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Measurement of the quark to photon fragmentation function through the inclusive production of prompt photons in hadronic Z^0 decays

The OPAL Collaboration

Abstract

The inclusive production of prompt photons with energy above 10 GeV is measured using the OPAL detector in hadronic Z^0 decays at LEP. In contrast to previous measurements, the prompt photons were not required to be isolated. The production rate and energy spectrum are found to be in agreement with QCD predictions for the quark-to-photon fragmentation function.

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1 Introduction

We gain a greater understanding of the properties of the elementary building blocks of matter and their interactions by studying the properties of hadrons, leptons and photons produced as a result of a primary interaction. The properties of hadrons and charged leptons produced in e^+e^- collisions have been studied in great detail at different centre-of-mass energies. In the case of photons radiated off quarks, prompt photons, in hadronic e^+e^- collisions, much less information is available due to the difficulty of separating these photons from those produced in the decays of other particles [1]-[5]. Both the shape and normalisation of the inclusive prompt photon energy spectrum in e^+e^- collisions are predicted, through the calculation of the quark-to-photon fragmentation function, by leading-order perturbative QCD [6, 7]. This asymptotic prediction has been parametrised in [8]. Non-perturbative effects can be included in the calculation through the vector-meson dominance ansatz as in [9, 10], where boundary terms missing in [8] were also accounted for. The higher-order terms were calculated, and seen to be small at the energies close to the Z^0 peak. Direct experimental study of these predictions is important in providing insight into the non-perturbative and higher order effects in the radiation of photons from quarks. This will also make theoretical predictions of photon production in other processes, such as those occuring at pp and $p\overline{p}$ colliders, more reliable, thus improving sensitivity to possible new phenomena.

At LEP the first measurement of prompt photon production in hadronic Z^0 decays was made by the OPAL Collaboration [11] for photons isolated from other particles in the event, as suggested in [12]. The production of isolated prompt photons was studied in great detail by all LEP experiments [13]-[16]. Following the suggestion of [17], the ALEPH Collaboration extracted the quark-to-photon fragmentation function from the study of non-isolated photons in jets containing a photon carrying more than 70% of the jet energy [18].

Here we present a measurement of the inclusive prompt photon energy spectrum in hadronic Z^0 decays at LEP. This method of studying the quark-to-photon fragmentation function was suggested in [10, 19]. To separate prompt photons from the photons from decays of other particles we use the following method. We selected clusters in the electromagnetic calorimeter not associated with charged tracks. A set of cuts were applied to reduce the background in the sample. The distribution of a variable characterising the transverse shape of the clusters in data was then fitted with a linear combination of the distributions for photons and for background to determine the fraction of prompt photons in the selected sample. The result was then corrected for the selection efficiencies, detector effects and initial state radiation. In the following sections we describe the OPAL detector, the event and electromagnetic cluster selection (sec. 3) and the determination of the number of photons in the selected sample (sec. 4). The efficiency and acceptance corrections are described (sec. 5) followed by the study of systematic effects (sec. 6). Finally the measured prompt photon energy spectrum is presented and discussed (sec. 7).

2 The OPAL detector

The OPAL detector operates at the LEP e^+e^- collider at CERN. A detailed description of the detector can be found in [20]. For this study, the most important components of OPAL were the central detector and the barrel electromagnetic calorimeter with its presampling detector. The central detector, measuring the momenta of charged particles, consists of a system of cylindrical tracking chambers surrounded by a solenoidal coil which produces a uniform axial magnetic field of 0.435 T along the beam axis[1]. The detection efficiency for charged particles is almost 100% within the polar angle range $|\cos\theta| < 0.95$.

The electromagnetic calorimeters completely cover the azimuthal range for polar angles satisfying $|\cos\theta| < 0.98$ providing excellent hermeticity. The barrel electromagnetic calorimeter covers the polar angle range $|\cos\theta| < 0.82$. It consists of 9440 lead glass blocks, each 24.6 radiation lengths deep, almost pointing towards the interaction region. Each block subtends an angular region of approximately 40×40 mrad^2. Half of the block width corresponds to 1.9 Moliére radii. Deposits of energy in adjacent blocks are grouped together to form clusters of electromagnetic energy. The intrinsic energy resolution of $\sigma_{\mathrm{E}}/\mathrm{E} = 0.2\% \oplus 6.3\%/\sqrt{\mathrm{E}}$ is substantially degraded (by a factor $\simeq 2$) due to the presence of two radiation lengths of material in front of the lead glass. For the intermediate region, $0.72 < |\cos\theta| < 0.82$, the amount of material increases up to eight radiation lengths causing further degradation in the energy resolution. The two endcap calorimeters, each made of 1132 lead glass blocks, 22 radiation lengths deep, cover the region of $0.81 < |\cos\theta| < 0.98$. In this study the measurement of inclusive photon production is restricted to the barrel part of the detector. Most of the electromagnetic showers start before the calorimeter and their position at the entrance of the calorimeter is measured by a barrel electromagnetic presampler made of limited streamer mode chambers. The presampler covers the polar angle range $|\cos\theta| < 0.81$ and its angular resolution for photons is approximately 2 mrad.

3 The selection of events and electromagnetic clusters

Our study was based on a sample of 2.5 million hadronic Z^0 decays selected as described in [21] from the data accumulated with the OPAL detector at LEP in 1992, 1993 and 1994 at an e^+e^- centre-of-mass energy of 91.2 GeV. We did not use the off-peak data to avoid additional complications in dealing with data collected at different e^+e^- centre-of-mass energies. We required that the central detector and the calorimeters were fully operational. Temporary, local inefficiencies in the presampler chambers were monitored and taken into account.

To study the properties of the background we used Monte Carlo events produced with the parton shower generators JETSET 7.4 [22] (3.9 million events) and HERWIG 5.8 [23](1.1 million events) with generator parameters given in [24]. We also used samples of events with only single photon or a π^0 meson present in the detector. The Monte Carlo (MC) samples

[1] In the OPAL coordinate system the x axis points towards the centre of the LEP ring, the y axis points upwards and the z axis points in the direction of the electron beam. The polar angle θ and the azimuthal angle ϕ are defined with respect to the z and x-axes, respectively, while r is the distance from the z-axis.

were passed through the full simulation of the OPAL detector [25] and subjected to the same reconstruction and analysis procedure as the data.

