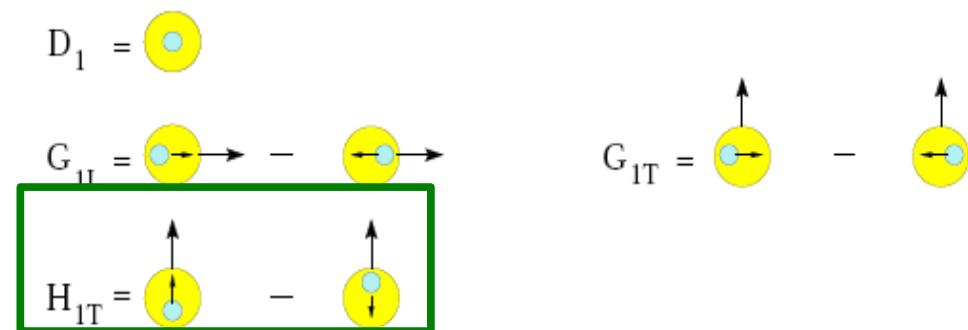
Distribution functions



transversity

<sup>q</sup> leading twist

Fragmentation functions



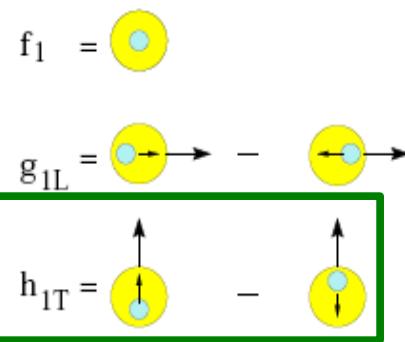
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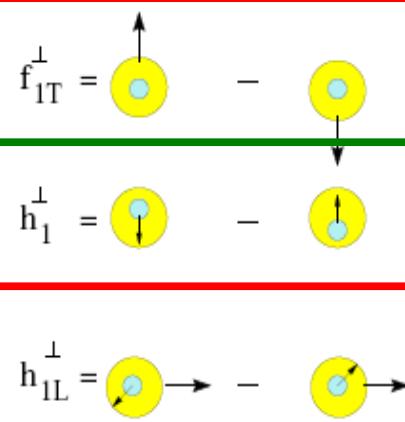
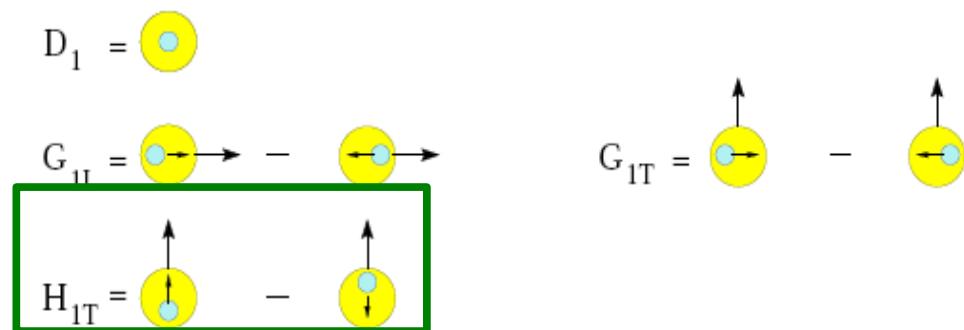
Distribution functions



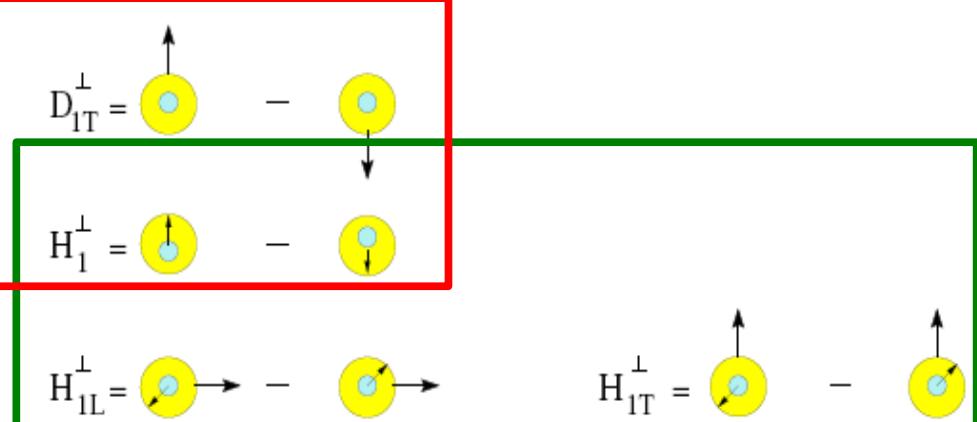
transversity

<sup>q</sup> leading twist

Fragmentation functions



Sivers



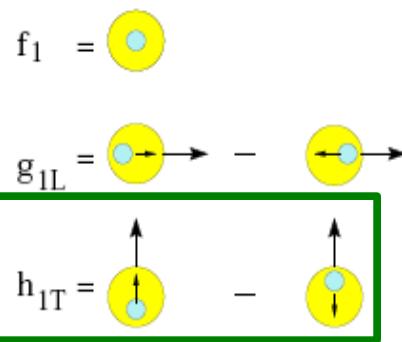
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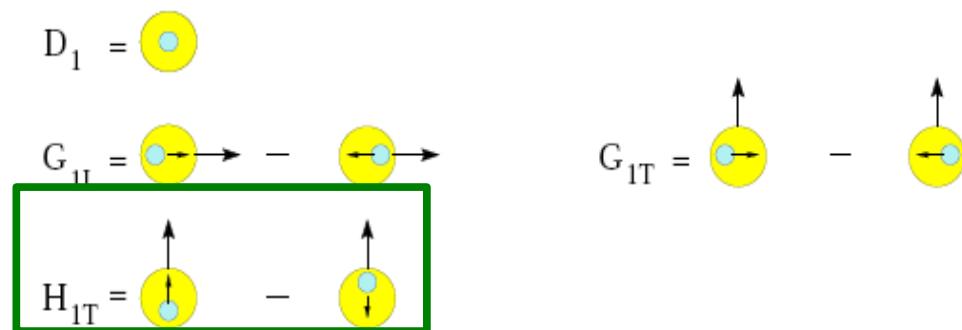
Distribution functions



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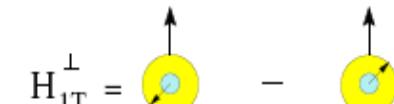
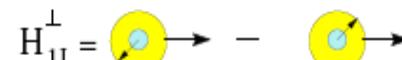
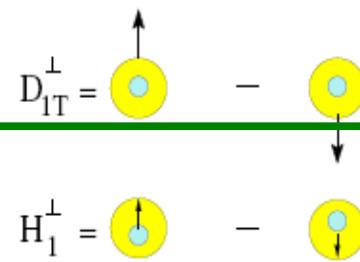
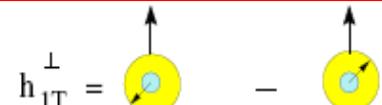
Fragmentation functions



