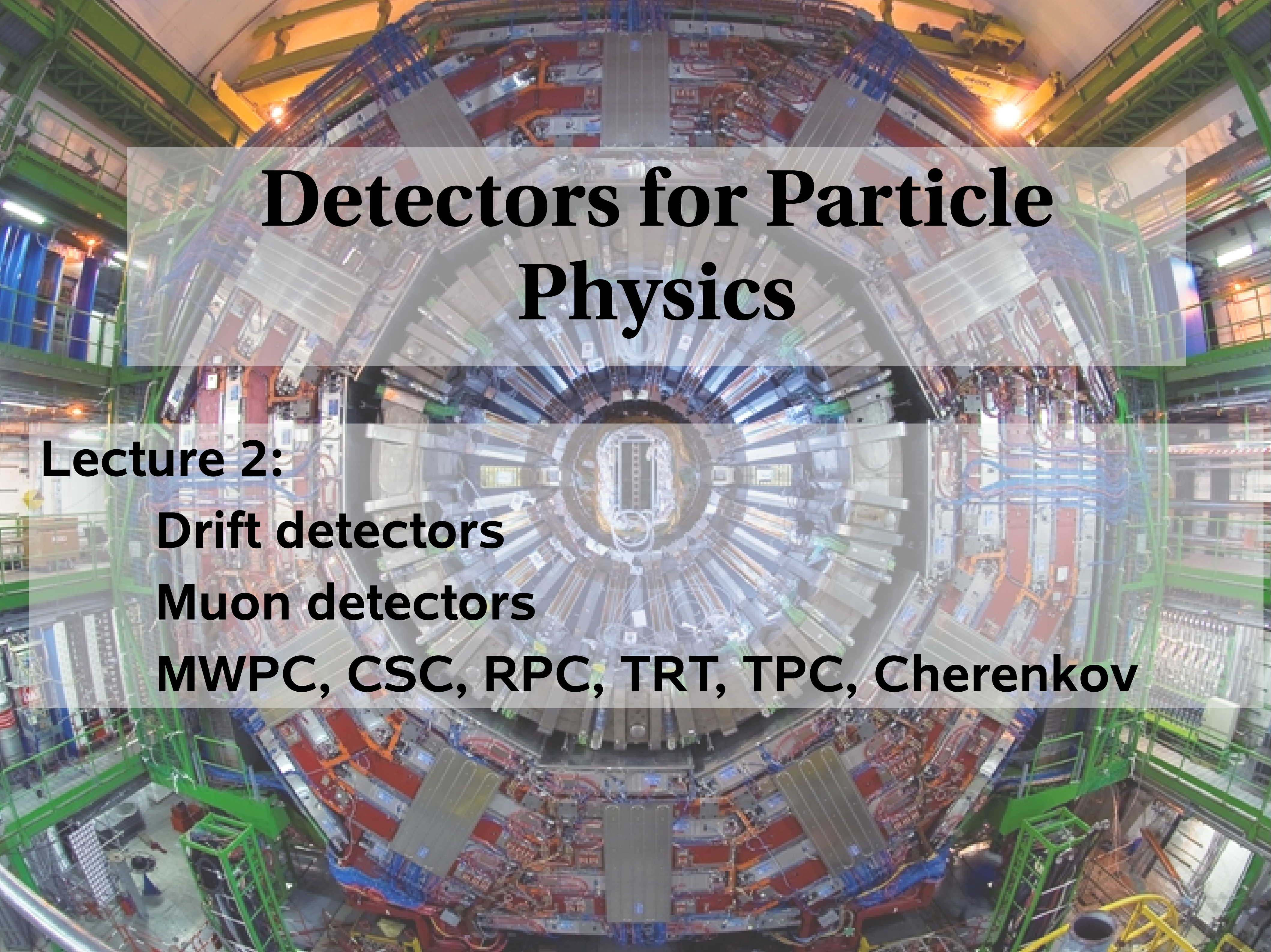


Detectors for Particle Physics

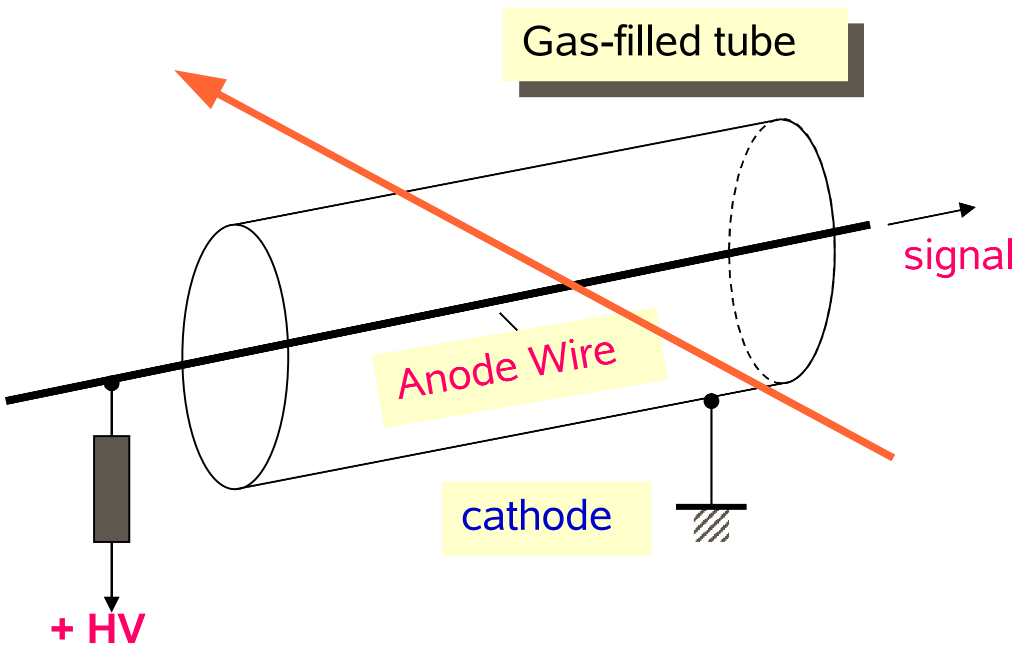
Lecture 2:

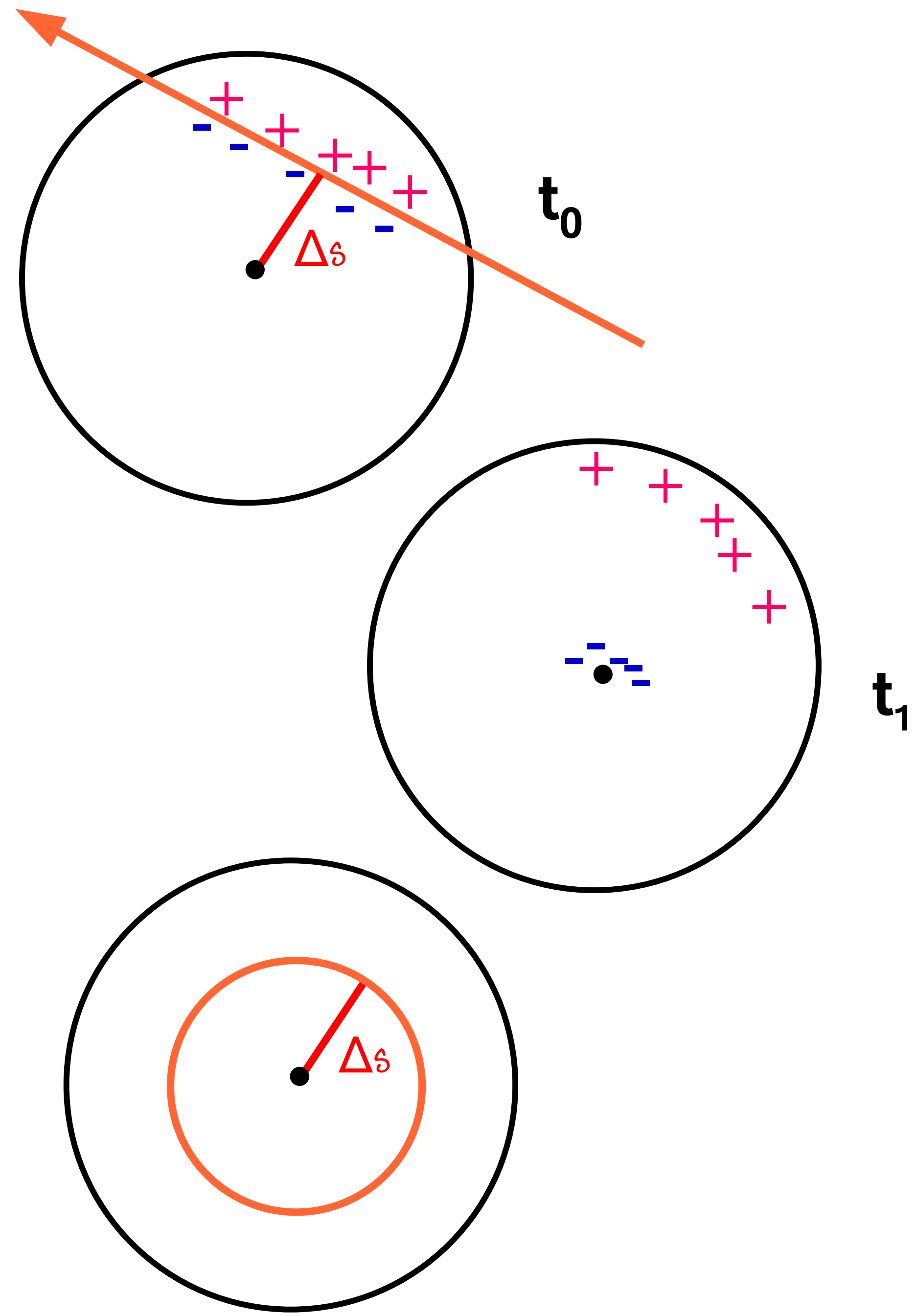
- Drift detectors
- Muon detectors
- MWPC, CSC, RPC, TRT, TPC, Cherenkov

Outline

- Lecture 1:
  - Collider detectors
  - Charged particles in a magnetic field
  - Silicon detectors
- Lecture 2:
  - Drift tubes
  - Muon systems
  - MWPCs, CSCs, RPCs, TRTs, TPCs, Cherenkovs
- Lecture 3:
  - Electromagnetic showers and calorimeters
  - Photon detectors
  - Hadronics showers and calorimeters
  - Particle flow technique
- Discussion session:
  - Your questions, please

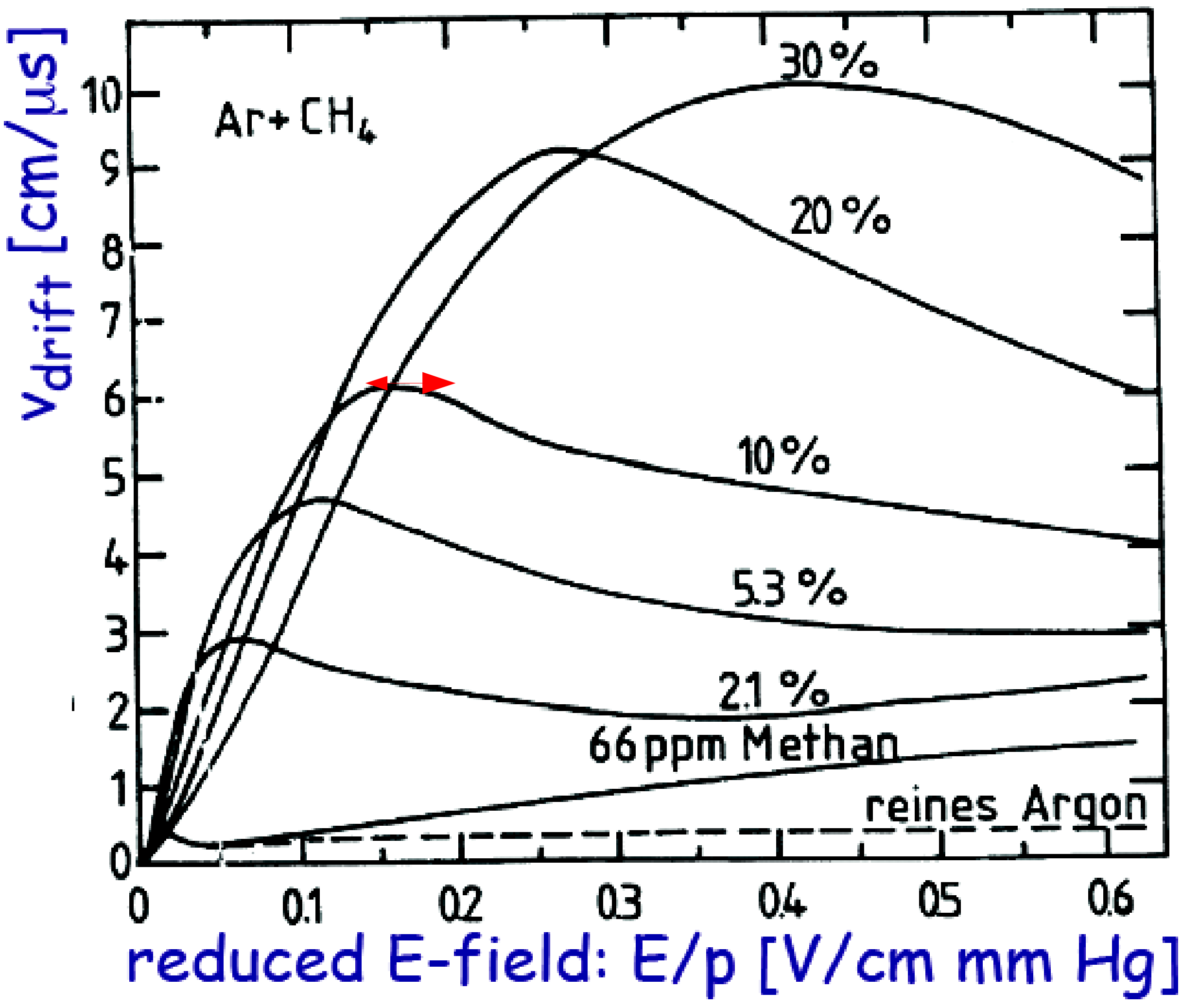
Principle of drift detectors





- Track ionises gas atoms
- electrons drift towards anode: v_D
- Amplify
- Measure drift time: $\Delta t = t_1 - t_0$
- Reconstruct radius: $\Delta s = v_D \, \Delta t$.
- vD depends on gas, voltage, pressure, temperature, field: need calibration.

Drift velocity

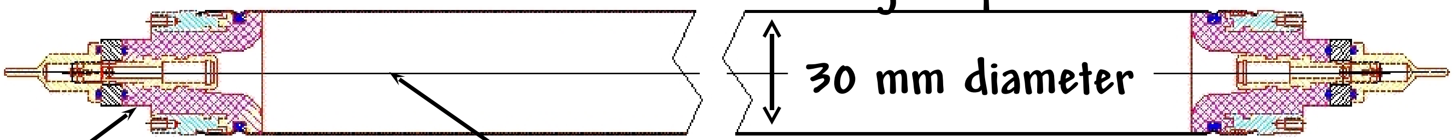


- Drift velocity depends on electric field, pressure, gas, temperature, magnetic field.
- Want stable operation point: just above maximum.
- trade-off:
  - slower gas = higher resolution.
  - faster gas better in a high-rate environment.

ATLAS drift tubes

tube wall: 0.4 mm Al

Length up to 6 m

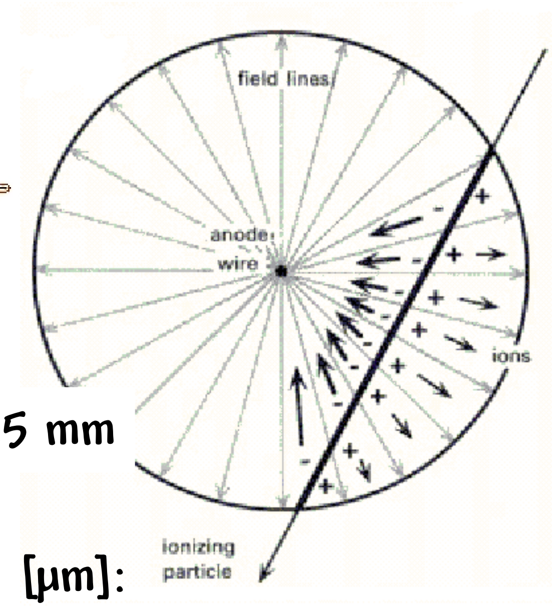


30 mm diameter

endplug

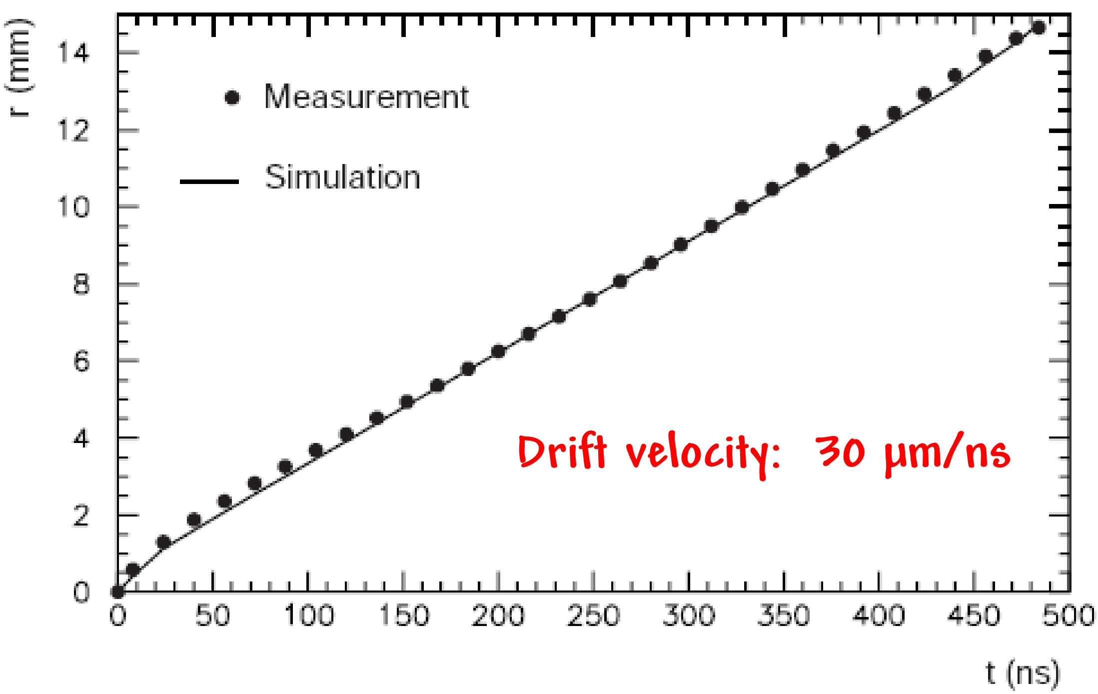
wire: 50 µm W-Re at +3300V

Gas: Ar (91%) CH4 (5%) N2 (4%) at 3 bar



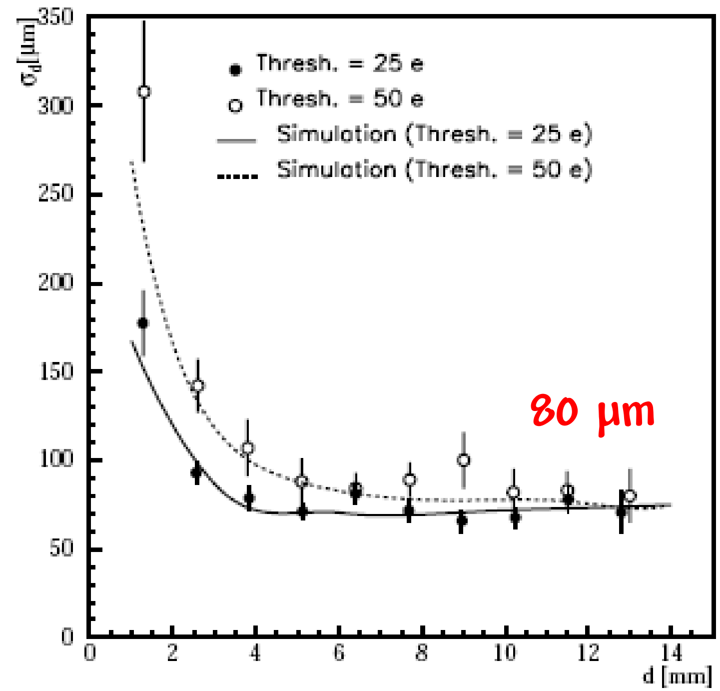
R=15 mm

Distance from time:



