

High Energy Physics - Theory

[Submitted on 11 Mar 2024]

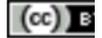
Holography and Regge Phases with $U(1)$ Charge

Giulia Fardelli, A. Liam Fitzpatrick, Wei Li

We use holography to study the large spin J limit of the spectrum of low energy states with charge Q under a $U(1)$ conserved current in CFTs in $d > 2$ dimensions, with a focus on $d = 3$ and $d = 4$. For $Q = 2$, the spectrum of such states is known to be universal and properly captured by the long-distance limit of holographic theories, regardless of whether the CFT itself is holographic. We study in detail the holographic description of such states at $Q > 2$, by considering the contribution to the energies of Q scalar particles coming from single photon and graviton exchange in the bulk of AdS; in some cases, scalar exchange and bulk contact terms are also included. For a range of finite values of Q and J , we numerically diagonalize the Hamiltonian for such states and examine the resulting spectrum and wavefunctions as a function of the dimension Δ of the charge-one operator and the central charges

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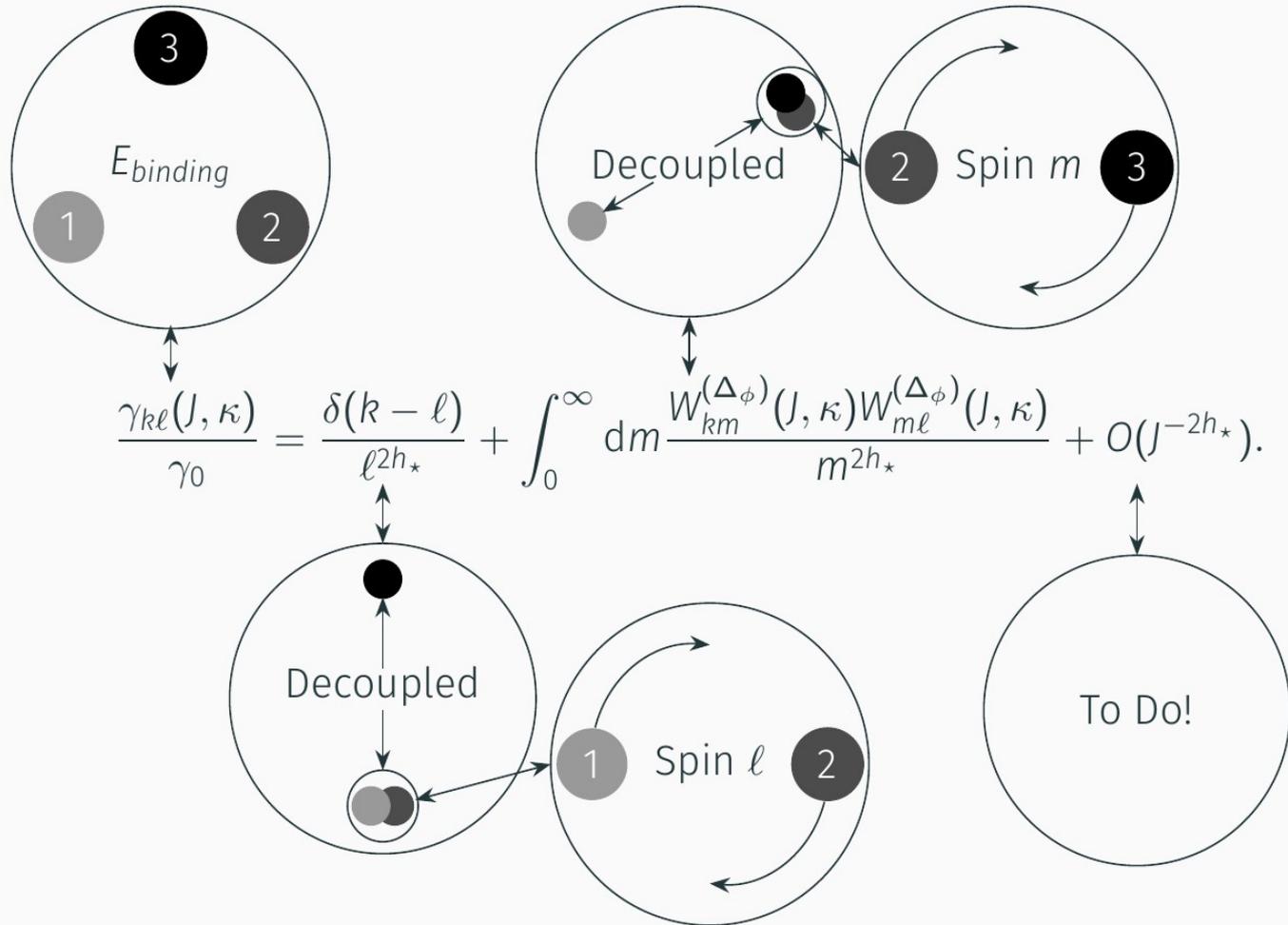
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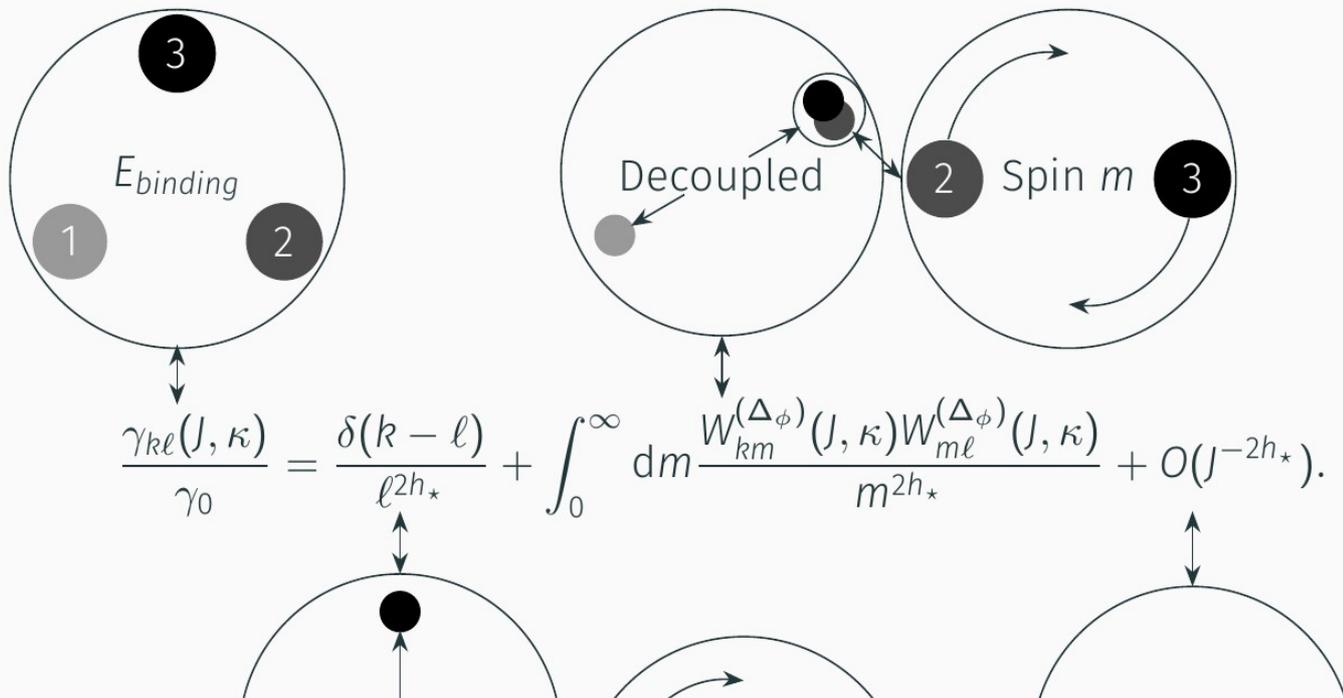
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Book

Perturbative Interactions between Three Bodies



Perturbative Interactions between Three Bodies



$$\Delta_{[\Phi, \Phi]_J} = 2\Delta + J + \frac{\Gamma(d)\Gamma^2(\Delta)}{\Gamma^2(\frac{d}{2})\Gamma^2(\Delta - \frac{d-2}{2})} \left(\frac{1}{c_{\mathcal{J}}} - \frac{2d(d+1)\Delta^2}{(d-1)^2 c_{\mathcal{T}}} \right) \frac{1}{J^{d-2}} + \dots$$

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Comments: 88 pages, many nice figures

Subjects: **High Energy Physics - Theory (hep-th)**

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(or [arXiv:2403.07079v1](https://arxiv.org/abs/2403.07079v1) [hep-th] for this version)

