

Determination of Higgs coupling uncertainties in the THDMS

Bachelor thesis colloquium

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Heading Agenda

Subheading, optional

1. Introduction
2. N2HDM
3. THDMS
4. Data generation
5. Evaluation

Introduction

LEP and CMS excesses

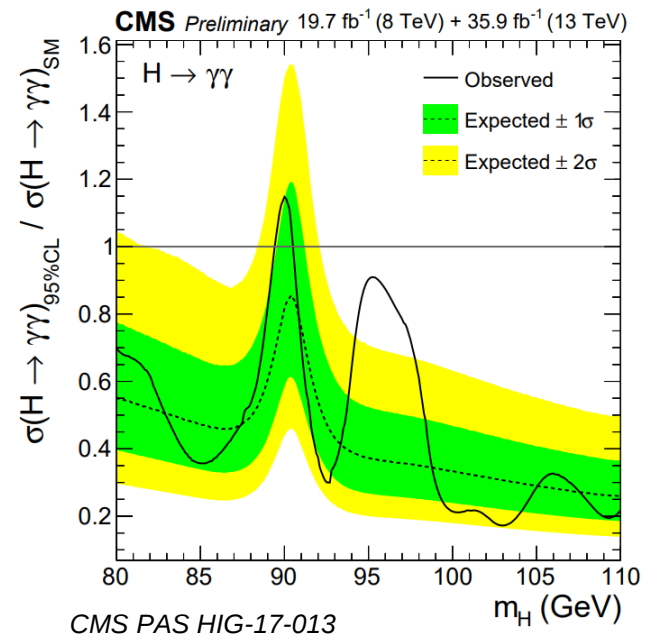
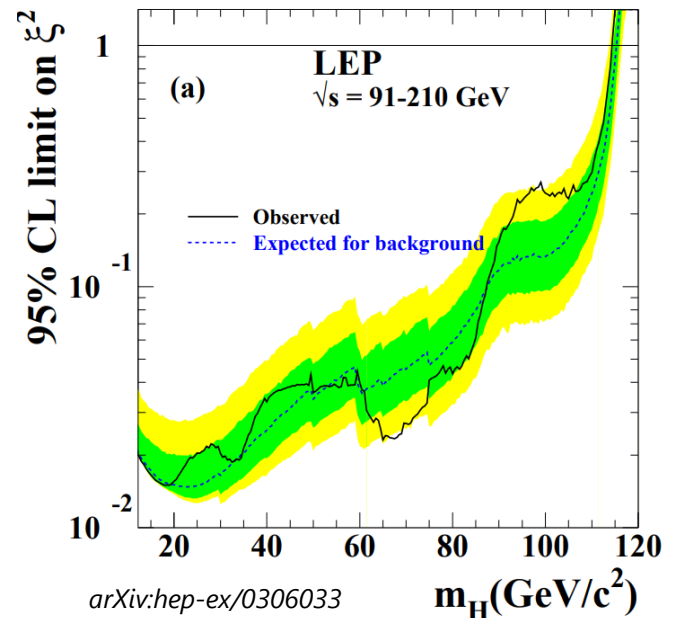
- 96 GeV excess at LEP and CMS

$$\mu_{LEP} = \frac{\sigma(e^+e^- \rightarrow Z\phi \rightarrow Zb\bar{b})}{\sigma^{\text{SM}}(e^+e^- \rightarrow ZH_{\text{SM}}^\phi \rightarrow Zb\bar{b})} = 0.117 \pm 0.057$$

$$\mu_{CMS} = \frac{\sigma(gg \rightarrow \phi \rightarrow \gamma\gamma)}{\sigma^{\text{SM}}(gg \rightarrow H_{\text{SM}}^\phi \rightarrow \gamma\gamma)} = 0.6 \pm 0.2$$

- Idea: fit both excesses with a 96 GeV Higgs boson.

$$\chi_{\text{CMS-LEP}}^2 = \frac{(\mu_{CMS} - 0.6)^2}{0.2^2} + \frac{(\mu_{LEP} - 0.117)^2}{0.057^2}$$



Introduction

Data generation

Scan

- Parameter scan using SARAH
- Create spectra with SPheno

Constraints

- Theoretical constraints
 - Tree-Level Perturbativity → SARAH & SPheno checks
 - Boundedness from below → SARAH & SPheno checks
- Experimental constraints
 - LEP, Tevatron & LHC Higgs searches → HiggsBounds
 - SM Higgs couplings → HiggsSignals
 - Electroweak precision observables → Fit for S, T, U parameters
 - Flavour physics $B \rightarrow X_s \gamma$ limit → Lower bounds for $m_{H^\pm}$

Introduction

Couplings

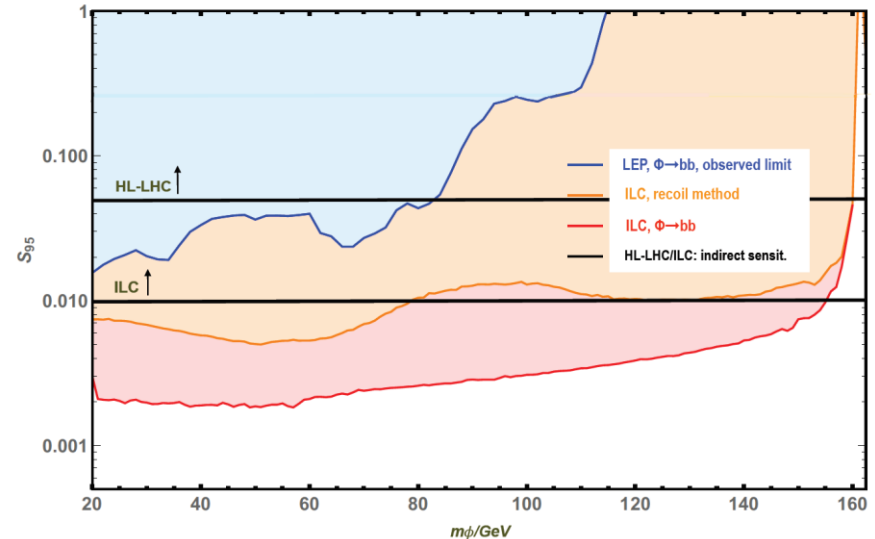
- S_{95} -method to calculate precisions at LEP, ILC and HL-LHC

$$S_{95} = \frac{\hat{\sigma}}{\sigma_{ref}} = \frac{\hat{n}}{n}$$

$$S_{95} \hat{=} \left| \frac{(g_{\phi ZZ})^2}{(g_{HZZ}^{\text{SM}})^2} \right|$$

- Upper bound for Higgs searches at HL-LHC

$$\frac{(g_{\phi ZZ})^2}{(g_{HZZ})^2} \leq 1 - (1 - |\Delta g_{HZZ}|)^2$$



arXiv:1801.09662 [hep-ph]

