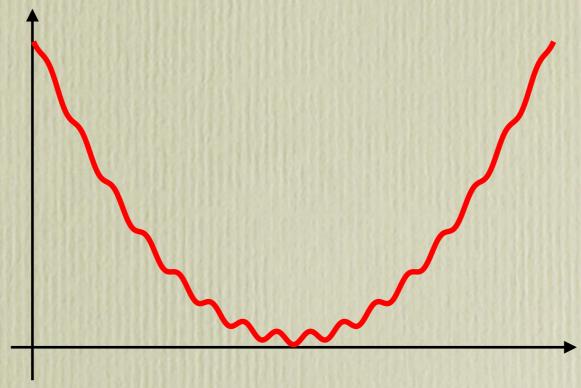
Back-reacted (rel)Axion Monodromy for Inflation & Relaxation in String Theory

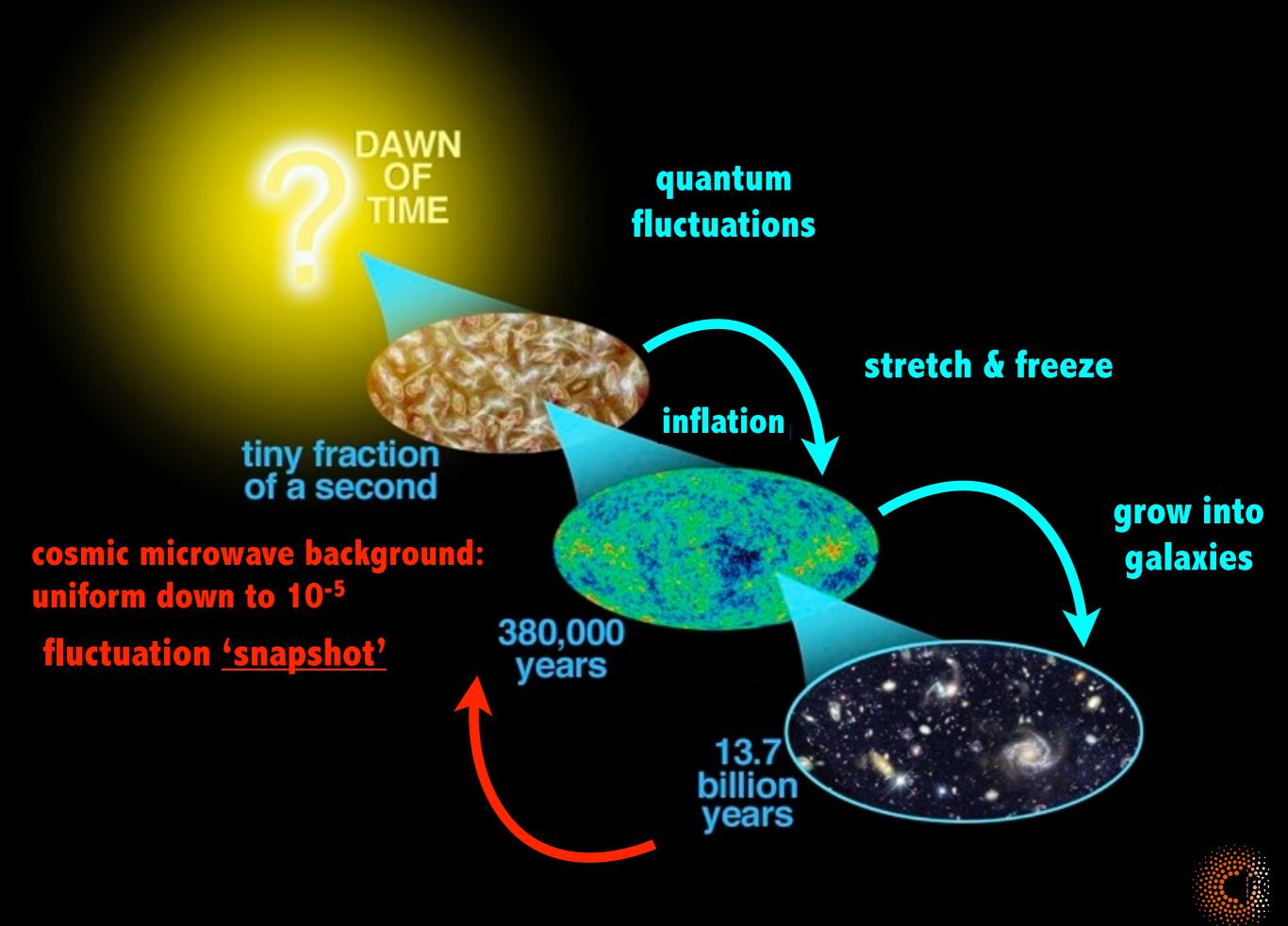
with Hebecker, Rompineve [1512.03768] with McAllister, Schwaller, Servant, Stout [1610.05320]



