#### THE TRANSITION FROM $\sigma(\gamma^* p)$ TO $\sigma(\gamma p)$

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### **1** Relation between $\sigma(\gamma^* p)$ and $F_2$

In electron-proton scattering the structure function  $F_2^{ep}$  and the cross section  $\sigma(\gamma^* p)$  are related. For small x one finds :

$$\sigma^{\gamma^* p}(W^2, Q^2) \approx \frac{4\pi^2 \alpha}{Q^2} \cdot F_2^{ep}(W^2, Q^2) \tag{1}$$

Empirically<sup>1</sup> the HERA  $F_2$ -data ( $F_2$  short for  $F_2^{ep}$ ) are well described by :

$$F_2(x, Q^2) = m \cdot \log (x_0/x) \cdot \log (1 + Q^2/Q_0^2)$$
(2)

 $(m=0.40, x_0=0.04 \text{ and } Q_0^2=0.50 \text{ GeV}^2)$  for x < 0.001 without restriction in  $Q^2$ . Therefore, it is convenient to use the quantity  $q = \log (1+Q^2/Q_0^2)$  rather than  $Q^2$  itself and cast eq. 1 in the form :

$$\sigma^{\gamma^* p}(W^2, Q^2) = \frac{4\pi^2 \alpha}{Q_0^2} \cdot \frac{q}{Q^2/Q_0^2} \cdot \frac{F_2(W^2, q)}{q}$$
(3)

Since  $Q^2 \to 0$  implies  $q \sim Q^2/Q_0^2$ , the observable  $F_2/q$  carries the information on the transition  $\sigma(\gamma^* p) \to \sigma(\gamma p)$ .

## 2 The data

For the purpose of this study the low- $Q^2 F_2$ -data<sup>2,4,5</sup> together with the recently published data<sup>3,6</sup> are applied requiring x < 0.001.

The two ZEUS samples called  $bpc^4$  and  $bpt^6$  have 34  $F_2$ -measurements in common allowing for the hypothesis test :

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<sup>\*</sup>Talk presented at DIS2001,Bologna, April 2001

 $F_2^{bpt} - F_2^{bpc} = 0$ . Fig. 1, 2 shows the differences grouped in  $Q^2$ bins. The errors are calculated from the uncorrelated uncertainties alone. The average being  $\langle F_2^{bpt} - F_2^{bpc} \rangle = 0.031 \pm 0.008$  and  $\chi^2/dof = 50/33$  imply that for obtaining consistency between the two samples the systematic uncertainties must be taken into account. The HERA Collaborations have performed a careful study of the systematics. With  $F_2/q$  linear in log 1/x (see fig. 3), the effect of each systematic source can be illustrated in a planar arrow diagram, each arrow representing the  $1\sigma$ - shift in the average  $u_0$ 



Figure 1. Difference of  $F_2^{bpt} - F_2^{bpc}$  versus log  $W^2$  for various  $Q^2$ -bins; each line corresponds to 0.

and slope  $u_1$ , as shown in fig. 2 for the *bpt*-data.

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Figure 2. Shifts induced by each systematic source for the ZEUS-bpt data; the ellipse refers to the statistical precision.

# **3** Properties of $F_2/q$

The measured  $F_2/q$  values are displayed in fig. 3 versus log  $x_0/x$  for all  $Q^2$  supporting again eq. 2. There is no observable  $Q^2$ -dependence.

For the study of  $\sigma(\gamma^* p) \to \sigma(\gamma p) F_2/q$  is reexpressed in terms of  $(q, W^2)$  instead of  $(x, Q^2)$ . In the selected phase space region  $W^2 = Q^2 \cdot \frac{1}{x}$ , thus  $F_2/q \sim \log 1/x$  implies  $F_2/q \sim \log W^2$ . The *q*-dependence of the data is therefore analysed in the form :

$$\frac{F_2(W^2, q)}{q} = u_0(q) + u_1(q) \cdot (\log W^2 - \langle \log W^2 \rangle)$$



The slopes of  $F_2/q$  for fixed values of  $\log W^2$  are plotted in fig. 4(left) and are q-independent. In fig. 4(right) the  $F_2/q$  points

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Figure 4. Left : Slopes for constant  $W^2$ . Right :  $F_2/q$  versus q for W = 200 GeV compared with extended parametrisation; also displayed are the two direct measurements of  $\sigma(\gamma p)$  by H1 and ZEUS.

are displayed for the value W = 200 GeV, since there the corresponding phase space region is populated and direct  $\gamma p$  measurements exist. The lowest measured  $Q^2$  being 0.05 GeV<sup>2</sup> is quite close to 0 and the trend of the data suggests to bridge the gap linearly. Physically, the transition has to be smooth.

The prediction using eq. 2 is consistent with the  $F_2$ -data. Its differential form is :  $dF_2/q = m d\log W^2 - m(1 + Q_0^2/Q^2) dq$ . The slopes  $u_1(q)$  are consistent with m (the line in fig. 4 left). The slopes of  $F_2/q$  for W=200 GeV are  $Q^2$ -independent for large q, while for smaller q an enhancement due to the  $Q_0^2/Q^2$ -term occurs. For yet smaller values, outside the measured region, the formula (eq. 2) is bound to fail, since  $F_2/q$  is a function of x alone in the measured region, whereas  $\sigma(\gamma p)$  is a function of  $W^2$  alone. This asks for a quantity turning from x to  $W^2$ , as q approaches 0, which is achieved by the replacement  $\frac{x_0}{x} \rightarrow \frac{x_0}{x} \frac{Q^2}{Q^2 + Q_w^2}$ ; for consistency  $0 < Q_w^2 \ll Q_0^2$ .

H1<sup>7</sup> and ZEUS<sup>8</sup> have measured  $\sigma(\gamma p)$  at W = 200 and 207 GeV with  $q = \mathcal{O}(10^{-4}) \approx 0$ . The two points expressed as  $F_2/q$  (see eq. 3)

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are also displayed in fig. 4(right). The line in the figure shows the prediction of the extended formula with  $Q_w^2 \approx 0.05 \text{ GeV}^2$ .

### 4 Conclusions

All  $F_2$ -data considered are consistent with each other. The systematic uncertainties of the H1 and ZEUS data samples are well estimated and need to be taken into account in detail.

 $F_2/q$ , the key quantity in the study of  $\sigma(\gamma^* p) \to \sigma(\gamma p)$ , is well described by eq. 2 and can be extended to the  $\gamma p$ -limit with a small modification in agreement with the direct  $\sigma(\gamma p)$  measurements of H1 and ZEUS. The  $W^2$ -dependence of  $\sigma(\gamma p)$  turns out to be *logarithmic*.

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