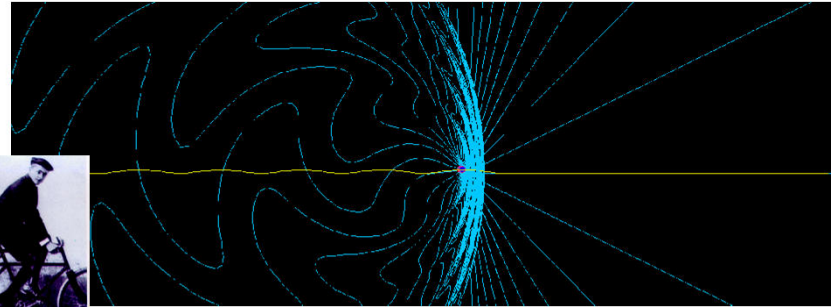
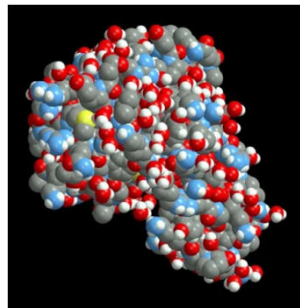
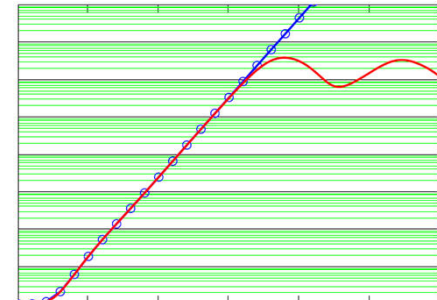
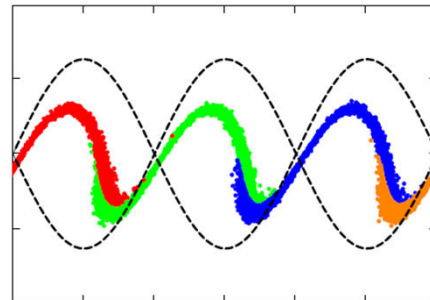


Free-Electron Laser

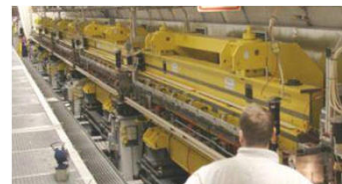
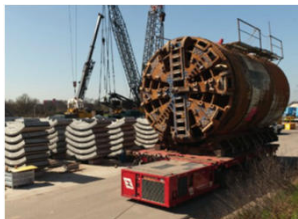
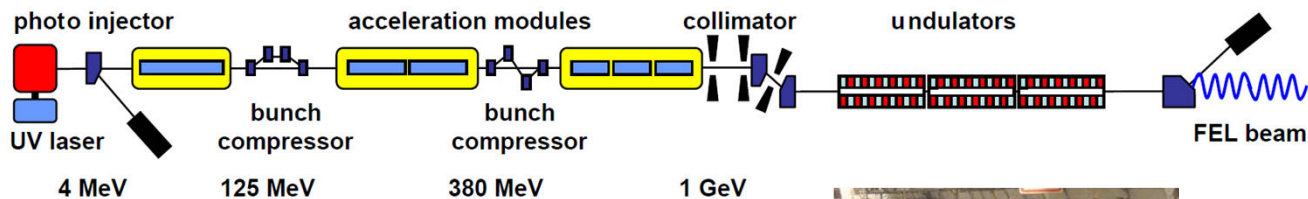
A) Motivation and Introduction



B) Theoretical Approach



C) Experimental Realization / Challenges



A) Motivation and Introduction

need for short wavelengths

why FELs?

free electron \leftrightarrow wave interaction

micro-bunching

amplifier and oscillator

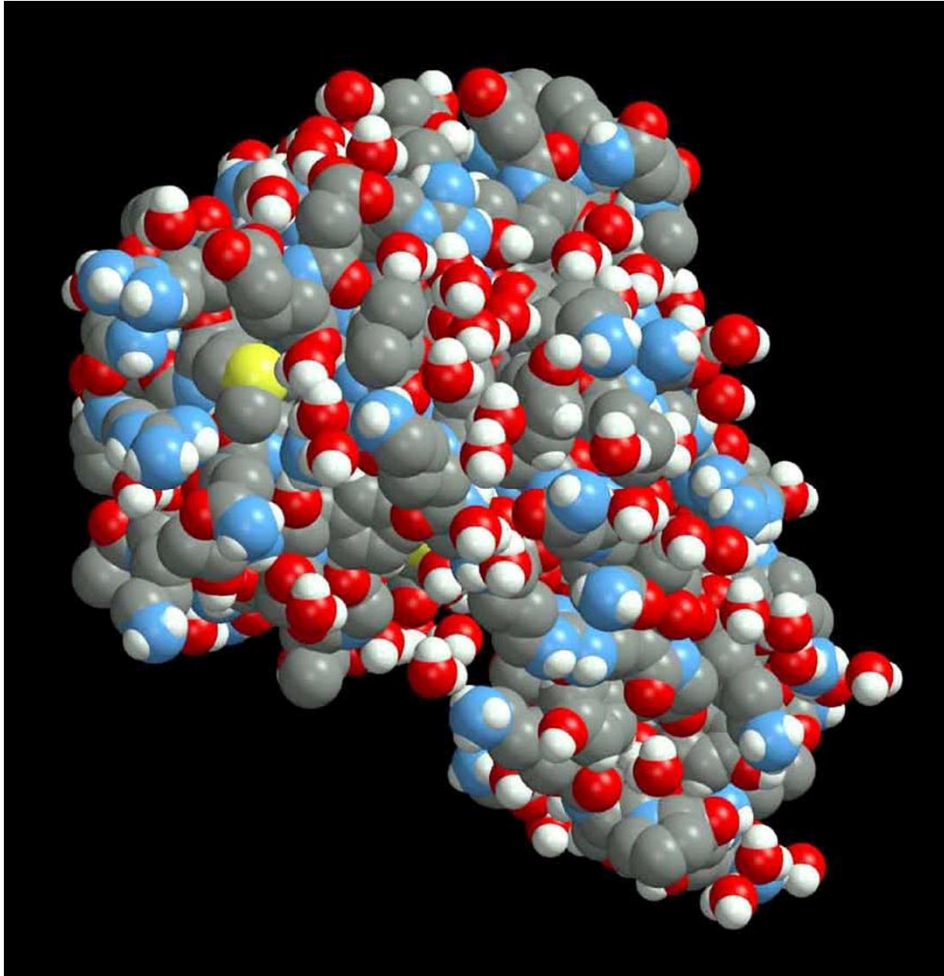
self amplifying spontaneous emission (SASE)

why SASE?

coherent radiation



need for short wavelengths

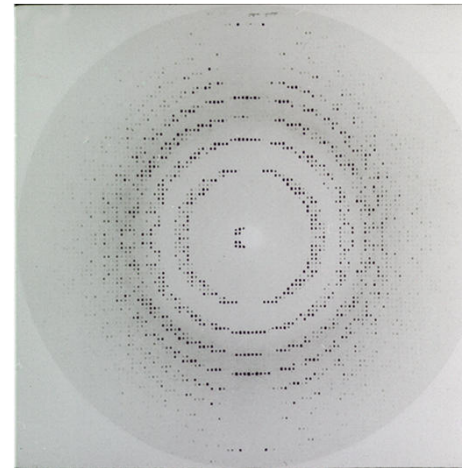


LYSOZYME MW=19,806

state of the art:

structure of biological macromolecule

reconstructed from diffraction
pattern of protein crystal:



needs $\approx 10^{15}$ samples

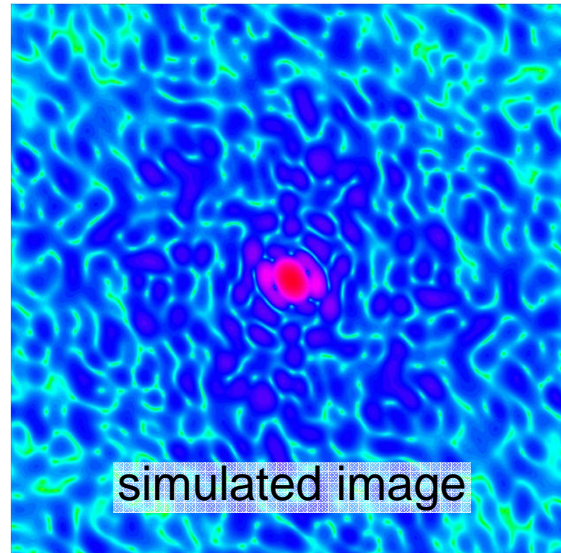
crystallized \rightarrow not in life environment

the crystal lattice imposes
restrictions on molecular motion

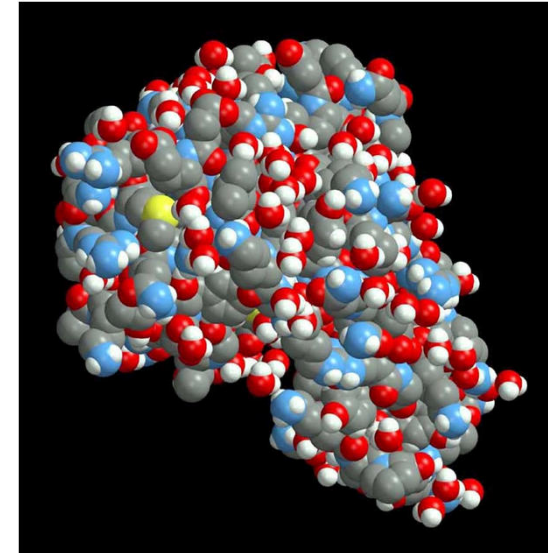


need for short wavelengths - 2

SINGLE MACROMOLECULE



courtesy Janos Hajdu



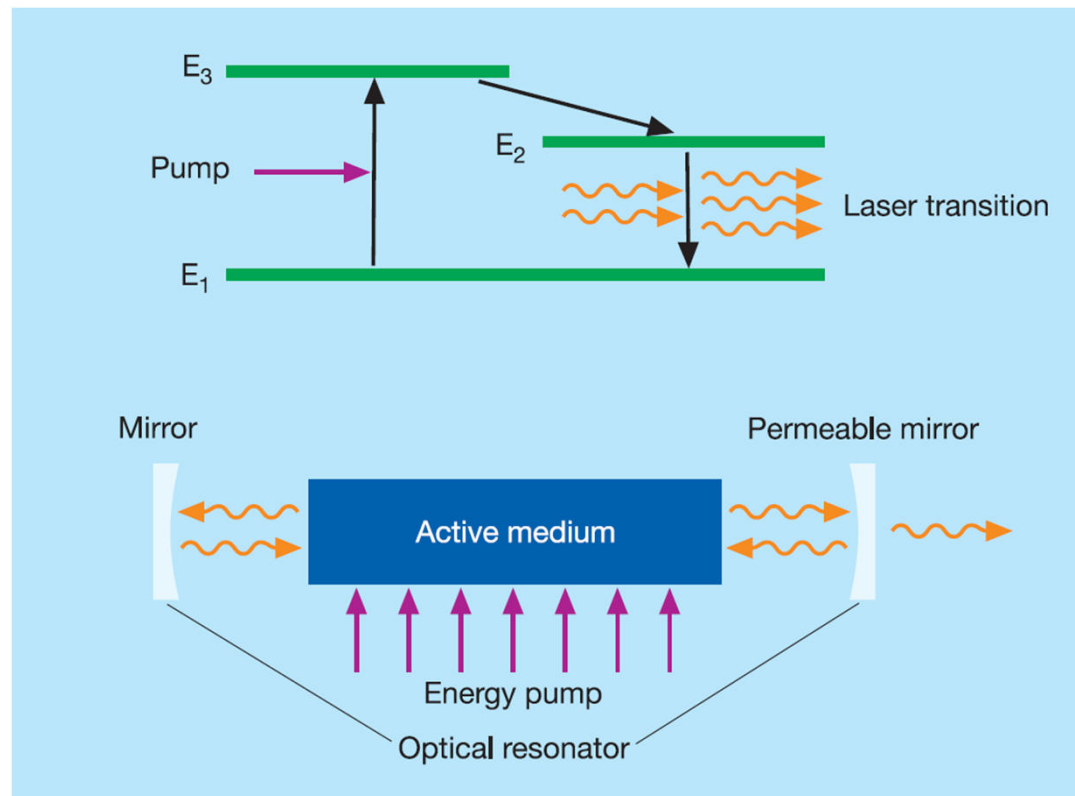
resolution does not depend on sample quality
needs very high radiation power @ $\lambda \approx 1\text{\AA}$
can see dynamics if pulse length < 100 fs