And work with: I. Ben-Dayan, W. Buchmüller, E. Dudas, K. Dutta, R. Flauger, R. Kallosh, A. Linde, J. Moritz, E. Pajer, F. Pedro, A. Retolaza, M. Rummel, F. Rühle, E. Silverstein, P. Vaudrevange, A. Uranga, C. Wieck, T. Wrase, G. Xu

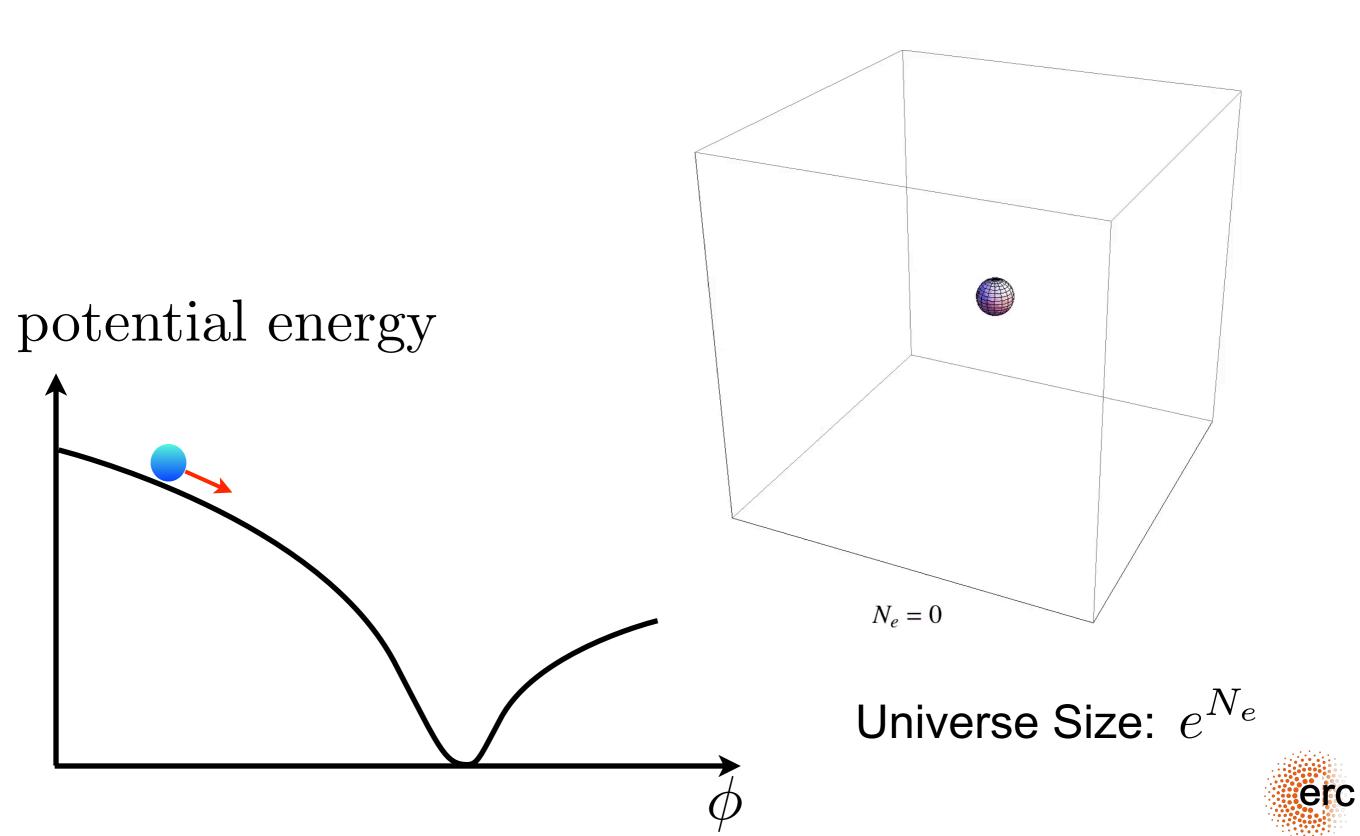
Alexander Westphal (DESY)