N2HDM

Potential and field conventions

$$\begin{aligned} V = & m_{11}^2 |\Phi_1|^2 + m_{22}^2 |\Phi_2|^2 - m_{12}^2 (\Phi_1^\dagger \Phi_2 + h.c.) + \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 \\ & + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) + \frac{\lambda_5}{2} [(\Phi_1^\dagger \Phi_2)^2 + h.c.] \\ & + \frac{1}{2} m_S^2 \Phi_S^2 + \frac{\lambda_6}{8} \Phi_S^4 + \frac{\lambda_7}{2} (\Phi_1^\dagger \Phi_1) \Phi_S^2 + \frac{\lambda_8}{2} (\Phi_2^\dagger \Phi_2) \Phi_S^2 . \end{aligned}$$

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{1}{\sqrt{2}}(v_1 + \rho_1 + i\eta_1) \end{pmatrix} \quad \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}}(v_2 + \rho_2 + i\eta_2) \end{pmatrix}$$

$$\Phi_S = v_S + \rho_S$$

$\mathbb{Z}_2$ -symmetry

$$\Phi_1 \rightarrow \Phi_1 , \quad \Phi_2 \rightarrow -\Phi_2 , \quad \Phi_S \rightarrow \Phi_S$$

N2HDM

Types of 2HDM

	u -type	d -type	leptons
type I	Φ_2	Φ_2	Φ_2
type II	Φ_2	Φ_1	Φ_1
lepton-specific	Φ_2	Φ_2	Φ_1
flipped	Φ_2	Φ_1	Φ_2

	u -type($c_{ht\bar{t}}$)	d -type($c_{hb\bar{b}}$)	leptons($c_{h\tau^+\tau^-}$)
type I	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$
type II	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$	$\frac{R_{i1}}{c_\beta}$
lepton-specific	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$
flipped	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$	$\frac{R_{i2}}{s_\beta}$

$$R = \begin{pmatrix} c_{\alpha_1} c_{\alpha_2} & s_{\alpha_1} c_{\alpha_2} & s_{\alpha_2} \\ -s_{\alpha_1} c_{\alpha_3} - c_{\alpha_1} s_{\alpha_2} s_{\alpha_3} & c_{\alpha_1} c_{\alpha_3} - s_{\alpha_1} s_{\alpha_2} s_{\alpha_3} & c_{\alpha_2} s_{\alpha_3} \\ s_{\alpha_1} s_{\alpha_3} - c_{\alpha_1} s_{\alpha_2} c_{\alpha_3} & -s_{\alpha_1} s_{\alpha_2} c_{\alpha_3} - c_{\alpha_1} s_{\alpha_3} & c_{\alpha_2} c_{\alpha_3} \end{pmatrix} \begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix} = R \begin{pmatrix} \rho_1 \\ \rho_2 \\ \rho_3 \end{pmatrix}$$

$$\begin{aligned}
 V = & m_{11}^2(\Phi_1^\dagger \Phi_1) + m_{22}^2(\Phi_2^\dagger \Phi_2) + \frac{\lambda_1}{2}(\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2}(\Phi_2^\dagger \Phi_2)^2 + \lambda_3(\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) \\
 & + \lambda_4(\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) + m_S^2(S^\dagger S) + \lambda'_1(S^\dagger S)(\Phi_1^\dagger \Phi_1) + \lambda'_2(S^\dagger S)(\Phi_2^\dagger \Phi_2) \\
 & + \frac{\lambda_3''}{4}(S^\dagger S)^2 + \left(-m_{12}^2 \Phi_1^\dagger \Phi_2 + \frac{\mu_{s1}}{6} S^3 + \mu_{12} S \Phi_1^\dagger \Phi_2 + \text{h.c.} \right)
 \end{aligned}$$

$$\begin{aligned}
 \Phi_1 = \begin{pmatrix} \chi_1^+ \\ \phi_1 \end{pmatrix} &= \begin{pmatrix} \chi_1^+ \\ v_1 + \frac{\rho_1 + i\eta_1}{\sqrt{2}} \end{pmatrix} \quad \Phi_2 = \begin{pmatrix} \chi_2^+ \\ \phi_2 \end{pmatrix} = \begin{pmatrix} \chi_2^+ \\ v_2 + \frac{\rho_2 + i\eta_2}{\sqrt{2}} \end{pmatrix} \\
 S &= v_S + \frac{\rho_S + i\eta_S}{\sqrt{2}}
 \end{aligned}$$

$$\tan \beta = \frac{v_2}{v_1}, \quad v = \sqrt{v_1^2 + v_2^2} = 174 \text{ GeV}$$

$\mathbb{Z}_2$ -symmetry

$$\Phi_1 \rightarrow \Phi_1, \quad \Phi_2 \rightarrow -\Phi_2, \quad \Phi_S \rightarrow \Phi_S$$

$\mathbb{Z}_3$ -symmetry

$$\begin{pmatrix} \Phi_1 \\ \Phi_2 \\ S \end{pmatrix} = \begin{pmatrix} 1 & & \\ & e^{i2\pi/3} & \\ & & e^{-i2\pi/3} \end{pmatrix} \begin{pmatrix} \Phi_1 \\ \Phi_2 \\ S \end{pmatrix}$$

$$\begin{aligned}
 V = & m_{11}^2(\Phi_1^\dagger \Phi_1) + m_{22}^2(\Phi_2^\dagger \Phi_2) + \frac{\lambda_1}{2}(\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2}(\Phi_2^\dagger \Phi_2)^2 + \lambda_3(\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) \\
 & + \lambda_4(\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) + m_S^2(S^\dagger S) + \lambda'_1(S^\dagger S)(\Phi_1^\dagger \Phi_1) + \lambda'_2(S^\dagger S)(\Phi_2^\dagger \Phi_2) \\
 & + \frac{\lambda_3''}{4}(S^\dagger S)^2 + \left(-m_{12}^2 \Phi_1^\dagger \Phi_2 + \frac{\mu_{s1}}{6} S^3 + \mu_{12} S \Phi_1^\dagger \Phi_2 + \text{h.c.} \right)
 \end{aligned}$$

12 free parameters

$$\lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda'_1, \lambda'_2, \lambda_3'', m_{12}, \mu_{s1}, \mu_{12}, v_S, \tan \beta$$

Mass input

$$m_{h_{1,2,3}}, m_{a_{1,2}}, m_{h^\pm}, \alpha_1, \alpha_2, \alpha_3, \alpha_4, v_S, \tan \beta$$

$\alpha_{1,2,3,4}$	→ mixing angles
$m_{h_{1,2,3}}$	→ scalar Higgs masses
$m_{a_{1,2}}$	→ pseudo-scalar Higgs masses
$m_{H^\pm}$	→ charged Higgs mass
$\tan \beta$	→ ratio of doublet vevs
v_S	→ singlet vev

