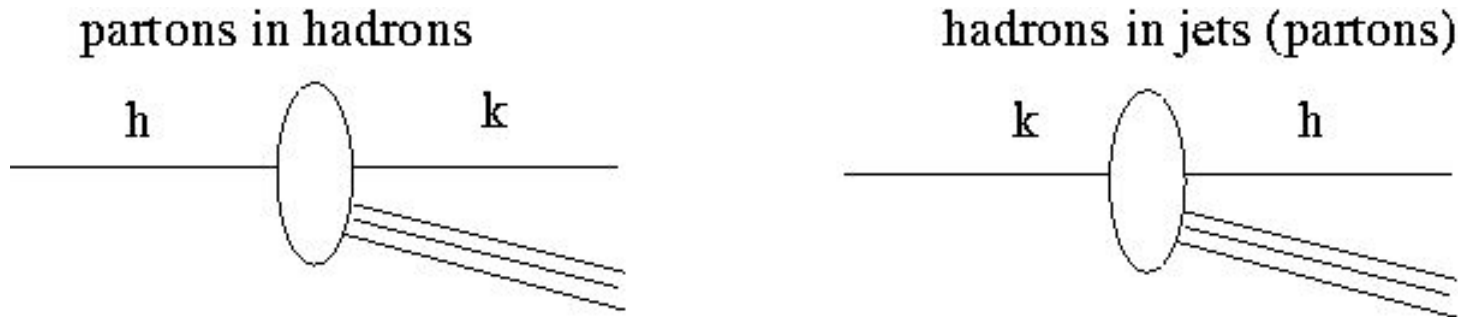


# QCD and Collider Physics III: Hadronization

- fragmentation function
  - heavy quarks
  - light quarks
- hadronization models
  - independent, cluster & string models
  - color flows in proton – proton interactions
- Literature:
  - Andersson et al., Parton fragmentation and string dynamics, PhysRep 97 (1983) 31
  - Andersson, The Lund Model
  - Barger/Phillips: *Collider Physics*
  - Dissertori, Knowles, Schmelling: *QCD - High Energy Exp and Theory*
  - Ellis, Stirling, Webber: *QCD and Collider Physics*
  - Field: *Applications of perturbative QCD*

[http://www-h1.desy.de/~jung/qcd\\_collider\\_physics\\_wise\\_2006](http://www-h1.desy.de/~jung/qcd_collider_physics_wise_2006)

# Scaling violations of Frag. Fcts.



- Similarity with evolution of parton density functions

$$t \frac{\partial}{\partial t} D_i(x, t) = \sum_j \int_x^1 \frac{dz}{z} \frac{\alpha_s}{2\pi} P_{ji}(z, \alpha_s) D_j\left(\frac{x}{z}, t\right)$$

- with splitting functions:  $P_{ji}(x, \alpha_s) = P_{ji}^{(0)} + \frac{\alpha_s}{2\pi} P_{ji}^{(1)}$
- lowest order splitting functions are the same as for PDF case
- higher order  $P_{gg}, P_{qg}$  are more singular than in PDFs
- resummation of small  $x$  enhanced terms have different behavior...

# Peterson fragmentation

PHYSICAL REVIEW D

VOLUME 27, NUMBER 1

1 JANUARY 1983

## Scaling violations in inclusive $e^+e^-$ annihilation spectra

C. Peterson,\* D. Schlatter, I. Schmitt,<sup>†</sup> and P. M. Zerwas<sup>‡</sup>

Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305

(Received 29 July 1982)

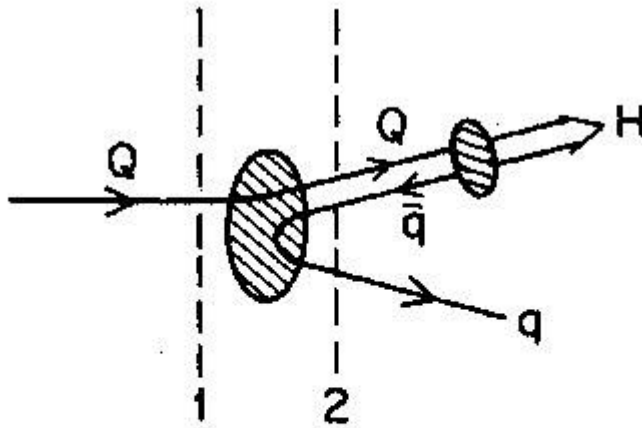


FIG. 3. The fragmentation of a heavy quark  $Q$  into a meson  $H(Q\bar{q})$ . Dashed lines are time slices used in the derivation of Eq. (3).

cussed in Ref. 18. The gross features of the amplitude for a fast moving heavy quark  $Q$  fragmentation into a hadron  $H=(Q\bar{q})$  and light quark  $q$  (Fig. 3) are determined by the value of the energy transfer  $\Delta E = E_H + E_q - E_Q$  in the breakup process,

$$\text{amplitude } (Q \rightarrow H + q) \propto \Delta E^{-1}. \quad (2)$$

Expanding the energies about the (transverse) particle masses ( $m_H \simeq m_Q$  for simplicity),

$$\begin{aligned} \Delta E &= (m_Q^2 + z^2 P^2)^{1/2} + (m_q^2 + (1-z)^2 P^2)^{1/2} \\ &\quad - (m_Q^2 + P^2)^{1/2} \\ &\propto 1 - (1/z) - (\epsilon_Q/1-z) \end{aligned} \quad (3)$$