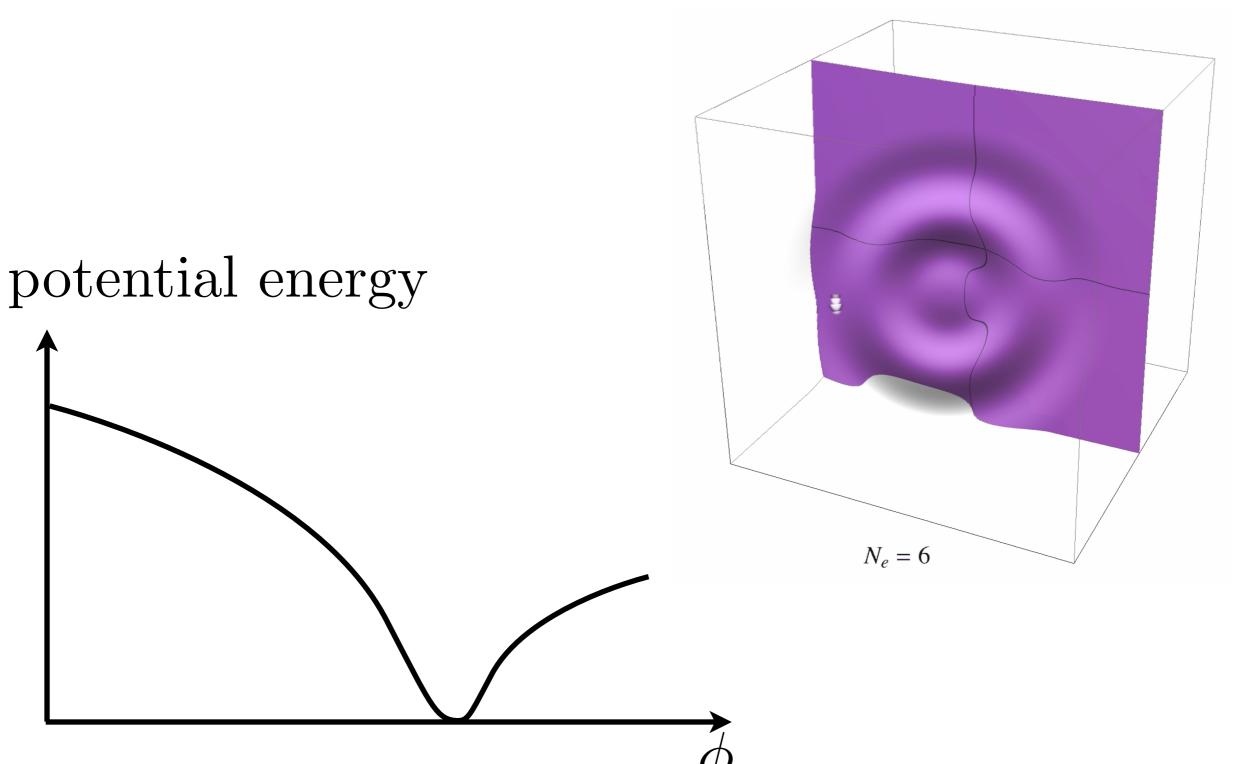
Inflation ...

(ideas by Guth, Linde & Steinhardt around 1980)



Inflation ...

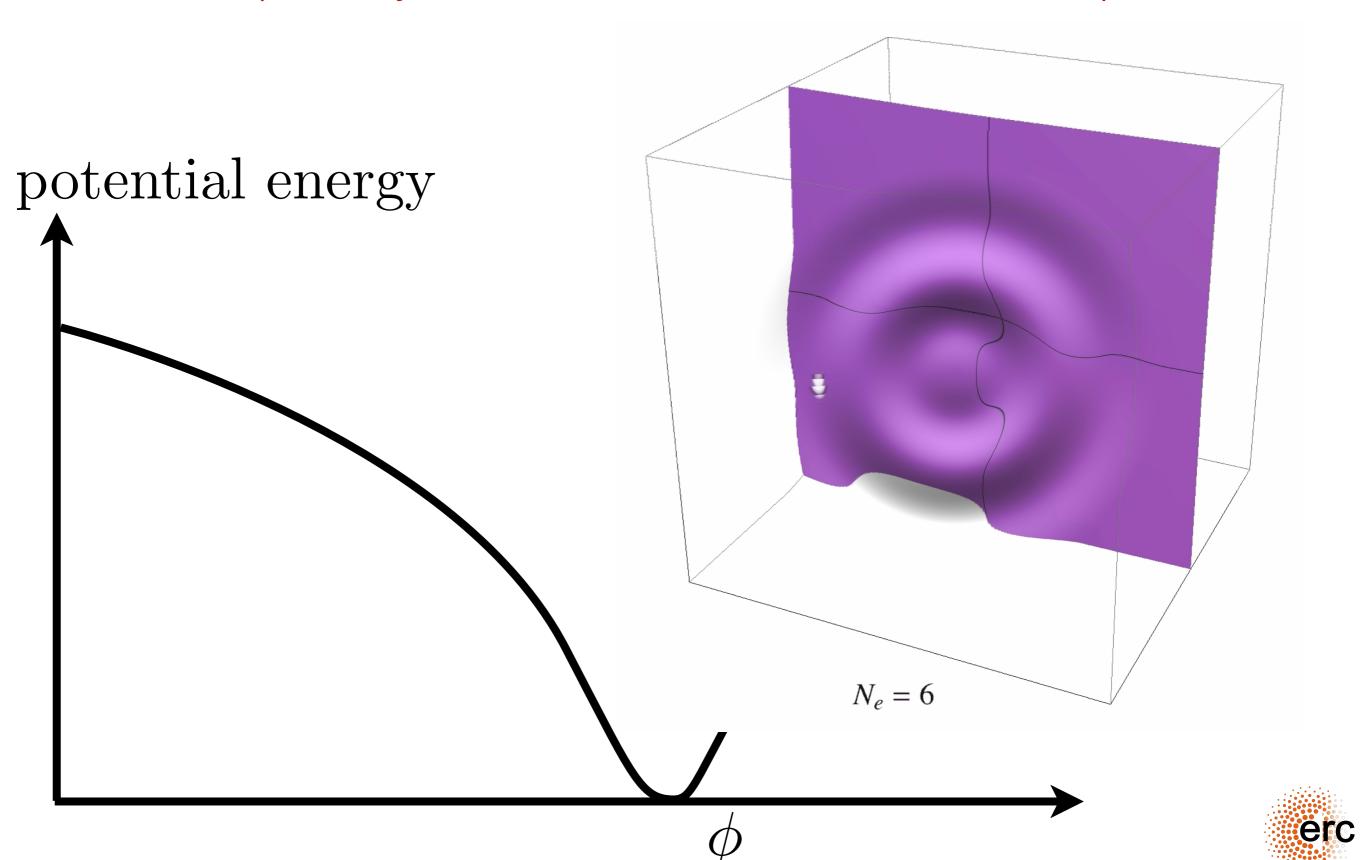
(ideas by Guth, Linde & Steinhardt around 1980)