Sivers



Boer-Mulders



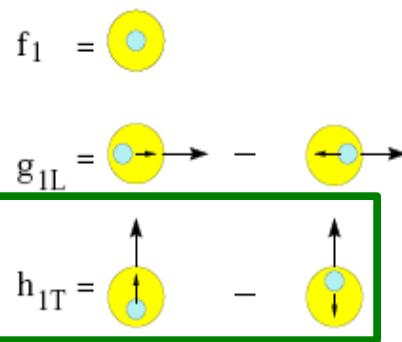
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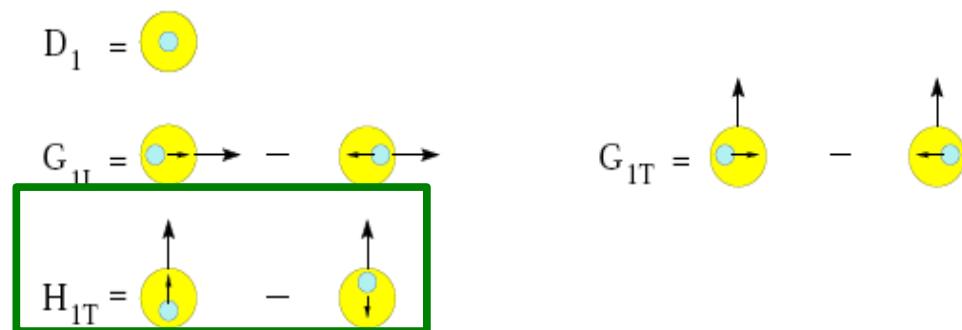
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transversity

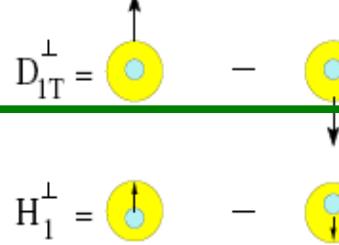
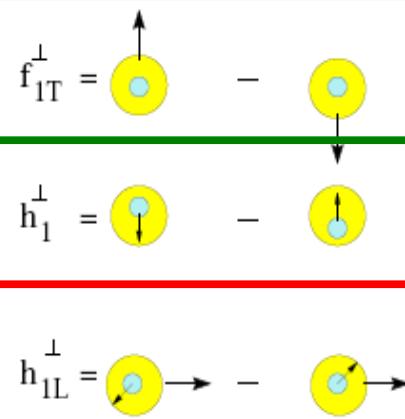
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Fragmentation functions



Sivers

Boer-Mulders

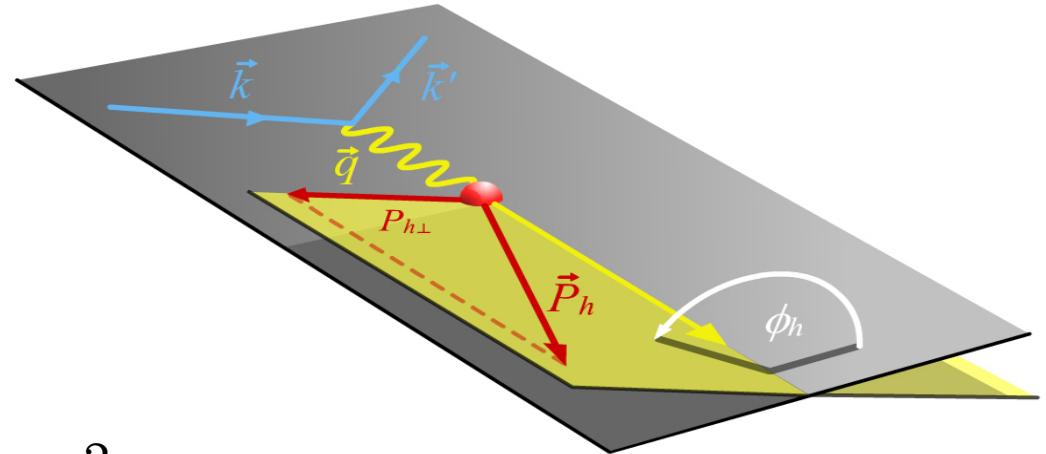


Collins

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# Spin-independent SIDIS cross section

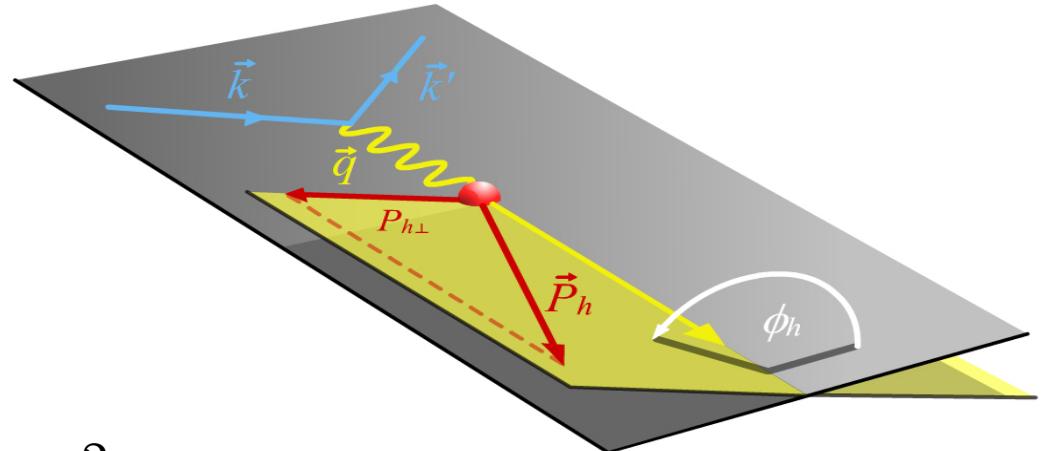


## Non-collinear cross section

$$\frac{d\sigma}{dxdydzdP_{h\perp}^2 d\phi_h} = \frac{\alpha^2}{xyQ^2} \left(1 + \frac{\gamma^2}{2x}\right) \{ A(y) F_{UU,T} + B(y) F_{UU,L} + C(y) \cos \phi_h \textcolor{blue}{F}_{UU}^{\cos \phi_h} + B(y) \cos 2\phi_h \textcolor{green}{F}_{UU}^{\cos 2\phi_h} \}$$

$$\gamma = \frac{2Mx}{Q}, \quad F = F(x, Q, z, \vec{P}_{h\perp})$$

# Spin-independent SIDIS cross section



## Non-collinear cross section

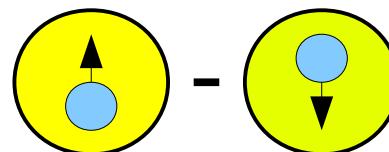
$$\frac{d\sigma}{dxdydzdP_{h\perp}^2 d\phi_h} = \frac{\alpha^2}{xyQ^2} \left(1 + \frac{\gamma^2}{2x}\right) \{ A(y) F_{UU,T} + B(y) F_{UU,L} + C(y) \cos \phi_h F_{UU}^{\cos \phi_h} + B(y) \cos 2\phi_h F_{UU}^{\cos 2\phi_h} \}$$

## leading twist term

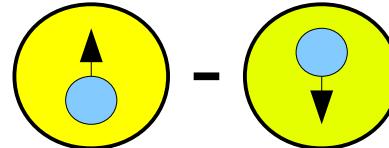
$$\gamma = \frac{2Mx}{Q}, \quad F = F(x, Q, z, \vec{P}_{h\perp})$$

$$F_{UU}^{\cos 2\phi_h} = \mathcal{I} \left[ -\frac{2(\hat{P}_{h\perp} \cdot \vec{k}_T)(\hat{P}_{h\perp} \cdot \vec{p}_T) - \vec{k}_T \cdot \vec{p}_T}{M_h M} h_1^\perp H_1^\perp \right]$$

