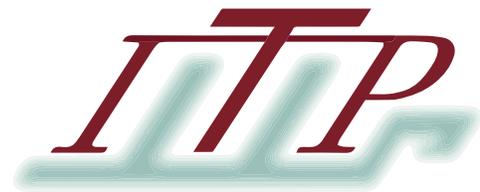


# Deciphering the Nature of the Higgs Boson

Henning Bahl



UNIVERSITÄT  
HEIDELBERG  
ZUKUNFT  
SEIT 1386

Hannover, 24.1.2024

# The Higgs mechanism in a nutshell

# The Higgs mechanism in a nutshell

- $U(1)$  model:  $\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} \Rightarrow$  massless gauge boson

# The Higgs mechanism in a nutshell

- $U(1)$  model:  $\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} \Rightarrow$  massless gauge boson
- Add complex scalar charged under the  $U(1) \rightarrow$  Abelian Higgs model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (\partial_\mu\phi^* - ieA_\mu\phi^*)(\partial^\mu\phi + ieA^\mu\phi) + m^2|\phi|^2 - \frac{\lambda}{4}|\phi|^4$$

# The Higgs mechanism in a nutshell

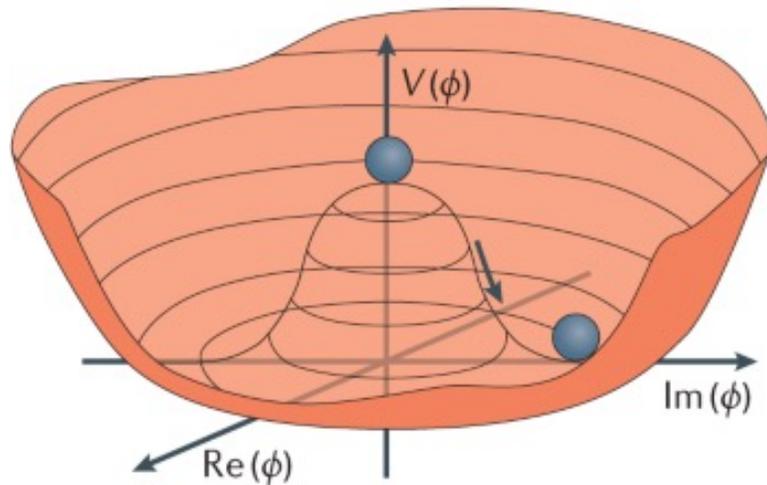
- $U(1)$  model:  $\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} \Rightarrow$  massless gauge boson
- Add complex scalar charged under the  $U(1) \rightarrow$  Abelian Higgs model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (\partial_\mu\phi^* - ieA_\mu\phi^*)(\partial^\mu\phi + ieA^\mu\phi) + \underbrace{m^2|\phi|^2 - \frac{\lambda}{4}|\phi|^4}_{-V(\phi)}$$

# The Higgs mechanism in a nutshell

- $U(1)$  model:  $\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} \Rightarrow$  massless gauge boson
- Add complex scalar charged under the  $U(1) \rightarrow$  Abelian Higgs model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (\partial_\mu\phi^* - ieA_\mu\phi^*)(\partial^\mu\phi + ieA^\mu\phi) + \underbrace{m^2|\phi|^2 - \frac{\lambda}{4}|\phi|^4}_{-V(\phi)}$$

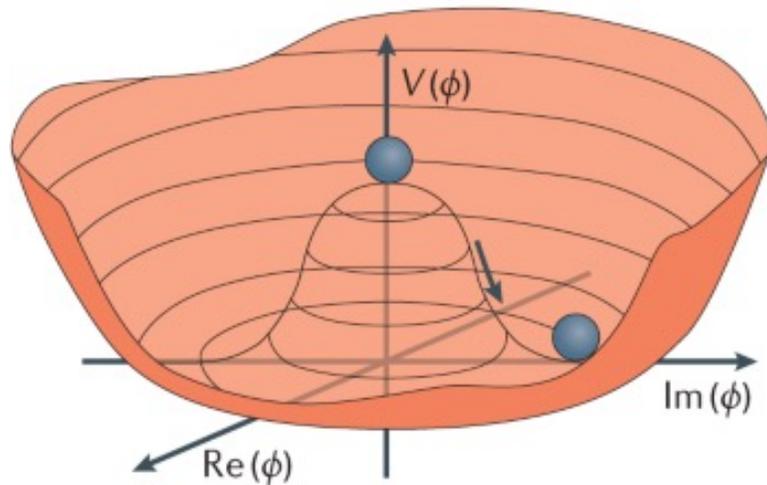


# The Higgs mechanism in a nutshell

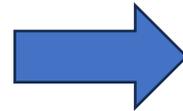
- $U(1)$  model:  $\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} \Rightarrow$  massless gauge boson

- Add complex scalar charged under the  $U(1) \rightarrow$  Abelian Higgs model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (\partial_\mu\phi^* - ieA_\mu\phi^*)(\partial^\mu\phi + ieA^\mu\phi) + \underbrace{m^2|\phi|^2 - \frac{\lambda}{4}|\phi|^4}_{-V(\phi)}$$



$$\phi(x) = \frac{1}{\sqrt{2}}(v + \sigma)e^{\frac{i\pi(x)}{F_\pi}}$$



$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{1}{2}e^2v^2A_\mu A^\mu + \dots$$

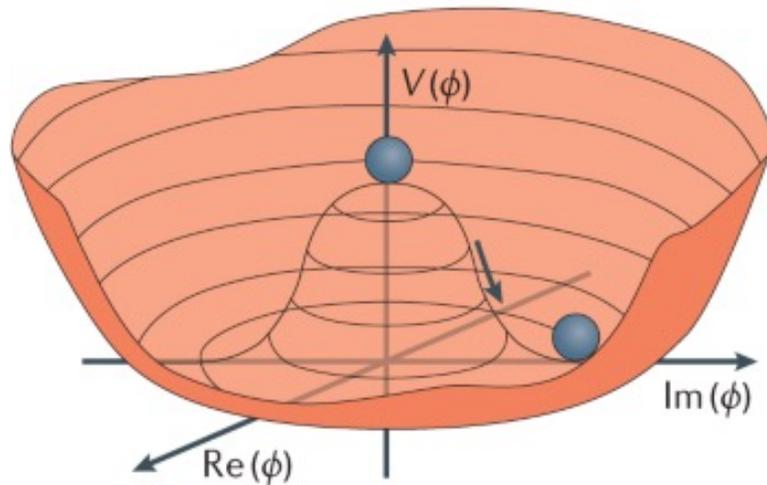
- generated mass term for gauge boson

# The Higgs mechanism in a nutshell

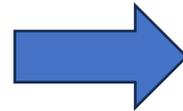
- $U(1)$  model:  $\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} \Rightarrow$  massless gauge boson

- Add complex scalar charged under the  $U(1) \rightarrow$  Abelian Higgs model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (\partial_\mu\phi^* - ieA_\mu\phi^*)(\partial^\mu\phi + ieA^\mu\phi) + \underbrace{m^2|\phi|^2 - \frac{\lambda}{4}|\phi|^4}_{-V(\phi)}$$



$$\phi(x) = \frac{1}{\sqrt{2}}(v + \sigma)e^{\frac{i\pi(x)}{F_\pi}}$$



$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{1}{2}e^2v^2A_\mu A^\mu + \dots$$

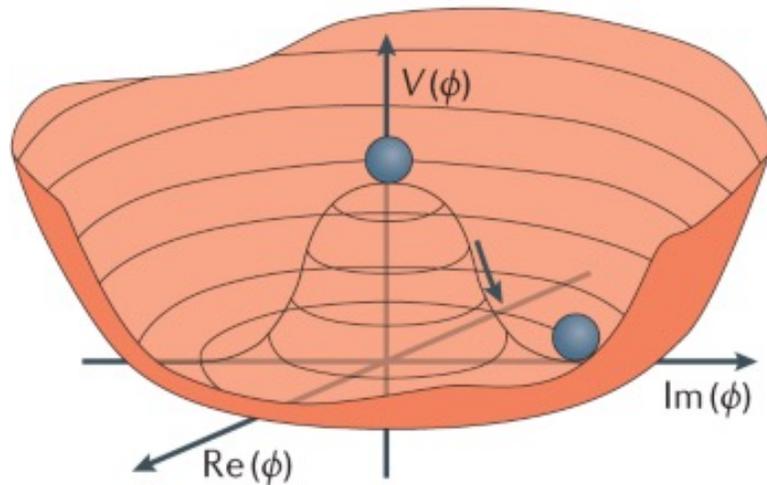
- generated mass term for gauge boson
- Can also generate fermion mass terms via Yukawa interaction:  $y_f\phi\bar{f}f$

# The Higgs mechanism in a nutshell

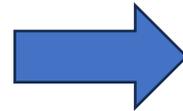
- $U(1)$  model:  $\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} \Rightarrow$  massless gauge boson

- Add complex scalar charged under the  $U(1) \rightarrow$  Abelian Higgs model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (\partial_\mu\phi^* - ieA_\mu\phi^*)(\partial^\mu\phi + ieA^\mu\phi) + \underbrace{m^2|\phi|^2 - \frac{\lambda}{4}|\phi|^4}_{-V(\phi)}$$



$$\phi(x) = \frac{1}{\sqrt{2}}(v + \sigma)e^{\frac{i\pi(x)}{F_\pi}}$$



$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{1}{2}e^2v^2A_\mu A^\mu + \dots$$

- generated mass term for gauge boson
- Can also generate fermion mass terms via Yukawa interaction:  $y_f\phi\bar{f}f$
- Model also used for describing superconductivity ( $\rightarrow$  Ginzburg-Landau theory)

# Constructing the Standard Model (SM)

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi}\not{D}\psi + \text{h.c.} \\ & + \chi_i Y_{ij} \chi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi)\end{aligned}$$

# Constructing the Standard Model (SM)

- Gauge group:  $SU(3)_C \times SU(2)_L \times U(1)_Y$ ,

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi} \not{D} \psi + \text{h.c.} \\ & + \chi_i Y_{ij} \chi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi)\end{aligned}$$

# Constructing the Standard Model (SM)

- Gauge group:  $SU(3)_C \times SU(2)_L \times U(1)_Y$ ,
- 3 quark and lepton generations,

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi} \not{D} \psi + \text{h.c.} \\ & + \sum_i y_{ij} \bar{\psi}_i \psi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi)\end{aligned}$$

# Constructing the Standard Model (SM)

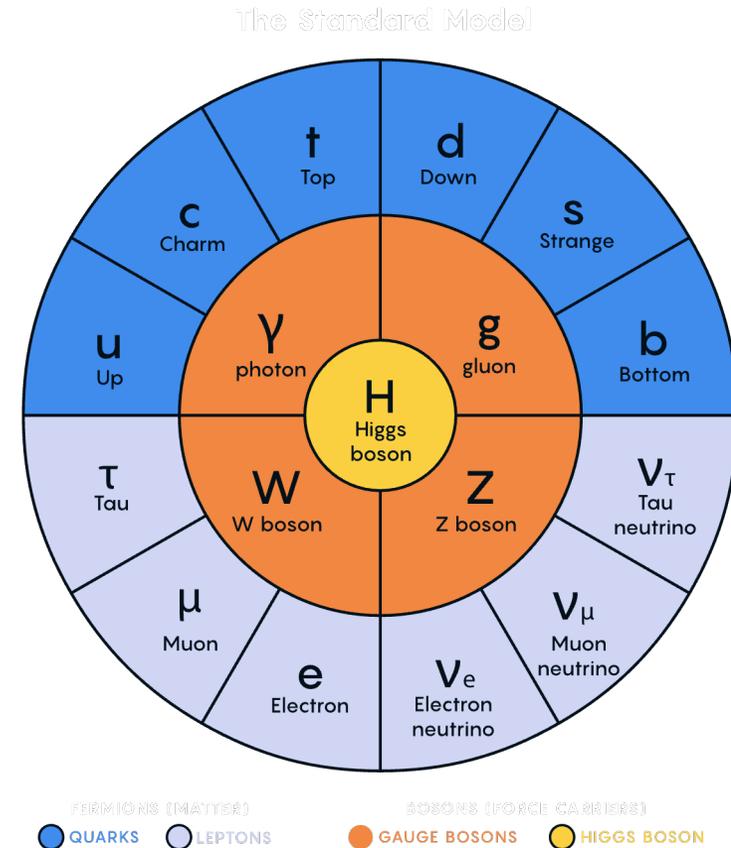
- Gauge group:  $SU(3)_C \times SU(2)_L \times U(1)_Y$ ,
- 3 quark and lepton generations,
- Higgs mechanism to break gauge group to  $SU(3)_C \times U(1)_{em} \rightarrow$  masses for fermions and electroweak gauge bosons.

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi} \not{D} \psi + \text{h.c.} \\ & + \chi_i y_{ij} \chi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi)\end{aligned}$$

# Constructing the Standard Model (SM)

- Gauge group:  $SU(3)_C \times SU(2)_L \times U(1)_Y$ ,
- 3 quark and lepton generations,
- Higgs mechanism to break gauge group to  $SU(3)_C \times U(1)_{em} \rightarrow$  masses for fermions and electroweak gauge bosons.

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi} \not{D} \psi + h.c. \\ & + \sum_i Y_{ij} \psi_i \phi + h.c. \\ & + |D_\mu \phi|^2 - V(\phi) \end{aligned}$$

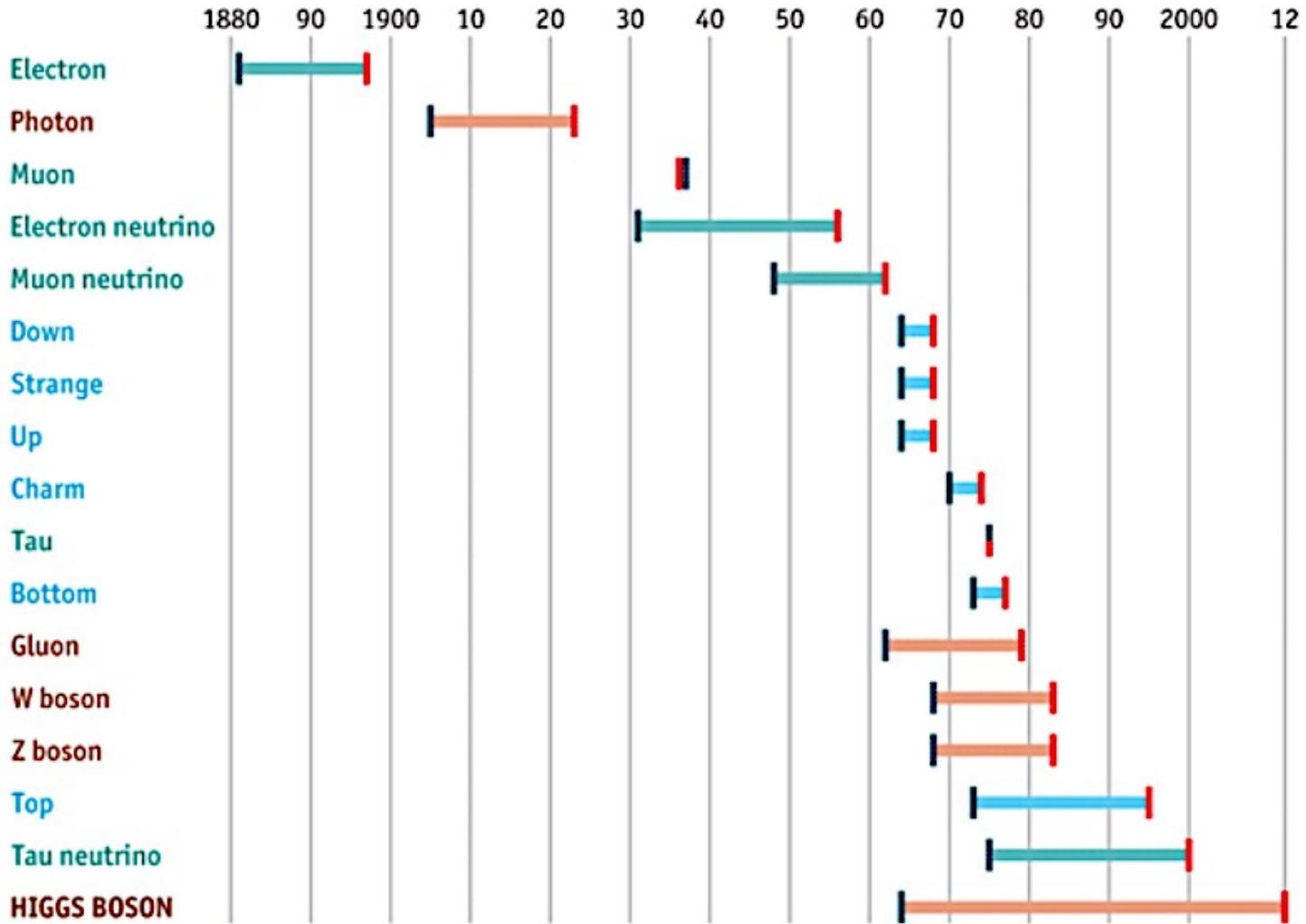


# The Standard Model of particle physics

Years from concept to discovery

Leptons  
Bosons  
Quarks

Theorised/explained  
Discovered



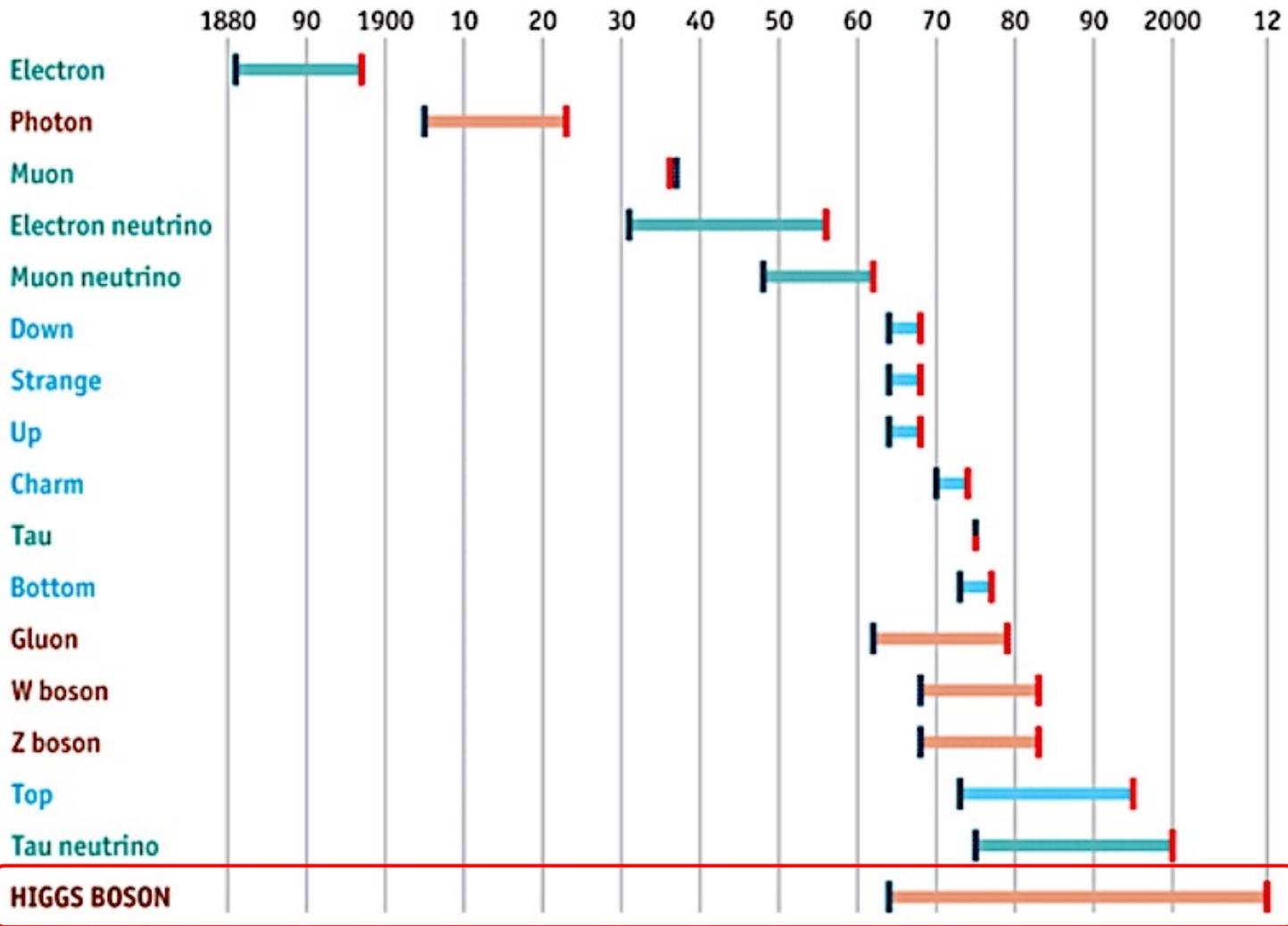
Source: *The Economist*

# The Standard Model of particle physics

Years from concept to discovery

Leptons  
Bosons  
Quarks

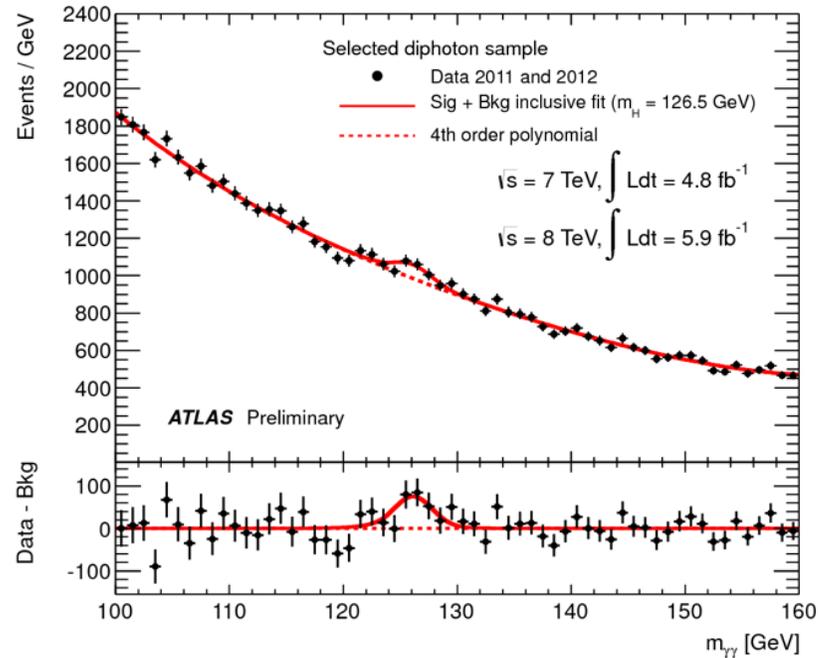
Theorised/explained  
Discovered



Source: *The Economist*

# The Higgs discovery

# 2012

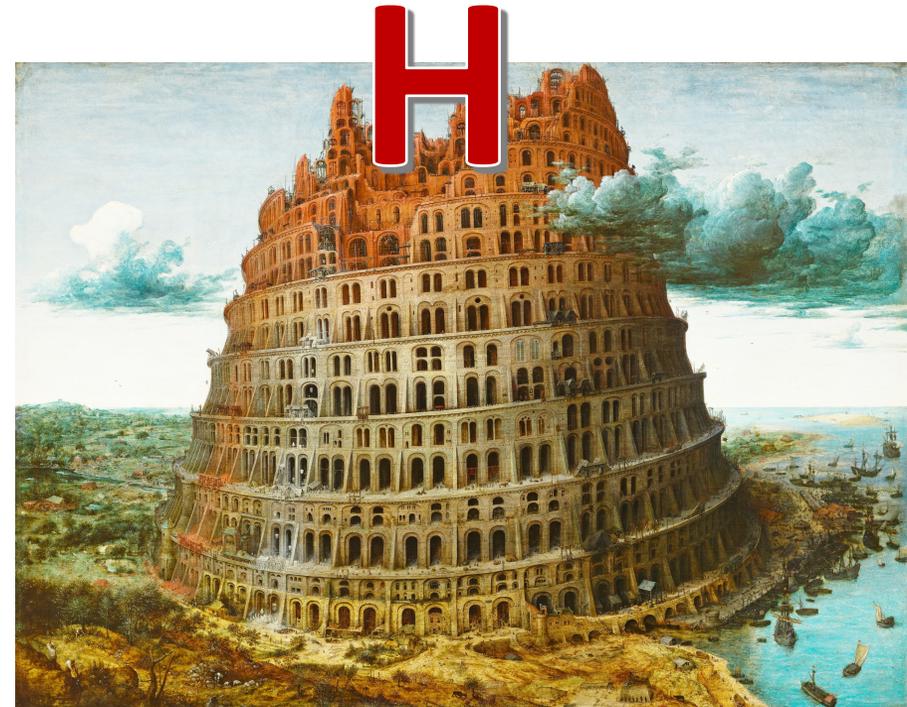
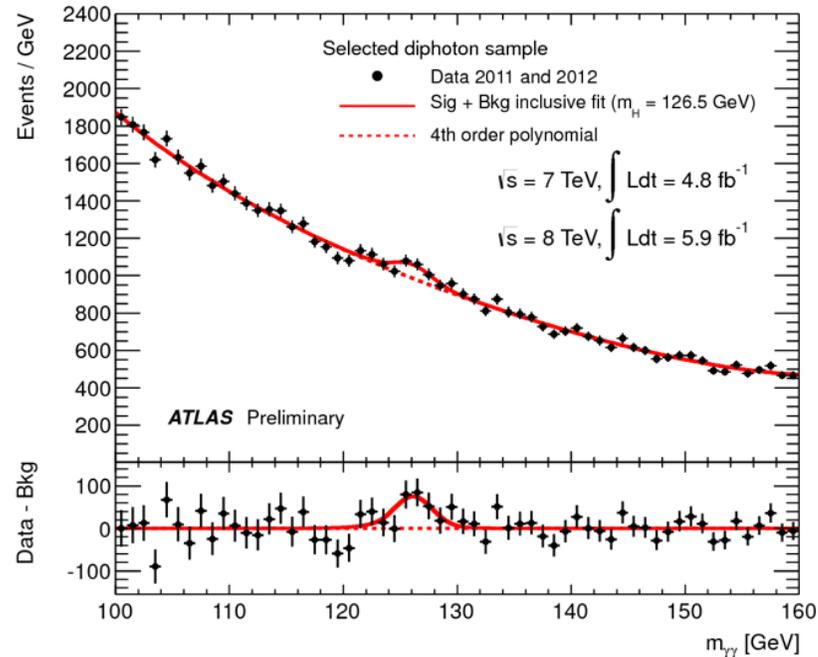


Tower of Babel

After decades of work, the Higgs discovery was a big success for particle physics ( $m_h = 125$  GeV).

# The Higgs discovery

# 2012

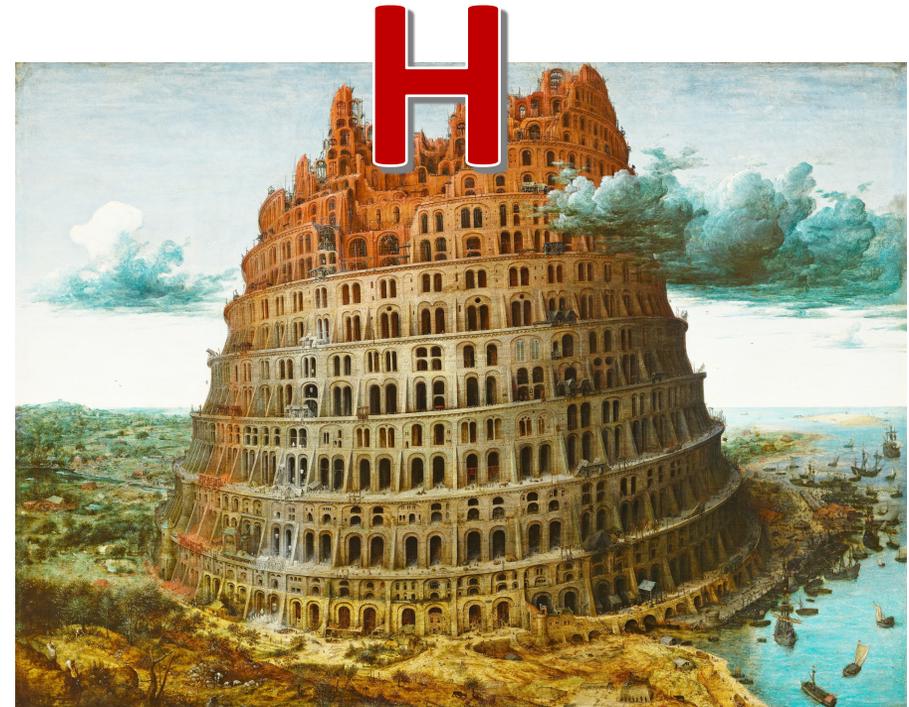
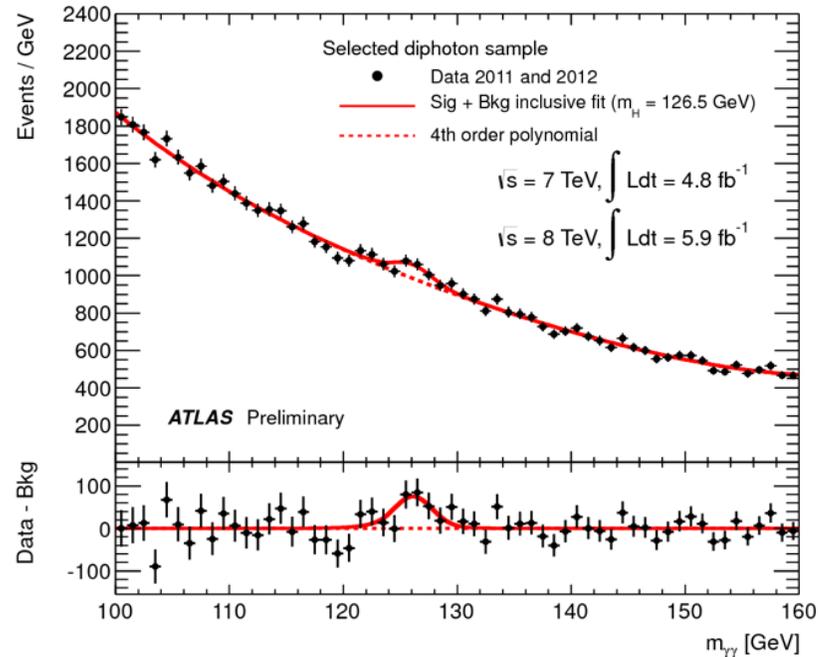


~~Tower of Babel~~ *the SM*

After decades of work, the Higgs discovery was a big success for particle physics ( $m_h = 125$  GeV).