Inflation ...

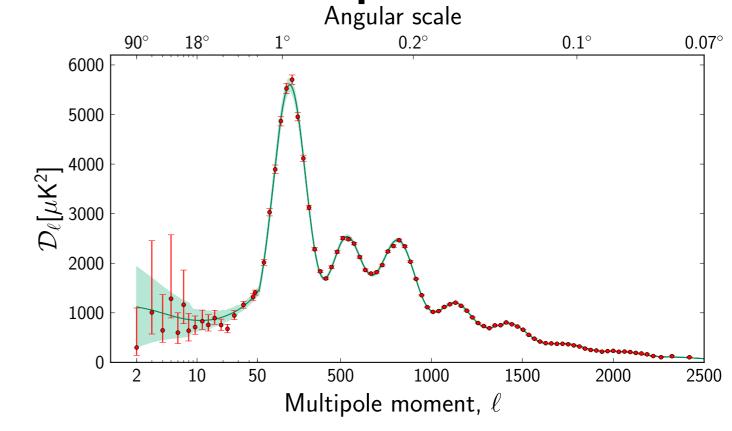
(ideas by Guth, Linde & Steinhardt around 1980)



Data!

PLANCK 2015 - most precise CMB data

BICEP2, Keck Array, BICEP3,, Spider ... since 2014: search for gravitational waves from inflation



LHC @ 8 TeV: Evidence for Scalars!

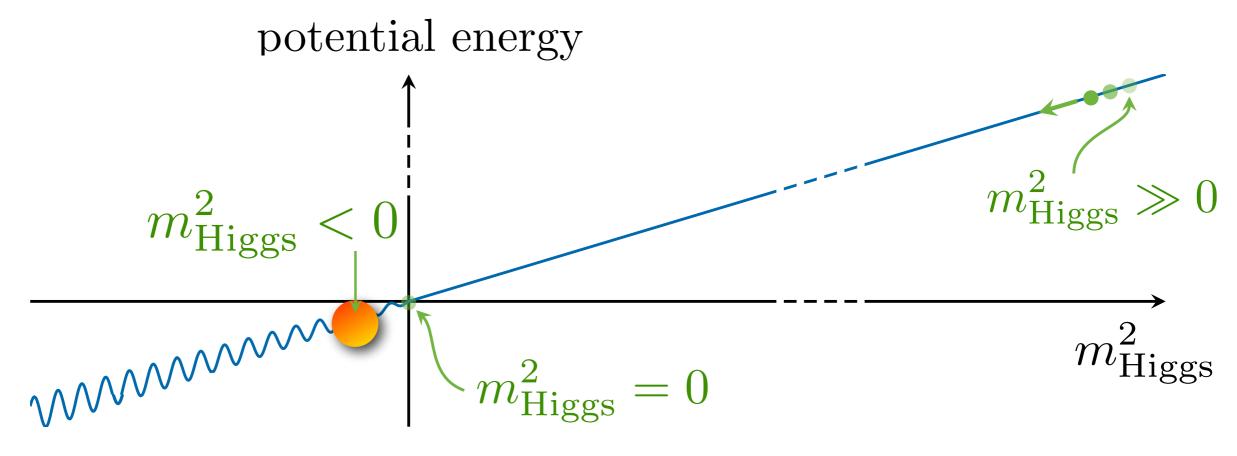
Higgs Discovery 2012!