<https://doi.org/10.48550/arXiv.2403.07079> 

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many nice figures

$$S = \frac{1}{\kappa^2} \int d^{d+1}x \sqrt{-g} \left(\frac{R - \Lambda}{2} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - (D^\mu \phi)^\dagger (D_\mu \phi) - m^2 \phi^\dagger \phi \right),$$

where we have defined the covariant derivative

$$D_\mu = \nabla_\mu - ig_{U(1)} A_\mu.$$

$$S = \frac{1}{\kappa^2} \int d^{d+1}x \sqrt{-g} \left(\frac{R - \Lambda}{2} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - (D^\mu \phi)^\dagger (D_\mu \phi) - m^2 \phi^\dagger \phi \right),$$

where we have defined the covariant derivative

$$D_\mu = \nabla_\mu - ig_{U(1)} A_\mu.$$

such that

$$S_{\text{eff}} = \int d^{d+1}x \sqrt{-\tilde{g}} \left(-\nabla^\mu \phi^\dagger \nabla_\mu \phi - m^2 \phi^\dagger \phi - V_{\text{eff}}[\phi, \phi^\dagger] \right)$$

$$V_{\text{eff}}[\phi, \phi^\dagger] = -\frac{1}{2} \kappa A^\mu[\phi, \phi^\dagger] \mathcal{J}_\mu[\phi, \phi^\dagger] - \frac{1}{4} \kappa h^{\mu\nu}[\phi, \phi^\dagger] \mathcal{T}_{\mu\nu}[\phi, \phi^\dagger],$$

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$$D_\mu = \nabla_\mu - ig_{U(1)} A_\mu.$$

$$\gamma^{(d,Q)}(J) = \int d^d x \sqrt{-\tilde{g}} \, J \langle \Psi | V_{\text{eff}}[\phi, \phi^\dagger] | \Psi \rangle_J \equiv J \langle \Psi | \mathcal{V}_{\text{eff}}[\phi, \phi^\dagger] | \Psi \rangle_J.$$

such that

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$$\langle \ell_2 \ell_4 | V_{\text{eff}}[\phi, \phi^\dagger] | \ell_1 \ell_3 \rangle = -2 \left(\frac{(\kappa g_{U(1)})^2}{2} \overbrace{\langle \ell_2 \ell_4 | A[\phi, \phi^\dagger] \cdot \mathcal{J}[\phi, \phi^\dagger] | \ell_1 \ell_3 \rangle} + \frac{\kappa^2}{4} \overbrace{\langle \ell_2 \ell_4 | h[\phi, \phi^\dagger] \cdot \mathcal{T}[\phi, \phi^\dagger] | \ell_1 \ell_3 \rangle} + \ell_2 \leftrightarrow \ell_4 \right),$$

$$S = \frac{1}{\kappa^2} \int d^{d+1}x \sqrt{-g} \left(\frac{R - \Lambda}{2} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - (D^\mu \phi)^\dagger (D_\mu \phi) - m^2 \phi^\dagger \phi \right),$$

where we have defined the covariant derivative

$$D_\mu = \nabla_\mu - ig_{U(1)} A_\mu.$$

$$\gamma^{(d,Q)}(J) = \int d^d x \sqrt{-\tilde{g}} \, J \langle \Psi | V_{\text{eff}}[\phi, \phi^\dagger] | \Psi \rangle_J \equiv J \langle \Psi | \mathcal{V}_{\text{eff}}[\phi, \phi^\dagger] | \Psi \rangle_J.$$

many nice figures

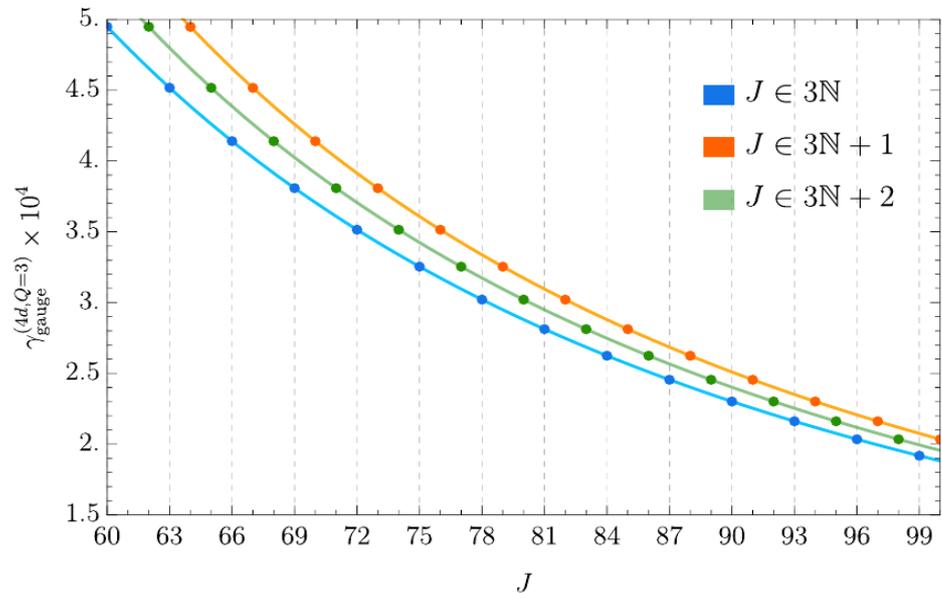
$$\begin{aligned} \langle \ell_2 \ell_4 | V_{\text{eff}}[\phi, \phi^\dagger] | \ell_1 \ell_3 \rangle &= -2 \left(\frac{(\kappa g_{U(1)})^2}{2} \langle \ell_2 \ell_4 | \overbrace{A[\phi, \phi^\dagger] \cdot \mathcal{T}[\phi, \phi^\dagger]} | \ell_1 \ell_3 \rangle \right. \\ &\quad \left. + \frac{\kappa^2}{4} \langle \ell_2 \ell_4 | \overbrace{h[\phi, \phi^\dagger] \cdot \mathcal{T}[\phi, \phi^\dagger]} | \ell_1 \ell_3 \rangle + \ell_2 \leftrightarrow \ell_4 \right), \end{aligned}$$

such that

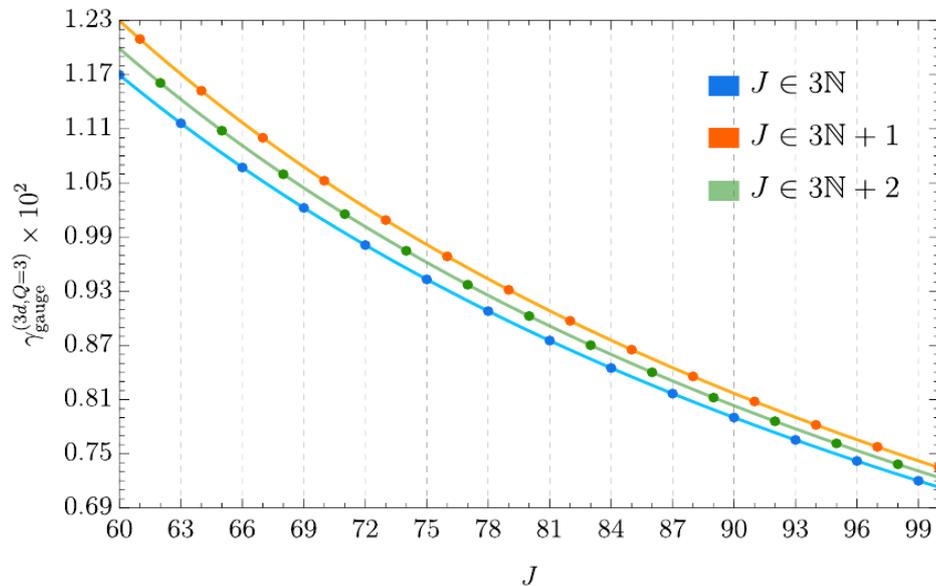
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Only Coulomb Force



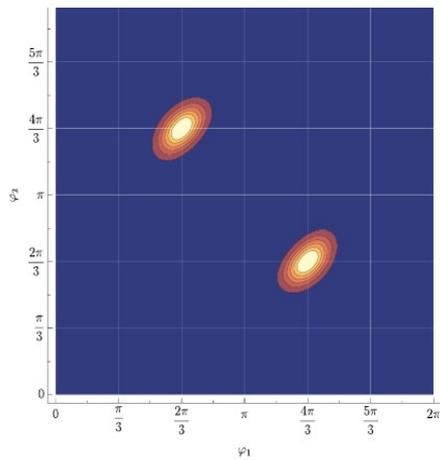
(a) $Q = 3$ $\Delta = 4$ in $4d$



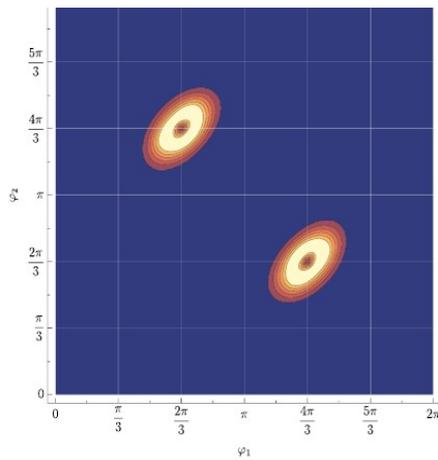
(b) $Q = 3$ $\Delta = \frac{5}{2}$ in $3d$

Figure 3. Anomalous dimensions corresponding to the lowest eigenvalue at each spin for

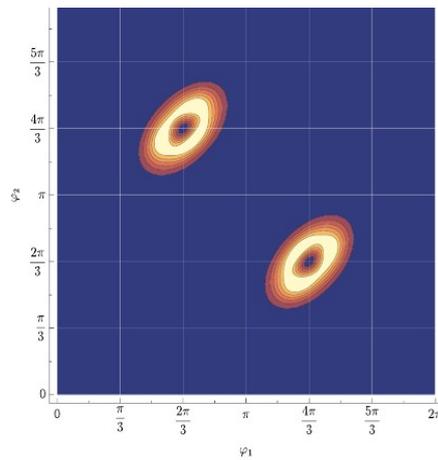
$|\Psi_{\min}(\varphi_1, \varphi_2)|$ for the 3-particle state with $\Delta = \frac{5}{2}$ in $3d$