Data generation

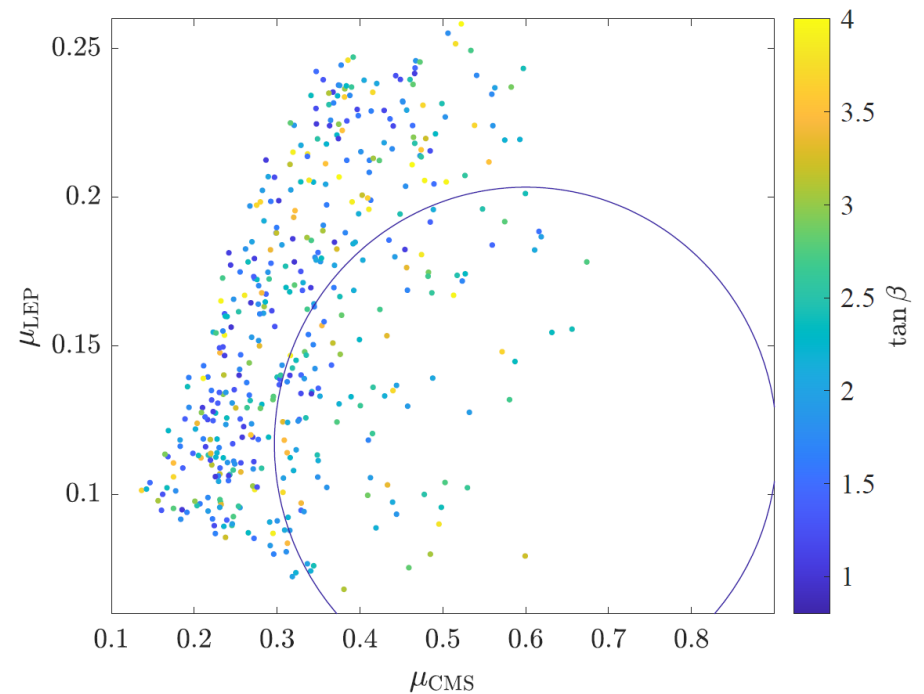
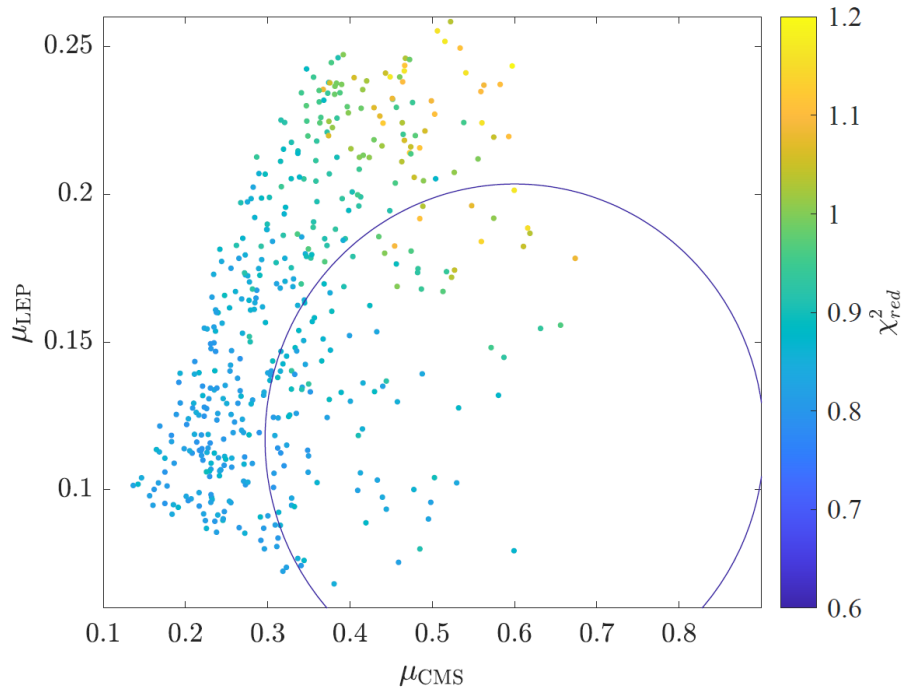
Scan using SARAH

Parameter	Scanning range	Parameter	Scanning range
α_1	(0.8, 1.45)	m_{h_3}	(650, 1000)
α_2	(0.93, 1.25)	m_a	(200, 500)
α_3	$(0.93, \frac{\pi}{2})$	m_A	(650, 1000)
α_4	$(1.565, \frac{\pi}{2})$	$m_{H^\pm}$	(650, 1000)
m_{h_1}	(80, 93)	$\tan \beta$	(1, 4)
m_{h_2}	(122, 126)	v_S	(100, 2000)