The difficulty in the measurement of prompt photon production lay in the separation of the signal from background. The QCD shower models predicted a signal-to-background ratio of approximately 1/200. Background clusters, with no charged track associated with them, are dominated by photons from decays of hadrons, particularly $\pi^0 \to \gamma\gamma$ ($\simeq$ 57%) and $\eta \to \gamma\gamma$ ($\simeq$ 10%). The other sources of background, like the interaction of neutral hadrons such as $\mathrm{K^0_L}$'s or neutrons in the material of the calorimeter contribute at the level of few per cent each. More than one particle can also contribute to a cluster. The transverse profile of the calorimeter cluster can be used to differentiate between clusters coming from different sources. The hardest to remove are clusters produced by $\pi^0 \to \gamma\gamma$ and $\eta \to \gamma\gamma$ meson decays because they can be very similar, especially for higher cluster energies, to those produced by one photon. An irreducible, but very well predicted, background comes from the initial state radiation (ISR) photons radiated by the beam particles before they interacted.

The tracks and calorimeter clusters were selected as described in [21]. In addition we required a cluster energy to be larger than 10 GeV and cluster polar angle such that $|\cos\theta| < 0.72$. We then applied three cuts motivated by studies with simulated events.

Cut 1. There was required to be no charged track associated with the cluster. Tracks were extrapolated to the calorimeter surface. A track was associated with a calorimeter cluster if it extrapolates to the calorimeter within 24 mrad (approximately half of a lead glass block width) of the centre of gravity of a cluster. In Figure 1a and 1b, the normalised distributions of the angle between the calorimeter cluster and the nearest track are shown for background clusters and prompt photons in the JETSET model, for small and large cluster energies. For the higher energy clusters a contribution from electrons and positrons from conversions of prompt photons in the beam pipe and central detector is seen for small angles δ. For lower energy clusters this effect is diluted both by the greater separation between the extrapolation of the track to the calorimeter surface and the centre of gravity of the calorimeter cluster, and by the presence of other tracks in the proximity of the prompt photon.

Cut 2. We required the presence of a presampler cluster within 24 mrad of the centre of gravity of the calorimeter cluster. The differences in azimuthal $|\Delta\phi|$ and polar $|\Delta\Theta|$ angles between the positions of the calorimeter cluster and the presampler cluster were required to satisfy $\Delta = \min(|\Delta\phi|, |\Delta\Theta|) < \Delta_{\mathrm{cut}}$. The distributions of Δ for single, isolated photons and π^0's in the detector as well as all background clusters in JETSET events are shown in Fig. 2a and 2b. The value of Δ tended to be larger for clusters produced by two overlapping photons, e.g. from a $\pi^0 \to \gamma\gamma$ than for clusters produced by a single photon, because the presampler measures the cluster position at an early stage of development of the electromagnetic shower.

Cut 3. The transverse profile of the cluster in the calorimeter was required to be compatible with that produced by an isolated photon. It was first assumed that the cluster was produced by a photon. The impact point of the photon was varied until the best description of the observed lateral shower profile by a reference profile was found. The reference profile was obtained by the parametrisation of the results of MC simulation of the isolated photon in the detector. The fast algorithm described in section 4.1 of [26] was used. The resulting variable, S, is proportional to the χ^2 for matching the measured and predicted energy sharing between the calorimeter blocks. The distributions of S for single photons, single π^0's and background clusters in JETSET events are shown in Fig. 2c and 2d.

The energy dependent values of Δ_{cut} and S_{cut} were chosen to optimise the separation of signal from background and are shown in Tab. 1.

In total 23106 clusters passed the selection procedure. The signal-to-background ratio, estimated from simulation, was improved from 1/200 to 1/130 after cut 1, then to 1/50 after cut 2 and finally to 1/6 after cut 3 with respectively 92%, 51% and 42% of the signal retained. According to the simulation the background clusters passing the selection were produced mostly by $\pi^0 \to \gamma\gamma$ ($\simeq$79%) and $\eta \to \gamma\gamma$ ($\simeq$11%) decays.

The cuts lead to a strong reduction of the background while prompt photons were much less affected. The fraction of photons rejected by the cuts can be corrected for and its knowledge will affect the systematic uncertainty of the measurement as detailed in Sections 5 and 6. The efficiency of cuts 2 and 3 for photon clusters well separated from other particles in the event was determined directly from the data using a sample of photons in radiative lepton pair events $\mathrm{e^+e^-} \to \ell^+\ell^-\gamma$ ($\ell = \mathrm{e}, \mu$). In addition, in hadronic events, further losses of prompt photons occur when other particles hit the calorimeter close to the photon. This can result in a cluster being associated to a track or being sufficiently distorted to fail the selection criteria. A small fraction ($\simeq$ 6%) of prompt photons also converted in the beam pipe or central detector. The correction for these effects was estimated using Monte Carlo as detailed in sec. 5. In Figure 1 we compare the data and MC distributions for the angle δ between the cluster and track closest to it. Differences between data and Monte Carlo are concentrated in the region of small angles, below our cut value of 24 mrad, a region containing only a small fraction of prompt photons. In Figure 3 we present data and MC distributions for the variables S and Δ used in cuts 2 and 3. We compare distributions for photons from radiative lepton events with MC for single, isolated photons (a and c). We also show distributions for data and MC clusters from the $\tau^+\tau^-$ events ($\tau^\pm \to \rho^\pm\nu_\tau$ and $\rho^\pm \to \pi^\pm\pi^0$, section 4.2 of [26]) (b and d). The differences between data and MC are small and concentrated mostly in the region below our cut values, so they will cause only small systematic effects.

4 The determination of the number of photons in the selected sample

We determined the fraction of photons in the sample remaining after cuts 1-3 above using the cluster shape fit variable C used in the previous OPAL studies of photon production [11, 16]. The fit algorithm applied was more sophisticated than that used in the the calculation of the variable S. The variable C had a better background rejection power than S. Due to the similar fit algorithms and shower parametrisations the C and S variables are correlated, although not fully. The definition of C is

$$C = \frac{1}{N_b} \sum_i \frac{(E_i^{\mathrm{pred}} - E_i^{\mathrm{obs}})^2}{(\sigma_i^{\mathrm{pred}})^2 + (\sigma_i^{\mathrm{obs}})^2}, \tag{1}$$

where: E_i^{obs} is the energy observed in calorimeter block number i; E_i^{pred} is the predicted energy in calorimeter block number i; σ_i^{pred} and σ_i^{obs} are the energy dependent errors on E_i^{pred} and E_i^{obs}, respectively; and N_b is the number of blocks in the cluster. E_i^{pred} was taken from the best fit of the shower profile parametrisation, assuming that the cluster was produced by a isolated photon, to the observed energy sharing between the calorimeter blocks. The reference profiles varied as a function of $\cos\theta$ because of the varying amount of material in front of the calorimeter.

