



Trilinear Higgs Couplings in 2HDMS and N2HDM Type II at e^+e^- Colliders

Master Thesis Project: Mid-term Talk

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26.10.2023



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N2HDM (real singlet)

$$\phi_{1,2} = \begin{pmatrix} \chi_i^\pm \\ \frac{v_i + \rho_i + i\eta_i}{\sqrt{2}} \end{pmatrix}, \quad S = v_S + \rho_S$$

$$\sqrt{v_1^2 + v_2^2} = 246.22 \text{ GeV}$$

symmetries:

- $\mathbb{Z}_2$ $\phi_2 \rightarrow -\phi_2$
- $\mathbb{Z}'_2$ $S \rightarrow -S$

2HDMS($\mathbb{Z}_3$) (complex singlet)

$$\phi_{1,2} = \begin{pmatrix} \chi_i^\pm \\ v_i + \frac{\rho_i + i\eta_i}{\sqrt{2}} \end{pmatrix}, \quad S = v_S + \frac{\rho_S + i\eta_S}{\sqrt{2}}$$

$$\sqrt{v_1^2 + v_2^2} = 174 \text{ GeV} = \frac{v_{\text{SM}}}{\sqrt{2}}$$

symmetries:

- $\mathbb{Z}_2$ $\phi_2 \rightarrow -\phi_2$
- $\mathbb{Z}_3$ $\begin{pmatrix} \phi_2 \\ S \end{pmatrix} \rightarrow \begin{pmatrix} e^{i\frac{2\pi}{3}} & 0 \\ 0 & e^{-i\frac{2\pi}{3}} \end{pmatrix} \begin{pmatrix} \phi_2 \\ S \end{pmatrix}$

N2HDM

$$\begin{aligned}
 V = & \underbrace{m_{11}^2 \phi_1^\dagger \phi_1 + m_{22}^2 \phi_2^\dagger \phi_2}_{\text{same}} - \underbrace{m_{12}^2 (\phi_2^\dagger \phi_1 + \phi_1^\dagger \phi_2)}_{\text{same with h.c.}} \\
 & + \underbrace{\frac{\lambda_1}{2} (\phi_1^\dagger \phi_1)^2 + \frac{\lambda_2}{2} (\phi_2^\dagger \phi_2)^2 + \lambda_3 (\phi_1^\dagger \phi_1) (\phi_2^\dagger \phi_2)}_{\text{same}} \\
 & + \underbrace{\lambda_4 (\phi_1^\dagger \phi_2) (\phi_2^\dagger \phi_1)}_{\text{same}} + \underbrace{\frac{\lambda_5}{2} ((\phi_1^\dagger \phi_2)^2 + (\phi_2^\dagger \phi_1)^2)}_{\text{same with h.c.}} \\
 & + \underbrace{\frac{m_S^2}{2} S^2 + \frac{\lambda_6}{8} S^4 + \frac{\lambda_7}{2} (\phi_1^\dagger \phi_1) S^2 + \frac{\lambda_8}{2} (\phi_2^\dagger \phi_2) S^2}_{\text{same}}
 \end{aligned}$$

same

same with h.c.

2HDMs

$$\begin{aligned}
 V = & \underbrace{m_{11}^2 \phi_1^\dagger \phi_1 + m_{22}^2 \phi_2^\dagger \phi_2}_{\text{same}} - \underbrace{m_{12}^2 (\phi_2^\dagger \phi_1 + \phi_1^\dagger \phi_2)}_{\text{same with h.c.}} \\
 & + \underbrace{\frac{\lambda_1}{2} (\phi_1^\dagger \phi_1)^2 + \frac{\lambda_2}{2} (\phi_2^\dagger \phi_2)^2 + \lambda_3 (\phi_1^\dagger \phi_1) (\phi_2^\dagger \phi_2)}_{\text{same}} \\
 & + \underbrace{\lambda_4 (\phi_1^\dagger \phi_2) (\phi_2^\dagger \phi_1)}_{\text{same}} + \underbrace{m_S^2 (S^\dagger S)}_{\text{same}} + \underbrace{\lambda'_1 (S^\dagger S) (\phi_1^\dagger \phi_1)}_{\text{same}} \\
 & + \underbrace{\lambda'_2 (S^\dagger S) (\phi_2^\dagger \phi_2)}_{\text{same}} + \underbrace{\lambda''_3 (S^\dagger S)^2}_{\text{same}} + \underbrace{\frac{\mu_{S1}}{6} (S^3 + S^{\dagger 3})}_{\text{same with h.c.}} \\
 & + \underbrace{\mu_{12} (S \phi_1^\dagger \phi_2 + S^\dagger \phi_2^\dagger \phi_1)}_{\text{same with h.c.}}
 \end{aligned}$$