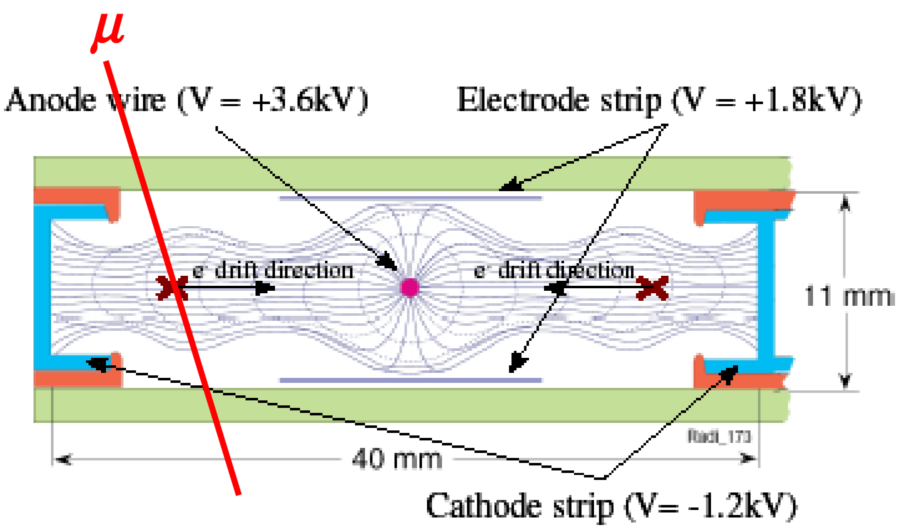
Drift velocity: 30 µm/ns

Resolution [µm]:



80 µm

CMS drift cells

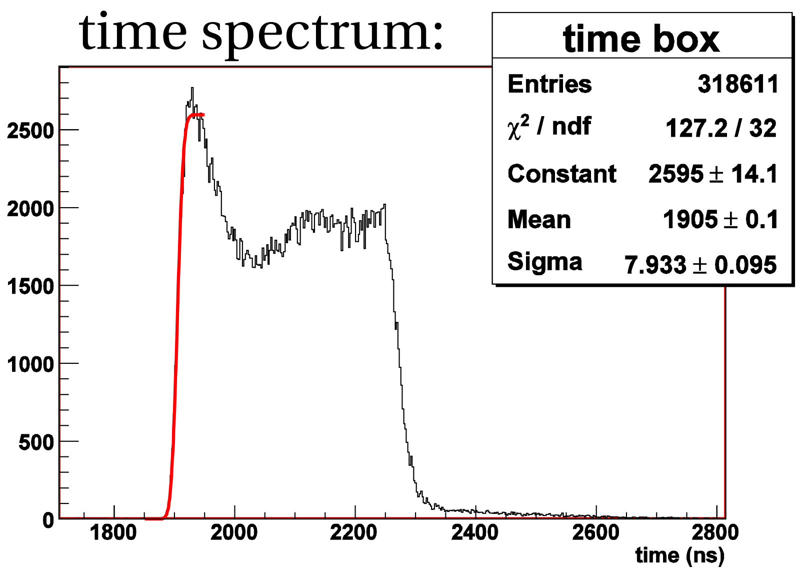


Extra electrode for field shaping:
more uniform drift.
Left/right ambiguity remains.

- **time synchronization:**

$$t_{meas} = \underbrace{t_{electr} + t.o.f. + t_{prop}}_{\text{time pedestal } (t_{trig})} + t_{drift}$$

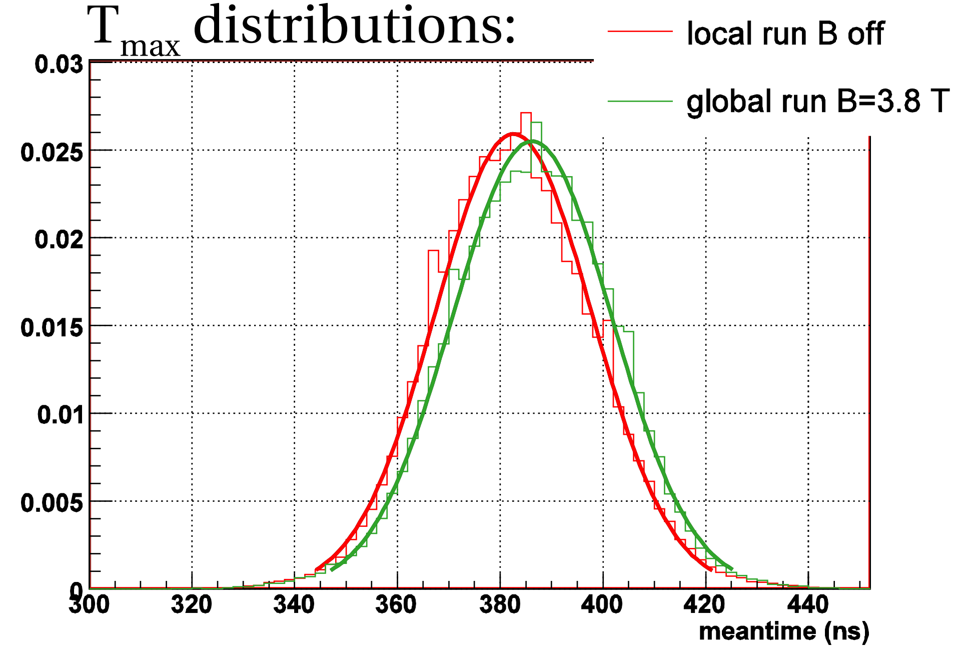
time spectrum:



t_{trig} t_{max}

- **drift velocity calibration:**

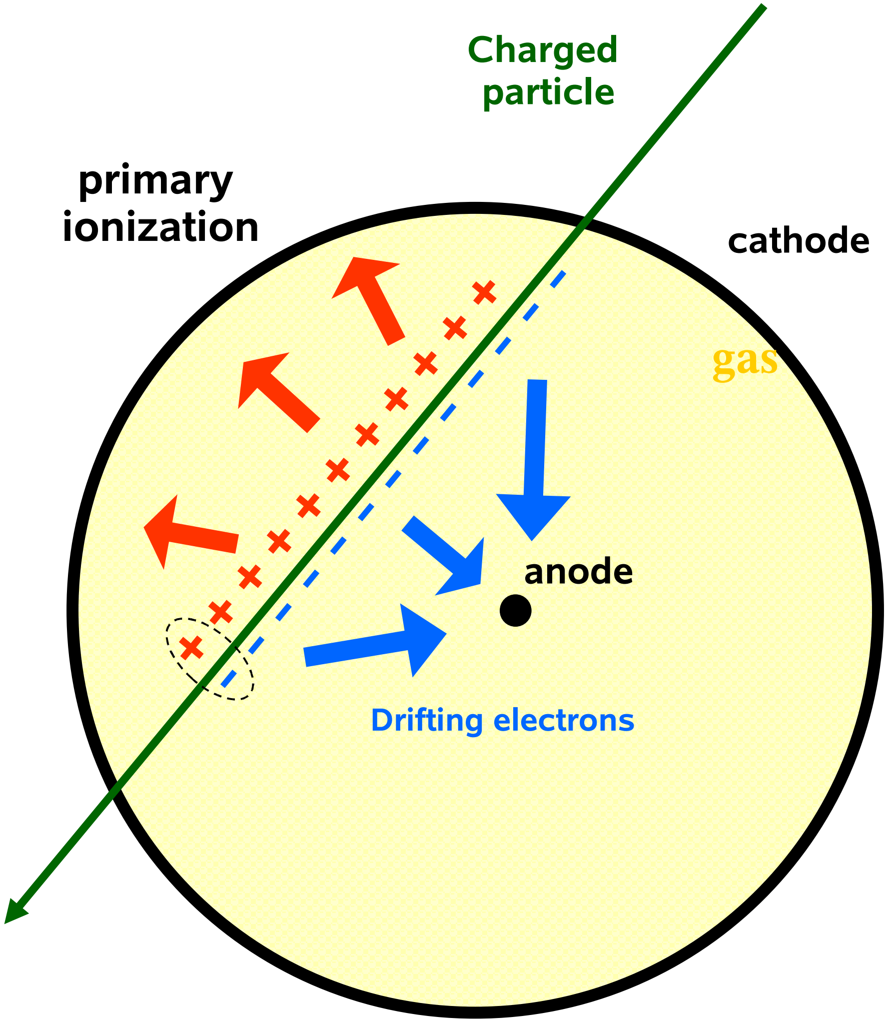
T_{max} distributions:



$$v_{drift} = L / (2 \times \langle T_{max} \rangle)$$

$$\text{resolution} = v_{drift} \times \langle \sigma_{Tmax} \rangle$$

Choice of Gas 1: ionization and drift



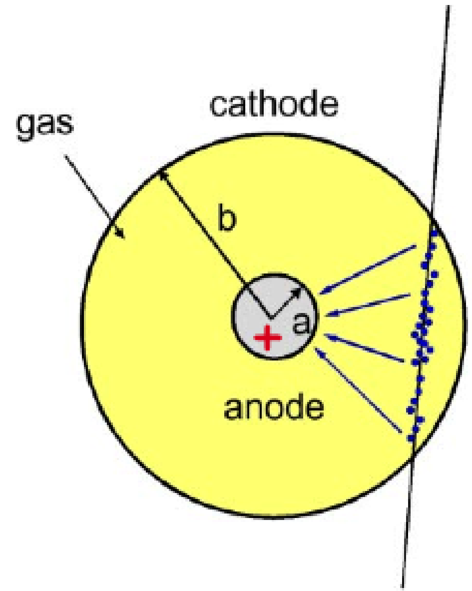
- Drifting electrons should not be trapped:
  - Use noble gas, e.g. Ar.
- Want large primary ionzation yield:
  - Ar gives 25 ions/cm at normal T, p for a minimm ionizing particle.
- The primary electrons may ionize further atoms:
  - $\times 3$ or $\times 4$ increase.
- Xe and/or higher pressure are even better (and more expensive).

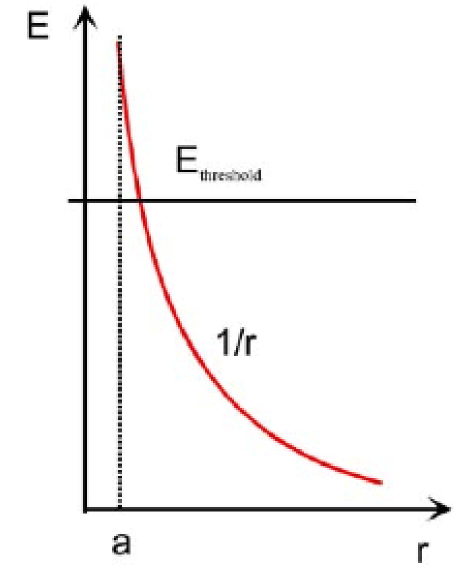
Signal amplification near the wire

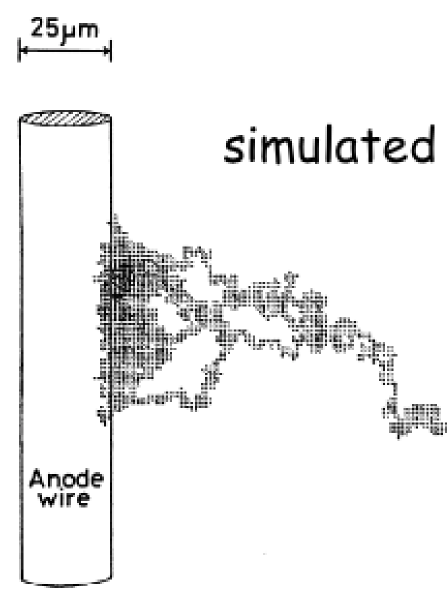
For cylindrical geometry:

$$E(r) \propto \frac{1}{r} \quad \text{and} \quad V(r) \propto \ln\frac{r}{a}$$

- the primary electrons drift towards the positive anode
- due to 1/r dependence the electric field close to very thin wires reaches values of E > kV/cm
- ⇒ in between collisions with atoms electrons gain enough energy to ionize further gas molecules
- ⇒ exponential increase in number of electron-ion pairs very close (few μm) to the wire



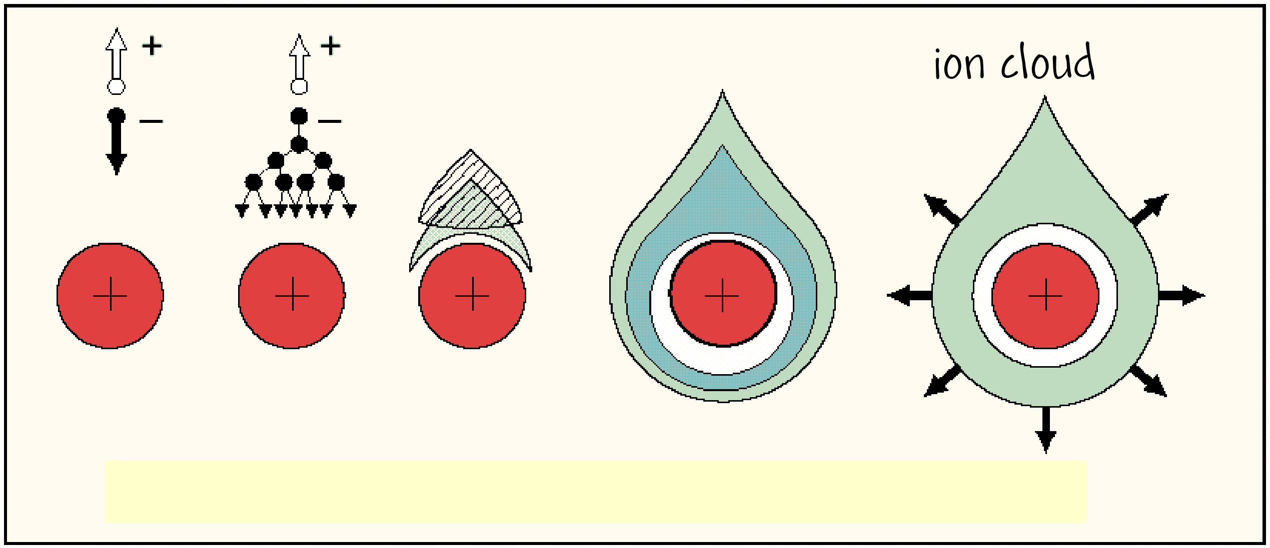




simulated avalanche

Amplification by 10^5 possible

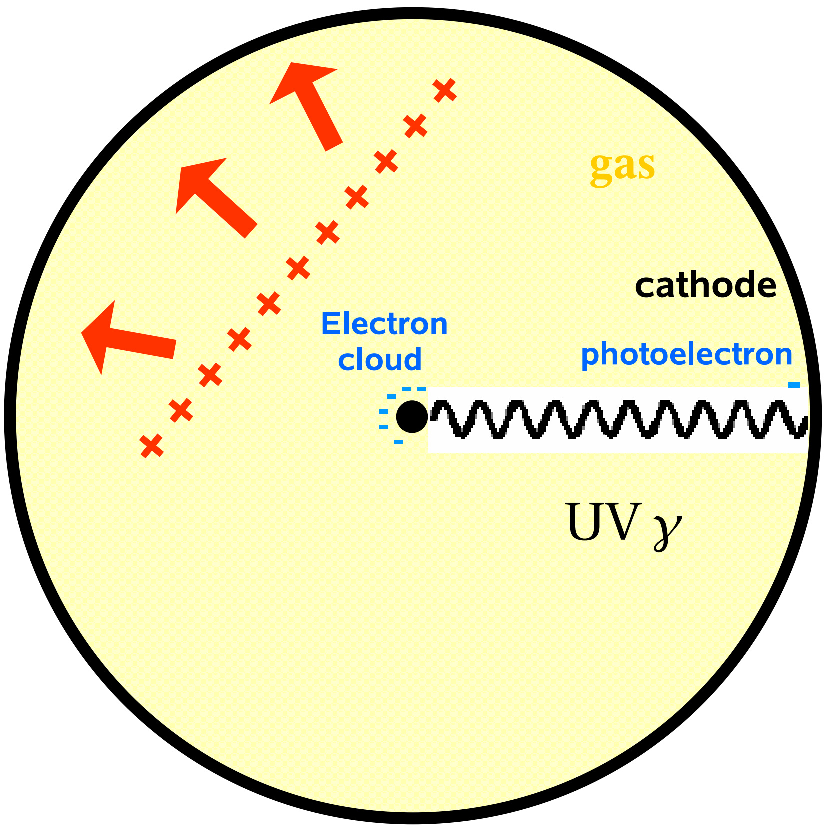
The Avalanche



- Amplification depends on:
  - Anode voltage
  - Wire radius
  - Gas composition
  - Pressure
  - temperature

Receding ion cloud induces signal on the wire

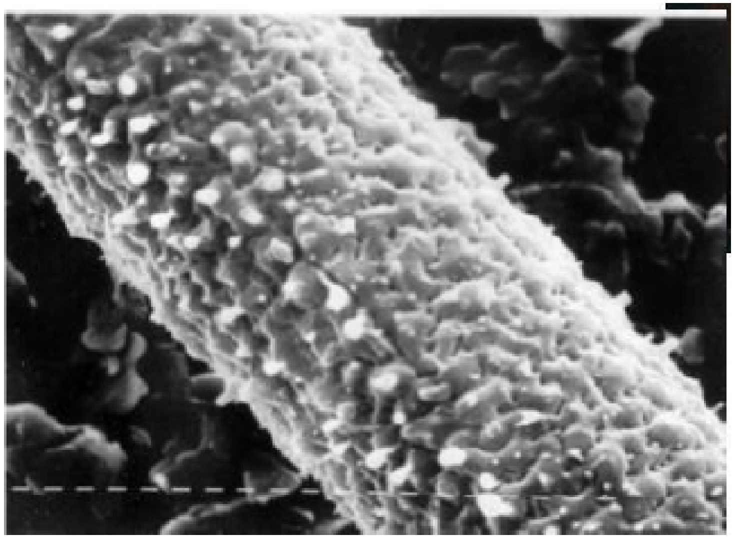
Choice of Gas 2: high gain, stable operation



- Want large gain at low voltage:
  - Ar is mono-atomic gas
  - No vibrational or rotational modes, only excitation and ionization.
- Excited Ar atoms may emit UV photons (11.6 eV):
  - UV photons may reach the cathode and produce photoelectrons.
  - Photoelectrons drift back towards the wire and may start a new avalanche
  - Continuous discharge!
- Need 'quencher' molecules that absorb UV photons without creating photoelectrons:
  - CH_4, C_2H_6, CO_2, ...