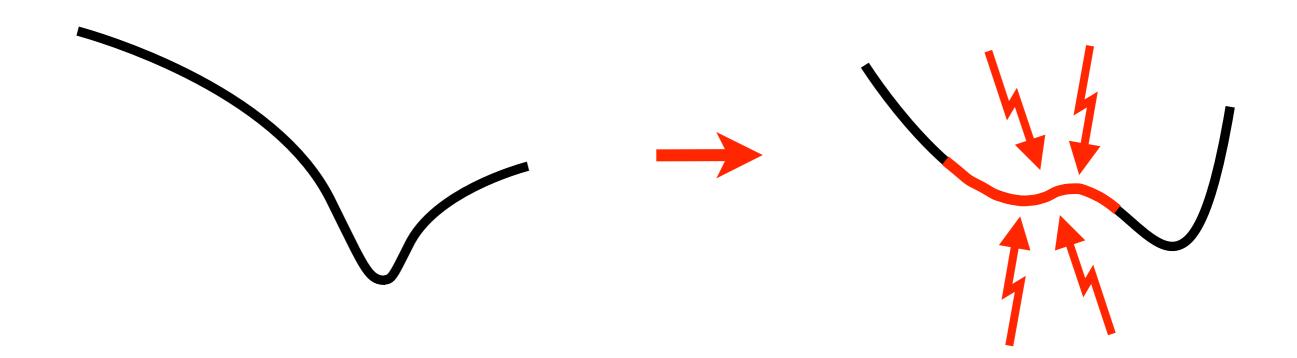
Light Higgs from Relaxation

$$m_{\text{Higgs}}^2 = M^2 - gM \left(\phi_{\text{init}} - \phi\right)$$



Quantum Theory ... Blessing & Bane

quantum theory can kill inflation / relaxation — quantum corrections!



need quantum gravity for control!



• use string theory to build inflation ...