softly breaks $\mathbb{Z}_2$

breaks $\mathbb{Z}_3 | \mathbb{Z}'_2$

N2HDM (11 dof)

- $\tan \beta = \frac{v_2}{v_1}$
- scalar: $m_{h_1}, m_{h_2}, m_{h_3}$
- pseudoscalar: m_A
- charged: $m_{h^\pm}$
- mixing: $\alpha_1, \alpha_2, \alpha_3$
- m_{12}
- v_S

basis change

$$\lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda_5, \lambda_6, \lambda_7, \lambda_8, \tan \beta, v_S, m_{12}$$

2HDMS (12 dof)

- $\tan \beta = \frac{v_2}{v_1}$
- scalar: $m_{h_1}, m_{h_2}, m_{h_3}$
- pseudoscalar: m_{a_1}, m_{a_2}
- charged: $m_{h^\pm}$
- mixing: $\alpha_1, \alpha_2, \alpha_3, \alpha_4$
- $v_S = v_S^{\text{N2HDM}} / \sqrt{2}$

basis change

$$\lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda'_1, \lambda'_2, \lambda''_3, \mu_{12}, \mu_{s_1}, \tan \beta, v_S, m_{12}$$

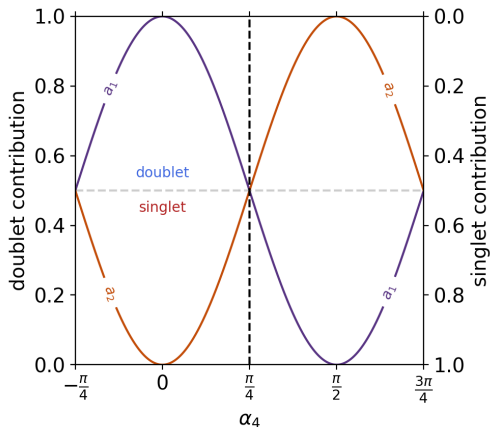


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Pseudo Scalar Mass Hierarchy

2 N2HDM Limit of the 2HDMS



$$\text{Doublet: } \begin{cases} \eta_1 = -a_1 \sin \beta \cos \alpha_4 + a_2 \dots \\ \eta_2 = a_1 \cos \beta \cos \alpha_4 + a_2 \dots \end{cases}$$

$$\text{Singlet: } \eta_S = a_1 \sin \alpha_4 + a_2 \dots$$

$$m_a^{\text{doublet}} = m_A^{N2HDM}$$

$$\alpha_4 \leq \frac{\pi}{4} :$$

$$m_{a_1} = m_a^{\text{doublet}}$$

$$m_{a_2} = m_a^{\text{singlet}}$$

$$\alpha_4 \geq \frac{\pi}{4} :$$

$$m_{a_1} = m_a^{\text{singlet}}$$

$$m_{a_2} = m_a^{\text{doublet}}$$

$$\left. \begin{array}{l} \alpha_4 \rightarrow 0 \\ \alpha_4 \rightarrow \frac{\pi}{2} \end{array} \right\} 2\text{HDMS} \rightarrow \text{N2HDM}$$

$$\tilde{\mu}_{2\text{HDMS}}^2 = m_{a_1}^2 \cos^2 \alpha_4 + m_{a_2}^2 \sin^2 \alpha_4$$

$$\rightarrow m_{a_i}^2 = \hat{\mu}_{\text{N2HDM}}^2 = \frac{m_{12}^2}{\sin \beta \cos \beta}$$

