



Doubly Heavy Baryons

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J.H., T.Mehen, PRD 73, 054003 (2006)



Outline:

- Motivation
- Heavy Quark-Diquark Symmetry
- Chiral Lagrangian of Doubly Heavy Baryons and Heavy Mesons
- Predictions for Doubly Charm Baryon
- Other Related Work
- Chiral Lagrangian for Doubly Heavy Baryon Semileptonic Decay
- Summary

Motivation

SELEX

PRL 89, 112001 (2002)
PLB 628, 18 (2005)

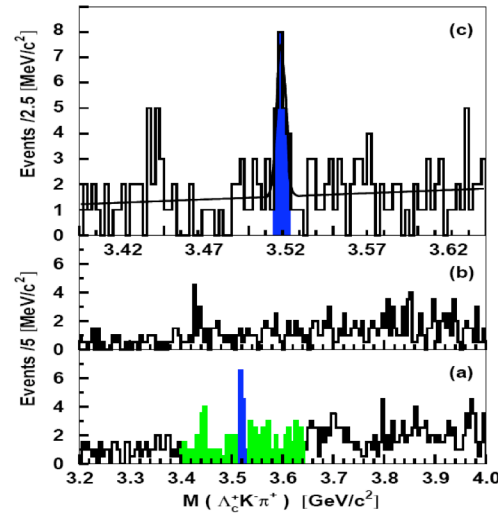


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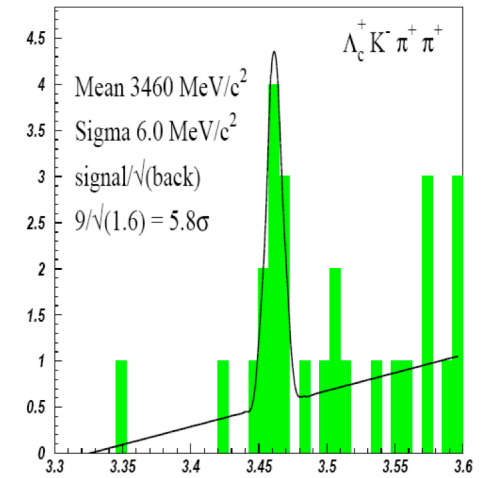


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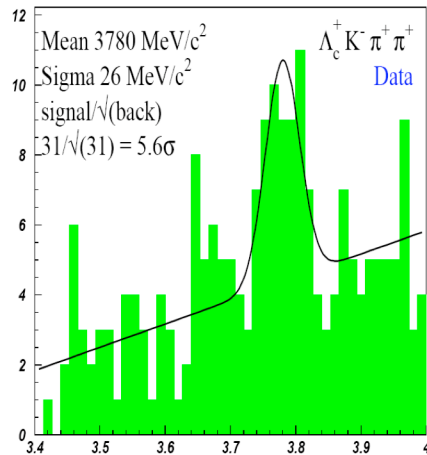


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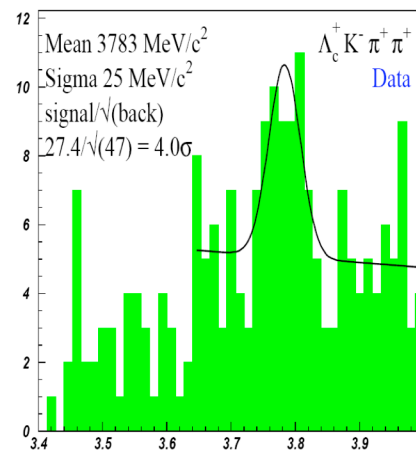


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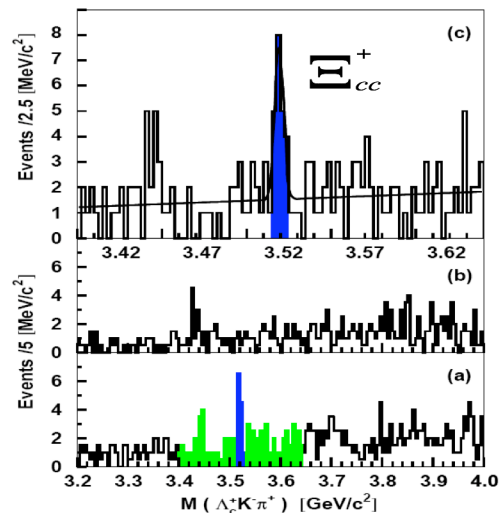


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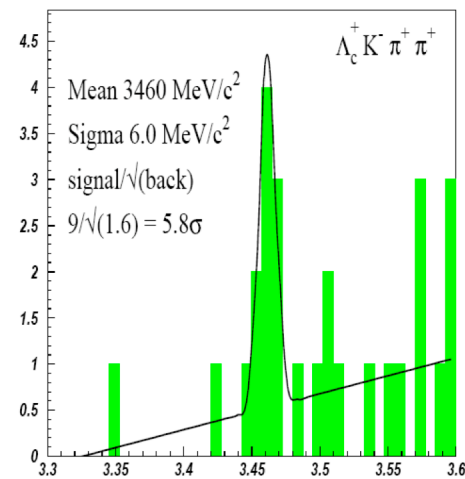