- we need a radiation source with
- very high peak and average **power**
 - **wavelengths** down to atomic scale $\lambda \sim 1\text{\AA}$
 - spatially coherent
 - monochromatic
 - fast tunability in wavelength & timing
 - sub-picosecond **pulse length**



why FELs?

principle of a quantum laser



problem & solution:

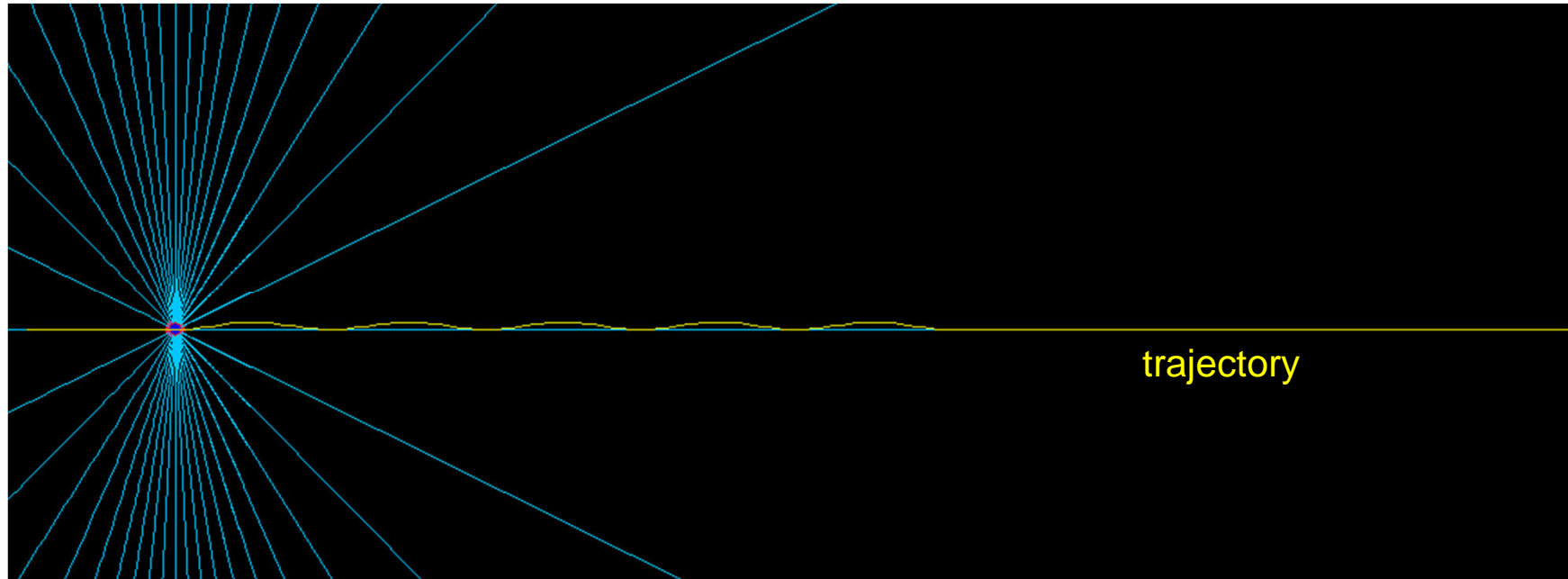
active medium \rightarrow free electron – EM wave interaction



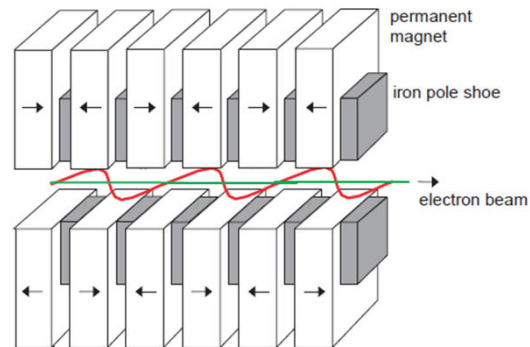
free electron \leftrightarrow wave Interaction

free electron in uniform motion + electric field lines

(before undulator)



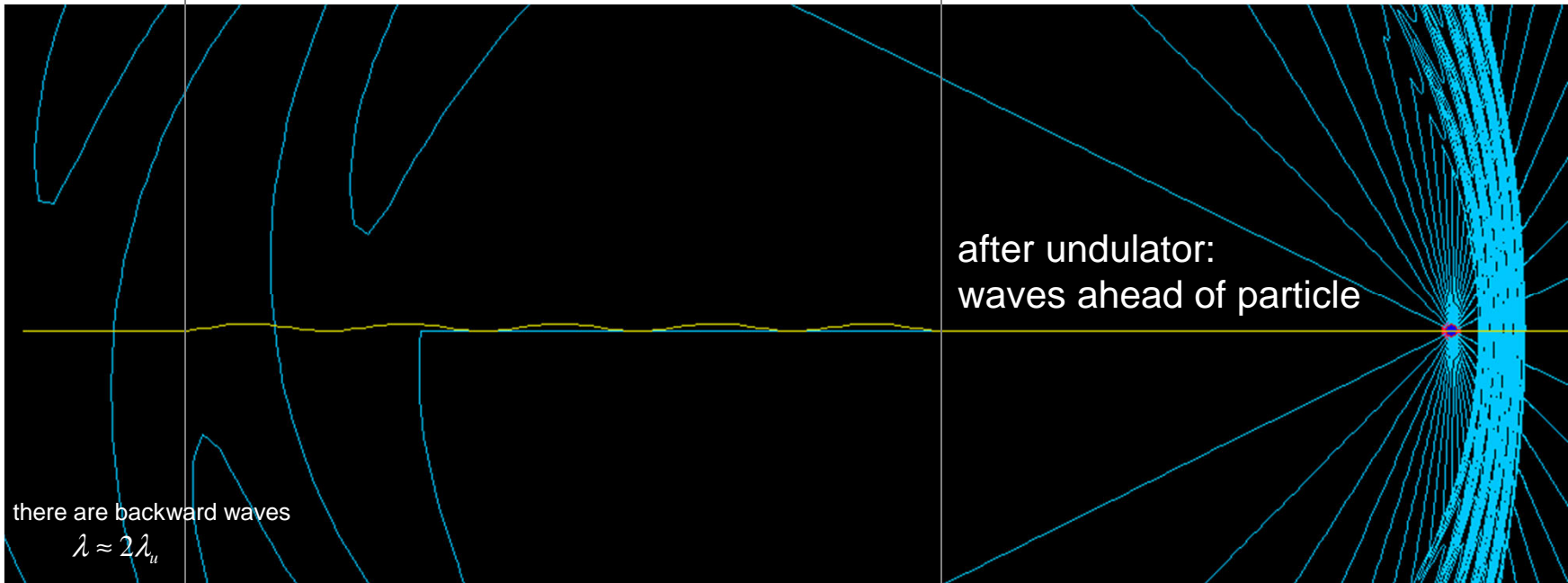
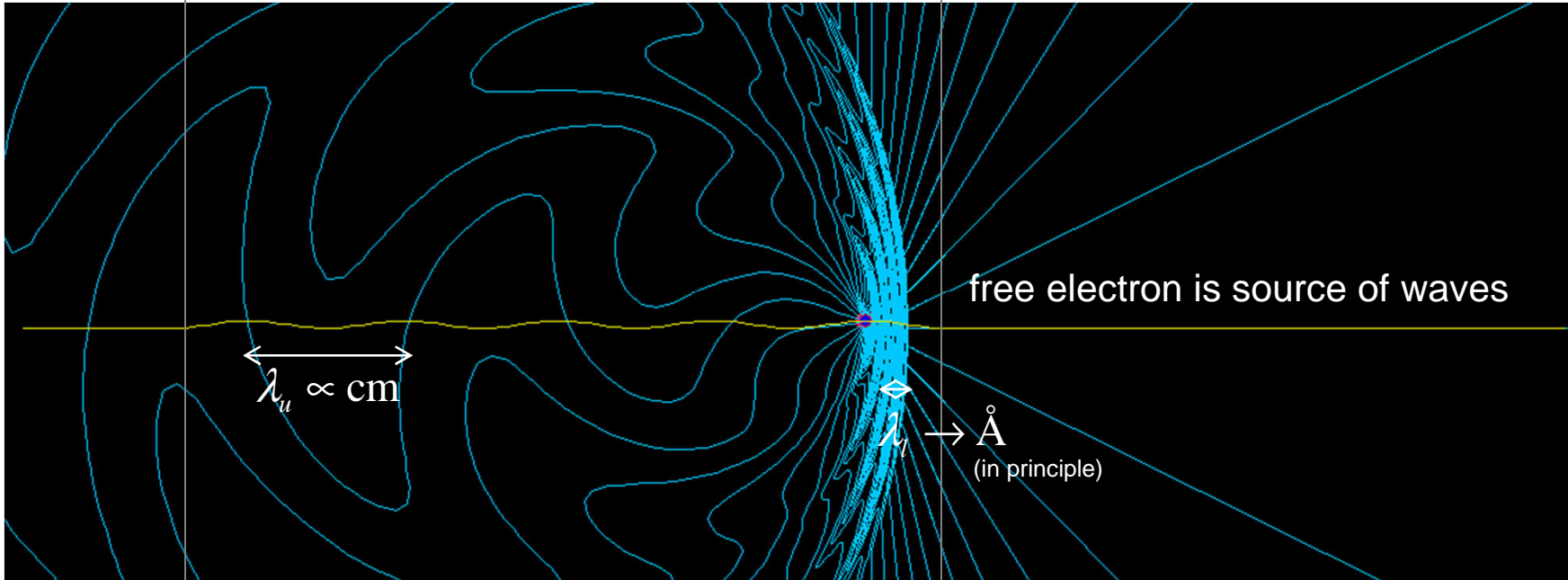
trajectory



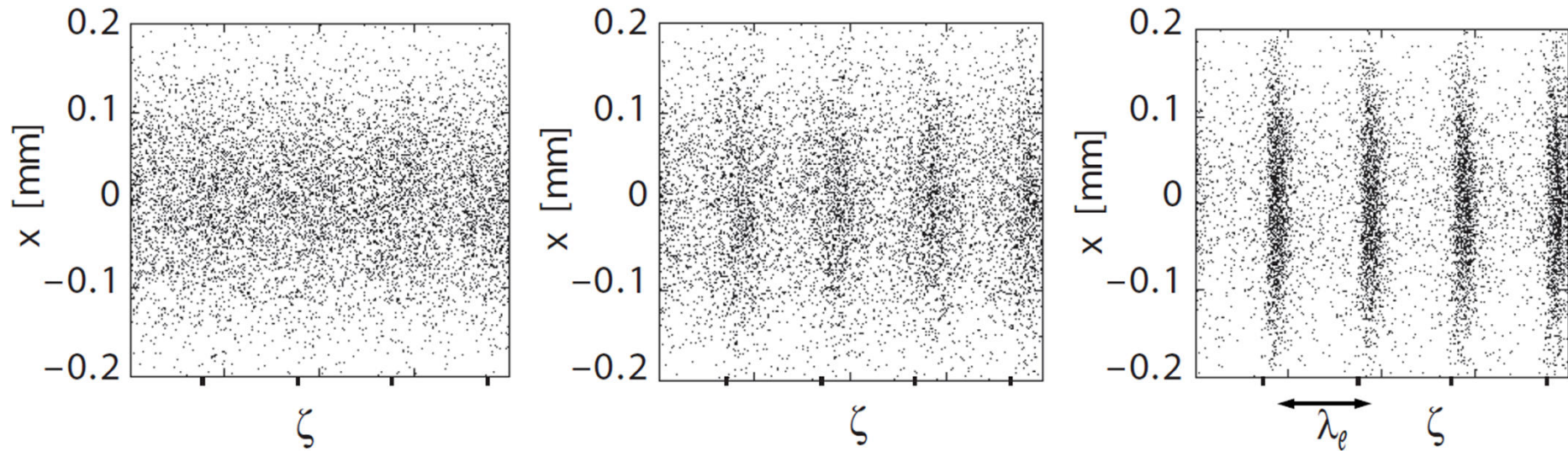
undulator

courtesy T. Shintake
<http://www.shintakelab.com/en/enEducationalSoft.htm>





micro-bunching



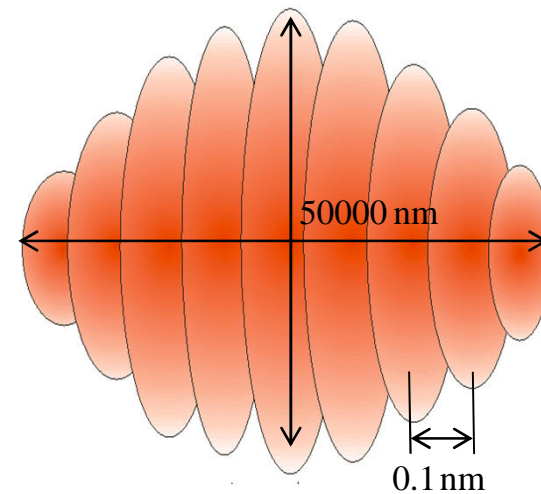
longitudinal motion to 1st order is trivial, **but**