N2HDM constraint:

$$\Rightarrow m_{12}^2 = m_A^2 \sin \beta \cos \beta$$

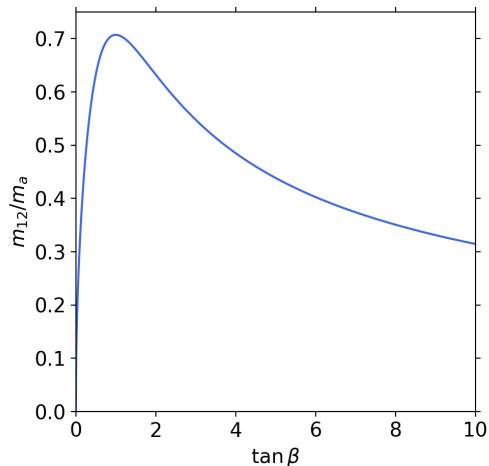
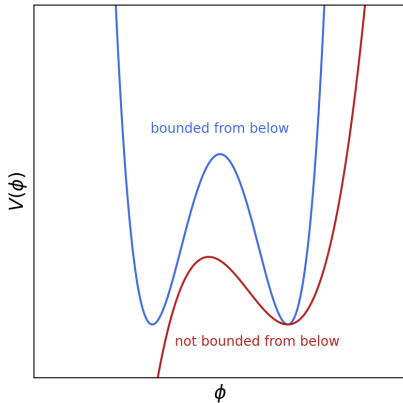




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- Tree level perturbative unitarity
 - $|\mathcal{M}| \leq 8\pi$
- Boundedness from below
 - $\phi_i, \mathcal{S} \rightarrow \pm\infty \Rightarrow V \rightarrow \infty$
- Conditions on $\lambda_1, \lambda_2, \lambda_3 \dots$
 - N2HDM [1]
 - 2HDMS [2]



- HiggsSignals[3]: h_2 agrees with measurements of 125 GeV Higgs boson (95 %CL)

$$\chi_{125;\text{SM}}^2 - \chi_{125;\text{BSM}}^2 \leq 5.99$$

- HiggsBounds[3]: exclusions on BSM searches (95 %CL)

- A hypothetical 95.4 GeV Higgs boson
- Fit experimental excesses (95 %CL):
 - ATLAS [4]

$$\mu_{\text{ATLAS}}^{\gamma\gamma} = 0.21 \pm 0.12$$

- CMS ¹ [5]

$$\mu_{\text{CMS}}^{\gamma\gamma} = 0.33_{-0.12}^{+0.19}$$

- LEP[6]

$$\mu_{\text{LEP}}^{b\bar{b}} = 0.117 \pm 0.057$$

¹ $\mu_{\text{CMS}}^{\tau\tau}$ also exists but not compatible with Type II



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Vector Boson and Yukawa Couplings

4 Type II Couplings

- $$R = \begin{pmatrix} c_{\alpha_1} c_{\alpha_2} & s_{\alpha_1} c_{\alpha_2} & s_{\alpha_2} \\ -(c_{\alpha_1} s_{\alpha_2} s_{\alpha_3} + s_{\alpha_1} c_{\alpha_3}) & -s_{\alpha_1} s_{\alpha_2} s_{\alpha_3} + c_{\alpha_1} c_{\alpha_3} & c_{\alpha_2} s_{\alpha_3} \\ -c_{\alpha_1} s_{\alpha_2} c_{\alpha_3} + s_{\alpha_1} s_{\alpha_3} & -(s_{\alpha_1} s_{\alpha_2} c_{\alpha_3} + c_{\alpha_1} s_{\alpha_3}) & c_{\alpha_2} c_{\alpha_3} \end{pmatrix}$$
- coupling ratio: $c_{h_i pp} = \frac{g_{h_i pp}}{g_{h_{SM} pp}}$

Vector Boson Coupling

- $c_{h_i VV} = R_{i1} \cos \beta + R_{i2} \sin \beta$
- alignment limit to SM
 - $c_{h_2 VV} \rightarrow 1$
 - $c_{h_{1,3} VV} \rightarrow 0$
 - $\beta - \alpha_1 - \text{sg}(\alpha_2) \alpha_3 \rightarrow \frac{\pi}{2}$

Yukawa Coupling

Type II: $t\bar{t} \sim \phi_2 \quad b\bar{b} \sim \phi_1 \quad \tau\bar{\tau} \sim \phi_1$