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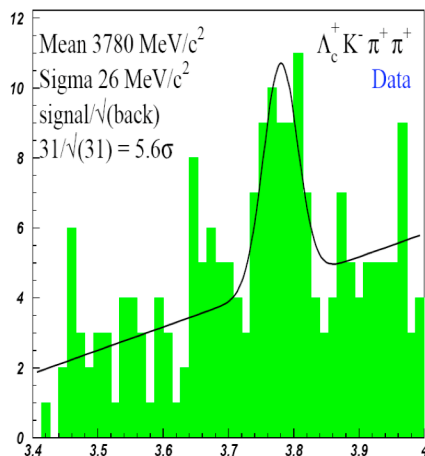


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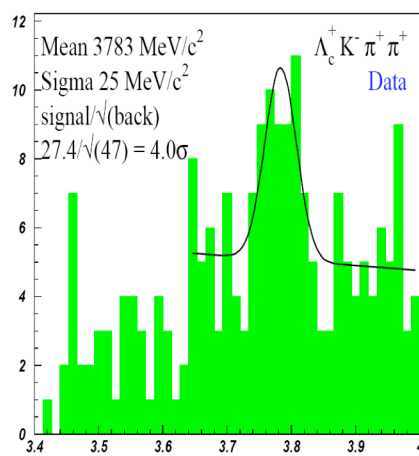


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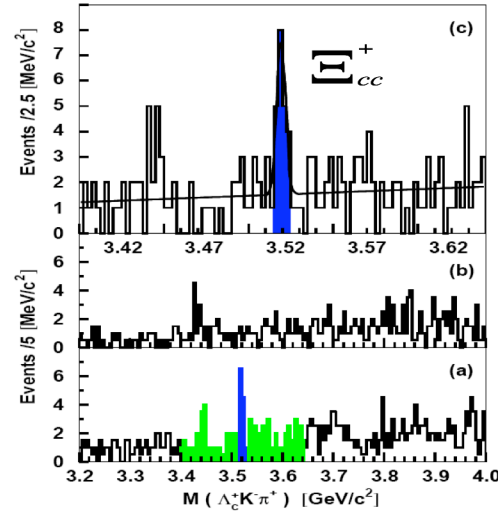


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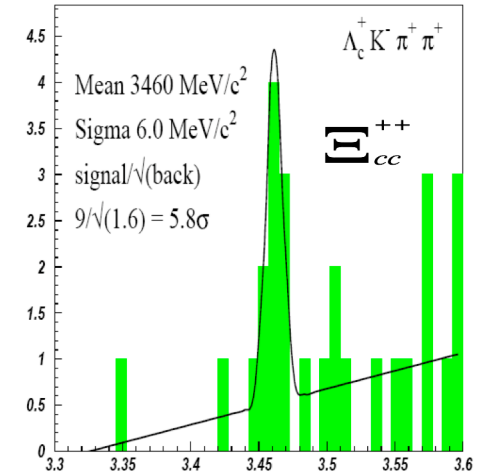


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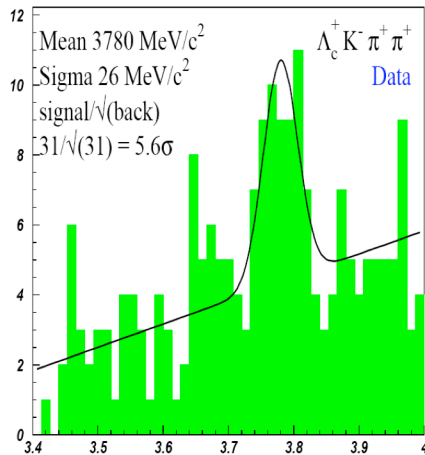


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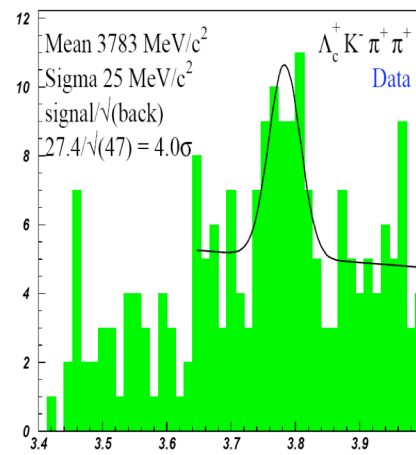


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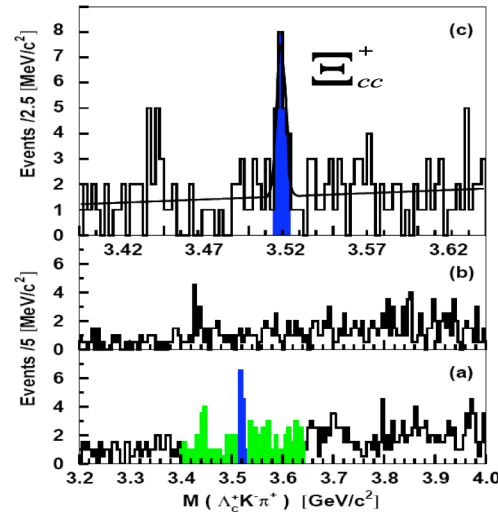


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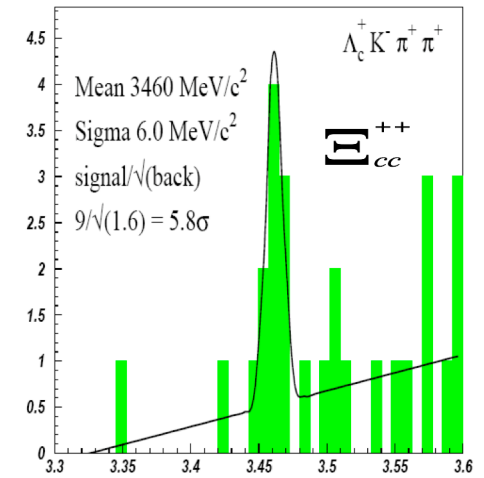