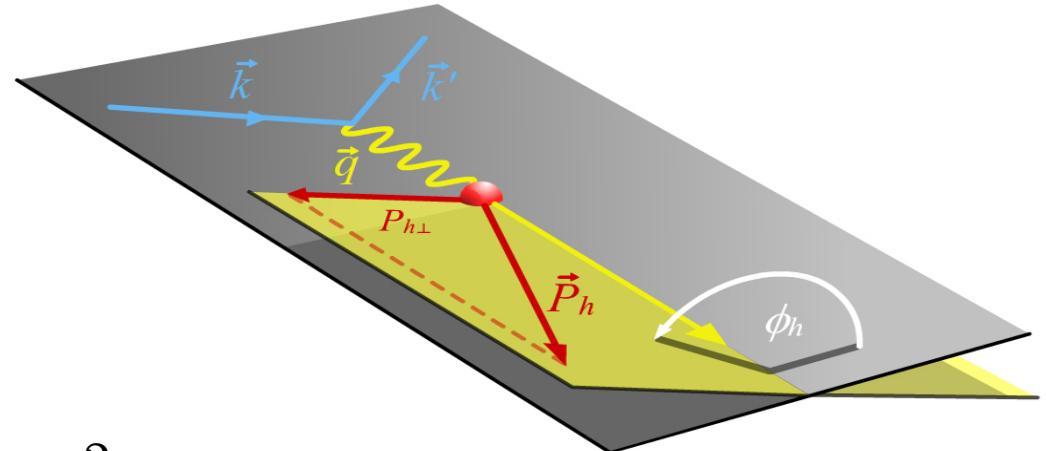
$h_1^\perp$  = Boer-Mulders distribution function



$H_1^\perp$  = Collins fragmentation function



# Spin-independent SIDIS cross section



Non-collinear cross section

$$\frac{d\sigma}{dxdydzdP_{h\perp}^2 d\phi_h} = \frac{\alpha^2}{xyQ^2} \left(1 + \frac{\gamma^2}{2x}\right) \{ A(y) F_{UU,T} + B(y) F_{UU,L} + C(y) \cos \phi_h F_{UU}^{\cos \phi_h} + B(y) \cos 2\phi_h F_{UU}^{\cos 2\phi_h} \}$$

sub-leading twist term

$$\gamma = \frac{2Mx}{Q}, \quad F = F(x, Q, z, \vec{P}_{h\perp})$$

$$F_{UU}^{\cos \phi_h} = \frac{2M}{Q} \mathcal{I} \left[ -\frac{\hat{P}_{h\perp} \cdot \vec{k}_T}{M} f_1 D_1 - \frac{\hat{P}_{h\perp} \cdot \vec{p}_T}{M_h} \frac{k_T^2}{M^2} h_1^\perp H_1^\perp + \dots \right]$$

Cahn effect



quark-gluon-quark correlations

# Extraction of the cosine moments

$$\langle \cos(n\phi_h) \rangle \stackrel{th.}{=} \frac{\int d\phi_h \cos(n\phi_h) \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}{\int d\phi_h \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}$$

↑  
↓

$$\langle \cos(n\phi_h) \rangle \stackrel{exp.}{=} \frac{\int d\phi_h \cos(n\phi_h) \epsilon_{acc}(\phi_h) \epsilon_{rad}(\phi_h) \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}{\int d\phi_h \epsilon_{acc}(\phi_h) \epsilon_{rad}(\phi_h) \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}$$

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Extraction is challenging!

Azimuthal modulations also possible due to

- detector geometrical acceptance
- higher-order QED effects

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fully differential analysis needed  
unfolding procedure with 400 x 12 bins \*

Variable	BINNING						#
	400 kinematic bins x 12 $\phi$ -bins						
x	0.023	0.042	0.078	0.145	0.27	1	5
y	0.3	0.45	0.6	0.7	0.85		4
z	0.2	0.3	0.45	0.6	0.75	1	5
$P_{hT}$	0.05	0.2	0.35	0.5	0.75		4

(\* )see F. Giordano,  
Proceedings of Transversity 2008 Workshop,  
May 28-31 2008, Ferrara, Italy,  
to be published by World Scientific

# Extraction of the cosine moments

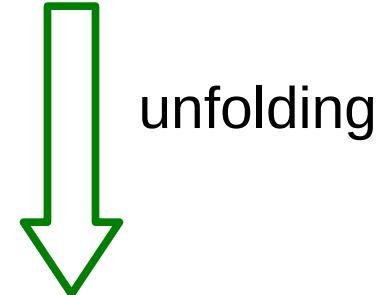
$$\langle \cos(n\phi_h) \rangle \stackrel{th.}{=} \frac{\int d\phi_h \cos(n\phi_h) \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}{\int d\phi_h \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}$$

$$\langle \cos(n\phi_h) \rangle \stackrel{exp.}{=} \frac{\int d\phi_h \cos(n\phi_h) \epsilon_{acc}(\phi_h) \epsilon_{rad}(\phi_h) \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}{\int d\phi_h \epsilon_{acc}(\phi_h) \epsilon_{rad}(\phi_h) \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}$$

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- higher-order QED effects



fully differential analysis needed  
unfolding procedure with  $400 \times 12$  bins \*

$$\langle \cos(n\phi_h) \rangle \Big|_{\text{bin } i} \approx \frac{\int d\phi_h \cos(n\phi_h) \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}{\int d\phi_h \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)} \Big|_{\text{bin } i}$$

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$$\langle \cos(n\phi_h) \rangle \stackrel{exp.}{=} \frac{\int d\phi_h \cos(n\phi_h) \epsilon_{acc}(\phi_h) \epsilon_{rc}(\phi_h) \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}{\int d\phi_h \epsilon_{acc}(\phi_h) \epsilon_{rc}(\phi_h) \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}$$

Extraction is challenging!

Azimuthal modulations also pose a challenge due to

- detector geometrical acceptance
- higher-order QED effects

unfolded

$$\langle \cos(n\phi_h) \rangle \Big|_{\text{bin } i} \approx$$

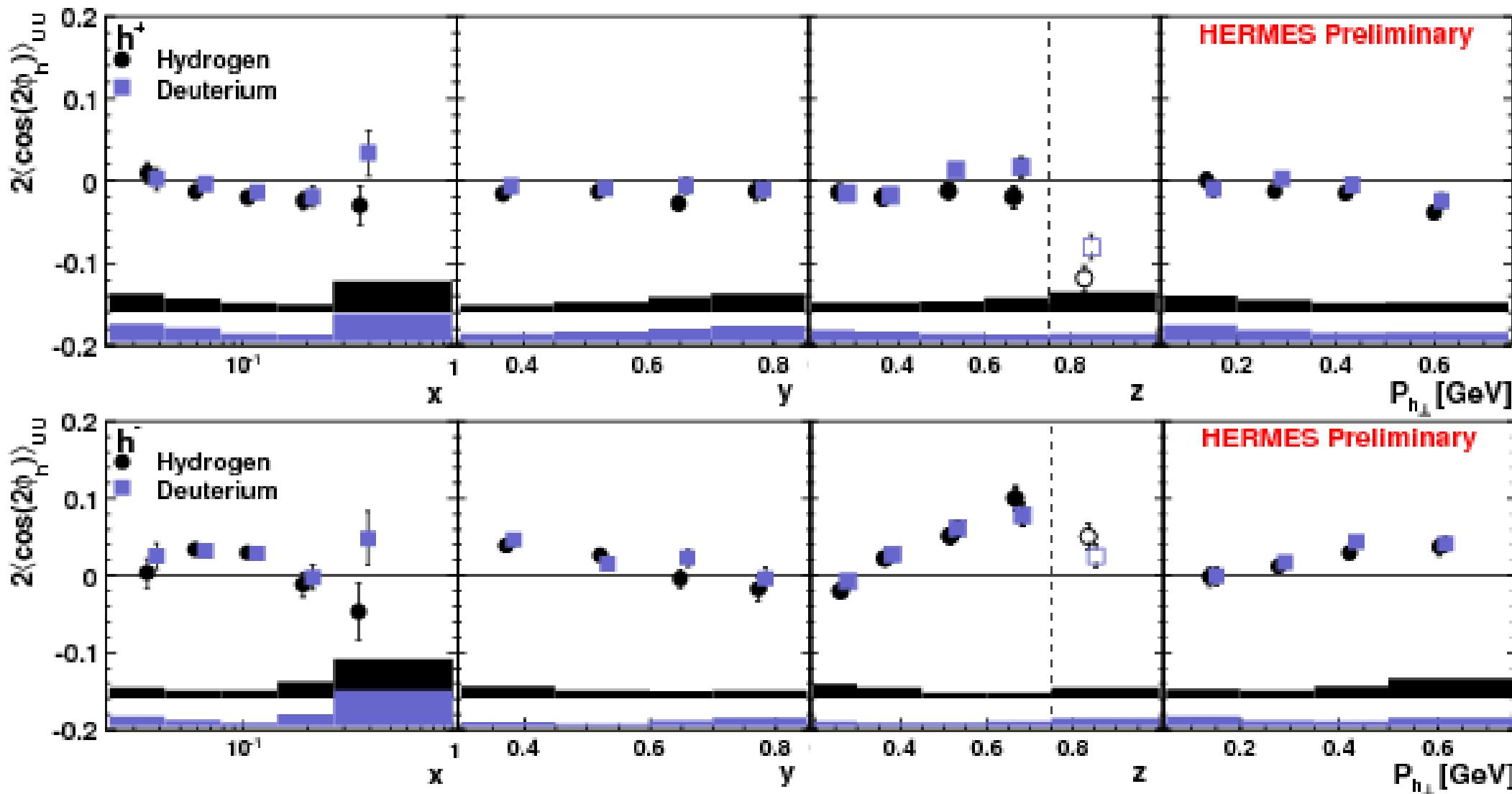
$$\frac{\int d\phi_h \cos(n\phi_h) \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)}{\int d\phi_h \sigma_{UU}(x, y, z, P_{h\perp}^2, \phi_h)} \Big|_{\text{bin } i}$$