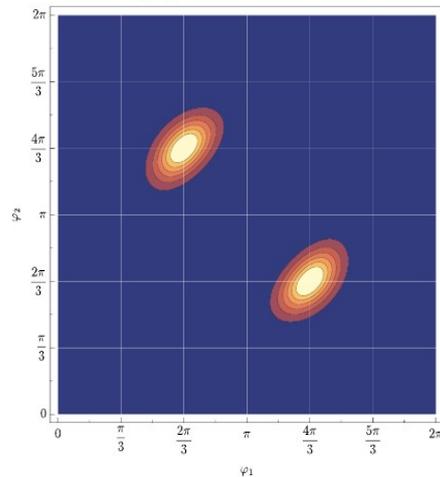
(a) $J = 99$



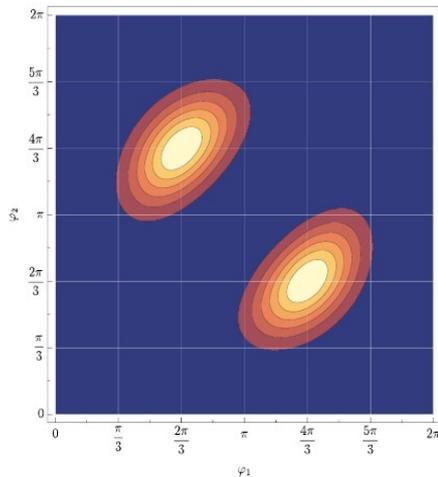
(b) $J = 98$



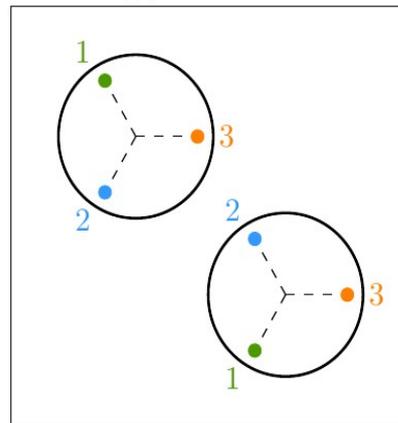
(c) $J = 97$



(d) $J = 57$



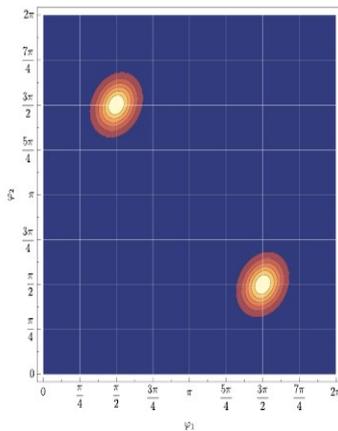
(e) $J = 18$



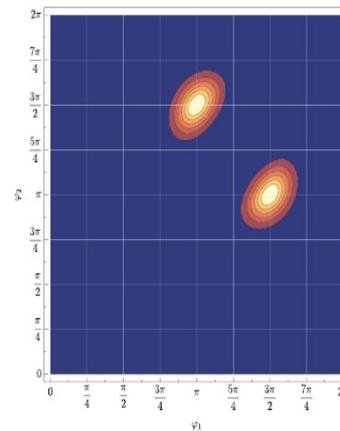
(f) Cartoon version

(Third particle is at angular position 0 on the equator)

The same pictures for four partons:

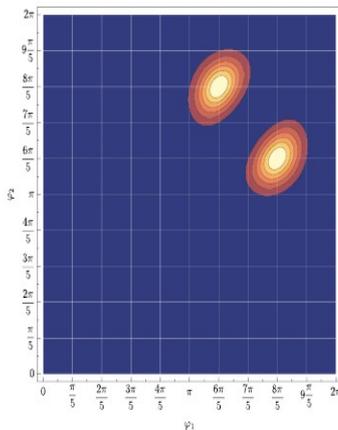


(a) $\varphi_3 = \pi$

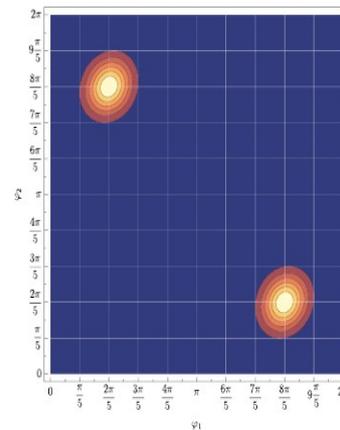


(b) $\varphi_3 = \frac{\pi}{2}$

Figure 7. $Q = 4 |\Psi_{\min}(\varphi_1, \varphi_2, \varphi_3, 0)|$ for $\Delta = \frac{5}{2}$ in $3d$ and $J = 60$.



(a) $(\varphi_3, \varphi_4) = (\frac{2\pi}{5}, \frac{4\pi}{5})$



(b) $(\varphi_3, \varphi_4) = (\frac{6\pi}{5}, \frac{4\pi}{5})$

Figure 8. $Q = 5 |\Psi_{\min}(\varphi_1, \varphi_2, \varphi_3, \varphi_4, 0)|$ for $\Delta = \frac{5}{2}$ in $3d$ and $J = 50$.

Excited states and radial resolution

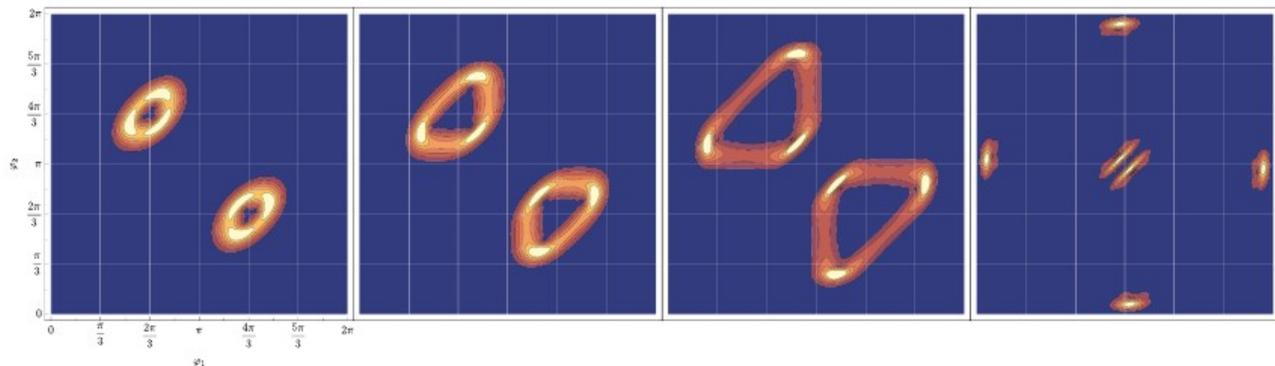
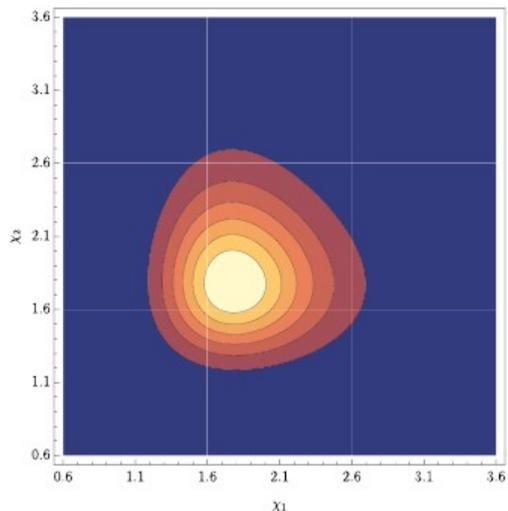
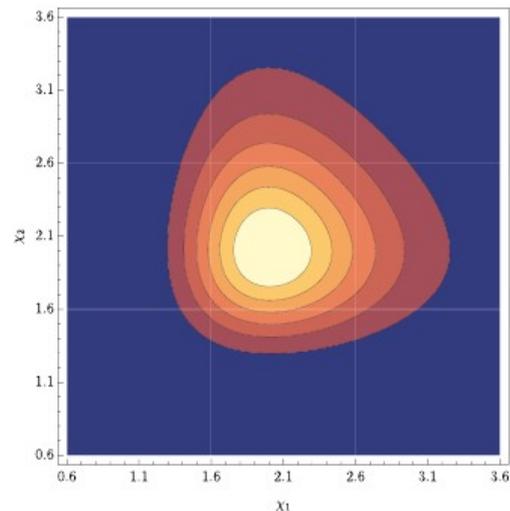


Figure 9. $|\Psi_r(\varphi_1, \varphi_2)|$ for the 3-particle state with $\Delta = \frac{5}{2}$ in $3d$ at $J = 99$ for various eigenvalues. Recall that r labels different primaries at a fixed spin, $r = 1, \dots, \mathcal{N}(3, 99) = 17$. We plot, from left to right, $r = 2, 4, 8, 16$.

