

Resummation of large-logarithms in MSSM Higgs-boson mass calculations

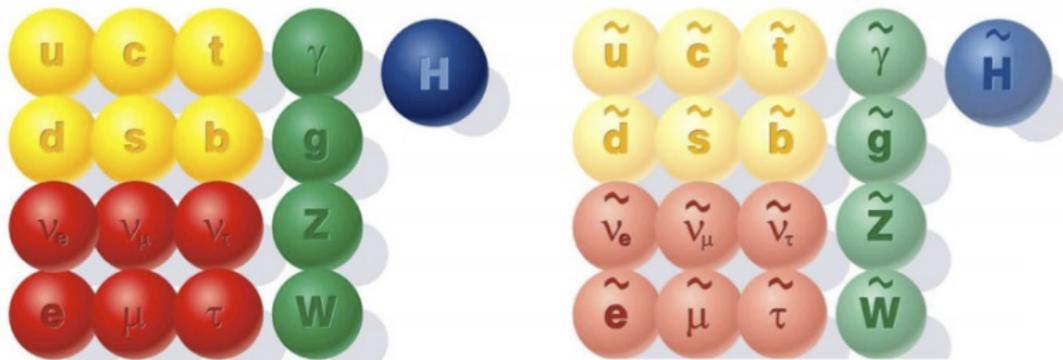
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- 1 MSSM (and the MSSM Higgs sector)
- 2 The lightest Higgs boson mass
- 3 Resummation of logarithms
- 4 Conclusion

Minimal Supersymmetric Standard Model:

- ▶ one of the most common models of BSM physics
- ▶ Supersymmetry relates bosons to fermions
- ▶ each SM particle gets a superpartner (e.g. stops \leftrightarrow top)
- ▶ able to address:
hierarchy problem, DM, gauge coupling unification,...

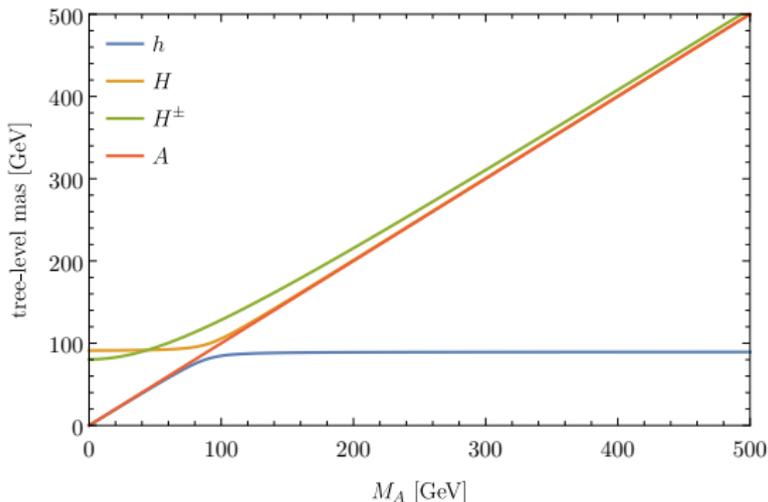


- ▶ Holomorphicity of superpotential and anomaly cancellation
⇒ need to introduce second Higgs doublet
- ▶ corresponds to type II THDM (remember Stephan's talk)
- ▶ five physical Higgses: h, H (CP-even); A (CP-odd); H^\pm
- ▶ SUSY restricts masses, couplings
- ▶ tree-level Higgs sector determined by 2 parameters:
 $M_A, \tan \beta = v_2/v_1$
(7 free parameters in general THDM models)

Distinct feature of the MSSM

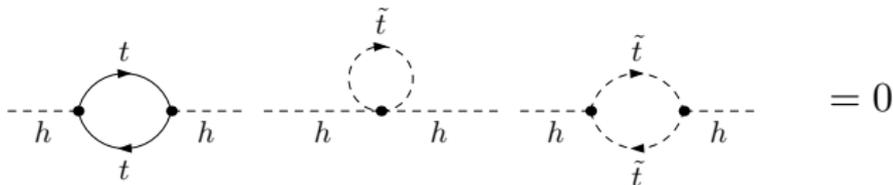
Mass of lightest Higgs boson is calculable in terms of model parameters.

