

Mirror symmetry and cluster algebras

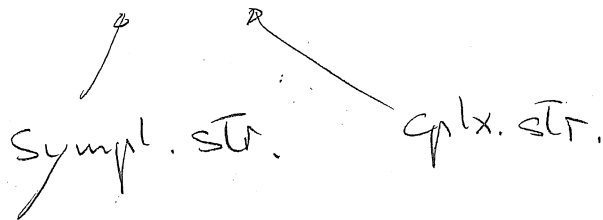
- an informal introduction -

(1)

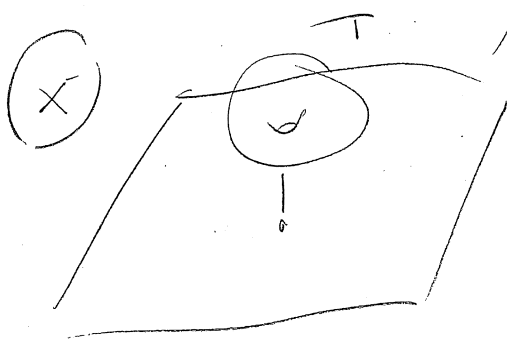
(1) A geometric constr. motivated by string th.

Conjecture: CY-mflds. come in pairs

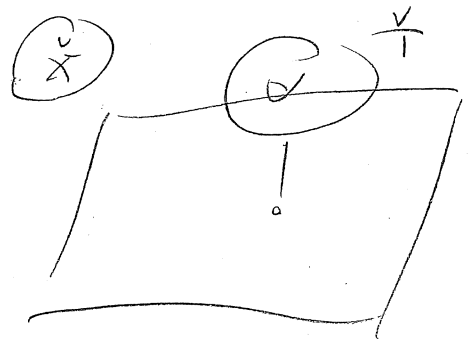
$$(X, \omega, I) \longleftrightarrow (\check{X}, \check{\omega}, \check{I})$$



Syz: Duality of Torus fibrations



B A-model



B B-model

Physics motivation: $X, \check{X} \sim$ moduli sp. of D-branes

Why? In B-model \exists D0-branes
 (\sim skyscraper sheaves, objects of $\mathcal{D}^b(\check{X})$)
 pointlike localized \rightsquigarrow probe (cplx.) geom.
 of \check{X} : Mod. sp. of D0 = \check{X}

CFT-version of MS suggests:

\exists Top. branes in A-model dual to D0

Branes in A-model: Objects in $Fuk(X)$

(L, ∇) ∇ : Unit. flat. Conn.

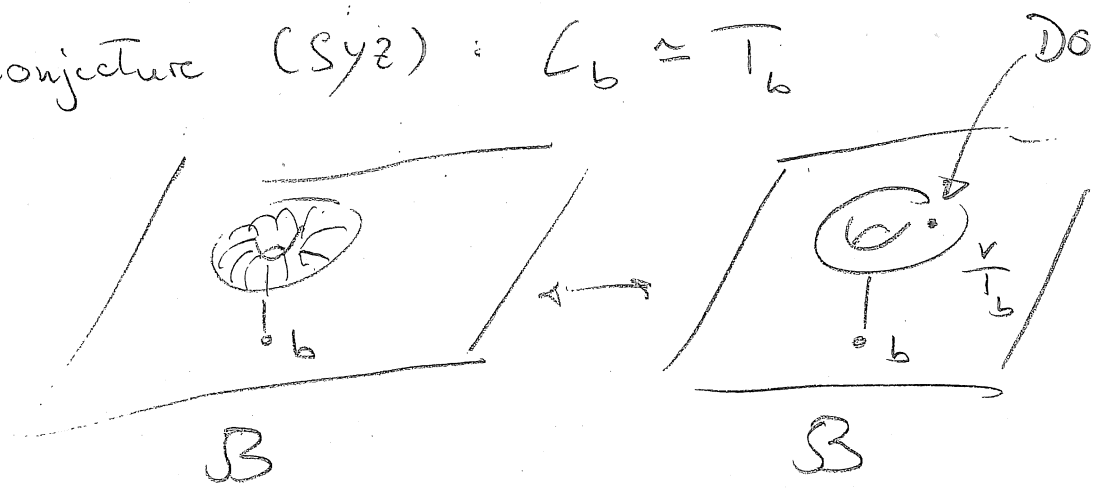
*) $\omega|_L = 0$, e.g.