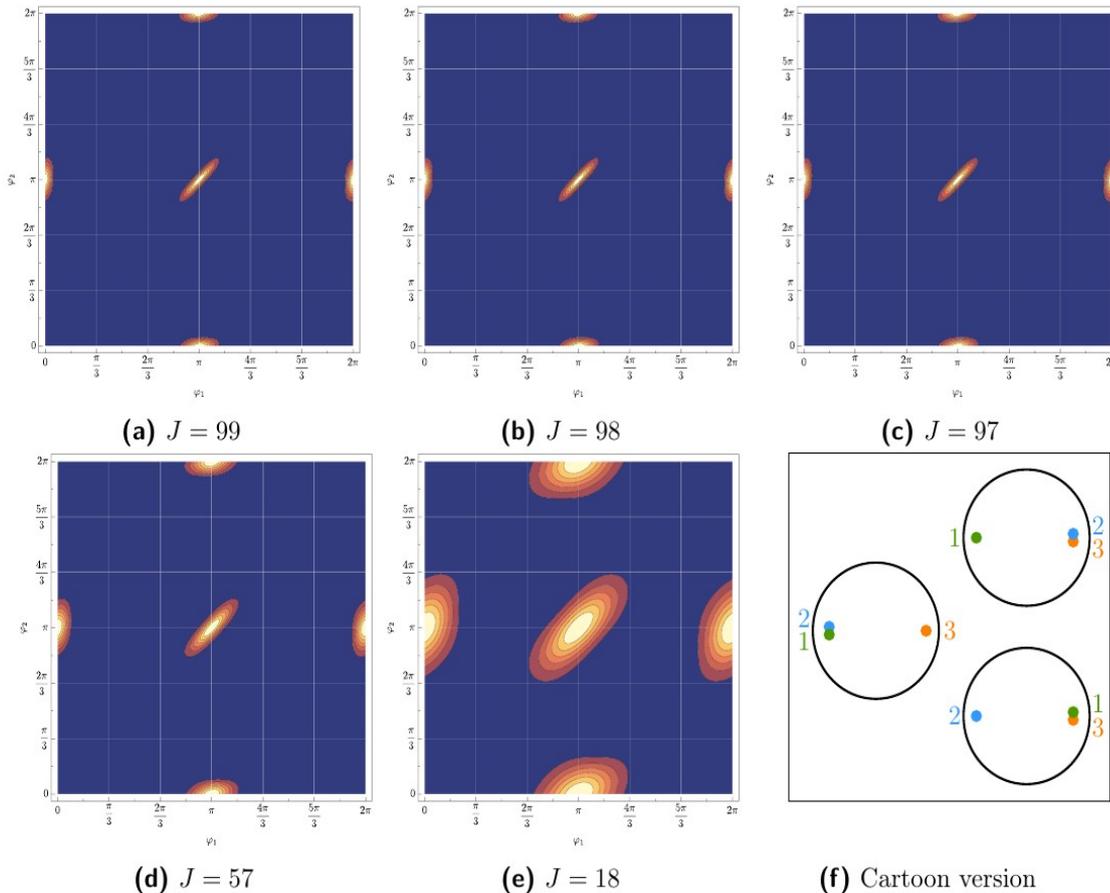


(a) $\Delta = 4$ and $\chi_3 \sim 1.78$



(b) $\Delta = \frac{5}{2}$ and $\chi_3 \sim 2$

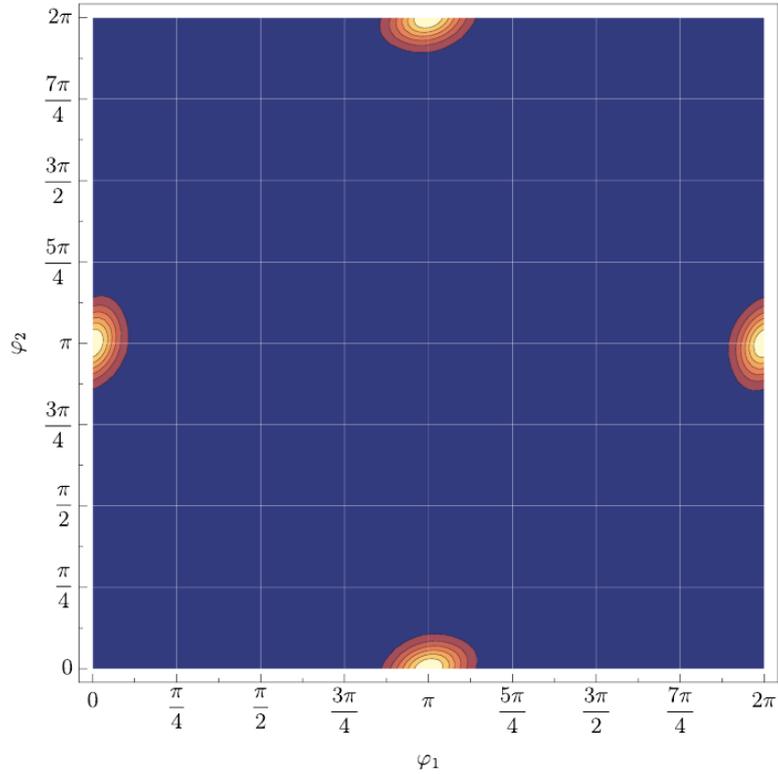
Only Gravitational Interaction



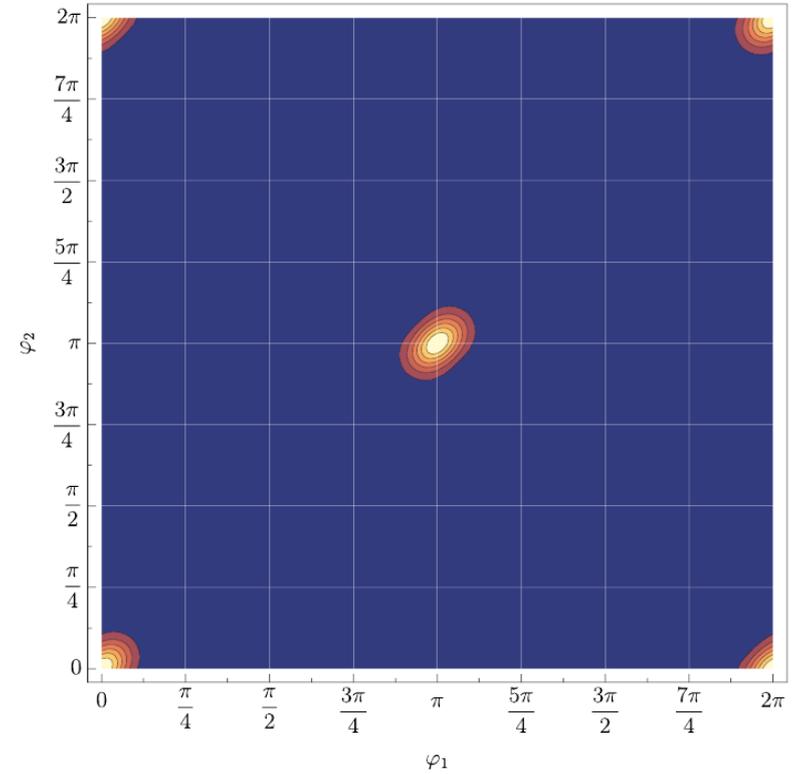
(Third particle is at angular position 0 on the equator)

Figure 16. $|\Psi_{\min}(\varphi_1, \varphi_2)|$ for the 3-particle state with $\Delta = \frac{5}{5}$ in $3d$ for different values