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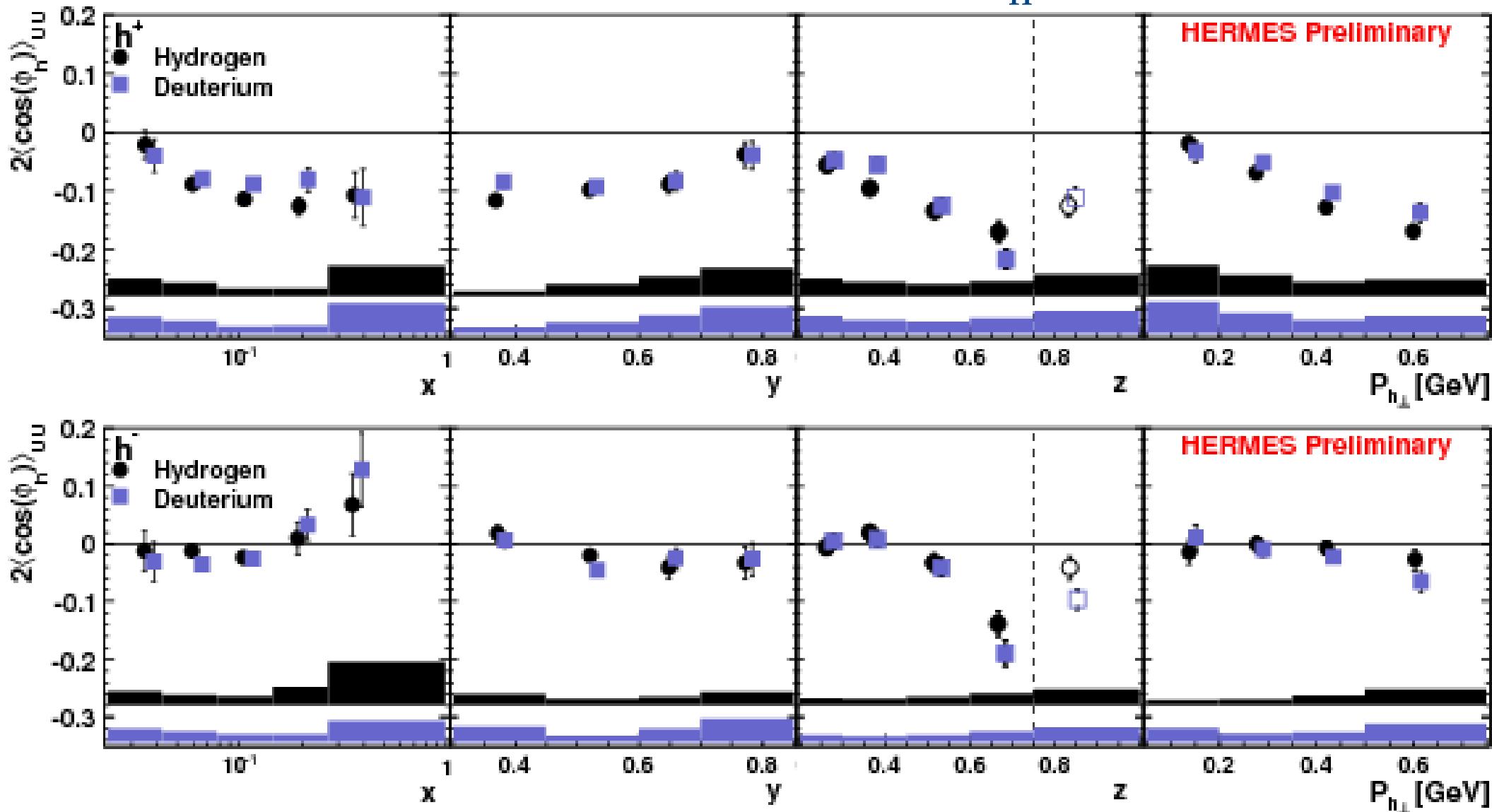
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# Results for $\langle \cos 2\phi_h \rangle$



- evidence for transversely polarized quarks in unpolarized nucleon!
- H<sup>+</sup> and H<sup>-</sup>: opposite sign in agreement with models

# Results for $\langle \cos \phi_h \rangle$



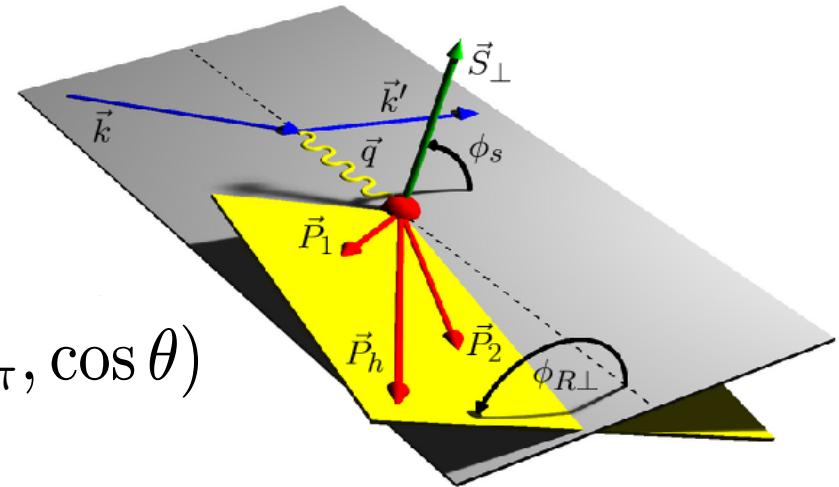
- predictions for Cahn effect: negative for  $h^+$  and  $h^-$
- → also other effects have to be taken into account

# Single Spin Asymmetries

$$A_{UT} = \frac{1}{\langle |S_T| \rangle} \frac{N^\uparrow(\phi, \phi_S) - N^\downarrow(\phi, \phi_S)}{N^\uparrow(\phi, \phi_S) + N^\downarrow(\phi, \phi_S)}$$

## 2-hadron production

$$\sim \sin(\phi_{R\perp} + \phi_S) \sum_q e_q^2 h_{1T}^{q\perp}(x) H_1^{\perp,q}(z, M_{\pi\pi}, \cos\theta)$$



