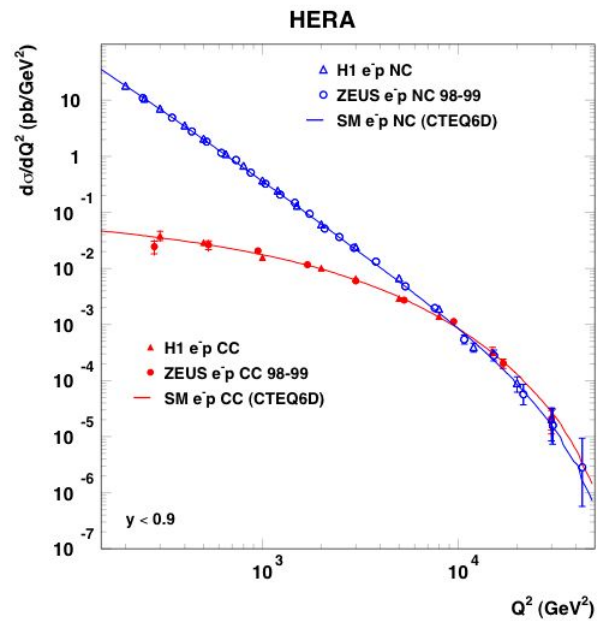


# Electroweak Physics at HERA

Joachim Meyer

DESY

The conference logo really fits this talk .....



# Outline

## Test of electroweak SM at high spacelike momentum transfers

### HERA I :

- Electroweak Unification
- Electroweak DIS cross section fits
- Propagator mass
- NC quark couplings

### HERA II :

- Polarized CC cross sections
- Polarized NC cross sections

### Outlook

# Inclusive Neutral Current (NC) & Charged Current (CC) Cross Sections

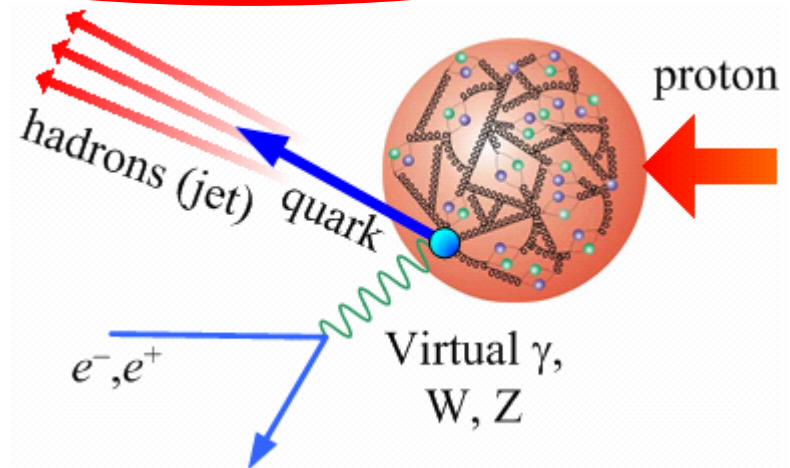
This talk

Structure Functions (SFs)

Electroweak Parameters

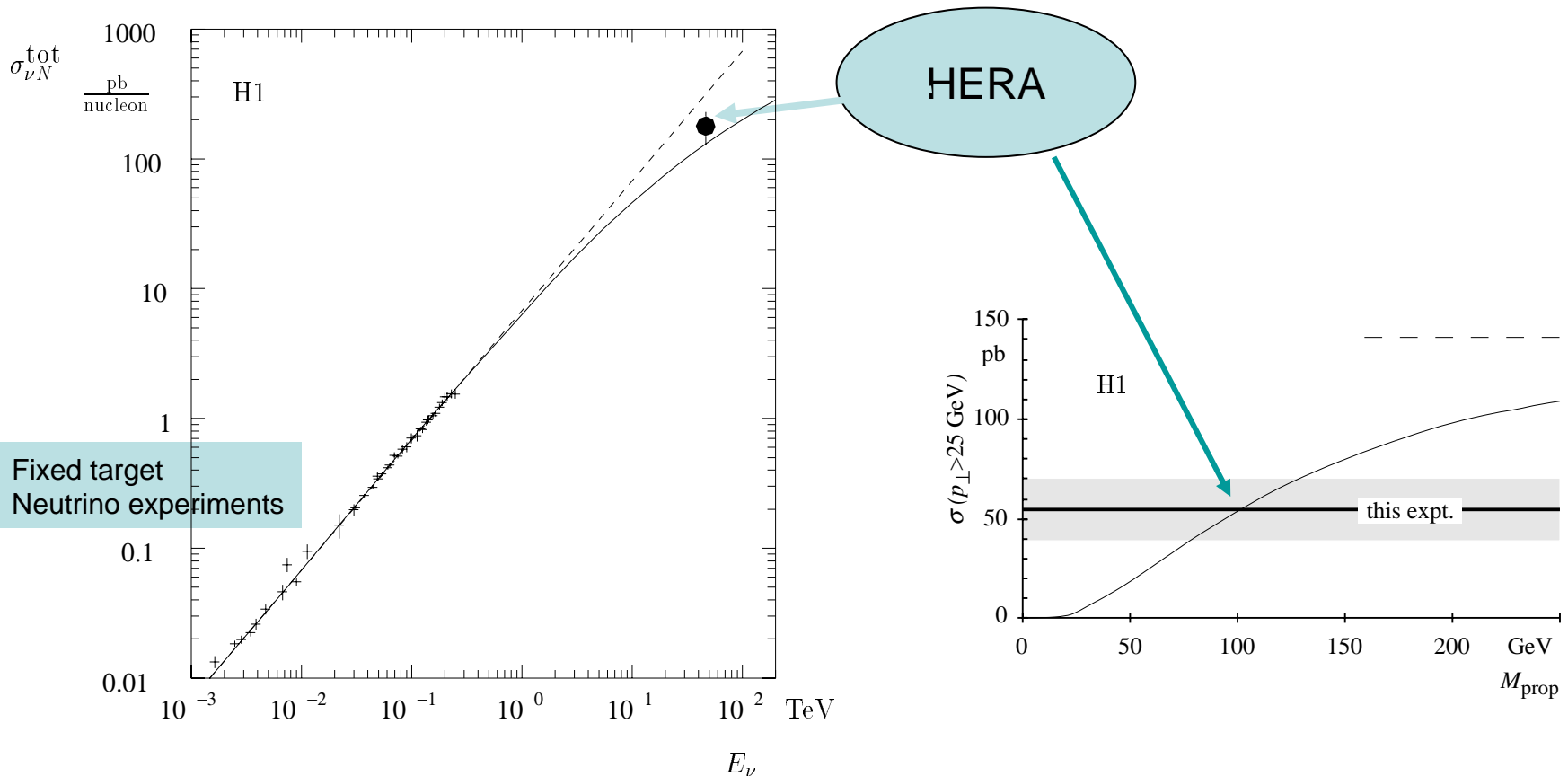
Polarization  $P_e$  (HERA-II only)

Parton Distribution Functions (PDFs)



1993 :Very first electroweak result at HERA from 0.3 pb<sup>-1</sup>

HERA total CC cross section converted to equivalent neutrino cross section



First evidence of W-Propagator effect in Charged Current DIS

# Measured NC and CC cross sections

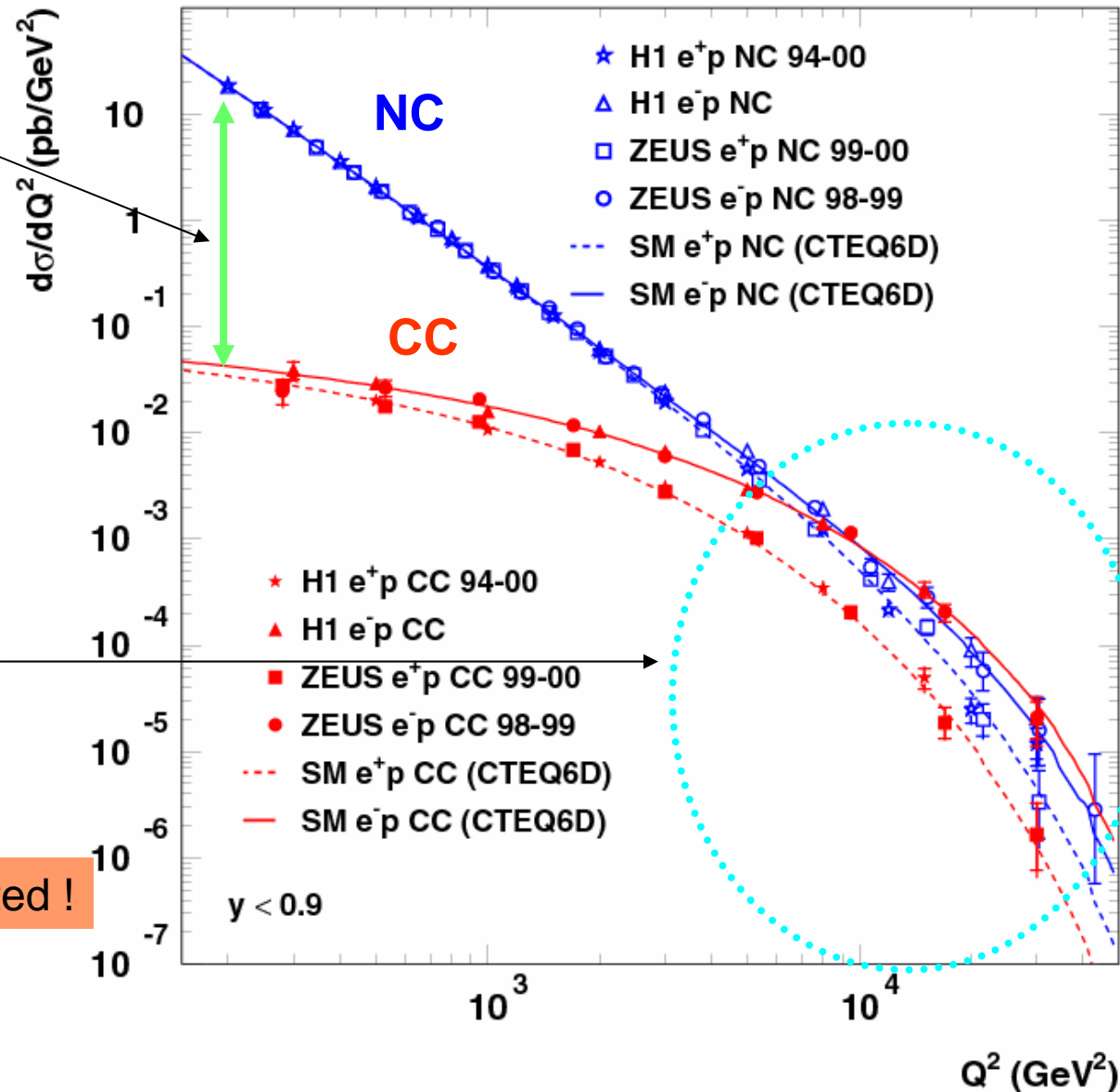
HERA I Data 1994 – 2000 (100 pb<sup>-1</sup> per experiment)