micro-bunching is a 2nd order effect

→ **coupled theory** of particle motion and wave generation

transverse bunch structure is much larger than longitudinal sub-structure

→ **1d theory** with plane waves

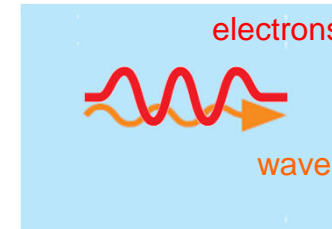
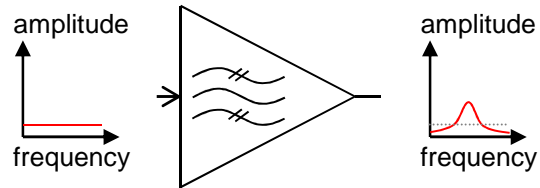


amplifier and oscillator

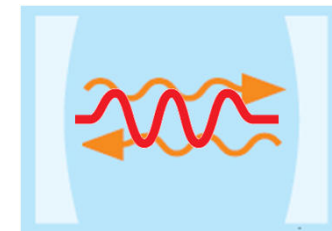
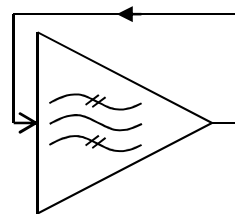
in principle

FEL

amplifier:

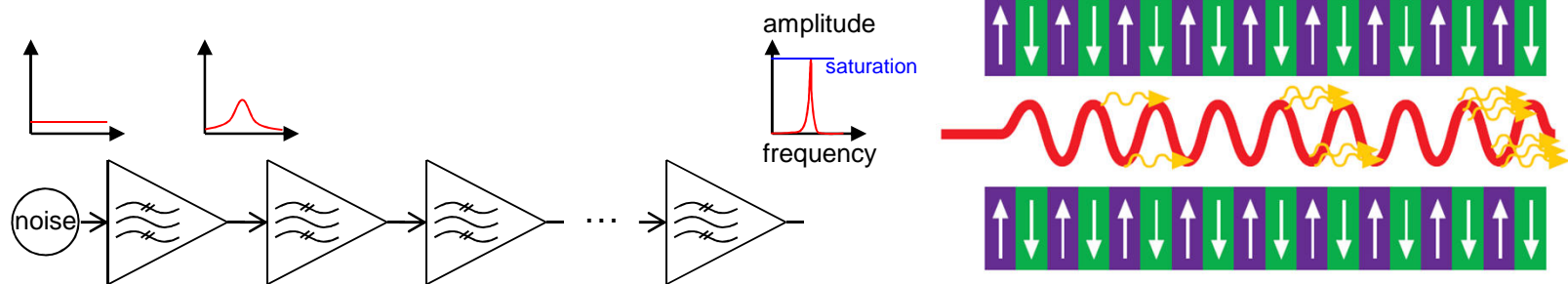


oscillator:

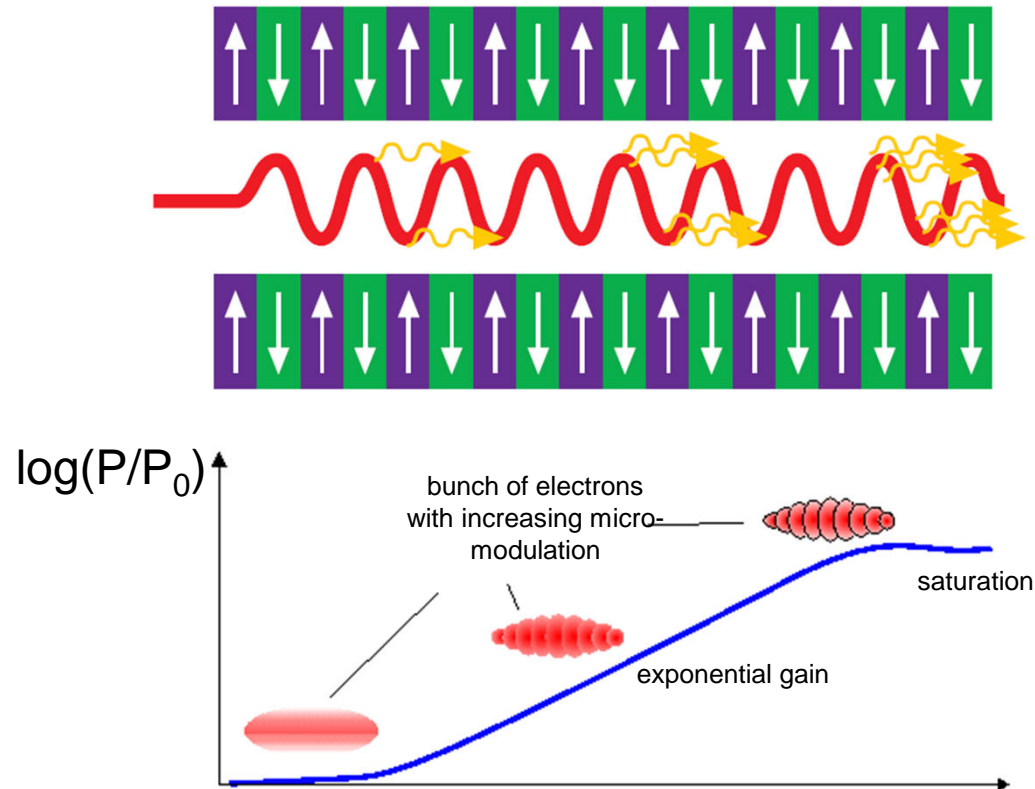


instability, driven by noise, growth until amplifier saturates

amplified noise:



self amplifying spontaneous emission (SASE)



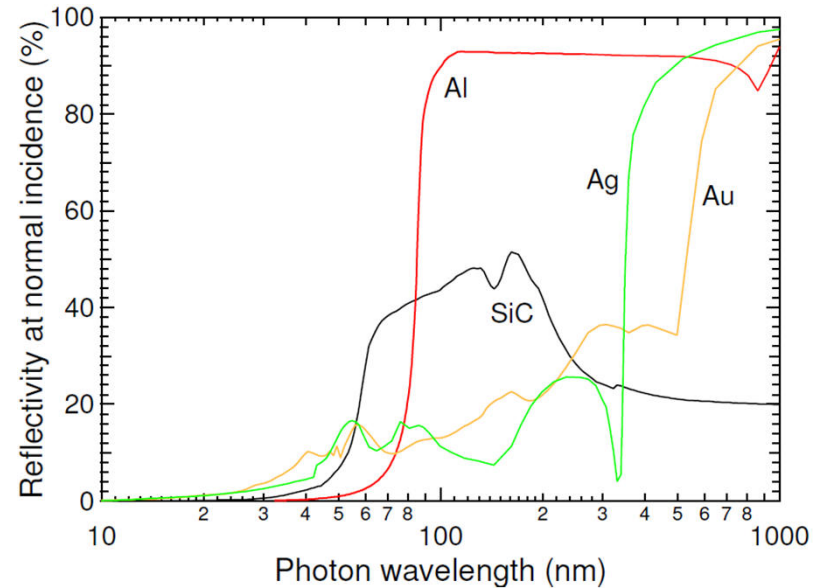
uniform **random** distribution of particles at entrance
incoherent emission of EM waves (noise, wide bandwidth)
amplification (\rightarrow resonant wavelength, micro-bunching)
saturation, full **micro modulation**, **coherent** radiation



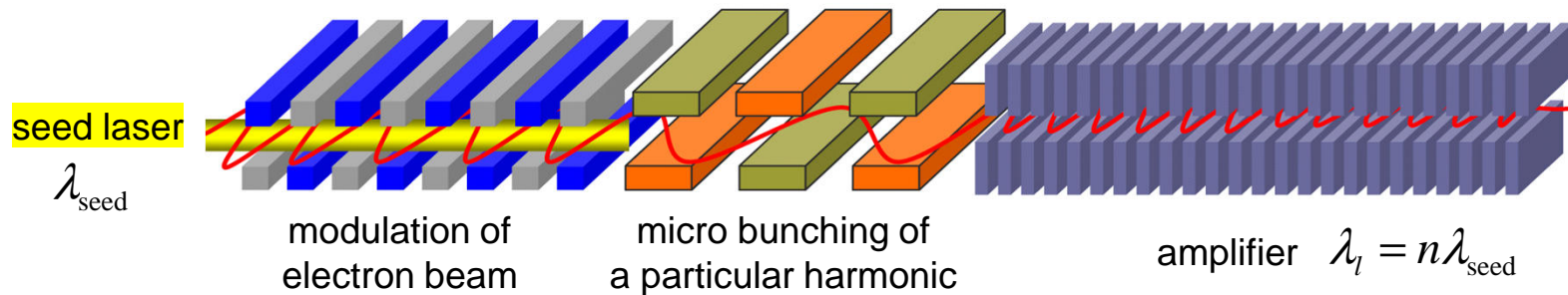
why SASE?

oscillator needs resonator

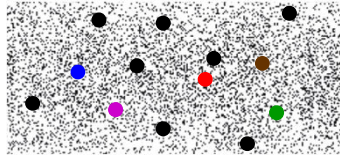
but there are no mirrors for wavelengths < 100 nm



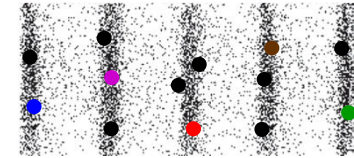
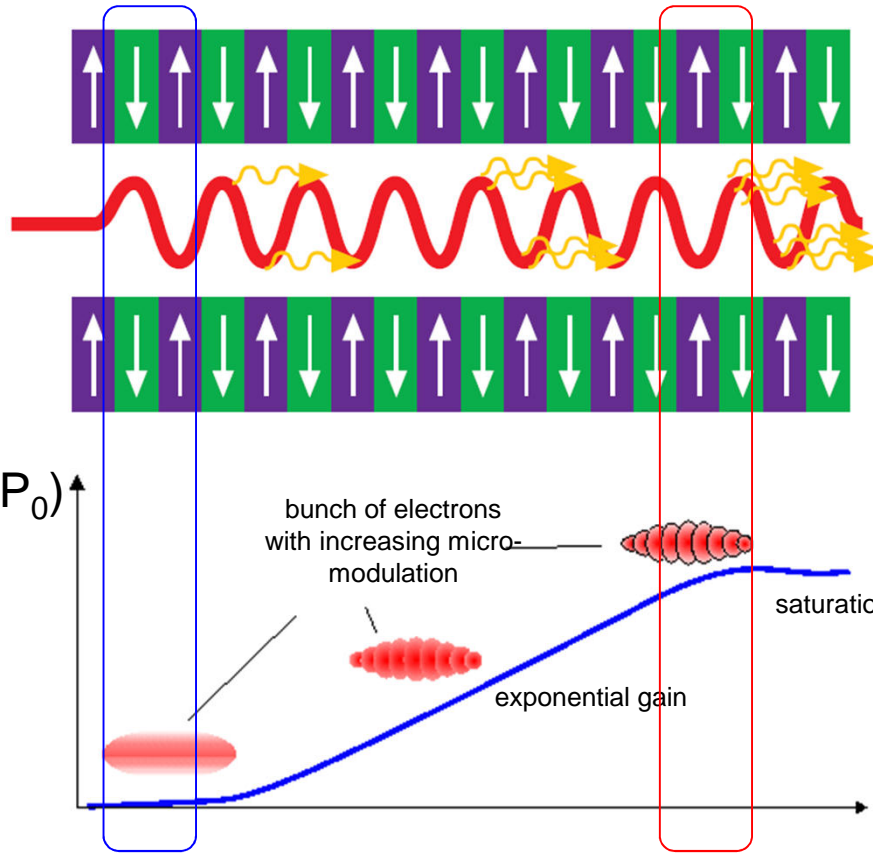
alternative: seed laser + harmonic generation + amplifier



coherent radiation



$\varphi_1 \varphi_2 \varphi_3$



$\varphi_1 \varphi_2 \varphi_3$

$\log(P/P_0)$

bunch of electrons
with increasing micro-
modulation

saturation

exponential gain

no micro-bunching
(only shot noise)

$$P_0 \propto N$$

saturation:
full micro-bunching

$$P \propto N^2$$

with $N =$ particles per λ_l



B) Theoretical Approach

particle dynamics: undulator motion

about particles

independent parameter: z

particle dynamics: interaction with EM wave

longitudinal equation of motion

phase space and pendulum

FEL low gain theory

micro bunching

electrodynamics (1D)

