



A model-independent analysis of interference effects in the $t\bar{t}$ final state at the LHC involving two CP-mixed Higgs bosons

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Outline

- **Introduction + Simplified model framework**
- **Summary of the Monte-Carlo implementation**
- **Results with mixing between the scalars**
- **(time permitting) Application to the C2HDM**

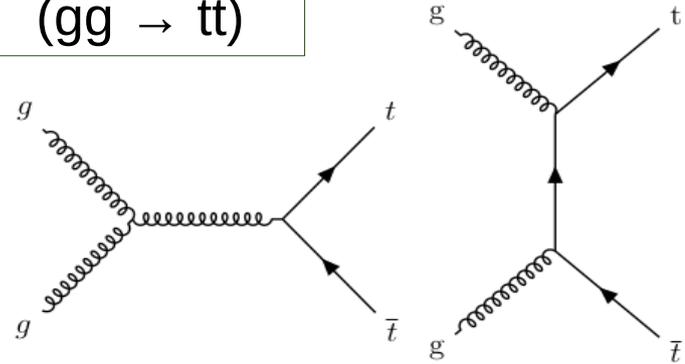
Di-top final state

- Total amplitude:

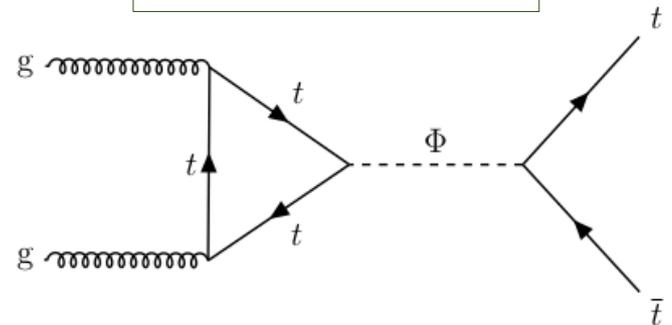
$$\mathcal{A} = \mathcal{A}(gg \rightarrow t\bar{t}) + \mathcal{A}(gg \rightarrow \Phi \rightarrow t\bar{t})$$

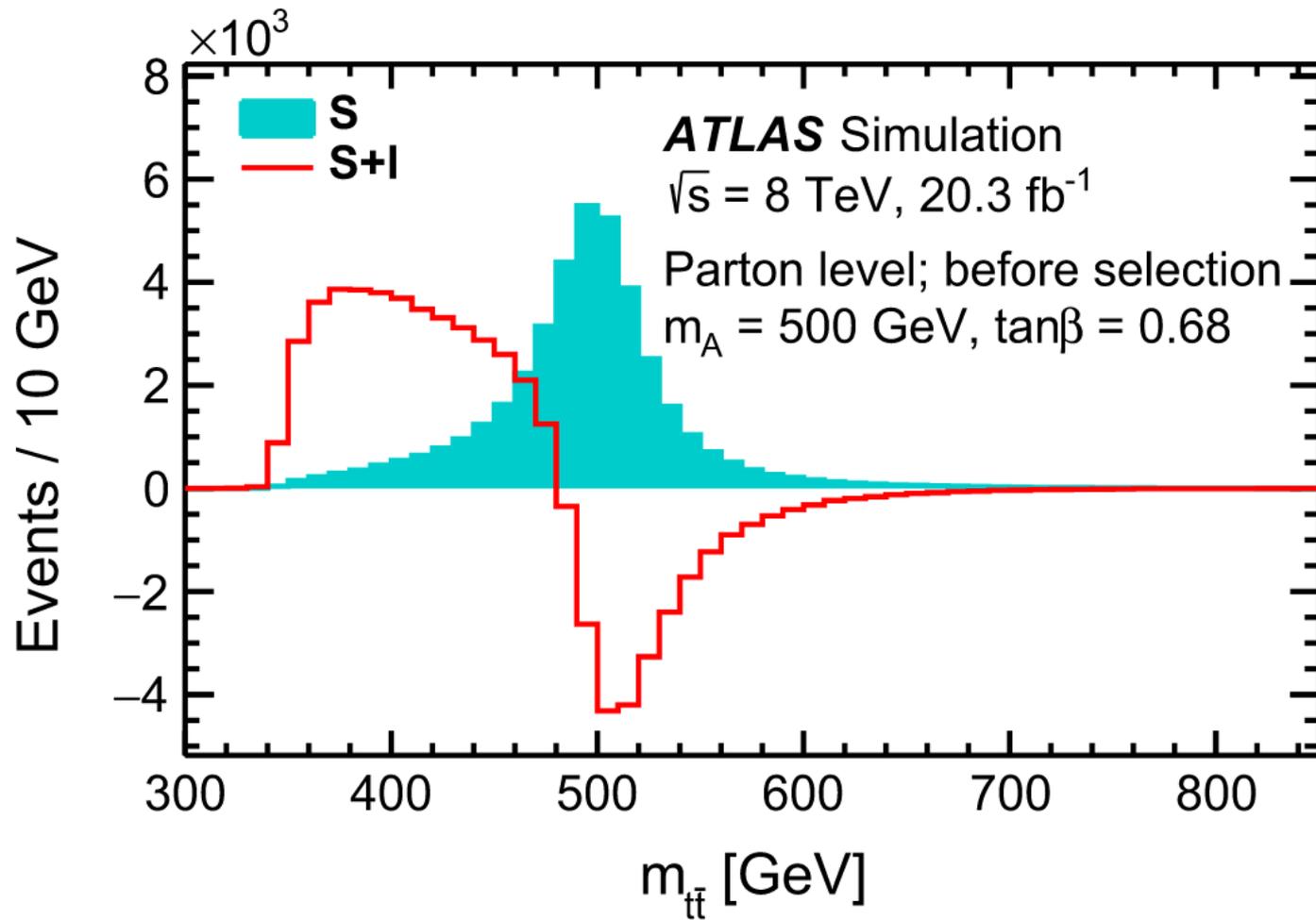
- Interference between the amplitudes in the cross-section!
- Signal-background interference: **large destructive contribution**
- Invariant mass distribution of the top quarks significantly distorted \rightarrow **peak-dip structure**

QCD
background
($gg \rightarrow t\bar{t}$)



Signal
($gg \rightarrow \Phi \rightarrow t\bar{t}$)





Simplified model framework

- Extended Higgs sector
- Consider two scalars (Φ_1, Φ_2) such that
 - massive than di-top threshold ($M_\Phi > 2m_t$)
 - produced via gluon fusion with top-triangle loop
 - CP-mixed character
 - decay to top quarks

- Analytical implementation (Mathematica)
- Monte-Carlo implementation (MadGraph 3.4.0)

$$\mathcal{L}_{\text{yuk}} = -\frac{y_t^{\text{SM}}}{\sqrt{2}} \bar{t} (c_t + i\gamma_5 \tilde{c}_t) t H$$

CP-even CP-odd

Yukawa-coupling modifiers



- Investigate impact of interference effects on the distribution profiles
- With two scalars: signal-signal interference possible
- Contribution of signal-signal interference
- Conditions? $|M_{H_1} - M_{H_2}| < \Gamma_{H_1} + \Gamma_{H_2}$
- Possible distribution profiles? Modifications to the usual peak-dip?
- Scalars mixing at loop-level \rightarrow Z-factors (complex numbers) \rightarrow modify distribution profiles?
- Followed by Monte-Carlo simulations

