# The Truth Is Out There

#### Line Shapes of the X(3872)

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## Line Shapes of the X(3872)

- What is the X(3872)?
- Line shapes of the X(3872)
- ... within  $\sim 1~{\rm MeV}$  of  $D^{*0}\bar{D}^0$  threshold
- ... within  $\sim 10~{\rm MeV}$  of  $D^*\bar{D}$  threshold

#### <u>References</u>

Braaten and Lu, arXiv:0709.2697 [hep-ph] Braaten and Lu, to appear soon on arXiv [hep-ph]

### What is the X(3872)?

### Two crucial experimental facts

• mass is extremely close to  $D^{*0}\overline{D}^{0}$  threshold

 $M_X - (M_{D^{*0}} + M_{D^0}) = -0.6 \pm 0.6 \text{ MeV}$ 

measured in  $J/\psi \pi^+\pi^-$  decay mode by Belle, CDF, Babar, D0 precise determination of  $D^0$  mass by CLEO

• quantum numbers  $J^{PC} = 1^{++}$  strongly preferred observation of  $X \to J/\psi \gamma$  by Belle analyses of  $X \to J/\psi \pi^+\pi^-$  by Belle, CDF observation of  $X \to D^0 \overline{D}{}^0 \pi^0$  by Belle, Babar

Two crucial experimental facts

- $J^{PC} = 1^{++}$
- $\implies$  S-wave coupling to  $D^{*0}\overline{D}^0$  (and  $D^0\overline{D}^{*0}$ )
- $M_X (M_{D^{*0}} + M_{D^0}) = -0.6 \pm 0.6 \text{ MeV}$  $\implies$  resonant interaction with  $D^{*0}\overline{D}^0$  (and  $D^0\overline{D}^{*0}$ )

Conclusion: X(3872) is either/or

• weakly-bound charm meson molecule

$$X = \frac{1}{\sqrt{2}} \left( D^{*0} \bar{D}^0 + D^0 \bar{D}^{*0} \right)$$

• virtual state of charm mesons

Nonrelativistic Quantum Mechanics

2-body system with short-range interactions and S-wave resonance sufficiently close to threshold has universal properties that

- depend only on the large scattering length *a*
- are insensitive to details of interactions at shorter distances structure of constituents mechanism for resonance fine-tuning of potential tuning of energy of molecule etc.

"Universality of Few-Body Systems with Large Scattering Length" Braaten and Hammer, arXiv:cond-mat/0410417 (Physics Reports)

Universal features

- large scattering length *a*
- $\bullet$  cross section at low energy E

$$\sigma(\underline{E}) = \frac{4\pi a^2}{1 + 2M_{D^*\bar{D}} a^2 E}$$

shallow S-wave bound states

*a* < 0: none

a>0: one binding energy:  $E_X=1/(2M_{D^*\bar{D}}a^2)$ mean separation:  $\langle r \rangle_X=a/2$ 

X(3872) has universal properties determined by large scattering length a in  $D^{*0}\overline{D}^{0}+D^{0}\overline{D}^{*0}$  channel insensitive to all shorter length scales of QCD

Universal results for a > 0:

$$E_X = 1/(2M_{D^*\bar{D}}a^2)$$
  
$$\langle r \rangle_X = a/2$$

measured binding energy:  $E_X = 0.6 \pm 0.6$  MeV

predicted mean separation:  $\langle r \rangle_X = 2.9^{+\infty}_{-0.9}$  fm

#### Beauty and the Beast

## $B \longrightarrow K + X(3872)$

# $\langle r \rangle_X = 2.9^{+\infty}_{-0.9}$ fm

LeFou: Tell us again, old man, just how big was the Beast? Maurice: It was enormous, I'd say at least 8, no, more like 10 fermis!

LeFou: Well, you don't get much crazier than that!

Belle: My father's not crazy and I can prove it!

## Line Shapes of the X(3872)

Line shape of X in decay mode C = invariant mass distribution of C:  $M_{D^{*0}} + M_{D^0} + E$ 

$$\frac{d\Gamma}{dE}[B \to K + C]$$

Mass measurements of X(3872) ...

