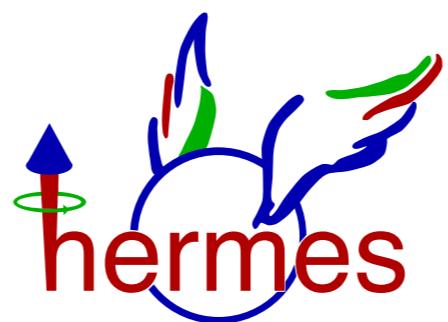


DVCS at HERMES

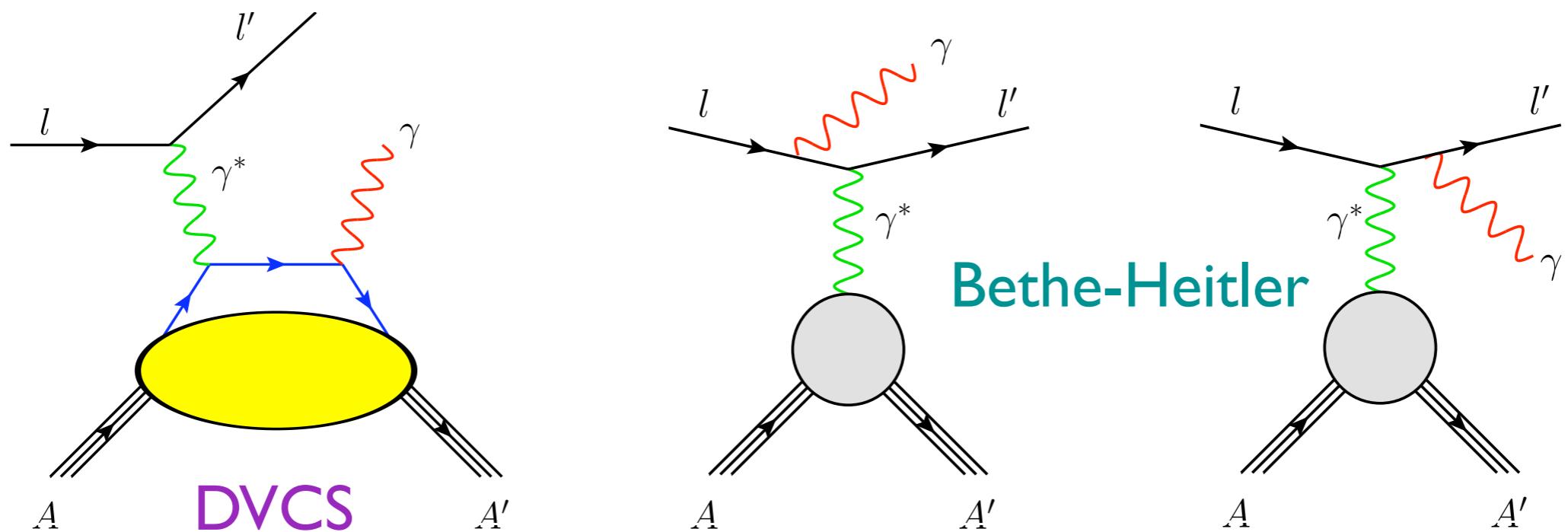
Aram Movsisyan

Yerevan Physics Institute



for the HERMES collaboration

DIS Newport News, 14.04.2011



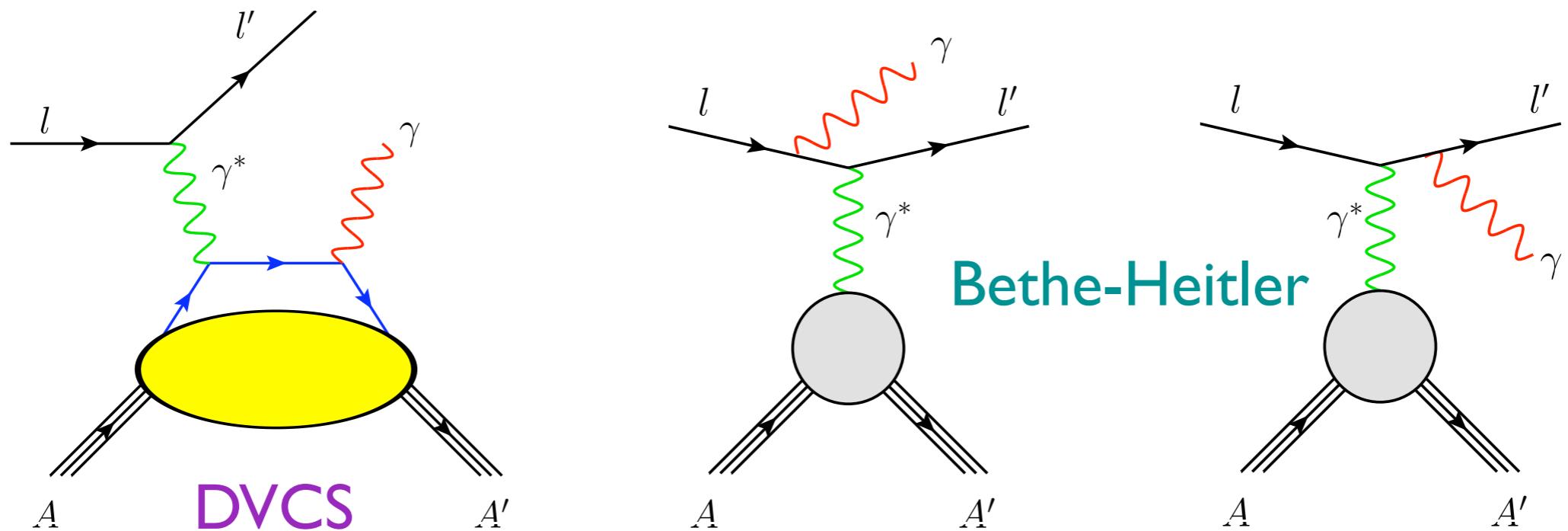
DVCS and **Bethe-Heitler** \Rightarrow Same final state \Rightarrow Interference

$$\frac{d\sigma}{dx_B dQ^2 d|t| d\phi} \propto |\mathcal{T}_{BH}|^2 + |\mathcal{T}_{DVCS}|^2 + \underbrace{\mathcal{T}_{DVCS} \mathcal{T}_{BH}^* + \mathcal{T}_{BH} \mathcal{T}_{DVCS}^*}_I$$

At HERMES kinematics $|\mathcal{T}_{DVCS}|^2 \ll |\mathcal{T}_{BH}|^2$

DVCS amplitudes can be accessed through **Interference**

Interference \Rightarrow non-zero azimuthal asymmetries



$$\frac{d\sigma}{dx_B dQ^2 d|t| d\phi} \propto |\mathcal{T}_{BH}|^2 + |\mathcal{T}_{DVCS}|^2 + \underbrace{\mathcal{T}_{DVCS} \mathcal{T}_{BH}^* + \mathcal{T}_{BH} \mathcal{T}_{DVCS}^*}_I$$

Bethe-Heitler is parametrized in terms of electromagnetic **Form-Factors**

$$\left. \begin{array}{ll} F_1, F_2 & \text{Nucleons} \\ G_1, G_2, G_3 & \text{Deuteron} \end{array} \right\} \Rightarrow \text{BH is calculable in QED}$$

DVCS is parametrized in terms of **Compton Form-Factors**

$$\left. \begin{array}{ll} \mathcal{H}, \mathcal{E}, \tilde{\mathcal{H}}, \tilde{\mathcal{E}} & \text{Nucleons} \\ \mathcal{H}, \mathcal{H}_1, \dots, \mathcal{H}_5, \tilde{\mathcal{H}}_1, \dots, \tilde{\mathcal{H}}_4 & \text{Deuteron} \end{array} \right\} =$$

= convolutions of hard scattering amplitudes and GPD's

- **Beam-Charge asymmetry**

$$\sigma(e^+, \phi) - \sigma(e^-, \phi) \propto \text{Re}[F_1 \mathcal{H}] (\text{Re}[G_1 \mathcal{H}_1])$$

- **Beam-Spin Asymmetry**

$$\sigma(\vec{e}, \phi) - \sigma(\overleftarrow{e}, \phi) \propto \text{Im}[F_1 \mathcal{H}] (\text{Im}[G_1 \mathcal{H}_1])$$

- **Longitudinal Target-Spin Asymmetry**

$$\sigma(\vec{P}, \phi) - \sigma(\overleftarrow{P}, \phi) \propto \text{Im}[F_1 \tilde{\mathcal{H}}] (\text{Im}[G_1 \tilde{\mathcal{H}}_1])$$

- **Longitudinal Double-Spin Asymmetry**

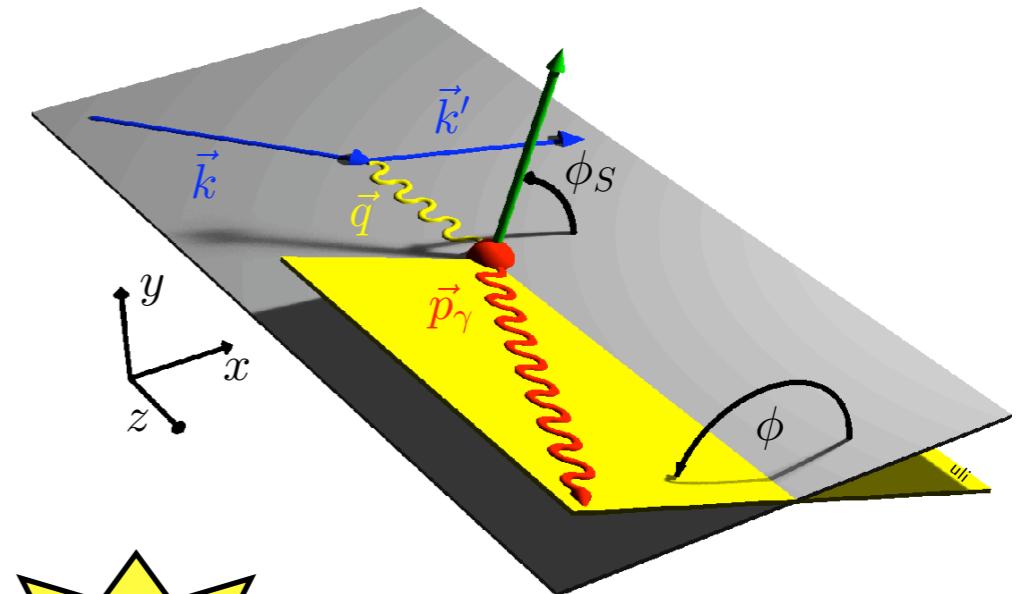
$$\sigma(\vec{P}, \vec{e}, \phi) - \sigma(\vec{P}, \overleftarrow{e}, \phi) \propto \text{Re}[F_1 \tilde{\mathcal{H}}] (\text{Re}[G_1 \tilde{\mathcal{H}}_1])$$

- **Transverse Target-Spin Asymmetry**

$$\sigma(\phi, \phi_S) - \sigma(\phi, \phi_S + \pi) \propto \text{Im}[F_2 \mathcal{H} - F_1 \mathcal{E}]$$

- **Transverse Double-Spin Asymmetry**

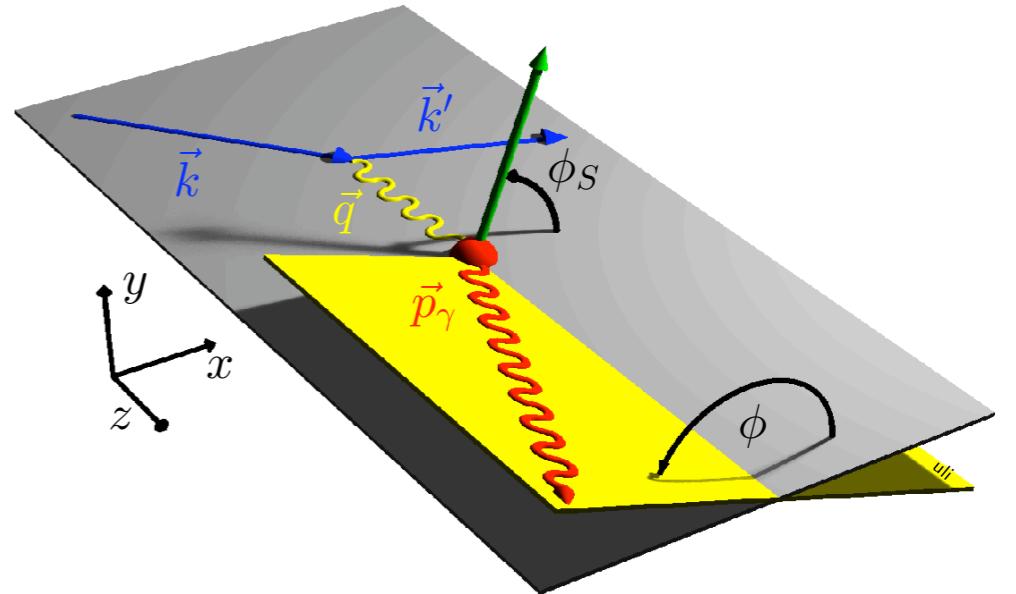
$$\sigma(\vec{e}, \phi, \phi_S) - \sigma(\overleftarrow{e}, \phi, \phi_S + \pi) \propto \text{Re}[F_2 \mathcal{H} - F_1 \mathcal{E}]$$



$$|\mathcal{T}_{\text{BH}}|^2 = \frac{K_{\text{BH}}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^2 c_n^{\text{BH}} \cos(n\phi) + s_1^{\text{BH}} \sin(\phi) \right\}$$

$$|\mathcal{T}_{\text{DVCS}}|^2 = K_{\text{DVCS}} \left\{ \sum_{n=0}^2 c_n^{\text{DVCS}} \cos(n\phi) + \sum_{n=1}^2 s_n^{\text{DVCS}} \sin(n\phi) \right\}$$

$$\mathcal{I} = -\frac{K_I e_\ell}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \left\{ \sum_{n=0}^3 c_n^{\text{I}} \cos(n\phi) + \sum_{n=1}^3 s_n^{\text{I}} \sin(n\phi) \right\}$$



Longitudinally polarized target:

$$\left. \begin{aligned} c_n &= c_{n,\text{unp}} + \lambda \Lambda c_{n,\text{LP}} \\ s_n &= \lambda s_{n,\text{unp}} + \Lambda s_{n,\text{LP}} \end{aligned} \right\} \text{Spin - 1/2}$$