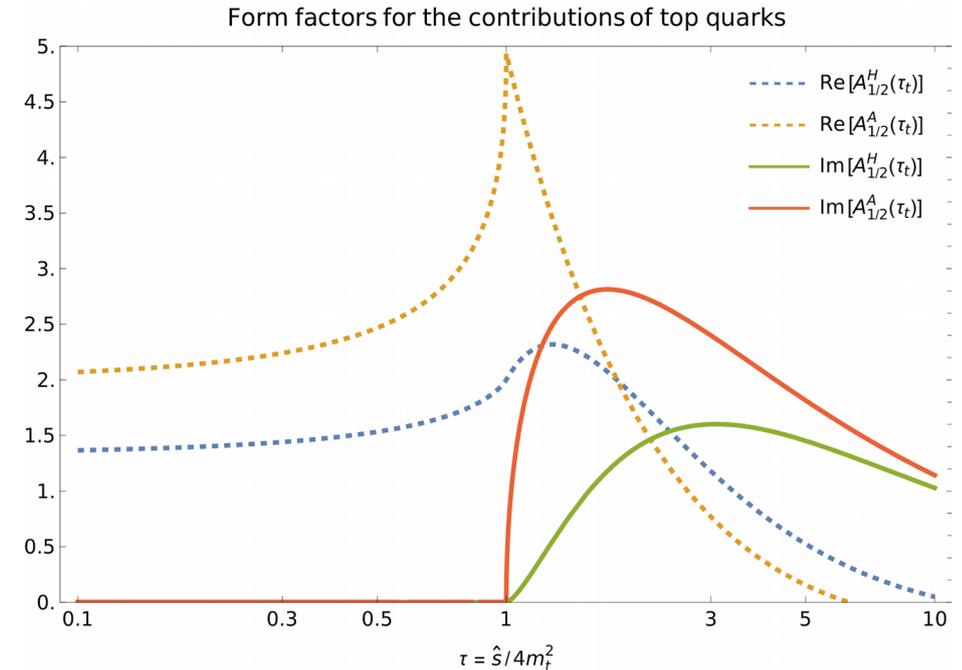


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- **Scalars mixing at loop-level → Z-factors (complex numbers) → modify distribution profiles?**
- **Followed by Monte-Carlo simulations**

Summary of the Monte-Carlo implementation

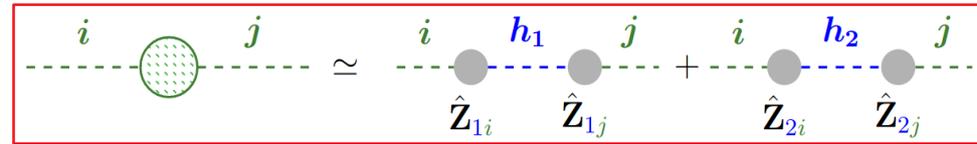
- Sizeable imaginary parts above the di-top threshold, effective coupling – poor approximation
- Incorporate the full top-triangle loop
- Two scalars, Yukawa-coupling modifiers, internal bookkeeping variables, and the Z-factors in the FeynRules file
- “hacked” python files in the FeynRules output files, Fortran routine for the top-loop



+ MadGraph caveats! (for e.g., internal parameters vs external parameters, etc)

Investigation with Z-factors

Allow the scalars to mix at the loop-level → Z-factors (complex numbers)
→ modify distribution profiles?



- Lowest-order interaction states H_1, H_2 → mix at loop-level → h_a, h_b loop-corrected mass eigenstates
- Following arxiv:1610.06193, for particles that mix, the total amplitude ($gg \rightarrow \Phi \rightarrow tt$) can be written using Z-factors and the Breit-Wigner propagators as

$$\mathcal{A} = \sum_{a=1,2} \left(\sum_{i=H_1, H_2} \hat{\mathbf{Z}}_{ai} \hat{\Gamma}_i^X \right) \Delta_a^{\text{BW}}(p^2) \left(\sum_{j=H_1, H_2} \hat{\mathbf{Z}}_{aj} \hat{\Gamma}_j^Y \right)$$

- $\hat{\mathbf{Z}}$ -matrix **complex elements!** [finite wave function normalisation factor (proper normalisation of the S-matrix)]

- In our case, we have,

$$\hat{\Gamma}_{H_1}^{(X,Y)} \propto (c_{t,1} + i\gamma_5 \tilde{c}_{t,1})$$

$$\hat{\Gamma}_{H_2}^{(X,Y)} \propto (c_{t,2} + i\gamma_5 \tilde{c}_{t,2})$$



$$\begin{aligned} \mathcal{A} \propto & \left(\hat{\mathbf{Z}}_{h_a H_1}(c_{t,1} + i\gamma_5 \tilde{c}_{t,1}) + \hat{\mathbf{Z}}_{h_a H_2}(c_{t,2} + i\gamma_5 \tilde{c}_{t,2}) \right) \Delta_{h_a}^{\text{BW}}(p^2) \\ & \left(\hat{\mathbf{Z}}_{h_a H_1}(c_{t,1} + i\gamma_5 \tilde{c}_{t,1}) + \hat{\mathbf{Z}}_{h_a H_2}(c_{t,2} + i\gamma_5 \tilde{c}_{t,2}) \right) \\ & + \left(\hat{\mathbf{Z}}_{h_b H_1}(c_{t,1} + i\gamma_5 \tilde{c}_{t,1}) + \hat{\mathbf{Z}}_{h_b H_2}(c_{t,2} + i\gamma_5 \tilde{c}_{t,2}) \right) \Delta_{h_b}^{\text{BW}}(p^2) \\ & \left(\hat{\mathbf{Z}}_{h_b H_1}(c_{t,1} + i\gamma_5 \tilde{c}_{t,1}) + \hat{\mathbf{Z}}_{h_b H_2}(c_{t,2} + i\gamma_5 \tilde{c}_{t,2}) \right) \end{aligned}$$

- With

$$\hat{\mathbf{Z}}_{h_a H_1} = Z_{11}, \hat{\mathbf{Z}}_{h_a H_2} = Z_{12}, \hat{\mathbf{Z}}_{h_b H_1} = Z_{21}, \hat{\mathbf{Z}}_{h_b H_2} = Z_{22}$$

- The amplitude can be written as

$$A \propto \underbrace{\left(Z_{11}(c_{t,1} + i\gamma_5 \tilde{c}_{t,1}) + Z_{12}(c_{t,2} + i\gamma_5 \tilde{c}_{t,2}) \right)}_{\text{Production side}} \Delta_{h_a}^{\text{BW}}(p^2) \underbrace{\left(Z_{11}(c_{t,1} + i\gamma_5 \tilde{c}_{t,1}) + Z_{12}(c_{t,2} + i\gamma_5 \tilde{c}_{t,2}) \right)}_{\text{Decay side}} \\ + \underbrace{\left(Z_{21}(c_{t,1} + i\gamma_5 \tilde{c}_{t,1}) + Z_{22}(c_{t,2} + i\gamma_5 \tilde{c}_{t,2}) \right)}_{\text{Production side}} \Delta_{h_b}^{\text{BW}}(p^2) \underbrace{\left(Z_{21}(c_{t,1} + i\gamma_5 \tilde{c}_{t,1}) + Z_{22}(c_{t,2} + i\gamma_5 \tilde{c}_{t,2}) \right)}_{\text{Decay side}}$$