HERA

Suppressed due to large mass of W boson compared to NC DIS

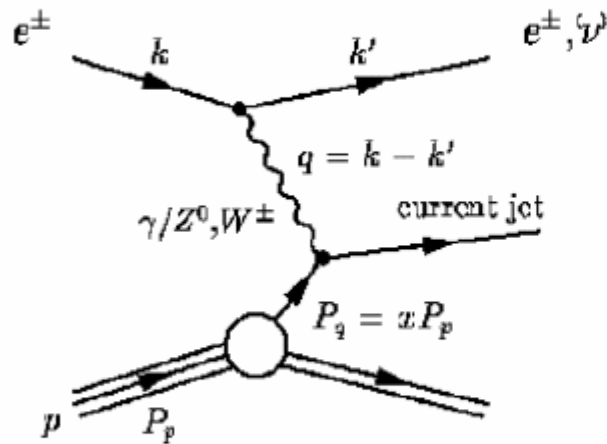
Electro-Weak unification at high Q<sup>2</sup>

High Q<sup>2</sup> results statistics limited !



# Kinematics :

## Deep Inelastic Scattering at HERA



- Neutral Current : exchange of  $\gamma$  or  $Z^0$
- Charged Current : exchange of  $W^\pm$

$$Q^2 = -q^2 = -(k - k')^2$$

Virtuality of exchanged boson  
spatial resolution :  $\lambda \approx \frac{1}{\sqrt{Q^2}}$

$$x = \frac{Q^2}{2p \cdot q} \quad \text{momentum fraction of the struck quark}$$

$$y = \frac{p \cdot q}{p \cdot k} \quad \text{inelasticity}$$

$$s = (p + k)^2 \quad Q^2 = s \cdot x \cdot y$$

Only two independent

# CC Cross Section

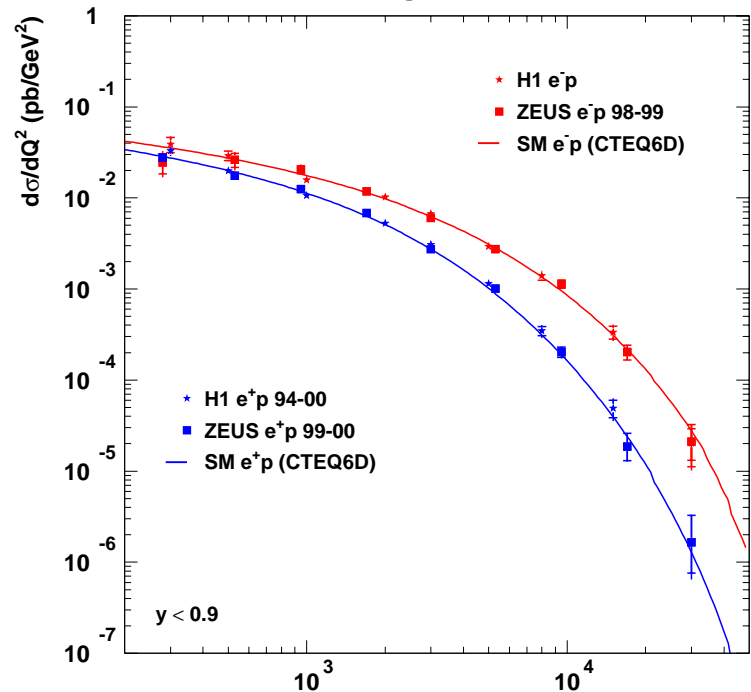
$e^+p$

$$\frac{d^2\sigma^{CC}(e^+p)}{dx dQ^2} = \frac{G_F^2}{4\pi x} \left( \frac{M_W^2}{M_W^2 + Q^2} \right)^2 [\bar{u} + \bar{c} + (1-y)^2(d+s)]$$

$e^-p$

$$\frac{d^2\sigma^{CC}(e^-p)}{dx dQ^2} = \frac{G_F^2}{4\pi x} \left( \frac{M_W^2}{M_W^2 + Q^2} \right)^2 [u + c + (1-y)^2(\bar{d} + \bar{s})]$$

HERA Charged Current





## NC Cross Section

$$\frac{d^2 \sigma^{\text{NC}}(e^\pm p)}{dx dQ^2} = \frac{2\pi \alpha^2}{xQ^4} [Y_+ F_2^{\text{NC}} \mp Y_- xF_3^{\text{NC}} - y^2 F_L^{\text{NC}}] \quad Y_\pm = 1 \pm (1-y)^2$$

Dominant contribution

Contribution only important at high  $Q^2$

Sizeable only at high  $y$

NC structure functions,  $F_2^{\text{NC}}$  and  $xF_3^{\text{NC}}$ , can be decomposed as

$\gamma$  exchange

$\gamma$ -Z interference

Z exchange

$$\begin{aligned} F_2^{\text{NC}} &= F_2^\gamma - v_e K_Z F_2^{\gamma Z} + (v_e^2 + a_e^2) K_Z^2 F_2^Z \\ xF_3^{\text{NC}} &= -a_e K_Z xF_3^{\gamma Z} + 2v_e a_e K_Z^2 xF_3^Z \end{aligned}$$

$$K_Z = \frac{1}{4 \sin^2 \theta_w \cos^2 \theta_w} \frac{Q^2}{Q^2 + M_Z^2}$$

$$[F_2, F_2^{\gamma Z}, F_2^Z] = x \sum_q [e_q^2, 2e_q v_q, v_q^2 + a_q^2](q + \bar{q})$$

$$[xF_3^{\gamma Z}, xF_3^Z] = 2x \sum_q [e_q a_q, v_q a_q](q - \bar{q})$$

Experiment measures **Cross-Sections** and extract **SFs**

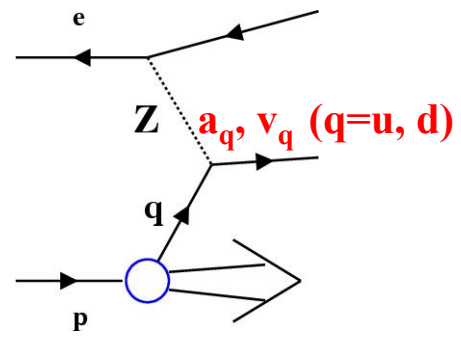
SFs : coupling constant  $\otimes$  Parton Distribution Functions (PDFs)

# HERA-I Results : Combined EW+PDF Fit (H1)

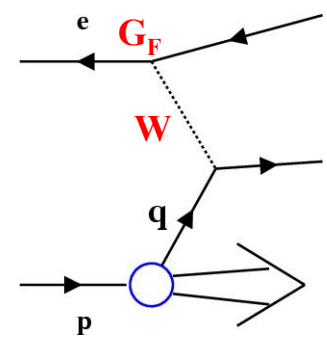
- The low  $Q^2$  precision cross section data +  
the high  $Q^2$  NC  $e^+p$  &  $e^-p$  data +  
the high  $Q^2$  CC  $e^+p$  &  $e^-p$  data } constrain 5 sets of PDFs:

→ gluon, up-type quark, down-type quark & their anti-quarks

- NC data at high  $Q^2$  also sensitive to quark couplings to the Z boson



- CC data also sensitive to
  - {  $G, W$  propagator mass [model independent]
  - {  $M_W, m_t$  [within the SM framework]



$$\frac{d^2\sigma_{CC}(e^\pm p)}{dx dQ^2} = \frac{G_F^2 M_W^4}{2\pi x (Q^2 + M_W^2)^2} \Phi(\text{PDFs})$$

# Combined EW+PDF Analysis Strategies

## • Model independent fits:

1) Fit  $a_u-v_u-a_d-v_d$ -PDF

to extract light quark couplings to the Z boson

2) Fit  $G-M_{\text{prop}}$ -PDF

to determine the normalization factor  $G$  and W propagator mass  $M_{\text{prop}}$

3) Fit  $M_{\text{prop}}$ -PDF [fix  $G$  to  $G_F$ ]

to determine the W propagator mass  $M_{\text{prop}}$

## • Fits within the SM framework:

$$G_F^2 = \frac{\pi\alpha}{\sqrt{2} M_W^2 \left(1 - \frac{M_W^2}{M_Z^2}\right)} \frac{1}{1-\Delta r}$$