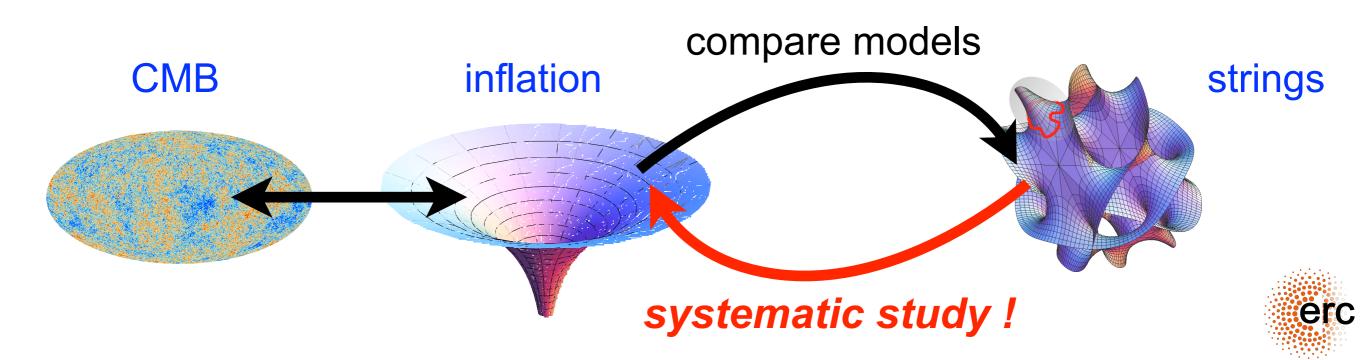


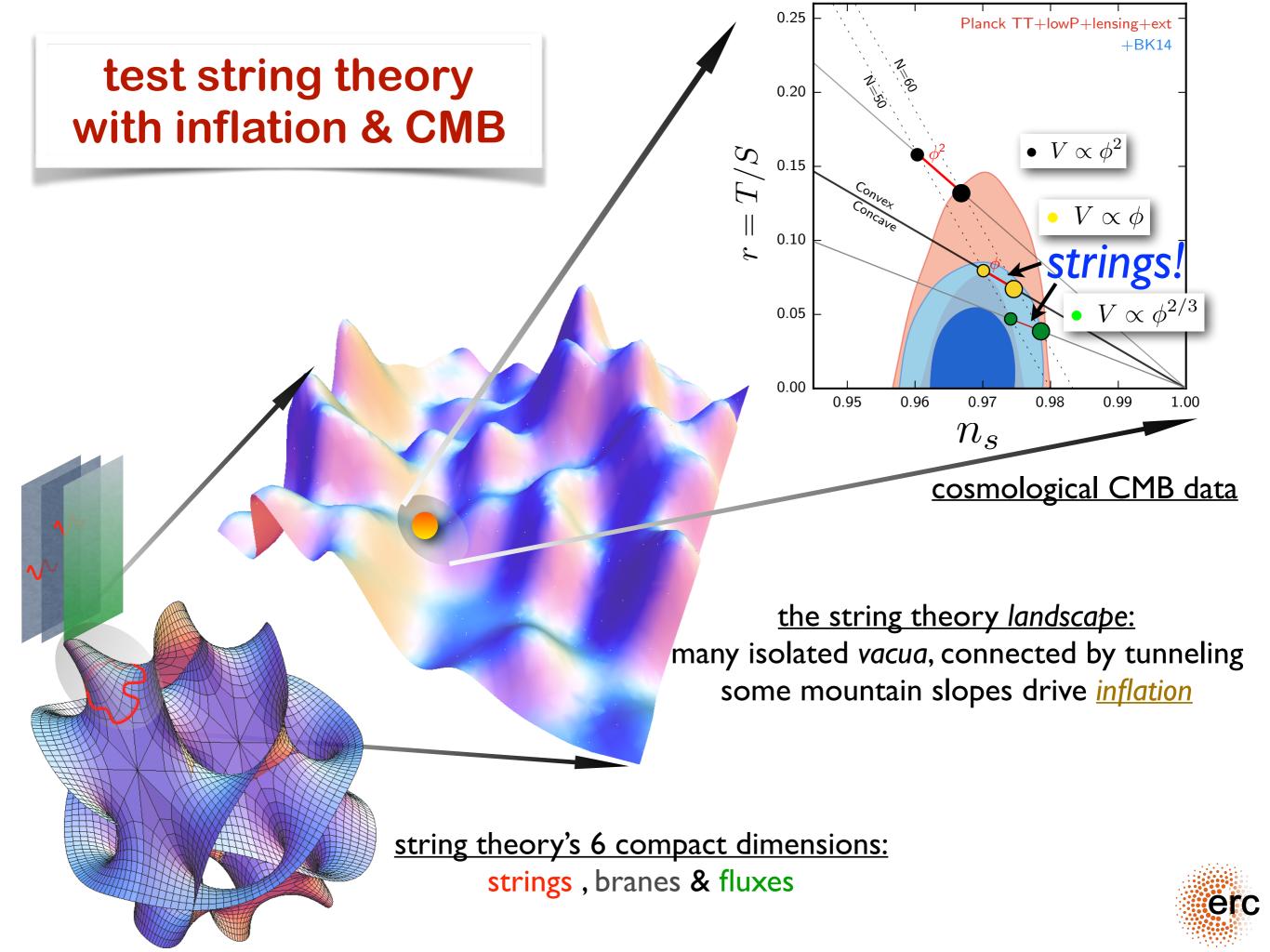
The Goal

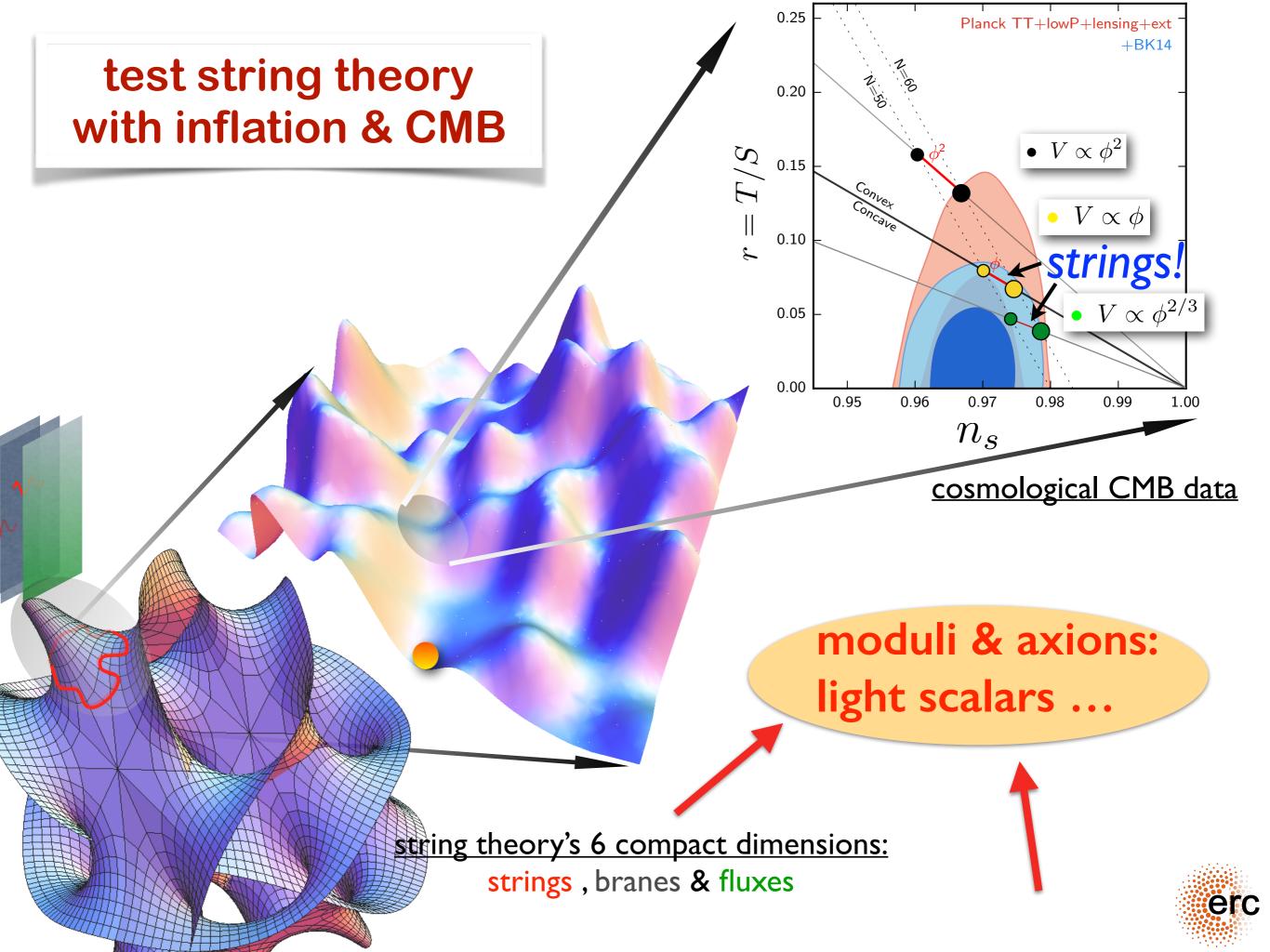
- energy scale of inflation 10^{13} x LHC close to **Planck Scale**!
 - → unique window to quantum gravity!