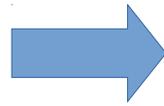
- Identifying

$$c_{t,1} \rightarrow Z_{11}c_{t,1} + Z_{12}c_{t,2}$$

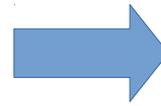
$$\tilde{c}_{t,1} \rightarrow Z_{11}\tilde{c}_{t,1} + Z_{12}\tilde{c}_{t,2}$$

$$c_{t,2} \rightarrow Z_{22}c_{t,2} + Z_{21}c_{t,1}$$

$$\tilde{c}_{t,2} \rightarrow Z_{22}\tilde{c}_{t,2} + Z_{21}\tilde{c}_{t,1}$$

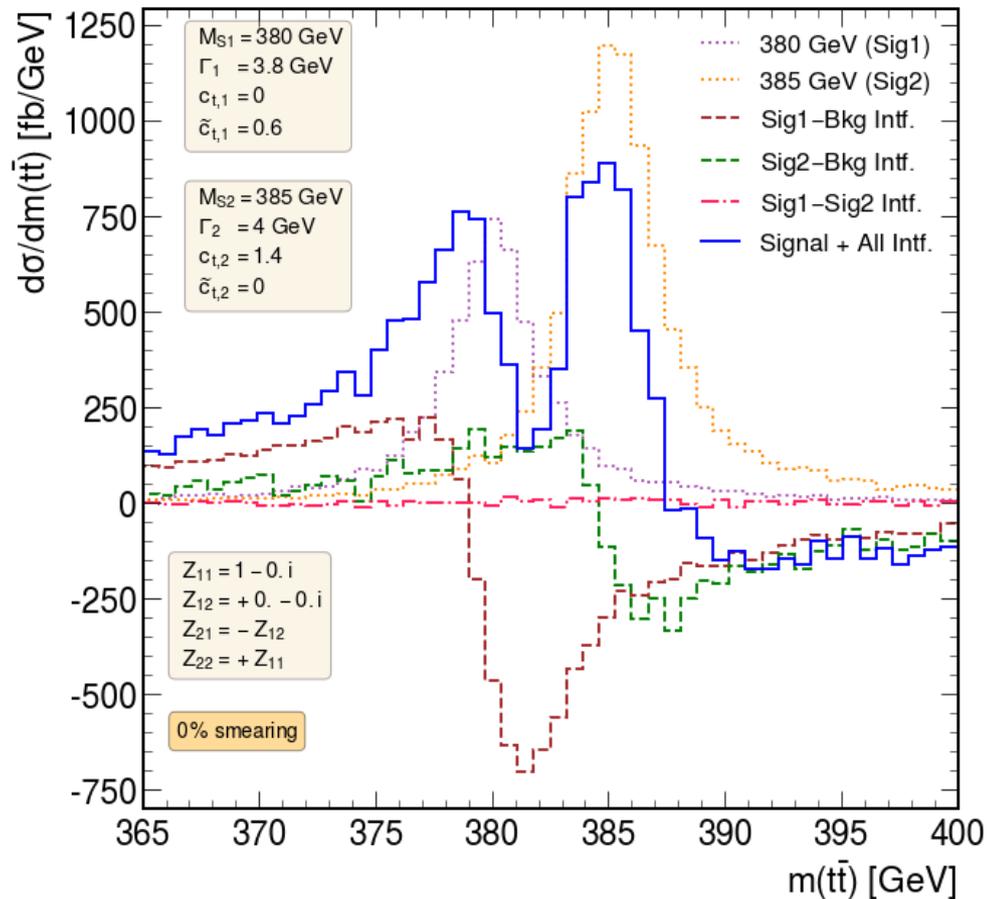


**Z-factors
can be
complex
numbers**

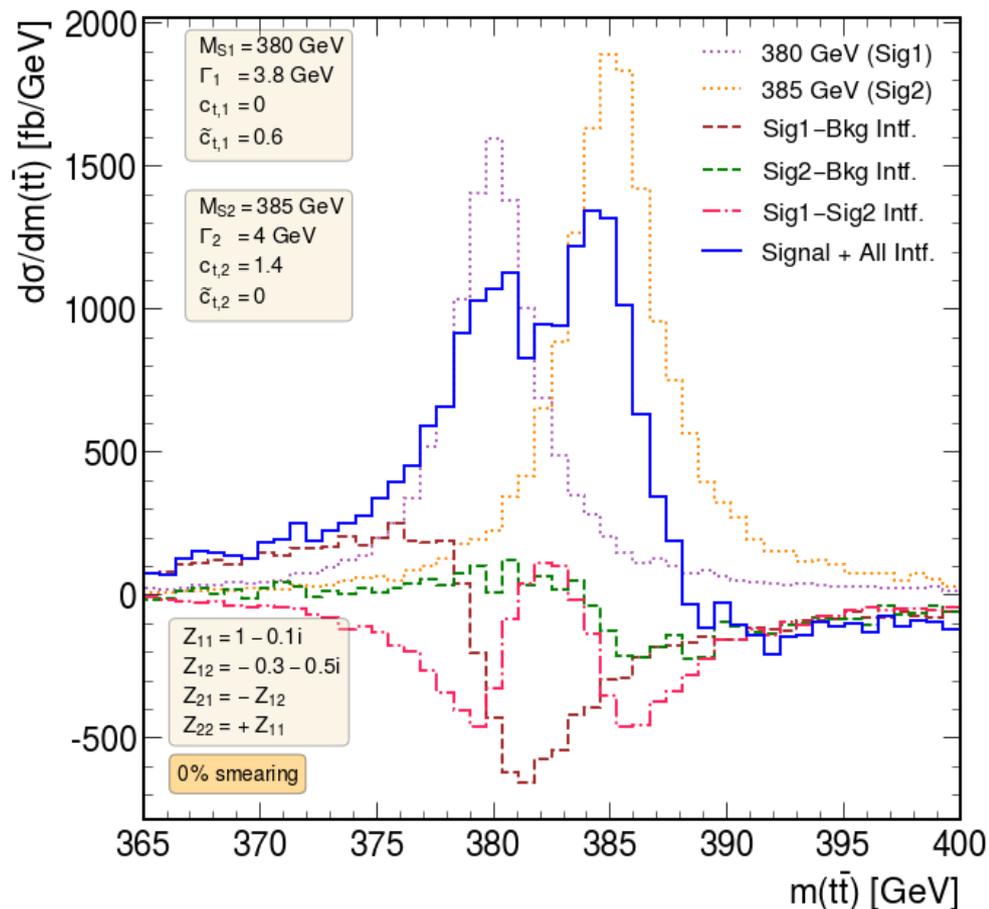


**Complex
Yukawa-
coupling
modifiers!**

**Additional
phases!**

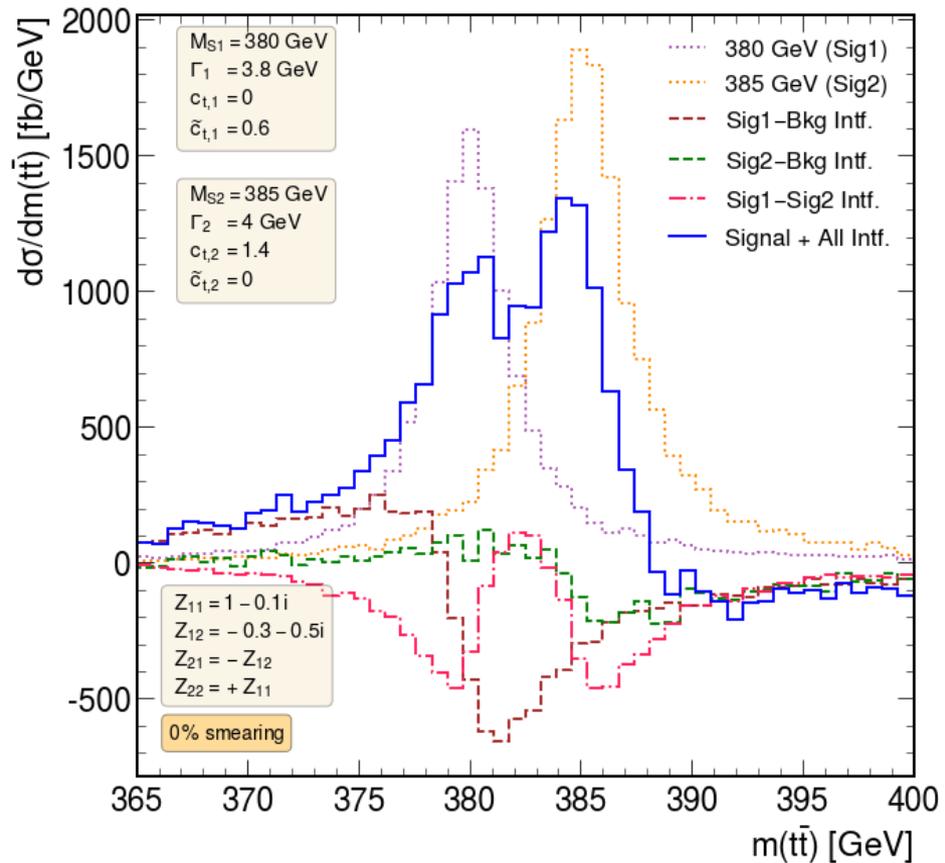
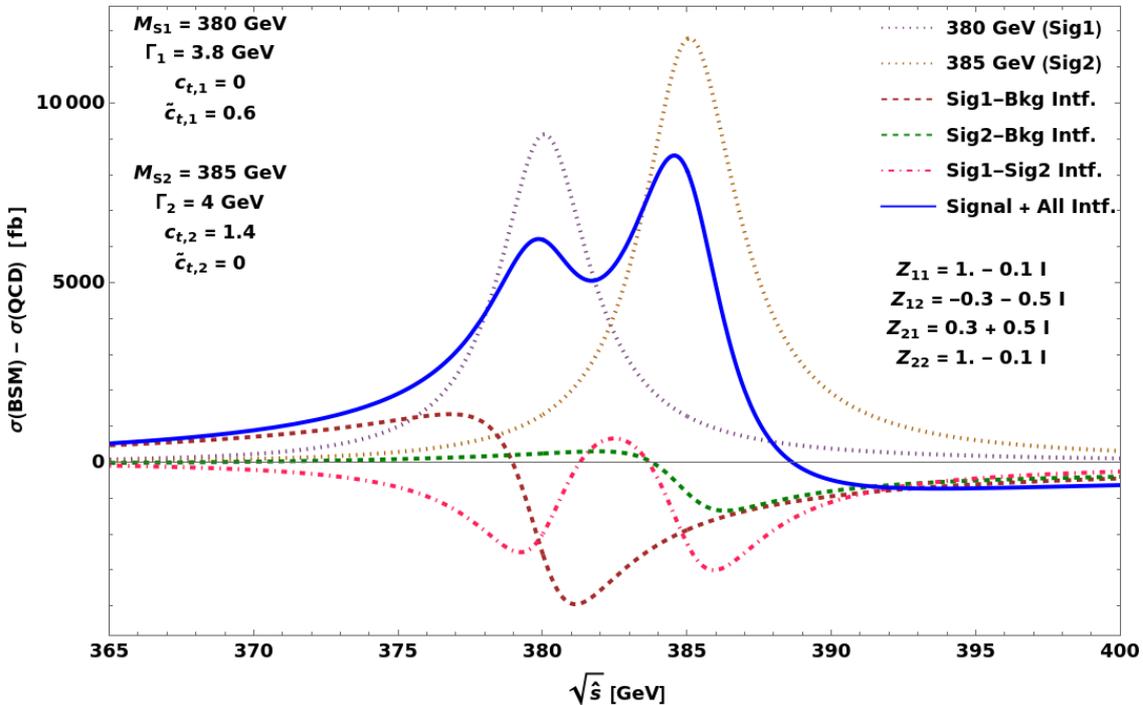