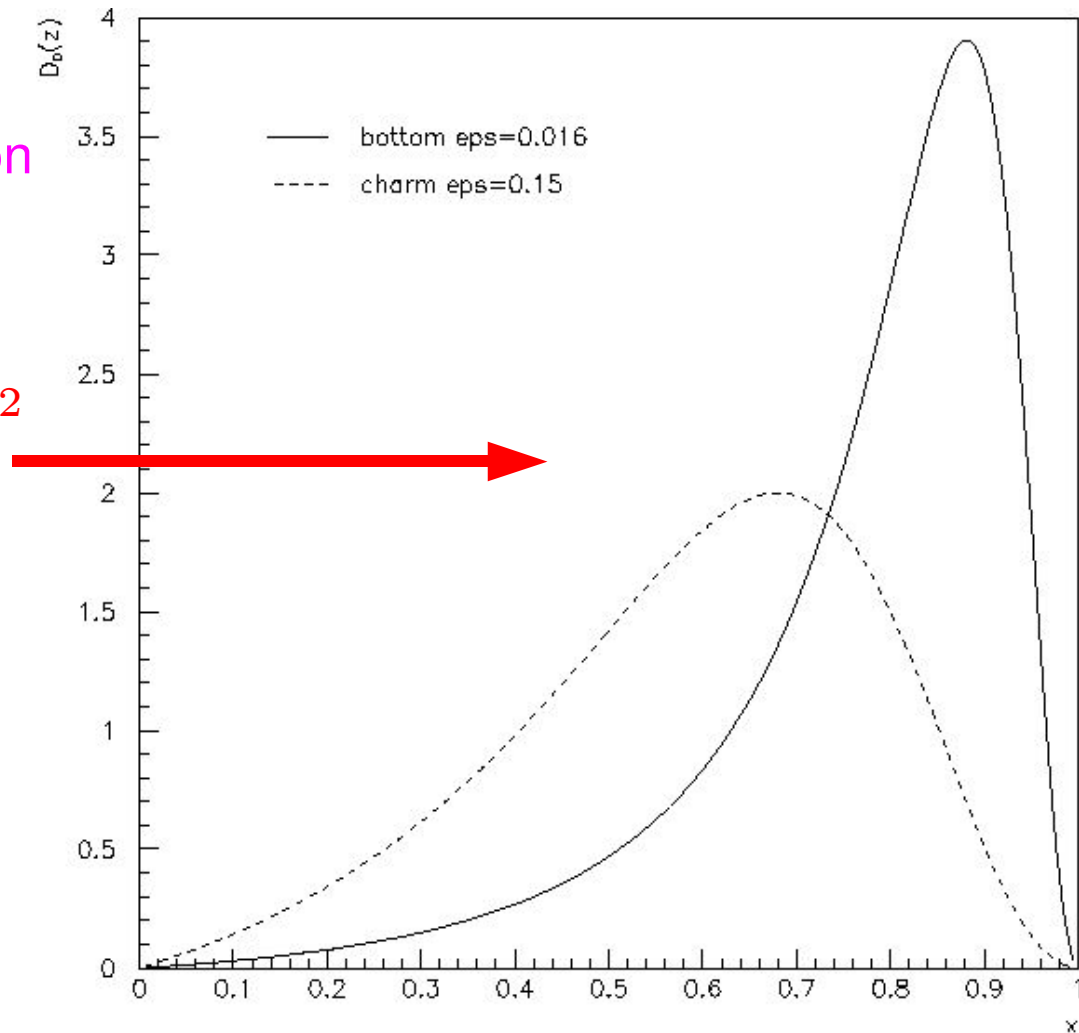
and taking a factor  $z^{-1}$  for longitudinal phase space, we suggest the following ansatz for the fragmentation function of heavy quarks  $Q$

$$D_Q^H(z) = \frac{N}{z[1 - (1/z) - \epsilon_Q/(1-z)]^2}. \quad (4)$$

# Heavy Quark Fragmentation

- transition from heavy quark to observable hadron by fragmentation function FF
- **Peterson FF:** (C. Peterson, D.Schlatter,I.Schmitt,P.Zerwas, PRD27 (1983) 105)

$$D_Q(z) = \frac{N}{z} \left[ 1 - \frac{1}{z} - \frac{\epsilon_Q}{1-z} \right]^{-2}$$



# Heavy Quark Fragmentation

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- **Peterson FF:** (C. Peterson, D.Schlatter,I.Schmitt,P.Zerwas, PRD27 (1983) 105)

$$D_Q(z) = \frac{N}{z} \left[ 1 - \frac{1}{z} - \frac{\epsilon_Q}{1-z} \right]^{-2}$$

- **Kartvelishvili** (V. Kartvelishvili,A.Likhoded,V.Petrov PLB78 (1978) 61)

$$D_Q(z) = N z^\alpha (1-z)$$

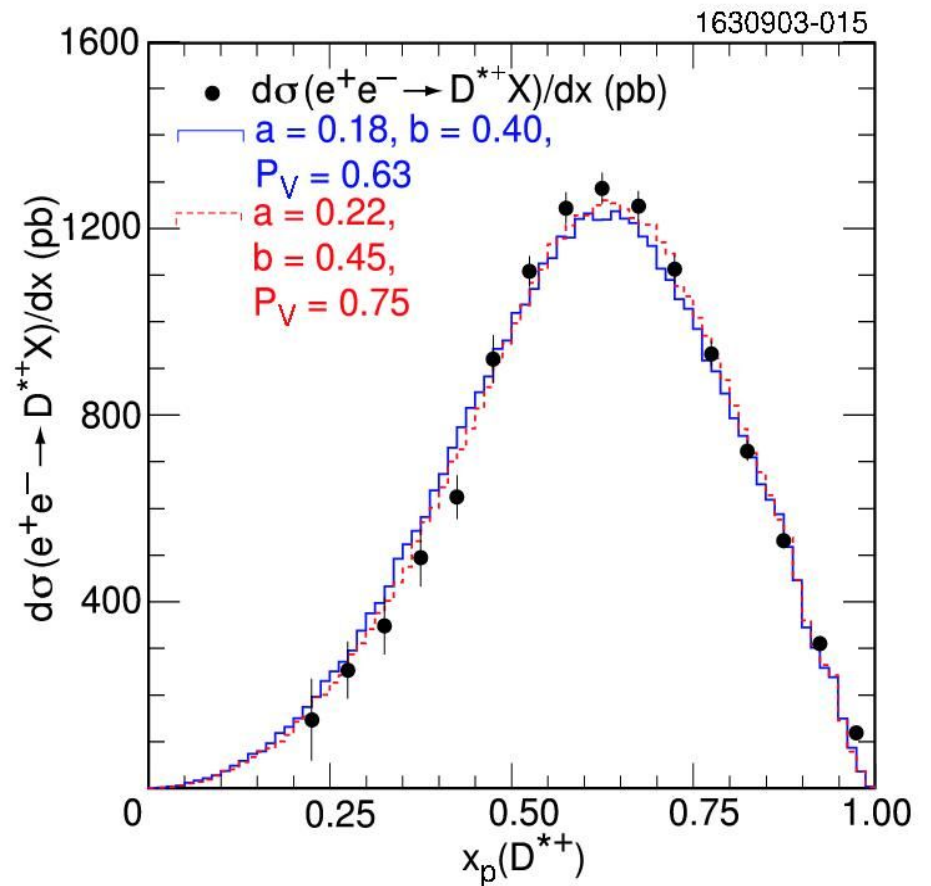
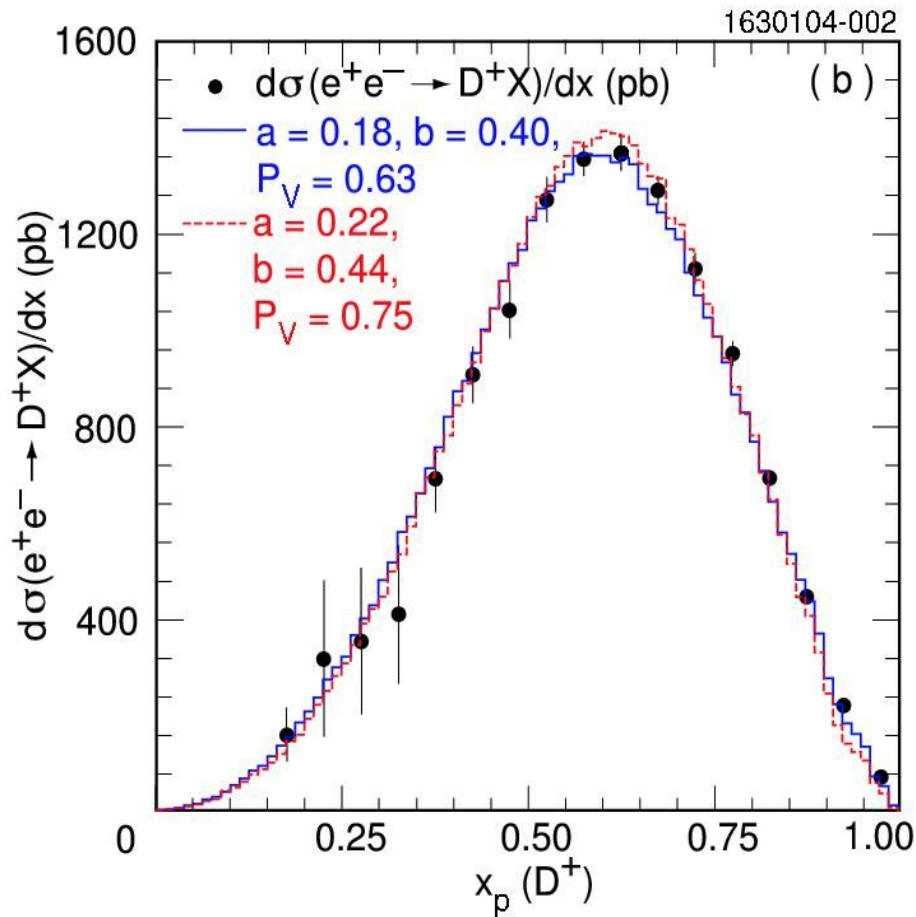
- **Lund FF** (B. Andersson et al, Phys.Rep 97, 31 (1983))