• proposal:

test string theory with inflation & CMB, or relaxation!



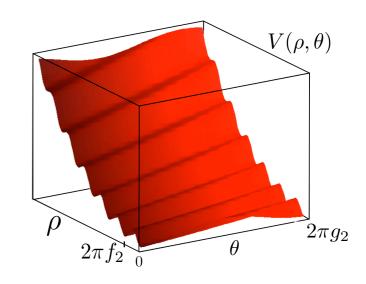


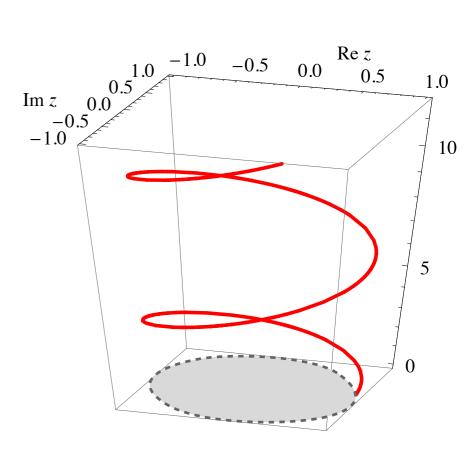


• string theory — use axions ...

crucial scale: f — not M_P !

backreaction at $\phi >> f$!







large fields need functional control

$$\mathcal{L} = (\partial \phi)^2 - V_0(\phi) - \sum_i g^{e_i} \left(\frac{\phi}{M}\right)^{d_i} , \quad d_i > 0 , \quad e_i \ge 0$$

 \Rightarrow $g \rightarrow 0$ limit: know your sequence $\{d_i, e_i\}$!

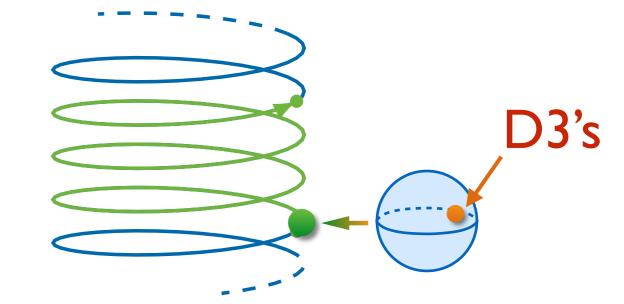
 $e_i = 0$ spells disaster at $\phi >> M \dots$



$$\frac{\phi}{f} = \mathcal{O}(100)$$

$$\frac{\phi}{f} \gg 10^6 \qquad \dots \qquad ?$$

$$\mathcal{L}_Q \sim q \int A_1$$



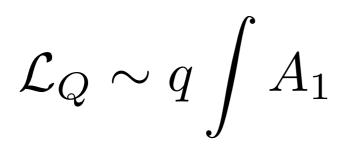
$$\mathcal{L}_{Q}^{NS5} \sim \int C_2 \wedge C_4 \sim \int C_2 \times \int C_4 \sim N_w \int C_4$$

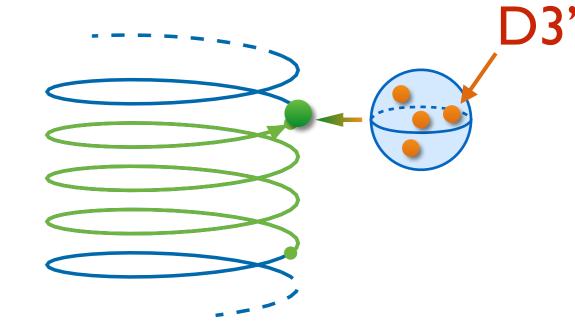
$$\underbrace{\int C_2 \times \int C_4}_{Q=N_w} C_4 \sim N_w \int C_4$$

$$\underbrace{\int C_4 \times \int C_4}_{Q^{D3}} C_4$$

charge backreacts on geometrycharged vs neutral BH!







$$\mathcal{L}_{Q}^{NS5} \sim \int C_2 \wedge C_4 \sim \int C_2 \times \int C_4 \sim N_w \int C_4$$

$$\underbrace{\int C_4}_{Q=N_w} \mathcal{L}_{Q}^{D3}$$

charge backreacts on geometrycharged vs neutral BH!