$h_{1T}$ : transversity

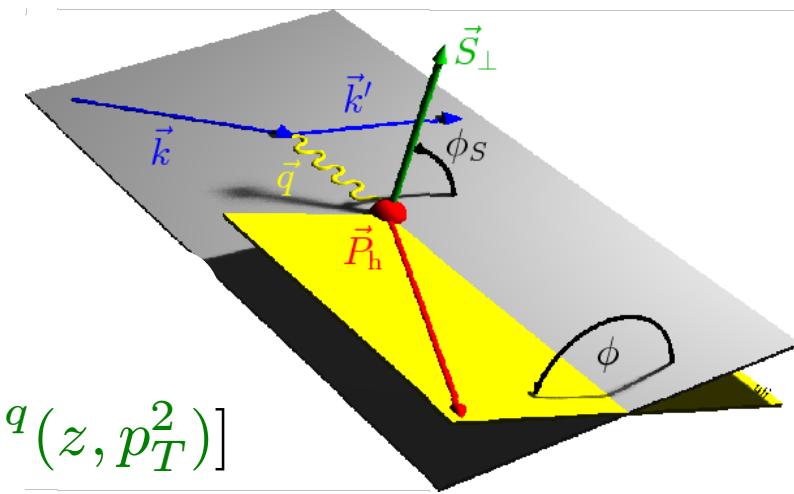
$H_1^\perp$ : Collins fragmentation function

$f_{1T}^{\perp,q}$ : Sivers distribution function

## 1-hadron production

$$\sim \sin(\phi + \phi_S) \sum_q e_q \mathcal{I} \left[ \frac{\vec{p}_T \cdot \hat{P}_{h\perp}}{M_h} h_{1T}^{q\perp}(x, k_T^2) \otimes H_1^{\perp,q}(z, p_T^2) \right]$$

$$+ \sin(\phi - \phi_S) \sum_q e_q \mathcal{I} \left[ \frac{\vec{k}_T \cdot \hat{P}_{h\perp}}{M_h} f_{1T}^{\perp,q}(x, k_T^2) \otimes D_1^q(z, p_T^2) \right]$$



# Single Spin Asymmetries

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## 2-hadron production

$$\sim \sin(\phi_{R\perp} + \phi_S) \sum_q e_q^2 h_{1T}^q(x) H_1^{\perp, q}(\gamma, \pi, \cos \theta)$$

$h_{1T}$ : transversity

$H_1^\perp$ : Collins fragmentation function

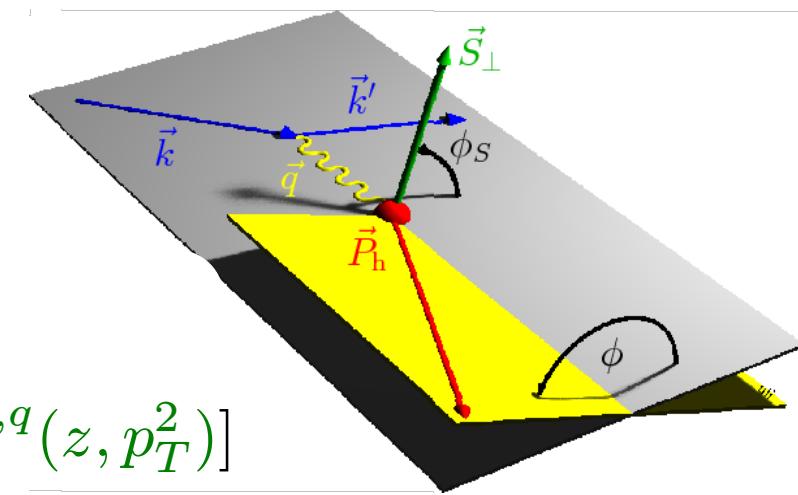
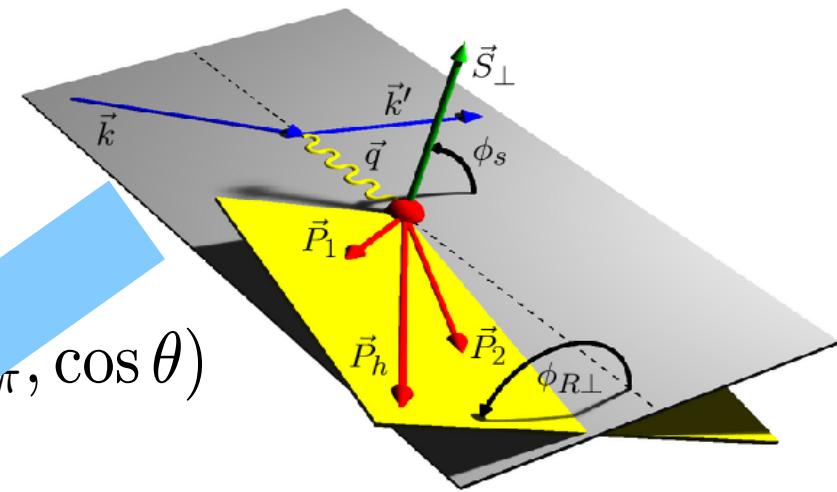
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see talk by M. Diefenthaler

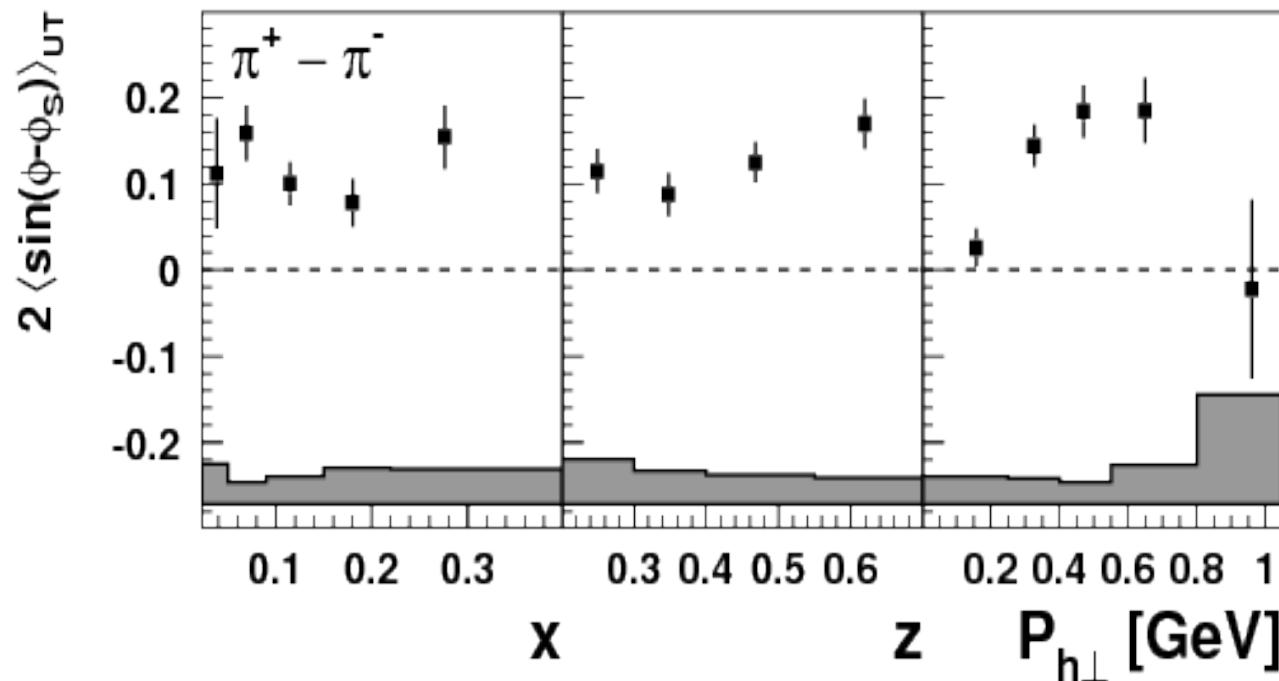


# Sivers distribution for valence quarks

$$A_{UT}^{\pi^+ - \pi^-} = \frac{1}{\langle |S_T| \rangle} \frac{(\sigma_{U\uparrow}^{\pi^+} - \sigma_{U\uparrow}^{\pi^-}) - (\sigma_{U\downarrow}^{\pi^+} - \sigma_{U\downarrow}^{\pi^-})}{(\sigma_{U\uparrow}^{\pi^+} - \sigma_{U\uparrow}^{\pi^-}) + (\sigma_{U\downarrow}^{\pi^+} - \sigma_{U\downarrow}^{\pi^-})}$$