λ - Beam helicity

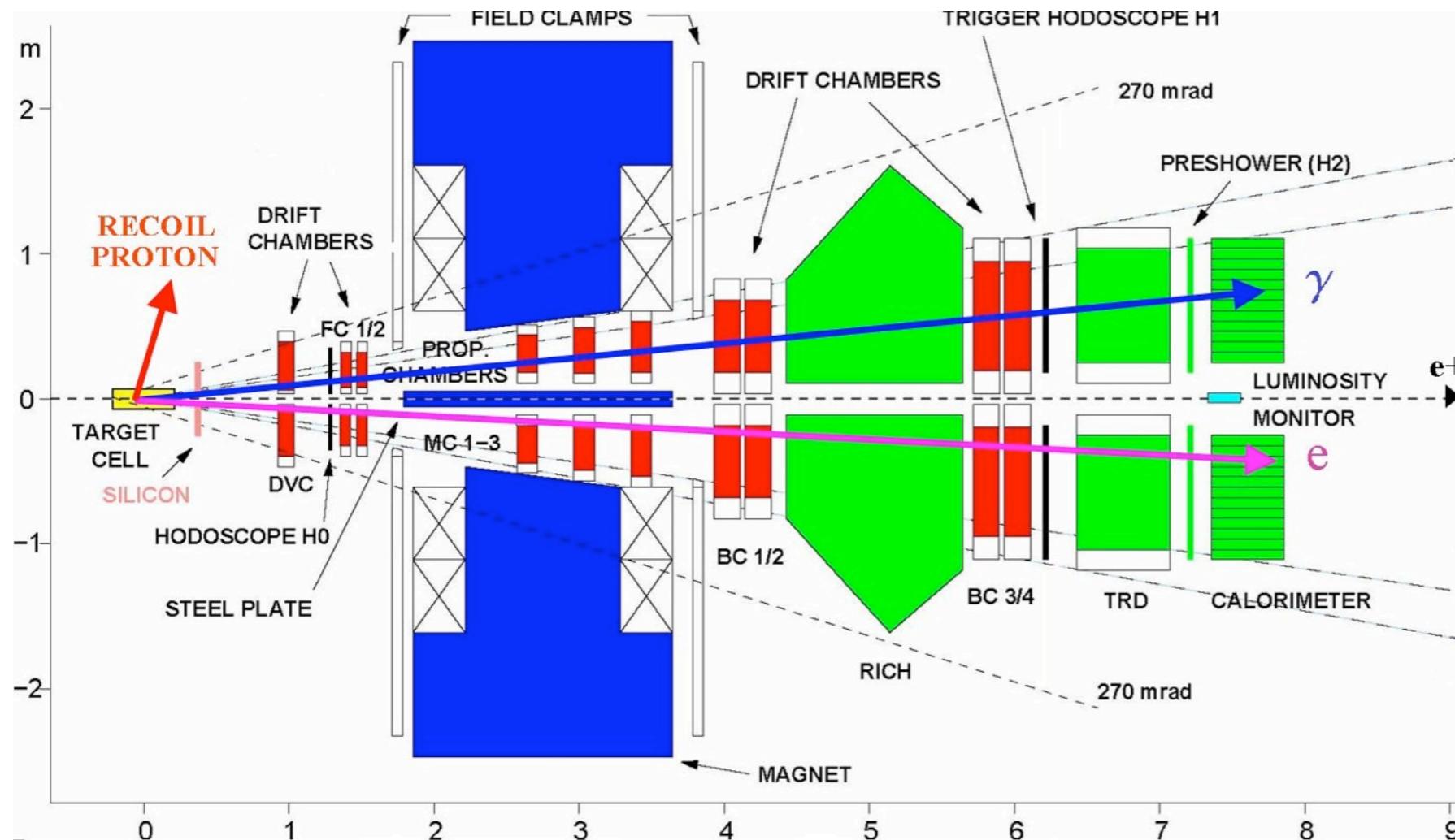
Λ - Target spin projection

e_ℓ - Beam charge

$$\left. \begin{aligned} c_n &= \frac{3}{2} \Lambda^2 c_{n,\text{unp}} + \lambda \Lambda c_{n,\text{LP}} + \left(1 - \frac{3}{2} \Lambda^2\right) c_{n,\text{LLP}} \\ s_n &= \frac{3}{2} \lambda \Lambda^2 s_{n,\text{unp}} + \Lambda s_{n,\text{LP}} + \left(1 - \frac{3}{2} \Lambda^2\right) \lambda s_{n,\text{LLP}} \end{aligned} \right\} \text{Spin - I}$$

Transversely polarized target:

$$\left. \begin{aligned} c_n &= c_{n,\text{unp}} + \Lambda c_{n,\text{UT}} + \lambda \Lambda c_{n,\text{LT}} \\ s_n &= \lambda s_{n,\text{unp}} + \Lambda s_{n,\text{UT}} + \lambda \Lambda s_{n,\text{LT}} \end{aligned} \right\} \text{Spin - 1/2}$$



Longitudinally polarized
 e^+/e^- Beam
 27.6 GeV

- 1996-1997 Longitudinally Polarized **Hydrogen** (e^+ Beam) ≈ 3 M DIS
- 2002-2005 Transversely Polarized **Hydrogen** (e^+/e^- Beam) ≈ 10 M DIS
- 1998-2000 Longitudinally Polarized **Deuterium** (e^+/e^- Beam) ≈ 6 M DIS
- 1996-2005 Unpolarized **Hydrogen** (e^+/e^- Beam) ≈ 17 M DIS
- 1996-2005 Unpolarized **Deuterium** (e^+/e^- Beam) ≈ 10 M DIS
- 2006-2007 Unpolarized **Hydrogen** (e^+/e^- Beam) ≈ 40 M DIS

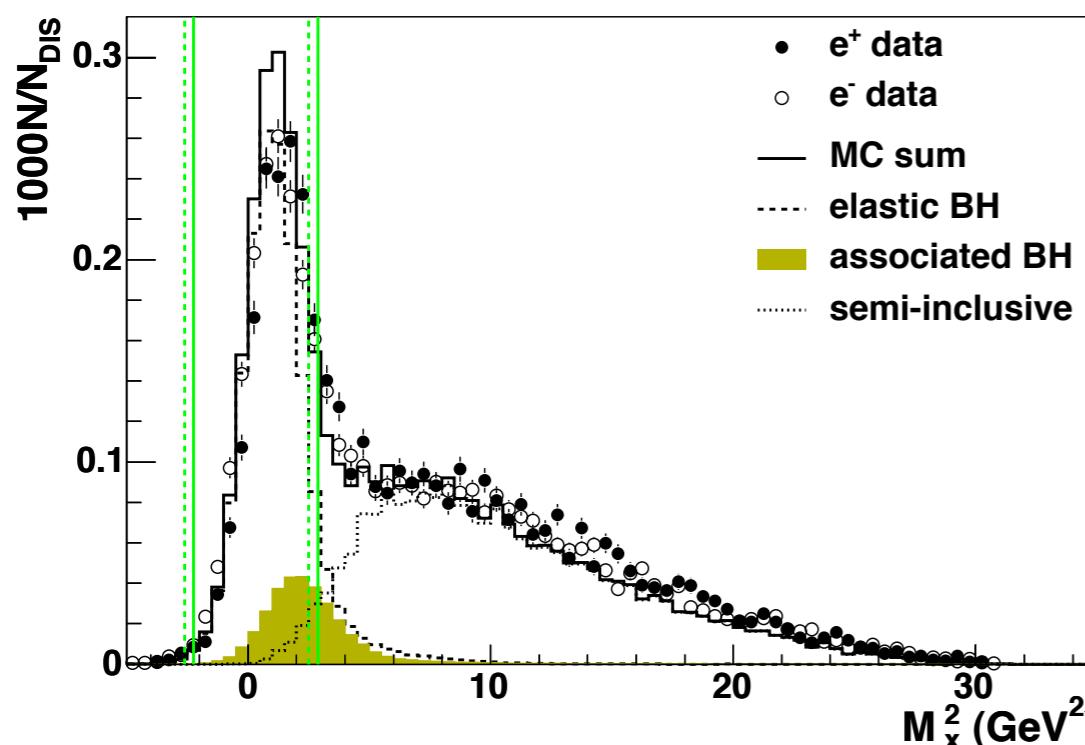
Event Selection

- Events with one DIS lepton and one trackless cluster in the calorimeter.
- Recoiling nucleon/nucleus was not detected
 \Rightarrow Exclusivity via missing mass technique: $M_x^2 = (P + q - q')^2$

$$W^2 > 9 \text{ GeV}^2, \quad \nu < 22 \text{ GeV}$$

$$0.03 < x_B < 0.35, \quad 1 < Q^2 < 10 \text{ GeV}^2$$

$$-t < 0.7 \text{ GeV}^2, \quad E_\gamma > 5 \text{ GeV}$$



Proton:

- Elastic; $ep \rightarrow ep\gamma$
- Associated; mainly $ep \rightarrow e\Delta^+\gamma$
- Semi-Inclusive; mainly $ep \rightarrow e\pi^0 X$

$$-2.25 \text{ GeV}^2 < M_x^2 < 2.89 \text{ GeV}^2$$

Associated cannot be resolved \rightarrow defined as a part of signal.

Event Selection

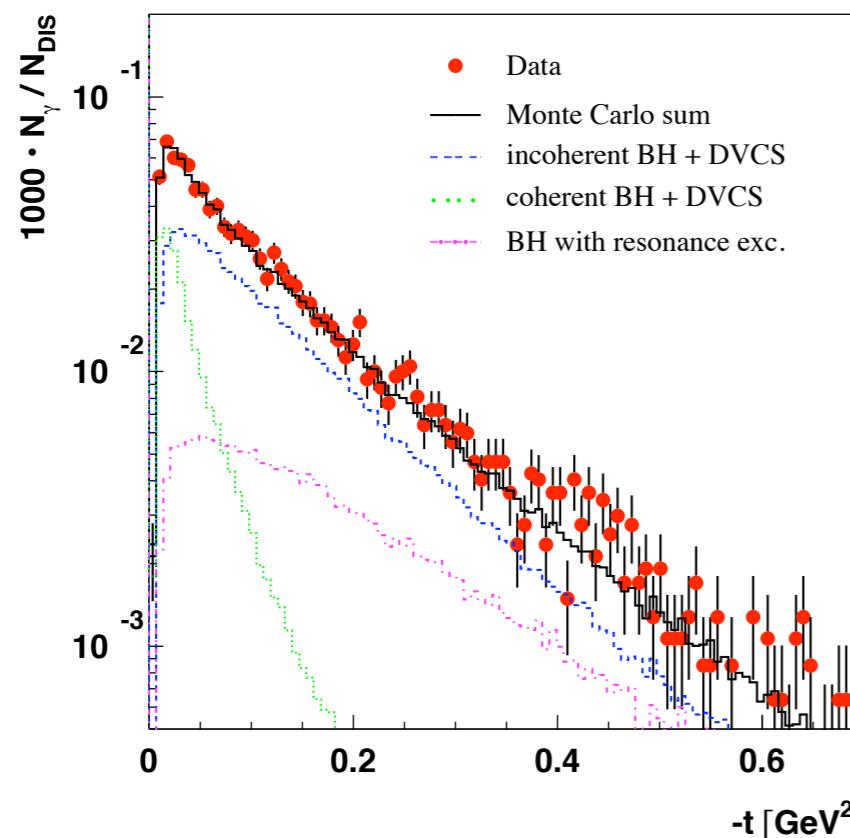
- Events with one DIS lepton and one trackless cluster in the calorimeter.
- Recoiling nucleon/nucleus was not detected

→ Exclusivity via missing mass technique: $M_x^2 = (P + q - q')^2$

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$$0.03 < x_B < 0.35, \quad 1 < Q^2 < 10 \text{ GeV}^2$$

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$$-2.25 \text{ GeV}^2 < M_x^2 < 2.89 \text{ GeV}^2$$

Associated cannot be resolved → defined as a part of signal.

Proton:

- Elastic; $ep \rightarrow ep\gamma$
- Associated; mainly $ep \rightarrow e\Delta^+\gamma$
- Semi-Inclusive; mainly $ep \rightarrow e\pi^0 X$

Deuteron:

- Elastic (Coherent); $ed \rightarrow ed\gamma$
- Quasi-Elastic; $ed \rightarrow epn\gamma$
- Associated; $eN \rightarrow eN^*\gamma$
- Semi-Inclusive; $eN \rightarrow e\pi^0 X$

Distribution in the expectation value of measured yield

$$\langle \mathcal{N}(e_\ell, P_l, S_t, \phi, \phi_S) \rangle \propto \sigma_{UU}(\phi) [1 + e_\ell \mathcal{A}_C + P_l \mathcal{A}_{LU}^{DVCS} + e_\ell P_l \mathcal{A}_{LU}^I + S_t \mathcal{A}_{UT}^{DVCS} + e_\ell S_t \mathcal{A}_{UT}^I + P_l S_t \mathcal{A}_{LT}^{BH+DVCS} + e_\ell P_l S_t \mathcal{A}_{LT}^I]$$

Combined e^+ and e^- data => allow to separate pure beam(target) polarization dependent parts of the cross section from that convoluted with beam charge .

e_ℓ - Beam Charge

P_l - Beam Polarization

S_t - Target Polarization

Distribution in the expectation value of measured yield

$$\langle \mathcal{N}(e_\ell, P_l, S_t, \phi, \phi_S) \rangle \propto \sigma_{UU}(\phi) [1 + e_\ell \mathcal{A}_C + \boxed{P_l \mathcal{A}_{LU}^{DVCS} + e_\ell P_l \mathcal{A}_{LU}^I} \\ + \boxed{S_t \mathcal{A}_{UT}^{DVCS} + e_\ell S_t \mathcal{A}_{UT}^I} \\ + \boxed{P_l S_t \mathcal{A}_{LT}^{BH+DVCS} + e_\ell P_l S_t \mathcal{A}_{LT}^I}]$$

Combined e^+ and e^- data => allow to separate pure beam(target) polarization dependent parts of the cross section from that convoluted with beam charge .

Expansion of the asymmetries:

$$\mathcal{A} \approx \sum_{n=0}^N A^{\cos(n\phi)} \cos(n\phi) \quad \text{or} \quad \mathcal{A} \approx \sum_{n=1}^N A^{\sin(n\phi)} \sin(n\phi)$$

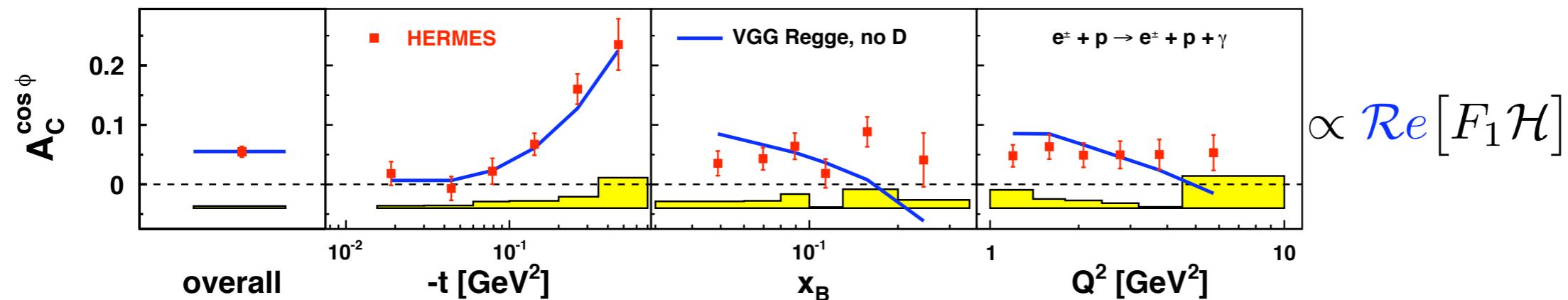
e_ℓ - Beam Charge

P_l - Beam Polarization

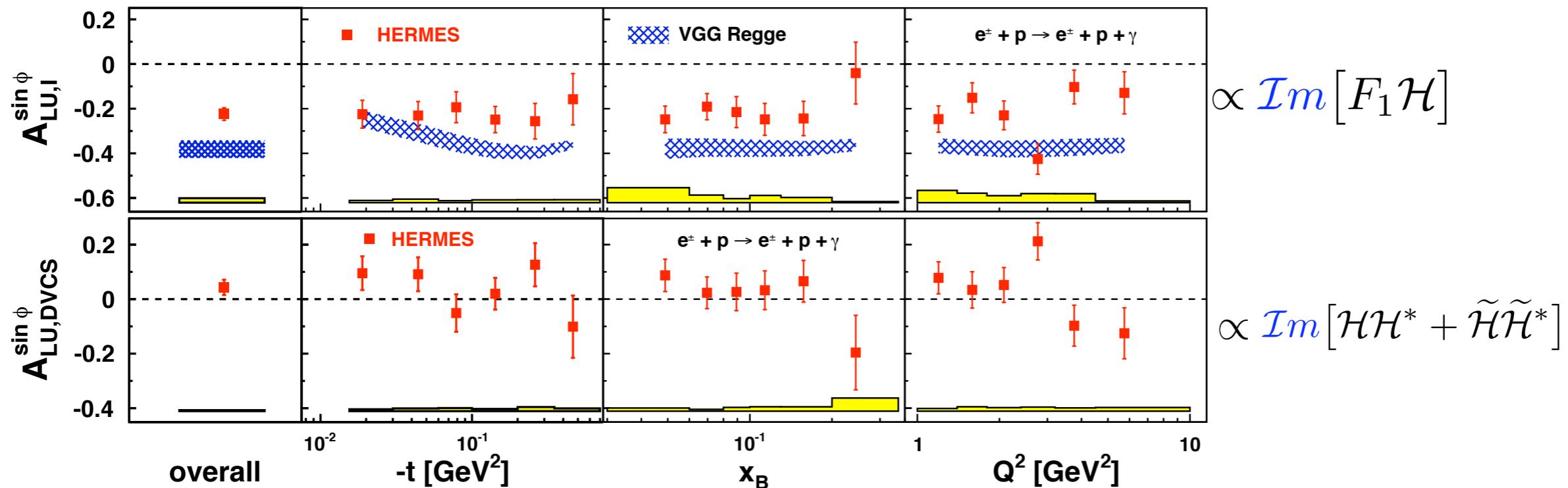
S_t - Target Polarization

Asymmetry amplitudes are extracted simultaneously with maximum likelihood method.