Values taken in $\overline{DR}$ -scheme and converted to on-shell scheme

Data generation

Scan using Sarah

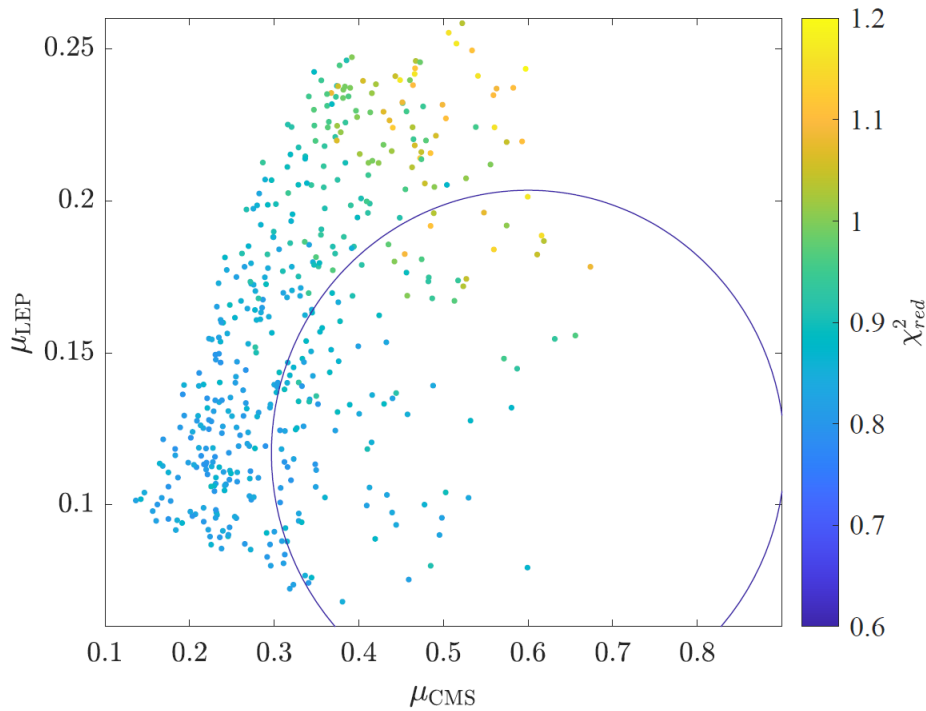


$$\chi^2_{\text{red}} = \frac{\chi^2}{n_{\text{obs}}}$$

χ^2 calculated by HiggsSignals with $n_{\text{obs}} = 107$

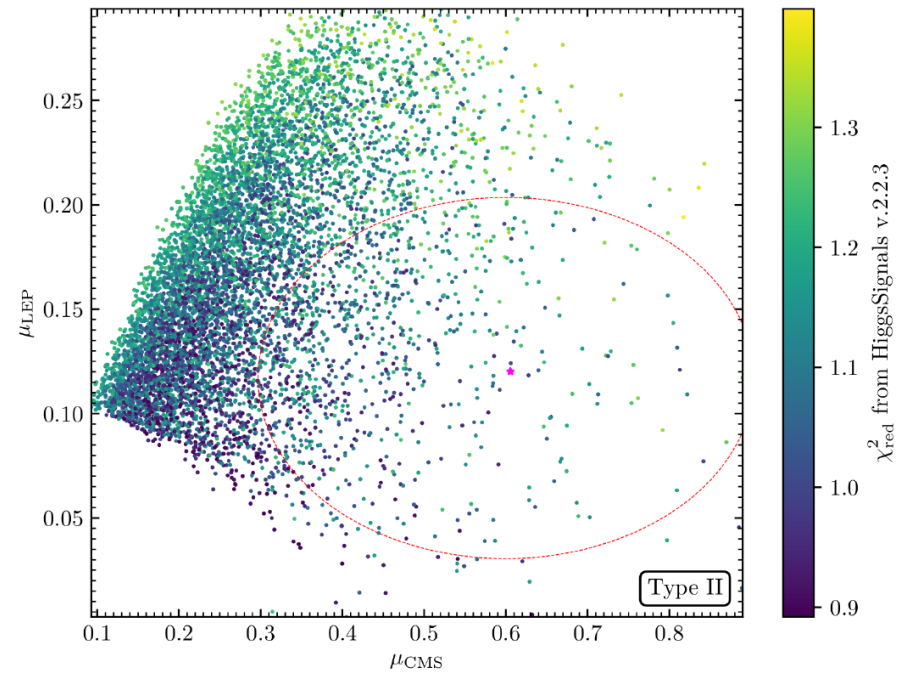
Evaluation

Comparison of N2HDM to THDMS data sets



$$n_{\text{obs}} = 107$$

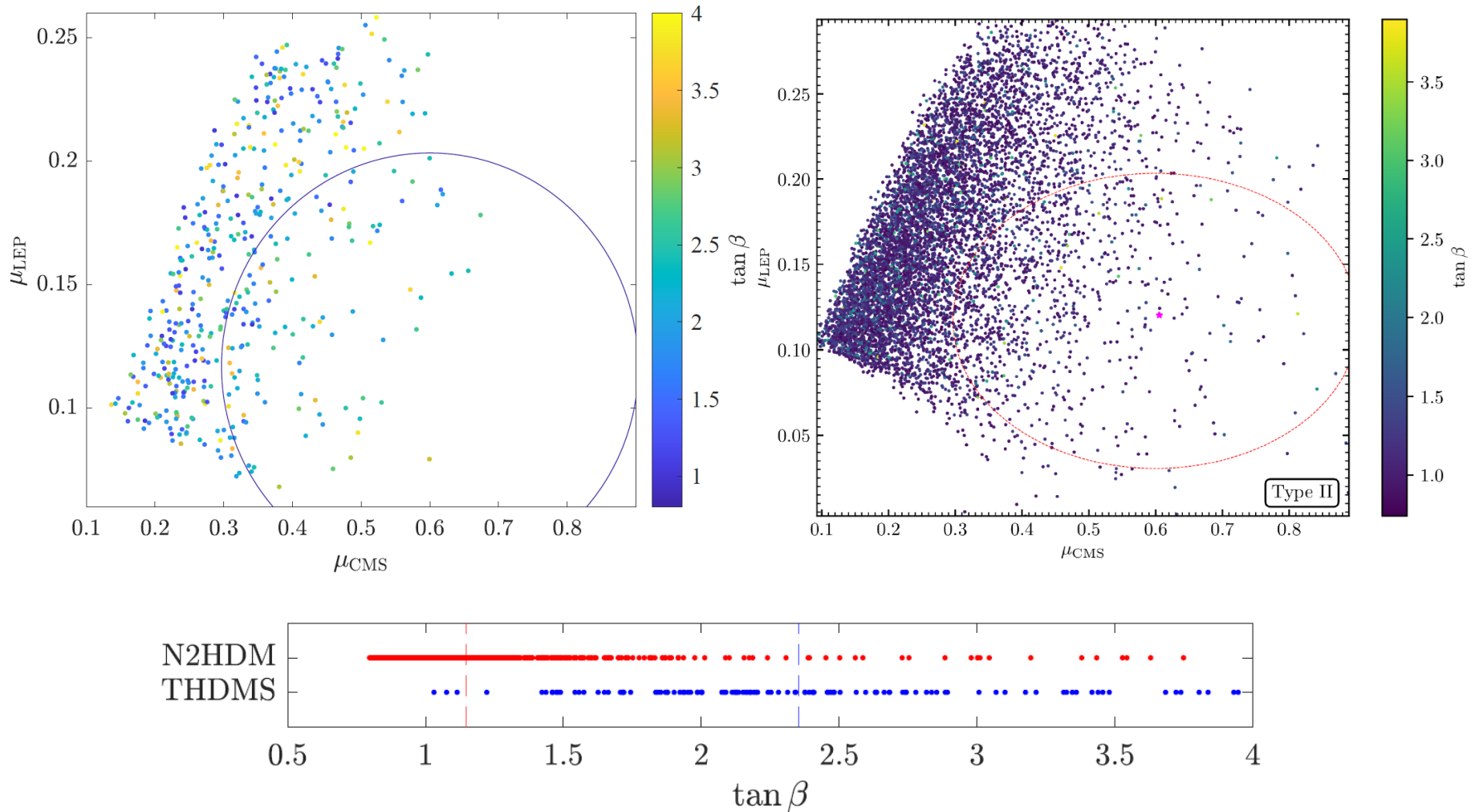
$$\chi_{\text{red}}^2 = \frac{\chi^2}{n_{\text{obs}}}$$