'gravity should be weak'

$$\frac{q}{m} > 1$$

[Arkani-Hamed et al. '06]

erc

magnetic WGC

dual coupling:

$$\tilde{e} = 1/e$$

monopoles:

$$m_{mon} \sim \tilde{e}^2 \int_{r_{min}} |F_2|^2 \sim \frac{M}{e^2} , \quad M \sim \frac{1}{r_{min}}$$

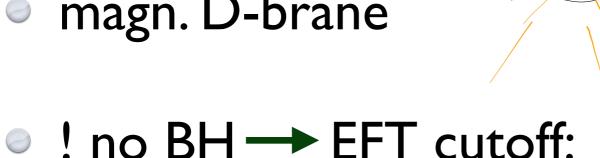


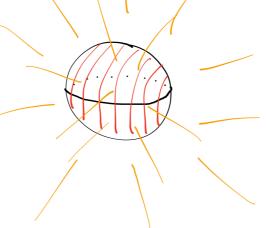
no BH!

$$r_{min} \sim \frac{1}{M} > R_{BH} \sim m_{mon} \sim \frac{M}{e^2}$$

implies EFT cutoff: M < e

- string theory:
 - magn. D-brane



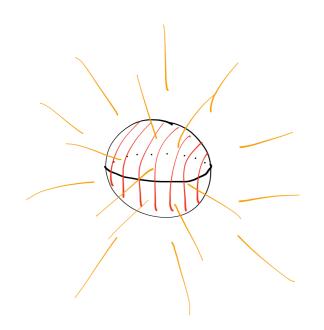


$$M < e^{1/(p+1)}$$

no BH!

$$r_{min} \sim \frac{1}{M} > R_{BH} \sim m_{mon} \sim \frac{M}{e^2}$$

implies EFT cutoff: M < e



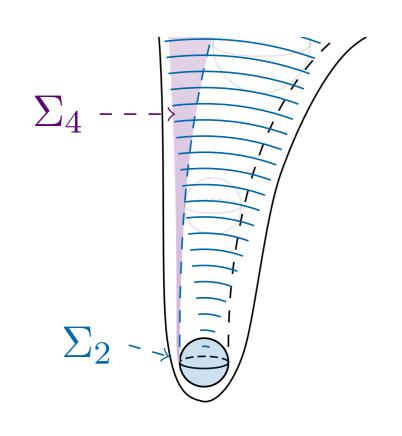
$$M < e^{1/(p+1)}$$

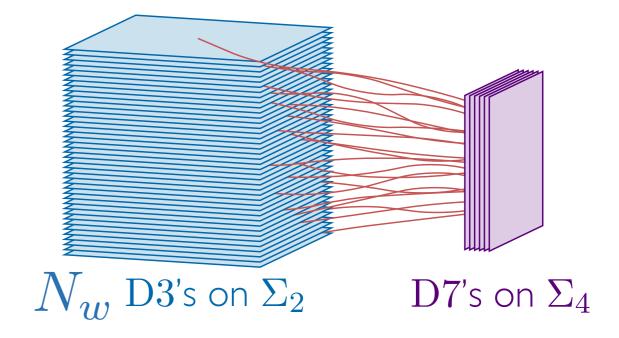
$$H \sim m\phi < M < e^{1//3} \sim (mf)^{1/3}$$

monodromy charge N_w is physical — backreacts

closed string picture:

open string picture:



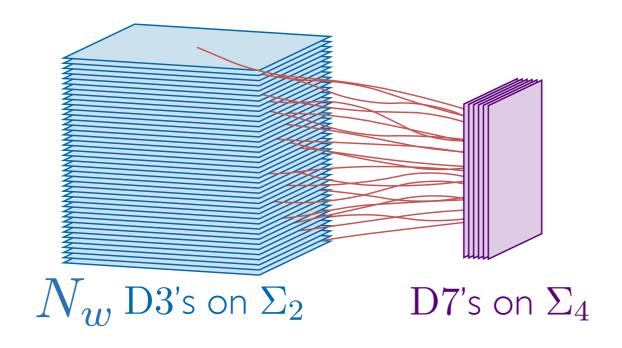




$$\delta g_{\rm YM}^{-2} \sim N_w$$



<u>light</u> & <u>charged</u> 3-7 string for every unit of N_w!



$$\delta \left(\frac{1}{g_{\rm YM}^2}\right)_{1-loop} \sim N_{3-7} = N_w$$



So ...

$$N_w \gtrsim 10^6 \ , \ c \sim 1/(2\pi)$$

$$\Lambda_c^3 \langle h \rangle \cos(\phi/f)$$
 : $\Lambda_c^3 \sim e^{-\frac{1}{g_{YM}^2}} \rightarrow e^{-\frac{1}{g_{YM}^2} - cN_w} \rightarrow 0$

Relaxions Run Away!



So ...

$$N_w \lesssim 100$$
 , $c \lesssim 1/(2\pi)$

$$\Lambda_c^3 \cos(\phi/f)$$
 : $\Lambda_c^3 \sim e^{-\frac{1}{g_Y^2 M}} \rightarrow e^{-\frac{1}{g_Y^2 M}} \stackrel{-cN_w}{\to} finite!$

Inflations proceeds - sometimes ...