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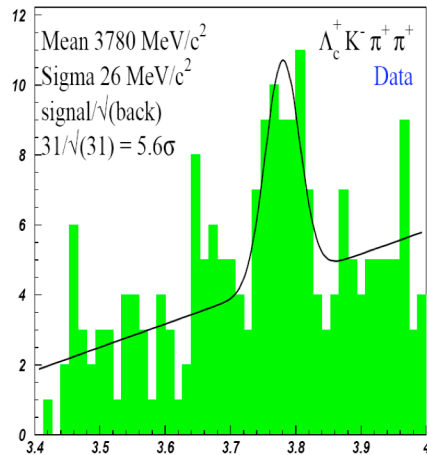


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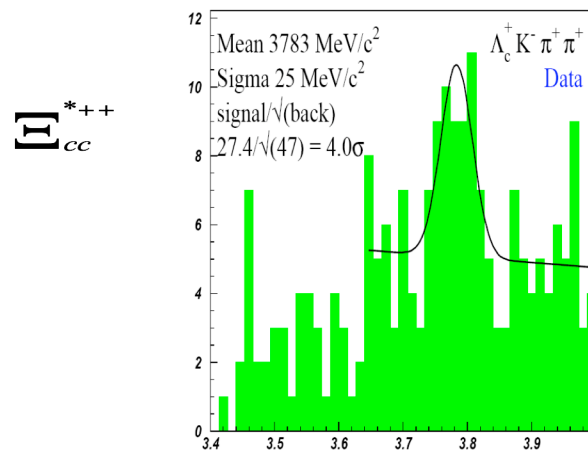
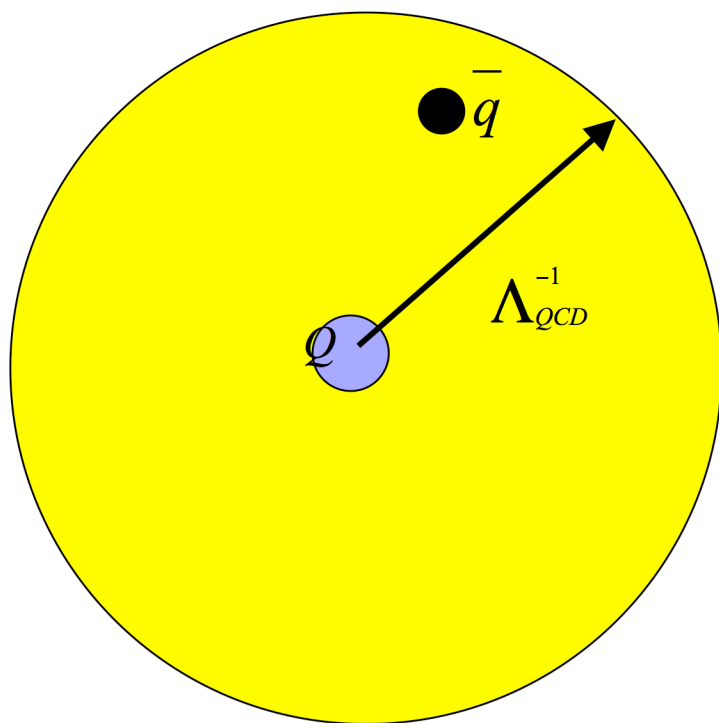


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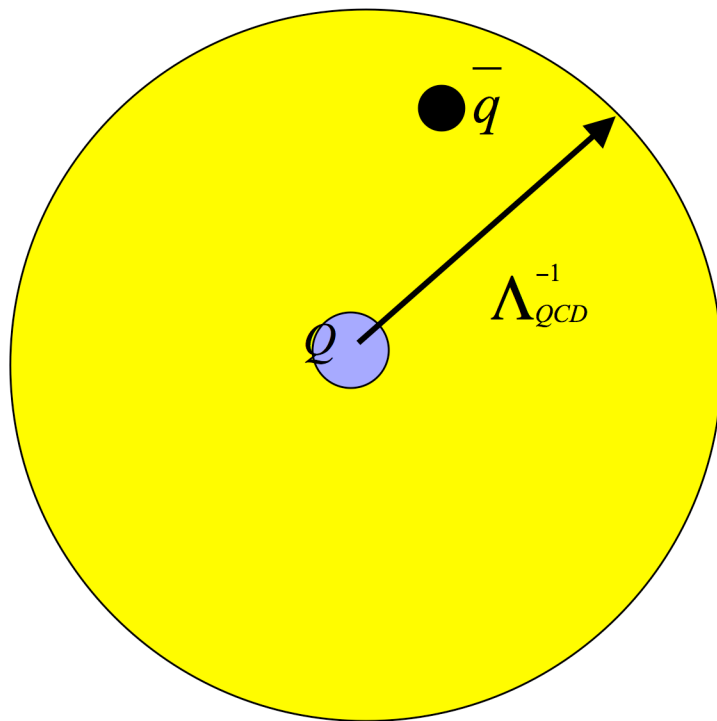