# The Higgs discovery

2012



~~Tower of Babel~~ *the SM*

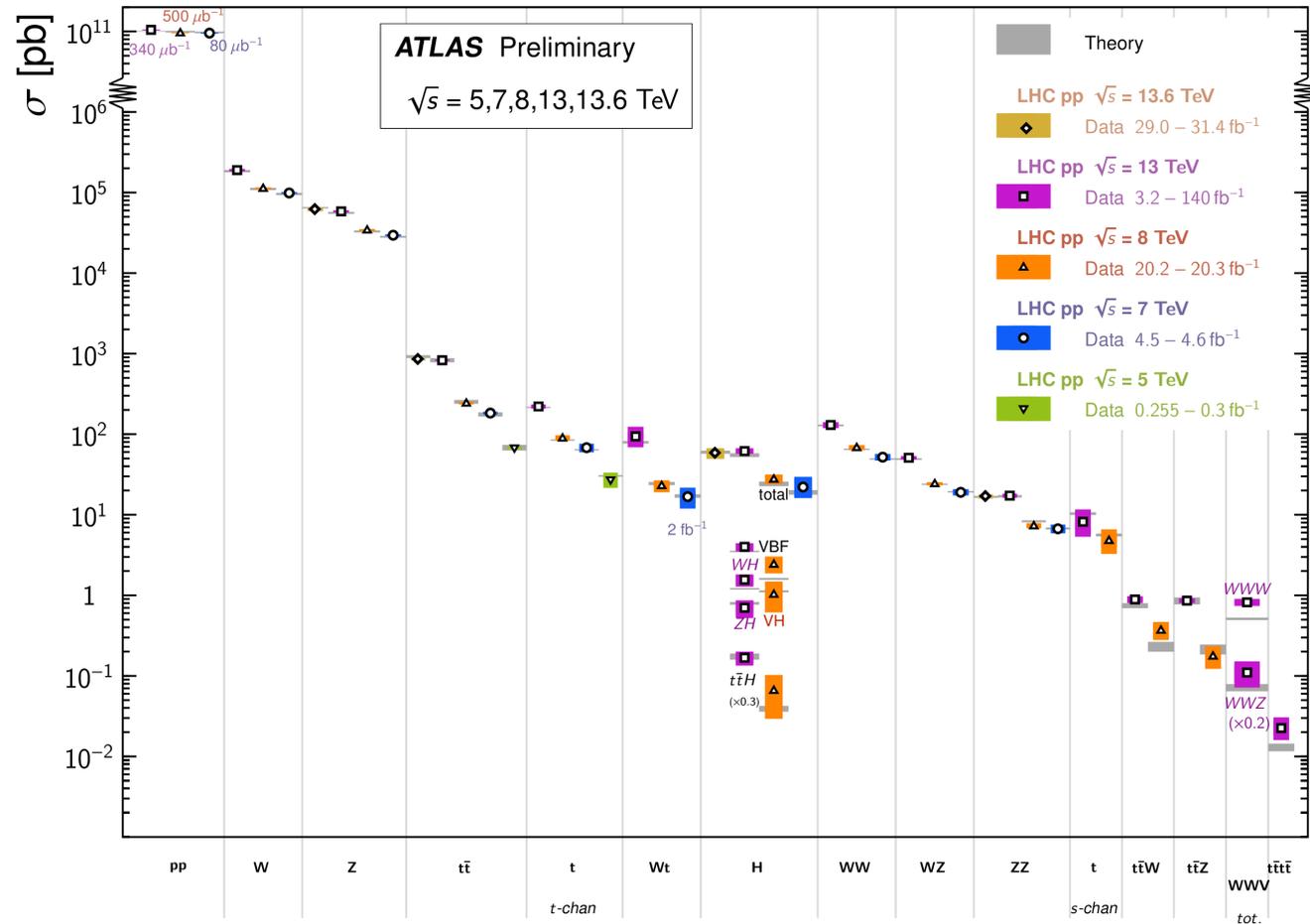
After decades of work, the Higgs discovery was a big success for particle physics ( $m_h = 125$  GeV).

→ So, have we *completed* the SM?

# The SM at the LHC

Standard Model Total Production Cross Section Measurements

Status: October 2023



# Going beyond the SM

So, everything left to do is to confirm the SM with even more precision?

# Going beyond the SM

So, everything left to do is to confirm the SM with even more precision? → **No!**

# Going beyond the SM

So, everything left to do is to confirm the SM with even more precision? → **No!**

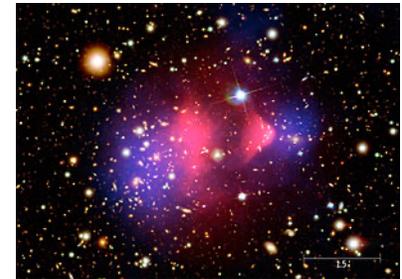
We still want to know about the

# Going beyond the SM

So, everything left to do is to confirm the SM with even more precision? → **No!**

We still want to know about the

- nature of dark matter,

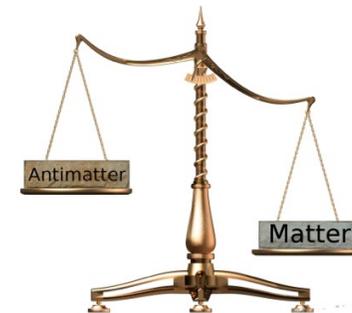
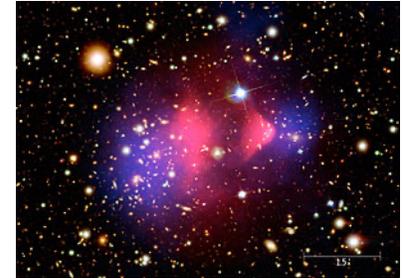


# Going beyond the SM

So, everything left to do is to confirm the SM with even more precision? → **No!**

We still want to know about the

- nature of dark matter,
- baryon asymmetry of the Universe,

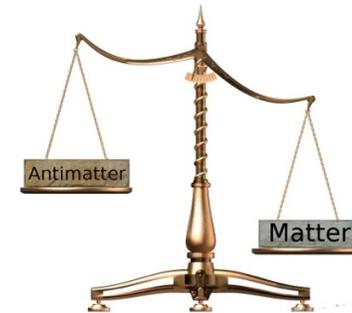
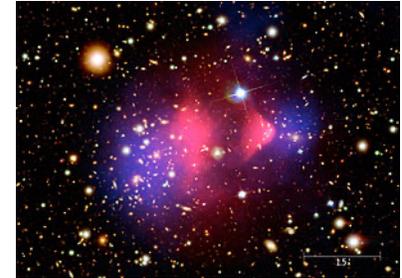


# Going beyond the SM

So, everything left to do is to confirm the SM with even more precision? → **No!**

We still want to know about the

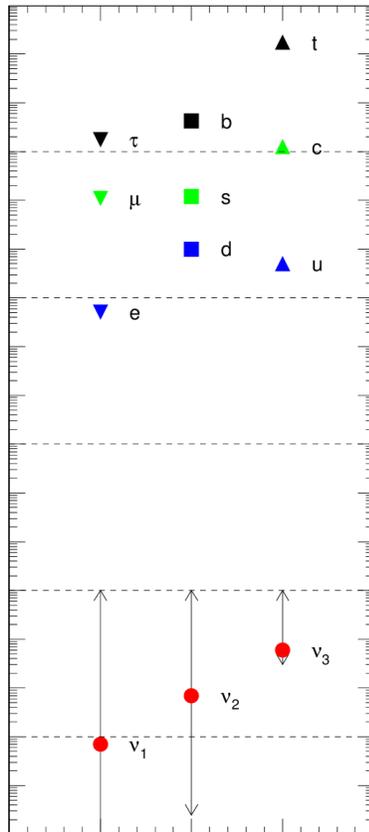
- nature of dark matter,
- baryon asymmetry of the Universe,
- origin of neutrino masses,



$m_\nu$

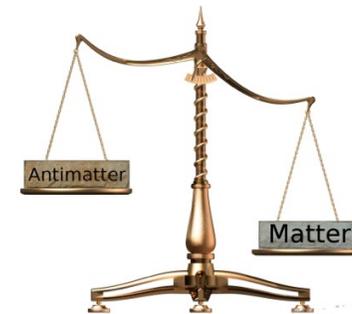
# Going beyond the SM

So, everything left to do is to confirm the SM with even more precision? → **No!**



We still want to know about the

- nature of dark matter,
- baryon asymmetry of the Universe,
- origin of neutrino masses,
- fermion mass hierarchy,

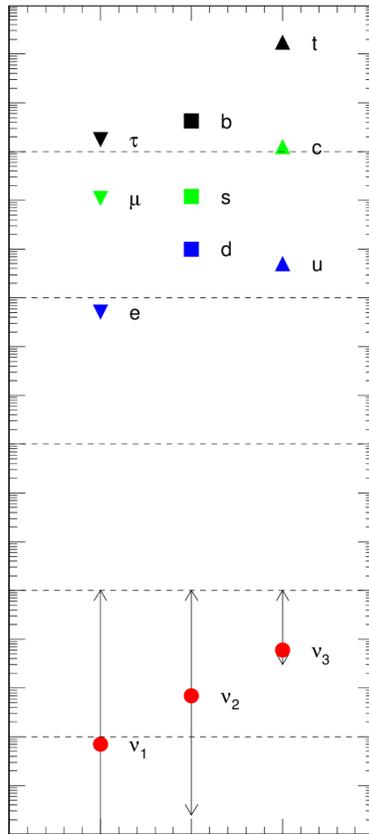


$m_\nu$



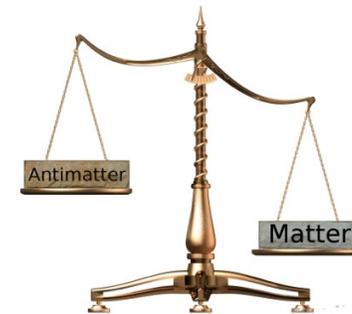
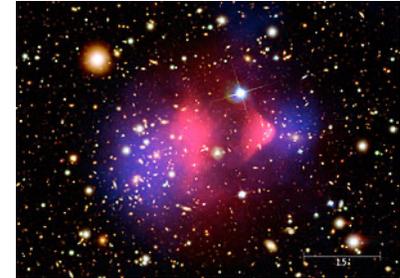
# Going beyond the SM

So, everything left to do is to confirm the SM with even more precision? → **No!**



We still want to know about the

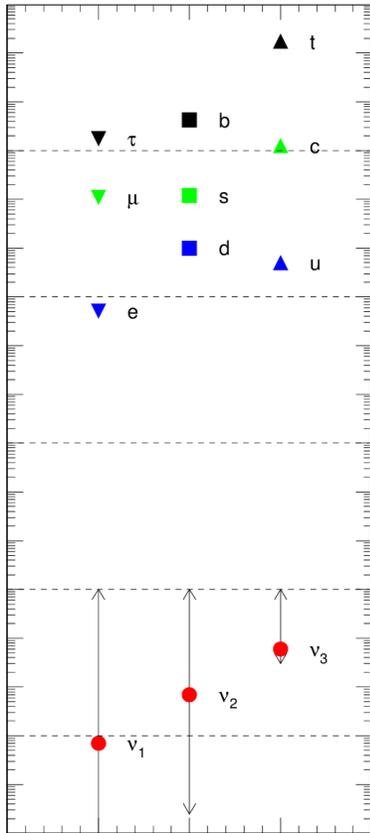
- nature of dark matter,
- baryon asymmetry of the Universe,
- origin of neutrino masses,
- fermion mass hierarchy,
- Higgs hierarchy problem,
- cosmological constant,



$m_\nu$

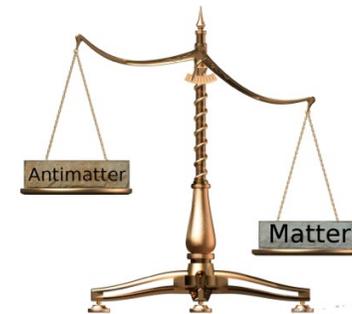
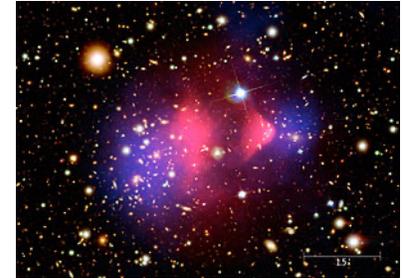
# Going beyond the SM

So, everything left to do is to confirm the SM with even more precision? → **No!**



We still want to know about the

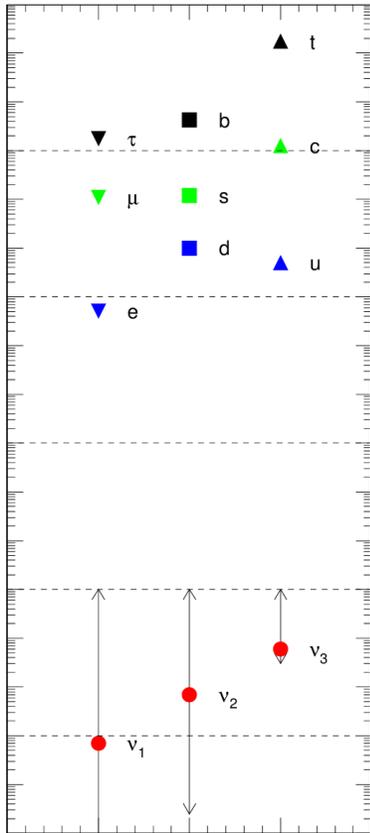
- nature of dark matter,
- baryon asymmetry of the Universe,
- origin of neutrino masses,
- fermion mass hierarchy,
- Higgs hierarchy problem,
- cosmological constant,
- thermal history of the Universe,



$m_\nu$

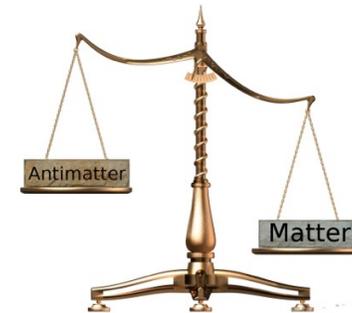
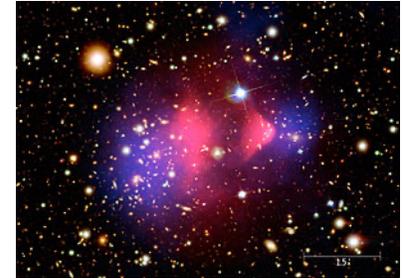
# Going beyond the SM

So, everything left to do is to confirm the SM with even more precision? → **No!**



We still want to know about the

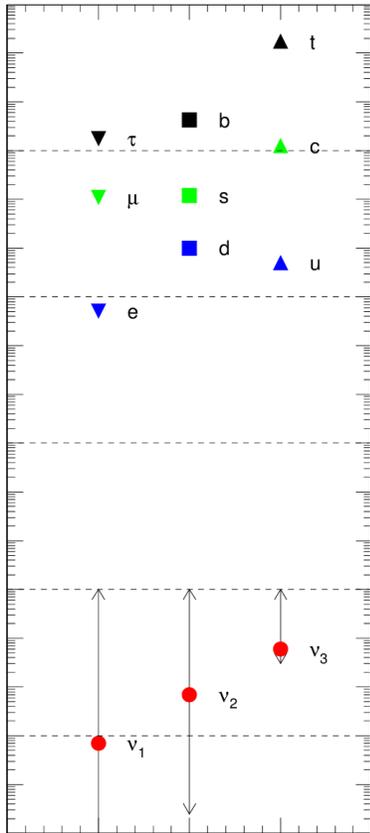
- nature of dark matter,
- baryon asymmetry of the Universe,
- origin of neutrino masses,
- fermion mass hierarchy,
- Higgs hierarchy problem,
- cosmological constant,
- thermal history of the Universe,
- stability of the Universe,



$m_\nu$

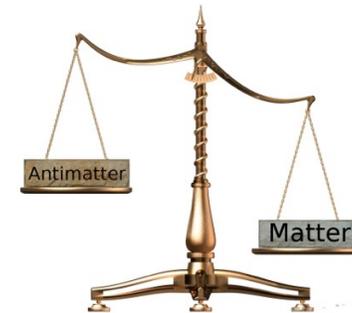
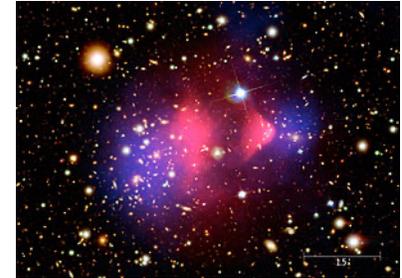
# Going beyond the SM

So, everything left to do is to confirm the SM with even more precision? → **No!**



We still want to know about the

- nature of dark matter,
- baryon asymmetry of the Universe,
- origin of neutrino masses,
- fermion mass hierarchy,
- Higgs hierarchy problem,
- cosmological constant,
- thermal history of the Universe,
- stability of the Universe,
- ...



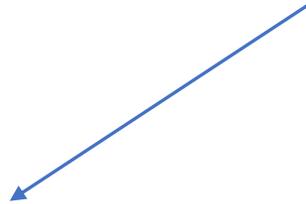
$m_\nu$

# The Higgs as a window to new physics

$$\mathcal{L}_{\text{Higgs}} = (D^\mu H)^\dagger (D_\mu H) + y_{ij} H f_i \bar{f}_j + \mu^2 |H|^2 - \lambda |H|^4 - V_0$$

# The Higgs as a window to new physics

$$\mathcal{L}_{\text{Higgs}} = (D^\mu H)^\dagger (D_\mu H) + y_{ij} H f_i \bar{f}_j + \mu^2 |H|^2 - \lambda |H|^4 - V_0$$



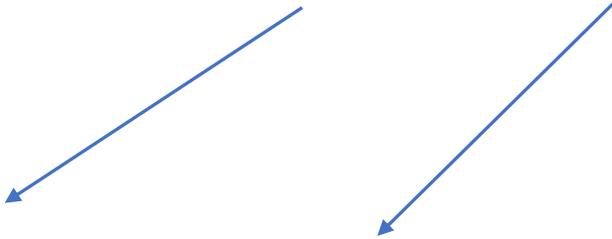
- fermion masses,
- neutrino masses,
- flavour structure.

# The Higgs as a window to new physics

$$\mathcal{L}_{\text{Higgs}} = (D^\mu H)^\dagger (D_\mu H) + y_{ij} H f_i \bar{f}_j + \mu^2 |H|^2 - \lambda |H|^4 - V_0$$

- fermion masses,
- neutrino masses,
- flavour structure.

hierarchy problem



# The Higgs as a window to new physics

$$\mathcal{L}_{\text{Higgs}} = (D^\mu H)^\dagger (D_\mu H) + y_{ij} H f_i \bar{f}_j + \mu^2 |H|^2 - \lambda |H|^4 - V_0$$

- fermion masses,
- neutrino masses,
- flavour structure.

hierarchy problem

- stability,
- thermal history.

# The Higgs as a window to new physics

$$\mathcal{L}_{\text{Higgs}} = (D^\mu H)^\dagger (D_\mu H) + y_{ij} H f_i \bar{f}_j + \mu^2 |H|^2 - \lambda |H|^4 - V_0$$

- fermion masses,
- neutrino masses,
- flavour structure.

hierarchy problem

- stability,
- thermal history.

cosmological  
constant

# The Higgs as a window to new physics

$$\mathcal{L}_{\text{Higgs}} = (D^\mu H)^\dagger (D_\mu H) + y_{ij} H f_i \bar{f}_j + \mu^2 |H|^2 - \lambda |H|^4 - V_0 - \lambda_{SH} |H|^2 S^2 + \dots$$

- fermion masses,
- neutrino masses,
- flavour structure.

hierarchy problem

- stability,
- thermal history.

cosmological constant

- portal to hidden sector,
- dark matter.

# The Higgs as a window to new physics

$$\mathcal{L}_{\text{Higgs}} = (D^\mu H)^\dagger (D_\mu H) + y_{ij} H f_i \bar{f}_j + \mu^2 |H|^2 - \lambda |H|^4 - V_0 - \lambda_{SH} |H|^2 S^2 + \dots$$

- fermion masses,
- neutrino masses,
- flavour structure.

hierarchy problem

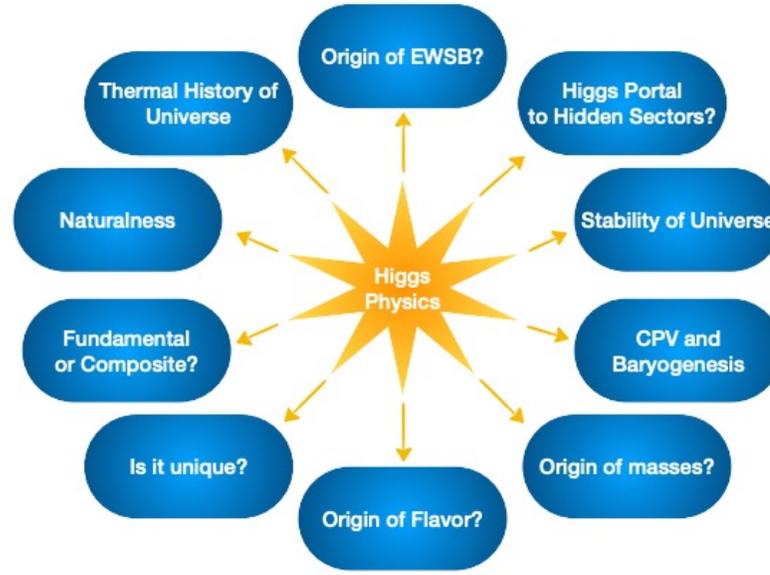
- stability,
- thermal history.

cosmological constant

- portal to hidden sector,
- dark matter.

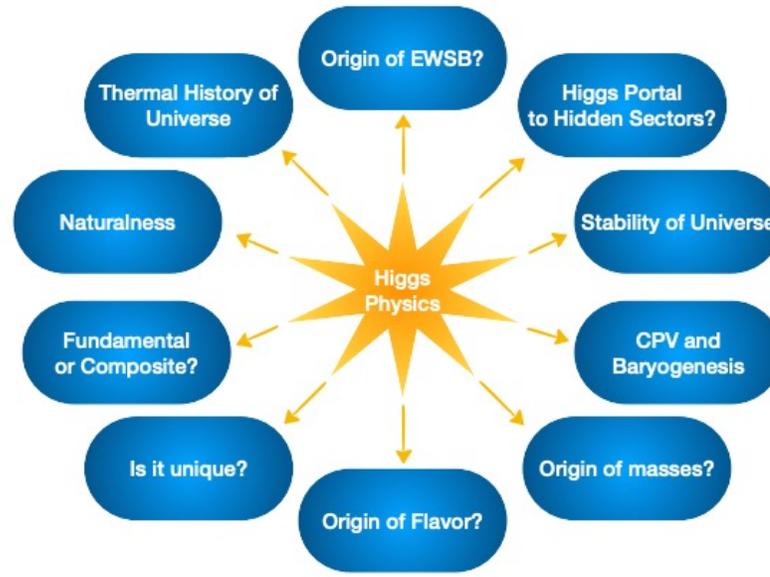
⇒ Strong motivation for on-going and future Higgs precision programs.

# The Higgs as a window to new physics



[Snowmass 2209.07510]

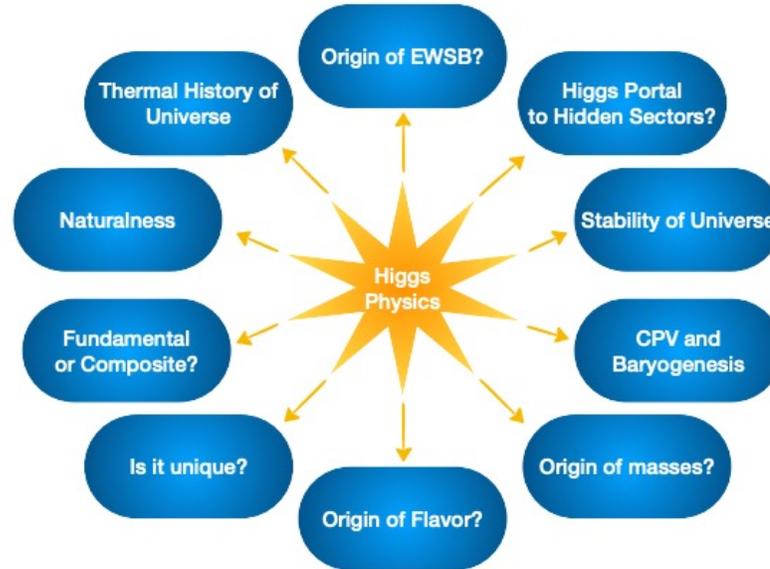
# The Higgs as a window to new physics



[Snowmass 2209.07510]

→ Where are we more than a decade after the Higgs discovery?

# The Higgs as a window to new physics

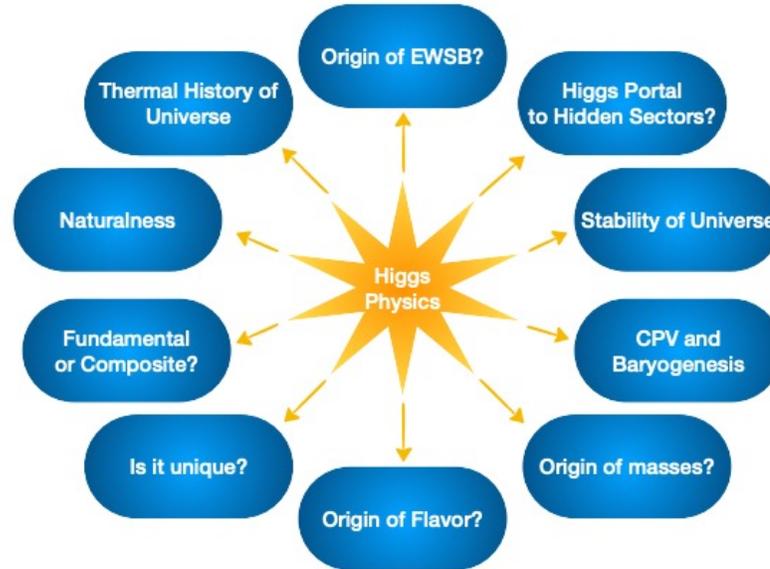


[Snowmass 2209.07510]

→ Where are we more than a decade after the Higgs discovery?

→ What have we learned about the Higgs in the mean time?

