

Precise Higgs mass predictions for multi-scale hierarchies with FeynHiggs

Henning Bahl

DESY, Hamburg



SUSY conference

24.8.2021

Introduction

Multi-scale hierarchies: low-mass BSM Higgs bosons

Multi-scale hierarchies: high-mass gluino

Conclusions

The SM-like Higgs mass as a precision observable

Special feature of the MSSM

Mass of lightest \mathcal{CP} -even Higgs, M_h , is calculable in terms of model parameters \Rightarrow can be used as a precision observable

- ▶ at tree-level $M_h^2 \simeq M_Z^2 \cos^2(2\beta) \leq M_Z^2$,
- ▶ M_h is, however, heavily affected by loop corrections,
- ▶ directly sensitive to the SUSY scale.

Experimentally measured mass: [Aad et al., 1503.07589]

$$M_h^{\text{exp}} = 125.08 \pm 0.21 \text{ (stat.)} \pm 0.11 \text{ (sys.) GeV}$$

To fully profit from experimental precision, higher order calculations are crucial!

FeynHiggs

Main purpose: precise calculation of Higgs mass spectrum in the MSSM.

[Authors: HB, Hahn, Heinemeyer, Hollik, Paßehr, Rzehak, Weiglein]

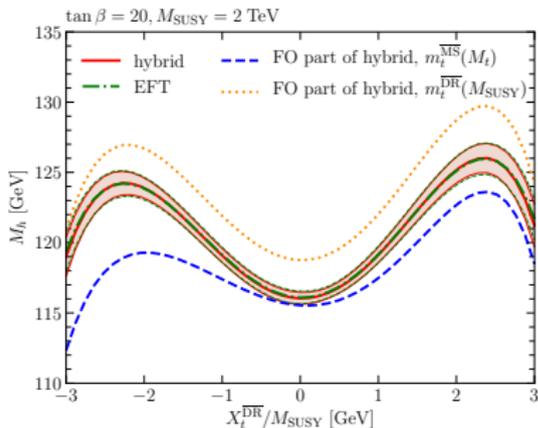
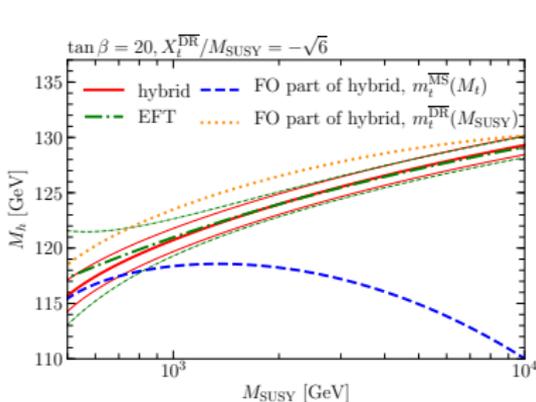
Higgs mass calculation II

Three approaches are used:

- ▶ Fixed-order (FO) approach:
 - + Precise for low SUSY scales,
 - but for high scales $\ln(M_{\text{SUSY}}^2/M_t^2)$ terms spoil convergence of perturbative expansion.
- ▶ effective field theory (EFT) approach:
 - + Precise for high SUSY scales (logs resummed),
 - but for low scales $\mathcal{O}(M_t/M_{\text{SUSY}})$ terms are missed if higher-dimensional operators are not included.
- ▶ hybrid approach combining FO and EFT approaches:
 - ++ Precise for low and high SUSY scales.

Current status for single-scale scenarios [HB, Heinemeyer, Hollik, Weiglein, 1912.04199]

Single-scale scenario with all non-SM particles at M_{SUSY} (SM as EFT)



“Rule of thumb”

Remaining theoretical uncertainties (for $\overline{\text{DR}}$ stop input parameter):

$$X_t/M_{\text{SUSY}} = 0 \rightarrow \Delta M_h \sim 0.5 \text{ GeV},$$

$$X_t/M_{\text{SUSY}} = \sqrt{6} \rightarrow \Delta M_h \sim 1 \text{ GeV}$$

Slightly higher for OS stop input parameters.