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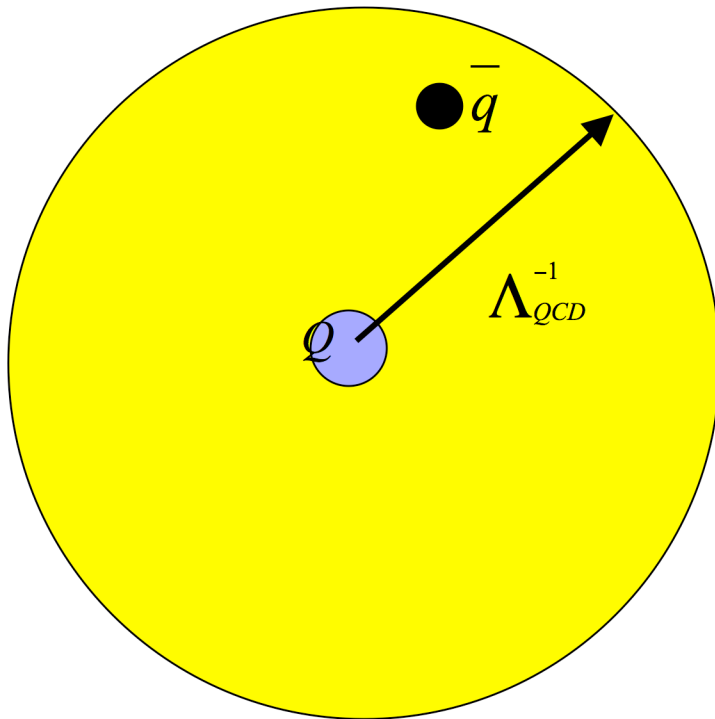
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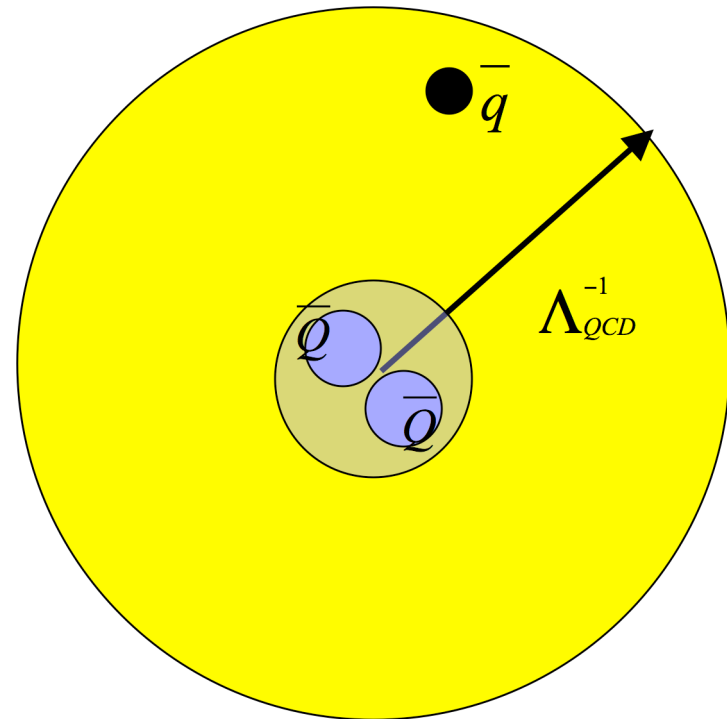


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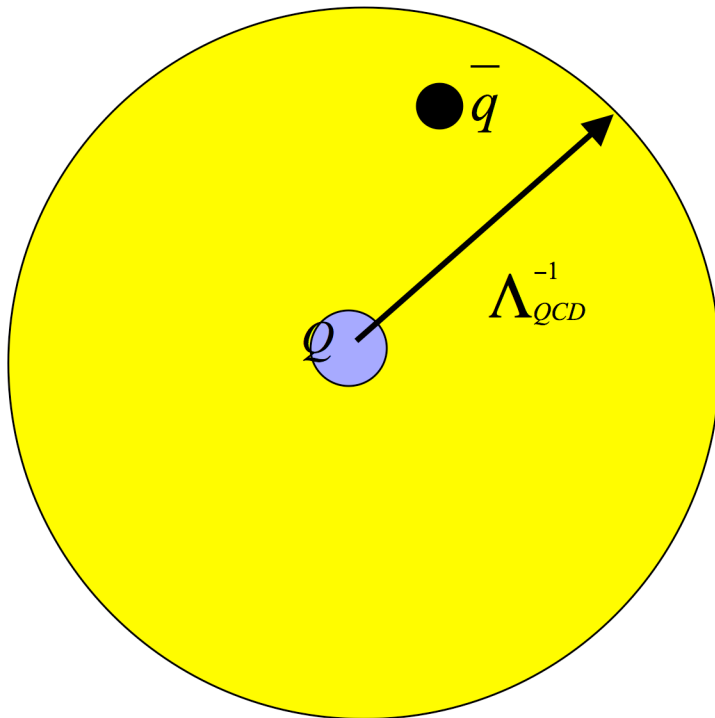


