What happens to atoms and molecules during XFEL pulses?

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XATOM and XMOLECULE developers



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Overview

- Introduction to XFEL—matter interaction
- > XATOM and x-ray multiphoton multiple ionization dynamics of Xe
- > XMOLECULE and x-ray-induced ultrafast explosion dynamics of CH₃I

> Summary



XFEL: X-ray free-electron laser







Why ultraintense and ultrafast?

- > Structural determination of biomolecules with x-rays
 → X-ray crystallography
- > Growing high-quality crystals is one of major bottlenecks
- Enough signals obtained from even single molecules by using ultraintense pulses
- Signals obtained before radiation damage by using *ultrafast* pulses



Gaffney & Chapman, Science 316, 1444 (2007).

How does matter interact with *ultraintense* and *ultrafast* pulses?





What high x-ray intensity means?



High x-ray intensity beyond one-photon absorption saturation

- synchrotron: at most one photon absorbed \rightarrow linear phenomena
- XFEL: at least one photon absorbed → nonlinear phenomena





Fundamental x-ray-matter interaction







Sequential multiphoton multiple ionization



Young et al., Nature 466, 56 (2010).

- > First LCLS experiment: fundamental atomic physics in XFEL
- Lots of x-ray photons: repeated K-shell ionization (P) followed by Auger (A)





Challenges for x-ray multiphoton ionization

- No standard quantum chemistry code available
- Theoretical challenges
 - tremendously many hole states by x-ray multiphoton absorption
 - highly excited system far from the ground state
 - electronic continuum states for ionization
 - complex inner-shell ionization dynamics, especially for heavy atoms





XATOM

- XATOM: describes
 dynamical behaviors of atoms interacting with
 XFEL pulses
- X-ray-induced atomic processes for any given element and configuration
- Sequential ionization model has been tested by a series of atomic XFEL experiments: Ne, Ar, Kr, Xe, …



Son, Young & Santra, *Phys. Rev. A* **83**, 033402 (2011). Jurek, Son, Ziaja & Santra, *J. Appl. Cryst.* **49**, 1048 (2016). Download executables: <u>http://www.desy.de/~xraypac</u>







XATOM: Electronic structure

- Efficient electronic structure calculation required
- Hartree-Fock-Slater method

$$\begin{bmatrix} -\frac{1}{2}\nabla^2 - \frac{Z}{r} + \int d^3r' \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} - \frac{3}{2} \left\{ \frac{3}{\pi} \rho(\mathbf{r}) \right\}^{1/3} \end{bmatrix} \psi(\mathbf{r}) = \varepsilon \psi(\mathbf{r})$$
spherically symmetric: $\psi_{nlm}(\mathbf{r}) = \frac{u_{nl}(r)}{r} Y_{lm}(\theta, \varphi)$

> Bound states → generalized pseudospectral method on *nonuniform* grid

➤ Continuum states: calculated with the same potential as used in bound states → 4th-order Runge-Kutta method on *uniform* grid

Son, Young & Santra, *Phys. Rev. A* 83, 033402 (2011).





XATOM: Cross sections and rates

> Calculate all cross sections and rates of x-ray-induced processes based on the perturbation theory

$$\Gamma_{FI} = 2\pi\delta \left(E_F - E_I \right) \left| \langle F | \hat{H}_{\text{int}} | I \rangle \right|^2$$

> Photoionization cross section

$$\sigma_{\rm P}(i,\omega) = \frac{4}{3} \alpha \pi^2 \omega N_i \sum_{l_j = |l_i - 1|}^{l_i + 1} \frac{l_{>}}{2l_i + 1} \left| \langle u_{n_i l_i}(r) | r | u_{\varepsilon l_j}(r) \rangle \right|^2$$

> Fluorescence rate

$$\Gamma_{\rm F}(i,j) = \frac{4}{3} \alpha^3 (I_i - I_j)^3 \frac{N_i^{\rm H} N_j}{4l_j + 2} \cdot \frac{l_{>}}{2l_i + 1} \left| \langle u_{n_i l_i}(r) | r | u_{n_j l_j}(r) \rangle \right|^2$$

> Auger rate

$$\Gamma_{A}(i,jj') = \pi \frac{N_{i}^{H}N_{jj'}}{2l_{i}+1} \sum_{L=|l_{j}-l_{j'}|}^{l_{j}+l_{j'}} \sum_{S=0}^{1} \sum_{l_{i'}}^{1} (2L+1)(2S+1)|M_{LS}(j,j',i,i')|^{2}$$





XATOM: Coupled rate equations

- > Electronic structure: calculated for every single configuration
- Cross sections and rates: calculated for every single configuration
- Solve coupled rate equations to simulate ionization dynamics

$$\frac{d}{dt}P_I(t) = \sum_{I' \neq I}^{\text{all config.}} \left[\Gamma_{I' \to I}P_{I'}(t) - \Gamma_{I \to I'}P_I(t)\right]$$

Tremendously large coupled rate equations (~millions configurations)
 solved by Monte Carlo approach



Son & Santra, *Phys. Rev. A* **85**, 063415 (2012).





