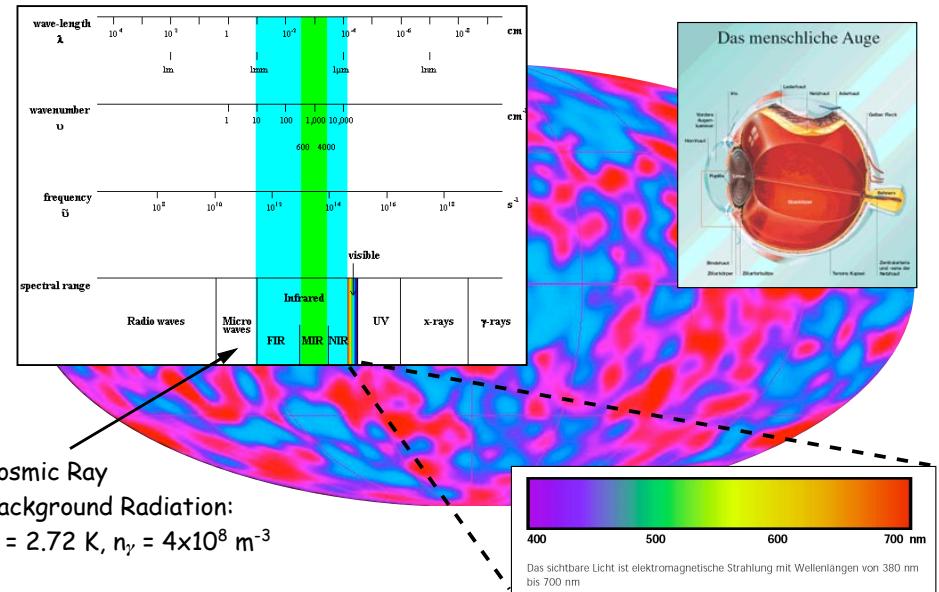


Detectors for Particle Physics

DESY Summer Student Lectures 2007

Carsten Niebuhr

Introduction



Cosmic Ray
Background Radiation:
 $T = 2.72 \text{ K}$, $n_\gamma = 4 \times 10^8 \text{ m}^{-3}$

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Topics of the Lecture

Part I

- Introduction
- Examples
- General Concepts
- Interaction of Charged Particles with Matter
 - Energy Loss: Bethe Bloch Formula
 - Multiple Scattering

Part II

- Use of Track Detectors for Momentum Measurement
- Gas Detectors
 - Proportional Chamber
 - Drift Chamber
 - TPC
 - MSGC, GEM
- Silicon Detectors
 - Strip Detectors
 - Pixel Detectors

Common lecture on Fr, 17.8.
by Georg Steinbrück

Part III

- Scintillation Counters
- Photodetectors
- Cherenkov Counters
- Transition Radiation
- Calorimeters
 - Shower Development
 - electromagnetic
 - hadronic

- not covered
 - Trigger
 - DAQ

Literatur

Text books:

K.Kleinknecht:
Teubner, 1992

Detektoren für Teilchenstrahlung

W.R. Leo:
Springer 1994

Techniques for Nuclear and Particle Physics Experiments

G.F.Knoll:
Wiley, 3rd edition

Radiation Detection and Measurement

C.Grunen:
BI Wissenschaftsverlag 1993

Teilchendetektoren

W.Bluem, L.Roland:
Springer, 1994

Particle Detection with Driftchambers

Review articles:
T.Ferbel:
Addison-Wesley 1987

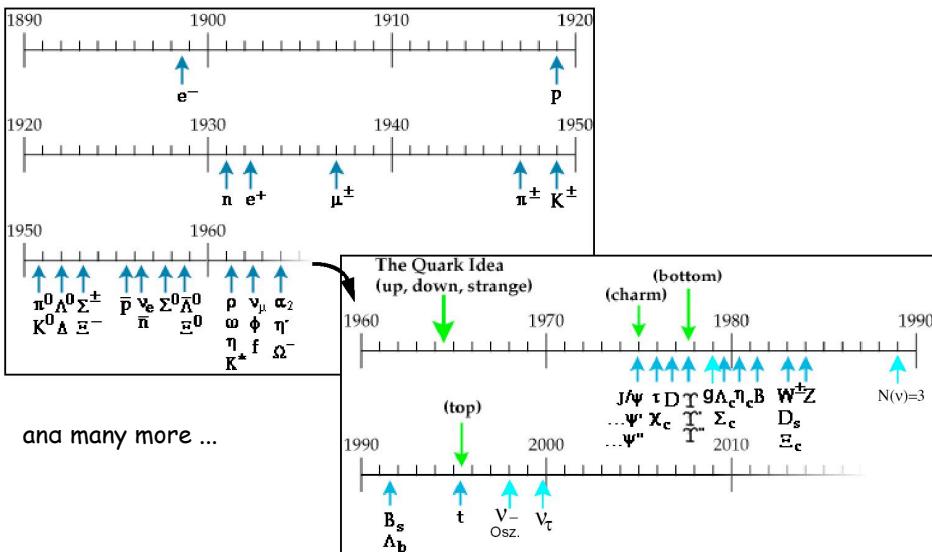
Experimental Techniques in High Energy Physics

Other sources:

Particle Data Group:
Review of Particle Physics
Eur. Phys. J. C15, 1-878 (2000)

R.K.Bock, A.Vasilescu:
The Particle Detector BriefBook
Springer, 1998 and //physics.web.cern.ch/Physics/ParticleDetector/BriefBook/

What are the Objects?