$$\mathcal{A}_C(\phi) = \frac{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) - (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}$$

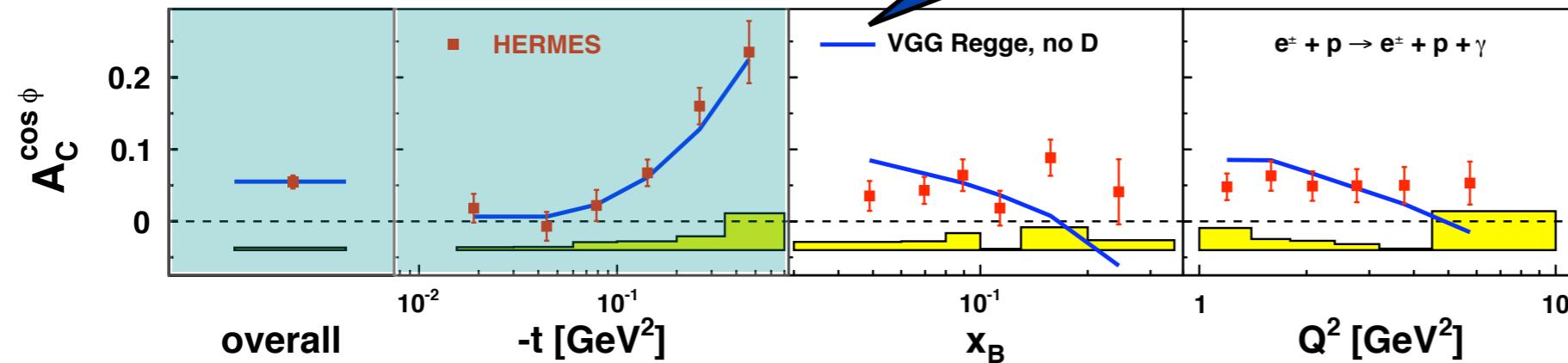


$$\mathcal{A}_{LU}^{I,DVCS}(\phi) = \frac{(\sigma^{+\rightarrow} - \sigma^{+\leftarrow})^+ - (\sigma^{-\rightarrow} - \sigma^{-\leftarrow})^-}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}$$



$$\mathcal{A}_C(\phi) = \frac{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) - (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}$$

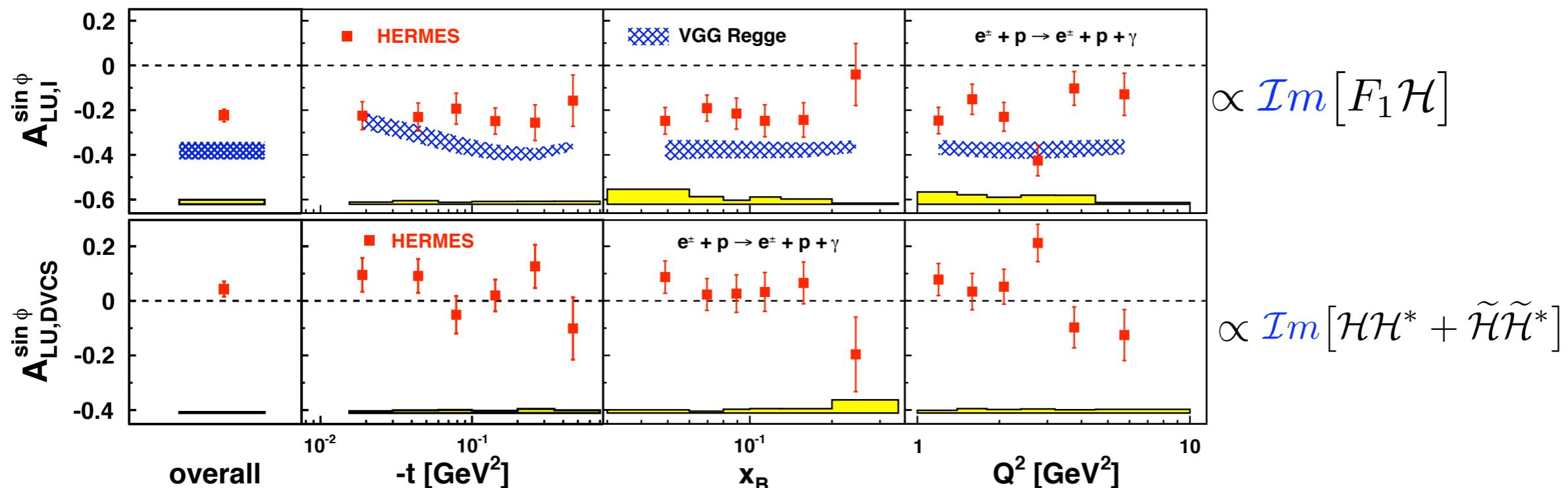
Model: VGG
 Phys..Rev.D (1999) 094017
 Prog. Nucl. Phys, 47 (2001) 401



- Beam charge asymmetry
- strong $-t$ dependence
 - no x_B and Q^2 dependencies
 - good agreement with model predictions

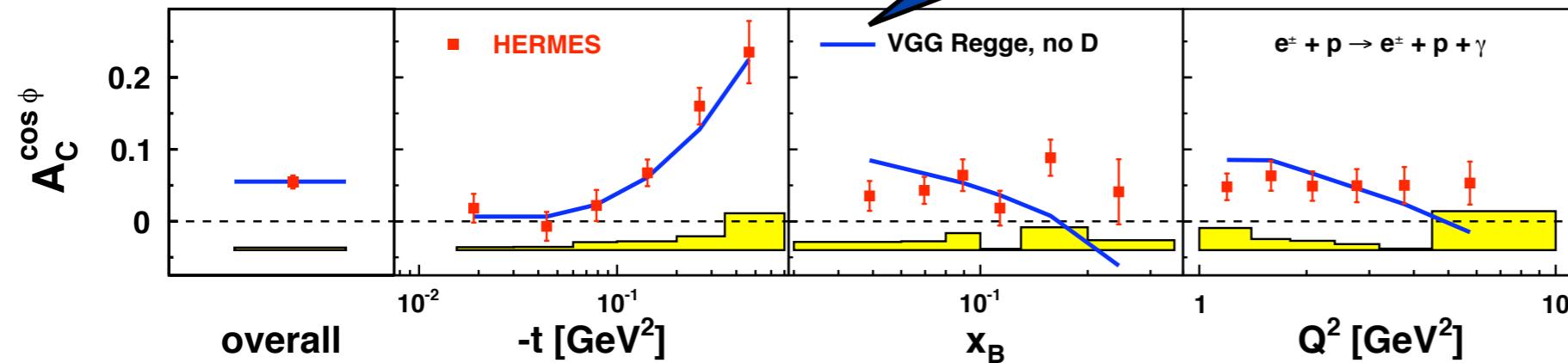
$$\propto \Re[F_1 \mathcal{H}]$$

$$\mathcal{A}_{LU}^{I,DVCS}(\phi) = \frac{(\sigma^{+\rightarrow} - \sigma^{+\leftarrow})^+ - (\sigma^{-\rightarrow} - \sigma^{-\leftarrow})^-}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}$$



$$\mathcal{A}_C(\phi) = \frac{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) - (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}$$

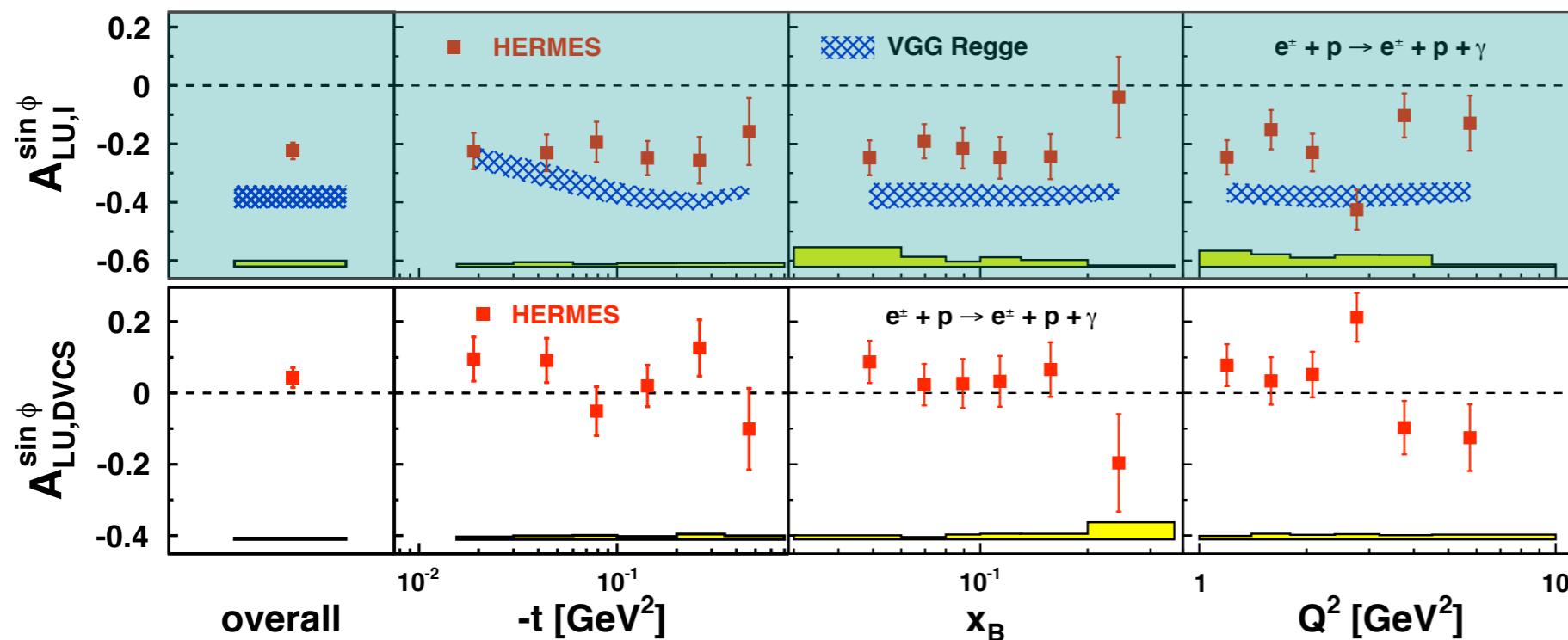
Model: VGG
 Phys..Rev.D (1999) 094017
 Prog. Nucl. Phys, 47 (2001) 401



- Beam charge asymmetry
- strong $-t$ dependence
 - no x_B and Q^2 dependencies
 - good agreement with model predictions

$$\propto \mathcal{R}e[F_1 \mathcal{H}]$$

$$\mathcal{A}_{LU}^{I,DVCS}(\phi) = \frac{(\sigma^{+\rightarrow} - \sigma^{+\leftarrow})^+ - (\sigma^{-\rightarrow} - \sigma^{-\leftarrow})^-}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}$$



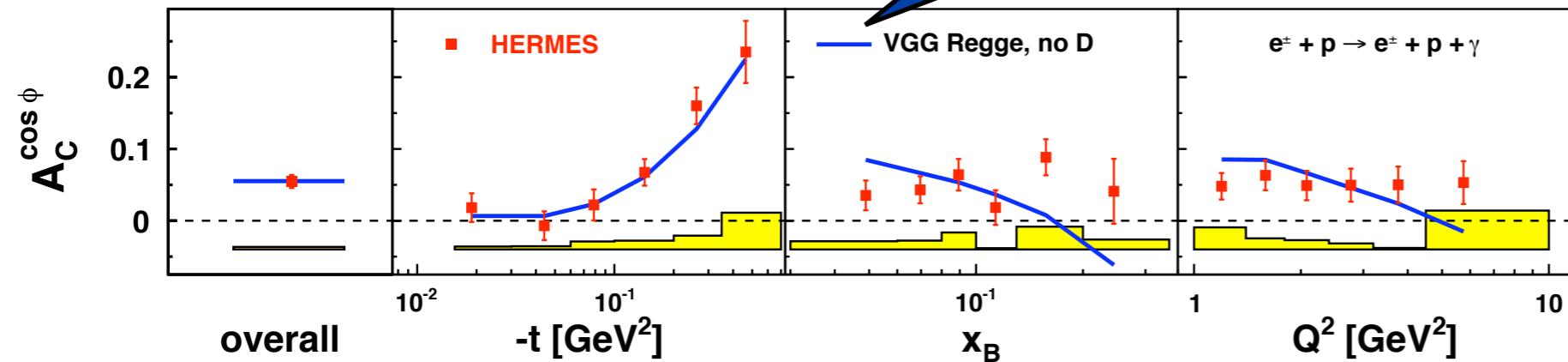
- Charge-difference beam-helicity asymmetry
- significant negative value
 - no kinematic dependencies
 - model predictions overshoot the data

$$\propto \mathcal{I}m[F_1 \mathcal{H}]$$

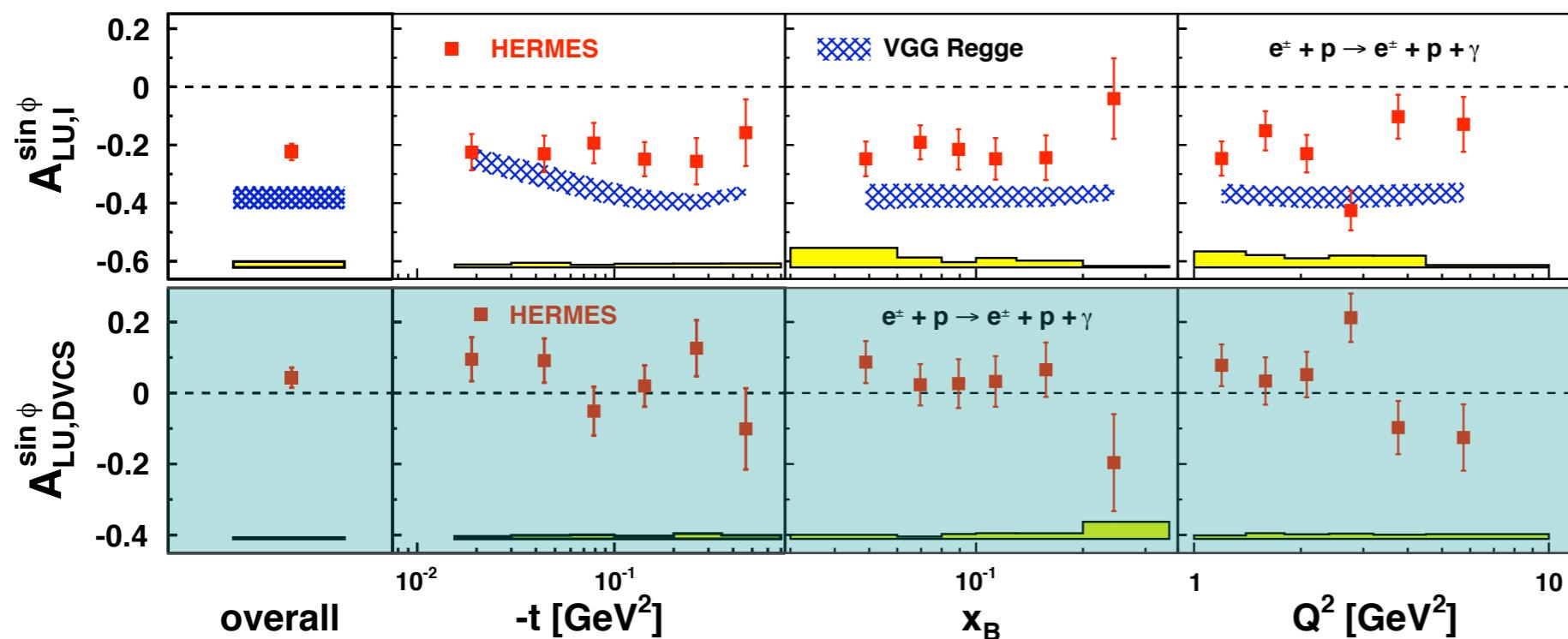
$$\propto \mathcal{I}m[\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*]$$

$$\mathcal{A}_C(\phi) = \frac{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) - (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}$$