Four Partons

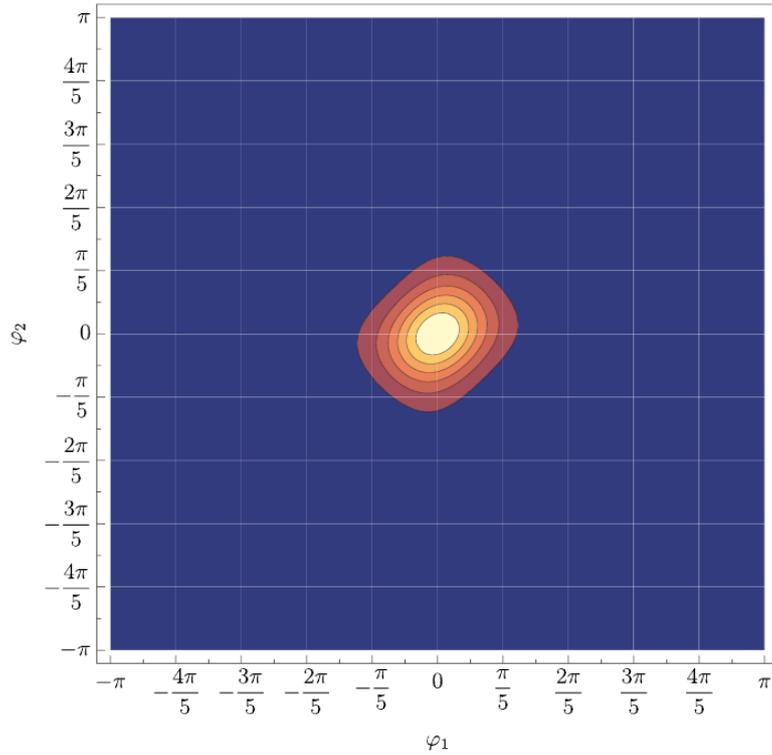


(a) $Q = 4$ $\Delta = 4$ in $4d$ and $\varphi_3 = 0$

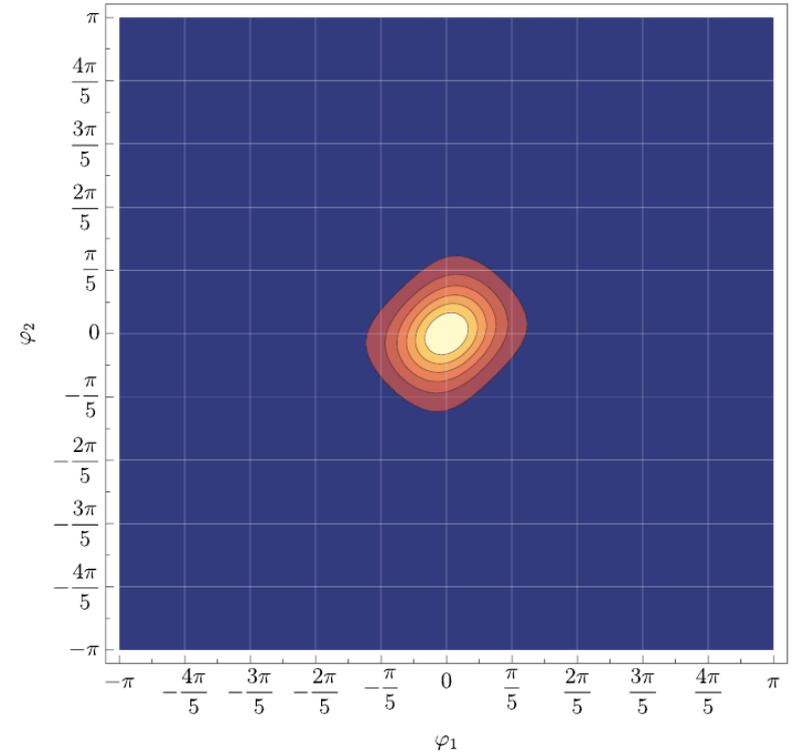


(b) $Q = 4$ $\Delta = 4$ in $4d$ and $\varphi_3 = \pi$

Five Partons



(a) $Q = 5$ $\Delta = 4$ in $4d$ and $J = 55$



(b) $Q = 5$ $\Delta = \frac{5}{2}$ in $3d$ and $J = 50$

Gauge + Gravity = Pretty Phase Diagrams!

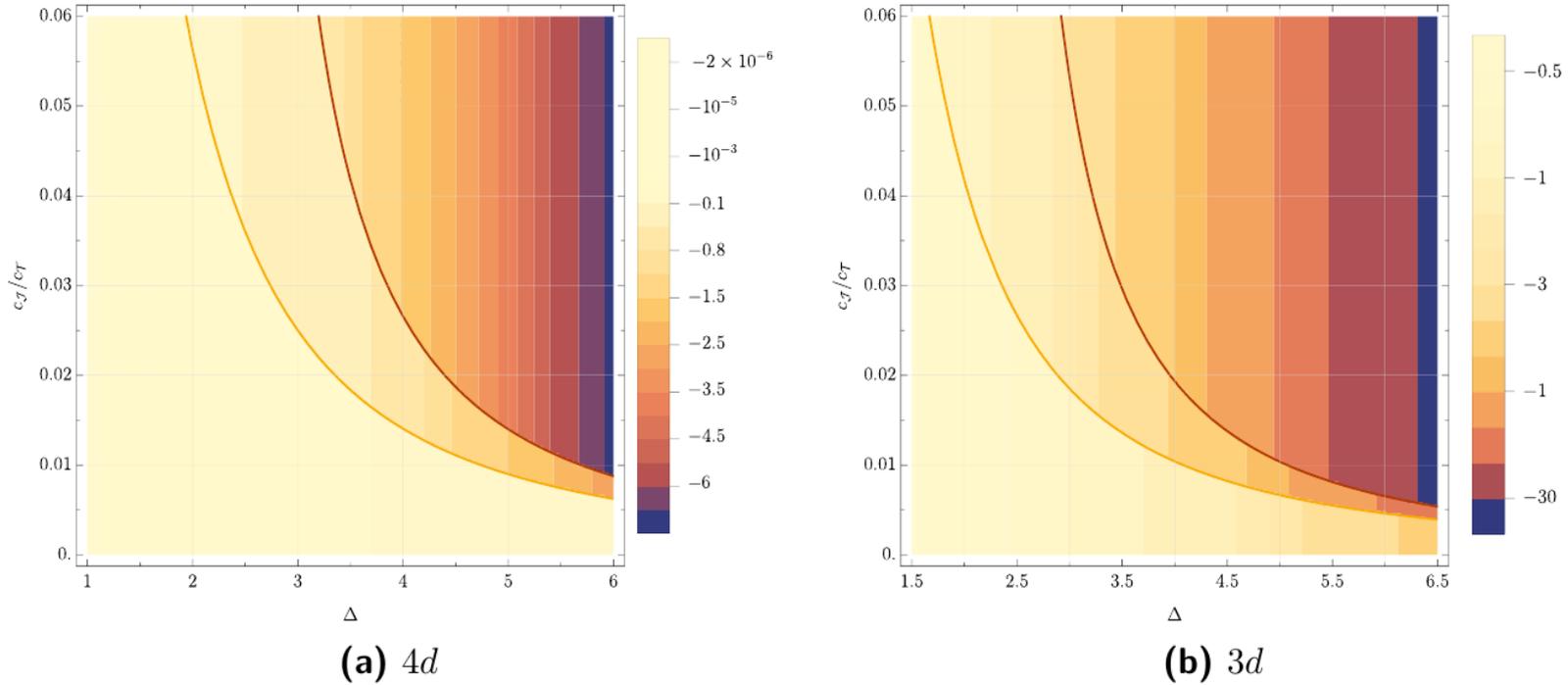


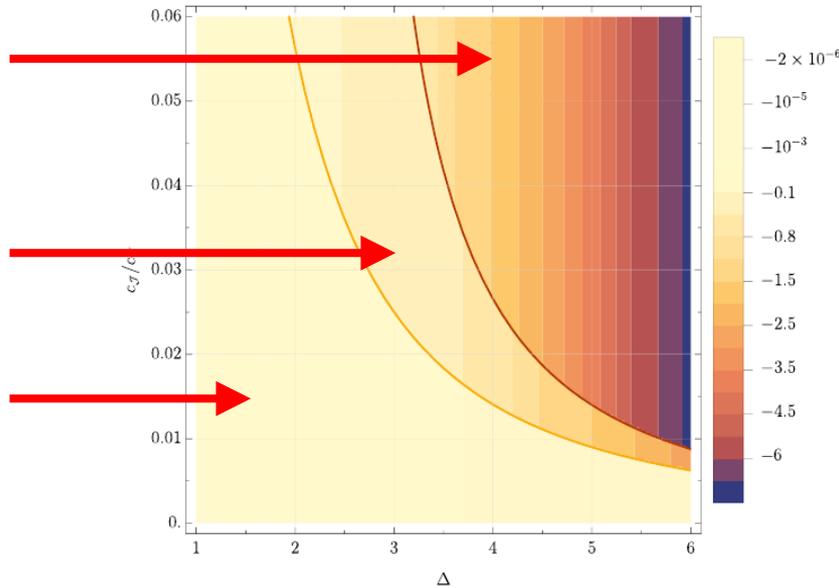
Figure 26. $\frac{\partial \gamma^{(d,Q)}}{\partial (c_J/c_T)}$ as a function of Δ and (c_J/c_T) at fixed $J = 100$ and $Q = 3$ in four and three dimensions. The lines correspond to a function $(c_J/c_T) = \beta(\Delta)$ such that the system passes from the repulsive \rightarrow attractive phase (orange line) and $a_0 = \gamma^{(d,Q=2)}(c_J/c_T)|_{\ell=2} \rightarrow a_0 = \gamma^{(d,Q=2)}(c_J/c_T)|_{\ell=0}$ (red line).

Gauge + Gravity = Pretty Phase Diagrams!