Choice of Gas 3: prevent ageing

deposits on the wire:

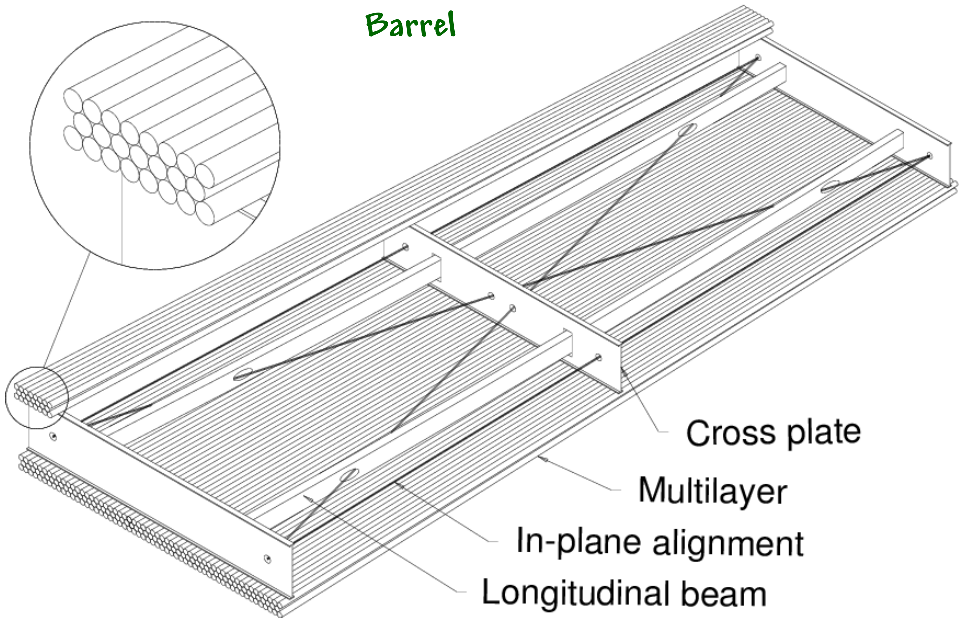


'whiskers':

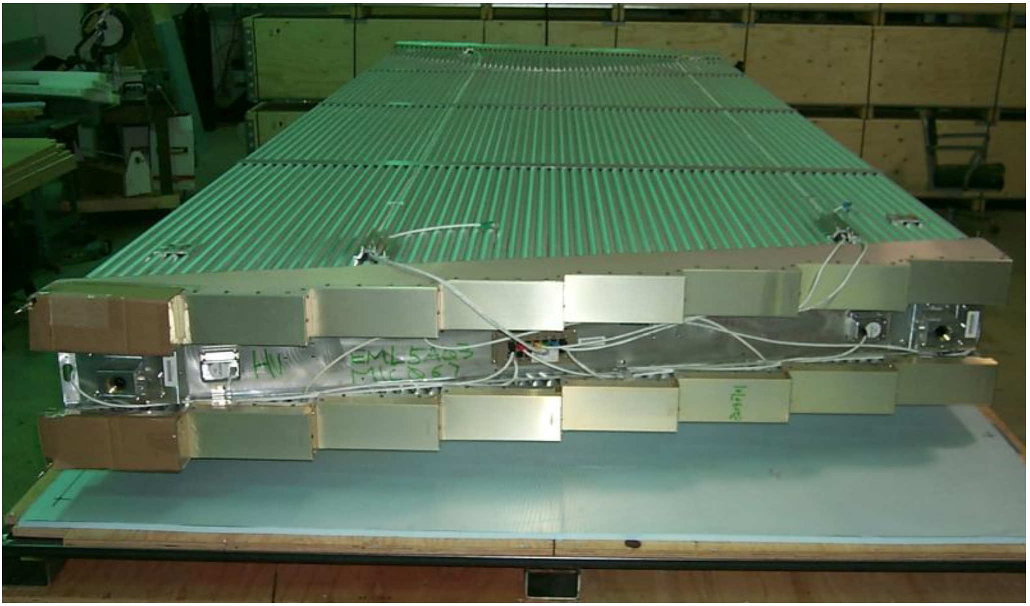
- Impurities in the gas (or in the chamber) may form deposits on the wire and reduce the gain.
- 'Whiskers' lead to HV instabilities.
- Prevention:
  - Build chamber in a clean room.
  - Use clean gas.
  - Add ~1% alcohol, water, or oxygen.

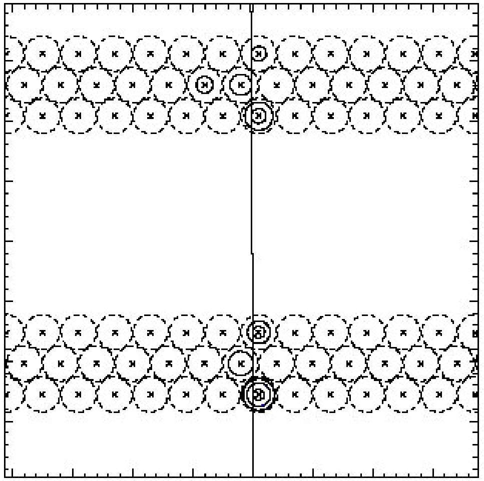
ATLAS Drift Tube Chambers

- 6 drift tube layers, arranged in 2 multilayers glued to a spacer frame
- length: 1 – 6 m, width: 1 – 2 m
- optical system to monitor chamber deformations
- gas: Ar:CO_2 (93:7) to prevent aging, 3 bar

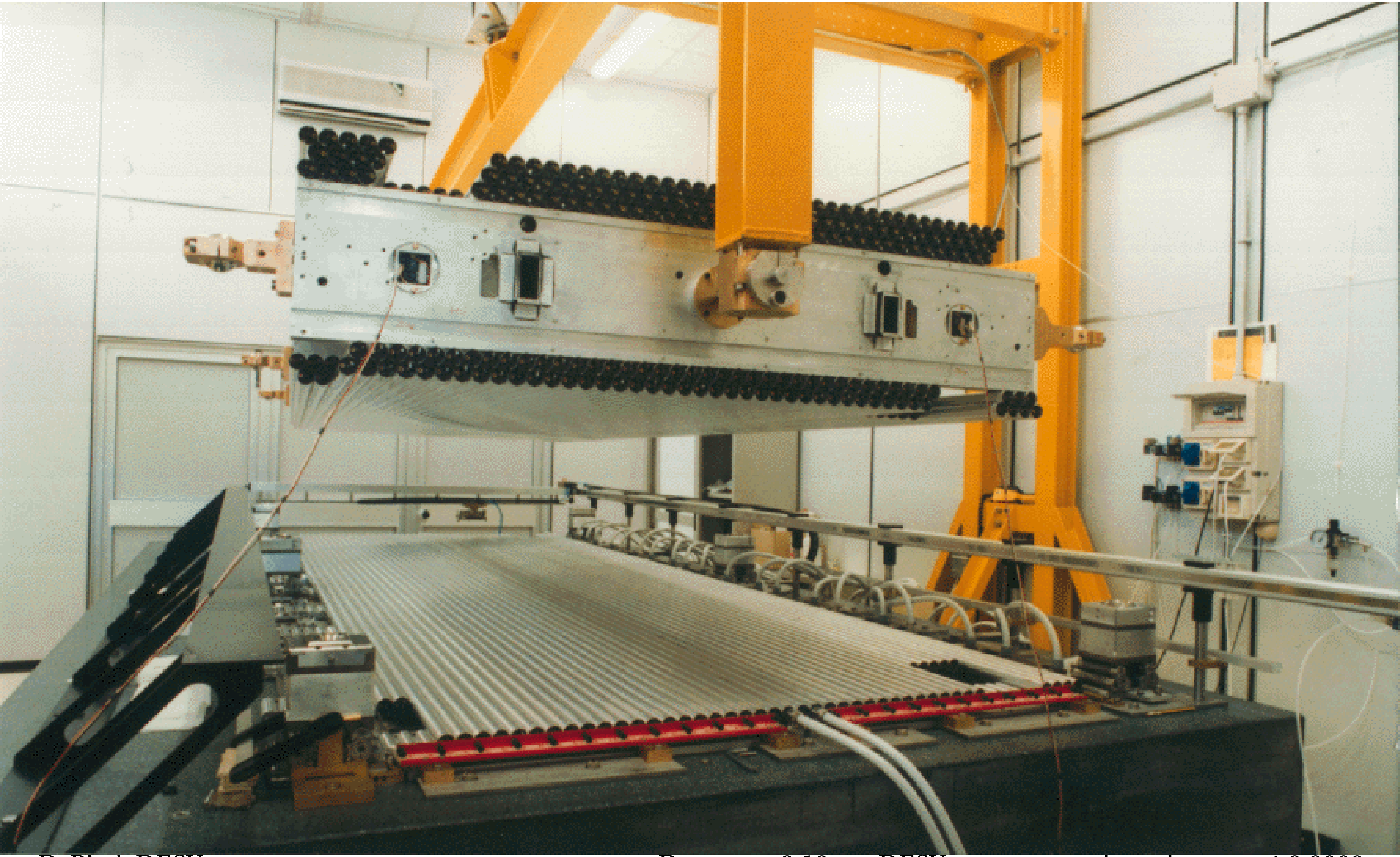


- chamber resolution: 50 μm
  - ➔ single tube resolution: 100 μm
  - ➔ required wire position accuracy: 20 μm

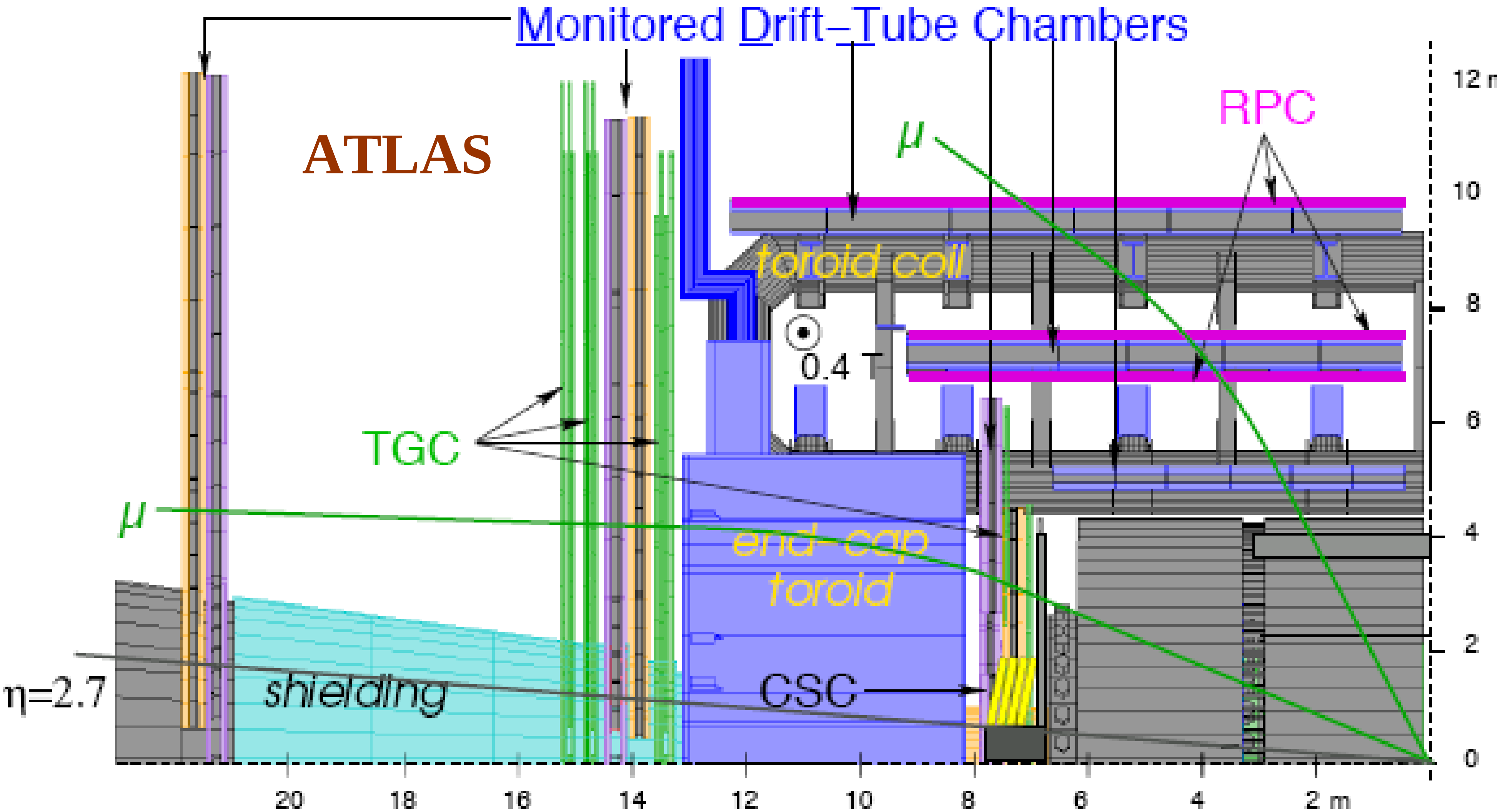




Assembly of MDT Chambers (Frascati, IT)



ATLAS muon system

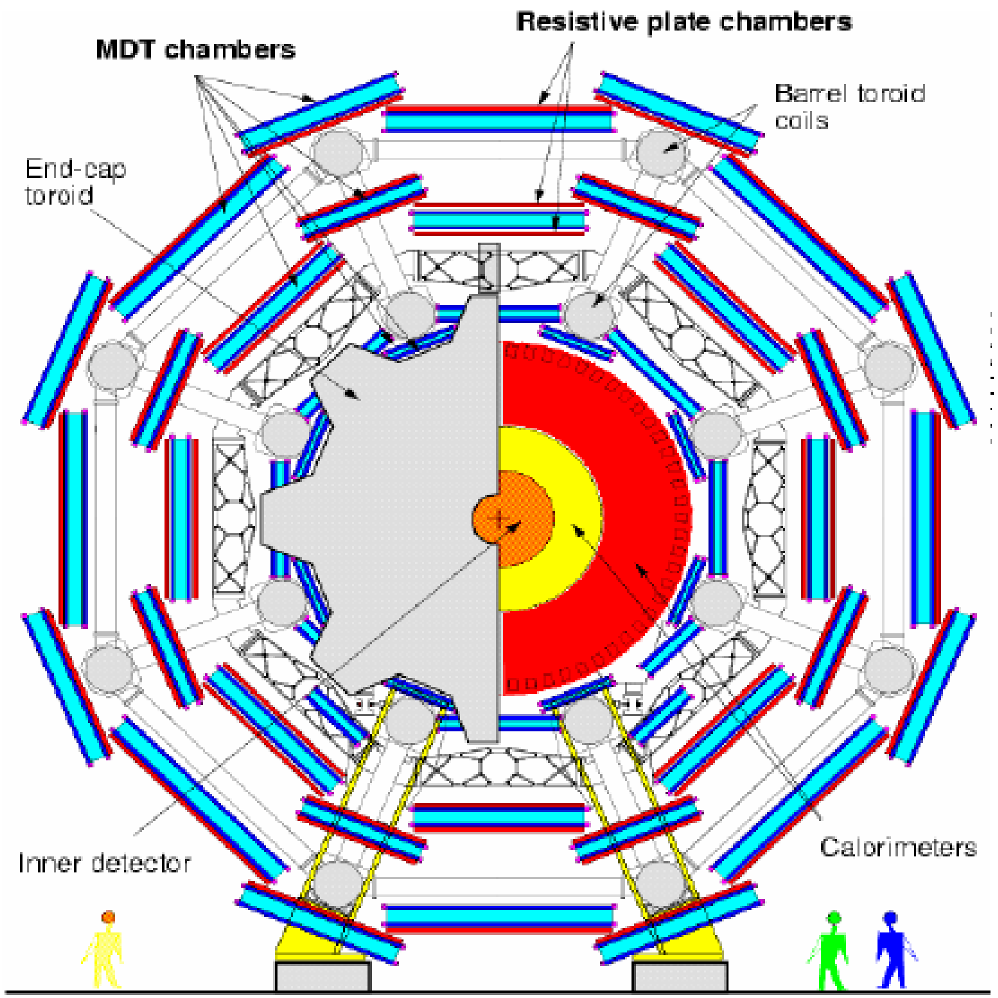


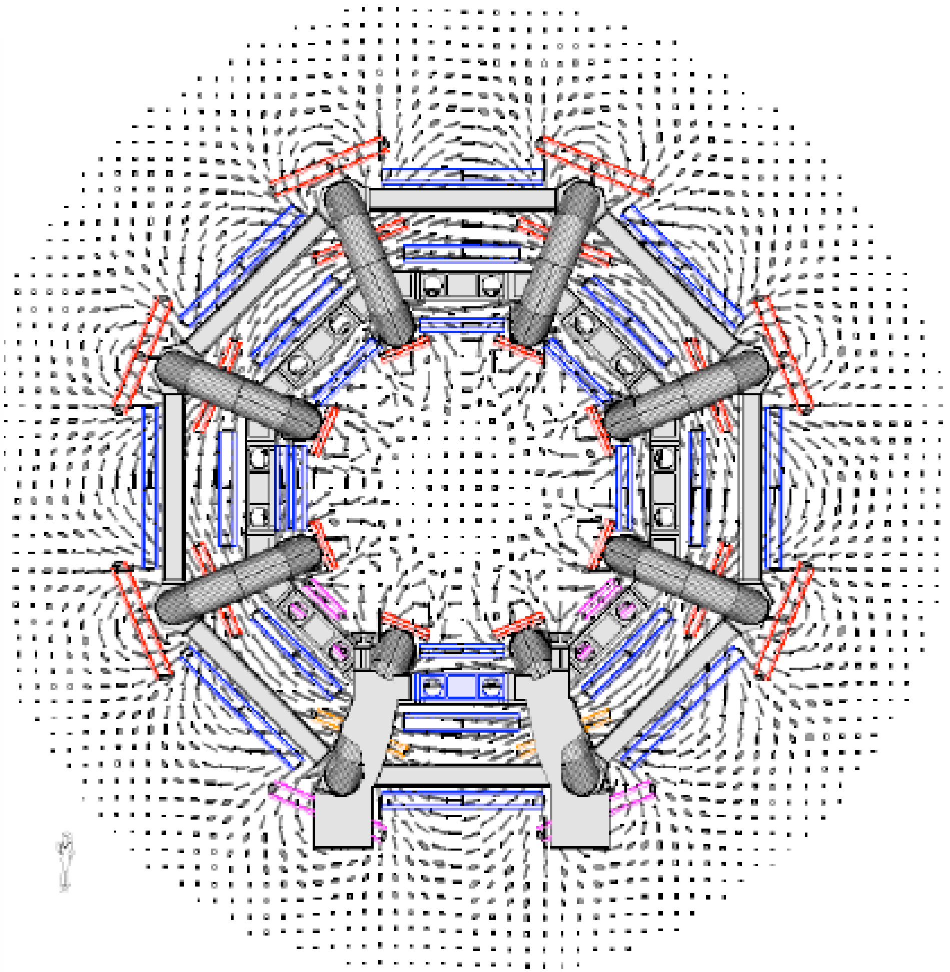
ATLAS muon spectrometer

" **Excellent stand-alone capabilities and coverage in open geometry**

" **Complicated geometry and field configuration (large fluctuations in acceptance and performance over full potential $\eta \times \varphi$ coverage ($|\eta| < 2.7$)**

ATLAS Barrel muon system in the toroid field

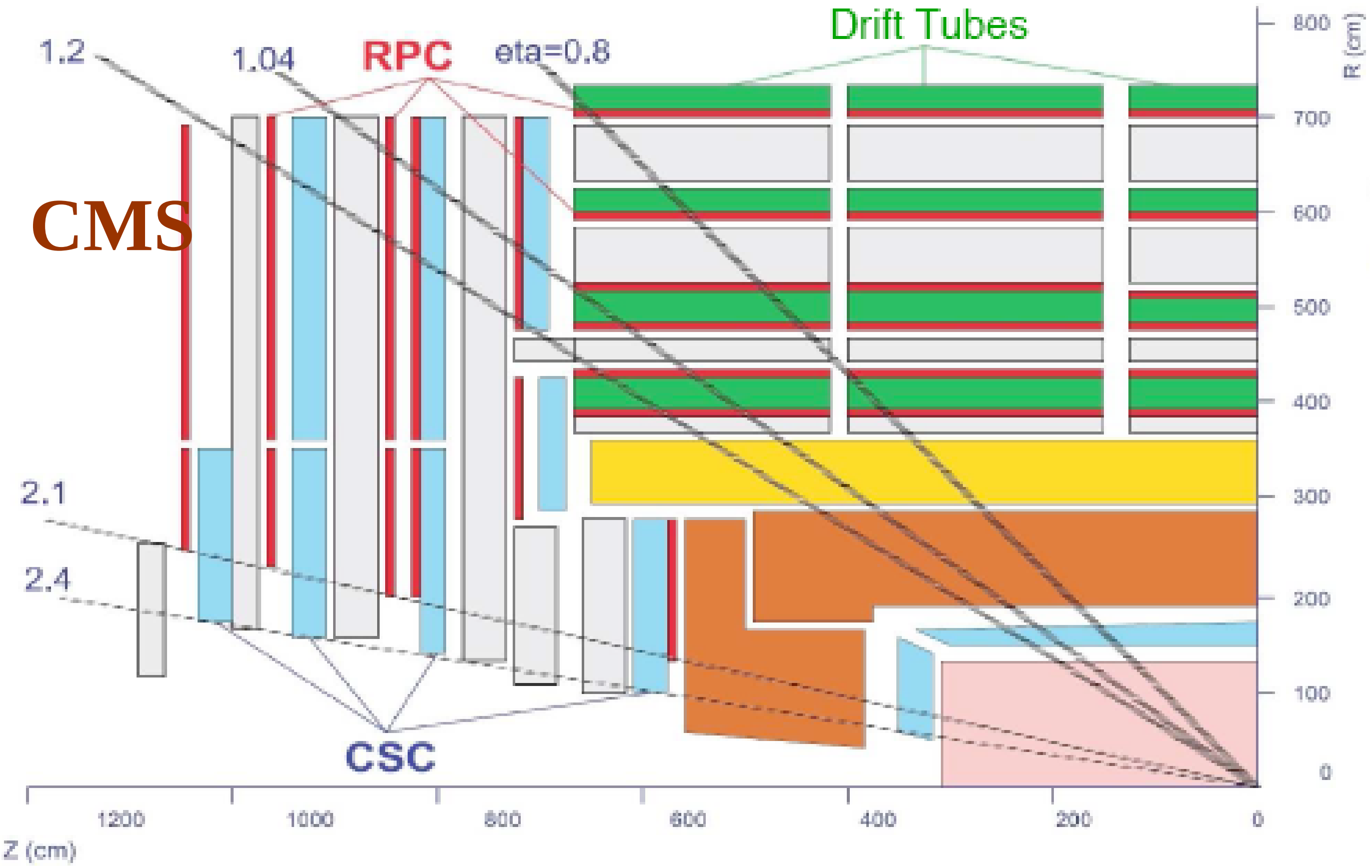


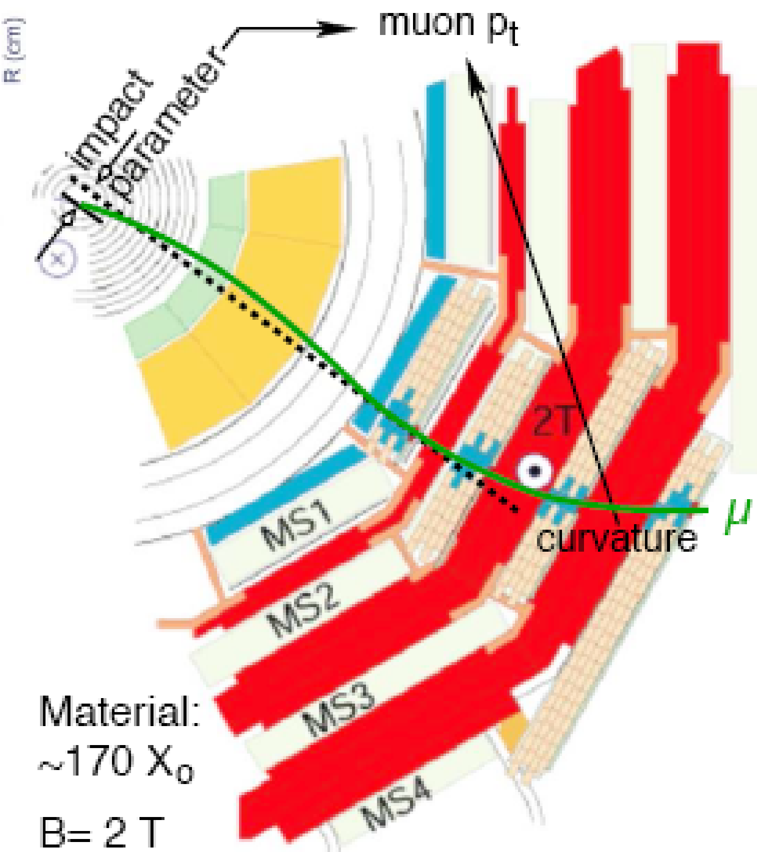


Detailed field map needed!