# The Higgs as a window to new physics

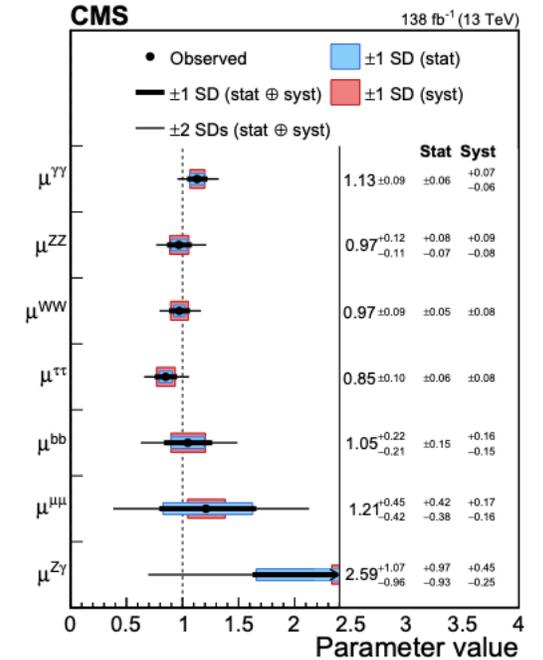
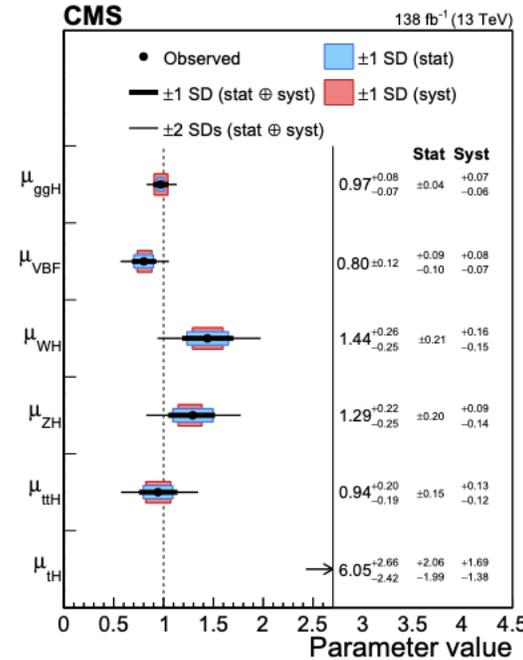
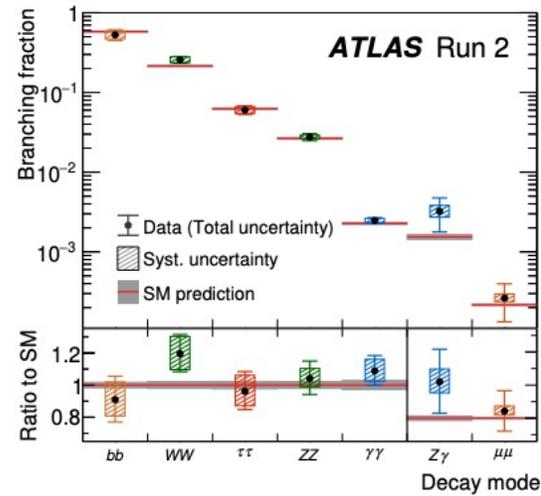
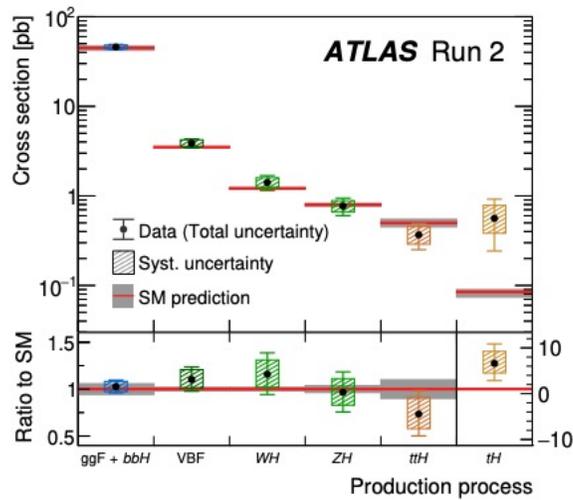


[Snowmass 2209.07510]

- Where are we more than a decade after the Higgs discovery?
- What have we learned about the Higgs in the mean time?
- What is still left to explore?

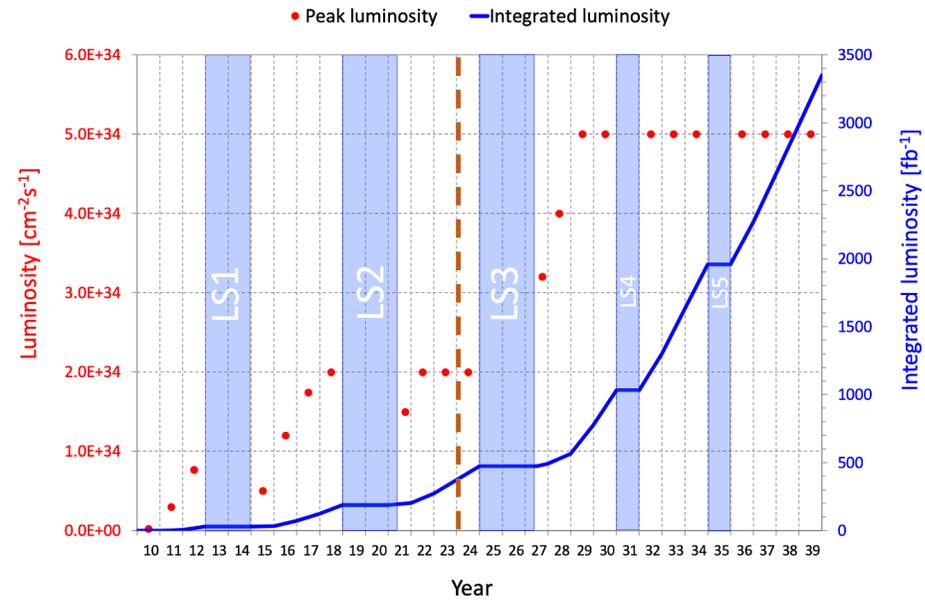
# The Higgs 10 years later

[ATLAS 2207.00092, CMS 2207.00043]



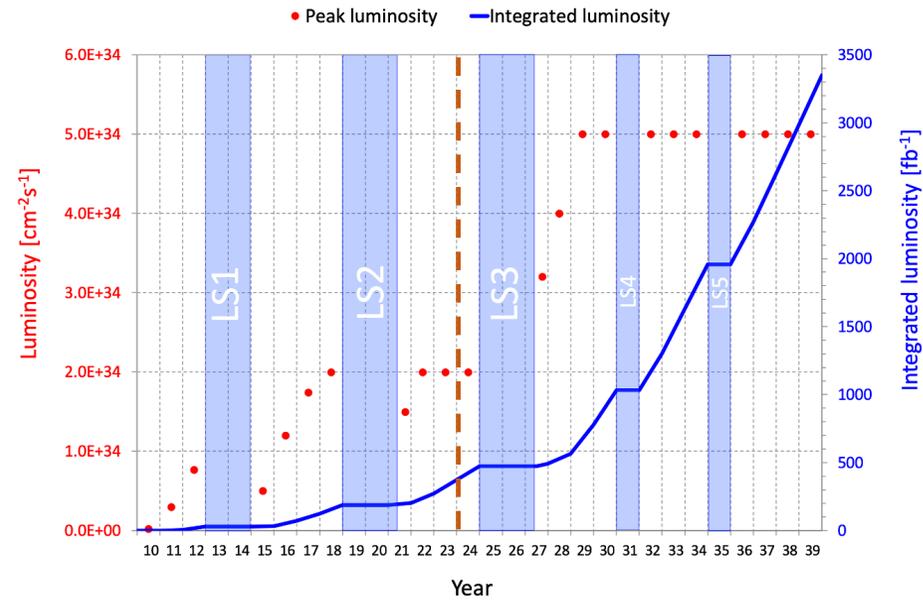
- Ten years later, we have entered the Higgs precision era.
- So far, all Higgs measurements agree with the SM predictions within the experimental and theoretical uncertainties.

# LHC Run-3 and beyond



Much more data will be collected in the next years.

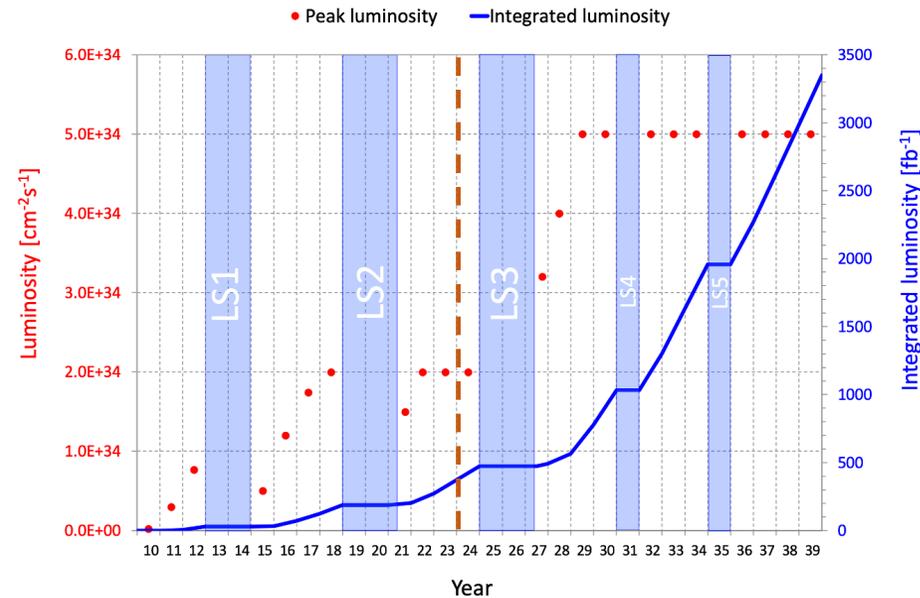
# LHC Run-3 and beyond



Much more data will be collected in the next years.

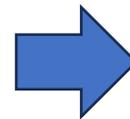
→ The LHC program has just started.

# LHC Run-3 and beyond



Much more data will be collected in the next years.

→ The LHC program has just started.



- What can we learn from existing measurements?
- What is still left to explore?
- Have we found **the** SM Higgs?

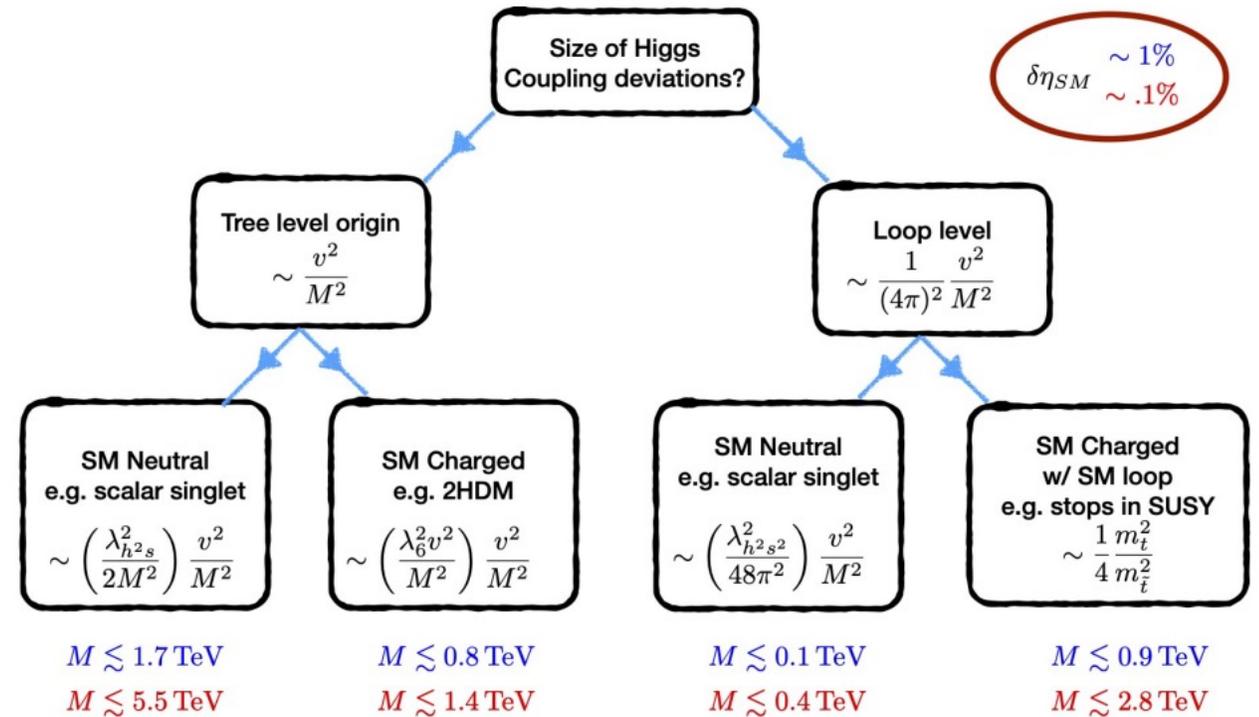
What can we learn from existing measurements?

# What can we learn from Higgs precision measurements?

- Higgs precision measurements put stringent constraints on many BSM scenarios.

Simplified scaling analysis:

- 1% precision level can constrain BSM particles with mass from 100 GeV to several TeV (within reach of the LHC or future colliders).



$\delta\eta_{SM} \sim 1\%$   
 $\sim .1\%$

Conservative Scaling for Upper Limit on Mass Scale Probed by Higgs Precision

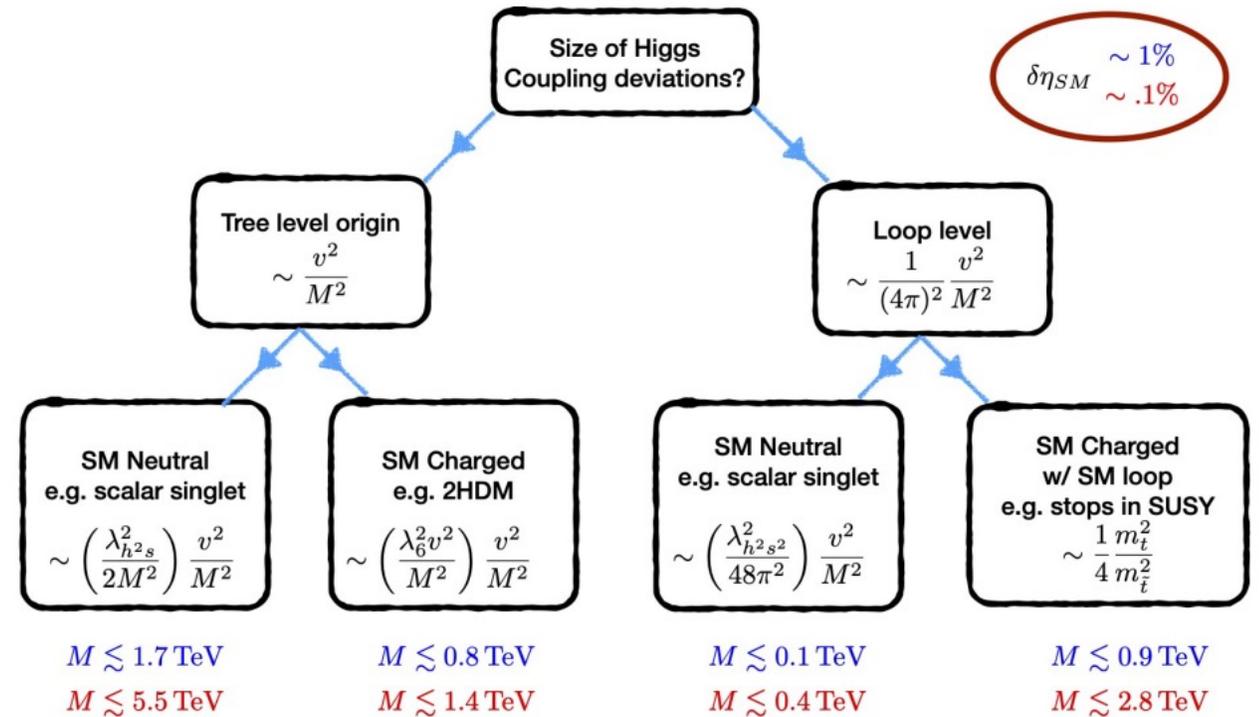
[Snowmass 2209.07510]

# What can we learn from Higgs precision measurements?

- Higgs precision measurements put stringent constraints on many BSM scenarios.

Simplified scaling analysis:

- 1% precision level can constrain BSM particles with mass from 100 GeV to several TeV (within reach of the LHC or future colliders).



Conservative Scaling for Upper Limit on Mass Scale Probed by Higgs Precision

[Snowmass 2209.07510]

We also shouldn't forget about the interplay with direct searches!

# Interplay with direct searches

- Important interplay between Higgs precision measurements and direct searches for BSM particles.

# Interplay with direct searches

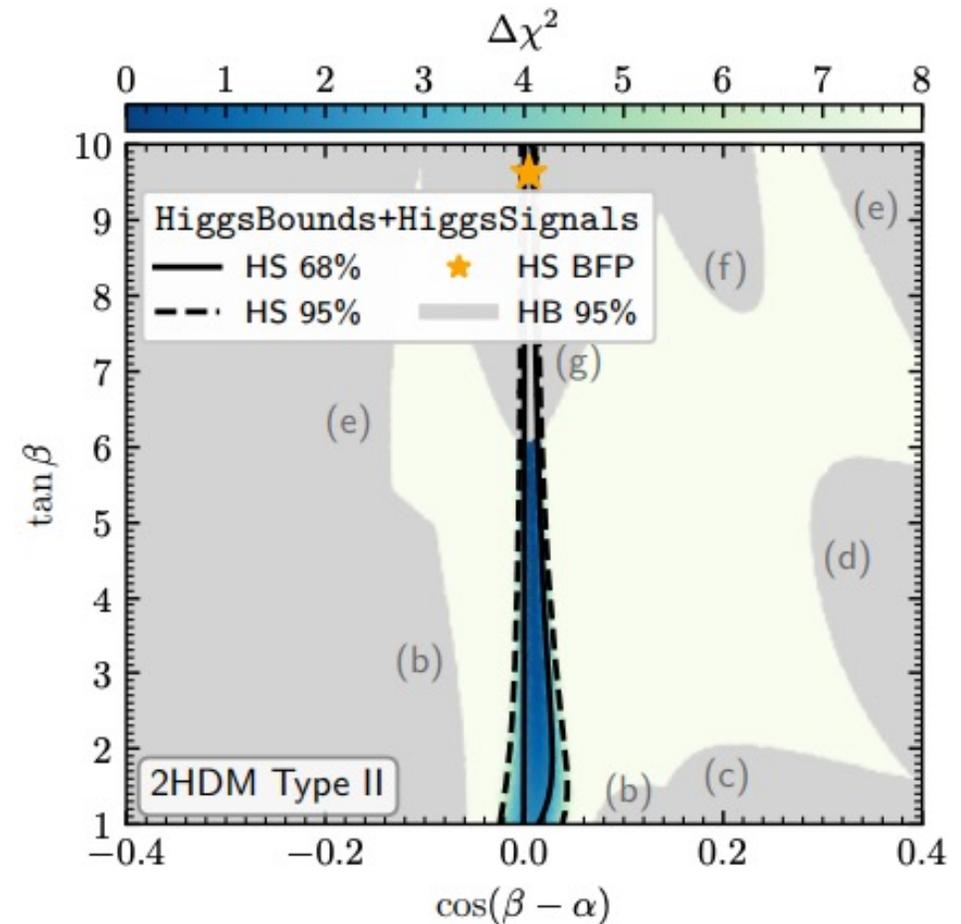
- Important interplay between Higgs precision measurements and direct searches for BSM particles.
- Cannot be captured in EFT framework  
→ use 2HDM here as a benchmark model.

# Interplay with direct searches

- Important interplay between Higgs precision measurements and direct searches for BSM particles.
- Cannot be captured in EFT framework  
→ use 2HDM here as a benchmark model.
- 2HDM: extend SM by 2<sup>nd</sup> Higgs doublet  
→ additional  $H, A, H^\pm$  BSM Higgs bosons

# Interplay with direct searches

- Important interplay between Higgs precision measurements and direct searches for BSM particles.
- Cannot be captured in EFT framework  
→ use 2HDM here as a benchmark model.
- 2HDM: extend SM by 2<sup>nd</sup> Higgs doublet  
→ additional  $H, A, H^\pm$  BSM Higgs bosons
- BSM searches:
  - b) CMS:  $pp \rightarrow \phi_1 \rightarrow h_{125} \phi_2 \rightarrow bb\tau\tau$
  - c) CMS:  $pp \rightarrow \phi \rightarrow Zh_{125}$
  - d) ATLAS:  $pp \rightarrow \phi \rightarrow WW, ZZ, WZ$
  - e) ATLAS:  $pp \rightarrow \phi \rightarrow h_{125} h_{125}$
  - f) ATLAS:  $pp \rightarrow \phi \rightarrow VV, Vh_{125}$
  - g) ATLAS:  $pp \rightarrow \phi \rightarrow \tau\tau$



$$m_H = m_A = m_{H^\pm} = \sqrt{m_{12}^2 / (\sin \beta \cos \beta)} = 800 \text{ GeV}$$

[HB et al. 2210.09332]

What is still left to explore?



# What we don't know about the Higgs (yet)

# What we don't know about the Higgs (yet)

Many Higgs properties only weakly constrained, e.g.:

- Higgs width/BSM decay channels,  
SM:  $\Gamma_h \simeq 4.1 \text{ MeV}$ ,
- light Yukawas,  
SM:  $y_f \propto m_f/v$ ,
- Higgs potential,  
SM:  $V(\Phi) = -\frac{1}{2}m_h^2\Phi^\dagger\Phi + \frac{m_h^2}{2v^2}(\Phi^\dagger\Phi)^2$ ,
- Higgs CP properties,  
SM: Higgs is  $\sim$  CP-even.

# What we don't know about the Higgs (yet)

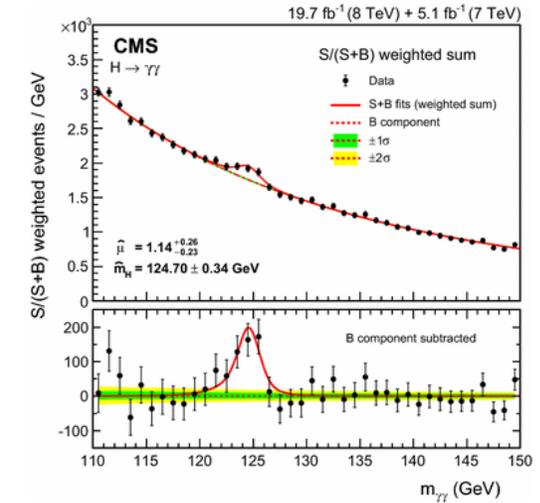
Many Higgs properties only weakly constrained, e.g.:

- Higgs width/BSM decay channels,  
SM:  $\Gamma_h \simeq 4.1 \text{ MeV}$ ,

- light Yukawas,  
SM:  $y_f \propto m_f/v$ ,

- Higgs potential,  
SM:  $V(\Phi) = -\frac{1}{2}m_h^2\Phi^\dagger\Phi + \frac{m_h^2}{2v^2}(\Phi^\dagger\Phi)^2$ ,

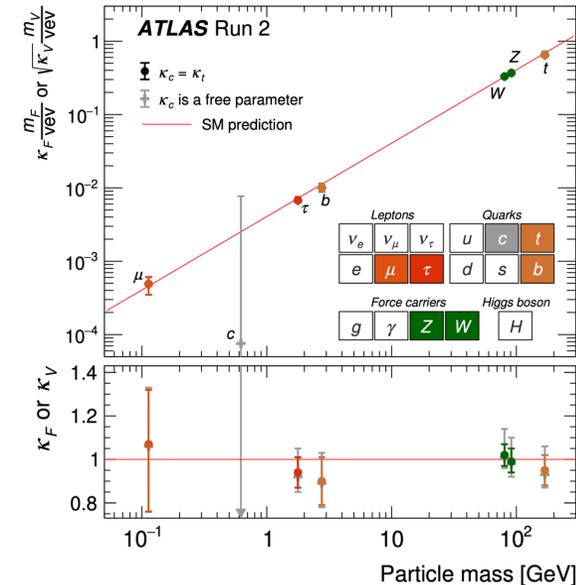
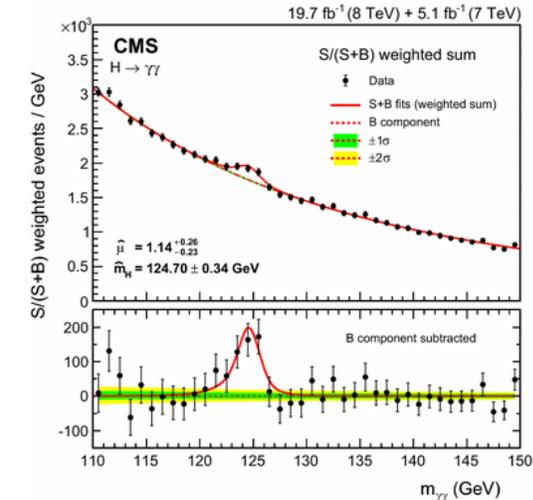
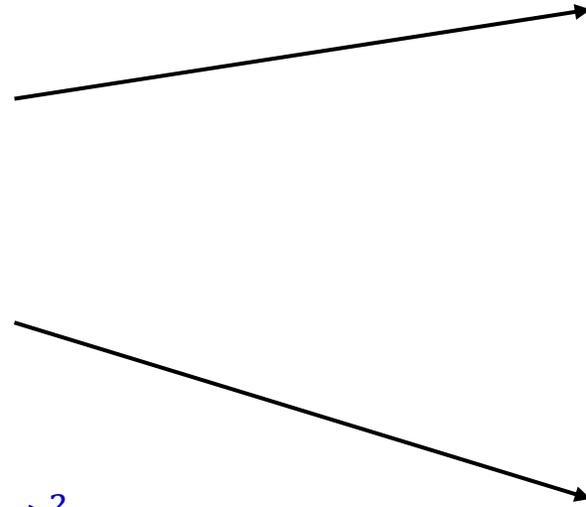
- Higgs CP properties,  
SM: Higgs is  $\sim$  CP-even.



# What we don't know about the Higgs (yet)

Many Higgs properties only weakly constrained, e.g.:

- Higgs width/BSM decay channels,  
SM:  $\Gamma_h \simeq 4.1 \text{ MeV}$ ,
- light Yukawas,  
SM:  $y_f \propto m_f/v$ ,
- Higgs potential,  
SM:  $V(\Phi) = -\frac{1}{2}m_h^2\Phi^\dagger\Phi + \frac{m_h^2}{2v^2}(\Phi^\dagger\Phi)^2$ ,
- Higgs CP properties,  
SM: Higgs is  $\sim$  CP-even.

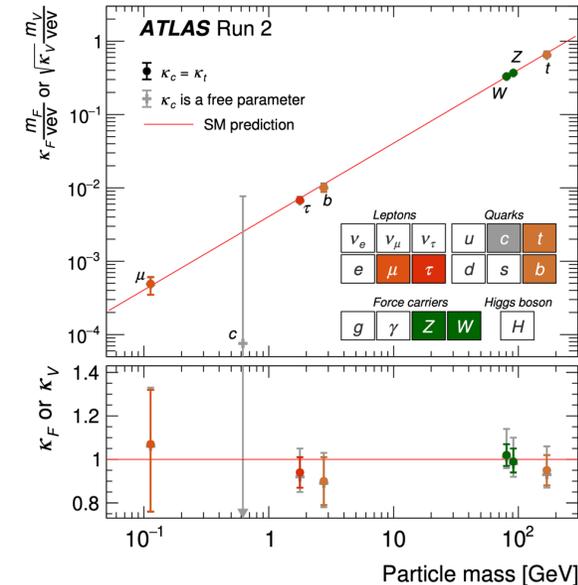
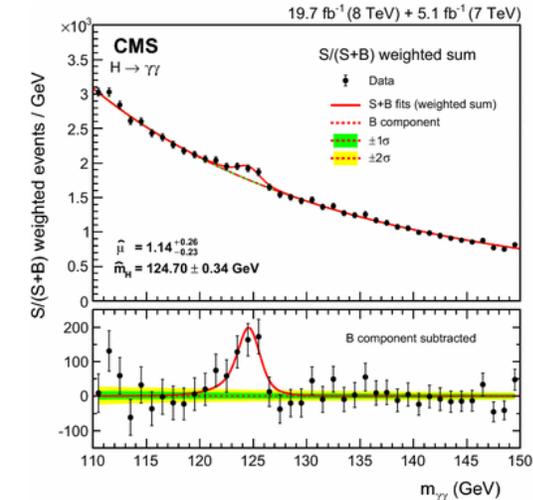


# What we don't know about the Higgs (yet)

Many Higgs properties only weakly constrained, e.g.:

- Higgs width/BSM decay channels,  
SM:  $\Gamma_h \simeq 4.1 \text{ MeV}$ ,
- light Yukawas,  
SM:  $y_f \propto m_f/v$ ,

- Higgs potential,  
SM:  $V(\Phi) = -\frac{1}{2}m_h^2\Phi^\dagger\Phi + \frac{m_h^2}{2v^2}(\Phi^\dagger\Phi)^2$ ,
- Higgs CP properties,  
SM: Higgs is  $\sim$  CP-even.