$\dim_{\mathbb{R}} L = 3$

$\omega = \sum db_i \wedge d\theta^i$

e.g. $L_b = \{x \in \bar{X}; b_i = \text{const.}\}$

Conjecture (SyZ): $L_b \cong T_b$

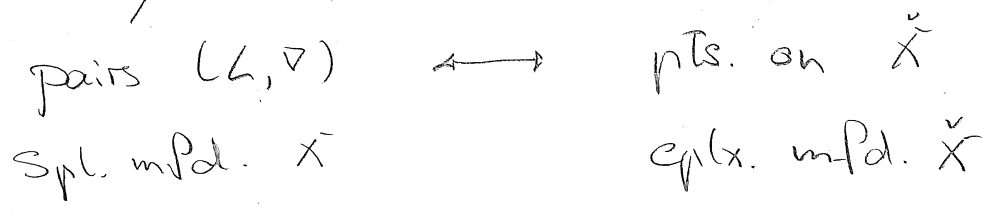


Mod. sp. of D0 on A-side:

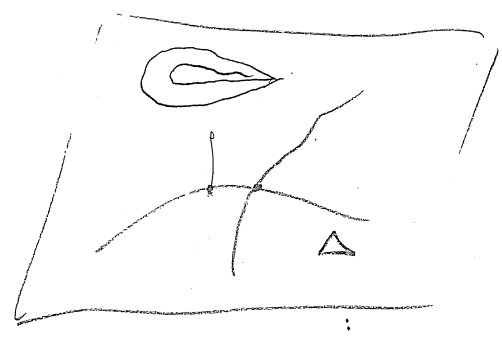
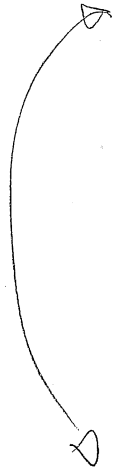
- Choice of $b_i \in \mathcal{B}$, $b_i \stackrel{(*)}{\sim} e^{-\int_{\beta_i} \omega}$
- Choice of $\nabla \sim \text{Hol}_{L_b}(\nabla) \sim (\Theta_i)$

(*) $\beta_i \in H_2(X, L_b)$: $\int_{\beta_i} \omega$ - Area of disk w/ boundary on Lagrangian L_b

Mirror symm. predicts corresp. between



Difficulty: Singular fibres, how to glue



$$\int_{\beta_i} \omega = 0$$

Divisor $\Delta \subset \mathcal{B}$

Ansatz: (Re-)construct gluing structure by defining (local?) holom. coords. for \mathcal{X} :

$$z_i = z_{\beta_i}(z, \nabla) = e^{-\int_{\beta_i} \omega} \text{Hol}_{\text{op}}(\nabla) \quad (2)$$

Problem: z_i not single valued around Δ

Refine ansatz:

- 1) Break $\mathcal{B}_0 = \mathcal{B} \setminus \Delta$ into local charts
- 2) Use (2) to define coords in local charts
- 3) Find transition fct. between chart making $\mathcal{M}_{(z, \nabla)}$ into cplx. mfd. \mathcal{X}



Very informal description of proposed answer to 1), 3): (Fukaya, Kontsevich-Schmidman, Auroux, Gross-Siebert)

Ad 1): Break \mathcal{B} into chambers separated by walls:

Wall: Locus in \mathcal{B} where holomorphic disks (with Maslov index $\mu=0$) exist

Example:

$$X = \mathbb{C}^2 \setminus D, \quad D = \{(x,y) \in \mathbb{C}^2, xy = \varepsilon\}$$

