



Midterm Talk

Parameter determination for dark matter scenarios in the MSSM

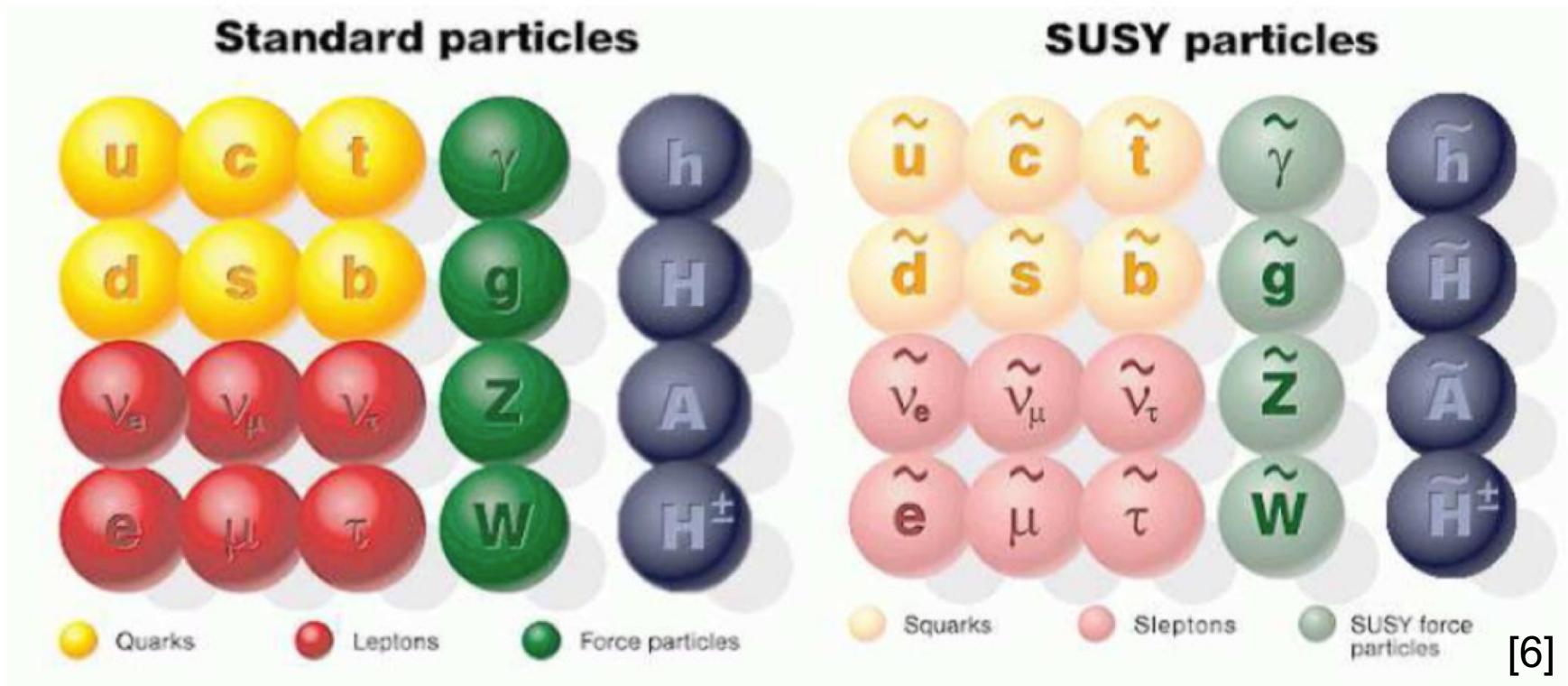
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MSSM

- SM-particles get supersymmetric partners
- Additionally:
 - Neutralinos: $\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$
 - Charginos: $\tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm$



Mass matrices in the MSSM

- Chargino and neutralino mass matrices
- We use the Bino-Wino basis $\{\tilde{B}, \tilde{W}^0, \tilde{H}_1^0, \tilde{H}_2^0\}$

- Parameters:

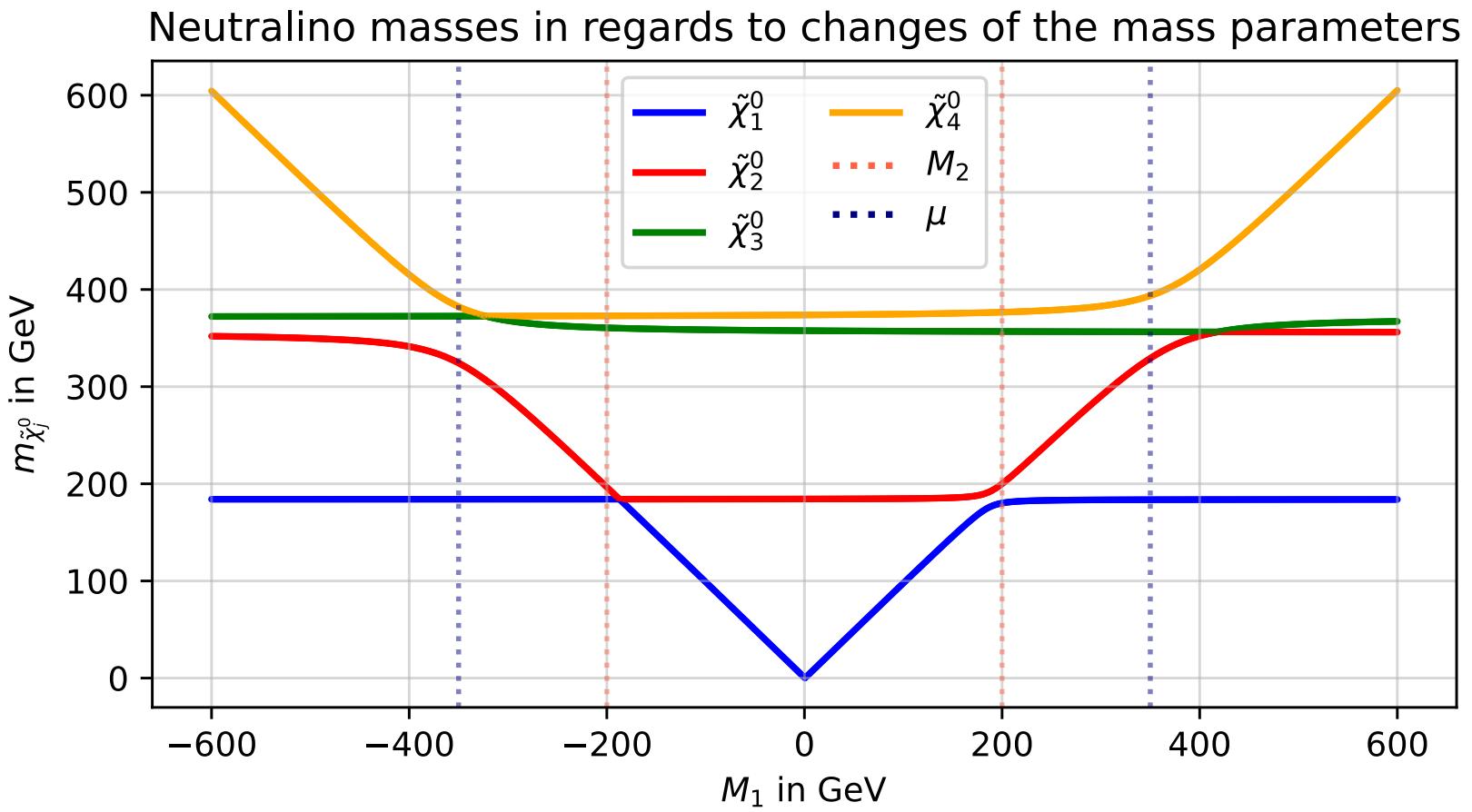
- M_1 : Bino mass parameter
- M_2 : Wino mass parameter
- μ : Higgsino mass parameter
- $\tan \beta = \frac{v_2}{v_1}$

$$M_C = \begin{pmatrix} M_2 & m_W \sqrt{2} \sin \beta \\ m_W \sqrt{2} \cos \beta & \mu \end{pmatrix}$$

$$[1] \quad M_N = \begin{pmatrix} M_1 & 0 & -m_Z \cos \beta \sin \theta_W & m_Z \sin \beta \sin \theta_W \\ 0 & M_2 & m_Z \cos \beta \cos \theta_W & -m_Z \sin \beta \cos \theta_W \\ -m_Z \cos \beta \sin \theta_W & m_Z \cos \beta \cos \theta_W & 0 & -\mu \\ m_Z \sin \beta \sin \theta_W & -m_Z \sin \beta \cos \theta_W & -\mu & 0 \end{pmatrix}^2$$

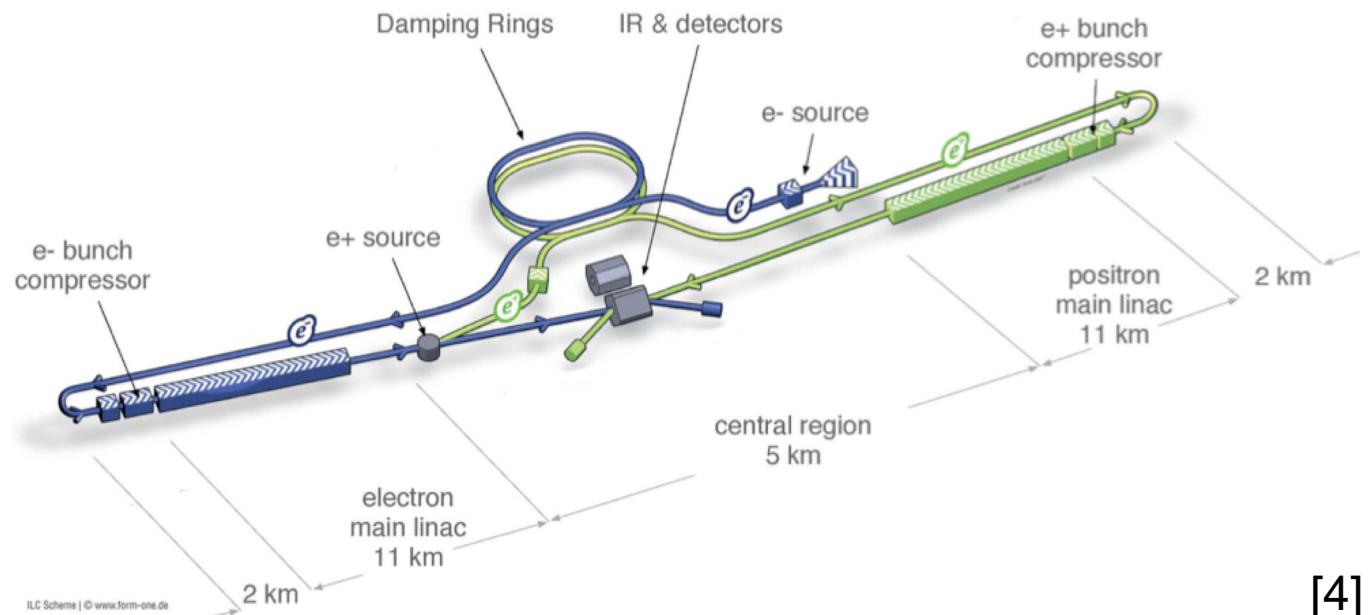
Mass hierarchy

- Particles are assigned by mass hierarchy
- Hierarchy determines sensitivity
- $\tilde{\chi}_1^0$ is a good candidate for DM



International Linear Collider

- Future e^+e^- linear collider with spin polarised beams
- Searches for signs of physics beyond the Standard Model
- Possible experiment:
 - $e^+e^- \rightarrow \tilde{\chi}_1^+\tilde{\chi}_1^-$
 - Luminosity $L = 500 \text{ fb}^{-1}$
 - $\sqrt{S} = 500 \text{ GeV}$



[4]