$$n_{\text{obs}} = 101$$

Evaluation

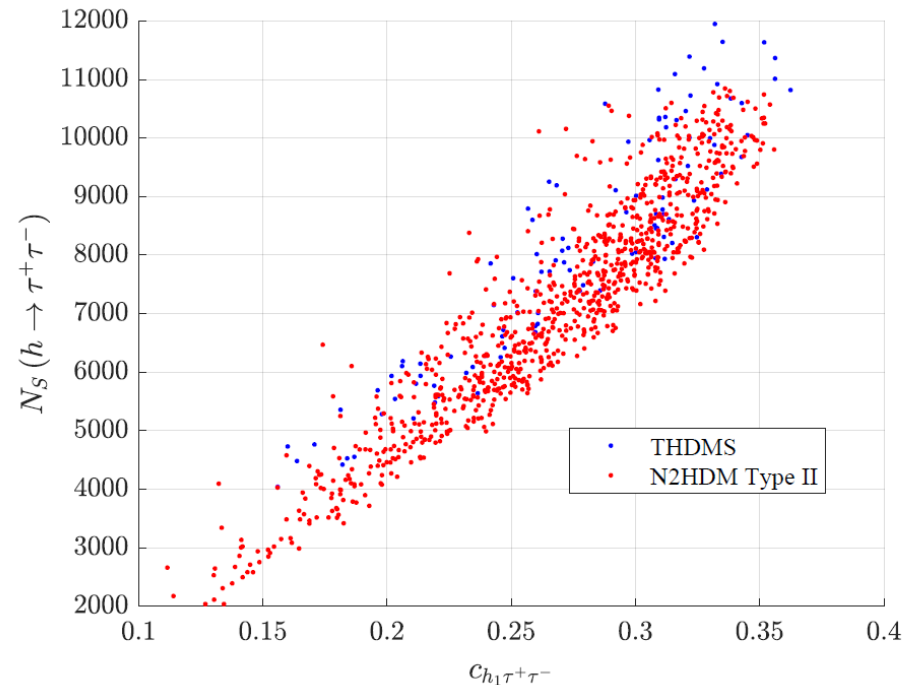
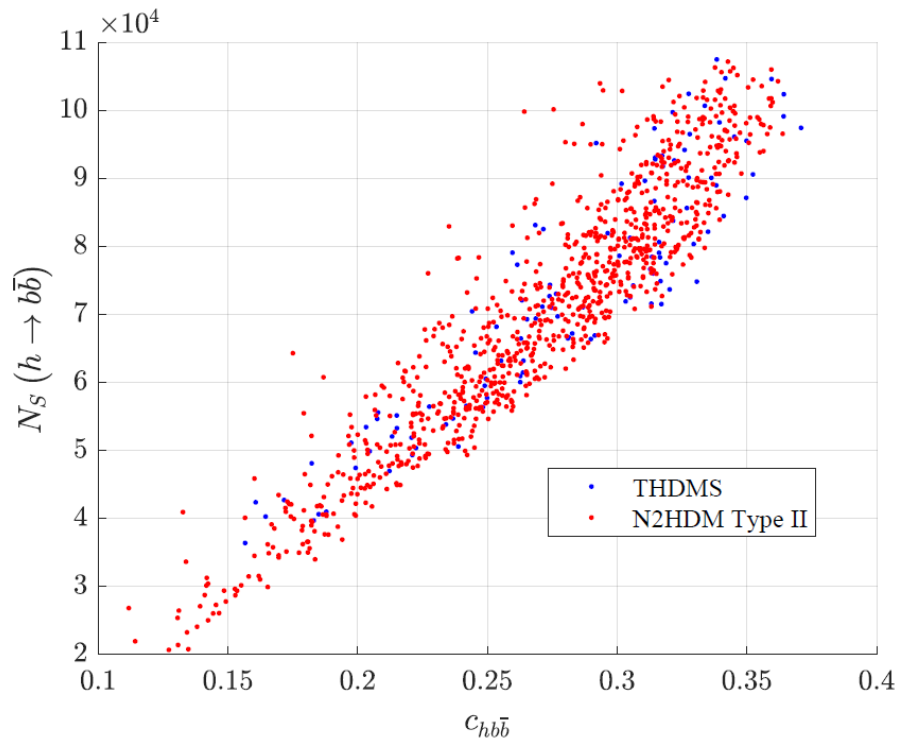
Comparison of N2HDM to THDMS data sets



Evaluation

Reproduction of N2HDM signal events

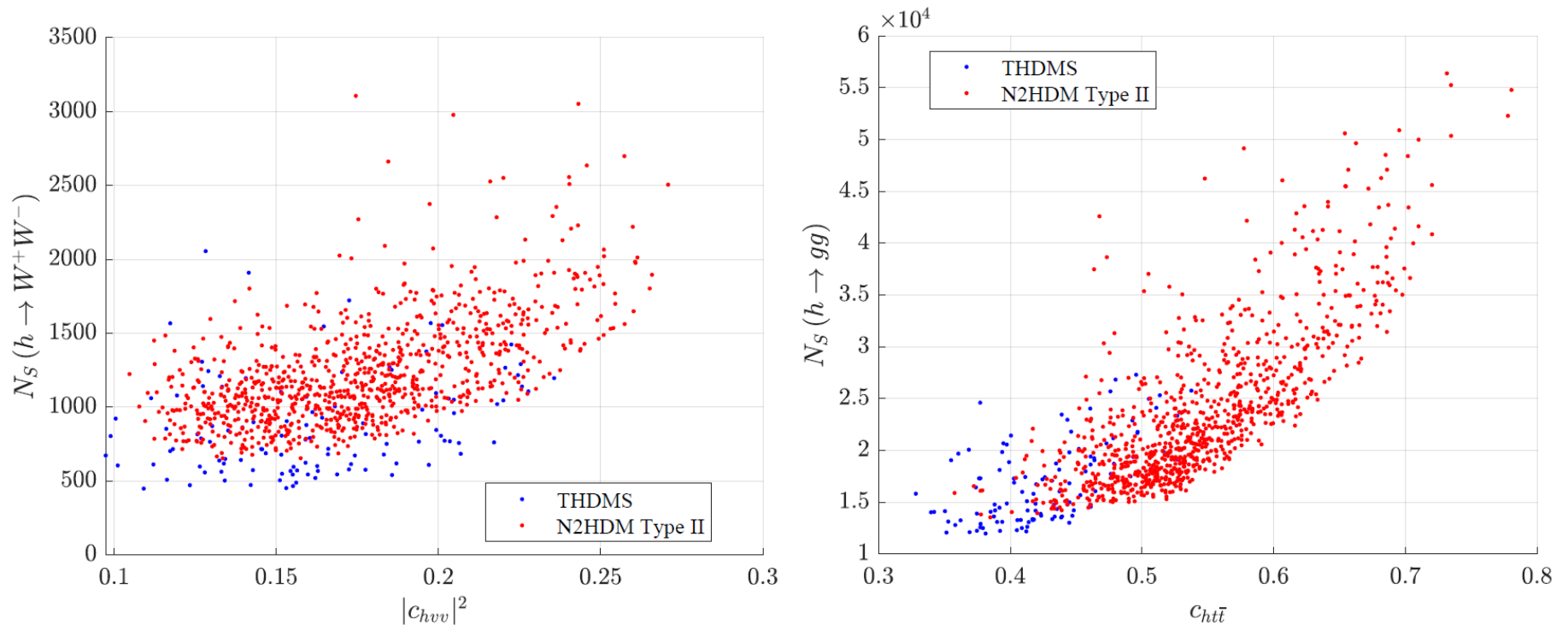
$$N_S(h \rightarrow xx) = \mathcal{L}_{int} \cdot \sigma(e^+e^- \rightarrow hZ) \cdot \text{BR}(h \rightarrow xx)$$



Evaluation

Reproduction of N2HDM signal events

$$N_S(h \rightarrow xx) = \mathcal{L}_{int} \cdot \sigma(e^+e^- \rightarrow hZ) \cdot \text{BR}(h \rightarrow xx)$$



Evaluation

Calculation of coupling uncertainties

Main equations:

$$\frac{\left(\frac{\Delta g_x}{g_x}\right)_h}{\left(\frac{\Delta g_x}{g_x}\right)_H} = \sqrt{\left(\frac{D + f_H}{1 + f_H}\right) \left(\frac{\sigma(e^+e^- \rightarrow ZH)}{\sigma(e^+e^- \rightarrow Zh)}\right) \left(\frac{\text{BR}(H \rightarrow xx)}{\text{BR}(h \rightarrow xx)}\right)} \cdot \frac{(1 - \text{BR}(H \rightarrow xx))}{(1 - \text{BR}(h \rightarrow xx))}$$