Model: VGG
 Phys..Rev.D (1999) 094017
 Prog. Nucl. Phys, 47 (2001) 401



$$\mathcal{A}_{LU}^{I,DVCS}(\phi) = \frac{(\sigma^{+\rightarrow} - \sigma^{+\leftarrow})^+ - (\sigma^{-\rightarrow} - \sigma^{-\leftarrow})^-}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}$$



Beam charge asymmetry

- strong $-t$ dependence
- no x_B and Q^2 dependencies
- good agreement with model predictions

$$\propto \Re[F_1 \mathcal{H}]$$

Charge-difference beam-helicity asymmetry

- significant negative value
- no kinematic dependencies
- model predictions overshoot the data

$$\propto \Im[F_1 \mathcal{H}]$$

Charge-averaged beam-helicity asymmetry

- consistent with zero

$$\propto \Im[\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*]$$

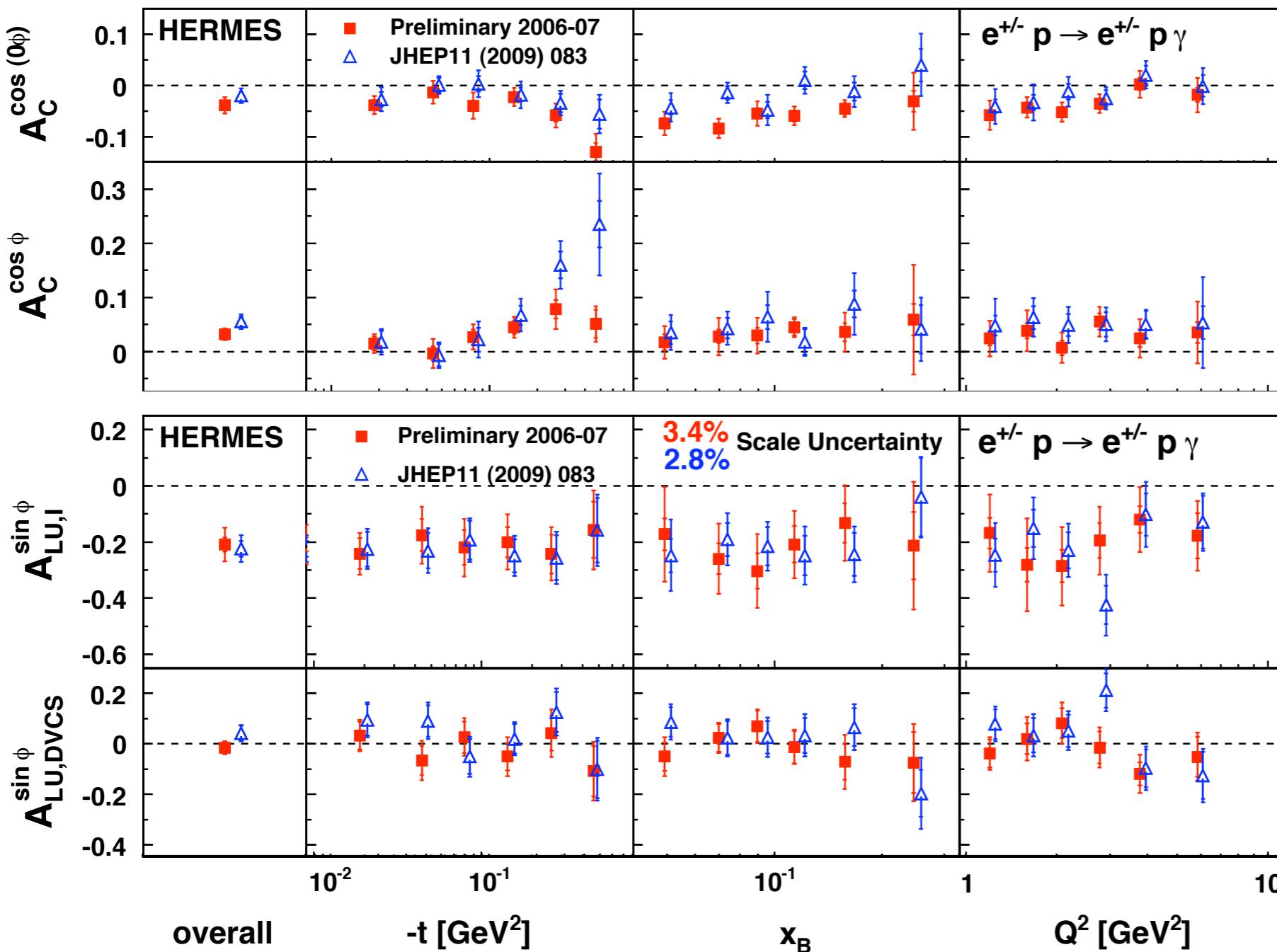
2 - dimensional (x_B , $-t$) binning also available

Beam-charge & Beam-helicity asymmetries (\mathbb{H})

1996-2005: JHEP 11 (2009) 083

2006-2007: Preliminary

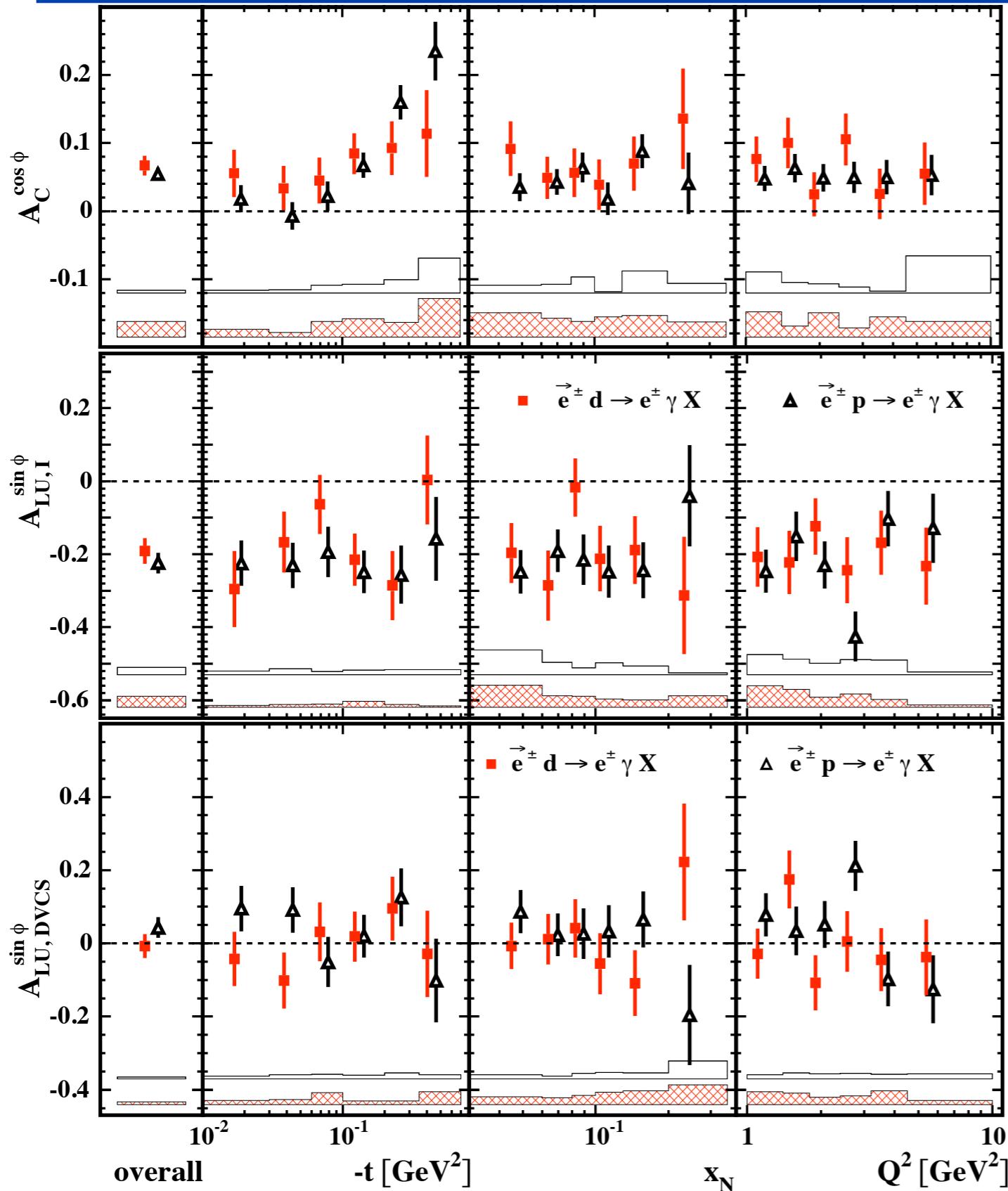
- improved precision
1996/2005 + 2006/2007 data



$$\mathcal{A}_C(\phi) = \frac{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) - (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}$$

$$\mathcal{A}_{LU}^I(\phi) = \frac{(\sigma^{+\rightarrow} - \sigma^{+\leftarrow}) - (\sigma^{-\rightarrow} - \sigma^{-\leftarrow})}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}$$

$$\mathcal{A}_{LU}^{DVCS}(\phi) = \frac{(\sigma^{+\rightarrow} - \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} - \sigma^{-\leftarrow})}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\rightarrow} + \sigma^{-\leftarrow})}$$



$$\propto \begin{aligned} & \mathcal{R}e[F_1 \mathcal{H}] \\ & \mathcal{R}e[G_1 \mathcal{H}_1] \end{aligned}$$

$$\propto \begin{aligned} & \mathcal{I}m[F_1 \mathcal{H}] \\ & \mathcal{I}m[G_1 \mathcal{H}_1] \end{aligned}$$

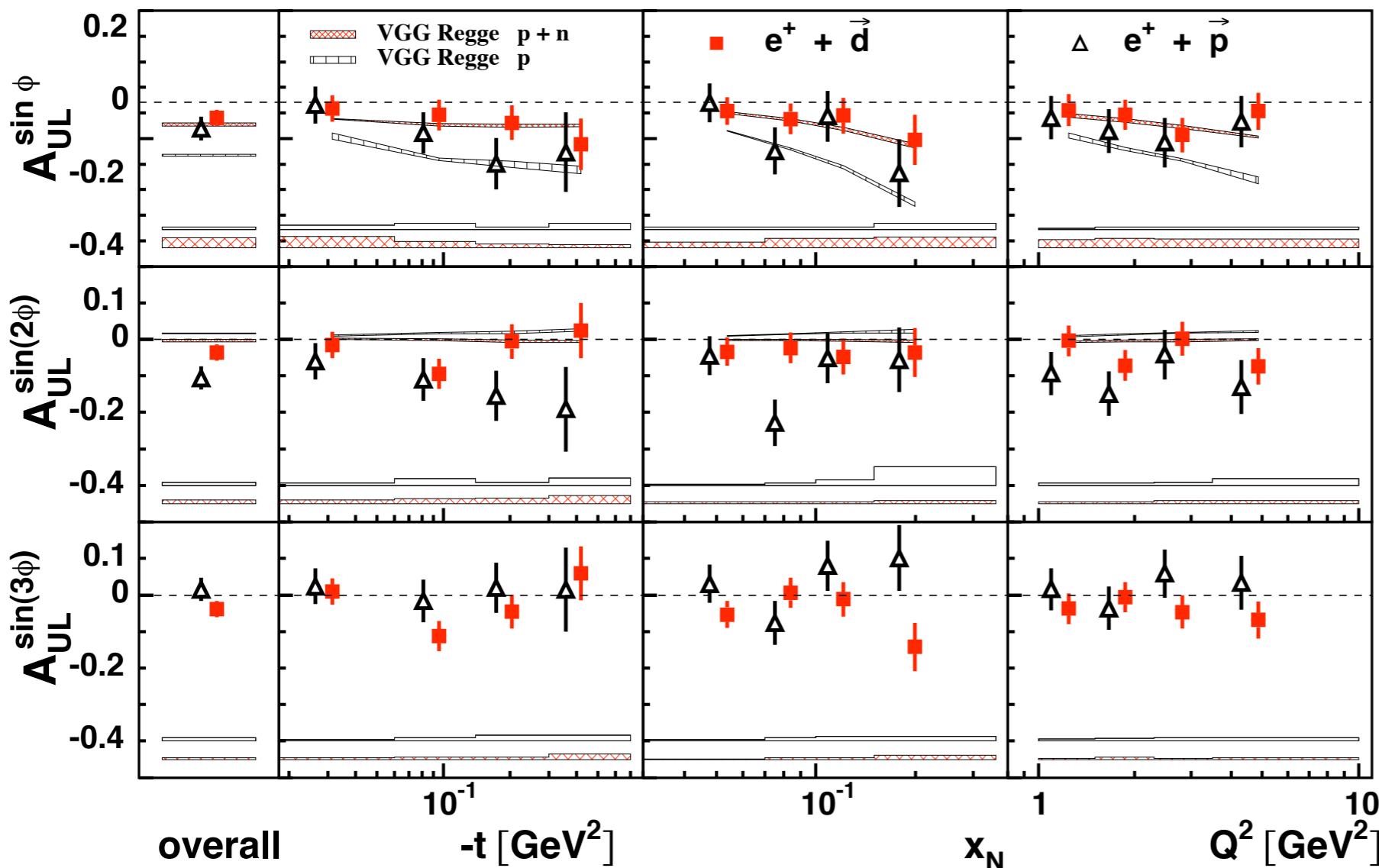
$$\propto \begin{aligned} & \mathcal{I}m[\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*] \\ & \mathcal{I}m[\mathcal{H}_1\mathcal{H}_1^* + \dots] \end{aligned}$$

Beam-charge & Beam-Helicity asymmetries

- Proton and Deuteron results are compatible at low and intermediate $-t$ regions.
- No clear signatures of 40% contribution from coherent process at low $-t$.
- no significant x_B and Q^2 dependencies.
- Difference at large $-t$ (for beam-charge asymmetry) might be caused by additional contributions from Neutron and its resonances.

Data collected with positron beam

$$\mathcal{A}_{UL}(\phi) = \frac{(\sigma^{\rightarrow\Rightarrow} + \sigma^{\leftarrow\Rightarrow}) - (\sigma^{\rightarrow\Leftarrow} + \sigma^{\leftarrow\Leftarrow})}{(\sigma^{\rightarrow\Rightarrow} + \sigma^{\leftarrow\Rightarrow}) + (\sigma^{\rightarrow\Leftarrow} + \sigma^{\leftarrow\Leftarrow})}$$



$$\propto \frac{\mathcal{Im}[F_1 \tilde{\mathcal{H}}]}{\mathcal{Im}[G_1 \tilde{\mathcal{H}}_1]}$$