# The Higgs potential

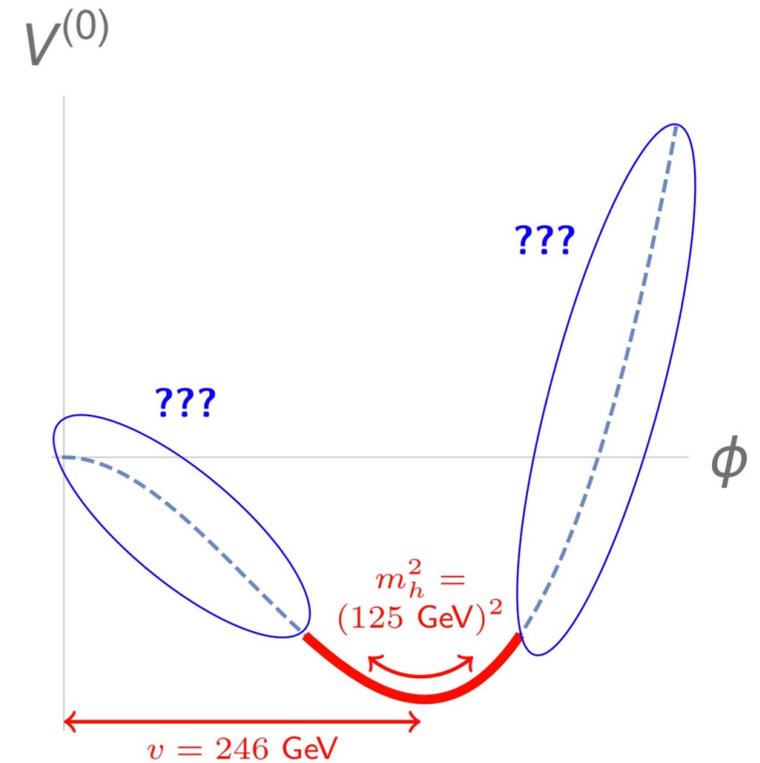
# The Higgs potential

*Is it really a Mexican hat?*



# What do we know about the Higgs potential?

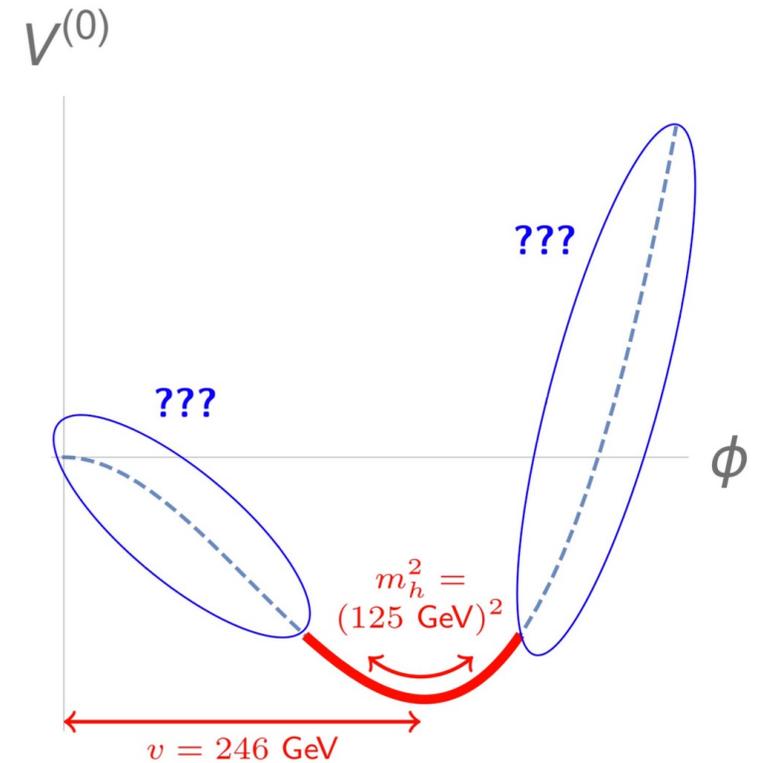
- After the Higgs discovery, we know
  - the location of the EW minimum:  $v = 246$  GeV,
  - the curvature of the potential close to the minimum:  $m_h = 125$  GeV.



[figure by J. Braathen]

# What do we know about the Higgs potential?

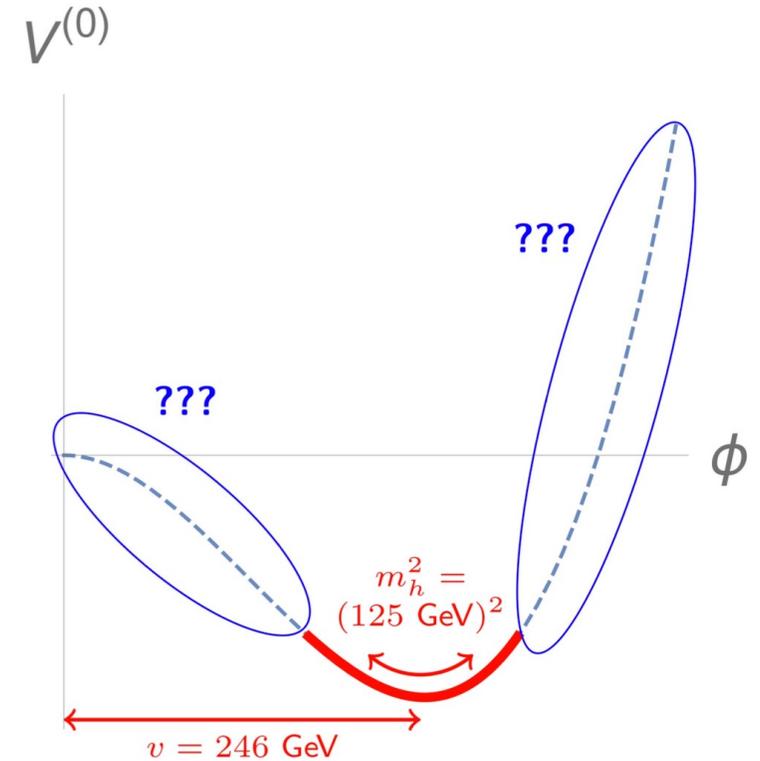
- After the Higgs discovery, we know
  - the location of the EW minimum:  $v = 246$  GeV,
  - the curvature of the potential close to the minimum:  $m_h = 125$  GeV.
- Away from the minimum, the shape of the potential is, however, unknown so far.
  - Determination of trilinear Higgs coupling  $\lambda_{hhh}$  crucial.



[figure by J. Braathen]

# What do we know about the Higgs potential?

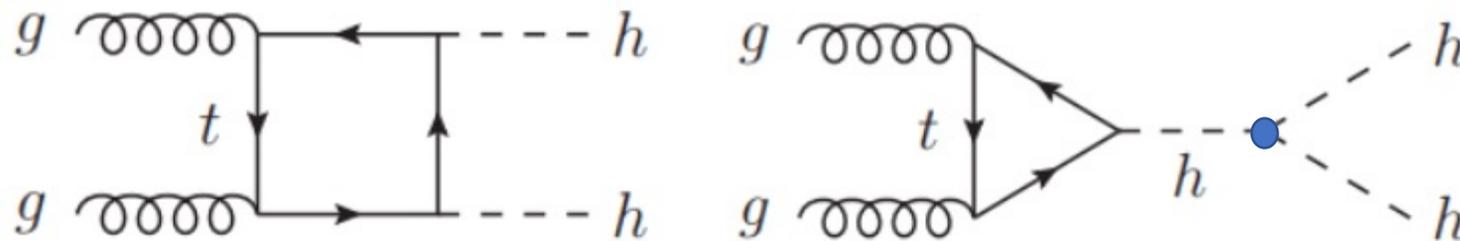
- After the Higgs discovery, we know
  - the location of the EW minimum:  $v = 246$  GeV,
  - the curvature of the potential close to the minimum:  $m_h = 125$  GeV.
- Away from the minimum, the shape of the potential is, however, unknown so far.
  - Determination of trilinear Higgs coupling  $\lambda_{hhh}$  crucial.
- $\lambda_{hhh}$  closely linked to
  - stability of EW vacuum,
  - nature of EW phase transition (→ EW baryogenesis?).



[figure by J. Braathen]

# Probing $\lambda_{hhh}$ via double-Higgs production

Most direct probe of trilinear Higgs coupling: double-Higgs production via gluon fusion.

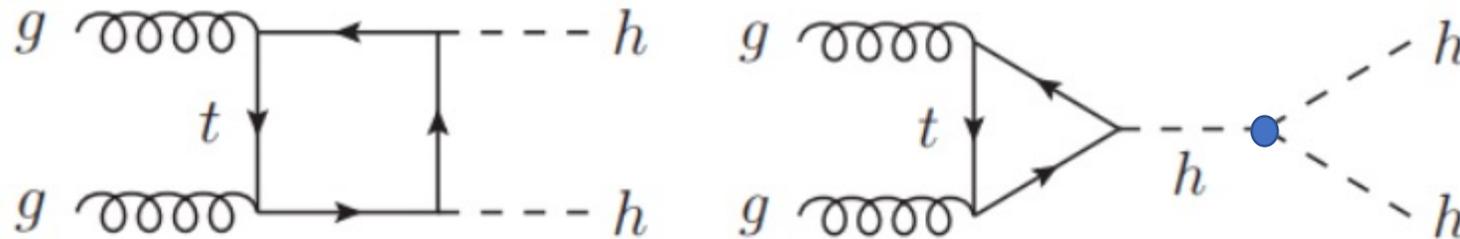


In the SM: large destructive interference between box and triangle contribution.

⇒ Deviations from SM trilinear Higgs coupling can significantly enhance the  $hh$  cross section.

# Probing $\lambda_{hhh}$ via double-Higgs production

Most direct probe of trilinear Higgs coupling: double-Higgs production via gluon fusion.



In the SM: large destructive interference between box and triangle contribution.

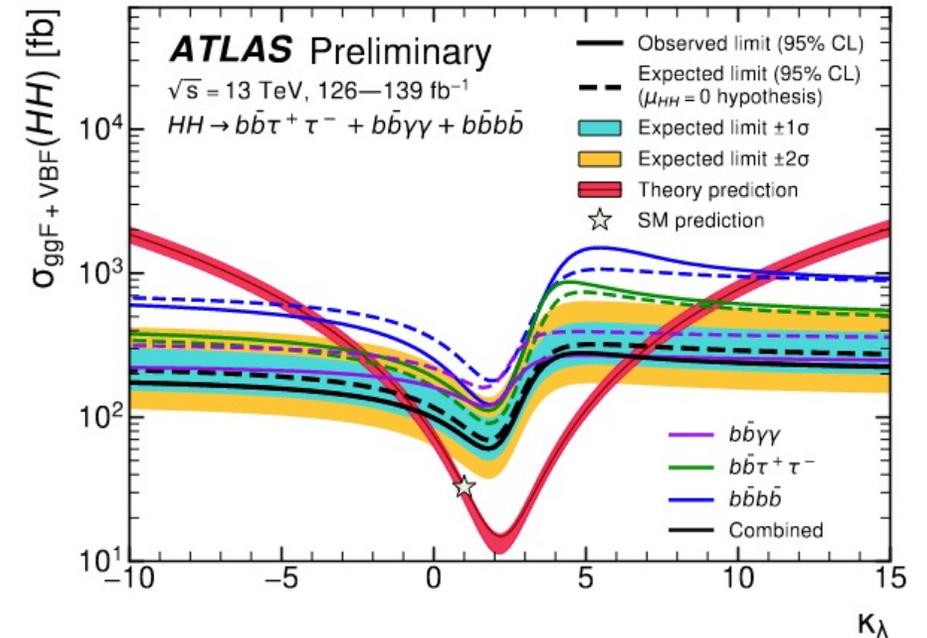
⇒ Deviations from SM trilinear Higgs coupling can significantly enhance the  $hh$  cross section.



Interpret experimental upper limits on  $hh$  cross section as limits on  $\kappa_\lambda$ .

# Experimental bound on $\kappa_\lambda \equiv \lambda_{hhh}/\lambda_{hhh}^{\text{SM}}$

Current strongest limit:  $-0.4 < \kappa_\lambda < 6.3$  at 95% CL [ATLAS-CONF-2022-050].

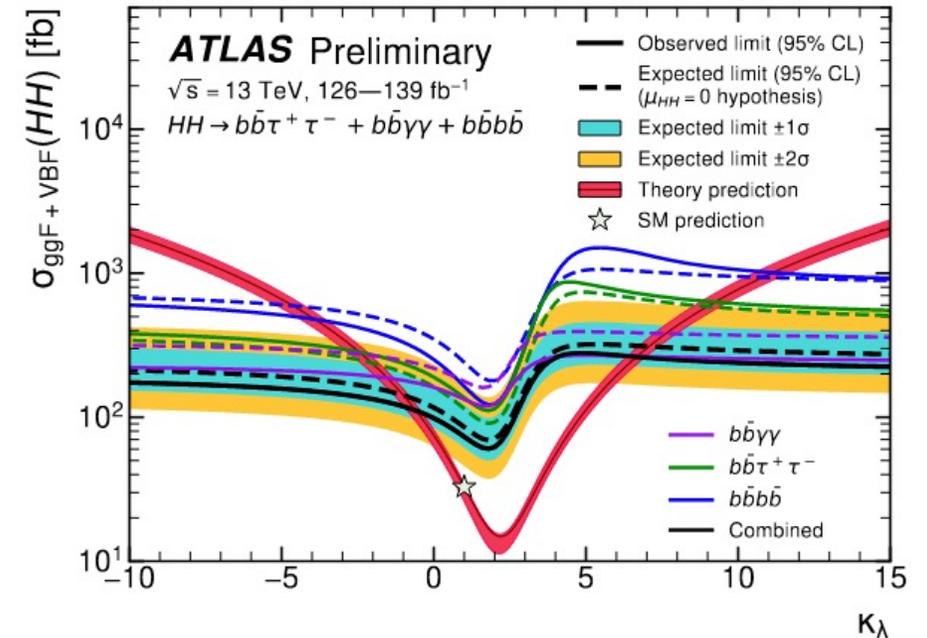


# Experimental bound on $\kappa_\lambda \equiv \lambda_{hhh}/\lambda_{hhh}^{\text{SM}}$

Current strongest limit:  $-0.4 < \kappa_\lambda < 6.3$  at 95% CL [ATLAS-CONF-2022-050].

## Assumptions:

- Simplest analysis assumes that all other Higgs couplings are SM-like.
- Non-resonant Higgs-boson pair production only deviates from the SM via a modified trilinear Higgs coupling (i.e., no heavy resonances).

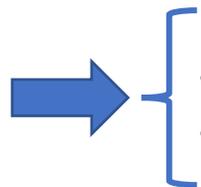
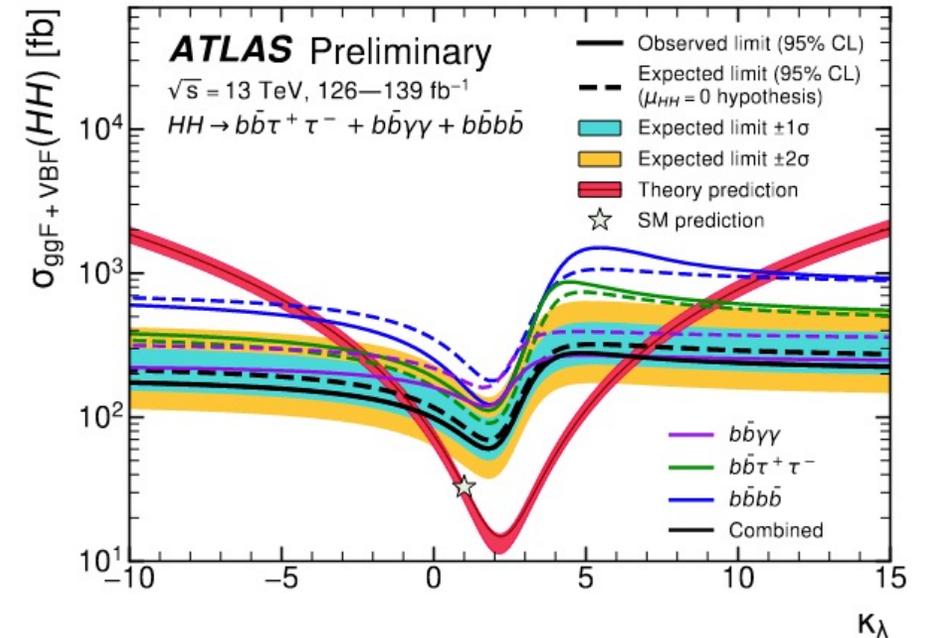


# Experimental bound on $\kappa_\lambda \equiv \lambda_{hhh}/\lambda_{hhh}^{\text{SM}}$

Current strongest limit:  $-0.4 < \kappa_\lambda < 6.3$  at 95% CL [ATLAS-CONF-2022-050].

## Assumptions:

- Simplest analysis assumes that all other Higgs couplings are SM-like.
- Non-resonant Higgs-boson pair production only deviates from the SM via a modified trilinear Higgs coupling (i.e., no heavy resonances).

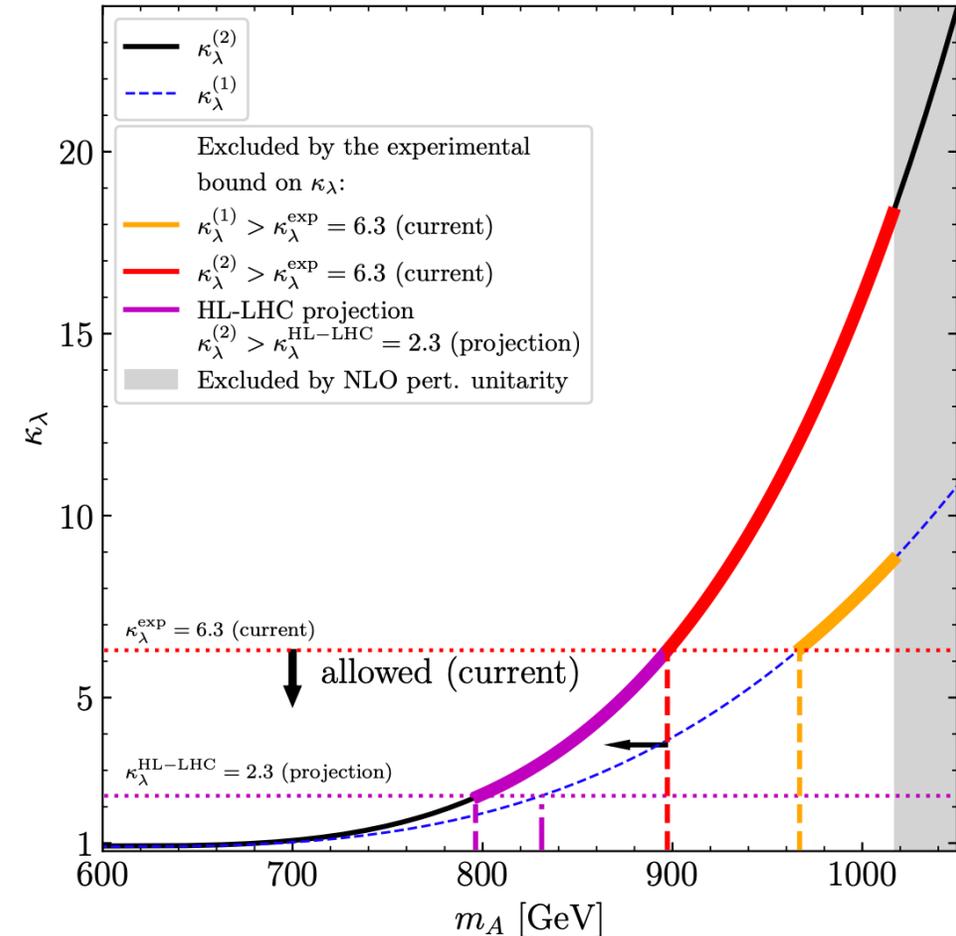


- Can we use this seemingly weak limit to constrain BSM models?
- Can large BSM deviations occur given other theoretical and experimental constraints?

# Trilinear Higgs coupling in the 2HDM

- Large deviations possible in the 2HDM without being in conflict with other measurements.
- Additional enhancement by 2L corrections.
- Maximal size bounded by perturbative unitarity.

2HDM type I,  $\alpha = \beta - \pi/2$ ,  $m_A = m_{H^\pm}$ ,  $M = m_H = 600$  GeV,  $\tan \beta = 2$



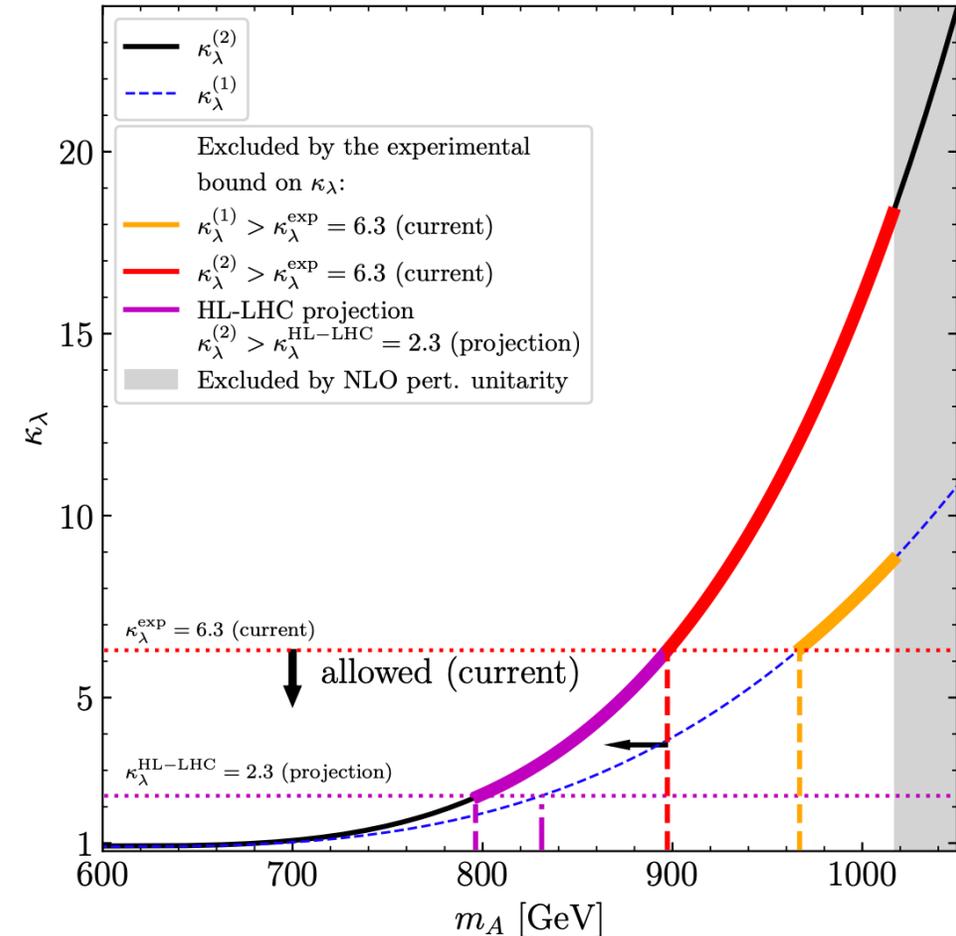
[HB,Braathen,Weiglein, 2202.03453]

# Trilinear Higgs coupling in the 2HDM

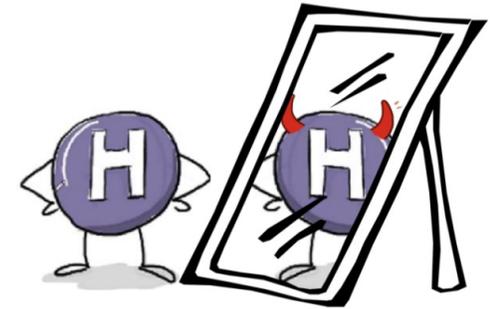
- Large deviations possible in the 2HDM without being in conflict with other measurements.
- Additional enhancement by 2L corrections.
- Maximal size bounded by perturbative unitarity.

Already current experimental limits on  $\kappa_\lambda$  probe so-far unconstrained BSM parameter space!

2HDM type I,  $\alpha = \beta - \pi/2$ ,  $m_A = m_{H^\pm}$ ,  $M = m_H = 600$  GeV,  $\tan \beta = 2$



[HB,Braathen,Weiglein, 2202.03453]



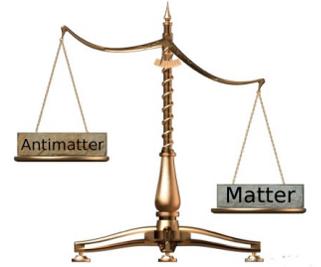
# The CP nature of the Higgs boson

*Does the Higgs sector provide additional sources of CP violation?*

# The CP nature of the Higgs boson

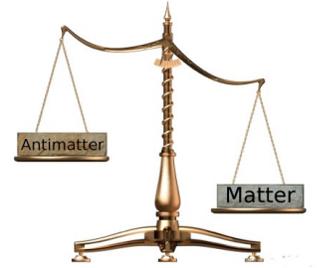


# The CP nature of the Higgs boson



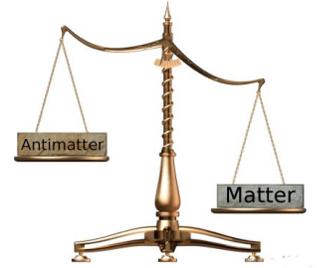
- Motivation: new sources of CP violation are necessary to explain the baryon asymmetry of the Universe.

# The CP nature of the Higgs boson



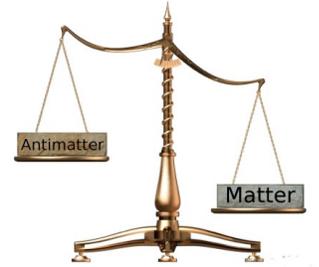
- Motivation: new sources of CP violation are necessary to explain the baryon asymmetry of the Universe.
- We know the Higgs boson is not a CP-odd state but it could be a CP-admixed state.

# The CP nature of the Higgs boson



- Motivation: new sources of CP violation are necessary to explain the baryon asymmetry of the Universe.
- We know the Higgs boson is not a CP-odd state but it could be a CP-admixed state.
- Parameterize CP-odd interactions using EFT framework by adding dimension-6 operators to the SM:
  - Gauge boson interactions:  $\Phi^\dagger \Phi W_{\mu\nu} \tilde{W}^{\mu\nu}, \Phi^\dagger \Phi B_{\mu\nu} \tilde{B}^{\mu\nu}, \Phi^\dagger \Phi W_{\mu\nu} \tilde{B}^{\mu\nu}, \Phi^\dagger \Phi G_{\mu\nu} \tilde{G}^{\mu\nu}$
  - Fermion interactions:  $\Phi^\dagger \Phi (Qu\tilde{\Phi}), \Phi^\dagger \Phi (Qd\Phi), \Phi^\dagger \Phi (Qe\Phi)$  with complex Wilson coefficients

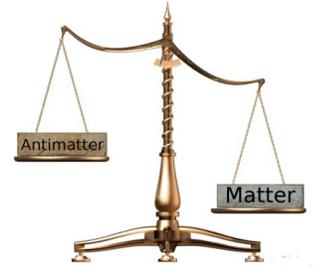
# The CP nature of the Higgs boson



- Motivation: new sources of CP violation are necessary to explain the baryon asymmetry of the Universe.
- We know the Higgs boson is not a CP-odd state but it could be a CP-admixed state.
- Parameterize CP-odd interactions using EFT framework by adding dimension-6 operators to the SM:
  - Gauge boson interactions:  $\Phi^\dagger \Phi W_{\mu\nu} \tilde{W}^{\mu\nu}, \Phi^\dagger \Phi B_{\mu\nu} \tilde{B}^{\mu\nu}, \Phi^\dagger \Phi W_{\mu\nu} \tilde{B}^{\mu\nu}, \Phi^\dagger \Phi G_{\mu\nu} \tilde{G}^{\mu\nu}$
  - Fermion interactions:  $\Phi^\dagger \Phi (Q_u \tilde{\Phi}), \Phi^\dagger \Phi (Q_d \Phi), \Phi^\dagger \Phi (Q_e \Phi)$  with complex Wilson coefficients

↳ Rewrite: 
$$\mathcal{L}_{\text{yuk}} = - \sum_{f=u,d,c,s,t,b,e,\mu,\tau} \frac{y_f^{\text{SM}}}{\sqrt{2}} \bar{f} (c_f + i\gamma_5 \tilde{c}_f) f H,$$

# The CP nature of the Higgs boson

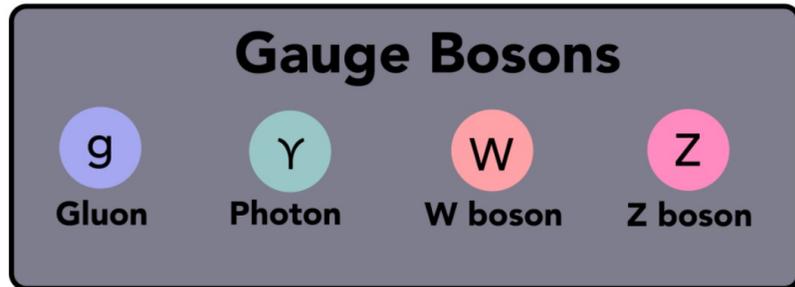


- Motivation: new sources of CP violation are necessary to explain the baryon asymmetry of the Universe.
- We know the Higgs boson is not a CP-odd state but it could be a CP-admixed state.
- Parameterize CP-odd interactions using EFT framework by adding dimension-6 operators to the SM:
  - Gauge boson interactions:  $\Phi^\dagger \Phi W_{\mu\nu} \tilde{W}^{\mu\nu}, \Phi^\dagger \Phi B_{\mu\nu} \tilde{B}^{\mu\nu}, \Phi^\dagger \Phi W_{\mu\nu} \tilde{B}^{\mu\nu}, \Phi^\dagger \Phi G_{\mu\nu} \tilde{G}^{\mu\nu}$
  - Fermion interactions:  $\Phi^\dagger \Phi (Qu\tilde{\Phi}), \Phi^\dagger \Phi (Qd\Phi), \Phi^\dagger \Phi (Qe\Phi)$  with complex Wilson coefficients

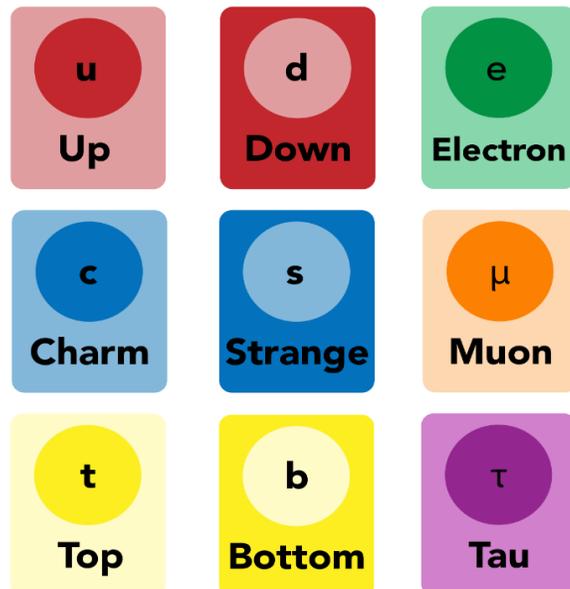
↳ Rewrite: 
$$\mathcal{L}_{\text{yuk}} = - \sum_{f=u,d,c,s,t,b,e,\mu,\tau} \frac{y_f^{\text{SM}}}{\sqrt{2}} \bar{f} (c_f + i\gamma_5 \tilde{c}_f) f H,$$

➔ What is the current status?

