

Holographic thermal propagator from modularity

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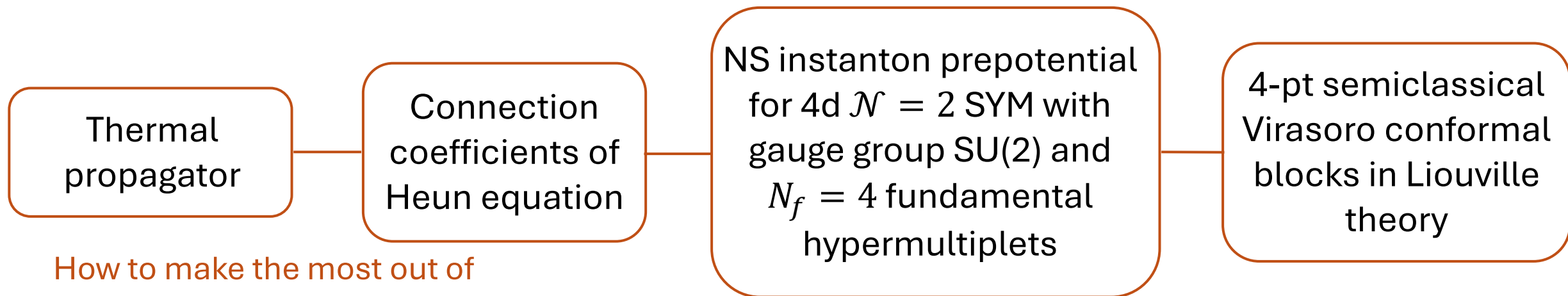
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Overview

Setting: Holographic 4d CFTs on the thermal geometry $S^1 \times \mathbb{R}^3$

Goal: Compute low T expansion of thermal two point function of scalars in momentum space

Strategy: Leverage connections with SUSY gauge theories and Virasoro conformal blocks



How to make the most out of these connections?

Thermal propagator from gauge theory

Gauge theory: $a \leftrightarrow u$, $\vec{a} = (a_0, a_t, a_1, a_\infty) \leftrightarrow \vec{m}$, $q \leftrightarrow t$

Propagator: $\omega, |\vec{k}|, \nu$ ($= 2 - \Delta_\phi$)

$$G_R = \pi^{4a_1} e^{-\partial_{a_1} F} \frac{M(a, a_1; a_\infty)}{M(a, -a_1; a_\infty)}, \quad \left| \quad q = e^{-\pi} \text{ (black brane)} \right.$$

Relation between parameters

$$\vec{a} = \left(0, \frac{i\omega}{4\pi}, \frac{\nu}{2}, \frac{\omega}{2\pi} \right) \quad \frac{\omega^2 - 2k^2}{8\pi^2} - \frac{\nu^2}{4} = u = -a^2 + a_t^2 + a_0^2 - \frac{1}{4} + t\partial_t F,$$

where $\omega, |\vec{k}|$ are in unit of T
 (at the end rescale $\omega \rightarrow \omega/T, |\vec{k}| \rightarrow |\vec{k}|/T$) $\Rightarrow$ low T = high $\omega, |\vec{k}|$ = high a, a_t, a_∞ 
 too many large parameters! So set $\omega = 0$

Computing the instanton prepotential

$$G_R = \pi^{4a_1} e^{-\partial_{a_1} F} \frac{M(a, a_1; a_\infty)}{M(a, -a_1; a_\infty)},$$

The prepotential F is an essential object, entering the propagator formula both directly and indirectly (to find a)

Typically, F computed via a *series expansion in q* (or t)

→ E.g., using the connection to conformal blocks

$$F(q, a, \vec{a}) = \dots + H(q, a, \vec{a}) \quad \text{where} \quad H(q, a, \vec{a}) \equiv \lim_{\epsilon_1 \rightarrow 1, \epsilon_2 \rightarrow 0} \epsilon_1 \epsilon_2 \log \mathcal{H}(q, \Delta, \vec{\mu}).$$

conformal block

$$\mathcal{H}(q, \Delta, \vec{\mu}) = 1 + \sum_{m,n=1}^{\infty} \frac{q^{mn} R_{m,n}(\vec{\mu})}{\Delta - \Delta_{m,n}} \mathcal{H}(q, \Delta_{m,n} + mn, \vec{\mu})$$

Zamolodchikov q -recursion

However, since low T = high a , we would like to have an expansion $H(a, q, \nu) = \sum_{n=1}^{\infty} \frac{\alpha_{2n}(q, \nu)}{a^{2n}}$

Main task: find such expansion! ↩

Prepotential as series in $1/a$

Main task: find an expansion $H(a, q, \nu) = \sum_{n=1}^{\infty} \frac{\alpha_{2n}(q, \nu)}{a^{2n}}$

