Wakefield calculation in the frequency domain



E. Gjonaj

Computational Elektromagnetics Lab, Technische Universität Darmstadt, Germany



DESY, Hamburg

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Motivation



- Long range wakefields
 - Low frequency, long bunches, bunch trains and/or high repetition rate, wall heating
- Approximation of geometry
 - Geometrical details smaller than bunch length, smooth tapering etc.
- Dispersive problems
 - Surface impedance, dielectrics
 - Free-space and waveguide boundary conditions
- Radiation effects
 - Curved beam trajectories and CSR
 - Beams with $\beta < 1$
- Coupler and waveguide signals





The frequency domain problem

$$\nabla \times \mu^{-1} \nabla \times E - k_0^2 \varepsilon E = -jk_0 Z_0 J_s \qquad J_s(x, y, z, \omega) = \rho(x, y) e^{-i\frac{\omega}{\nu}z}$$

• Weak FE formulation: find $E \in H(curl)$ such that:

$$\int dV \mu^{-1} \nabla \times E \cdot \nabla \times v_h - k_0^2 \int dV \varepsilon E \cdot v_h =$$

$$-jk_0 Z_0 \int dV J_s \cdot v_h + \oint_S dS \ n \cdot [v_h \times \mu^{-1} \nabla \times E] \qquad \forall v_h \in H(curl)$$







High-order hierarchic basis functions*



- Allows for simple hp-adaption
- Supports mesh elements of different type + hybrid meshes

*M. Ainsworth, J. Coyle: Int. J. of Numerical Methods in Eng., 2003. *J. Schöberl, S. Zaglmayr: Int. J. Comp. and Math. in Electrical and Electronic Eng., 2005.





Treatment of boundary surfaces

$$\int dV \mu^{-1} \nabla \times E \cdot \nabla \times v_h - k_0^2 \int dV \varepsilon E \cdot v_h =$$

$$-jk_0 Z_0 \int dV J_s \cdot v_h + \int_{S_{SIBC}} dS n \cdot [v_h \times \mu^{-1} \nabla \times E] + \int_{S_{WG}} dS n \cdot [v_h \times \mu^{-1} \nabla \times E]$$
resistive wall in & outgoing pipes

SIBC boundaries

$$\oint_{S_{SIBC}} dS \ n \cdot [v_h \times \mu^{-1} \nabla \times E] = \dots = j \omega \mathbf{Y}_{\mathcal{S}}(\boldsymbol{\omega}) \oint_{S_{SIBC}} dS \ v_h \cdot E$$

Simple modification of the system matrix on SIBC surfaces No fitting of the surface impedance function or ADE/convolution is needed





Treatment of boundary surfaces

$$\int dV \mu^{-1} \nabla \times E \cdot \nabla \times v_h - k_0^2 \int dV \varepsilon E \cdot v_h =$$

$$-jk_0 Z_0 \int dV J_s \cdot v_h + \int_{S_{SIBC}} dS n \cdot [v_h \times \mu^{-1} \nabla \times E] + \int_{S_{WG}} dS n \cdot [v_h \times \mu^{-1} \nabla \times E]$$
resistive wall in & outgoing pipes

Beam pipe boundaries

$$\begin{split} n \times \nabla \times E &= n \times \nabla \times E^{inc} + \sum_{m} a_{m}^{TE} \gamma_{m}^{TE} e_{m}^{TE} + \sum_{m} a_{m}^{TM} \frac{-k_{0}^{2}}{\gamma_{m}^{TM}} e_{m}^{TM} \\ a_{m}^{TE} &= \int_{S_{WG}} dS e_{m}^{TE} \cdot \left[E - E^{inc} \right] \\ \text{Reflection coefficients for each mode} \\ a_{m}^{TM} &= \int_{S_{WG}} dS e_{m}^{TM} \cdot \left[E - E^{inc} \right] \end{split}$$





Beam pipe boundary conditions

$$\int dV \mu^{-1} \nabla \times E \cdot \nabla \times v_h - k_0^2 \int dV \varepsilon E \cdot v_h + \sum_m P_m^{TE}(E) + \sum_m P_m^{TM}(E) = -jk_0 Z_0 \int dV J_s \cdot v_h + \oint_{S_{WG}} dS n \cdot [v_h \times \mu^{-1} \nabla \times E^{inc}] + \sum_m U_m^{TE} + \sum_m U_m^{TM}$$

with $P_m^{TE}(E) = -\gamma_m^{TE} \left(\int_{S_{WG}} dS v_h \cdot e_m^{TE} \right) \left(\int_{S_{WG}} dS e_m^{TE} \cdot E \right), \qquad P_m^{TM}(E) = \cdots$

and matrix representation (TE):

$$P_m^{TE}(E) \rightarrow P_m^{TE} \cdot \mathbf{e} = -\gamma_m^{TE} \mathbf{R}^T \cdot \mathbf{M}_m^{TE} \cdot \mathbf{R} \cdot \mathbf{e} \qquad [\mathbf{R}]_{ij} = \int_{S_{WG}} dS \, \varphi_i^{2D} \cdot \varphi_j^{3D}$$
$$\underbrace{\mathbf{M}_m^{TE} = \mathbf{e} \, {}_m^{TE} \otimes \mathbf{e} \, {}_m^{TE}}_{m} \text{ dense modal dyadic} \qquad 3D\text{-to-2D projection matrix}$$





- Beam pipe boundary excitation
- For an ultra-relativistic bunch (same idea for $\beta < 1$):

2D-electrostatic problem at both ends of the pipe

- Modal contribution to the RHS

$$U_m^{TE}(E^{inc}) = -\gamma_m^{TE} \left(\int_{S_{WG}} dS \, v_h \cdot e_m^{TE} \right) \left(\int_{S_{WG}} dS \, e_m^{TE} \cdot E^{inc} \right)$$

$$U_m^{TE}(E^{inc}) \rightarrow U_m^{TE} \cdot \mathbf{e}^{inc} = -\gamma_0^{TE} \mathbf{R}^T \cdot \mathbf{M}_m^{TE} \cdot \mathbf{R}^{2D} \cdot \mathbf{e}^{inc}$$

...do this for all waveguide modes supported in the pipe





Collimator – hybrid meshes







Collimator – impedance



































XFEL bunch compressor







XFEL bunch compressor





Discussion & Outlook



- The frequency domain approach
 - Purpose: Fill the gap for some important wakefield/impedance problems
 - Complicated geometry
 - Long range wakefields Joule losses
 - Resistive, rough surfaces, dispersive materials, waveguide openings
 - Curved beam trajectories and CSR. Validation of other CSR approaches
 - Status: Implementation of main code finished
 - Mixed, high-order elements, parallelization,...
 - Waveguide operators
 - ToDo: Enable larger problems and faster solutions
 - Domain decomposition based solver, multigrid, ...
 - Fast spectral evaluation by model order reduction
 - Limitation: Huge size of discrete problem for ultra-high frequencies
 - Estimated upper limit using proper solvers and computing power ~100 GHz



Thank You for your attention