- ▶ measured M_h can be used as a precision observable to constrain parameter space of the model (\rightarrow Stephan's talk)
- ▶ at tree-level: $M_h^2 \leq M_Z^2 \cos(2\beta)^2 \leq M_A^2$



What's about quantum correction?

- ▶ exact SUSY \rightarrow all quantum corrections cancel



Reasons?

- ▶ exact SUSY forces sparticles to have the same mass as the corresponding SM particle ($m_t = m_{\tilde{t}_1} = m_{\tilde{t}_2}$)
- ▶ SUSY forces sparticles to have the same couplings as the corresponding SM particle
- ▶ fermion (scalar) loop has opposite sign as scalar (fermion) loop

But M_h was measured to be ~ 125 GeV ($> M_Z \sim 90$ GeV)!

And no SUSY particles have been found yet!

\Rightarrow **SUSY must be broken**

- ▶ We don't know breaking mechanism
- ▶ Parametrize our ignorance by adding additional mass terms for sparticles to Lagrangian (soft-breaking terms)

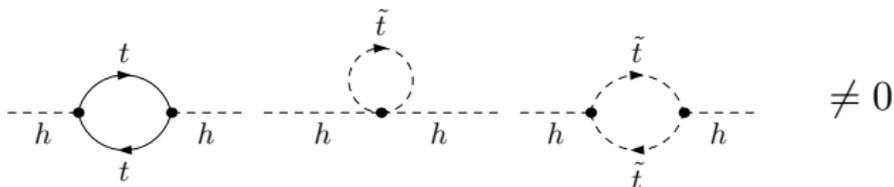
$$m_{\tilde{t}_1}^2 = m_t^2 + \Delta m_1^2$$

$$m_{\tilde{t}_2}^2 = m_t^2 + \Delta m_2^2$$

- ▶ We don't change couplings (\rightarrow no quadratic divergences)

SUSY broken

Loop corrections contribute to Higgs mass



but

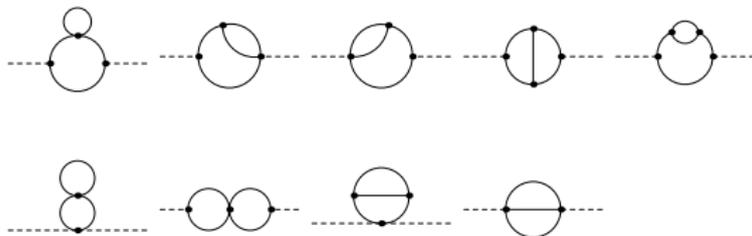
$$M_h^2 = M_Z^2 \cos(2\beta)^2 + \frac{3}{4\pi^2} m_t^2 h_t^2 \ln \left(\frac{M_S^2}{m_t^2} \right) + (\text{non logarithmic terms})$$

with $M_S^2 = m_{\tilde{t}_1} m_{\tilde{t}_2}$.

- ▶ 1-loop correction can raise tree-level mass by more than $\sim 80\%$

\Rightarrow **Higher order corrections necessary to obtain accurate result!**

→ 2-loop calculation (already quite complicated, but doable):



- ▶ calculation yields terms $\propto \ln^2 \left(\frac{M_S^2}{m_t^2} \right)$
- ▶ for $M_S \gg m_t$ (heavy stops ← LHC stop searches):

logarithms get large spoiling convergence of perturbative expansion

→ 3-loop calculation ???

Not feasible, but large logarithmic terms ($\ln^3(\dots)$) expected!!



**Alternative method needed to control large
logarithms!**

Origin of the problem

Huge hierarchy between scales m_t and M_S



We need to separate the effects of high- and low-energy physics!