We fitted the distribution of C in the data with a linear combination of MC distributions for photons and background for clusters passing the same selection criteria as the data:

$$C_{\mathrm{fit}} = fC_\gamma + (1-f)C_{\mathrm{bkg}}. \tag{2}$$

The fraction f of photons in the selected sample was the fit parameter. For the distribution of the background, C_{bkg}, clusters from JETSET hadronic Z^0 events were used, where initial state radiation and prompt photons were removed from the sample. For the distribution of the photons, C_γ, we used a simulated sample of isolated photons. The fit for f to the C variable distribution in the data was performed separately in seven bins of cluster energy as shown in first column of Tables 1 and 2. The C_γ and C_{bkg} distributions had only a small dependence on cluster energy within a given bin. A binned maximum likelihood method [27] was used to fit the C variable distribution between 0 and 5.

Since we used simulated distributions for C_γ and C_{bkg} in eq. 2, it is crucial to check that the C variable is well described in the simulation. We show data and MC distributions for photons (from radiative lepton pair events and single, isolated photon MC) in Fig. 3e and clusters from data and MC τ decays ($\tau^\pm \to \rho^\pm \nu_\tau$ and $\rho^\pm \to \pi^\pm \pi^0$, sec. 4.2 in [26]) in Fig. 3f. The simulation describes well the C variable distributions for clusters produced by isolated photons as well as by π^0's. In Fig. 5 we show distributions of the C variable for background clusters from different sources in the JETSET simulation. The shapes of the distributions are similar, although $\pi^0 \to \gamma\gamma$ and $\eta \to \gamma\gamma$ decays tend to produce less clusters with higher values of C than other sources of background.

The fitted fraction of photons f is shown in Tab. 2 for different cluster energy ranges. The comparison of the data and fit results is given in Fig. 4. The contributions from prompt photons and background to the fit are shown. The χ^2, taking into account the statistical errors on the C_γ and C_{bkg} distributions, were between 16 and 37 for 23 degrees-of-freedom.

5 Corrections for efficiency, acceptance and initial state radiation

The energy spectrum of prompt photons obtained in the previous section was corrected for photons lost in the selection process and outside the geometrical acceptance. To ensure that the energy spectrum of photons was not biased, efficiency corrections were determined separately for each energy bin. Then the contribution due to initial state radiation was subtracted and the corrected photon energy spectrum was normalised to the total number of hadronic events.

We applied corrections for the following effects.

1. Local, temporary inefficiencies in the presampler system. This factor was determined from data to be 1.23 with negligible statistical error.
2. Rejection of real photons by cuts 2 and 3. A correction was determined using the data sample of photons in radiative lepton pair events.
3. Additional rejection, by the combined effects of cuts 1, 2 and 3, of photons that formed calorimeter clusters with other particles in the event or converted in the beam-pipe or central detector. This correction, called photon environment in Tab. 1, was determined with the JETSET Monte Carlo as the ratio of the combined cut efficiency for prompt photons where no other particle contributed to the calorimeter cluster, to the efficiency for all prompt photons.
4. The contribution of initial state radiation, estimated using the KORALZ program version 4.0, was subtracted.

Values of the energy dependent corrections are shown in Tab. 1. To compare our result with theoretical predictions we applied an additional correction for the rejection of photons by the cut on the polar angle θ ($|\cos\theta| < 0.72$). The correction was 1.58. We assumed the leading-order $1+\cos^2\theta$ dependence of the photon production cross-section, which was consistent with the polar angle distribution of the selected calorimeter clusters. The fully corrected energy spectrum of prompt photons in hadronic Z^0 decays is shown in Tab. 2 and in Fig. 6.

6 Systematic effects

We checked the dependence of our result on possible deficiencies in the simulation of the detector and on the particle composition of the background. The main sources of the systematic uncertainties were estimated as follows:

1. The sensitivity of the fit result to the quality of the MC reproduction of the C variable was determined as follows, separately for each cluster energy bin. The C_γ and C_{bkg} distributions used in the fit were simultaneously scaled by $(1 \pm \alpha)$, where α ($\simeq 4\%$) was

the error on the mean value of C from the Monte Carlo added in quadrature to that from the data. These modified distributions were then used in the fit. The differences between these results and those obtained with unmodified distributions were assigned as the systematic error. These uncertainties, ranging from 10-23% and partly correlated between different photon energy points, were the dominant source of systematic error in our measurement.

2. The fractions of photons in the data obtained from the fit using the background spectrum predicted by the HERWIG MC were in agreement with those obtained using the JETSET MC background. This showed that, within our statistical precision, the result did not depend on the details of the model implementation of the parton shower development and hadronisation processes in the MC generators. The difference between the fraction of prompt photons obtained from the data fitted with the $C_{\rm bkg}$ predicted by the JETSET and HERWIG was assigned as a systematic error.

3. Although we relied on the data as much as possible in the determination of the efficiency corrections we had to resort to MC to estimate how the efficiencies for photons passing the selection cuts were modified by the presence of other particles in the event (correction 3 in sec. 5). The corrections obtained with JETSET and HERWIG models gave results consistent within statistical errors despite the differences in the modelling of the prompt photon radiation, parton shower and hadronisation for the two models. The difference between corrections obtained with the JETSET and HERWIG models was assigned as a systematic error.

In addition several further checks were performed, none of which produced a statistically significant, at the one standard deviation level, difference from the result of the default procedure and were not included in the systematic error.

- In equation 2 the amount of the background in the sample was the fit parameter. Therefore, the fitted value of f was not, to first order, sensitive to the background flux as incorporated in the MC generator. In principle, some sensitivity to the relative fluxes of different background sources remained, since it could change the shape of the C variable distribution of the background. This sensitivity was estimated by repeating the fit with the number of clusters in the background produced by η mesons, the second largest source of background, adjusted by the uncertainty on the η yield of $\pm 20\%$ [29]. This resulted in 1.5% change of the prompt photon yield.

- The systematic error from the modelling of the C variable (point 1) was consistent with an estimate which assumed a dependence of the factor α on the value of C variable. For C greater than 2 the factor α was put to zero. For lower C values a linear dependence was assumed, such that the mean value of α for C between 0 and 2 was equal to the value of α used in the default procedure.