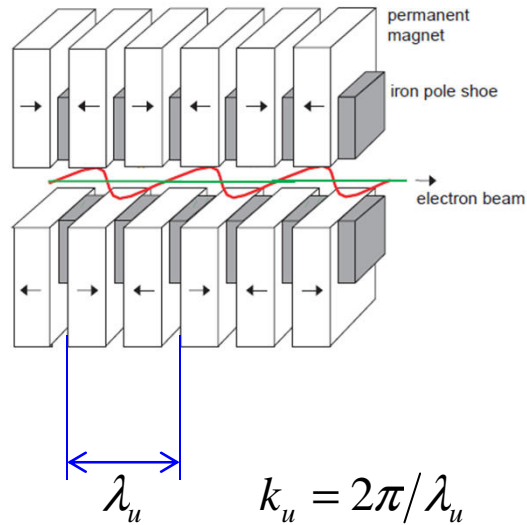
FEL high gain theory (1D)

continuous phase space: Vlasov equation

FEL third order equation



particle dynamics: undulator motion



field of planar undulator $\mathbf{B}_u = -B_0 \mathbf{e}_y \sin k_u z$

equation of motion $\gamma m_e \dot{\mathbf{v}} = -e \mathbf{v} \times \mathbf{B}$

approach: $z \approx \bar{v}t - \hat{z} \cos 2\omega_u t$ with $\omega_u = 2\pi\bar{v}/\lambda_u$

$x \approx \hat{x} \sin \omega_u t$

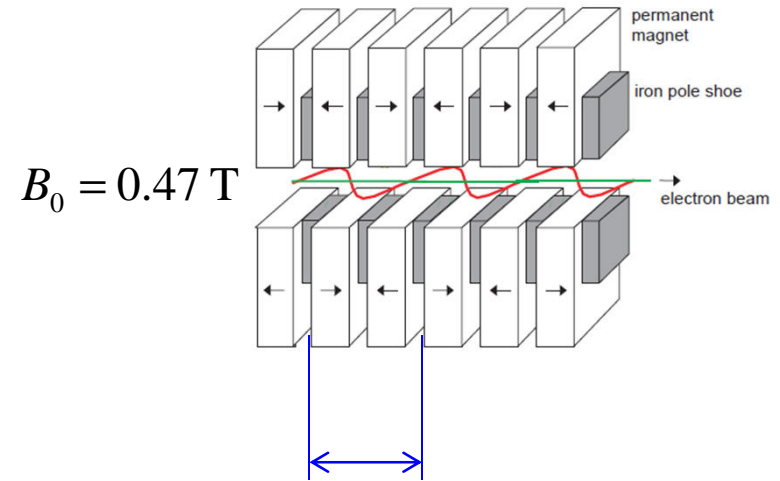
$$\hat{x} = \frac{e}{v\gamma m_e k_u^2} B_0 = \frac{K}{\beta\gamma k_u} \quad \text{with undulator parameter} \quad K = \frac{eB_0}{m_e c k_u} \ll 1$$

$$\bar{v} = v - v \left(\frac{K}{2\beta\gamma} \right)^2 \approx c - \frac{c}{2\gamma^2} \left(1 + \frac{K^2}{2} \right)$$

$$\hat{z} = \frac{K^2}{8\gamma^2 k_u}$$



example: FLASH



$$\lambda_u = 27 \text{ mm}$$

$$K \approx \frac{0.934}{\text{cmT}} B_0 \lambda_u \approx 1.2$$

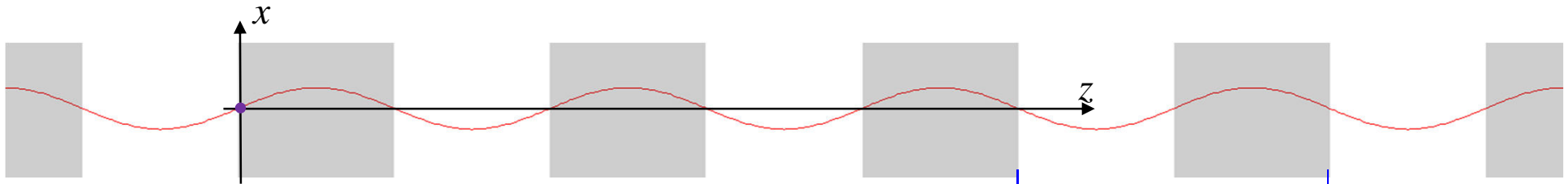
$$W \approx 1 \text{ GeV} \rightarrow \gamma \approx 1957$$

$$\hat{x} = \frac{K}{\gamma k_u} = 2.6 \text{ } \mu\text{m}$$

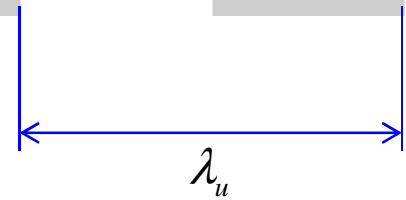
$$\hat{z} = \frac{K^2}{8\gamma^2 k_u} = 0.2 \text{ nm}$$



about particles

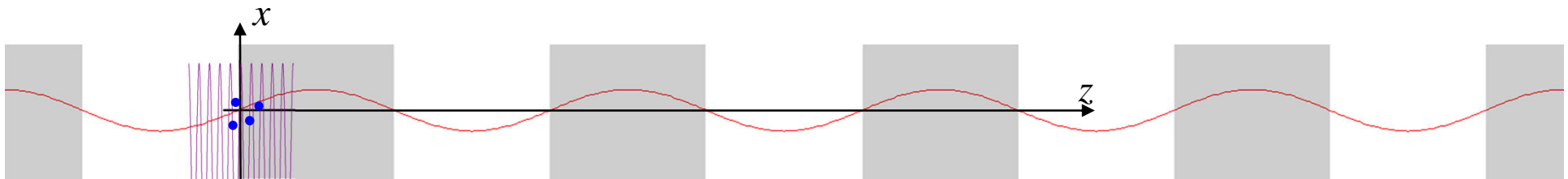


“reference particle”: only interaction with undulator field
constant energy, averaged velocity, ...
index “0”

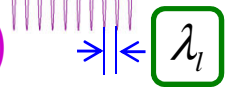


resonance
condition

$$\boxed{\gamma_0} \quad v_0 = c\sqrt{1-\gamma_0^2} \quad \bar{v}_0 = v_0 - v_0 \left(\frac{K}{2\gamma_0} \right)^2 \quad x_0 \approx \hat{x}_0 \sin \omega_u t \quad z_0 \approx \dots$$



wave
 $E(\mathbf{r}, t)$



ordinary particles:

in interaction with undulator, **external waves, self fields**, ...
slowly variation of energy, velocity (compared to λ_u)
index “ ν ” or skipped

$$\gamma_\nu \quad v_\nu \quad \bar{v}_\nu \quad \dots$$



new independent parameter: z

“reference particle”: $\gamma_0 \quad v_0 \quad \bar{v}_0$

$$t_0(z) = \frac{z}{\bar{v}_0} + \frac{\hat{z}_0}{\bar{v}_0} \cos 2k_u z$$

$$x_0(z) = \hat{x}_0 \sin k_u z$$

ordinary particles:

$\gamma_v(z) \quad v_v(z) \quad \bar{v}_v(z)$ approach: nearly constant on one period λ_u

$$t_v(z) \approx \underbrace{t_{v,i}}_{\text{initial condition}} + \underbrace{\int_0^z \frac{dz}{\bar{v}_v(z)} + \frac{\hat{z}_0}{\bar{v}_0} \cos 2k_u z}_{\text{slippage effects}} = T_v(z) + t_0(z)$$

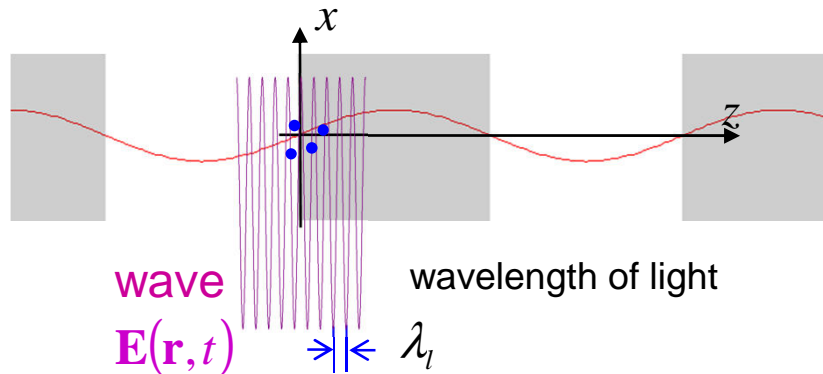
initial condition slippage effects

$$x_v(z) \approx \underbrace{x_{v,i}}_{\text{initial condition}} + \hat{x}_0 \sin k_u z = x_{v,i} + x_0(z)$$

energy parameter: $\eta_v = \frac{\gamma_v - \gamma_0}{\gamma_0}$



particle dynamics: interaction with EM wave



$$dW = -e\mathbf{E}(\mathbf{r}, t) \cdot d\mathbf{r}$$

plane wave with $k_l = 2\pi/\lambda$
x polarization

dx/dz

$$\frac{dW_v}{dz} = -e \left[E_0 \cos\{k_l(z - ct_v(z))\} \right] \cdot \left[\frac{eK}{\gamma} \cos k_u z \right]$$

$$t_v = T_v + \frac{z}{v_0} + \frac{\hat{z}_0}{v_0} \cos 2k_u z$$

$$\frac{dW_v}{dz} = -\frac{eE_0K}{\gamma} \cos\left\{ \left[k_l(1 - c/\bar{v}_0) \right] z - k_l(cT_v(z) + \hat{z}_0 \cos 2k_u z) \right\} \cdot \cos k_u z$$

resonance condition (permanent energy transfer): $k_l(1 - c/\bar{v}_0) = k_u \rightarrow \lambda_l = \frac{\lambda_u}{2\gamma_0^2} \left(1 + \frac{K}{2} \right)$

it is a condition for the energy of the reference particle or the wavelength λ_l



particle dynamics: interaction with EM wave

averaged vs. undulator period $\left\langle \frac{dW_\nu}{dz} \right\rangle = ?$ $T_\nu(z) \approx \text{const}$

estimation **without** longitudinal oscillation $\langle \cos(k_u z + \psi + \hat{z}_0 \cos 2k_u z) \cos k_u z \rangle = \frac{1}{2} \cos \psi$

$$\frac{dW_\nu}{dz} = -\frac{eE_0 K}{2\gamma} \cos \psi_\nu$$

$$\psi_\nu = k_l c T_\nu(z)$$

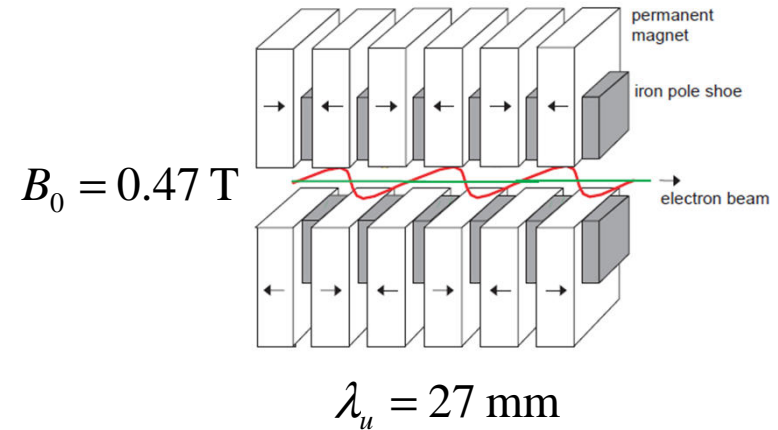
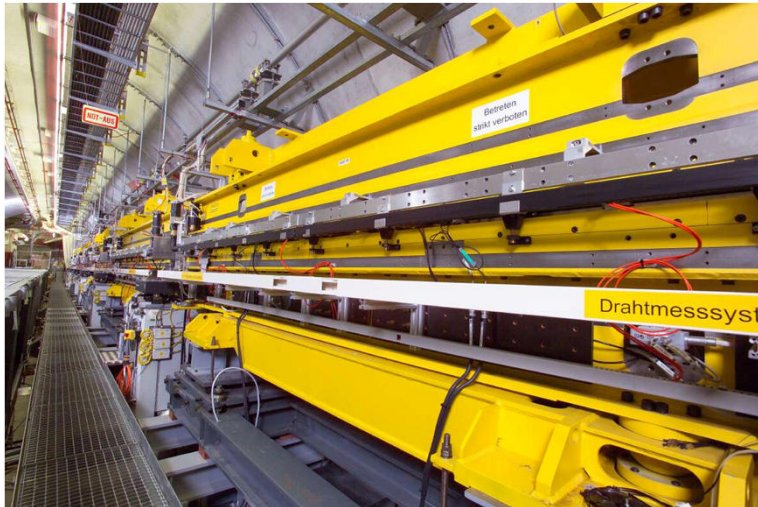
ponderomotive phase

with longitudinal oscillation: replace K by $\hat{K} = K \left[J_0 \left(\frac{K^2}{4+2K^2} \right) - J_1 \left(\frac{K^2}{4+2K^2} \right) \right]$