Without mixing

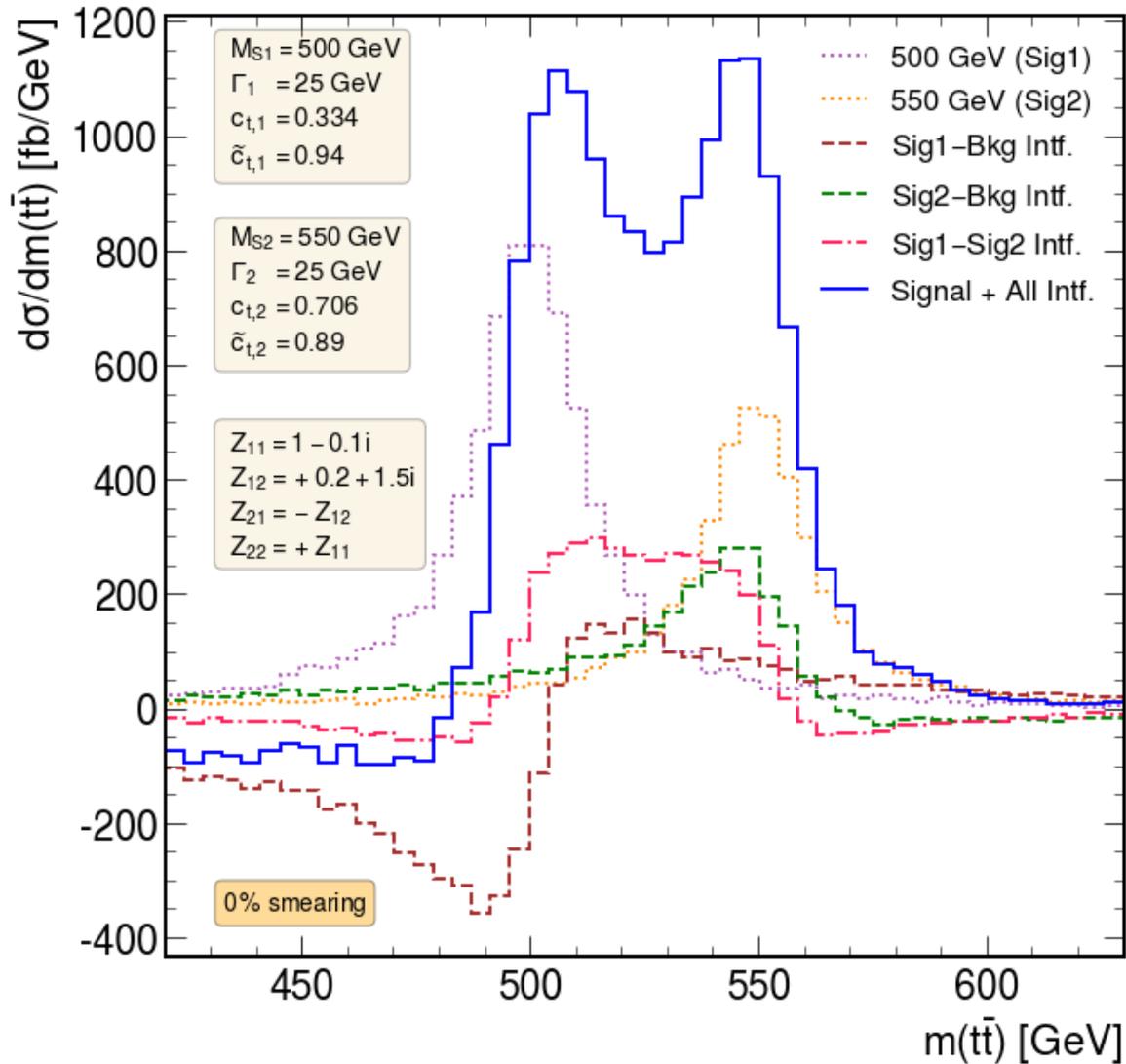


With mixing

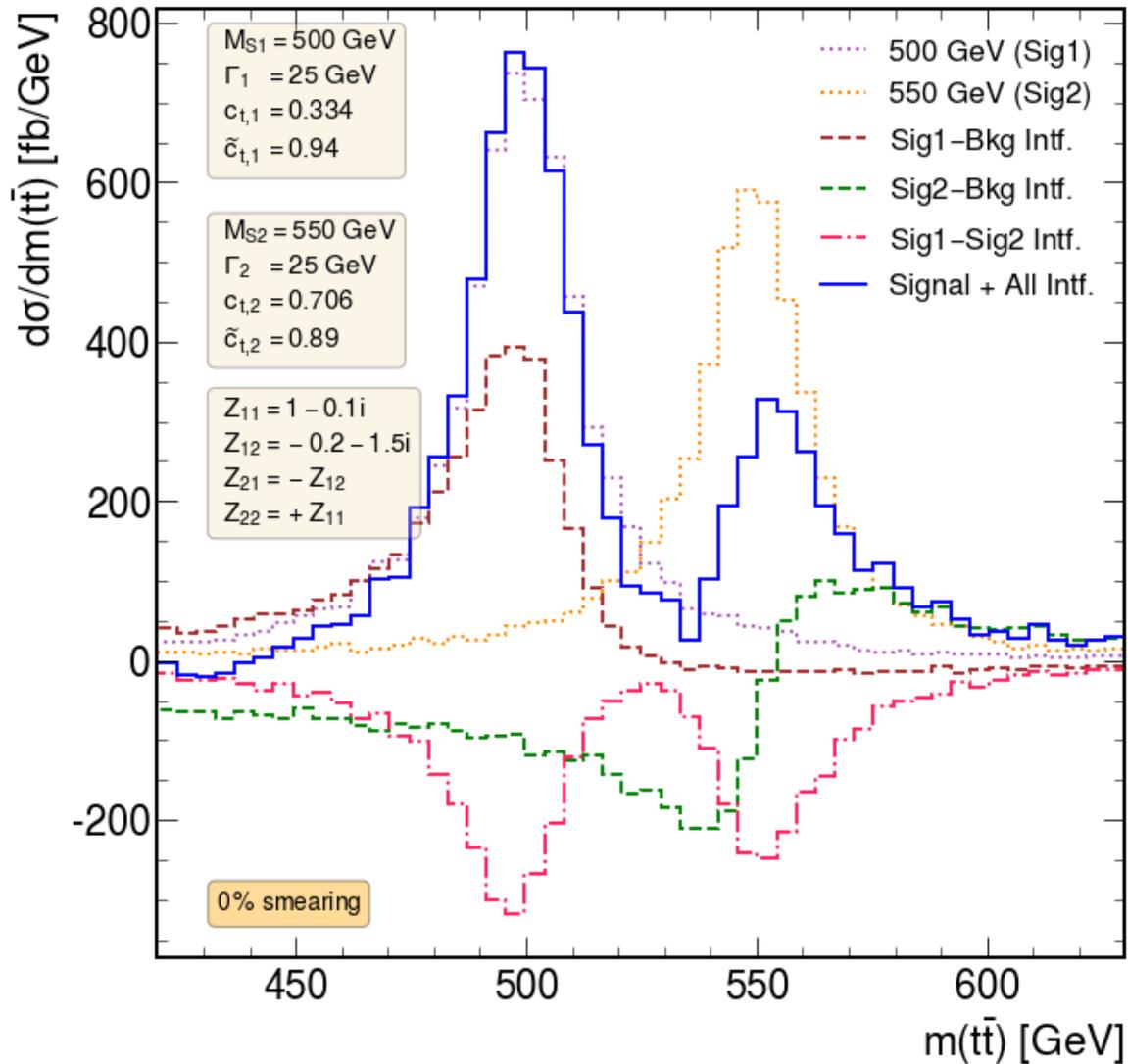
With mixing



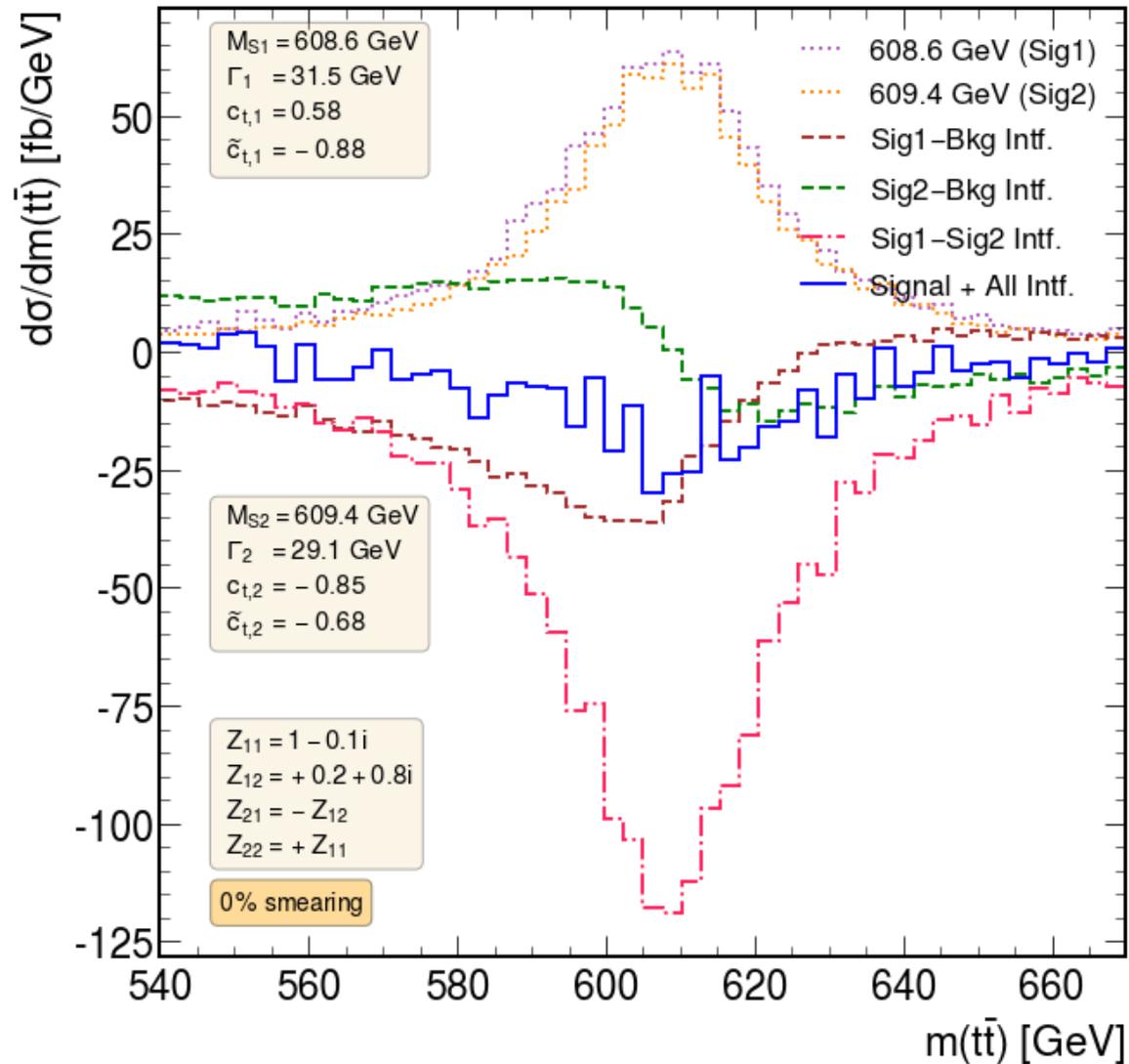
Good agreement between analytical and Monte-Carlo results