- The fit of equation 2 was repeated with the distribution of C for prompt photons in JETSET used as C_γ instead of the default of isolated photons.

- The values of the cuts $\Delta_{\rm cut}$, $S_{\rm cut}$ and for the association cut were changed by ± 20 %.

- The fit result (section 4) did not depend on the bin size of the fitted distribution (changed by a factor of 2) or the upper fit boundary (moved between 5 and 10 in C).

- The effect of the energy resolution was negligible for the size of the photon energy bins used. The absolute energy scale for the electromagnetic calorimeters is very well calibrated using Bhabha scattering events and cross-checked with, for example, the π^0 mass peak position.

- The background of calorimeter clusters from $\tau^+\tau^-$ events in the selected sample (estimated with KORALZ 4.0 [28]) was less than 0.5% of the observed photon signal.

- The contamination from clusters produced by the LEP accelerator background or cosmic rays, estimated with the events collected using a random beam-crossing trigger, was below the level of 2 clusters per million events and no cluster passed the selection criteria.

To summarise, the principal systematic error comes from the uncertainty on the fitted fraction of photons in the data due to the uncertainty on the quality of the MC reproduction of the C variable (point 1). The uncertainties due to the possible dependence of our result on the modelling of the prompt photon radiation, parton shower and hadronisation in the MC generators were estimated in points 2 and 3. The total systematic uncertainty was calculated as the sum in quadrature of the propagated errors listed in points 1, 2 and 3, as shown in Tab. 3. The systematic errors were partly correlated between the different photon energy ranges.

7 Summary and discussion

The energy spectrum of prompt photons in hadronic Z^0 decays, corrected for the geometrical acceptance as discussed in section 5, is shown in Tab. 2 and in Fig. 6. The inner error bars are statistical, and the outer combined statistical and systematic (added in quadrature). The breakdown of the systematic uncertainties is given in Tab. 3.

To compare our result with leading-order theoretical predictions for the quark-to-photon fragmentation function we use the leading-order cross-section for prompt photon production in $\mathrm{e^+e^-}$ annihilation given by formula (12) in [19]:

$$\frac{1}{\sigma_{\mathrm{had}}}\frac{d\sigma(E_\gamma)}{dE_\gamma} = \frac{4}{\sqrt{s}}\sum_{\mathrm{q}} w_{\mathrm{q}} D_{\gamma/\mathrm{q}}, \tag{3}$$

where E_γ is the photon energy; $\sqrt{s}$ is the $\mathrm{e^+e^-}$ centre of mass energy; w_{q} is the relative contribution from quark flavour q ($w_{\mathrm{q}} = \Gamma_{\mathrm{Z}\to\mathrm{q\overline{q}}}/\Gamma_{\mathrm{Z}\to\mathrm{hadrons}}$); σ_{had} is the production cross-section for $\mathrm{e^+e^-}$ hadronic events, and $D_{\gamma/\mathrm{q}}$ is the quark q to photon fragmentation function. In Fig. 6 we plot the QCD prediction using an asymptotic leading-order $D_{\gamma/\mathrm{q}}$ [6, 7] as parametrised by Duke and Owens [8] with $Q^2 = M_Z^2$ and $\Lambda = 0.2$ GeV. In Fig. 6 we also plot the QCD predictions using $D_{\gamma/\mathrm{q}}$ from leading-order (LO) calculations by Glück, Reya and Vogt [9]. We also show the higher-order (HO) prediction for the prompt photon production cross-section by Glück, Reya and Vogt [9] and the beyond leading logarithm (BLL) prediction by Bourhis, Fontannaz and Guillet [10]. The factorisation schemes used were DIS_γ and $\overline{\mathrm{MS}}$ respectively. The authors of [9] and [10] included in their calculations non-perturbative effects through the vector-meson dominance ansatz, although using different experimental inputs.

Our data are in agreement with these theoretical predictions. The experimental precision is not sufficient to discriminate between them. The ALEPH data on the production of jets containing a photon carrying a substantial fraction (above 70%) of the jet energy (Fig. 4 in [18]) show clear disagreement with the Duke-Owens parametrisation [8]. It was noted in [10] that a possible reason is that the ALEPH measurement is restricted to only part of the phase-space for photon production in hadronic Z^0 decays.

To summarise, we have measured the inclusive production of prompt photons with energy above 10 GeV in hadronic Z^0 decays. Good agreement is found with current QCD predictions for the quark-to-photon fragmentation function.

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Table 1: Energy dependent cuts and corrections to the photon energy spectrum. Corrections for photon efficiency and environment are multiplicative, ISR is subtracted.

E, GeV	$\Delta_{\rm cut}$ mrad	$S_{\rm cut}$	isolated photon efficiency	photon environment	N_γ^{ISR}/GeV/10^6 events
$10-15$	3.	1.5	2.51 ± 0.13	1.34 ± 0.10	6.71 ± 0.13
$15-20$	4.	2.5	2.06 ± 0.12	1.36 ± 0.11	2.62 ± 0.08
$20-25$	5.	4.	1.51 ± 0.07	1.37 ± 0.10	1.52 ± 0.06
$25-30$	5.	5.	1.47 ± 0.08	1.25 ± 0.09	1.27 ± 0.06
$30-35$	5.	5.	1.42 ± 0.08	1.20 ± 0.09	1.28 ± 0.06
$35-40$	5.	5.	1.60 ± 0.09	1.14 ± 0.10	1.20 ± 0.06
$40-45.6$	5.	5.	1.57 ± 0.05	1.03 ± 0.13	2.87 ± 0.09

Table 2: The fraction f of prompt photons in the selected sample and the corrected energy spectrum of prompt photons in hadronic Z^0 events (the first error is statistical, the second systematic).

E, GeV	Mean E, GeV	f, %	N_γ/GeV/10^6 events
$10-15$	12.0	32.7 ± 2.1	$625 \pm 70 \pm 99$
$15-20$	17.3	18.2 ± 1.7	$266 \pm 36 \pm 63$
$20-25$	22.3	12.9 ± 1.4	$142 \pm 20 \pm 37$
$25-30$	27.3	27.9 ± 1.9	$212 \pm 24 \pm 39$
$30-35$	32.3	25.7 ± 2.8	$118 \pm 17 \pm 28$
$35-40$	37.1	29.9 ± 5.0	$75 \pm 15 \pm 20$
$40-45.6$	42.1	80.4 ± 10.2	$48 \pm 9 \pm 8$

Table 3: Contributions to the systematic error on the photon energy spectrum shown in Tab. 2.