$\Delta r$  contains - quadratic dependence on  $m_t$   
- logarithmic dependence on  $M_H$

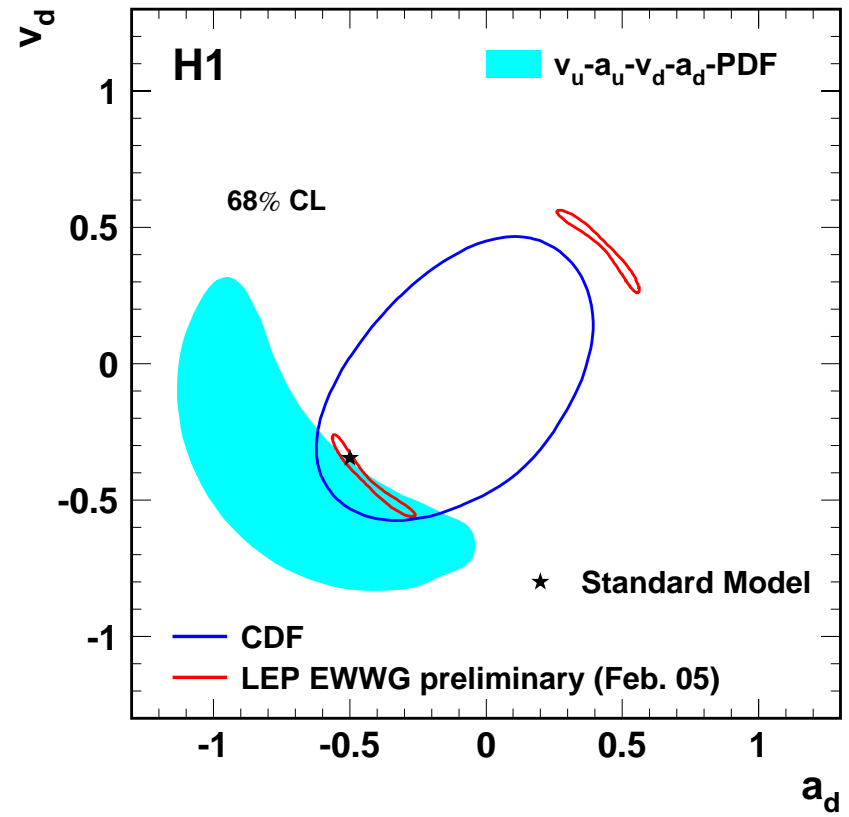
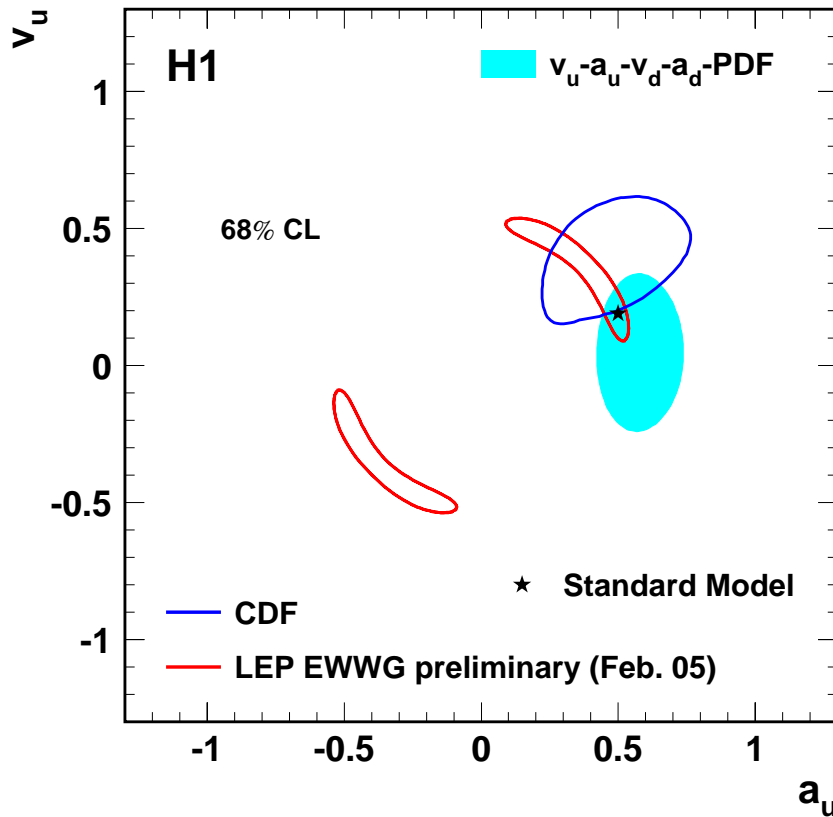
4) Fit  $M_W$ -PDF [fix  $m_t$  to 178GeV,  $M_H$  to 120GeV]

to determine the SM W mass  $M_W$

5) Fit  $m_t$ -PDF [fix  $M_W$  to 80.425GeV,  $M_H$  to 120GeV]