Introduction

Multi-scale hierarchies: low-mass BSM Higgs bosons

Multi-scale hierarchies: high-mass gluino

Conclusions

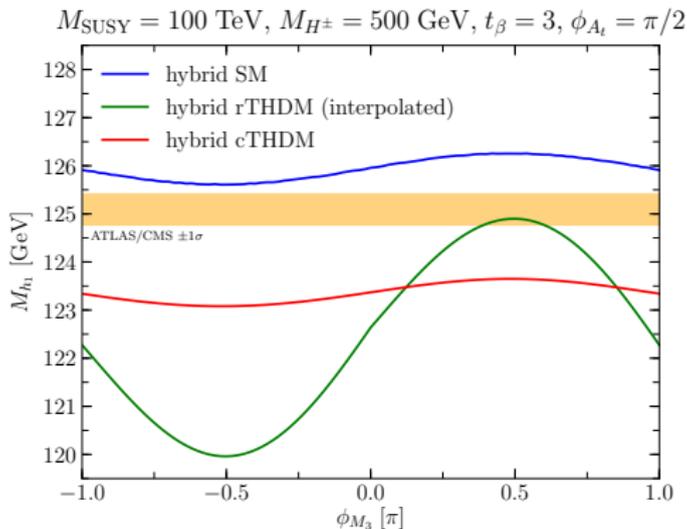
THDM as EFT

- ▶ For low M_A , the EFT of the MSSM is not the THDM type-II,
→ both Higgs doublets couple to e.g. top quarks,
 - ▶ loop corrections induce non-zero (potentially complex) values for $\lambda_{5,6,7}$
- ⇒ Large number of EFT parameters complicating the calculation.

Recent progress:

- ▶ complex THDM as EFT [HB,Murphy,Rzehak,1909.00726,2010.04711],
- ▶ calculation of $\mathcal{O}(\alpha_t^2)$ threshold corrections [HB,Sobolev, 2010.01989].

Complex THDM as EFT



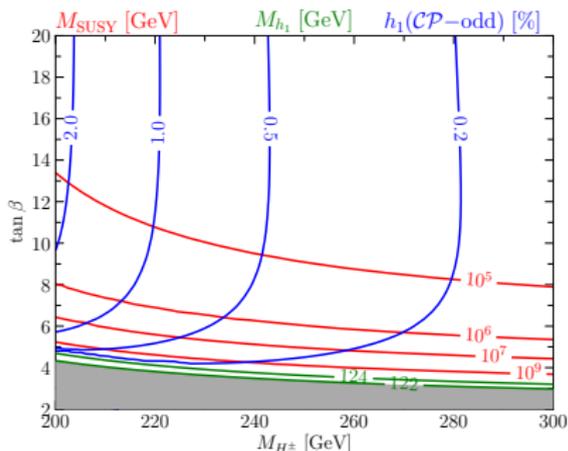
Including phase dependence fully in

- ▶ 2L RGEs,
- ▶ one-loop threshold corrections,
- ▶ $\mathcal{O}(\alpha_t \alpha_s)$ λ_i -threshold corrections.

Application: the \mathcal{CP} -odd component of the SM-like Higgs boson

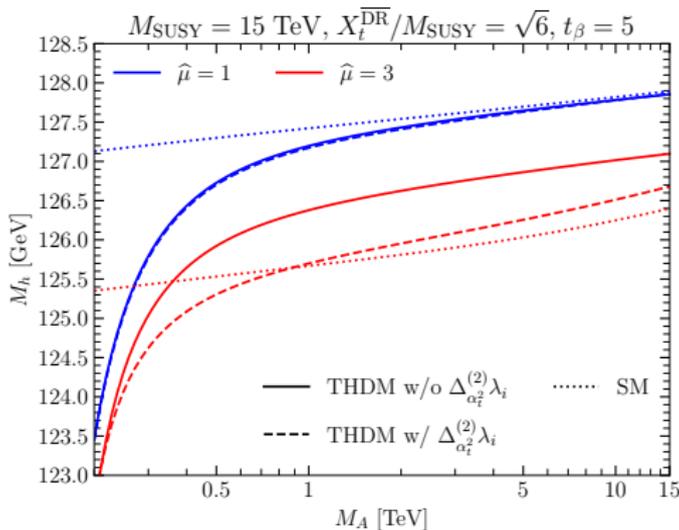
Sizeable \mathcal{CP} -odd component requires

- ▶ large mixing with \mathcal{CP} -odd A boson
 - imaginary parts of couplings have to be large ($\phi_{A_t} = 2\pi/3, \phi_{M_3} = \pi/4$)
 - $\tan \beta$ and M_{H^\pm} must be small
- ▶ large SUSY scale required to ensure $M_h \sim 125$ GeV
→ \mathcal{CP} -mixing decouples



Potential discovery of \mathcal{CP} -odd component at the LHC has potential to exclude the MSSM.

$\mathcal{O}(\alpha_t^2)$ threshold corrections to λ_i



- ▶ compared different calculation methods,
- ▶ easiest method: calculate 2L 4-point functions in the unbroken phase,
- ▶ calculation fully includes CP-violating phases.

⇒ similar precision level as for single-scale hierarchy: theoretical uncertainty for low M_A should be under control!