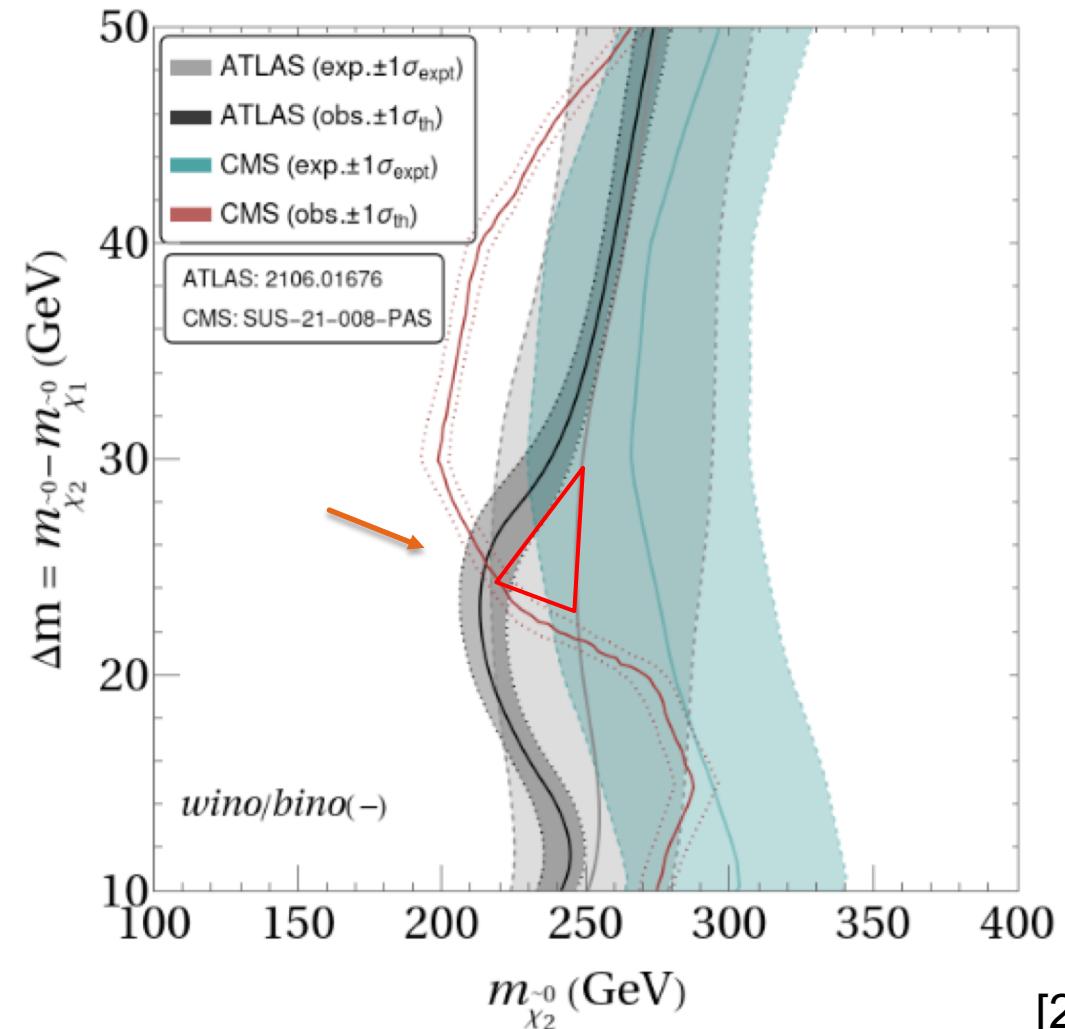
Scenario

- Bino/Wino(-): Bino-wino DM with $\mu \times M_1 < 0$

$$100 \text{ GeV} \leq -M_1 \leq 400 \text{ GeV}, \quad |M_1| \leq M_2 \leq 1.4 |M_1|$$
$$1.2 |M_1| \leq \mu \leq 2 \text{ TeV}, \quad 2 \leq \tan \beta \leq 60$$

$$100 \text{ GeV} \leq m_{\tilde{l}_L} = m_{\tilde{l}_R} \leq 1.5 \text{ TeV}, \quad M_A = 3 \text{ TeV}$$

- Some excess of events → Search for points in this excess



[2]

Data

- SLHA files are taken from „Consistent Excesses in the Search for $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm$: Wino/bino vs. Higgsino Dark Matter“[2]
 - Florian Lika did the same analysis before but with other data
- Constraints:
 - Muon g-2 contribution
 - Vacuum stability constraints
 - DM relic density constraints
 - Constraints from LHC measurements
 - DM direct detection constraints

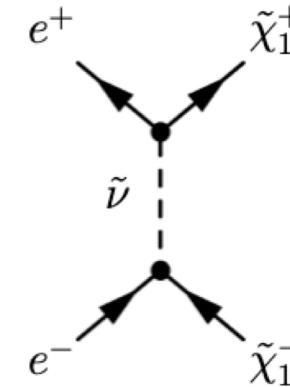
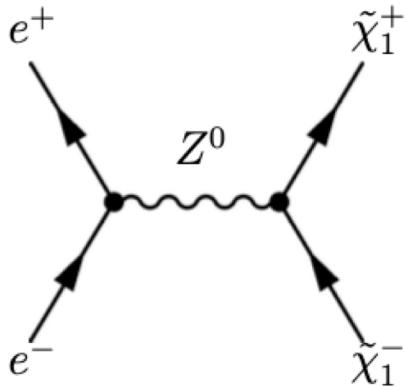
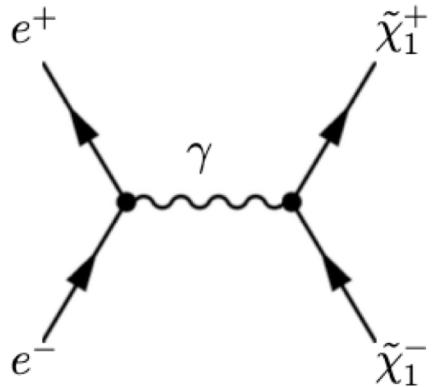
Steps of the analysis

- Chargino production crosssection $\sigma(e^+e^- \rightarrow \tilde{\chi}_1^+\tilde{\chi}_1^-)$ and chargino and neutralino masses ($m_{\tilde{\chi}_1^\pm}, m_{\tilde{\chi}_1^0}$) are measured
- Left and right chargino mixing angles $\phi_{L,R}$ are determined
- The parameters (M_1, M_2, μ) are reconstructed with uncertainties via a scan over the parameter space
- DM relic density (Ωh^2) is calculated and compared to the relic density gained from the true values

Chargino pair production

- Contributing first order Feynman diagrams for chargino pair production:

[1]



- Crosssection calculated with mixing angles

$$\sigma^\pm\{ij\} = c_1 \cos^2 2\Phi_L + c_2 \cos 2\Phi_L + c_3 \cos^2 2\Phi_R + c_4 \cos 2\Phi_R + c_5 \cos 2\Phi_L \cos 2\Phi_R + c_6$$

[5]

Configurations

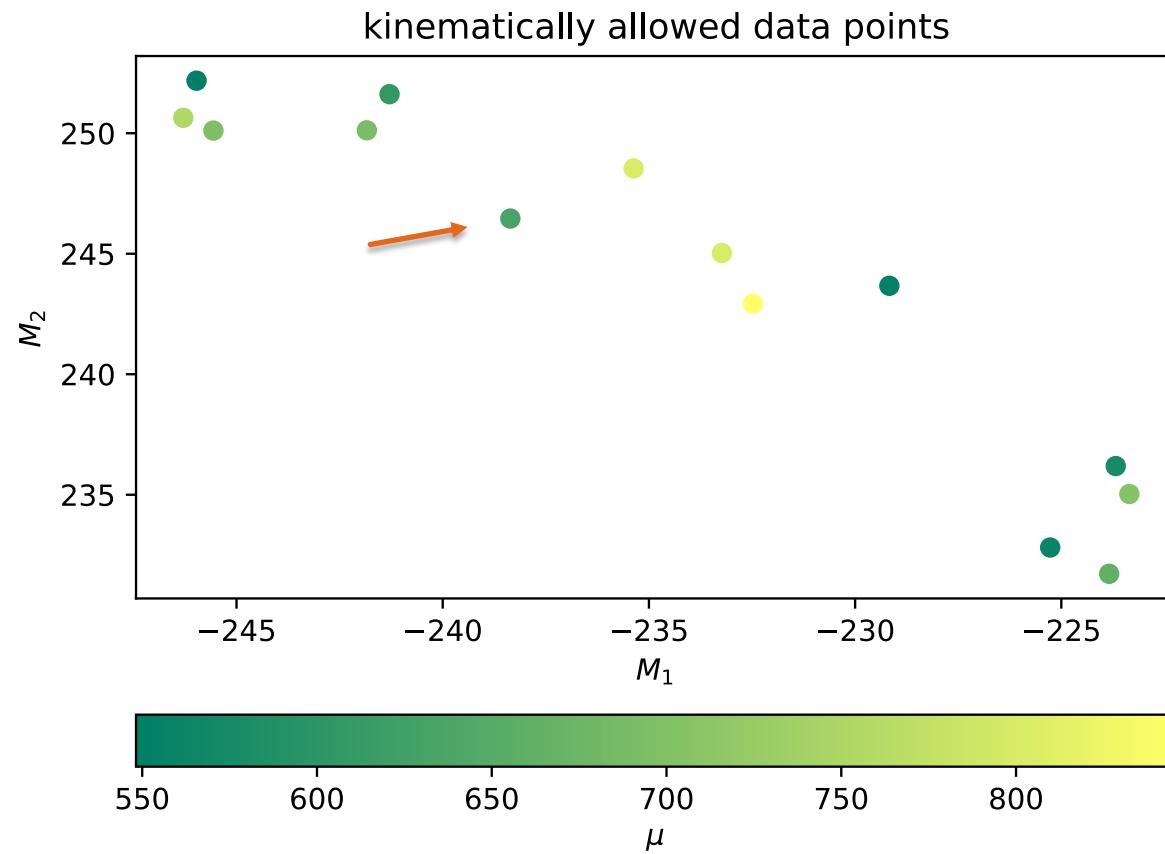
- We use four beam configurations to increase precision
- Consider the following uncertainties:

statistical uncertainty	$\sigma_{\text{stat}} = \frac{\sqrt{\sigma L}}{L}$
polarisation uncertainty	$P_{e^-} = \pm 0.804, P_{e^+} = \mp 0.603$ $P_{e^-} = \pm 0.796, P_{e^+} = \mp 0.597$
mass uncertainty	$\delta m_{\tilde{\chi}_1^\pm} = 0.005m_{\tilde{\chi}_1^\pm}$

\sqrt{S}	P_{e^-}	P_{e^+}
500 GeV	-0.8	+0.6
500 GeV	+0.8	-0.6
550 GeV	-0.8	+0.6
550 GeV	+0.8	-0.6