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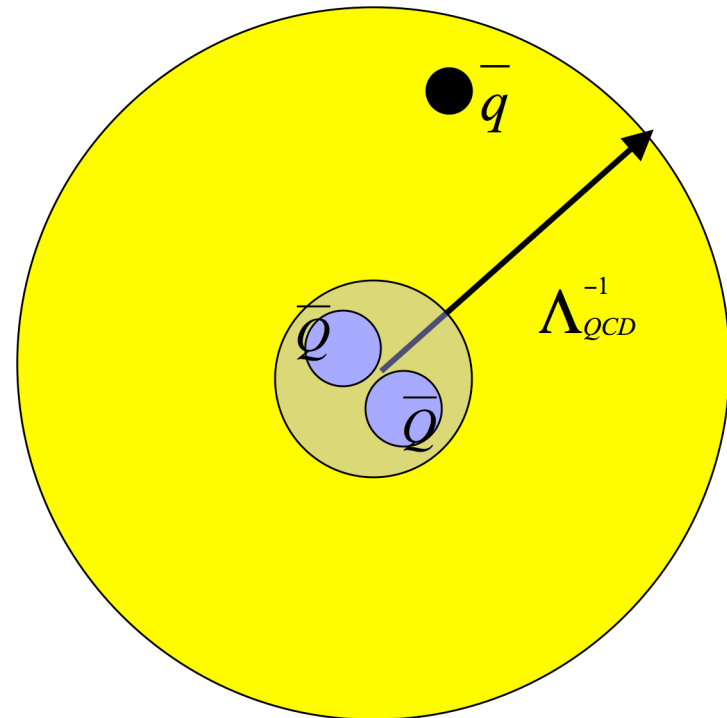
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Heavy Quark-Diquark Symmetry

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HQET Lagrangian

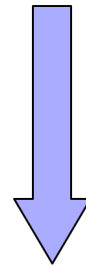
(Savage, Wise, Fleming, Mehen)

$$\begin{aligned}\mathcal{L} = & h^\dagger \left(iD_0 - \frac{\vec{D}^2}{2m_Q} \right) h + \vec{V}^\dagger \cdot \left(iD_0 + \delta - \frac{\vec{D}^2}{4m_Q} \right) \vec{V} \\ & + \frac{g_s}{2m_Q} h^\dagger \vec{\sigma} \cdot \vec{B}^a \frac{\lambda^a}{2} h + \frac{ig_s}{2m_Q} \vec{V}^\dagger \cdot \vec{B}^a \frac{\lambda^a}{2} \times \vec{V}.\end{aligned}$$

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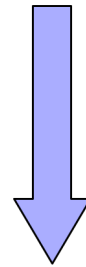


$$\mathcal{Q}_\mu = \begin{pmatrix} h_\alpha \\ V_i \end{pmatrix}$$

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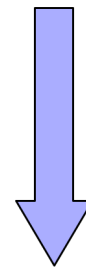
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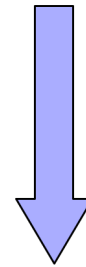
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$$\vec{\Sigma}_{\mu\nu} = \begin{pmatrix} \vec{\sigma}_{\alpha\beta} & 0 \\ 0 & \vec{T}_{jk} \end{pmatrix} \quad (\mathcal{T}^i)_{jk} = -i\epsilon_{ijk}$$



Chiral Lagrangian with Heavy Quark-Diquark Symmetry

$$H_v = \left(\frac{1 + \not{v}}{2} \right) (P_v^{*\mu} \gamma_\mu - \gamma_5 P_v)$$



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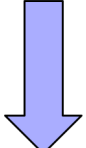
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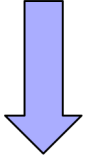
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- Chiral corrections



$$\delta m_{\Xi_{cc}^*} - \delta m_{\Xi_{cc}} - \frac{3}{4}(\delta m_{D^*} - \delta m_D) = \begin{cases} 3.9 \text{ MeV} & \mu = 500 \text{ MeV} \\ 5.3 \text{ MeV} & \mu = 1000 \text{ MeV} \\ 6.1 \text{ MeV} & \mu = 1500 \text{ MeV} \end{cases}$$

Electromagnetic Decay

$$\mathcal{L} = \frac{e\beta}{2} \text{Tr}[\mathcal{H}_a^\dagger \mathcal{H}_b \vec{\sigma} \cdot \vec{B} Q_{ab}] \\ + \frac{e}{2m_Q} Q' \text{Tr}[\mathcal{H}_a^\dagger \vec{\Sigma}' \cdot \vec{B} \mathcal{H}_b]$$

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■ Electromagnetic widths

$$\Gamma[P_a^* \rightarrow P_a \gamma] = \frac{\alpha}{3} \left(\beta Q_{aa} + \frac{Q'}{m_Q} \right)^2 \frac{m_P}{m_{P^*}} E_\gamma^3 \\ \Gamma[\Xi_a^* \rightarrow \Xi_a \gamma] = \frac{4\alpha}{9} \left(\beta Q_{aa} - \frac{Q'}{m_Q} \right)^2 \frac{m_\Xi}{m_{\Xi^*}} E_\gamma^3$$

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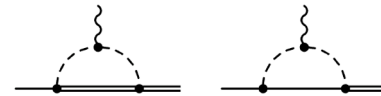
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TABLE I. Predictions for the electromagnetic widths of the Ξ_{cc}^{*+} and Ξ_{cc}^{*++} . The fits are explained in the text.