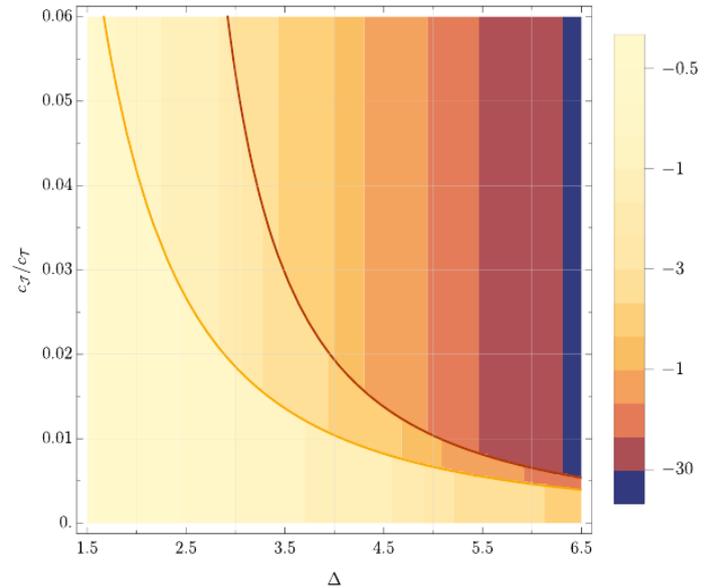
Attractive
spin 0

Attractive
spin 2

Repulsive



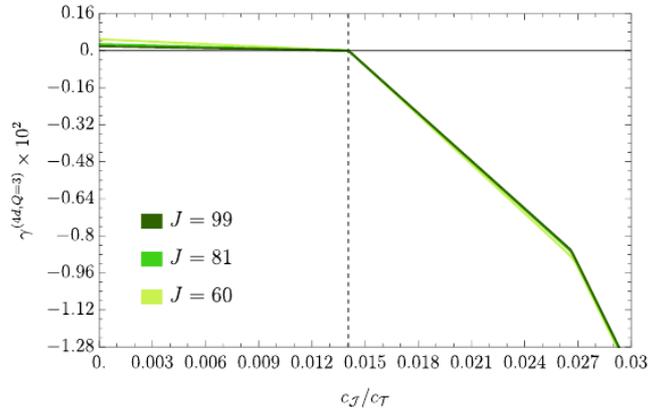
(a) $4d$



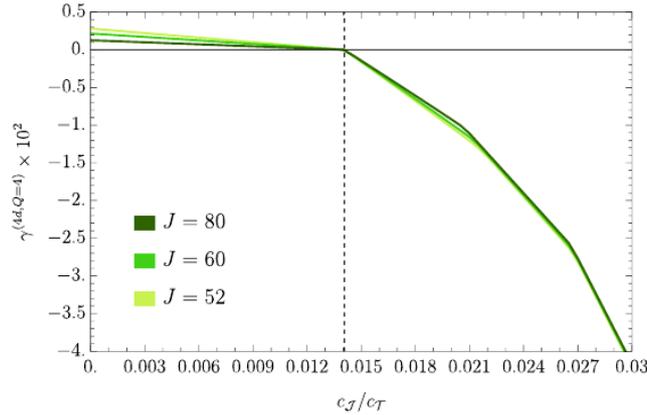
(b) $3d$

Figure 26. $\frac{\partial \gamma^{(d,Q)}}{\partial (c_{\mathcal{J}}/c_{\mathcal{T}})}$ as a function of Δ and $(c_{\mathcal{J}}/c_{\mathcal{T}})$ at fixed $J = 100$ and $Q = 3$ in four and three dimensions. The lines correspond to a function $(c_{\mathcal{J}}/c_{\mathcal{T}}) = \beta(\Delta)$ such that the system passes from the repulsive \rightarrow attractive phase (orange line) and $a_0 = \gamma^{(d,Q=2)}(c_{\mathcal{J}}/c_{\mathcal{T}})|_{\ell=2} \rightarrow a_0 = \gamma^{(d,Q=2)}(c_{\mathcal{J}}/c_{\mathcal{T}})|_{\ell=0}$ (red line).

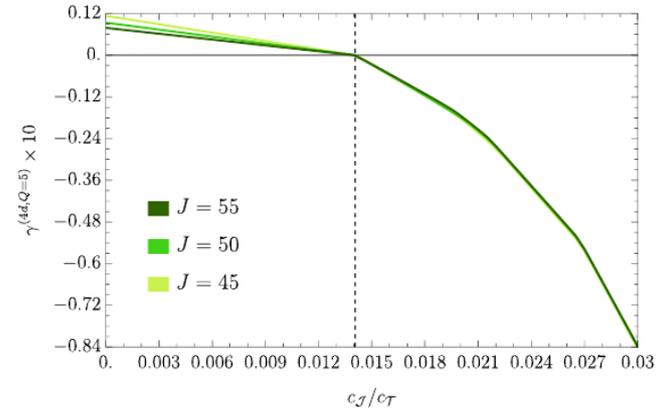
For people who don't like derivatives:



(a) $Q = 3$



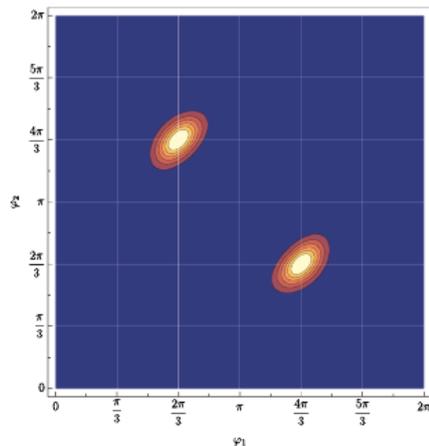
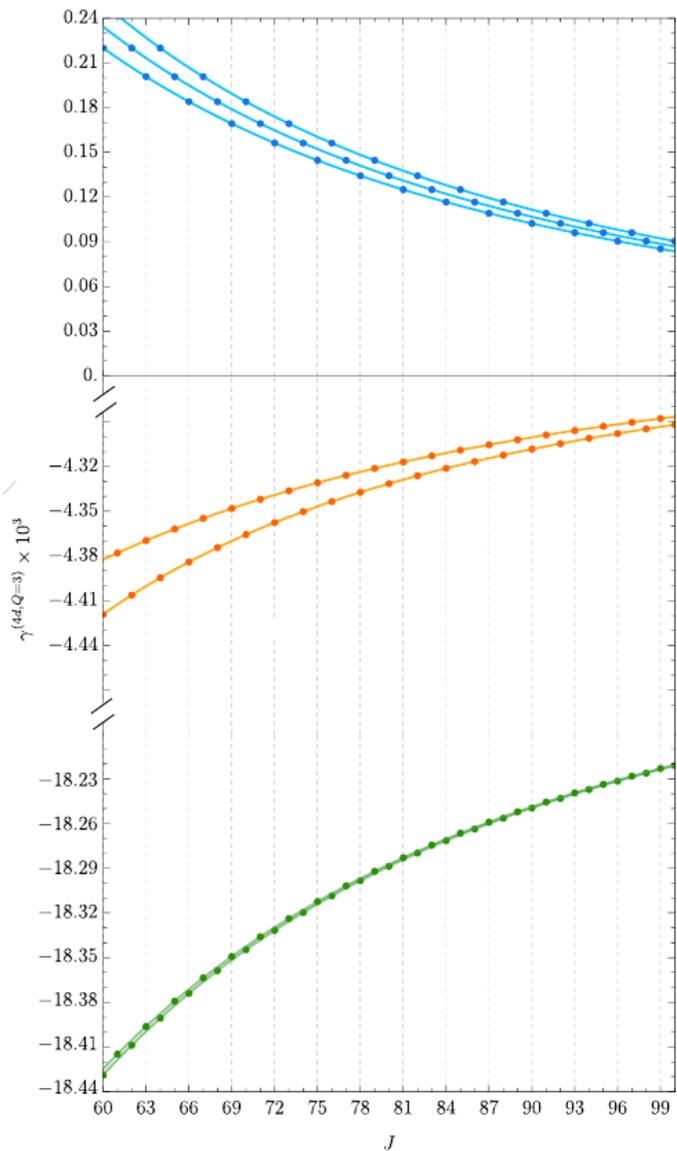
(b) $Q = 4$



(c) $Q = 5$

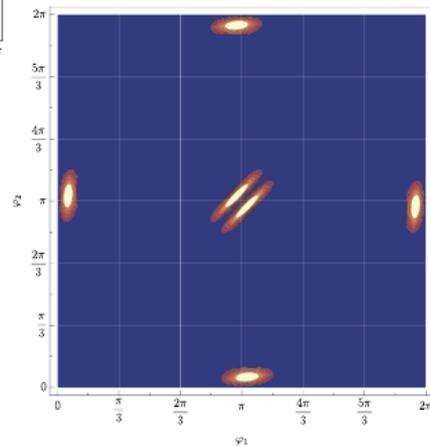
Figure 27. $4d$ anomalous dimensions for Q -particle states due the photon and graviton exchange combined at $\Delta = 4$ as a function of c_J/c_T . At fixed Q , every line corresponds to a different value of the spin. There are Q different slopes corresponding to different configurations of the partons.

Gravity gets stronger

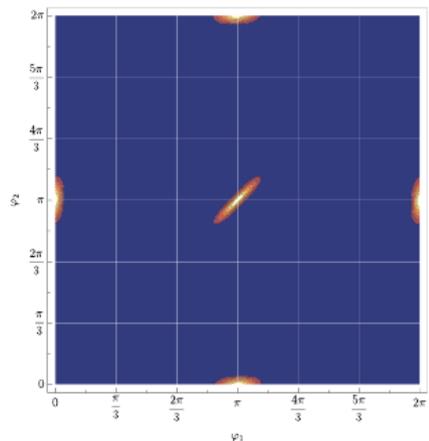


Repulsive

Attractive
spin 2



Attractive
spin 0



More Partons, more phases.