$$D_Q(z) = N \frac{(1-z)^a}{z} \exp \left[ \frac{-bm_\perp^2}{z} \right]$$

# Heavy Quark FF

• Results from CLEO (hep-ex/0402040)

➔ compared to Lund FF



# Light Quark FF

- parametrisations by:  
Kniehl, Kramer & Poetter NPB582  
(2000) 514

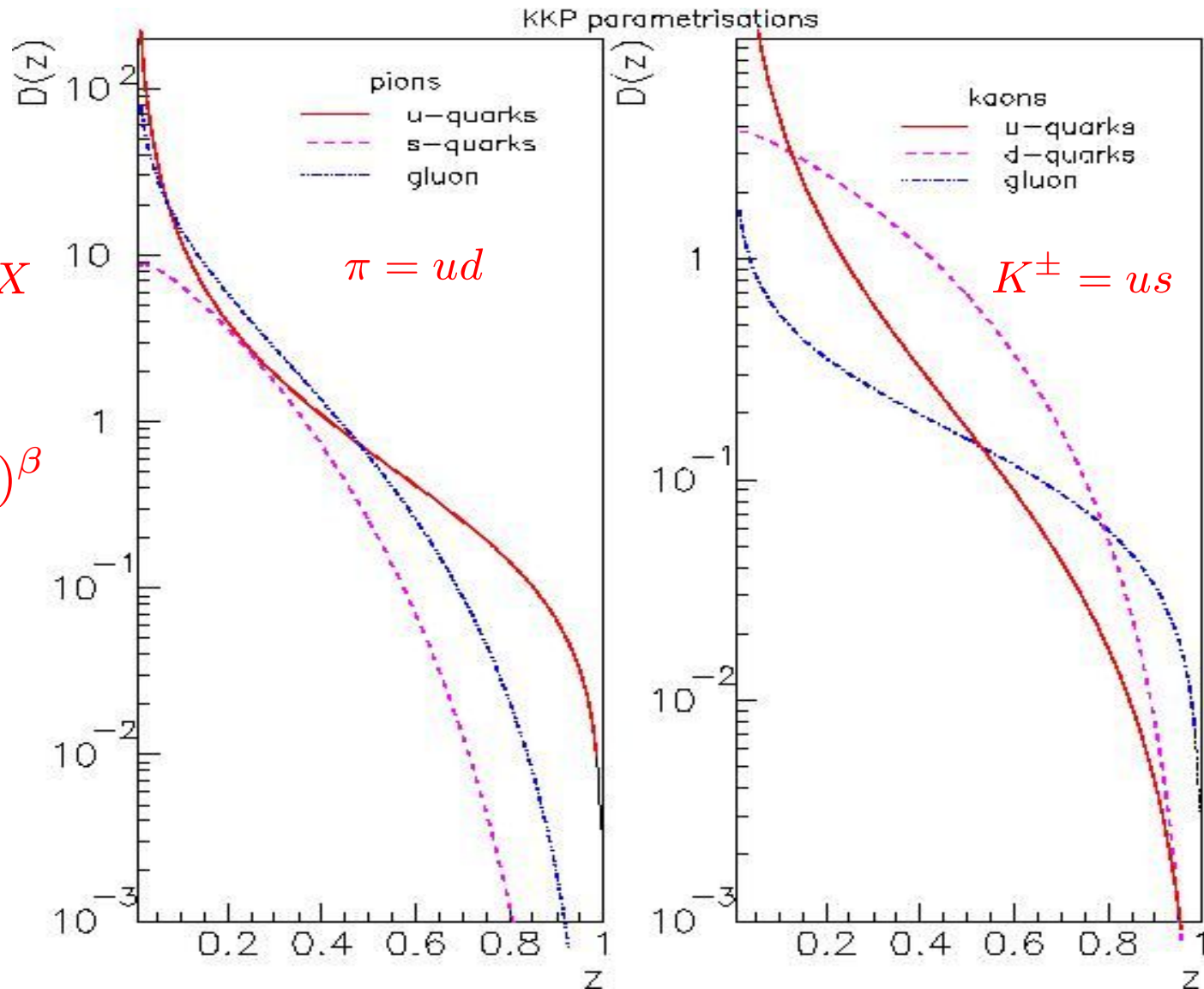
- use

$$e^+e^- \rightarrow (\gamma, Z) \rightarrow h + X$$

- starting distribution:

$$D_Q(z) = N z^\alpha (1-z)^\beta$$

- scaling violations...  
evolution in  $Q^2$



# Fragmentation Models

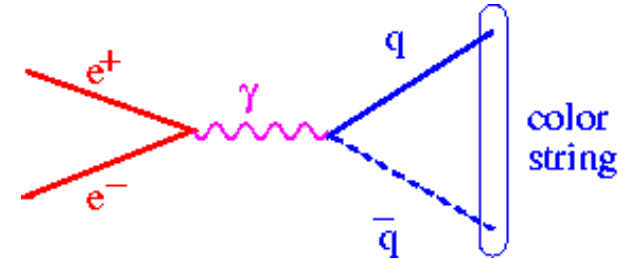
- describe transition from quarks to hadrons
- **Independent Fragmentation** (Feynman & Field: Phys. Rev D15 (1977)2590, NPB 138 (1978) 1)
  - quarks fragment independently
  - gluon are split:  $g \rightarrow q\bar{q}$
  - fragmentation depends on momentum (energy), **but not** on virtuality
  - **not** Lorentz invariant
  - with 4 parameters can describe broad features of 2-jet and 3-jet
- **Lund String Fragmentation** (Andersson, Gustafson, Ingelman, Sjostrand Phys. Rep 97 (1983) 33)
  - for qq is similar to independent fragmentation
    - **BUT** is covariant and has no leftover
    - constraints on fragmentation function:  $q\bar{q}$  symmetric
    - transverse momentum distribution from tunneling effect
- **Cluster Fragmentation** (Webber NPB 238 (1984) 492)
  - pre-confinement of color
  - gluon split  $g \rightarrow q\bar{q}$