Mean E, GeV	N_γ/GeV/10^6 events	Systematic errors on N_γ/GeV/10^6 events		
		description of C variable	model dependence of background distribution	photon acceptance
12.0	$625 \pm 70 \pm 99$	± 60	± 73	± 29
17.3	$266 \pm 36 \pm 63$	± 54	± 31	± 12
22.3	$142 \pm 20 \pm 37$	± 32	± 17	± 6
27.3	$212 \pm 24 \pm 39$	± 28	± 25	± 10
32.3	$118 \pm 17 \pm 28$	± 23	± 14	± 6
37.1	$75 \pm 15 \pm 20$	± 17	± 9	± 4
42.1	$48 \pm 9 \pm 8$	± 5	± 6	± 3

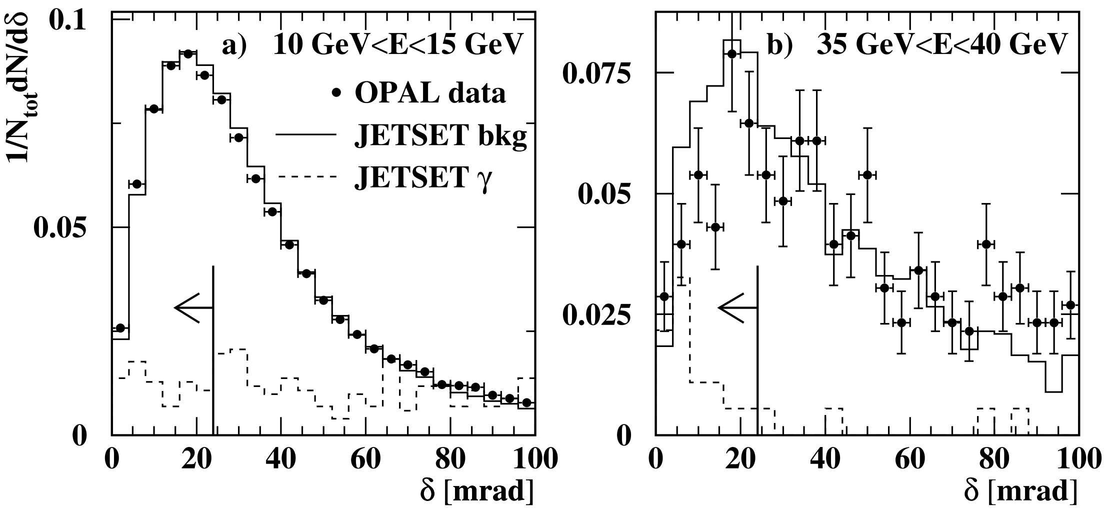


Figure 1: The angle, δ, between the calorimeter cluster and the nearest track for hadronic events: for the data, background clusters in the JETSET simulated event sample and prompt photons in the JETSET sample, for different cluster energies: $10 < E < 15$ GeV a), and $35 < E < 40$ GeV b). Distributions are normalised such that the total yield for each curve is unity. Most of the photon signal is off scale, with $\delta > 100$ mrad. The regions removed by the cut are shown by arrows.

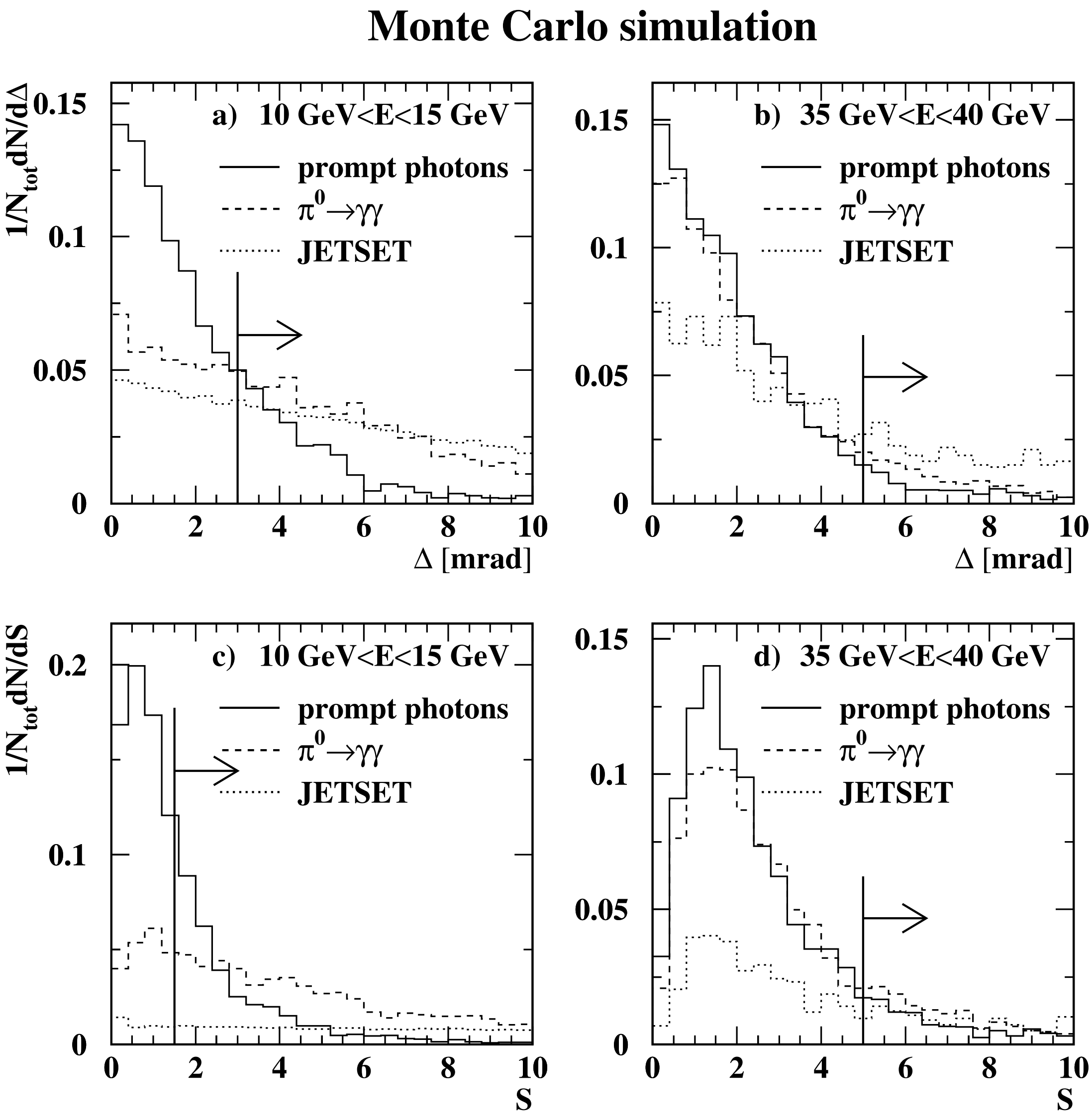


Figure 2: The distributions of the variables Δ (a and b), and S (c and d) for clusters from simulation of single isolated photons and π^0's and for all background clusters in JETSET events (after cut 1), for different cluster energies: $10 < E < 15$ GeV a) and c), and $35 < E < 40$ GeV b) and d). Distributions are normalised such that the total yield for each curve is unity. The regions removed by the cuts are shown by arrows.