to determine the top quark mass  $m_t$

# First Results on Light Quark Couplings to Z at HERA

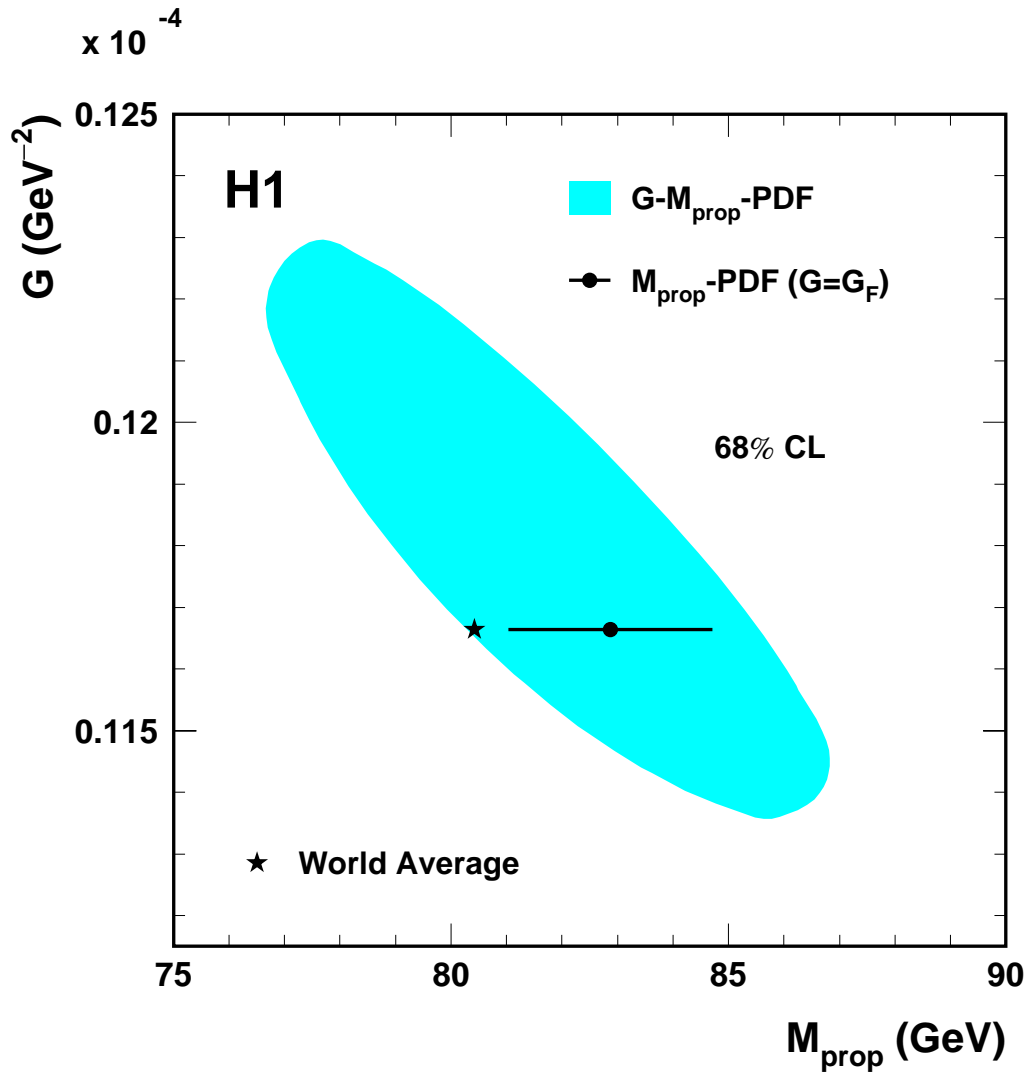


SM :  $v_q = I_q^3 - 2e_q \sin^2 \theta_w$        $a_q = I_q^3$

**Tevatron:**  $qq \rightarrow ee$  Drell-Yan,  $A_{FB}$ : CDF Collab., Phys. Rev. D71(2005)052002, hep-ex/0411059

**LEP:**  $ee \rightarrow qq(\gamma)$  [ $a_q^2 + v_q^2$ ]: <http://lepewwg.web.cern.ch/LEPEWWG/plots/winter2005/>

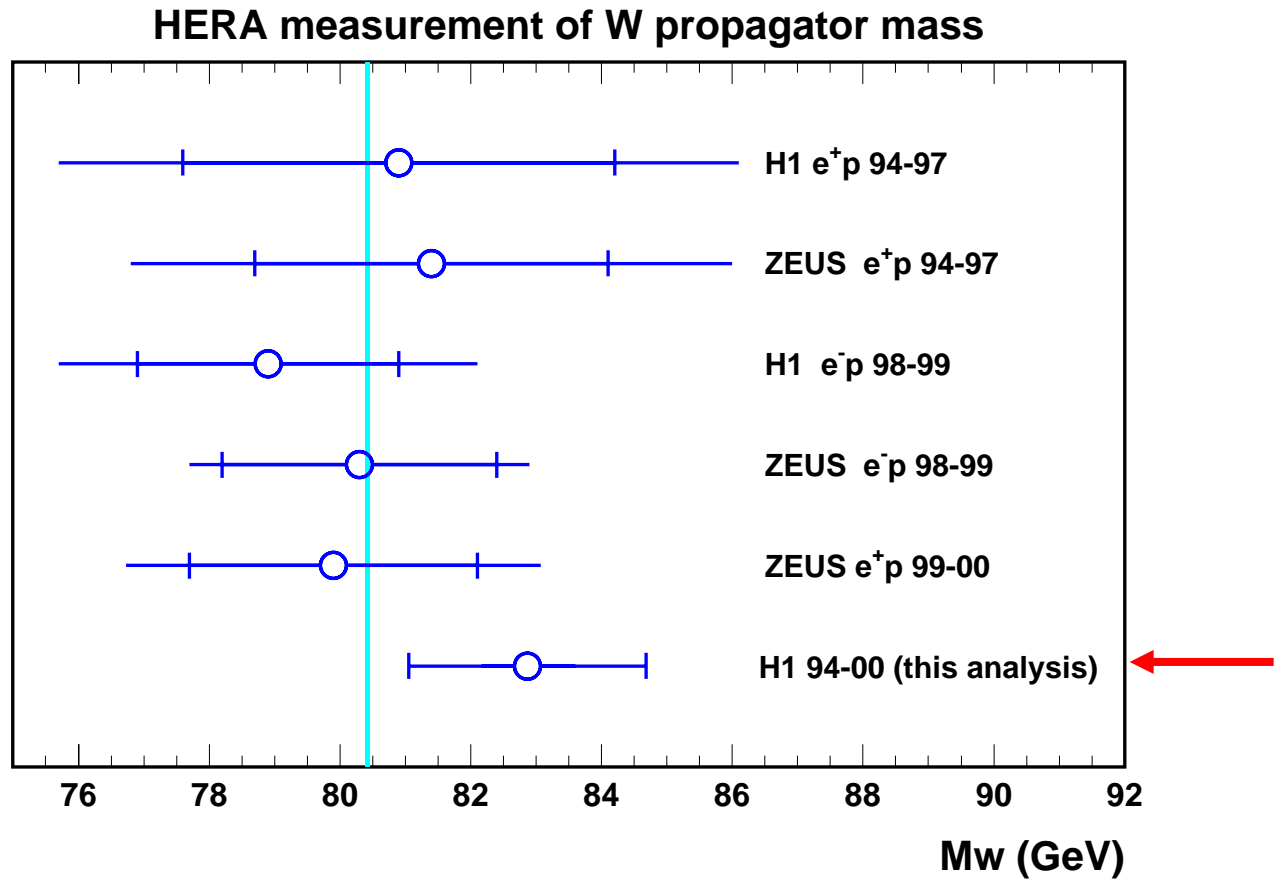
# W Propagator Mass and Coupling



With  $G_F$  from PDG :

$$M_{\text{Prop}} = 82.9 \pm 1.9 \text{ GeV}$$

# Improved Precision on W Propagator Mass



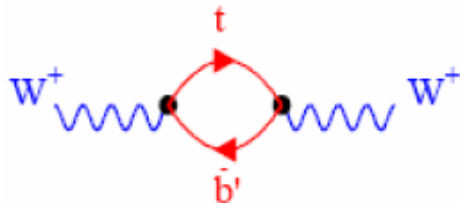
# Fits Imposing the SM Constraints

$$\frac{d^2 \sigma_{cc}^\pm}{dx dQ^2} = \frac{G^2}{2\pi} \cdot \left( \frac{M_W^2}{Q^2 + M_W^2} \right)^2 \cdot \Phi^\pm(pdfs)$$

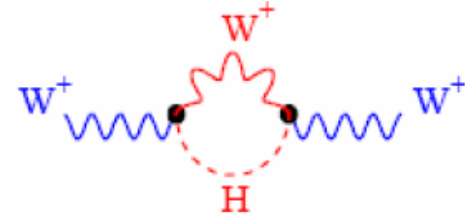
Introducing SM  $G_F$ - $M_W$  relation  
in On-Mass-Shell (OMS) scheme

$$\frac{d^2 \sigma_{cc}^\pm}{dx dQ^2} = \frac{\pi \alpha^2}{4M_W^4 \left(1 - \frac{M_W^2}{M_Z^2}\right)^2} \cdot \frac{1}{(1 - \Delta r)^2} \cdot \left( \frac{M_W^2}{Q^2 + M_W^2} \right)^2 \cdot \Phi^\pm(pdfs)$$

Quadratic dependence on  $m_t$



Logarithmic dependence on  $M_H$



# Determination of SM Parameters

## W mass value:

$$M_W = 80.786 \pm 0.205_{\text{exp}} \text{ GeV}$$

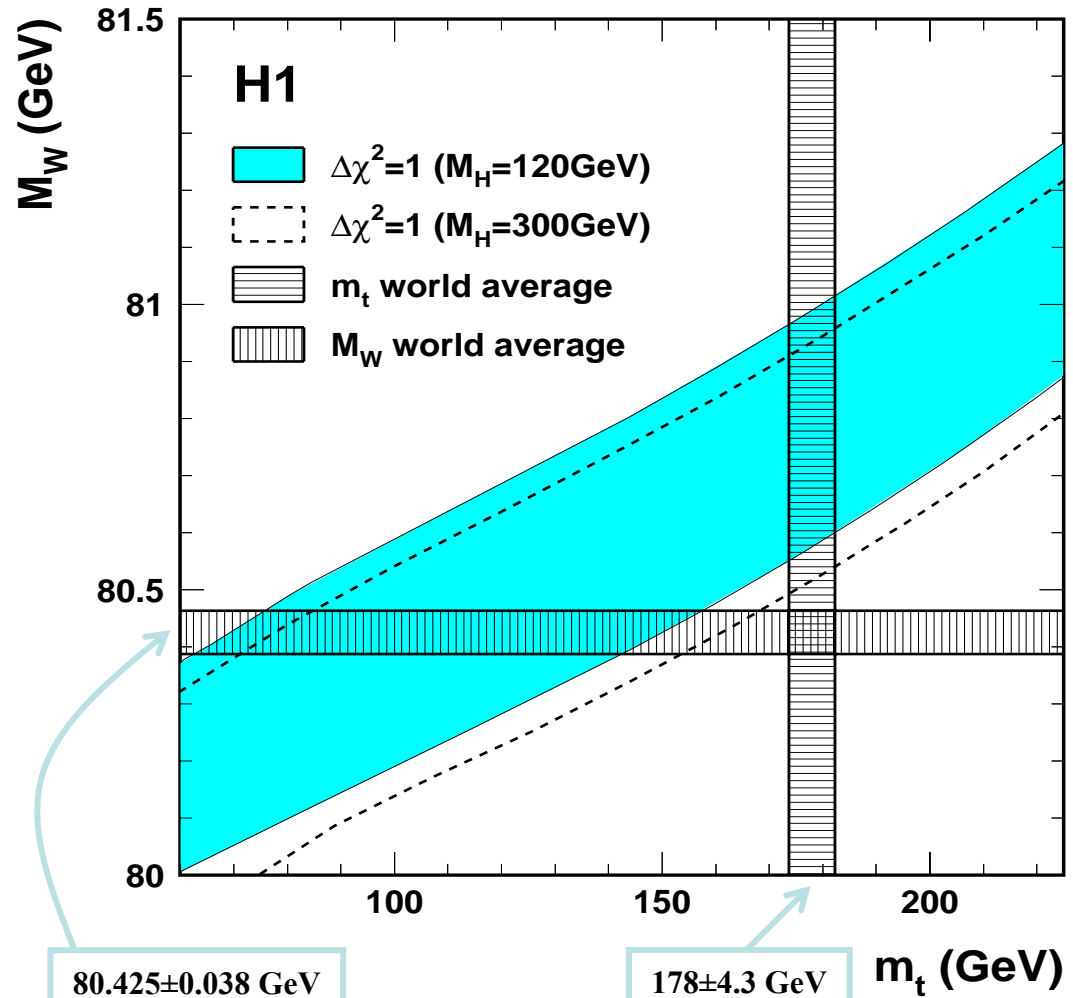
+ the world average  $M_Z$   
 → indirect determination of