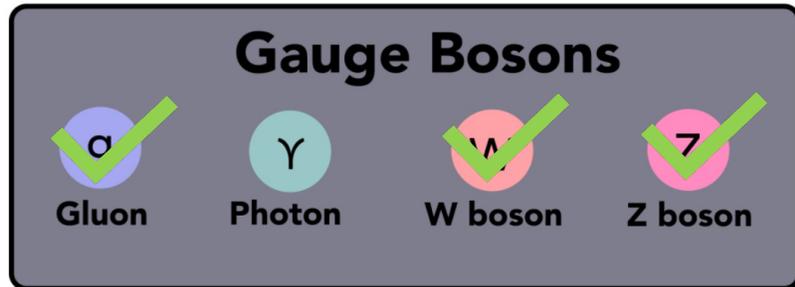
# Which CP structures are accessible at the LHC?



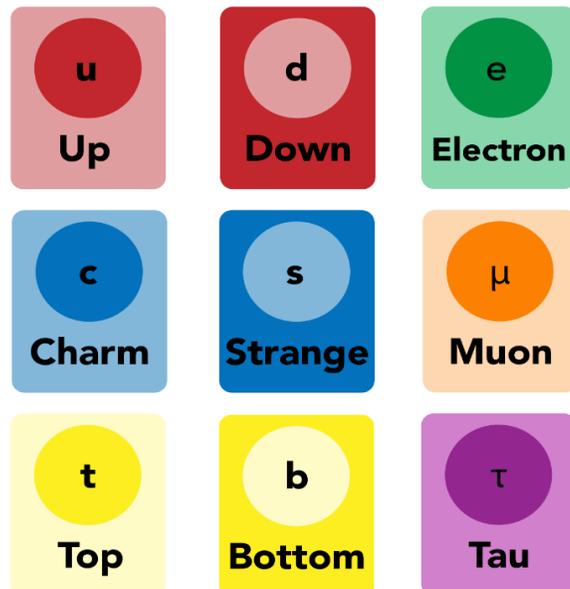
## Fermions



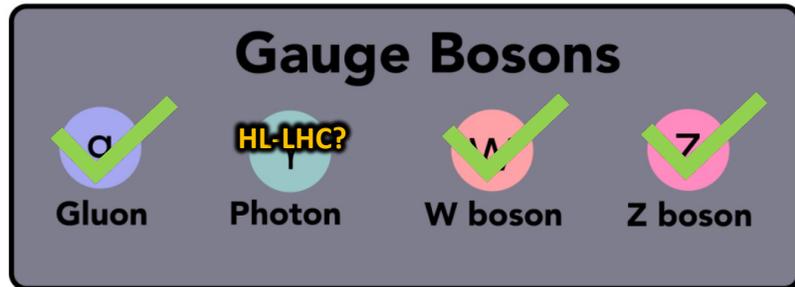
# Which CP structures are accessible at the LHC?



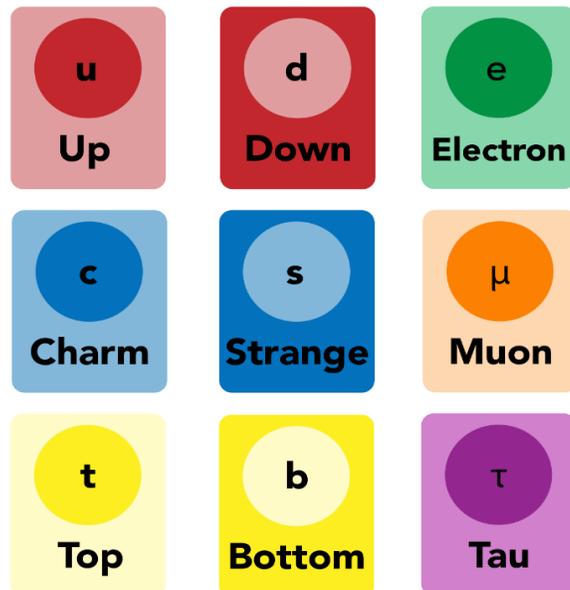
## Fermions



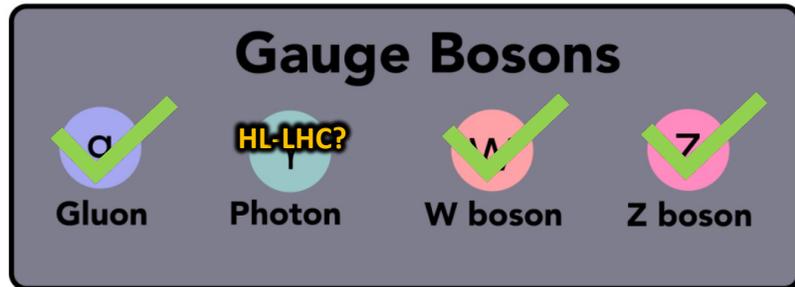
# Which CP structures are accessible at the LHC?



## Fermions

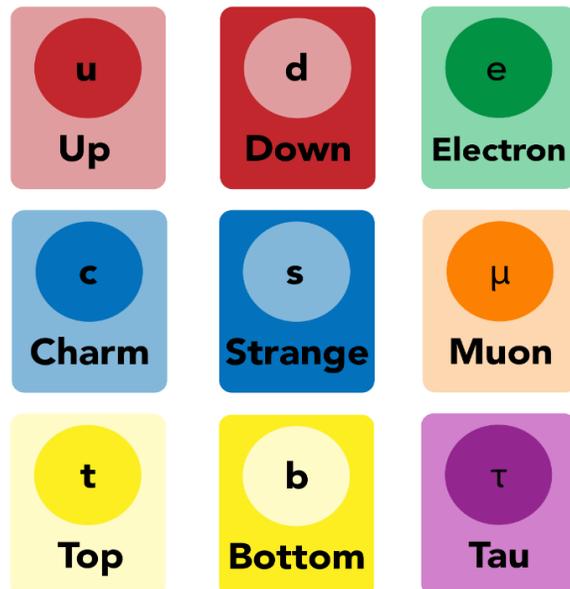


# Which CP structures are accessible at the LHC?

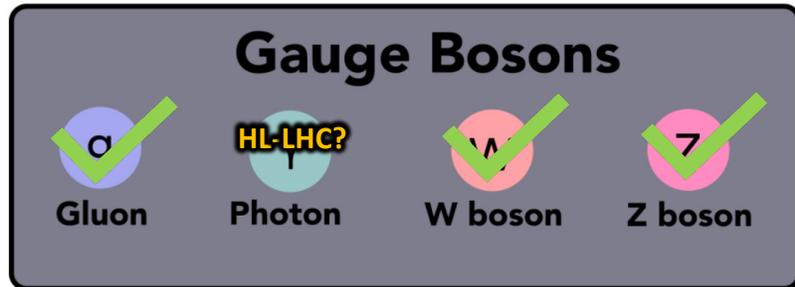


- CP structure of  $HWW$ ,  $HZZ$  interactions is comparably well-constrained. [ATLAS,CMS:...,2002.05315, 2104.12152,2109.13808,2202.06923,2205.05120]

## Fermions

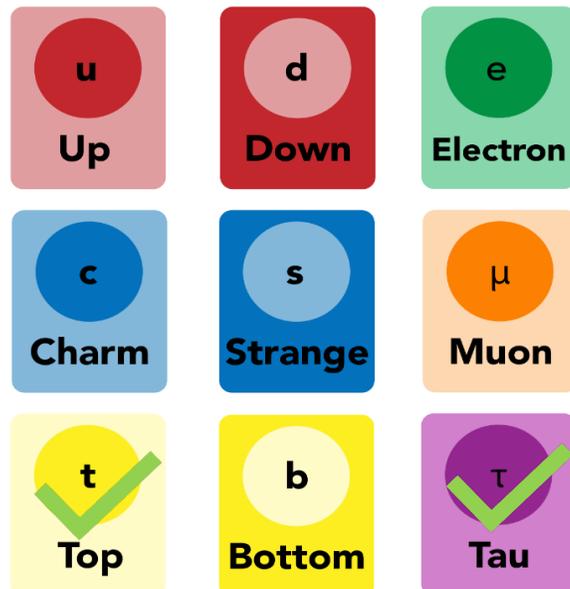


# Which CP structures are accessible at the LHC?

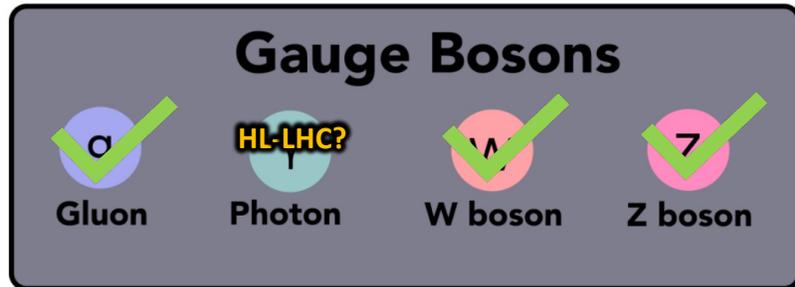


- CP structure of  $HWW$ ,  $HZZ$  interactions is comparably well-constrained. [ATLAS,CMS:...,2002.05315, 2104.12152,2109.13808,2202.06923,2205.05120]

## Fermions

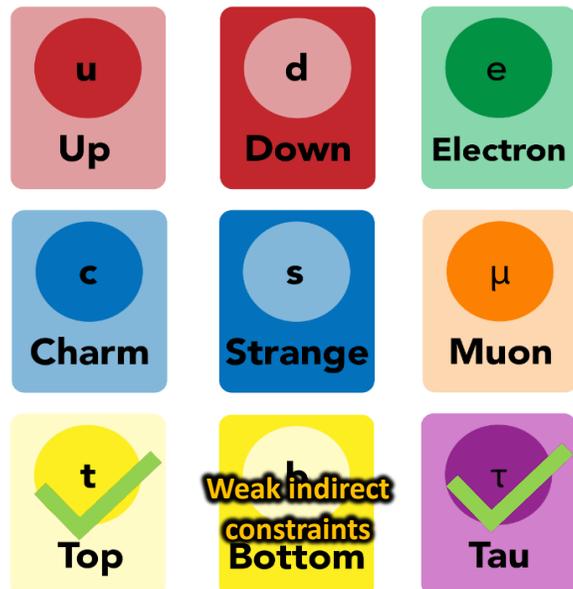


# Which CP structures are accessible at the LHC?

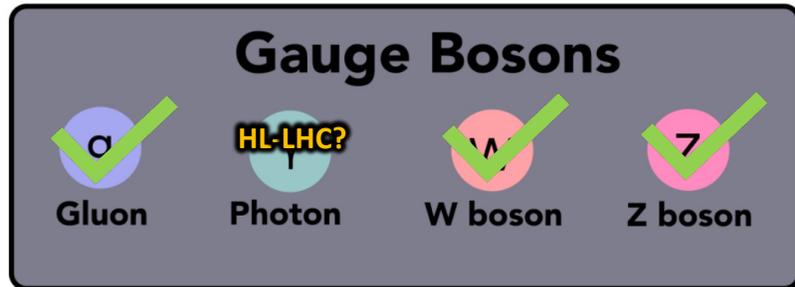


- CP structure of  $HWW$ ,  $HZZ$  interactions is comparably well-constrained. [ATLAS,CMS:...,2002.05315, 2104.12152,2109.13808,2202.06923,2205.05120]

## Fermions



# Which CP structures are accessible at the LHC?

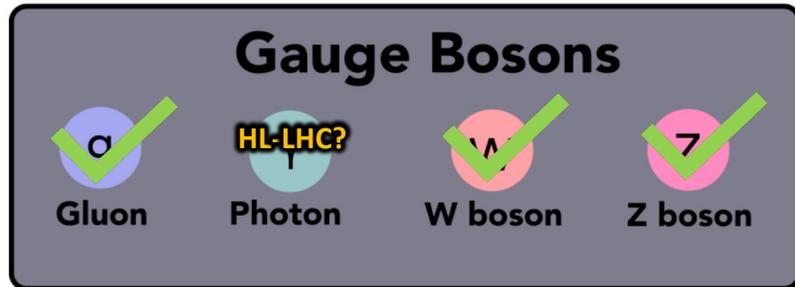


- CP structure of  $HWW$ ,  $HZZ$  interactions is comparably well-constrained. [ATLAS,CMS:...,2002.05315, 2104.12152,2109.13808,2202.06923,2205.05120]

## Fermions



# Which CP structures are accessible at the LHC?

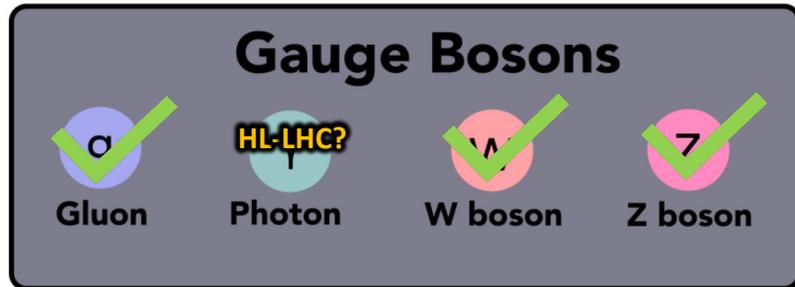


## Fermions



- CP structure of  $HWW$ ,  $HZZ$  interactions is comparably well-constrained. [ATLAS,CMS:...,2002.05315, 2104.12152,2109.13808,2202.06923,2205.05120]
- The CP structure of the  $Hf\bar{f}$ ,  $H\gamma\gamma$ ,  $Hgg$  interactions is far less known.

# Which CP structures are accessible at the LHC?

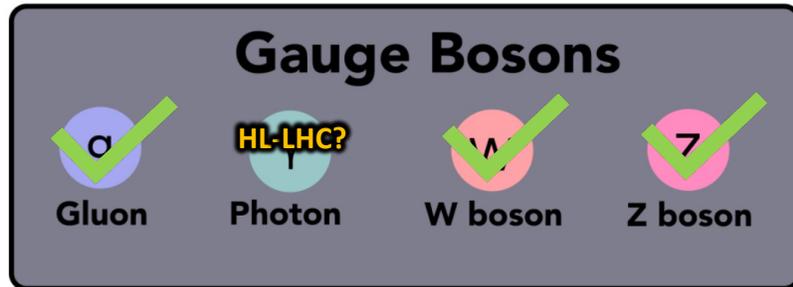


## Fermions



- CP structure of  $HWW$ ,  $HZZ$  interactions is comparably well-constrained. [ATLAS,CMS:...,2002.05315, 2104.12152,2109.13808,2202.06923,2205.05120]
- The CP structure of the  $Hf\bar{f}$ ,  $H\gamma\gamma$ ,  $Hgg$  interactions is far less known.
- Most BSM theories predict largest CP violation in  $Hf\bar{f}$ ,  $H\gamma\gamma$ ,  $Hgg$  couplings.

# Which CP structures are accessible at the LHC?



## Fermions



- CP structure of  $HWW$ ,  $HZZ$  interactions is comparably well-constrained. [ATLAS,CMS:...,2002.05315, 2104.12152,2109.13808,2202.06923,2205.05120]
- The CP structure of the  $Hf\bar{f}$ ,  $H\gamma\gamma$ ,  $Hgg$  interactions is far less known.
- Most BSM theories predict largest CP violation in  $Hf\bar{f}$ ,  $H\gamma\gamma$ ,  $Hgg$  couplings.



New ideas/techniques are needed to make the most of current and future data!

# Constraining CP violation

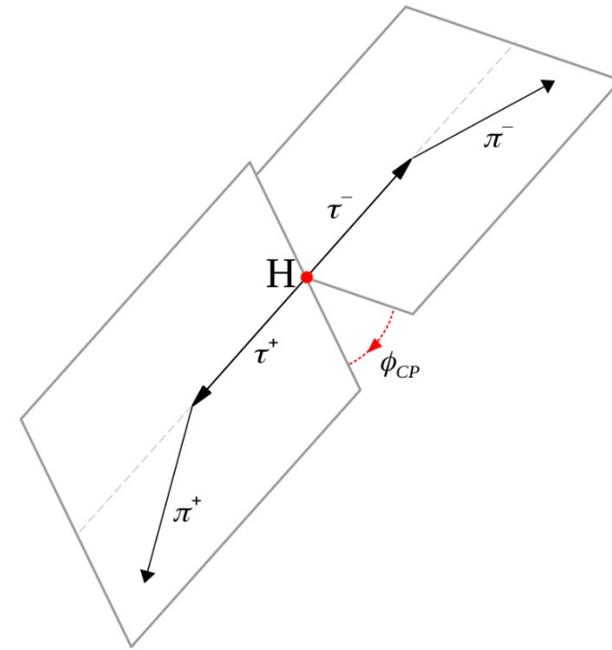
CP violation in the Higgs sector can be constrained using:

- **Pure CP-odd observables:**
  - Unambiguous markers for CP violation: e.g.
    - EDM measurements,
    - decay angle in  $H \rightarrow \tau^+ \tau^-$ .
  - Typically requires to access polarization of particles coupling to the Higgs.
  - Experimentally difficult for many LHC processes (i.e., top-associated Higgs production).
  - Almost impossible for  $H \rightarrow b\bar{b}$  or  $H \rightarrow \mu^+ \mu^-$

# Constraining CP violation

CP violation in the Higgs sector can be constrained using:

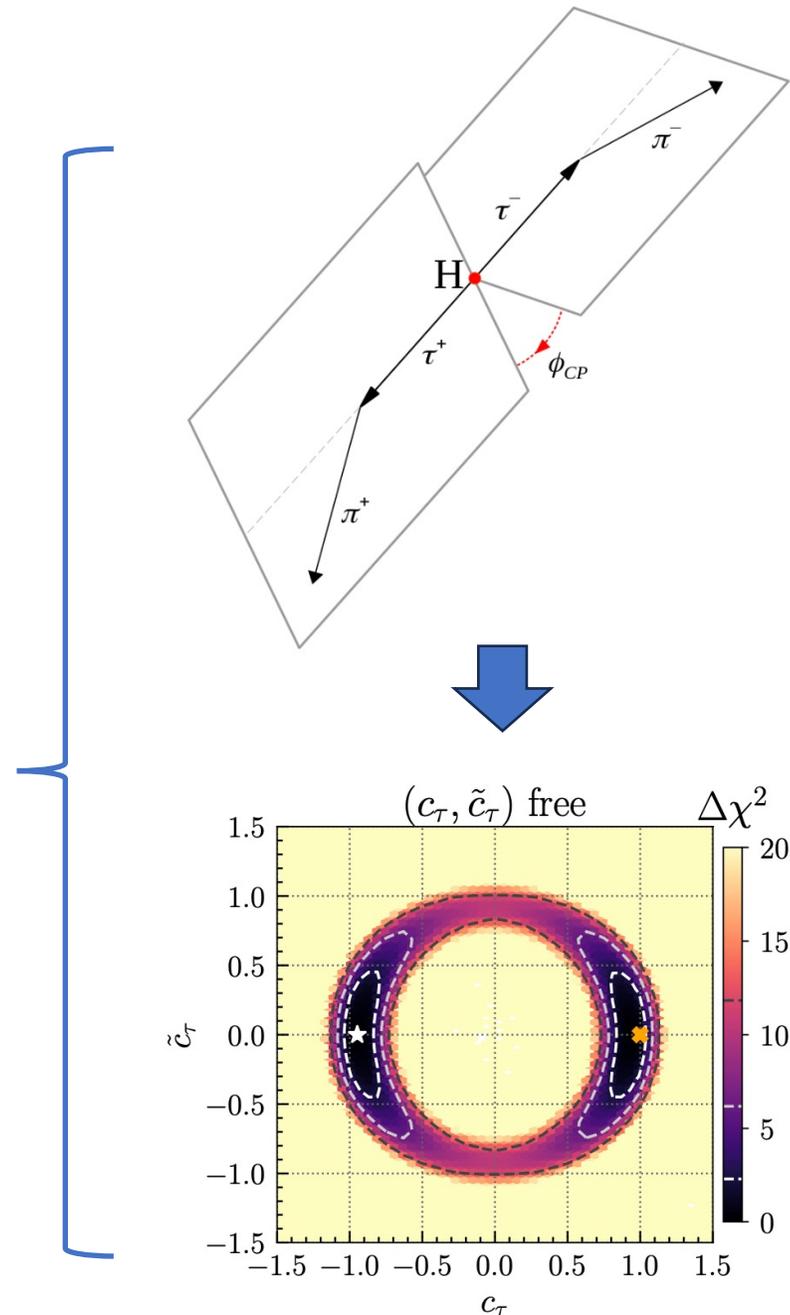
- **Pure CP-odd observables:**
  - Unambiguous markers for CP violation: e.g.
    - EDM measurements,
    - decay angle in  $H \rightarrow \tau^+ \tau^-$ .
  - Typically requires to access polarization of particles coupling to the Higgs.
  - Experimentally difficult for many LHC processes (i.e., top-associated Higgs production).
  - Almost impossible for  $H \rightarrow b\bar{b}$  or  $H \rightarrow \mu^+ \mu^-$



# Constraining CP violation

CP violation in the Higgs sector can be constrained using:

- **Pure CP-odd observables:**
  - Unambiguous markers for CP violation: e.g.
    - EDM measurements,
    - decay angle in  $H \rightarrow \tau^+ \tau^-$ .
  - Typically requires to access polarization of particles coupling to the Higgs.
  - Experimentally difficult for many LHC processes (i.e., top-associated Higgs production).
  - Almost impossible for  $H \rightarrow b\bar{b}$  or  $H \rightarrow \mu^+ \mu^-$



# Constraining CP violation

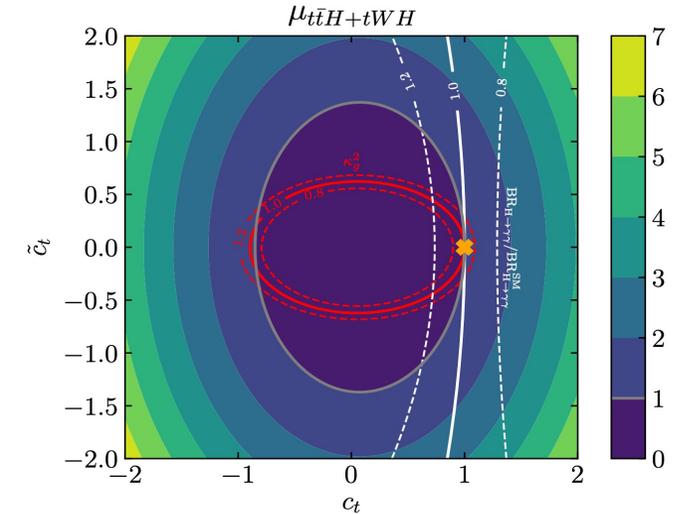
CP violation in the Higgs sector can be constrained using:

- **Pure CP-even observables:**
  - Many rate measurements are indirectly sensitive: e.g.  $ggH$ .
  - Subtle effects in kinematic distributions of CP-even observables (e.g.  $p_{T,H}$  in  $t\bar{t}H$ ).
- Deviations from SM need not be due to CP violation  
→ degeneracies with non-CPV BSM effects.

# Constraining CP violation

CP violation in the Higgs sector can be constrained using:

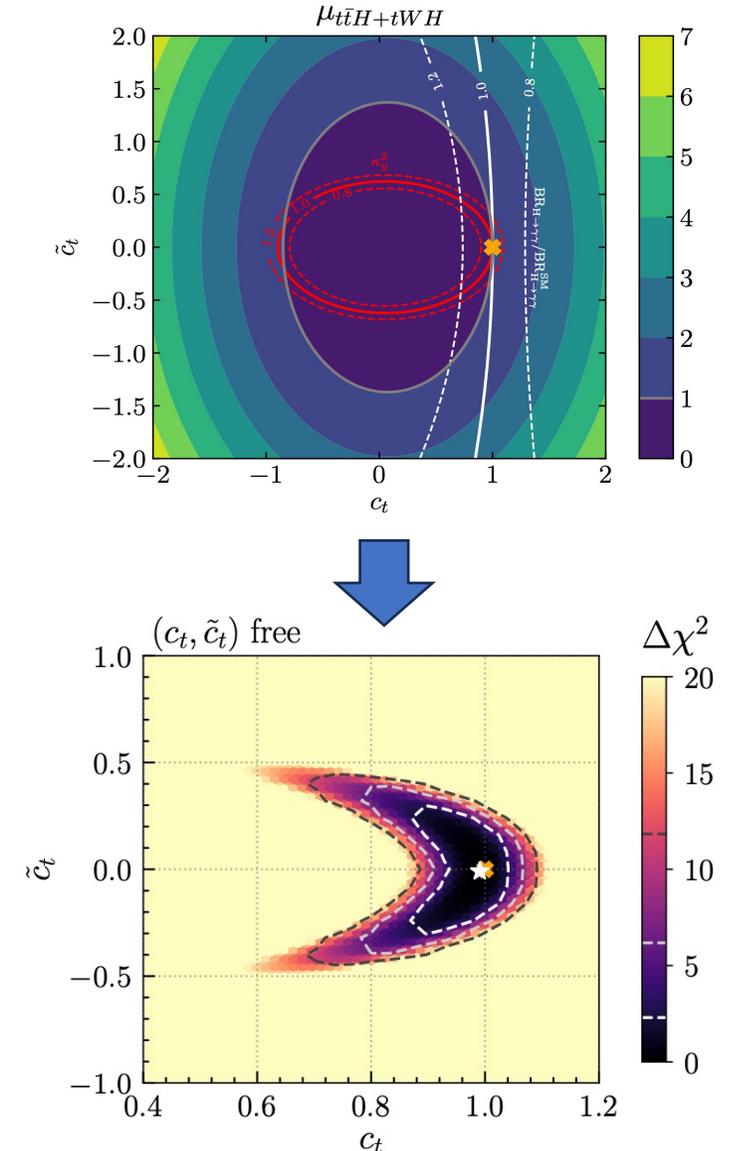
- **Pure CP-even observables:**
  - Many rate measurements are indirectly sensitive: e.g.  $ggH$ .
  - Subtle effects in kinematic distributions of CP-even observables (e.g.  $p_{T,H}$  in  $t\bar{t}H$ ).
  - Deviations from SM need not be due to CP violation  
→ degeneracies with non-CPV BSM effects.



# Constraining CP violation

CP violation in the Higgs sector can be constrained using:

- **Pure CP-even observables:**
  - Many rate measurements are indirectly sensitive: e.g.  $ggH$ .
  - Subtle effects in kinematic distributions of CP-even observables (e.g.  $p_{T,H}$  in  $t\bar{t}H$ ).
- Deviations from SM need not be due to CP violation  
→ degeneracies with non-CPV BSM effects.



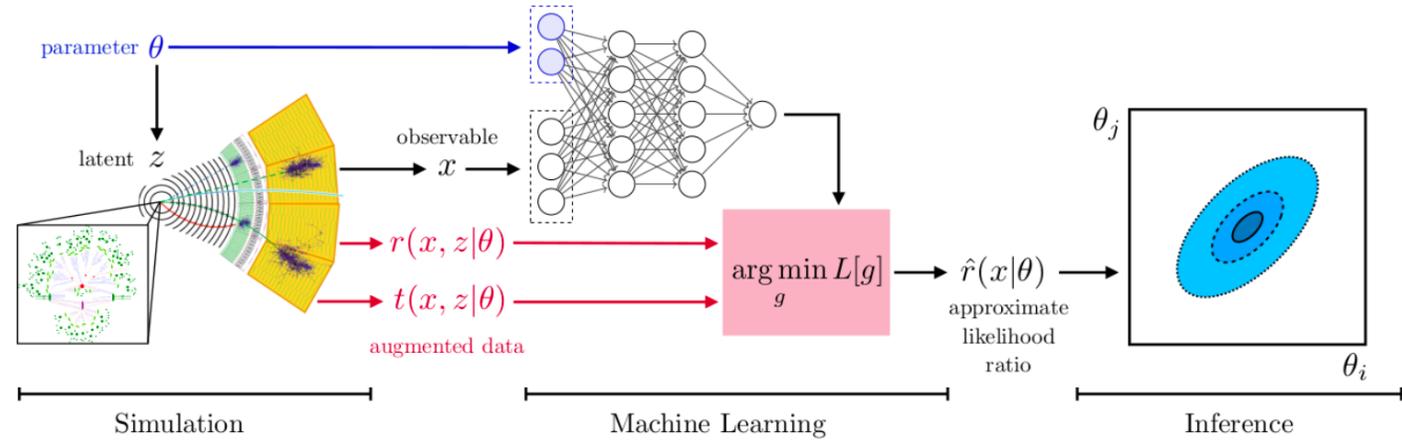
# Constraining CP violation

- **Multivariate analyses:**
  - Exploit full kinematic information using machine learning.
  - Often mixes CP-even and CP-odd observables.
  - High sensitivity.
  - Can be difficult to reinterpret.

# Constraining CP violation

- **Multivariate analyses:**

- Exploit full kinematic information using machine learning.
- Often mixes CP-even and CP-odd observables.
- High sensitivity.
- Can be difficult to reinterpret.