\propto Higher twist

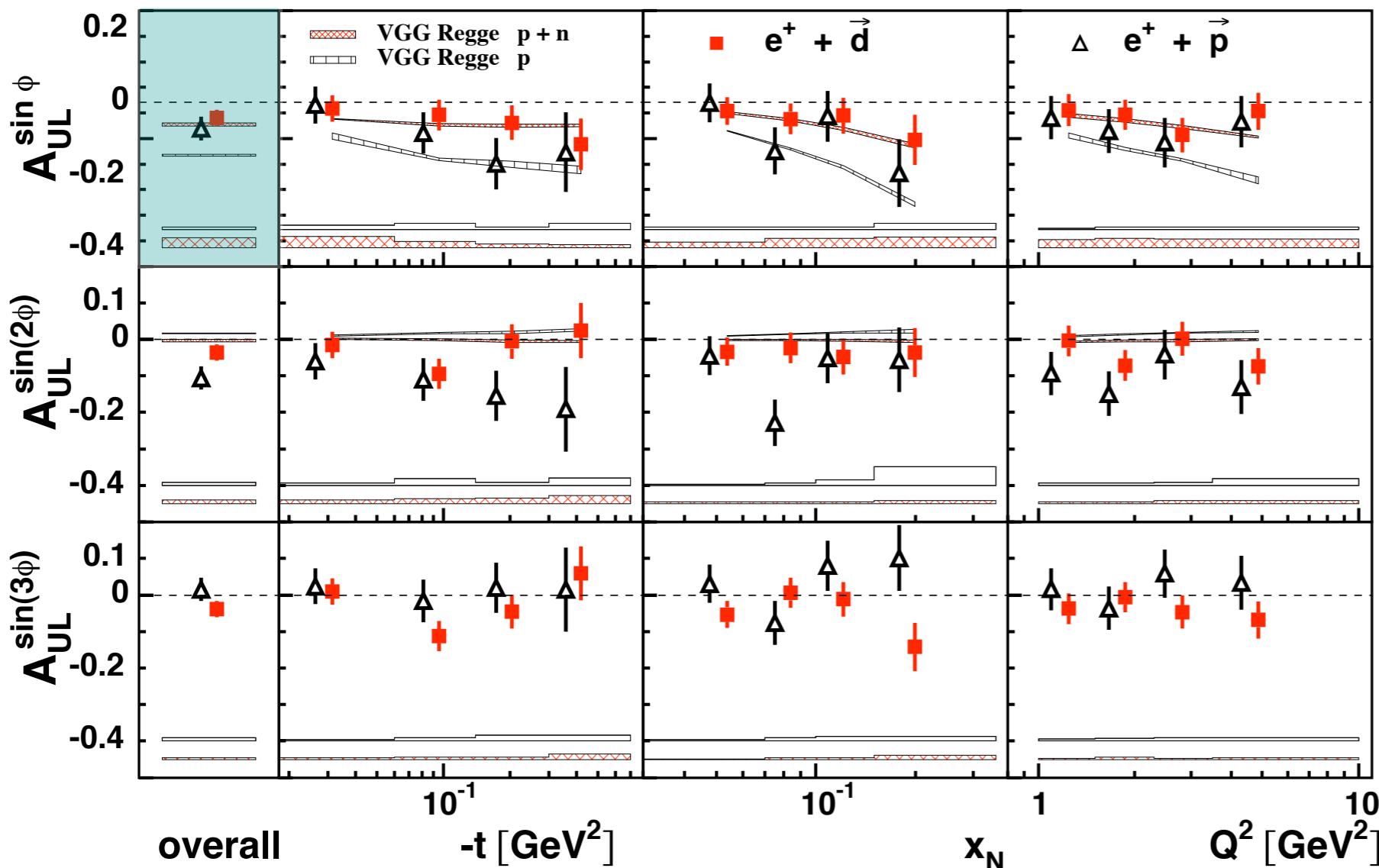
Model predictions for deuteron: Sum of incoherent processes on the proton and neutron

Data collected with positron beam

$$\mathcal{A}_{UL}(\phi) = \frac{(\sigma^{\rightarrow\Rightarrow} + \sigma^{\leftarrow\Rightarrow}) - (\sigma^{\rightarrow\Leftarrow} + \sigma^{\leftarrow\Leftarrow})}{(\sigma^{\rightarrow\Rightarrow} + \sigma^{\leftarrow\Rightarrow}) + (\sigma^{\rightarrow\Leftarrow} + \sigma^{\leftarrow\Leftarrow})}$$

Longitudinal Target-Spin asymmetry

- Non-zero negative value of leading $\sin(\phi)$ amplitude on both targets.



$$\propto \frac{\mathcal{Im}[F_1 \tilde{\mathcal{H}}]}{\mathcal{Im}[G_1 \tilde{\mathcal{H}}_1]}$$

\propto Higher twist

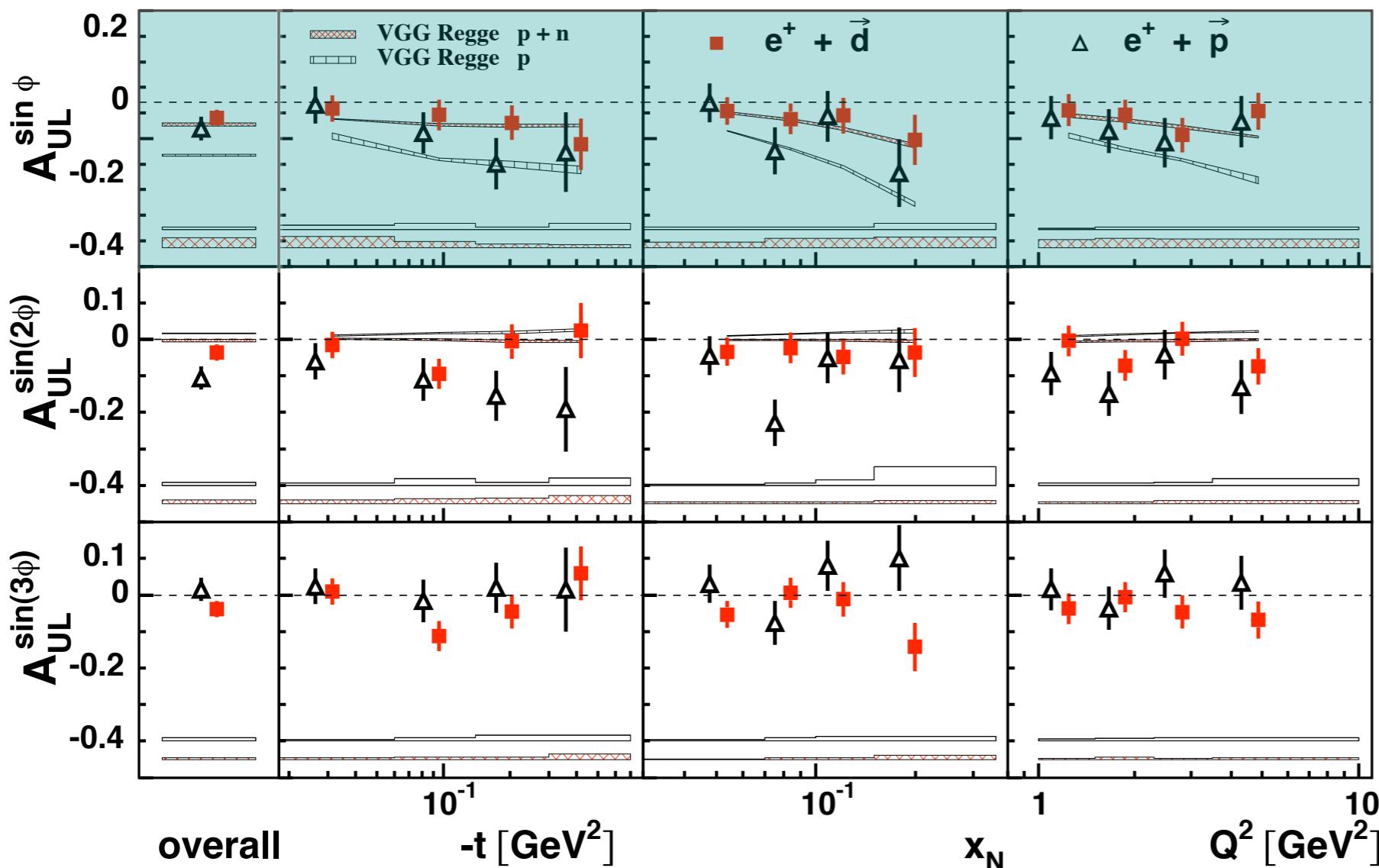
Model predictions for deuteron: Sum of incoherent processes on the proton and neutron

Data collected with positron beam

$$\mathcal{A}_{UL}(\phi) = \frac{(\sigma^{\rightarrow\Rightarrow} + \sigma^{\leftarrow\Rightarrow}) - (\sigma^{\rightarrow\Leftarrow} + \sigma^{\leftarrow\Leftarrow})}{(\sigma^{\rightarrow\Rightarrow} + \sigma^{\leftarrow\Rightarrow}) + (\sigma^{\rightarrow\Leftarrow} + \sigma^{\leftarrow\Leftarrow})}$$

Longitudinal Target-Spin asymmetry

- Non-zero negative value of leading $\sin(\phi)$ amplitude on both targets.



$$\propto \frac{\mathcal{Im}[F_1 \tilde{\mathcal{H}}]}{\mathcal{Im}[G_1 \tilde{\mathcal{H}}_1]}$$

- Results on deuteron neither support nor disfavor large contribution from neutron, predicted by the model.
- Results on proton and deuteron targets are compatible.

\propto Higher twist

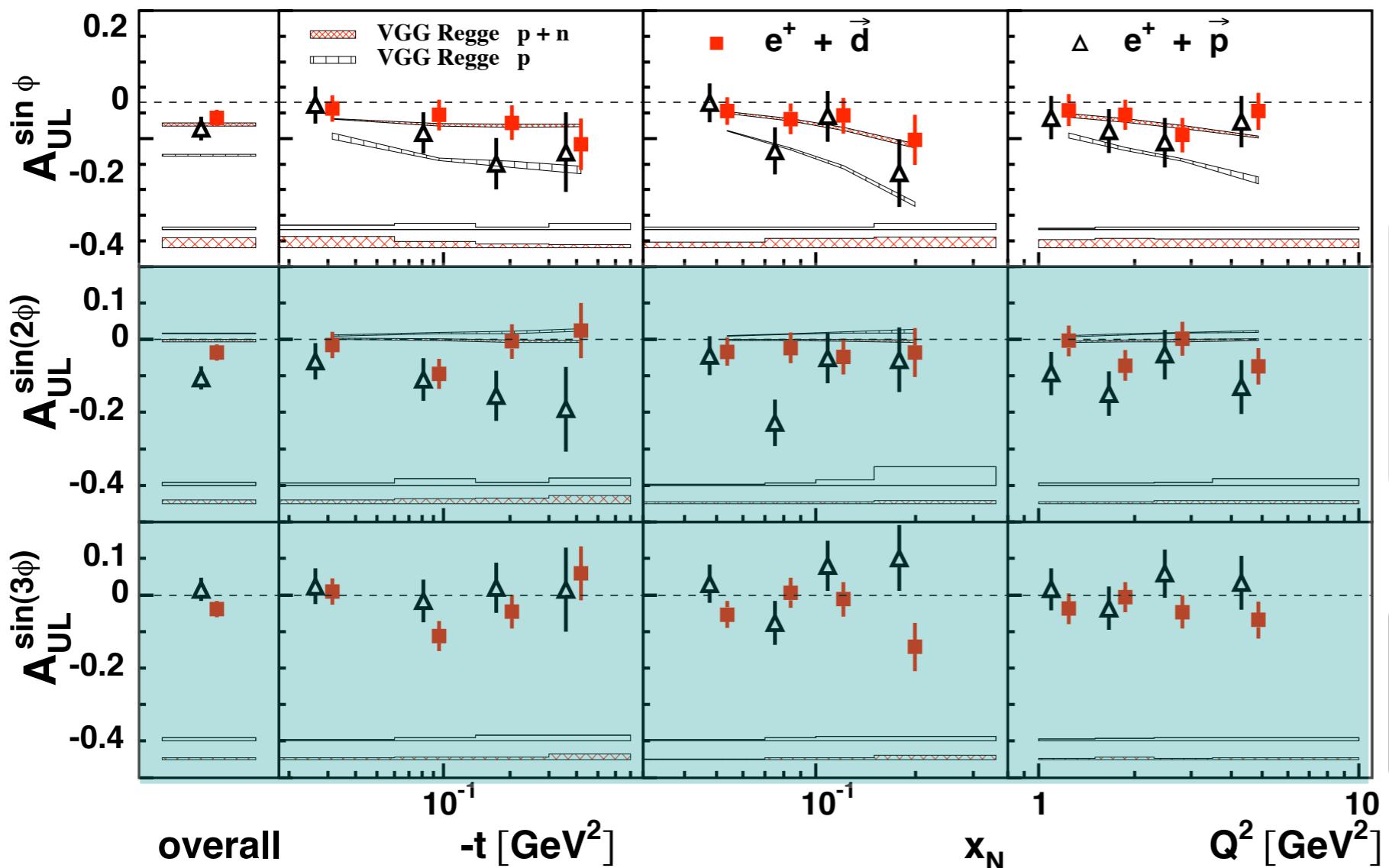
Model predictions for deuteron: Sum of incoherent processes on the proton and neutron

Data collected with positron beam

$$\mathcal{A}_{UL}(\phi) = \frac{(\sigma^{\rightarrow\Rightarrow} + \sigma^{\leftarrow\Rightarrow}) - (\sigma^{\rightarrow\Leftarrow} + \sigma^{\leftarrow\Leftarrow})}{(\sigma^{\rightarrow\Rightarrow} + \sigma^{\leftarrow\Rightarrow}) + (\sigma^{\rightarrow\Leftarrow} + \sigma^{\leftarrow\Leftarrow})}$$

Longitudinal Target-Spin asymmetry

- Non-zero negative value of leading $\sin(\phi)$ amplitude on both targets.



$$\propto \frac{\mathcal{I}m[F_1 \tilde{\mathcal{H}}]}{\mathcal{I}m[G_1 \tilde{\mathcal{H}}_1]}$$

- Results on deuteron neither support nor disfavor large contribution from neutron, predicted by the model.
- Results on proton and deuteron targets are compatible.

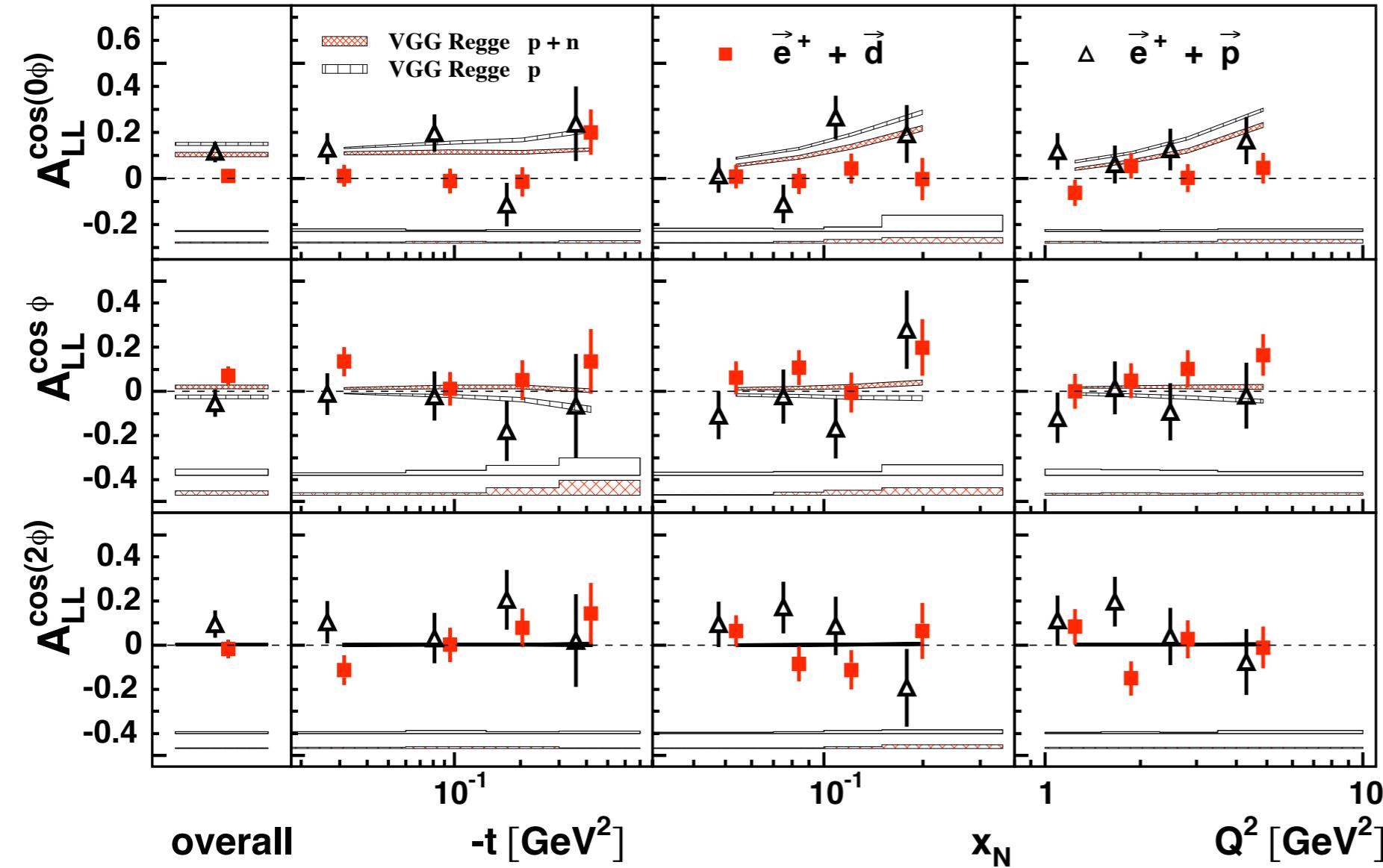
\propto Higher twist

- Amplitudes related to the Higher twist contributions are consistent with zero.