- $c_{h_i tt} = \frac{R_{i2}}{\sin \beta}$
- $c_{h_i bb} = \frac{R_{i1}}{\cos \beta}$
- $c_{h_i \tau\tau} = \frac{R_{i1}}{\cos \beta}$

Trilinear Higgs Couplings

4 Type II Couplings

- $g_{h_i h_j h_k} = \left. \frac{\partial^3 V}{\partial h_i \partial h_j \partial h_k} \right|_{h_{i,j,k}=0}$

- $\Lambda(a, b, c) = \sum_{\sigma \in S_3} R_{i\sigma(a)} R_{j\sigma(b)} R_{k\sigma(c)}$

$$\begin{aligned}
 g_{h_i h_j h_k} = & \underbrace{\left(m_i^2 + m_j^2 + m_k^2 \right) \left[\frac{R_{i1} R_{j1} R_{k1}}{v \cos \beta} + \frac{R_{i2} R_{j2} R_{k2}}{v \sin \beta} + \frac{R_{i3} R_{j3} R_{k3}}{v_S} \right]} \\
 & + \underbrace{\frac{\tilde{\mu}^2}{v} \left[- \frac{3 R_{i1} R_{j1} R_{k1} \sin^2 \beta}{\cos \beta} - \frac{3 R_{i2} R_{j2} R_{k2} \cos^2 \beta}{\sin \beta} + \frac{1}{2} \cos \beta \Lambda(1, 2, 2) + \frac{1}{2} \sin \beta \Lambda(1, 1, 2) \right]} \\
 & + \underbrace{\left(\frac{m_{a_2}^2 - m_{a_1}^2}{v} \cos \alpha_4 \sin \alpha_4 \right) \left[\frac{5 R_{i3} R_{j3} R_{k3} v^2 \sin 2\beta}{3 v_S^2} - \frac{1}{2} \tan \beta \Lambda(1, 1, 3) - \frac{v \sin \beta}{2 v_S} \Lambda(1, 3, 3) \right]} \\
 & - \underbrace{\frac{1}{2} \cot \beta \Lambda(2, 2, 3) - \frac{v \cos \beta}{2 v_S} \Lambda(2, 3, 3) + \Lambda(1, 2, 3)} + \underbrace{\frac{R_{i3} R_{j3} R_{k3}}{3 v_S} \left(m_{a_1}^2 \sin^2 \alpha_4 + m_{a_2}^2 \cos^2 \alpha_4 \right)}
 \end{aligned}$$

same part

$\tilde{\mu}_{2\text{HDMS}} \neq \hat{\mu}_{\text{N2HDM}}$

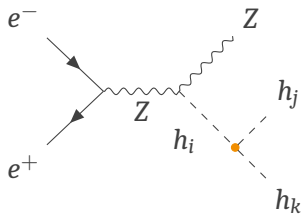
only in 2HDMS



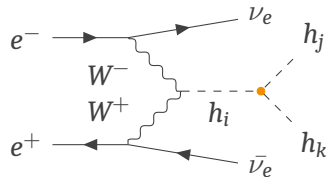
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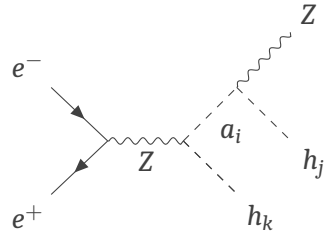
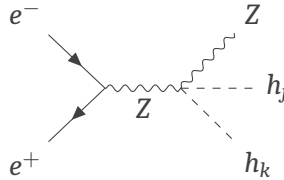
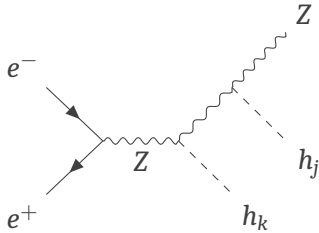
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Higgs Strahlung