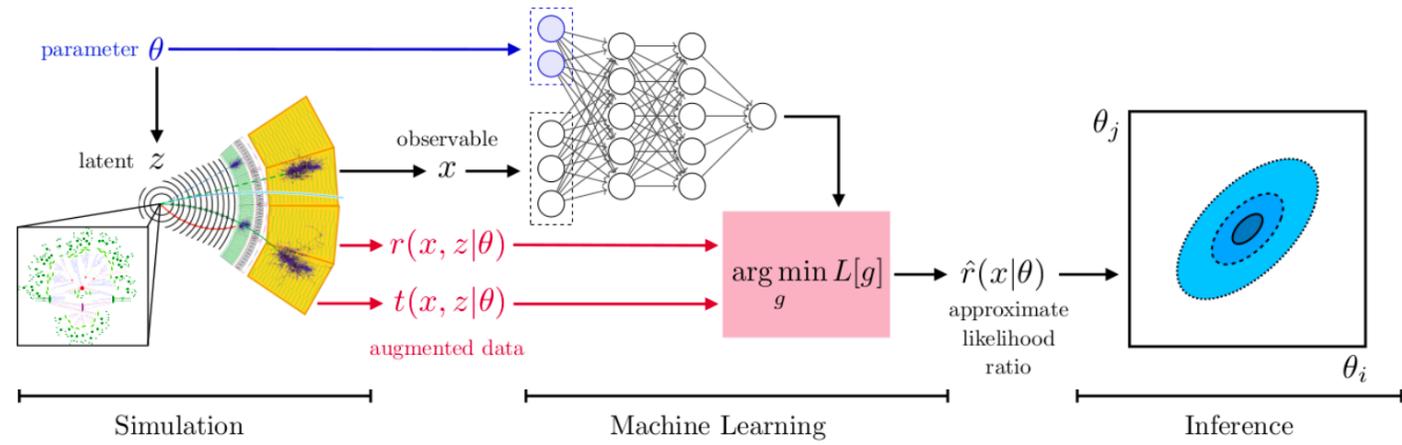


[e.g. simulation-based inference, Brehmer et al., 1805.00013, ...]

# Constraining CP violation

- **Multivariate analyses:**

- Exploit full kinematic information using machine learning.
- Often mixes CP-even and CP-odd observables.
- High sensitivity.
- Can be difficult to reinterpret.

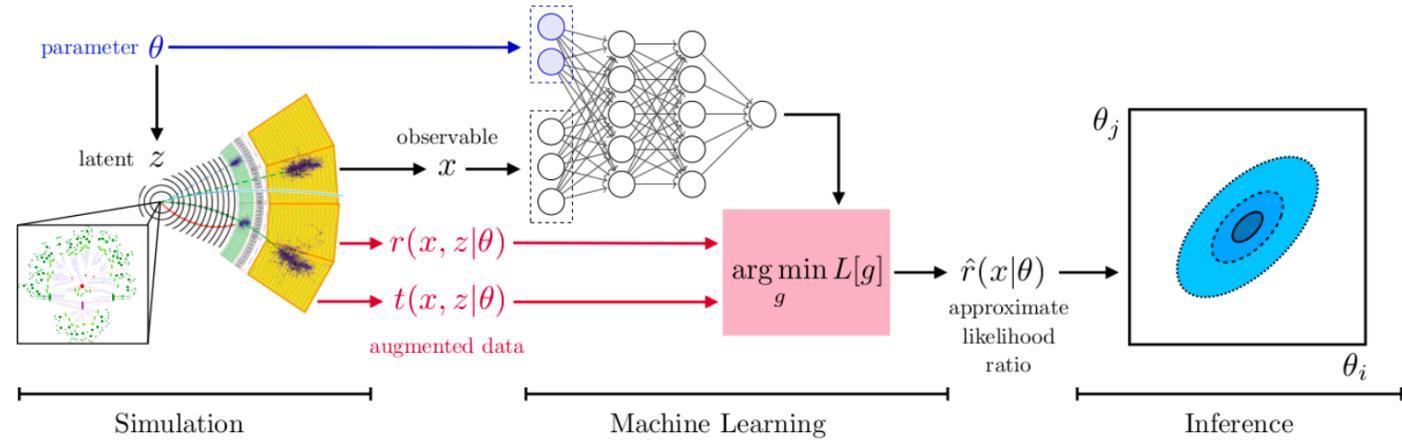


[e.g. simulation-based inference, Brehmer et al., 1805.00013, ...]

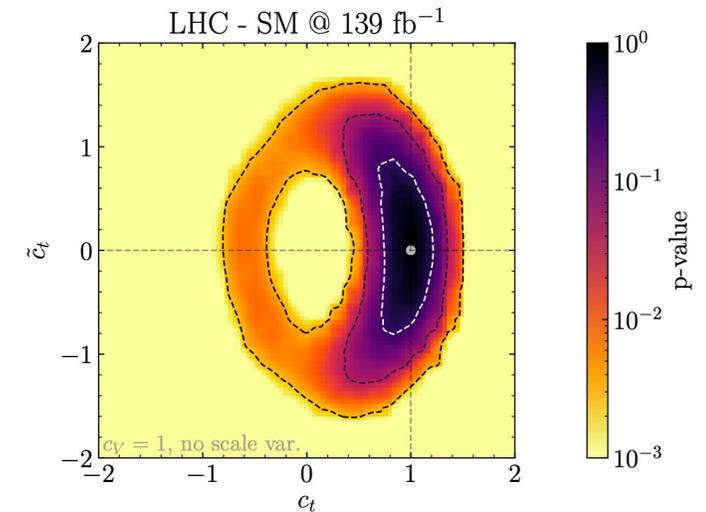
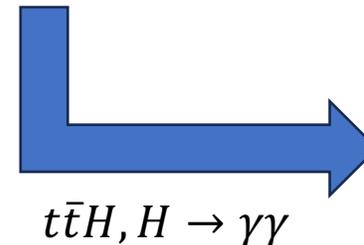
# Constraining CP violation

- **Multivariate analyses:**

- Exploit full kinematic information using machine learning.
- Often mixes CP-even and CP-odd observables.
- High sensitivity.
- Can be difficult to reinterpret.



[e.g. simulation-based inference, Brehmer et al.,1805.00013, ...]

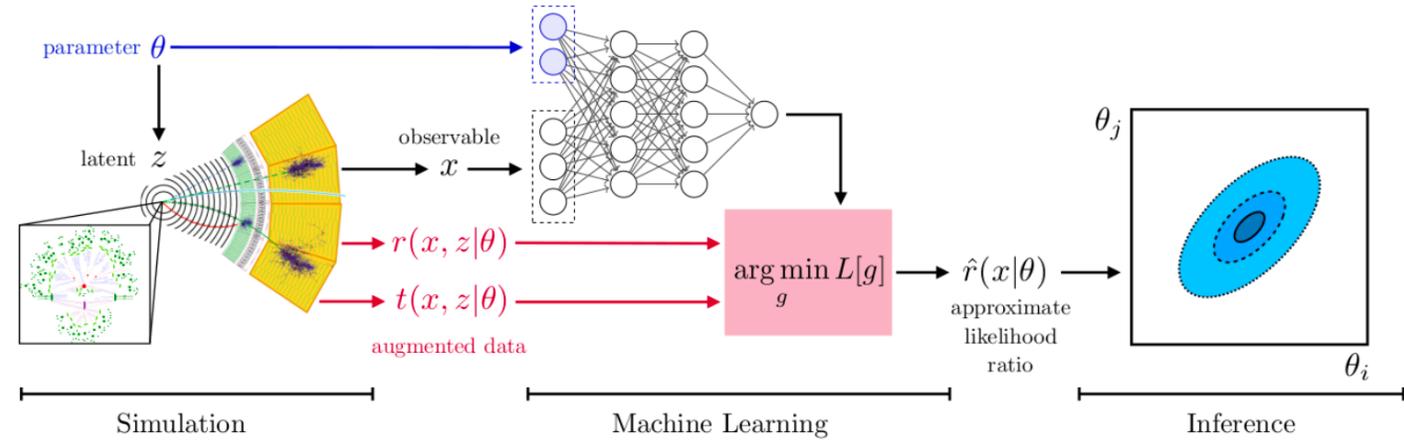


[HB&Brass,2110.10177]

# Constraining CP violation

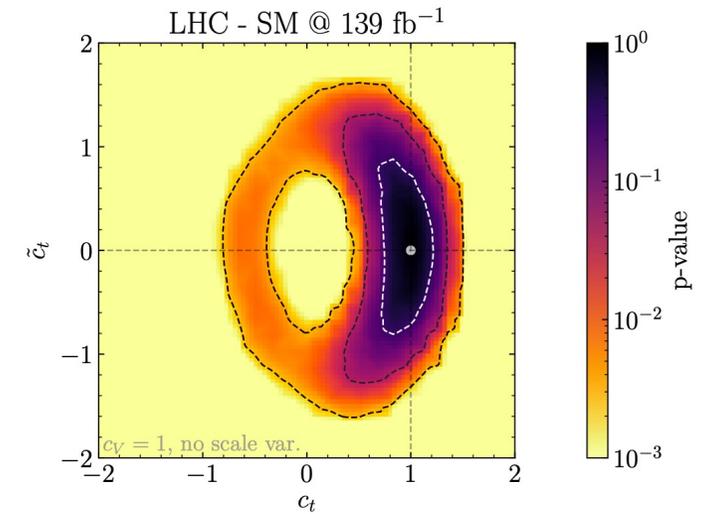
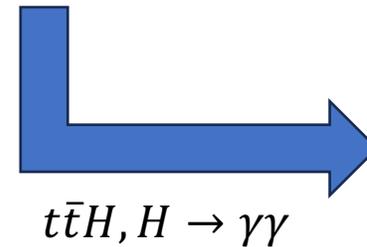
- **Multivariate analyses:**

- Exploit full kinematic information using machine learning.
- Often mixes CP-even and CP-odd observables.
- High sensitivity.
- Can be difficult to reinterpret.



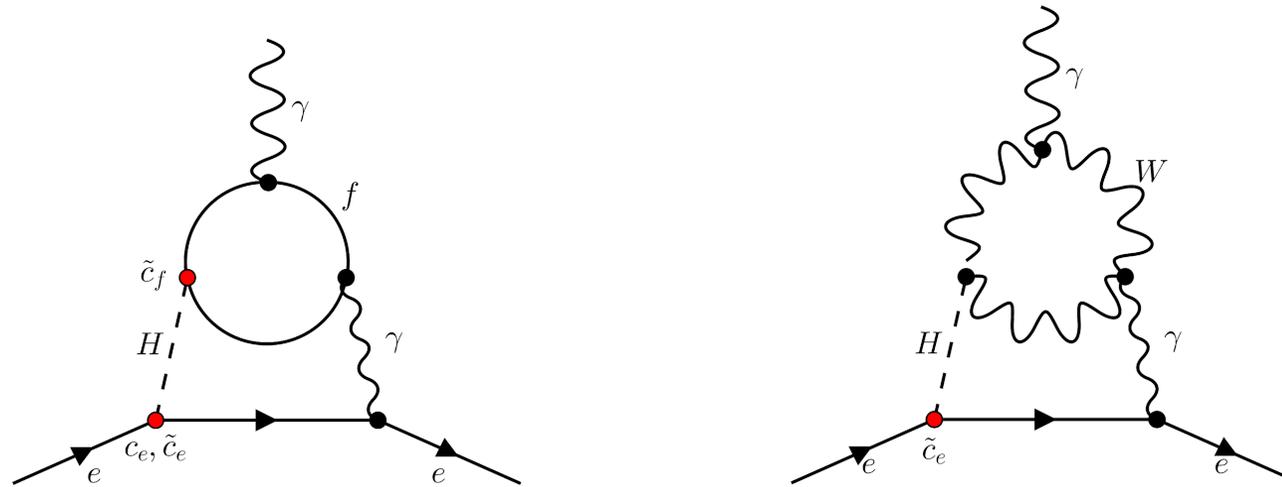
[e.g. simulation-based inference, Brehmer et al.,1805.00013, ...]

Exploit and combine all three complementary approaches to learn as much as possible!



[HB&Brass,2110.10177]

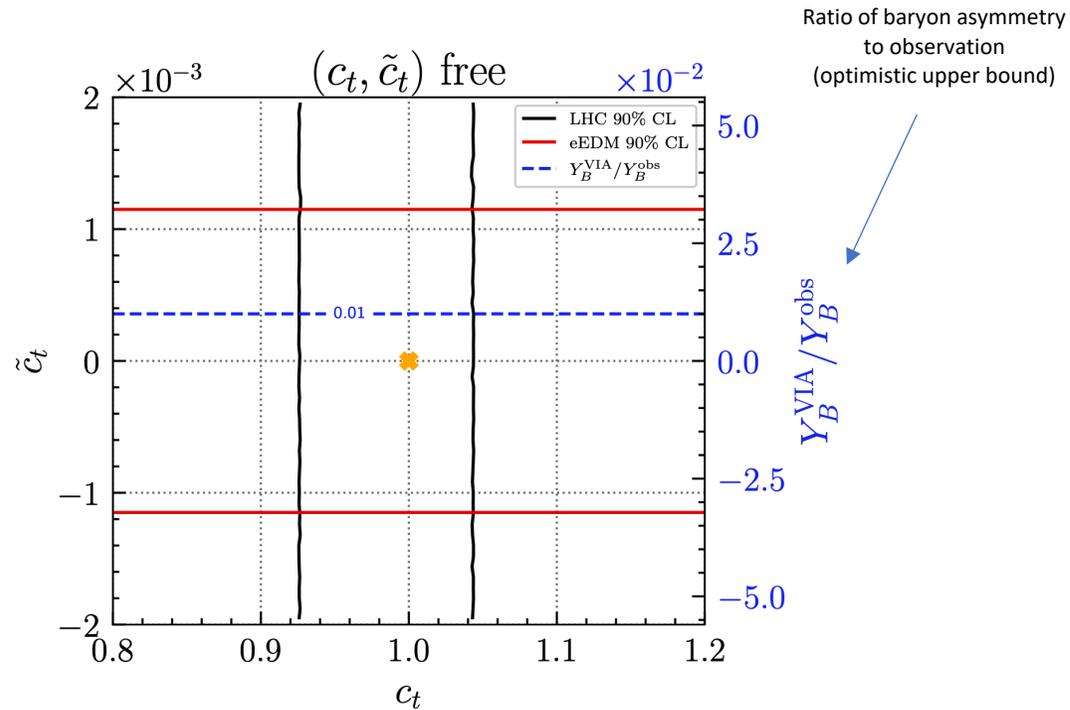
# Complementarity with EDM constraints



- Several EDMs are sensitive to CP violation in the Higgs sector.
- Consider here only constraints from theoretically cleanest EDM: the electron EDM.
- Limit by ACME collaboration:  $d_e^{\text{ACME}} = 1.1 \cdot 10^{-29} e \text{ cm}$  at 90% CL.

# Complementarity with EDM constraints: $t$ and $\tau$

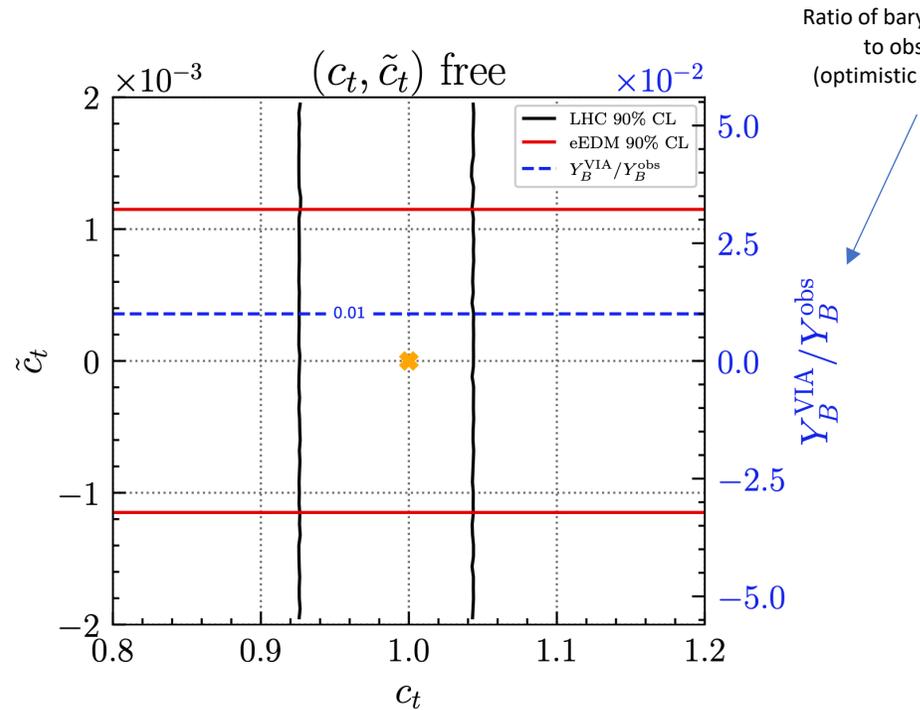
[HB et al., 2202.11753]



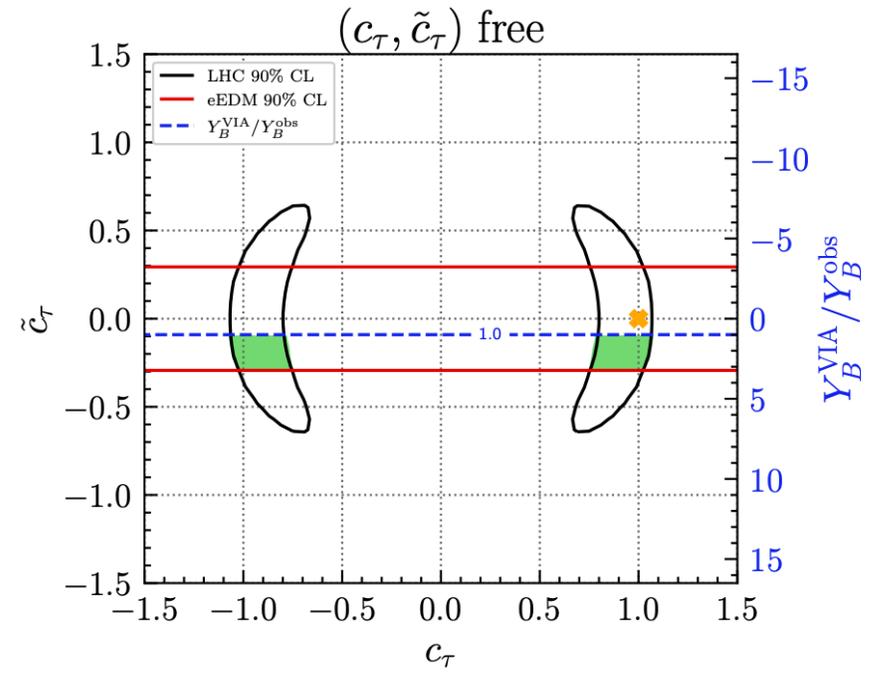
Very strong constraints on CP-odd top-Yukawa coupling.

# Complementarity with EDM constraints: $t$ and $\tau$

[HB et al., 2202.11753]



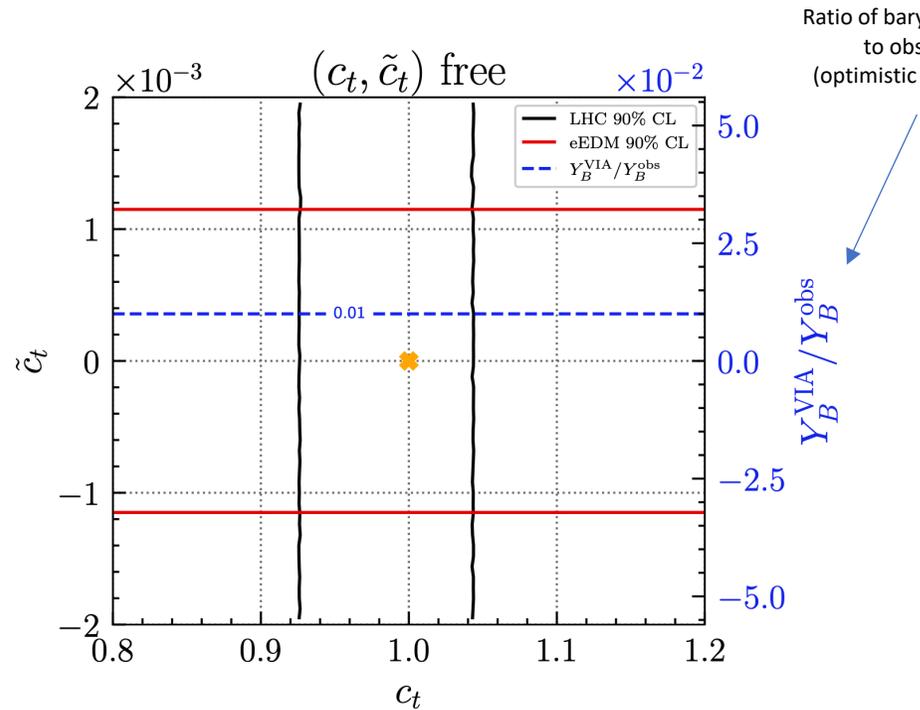
Very strong constraints on CP-odd top-Yukawa coupling.



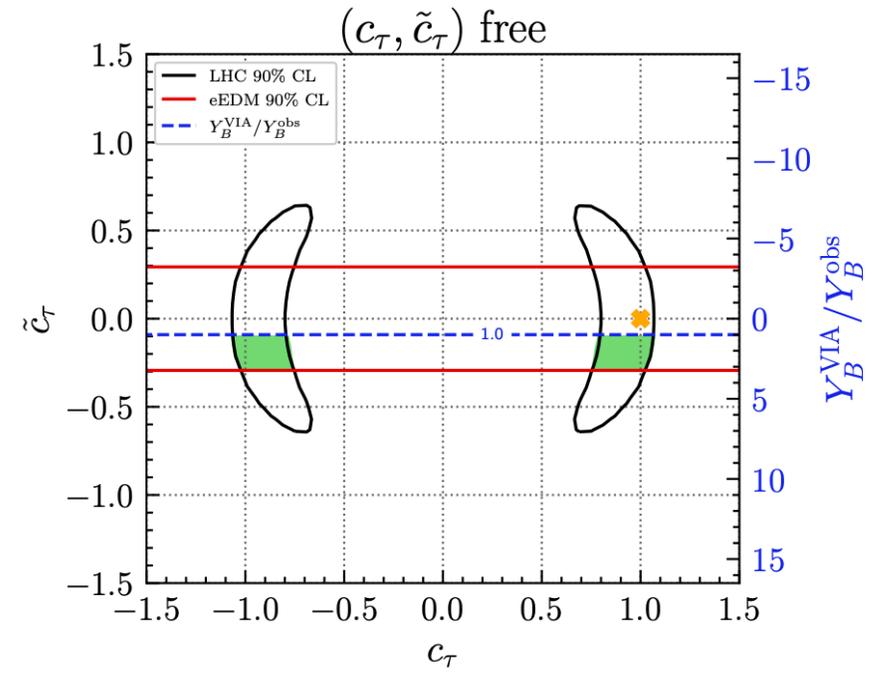
CP-odd  $\tau$  coupling can contribute significantly to baryon asymmetry.

# Complementarity with EDM constraints: $t$ and $\tau$

[HB et al., 2202.11753]



Very strong constraints on CP-odd top-Yukawa coupling.



CP-odd  $\tau$  coupling can contribute significantly to baryon asymmetry.

→ updated EDM measurement almost completely excludes green area

# EDM > LHC?

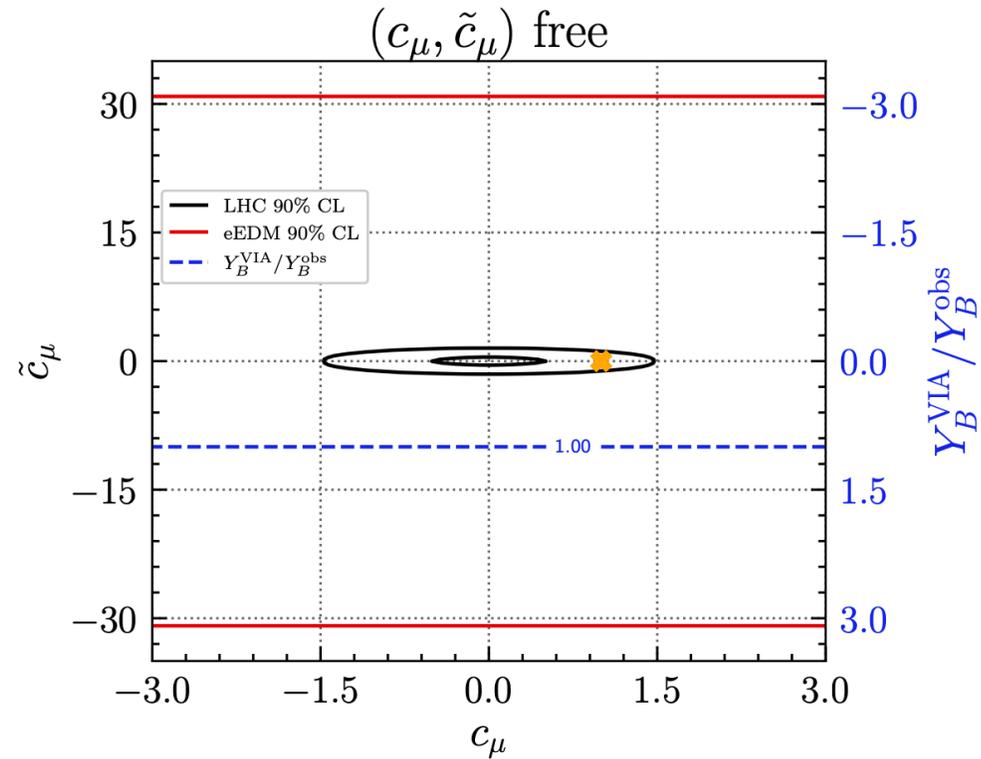
[see also Fuchs et al.,1911.08495]

EDM > LHC? **No.**

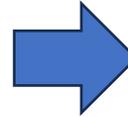
[see also Fuchs et al.,1911.08495]

# EDM > LHC? No.

[see also Fuchs et al.,1911.08495]

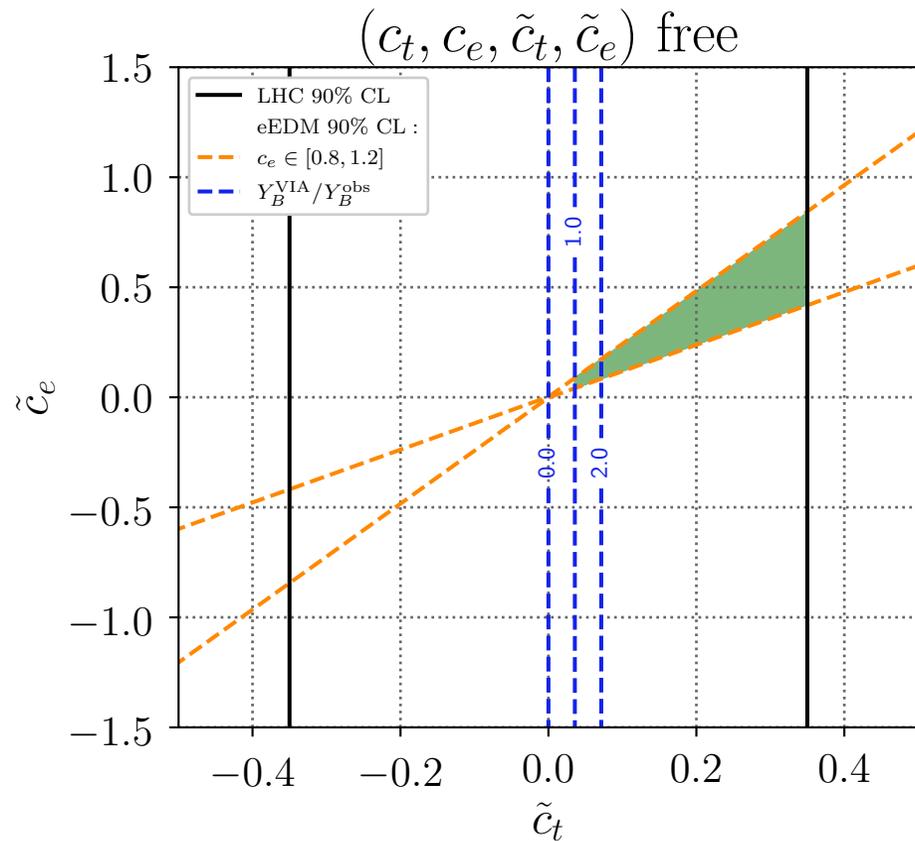


$$\Gamma_{H \rightarrow \mu^+ \mu^-} \propto c_\mu^2 + \tilde{c}_\mu^2$$



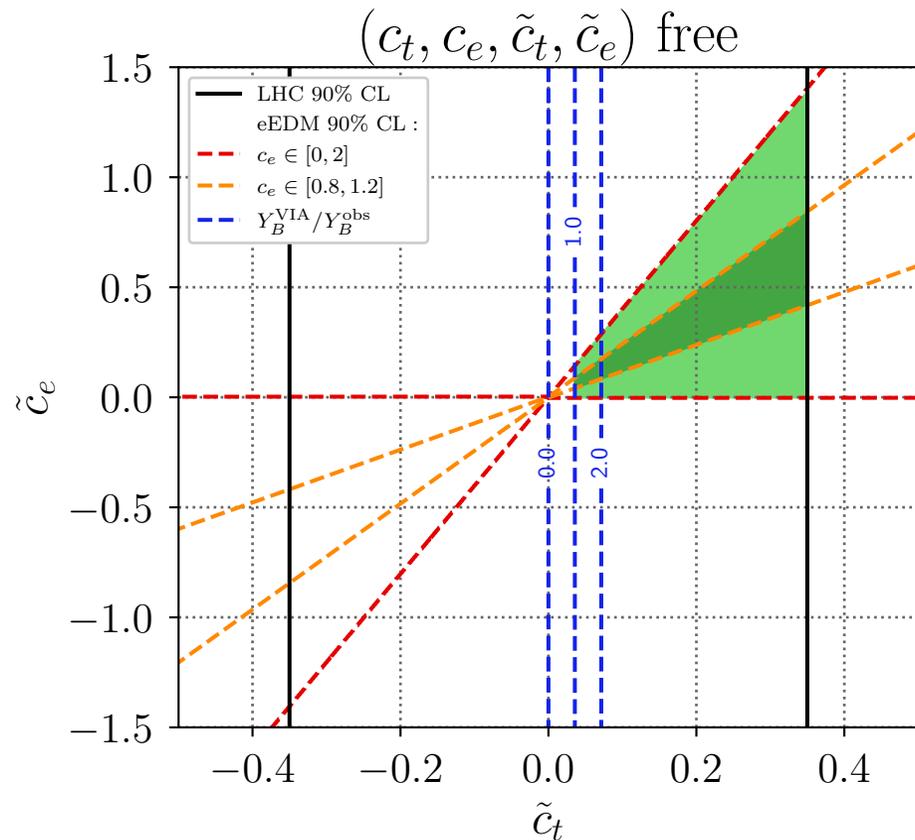
CP-insensitive  $H \rightarrow \mu^+ \mu^-$  rate measurement outperforms EDM constraint.