→  $\langle \sin(\phi - \phi_S) \rangle_{UT}^{\pi^+ - \pi^-} = -\frac{4f_{1T}^{\perp, u_v} - f_{1T}^{\perp, d_v}}{4f_1^{u_v} - f_1^{d_v}}$

**HERMES PRELIMINARY 2002-2005**  
lepton beam amplitudes, 8.1% scale uncertainty



- Sivers distribution for d-valence >> u-valence or
- Sivers distribution for u-valence is large & <0 (more likely)

# Generalized parton distributions and Deeply virtual Compton scattering

# Spin decomposition of the nucleon

X. Ji, *Phys. Rev. Lett.* **78**, 610-613 (1997)

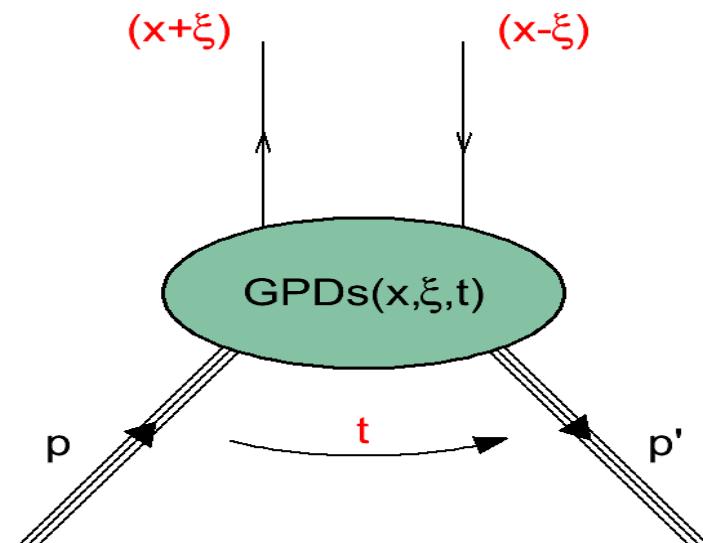
$$\frac{1}{2} = \sum_q J^q + J^g = \frac{1}{2} \Delta \Sigma + \sum_q L^q + J^g$$

$$J^q = \lim_{t \rightarrow 0} \frac{1}{2} \int_{-1}^1 dx x [H^q(x, \xi, t) + E^q(x, \xi, t)]$$

→ possibility to access quark orbital angular momentum

## Generalized Parton Distributions (GPDs)

- 4 twist-2 quark helicity conserving GPDs:  
 $H^q, \tilde{H}^q, E^q, \tilde{E}^q$



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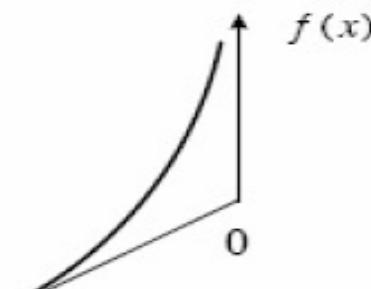
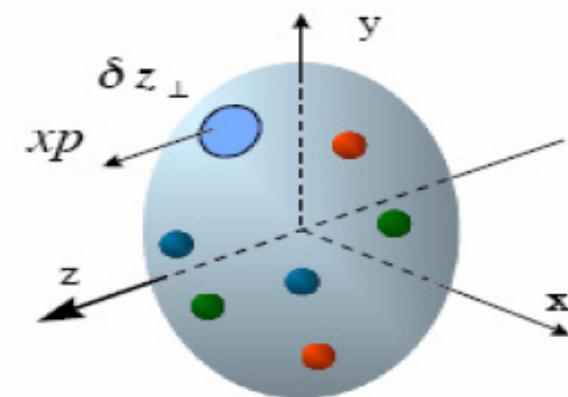
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- forward limit → parton distribution functions

$$H^q(x, 0, 0) = q(x)$$

$$\tilde{H}^q(x, 0, 0) = \Delta q(x)$$



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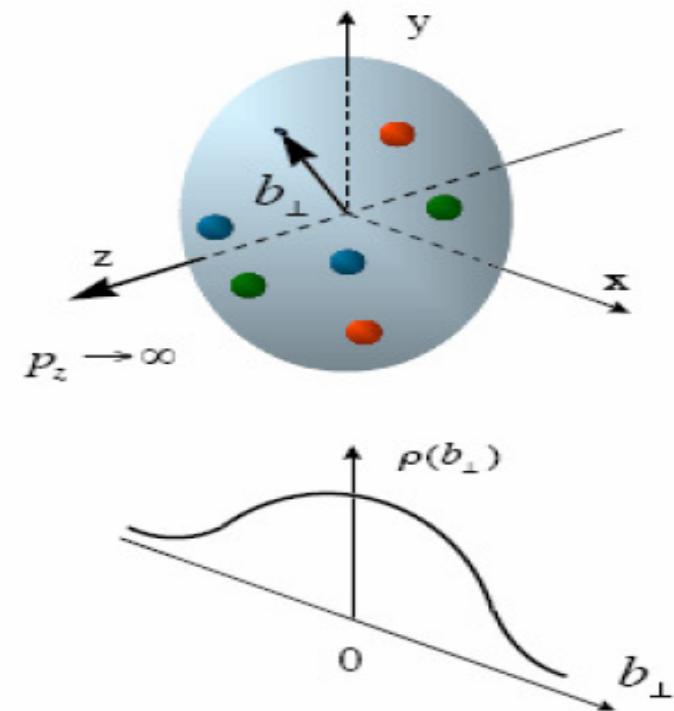
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- forward limit → parton distribution functions
- moments → form factors