Vector Boson Fusion





- $\sqrt{s} = 500 \text{ GeV}$
- $\mathcal{L}_{\text{int}} = 4 \text{ ab}^{-1}$
- Pol. $e^+ = +0.3$
- Pol. $e^- = -0.8$
- Expected $\sigma_{\text{SM}}(ZHH) = 0.232 \text{ fb}$

uncertainty estimation

- ILC₅₀₀ significance: 8σ at 2 ab^{-1} [7]
- $\text{unc} = \frac{\sigma_{\text{SM}}}{8} \sqrt{\frac{2}{4} \frac{\sigma_{\text{SM}}}{\sigma(Zh_i h_j)}}$
- $\text{dev} = \frac{|\sigma_{2\text{HDMS}} - \sigma_{\text{N2HDM}}|}{\text{unc}_{\text{N2HDM}}}$

Cross-sections with MadGraph5_aMC@NLO[8]



Analysis

α_4 Limits



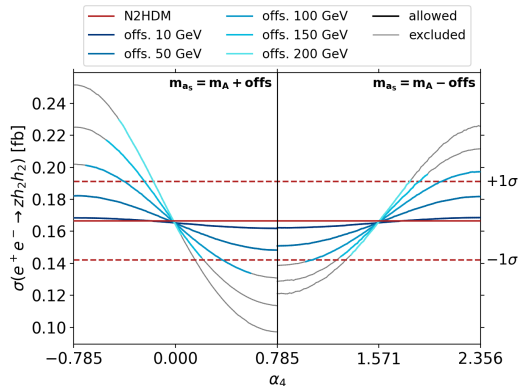
- differences only in α_4 and m_{A_S}
- hierarchy: $\alpha_4 \lesssim \pi/4$
- angles by coupling:

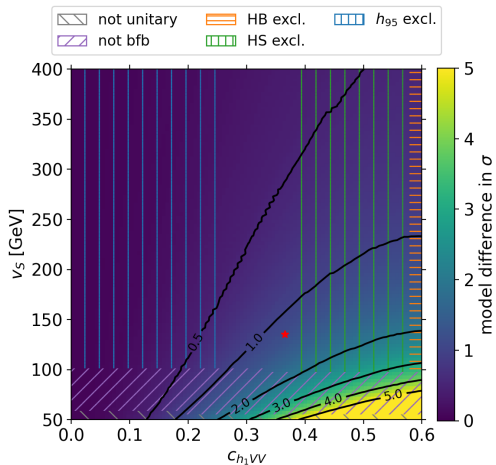
$$c_{h_1 VV}, \frac{c_{h_1 bb}}{c_{h_1 tt}}, \varepsilon \rightarrow \alpha_1, \alpha_2, \alpha_3$$

- alignment limit offset:

$$\beta - \alpha_1 - \text{sg}(\alpha_2)\alpha_3 = \frac{\pi}{2} - \varepsilon$$

⇒ 1σ limits on α_4 for $\tan \beta$ vs m_A



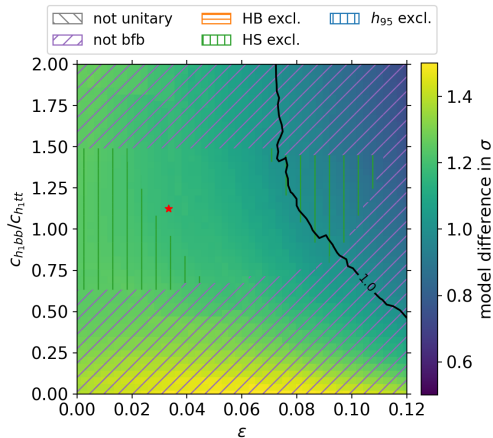


- set $\alpha_4 = 2.2$
- set $m_{a_1} = m_A - 120 \text{ GeV}$

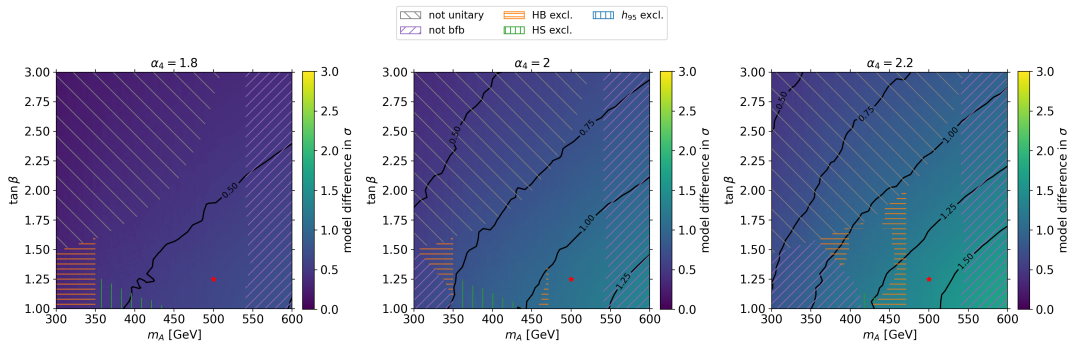


- $v_S \leq 150 \text{ GeV}$
- $c_{h_1 VV} \in (0.23, 0.4)$
- ★ benchmark point

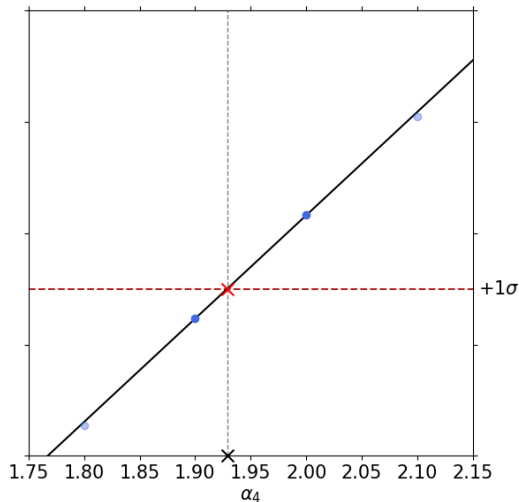
- small impact
- $\varepsilon \approx 0.05^1$
- $\frac{c_{h_1 bb}}{c_{h_1 tt}} \approx 1$



¹further investigation needed!



- precision vs running time
- 50 α_4 steps
- linear interpolation in between



$$m_{h_1} = 95.4 \text{ GeV}$$

$$m_{h_2} = 125.09 \text{ GeV}$$

$$m_{h_3} = 539.662 \text{ GeV}$$

$$m_{h^\pm} = 539.662 \text{ GeV}$$

$$v_S = 135.31 \text{ GeV}$$

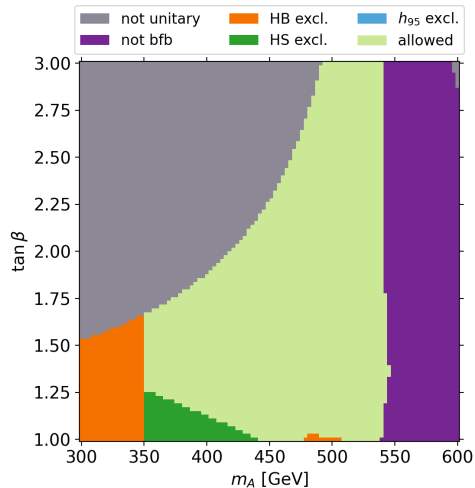
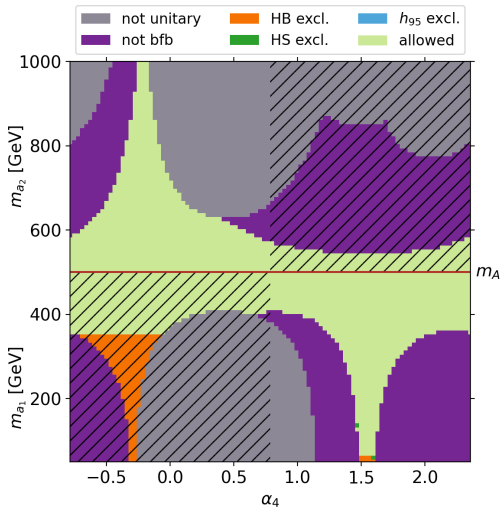
$$m_{a_S} = m_{A_D} \pm 120 \text{ GeV}$$

$$c_{h_1 VV} = 0.36$$

$$c_{h_1 bb}/c_{h_1 tt} = 1.124$$

$$\varepsilon = 0.033$$

$$\text{scan: } \begin{array}{ll} \tan \beta & \in (1, 3) \\ m_A & \in (300, 600) \text{ GeV} \end{array}$$