Fit	β^{-1} (MeV)	m_c (MeV)	$\Gamma[\Xi_{cc}^{*++}]$ (keV)	$\Gamma[\Xi_{cc}^{*+}]$ (keV)
QM 1	379	1863	$3.3 \left(\frac{E_\gamma}{80 \text{ MeV}} \right)^3$	$2.6 \left(\frac{E_\gamma}{80 \text{ MeV}} \right)^3$
QM 2	356	1500	$3.4 \left(\frac{E_\gamma}{80 \text{ MeV}} \right)^3$	$3.2 \left(\frac{E_\gamma}{80 \text{ MeV}} \right)^3$
χ PT 1	272	1432	$2.3 \left(\frac{E_\gamma}{80 \text{ MeV}} \right)^3$	$3.5 \left(\frac{E_\gamma}{80 \text{ MeV}} \right)^3$
χ PT 2	276	1500	$2.3 \left(\frac{E_\gamma}{80 \text{ MeV}} \right)^3$	$3.3 \left(\frac{E_\gamma}{80 \text{ MeV}} \right)^3$



Excited States

■ P-wave excited diquark

$$\Xi_{cc}^{\mathcal{P}} \quad J^P = \frac{1}{2}^+ \qquad \Xi_{cc}^{\mathcal{P}*} \quad J^P = \frac{3}{2}^+$$

$$\Gamma[\Xi_{cc}^{\mathcal{P}*} \rightarrow \Xi_{cc}^* \pi] = \lambda_{3/2}^2 111 \text{ MeV}$$

$$\Gamma[\Xi_{cc}^{\mathcal{P}} \rightarrow \Xi_{cc} \pi] = \lambda_{1/2}^2 111 \text{ MeV}$$

■ Radially excited diquark

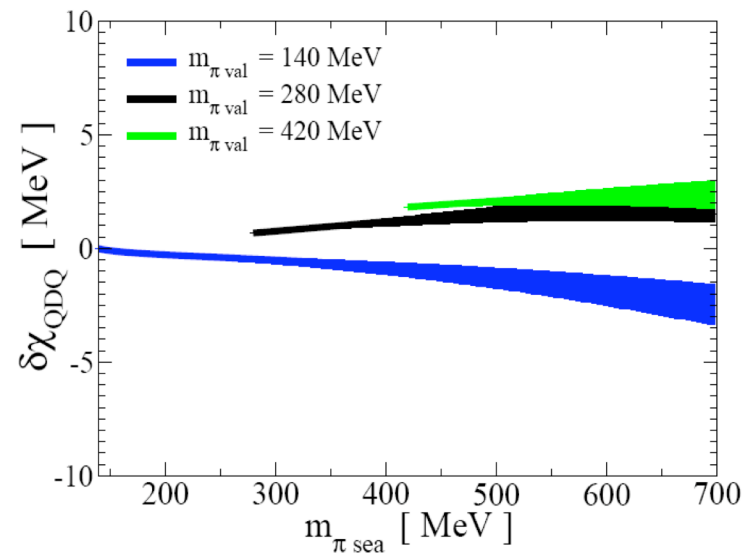
$$\Xi_{cc}'^* \quad J^P = \frac{3}{2}^- \qquad \Xi_{cc}' \quad J^P = \frac{1}{2}^-$$

$$\frac{\Gamma[\Xi_{cc}'^* \rightarrow \Xi_{cc}^* \pi]}{\Gamma[\Xi_{cc}'^* \rightarrow \Xi_{cc} \pi]} = 0.56 \qquad \frac{\Gamma[\Xi_{cc}' \rightarrow \Xi_{cc}^* \pi]}{\Gamma[\Xi_{cc}' \rightarrow \Xi_{cc} \pi]} = 2.3$$

In Quenched and Partially Quenched ChiPT

T.M., B.Tiburzi, PRD 74, 054505

- Partially quenched corrections to Quark-diquark symmetry prediction for hyperfine splittings



$$\delta\chi_{QDQ} = \left[M_{\Xi^{(3/2)}}^{(3/2)} - M_{\Xi}^{(3/2)} - \frac{3}{4} \left(M_{P^*}^{(3/2)} - M_P^{(3/2)} \right) \right]_{\text{PQ}\chi\text{PT}} - \left[M_{\Xi^{(3/2)}}^{(3/2)} - M_{\Xi}^{(3/2)} - \frac{3}{4} \left(M_{P^*}^{(3/2)} - M_P^{(3/2)} \right) \right]_{\chi\text{PT}}$$

Diquark Effective Action from vNRQCD

S. Fleming, T.Mehen, PRD 73, 034502

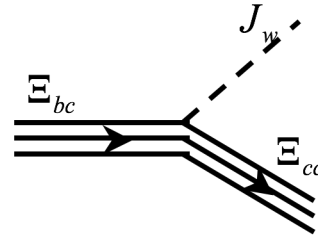
■ vNRQCD Lagrangian (Luke, Manohar, Rothstein)

$$\mathcal{L} = \sum_{\mathbf{p}} \chi_{\mathbf{p}}^{\dagger} \left(iD^0 - \frac{(\mathbf{p} - i\mathbf{D})^2}{2m_Q} + \frac{g}{2m_Q} \boldsymbol{\sigma} \cdot \mathbf{B} \right) \chi_{\mathbf{p}} \\ - \frac{1}{2} \sum_{\mathbf{p}, \mathbf{q}} \frac{g_s^2}{(\mathbf{p} - \mathbf{q})^2} \chi_{\mathbf{q}}^{\dagger} \bar{T}^A \chi_{\mathbf{p}} \chi_{-\mathbf{q}}^{\dagger} \bar{T}^A \chi_{-\mathbf{p}} + \dots$$