Model predictions for deuteron: Sum of incoherent processes on the proton and neutron

Data collected with positron beam

$$\mathcal{A}_{LL}(\phi) = \frac{(\sigma^{\rightarrow\rightarrow} + \sigma^{\leftarrow\leftarrow}) - (\sigma^{\rightarrow\leftarrow} + \sigma^{\leftarrow\rightarrow})}{(\sigma^{\rightarrow\rightarrow} + \sigma^{\leftarrow\leftarrow}) + (\sigma^{\rightarrow\leftarrow} + \sigma^{\leftarrow\rightarrow})}$$



$$\propto -A_{LL}^{\cos(\phi)}$$

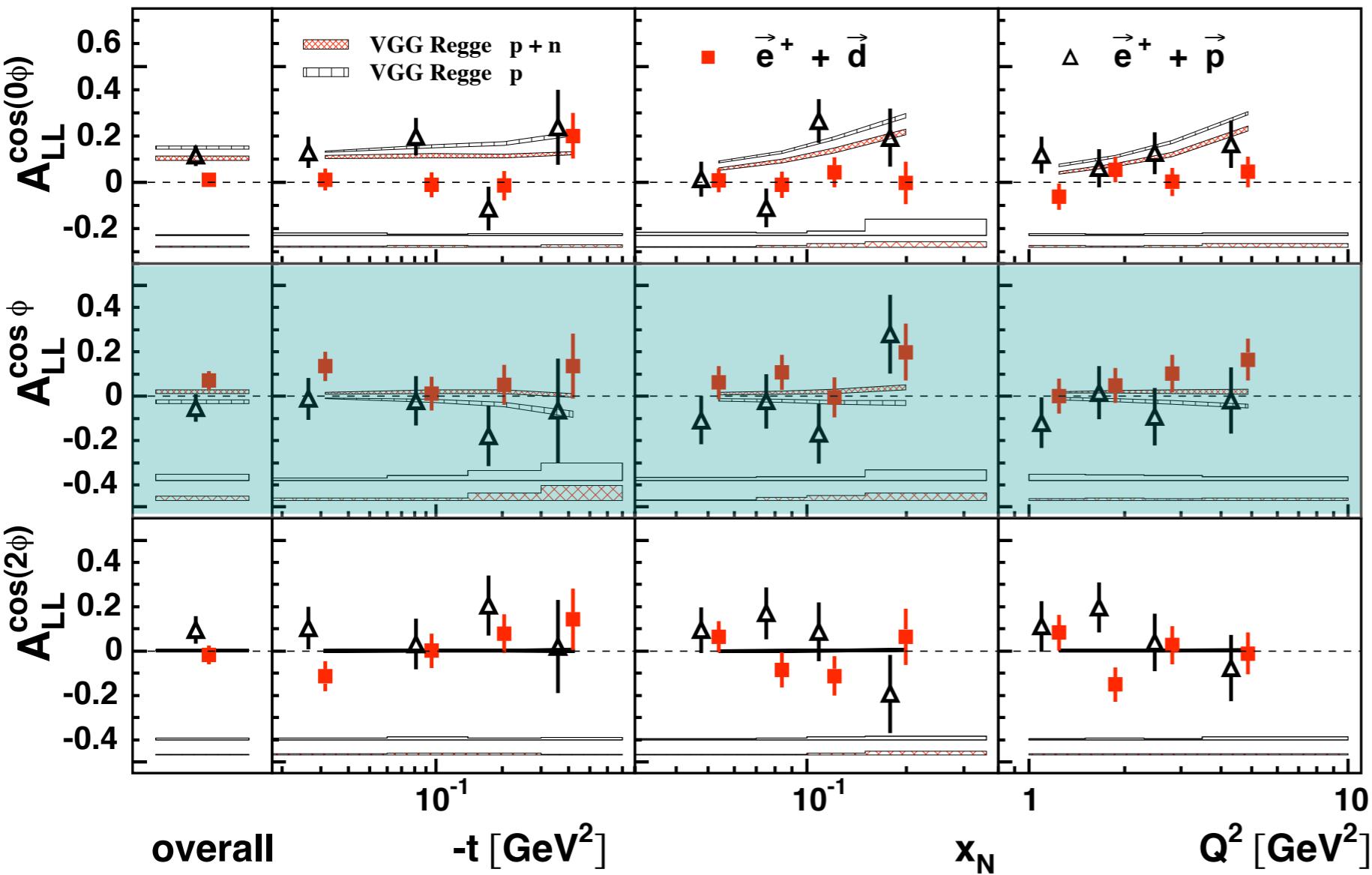
$$\begin{aligned} &\propto \text{Re}[F_1 \tilde{\mathcal{H}}] \\ &\propto \text{Re}[G_1 \tilde{\mathcal{H}}_1] \end{aligned}$$

$$\propto \text{Higher twist}$$

Asymmetry amplitudes are attributed not only to squared DVCS or Interference term, but also to squared Bethe-Heitler term.

Data collected with positron beam

$$\mathcal{A}_{LL}(\phi) = \frac{(\sigma^{\rightarrow\rightarrow} + \sigma^{\leftarrow\leftarrow}) - (\sigma^{\rightarrow\leftarrow} + \sigma^{\leftarrow\rightarrow})}{(\sigma^{\rightarrow\rightarrow} + \sigma^{\leftarrow\leftarrow}) + (\sigma^{\rightarrow\leftarrow} + \sigma^{\leftarrow\rightarrow})}$$



Longitudinal Double-Spin asymmetry

- Leading $\cos(\phi)$ amplitude is compatible with zero for both targets.

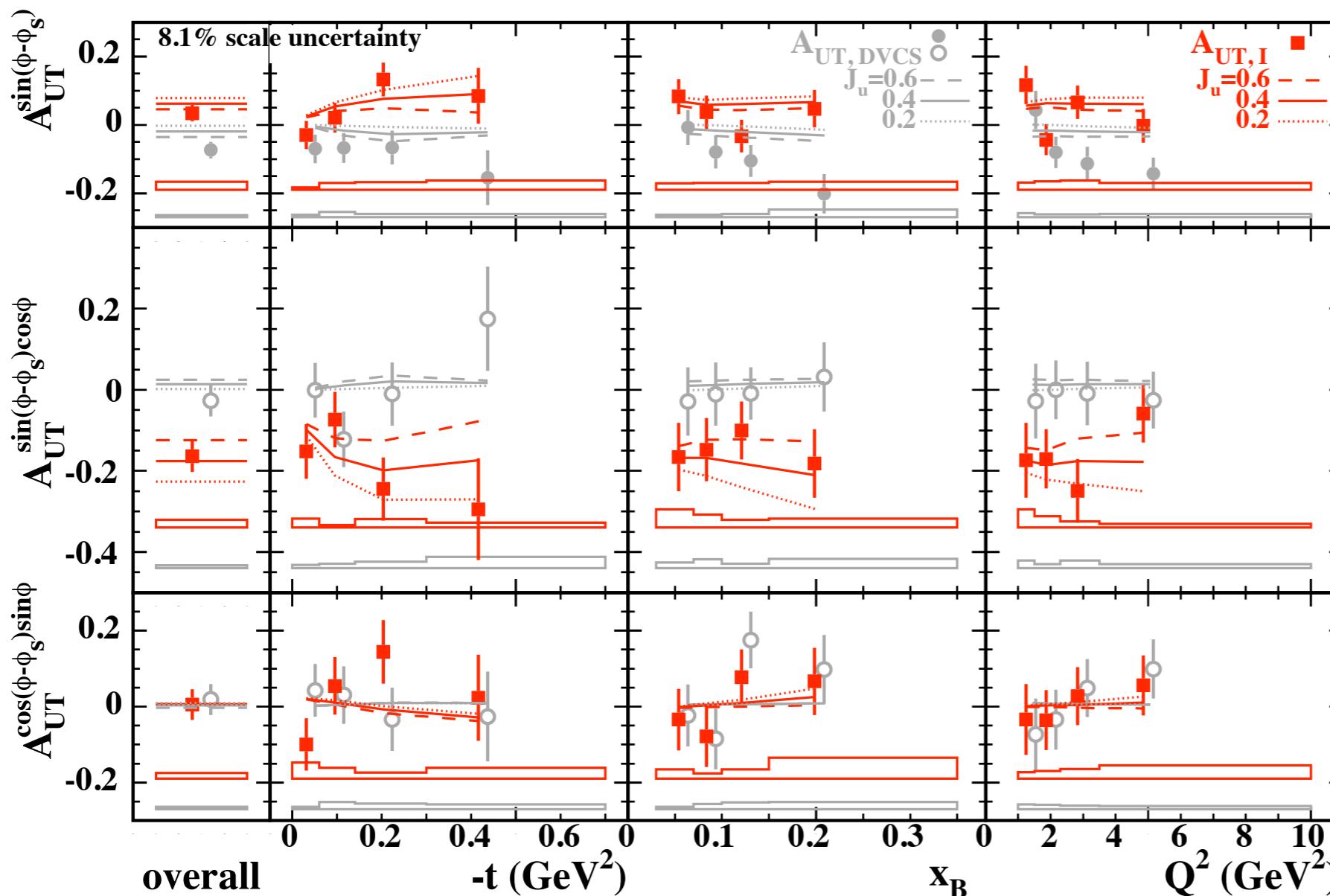
$$\propto -A_{LL}^{\cos(\phi)}$$

$$\begin{aligned} &\propto \text{Re}[F_1 \tilde{\mathcal{H}}] \\ &\propto \text{Re}[G_1 \tilde{\mathcal{H}}_1] \end{aligned}$$

\propto Higher twist

Asymmetry amplitudes are attributed not only to squared DVCS or Interference term, but also to squared Bethe-Heitler term.

$$\mathcal{A}_{UT}^{I,DVCS}(\phi, \phi_S) = \frac{(\sigma^{+\uparrow} - \sigma^{+\downarrow})^+ (\sigma^{-\uparrow} - \sigma^{-\downarrow})}{(\sigma^{+\uparrow} + \sigma^{+\downarrow}) + (\sigma^{-\uparrow} + \sigma^{-\downarrow})}$$

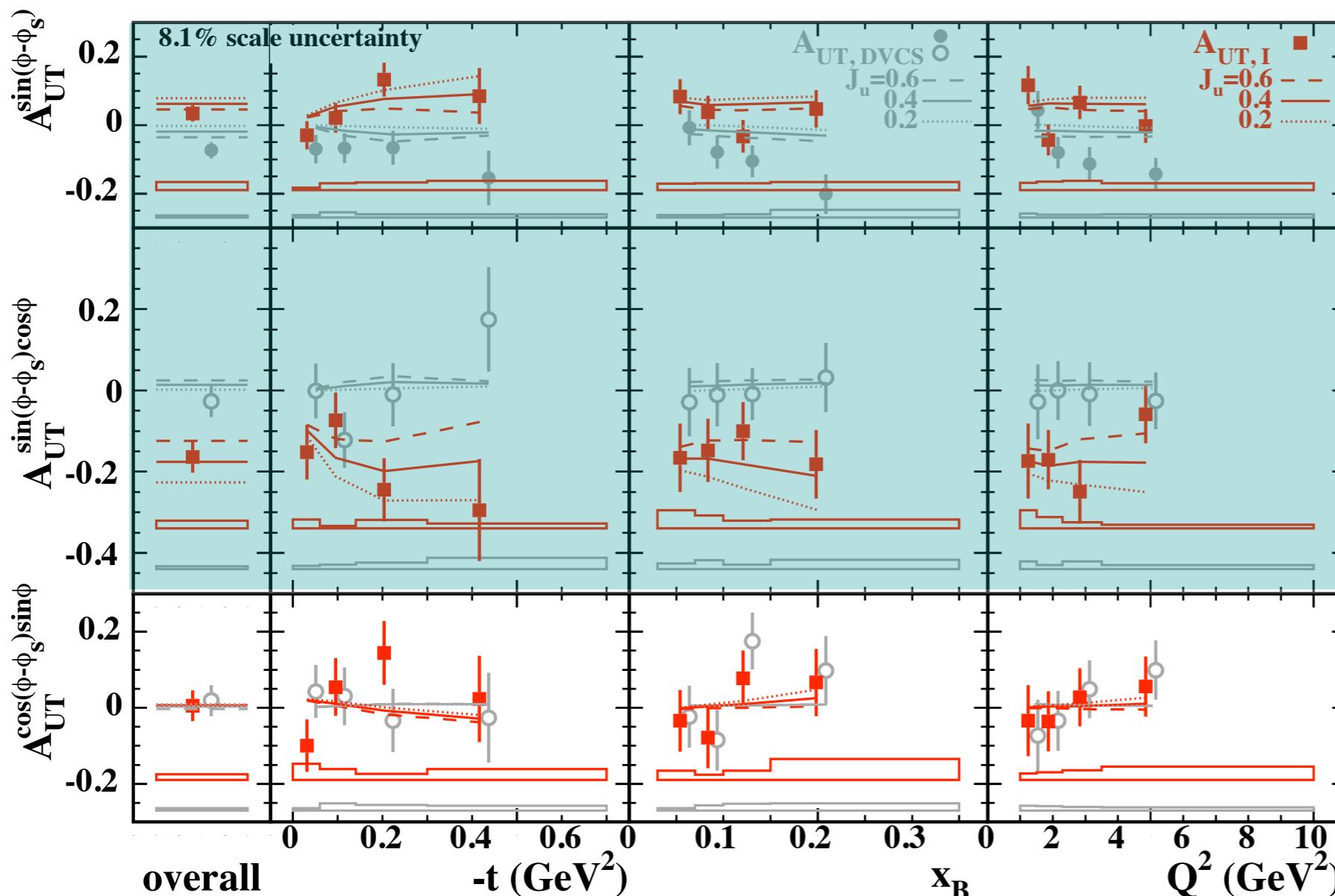


$$\propto A_{UT}^{\sin(\phi-\phi_S)\cos(\phi)}$$

$$\propto \begin{aligned} & \mathcal{Im}[F_2 \mathcal{H} - F_1 \mathcal{E}] \\ & \mathcal{Im}[\mathcal{H}\mathcal{E}^* - \mathcal{E}\mathcal{H}^* - \xi(\tilde{\mathcal{H}}\tilde{\mathcal{E}}^* - \tilde{\mathcal{E}}\tilde{\mathcal{H}}^*)] \end{aligned}$$

$$\propto \begin{aligned} & \mathcal{Im}[F_2 \tilde{\mathcal{H}} - (F_1 + \xi F_2) \tilde{\mathcal{E}}] \\ & \mathcal{Im}[-\tilde{\mathcal{H}}\mathcal{E}^* - \tilde{\mathcal{H}}^*\mathcal{E} + \xi(\mathcal{H}\tilde{\mathcal{E}}^* + \tilde{\mathcal{E}}\tilde{\mathcal{H}}^*)] \end{aligned}$$

$$\mathcal{A}_{UT}^{I,DVCS}(\phi, \phi_S) = \frac{(\sigma^{+\uparrow} - \sigma^{+\downarrow})^+ (\sigma^{-\uparrow} - \sigma^{-\downarrow})^-}{(\sigma^{+\uparrow} + \sigma^{+\downarrow})^+ + (\sigma^{-\uparrow} + \sigma^{-\downarrow})^-}$$



Charge-difference Transverse Target-Spin asymmetry

- Non-zero leading $\cos(n\phi)$ amplitudes.