Example point

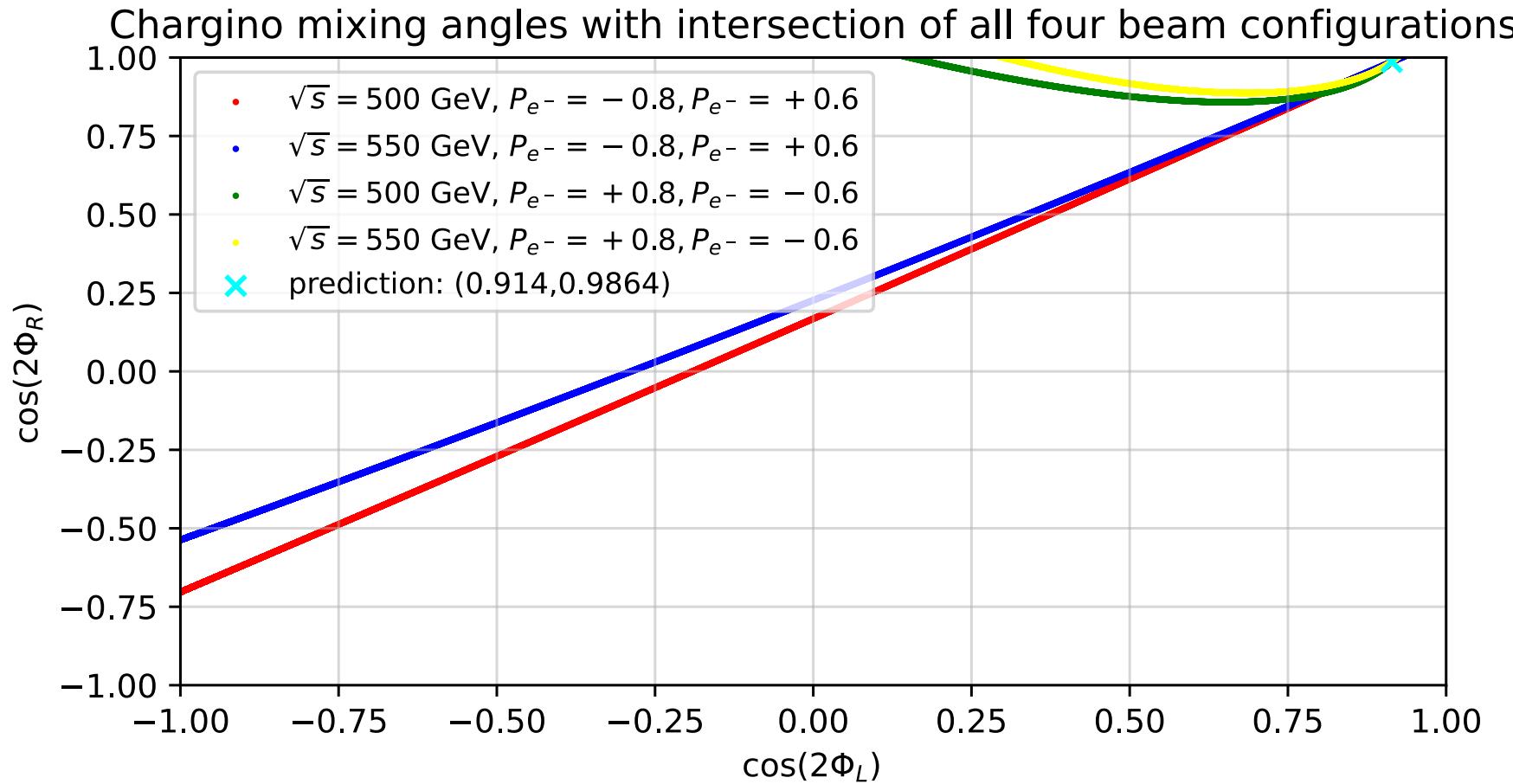
- Out of 14 kinematically allowed data points, one is taken as example point for following explanations



M_1	-238.36 GeV	$\sigma_{-8,+6}^{500}$	394.9612 fb
M_2	246.46 GeV	$\sigma_{+8,-6}^{500}$	11.0373 fb
μ	632.64 GeV	$\sigma_{-8,+6}^{550}$	519.8185 fb
$\tan \beta$	51.97	$\sigma_{+8,-6}^{550}$	14.4873 fb
$M_{\tilde{\nu}}$	496.78 GeV	$m_{\tilde{\chi}_1^\pm}$	241.48 GeV

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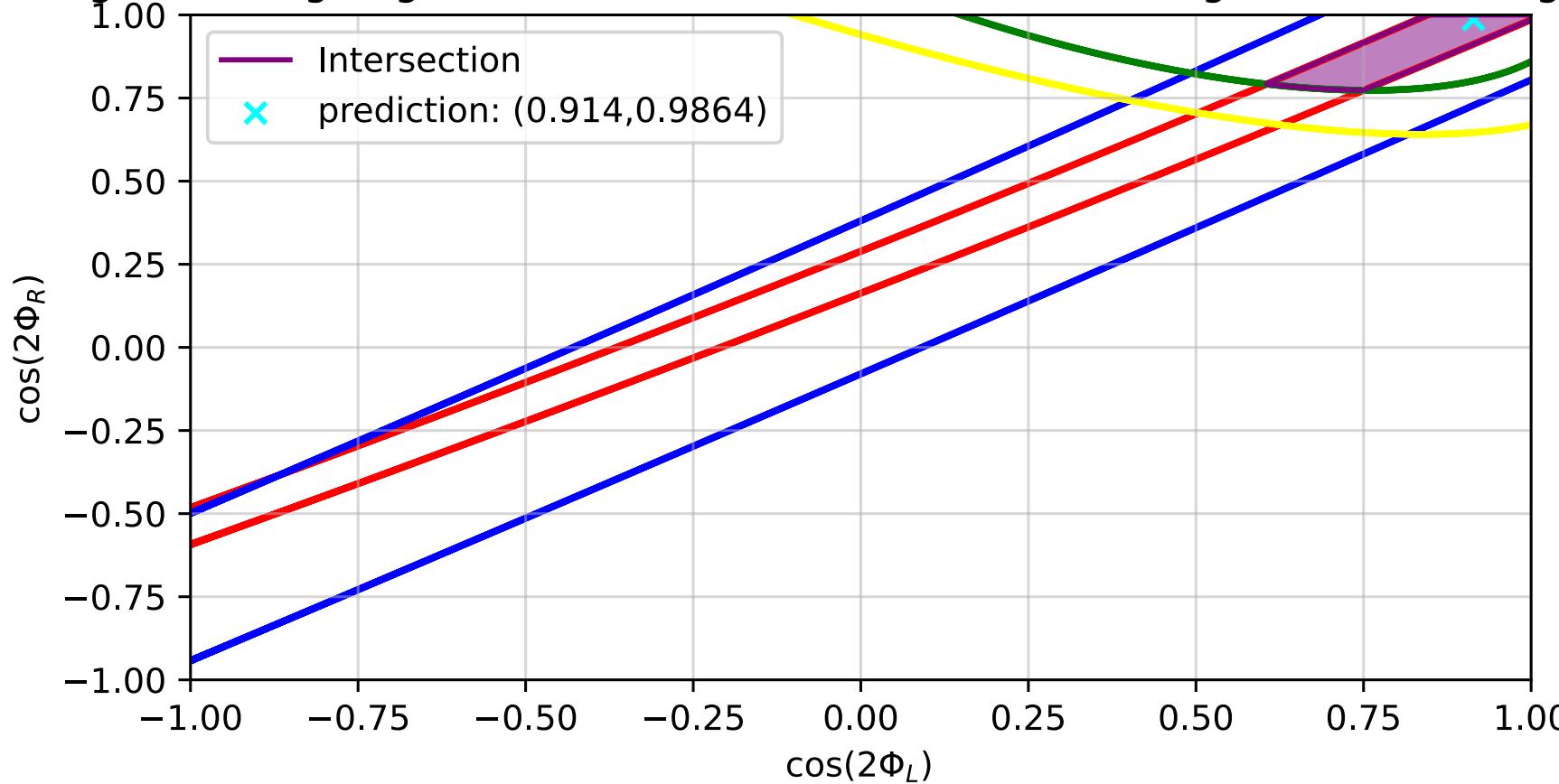
Mixing angles without uncertainties



$$\sigma^\pm\{ij\} = c_1 \cos^2 2\Phi_L + c_2 \cos 2\Phi_L + c_3 \cos^2 2\Phi_R + c_4 \cos 2\Phi_R + c_5 \cos 2\Phi_L \cos 2\Phi_R + c_6$$

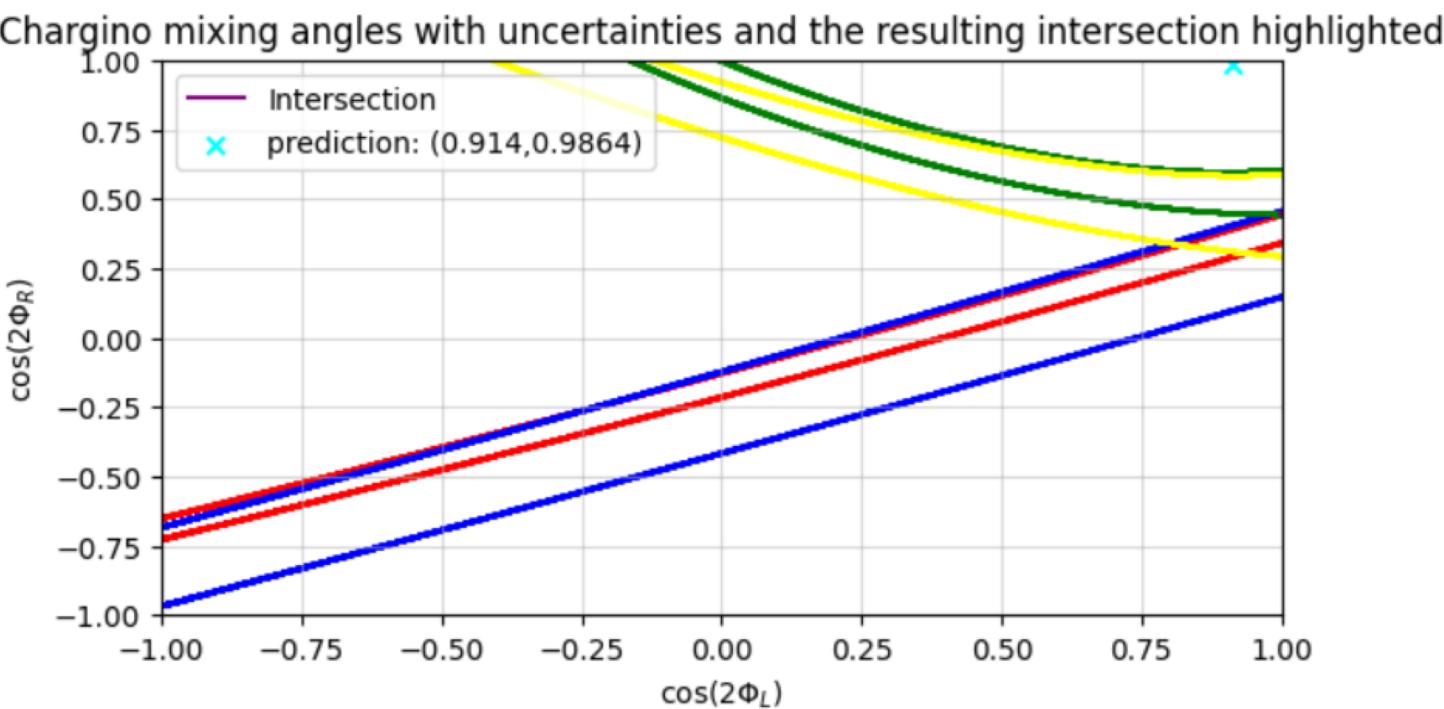
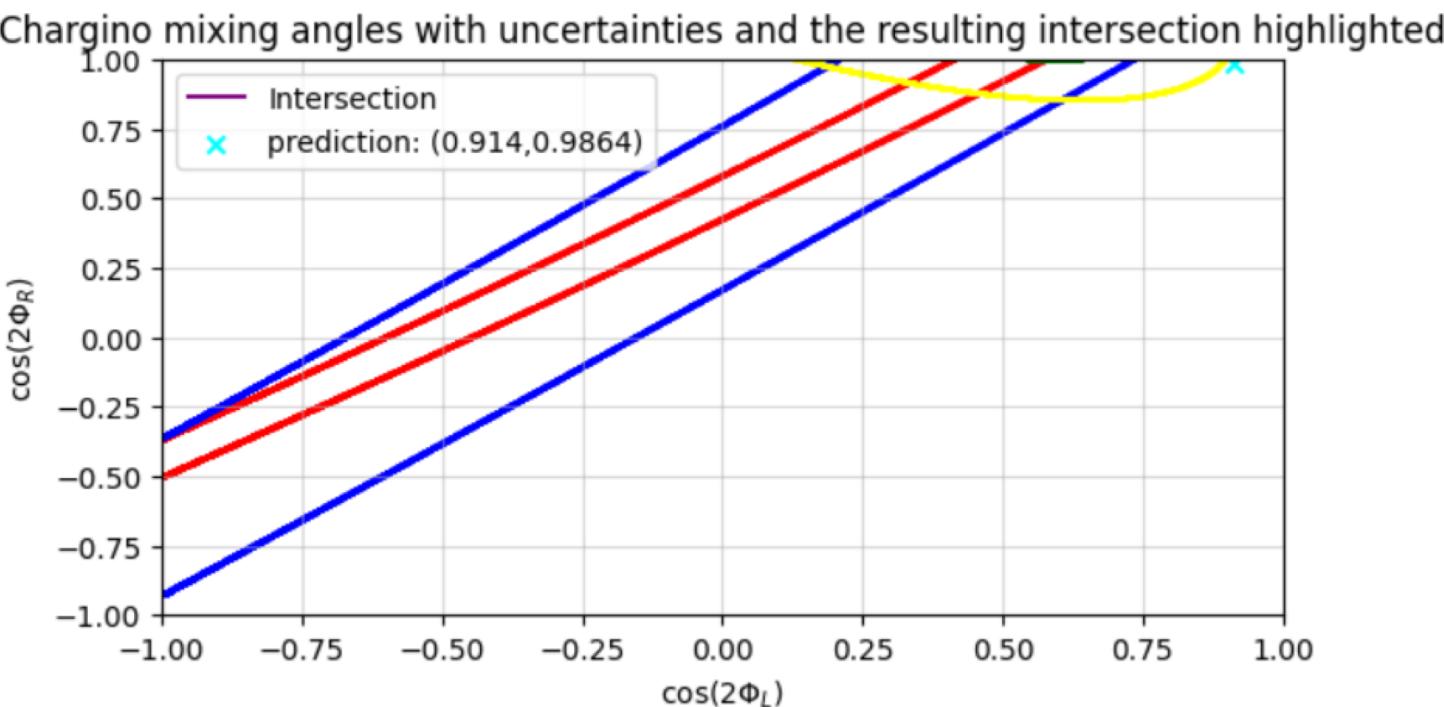
Valid mixing angles