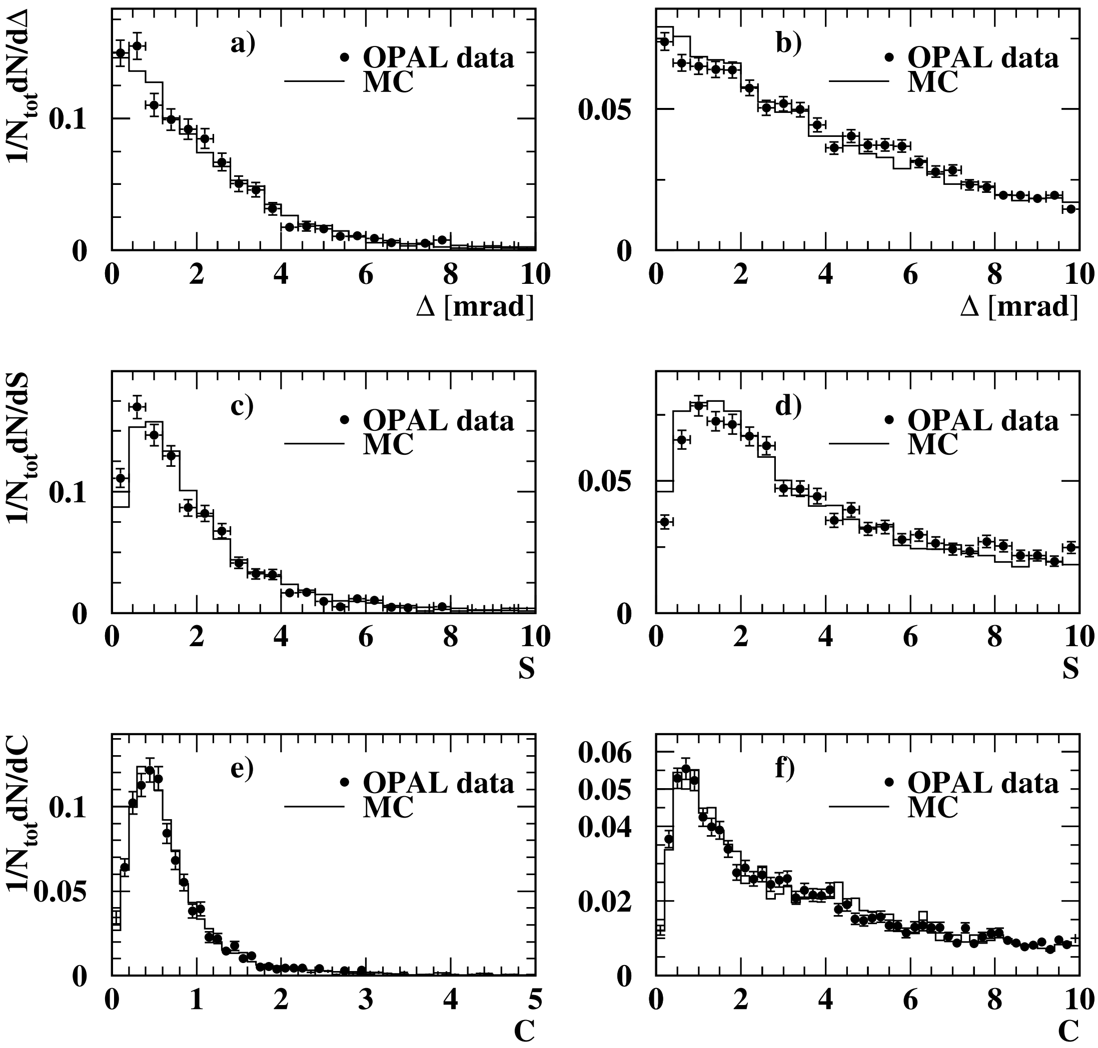
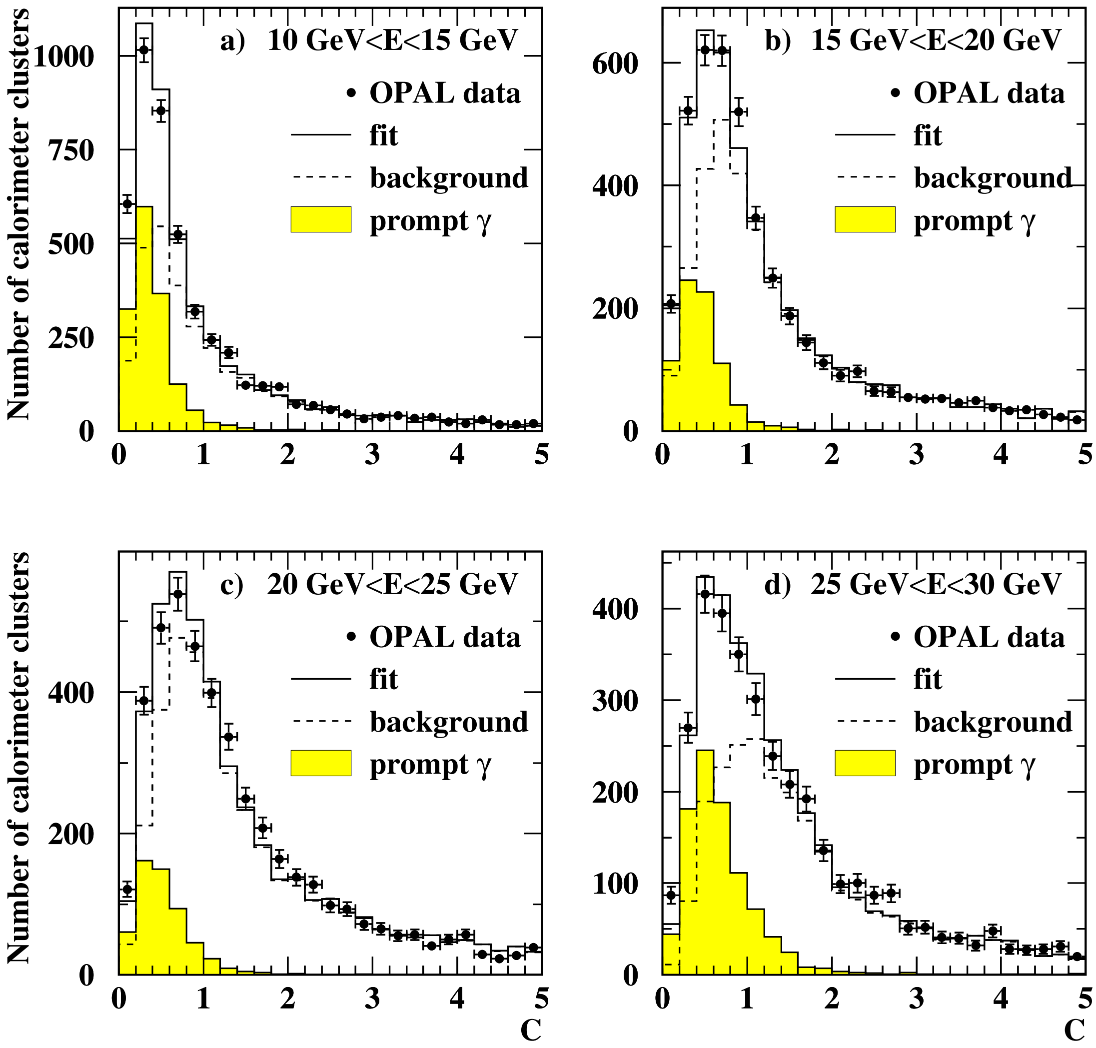