Complex inner-shell decay cascade





Multiphoton absorption after/during decay cascade

- More than 20 million multiple-hole configurations
- More than 2 billion x-ray-induced processes





X-ray multiphoton ionization dynamics



X-ray multiphoton ionization dynamics



X-ray multiphoton ionization dynamics



Ionization mechanism described by theory

To reach Xe²⁴⁺: 5 photons absorbed, 24 electrons ejected







Charge-state distributions of Xe



Comparison with LCLS experiment



LCLS experiment







Artem Rudenko at KSU



- Xe M-shell ionization
- 2 keV: excellent agreement between theory and experiment
- 1.5 keV: further ionization via resonance

Rudek et al., Nature Photon. 6, 858 (2012).





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Ultra-efficient ionization by XFEL

- > REXMI: <u>Resonance-Enabled X-ray Multiple Ionization</u>
- > Broad bandwidth (~15 eV): resonances for many charge states



Rudek et al., Nature Photon. 6, 858 (2012).



Comparison with SACLA experiment



SACLA experiment



Kiyoshi Ueda at Tohoku Univ.

- Hironobu Fukuzawa
- Koji Motomura

Fukuzawa et al., *Phys. Rev. Lett.* **110**, 173005 (2013).

- Xe L-shell ionization: good agreement
- underestimation in theory: lack of relativity, shake-off, and resonance





Improving theory

> One resonance channel: Ne@1050 eV

- Xiang et al., Phys. Rev. A 86, 061401 (2012).
- > Multiple resonance channels: Xe@1.5 keV
 - Ho et al., Phys. Rev. Lett. **113**, 253001 (2014).
 - Ho, Kanter, and Young, *Phys. Rev. A* **92**, 063430 (2015).
- > What about both resonance effect and relativistic effects?
 - resonance effect: Xe M-shell ionization at 1.5 keV
 - relativistic effects: Xe L-shell ionization at 5.5 keV
 - Interplay between resonant excitations and relativistic corrections





Inclusion of relativistic effects

Relativistic energy correction within first-order perturbation theory

$$\hat{H} = \hat{H}_0 - \frac{\alpha^2}{8}\hat{p}^4 - \frac{\alpha^2}{4}\frac{dV}{dr}\frac{d}{dr} + \frac{\alpha^2}{2}\frac{1}{r}\frac{dV}{dr}\hat{l}\cdot\hat{s}$$
$$E_{nlj} = \varepsilon_{nl} + \Delta\varepsilon_{nl}^{\text{mass}} + \Delta\varepsilon_{nl}^{\text{Darwin}} + \Delta\varepsilon_{nlj}^{\text{so}}$$

Toyota, Son & Santra, *Phys. Rev. A* **95**, 043412 (2017).

Open new Coster-Kronig decay channels due to spin-orbit splitting

group	non-rel (a.u.)	rel	McGuire [37]
M_2 -X	1.61×10^{-4}	2.20×10^{-4}	1.60×10^{-4}
M_3 -X		1.75×10^{-4}	1.37×10^{-4}
M_2 - XY	2.09×10^{-2}	1.97×10^{-2}	2.27×10^{-2}
M_3 - XY		2.06×10^{-2}	2.13×10^{-2}
M_2 - M_3X	forbidden	5.38×10^{-3}	$5.50 imes 10^{-3}$
$M_2 - M_{45}X$	2.05×10^{-1}	1.54×10^{-1}	1.50×10^{-1}
$M_3 - M_{45}X$		1.88×10^{-1}	1.80×10^{-1}