# Fragmentation: simple example

- process  $e^+e^- \rightarrow q\bar{q}$

- $$\frac{d\sigma}{d\cos\theta d\phi} = \frac{\alpha_{em}^2}{4s} (1 + \cos^2\theta)$$



- **BUT** what about fragmentation/ hadronization ???
- use concept of **local parton-hadron duality**

linear confinement potential:  $V(r) \sim -1/r + \kappa r$   
with  $\kappa \sim 1 \text{ GeV/fm}$

qq connected via color flux tube of transverse size of hadrons ( $\sim 1 \text{ fm}$ )  
color tube: uniform along its length  $\rightarrow$  linearly rising potential

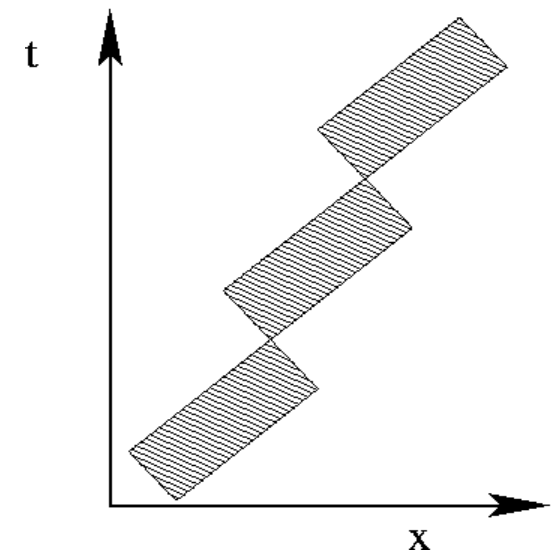
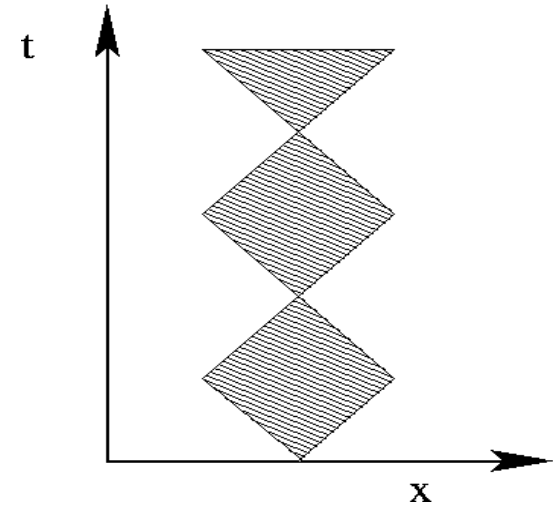
**$\rightarrow$  Lund string fragmentation**

# Lund string fragmentation

- in a color neutral qq-pair, a color force is created in between
- color lines of the force are concentrated in a narrow tube connecting q and  $\bar{q}$ , with a string tension of:

$$\kappa \sim 1 \text{ GeV/fm} \sim 0.2 \text{ GeV}^2$$

- as q and  $\bar{q}$  are moving apart in qq rest frame, they are de-accelerated by string tension, accelerated back etc ... (periodic oscillation)
- viewed in a moving system, the string is boosted



# Fragmentation in the String Model

- hadronization: iterative process
- string breaks in  $q\bar{q}$  pairs (still respecting color flow)
- select transverse motion with  $m=m_{qq}$  (and flavor)

$$P \sim \exp\left(-\frac{\pi m_t^2}{\chi}\right) = \exp\left(-\frac{\pi m^2}{\chi}\right) \exp\left(-\frac{\pi p_t^2}{\chi}\right)$$

- suppression of heavy quark production

$$u : d : s : c \sim 1 : 1 : 0.37 : 10^{-10}$$

actually leave it as a free parameter

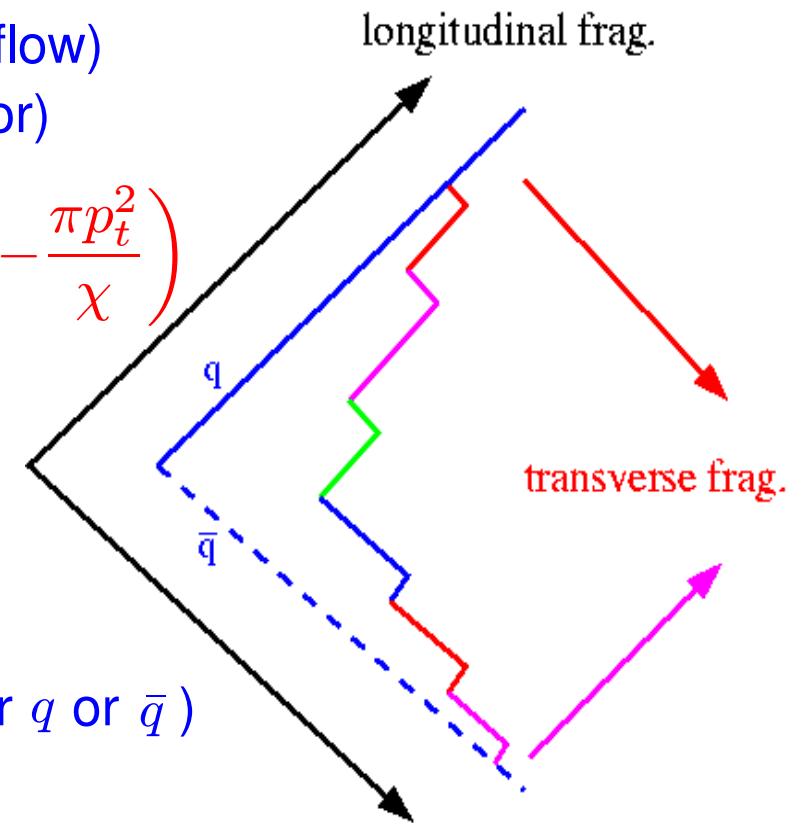
- longitudinal fragmentation

symmetric fragmentation function (from either  $q$  or  $\bar{q}$ )

$$f(z) \sim \frac{1}{z}(1-z)^a \exp\left[-\frac{bm_{\perp}^2}{z}\right]$$

harder spectrum for heavy quarks

- start from  $q$  or  $\bar{q}$
- repeat until cutoff is reached
- heavy use of random numbers and importance sampling method



# Hadronization: particle mass and decays

- particle masses

- taken from PDG, where known, otherwise from constituent masses

- particle widths

- in hard scattering production process short lived particles ( $\rho, \Delta$ ) have nominal mass, without mass broadening

- in hadronization use Breit-Wigner:

$$\mathcal{P}(m)dm \propto \frac{1}{(m - m_0)^2 + \Gamma^2/4}$$

- lifetimes

- related to widths ... but for practical purpose separated

- $P(\tau)d\tau \sim \exp(-\tau/\tau_0) d\tau$

- calculate new vertex position  $v' = v + \tau p/m$

- decays

- taken from PDG, where known

- assume momentum distribution given by phase space only

- exceptions, like  $\omega, \phi \rightarrow \pi^+ \pi^- \pi^0$ , or  $D \rightarrow K\pi$ ,  $D^* \rightarrow K\pi\pi$   
and some semileptonic decays use matrix elements