# Dependence on electron-Yukawa coupling



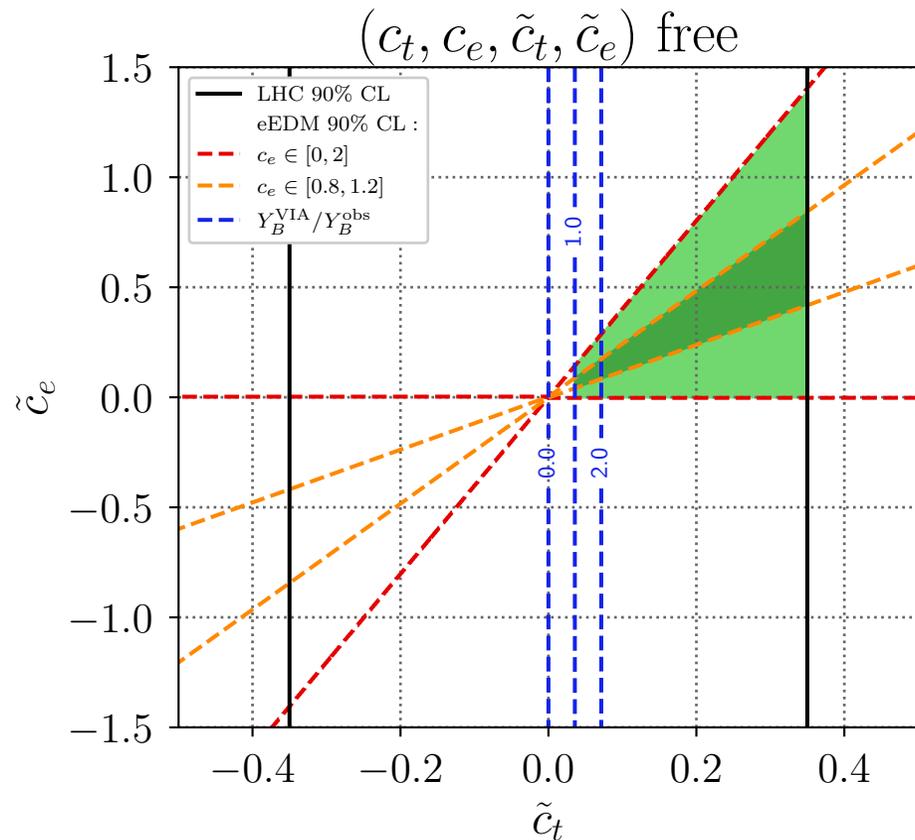
- Electron Yukawa-coupling only very weakly constrained ( $g_e \leq 268$  at 95% CL).
- If  $c_e$  smaller, eEDM significantly weakened.
- Moreover, we can fine-tune CP-odd electron-Yukawa coupling such that  $d_e < d_e^{\text{ACME}}$ .
- Neutron EDM has similar dependence on first-generation quark-Yukawa couplings.

# Dependence on electron-Yukawa coupling

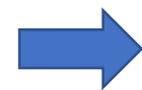


- Electron Yukawa-coupling only very weakly constrained ( $g_e \leq 268$  at 95% CL).
- If  $c_e$  smaller, eEDM significantly weakened.
- Moreover, we can fine-tune CP-odd electron-Yukawa coupling such that  $d_e < d_e^{\text{ACME}}$ .
- Neutron EDM has similar dependence on first-generation quark-Yukawa couplings.

# Dependence on electron-Yukawa coupling



- Electron Yukawa-coupling only very weakly constrained ( $g_e \leq 268$  at 95% CL).
- If  $c_e$  smaller, eEDM significantly weakened.
- Moreover, we can fine-tune CP-odd electron-Yukawa coupling such that  $d_e < d_e^{\text{ACME}}$ .
- Neutron EDM has similar dependence on first-generation quark-Yukawa couplings.

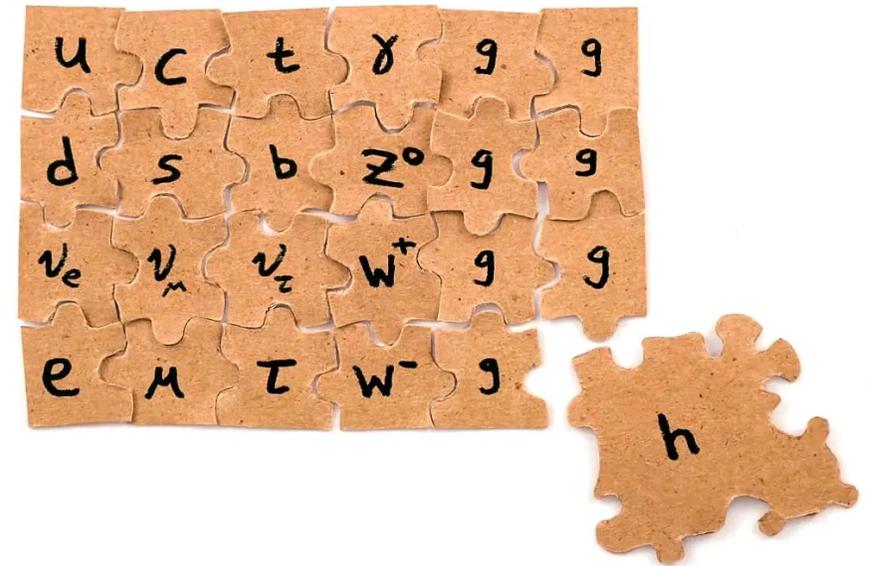


LHC bounds important since they do not depend on 1<sup>st</sup> gen. Yukawa couplings.

# Conclusions

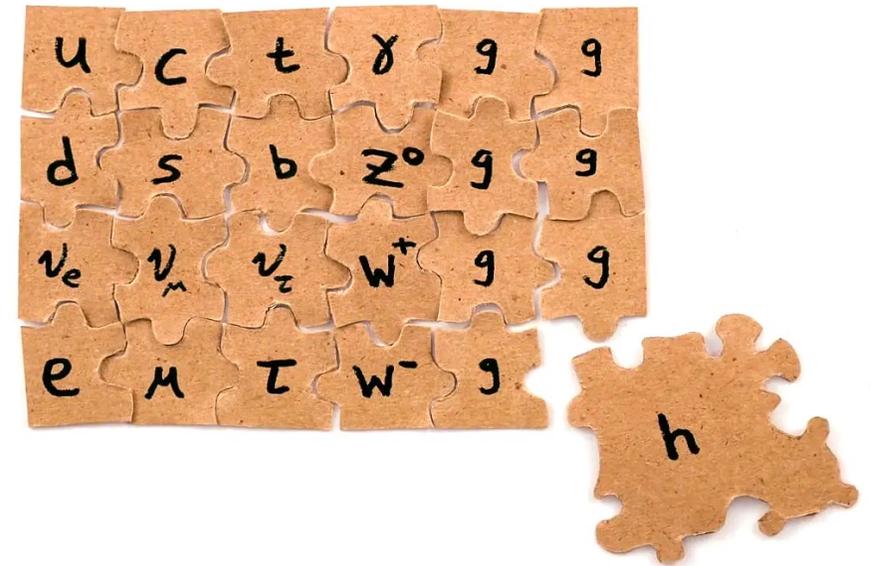
# Conclusions

- The Higgs is not the last missing puzzle piece of the SM but could be the link to many BSM scenarios.
- Higgs precision measurements and precision predictions are crucial to understand electroweak symmetry breaking.
- Existing measurements already teach us a lot about possible BSM extensions.
- Much work still left to do:
  - Light Yukawas,
  - Higgs CP structure,
  - Higgs potential,
  - Higgs width,
  - ...



# Conclusions

- The Higgs is not the last missing puzzle piece of the SM but could be the link to many BSM scenarios.
- Higgs precision measurements and precision predictions are crucial to understand electroweak symmetry breaking.
- Existing measurements already teach us a lot about possible BSM extensions.
- Much work still left to do:
  - Light Yukawas,
  - Higgs CP structure,
  - Higgs potential,
  - Higgs width,
  - ...



**The Higgs will keep us busy for many decades to come!**

# Appendix

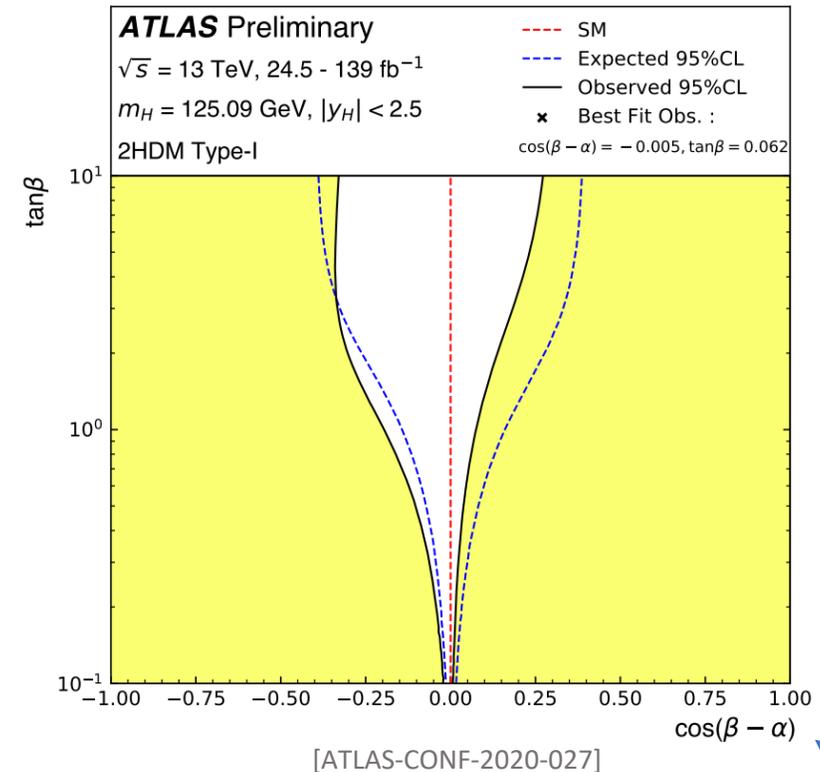
# Impact of Higgs precision measurements on 2HDM

- BSM benchmark model: 2HDM type-I
- Two Higgs doublets  $\rightarrow$  CP-even  $h_1, h_2$  (and  $A, H^\pm$ )
  - $\tan\beta$ : ratio of vevs
  - $\alpha$ : mixing angle
  - $m_{h_1} < m_{h_2}$
- Scaling of vector boson couplings

$$c(h_1 VV) \propto \sin(\beta - \alpha)$$

$$c(h_2 VV) \propto \cos(\beta - \alpha)$$

$\rightarrow$  Measurements enforce approximate alignment of the SM-like Higgs with the electroweak vacuum.



assumes that  $h_1$  is  $h_{125}$

# Impact of Higgs precision measurements on 2HDM

- BSM benchmark model: 2HDM type-I
- Two Higgs doublets  $\rightarrow$  CP-even  $h_1, h_2$  (and  $A, H^\pm$ )
  - $\tan\beta$ : ratio of vevs
  - $\alpha$ : mixing angle
  - $m_{h_1} < m_{h_2}$
- Scaling of vector boson couplings

$$c(h_1 VV) \propto \sin(\beta - \alpha)$$

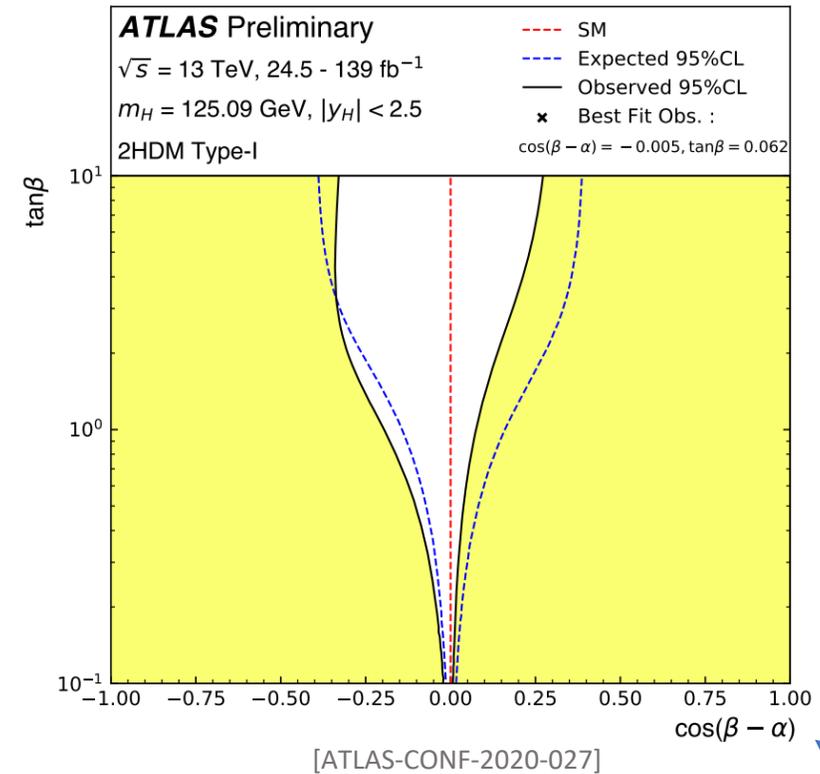
$$c(h_2 VV) \propto \cos(\beta - \alpha)$$

$\rightarrow$  Measurements enforce approximate alignment of the SM-like Higgs with the electroweak vacuum.



$$\sin(\beta - \alpha) \sim 1 \text{ if } h_1 \text{ is } h_{125} \text{ or}$$

$$\cos(\beta - \alpha) \sim 1 \text{ if } h_2 \text{ is } h_{125}$$



assumes that  $h_1$  is  $h_{125}$

# Impact of Higgs precision measurements on 2HDM

- BSM benchmark model: 2HDM type-I
- Two Higgs doublets  $\rightarrow$  CP-even  $h_1, h_2$  (and  $A, H^\pm$ )
  - $\tan\beta$ : ratio of vevs
  - $\alpha$ : mixing angle
  - $m_{h_1} < m_{h_2}$
- Scaling of vector boson couplings

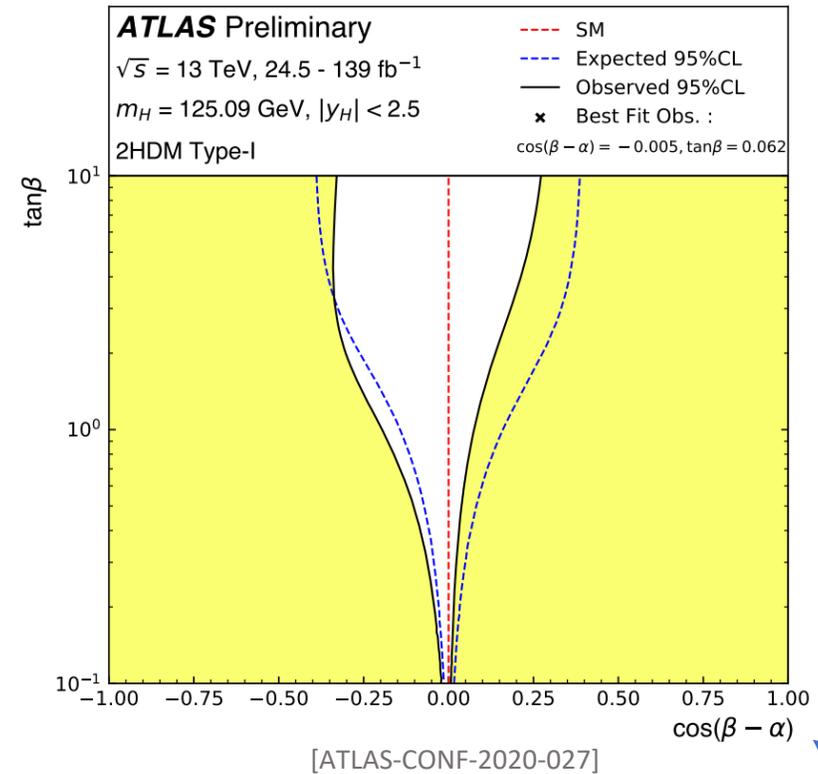
$$c(h_1 VV) \propto \sin(\beta - \alpha)$$

$$c(h_2 VV) \propto \cos(\beta - \alpha)$$

$\rightarrow$  Measurements enforce approximate alignment of the SM-like Higgs with the electroweak vacuum.



$\sin(\beta - \alpha) \sim 1$  if  $h_1$  is  $h_{125}$  or  
 $\cos(\beta - \alpha) \sim 1$  if  $h_2$  is  $h_{125}$

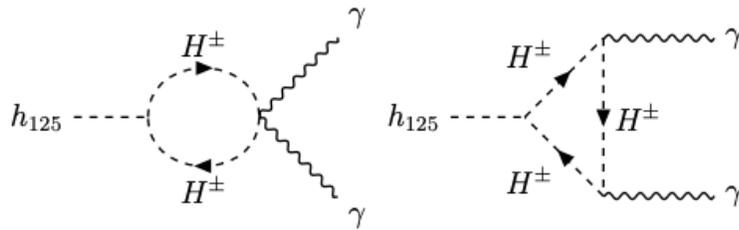


assumes that  $h_1$  is  $h_{125}$

How can we distinguish the two cases?

# Interplay with $H\gamma\gamma$ coupling

- Also loop effects can be important as evident in the di-photon decay channel.
- Charged Higgs yields sizeable contribution:



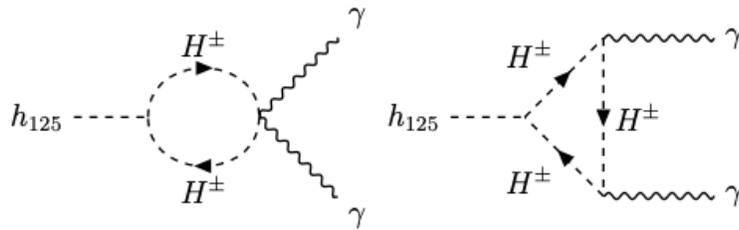
Coupling:  $g_{h_{125}H^+H^-} \rightarrow -\frac{1}{v} [m_{h_{125}}^2 + 2(m_{H^\pm}^2 - \bar{m}^2)]$

Loop suppression:  $v^2/m_{H^\pm}^2$

Higgs potential parameter.  
 Pert. unitarity enforces  $m_{h_2}^2 \sim m_{H^\pm}^2 \sim \bar{m}^2$

# Interplay with $H\gamma\gamma$ coupling

- Also loop effects can be important as evident in the di-photon decay channel.
- Charged Higgs yields sizeable contribution:

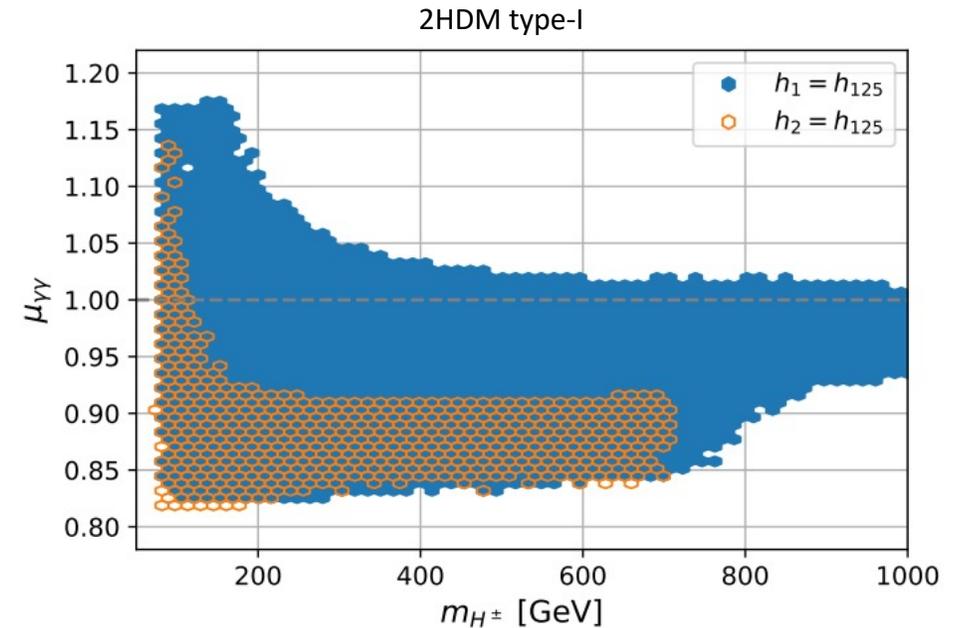


Coupling:  $g_{h_{125}H^+H^-} \rightarrow -\frac{1}{v} [m_{h_{125}}^2 + 2(m_{H^\pm}^2 - \bar{m}^2)]$

Loop suppression:  $v^2/m_{H^\pm}^2$

Higgs potential parameter.  
Pert. unitarity enforces  $m_{h_2}^2 \sim m_{H^\pm}^2 \sim \bar{m}^2$

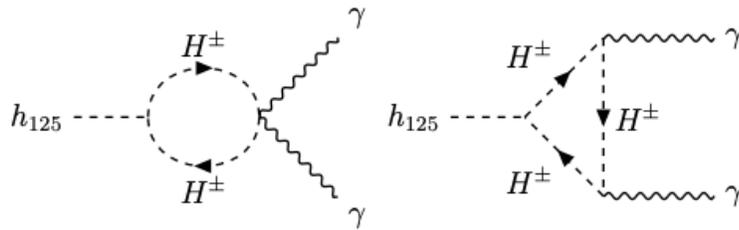
⇒ Lower di-photon signal rate predicted if heavier CP-even Higgs  $H$  is  $h_{125}$



[HB et al. 2103.07484, see also Bernon et al 1511.03682]

# Interplay with $H\gamma\gamma$ coupling

- Also loop effects can be important as evident in the di-photon decay channel.
- Charged Higgs yields sizeable contribution:

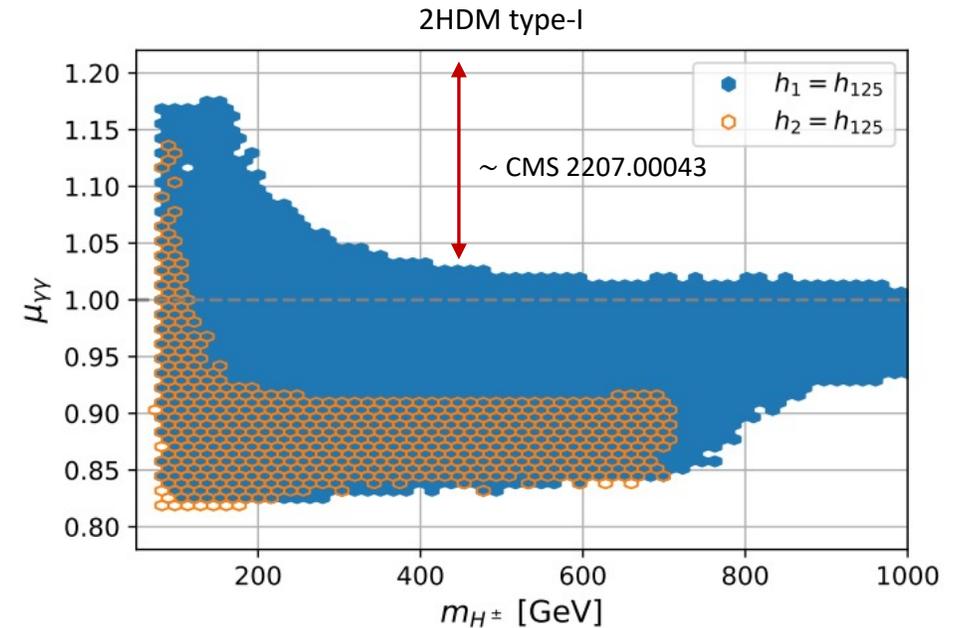


Coupling:  $g_{h_{125}H^+H^-} \rightarrow -\frac{1}{v} [m_{h_{125}}^2 + 2(m_{H^\pm}^2 - \bar{m}^2)]$

Loop suppression:  $v^2/m_{H^\pm}^2$

Higgs potential parameter.  
Pert. unitarity enforces  $m_{h_2}^2 \sim m_{H^\pm}^2 \sim \bar{m}^2$

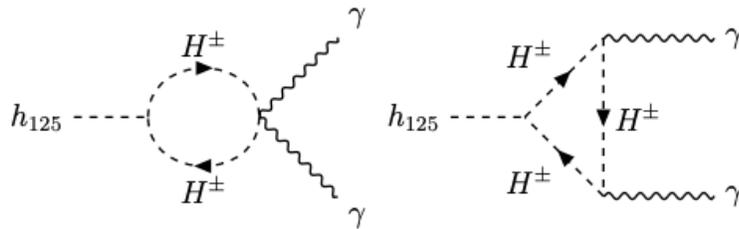
⇒ Lower di-photon signal rate predicted if heavier CP-even Higgs  $H$  is  $h_{125}$



[HB et al. 2103.07484, see also Bernon et al 1511.03682]

# Interplay with $H\gamma\gamma$ coupling

- Also loop effects can be important as evident in the di-photon decay channel.
- Charged Higgs yields sizeable contribution:

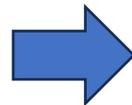


Coupling:  $g_{h_{125}H^+H^-} \rightarrow -\frac{1}{v} [m_{h_{125}}^2 + 2(m_{H^\pm}^2 - \bar{m}^2)]$

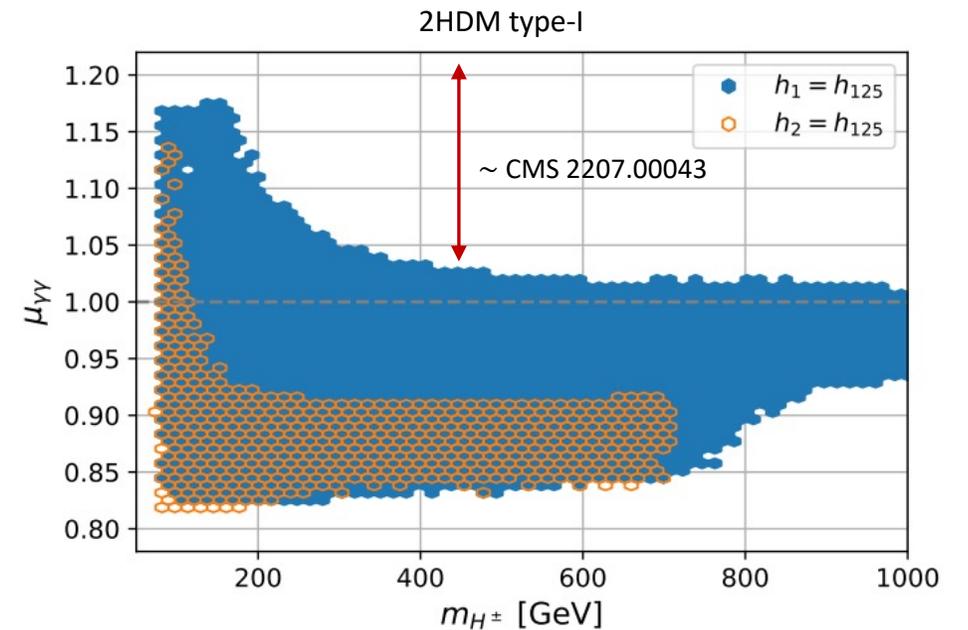
Loop suppression:  $v^2/m_{H^\pm}^2$

Higgs potential parameter.  
Pert. unitarity enforces  $m_{h_2}^2 \sim m_{H^\pm}^2 \sim \bar{m}^2$

⇒ Lower di-photon signal rate predicted if heavier CP-even Higgs  $H$  is  $h_{125}$



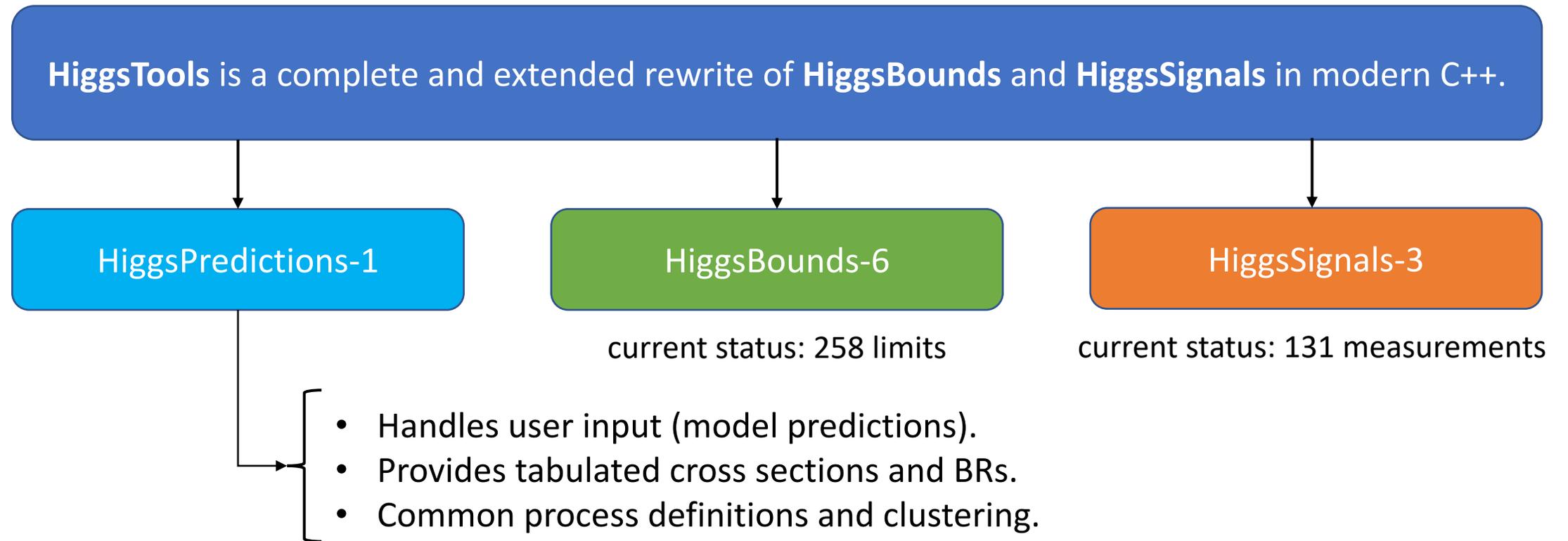
Important interplay between different Higgs couplings!