Close photoionization earlier due to relativistic energy corrections



Closer to photon energy due to energy corrections: higher cross sections





Relativistic heavy atom

> Number of all possible electronic configurations

non-relativistic

Xe: 1s² 2s² 2p⁶ 3s² 3p⁶ 3d¹⁰ 4s² 4p⁶ 4d¹⁰ 5s² 5p⁶

M-shell: 1,120,581 *L*-shell: 23,532,201

relativistic

Xe: $1s_{\frac{1}{2}}^{2} 2s_{\frac{1}{2}}^{2} 2p_{\frac{1}{2}}^{2} 2p_{\frac{3}{2}}^{4} 3s_{\frac{1}{2}}^{2} 3p_{\frac{1}{2}}^{2} 3p_{\frac{3}{2}}^{4} 3d_{\frac{3}{2}}^{4} 3d_{\frac{5}{2}}^{6}$ $4s_{\frac{1}{2}}^{2} 4p_{\frac{1}{2}}^{2} 4p_{\frac{3}{2}}^{4} 4d_{\frac{3}{2}}^{4} 4d_{\frac{5}{2}}^{6} 5s_{\frac{1}{2}}^{2} 5p_{\frac{1}{2}}^{2} 5p_{\frac{3}{2}}^{4}$ *M*-shell: 111,628,125 *L*-shell: 5,023,265,625





Inclusion of resonant excitations

> Resonant photoexcitation cross section

SCIENCE

$$\sigma_{\rm R}(i, f, \omega_{\rm in}) = \frac{4}{3}\pi^2 \alpha \omega_{\rm in} l_> N_i N_f^H \left\{ \begin{array}{c} l_i & s & j_i \\ j_f & 1 & l_f \end{array} \right\}^2 |\langle u_{n_f l_f}(r) | r | u_{n_i l_i}(r) \rangle|^2 \\ \times \delta(E - E_{n_f l_f j_m} + E_{n_i l_i j_i}) \\ \times \delta(E - E_{n_f l_f j_m} + E_{n_i l_i j_i}) \\ \end{array} \right\}$$
Number of possible electron configurations explodes depending on n_{max} (=15~20) and l_{max} (=6~7).
Toyota, Son & Santra, Phys. Rev. A **95**, 043412 (2017).

Interplay between resonance and relativity



SCIENCE

Xe atom at higher x-ray intensity

- New experimental setup:
 LCLS CXI using nano-focus
 new realm of intensity approaching ~10²⁰ W/cm²
- > Various photon energies: 5.5 keV ~ 8.3 keV
 - Trend of REXMI examined
 - Peak structure in the REXMI region
 - L-shell initiated ionization → large relativistic effects
- Improved experimental analysis and improved calibration procedure with Ar data XCALIB: developed by Zoltan Jurek and Koudai Toyota





Benedikt Rudek at PTB

XCALIB development







Zoltan Jurek





New data for Xe at 5.5 keV



SCIENCE

Preliminary results

Rudek, Toyota, et al., (in preparation).

- ~40× higher fluence than before
- REXMI region: accurate electronic structure of bound and Rydberg states
- > Only relativistic, resonance calculation reproduces well the peak structure of experiment.
- Importance of interplay between resonance and relativistic effects





Challenges for molecular dynamics at XFEL

No ab initio theoretical tools available for high x-ray intensity

- coupled ionization and nuclear dynamics in the same time scales
- formidable task: e.g. CH₃I ~ 200 trillion rate eqs at single geometry
- highly excited molecular electronic structure

XMOLECULE

- quantum electrons, classical nuclei
- efficient electronic structure calculation: core-hole adapted basis functions calculated by XATOM
- Monte Carlo on the fly

Hao, Inhester, Hanasaki, Son & Santra, *Struc. Dyn.* **2**, 041707 (2015). Inhester, Hanasaki, Hao, Son & Santra, *PRA* **94**, 023422 (2016).





XMOLECULE: Electronic structure

- Efficient electronic structure calculation required
- Hartree-Fock-Slater method
- Bound states: LCAO-MO

$$\psi_i(\mathbf{r}) = \sum_{\mu} C_{\mu i} \phi_{\mu}(\mathbf{r})$$

- Core-hole-adapted NAO

$$H_{\mu\nu} = \int d^3 r \,\phi_\mu(\mathbf{r}) \left[-\frac{1}{2} \nabla^2 + V_{\text{eff}}(\mathbf{r}) \right] \phi_\nu(\mathbf{r}), \quad S_{\mu\nu} = \int d^3 r \,\phi_\mu(\mathbf{r}) \phi_\nu(\mathbf{r})$$

Cross sections and rates calculated on the fly for given elec. structure

Stradients calculated on the fly for given electronic structure





N^{2+:} 1s⁰2s²2p³

N₂³⁺

XMOLECULE: Core-hole-adapted NAOs



> NAO: numerical atomic orbitals calculated by XATOM

$$\psi_i(\mathbf{r}) = \sum_{\mu} C_{\mu i} \phi_{\mu}(\mathbf{r}) \qquad \qquad \phi_{nlm}(\mathbf{r}) = \frac{u_{nl}(r)}{r} Y_{lm}(\theta, \varphi)$$

Molecular core-hole states with corresponding atomic core-hole states

Sood treatment for molecular core-hole states with a minimal basis set Hao, Inhester, Hanasaki, Son & Santra, Struc. Dyn. 2, 041707 (2015).