Evaluation via decay

$$D := \frac{f_H}{f_h} \quad f_\phi := \left(\frac{S}{B}\right)_\phi = \left(\frac{N_S}{N_B}\right)_\phi$$

Evaluation via production

$$\frac{\left(\frac{\Delta g_Z}{g_Z}\right)_h}{\left(\frac{\Delta g_Z}{g_Z}\right)_H} = \sqrt{\frac{\sigma(e^+e^- \rightarrow ZH)}{\sigma(e^+e^- \rightarrow Zh)}}$$

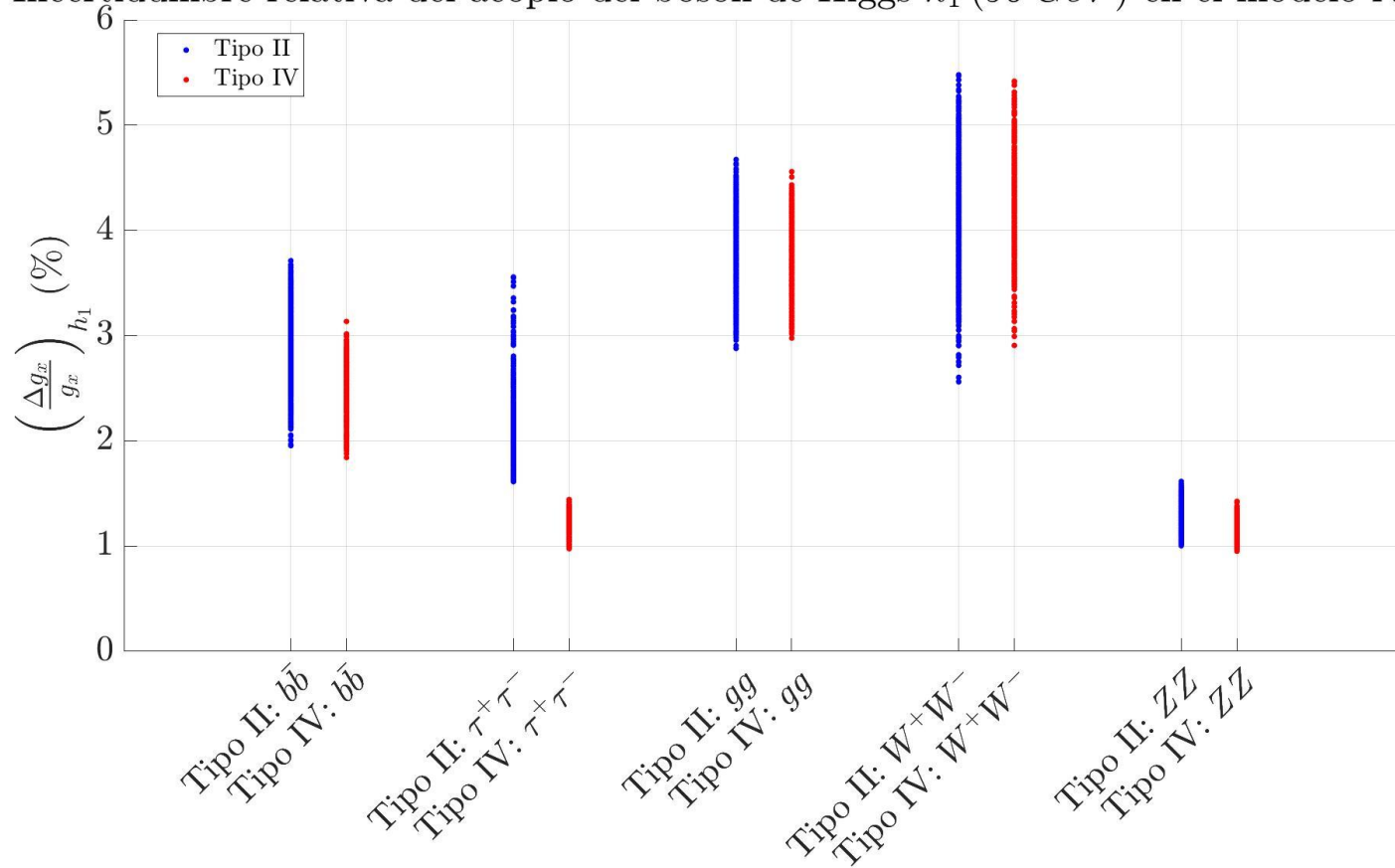
- D parameter is the only unknown value

Previous studies of Higgs between 20 and 100 GeV suggest $D \sim 2 - 3$

Evaluation

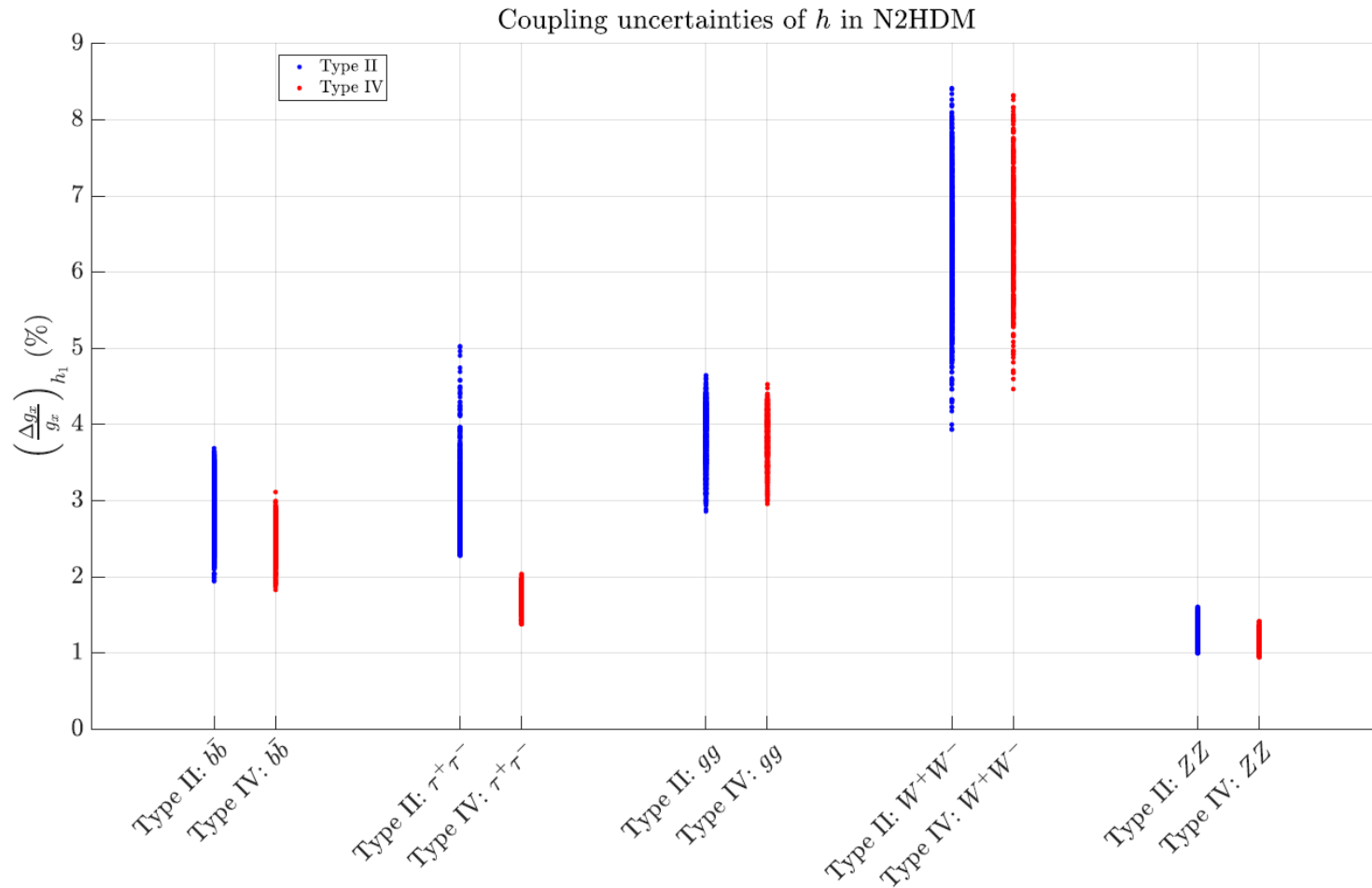
Reproduction

Incertidumbre relativa del acoplo del bosón de Higgs h_1 (96 GeV) en el modelo N2HDM