Effective Field Theory

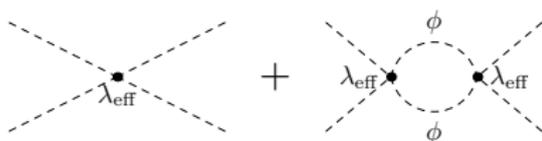
$$\mathcal{L}_{\text{toy}} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi + \frac{1}{2} \partial_\mu \Phi \partial^\mu \Phi - m^2 \phi^2 - M^2 \Phi^2 - V(\phi, \Phi)$$
$$V(\phi, \Phi) = \lambda_1 \phi^4 + \lambda_2 \phi^2 \Phi^2 + \lambda_3 \Phi^4$$

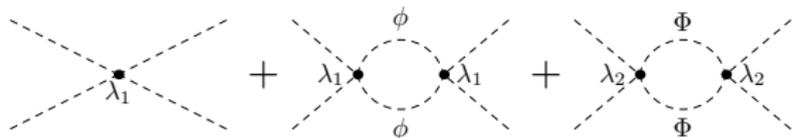
Effective Lagrangian for energies $Q \sim m \ll M$:

$$\mathcal{L}_{\text{eff}} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - m^2 \phi^2 - \frac{\lambda_{\text{eff}}}{4!} \phi^4$$

→ 'high-energy physics' encoded in λ_{eff}

Determine λ_{eff} by matching with full theory at $Q \sim M$ via $\phi\phi \rightarrow \phi\phi$ scattering:

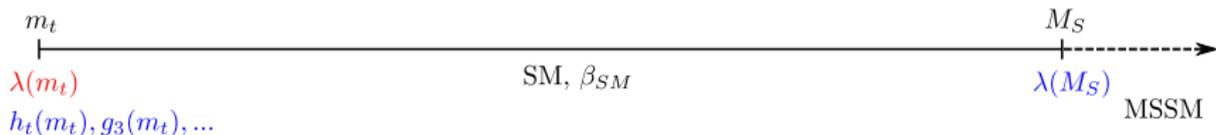
EFT result = 

full theory result = 

EFT result $\stackrel{!}{=}$ full theory result

Regard SM as EFT

$$M_h^2 = 2\lambda(Q = m_t)v^2 \quad (\lambda: \text{SM Higgs self-coupling, } v: \text{SM vev})$$



- ▶ Match SM to MSSM at $Q = M_S$: $\lambda(M_S)$ fixed in MSSM

Calculate Higgs mass in SM:

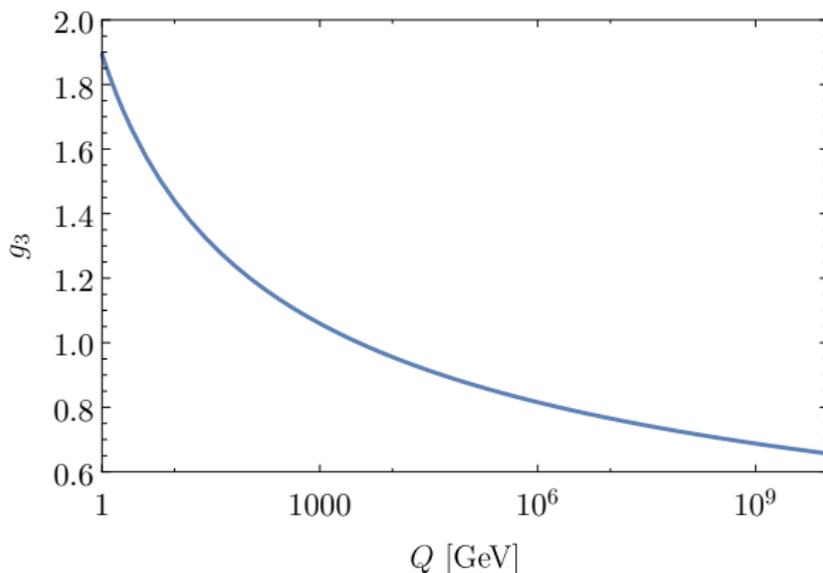
- ▶ In EFT (SM) all SUSY-particles are integrated out
 → no large logarithms
- ▶ All effects of SUSY-particles are absorbed into $\lambda(m_t)$

How to get $\lambda(m_t)$? We only have $\lambda(M_S)$.

