



Domain Walls in the N2HDM: Exploring Vacuum Trapping and Scalar Potential Evolution

Master Thesis Project

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Vacuum Trapping

- Universe finds itself in a local vacuum state rather than the global one.
- Universe remains trapped in an EW symmetric phase down to $T \rightarrow 0$.

False vacuum
Metastable vacuum state

- Prevents EWSB
- Doublet VEVs remain zero
- Unphysical parameter points



Domain Walls

SSB

Discrete Symmetry

• Topological defects.

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- Act as Boundaries for distinct vacuum states.
- Inside the Domain Wall

Energy of Metastable state $> {\rm True} \; {\rm EW} \; {\rm Vacuum}$

- Problematic Dominate the energy of the universe at some point in time.
- Permitted → For approx. discrete symmetries
 If Annihilation occurs before energy domination





Motivation

- Vacuum Trapping Provides insights into the early Universe's phase structure.
- Domain Walls may seed the nucleation process during phase transitions, affecting vacuum decay.
- Studying vacuum trapping is essential to understand the correct evolution of the universe despite vacuum trapping.

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N2HDM = SM(ϕ_1) + Second Higgs Doublet(ϕ_2) + Real Scalar Singlet(ϕ_s) = 2HDM(ϕ_1, ϕ_2) + Real Scalar Singlet(ϕ_s)

 $\begin{aligned} \mathbf{V} &= m_{11}^2 |\Phi_1|^2 + m_{22}^2 |\Phi_2|^2 - m_{12}^2 (\Phi_1^{\dagger} \Phi_2 + h.c.) + \frac{\lambda_1}{2} (\Phi_1^{\dagger} \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^{\dagger} \Phi_2)^2 \\ &+ \lambda_3 (\Phi_1^{\dagger} \Phi_1) (\Phi_2^{\dagger} \Phi_2) + \lambda_4 (\Phi_1^{\dagger} \Phi_2) (\Phi_2^{\dagger} \Phi_1) + \frac{\lambda_5}{2} [(\Phi_1^{\dagger} \Phi_2)^2 + h.c.] \\ &\left(+ \frac{1}{2} m_5^2 \Phi_5^2 + \frac{\lambda_6}{8} \Phi_5^4 + \frac{\lambda_7}{2} (\Phi_1^{\dagger} \Phi_1) \Phi_5^2 + \frac{\lambda_8}{2} (\Phi_2^{\dagger} \Phi_2) \Phi_5^2 \right) \end{aligned}$

Symmetries: $Z_2: \phi_1 \to \phi_1, \phi_2 \to -\phi_2$ and $\phi_s \to \phi_s$, only softly broken by m_{12}^2 $Z'_2: \phi_1 \to \phi_1, \phi_2 \to \phi_2$ and $\phi_s \to -\phi_s$, spontaneously broken by v_s

Next-to 2 Higgs Doublet Model (N2HDM)

$$\bigvee m_{11}^{2} |\phi_{1}|^{2} + \frac{\lambda_{1}}{2} (\phi_{1}^{\dagger} \phi_{1})^{2} + m_{22}^{2} |\phi_{2}|^{2} + \frac{\lambda_{2}}{2} (\phi_{2}^{\dagger} \phi_{2})^{2} + \lambda_{3} (\phi_{1}^{\dagger} \phi_{1}) (\phi_{2}^{\dagger} \phi_{2}) + \lambda_{4} (\phi_{1}^{\dagger} \phi_{2}) (\phi_{2}^{\dagger} \phi_{1}) + \frac{\lambda_{5}}{2} [(\phi_{1}^{\dagger} \phi_{2})^{2} + h.c]$$

$$\frac{+1}{2} m_{s}^{2} \phi_{s}^{2} + \frac{\lambda_{6}}{8} \phi_{s}^{4} + \frac{\lambda_{7}}{2} (\phi_{1}^{\dagger} \phi_{1}) \phi_{s}^{2} + \frac{\lambda_{8}}{2} (\phi_{2}^{\dagger} \phi_{2}) \phi_{s}^{2}$$
Symmetries: $\phi_{2} \rightarrow -\phi_{2}$
 $\vdots \phi_{s} \rightarrow -\phi_{s}$