Figure 3: The distributions of the variables Δ (a and b), S (c and d) and C (e and f) for photons in radiative lepton pair data events and simulated single, isolated photons (a, c and e); and π^0's in $\tau^\pm \to \rho^\pm \nu_\tau$, $\rho^\pm \to \pi^\pm \pi^0$ decays (b, d and f). Distributions are normalised such that the total yield for each curve is unity


a) 10 GeV<E<15 GeV
b) 15 GeV<E<20 GeV
c) 20 GeV<E<25 GeV
d) 25 GeV<E<30 GeV
OPAL data
fit
background
prompt γ
Number of calorimeter clusters
C

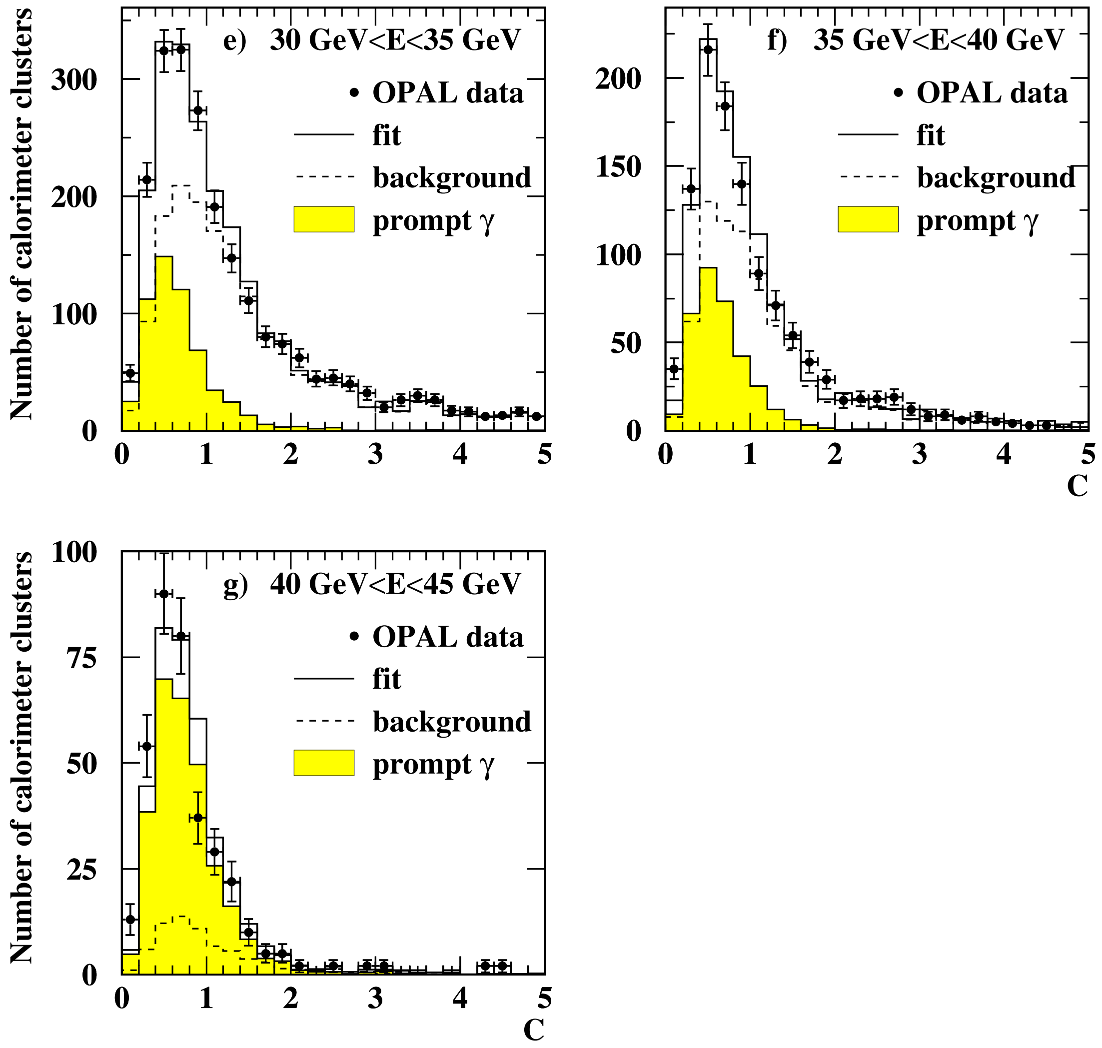


Figure 4: Comparison of the data and fit for different cluster energy ranges. Contributions from prompt photons and background in the fit are shown.