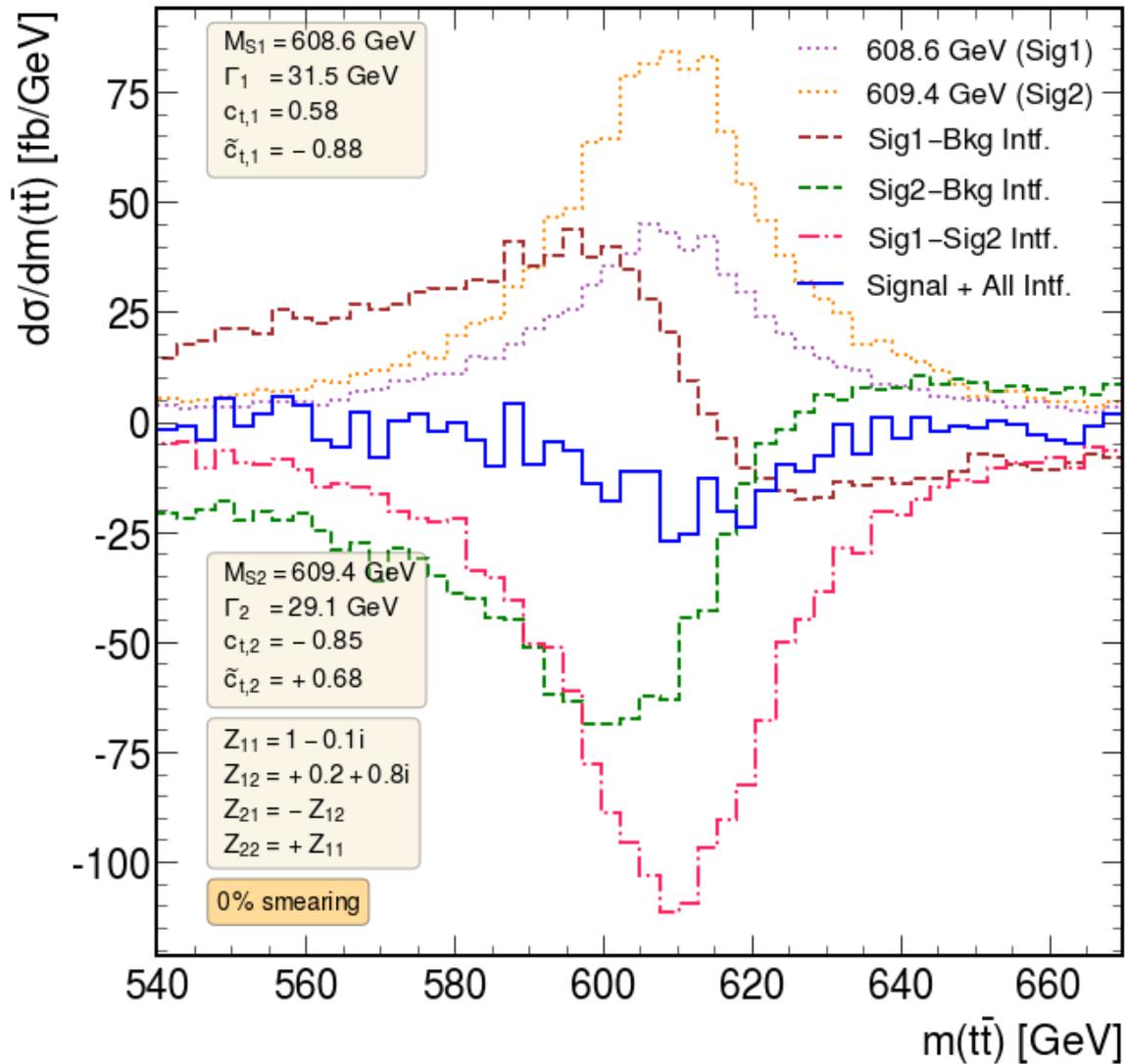
Large mixing between the scalars



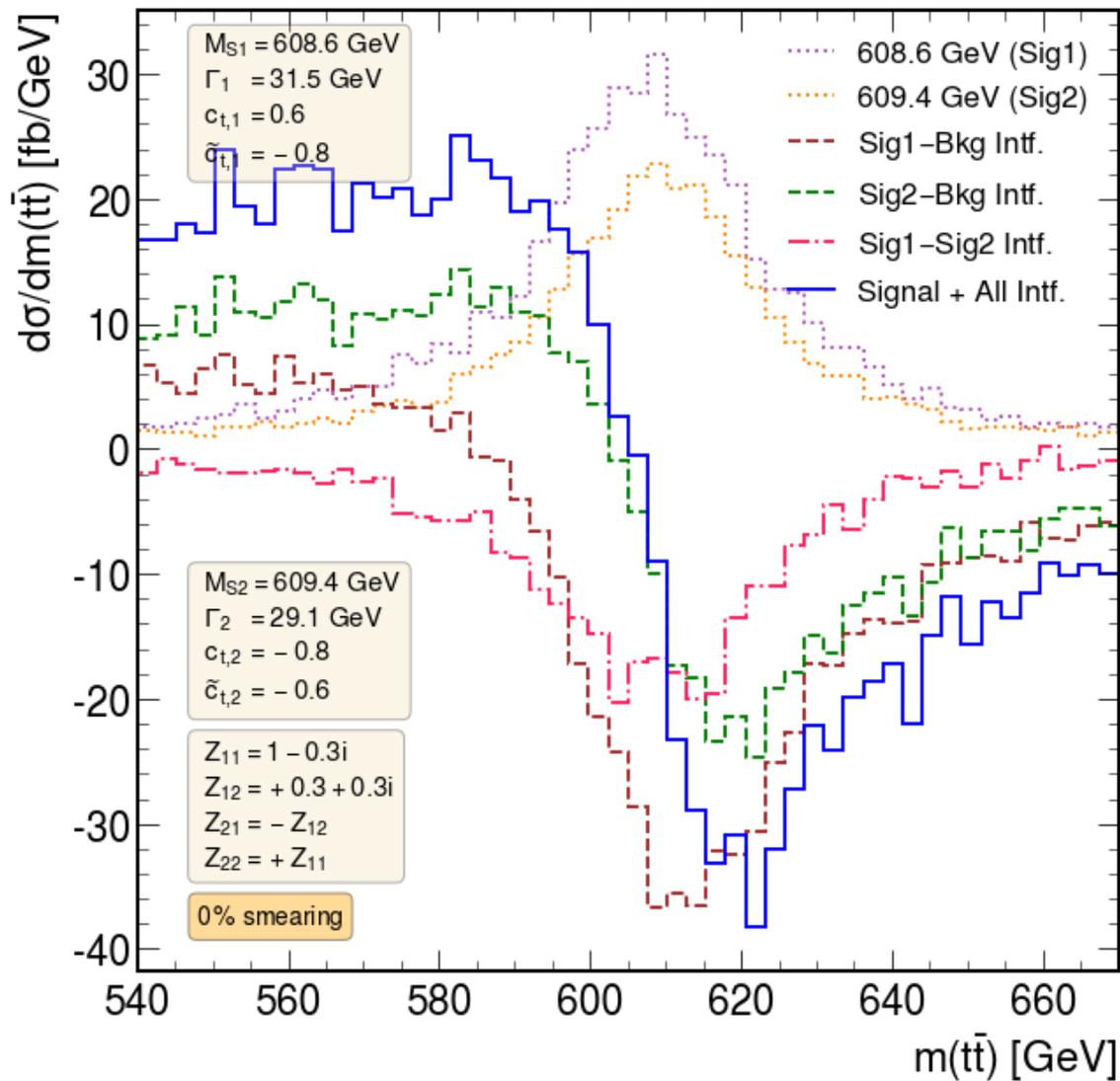
Large mixing between the scalars,
different sign for off-diagonal term



Large destructive signal-signal interference



ct2 flipped, different behaviour of signal-background interferences



Slightly smaller values of Yukawa-coupling modifiers and mixing



Takeaways!

- Monte-Carlo (and analytical) implementation of individual signal and interference contributions considering mixing between the scalars
