

Matching Coefficients in NRQCD and HQET

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in collaboration with

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Outline

- 1 Framework
 - Introduction
 - Calculation
- 2 Results
 - NRQCD
 - HQET
- 3 Summary

Introduction

- matching coefficients: Connection between full theory (QCD) and an effective theory.

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NRQCD

- describes a non-relativistic quark-antiquark pair
- top production at threshold at a future linear collider $\Rightarrow$ determination of top quark mass $\rightarrow$ [talk by K. Schuller](#)
- decay of heavy mesons

Introduction

- matching coefficients: Connection between full theory (QCD) and an effective theory.

NRQCD

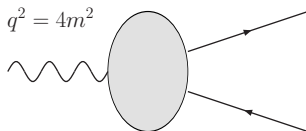
- describes a non-relativistic quark-antiquark pair
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HQET

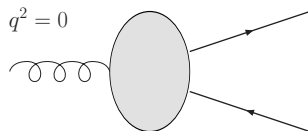
- describes a single heavy quark with small residual momentum
- breaking of the spin symmetry $\Rightarrow B - B^*$ mass splitting cmp talk by R. Sommer

Calculation

NRQCD



HQET



- vector current

$$j_v^k = c_v(\mu) \tilde{j}^k + \frac{d_v(\mu)}{6m_Q^2} \phi^\dagger \sigma^k \vec{D}^2 \chi + \dots$$

- consider diagrams with at least one closed quark loop

- chromomagnetic operator

$$O_{cm} = \frac{1}{2} \bar{Q}_v G_{\mu\nu} \sigma^{\mu\nu} Q_v$$

- reduces to propagator type diagrams

What's known?

NRQCD

- one-loop MC
- two-loop MC

Källen et al '55

Czarnecki et al '97; Beneke et al '97

Kniehl et al '06

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HQET

- one-loop MC
- one-loop AD

Eichten, Hill '90

Falk, Grinstein, Luke '91

- two-loop MC
- two-loop AD

Czarnecki, Grozin '97

Amorós et al '97; Czarnecki et al '97

- large- β_0 limit

Grozin, Neubert '97

- m_c effects at two-loops

Davydychev, Grozin '99

General Approach

- diagrams generated with qgraf Nogueira '91
- topologies identified and mapped with the help of q2e and exp Harlander, Seidensticker, Steinhauser '96/'97
- calculation done with form Vermaseren '00
- reduction to master integrals using IBP identities with crusher PM, Seidel
- master integrals partially known Melnikov, v. Rittbergen '00; Laporta, Remiddi '96
- master integrals calculated/checked using Mellin-Barnes representations (MB.m) Czakon '05

Details: NRQCD

Consider renormalized quark-antiquark-photon vertex at threshold

$$Z_2 \Gamma_V = c_V \tilde{Z}_2 \tilde{Z}_V^{-1} \tilde{\Gamma}_V + \dots$$

we want to extract c_V and $\tilde{Z}_V \rightarrow$ we need:

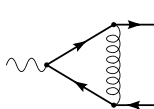
- wave function renormalization constant Z_2

Broadhurst et al '91; Melnikov, v. Rittbergen '00

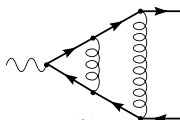
$\hookrightarrow$ we recalculated the necessary parts of Z_2

- $\tilde{\Gamma}_V$ only needed at tree level
- $\tilde{Z}_2 = 1$

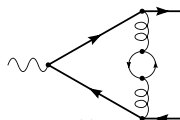
Details: NRQCD: Graphs



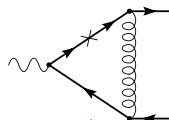
(a)



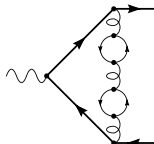
(b)



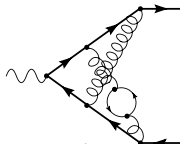
(c)



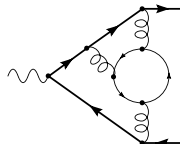
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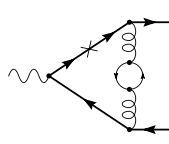
(e)



(f)



(g)



(h)

Details: HQET

on QCD side calculate quark-antiquark-gluon vertex with onshell quarks and zero momentum transfer

- decompose into form factors

$$\Gamma^\mu = \gamma^\mu F_1(q^2) - \frac{i}{2m_Q} \sigma^{\mu\nu} q_\nu F_2(q^2)$$

- use projectors to obtain form factors
- colour charge

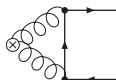
$$Z_2^{OS} F_1(0) = 1 \quad \Rightarrow \quad F_1(0)^{-1} = Z_2^{OS}$$