Chargino mixing angles with uncertainties and the resulting intersection highlighted



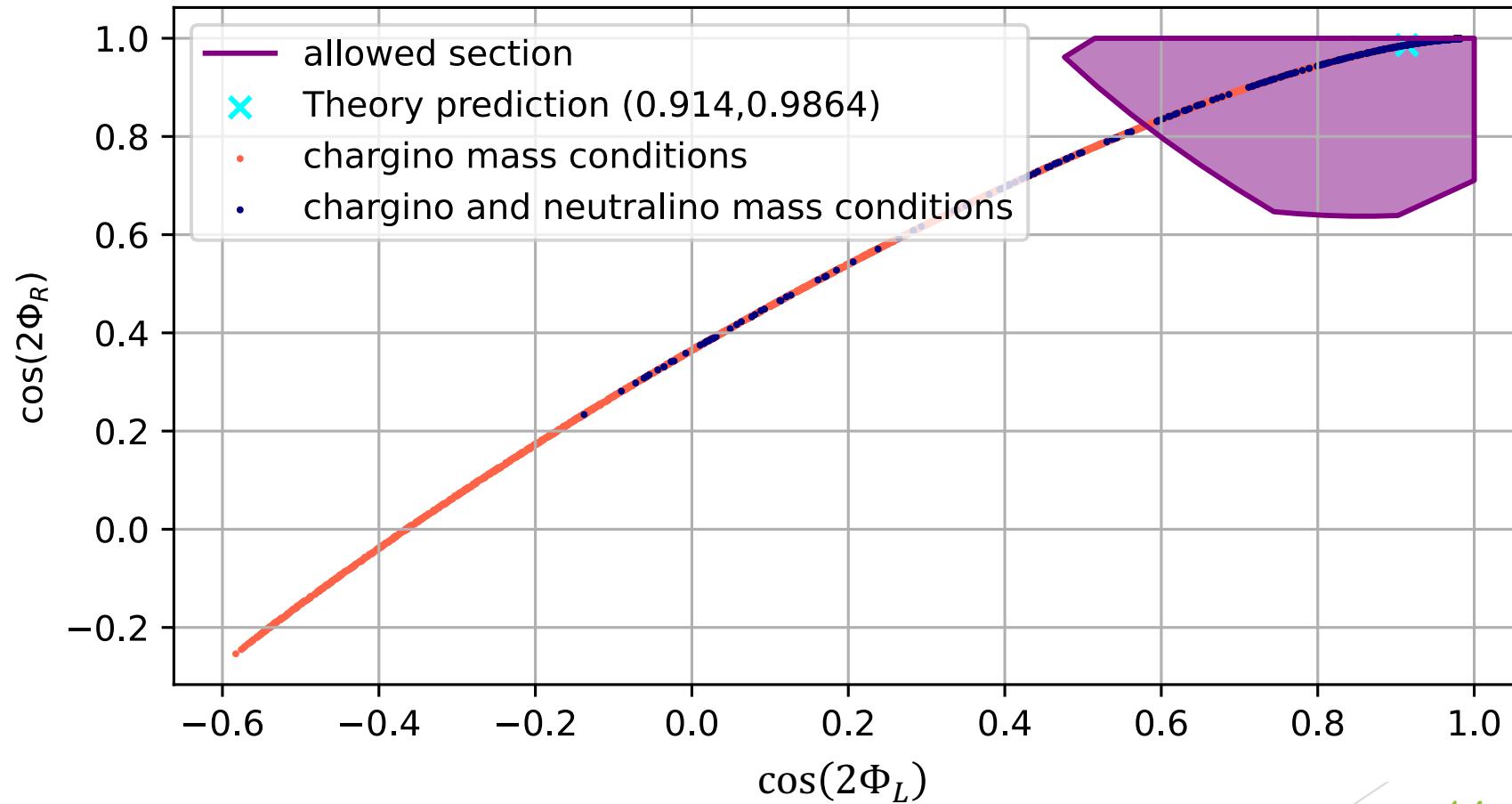
Variation of $M_{\tilde{\chi}}$

- In steps of 10 GeV
- Variation for the whole data set: -140 to +110 GeV
- Variation for this point: -40 to +80 GeV



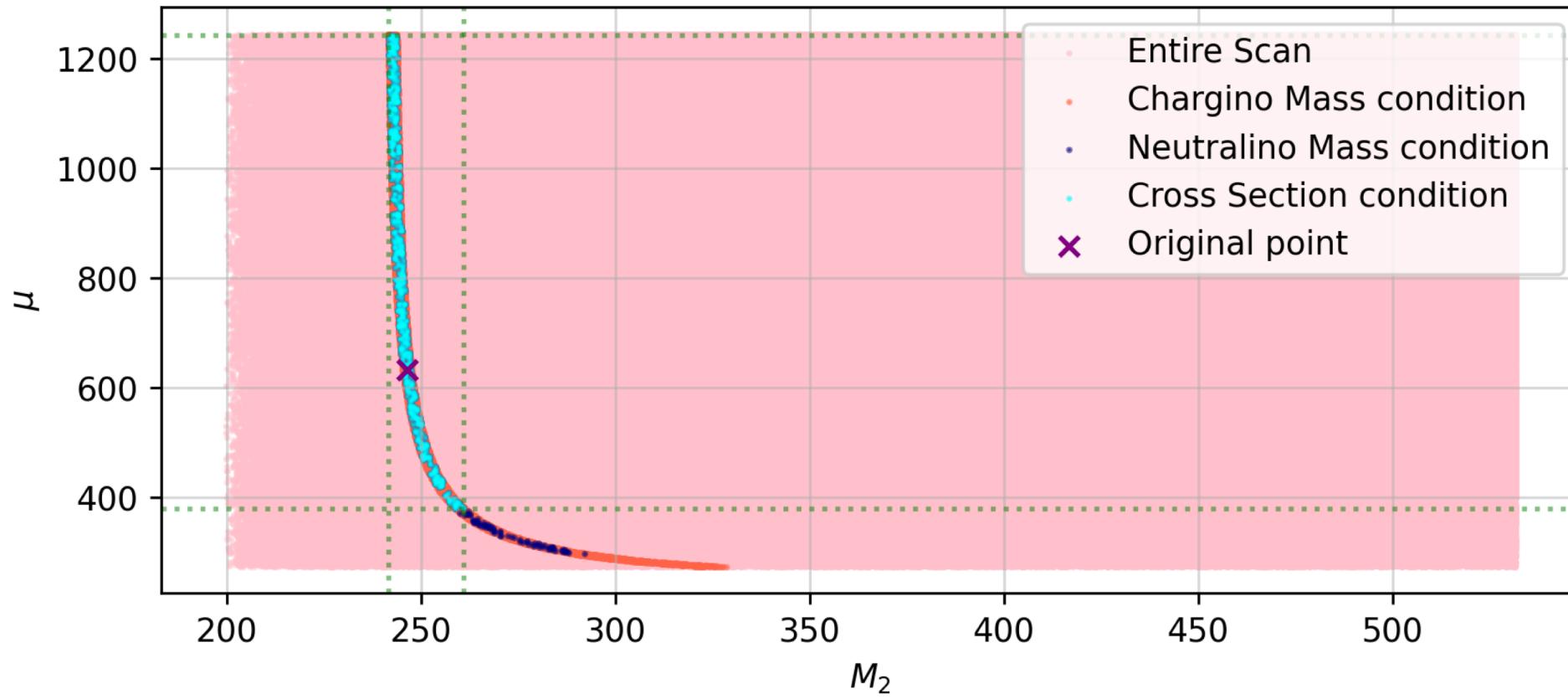
Scanning for valid points

Scan results with cross-section condition visualised using the intersection surface.