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Example for Resonance at HERA

Pentaquark candidate:

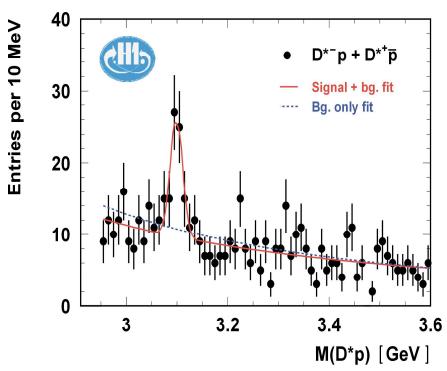
$$\theta_c^0 \rightarrow D^{*-} p \rightarrow K^- \pi^+ \pi^- p$$

minimal quark content: uuddc̄
so far only „seen“ by H1 ...

- is this a real signal ??
- or just a statistical fluctuation ?
- or a detector effect ?

=> very good understanding of detector response is required

- significance: signal/background
- resolution
- efficiency / acceptance



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Fundamental Interactions

Forces	Strong force	Electro-weak force	Gravity
Exchanged particles	Gluon	Photon	W,Z bosons
Magnitude	1	0.01	10^{-5}
	Nuclei Hadron Nuclear fusion Solar energy	Molecule, Atom Electronics Synchrotron rad. Aurora	Neutron decay Nuclei decay Neutrino Geothermy
Example	$\rho^0 \rightarrow \pi^+ \pi^-$ $\approx 10^{-24}$	$\pi^0 \rightarrow \gamma\gamma$ $\approx 10^{-16}$	$K^0 \rightarrow \pi^+ \pi^-$ $\approx 10^{-10}$
Lifetime [s]	$\approx 3 \times 10^{-13}$	$\approx 3 \times 10^{-5}$	≈ 30
cτ [mm]			

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Detection of Particles and Radiation

The goal of experimental particle physics: measurement of

- particle properties
- reaction probabilities (\rightarrow cross sections)

This requires determination of:

- particle type (mass, charge, spin etc)
- momentum / energy of particle
- emission angles

Elements contributing to such measurements :

- position sensitive detectors
- deflection in magnetic field
- calorimetry: total energy absorption and measurement
- mass determination
- Cherenkov radiation or time of flight
- transition radiation

\rightarrow position, direction
 $\rightarrow |\vec{p}|$
 $\rightarrow E_{\text{tot}}$
 $\rightarrow m$
 $\rightarrow \beta$
 $\rightarrow \gamma$

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Criteria for an ideal Detector

Because in general there can be very complex event topologies one often aims at reconstruction of full event kinematics (background rejection)

Most important:

- high **efficiency**
- high **resolution**
- high **acceptance** → try to cover full solid angle (4π)

also very important (partly conflicting demands):

- particle identification capability
- fast response
- high rate capability
- small dead time
- hermeticity
- longevity of detector components
- high reliability
- good accessibility (for repairs)
- low cost

Example of Particle Physics Experiments

Besides the large collider detectors there are many other expts: www.hep.net/experiments/all_sites.html

Bates Linear Accelerator (MIT)
BLAST , OOPS , SAMPLE

Beijing IHEP
ARGO-YBJ , BES , Tibet Asgamma

Brookhaven
BRAHMS , Crystal Ball (E913/914) , E787 , E821/muon g-2 , E850 , E852 , E863/EMU01 , E864 , E865 , E869 , E877 , E881 , E885 , E890 , E891 , E895 , E905 , E906 , E907 , E909 , E910 , E913/914 (Crystal Ball) , E917 , E923 , E926 , E927 , E949 , E953 , EIC , EMU01/E863 , High Gain Harmonic Generation FEL , ICAE , IFEL , IMB , LEGS , MECO , Microundulator FEL , NuMass/E952 , PHENIX , PHOBOS , pp2pp , Smith-Purcell , STAR , Zero Degree Calorimeter