[HB et al. 2103.07484, see also Bernon et al 1511.03682]

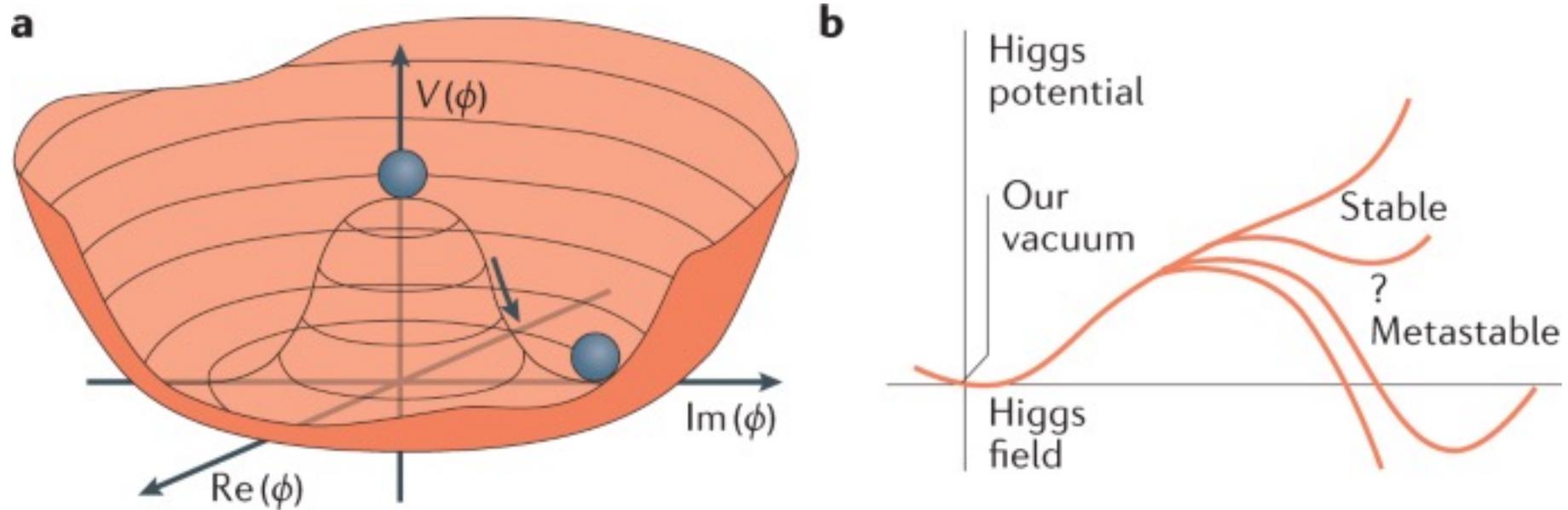
# Interlude: HiggsTools

[HB et al., 2210.09332]



➔ C++ interface for high performance; Python and Mathematica interfaces for ease of use.

# The stability of the Universe



But could we see a BSM effects first in the Higgs potential?

# Baryon asymmetry of the Universe

- Different techniques used in the literature to calculate BAU  $Y_B$ :

- Vev-insertion approach (VIA),

[Huet&Nelson,9504427,9506477;Carena et al., 9603420;Riotto, 9712221;Lee et al.,0412354;Postma et al.,2206.01120]

- WKB (or FH) approximation.

[Joecy et al.,9410282;Kainulainen et al.,0105295, 0202177;Prokopec et al., 0312110, 0406140;Konstandin et al.,1302.6713, 1407.3132]

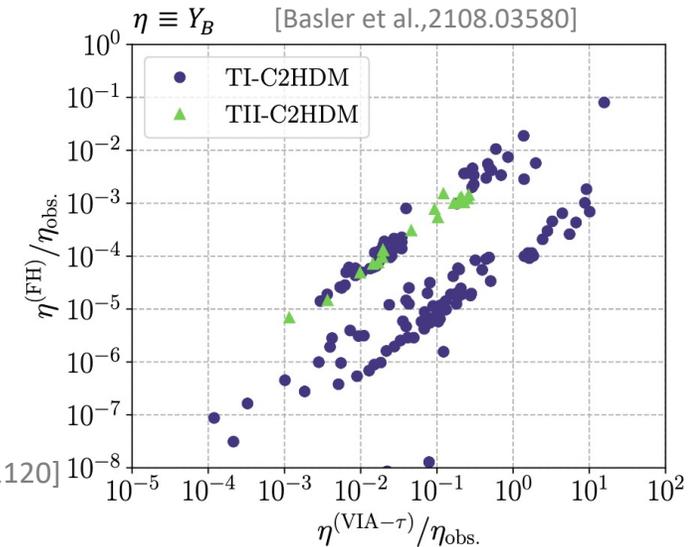
- VIA approach yields consistently higher results by orders of magnitude.
- We use VIA approach with bubble wall parameters close to optimal values for  $Y_B$ :

[de Vries,1811.11104;Fuchs et al.,2003.00099,2007.06940;Shapira,2106.05338]

$$\frac{Y_B}{Y_B^{\text{obs}}} \simeq 28\tilde{c}_t - 0.2\tilde{c}_b - 11\tilde{c}_\tau + \dots$$



$Y_B$  values should be regarded as **upper bound** on what is theoretically achievable.



# Case study: real singlet extension of the SM

$$V(\Phi, S) = V_{\text{SM}}(\Phi) + \frac{1}{2} \mu_S^2 S^2 + \frac{1}{4!} \lambda_S S^4 + \lambda_{S\Phi} S^2 \Phi^\dagger \Phi$$

If  $S$  does not get a vev,  $\lambda_{HHH} = \lambda_{HHH}^{\text{SM}}$  at the tree-level ( $m_S^2 = \mu_S^2 + \lambda_{S\Phi} v^2$ ).

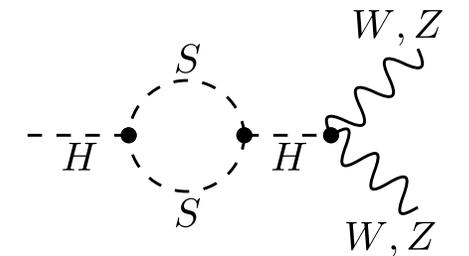
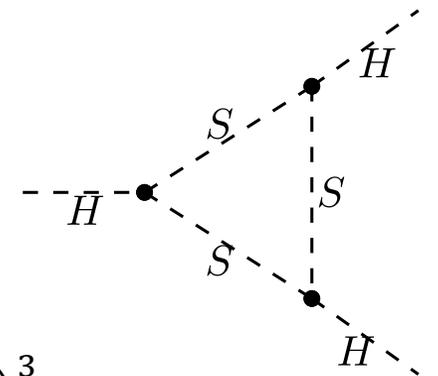
The 1L correction to  $\lambda_{HHH}$  scales like

$$\lambda_{HHH}^{1L} \propto \frac{g_{HSS}^3}{(4\pi)^2} C_0(\dots) \propto \frac{g_{HSS}^3}{(4\pi)^2} \frac{1}{m_S^2} \propto \frac{1}{(4\pi)^2} \frac{m_S^4}{v^3} \left(1 - \frac{\mu_S^2}{m_S^2}\right)^3 \Rightarrow \kappa_\lambda \equiv \frac{\lambda_{HHH}}{\lambda_{HHH}^{\text{SM}}} = 1 + \frac{1}{(4\pi)^2} \frac{m_S^4}{v^4 \lambda_\Phi^{\text{SM}}} \left(1 - \frac{\mu_S^2}{m_S^2}\right)^3$$

whereas the dominant correction to other Higgs couplings scale like

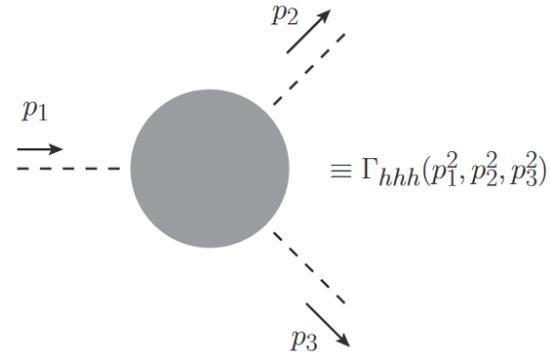
$$g^{1L} \propto \frac{g_{HSS}^2}{(4\pi)^2} B'_0(\dots) \cdot g_{\text{tree}} \propto \frac{1}{(4\pi)^2} \frac{m_S^2}{v^2} \left(1 - \frac{\mu_S^2}{m_S^2}\right)^2 \Rightarrow \kappa_g \equiv \frac{g}{g^{\text{SM}}} = 1 + \frac{1}{(4\pi)^2} \frac{m_S^2}{v^2} \left(1 - \frac{\mu_S^2}{m_S^2}\right)^2$$

➡ Deviation in  $\lambda_{HHH}$  enhanced by a factor  $\frac{m_S^2}{v^2 \lambda_\Phi^{\text{SM}}} \left(1 - \frac{\mu_S^2}{m_S^2}\right)$  w.r.t. to other Higgs couplings!



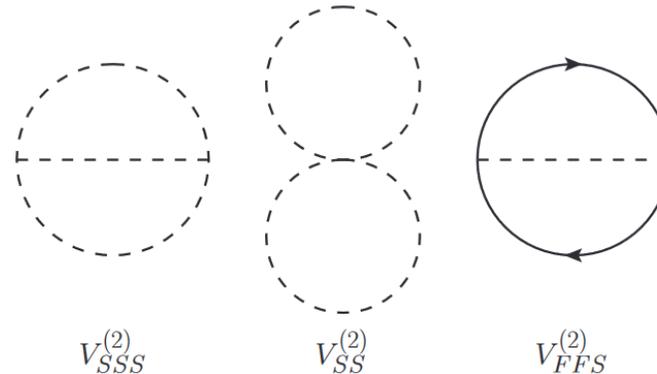
# Calculating BSM corrections to $\kappa_\lambda$

- Need to calculate Higgs three-point function:



- Alternatively, employ zero momentum approximation and then use effective potential:

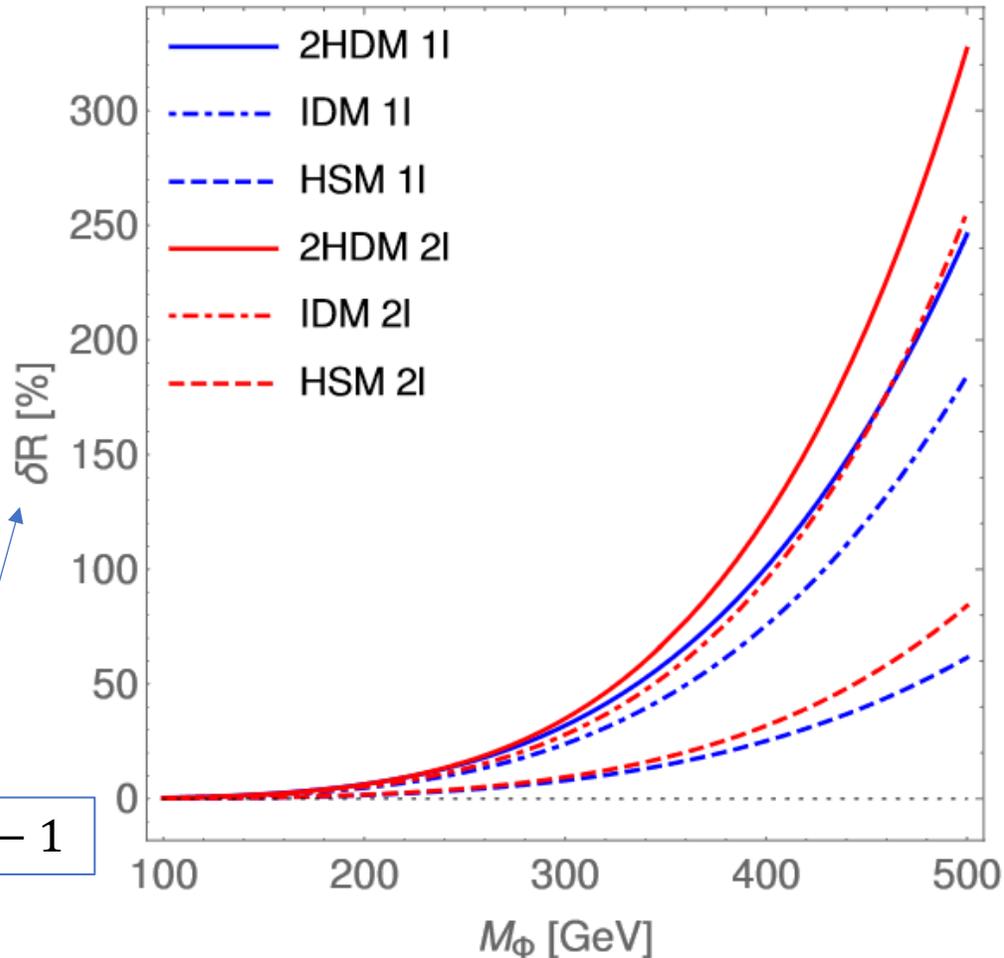
$$\lambda_{hhh} \equiv \left. \frac{\partial^3 V_{\text{eff}}}{\partial h^3} \right|_{\text{min}} \equiv \lambda_{hhh}^{(0)} + \kappa \delta^{(1)} \lambda_{hhh} + \kappa^2 \delta^{(2)} \lambda_{hhh}$$



- Using  $V_{\text{eff}}$ , 1L and 2L corrections have been calculated in various BSM Higgs models (see e.g. [Braathen, Kanemura, 1911.11507]).

# Calculating BSM corrections to $\kappa_\lambda$

[Braathen,Kanemura,1911.11507]



- Large non-decoupling corrections found in several BSM models.
- Analysis assumed that all BSM masses are equal  $M_\Phi$ .
- No phenomenological analysis has been performed.



**Idea of this work:**

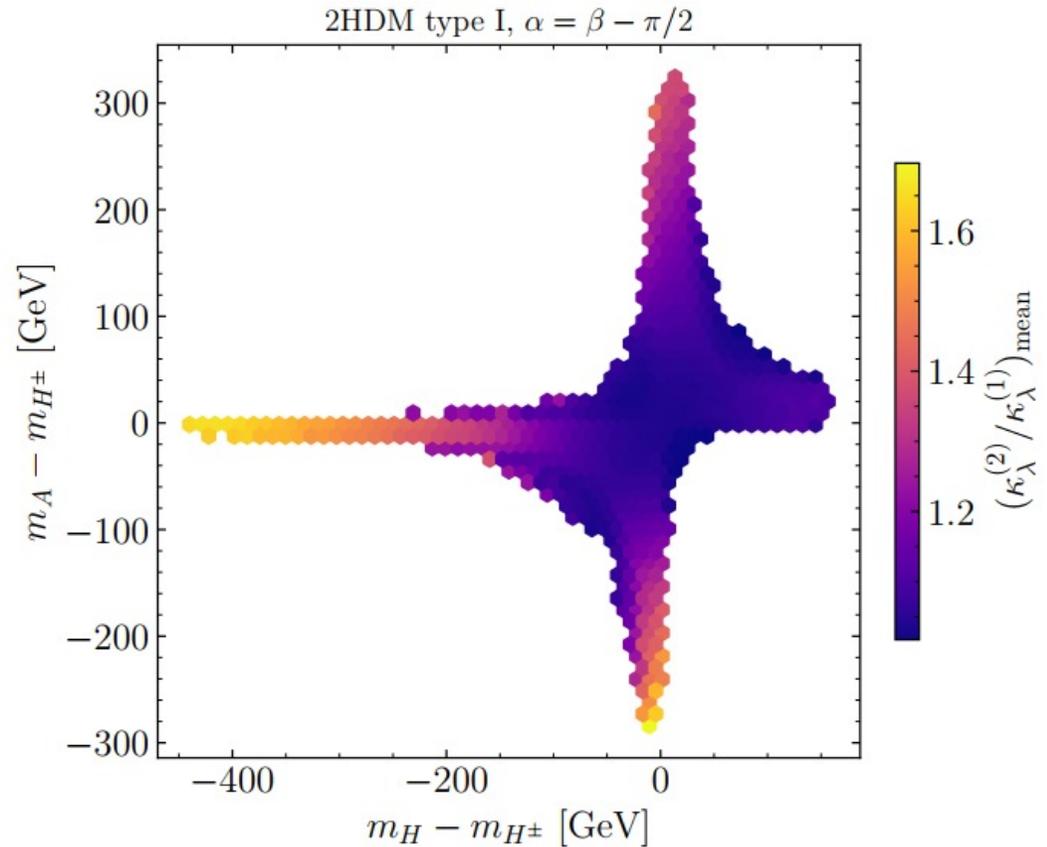
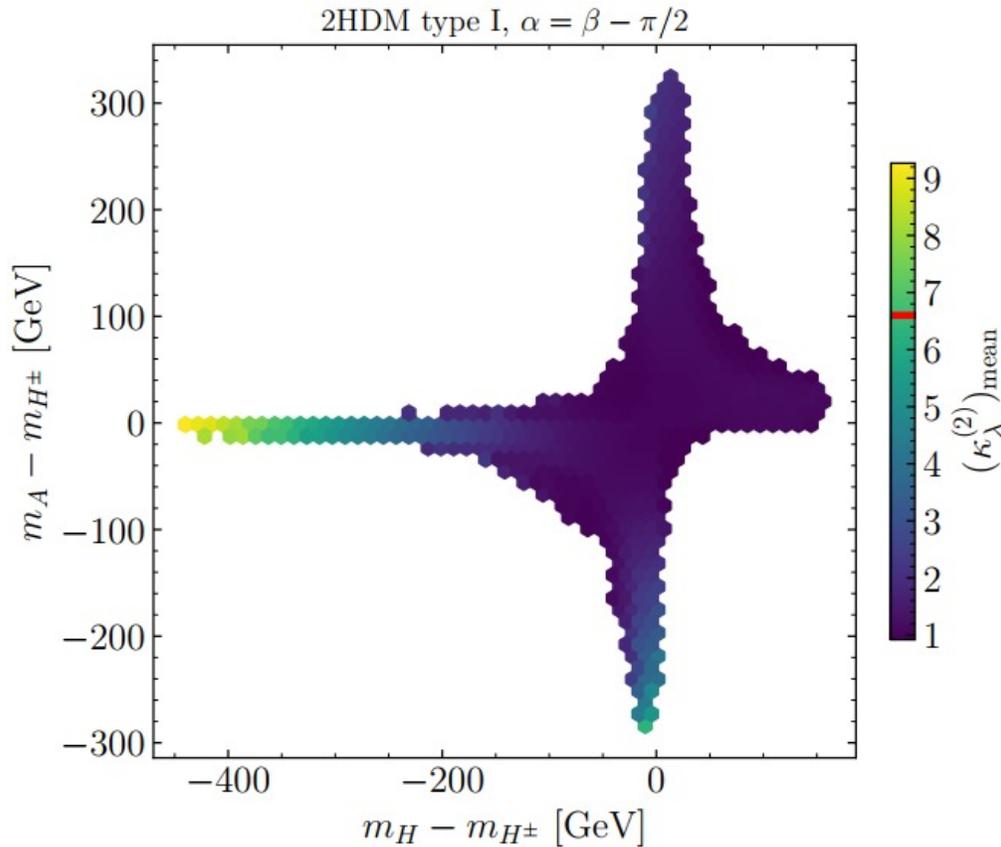
Can we constrain these models based on the large corrections to  $\kappa_\lambda$ ?

# 2HDM parameter scan

- We checked for
  - vacuum stability and boundedness-from-below,
  - NLO perturbative unitarity, [Grinstein et al., 1512.04567; Cacchio et al., 1609.01290]
  - electroweak precision observables (calculated at the 2L level using THDM\_EWPOS), [Hessenberger & Hollik, 1607.04610, 2207.03845]
  - SM-like Higgs measurements via HiggsSignals, [Bechtle et al., 2012.09197]
  - direct searches for BSM scalars via HiggsBounds, [Bechtle et al., 2006.06007]
  - b-physics constraints.
- Most constraints checked using ScannerS. [Mühlleitner et al., 2007.02985]
- For each point passing the constraints, we calculate  $\kappa_\lambda$  at the 1L and 2L level ( $\kappa_\lambda^{(1)}$  and  $\kappa_\lambda^{(2)}$ ). [Braathen, Kanemura, 1911.11507]

# 2HDM parameter scan — results

(showing only points passing all constraints mentioned on previous slide)

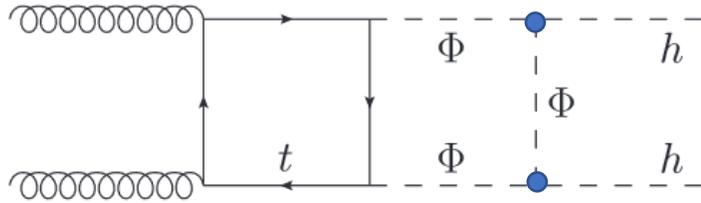


- Largest corrections for  $m_A \simeq m_{H^\pm}$ ,  $m_H < m_{H^\pm}$  and  $m_H \simeq m_{H^\pm}$ ,  $m_A < m_{H^\pm}$  ( $\kappa_\lambda$  of up to 9).
- 2L corrections have sizeable impact (up to 70%).

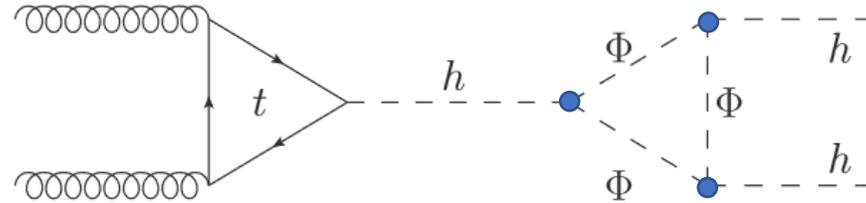
# Can we apply the experimental constraints on $\kappa_\lambda$ ?

Assumptions of experimental bound:

- All other Higgs couplings are SM-like.
  - 2HDM in the alignment limit with heavy BSM masses. ✓
- Higgs-boson pair production only deviates from the SM via a modified trilinear Higgs coupling.
  - No resonant contribution because  $Hhh$  coupling is zero in alignment limit. ✓
  - Other BSM contributions to  $hh$  production?



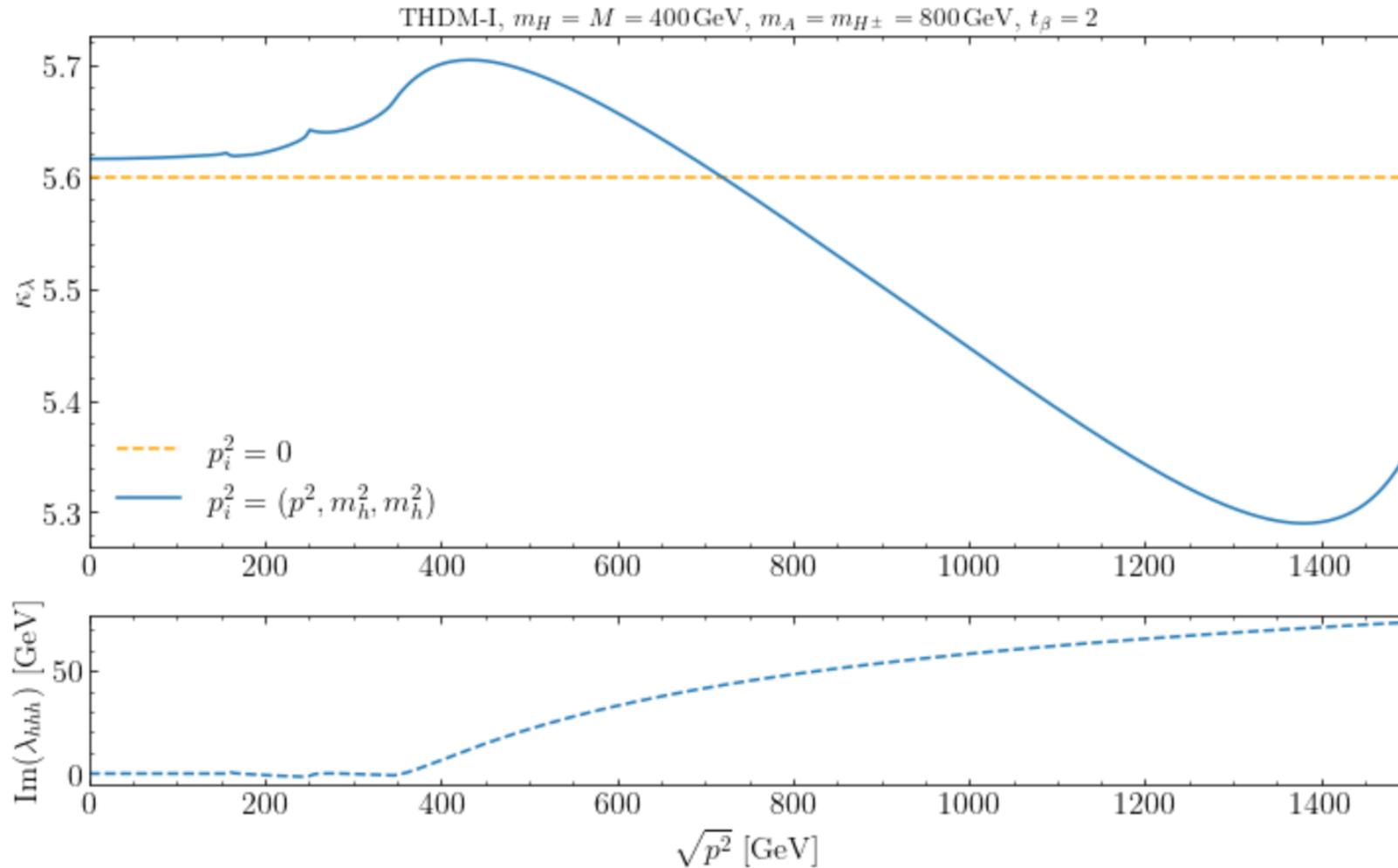
$$\propto \mathcal{O}(y_t^2 g_{hh\Phi\Phi}^2) \text{ (not included)}$$



$$\propto \mathcal{O}(y_t g_{hh\Phi\Phi}^3) \text{ (included)}$$

- We include the all corrections leading in the large coupling  $g_{hh\Phi\Phi}$  at the NLO and NNLO level. ✓

# Momentum dependence



# The Higgs mass as a precision observable

- Also the Higgs mass is a precision observable useful for BSM phenomenology.
- In SUSY models, the Higgs mass can be predicted in terms of the model parameters.
- MSSM:  $M_h \sim 125 \text{ GeV} \Rightarrow$  stop masses  $\gtrsim 2 \text{ TeV}$ .
- Experimental precision significantly better than remaining theoretical uncertainty.  
( $\sim 0.5 \text{ GeV}$  for  $X_t/M_S = 0$  and  $\sim 1 \text{ GeV}$  for  $X_t/M_S = \sqrt{6}$ )