# Cluster Fragmentation

Volume 83B, number 1

PHYSICS LETTERS

23 April 1979

## PRECONFINEMENT AS A PROPERTY OF PERTURBATIVE QCD

D. AMATI and G. VENEZIANO

*CERN, Geneva, Switzerland*

Received 2 February 1979

The first important point to realize is that, in the axial gauge and at the leading log level we are working in all relevant graphs are planar [2]. It follows that the final quanta can be *ordered*, as shown in fig. 1. Furthermore, there is a natural way to group them (fig. 1) into sets  $C_i$  of adjacent partons each consisting of a quark, an antiquark and a number of gluons. These systems contain a dominant singlet component and,

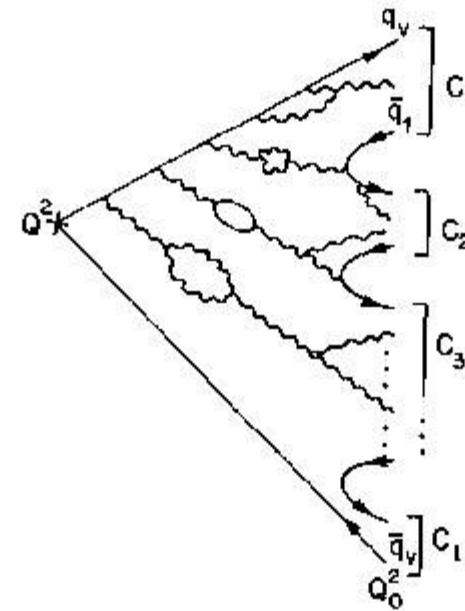
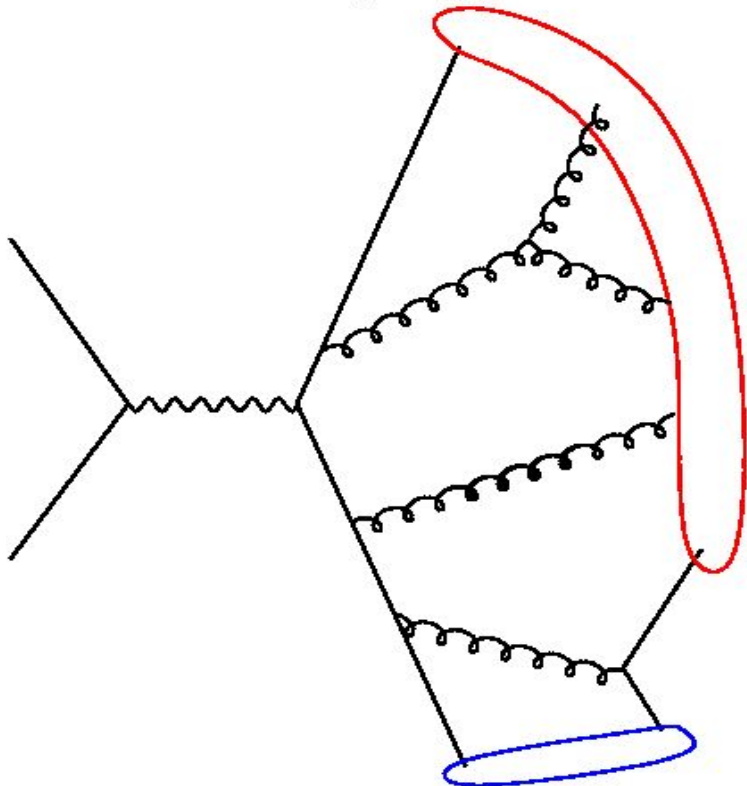


Fig. 1. Planar diagrams contributing to the leading log evolution of  $e^+e^-$  jets in the axial gauge. Final partons are  $O(Q_0^2)$  off-shell and naturally group into ordered colour singlets  $C_i$ .