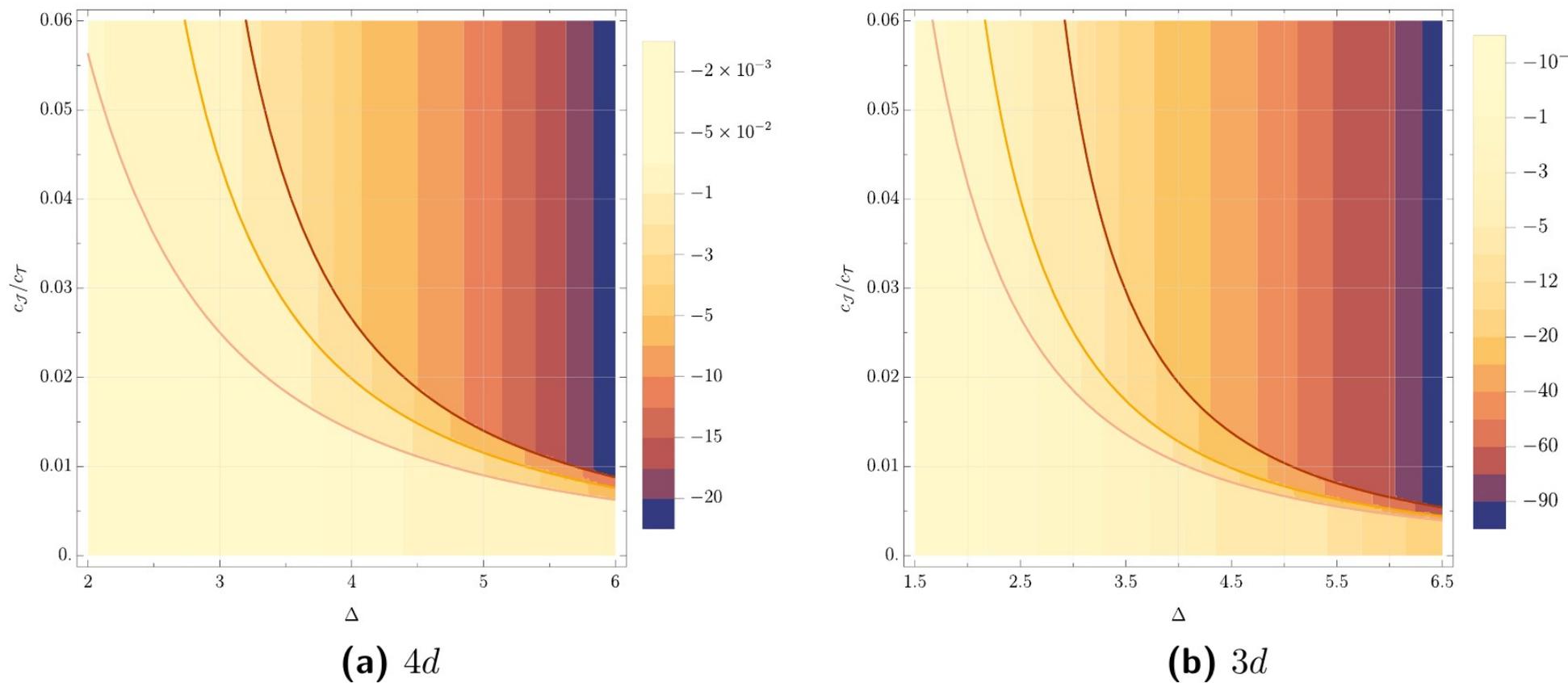


Figure 30. $\frac{\partial \gamma^{(d,Q)}}{\partial (c_J/c_T)}$ as a function of Δ and (c_J/c_T) at fixed $J = 100$ and $Q = 4$:

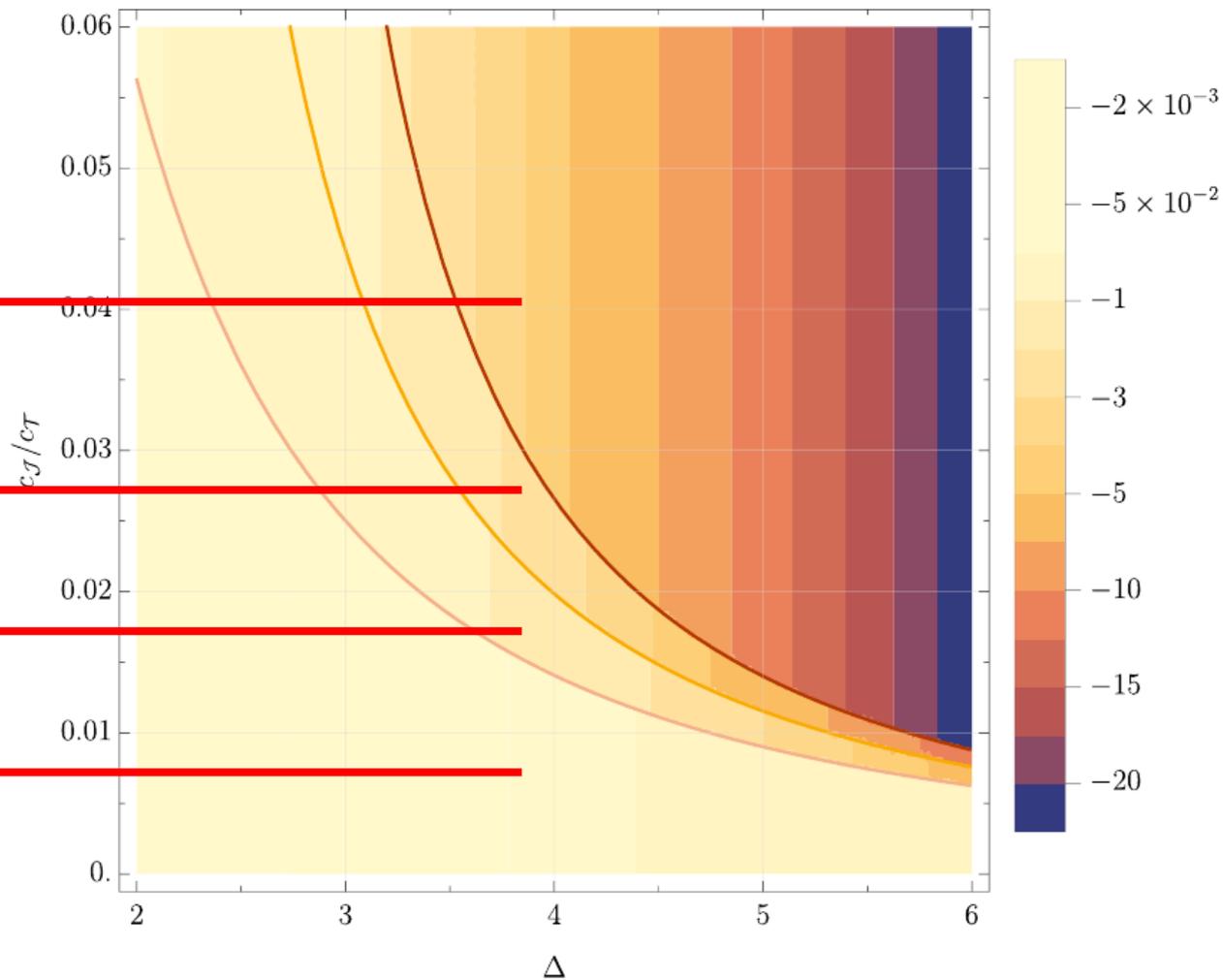
Two blobs of two partons?!