$$\sin^2\theta_W = 1 - M_W^2/M_Z^2 = 0.2151 \pm 0.0040_{\text{exp}}$$

Using  $M_W$  (PDG) restricts top quark mass

$$m_t = 104 \pm 44_{\text{exp}} \text{ GeV}$$

→ First determination in DIS at EW scale



Ringberg 2005

Joachim Meyer

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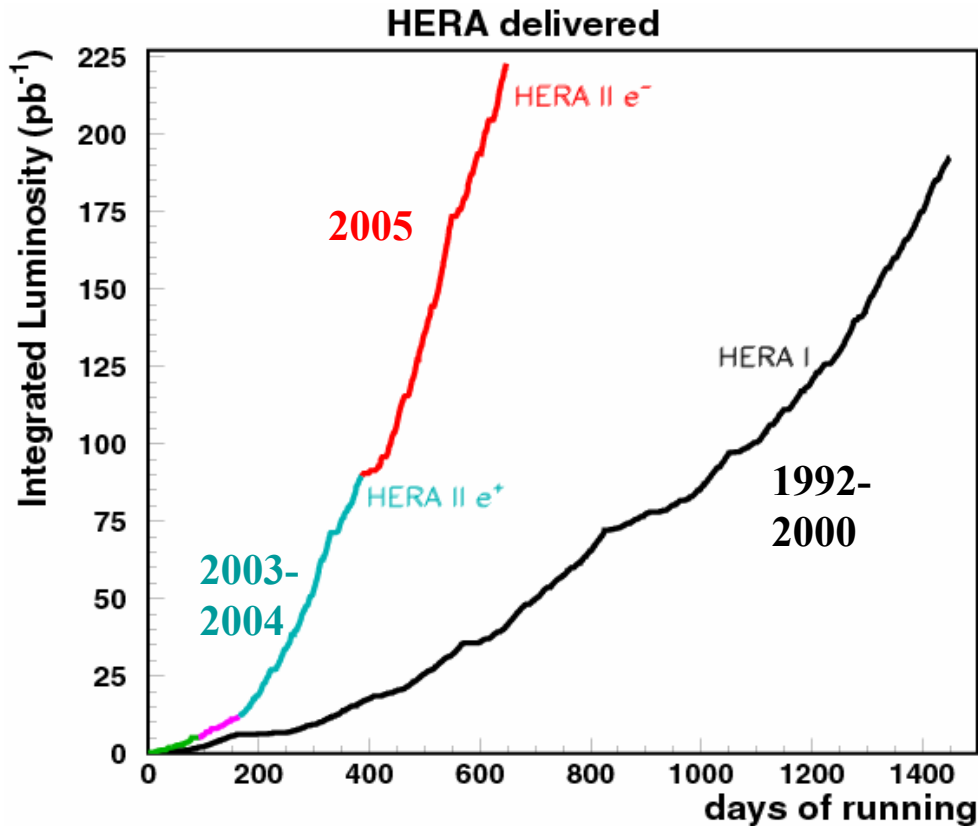


# HERA II

- Substantial increase in luminosity
- Longitudinally polarized lepton beams

# Luminosity at HERA-II

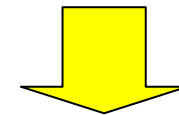
High Luminosity → sensitivity in High- $Q^2$  region



Luminosity used for physics analysis per experiment :

HERA-I :  $100\text{pb}^{-1}(\text{e}^+\text{p})$ ,  $20\text{pb}^{-1}(\text{e}^-\text{p})$

HERA-II :  $40\text{pb}^{-1}(\text{e}^+\text{p})$ ,  $100\text{pb}^{-1}(\text{e}^-\text{p})$



By the end of the HERA-II in July 2007,

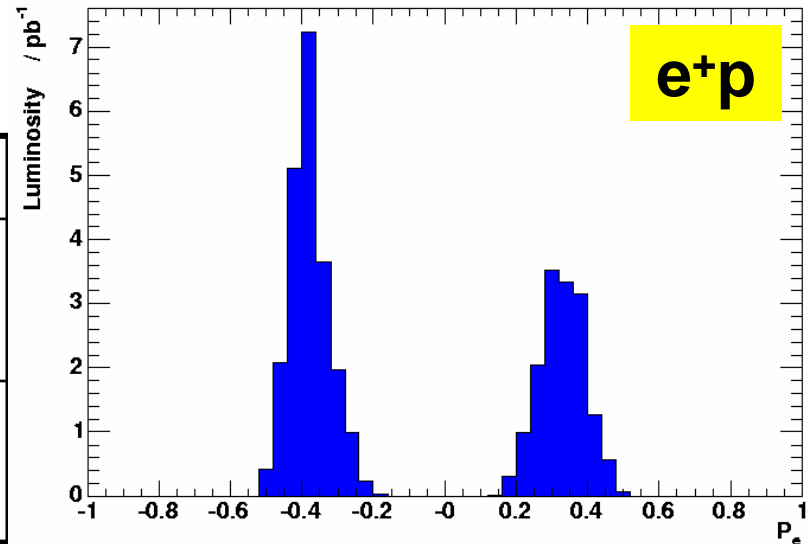
expect  $\sim 700\text{pb}^{-1}$

per experiment

# Data samples

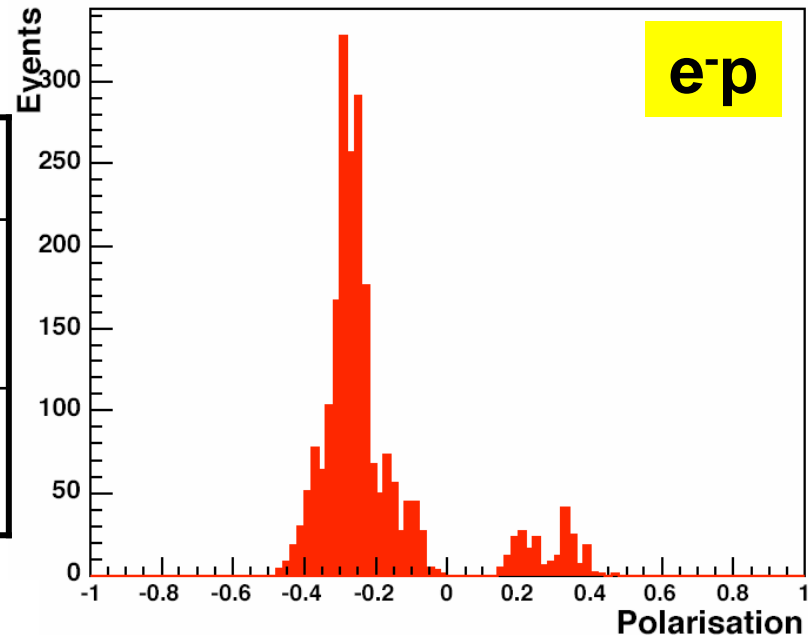
## H1 data samples

	P < 0 (LH)	P > 0 (RH)
e <sup>+</sup> p data	L = 21.7 pb <sup>-1</sup> P = - 40.2 %	L = 15.3 pb <sup>-1</sup> P = + 33.0 %
e <sup>-</sup> p data	L = 17.8 pb <sup>-1</sup> P = - 25.4 %	



## ZEUS data samples

	P < 0 (LH)	P > 0 (RH)
e <sup>+</sup> p data	L = 16.4 pb <sup>-1</sup> P = - 40.2 %	L = 14.1 pb <sup>-1</sup> P = + 31.8 %
e <sup>-</sup> p data	L = 35.3 pb <sup>-1</sup> P = - 25.9 %	L = 6.5 pb <sup>-1</sup> P = + 29.2 %



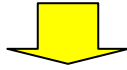
# Charged current physics at HERA II

# CC Total Cross-Section (H1)

$Q^2 > 400 \text{ GeV}^2, y < 0.9$

Remind : CC is pure weak

$$\sigma_{\text{CC}}(P_{e^\pm}) = (1 \pm P_{e^\pm}) \sigma_{\text{CC}}(P_{e^\pm} = 0)$$



Direct observation of chiral structure of weak interaction

- A clear linear dependence is observed both  $e^+$  and  $e^-$
- Data are in agreement with the SM prediction

$e^+p$

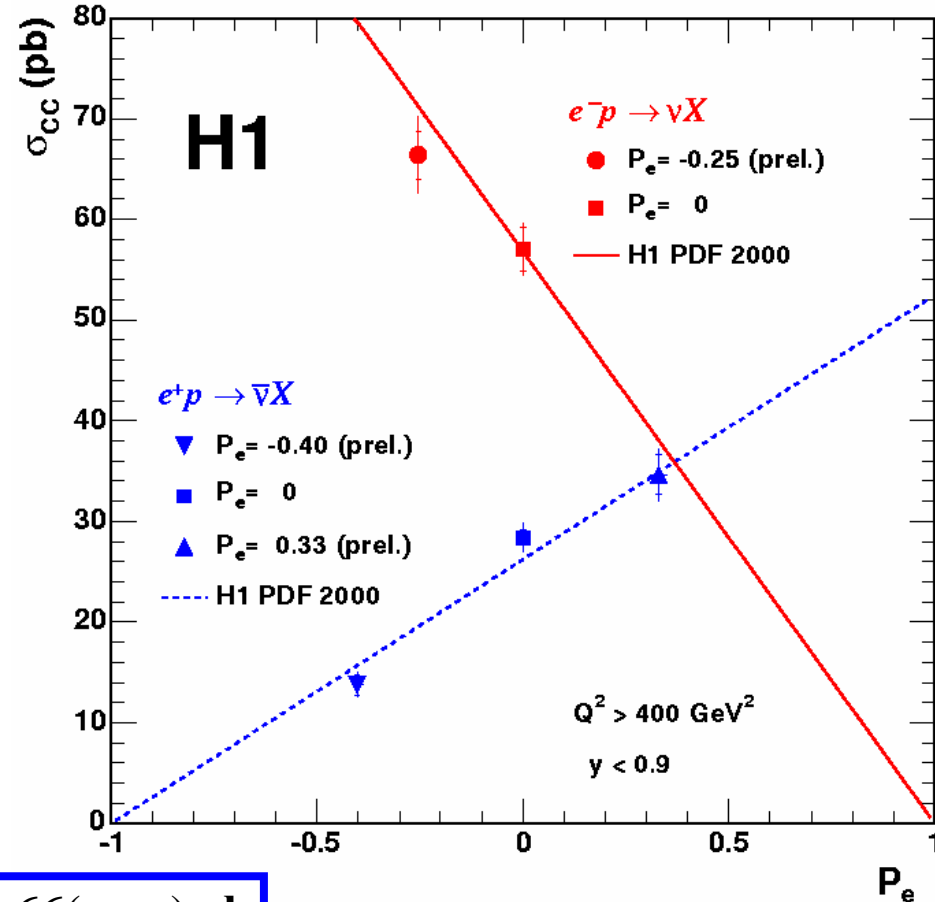
$$\sigma_{\text{CC}}(P_e = +33\%) = 34.67 \pm 1.94(\text{stat.}) \pm 1.66(\text{syst.}) \text{ pb}$$

$$\sigma_{\text{CC}}(P_e = -40\%) = 13.80 \pm 1.04(\text{stat.}) \pm 0.94(\text{syst.}) \text{ pb}$$