(modified undulator parameter)



example: FLASH



$$K \approx \frac{0.934}{\text{cmT}} B_0 \lambda_u \approx 1.2$$

$$W \approx 1 \text{ GeV} \rightarrow \gamma \approx 1957$$

$$\hat{x} = \frac{K}{\gamma k_u} = 2.6 \mu\text{m}$$

$$\hat{z} = \frac{K^2}{8\gamma^2 k_u} = 0.2 \text{ nm}$$

$$\lambda_l = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} \right) \approx 6 \text{ nm}$$

$$\hat{K} = K \left[J_0 \left(\frac{K^2}{4 + 2K^2} \right) - J_1 \left(\frac{K^2}{4 + 2K^2} \right) \right]$$

$$\hat{K} = 1.06$$



longitudinal equation of motion

new longitudinal parameters $t_v(z) = T_v(z) + t_0(z) = \frac{\psi_v(z)}{k_l c} + t_0(z)$ ponderomotive phase

$$\eta_v(z) = \frac{W_v(z)}{W_0} - 1 \quad \text{relative energy deviation}$$

with $\frac{d\psi_v}{dz} = k_l c \frac{dT_v(z)}{z} = \frac{k_l c}{\bar{v}_v} - \frac{k_l c}{\bar{v}_0} \approx 2k_u \eta_v$ and $W_0 = m_e c^2 \gamma_0^2$ follow

particle equations

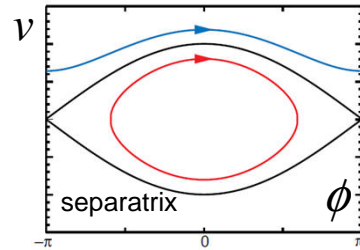
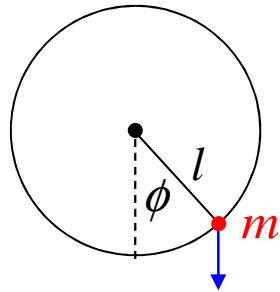
$$\frac{d\eta_v}{dz} \approx -\frac{eE_0 \hat{K}}{2m_e c^2 \gamma_0^2} \cos \psi_v$$

$$\frac{d}{dz} \psi_v \approx 2k_u \eta_v$$

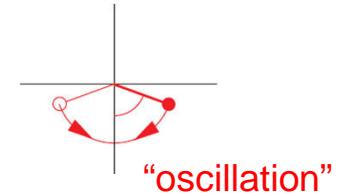
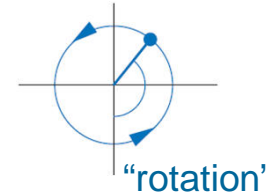


phase space and pendulum

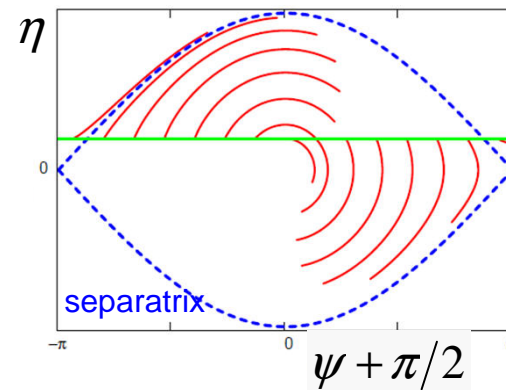
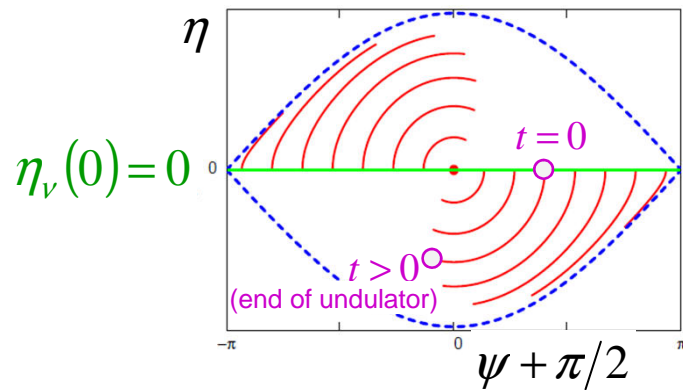
equations are analog to mathematical pendulum



two types of solution:



trajectories in phase space



$\eta_v(0) > 0$
or
 $\gamma_v(0) > \gamma_0$

15 particles with different initial conditions



FEL low gain theory

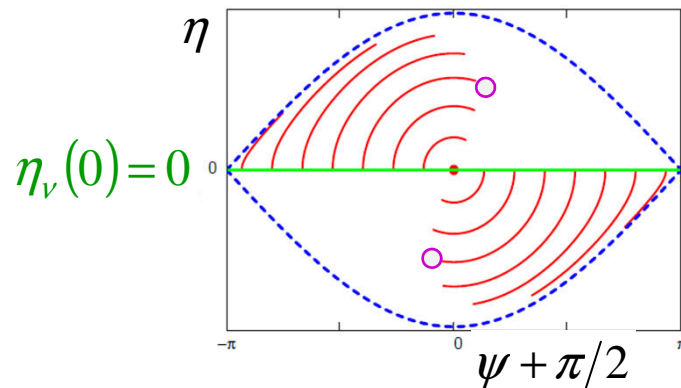
neglect change of field amplitude

indirect gain calculation

$$G = \frac{\text{gain of field energy}}{\text{initial field energy}} = \frac{\text{loss of particle energy}}{\text{initial field energy}} = \frac{W_{\Sigma}(\text{in}) - W_{\Sigma}(\text{out})}{\text{initial field energy}}$$

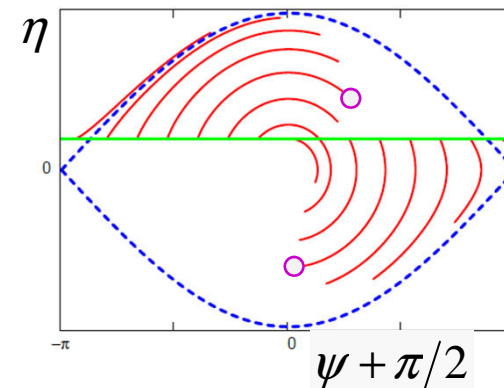
$$W_{\Sigma}(\text{in}) = W_{\Sigma}(\text{out})$$

$$G = 0$$



$$W_{\Sigma}(\text{in}) > W_{\Sigma}(\text{out})$$

$$G > 0$$



$$\eta_v(0) > 0$$

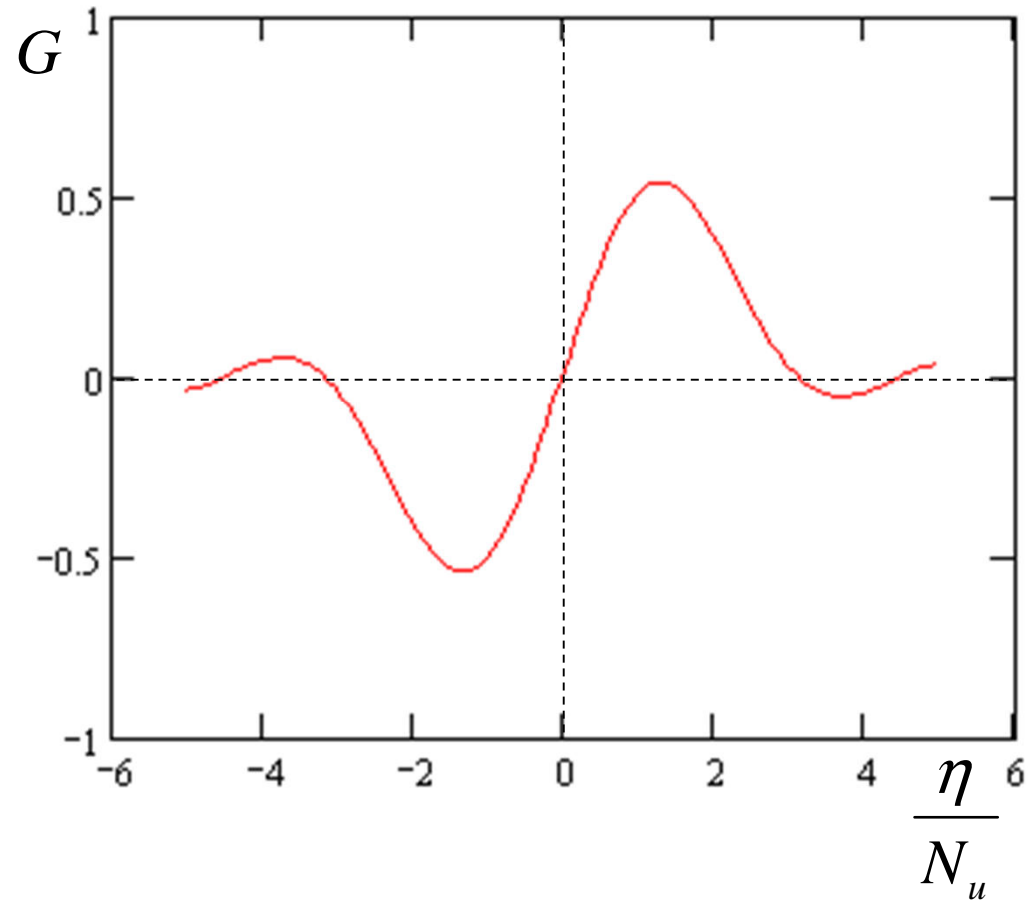
or

$$\gamma_v(0) > \gamma_0$$



FEL low gain theory

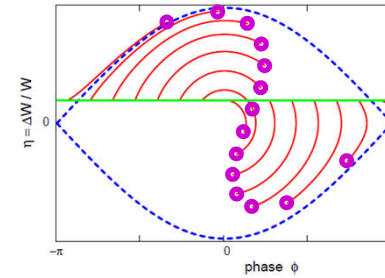
neglect change of field amplitude



$N_u =$ periods of undulator



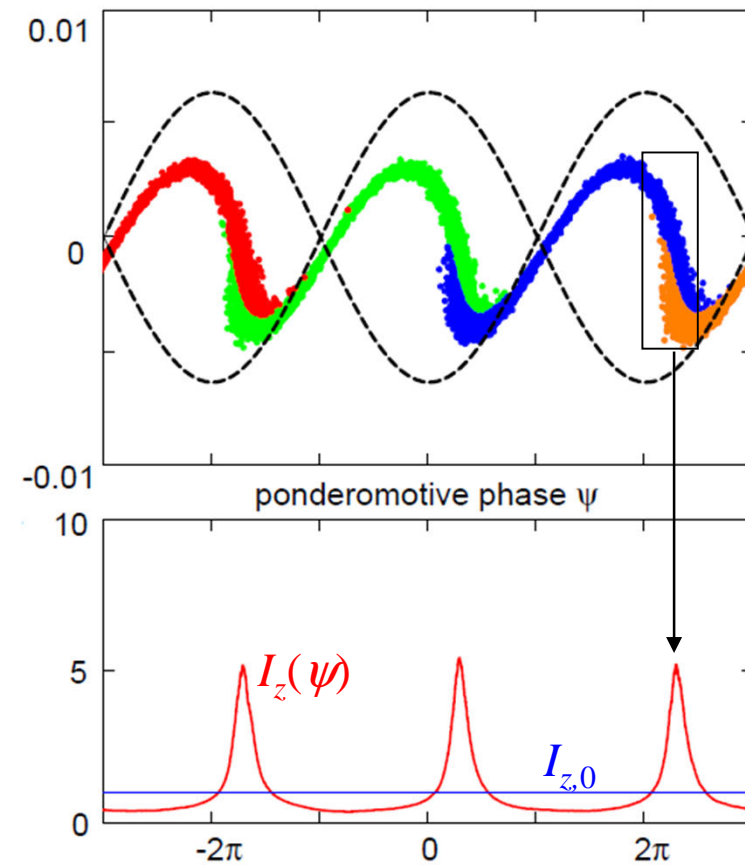
micro-bunching



Fourier analysis
→ amplitude of micro modulation

$$\hat{I} \propto \sum \exp(-i\psi_v)$$

(fundamental mode)



electrodynamics (1D)