They tackle the task in two steps

1) Expand Zamolodchikov q –series in a to find a (convoluted) representation of $\alpha_{2n}(q, \nu)$ as a series in q

$$\alpha_{2n} = - \lim_{\substack{\epsilon_1 \rightarrow 1 \\ \epsilon_2 \rightarrow 0}} \left[\sum_{m=1}^n \frac{(\epsilon_1 \epsilon_2)^{m+1}}{m} \sum_{\substack{\sum_{l=0}^{\infty} k_l = m \\ \sum_{l=0}^{\infty} k_l(l+1) = n}} \frac{m!}{\prod_{l=0}^{\infty} k_l!} \prod_{l=0}^{\infty} \left(\sum_{\chi=1}^{\infty} \tilde{c}_{\chi}^l q^{\chi} \right)^{k_l} \right]$$

$$\tilde{c}_{\chi}^l = \sum_{k=1}^{\chi} \sum_{\sum_{i=1}^k m_i n_i = \chi} \left(\frac{m_1 \epsilon_1 + n_1 \epsilon_2}{2} \right)^{2l} \frac{\prod_{i=1}^k R_{m_i n_i}}{\prod_{i=1}^{k-1} \Delta_{m_i n_i + m_i n_i - \Delta_{m_{i+1} n_{i+1}}}}$$

This allows for the computation of first terms in the q series for a generic α_{2n}

e.g. $\alpha_4 = - \frac{(\nu^2 - 1)^2}{8} q^2 (2 + 3q^2(7 - \nu^2) + 8q^4(9 - 2\nu^2) + q^6(191 - 45\nu^2) + \dots)$

But no hope in resumming!

Prepotential as series in $1/a$

Main task: find an expansion $H(a, q, \nu) = \sum_{n=1}^{\infty} \frac{\alpha_{2n}(q, \nu)}{a^{2n}}$

2) Use constraints coming from gauge theory [Billò, Frau, ... ; 2013-2015]

→ The coefficients α_{2n} , paired with some known q -independent functions, are quasi-modular forms

$$\alpha_{2n}(q, \nu) + f_{2n}(\nu) = \sum_{2k+4l+6m=2n} c_{k,l,m}(\nu) E_2^k(q) E_4^l(q) E_6^m(q) \equiv \tilde{\alpha}_{2n}(q, \nu)$$

↙ Eisenstein series

finite number of coefficients to be determined!

→ S-duality $\mathcal{S} : \tau \rightarrow -\frac{1}{\tau}$ and also $\mathcal{S} : a \rightarrow a_D, a_D \rightarrow -a$ $q = e^{i\pi\tau}$

$$F_{\text{NS}}(-1/\tau, a_D) = F_{\text{NS}}(a) - a \frac{\partial F_{\text{NS}}}{\partial a} \Rightarrow \frac{\partial \tilde{\alpha}_{2l}}{\partial E_2} = \frac{l-1}{12} (\nu^2 - 1) \tilde{\alpha}_{2(l-1)} + \frac{1}{6} \sum_{i=1}^{l-2} i(l-i-1) \tilde{\alpha}_{2i} \tilde{\alpha}_{2l-2i-2}$$

recursion relation for the coefficients $c_{k,l,m}$ with $k \neq 0$!

Summary of the strategy

$$G_R = \pi^{4a_1} e^{-\partial_{a_1} F} \frac{M(a, a_1; a_\infty)}{M(a, -a_1; a_\infty)}, \quad \Bigg|_{q = e^{-\pi} \text{ (black brane)}}$$

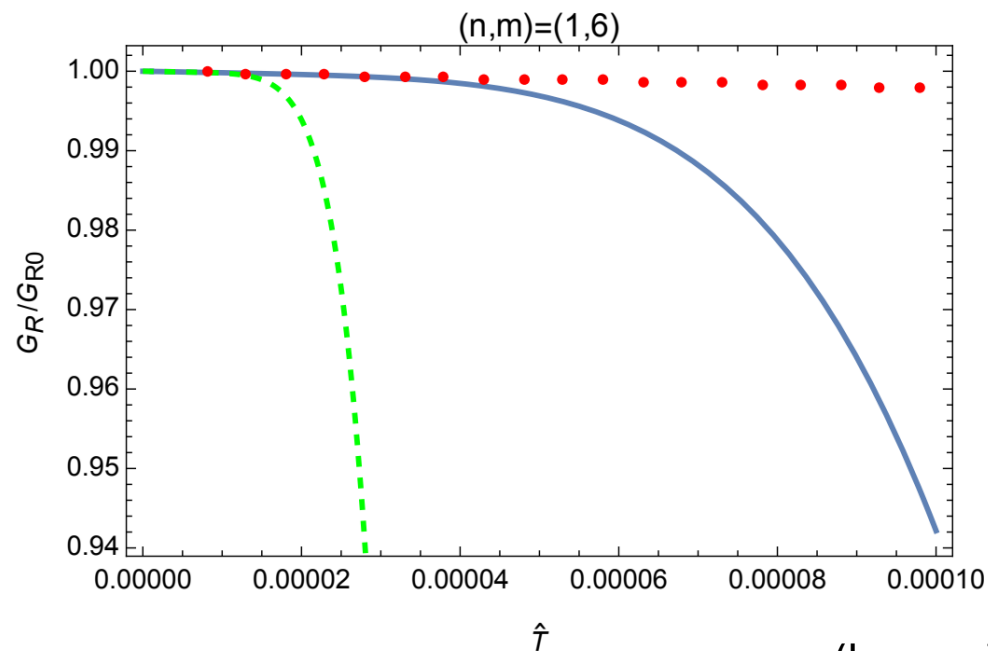
1. Compute the coefficients $\alpha_{2n}(q, \nu)$ of the large a expansion of the prepotential (also of M , but that is simpler):
 - Use Zamolodchikov q -relation to determine the leading terms in their q series
 - Feed these into the recurrence relation + matching to the expansion in the Eisenstein basis to fully fix the $\alpha_{2n}(q, \nu)$
2. Invert the Matone relation and find a in terms of $|\vec{k}|$