CMS muon system

CMS





CMS muon spectrometer

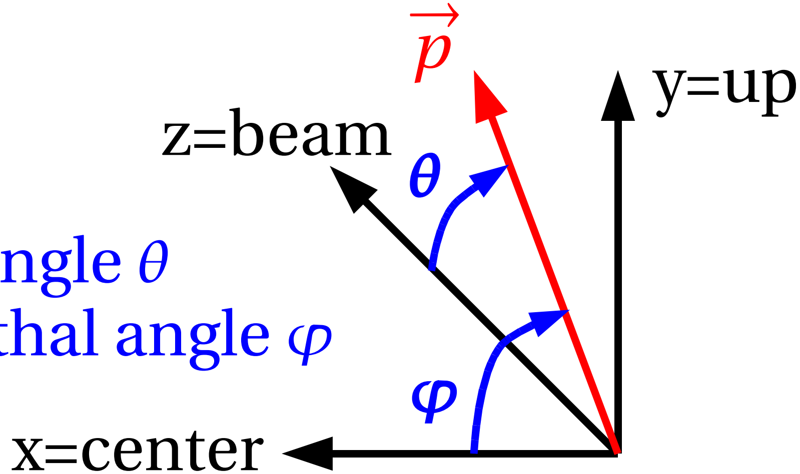
" Superior combined momentum resolution in the central region with silicon tracker.

" Limited stand-alone resolution and trigger (at very high luminosities) due to multiple scattering in iron

" Degraded overall resolution in the forward regions (|η| > 2.0) where solenoid bending power becomes insufficient

Pseudo-rapidity

Polar angle θ
Azimuthal angle φ



$$p_t = \sqrt{p_x^2 + p_y^2}$$

$$p = \sqrt{p_x^2 + p_y^2 + p_z^2}$$

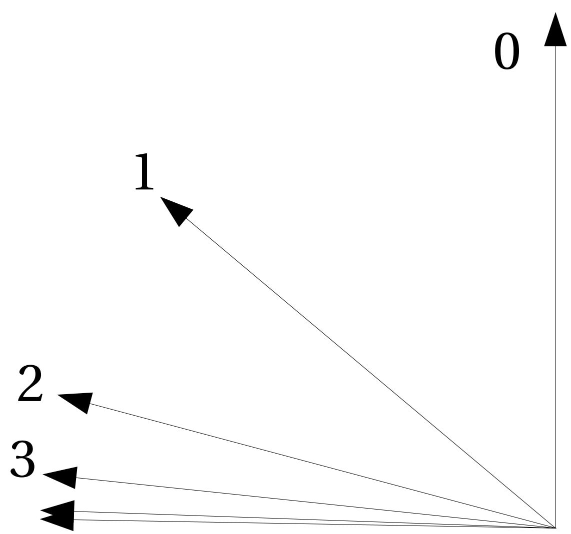
$$\varphi = \arctan(p_y / p_x) \in [-\pi, \pi]$$

$$\theta = \arccos(p_z / p) \in [0, \pi]$$

pseudo-rapidity: $\eta = -\ln(\tan(\theta/2))$

(soft hadron production: dn/dη ≈ const. 'central rapidity plateau')

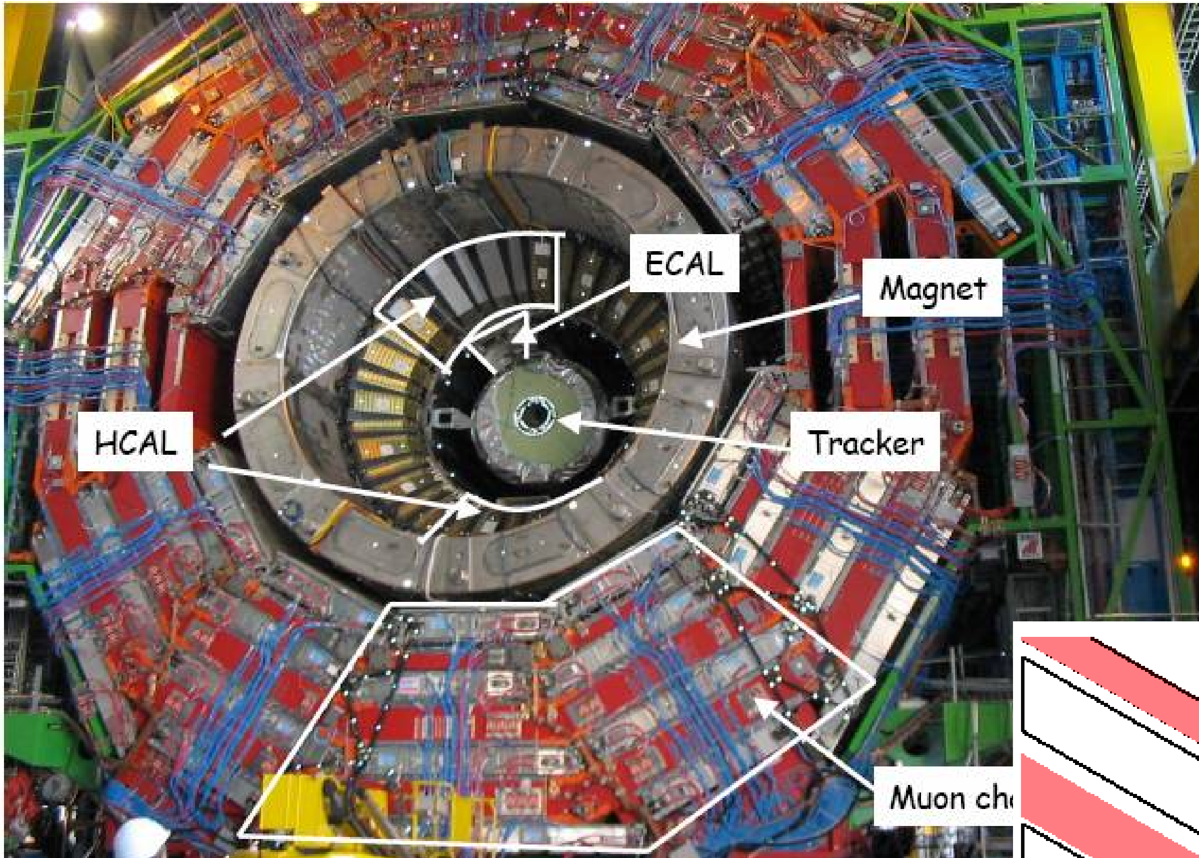
η	θ
0	90°
±1	±40°
±2	±15°
±3	±6°
±4	±2°
±5	±1°

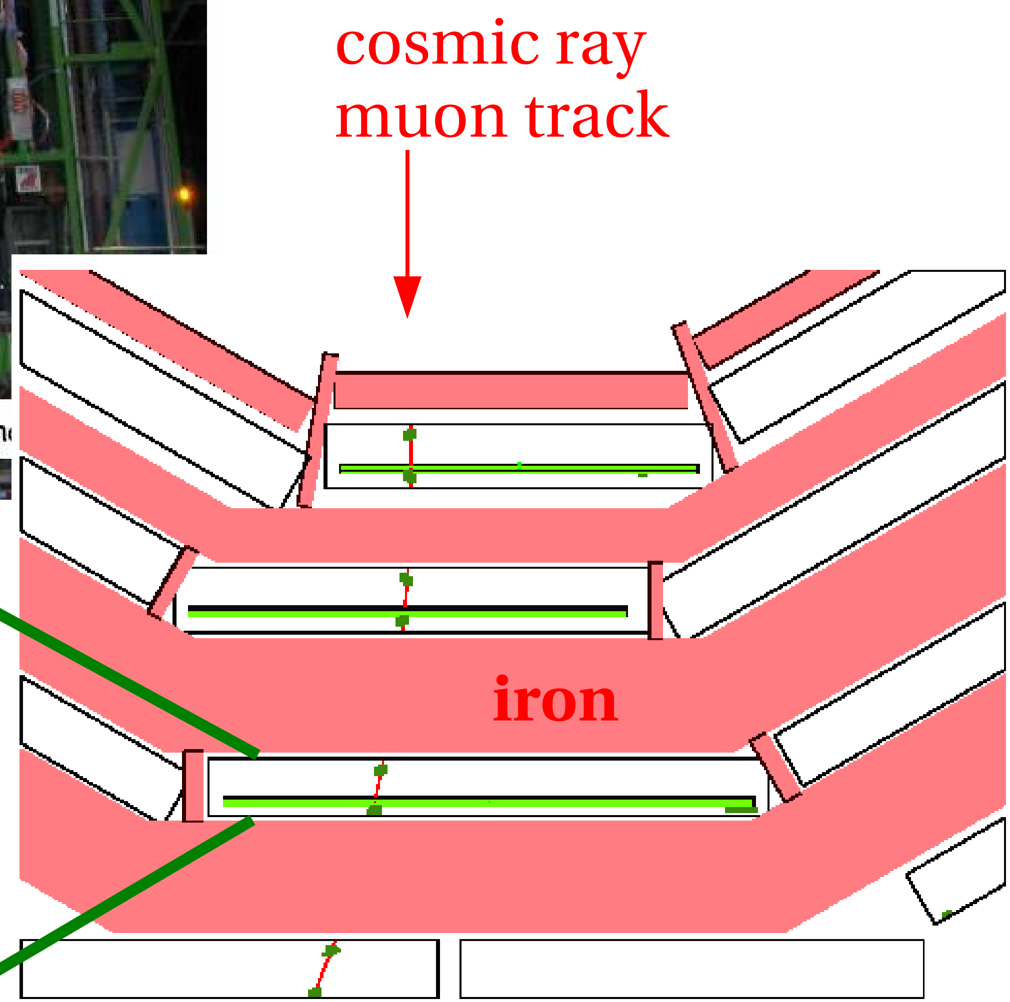


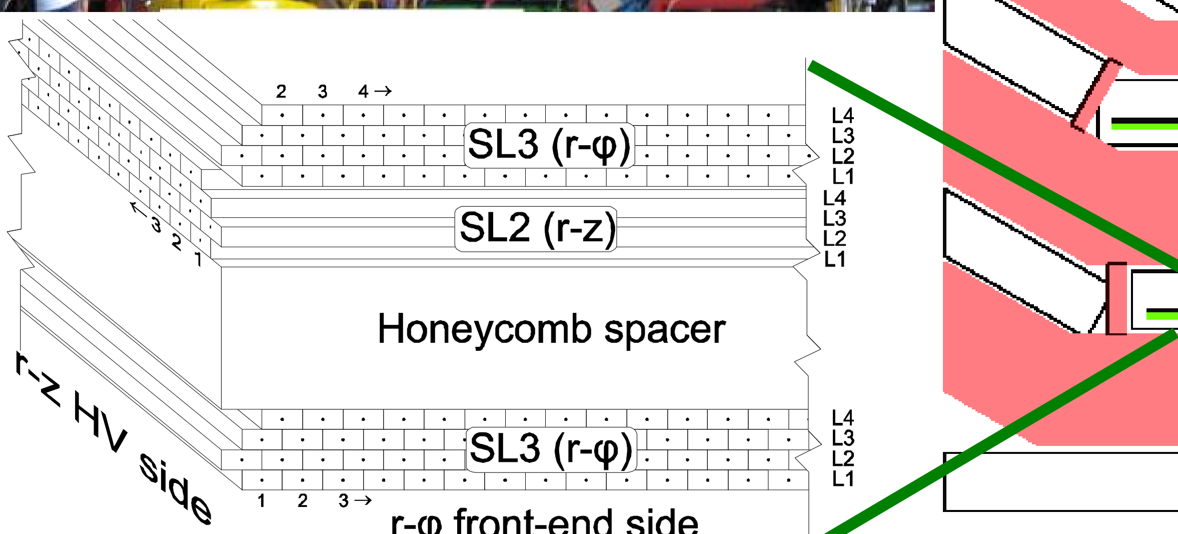
rapidity: $y = \frac{1}{2} \ln\left(\frac{E + p_z}{E - p_z}\right)$