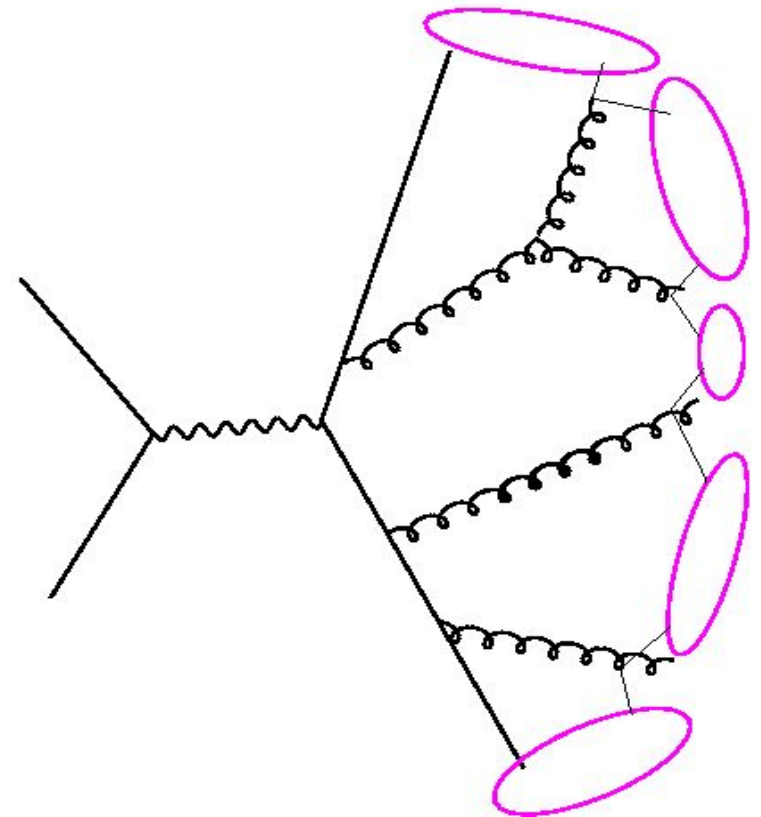
# Cluster Fragmentation

- Pre-confinement of color
- Gluon split  $g \rightarrow q\bar{q}$

String Fragmentation

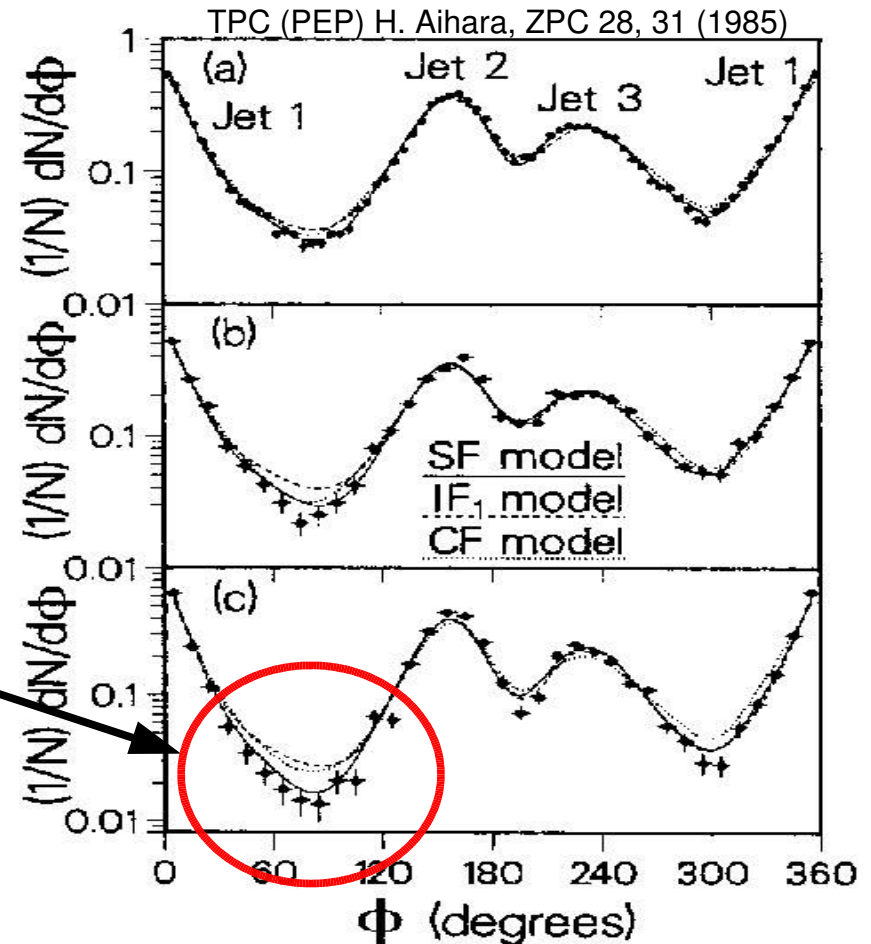
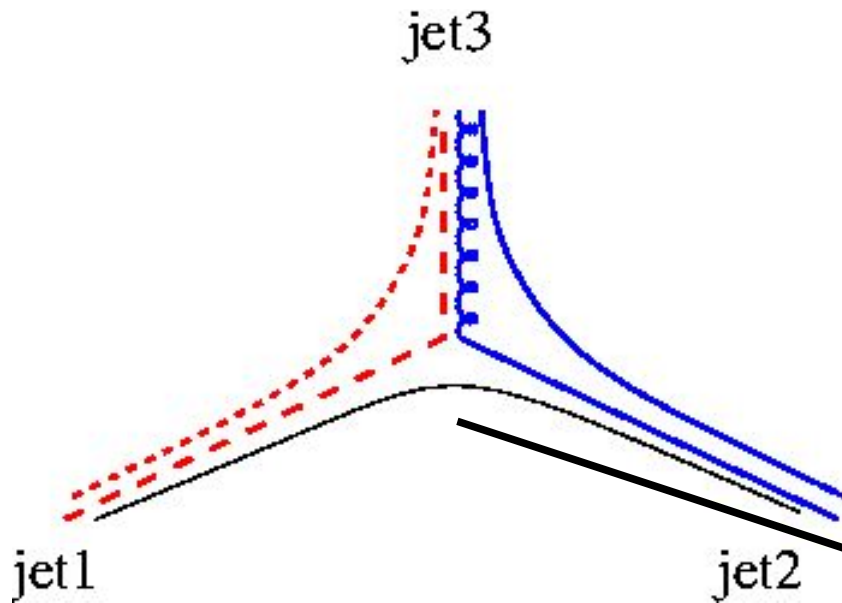
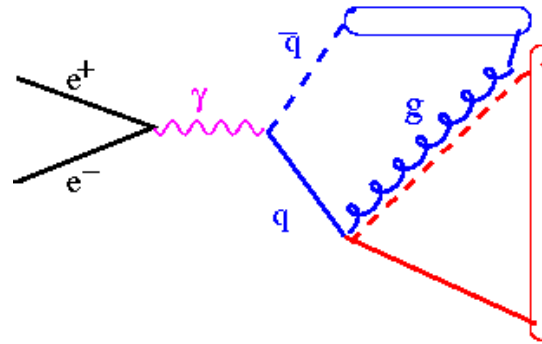


Cluster Fragmentation



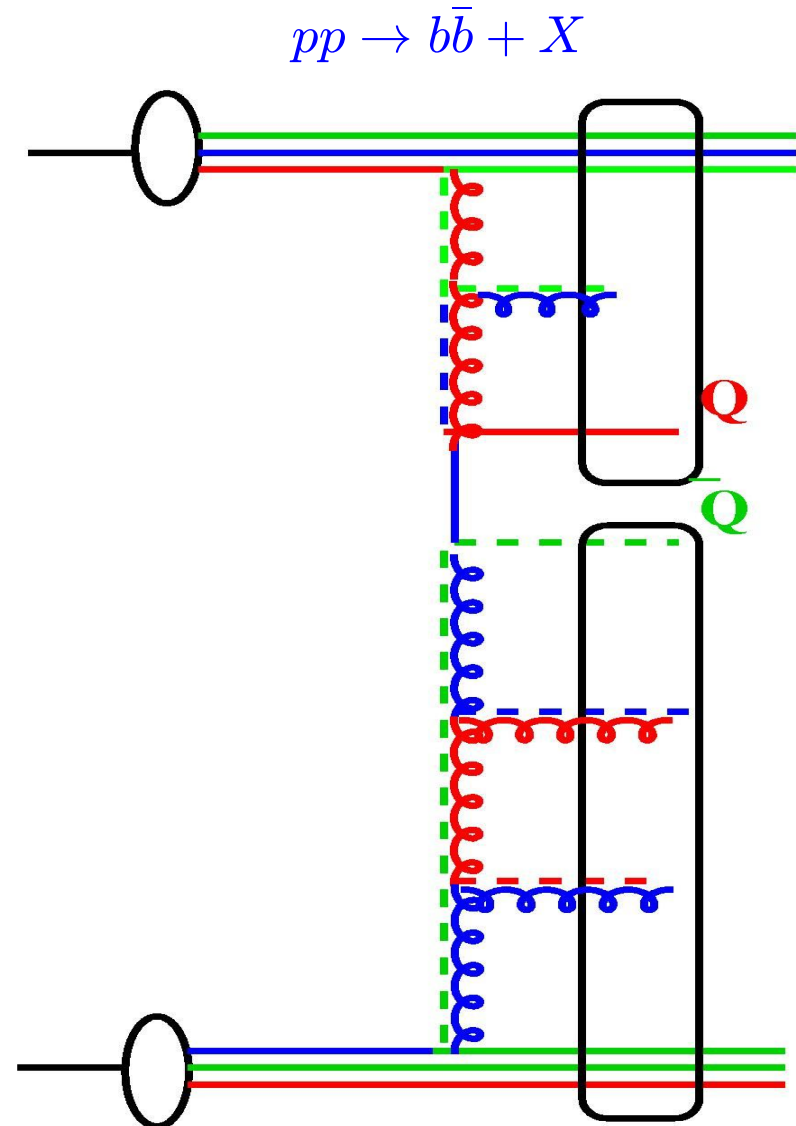
# Gluons in string fragmentation

- process  $e^+e^- \rightarrow q\bar{q}g$
- watch out color flow !!!
- gluons act as kinks on strings
- string effect seen in experiment



# Color Flow in String Fragmentation

- quarks carry color
- anti-quarks carry anticolor
- gluons carry color – anticolor
- connect to color singlet systems
- watch out  $pp$  or  $p\bar{p}$



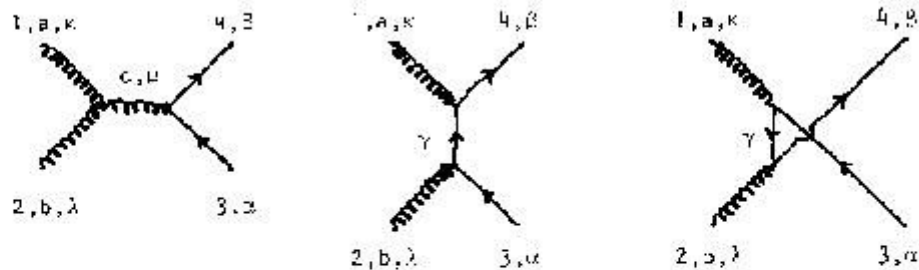


# Color Flow in pp

The Lund Monte Carlo For High P(T) Physics H.U. Bengtsson  
Comput.Phys.Commun.31:323,1984.