$e^-p$

$$\sigma_{\text{CC}}(P_e = -25\%) = 66.42 \pm 2.39(\text{stat.}) \pm 2.99(\text{syst.}) \text{ pb}$$

CC cross sections



# CC Total Cross-Section : H1 and ZEUS

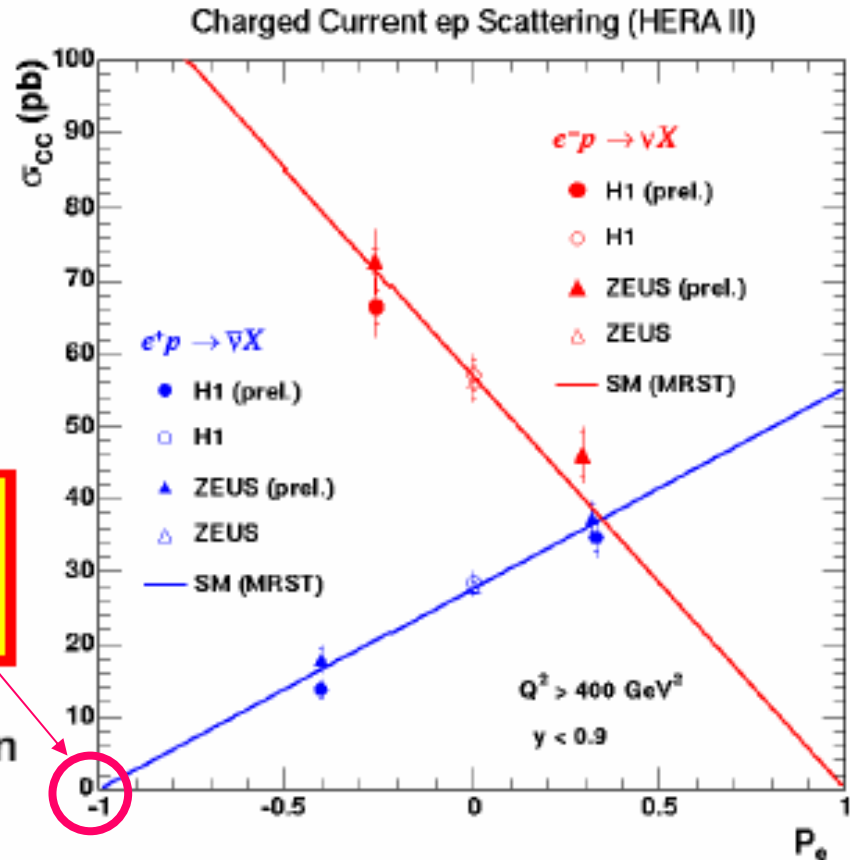
$Q^2 > 400 \text{ GeV}^2, y < 0.9$

Right Handed CC cross section is extrapolated by linear fit to H1+ZEUS  $e^+p$  data



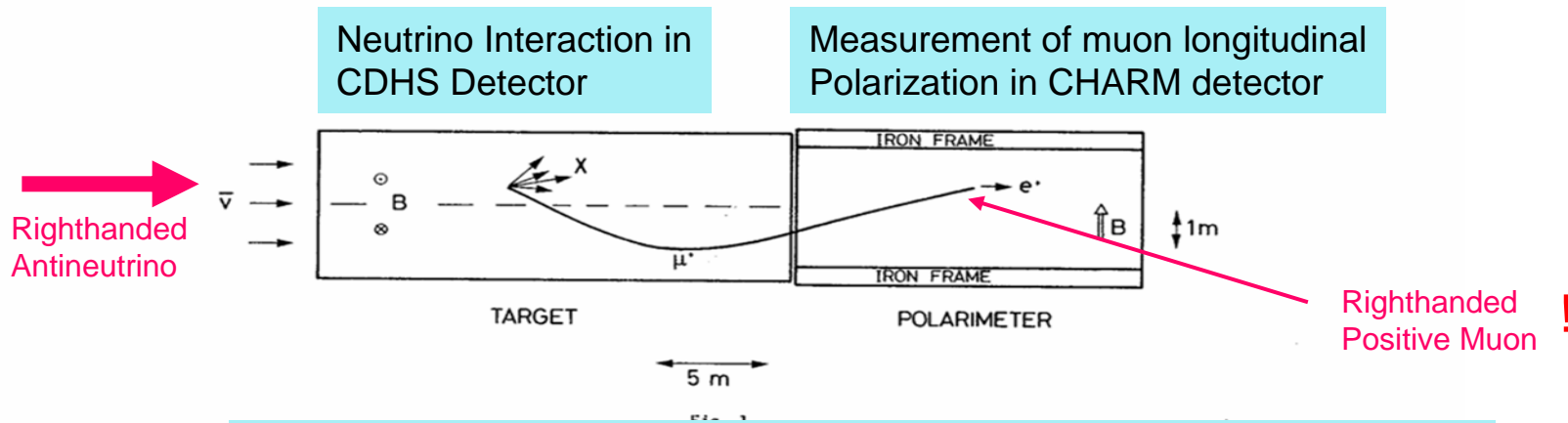
$$\sigma_{e^+p \rightarrow \bar{\nu}X}(P_{e^+} = -100\%) = 0.2 \pm 1.8(\text{stat.}) \pm 1.6(\text{syst.}) \text{ pb}$$

Consistent with the SM prediction of:  $\sigma_{CC}(RH) = 0$



# History:

The **polarization dependence of CC – DIS scattering** has already been measured in 1979 at the CERN Neutrino beam. (Phys.Lett. B86 ; 222 (1979))



Result :

The longitudinal Polarisation of the positive muons is

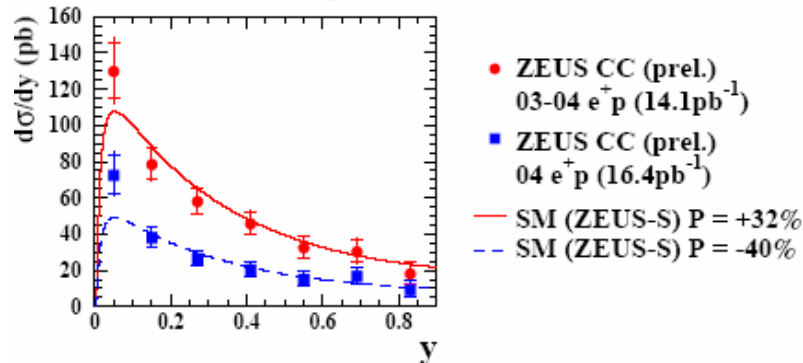
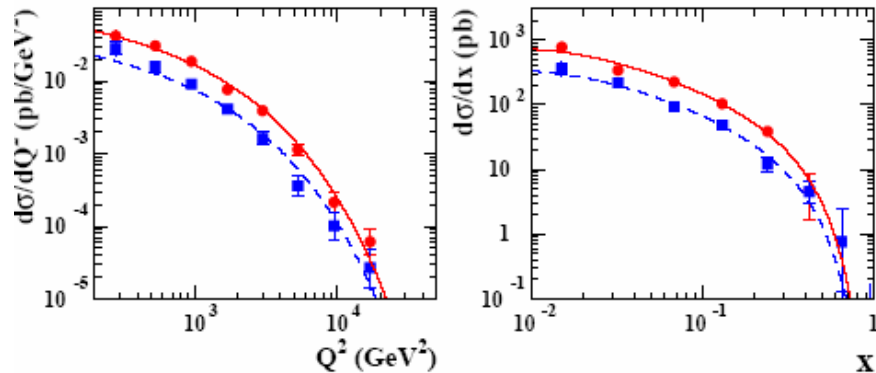
$$P = + 1.09 \pm 0.22$$

(at an average momentum transfer  $3.2 \text{ GeV}^2$  !!!)