Introduction

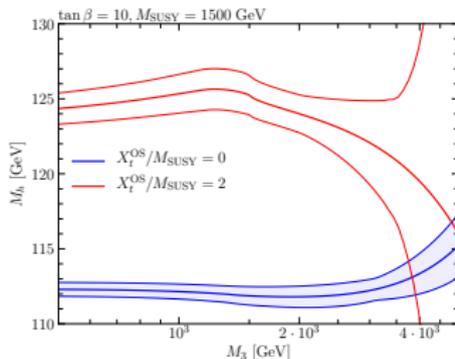
Multi-scale hierarchies: low-mass BSM Higgs bosons

Multi-scale hierarchies: high-mass gluino

Conclusions

The heavy gluino limit: $M_{\tilde{g}} \gg M_{\tilde{t}}$

Increasingly relevant due to tightening LHC gluino limits.



Large uncertainty due to M_3 power-enhanced terms appearing at the two-loop level in $\overline{\text{DR}}$ EFT calculation (do not appear in OS scheme).

Needed EFT: MSSM without gluino

Has not been worked out yet...

(available in FeynHiggs since version 2.18.0)

Solution: Absorb power-enhanced terms into renormalization scheme

[HB,Sobolev,Weiglein,1912.10002]

Use $\overline{\text{MDR}}$ instead of $\overline{\text{DR}}$ in EFT ($\overline{\text{DR}}$ ill-defined for $Q < |M_3|$),

$$\left(m_{\tilde{t}_{L,R}}^{\overline{\text{MDR}}}\right)^2 = \left(m_{\tilde{t}_{L,R}}^{\overline{\text{DR}}}\right)^2 \left[1 + \frac{\alpha_s}{\pi} C_F \frac{|M_3|^2}{m_{\tilde{t}_{L,R}}^2} \left(1 + \ln \frac{Q^2}{|M_3|^2} \right) \right],$$

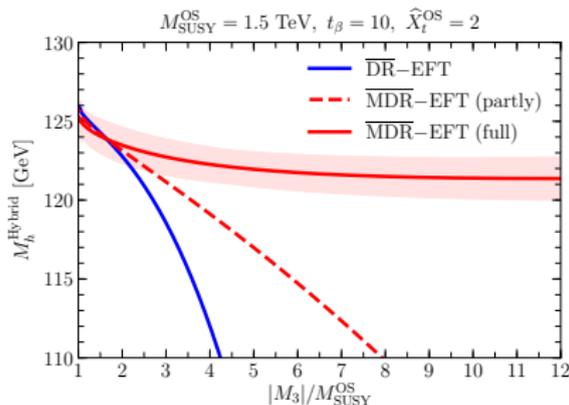
$$X_t^{\overline{\text{MDR}}}(Q) = X_t^{\overline{\text{DR}}}(Q) - \frac{\alpha_s}{\pi} C_F M_3 \left(1 + \ln \frac{Q^2}{|M_3|^2} \right),$$

resums all $\mathcal{O}(\alpha_s^n M_3^{2n}, \alpha_s^n M_3^n)$ terms.

⇓

Drastically reduced uncertainty.

(coming to FeynHiggs soon)



Introduction

Multi-scale hierarchies: low-mass BSM Higgs bosons

Multi-scale hierarchies: high-mass gluino

Conclusions

Conclusions

- ▶ The SM-like Higgs mass is a unique observable in the MSSM directly sensitive to the SUSY scale,
- ▶ most precise prediction available for single-scale scenario (SM as EFT).

What about multi-scale hierarchies?

- ▶ Low-mass BSM Higgs bosons (THDM as EFT):
 - have now reached similar precision level as for single-scale hierarchy,
 - \mathcal{CP} violating phases are fully taken into account.
- ▶ High-mass gluino:
 - large uncertainty if gluino is heavier than stops in the $\overline{\text{DR}}$ scheme,
 - large corrections can be absorbed by modifying the renormalization scheme.

Conclusions

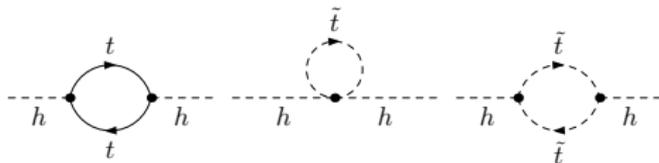
- ▶ The SM-like Higgs mass is a unique observable in the MSSM directly sensitive to the SUSY scale,
- ▶ most precise prediction available for single-scale scenario (SM as EFT).

What about multi-scale hierarchies?

- ▶ Low-mass BSM Higgs bosons (THDM as EFT):
 - have now reached similar precision level as for single-scale hierarchy,
 - \mathcal{CP} violating phases are fully taken into account.
- ▶ High-mass gluino:
 - large uncertainty if gluino is heavier than stops in the $\overline{\text{DR}}$ scheme,
 - large corrections can be absorbed by modifying the renormalization scheme.