Process:  $gg \rightarrow q_i \bar{q}_i$

Diagrams:



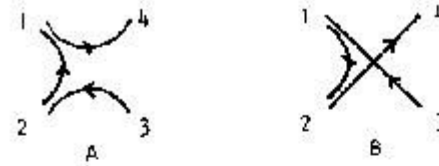
Amplitudes:

$$s: g^2 f^{abc} T_{\alpha\beta}^c \bar{u}_i^\beta(q_4) \frac{\epsilon_1^\nu \epsilon_2^\lambda \gamma^\mu}{s} C_{\alpha\lambda\mu}(q_1, q_2, -q_1 - q_2) v_i^\alpha(q_3)$$

$$t: -ig^2 T_{\alpha\gamma}^b T_{\gamma\beta}^a \bar{u}_i^\beta(q_4) \epsilon_1 \frac{q_1 - q_4}{t} \epsilon_2 v_i^\alpha(q_3)$$

$$u: -ig^2 T_{\alpha\gamma}^a T_{\gamma\beta}^b \bar{u}_i^\beta(q_4) \epsilon_2 \frac{q_1 - q_3}{u} \epsilon_1 v_i^\alpha(q_3)$$

Colour flows:



String configurations:



Colour factors: A:  $T_{\alpha\gamma}^b T_{\gamma\beta}^a$ ; B:  $T_{\alpha\gamma}^a T_{\gamma\beta}^b$

Amplitudes:

$$A: -ig^2 \bar{u}_i^\beta(q_4) \left[ \epsilon_1 \frac{q_1 - q_4}{t} \epsilon_2 - \frac{\epsilon_1^\nu \epsilon_2^\lambda \gamma^\mu}{s} C_{\alpha\lambda\mu}(q_1, q_2, -q_1 - q_2) \right] v_i^\alpha(q_3)$$

$$B: -ig^2 \bar{u}_i^\beta(q_4) \left[ \epsilon_2 \frac{q_1 - q_3}{u} \epsilon_1 + \frac{\epsilon_1^\nu \epsilon_2^\lambda \gamma^\mu}{s} C_{\alpha\lambda\mu}(q_1, q_2, -q_1 - q_2) \right] v_i^\alpha(q_3)$$

Cross-sections:

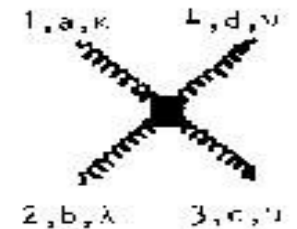
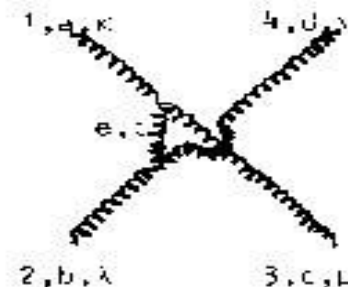
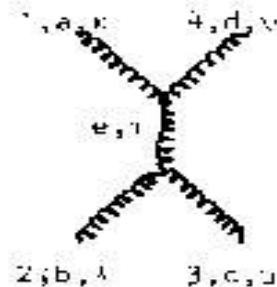
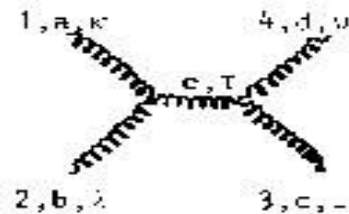
$$A: \frac{\pi\alpha_s^2}{s^2} \frac{1}{6} \left( \frac{\hat{u}}{t} - 2 \frac{\hat{u}^2}{s^2} \right); \quad B: \frac{\pi\alpha_s^2}{s^2} \frac{1}{6} \left( \frac{\hat{t}}{u} - 2 \frac{\hat{t}^2}{s^2} \right)$$

# Color Flow in pp

Process:  $gg \rightarrow gg$

The Lund Monte Carlo For High P(T) Physics H.U. Bengtsson  
Comput.Phys.Commun.31:323,1984.

Diagrams:



Amplitudes:

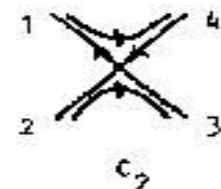
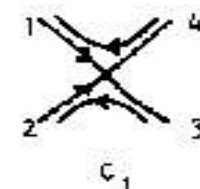
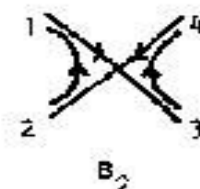
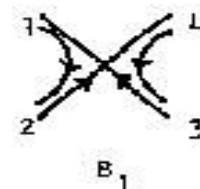
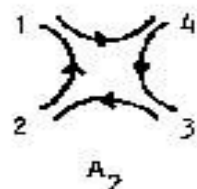
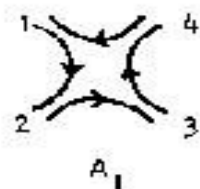
$$s: -ig^2 \frac{1}{s} f^{abc} f^{ode} \epsilon_1^\mu \epsilon_2^\lambda \epsilon_3^\nu \epsilon_4^\rho C_{\kappa\lambda\rho}(-q_1, -q_2, q_1 + q_2) C_{\mu\nu\sigma}(q_3, q_4, -q_3 - q_4)$$

$$t: -ig^2 \frac{1}{t} f^{dae} f^{bhc} \epsilon_1^\mu \epsilon_2^\lambda \epsilon_3^\nu \epsilon_4^\rho C_{\nu\kappa\rho}(q_4, -q_1, q_1 - q_4) C_{\lambda\mu\sigma}(-q_2, q_3, q_2 - q_3)$$

$$u: -ig^2 \frac{1}{u} f^{cae} f^{bhd} \epsilon_1^\mu \epsilon_2^\lambda \epsilon_3^\nu \epsilon_4^\rho C_{\mu\kappa\rho}(q_3, -q_1, q_1 - q_3) C_{\lambda\nu\sigma}(-q_2, q_4, q_2 - q_4)$$

$$4: -ig^2 f^{abc} f^{ode} (\delta_{\kappa\mu} \delta_{\lambda\nu} - \delta_{\kappa\nu} \delta_{\lambda\mu}) + ig^2 f^{dae} f^{bhc} \times (\delta_{\kappa\mu} \delta_{\lambda\nu} - \delta_{\kappa\nu} \delta_{\lambda\mu})$$