- Mixing between scalars can lead to highly non-trivial distribution profiles, rich phenomenology
- One needs precise estimation of the Z-factors



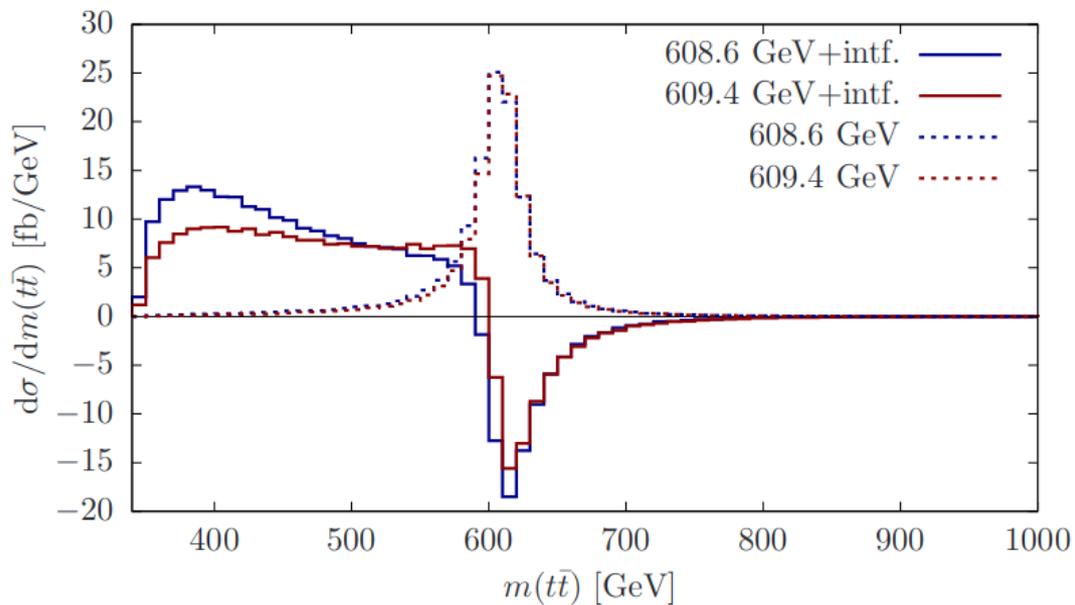
Backup/Extra slides

Application to the C2HDM

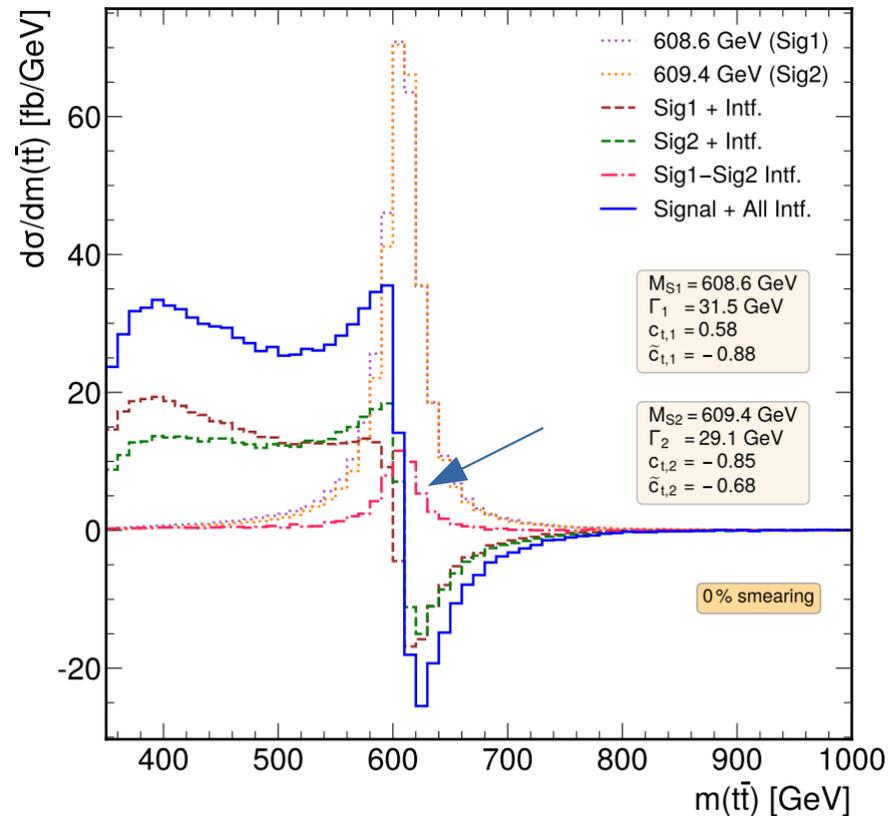
- Compare results with existing literature in arXiv:1909.09987v2
- 2HDM with a CP-violating scalar sector
- The Yukawa-coupling modifiers can be calculated using the elements of the rotation matrix that diagonalizes the 3x3 mass matrix to give a diagonal matrix with mass eigenstates
- We consider the lower-right 2x2 submatrix