Melnikov, v. Rittbergen '00; PM, Mihaila, Piclum, Steinhauser '07

- chromomagnetic moment

$$\gamma_{cm} = Z_2^{OS} F_2(0) \quad \Rightarrow \quad C_{cm}(\mu) = Z_{cm}(\alpha_s(\mu)) [1 + \mu_{cm}(\alpha_s(\mu))]$$

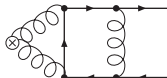
Details: HQET: Graphs



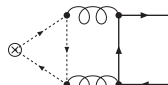
(a)



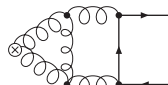
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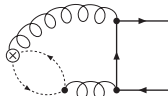
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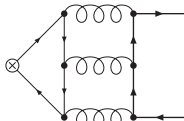
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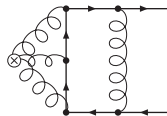
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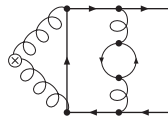
(f)



(g)



(h)



(i)

Results for NRQCD

- bottom system
- top system

Bottom System

Decay rate of $\Upsilon(1S)$ into leptons

$$\Gamma(\Upsilon(1S) \rightarrow l^+l^-) = \Gamma_{LO} \rho_1 \left[c_V^2(m_b) + \frac{C_F^2 \alpha_S^2(\mu_S)}{12} c_V(m_b) (d_V(m_b) + 3) \right]$$

with $\alpha_S(M_Z) = 0.118$, $m_b = 5.3$ GeV and $\mu_S = 2.0967$ GeV

$$\Gamma_1 \approx \Gamma_1^{LO} (1 - 0.446_{NLO} + 1.75_{NNLO} - 1.67_{NNNLO})$$

changes to

$$\Gamma_1 \approx \Gamma_1^{LO} (1 - 0.446_{NLO} + 1.75_{NNLO} - 1.20_{NNNLO})$$

convergence improved

Top System

normalized $t\bar{t}$ peak cross section at a linear collider

$$\begin{aligned} R(e^+e^- \rightarrow t\bar{t}) &= \frac{\sigma(e^+e^- \rightarrow t\bar{t})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} \\ &= R_{LO\rho_1} \left[c_V^2(m_b) + \frac{C_F^2 \alpha_S^2(\mu_S)}{12} c_V(m_b) (d_V(m_b) + 3) \right] \end{aligned}$$

with $\alpha_S(M_Z) = 0.118$, $m_t = 175$ GeV and $\mu_S = 32.625$ GeV

$$R_1 \approx \Gamma_1^{LO} (1 - 0.243_{NLO} + 0.435_{NNLO} - 0.268_{NNNLO})$$

changes to

$$R_1 \approx \Gamma_1^{LO} (1 - 0.243_{NLO} + 0.435_{NNLO} - 0.195_{NNNLO})$$

Sum of all corrections ~ 0.003

Results for HQET

- mass splitting
- magnetic moments

Results

matching coefficient

$$C_{cm}(m_c) = 1 + 0.2309 + 0.1835 + 0.2362 = 1.6506$$

$$C_{cm}(m_b) = 1 + 0.1492 + 0.0676 + 0.0497 = 1.2664$$

$$C_{cm}(m_t) = 1 + 0.0741 + 0.0144 + 0.0045 = 1.0930$$

anomalous dimension

$$\gamma_{cm}(m_c) = 0.1599 + 0.0298 + 0.0163 = 0.2060$$

$$\gamma_{cm}(m_b) = 0.1033 + 0.0099 + 0.0029 = 0.1160$$

$$\gamma_{cm}(m_t) = 0.0513 + 0.0018 + 0.0002 = 0.0533$$

Heavy Quark Mass Splittings: Lattice

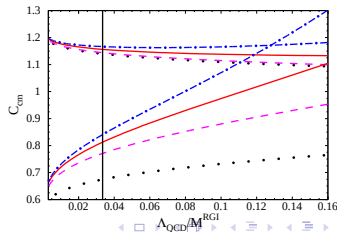
mass splitting of heavy mesons

$$m_{B^*}^2 - m_B^2 = \frac{4}{3} C_{cm}^{(4)}(\mu) \mu_{G(4)}^2(\mu) + \mathcal{O}\left(\frac{\Lambda_{QCD}}{m_b}\right)$$