- ▶ Higgs self-coupling is a running parameter
→ changes value depending on the scale (energy) of process
- ▶ running is governed by renormalization group equations (RGEs)

- ▶ typical example: running gauge coupling of QCD:

$$\frac{dg_3}{d \ln Q^2} = -\frac{7}{2}k g_3^3 \text{ with } (k = 1/(4\pi)^2)$$



$$\frac{d\lambda}{d \ln Q^2} = 6k \left(\lambda^2 + \lambda h_t^2 - h_t^4 \right)$$

Solve iteratively:

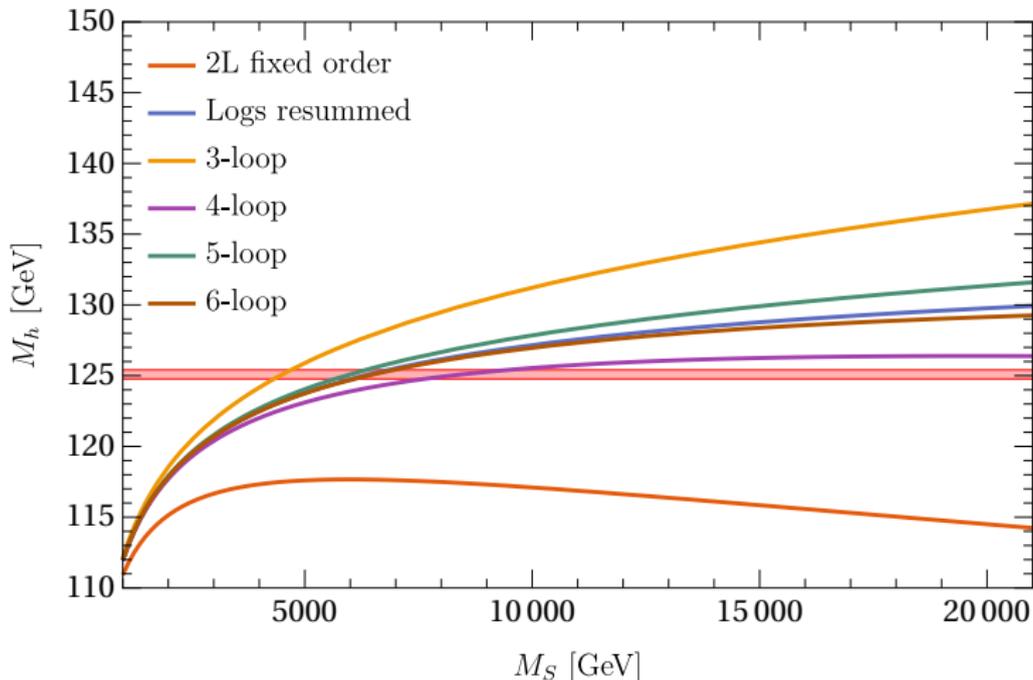
$$\begin{aligned} \lambda(m_t) &\approx \lambda(M_S) + \int_{Q=M_S}^{Q=m_t} \frac{d\lambda}{d \ln Q^2} d \ln Q^2 \approx \\ &\approx \lambda(M_S) - 6k \left(\lambda^2(M_S) + \lambda(M_S)h_t^2(m_t) - h_t^4(m_t) \right) \ln \left(\frac{M_S^2}{m_t^2} \right) \approx \\ &\approx 6kh_t^4(m_t) \ln \left(\frac{M_S^2}{m_t^2} \right) \end{aligned}$$

Multiply by $2v^2 \rightarrow$ 1-loop large logarithm reproduced ($h_t v = m_t$)

\Rightarrow large logarithms originate from RGE running

Solve system of RGEs ($dg_i/d\ln Q^2 = \dots$) numerically:

\Rightarrow Resummation of large logarithms up to all orders



- ▶ In the MSSM lightest, the Higgs mass M_h is calculable in terms of model parameters
- ▶ M_h can be used as a precision observable to constrain parameter space
- ▶ Calculation suffers from large logarithmic contributions spoiling the convergence of the perturbative expansion
- ▶ Effective field theory techniques provide tools to resum these contributions up to all orders