$y \approx \eta$ for $p \gg m$

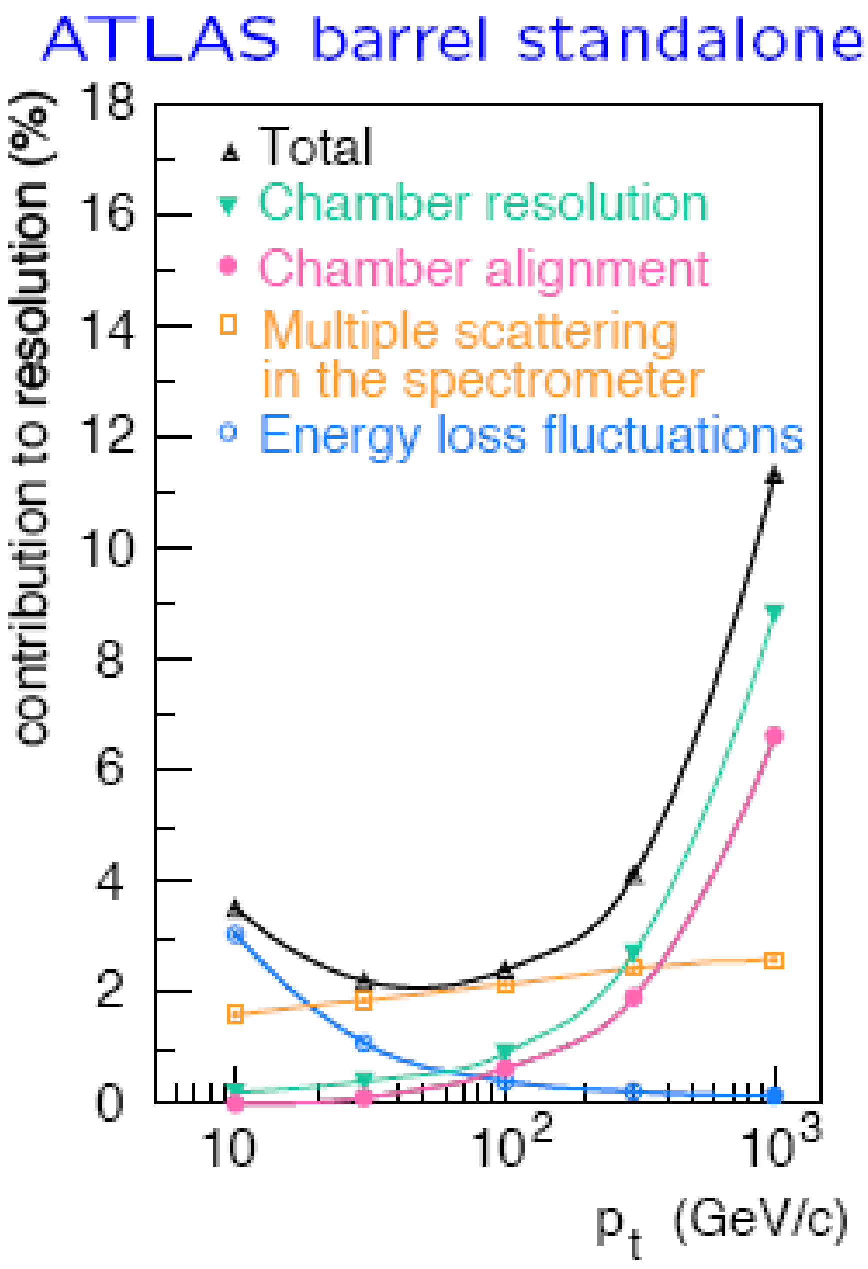
CMS Muon chambers

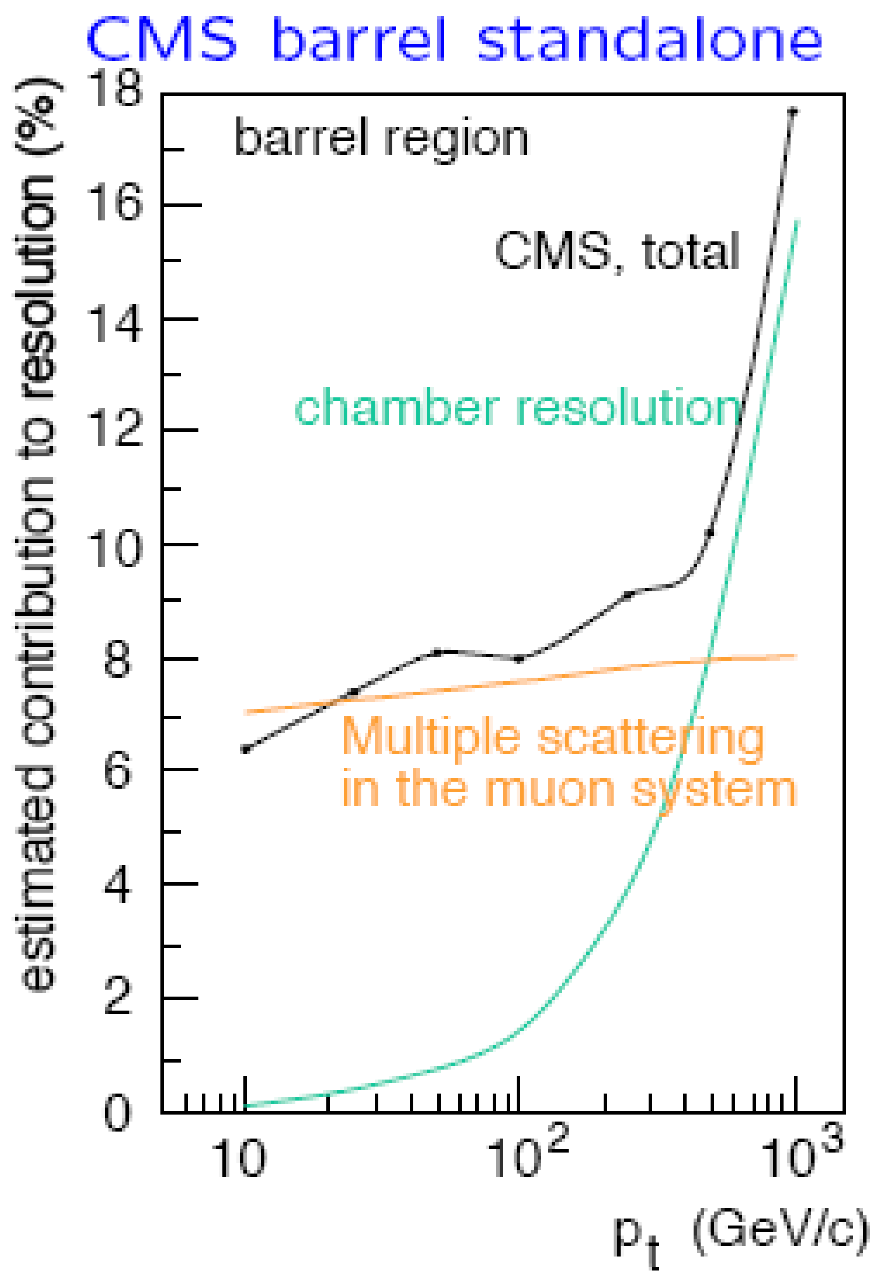






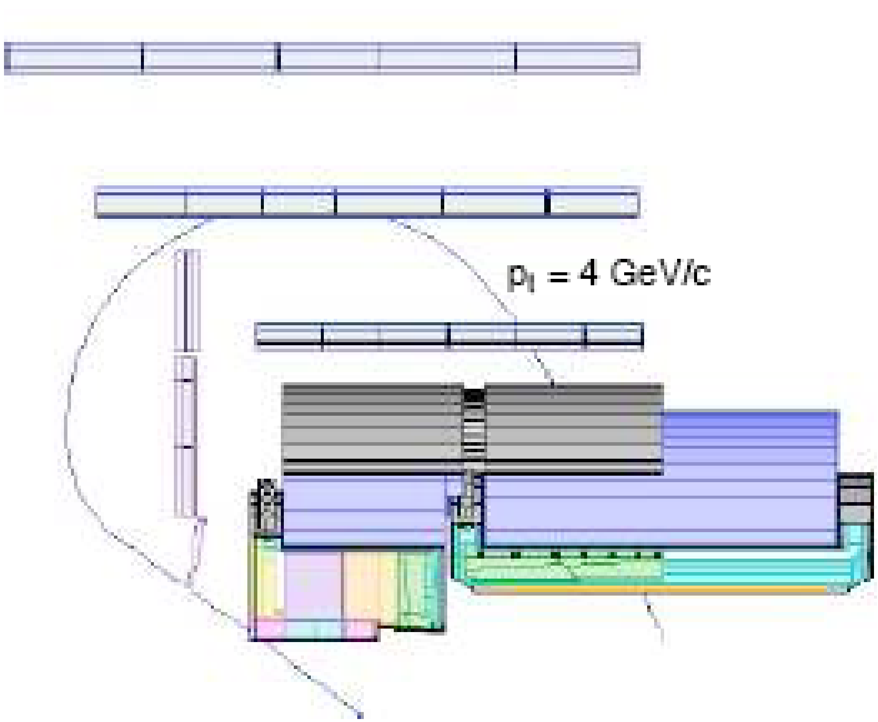
ATLAS and CMS: muon momentum resolution



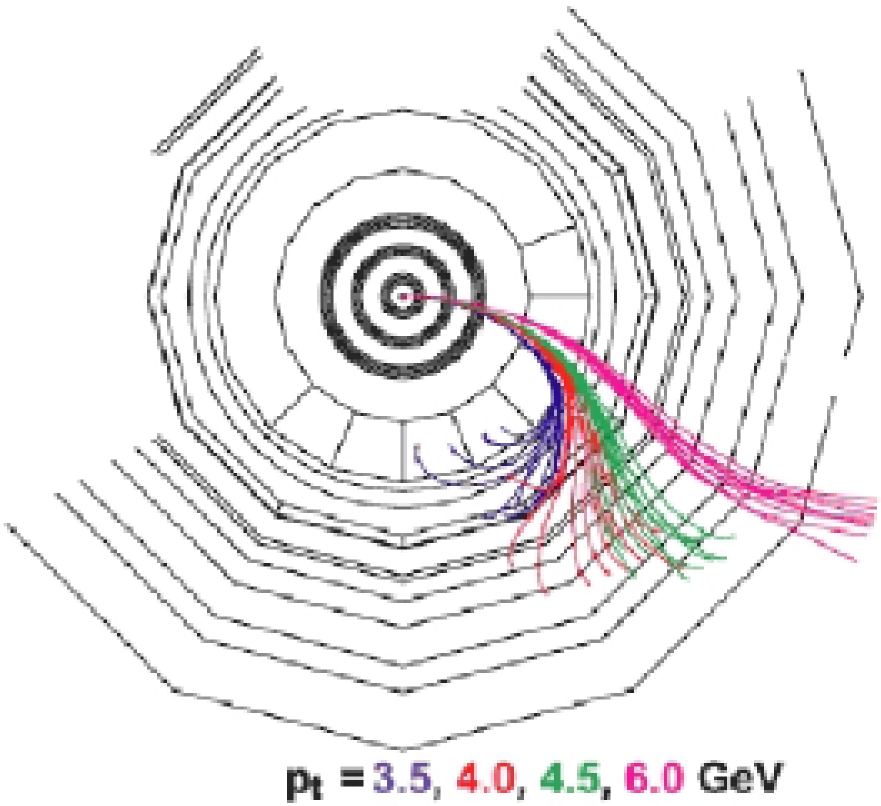


low-p_T muons

ATLAS



CMS



Requirements for muon identification and reconstruction at low p_T

" Identify track stub in first layer of muon system

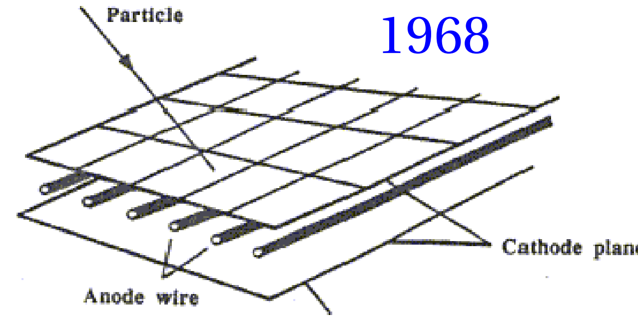
" Check for minimum ionising signals in last layers of hadron calorimeter

" Match as precisely as feasible (within limitations due to large MS and energy loss in calorimetry) measured track in inner detector with track stub in muon system

Multi-wire proportional chambers MWPCs

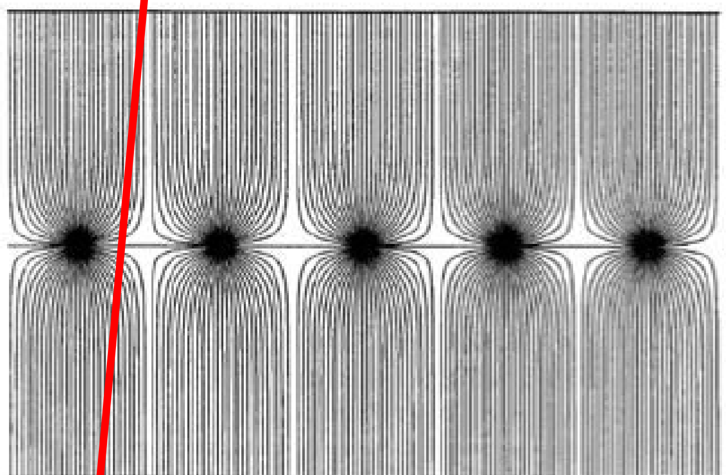
string many anode wires between cathode planes:

1968



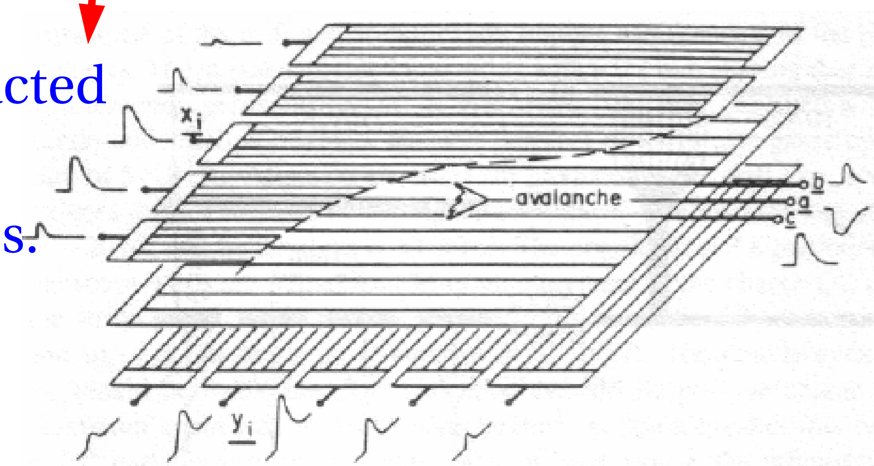
field lines:

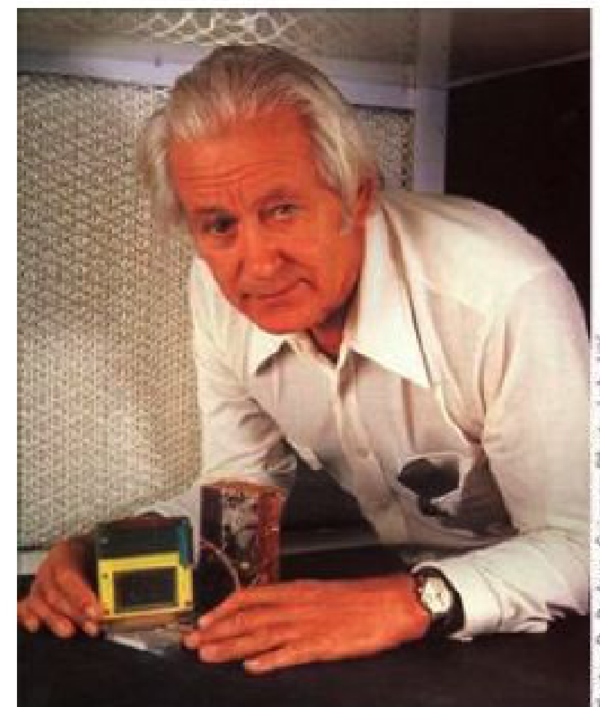
±3 mm. max. drift time can be < 50 ns. Fast!



2 mm wire spacing
resolution ~ 0.5 mm

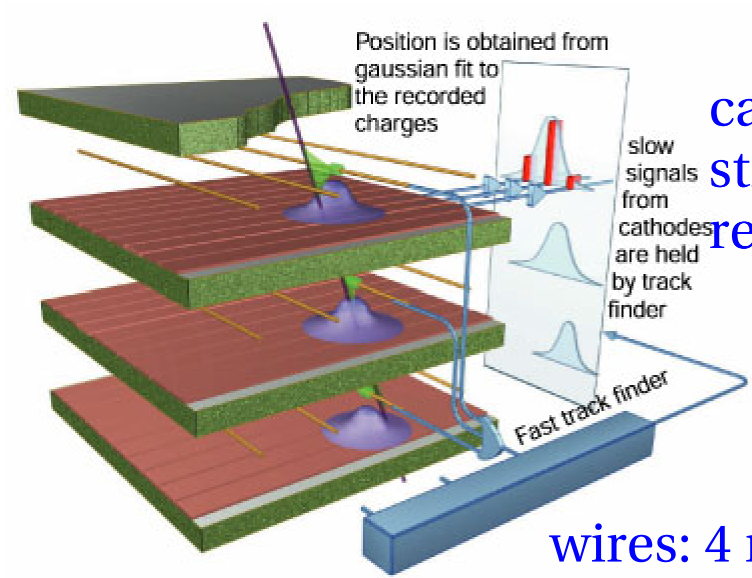
signals can be extracted from the wire and both cathode planes. Cathodes can be segmented.





Georges Charpak, CERN
Nobel prize 1992

CMS Cathode Strip Chambers

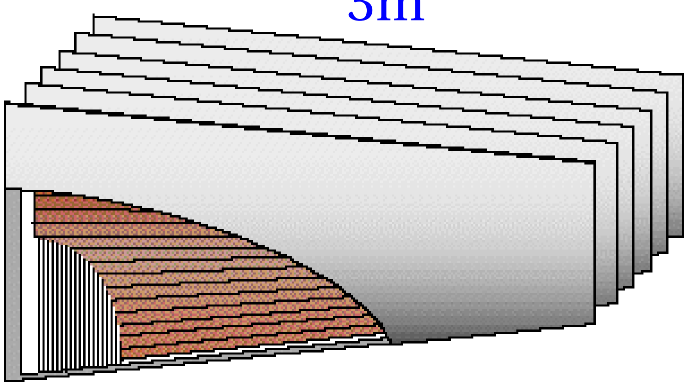


cathode strips: 0.1 mm resolution.

wires: 4 ns timing for the trigger

3m

1m

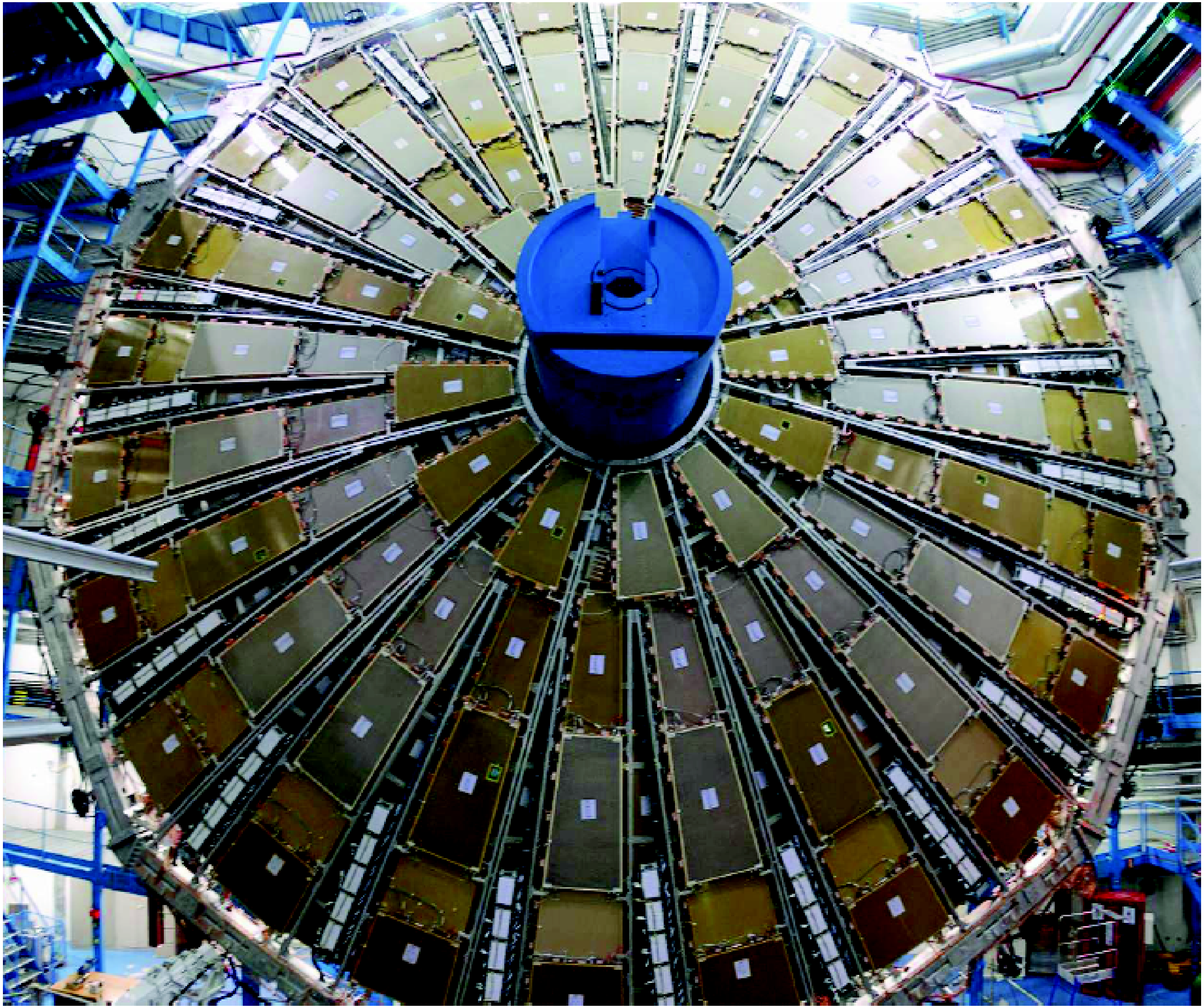


6 planes/chamber
468 chambers

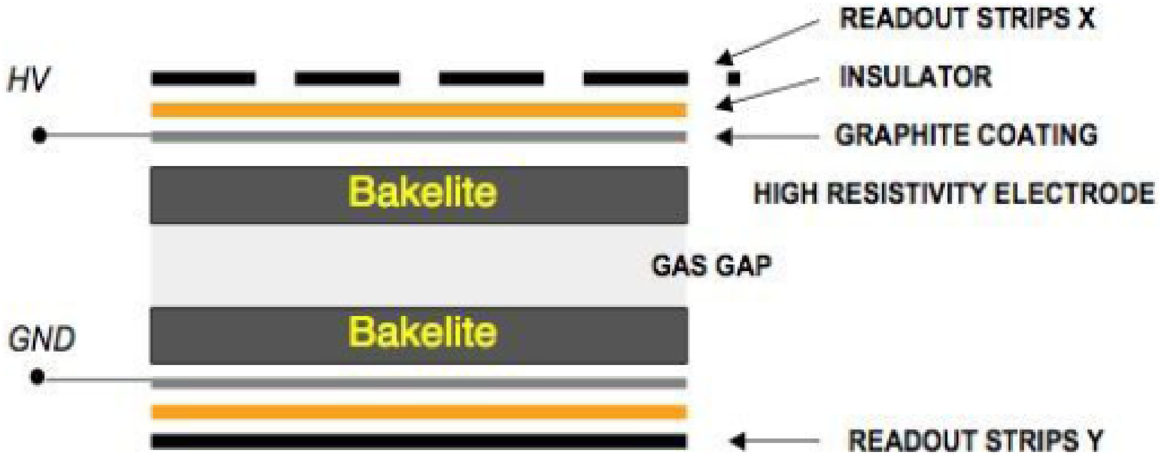
6000 m^2
450k channels

ATLAS muon chamber wheel

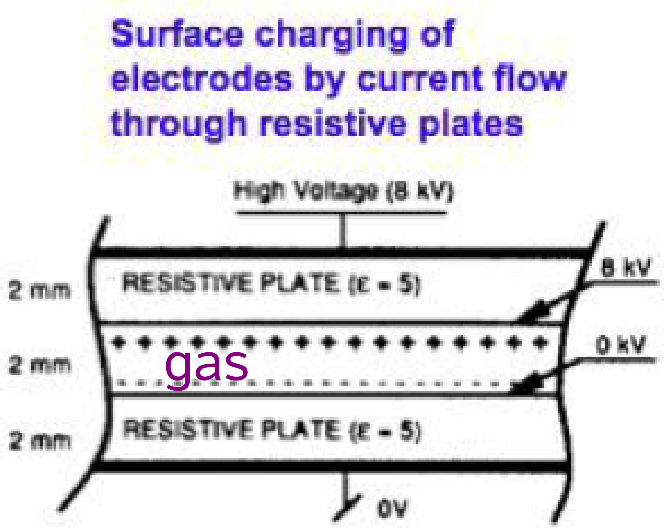
Ø25m



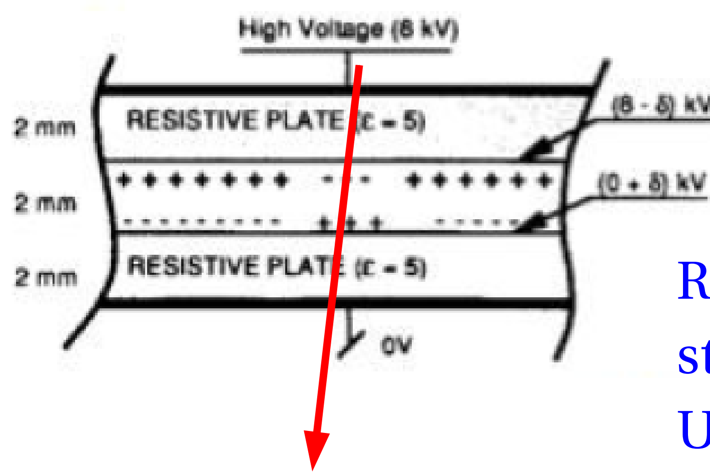
RPC Resistive plate chambers



no wires!
2 Bakelite plates separated by a gas gap and 8kV.



a particles causes local discharge which induces signals in the readout strips.