# CC Differential Cross Sections (ZEUS)

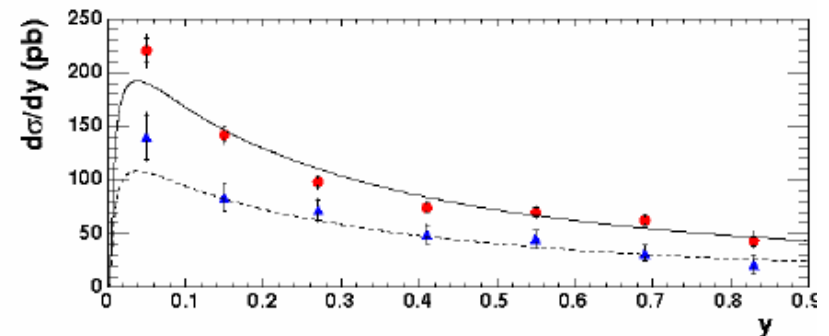
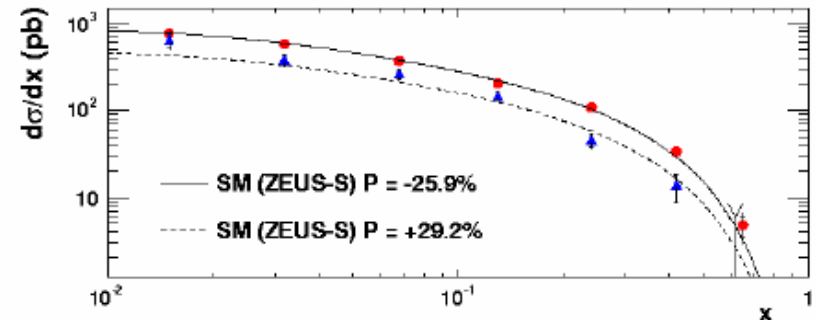
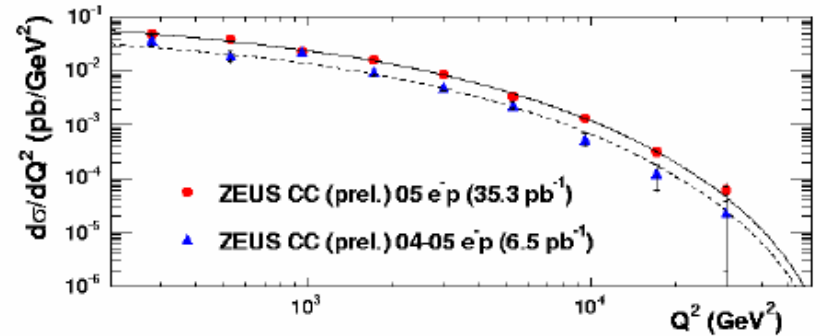
$e^+$

ZEUS



$e^-$

ZEUS



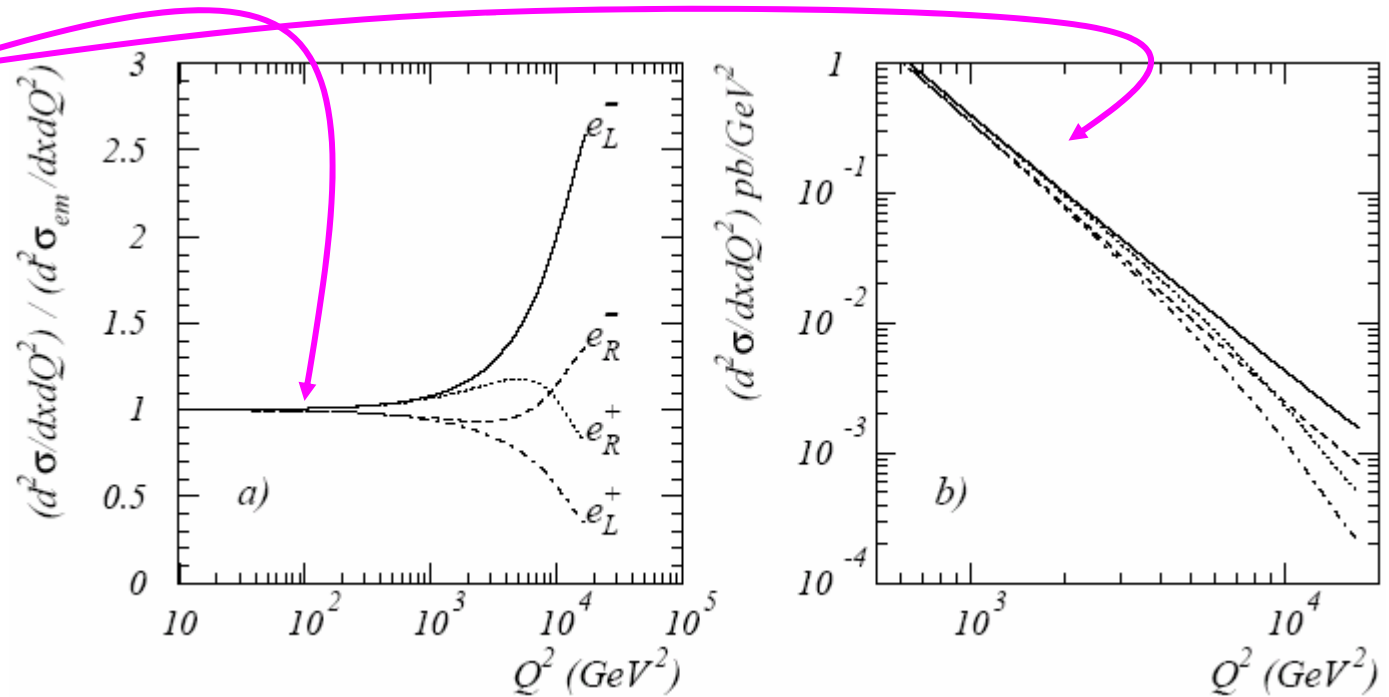
polarisation dependence seen not only total cross section but differential cross section in all kinematic regions



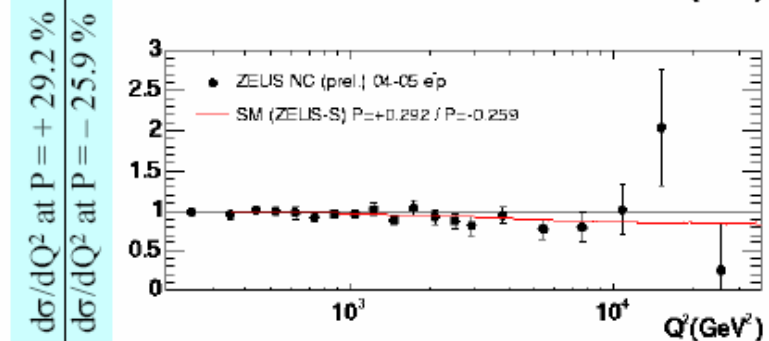
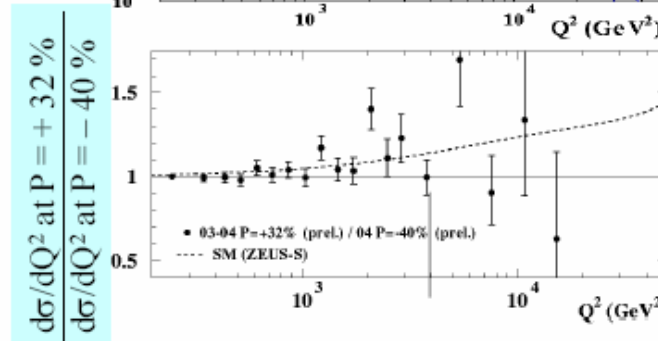
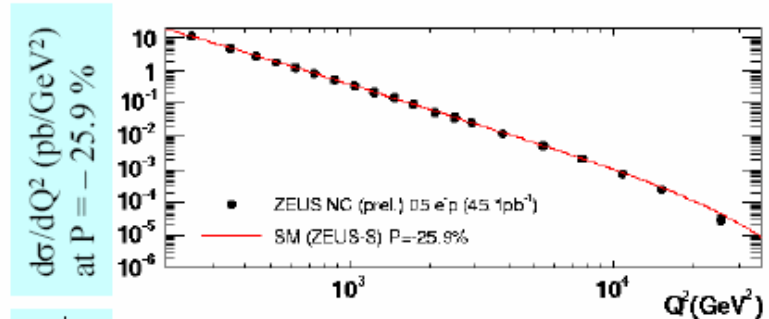
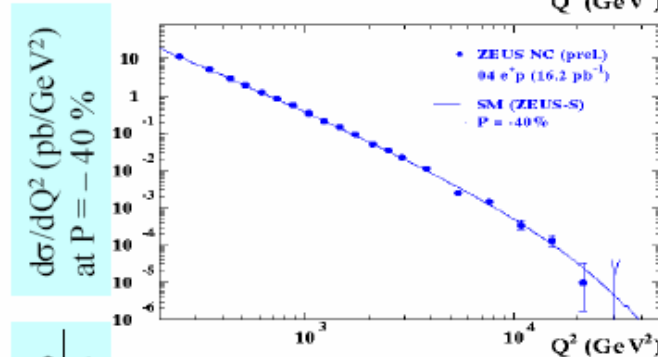
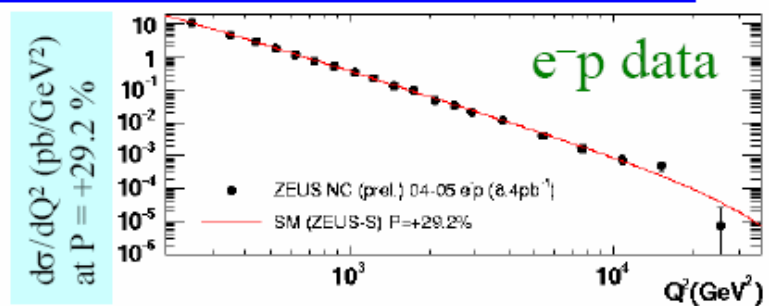
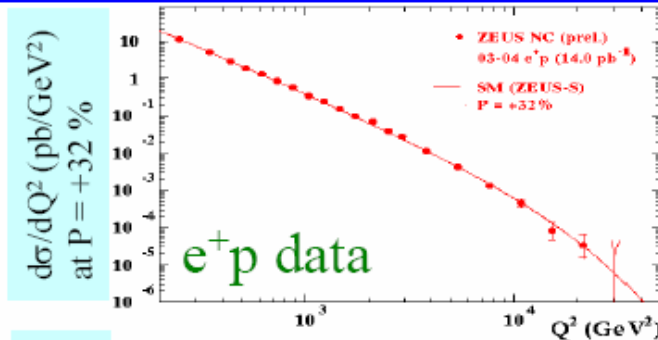
# Neutral current physics At HERA II