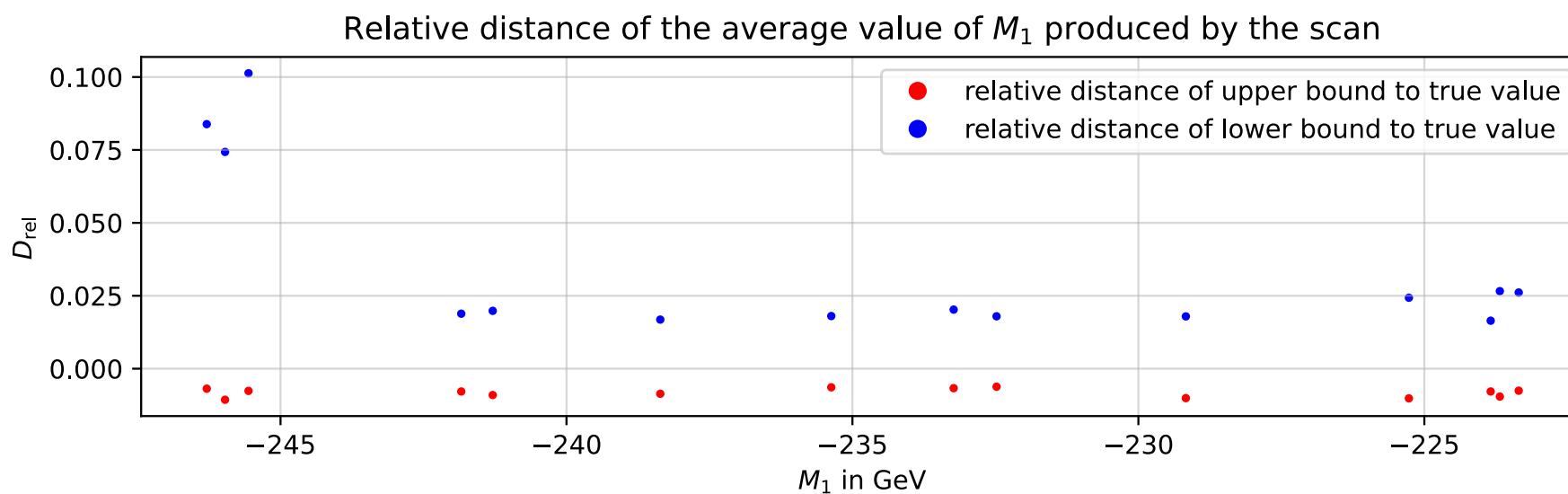
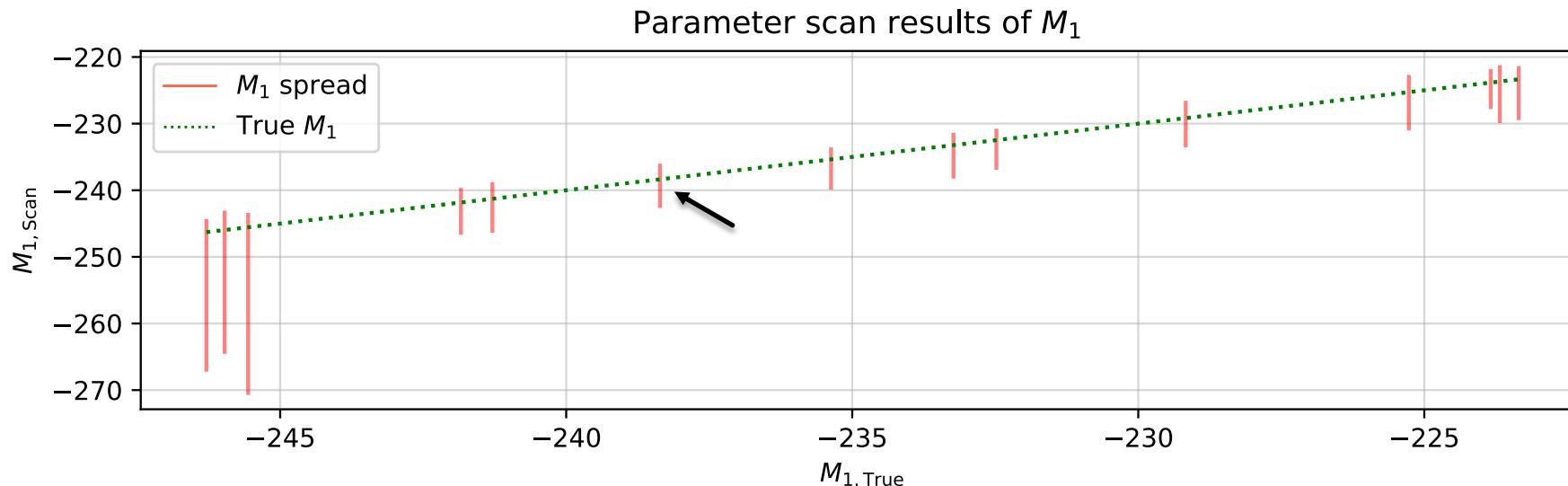


Scanning for valid points

Scan results in the M_2, μ plane.



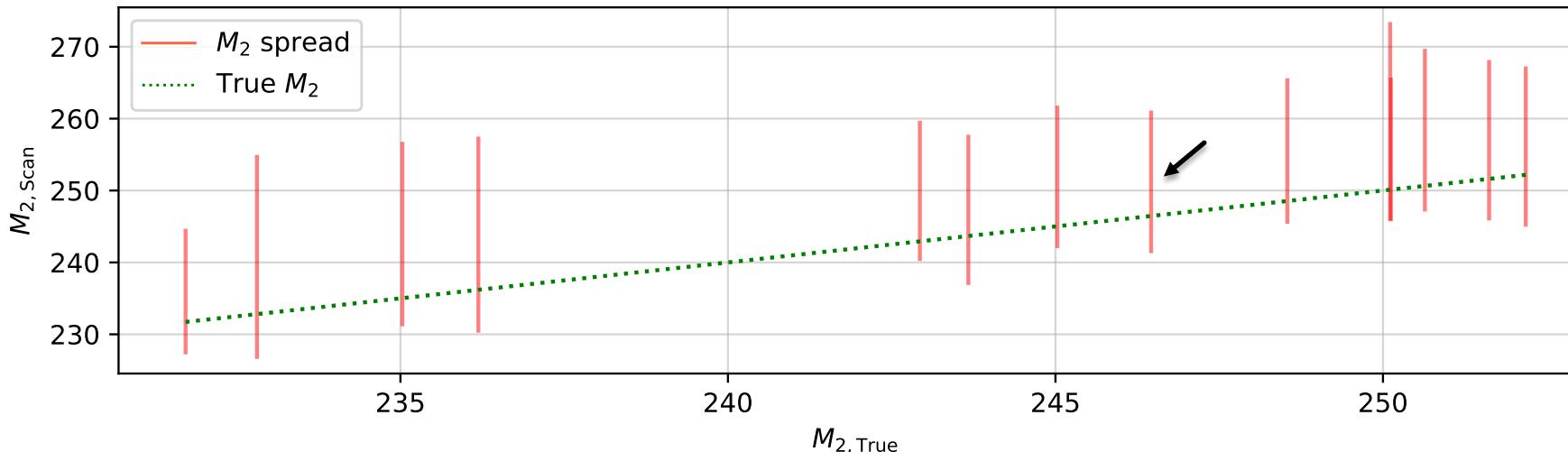
Results for M_1



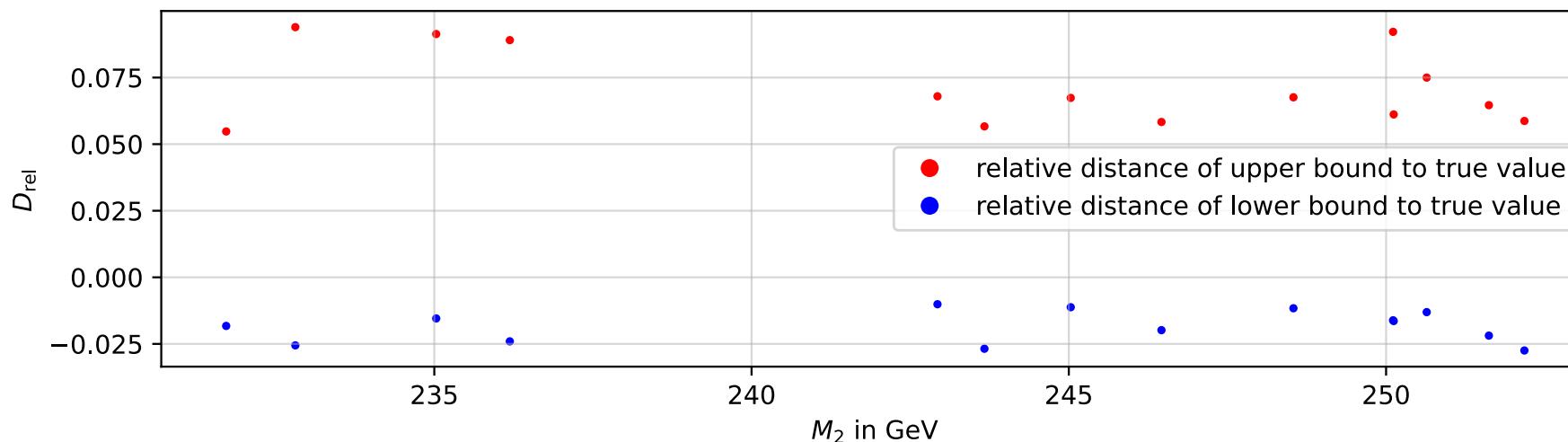
$$D_{\text{rel}} = \frac{M_{1,\text{scan}} - M_{1,\text{true}}}{M_{1,\text{true}}}$$

Results for M_2

Parameter scan results of M_2



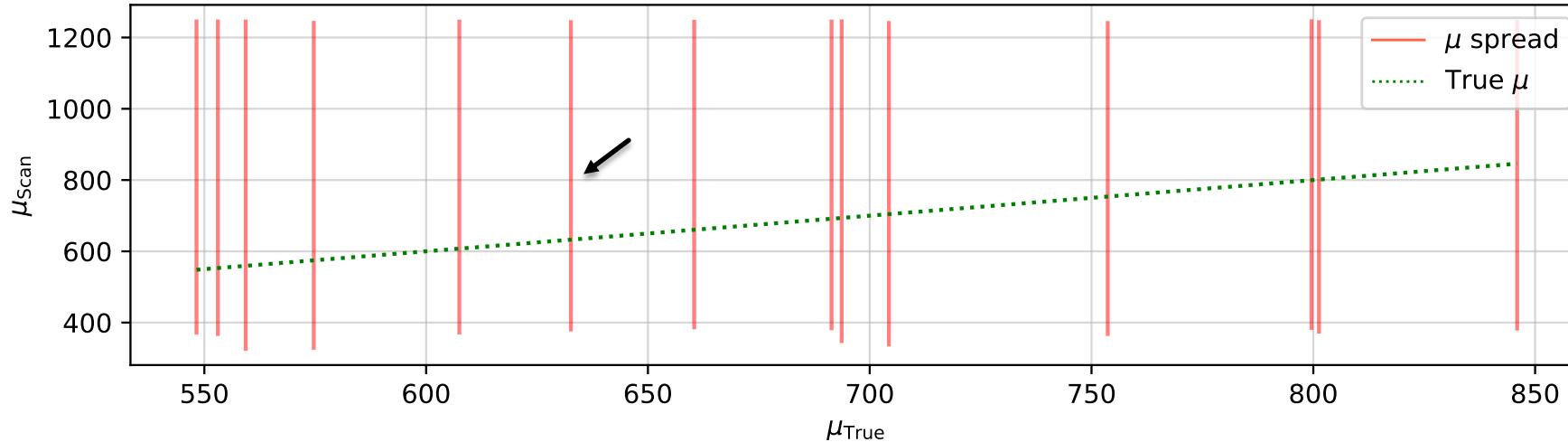
Relative distance of the average value of M_2 produced by the scan



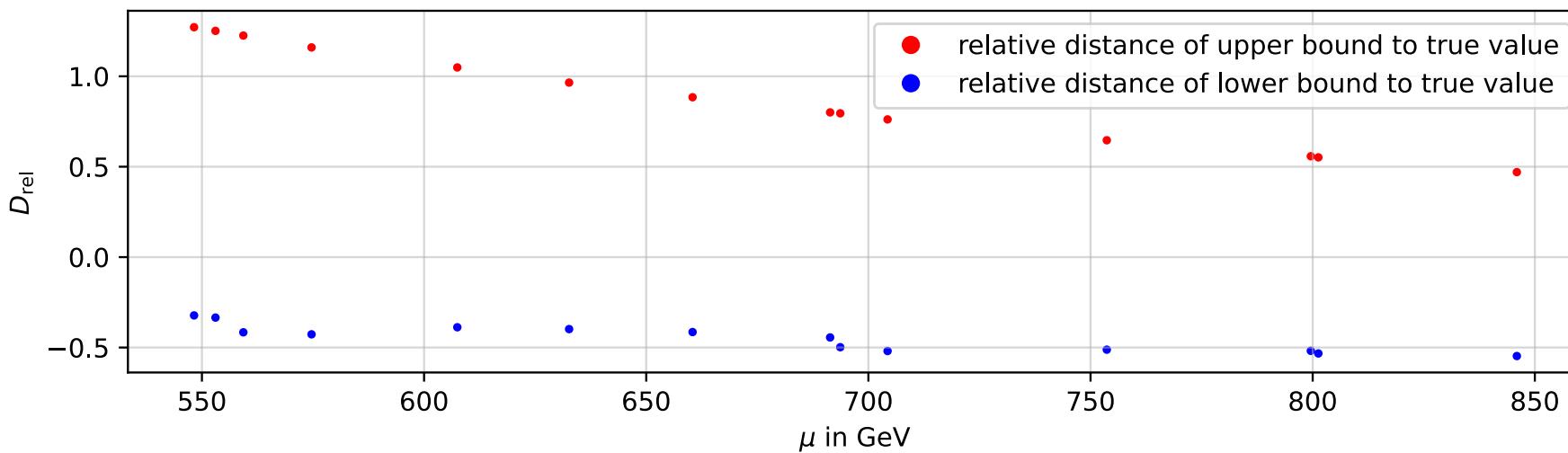
$$D_{\text{rel}} = \frac{M_{2,\text{scan}} - M_{2,\text{true}}}{M_{2,\text{true}}}$$

