

Proposal for a micromagnetic standard problem for materials with Dzyaloshinskii-Moriya interaction

Marijan Beg^{1,2}, Leoni Breth², David Cortés-Ortuño², Ryan A. Pepper², Thomas Kluyver², Gary Downing², Thorsten Hesjedal³, Peter Hatton⁴, Tom Lancaster⁴, Geetha Balakrishnan⁵, Ondrej Hovorka², and Hans Fangohr^{1,2,*}



¹European XFEL GmbH, Holzkoppel 4, 22869 Schenefeld, Germany

²Faculty of Engineering and the Environment, University of Southampton, Southampton SO17 1BJ, United Kingdom

³Department of Physics, University of Oxford, Oxford OX1 3PU, United Kingdom

⁴Department of Physics, University of Durham, Durham DH1 3LE, United Kingdom

⁵Department of Physics, University of Warwick, Coventry CV4 7AL, United Kingdom

*email: hans.fangohr@xfel.eu



Introduction

- Simulations emerge as the **third pillar** of research and development in academia and industry, and correctness of simulation tools needs to be ensured.

- Magnetic **simulations with Dzyaloshinskii-Moriya interaction** (DMI) are becoming increasingly popular after the discovery of magnetic skyrmions and their promising features.

- **Standard problems** [1, 2, 3] are used to **test and provide confidence for the correctness** of newly developed and existing micromagnetic simulation tools.

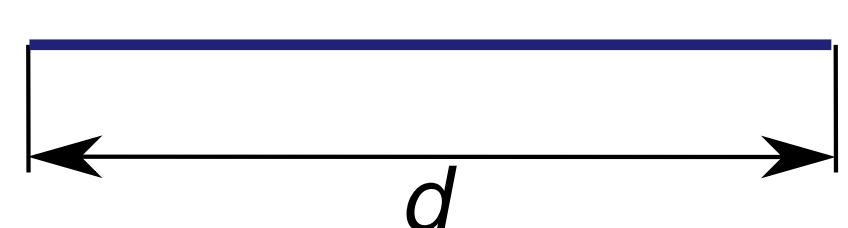
- There is **no set of standard problems** that can be used to test micromagnetic simulation tools that include Dzyaloshinskii-Moriya interaction.

- In this work, **we present simple 1D and 2D problems** where their solutions can be compared to the simulation results to support practical computational micromagnetics.

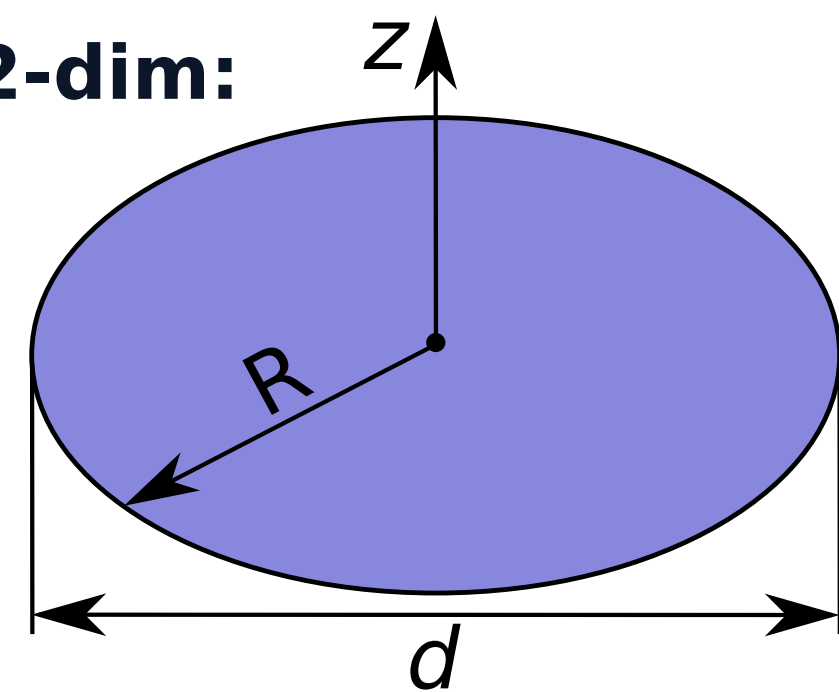
Methods

- **Geometries and material parameters:**

1-dim:



2-dim:



FeGe [4]:

$M_s = 384$ kA/m

$A = 8.78$ pJ/m

$D = 1.58$ mJ/m²

- **Hamiltonian:**

$$w = A(\nabla \mathbf{m})^2 + D \mathbf{m} \cdot (\nabla \times \mathbf{m}) - K(\mathbf{m} \cdot \mathbf{u})^2$$

↖ exchange
↖ Dzyaloshinskii-Moriya
↖ uniaxial anisotropy

- We run simulations using **OOMMF** [5] via our Python interface **JOOMMF** [6, 7].

- We use **our implementation of bulk DMI** extension for OOMMF [8].

- **No assumption about translational invariance** of magnetisation in any direction.

References

- [1] <http://www.ctcms.nist.gov/~rdm/mumag.org.html>
- [2] M. Najafi et al. *J. Appl. Phys.* **11**, 113914 (2009).
- [3] A. Baker et al. *J. Magn. Magn. Mater.* **421**, 428 (2017).
- [4] M. Beg et al. *Sci. Rep.* **5**, 17137 (2015).
- [5] M. J. Donahue and D. G. Porter, OOMMF User's Guide, Version 1.0, Interag. Rep. NISTIR 6376, Natl. Inst. Stand. Technol. Gaithersburg, MD, 1999.
- [6] M. Beg et al. *AIP Advances* **7**, 056025 (2017).
- [7] <http://joommf.github.io>
- [8] <https://github.com/joommf/oommf-bulk-dmi>
- [9] S. Rohart and A. Thiaville. *Phys. Rev. B* **88**, 184422 (2013)
- [10] <https://github.com/fangohr/paper-2017-dmi-standard-problem>

1D problems

Quasi-ferromagnetic state

- Due to the specific boundary conditions, magnetisation tilts at the edges of the one-dimensional sample (quasi-ferromagnetic).

- We solve the boundary value problem [9] using the shooting method.

- We include uniaxial anisotropy in the (0, 0, 1) direction.

$$\frac{d^2\theta}{dx^2} = \frac{\cos\theta \sin\theta}{\Delta^2}$$

$$\frac{d\theta}{dx} = -\frac{D}{2A}$$

Helical state

- We have only symmetric exchange and DMI in the Hamiltonian.