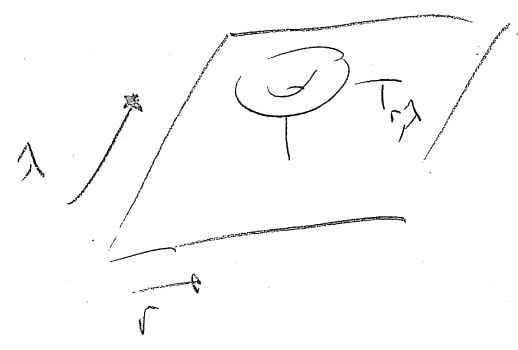
$$\omega = \frac{i}{2} (dx \wedge d\bar{x} + dy \wedge d\bar{y})$$

$$\mathbb{I}: \Omega = \frac{dx \wedge dy}{xy - \varepsilon} \quad \text{holom. (in } x,y) \text{ vol. form}$$

Lemma:

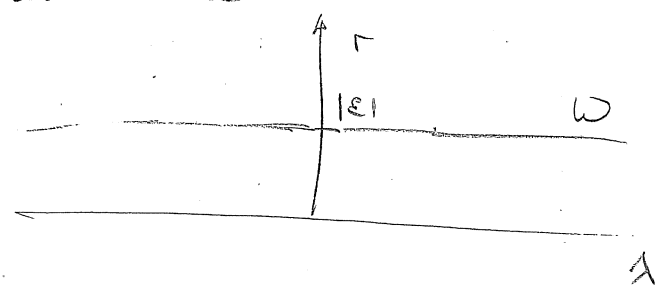
$$T_{r,\lambda} = \{(x,y) \in \mathbb{C}^2, |xy - \varepsilon| = r, \mu_S(x,y) = \lambda\}$$

are (special) Lagrangian tori $\frac{1}{2}(|x|^2 - |y|^2)$



Singular at $(r,\lambda) = (|\varepsilon|, 0)$

Wall in \mathcal{B} at $r = |\varepsilon|$



Proof uses projection $p: (x,y) \rightarrow xy$

Assume there exists holom. map

$$u: (D, \partial D) \rightarrow (X, T_{r, \lambda})$$

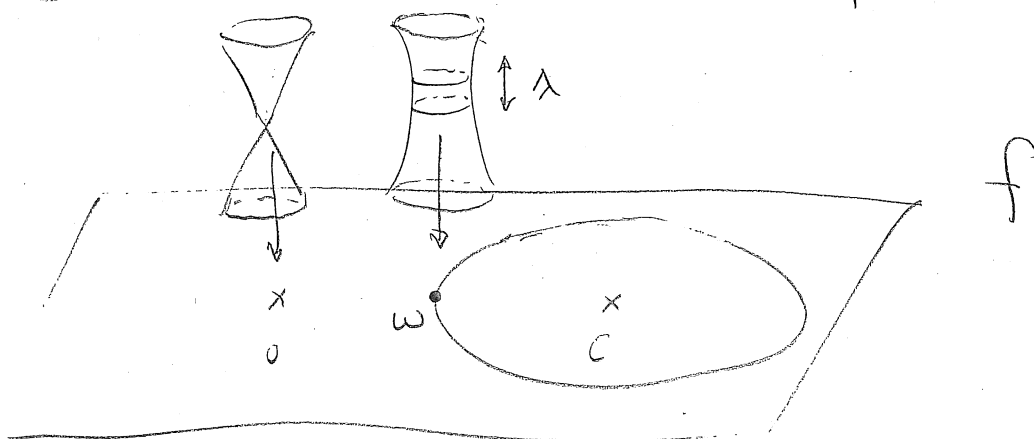
Consider $g := f \circ u - \varepsilon : D \rightarrow \mathbb{C} \setminus \{0\}$ holom.

$$\text{on } \partial D : |g| = r = \text{const.}$$

Maximum principle: Since g does not vanish on D , Max and Min of g on ∂D

$$\Rightarrow g = \text{const.}$$

So $D \sim$ fiber of f over pt. w



Fibres of $f \sim S^1 \times \mathbb{R}$ for $r \neq |\varepsilon|$

" " " \sim union of discs $r = |\varepsilon|$.