RPCs are fast and stable at high rate. Used as trigger chambers.

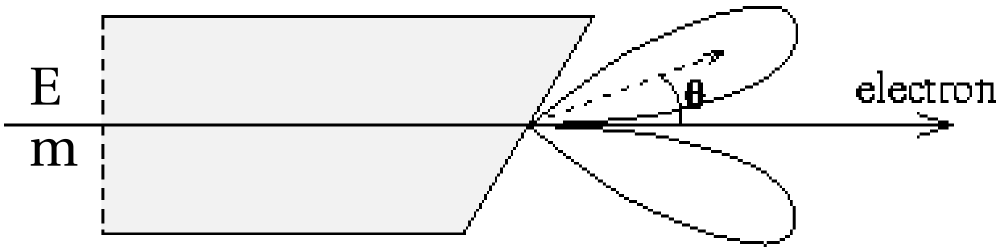
Transition Radiation

- Relativistic particles passing through an interface radiate.

polypropylene
$(C_3H_6)_x$
$\hbar\omega_P = 20$ eV

air
$\hbar\omega_P = 0.7$ eV

$\omega_p^2 = n_e e^2 / m_e \varepsilon_0$

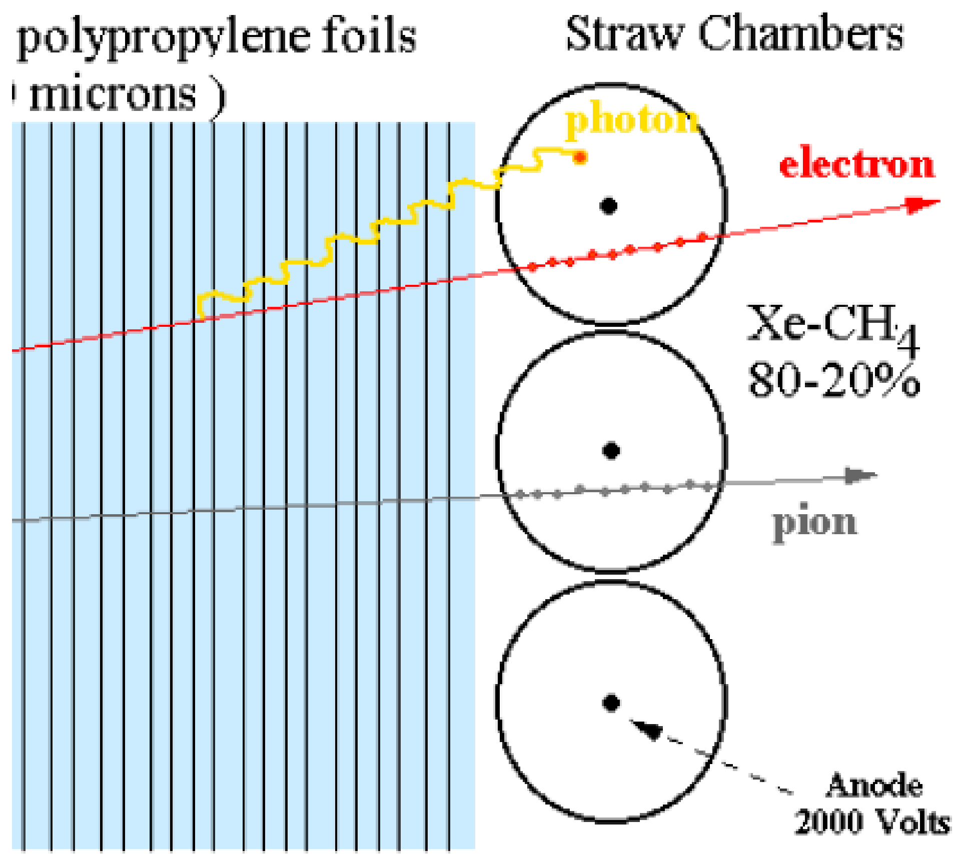


$\gamma = E/m$

$\theta \approx 1/\gamma$

$\hbar\omega \approx$ 10-30 keV
for $\gamma \approx$ 10'000
$N \sim \gamma \;\; \hbar\omega_P$
depends on particle!

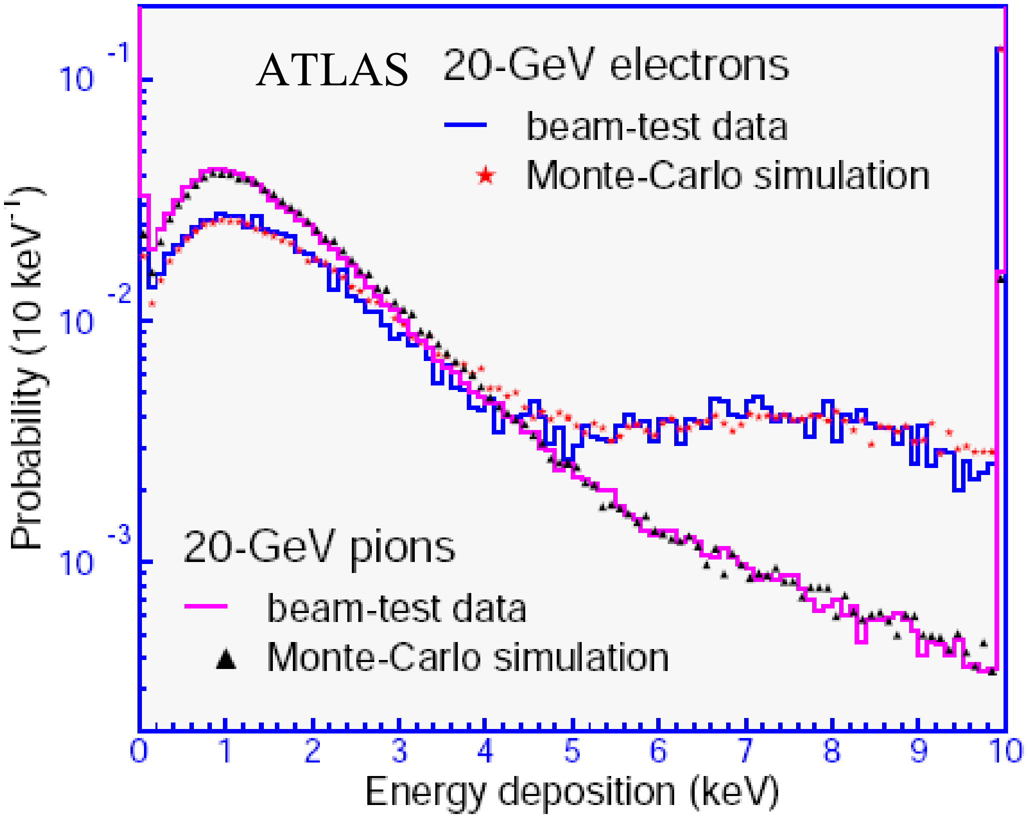
$$\frac{d^2W}{d\omega d\theta} = \frac{2\alpha\theta^3}{\pi}\left(\frac{1}{\frac{1}{\gamma^2}+\theta^2+\frac{\omega_1^2}{\omega^2}} - \frac{1}{\frac{1}{\gamma^2}+\theta^2+\frac{\omega_2^2}{\omega^2}}\right)^2$$

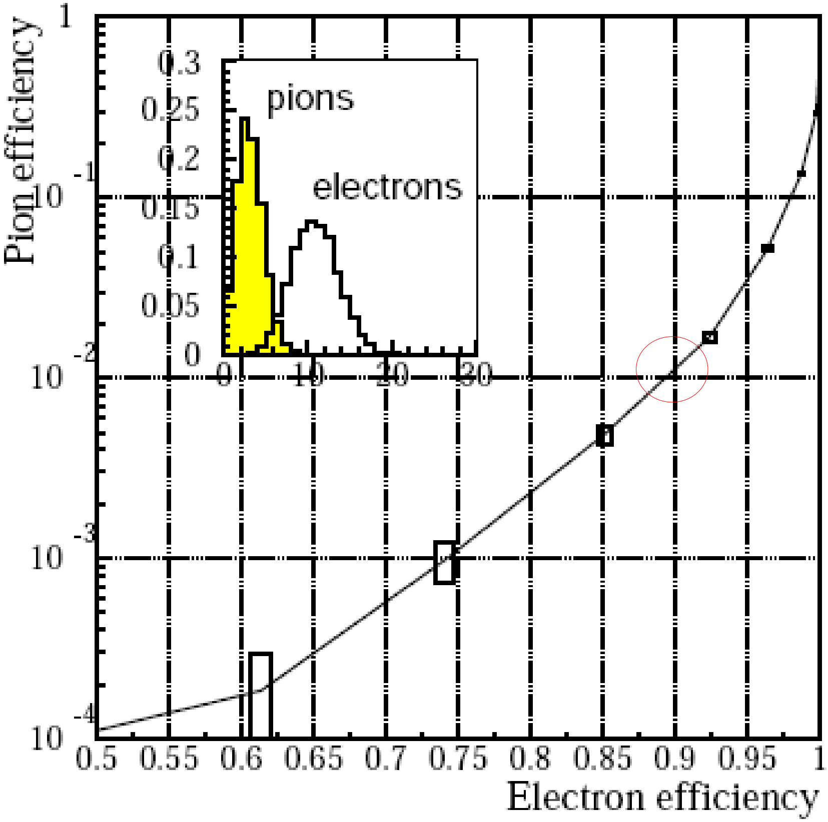


1% probability per foil

Xe gas to absorb X-ray.

e/π separation using TRT

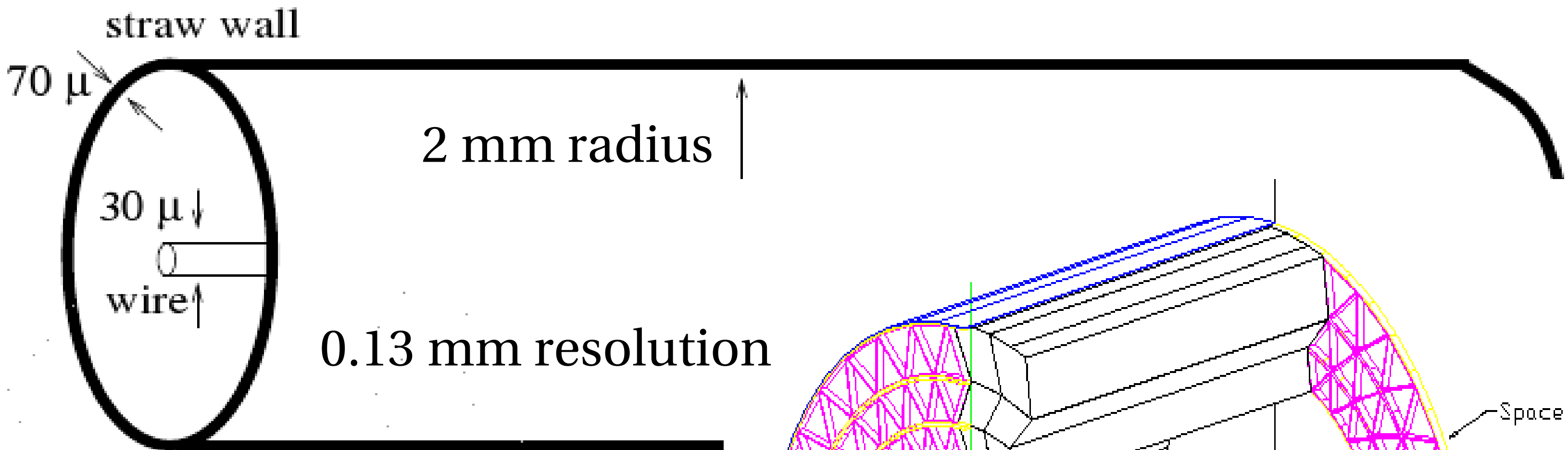


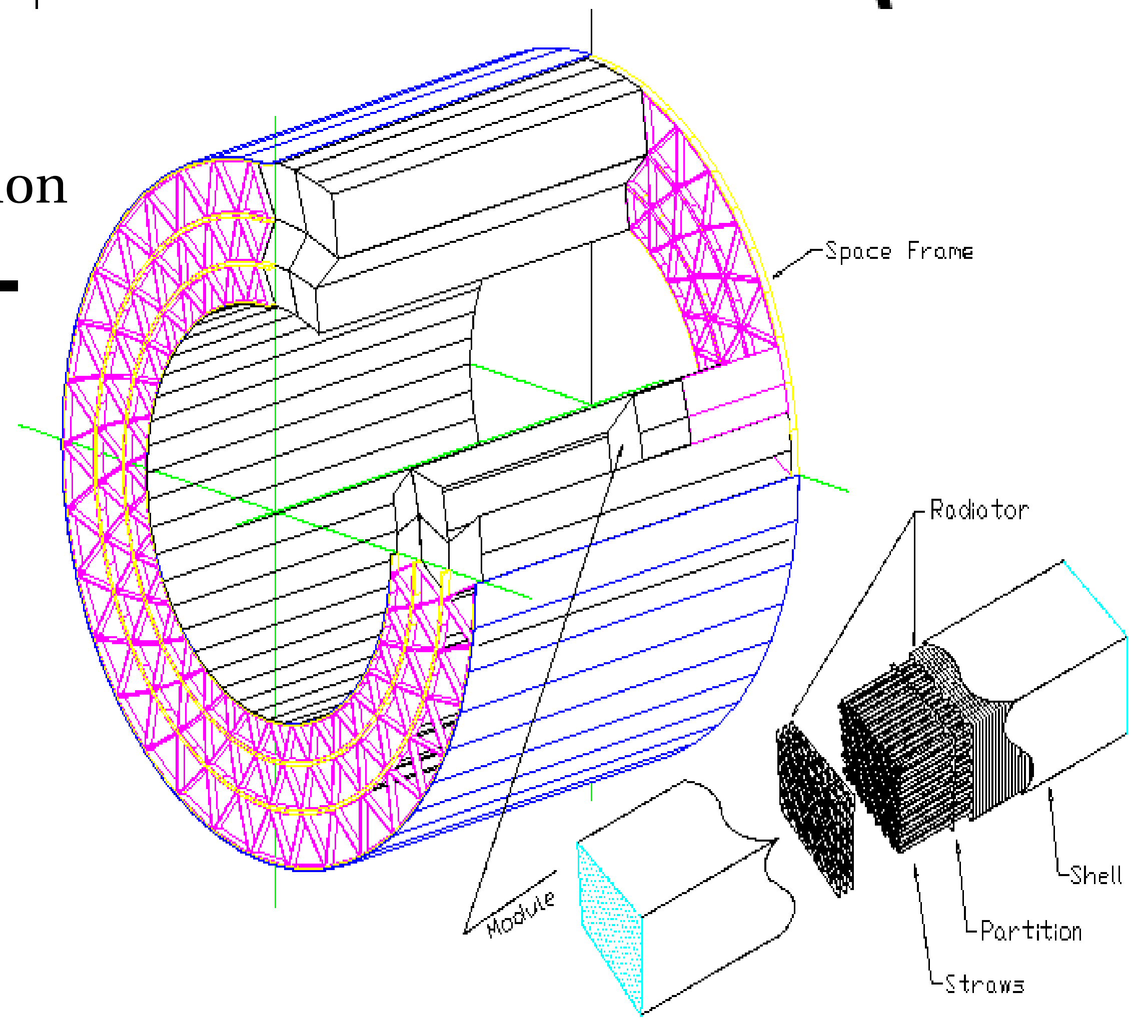


pion reduction by factor 75
for 90% electron efficiency

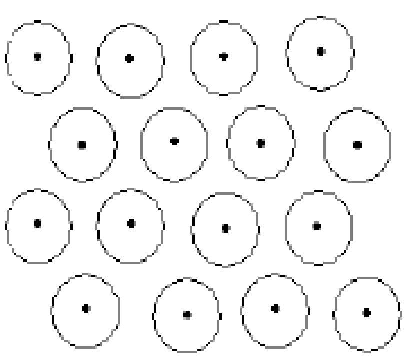
ATLAS Transition Radiation Tracker

built from straw tubes:



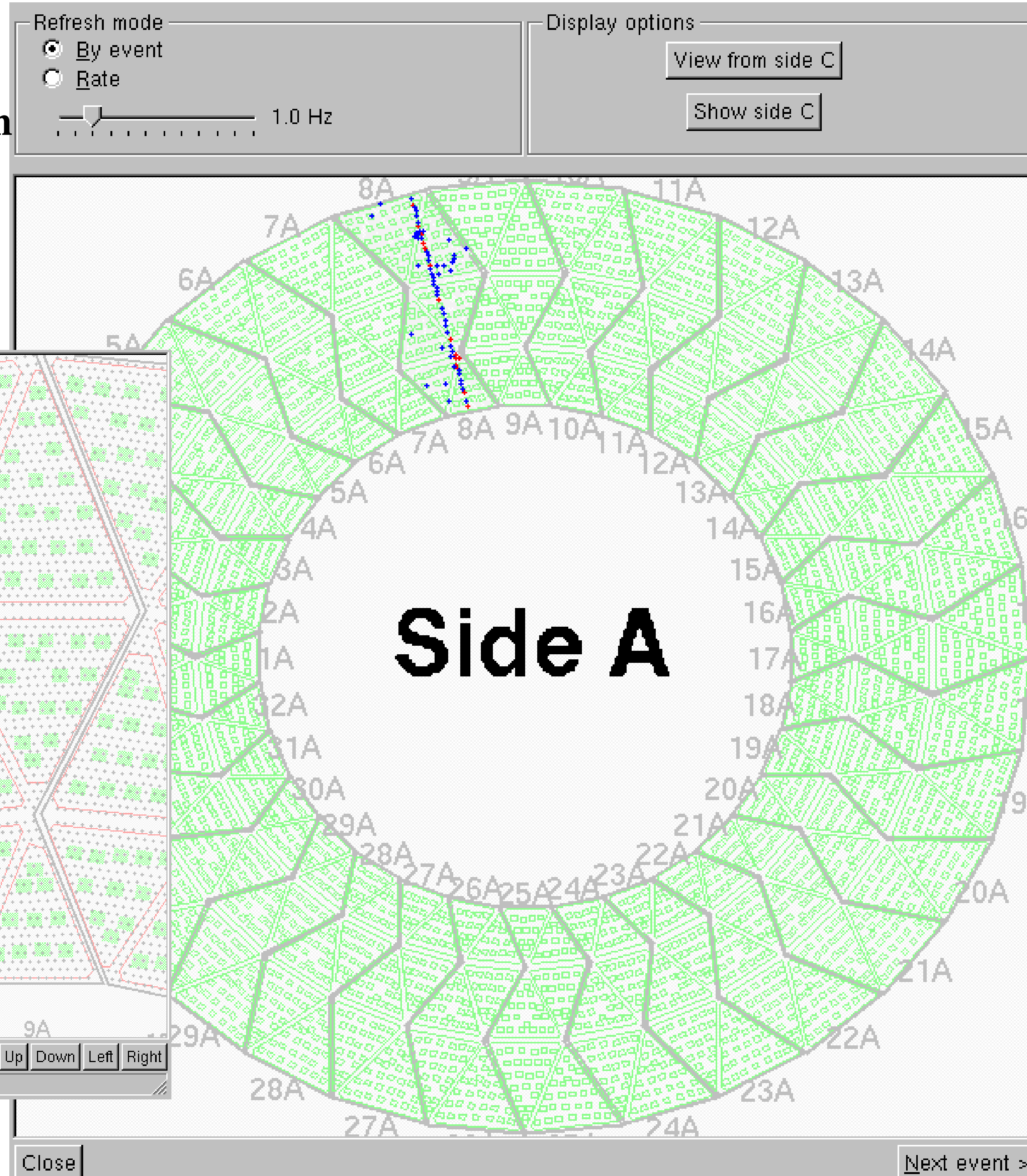


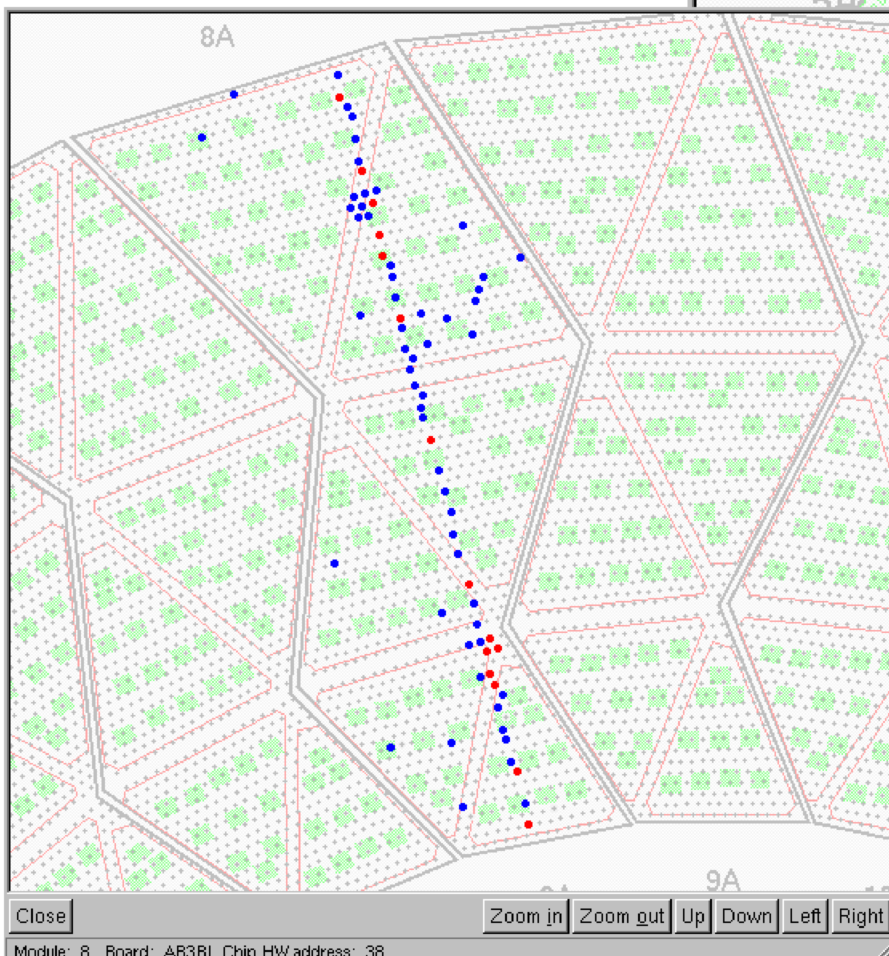
Straw gas mixture :
Xe(70%) CO_2(27%) O_2(3%)