1 σ Limits $\alpha_4 \geq \pi/4$ ($m_A \geq m_{a_s}$)

6 Analysis

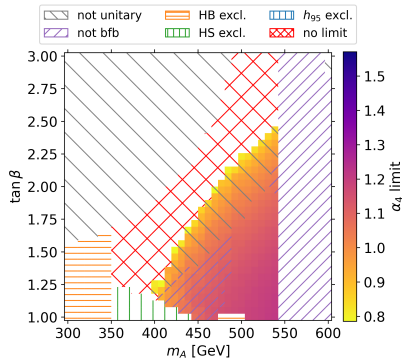


Figure: $\alpha_4 \in \left(\frac{\pi}{4}, \frac{\pi}{2}\right)$

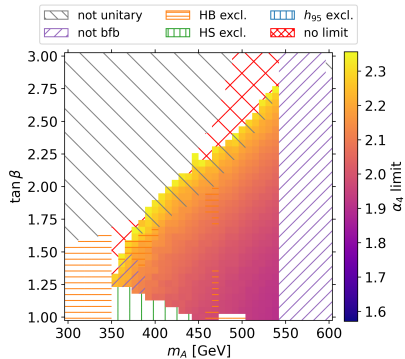


Figure: $\alpha_4 \in \left(\frac{\pi}{2}, \frac{3\pi}{4}\right)$

1σ Limits $\alpha_4 \leq \pi/4$ ($m_A \leq m_{a_s}$)

6 Analysis

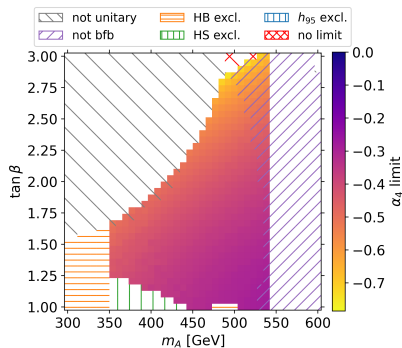


Figure: $\alpha_4 \in (-\frac{\pi}{4}, 0)$

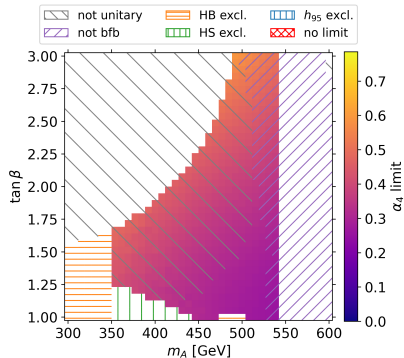


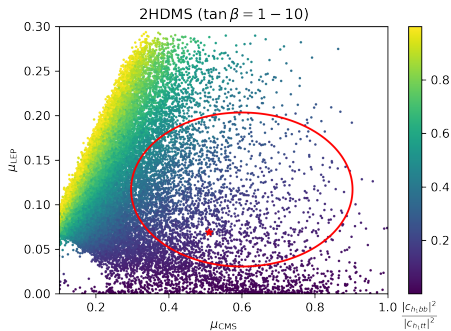
Figure: $\alpha_4 \in (0, \frac{\pi}{4})$



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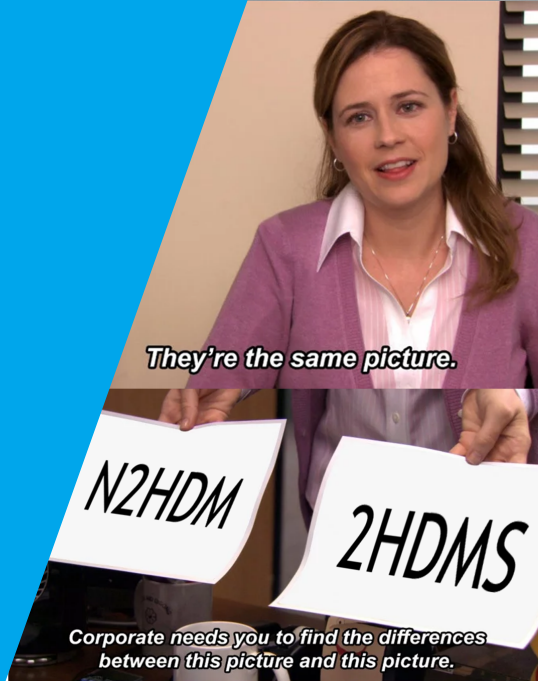
- Better choice of $c_{h_1 bb}/c_{h_1 tt}$
- mass limits for fixed α_4
- ML Limit enhancement?
- different final states ($h_1 h_2, h_1 h_1 \dots$)
- cross-check with different processes
 $pp \rightarrow t\bar{t}t\bar{t}$ from Cheng Li
- TeV scale m_{A_S} in decoupling limit