... in  $J/\psi \pi^+\pi^-$ 

$$M_X - (M_{D^{*0}} + M_{D^0}) = -0.6 \pm 0.6$$
 MeV

... in  $D^0 \bar{D}^0 \pi^0$ 

$$\begin{split} M - (M_{D^{*0}} + M_{D^0}) &= +4.1 \pm 0.7^{+0.3}_{-1.6} \text{ MeV} \quad \text{(Belle)} \\ &= +4.3 \pm 1.1 \pm 0.5 \text{ MeV} \quad \text{(Babar)} \end{split}$$

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Line Shapes of X (cont.)
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short distances: \ll |a| \sim 6 fm
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large momenta: \gg 1/|a| \sim 30 MeV
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Qualitative difference between decay modes

 $X \longrightarrow J/\psi \, \pi^+ \pi^-$ 

 $X \longrightarrow D^0 \overline{D}{}^0 \pi^0$ 

involves decay of constituent  $D^{*0} \rightarrow D^0 \pi^0$  $\bar{D}^{*0} \rightarrow \bar{D}^0 \pi^0$ 

decay products have small momenta  $\sim 1/|a|$ 

Qualitative difference between bound state and virtual state



Quantitative behavior of line shapes may depend on

- $D^*$  widths  $\Gamma[D^{*0}] = 65.5 \pm 15.4 \text{ keV}$
- inelastic scattering channels of charm mesons  $J/\psi \pi^+\pi^-$ ,  $J/\psi \pi^+\pi^-\pi^0$ , ...
- charged charm mesons  $D^{*+}D^{-}$  threshold: +8.1 MeV
- 3-body channels:  $D\bar{D}\pi$   $D^0\bar{D}^0\pi^0$  threshold: -7.1 MeV  $D^+\bar{D}^0\pi^-$ ,  $D^0D^-\pi^+$  threshold: +2.3 MeV see talk by Tom Mehen

Recent analyses of data from Belle and Babar on  $B^+ \longrightarrow K^+ + X(3872)$ in decay channels  $J/\psi \pi^+\pi^-$ ,  $D^0 \overline{D}{}^0 \pi^0$ 

 Hanhart, Kalashnikova, Kudryavtsev, Nefediev [arXiv:0704.0605] included effects of D\*±D<sup>∓</sup> channel included effects of inelastic channels J/ψ π<sup>+</sup>π<sup>-</sup>, J/ψ π<sup>+</sup>π<sup>-</sup>π<sup>0</sup> ignored effects of D\* widths assumed bound state cannot decay into D<sup>0</sup>D¯<sup>0</sup>π<sup>0</sup>
 Conclusion: X(3872) must be a virtual state See talk by Hanhart

#### • Braaten and Lu [arXiv:0709.2697]

neglected effects of  $D^{*\pm}D^{\mp}$  channel included effects of inelastic channels  $J/\psi \pi^+\pi^-$ ,  $J/\psi \pi^+\pi^-\pi^0$ included effects of  $D^{*0}$  width Conclusion: data prefers X(3872) to be a bound state but virtual state not excluded

#### Analysis of Braaten and Lu [arXiv:0709.2697]

Belle data on  $B^+ \rightarrow K^+ + X$  with total experimental errors subtracted line shapes with 3 adjustable parameters two local minima of  $\chi^2$  with Re $\gamma = +17.3$  MeV and +57.8 MeV



#### Conclusions

data prefers X(3872) to be a bound state, but virtual state not excluded can explain difference between measured mass in  $J/\psi \pi^+\pi^-$  and in  $D^0 \overline{D}{}^0 \pi^0$ 

## Line Shapes of X(3872)within $\sim 1$ MeV of $D^{*0}\bar{D}^0$ Threshold

S-wave resonance in 1<sup>++</sup> channel:  $D^{*0}\overline{D}^0 + D^0\overline{D}^{*0}$ 

1. begin with scattering amplitude f(E) that satisfies unitarity exactly:

 $\operatorname{Im} f(E) = |f(E)|^2 \sqrt{2M_{D^*\bar{D}}E}$ 

- 2. apply deformations that take into account  $D^{*0}$  width inelastic scattering channels
- 3. insert into factorization formulas for the line shapes

Line Shapes within  $\sim 1$  MeV of Threshold (cont.)

1. scattering amplitude that satisfies unitarity exactly:

$$f(E) = \frac{1}{-\gamma + \kappa(E)}$$
  

$$\gamma = 1/a$$
  

$$\kappa(E) = (-2M_{D^*\bar{D}}E - i\varepsilon)^{1/2}$$

- a > 0: bound state pole at  $E = -\gamma^2/(2M_{D^*\bar{D}})$
- *a* < 0: virtual state

pole on second sheet of complex energy E

Line Shapes within  $\sim 1$  MeV of Threshold (cont.)

2. Apply deformations to unitary scattering amplitude

$$f(E) = \frac{1}{-\gamma + \kappa(E)}$$

• take into account  $D^{*0}$  width:  $M_{D^{*0}} \rightarrow M_{D^{*0}} - i\Gamma[D^{*0}]/2$ 

$$\kappa(E) = \left(-2M_{D^*\bar{D}}(E + i\Gamma[D^{*0}]/2)\right)^{1/2}$$

• take into account inelastic scattering channels

$$\gamma \longrightarrow \operatorname{Re} \gamma + i \operatorname{Im} \gamma, \quad \operatorname{Im} \gamma > 0$$

Optical theorem:

 $\operatorname{Im} f(E) = |f(E)|^2 \left( \operatorname{Im} \gamma - \operatorname{Im} \kappa(E) \right)$ consistent with multi-channel unitarity Line Shapes within  $\sim 1$  MeV of Threshold (cont.)