Maxwell equations



wave equation



1D wave equation

$$\left(\frac{\partial^2}{\partial z^2} - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) E_x = \mu_0 \frac{\partial}{\partial t} J_x$$

$$J_x(z, t) = \frac{\hat{x}_0 k_u \cos k_u z}{v_z(z)} J_z(z, t)$$

approach with slowly varying amplitude

$$J_z(z, t) \sim \text{Re} \left\{ \hat{I}_1(z) \exp(ik_l(z - t/c)) \right\} + \dots$$

$$E_x(z, t) \sim \text{Re} \left\{ \hat{E}_x(z) \exp(ik_l(z - t/c)) \right\} + \dots$$

$$\frac{d}{dz} \hat{E}_x(z) = -\frac{\mu_0 c \hat{K}}{4\gamma_0} \hat{I}_1$$



FEL high gain theory (1D)

$$\frac{d\eta_\nu}{dz} \approx -\frac{e\hat{K}}{2m_e c^2 \gamma_0^2} \operatorname{Re}\{\hat{E}_x \exp(i\psi_\nu)\}$$

$$\frac{d}{dz}\psi_\nu \approx 2k_u \eta_\nu$$

$$\hat{I} \propto \sum \exp(-i\psi_\nu)$$

$$\frac{d}{dz}\hat{E}_x \approx -\frac{\mu_0 c \hat{K}}{4\gamma_0} \hat{I}_1$$

particle equations

micro modulation

electrodynamics



FEL high gain theory (1D)

numerical solution
(Mathcad)

state vector (initial)

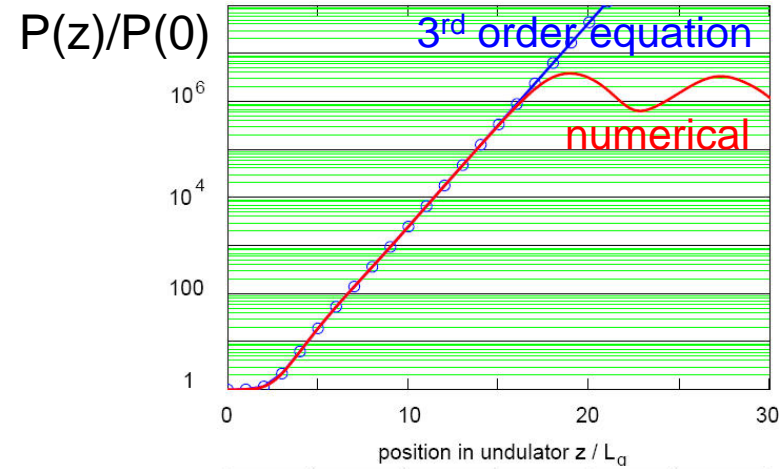
$$x0 := \text{stack} \left[\text{stack}(\eta_i, \psi_i), \begin{pmatrix} \text{Re}(Ex_i) \\ \text{Im}(Ex_i) \end{pmatrix} \right]$$

first derivative

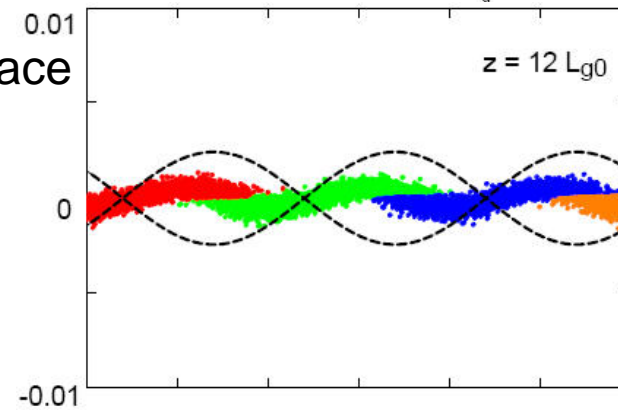
$$D(Z, x) := \begin{cases} \eta \leftarrow \text{submatrix}(x, 0, N-1, 0, 0) \\ \psi \leftarrow \text{submatrix}(x, N, 2 \cdot N-1, 0, 0) \\ Ex \leftarrow x_{2 \cdot N} + i \cdot x_{2 \cdot N+1} \\ J1 \leftarrow -c_0 \cdot N_e \cdot q_e \cdot \frac{2}{N} \cdot \sum_{n=0}^{N-1} \exp(-i \cdot \psi_n) \\ \text{for } n \in 0..N-1 \\ \left| \begin{aligned} d\eta_n &\leftarrow \frac{-q_e \cdot K}{2 \cdot m_e \cdot c_0^2 \cdot \gamma^2} \cdot \text{Re}(Ex \cdot e^{i \cdot \psi_n}) \\ d\psi_n &\leftarrow 2 \cdot k_u \cdot \eta_n \end{aligned} \right. \\ \text{stack} \left[\text{stack}(d\eta, d\psi), \frac{-\mu_0 \cdot c_0 \cdot K}{4 \cdot \gamma} \cdot \begin{pmatrix} \text{Re}(J1) \\ \text{Im}(J1) \end{pmatrix} \right] \end{cases}$$

runge kuta integration

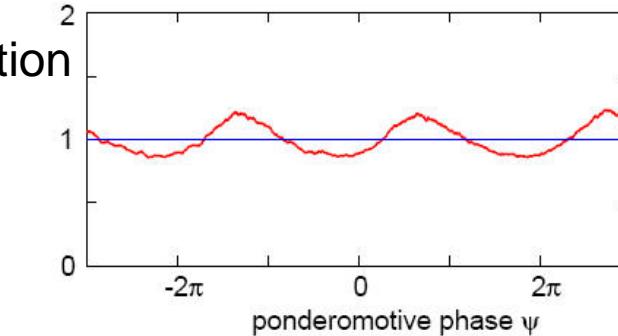
$$RK := \text{rkfixed}(x0, 0, z_s, N_s, D)$$



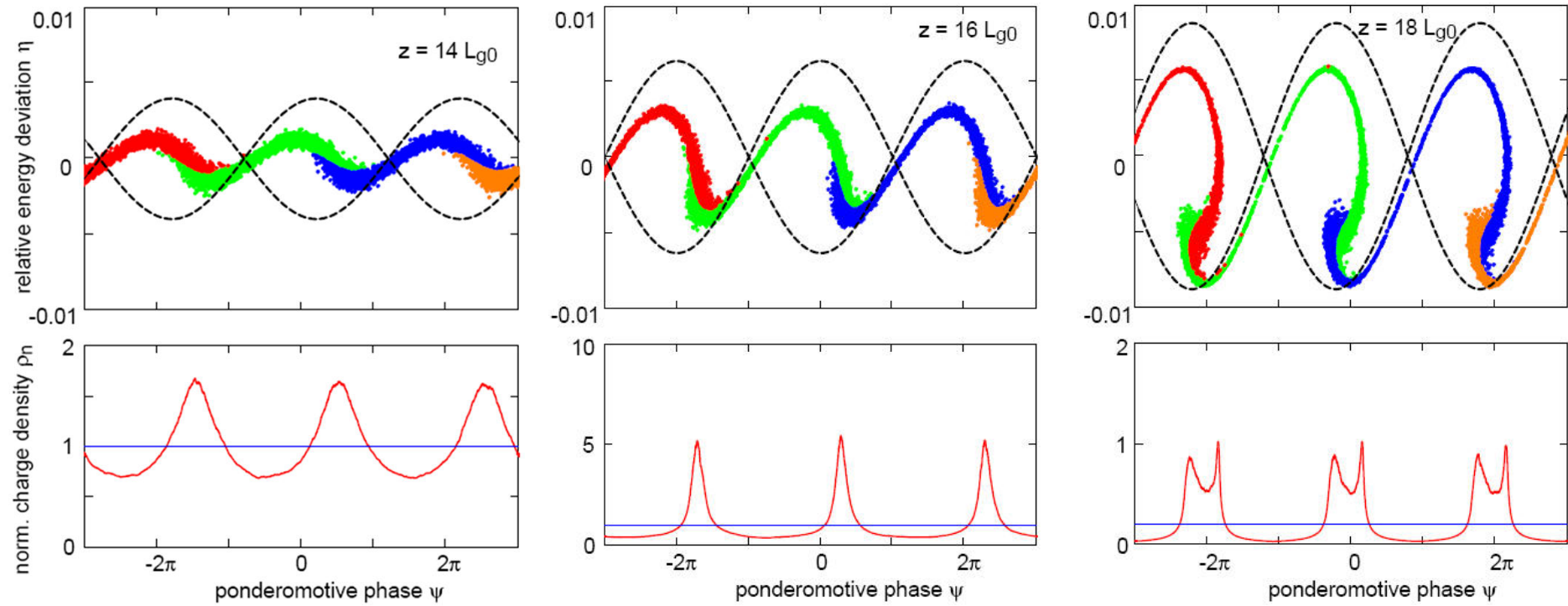
phase space



modulation



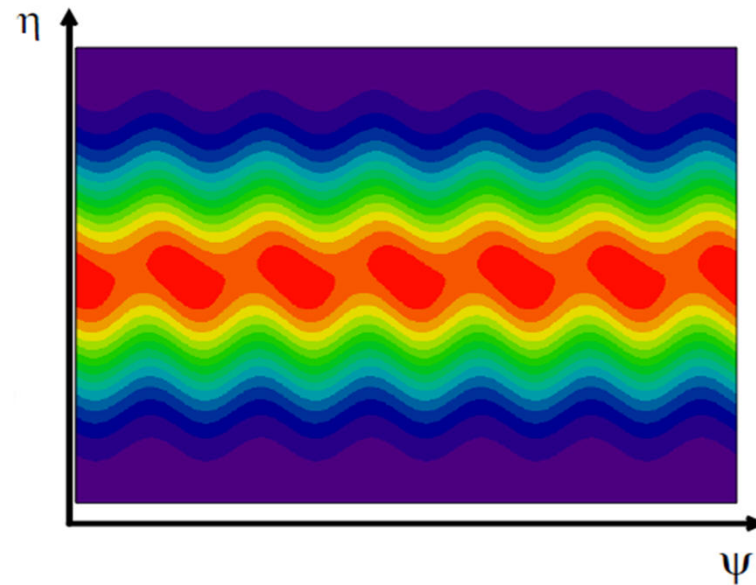
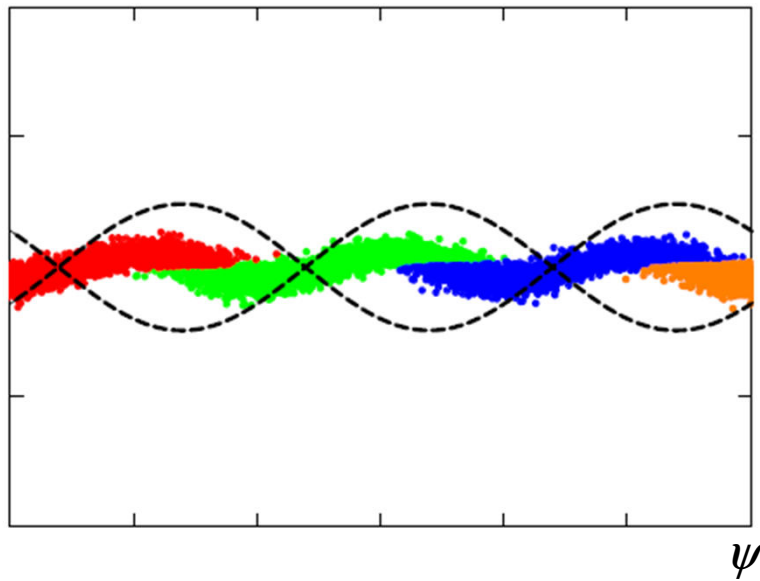
FEL high gain theory (1D)



continuous phase space, Vlasov equation

many point particles $\psi_v, \eta_v \rightarrow$ continuous density distribution $F(\psi, \eta, z)$

$$\eta = \gamma/\gamma_0 - 1$$



continuous phase space, Vlasov equation

phase space density: $F(\psi, \eta, z)$

continuity equation: $\nabla \cdot \mathbf{J}(\psi, \eta, z) + \frac{\partial F}{\partial z} = 0$

with “current density” $\mathbf{J}(\psi, \eta, z) = \begin{pmatrix} \psi' \\ \eta' \end{pmatrix} F(\psi, \eta, z)$

therefore $\frac{\partial(\psi'F)}{\partial \psi} + \frac{\partial(\eta'F)}{\partial \eta} + \eta' \frac{\partial F}{\partial \eta} + \frac{\partial F}{\partial z} = 0$

with particle equations $\psi' = \frac{\partial \psi}{\partial z} = f(\eta, z)$

$\eta' = \frac{\partial \eta}{\partial z} = g(\psi, z)$

Vlasov equation:

$$\psi' \frac{\partial F}{\partial \psi} + \eta' \frac{\partial F}{\partial \eta} + \frac{\partial F}{\partial z} = 0 \quad \text{or} \quad \frac{dF}{dz} = 0$$