$$\int_{-1}^1 dx H^q(x, \xi, t) = F_1^q(t)$$

■  
■  
■



# Spin decomposition of the nucleon

X. Ji, *Phys. Rev. Lett.* **78**, 610-613 (1997)

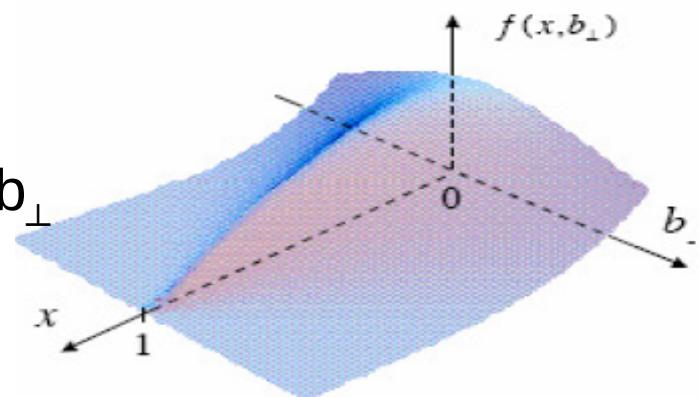
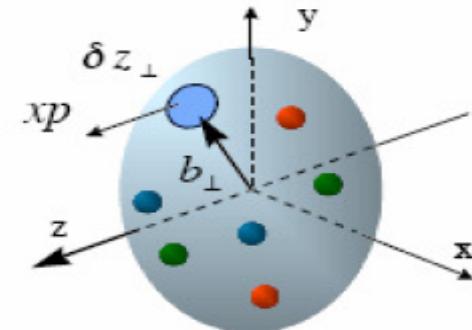
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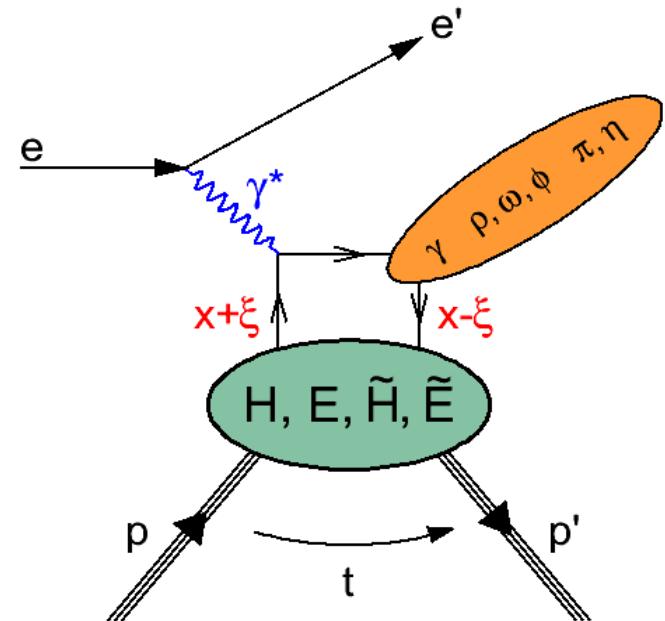
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possibility to access quark orbital angular momentum

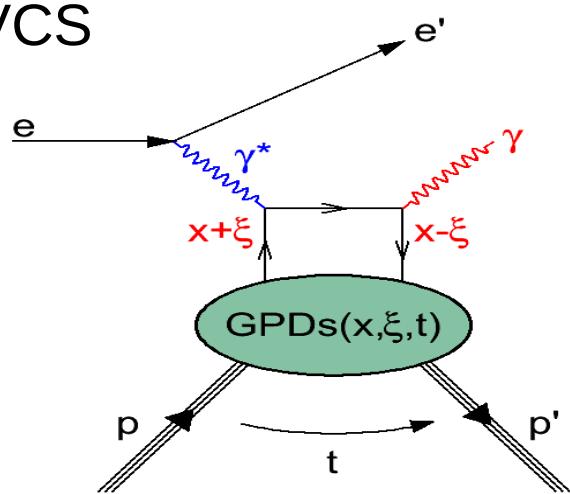
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- access via exclusive meson production and deeply virtual Compton scattering

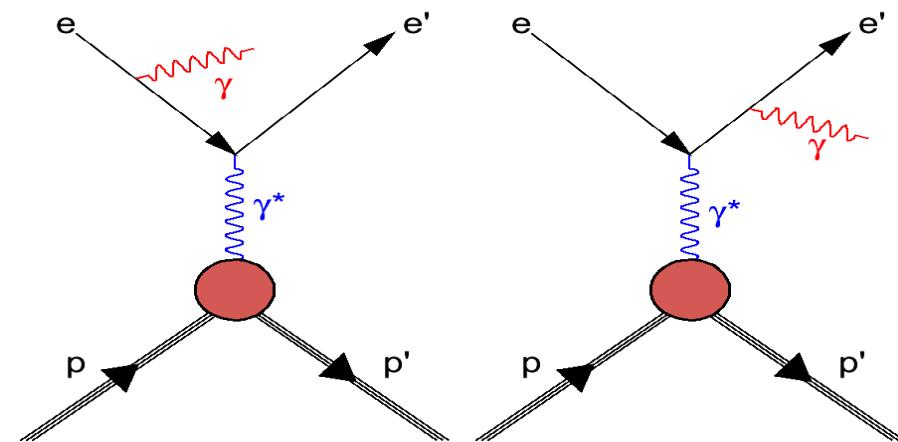


# Deeply virtual Compton scattering

DVCS



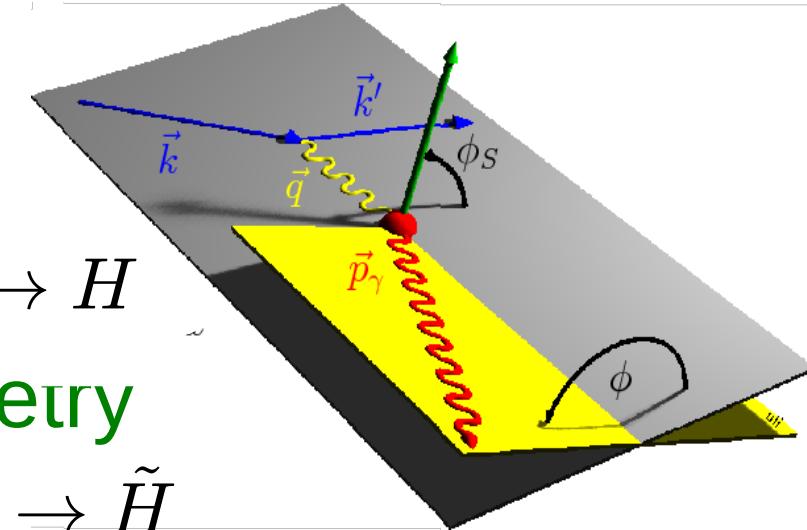
Bethe-Heitler



$$d\sigma \sim |T_{DVCS}|^2 + |T_{BH}|^2 + \underbrace{T_{BH} T_{DVCS}^* + T_{BH}^* T_{DVCS}}_{\text{interference term}}$$

- theoretically cleanest process to access GPDs
- DVCS and BH: indistinguishable → interference
- $T_{BH} \gg T_{DVCS}$
- $T_{BH}$ : calculable (form factors)
- access interference term by measuring azimuthal asymmetries

# Azimuthal asymmetries



- Beam-spin asymmetry

$$d\sigma(\vec{e}, \phi) - d\sigma(\overleftarrow{e}, \phi) \sim \Im[F_1 \mathcal{H}] \sin \phi \rightarrow H$$

- Longitudinal target-spin asymmetry

$$d\sigma(\vec{P}, \phi) - d\sigma(\overleftarrow{P}, \phi) \sim \Im[F_1 \tilde{\mathcal{H}}] \sin \phi \rightarrow \tilde{H}$$

- Beam-charge asymmetry

$$d\sigma(e^+, \phi) - d\sigma(e^-, \phi) \sim \Re[F_1 \mathcal{H}] \cos \phi \rightarrow H$$

- Transverse target-spin asymmetry

$$\begin{aligned} d\sigma(\phi, \phi_S) - d\sigma(\phi, \phi_S + \pi) \sim \\ \Im[F_2 \mathcal{H} - F_1 \mathcal{E}] \sin(\phi - \phi_S) \cos \phi \\ \Im[F_2 \tilde{\mathcal{H}} - F_1 \xi \tilde{\mathcal{E}}] \cos(\phi - \phi_S) \sin \phi \end{aligned}$$