Results for μ

Parameter scan results of μ

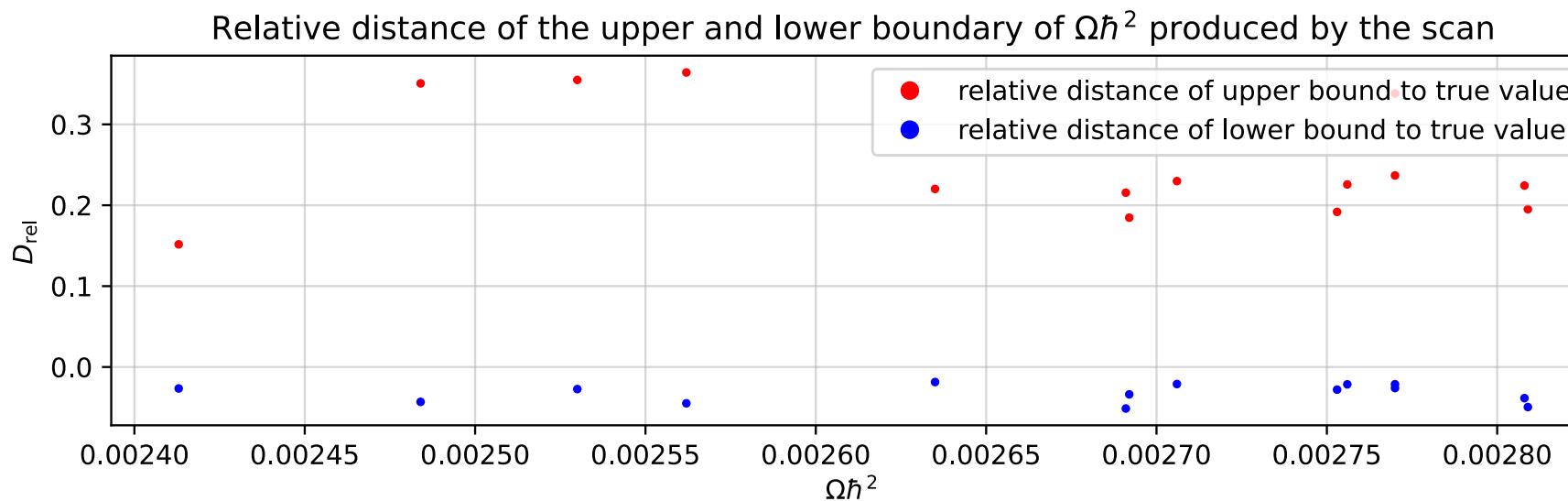
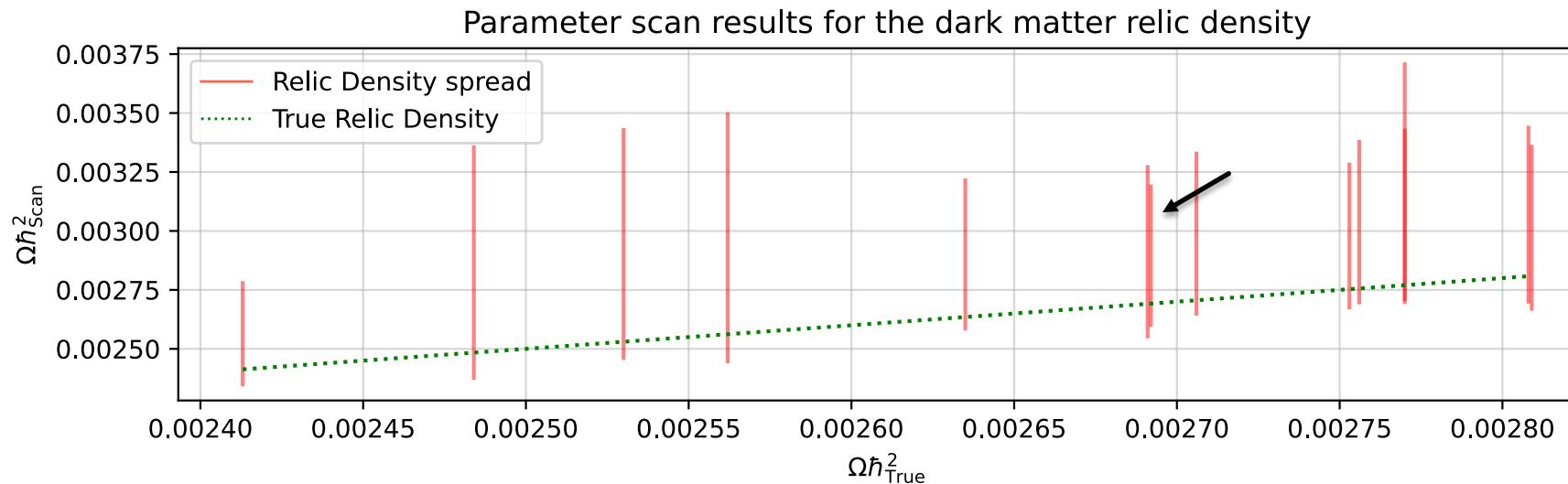


Relative distance of the average value of μ produced by the scan



$$D_{\text{rel}} = \frac{\mu_{\text{scan}} - \mu_{\text{true}}}{\mu_{\text{true}}}$$

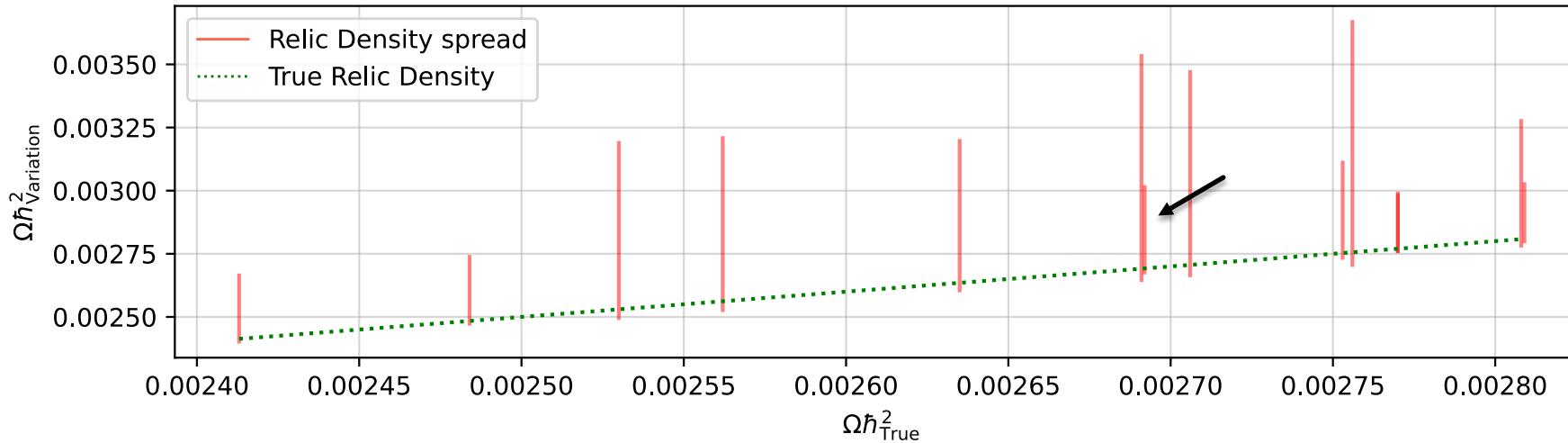
Results for DM relic density



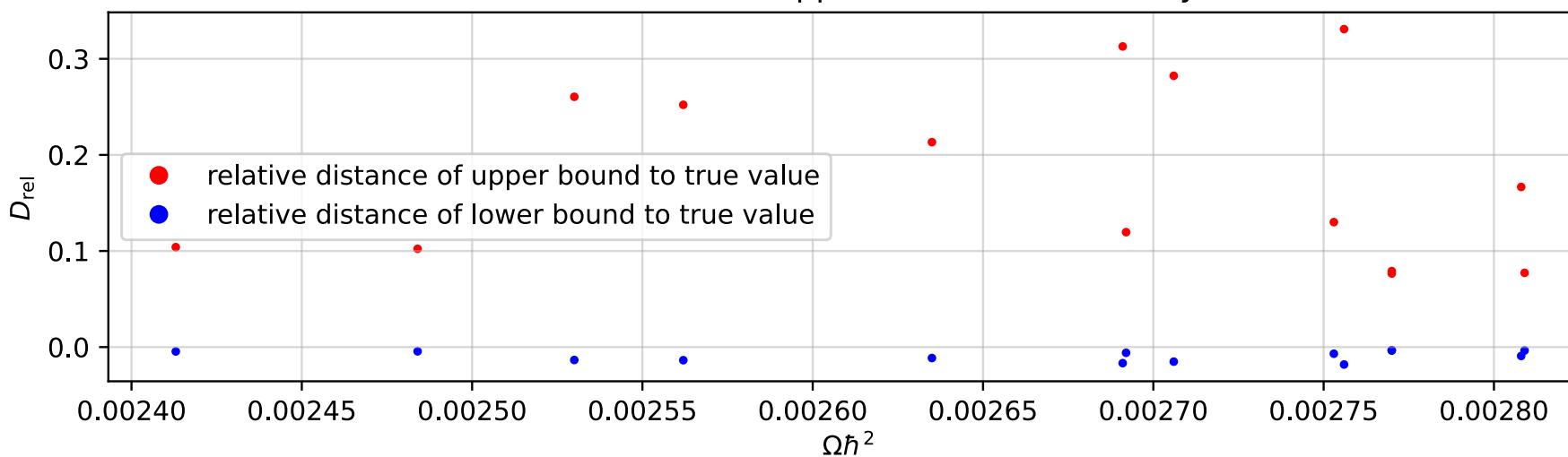
$$D_{\text{rel}} = \frac{\Omega h^2_{\text{scan}} - \Omega h^2_{\text{true}}}{\Omega h^2_{\text{true}}}$$