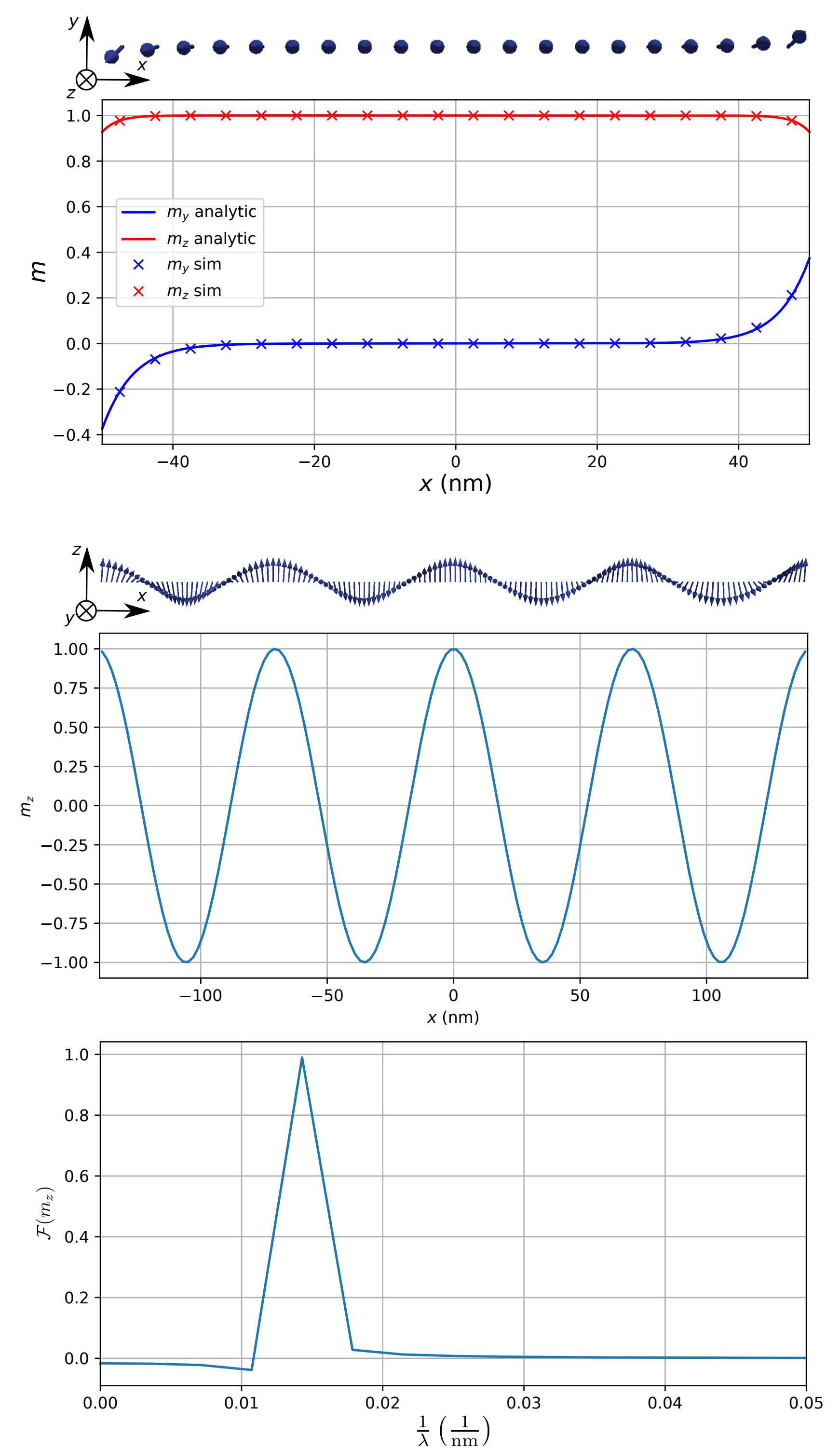
- Due to a mutual competition, the helical state is formed.

- We know that the helical period should be:

$$\lambda = 4\pi\sqrt{\frac{A}{D}} = 70 \text{ nm}$$

- By computing the Fourier transform, we can obtain the helical period.

- The helical period strongly depends on the length of the one-dimensional sample because the magnetisation configuration must always satisfy the specific boundary conditions.



2D problem

Isolated skyrmion

- We have only symmetric exchange, DMI, and uniaxial anisotropy energy terms in the Hamiltonian.

- In a two-dimensional disk sample with 50 nm radius, an isolated skyrmion is formed.

- We solve the boundary value problem [9] using the shooting method:

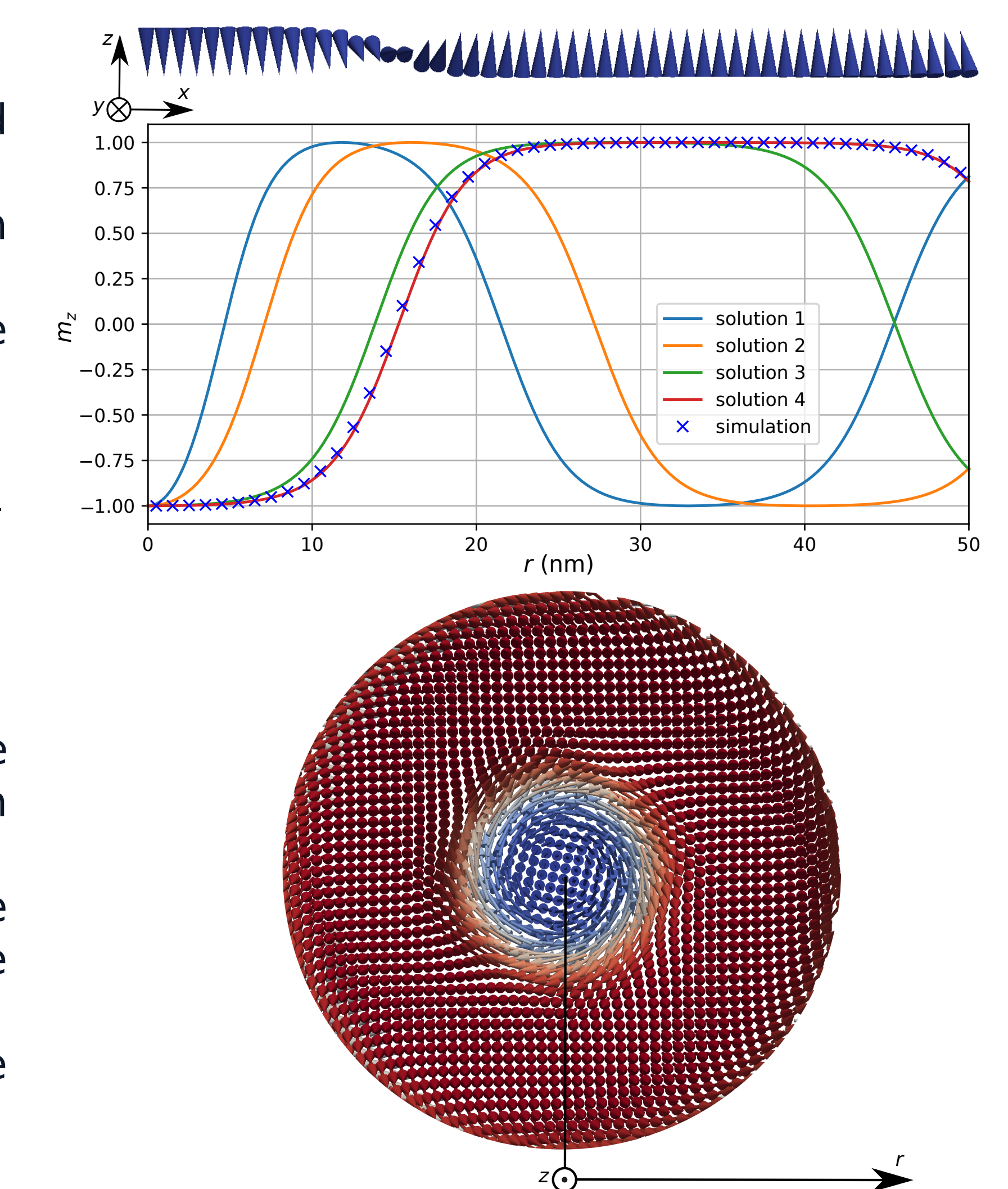
$$\frac{d^2\theta}{dr^2} = -\frac{1}{r} \frac{d\theta}{dr} + \left(\frac{1}{r^2} + \frac{1}{\Delta^2}\right) \frac{\sin 2\theta}{2} - \frac{2 \sin^2 \theta}{\xi r}$$

$$\frac{d\theta}{dr} = -\frac{D}{2A}$$

- We perform shooting at $m_z = -1$ and $R = 0$ because we assume the skyrmion orientation to point down in core.

- The magnetisation at the core points in the negative z direction and then rotates in a Bloch-type wall configuration to the periphery.

- There is additional tilting of magnetisation at the boundary due to the specific boundary conditions.



Summary

- We implemented the OOMMF extension module for simulating bulk Dzyaloshinskii-Moriya interactions [8].

- We collect a set of simple problems that can be used to test new and existing micromagnetic simulation tools with DMI tools effectively.

- We provide the full calculation of the semi-analytical solution and the numerical solution (computed with OOMMF) in public Jupyter Notebooks [10].