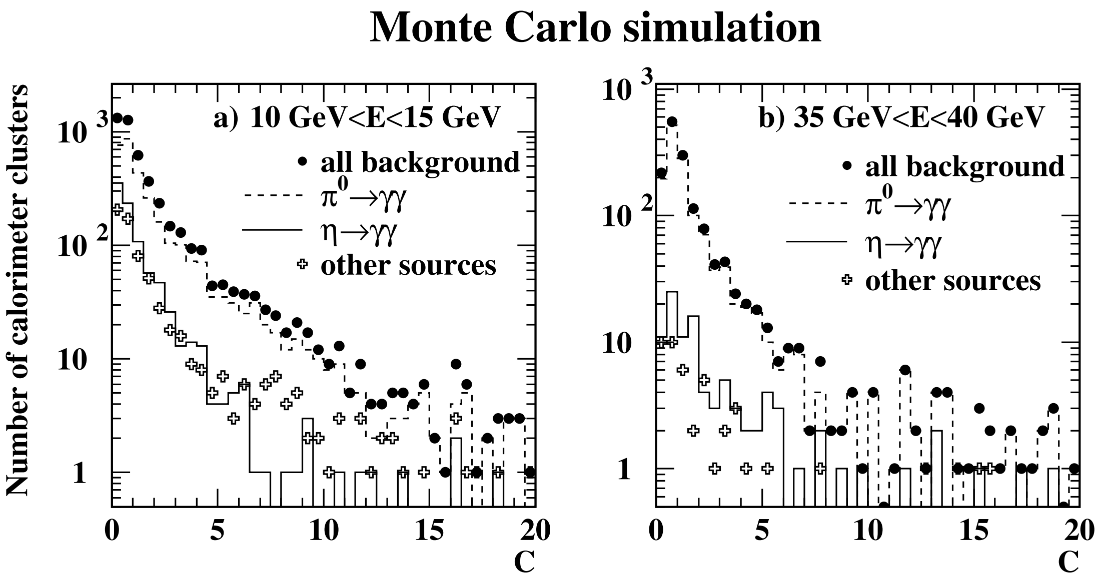


Figure 5: C variable distributions for background clusters from different sources in the JETSET simulation for different cluster energy ranges.

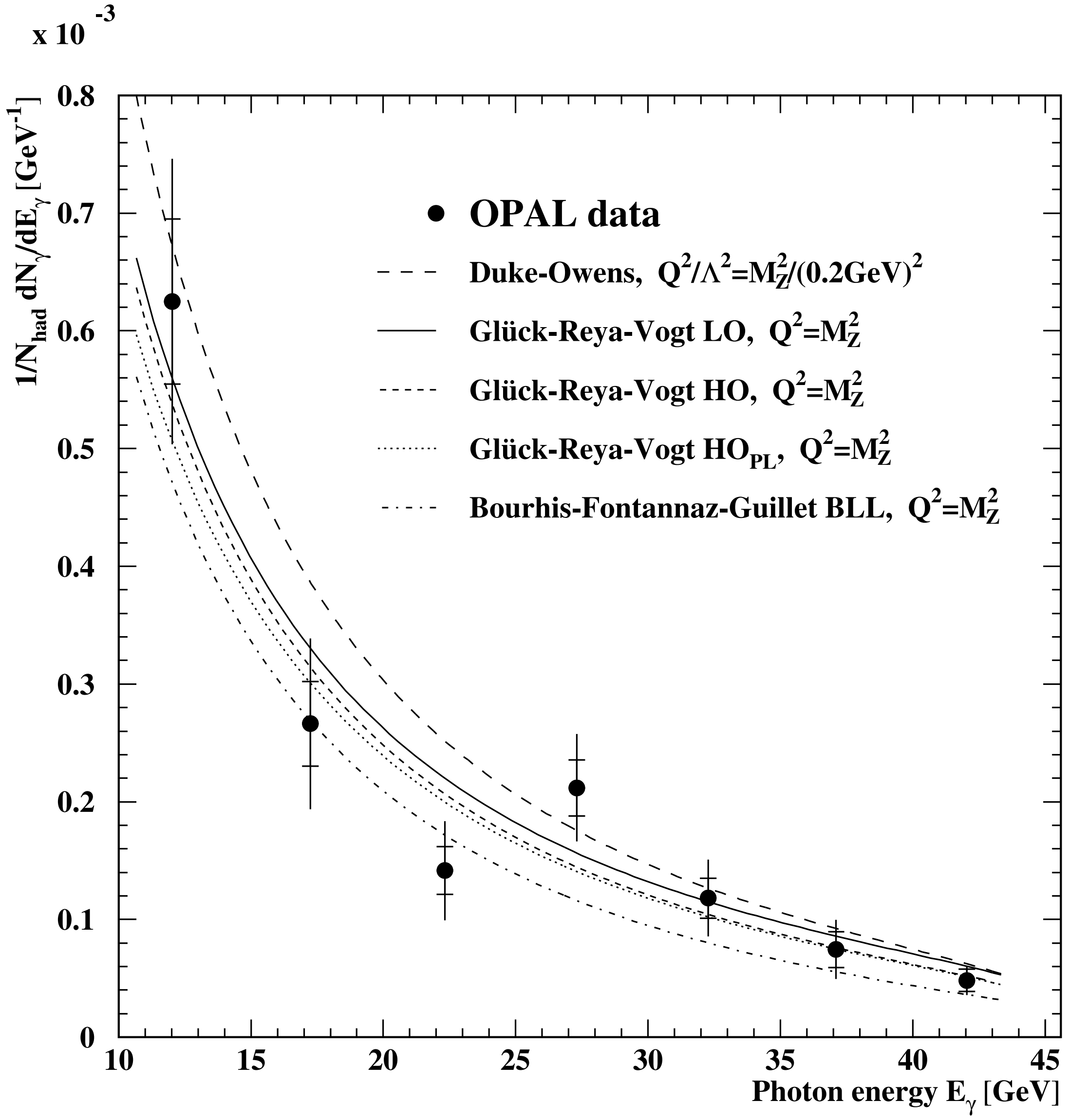


Figure 6: The photon energy spectrum in hadronic Z^0 decays compared to various theoretical predictions: the Duke-Owens parametrisation [8], the Glück, Reya and Vogt predictions including leading-order (LO), higher-order (HO) and higher-order without the non-perturbative corrections (HO_{PL}) [9]. The Bourhis, Fontannaz and Guillet prediction shown include effects beyond leading logarithms (BLL) [10].