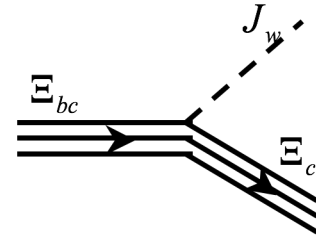
■ Hubbard-Stratonovich trans. Integrate out χ

$$\mathcal{L}_{\mathbf{T}} = \int d^3\mathbf{r} \mathbf{T}_{\mathbf{r}}^{\dagger} \left(iD_0 + \frac{\nabla_{\mathbf{r}}^2}{m_Q} - V^{(3)}(r) \right) \mathbf{T}_{\mathbf{r}} + \frac{g}{2m_Q} \int d^3\mathbf{r} i \mathbf{T}_{\mathbf{r}}^{\dagger} \cdot \mathbf{B} \times \mathbf{T}_{\mathbf{r}} \\ = \sum_n \mathbf{T}_n^{\dagger} (iD_0 + \delta_n) \mathbf{T}_n + \frac{g}{2m_Q} i \sum_n \mathbf{T}_n^{\dagger} \cdot \mathbf{B} \times \mathbf{T}_n \\ \mathbf{T}_{\mathbf{r}}^i = \sum_n \mathbf{T}_n^i \phi_n(\mathbf{r}) \quad \left(-\frac{\nabla_{\mathbf{r}}^2}{m_Q} + V^{(3)}(r) \right) \phi_n(\mathbf{r}) = -\delta_n \phi_n(\mathbf{r})$$

Semileptonic Decay



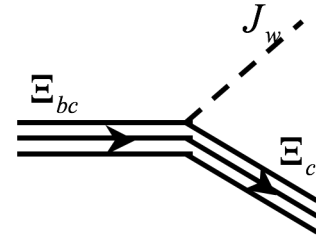
Semileptonic Decay



- Quark current $\bar{c}\gamma^\mu(1-\gamma^5)b$

$$\downarrow$$
$$c^+(\delta^{\mu 0} - \delta^{\mu i}\sigma^i)b$$

Semileptonic Decay

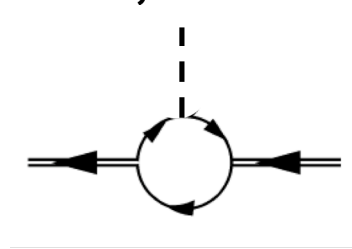


- Quark current $\bar{c}\gamma^\mu(1-\gamma^5)b$

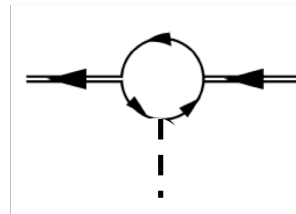
$$\Downarrow$$

$$c^+(\delta^{\mu 0} - \delta^{\mu i}\sigma^i)b$$

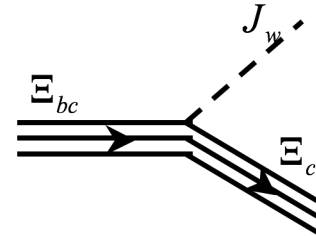
- Integrate out b,c



S. Fleming, T.Mehen, PRD 73, 034502



Semileptonic Decay

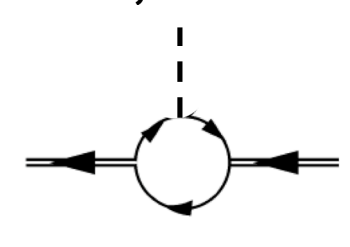


- Quark current $\bar{c}\gamma^\mu(1-\gamma^5)b$

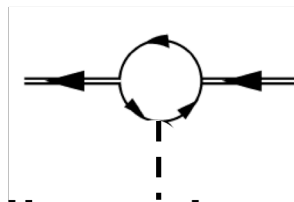


$$c^+(\delta^{\mu 0} - \delta^{\mu i}\sigma^i)b$$

- Integrate out b,c



S. Fleming, T.Mehen, PRD 73, 034502



- Effective Lagrangian of diquark coupled to weak current

$$L \propto J_w^0 T_{cc}^{+i} T_{bc}^i - i J_w^i T_{cc}^{+j} \epsilon_{ijk} T_{bc}^k - J_w^i T_{cc}^{+i} T_{bc}^i$$

Semileptonic Decay

$$T_{a,i\beta} = \sqrt{2} \left(\Xi_{a,i\beta}^* + \frac{1}{\sqrt{3}} \Xi_{a,\gamma} \sigma_{\gamma\beta}^i \right) \quad T' = \sqrt{2} \Xi'$$

- Reproduce zero recoil Cascade bc to cascade cc transition matrix elements.

$$\Xi_{bc} \rightarrow \Xi_{cc} \quad \eta \bar{u}_{cc} \left(2\gamma^\mu - \frac{4}{3} \gamma^\mu \gamma_5 \right) u_{bc}$$

$$\Xi'_{bc} \rightarrow \Xi_{cc} \quad \frac{-2}{\sqrt{3}} \eta \bar{u}_{cc} (-\gamma^\mu \gamma_5) u_{bc}$$

$$\Xi_{bc} \rightarrow \Xi_{cc}^* \quad \frac{-2}{\sqrt{3}} \eta \bar{u}_{cc}^\mu u_{bc}$$

$$\Xi'_{bc} \rightarrow \Xi_{cc}^* \quad -2\eta \bar{u}_{cc}^\mu u_{bc}$$

$$\Xi_{bc}^* \rightarrow \Xi_{cc} \quad \frac{-2}{\sqrt{3}} \eta \bar{u}_{cc} u_{bc}^\mu$$

$$\Xi_{bc}^* \rightarrow \Xi_{cc}^* \quad -2\eta \bar{u}_{cc}^\lambda (\gamma^\mu - \gamma^\mu \gamma_5) u_{bc\lambda}$$