Radiator foils are placed between the straws

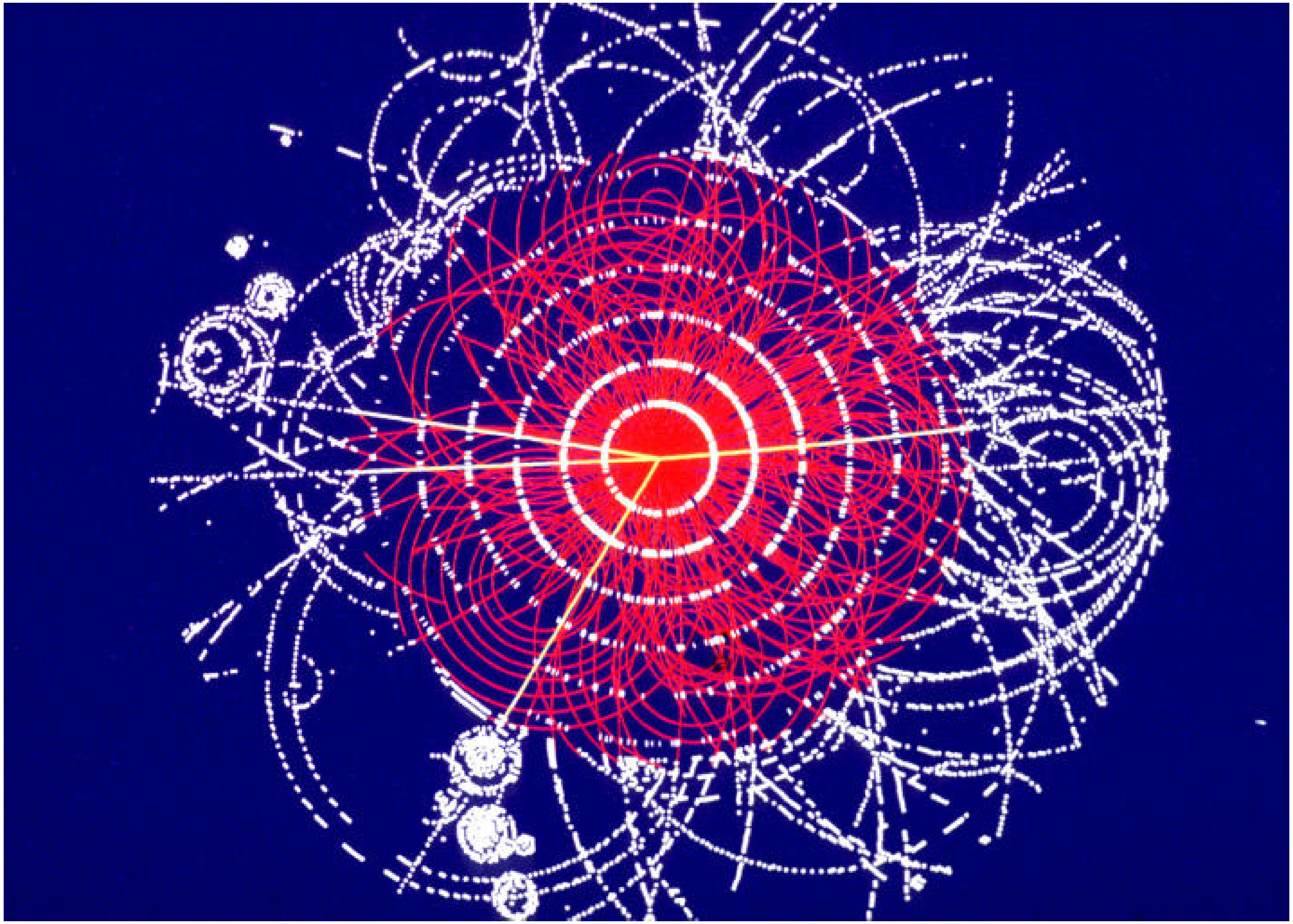
First cosmic ray event seen in the Barrel TRT!





D. Pitzl, DESY

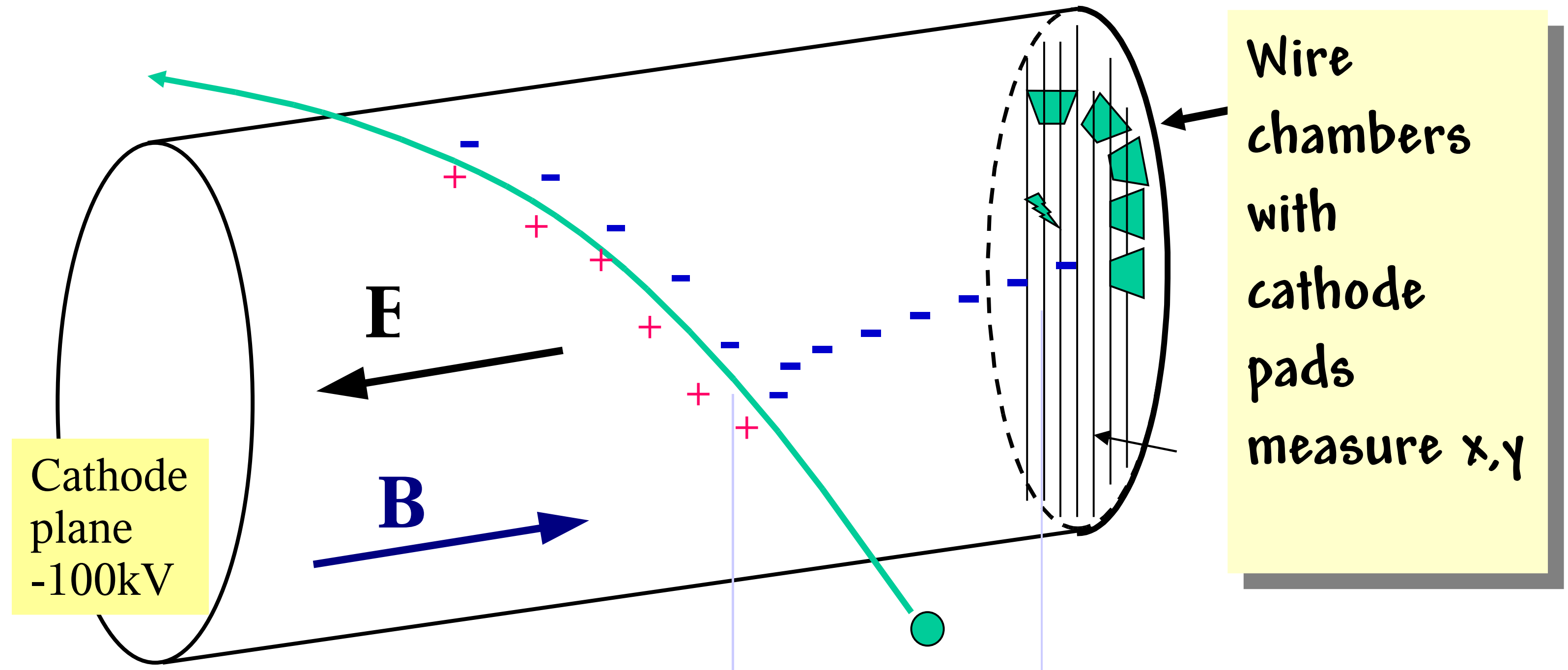
ATLAS inner tracking: silicon and straws

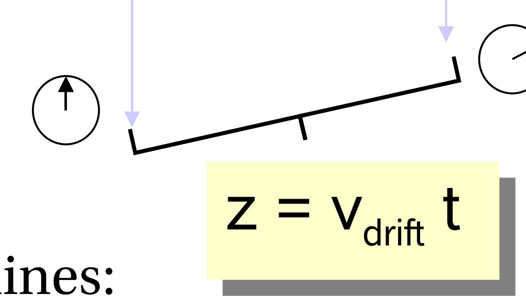


red = tracks found in the silicon layers.

white = hits in the straw tubes. 36 hits/track: good for pattern recognition.

Time Projection Chamber in a solenoid field





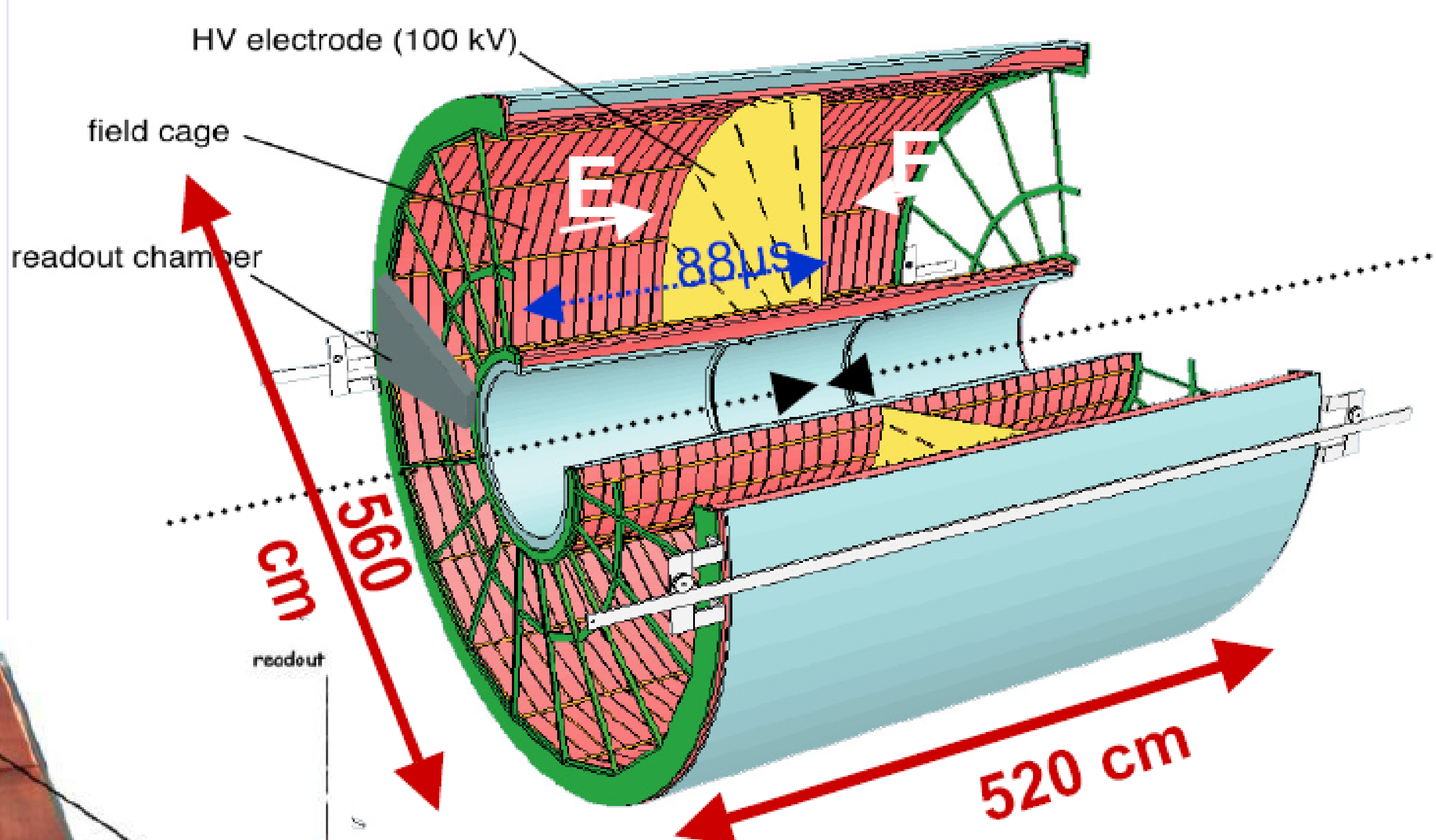
Separate two regions:

- ❑ Drift along z: 20-30 μs/m.
- ❑ Amplify at the end plate

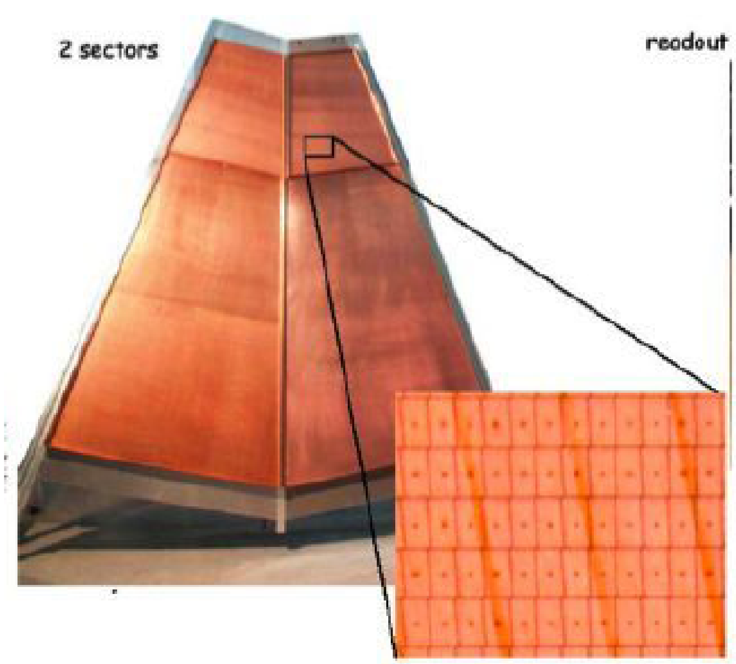
No material inside drift volume!

E||B: drifting electrons curl around B field lines: limited spread.

ALICE TPC

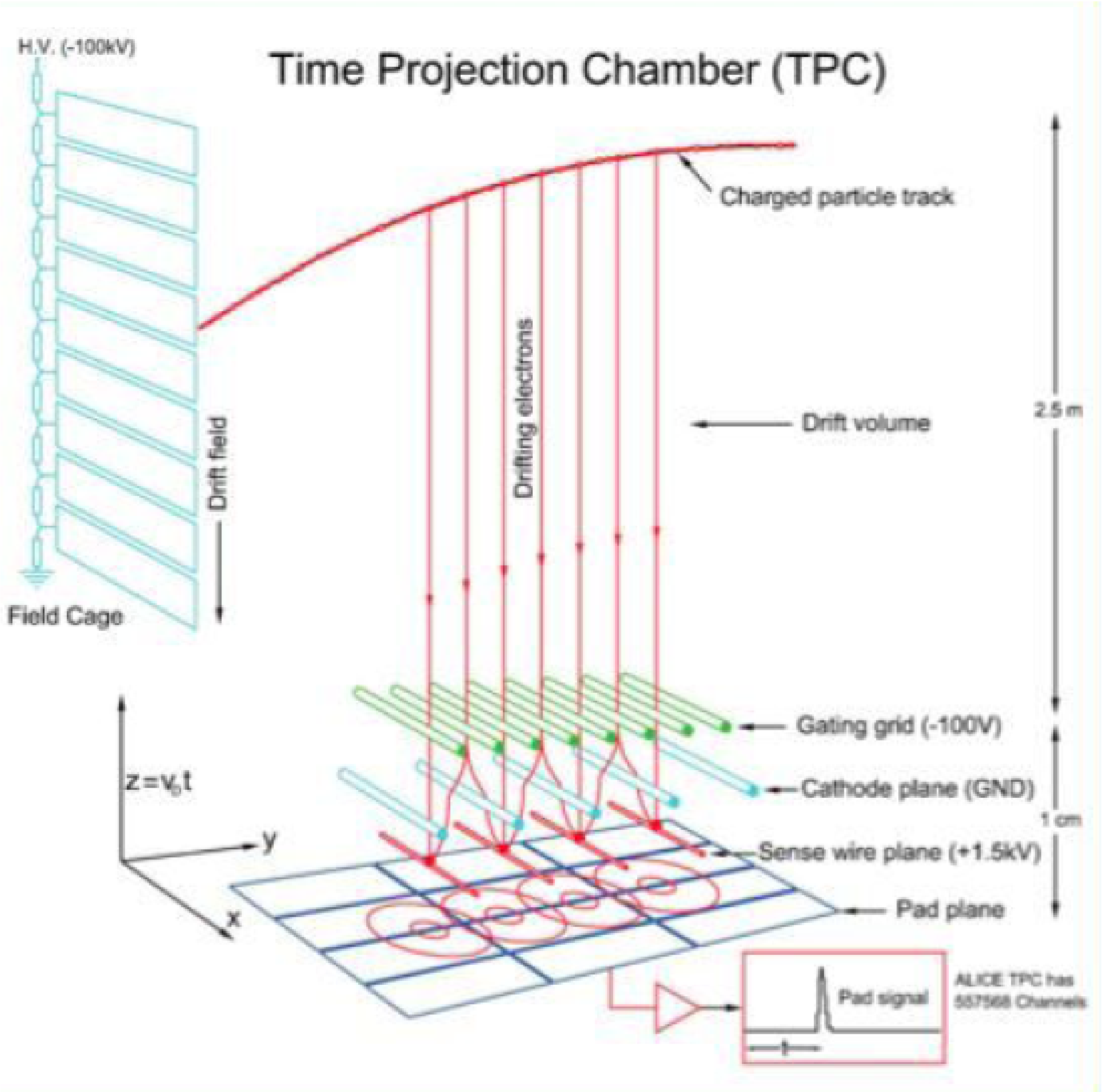


readout:



max. drift time 88 µs.
LHC operates with low luminosity
at ALICE (pp and PbPb).