$$\frac{\omega^2 - 2k^2}{8\pi^2} - \frac{\nu^2}{4} = -a^2 + a_t^2 + a_0^2 - \frac{1}{4} + t\partial_t F,$$
3. Reinstate the temperature as $|\vec{k}| \rightarrow |\vec{k}|/T$ and finally expand everything in T

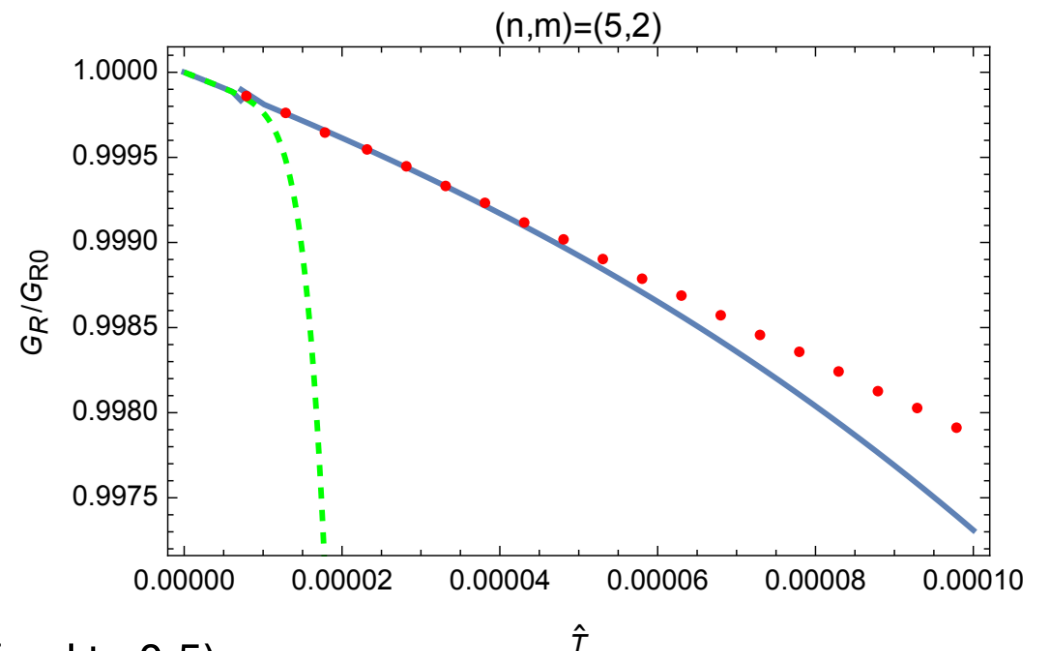
↪ In this way they get an expansion for G_R up to T^{40}

Check of the result

- The expansion is tested against a numerical solution of the propagator, computing it from the Heun equation
- As the low T expansion obtained from the propagator seems asymptotic, they have also tested extracting some Padé approximants



(here ν is fixed to 0.5)



Conclusions

- The authors have derived a **low T expansion** for the momentum space two point function of scalars in 4d holographic CFTs on $S^1 \times \mathbb{R}^3$
- In order to obtain it, they have used tools coming from **SUSY gauge theories** (quasi-modularity and S-duality) and **Virasoro blocks** (q-recursion)
- The method is efficient, and can be used to push to high orders. The series appears to be **asymptotic**, confirming previous predictions [Čeplak, Liu, Parnachev, Valach; 2024]

Outlooks

- Study the case of $\omega \neq 0$
- Perform a resurgent analysis of the divergent series
- Interpret S-duality constraints as coming from crossing in the four-point block