3. Factorization formulas for the line shapes factor rates into long-distance factor (depends on  $E, \gamma$ ) × short-distance factors (insensitive to  $E, \gamma$ )

$$B^{+} \to K^{+} + C, \text{ where } C = J/\psi \pi^{+} \pi^{-}, J/\psi \pi^{+} \pi^{-} \pi^{0}, \dots$$
$$\frac{d\Gamma}{dE} = 2\Gamma_{B^{+}}^{K^{+}} \times |f(E)|^{2} \times \Gamma^{C}$$
$$B^{+} \to K^{+} + D^{0} \bar{D}^{0} \pi^{0}$$
$$\frac{d\Gamma}{dE} = 2\Gamma_{B^{+}}^{K^{+}} \times |f(E)|^{2} \left[ M_{D^{*}\bar{D}} \left( \sqrt{E^{2} + \Gamma[D^{*0}]^{2}/4} + E \right) \right]^{1/2}$$
$$\times \text{Br}[D^{*0} \to D^{0} \pi^{0}]$$

#### Short-distance factors

- $\Gamma_B^K$  different for  $B^+ \to K^+$  and  $B^0 \to K^0$
- $\Gamma^{C}$  different for  $J/\psi \pi^{+}\pi^{-}$  and  $J/\psi \pi^{+}\pi^{-}\pi^{0}$

# Line Shapes of X(3872)within $\sim 10$ MeV of $D^{*0}\bar{D}^0$ Threshold

S-wave resonance in two coupled  $1^+$  + channels:  $D^{*0}\overline{D}^0 + D^0\overline{D}^{*0}$ ,  $D^{*+}D^- + D^+D^{*-}$ 

- 1. begin with scattering amplitudes  $f_{00}(E)$ ,  $f_{01}(E)$ ,  $f_{11}(E)$ that satisfy two-channel unitarity exactly with isospin symmetry at high energy  $\implies$  2 scattering parameters:  $\gamma_{I=0}$ ,  $\gamma_{I=1}$
- 2. apply deformations that take into account  $D^{*0}$ ,  $D^{*+}$  widths inelastic scattering channels
- 3. insert into factorization formulas for the line shapes, with short-distance factors constrained by isospin symmetry

Line Shapes within  $\sim 10$  MeV of Threshold (cont.)

### Implications

1. Conceptual error by Braaten and Kusunoki [hep-ph/0412268]

prediction:  $B^0 \to K^0 + X$  is suppressed compared to  $B^+ \to K^+ + X$ Belle, Babar: no indication of strong suppression error: implicitly assumed  $|\gamma_0|, |\gamma_1| \ll \sqrt{2M_{D^*\bar{D}}\nu} = 125$  MeV

2. Conceptual error by Voloshin [arXiv:0704.3029]

prediction: line shapes of X from  $B^0 \to K^0$  same as from  $B^+ \to K^+$ errors: did not allow for resonant scattering between neutral and charged  $D^*\overline{D}$  channels results inconsistent with isospin symmetry in short-distance factors for  $B \to K$  Line Shapes within  $\sim 10$  MeV of Threshold (cont.)

3. Interpretation of scattering amplitude of Hanhart et al.

"generalization of Flatté parametrization for near-threshold resonance" limit  $|\gamma_1| \gg |\gamma_0|, \sqrt{2M_{D^*\bar{D}}\nu}$  gives essentially same scattering amplitude

4. Line shapes depend on production process, decay channel

determined by  $\gamma_0$ ,  $\gamma_1$ different for  $B^+ \to K^+ + X$  and  $B^0 \to K^0 + X$ different for  $J/\psi \ \pi^+\pi^-$ ,  $J/\psi \ \pi^+\pi^-\pi^0$ ,  $D^0 \overline{D}^0 \pi^0$ zeroes in line shapes of  $J/\psi \ \pi^+\pi^ B^+ \to K^+ + X$ : zero near +6 MeV  $B^0 \to K^0 + X$ : zero near -2 MeV no zeroes in line shapes of  $J/\psi \ \pi^+\pi^-\pi^0$ ,  $D^0 \overline{D}^0 \pi^0$ 

5. Ratios of production rates from  $B^0 \to K^0$  and  $B^+ \to K^+$ determined by  $\gamma_0, \gamma_1$ 