TPC


Time Projection Chamber (TPC)
H.V. (-100kV)
Field Cage
Drift field
Charged particle track
Drifting electrons
Drift volume
2.5 m
Gating grid (-100V)
Cathode plane (GND)
Sense wire plane (+1.5kV)
Pad plane
1 cm
z=v_D t
y
x
Pad signal
t
ALICE TPC has
557568 Channels

Energy loss of charged particles in matter

- Charged particles loose energy in collisions with electrons in matter:

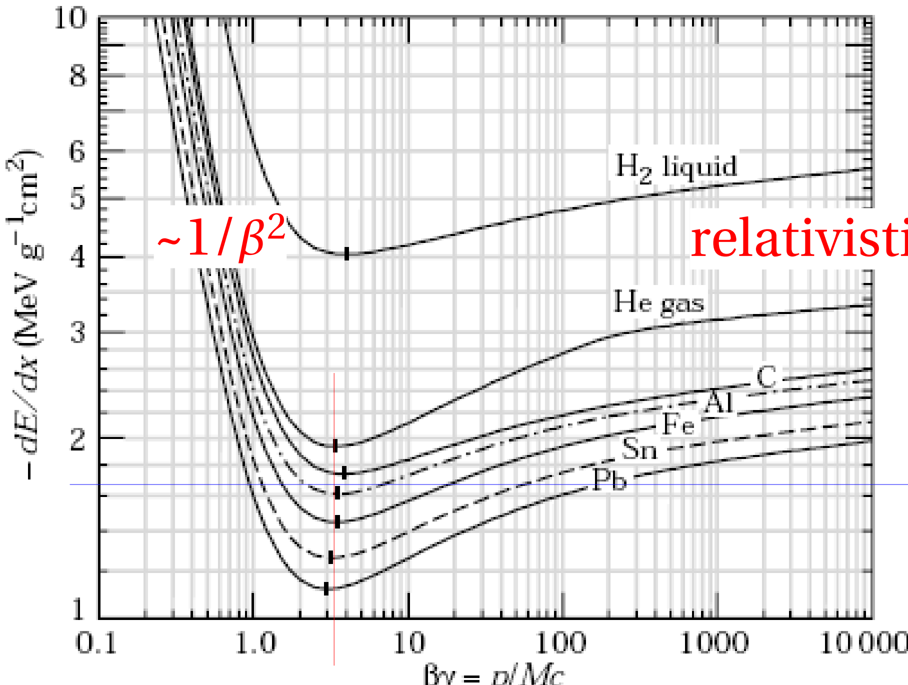
Bethe-Bloch:

$$\frac{1}{\rho}\frac{dE}{dx} = -4\pi N_A r_e^2 m_e c^2 \frac{Z}{A}\frac{z^2}{\beta^2}\left[\ln\left(\frac{2 m_e c^2 \beta^2\gamma^2}{\langle I\rangle}\right) - \beta^2 - \frac{\delta}{2}\right]$$

medium | particle | ionisation potential | density correction

$$\beta = \frac{v}{c} = \frac{cp}{E}, \quad \gamma = \frac{E}{mc^2}, \quad \beta\gamma = \frac{cp}{mc^2}$$

dE/dρx



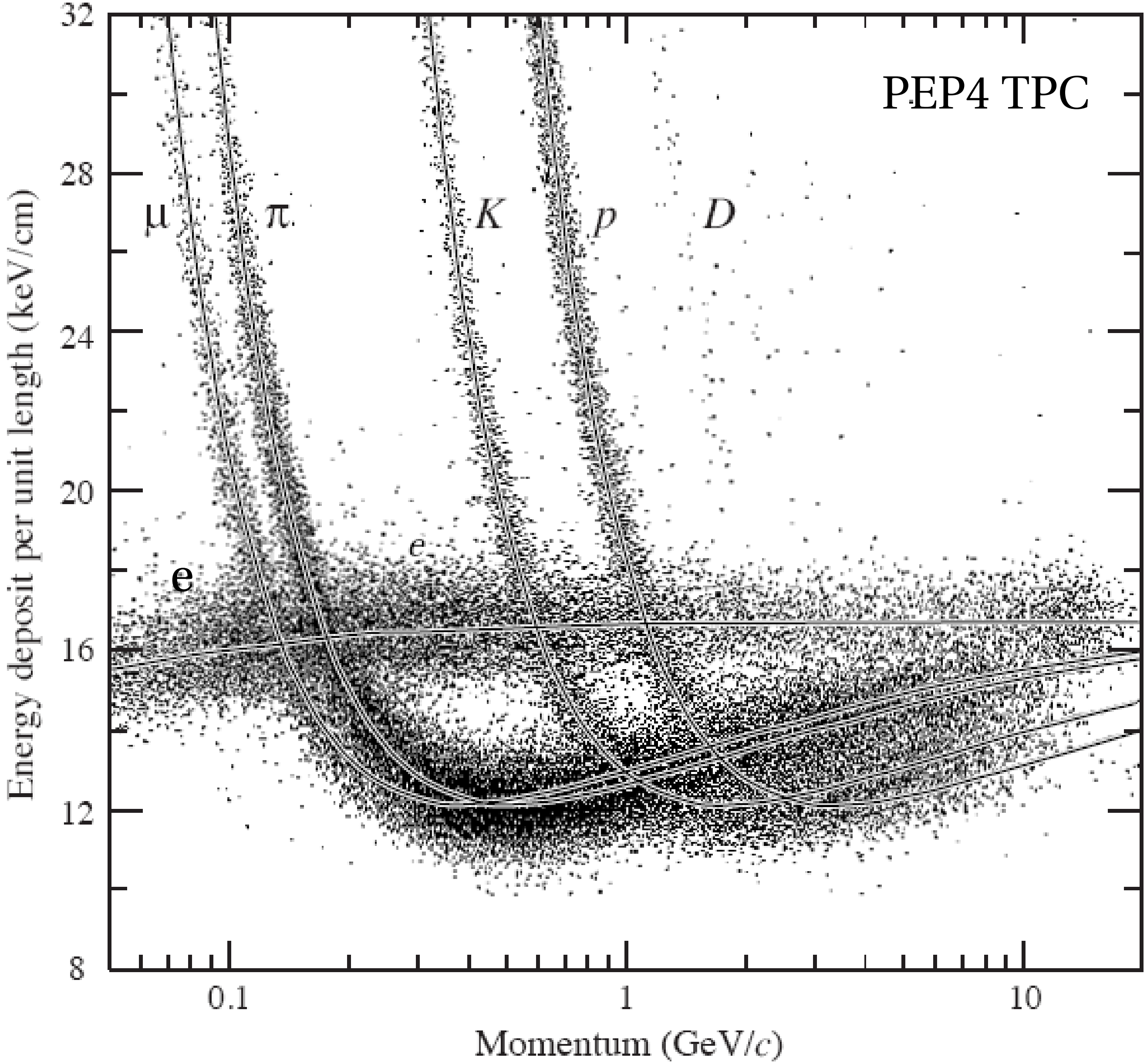
$\sim 1/\beta^2$

relativistic rise $\sim \ln(\gamma)$

in iron (ρ=7.9 g/cm^3):
dE/dx = 1.3 GeV/m.

minimum at
p/m ≈ 3.5 for all particles

dE/dx Data



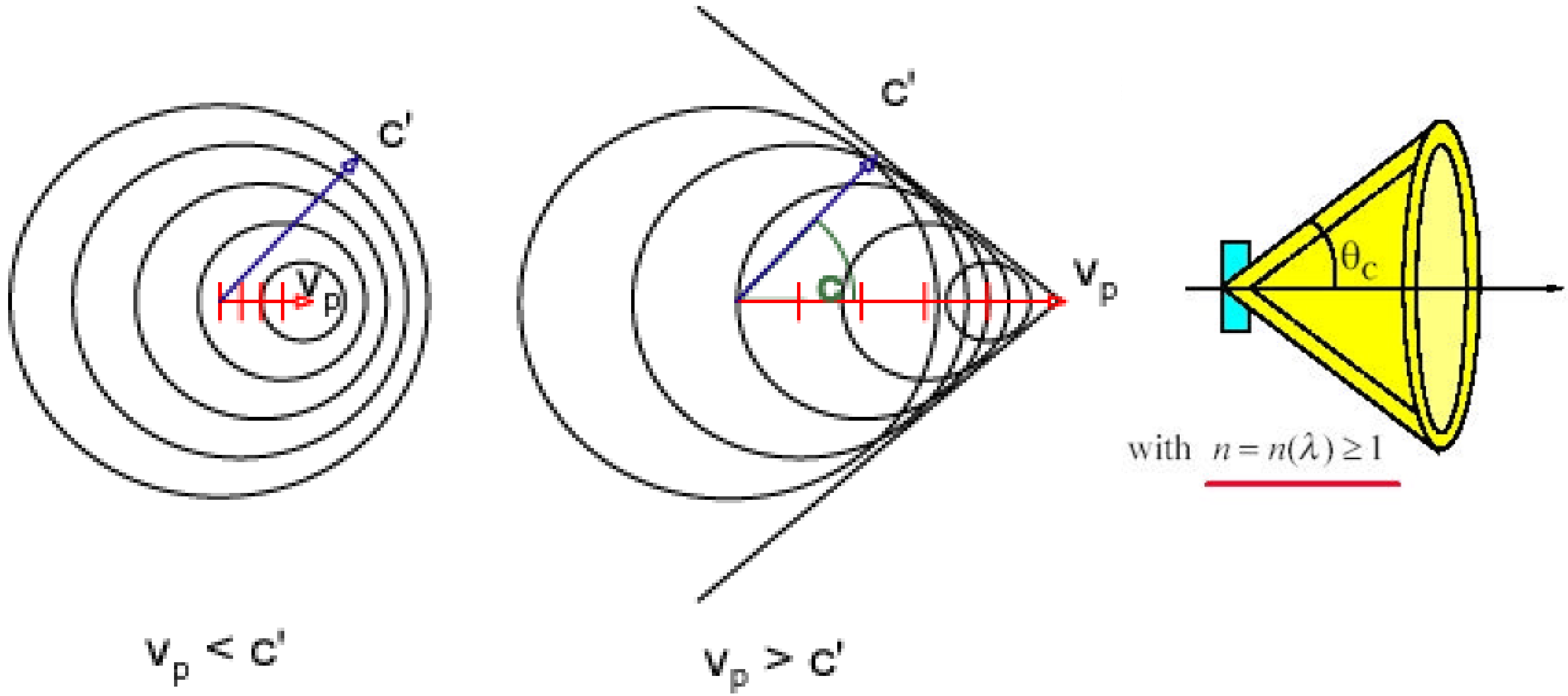
<dE/dx> averaged over many samplings: $\sigma \sim \sqrt{N}$.

good for particle identification at low momenta.

Cerenkov Radiation

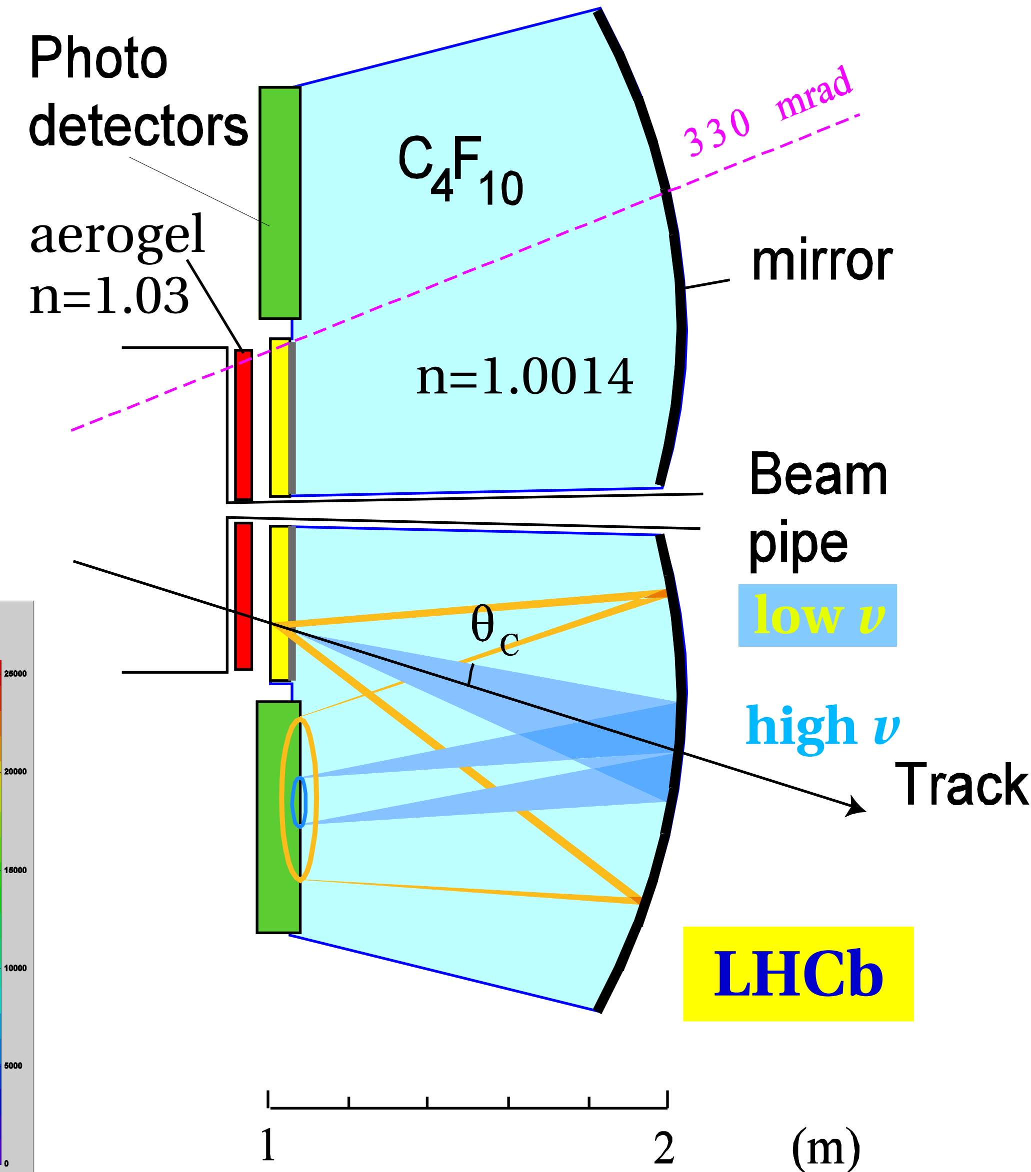
Cherenkov-Effect:
A charged particle moving faster than the speed of light in a medium $v > c/n$ emits Cherenkov radiation.

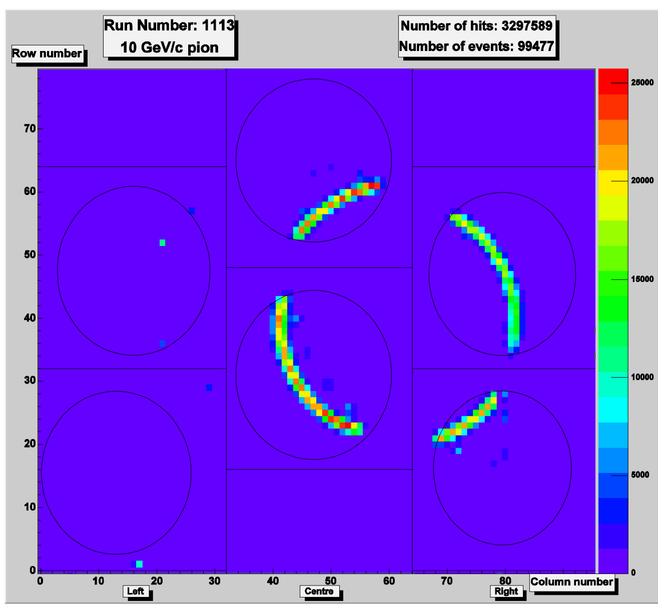
Emission of a coherent wave front: $\cos\theta_c = 1/(\beta n)$



LHCb Cerenkov Detector

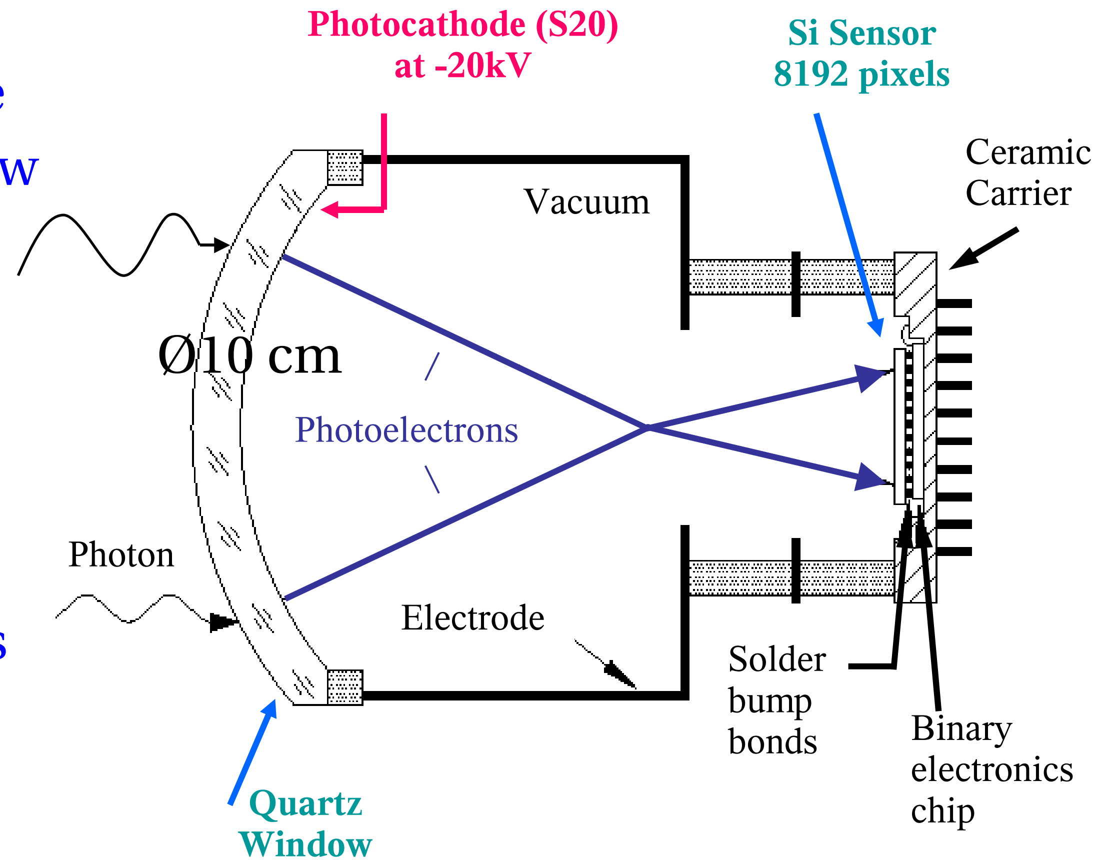
The Cherenkov cone is imaged into a ring at a position-sensitive photon detector.
Ring radius -> Cherenkov angle
-> particle velocity.
Together with momentum measurement:
determine particle mass.
Good for π - K- p separation.

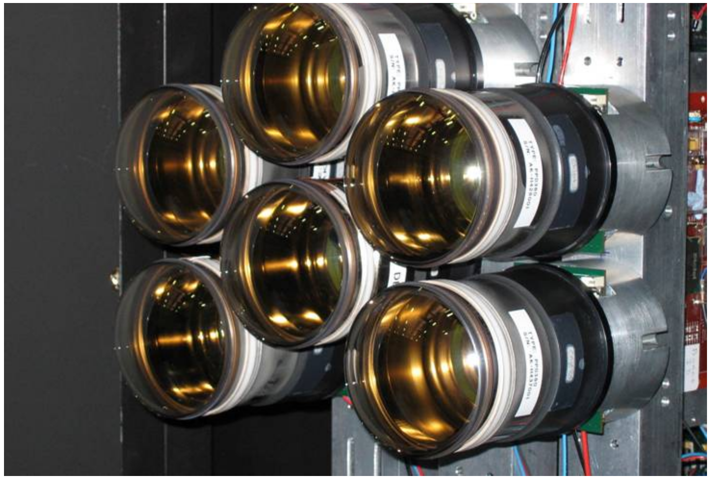




Hybrid photon detector HPD

- A Cherenkov photon reaches the multialkali-coated quartz window of an HPD.
- A photoelectron is released, and is accelerated toward the silicon array by a 20kV potential.
- The photoelectron creates around 5,000 electron-hole pairs in the silicon.
- The silicon array has 1024 pixels for position measurement.
- Readout is by a bump-bonded amplifier and discriminator chip.





Summary

- Drift detectors
- muons systems
- MWPC, CSC, RPC
- transition radiation
- TPCs
- Bethe-Bloch dE/dx
- Cherenkov and HPDs