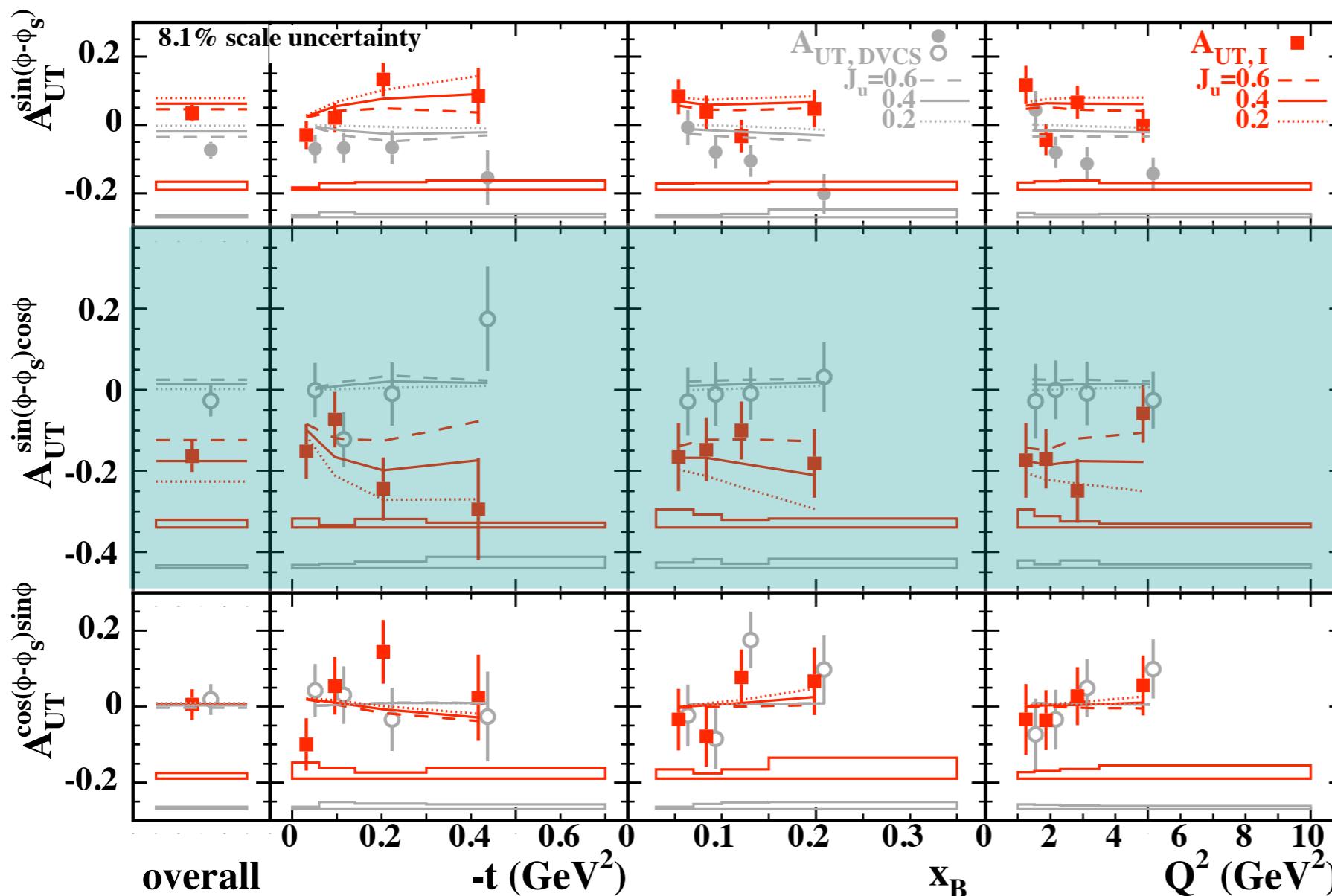
$$\propto A_{UT}^{\sin(\phi-\phi_S) \cos(\phi)}$$

$$\propto \begin{aligned} & \text{Im}[F_2 \mathcal{H} - F_1 \mathcal{E}] \\ & \text{Im}[\mathcal{H} \mathcal{E}^* - \mathcal{E} \mathcal{H}^* - \xi(\tilde{\mathcal{H}} \tilde{\mathcal{E}}^* - \tilde{\mathcal{E}} \tilde{\mathcal{H}}^*)] \end{aligned}$$

$$\propto \begin{aligned} & \text{Im}[F_2 \tilde{\mathcal{H}} - (F_1 + \xi F_2) \tilde{\mathcal{E}}] \\ & \text{Im}[-\tilde{\mathcal{H}} \mathcal{E}^* - \tilde{\mathcal{H}}^* \mathcal{E} + \xi(\mathcal{H} \tilde{\mathcal{E}}^* + \tilde{\mathcal{E}} \mathcal{H}^*)] \end{aligned}$$

Transverse Target-Spin Asymmetries (\mathcal{H})

$$\mathcal{A}_{UT}^{I,DVCS}(\phi, \phi_S) = \frac{(\sigma^{+\uparrow} - \sigma^{+\downarrow})^+ (\sigma^{-\uparrow} - \sigma^{-\downarrow})^-}{(\sigma^{+\uparrow} + \sigma^{+\downarrow})^+ + (\sigma^{-\uparrow} + \sigma^{-\downarrow})^-}$$



Charge-difference Transverse Target-Spin asymmetry

- Non-zero leading $\cos(n\phi)$ amplitudes.

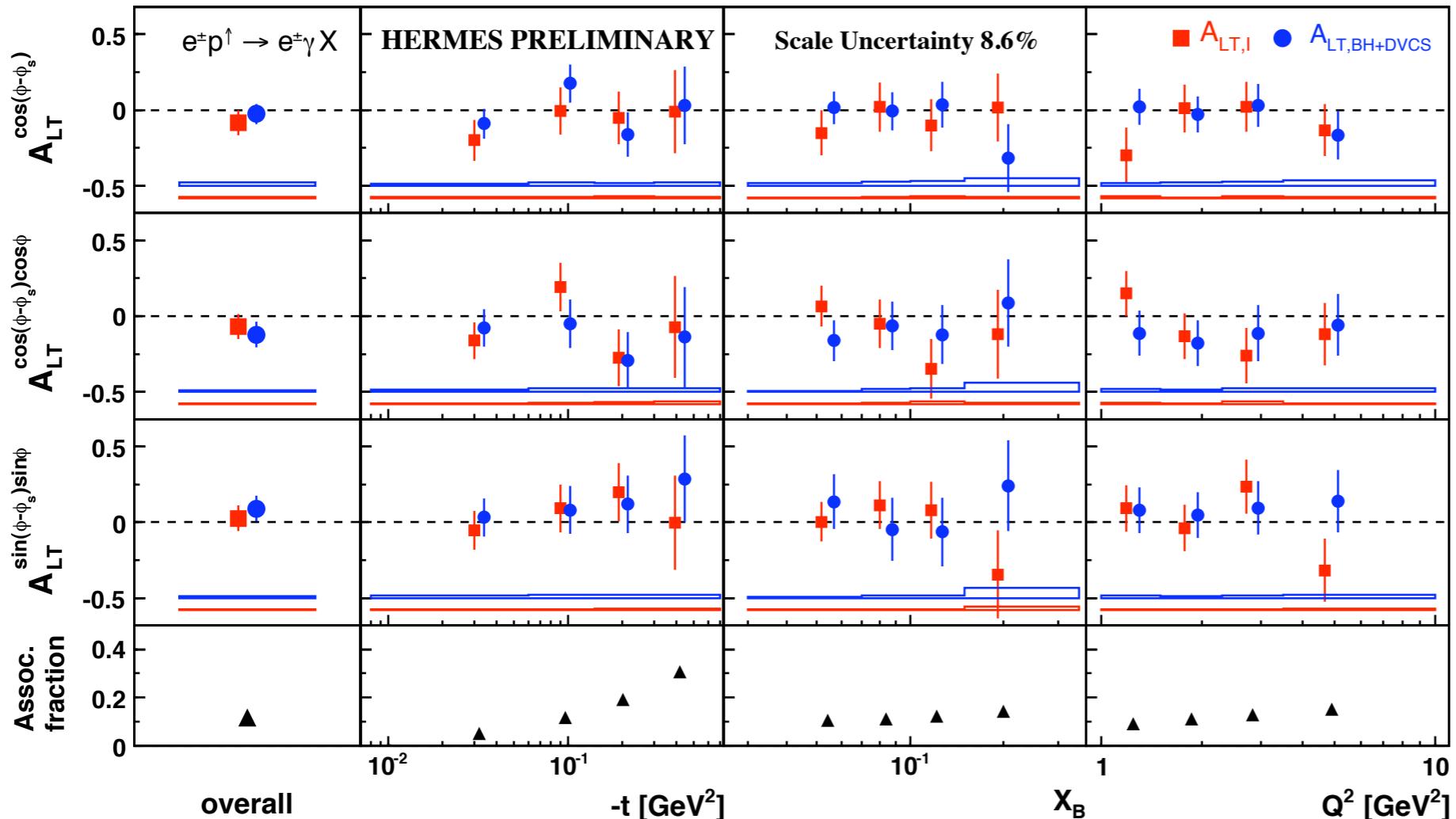
$$\propto A_{UT}^{\sin(\phi-\phi_S) \cos(\phi)}$$

$$\propto \begin{aligned} & \text{Im}[F_2 \mathcal{H} - F_1 \mathcal{E}] \\ & \text{Im}[\mathcal{H} \mathcal{E}^* - \mathcal{E} \mathcal{H}^* - \xi(\tilde{\mathcal{H}} \tilde{\mathcal{E}}^* - \tilde{\mathcal{E}} \tilde{\mathcal{H}}^*)] \end{aligned}$$

$$\propto \begin{aligned} & \text{Im}[F_2 \tilde{\mathcal{H}} - (F_1 + \xi F_2) \tilde{\mathcal{E}}] \\ & \text{Im}[-\tilde{\mathcal{H}} \mathcal{E}^* - \tilde{\mathcal{H}}^* \mathcal{E} + \xi(\mathcal{H} \tilde{\mathcal{E}}^* + \tilde{\mathcal{E}} \mathcal{H}^*)] \end{aligned}$$

Leading $\cos(\phi)$ amplitude of charge-difference target-spin asymmetry A_{UT}^I
is sensitive to CFF \mathcal{E} , therefore J_u .

$$\mathcal{A}_{LT}^{I, BH+DVCS}(\phi, \phi_S) = \frac{(\vec{\sigma}^{+\uparrow} + \vec{\sigma}^{+\downarrow} - \vec{\sigma}^{+\downarrow} - \vec{\sigma}^{+\uparrow})^+ (\vec{\sigma}^{-\uparrow} + \vec{\sigma}^{-\downarrow} - \vec{\sigma}^{-\downarrow} - \vec{\sigma}^{-\uparrow})^-}{(\vec{\sigma}^{+\uparrow} + \vec{\sigma}^{+\downarrow} + \vec{\sigma}^{+\downarrow} + \vec{\sigma}^{+\uparrow}) + (\vec{\sigma}^{+\uparrow} + \vec{\sigma}^{+\downarrow} + \vec{\sigma}^{+\downarrow} + \vec{\sigma}^{+\uparrow})}$$



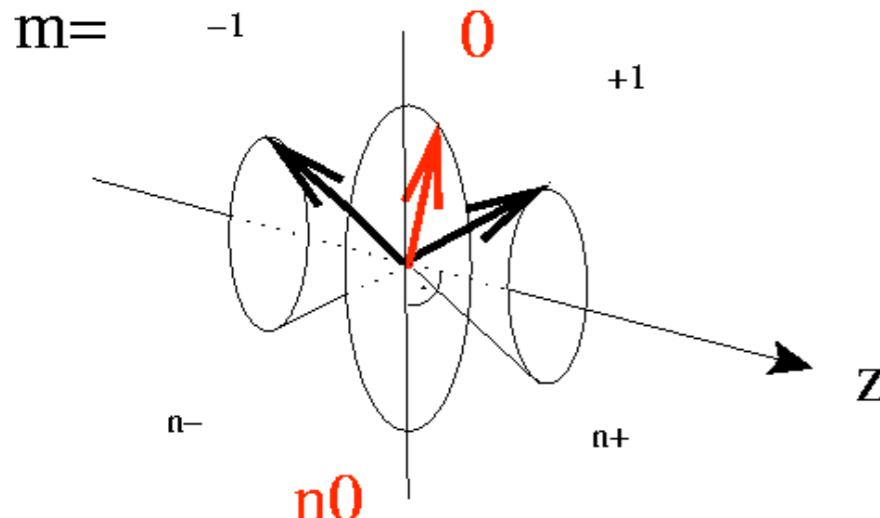
$$\propto A_{LT}^{\cos(\phi-\phi_S)\cos(\phi)}$$

$$\propto \frac{\text{Re}[F_2 \tilde{\mathcal{H}} - (F_1 + \xi F_2) \tilde{\mathcal{E}}]}{\text{Re}[\mathcal{H}\mathcal{E}^* - \mathcal{E}\mathcal{H}^* - \xi(\tilde{\mathcal{H}}\tilde{\mathcal{E}}^* - \tilde{\mathcal{E}}\tilde{\mathcal{H}}^*)]}$$

$$\propto \frac{\text{Re}[F_2 \mathcal{H} - F_1 \mathcal{E}]}{\text{Re}[-\tilde{\mathcal{H}}\mathcal{E}^* - \tilde{\mathcal{H}}^*\mathcal{E} + \xi(\mathcal{H}\tilde{\mathcal{E}}^* + \tilde{\mathcal{E}}\tilde{\mathcal{H}}^*)]}$$

Leading amplitudes of charge-difference and charge-averaged transverse double-spin asymmetries are compatible with zero over all kinematic regions.
Sensitivity to J_u is suppressed by kinematic pre-factor.

Beam-Spin and Beam-Charge Asymmetries



Beam-Spin and Beam-Charge Asymmetries
on longitudinally polarized Deuterium
with vanishing vector polarization
(non-vanishing tensor polarization)

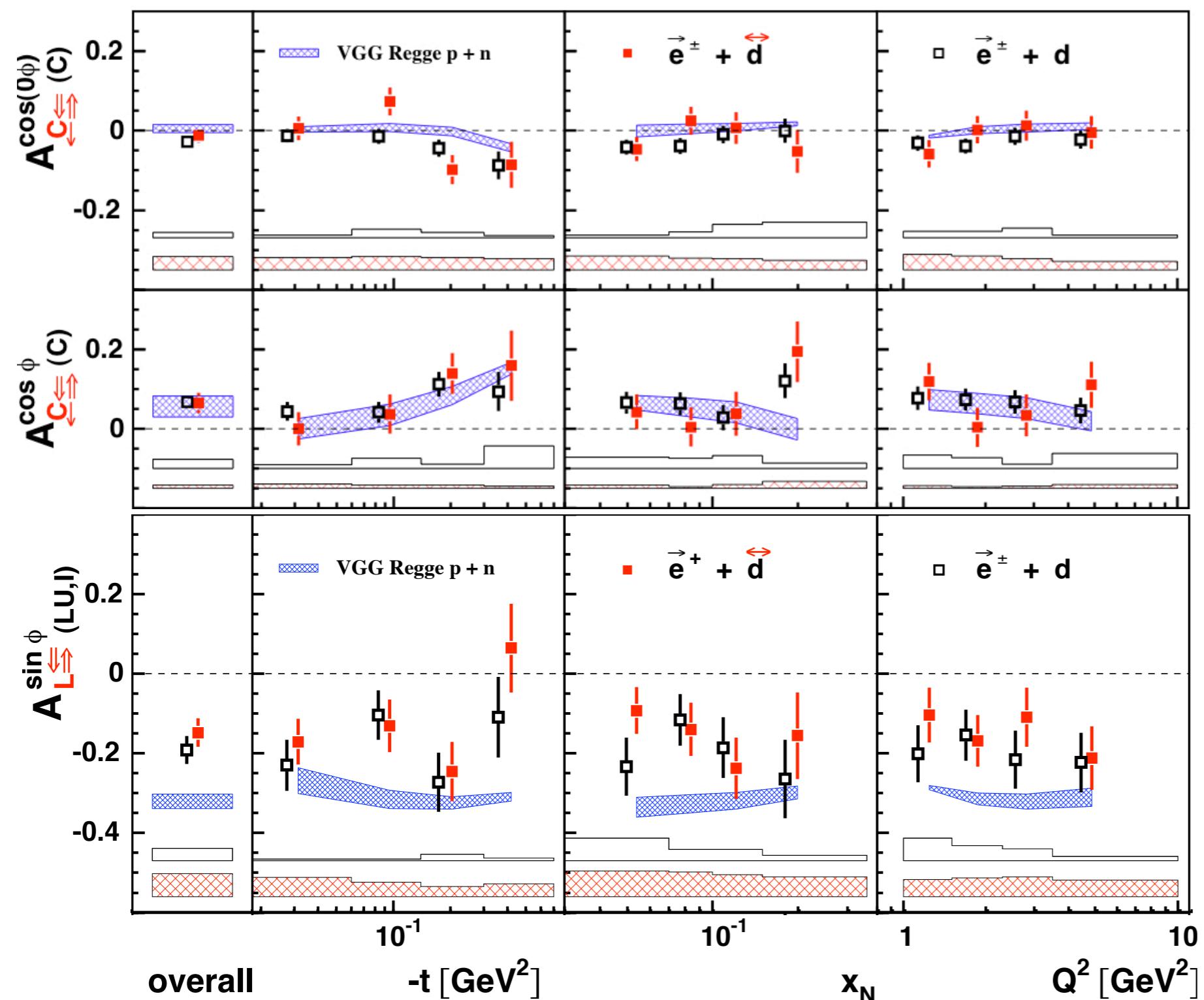
Coherent contribution :
small $-t \simeq 40\%$

$$\mathcal{A}_{L\rightleftarrows}(\phi) = \frac{(\sigma^{\rightarrow\rightarrow} + \sigma^{\rightarrow\leftarrow}) - (\sigma^{\leftarrow\leftarrow} + (\sigma^{\leftarrow\rightarrow}))}{(\sigma^{\rightarrow\rightarrow} + \sigma^{\rightarrow\leftarrow}) + (\sigma^{\leftarrow\leftarrow} + (\sigma^{\leftarrow\rightarrow}))} \quad \left. \begin{array}{l} \\ \end{array} \right\} \mathcal{Im}[G_1(\mathcal{H}_1 - \frac{1}{3}\mathcal{H}_5)]$$

$$\mathcal{A}_{LU}(\phi) = \frac{\sigma^{\rightarrow} - \sigma^{\leftarrow}}{\sigma^{\rightarrow} + \sigma^{\leftarrow}} \quad \left. \begin{array}{l} \\ \end{array} \right\} \mathcal{Im}[G_1\mathcal{H}_1]$$

$$\mathcal{A}_{C\rightleftarrows}(\phi) = \frac{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) - (\sigma^{-\leftarrow} + (\sigma^{-\rightarrow}))}{(\sigma^{+\rightarrow} + \sigma^{+\leftarrow}) + (\sigma^{-\leftarrow} + (\sigma^{-\rightarrow}))} \quad \left. \begin{array}{l} \\ \end{array} \right\} \mathcal{Re}[G_1(\mathcal{H}_1 - \frac{1}{3}\mathcal{H}_5)]$$

$$\mathcal{A}_C(\phi) = \frac{\sigma^+ - \sigma^-}{\sigma^+ - \sigma^-} \quad \left. \begin{array}{l} \\ \end{array} \right\} \mathcal{Re}[G_1\mathcal{H}_1]$$