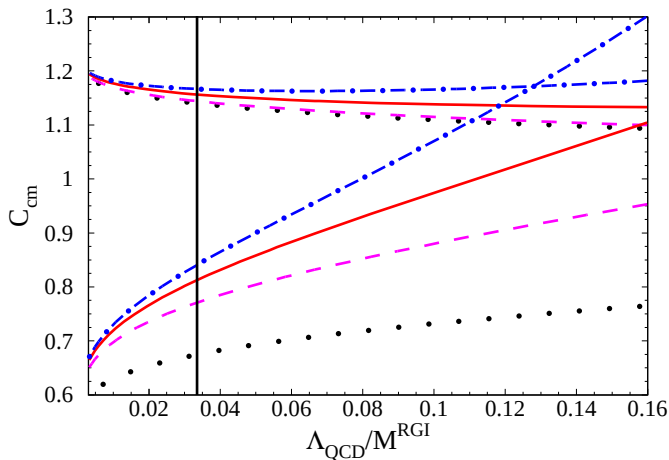
- transform matching coefficient to RGI scheme and rewrite it in terms of Λ_{QCD}/M^{RGI}

$$M^{RGI} = \bar{m}_* \left(2\beta_0 \frac{\alpha_s(\bar{m}_*)}{\pi} \right)^{-\frac{\gamma_{m,0}}{\beta_0}} \exp \left[- \int_0^{\alpha_s(\bar{m}_*)} \left(\frac{\gamma_m}{\beta} - \frac{\gamma_{m,0}}{\beta_0} \right) \frac{d\alpha_s}{\alpha_s} \right]$$

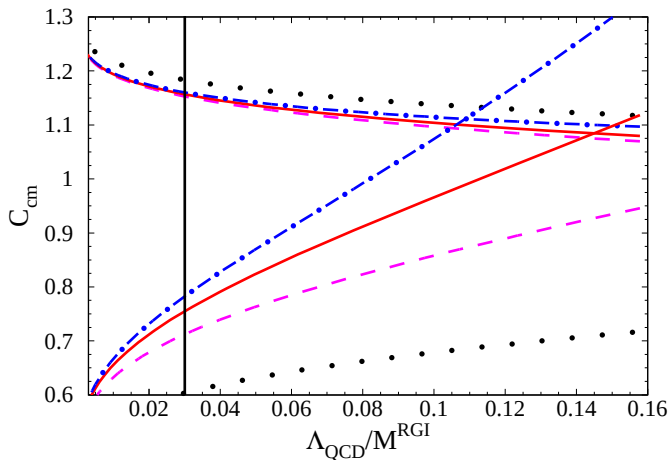
- improved convergence of the matching coefficient
 $\hookrightarrow$ error $\approx 1\%$ [Grozin, Guazzini, PM, Meyer, Piclum, Sommer, Steinhauser '07](#)



Heavy Quark Mass Splittings: Lattice: $N_f=0$



Heavy Quark Mass Splittings: Lattice: $N_f=4$



Heavy Quark Mass Splittings: Perturbative

- matrix element perturbatively not accessible
↪ consider ratio
- matrix elements related by renormalization and decoupling relations

Grozin '00

$$R = \frac{m_{B^*}^2 - m_B^2}{m_{D^*}^2 - m_D^2} = 0.88 \quad (\text{exp.})$$

resummed logs

$$R = 0.8517 - 0.0696 - 0.0908 + [-0.1285] = 0.6914 + [-0.1285]$$

fixed order

$$R = 1 - 0.1113 - 0.0780 - 0.0755 + \dots = 0.7352 + \dots$$

Magnetic moments

replace external gluon by photon

$$\begin{aligned} \frac{a_Q}{Q_Q} = & 0.2122 \alpha_s^{(n_f)}(m_Q) + (0.8417 - 0.0469 n_l) \left[\alpha_s^{(n_f)}(m_Q) \right]^2 \\ & + \left(4.5763 - 0.5856 n_l + 0.0145 n_l^2 \right) \left[\alpha_s^{(n_f)}(m_Q) \right]^3 + \mathcal{O}(\alpha_s^4) \end{aligned}$$

$$a_c(m_c) = 0.0478 + 0.0533 + 0.0758 = 0.1770$$

$$a_b(m_b) = -0.0153 - 0.0103 - 0.0084 = -0.0340$$

$$a_t(m_t) = 0.0152 + 0.0047 + 0.0017 = 0.0215$$

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$$a_b(m_Z) = -0.0083 - 0.0066 - 0.0056 = -0.0206$$

exp: $a_b > -0.135$ (68% CL)

Escribano, Massó '94

Magnetic moments

replace external gluon by photon

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- no measurement of the magnetic moment of the top quark
- sensitive to new physics

Conclusion

- NRQCD
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 - generally improves convergence
 - still incomplete, full calculation in progress

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 - non-singlet three loop contributions calculated
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- HQET
 - chromomagnetic moment calculated
 - important input for lattice calculations
 - convergence improved in RGI scheme
 - pert. results for mass splittings not in good agreement with experimental value
 - byproduct: magnetic moments