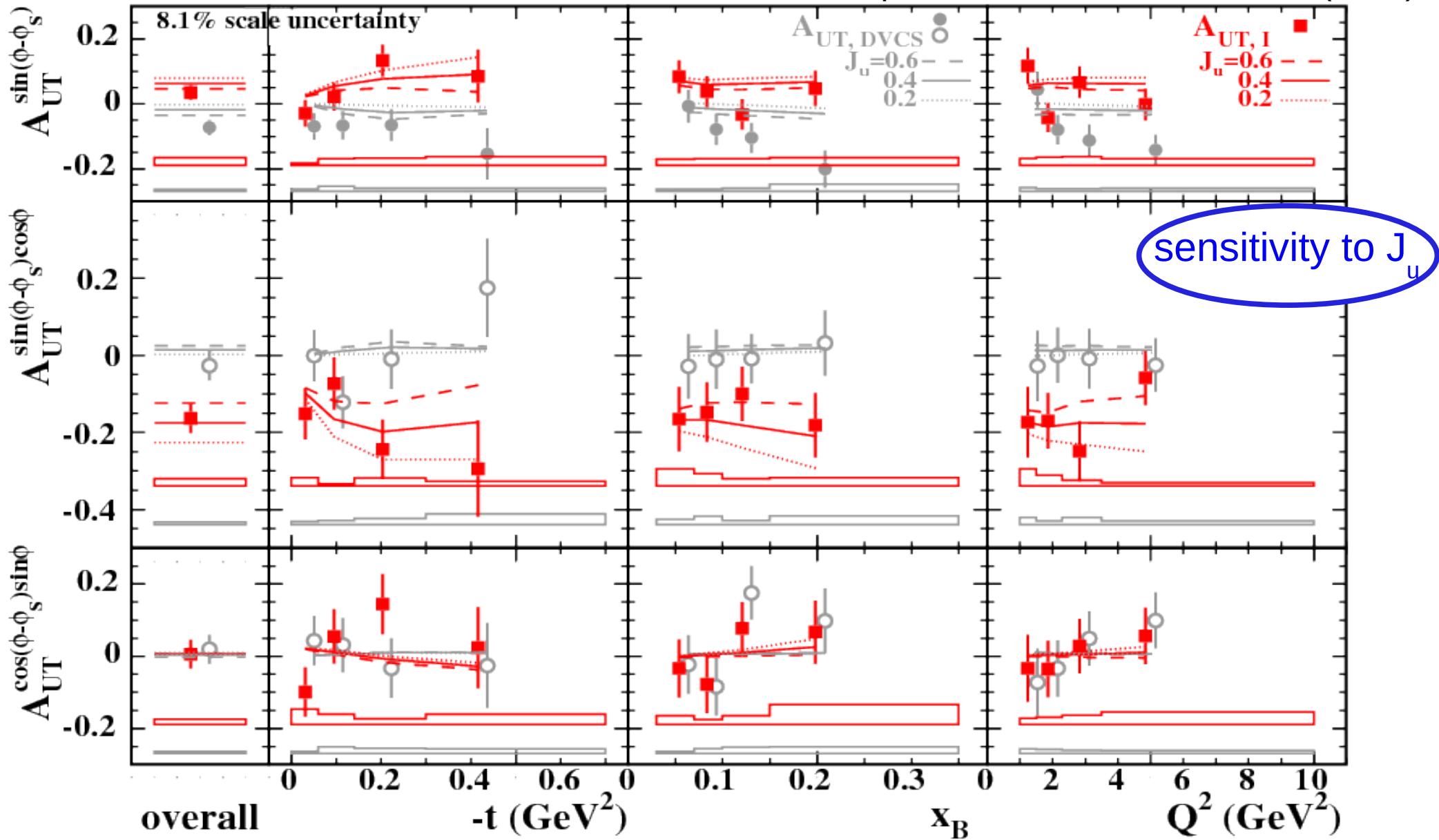


only access to  $E!$

$\mathcal{H}, \mathcal{E}, \tilde{\mathcal{H}}, \tilde{\mathcal{E}}$ : Compton form factors

# Transverse target-spin asymmetry

A. Airapetian et al., JHEP 0806, 066 (2008)

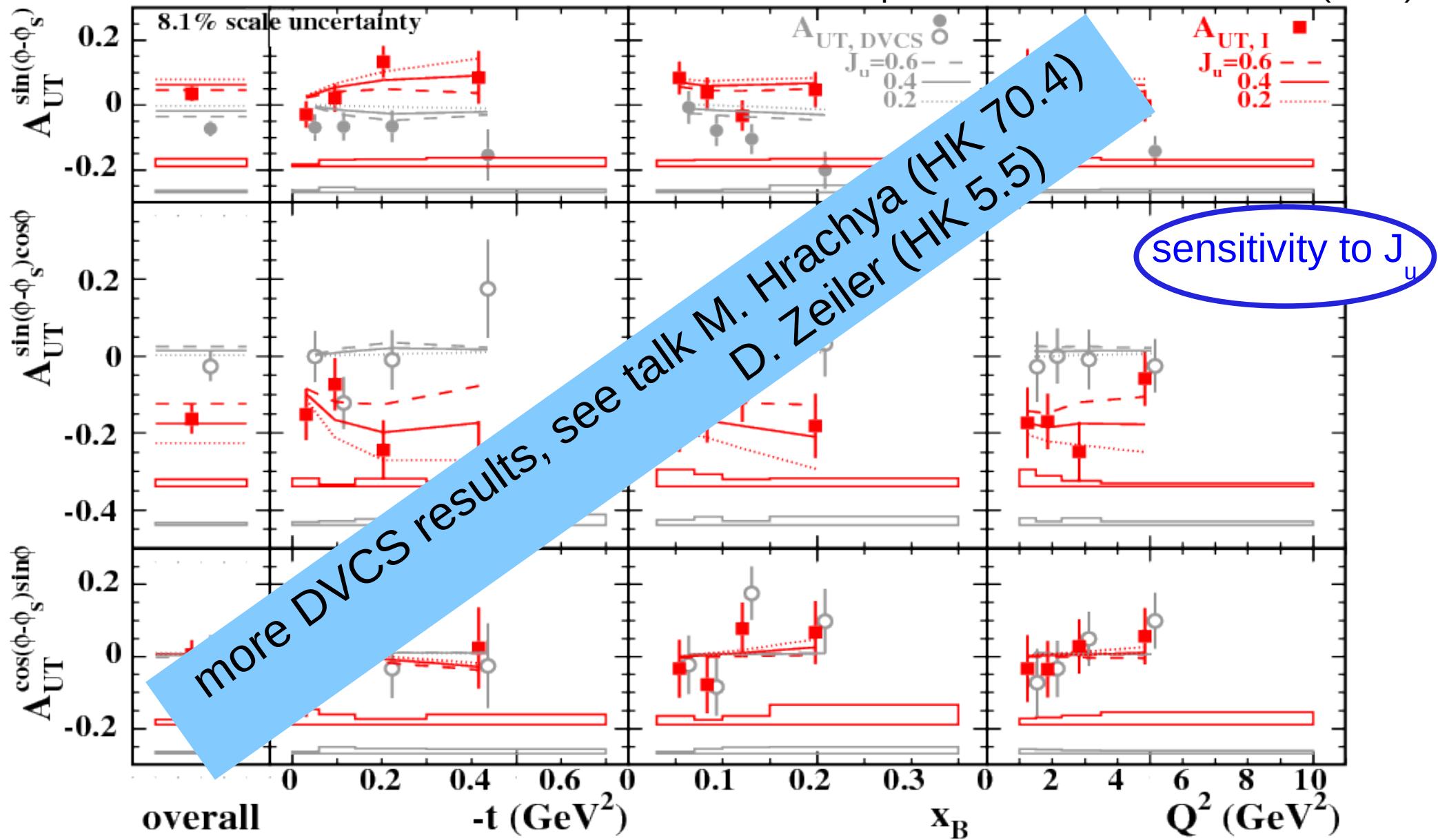


VGG model: Phys. Rev. D60, 094017 (1999)

Prog. Part. Nucl. Phys. 47, 401 (2001)

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# Summary

- strange quark distributions
  - SU(3) symmetry appears to be violated
  - $S(x)$  much softer than light isoscalar
  - $\Delta S(x)$  consistent with 0
- accounting for transverse momentum of parton → azimuthal dependence of unpolarized cross-section  
Boer-Mulders effect: non zero!
- Sivers distribution for valence quarks:  
likely, large and negative for  $u_v$
- generalized parton distributions → access to quark orbital angular momentum:  
transverse target-spin asymmetry sensitive to  $J_u$