FEL 3rd order equation

perturbation approach: $F(\psi, \eta, z) \approx F_0(\eta - \eta_{\text{off}}) + \text{Re}\{\hat{F}_1(\eta, z)\exp(i\psi)\}$

$$\frac{d^3 \hat{E}_x}{dz^3} + 4ik_u \eta_{\text{off}} \frac{d^2 \hat{E}_x}{dz^2} - 4k_u^2 \eta_{\text{off}}^2 \frac{d\hat{E}_x}{dz} - i\Gamma^3 \hat{E}_x = 0$$

with energy offset $\eta_{\text{off}} = \langle \eta_v \rangle = \frac{\langle \gamma_v \rangle - \gamma_0}{\gamma_0}$

gain parameter $\Gamma = \sqrt[3]{\frac{1}{4} \frac{e}{m_e} \frac{\mu_0}{c} \frac{I}{A} \frac{k_u \hat{K}^2}{\gamma_0^3}}$

beam current I
beam cross-section A

solution $\hat{E}_x(z) = A_1 \exp \alpha_1 z + A_2 \exp \alpha_2 z + A_3 \exp \alpha_3 z$



FEL 3rd order equation

$$\hat{E}_x(z) = A_1 \exp \alpha_1 z + A_2 \exp \alpha_2 z + A_3 \exp \alpha_3 z$$

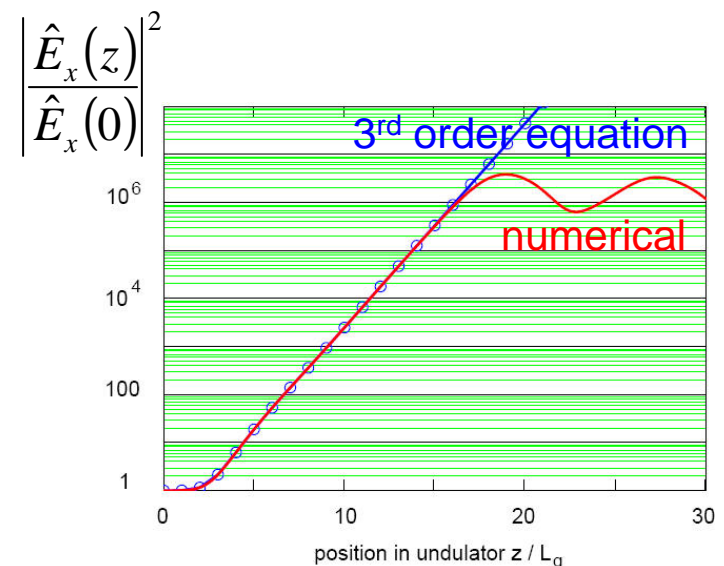
no energy offset: $\eta_{\text{off}} = 0$ or $\langle \gamma_v \rangle = \gamma_0$

$$\alpha_1 = \frac{i + \sqrt{3}}{2} \Gamma$$

$$\alpha_2 = \frac{i - \sqrt{3}}{2} \Gamma$$

$$\alpha_3 = -i\Gamma$$

positive real part \rightarrow exponential growth !



power gain length: $P(z) \rightarrow |A_1^2 \exp 2\alpha_1 z| \propto \exp\left(z \sqrt{3}\Gamma\right)^{1/L_g}$



C) Experimental Realization / Challenges

Linac Coherent Light Source - LCLS

scales

challenges

rf gun

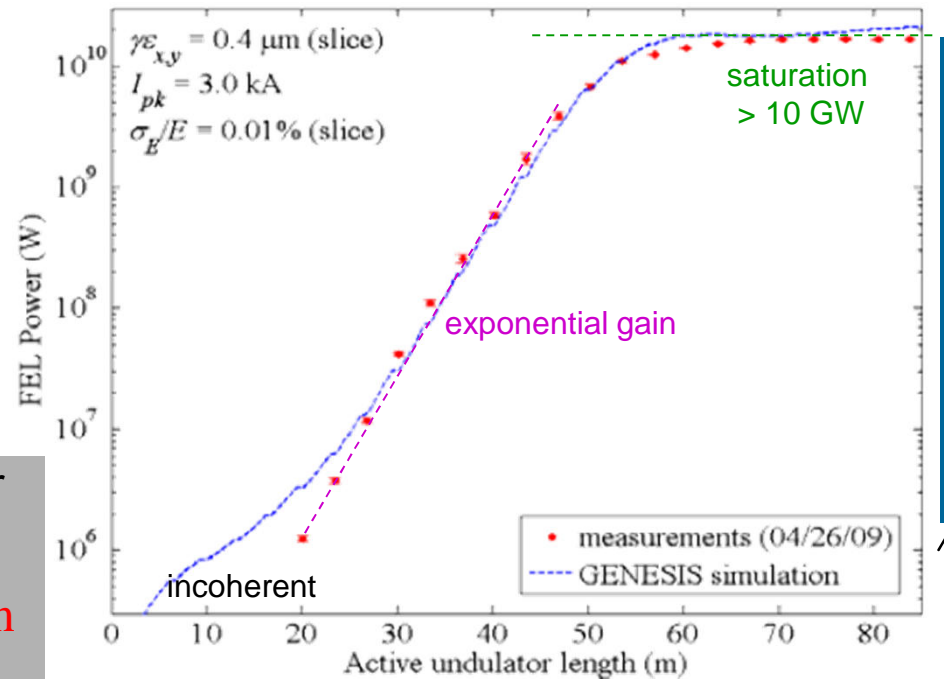
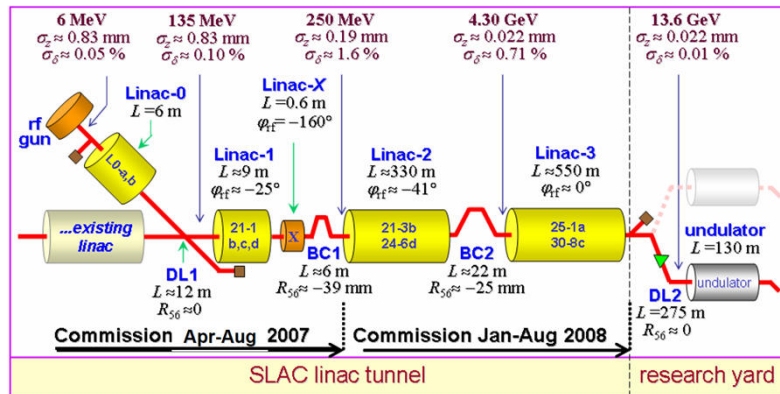
bunch compression

European X-FEL



Linac Coherent Light Source- LCLS

SLAC mid-April 2009 – first lasing at 1.5 Å



injector	linac & bunch compression	undulator
$q \approx 0.25 \text{ nC}$	$L \propto 10^3 \text{ m}$ $E = 13.6 \text{ GeV}$ $I = 3 \text{ kA}$	$L \propto 100 \text{ m}$ $\lambda_u \approx 3 \text{ cm}$ $L_g \approx 3.3 \text{ m}$
		$\lambda_l = 1.5 \text{ \AA}$

particles per λ_l

$$N = \frac{I\lambda_l}{ec} \approx 10^4$$



scales

photon wavelength $\lambda_l \propto 10^{-10} \text{ m} \propto \lambda_u / \gamma^2$

cooperation length $L_l \propto 10^{-8} \text{ m}$

transverse oscillation $\hat{x} \propto 10^{-6} \text{ m}$ (undulator trajectory)

bunch length $L_b \propto 10^{-5} \text{ m}$

bunch width $\sigma_w^{\text{bunch}} \propto 10^{-5} \text{ m}$ width of photon beam $\sigma_w^{\text{wave}} \propto \sqrt{\lambda_l L_R}$

undulator period $\lambda_u \propto 10^{-2} \text{ m}$

overlap of particle beam
with photon beam

power gain length $L_g \approx 1..10 \text{ m}$

Rayleigh length L_R (scale of widening of photon beam)

saturation length $L_s \approx 10L_g .. 20L_g < L_u$

undulator length $L_u \propto 100 \text{ m}$

total length $L \propto 10^3 \text{ m}$



challenges

$$\lambda_l = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} \right)$$

$$L_g = \frac{1}{\sqrt{3}} \left(\frac{4mc}{\mu e} \frac{\gamma^3 \lambda_u}{K^2} \frac{\sigma_r^2}{I} \right)^{1/3}$$

$$\sigma_r^2 \propto \lambda_l L_g$$

- $\lambda_l \rightarrow \text{\AA}$
- Energy $\rightarrow 10 \dots 20 \text{ GeV}$
- gain length $L_g < \sim 10 \text{ m}$
- high peak current $> \sim \text{kA}$
- transverse beam size $\sigma_r \propto 10 \mu\text{m}$
- energy spread
- overlap electron-photon beam

(undulator parameter $K \propto 1$)

transverse: generate low emittance beam
preservation of emittance

longitudinal: compression

acceleration

diagnostic and steering

undulator alignment

$$\frac{\sigma_r^2}{I} = \frac{\sigma_r^2 L_b}{qc}$$

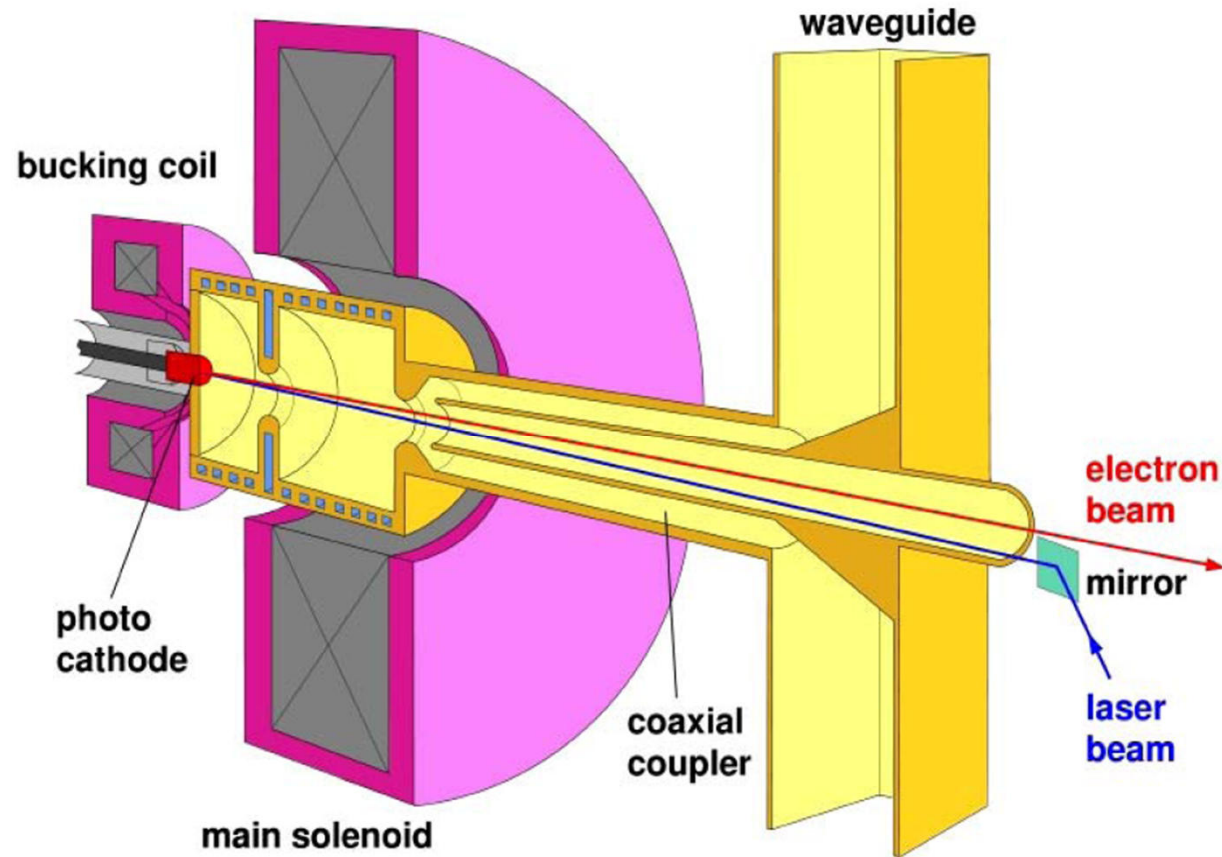
volume
bunch charge

space charge forces:

$$E_{sq} \propto \frac{1}{\gamma^2} \frac{q}{\sigma_r^2}$$



rf gun



typical parameters of FLASH & European XFEL:

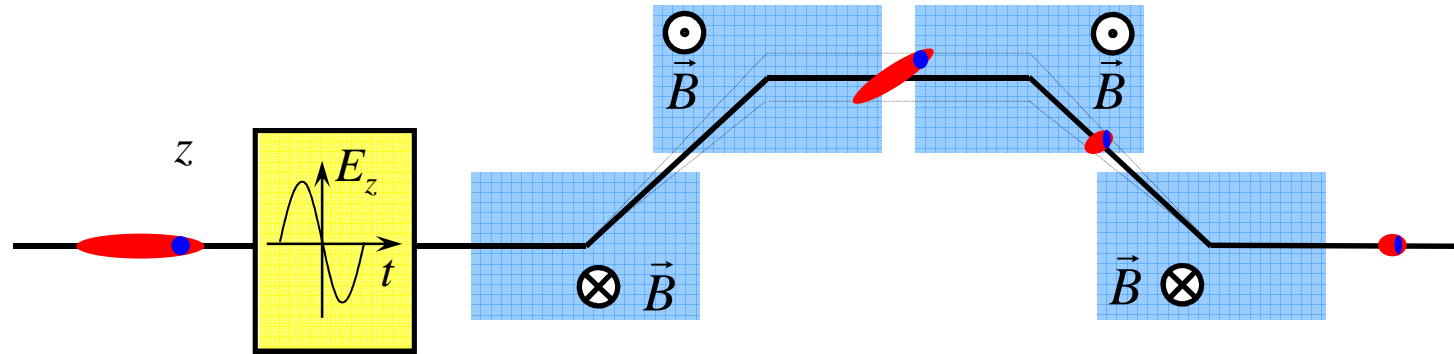
$$q \propto 0.1 \text{ nC} \quad E \propto 5 \text{ MeV} \quad I \propto 5 \text{ A}$$
$$\gamma \propto 10$$

longitudinal compression $1 \rightarrow 0.001$ needed ! (5 kA)



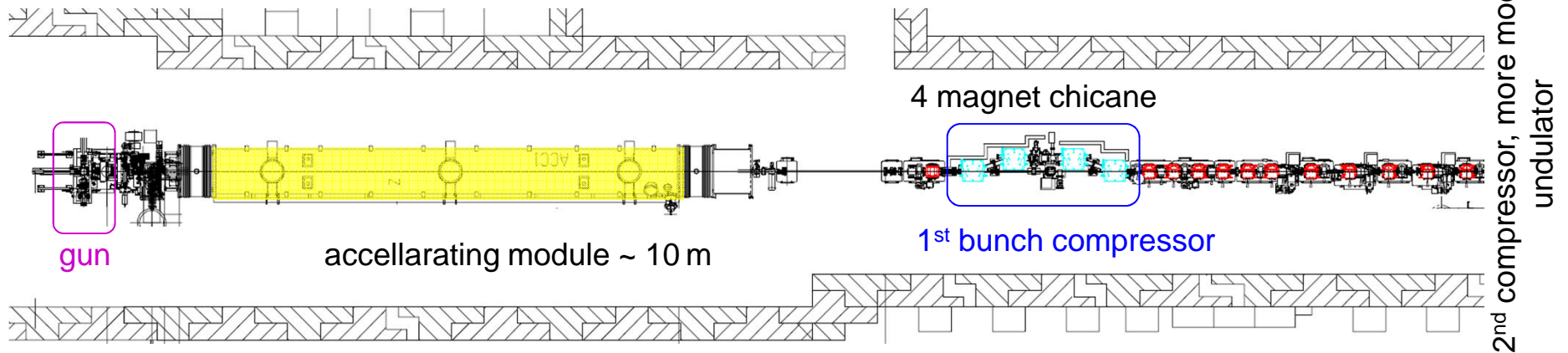
magnetic bunch compression

$\gamma \gg 1 \rightarrow$ velocity differences are too small for effective compression
magnetic compression: path length depends on energy



acceleration "off crest" \rightarrow
head particle with less energy than tail

FLASH:

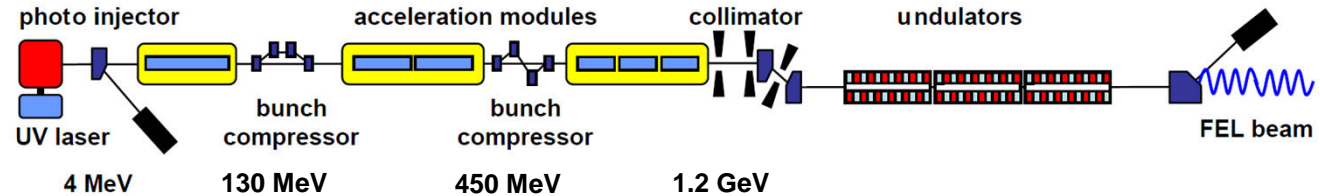


beam dynamics with space charge and CSR effects

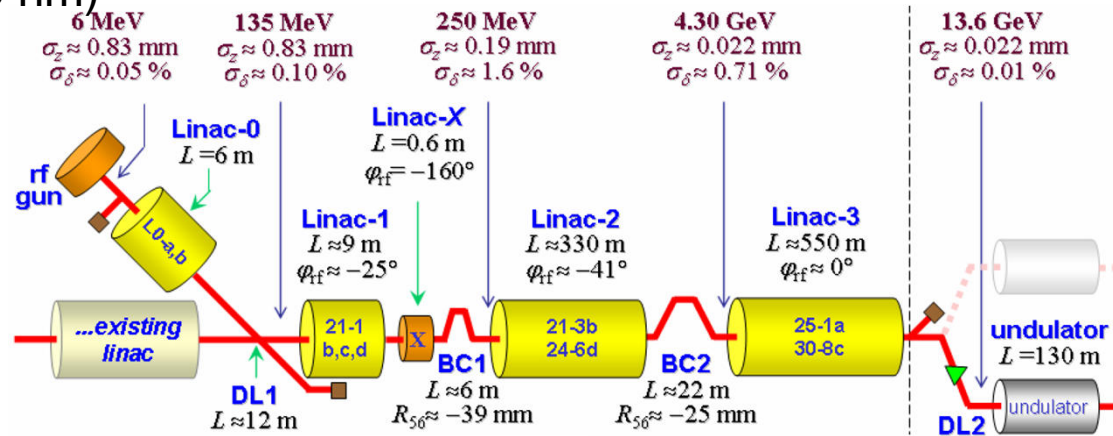


magnetic bunch compression - 2

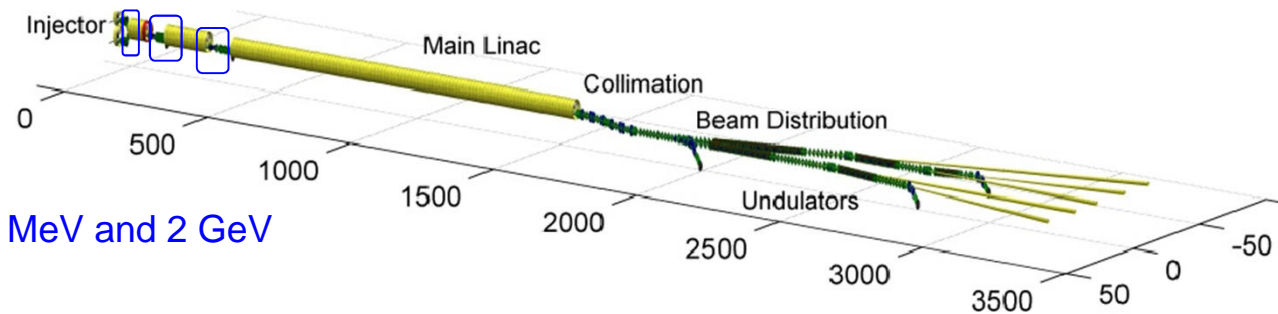
FLASH (1.2 GeV, 4 nm)



LCLS (14 GeV, 0.15 nm)



European XFEL (0.1 nm)



BCs at 130 MeV, 500 MeV and 2 GeV

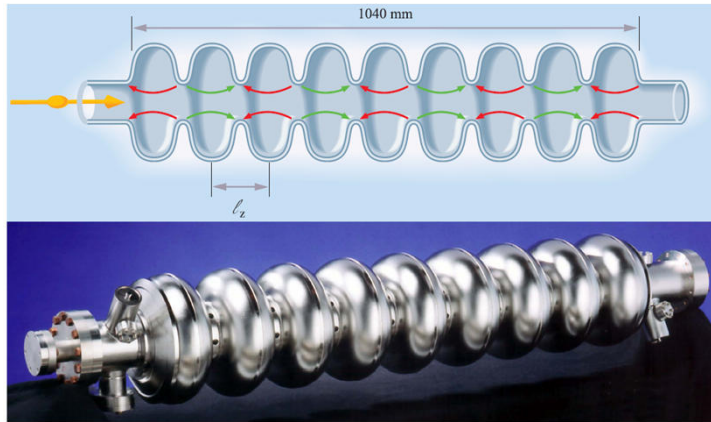


European XFEL

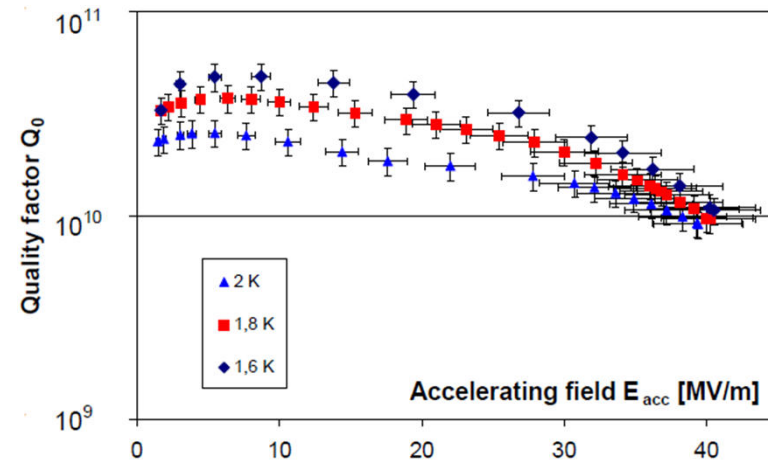


European X-FEL - 2

superconducting cavity, 1.3 GHz $E_{\text{acc}} \rightarrow 40 \text{ MeV/m}$
23.5 MeV/m are needed



FLASH tunnel: cryo module



undulator



European X-FEL - 3

	LCLS	SCSS	European XFEL
Abbreviation for	Linac Coherent Light Source	Spring-8 Compact SASE Source	European X-Ray Free-Electron Laser
Location	California, USA	Japan	Germany
Start of commissioning	2009	2010	2014
Accelerator technology	normal conducting	normal conducting	superconducting
Number of light flashes per second	120	60	30 000 multi bunch operation
Minimum wavelength of the laser light	0.15 nanometres	0.1 nanometres	0.1 nanometres
Maximum electron energy	14.3 billion electron volts (14.3 GeV)	6-8 billion electron volts (6-8 GeV)	17.5 billion electron volts (17.5 GeV)
Length of the facility	3 Kilometer	750 Meter	3.4 Kilometer
Number of undulators (magnet structures for light generation)	1	3	5
Number of experiment stations	3-5	4	10
Peak brilliance [photons / s / mm ² / mrad ² / 0.1% bandwidth]	8.5 · 10 ³²	5 · 10 ³³	5 · 10 ³³

beamlines