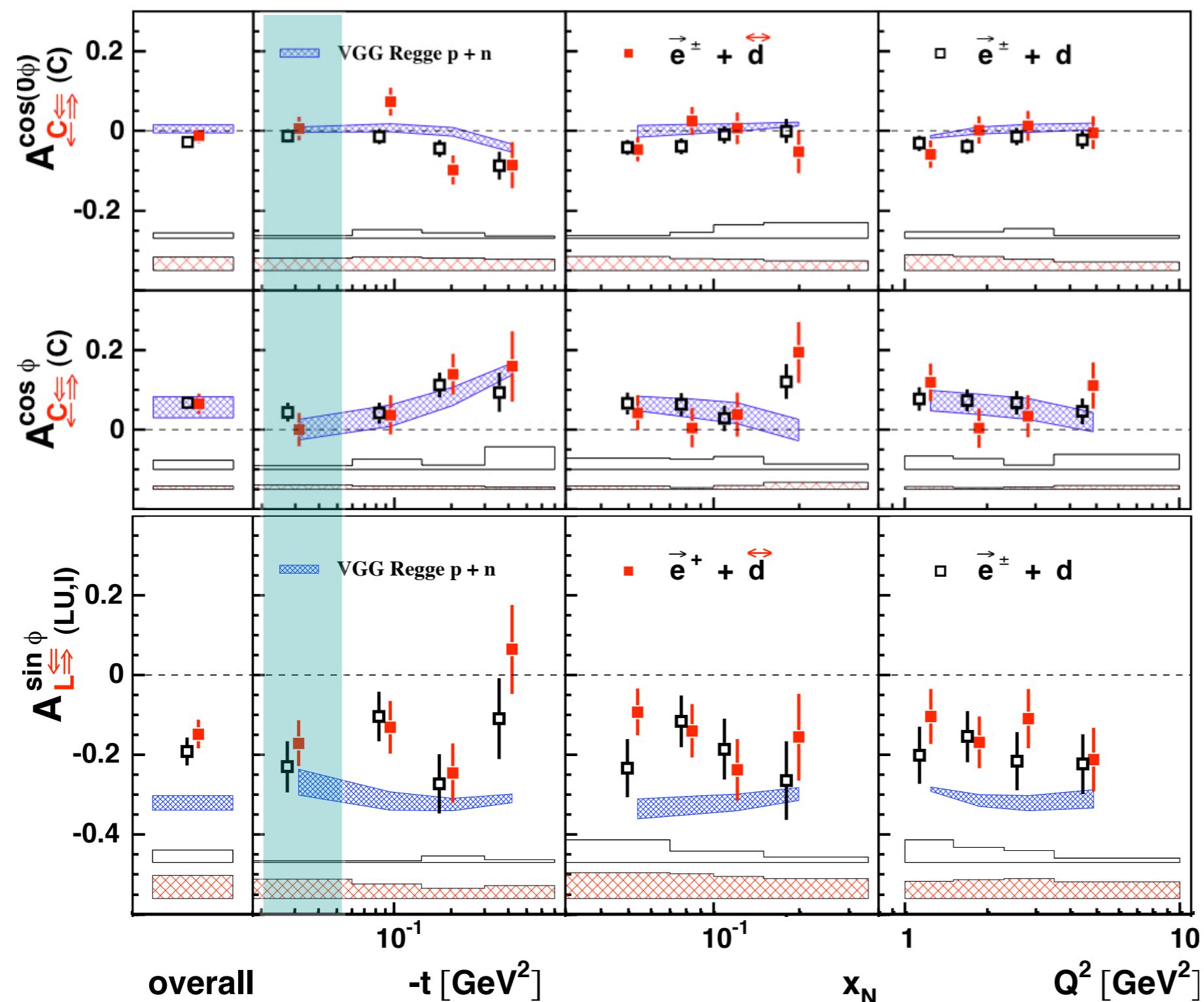


$$\propto -A \cos(\phi)$$

$$\propto \begin{aligned} & Re[G_1(\mathcal{H}_1 - \frac{1}{3}\mathcal{H}_5)] \\ & Re[G_1\mathcal{H}_1] \end{aligned}$$

$$\propto \begin{aligned} & Im[G_1(\mathcal{H}_1 - \frac{1}{3}\mathcal{H}_5)] \\ & Im[G_1\mathcal{H}_1] \end{aligned}$$

- Results on an unpolarized and longitudinally polarized deuterium targets are consistent over all kinematic region.
- Comparison at low $-t$ does not reveal signatures of tensor effects.
Small contribution from CFF \mathcal{H}_5



$$\propto -A \cos(\phi)$$

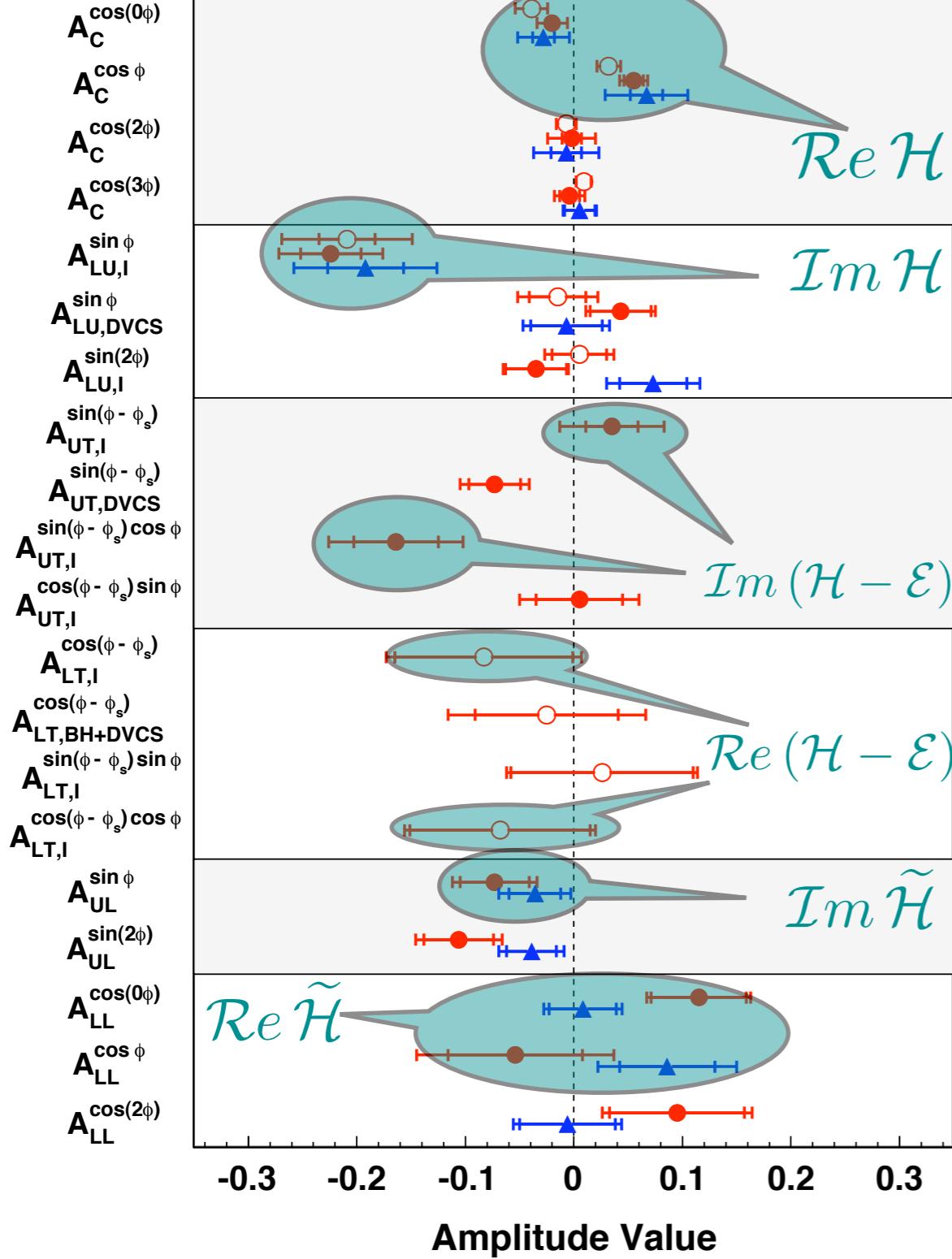
$$\propto \begin{aligned} & Re[G_1(\mathcal{H}_1 - \frac{1}{3}\mathcal{H}_5)] \\ & Re[G_1\mathcal{H}_1] \end{aligned}$$

$$\propto \begin{aligned} & Im[G_1(\mathcal{H}_1 - \frac{1}{3}\mathcal{H}_5)] \\ & Im[G_1\mathcal{H}_1] \end{aligned}$$

- Results on an unpolarized and longitudinally polarized deuterium targets are consistent over all kinematic region.
- Comparison at low $-t$ does not reveal signatures of tensor effects.
Small contribution from CFF \mathcal{H}_5

HERMES DVCS

● Hydrogen
▲ Deuterium
○ Hydrogen Preliminary



Beam-Charge Asymmetry

Beam-Spin Asymmetry

Transverse Target-Spin Asymmetry

Transverse Double-Spin Asymmetry

Longitudinal Target-Spin Asymmetry

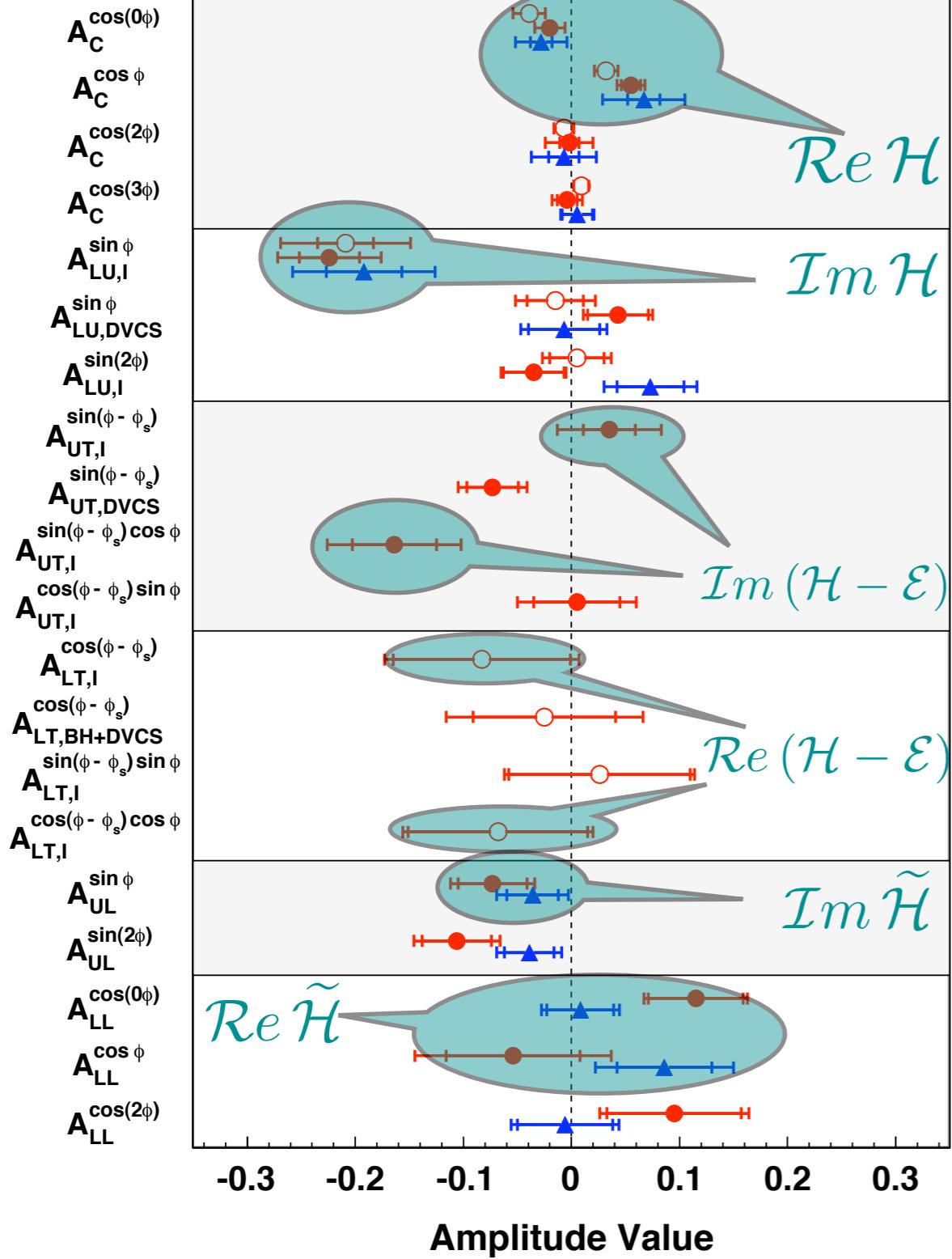
Longitudinal Double-Spin Asymmetry

+ BCA and BSA on nuclear targets

First results with Recoil measurement: See talk by S.Yaschenko

HERMES DVCS

● Hydrogen
▲ Deuterium
○ Hydrogen Preliminary



Beam-Charge Asymmetry

JHEP 11 (2009) 083

Nucl. Phys. B829 (2010) 1

Beam-Spin Asymmetry

JHEP 11 (2009) 083

Nucl. Phys. B829 (2010) 1

Transverse Target-Spin Asymmetry

JHEP 06 (2008) 066

Transverse Double-Spin Asymmetry

Preliminary

Longitudinal Target-Spin Asymmetry

JHEP 06 (2010) 019

Nucl. Phys. B842 (2011) 265

Longitudinal Double-Spin Asymmetry

JHEP 06 (2010) 019

Nucl. Phys. B842 (2011) 265

+ BCA and BSA on nuclear targets

Phys. Rev. C81, 035202 (2010)

First results with Recoil measurement: See talk by S.Yaschenko