Variation of M_1

Dark matter relic density with M1 variation



Relative distance of the upper and lower boundary of Ωh^2

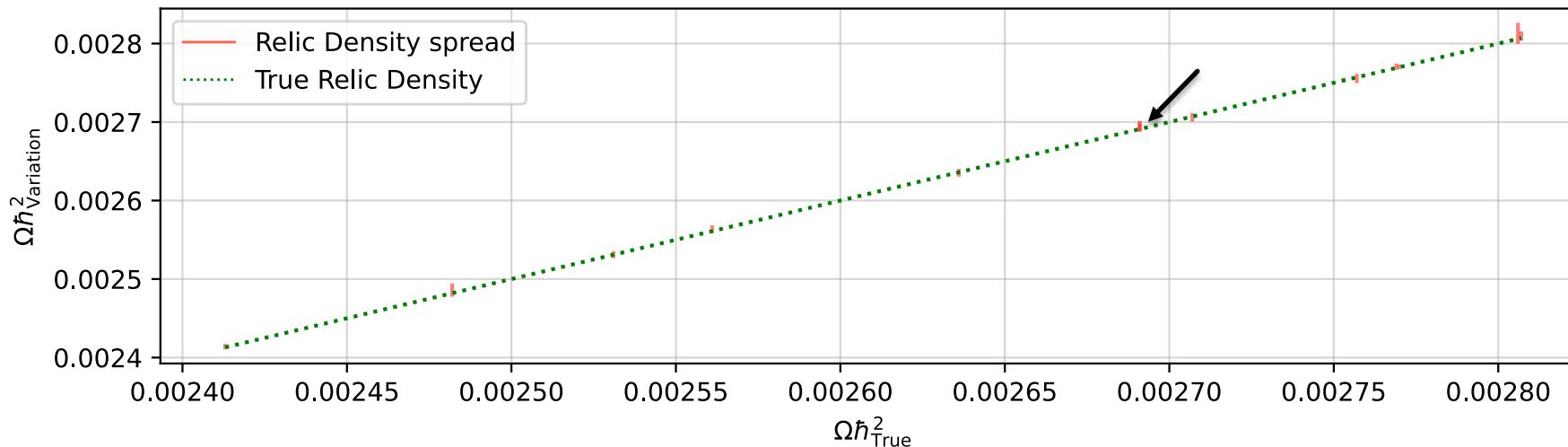


$$D_{\text{rel}} = \frac{M_{1,\text{var}} - M_{1,\text{true}}}{M_{1,\text{true}}}$$

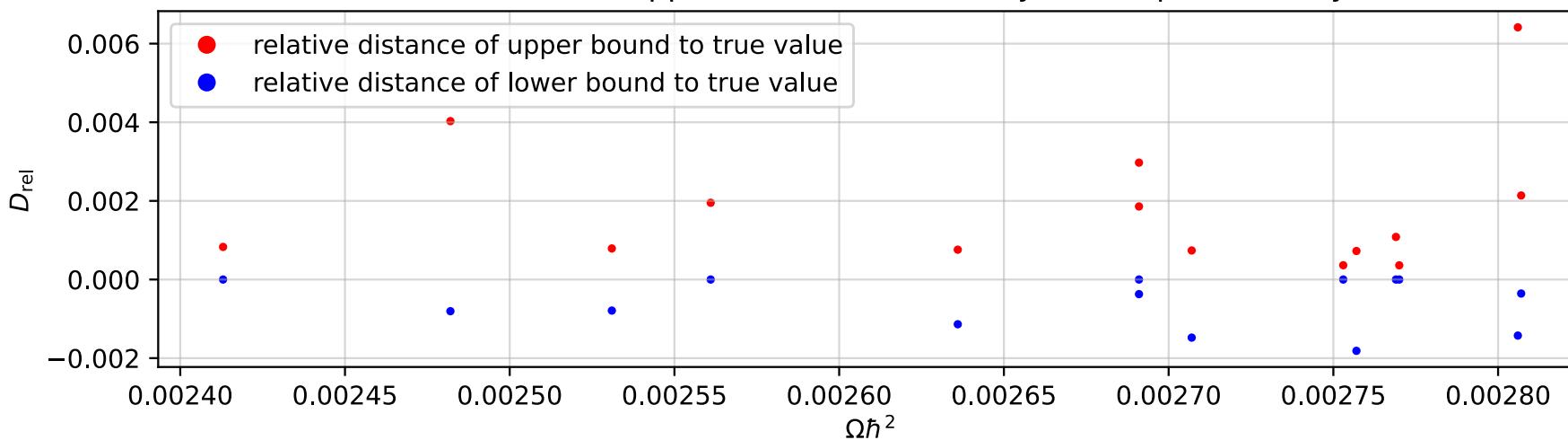
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Variation of μ

Parameter scan results for the dark matter relic density



Relative distance of the upper and lower boundary of Ωh^2 produced by the scan



$$D_{\text{rel}} = \frac{\mu_{\text{var}} - \mu_{\text{true}}}{\mu_{\text{true}}}$$

Conclusion & Outlook

- Florians code was debugged and improved
- Different data set was studied
- Next steps:
 - Include forward-backward asymmetry (A_{FB}) as additional parameter to improve accuracy
 - Find alternative method of getting $M_{\tilde{\nu}}$ (via A_{FB}) to get a more reliable uncertainty and thus improve the accuracy
 - Optimize beam energy to get higher crosssections and better accuracy

Thank you for your
attention! ☺

Questions?

Sources

- [1] Combined LHC/ILC analysis of a SUSY scenario with heavy sfermions, 2018, arXiv:hep-ph/0607104v2
- [2] Consistent Excesses in the Search for $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm$: Wino/bino vs. Higgsino Dark Matter, 2024, arXiv:2403.14759v1
- [3] Master thesis of Florian Lika: SUSY Parameter determination within Dark Matter Phenomenology at future e+e- Colliders, 2023
- [4] On the physics potential of ILC and CLIC, A.F.Zarnecki, 2020
- [5] SUSY Parameter Determination in Combined Analyses at LHC/LC, 2003, arXiv:hep-ph/0312069v1
- [6] Beyond the standard model seraches at the LHC, 2016,
https://indico.cern.ch/event/555909/contributions/2265363/attachments/1325399/1989408/deroeck_Islamabad2016_2_v1.pdf

Backup

- Coefficients from crosssection formula:

$$c_1 = \int_C |Z|^2 \{ c_{LR} L^2 f_2 + c_{RL} R^2 f_1 \}$$

$$\begin{aligned} c_2 = & \int_C |Z|^2 \{ c_{LR} L^2 (1 - 4L)(2f_2 + f_3) + c_{RL} R^2 (1 - 4R)(2f_1 + f_3) \} \\ & - \int_C G \tilde{N} 4 \{ c_{LR} L (2f_2 + f_3) + c_{RL} R (2f_1 + f_3) \} - \int_C Re(Z) \tilde{N} c_{LR} L f_3 \end{aligned}$$

$$c_3 = \int_C |Z|^2 (c_{LR} L^2 f_1 + c_{RL} R^2 f_2) - \int_C Z \tilde{N} 2 c_{LR} L f_1 + \int_C \tilde{N}^2 c_{LR} f_1$$

$$\begin{aligned} c_4 = & \int_C |Z|^2 (1 - 4L) \{ c_{LR} L^2 (2f_1 + f_3) + c_{RL} R^2 (2f_2 + f_3) \} + \int_C \tilde{N}^2 2 c_{LR} f_1 \\ & + \int_C Re(Z) \tilde{N} c_{LR} L \{ -4f_1 - f_3 + 4L(2f_1 + f_3) \} + \int_C G \tilde{N} 4 c_{LR} (2f_1 + f_3) \\ & - \int_C G Re(Z) 4 \{ c_{LR} L (2f_1 + f_3) + c_{RL} R (2f_2 + f_3) \} \end{aligned}$$

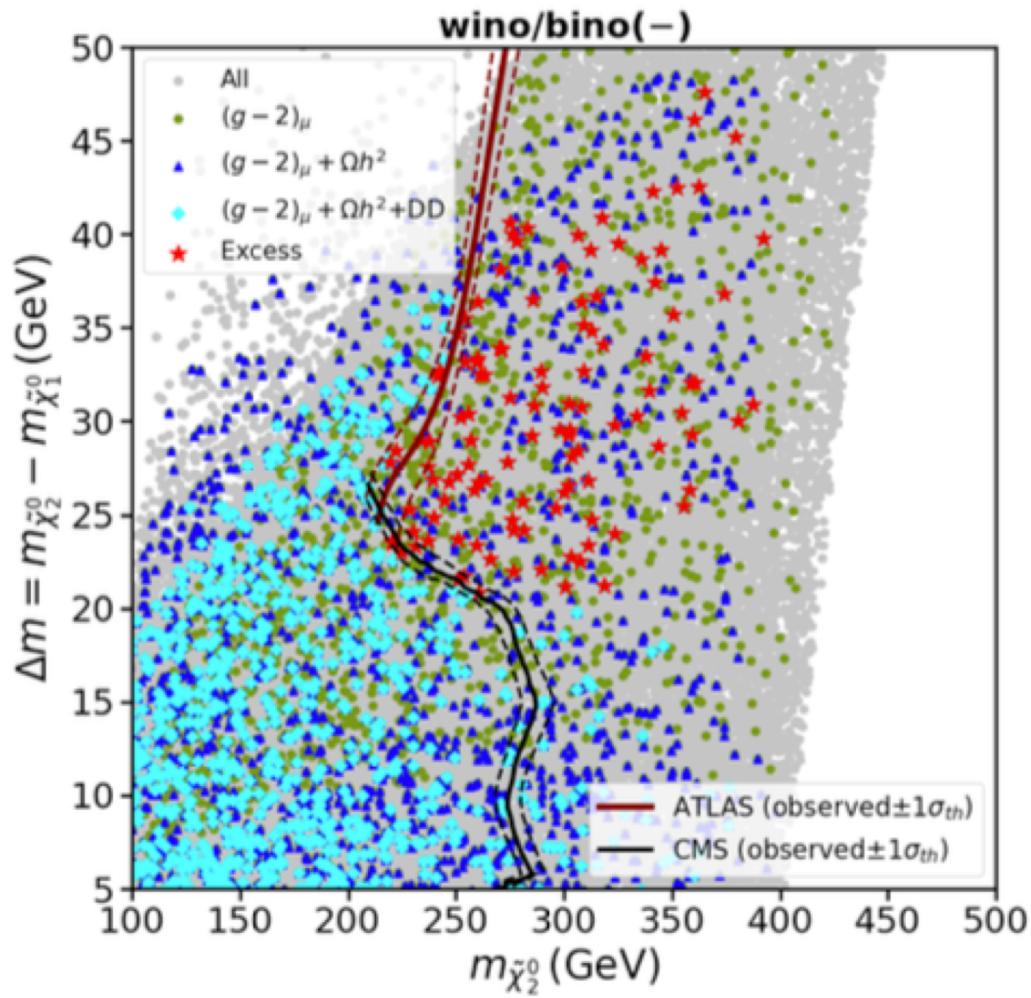
$$\begin{aligned} c_5 = & \int_C |Z|^2 (c_{LR} L^2 + c_{RL} R^2) f_3 - \int_C Re(Z) \tilde{N} c_{LR} L f_3 \\ c_6 = & \int_C |Z|^2 \{ c_{LR} L^2 (1 - 8L) + c_{RL} R^2 (1 - 8R) + 16L^2 (c_{LR} L^2 + c_{RL} R^2) \} (f_1 + f_2 + f_3) \\ & - \int_C Re(Z) \tilde{N} c_{LR} L (1 - 4L) (2f_1 + f_3) + \int_C G^2 (c_{LR} + c_{RL}) (f_1 + f_2 + f_3) \\ & - \int_C Re(Z) G 8 \{ c_{RL} R + c_{LR} L (1 - 4L) \} (f_1 + f_2 + f_3) + \int_C \tilde{N}^2 c_{LR} f_1 \\ & + \int_C G \tilde{N} 4 c_{LR} (2f_1 + f_3) \end{aligned}$$

Backup

$$m_{\tilde{\chi}_{1,2}^\pm}^2 = \frac{1}{2}(M_2^2 + \mu^2 + 2m_W^2 \mp \Delta_C)$$