Spin one three parton blob

Spin two three parton blob

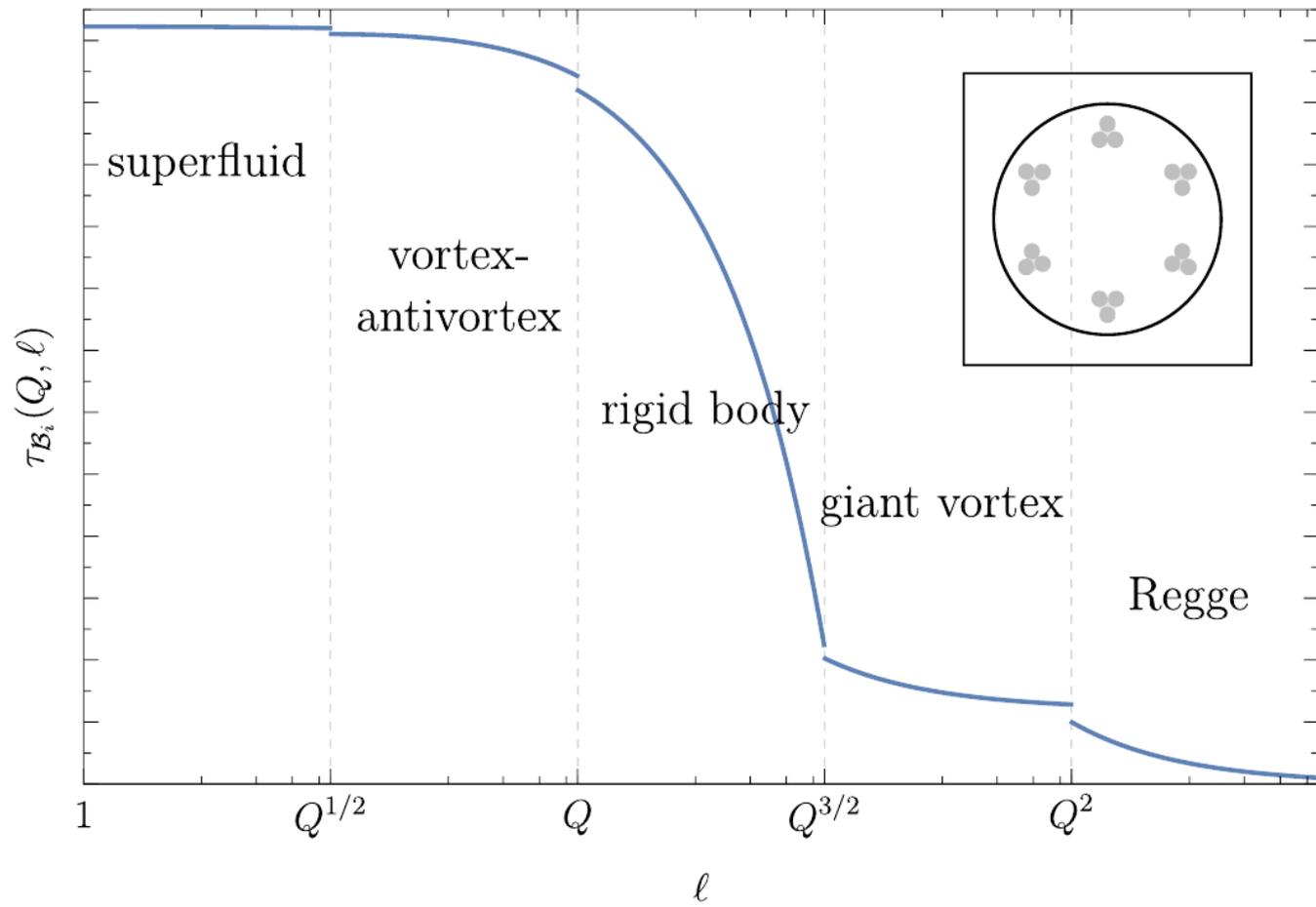
Spin three two parton blobs

Repulsive Phase

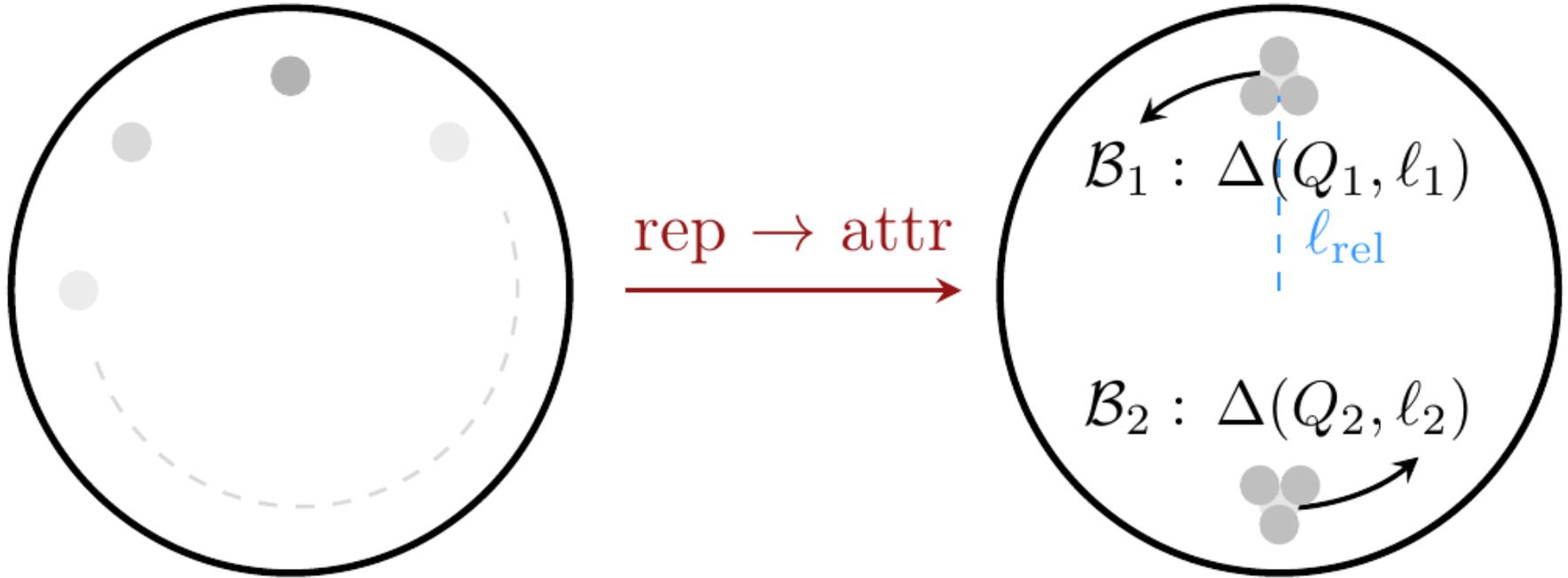


(a) $4d$

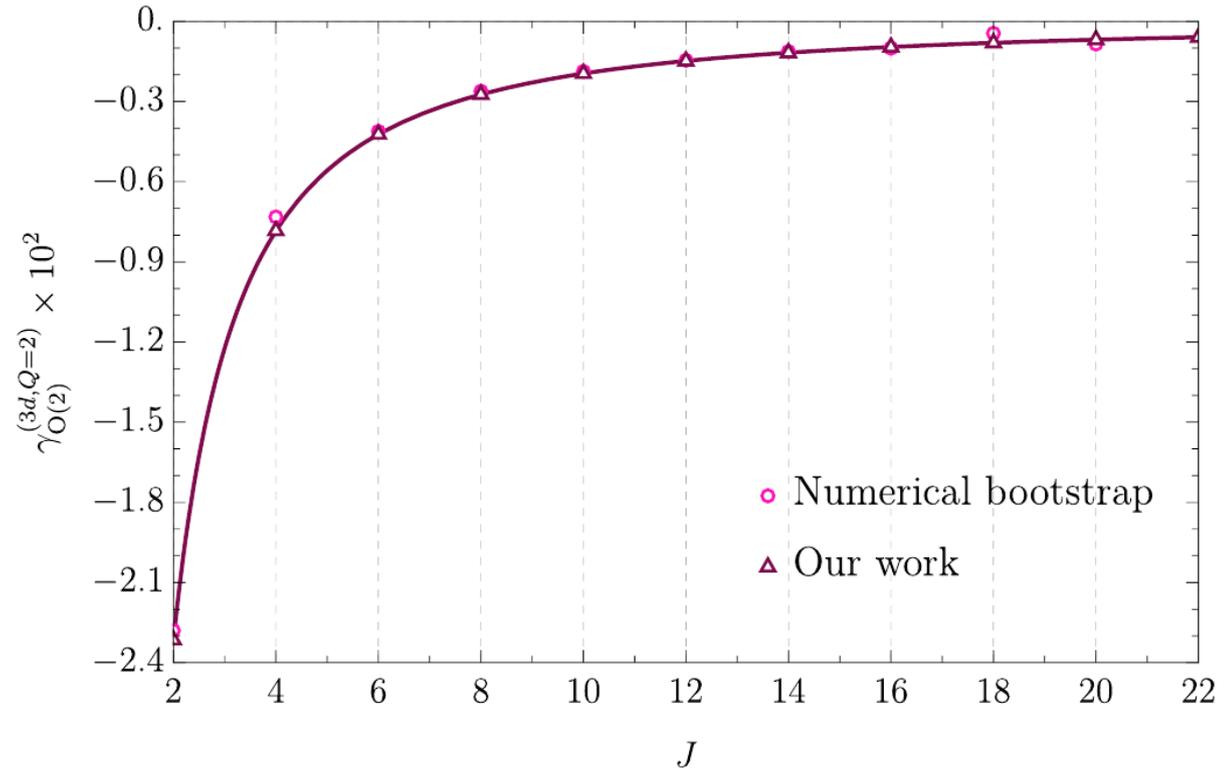
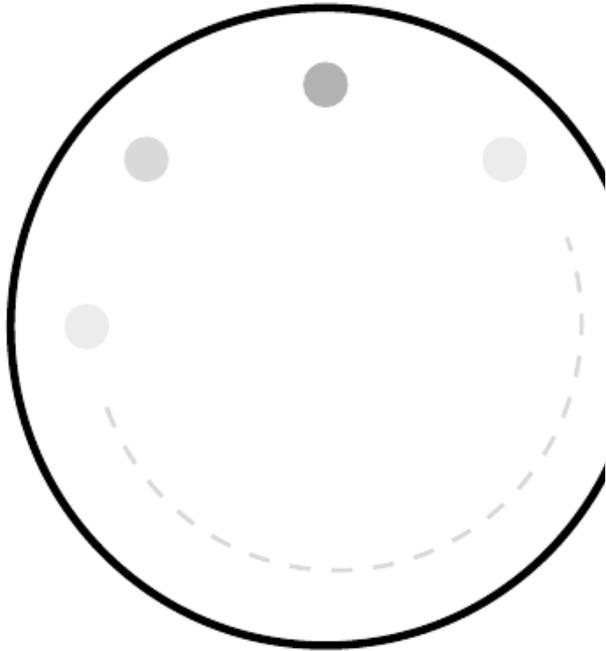
Large Charge EFT



Feeding in non perturbative data and matching known $O(2)$ model results.



Feeding in non perturbative data and matching known $O(2)$ model results.



Feeding in non p matching known