Evaluation

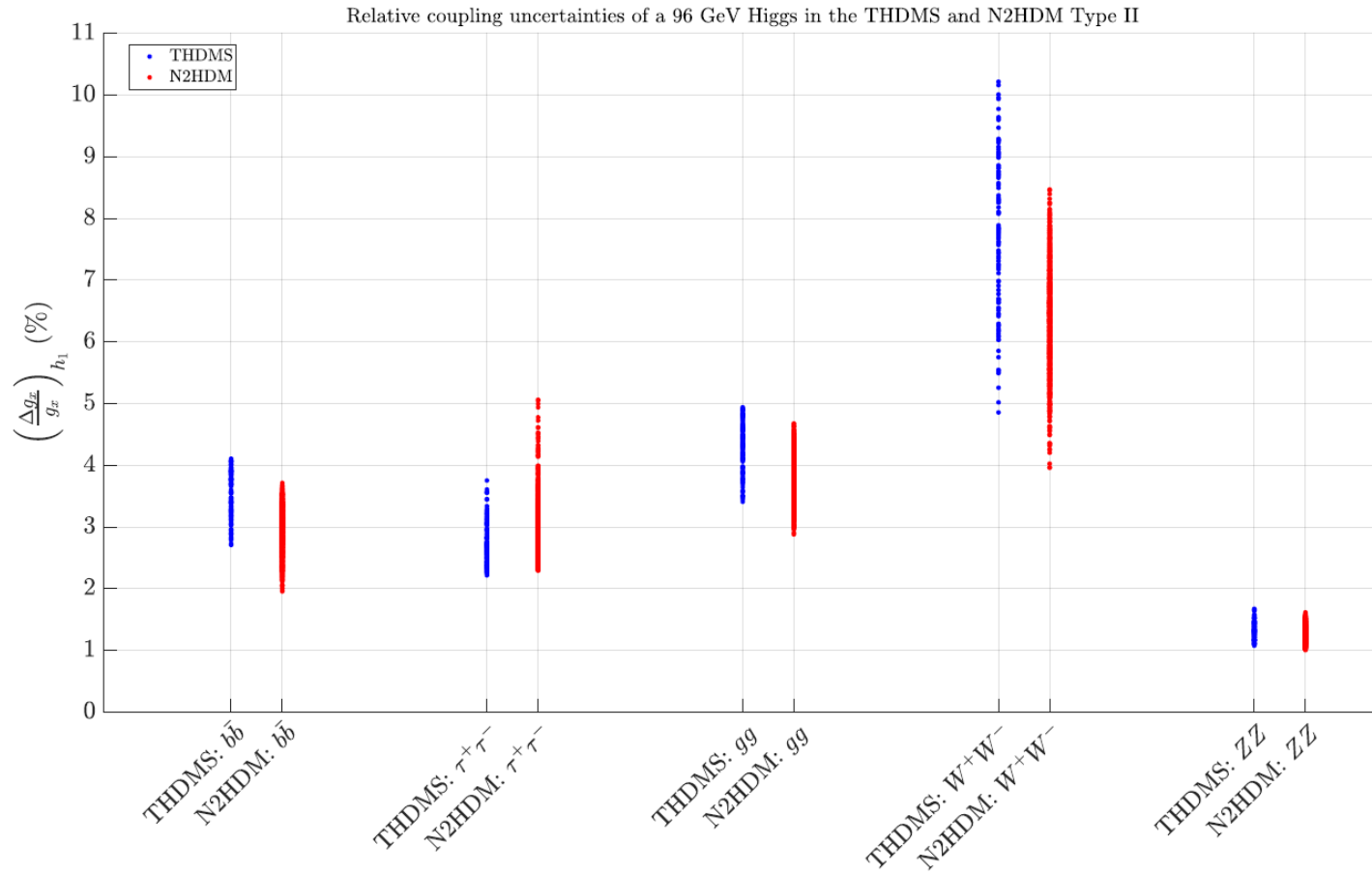
Reproduction



Evaluation

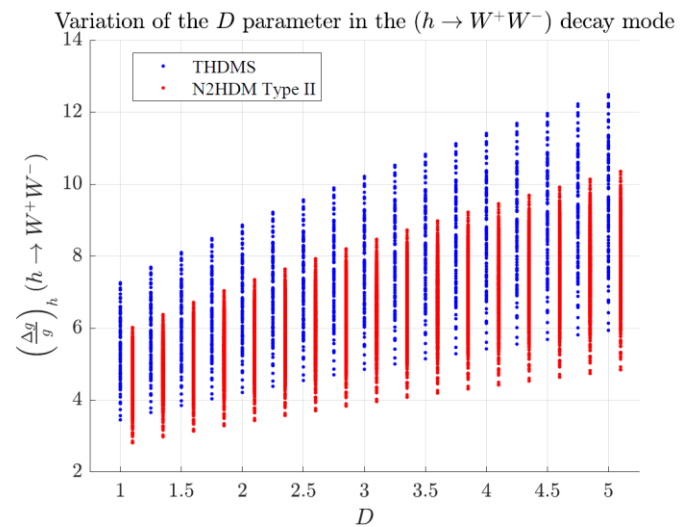
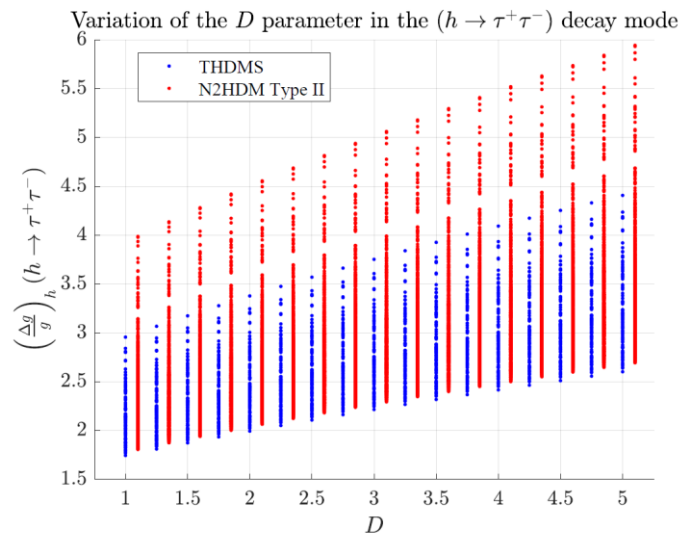
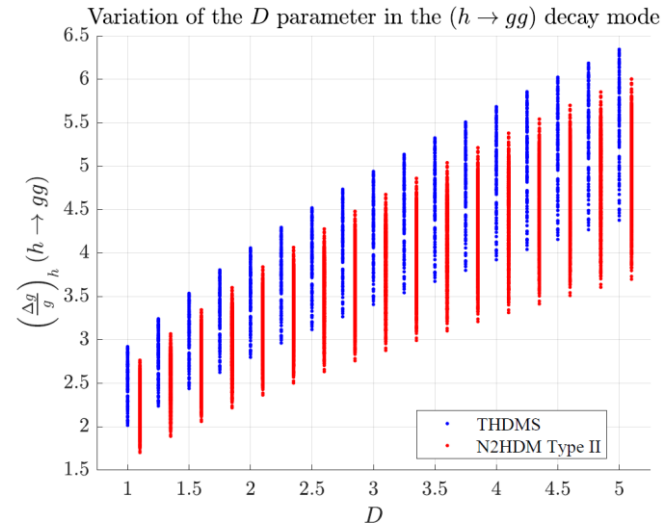
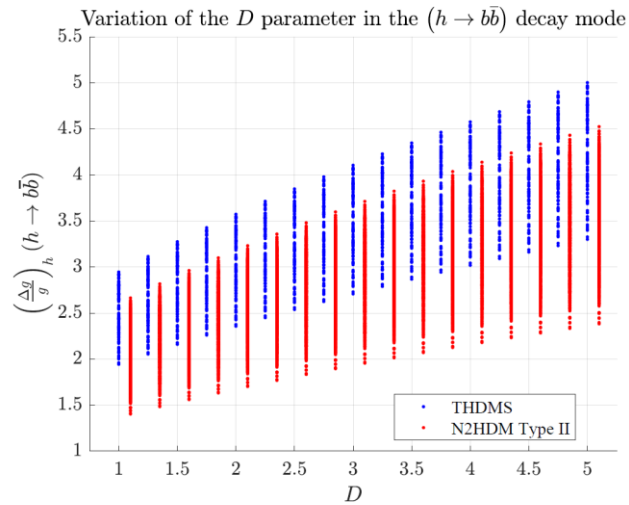
Results

$\left(\frac{\Delta g_x}{g_x}\right)_{h_1}$	range in THDMS [%]	range in N2HDM [%]
bb	(2.690, 4.08)	(1.966, 3.733)
$\tau^+\tau^-$	(2.200, 3.728)	(2.299, 5.075)
gg	(3.385, 4.907)	(2.906, 4.718)
W^+W^-	(4.824, 10.149)	(3.983, 8.518)
ZZ	(1.069, 1.663)	(0.997, 1.604)



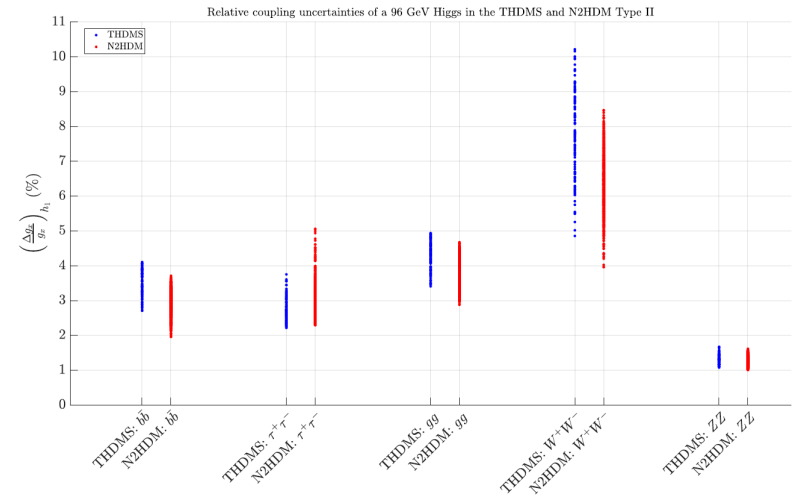
Evaluation

Dependence on D



Conclusion

- THDMS shows similar behaviour as the N2HDM