J. Flynn, J. Nieves, arXiv:0706.2805

Semileptonic Decay

$$L \propto J_w^0 T_{cc}^{+i} T_{bc}^i - i J_w^i T_{cc}^{+j} \varepsilon_{ijk} T_{bc}^k - J_w^i T_{cc}^{+i} T_{bc}'$$

$$T_{a,i\beta} = \sqrt{2} \left(\Xi_{a,i\beta}^* + \frac{1}{\sqrt{3}} \Xi_{a,\gamma} \sigma_{\gamma\beta}^i \right) \quad T' = \sqrt{2} \Xi'$$

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$$\Xi_{bc} \rightarrow \Xi_{cc} \quad \eta \bar{u}_{cc} \left(2\gamma^\mu - \frac{4}{3} \gamma^\mu \gamma_5 \right) u_{bc}$$

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$$\Xi_{bc}^* \rightarrow \Xi_{cc}^* \quad -2\eta \bar{u}_{cc}^\lambda (\gamma^\mu - \gamma^\mu \gamma_5) u_{bc\lambda}$$

J. Flynn, J. Nieves, arXiv:0706.2805



Summary

- Chiral Lagrangian of doubly heavy baryons and heavy mesons with quark-diquark symmetry
- Predictions:
Chiral mass correction,
em decays,
strong decays of excited states
- In progress:
Chiral Lagrangian of doubly heavy baryon semileptonic decay
Chiral correction



Electromagnetic couplings

$$\begin{aligned}\mathcal{L}_{em} &= \frac{e}{2m_Q} Q' h^\dagger \vec{\sigma} \cdot \vec{B} h - \frac{ie}{m_Q} Q' \vec{V}^\dagger \cdot \vec{B} \times \vec{V} \\ &= \frac{e}{2m_Q} Q' \mathcal{Q}_\mu^\dagger \vec{\Sigma}'_{\mu\nu} \cdot \vec{B} \mathcal{Q}_\nu,\end{aligned}$$

$$\vec{\Sigma}'_{\mu\nu} = \begin{pmatrix} \vec{\sigma}_{\alpha\beta} & 0 \\ 0 & -2\vec{T}_{jk} \end{pmatrix}$$

Electromagnetic couplings

$$\begin{aligned}\mathcal{L}_{em} &= \frac{e}{2m_Q} Q' h^\dagger \vec{\sigma} \cdot \vec{B} h - \frac{ie}{m_Q} Q' \vec{V}^\dagger \cdot \vec{B} \times \vec{V} \\ &= \frac{e}{2m_Q} Q' Q_\mu^\dagger \vec{\Sigma}'_{\mu\nu} \cdot \vec{B} Q_\nu,\end{aligned}$$

$$\vec{\Sigma}'_{\mu\nu} = \begin{pmatrix} \vec{\sigma}_{\alpha\beta} & 0 \\ 0 & -2\vec{T}_{jk} \end{pmatrix}$$

$$\mathcal{L} = \frac{e\beta}{2} \text{Tr}[H_a^\dagger H_b \vec{\sigma} \cdot \vec{B} Q_{ab}] + \frac{e}{2m_Q} Q' \text{Tr}[H_a^\dagger \vec{\sigma} \cdot \vec{B} H_a]$$



Electromagnetic couplings

$$\begin{aligned}\mathcal{L}_{em} &= \frac{e}{2m_Q} Q' h^\dagger \vec{\sigma} \cdot \vec{B} h - \frac{ie}{m_Q} Q' \vec{V}^\dagger \cdot \vec{B} \times \vec{V} \\ &= \frac{e}{2m_Q} Q' Q_\mu^\dagger \vec{\Sigma}'_{\mu\nu} \cdot \vec{B} Q_\nu,\end{aligned}$$

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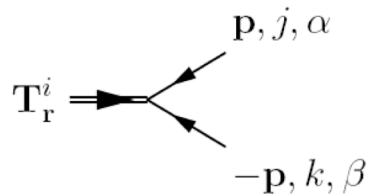
$$\mathcal{L} = \frac{e\beta}{2} \text{Tr}[H_a^\dagger H_b \vec{\sigma} \cdot \vec{B} Q_{ab}] + \frac{e}{2m_Q} Q' \text{Tr}[H_a^\dagger \vec{\sigma} \cdot \vec{B} H_a]$$



$$\begin{aligned}\mathcal{L} &= \frac{e\beta}{2} \text{Tr}[\mathcal{H}_a^\dagger \mathcal{H}_b \vec{\sigma} \cdot \vec{B} Q_{ab}] \\ &\quad + \frac{e}{2m_Q} Q' \text{Tr}[\mathcal{H}_a^\dagger \vec{\Sigma}' \cdot \vec{B} \mathcal{H}_b]\end{aligned}$$

Coupled to Weak Current

	S	J^P	I	S_{hh}^π	
Ξ_{cc}	0	$\frac{1}{2}^+$	$\frac{1}{2}$	1^+	ccl
Ξ_{cc}^*	0	$\frac{3}{2}^+$	$\frac{1}{2}$	1^+	ccl
Ξ'_{bc}	0	$\frac{1}{2}^+$	$\frac{1}{2}$	0^+	bcl
Ξ_{bc}	0	$\frac{1}{2}^+$	$\frac{1}{2}$	1^+	bcl
Ξ_{bc}^*	0	$\frac{3}{2}^+$	$\frac{1}{2}$	1^+	bcl



$$-ie^{-i\mathbf{p}\cdot\mathbf{r}} V^{(3)}(r) \frac{\epsilon_{ijk}}{2} (\sigma\epsilon)_{\alpha\beta}$$