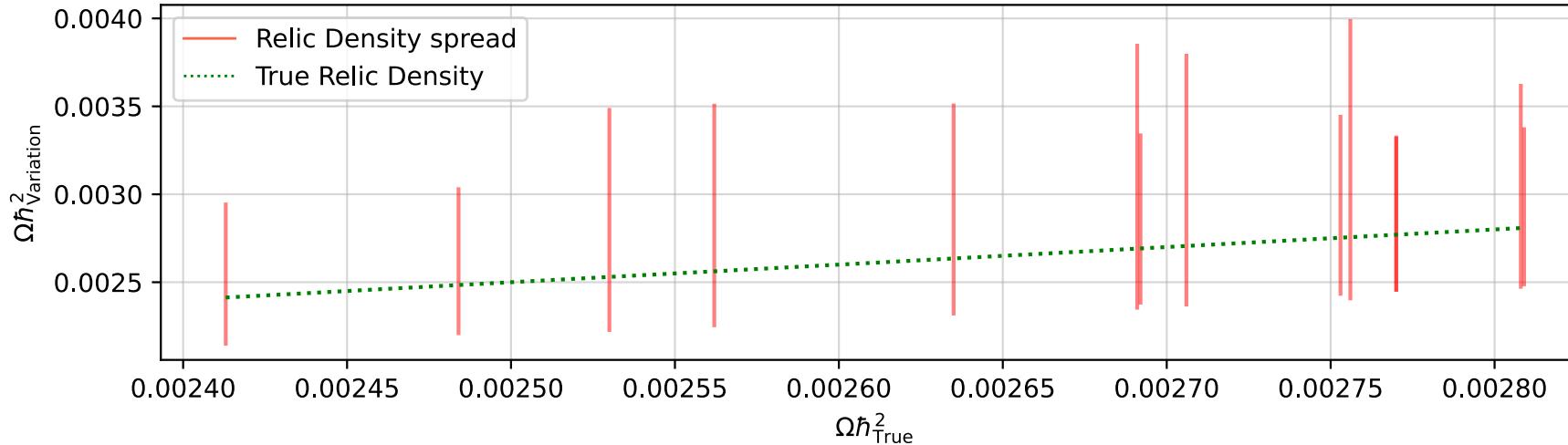
$$\Delta c = \left[(M_2^2 - \mu^2)^2 + 4m_W^4 \cos^2 2\beta + 4m_W^2(M_2^2 + \mu^2) + 8m_W^2 M_2 \mu \sin 2\beta \right]^{\frac{1}{2}}$$

Backup

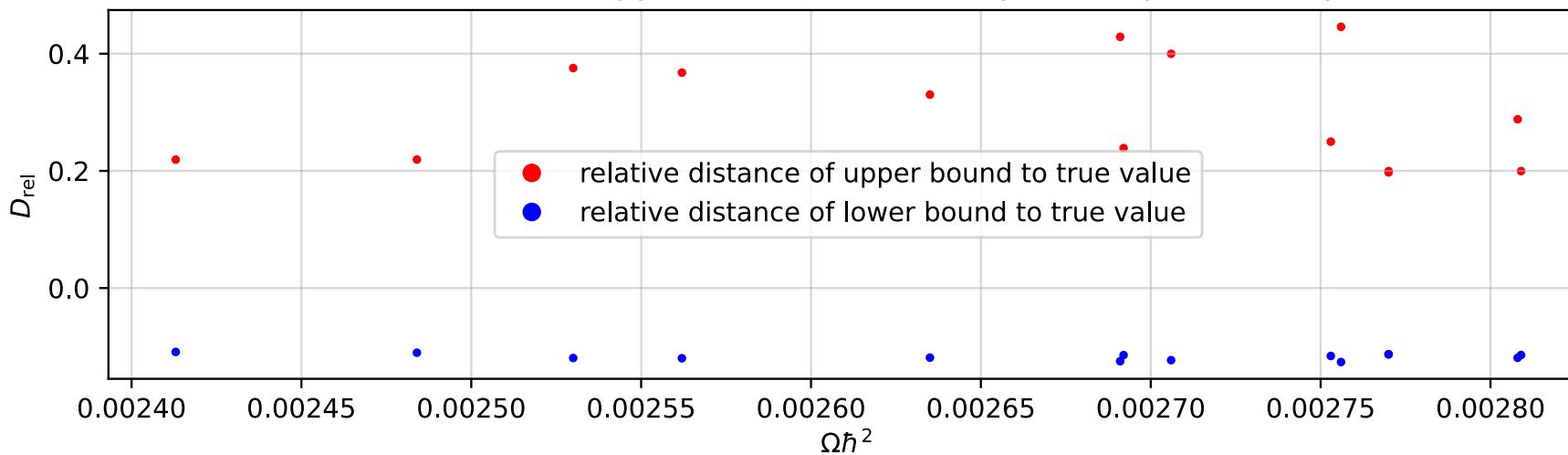


Variation of M_2

Parameter scan results for the dark matter relic density

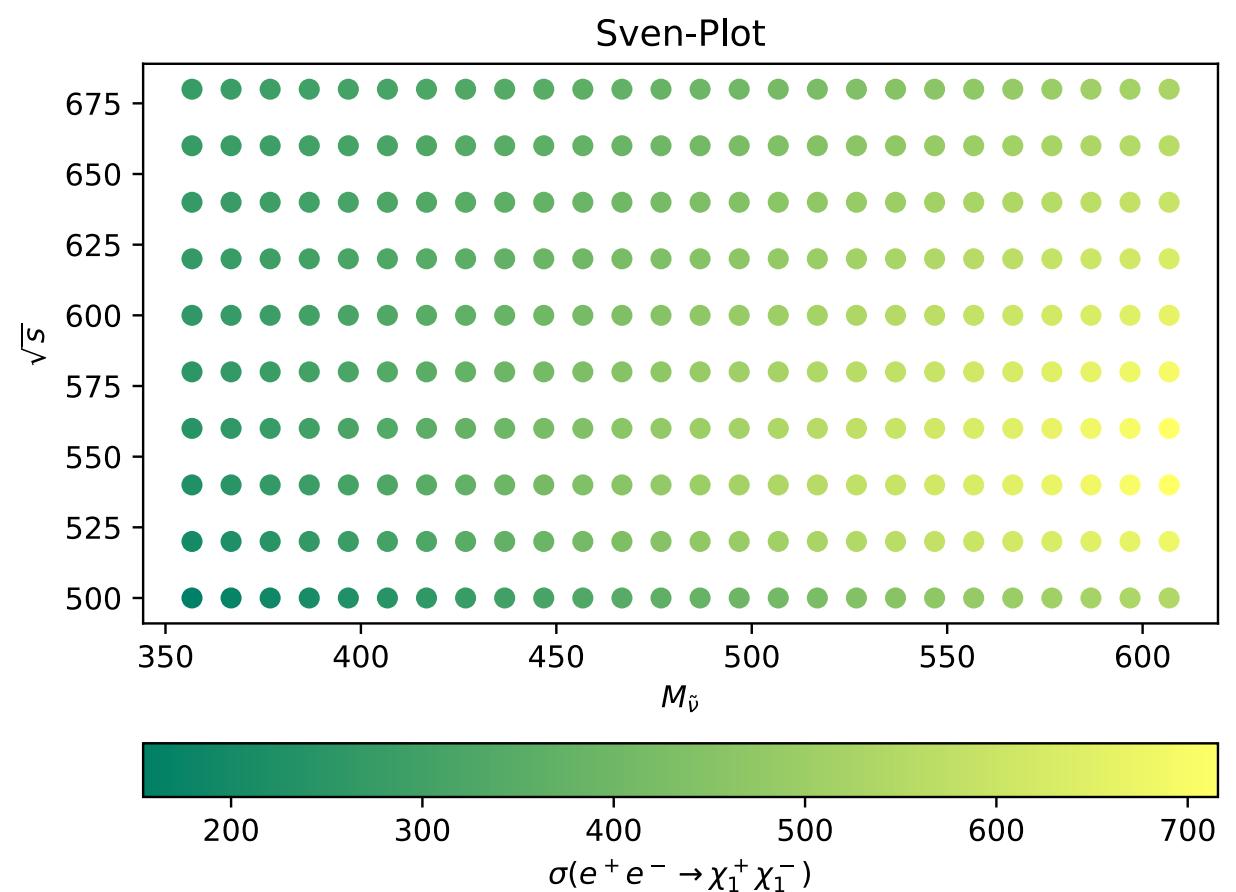
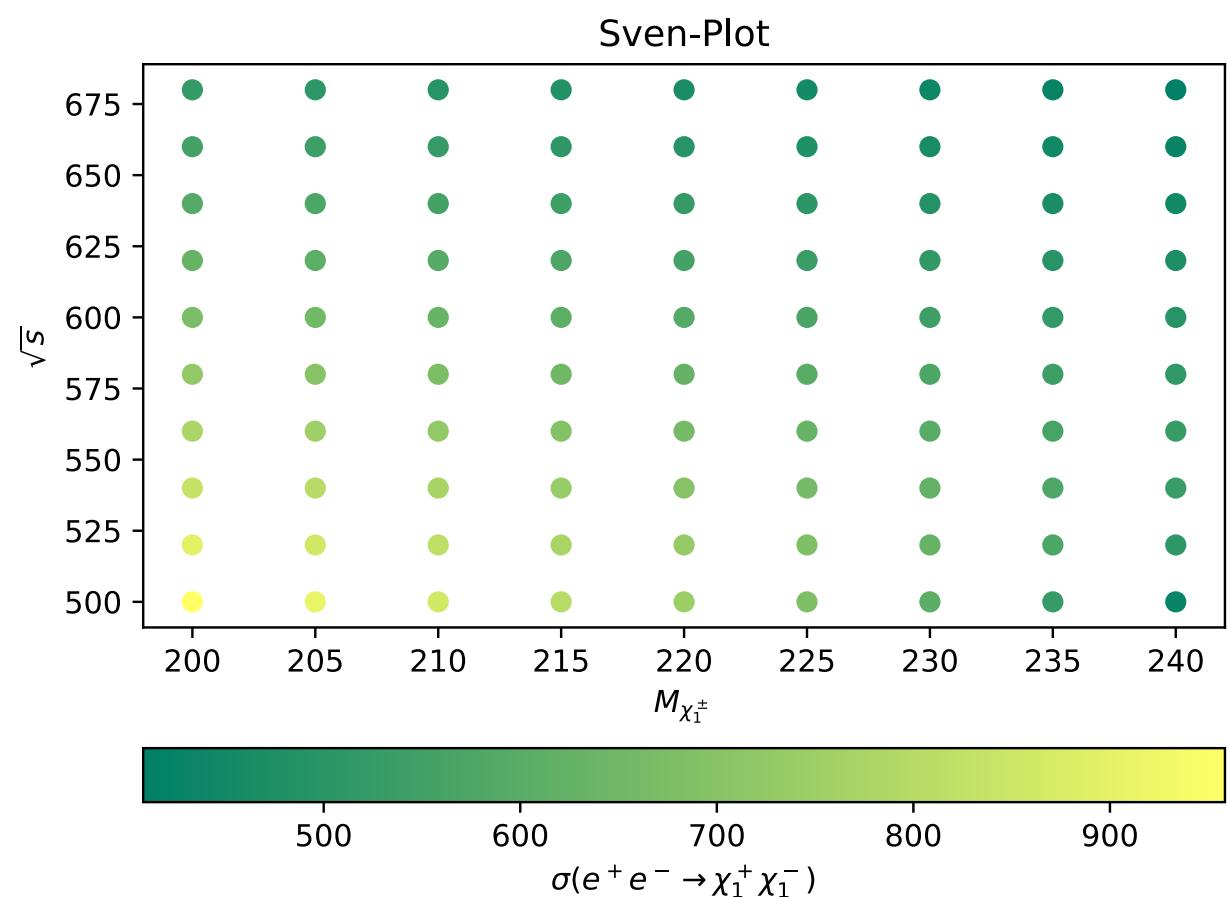


Relative distance of the upper and lower boundary of Ωh^2 produced by the scan



$$D_{\text{rel}} = \frac{M_{2,\text{var}} - M_{2,\text{true}}}{M_{2,\text{true}}}$$

Backup



Muster

- blabla