- THDMS shows slight tendency for higher uncertainties
- $\tan \beta$ value helps in differentiating between THDMS and N2HDM
- Fewer points compared to the N2HDM may shift the results

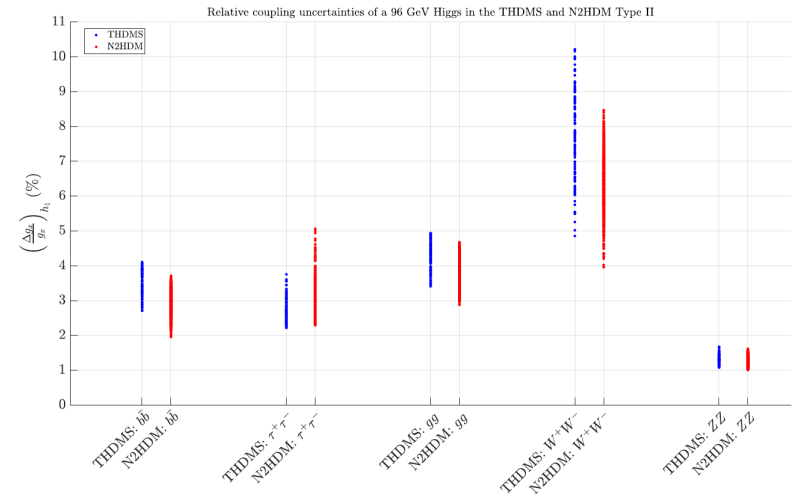


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Conclusion

Next steps

- More thorough scan
 - Greater ranges
 - Higher values
- Inclusion of vacuum stability constraints
- Find a way to clearly differentiate between THDMS and N2HDM



$\left(\frac{\Delta g_x}{g_x}\right)_{h_1}$	range in THDMS [%]	range in N2HDM [%]
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Thank you