Molecular cross sections and rates

> Molecular continuum approximated by atomic continuum: $\psi_{\varepsilon} \rightarrow \phi_{A\varepsilon}$

Photoionization cross section

$$\sigma_{\rm P} \propto \left| \langle \psi_i | d | \psi_{\varepsilon} \rangle \right|^2 \approx \sum_A \left| \sum_{\mu}^{\text{on atom } A} C_{\mu i} \langle \phi_{\mu} | d | \phi_{A\varepsilon} \rangle \right|^2$$

> Fluorescence rate $\Gamma_{\rm F} \propto |\langle \psi_i | d | \psi_j \rangle|^2$ evaluated by multicenter integration

> Auger rate: one-center approximation

$$\begin{split} \langle jj'|i\varepsilon\rangle &= \int d^3r_1 \int d^3r_2 \,\psi_j(\mathbf{r}_1)\psi_{j'}(\mathbf{r}_2) \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \psi_i(\mathbf{r}_1)\psi_\varepsilon(\mathbf{r}_2) \\ \langle jj'|i\varepsilon\rangle &\approx \sum_{\mu,\nu,\lambda}^{\text{on atom }A} C_{\mu j} C_{\nu j'} C_{\lambda i} \langle \mu \nu | \lambda \varepsilon \rangle \quad \text{ intra-atomic process included} \end{split}$$

Inhester, Hanasaki, Hao, Son & Santra, Phys. Rev. A 94, 023422 (2016).





Various multiple-hole states of CO



All possible multiple-hole configurations (*N*=2187) formed by x-ray multiphoton ionization (at single molecular geometry)





Molecules at low x-ray intensity



Erk *et al.*, *PRL* **110**, 053003 (2013).



CH₃I: charge rearrangement as a function of bond distance Erk *et al., Science* **345**, 288 (2014).

Total charge of molecule is similar to atomic charge. Heavy atom charges are reduced after charge rearrangement. Still valid for high x-ray intensity?





Iodomethane at higher x-ray intensity

- New experimental setup: LCLS CXI using nano-focus
 → new realm of intensity approaching ~10²⁰ W/cm²
- Selective ionization on heavy atom

LCLS experiment





Daniel Rolles at KSU

Artem Rudenko at KSU

CH₃I @ 8.3 keV



σ(I)~50 kbarn σ(C)~80 barn σ(H)~8 mbarn

- X-ray multiphoton ionization occurs at high intensity
- > Charge imbalance induces charge rearrangement
- > Coulomb explosion after/during ionization & charge rearrangement





Coulomb explosion of iodomethane

CH₃I (t = 0 fs)







Coulomb explosion of iodomethane







What happened?







Comparison of CSD and KER



Capturing ultrafast ionization and fragmentation dynamics

- CSD (charge-state distribution): direct outcome of ionization dynamics
- KER (kinetic energy release): molecular information when it breaks apart, influenced by detailed dynamical behaviors





Molecular ionization enhancement



molecular charge > ∑(atomic charges) theoretically predicted and experimentally confirmed





Ionization enhanced by charge rearrangement



- Electrons from light atoms become available for further ionization on heavy atoms after charge rearrangement.
- CREXIM: <u>Charge-Rearrangement-Enhanced X-ray Ionization of Molecules</u>
- Impact on molecular imaging: not reducing partial charges of heavy atoms due to charge rearrangement, but inducing more ionization overall





The bigger molecule, the larger effect



- > Coincidence measurement at intermediate intensity: higher iodine charge always along with highest carbon charge
- > Estimated molecular charge: Xe^{48+} , CH_3I^{54+} , and $C_6H_5I^{>54+}$





Molecular black hole

DESY news: http://www.desy.de/news/news_search/ index_eng.html?openDirectAnchor=1232

The extremely intense X-ray flash knocks so many electrons out of the iodine atom (right) such that it pulls in the electrons of the methyl group (left) like an electromagnetic version of a black hole, before finally spitting them out.

"Femtosecond response of polyatomic molecules to ultra-intense hard X-rays," Rudenko *et al., Nature* **546**, 129 (2017).







Summary



- Enabling tools to investigate x-ray multiphoton physics of atoms, molecules, and clusters exposed to high intensity x-ray pulses
- X-ray multiphoton ionization dynamics of Xe
 - interplay between resonance and relativistic effects
- X-ray-induced ultrafast explosion dynamics of CH₃I
 - first quantitative comparison for molecules under XFEL irradiation
 - molecular ionization enhancement via CREXIM





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Thank you for your attention!