Plot taken from [2]

Any Questions?

- Daniel.Schieber@desy.de



- [1] Margarete Mühlleitner, Marco O. P. Sampaio, Rui Santos, and Jonas Wittbrodt.
The n2hdm under theoretical and experimental scrutiny.
[Journal of High Energy Physics](#), 2017(3), mar 2017.
- [2] S. Heinemeyer, C. Li, F. Lika, G. Moortgat-Pick, and S. Paasch.
A 96 gev higgs boson in the 2hdm plus singlet, 2021.
- [3] Henning Bahl, Thomas Biekötter, Sven Heinemeyer, Cheng Li, Steven Paasch, Georg Weiglein, and Jonas Wittbrodt.
HiggsTools: BSM scalar phenomenology with new versions of HiggsBounds and HiggsSignals.
[Computer Physics Communications](#), 291:108803, oct 2023.
- [4] Search for resonances in the 65 to 110 GeV diphoton invariant mass range using 80 fb^{-1} of pp collisions collected at $\sqrt{s} = 13 \text{ TeV}$ with the ATLAS detector.
2018.
- [5] Search for a standard model-like Higgs boson in the mass range between 70 and 110 GeV in the diphoton final state in proton-proton collisions at $\sqrt{s}=13 \text{ TeV}$.
2023.
- [6] Search for the standard model higgs boson at LEP.
[Physics Letters B](#), 565:61–75, jul 2003.
- [7] ILC-Collaboration.
The international linear collider: Report to snowmass 2021, 2023.
- [8] J. Alwall, R. Frederix, S. Frixione, V. Hirschi, F. Maltoni, O. Mattelaer, H.-S. Shao, T. Stelzer, P. Torrielli, and M. Zaro.
The automated computation of tree-level and next-to-leading order differential cross sections, and their matching to parton shower simulations.
[Journal of High Energy Physics](#), 2014(7), jul 2014.



Back up



scattering matrix unitary! Optical Theorem:

$$\mathbb{1} = S^\dagger S = (\mathbb{1} - iT^\dagger)(\mathbb{1} + iT) = \mathbb{1} + i(T - T^\dagger) + T^\dagger T \Rightarrow T^\dagger T = -i(T - T^\dagger) = \Im(T)$$

Decompose matrix element $\mathcal{M}$ with partial wave expansion:

$$\mathcal{M}(s, \theta) = 16\pi \sum_J (2J+1) a_J P_J(\cos \theta) \text{ with } |a_J|^2 = \Im(a_J) \Rightarrow \Re(a_J) \leq \frac{1}{2}$$

In high energy limit leading contribution $\mathcal{M}_0$ to tree level matrix element

$$\Rightarrow |\mathcal{M}| \leq 8\pi$$

$\mathcal{M}$ from interaction basis all possible initial and final states: $\chi_{1,2}^\pm, \rho_{1,2,S}, \eta_{1,2,S}$

Eigenvalues smaller than 8π

$\alpha_1, \alpha_2, \alpha_3$ calculation

8 Back up

$$\alpha_1 = \arctan \left(\frac{\tan \beta}{c_{h_1 bb}/c_{h_1 tt}} \right) \quad (1)$$

$$\alpha_2 = \arccos \left(\frac{c_{h_1 VV}}{\cos \beta \cos \alpha_1 + \sin \beta \sin \alpha_1} \right) \quad (2)$$

$$\alpha_3 = \text{sg}(\alpha_2) \left(\beta - \alpha_1 - \frac{\pi}{2} + \varepsilon \right) \quad (3)$$