# NC & CC Cross Sections Dependence on Polarization $P_e$

For NC: **em** contribution dominating at low  $Q^2$  is independent of  $P_e$ .  
weak NC only significant at high  $Q^2$



# NC Differential Cross Section (ZEUS)



- $\frac{d\sigma}{dQ^2}$  for NC : Polarisation dependence is not observed conclusively with the current limited statistics

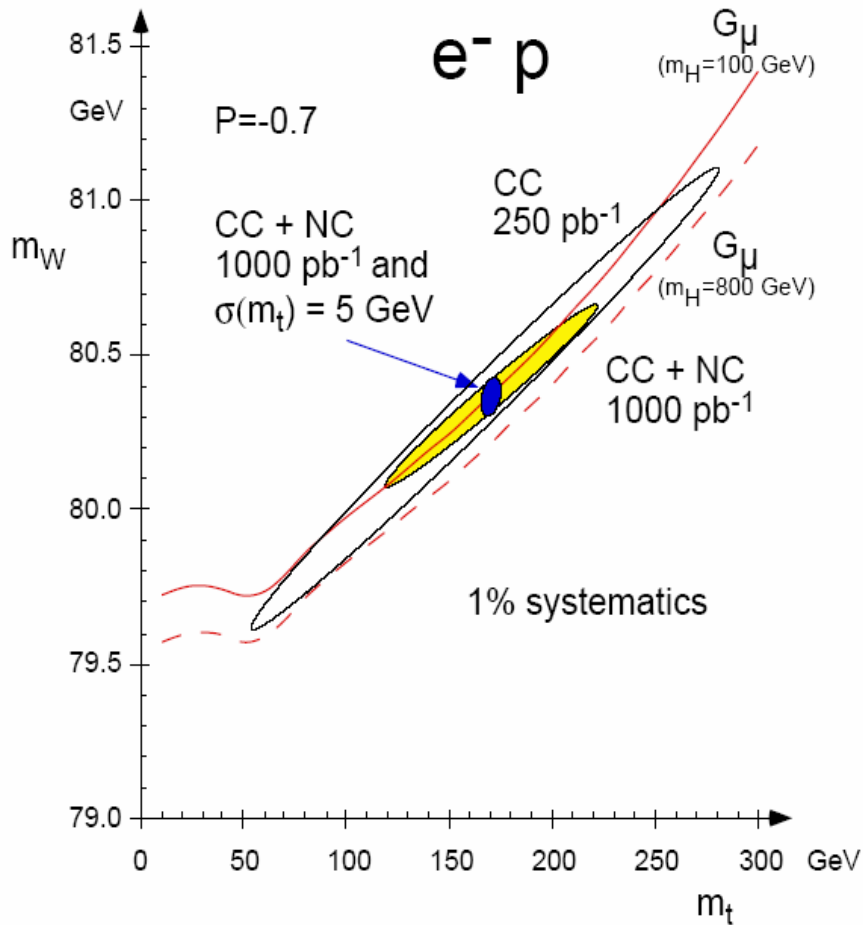
# Outlook

HERA II is expected to deliver approx  $700 \text{ pb}^{-1}$  to each experiment. There will be an equal share of  $e^+$  and  $e^-$ , left- and righthanded polarized.

More precise checks of electroweak SM possible :

- Better  $M_W$  determination (spacelike Propagatormass)
- More precise  $Z_0$  couplings to light quarks

# HERA Physics workshop studies : W - Propagator Mass



## H1 Result :

HERA I (100 pb<sup>-1</sup> , unpol.) :

$$M_{\text{prop}} = 82.9 \pm 1.9 \text{ GeV}$$

$$M_W = 80.8 \pm 0.2 \text{ GeV}$$

Precise check of EW theory when combined with  $M_t$  from Tevatron (LHC)

## Polarised NC DIS cross section

$$\frac{d^2\sigma^{\text{NC}}(e^\pm p)}{dx dQ^2} = \frac{2\pi a^2}{xQ^4} [H_0^\pm + P_e H_P^\pm]$$

Unpolarised contribution

Polarised contribution : only includes  $\gamma$ -Z and Z terms

$$F_2^{\text{NC}} = F_2^\gamma - (v_e - \underline{P_e a_e}) K_Z F_2^{\gamma Z} + (v_e^2 + a_e^2 - 2\underline{P_e v_e a_e}) K_Z^2 F_2^Z$$

$$xF_3^{\text{NC}} = -(a_e - \underline{P_e v_e}) K_Z xF_3^{\gamma Z} + [2v_e a_e - \underline{P_e}(v_e^2 + a_e^2)] K_Z^2 xF_3^Z$$

$$K_Z = \frac{1}{4 \sin^2\theta_W \cos^2\theta_W} \frac{Q^2}{Q^2 + M_Z^2}$$

$$[F_2, F_2^{\gamma Z}, F_2^Z] = x \sum_q [e_q^2, \underline{2e_q v_q}, \underline{v_q^2 + a_q^2}] (q + \bar{q})$$

$$[xF_3^{\gamma Z}, xF_3^Z] = 2x \sum_q [\underline{e_q a_q}, \underline{v_q a_q}] (q - \bar{q})$$



Polarised  $e^\pm$  beam helps to constrain  $v_q$

## Polarisation physics

- Quark couplings can be accurately measured, e.g. light quark couplings by looking at differences between  $\sigma(L,R)$ .
- Great improvement over unpolarised case.
- $e_{L,R}^{\pm}, P = \pm 70\%$   
250 pb<sup>-1</sup> per beam

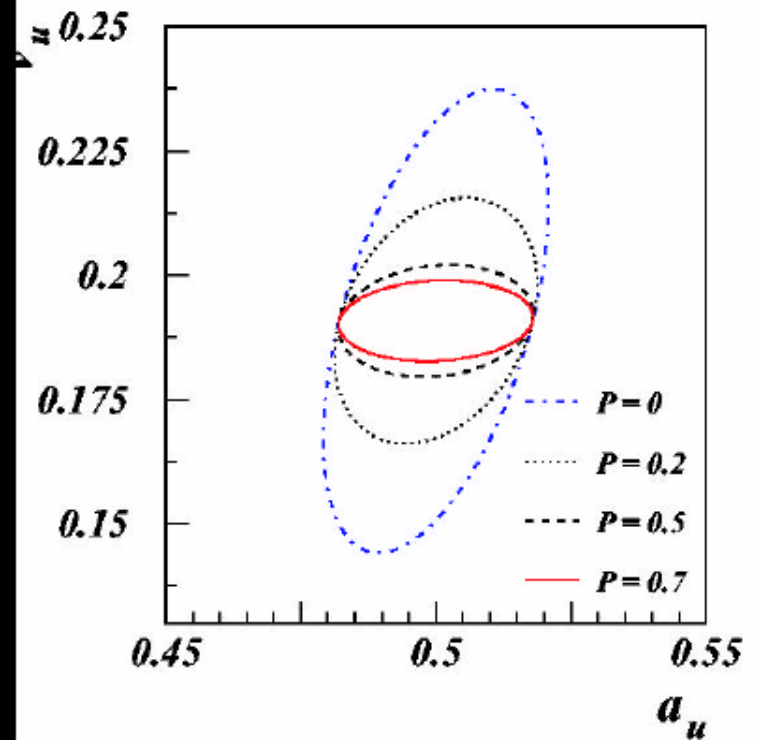
### H1 RESULT :

Hera I (unpol):

$$a_u = 0.56 \pm 0.10$$

$$v_u = 0.04 \pm 0.19$$

	$v$	$a$
$u$	13%	6%
$d$	17%	17%

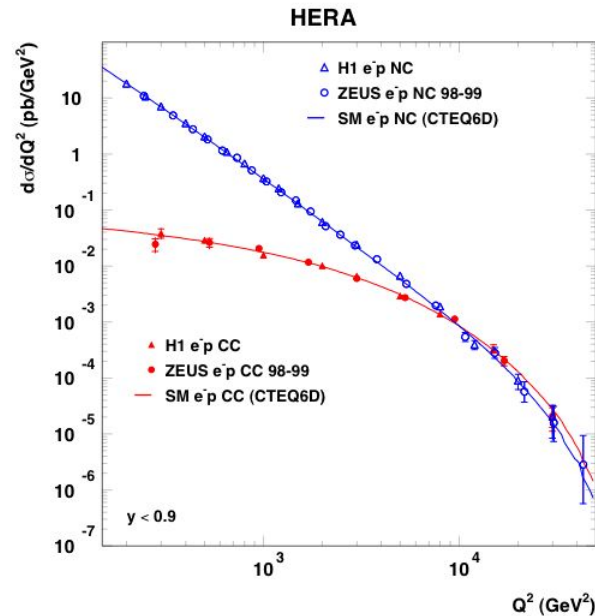


# Conclusion

- With approx. 200 pb<sup>-1</sup> significant tests of the SM electroweak sector at high spacelike momentum transfers have been performed :
  - Electroweak unification, CC-Propagator mass,
  - Light quark Z<sub>0</sub> couplings, CC-chiral structure,
  - NC polarisation dependences
- A factor 3 more luminosity still to come at HERA II and the availability of longitudinally polarized electrons and positrons will significantly enhance the sensitivity of these SM tests.



Despite 'HERA is a QCD machine' there are also interesting electroweak results which I hope justify the conference logo ...



Many thanks for  
your attention

Thanks to my colleagues from Zeus and H1 for supplying some of these slides