Thanks for your attention!

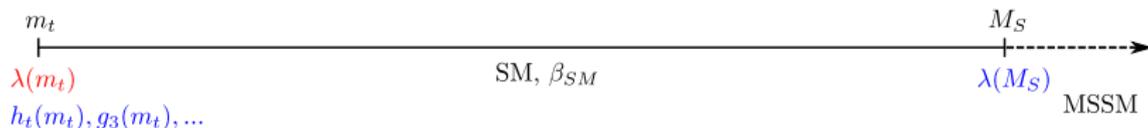
Fixed-order techniques



$$M_h^2 = m_h^2 + \frac{6y_t^4}{(4\pi)^2} v^2 \left[\ln \frac{M_{\tilde{t}}^2}{M_t^2} + \left(\frac{X_t}{M_{\tilde{t}}} \right)^2 - \frac{1}{12} \left(\frac{X_t}{M_{\tilde{t}}} \right)^4 \right] + \dots$$

- ▶ Stop mass scale $M_{\tilde{t}} = \sqrt{M_{\tilde{t}_1} M_{\tilde{t}_2}}$,
- ▶ large logarithms spoil perturbative convergence if $M_{\tilde{t}} \gg M_t$,
- ▶ status in FeynHiggs: $\mathcal{O}(\text{full 1L}, \alpha_s(\alpha_b + \alpha_t), (\alpha_b + \alpha_t)^2)$.

EFT calculation (simplest hierarchy)



- ▶ Integrate out all SUSY particles \rightarrow SM as EFT,
- ▶ Higgs self-coupling fixed at matching scale

$$\lambda(M_{\text{SUSY}}) = \frac{1}{4}(g^2 + g_y^2) + \frac{6y_t^4}{(4\pi)^2} \left[\left(\frac{X_t}{M_{\text{SUSY}}} \right)^2 - \frac{1}{12} \left(\frac{X_t}{M_{\text{SUSY}}} \right)^4 \right] + \dots,$$

- ▶ run Higgs self-coupling down to electroweak scale,
- ▶ calculate Higgs mass: $M_h^2 = \lambda(M_t)v^2 + \dots$,
- ▶ status in FeynHiggs: full LL + NLL resummation, NNLL resummation in gaugeless limit, partial N³LL resummation; similar precision for multi-scale hierarchies.

How to deal with intermediary SUSY scales?

For sparticles in the LHC range, both logs and suppressed terms might be relevant. We could try to improve

- ▶ fixed-order calculation → need to calculate more three- and two-loop corrections,
- ▶ EFT calculation → need to include higher-dimensional operators into calculation.

or ...



Hybrid approach

Combine both approaches to get precise results for both regimes!

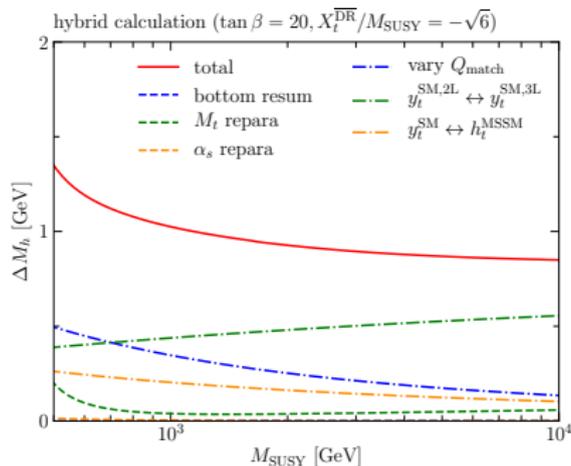
Procedure in FeynHiggs

1. Calculation of diagrammatic fixed-order self-energies $\hat{\Sigma}_{hh}$
2. Calculation of EFT prediction $\lambda(M_t)v^2$
3. Add non-logarithmic terms contained in fixed-order result and the logarithms contained in EFT result

$$\hat{\Sigma}_{hh}(m_h^2) \longrightarrow [\hat{\Sigma}_{hh}(m_h^2)]_{\text{nolog}} - [v^2\lambda(M_t)]_{\text{log}}$$

In practice, this is achieved by using subtraction terms.

Remaining uncertainties – individual sources



Uncertainty estimate dominated by:

- ▶ Uncertainty from higher order threshold corrections:
 - vary matching scale between SM and MSSM,
 - reexpress threshold correction in terms of h_t^{MSSM} instead of y_t^{SM} .
- ▶ Uncertainty of SM input couplings:
 - $y_t(M_t)$ extracted at the 2- or 3-loop level out of OS top mass.

→ FeynHiggs provides point-by-point uncertainty estimate.