Dr. Till Bargheer Due: 09.07.2019

→

5.1. Heisenberg Spin Chain: Direct Diagonalization (4 points)

Consider the Heisenberg spin chain with periodic boundary conditions and Hamiltonian

$$\mathcal{H} = \sum_{j=1}^{L} (\mathcal{I}_{j,j+1} - \mathcal{P}_{j,j+1}).$$
(5.1)

Compute the spectrum of eigenvalues of \mathcal{H} (energies) by direct diagonalization for the cases specified below. M denotes the number of up spins.

- a) Compute the spectrum for a spin chain of length L = 3 and arbitrary number M of spin flips. How do the eigenstates organize into $\mathfrak{su}(2)$ multiplets? The generators of $\mathfrak{su}(2)$ are $Q^{\alpha} = \sum_{i} \sigma_{i}^{\alpha}/2$, $\alpha = x, y, z$, and $Q^{\pm} = Q^{x} \pm iQ^{y}$.
- b) Restrict to cyclic states, i.e. identify all states that are equivalent under cyclic permutations of the spin chain sites. Compute the spectrum for the states with L = 4, M = 2, and for L = 6, M = 2, 3.

5.2. Heisenberg Spin Chain: Bethe Equations (4 points)

The Bethe equations for the $XXX_{1/2}$ Heisenberg spin chain read

$$\left(\frac{u_k + i/2}{u_k - i/2}\right)^L = \prod_{\substack{j=1\\j \neq k}}^M \frac{u_k - u_j + i}{u_k - u_j - i} , \qquad k = 1, \dots, M.$$
(5.2)

Each solution to these equations (such that all finite u_k are distinct) defines an eigenstate of the Heisenberg Hamiltonian with M up spins (magnons). The energy E and momentum P are given by

$$E = \sum_{k=1}^{M} \left(\frac{i}{u_k + i/2} - \frac{i}{u_k - i/2} \right), \qquad e^{iP} = \prod_{k=1}^{M} \frac{u_k + i/2}{u_k - i/2}.$$
 (5.3)

a) Use the Bethe equations to compute the energy spectrum for L = 3 and $M \le 1$. States with M > L/2 are obtained from states with $M \le L/2$ by flipping all spins. Compare to the results of problem 5.1 a). How are the $\mathfrak{su}(2)$ multiplets realized?

In the following, restrict to cyclic states, i.e. require $e^{iP} = 1$.

- **b)** Compute the energy spectrum for L = 4, M = 2, and for L = 6, M = 2. Compare to the results of problem 5.1 b).
- c) Compute the energy spectrum from the Bethe equations for any L and M = 2.
- d) The solution for L = 6, M = 3 is singular. Show that the regularized rapidities

$$u_1 = \frac{i}{2} + \varepsilon + c \varepsilon^6$$
, $u_2 = -\frac{i}{2} + \varepsilon$, $u_3 = \frac{1 - 4u_1 u_2}{4(u_1 + u_2)} + d(\varepsilon)$ (5.4)

solve the Bethe equations and the condition $e^{iP} = 1$ in the limit $\varepsilon \to 0$ for a suitable constant c and function $d(\varepsilon)$. Compare to your result of problem 5.1 b).

5.3. Coordinate Bethe Ansatz for the XXZ Spin Chain (4 points)

Consider the Hamiltonian \mathcal{H} of the XXZ spin chain with periodic boundary conditions:

$$\mathcal{H} = \sum_{j=1}^{L} \mathcal{H}_{j,j+1}, \qquad \mathcal{H}_{j,k} = \frac{1}{2} \left[\sigma_j^x \sigma_k^x + \sigma_j^y \sigma_k^y + \Delta (\sigma_j^z \sigma_k^z - \mathbf{1}_j \mathbf{1}_k) \right], \qquad \sigma_{L+1} \equiv \sigma_1, \quad (5.5)$$

which acts on a spin chain of length L with spins $|\downarrow\rangle = (0,1)^{\mathsf{T}}$ and $|\uparrow\rangle = (1,0)^{\mathsf{T}}$. Here, $\sigma_j^i, i \in \{x, y, z\}$ are the Pauli matrices acting on the spin state at site j.

a) Consider a general state with a single up spin

$$|\psi_1\rangle = \sum_{k=1}^{L} f(k) |k\rangle, \qquad |k\rangle = |\downarrow\downarrow\downarrow\dots\downarrow\downarrow\uparrow\downarrow\downarrow\downarrow\dots\downarrow\downarrow\rangle.$$
(5.6)

Convert the eigenvalue equation $\mathcal{H}|\psi_1\rangle = e_1|\psi_1\rangle$ to a finite difference equation for f(k). Show that the one-magnon ansatz $f(k) = e^{ipk}$ solves the equation, and that the dispersion relation becomes $e_1(p) = 2(\cos(p) - \Delta)$.

b) Now consider states with two up spins:

$$|\psi_2\rangle = \sum_{1 \le k < \ell \le L} f(k,\ell) |k,\ell\rangle, \qquad |k,\ell\rangle = |\downarrow \dots \downarrow^k \downarrow \dots \downarrow^\ell \downarrow \dots \downarrow\rangle.$$
(5.7)

Starting with the eigenvalue equation $\langle k, \ell | \mathcal{H} | \psi_2 \rangle = e_2 f(k, \ell)$, derive two difference equations for $f(k, \ell)$ by considering the two cases $k + 1 < \ell$ and $k + 1 = \ell$. Using the two-magnon ansatz

$$f(k,\ell) = e^{ipk+iq\ell} + S(p,q) e^{iqk+ip\ell} , \qquad (5.8)$$

show that the dispersion relation is $e_2(p,q) = e_1(p) + e_1(q)$, and that the scattering phase S(p,q) must satisfy

$$S(p,q) = -\frac{1 + e^{i(p+q)} - 2\Delta e^{iq}}{1 + e^{i(p+q)} - 2\Delta e^{ip}}.$$
(5.9)

Hint: First compute the action of \mathcal{H} on neighboring spins $|\downarrow\downarrow\rangle$, $|\uparrow\uparrow\rangle$, $|\downarrow\uparrow\rangle$, and $|\uparrow\downarrow\rangle$.

c) Express S(p,q) in terms of rapidities u, v, which are related to the momenta p, q via

$$e^{ip} = \frac{u+i/2}{u-i/2} . (5.10)$$

Taking the limit $\Delta \to 1$, show that the Bethe equations $e^{ip_k L} = \prod_{j=1, j \neq k}^M S(p_j, p_k)$ for an *M*-magnon state become

$$\left(\frac{u_k + i/2}{u_k - i/2}\right)^L = \prod_{\substack{j=1\\j \neq k}}^M \frac{u_k - u_j + i}{u_k - u_j - i} .$$
(5.11)