$$\{R_{i,j}\} \equiv R = \begin{pmatrix} c_1 c_2 & s_1 c_2 & s_2 \\ -(c_1 s_2 s_3 + s_1 c_3) & c_1 c_3 - s_1 s_2 s_3 & c_2 s_3 \\ -c_1 s_2 c_3 + s_1 s_3 & -(c_1 s_3 + s_1 s_2 c_3) & c_2 c_3 \end{pmatrix}$$

$$\begin{aligned} \text{Type I : } & c_t = \frac{R_{i2}}{s_\beta}, \quad \tilde{c}_t = -i \frac{R_{i3}}{t_\beta} \\ \text{Type II : } & c_t = \frac{R_{i2}}{s_\beta}, \quad \tilde{c}_t = -i \frac{R_{i3}}{t_\beta} \end{aligned}$$

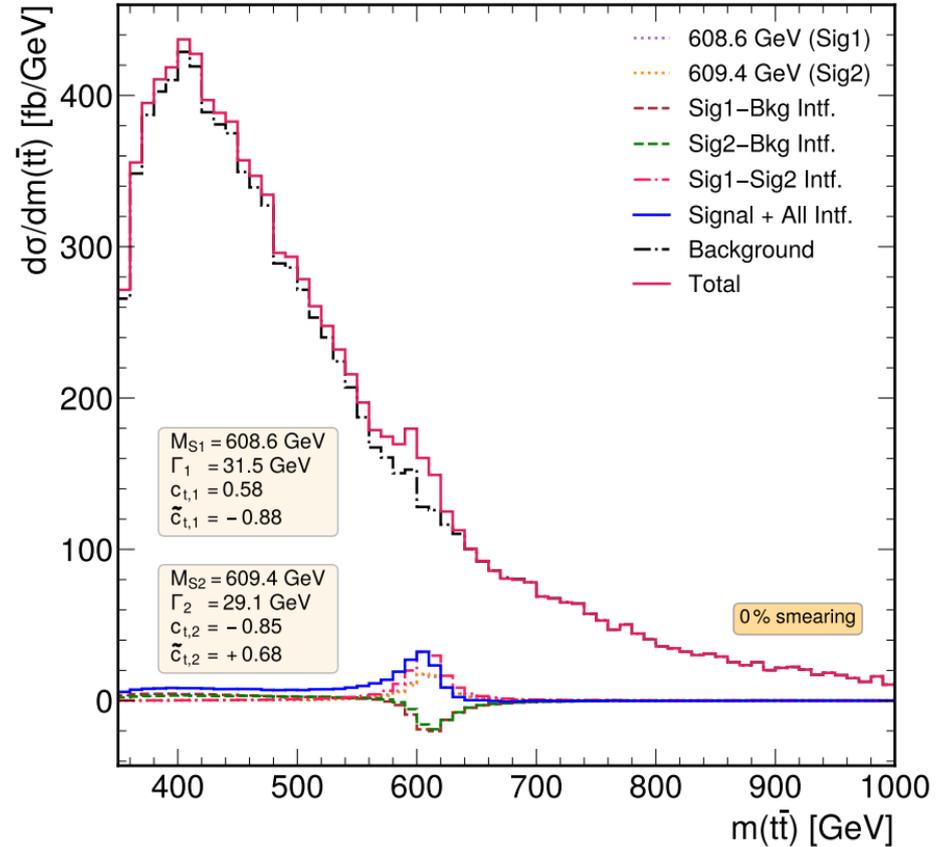
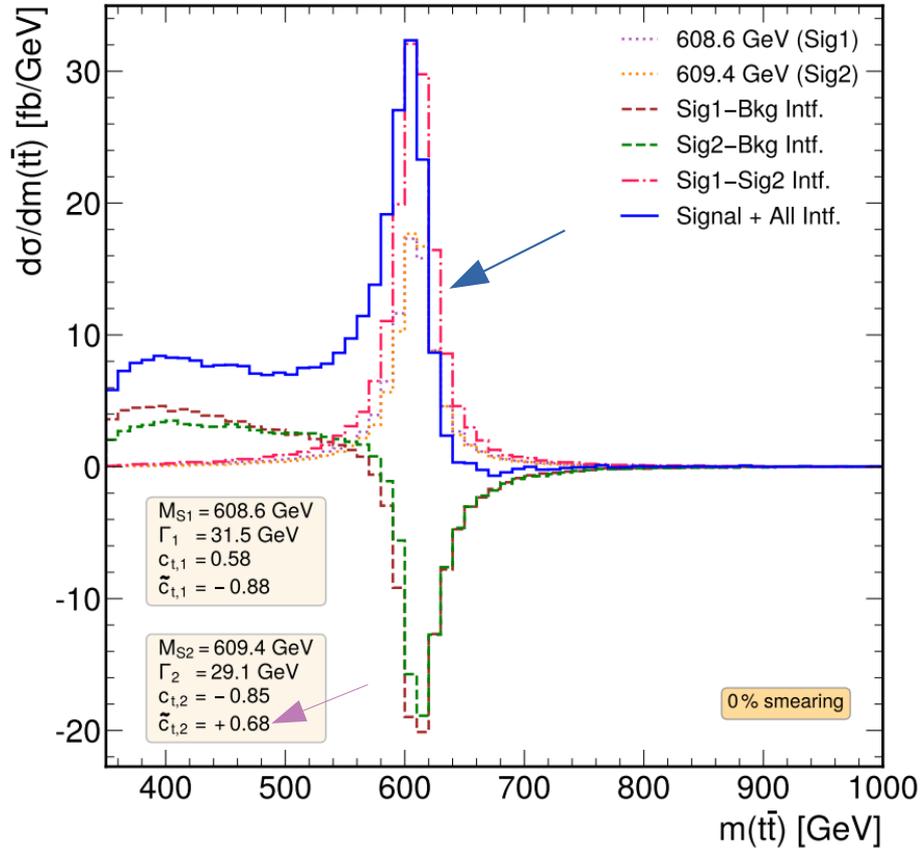


arXiv:1909.09987v2



- Signal-Signal interference can be significant (not considered previously)

Sign of ctt_2 flipped



- Signal-Signal could be as large as one of the pure signals!