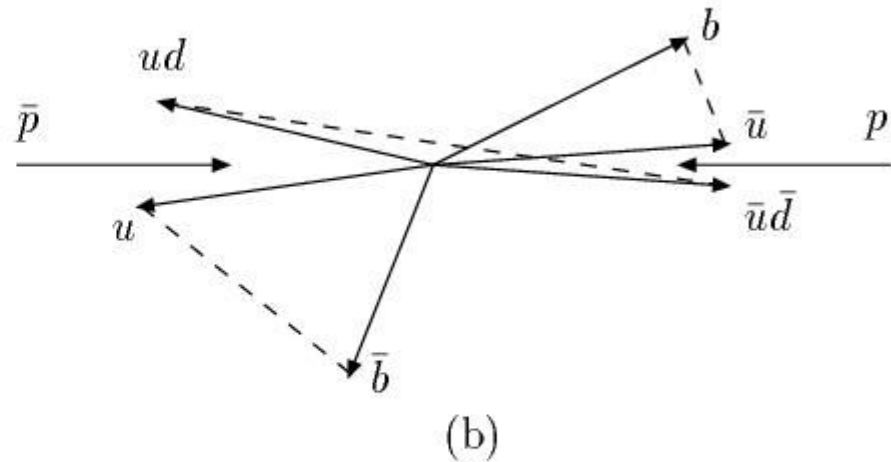
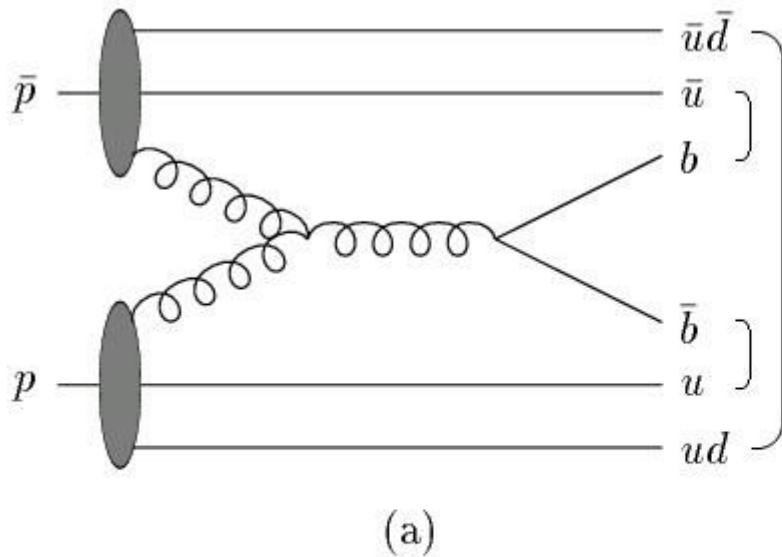
Colour flows:



# Color Flow in $p\bar{p} \rightarrow b\bar{b} + X$

B physics at the Tevatron: Run II and beyond  
E. Norrbin, hep-ph/0201071,p522

$$p\bar{p} \rightarrow b\bar{b} + X$$

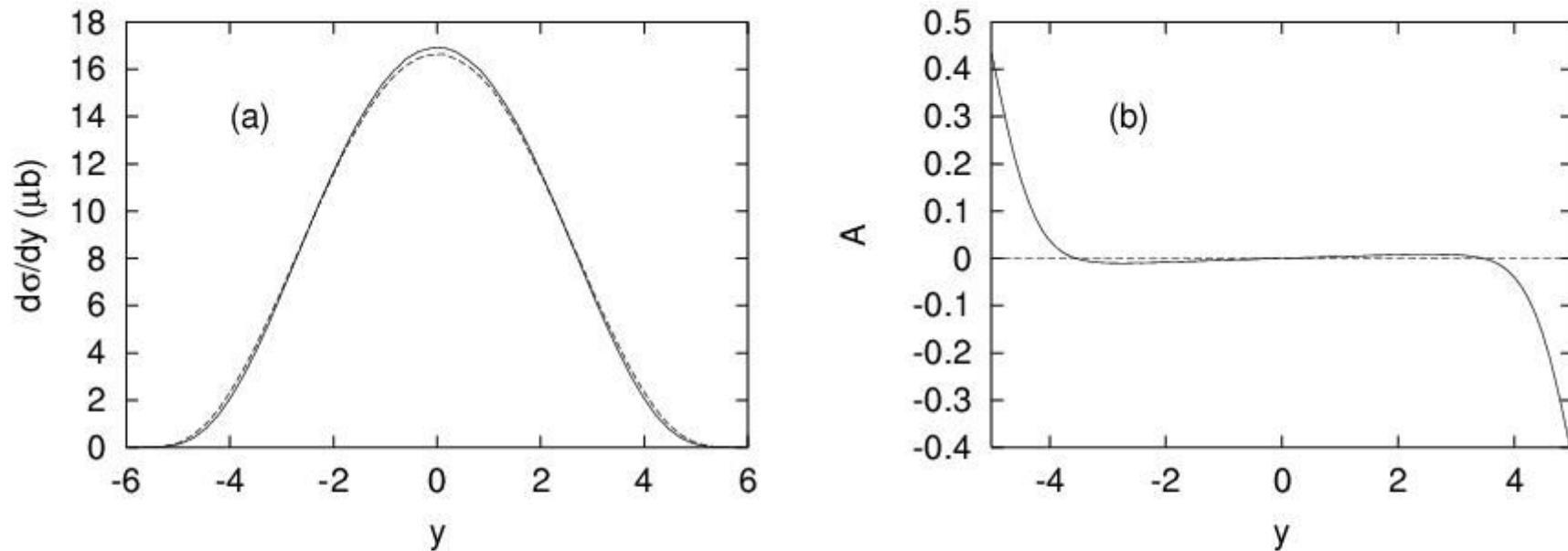


**Figure 9.45:** Example of a string configuration in a  $p\bar{p}$  collision. (a) Graph of the process, with brackets denoting the final color singlet subsystems. (b) Corresponding momentum space picture, with dashed lines denoting the strings.

# Beam - drag effect

- Due to color connection of produced b-quark with beam remnants, the rapidity distribution of b-quarks and B-hadrons is different.
- Asymmetry of  $B^0\bar{B}^0$

B physics at the Tevatron: Run II and beyond  
E. Norrbin, hep-ph/0201071,p525



**Figure 9.47:** Bottom production at the Tevatron. (a) Rapidity distribution of bottom quarks (full) and the B hadrons produced from them (dashed). (b) The asymmetry  $A = \frac{\sigma(B^0) - \sigma(\bar{B}^0)}{\sigma(B^0) + \sigma(\bar{B}^0)}$  as a function of rapidity. For simplicity, only pair production is included.

- **HowTo connect this to factorised fragmentation functions ?**

# Summary

- Fragmentation functions from longitudinal phase space model
  - different behavior for light and heavy quarks
- Hadronization models (iterative procedure)
  - **NEEDED** if more than single inclusive quantities are investigated
  - different models available
    - independent fragmentation (**do not use anymore !!!**)
    - LUND string fragmentation (PYTHIA / JETSET)
    - cluster fragmentation (HERWIG)
      - **dedicated effects observable:**
        - string effect
        - beam drag effect
- **Important to respect all color informations !!!!!**