Extension of Z_2 to Yukawa sector \Rightarrow 4 types of the (N)2HDM

EW vacuum:

$$\left\langle \Phi_1 \right\rangle \left| \tau_{=0} = \frac{1}{\sqrt{2}} \left(\begin{array}{c} 0 \\ \boldsymbol{v}_1 \end{array} \right), \quad \left\langle \Phi_2 \right\rangle \left| \tau_{=0} = \frac{1}{\sqrt{2}} \left(\begin{array}{c} 0 \\ \boldsymbol{v}_2 \end{array} \right), \quad \left\langle \Phi_S \right\rangle |_{\mathcal{T}=0} = \boldsymbol{v}_s \right\rangle$$

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Physical Input Parameters:

- CP-even Higgs bosons
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- CP-odd Higgs bosons
- Charged Higgs bosons
- Soft breaking Z₂ term

- Singlet vev
- tan
- EW scale
- Mixing angles

125.09	400	[30, 1000]	650	2	1.10714	1.82*e-8	[-1.5, 1.5]	255 ²	[1, 1000]

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Evolution of VEVs

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125.09	400	550.25	650	2	1.10714	1.82e-8	0.50065	255 ²	544.63



For initialization of the fields, v_1 and v_2 were both taken as a gaussian function with a small value.



Evolution of VEVs

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125.09 400 550.25 650 2 1.10714 1.82e-8 0.50065 255² 544.63



Phase Tracer: Electroweak VEVs and Singlet VEV in dependence of the temperature

EW symmetry is broken at high temperatures and becomes unbroken at T \approx 500 GeV.

vs decreasing for T > 500 GeV, suggesting an EW-breaking but symm. conserving vacuum configuration.

650

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Evolution of VEVs 400 125.09 451.997 1000 1000 Phase 1 Phase 2 F OPT 800 800 800 Temperature 600 600 600 400 400 400 200 200 200 0 -200 -200 ò 200 400 ò 200 400 -200 Ó 200 400 **v**1 **v** 2 VS

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One of the Parameter points featuring Vacuum Trapping, obtained through ScannerS.

1.82e-8

-1.18487

255²

1.10715

EW symmetry is broken at high temperatures and becomes unbroken at T \approx 570 GeV.

Co-existence of two vacuum down to T=0.

184.019

Preventing Vacuum Trapping through Domain Walls

- Bounce Action and Tunneling Rate
- Reduced Bounce Action for tunneling
- Lesser Energy barrier for Nucleation between the metastable and true vacuum states.
- Smaller Critical radius for a stable bubble to form.
- In N2HDM, the tunneling, for field moving from the meta-stable state, occurs through the Domain wall to the EW vacuum.

$$S_{3} = 4\pi \int r^{2} \mathrm{d}r \, \left[\frac{1}{2} \left(\frac{\mathrm{d}\phi_{\mathrm{B}}}{\mathrm{d}r} \right)^{2} + V \left(\phi_{\mathrm{B}}, T \right) \right]$$

$$\Gamma(T) = A(T) e^{-S_3(T)/T}$$

 $A(T) \approx T^4 (S_3/2\pi T)^{3/2}$



S. Blasi and A. Mariotti, Domain Walls Seeding the Electroweak Phase Transition

Further Research: Finite Temperature effects

- The strategy is to look for parameters featuring vacuum trapping by finding parameter points where the symmetric and broken phases coexist at temperature T = 0 and no nucleation occurs.
- Calculating the bounce action inside the Domain Wall which leads to the tunneling probability.
- To check which parameter points can be rescued from Vacuum Trapping through Domain Walls.

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Summary and Outlook

• Vacuum Trapping in N2HDM

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• Breaking of discrete symmetry related to

singlet — Domain Walls

- DWs seed the nucleation process
- Evolution of VEVs with temperature
- Finite Temperature effects
- Calculating Bounce Action and Tunneling Probability



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doublet Mocielse Millio 2Esploring decyun 02pp (2020) Scalar Potential Evolution