\Rightarrow Wall (\exists hole disk with μ -index 0) at $r = |\varepsilon|$.

Coords. in chamber I ($r > |\varepsilon|$) ?

Consider disks Touching $xy = \varepsilon$ ($\rightsquigarrow \mu=2$)

\exists Two types of such disks

$$(\sim D(r_1) \times \{y\}, \{x\} \times D(r_2))$$

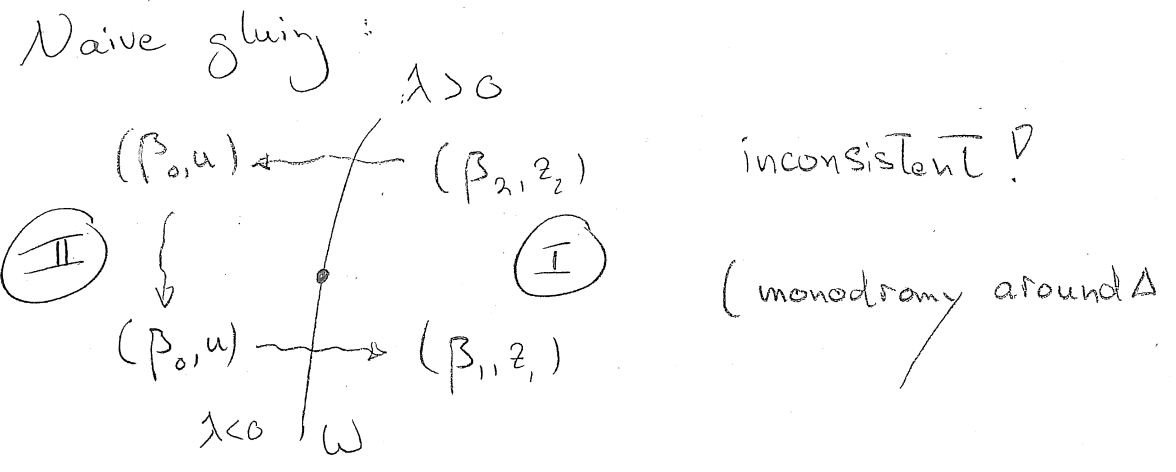
defining classes β_1, β_2 in $H_2(X, \mathbb{Z})$
 \leadsto coords. $z_i = z_{\beta_i} = e^{-\int_{\beta_i} \omega} \text{Hol}_{\beta_i}(\nabla)$

*) ω : coordinate associated to $\mu=0$ class α

Chamber II ($r < |z|$): Supplement by ω

Only one such class $\beta_0 \leadsto$ coord. u ✓

In order to define gluing, continue disks / coordinates across wall:



Way out: Define corrected gluing:

$$\lambda > 0 \quad u = z_2 (1 + \omega)$$

$$\lambda < 0 \quad u = z_1 (1 + \bar{\omega}')$$

Noting that $\omega = z_1 / z_2 \leadsto$ consistency ✓

Tip of an iceberg:

- Corrected gluing clusters-like ✓
- More generally:

$$z_0^{\text{I}} = z_0^{\text{II}} f_{\nu, \alpha}(z_\alpha) \quad \left(\begin{array}{l} \text{count of} \\ (\mu=0)\text{-disks} \end{array} \right)$$

More conceptual approach: Theta-fcts.
(here: superpotential)

$$\mathcal{J}_L(L, \mathcal{D}) := \sum_{\beta \in H_2(X, L)} N_\beta z_\beta(L, \mathcal{D})$$

where/how u allowed to touch singularities at \mathcal{D}

(virtual) count of holom. discs $u: (D, \partial D) \rightarrow (X, L)$

Claim: \mathcal{J} : globally defined function on cluster variety defd. by corrected gluing.

\rightsquigarrow Ring of alg fcts. on \check{X} ,
 \check{X} as alg. variety \check{V}