CERN
ALEPH , ALICE , AMS , ANTARES , ASACUSA , ATHENA , Atlas (European) , ATRAP , CDHS neutrino experiment/WA1 , CERES/NA45 , CHORUS , CMS , CosmoLEP , CPLEAR/PS195 , Crystal Barrel/PS 197 , Crystal Clear/RD18 , DELPHI , EMU01 , FELIX , HARP , ICANOE , ISOLDE , L3 , LHC-B , MISTRAL , NTOF1 , NTOF2 , NTOF3 , NA45.2/IONS/EL_PAR , NA47/SMC , NA48 , NA48.1 , NA48.2 , NA49 , NA50 , NA51 , NA52/ Newmass , NA56/SPY , NA57 , NA58/COMPASS , NA59 , NA60 , NOMAD , OBELIX/PS201 , OPAL ,

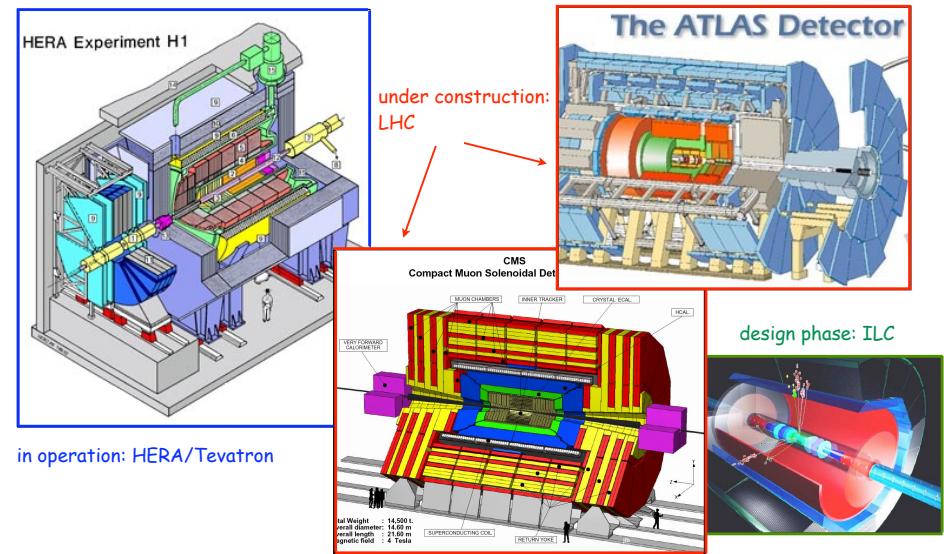
OPERA , PAMELA , PS185 , PS205/HELIUMTRAP , PS210 , PS212/DIRAC , PS214/HARP , RD8 , RD11 , RD12/TTC , RD13 , RD27 , RD39/SMSD , RD41/ MOOSE , RD42 , RD44/Geant 4 , RD45 , RD46 , RD48/ ROSE , RD49/RADTOL , TOSCA , TOTEM , WA85 , WA92 (Beatrice) , WA94 , WA97 , WA98 , WA102

DESY
H1 , HERA-B , Hermes , TESLA , ZEUS

Fermilab
Antihydrogen/E862 , APEX/E868 , Auger Project , BooNE/E898 , BTEV/C0 , CDF/E830 , CDMS/E981 , CEX/E853 , Charmonium/E835 , CMS (US Server) , COSMOS/E803 , D0 (DZero)/E823 , Donut/E872 , E665 , E771 , E789 , Fermi III Project , FOCUS/E831 , HyperCP/E871 , KTEV/E799/E832 , MINOS/E875 , NuMI , NUSEA/E866 , NuTeV/E815 , SDSS , SELEX/E781 , Zero Degrees/C0

Gran Sasso
BOREXino , CRESST , CUORICINO , DAMA , EASTOP , GALLEX(finished) , GENIUS , GNO , Heidelberg Dark Matter Search (HDMS) , Heidelberg-Moscow Experiment , ICARUS , LUNA , LVD , MACRO , MONOLITH , NOE , OPERA , USA

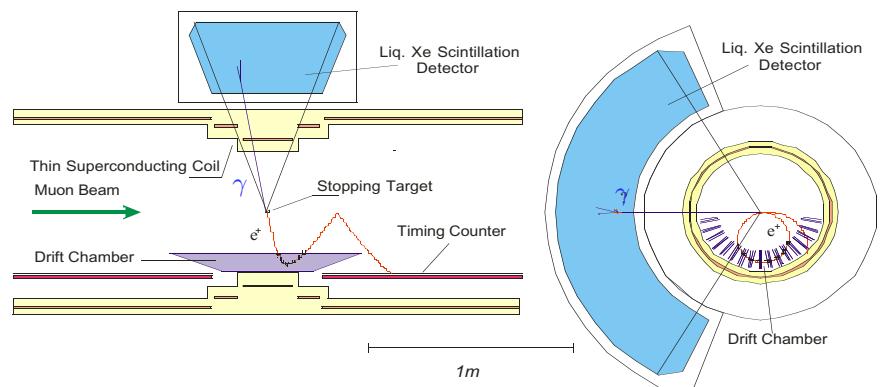
Modern Collider Detectors



Search for Rare/Forbidden Decays

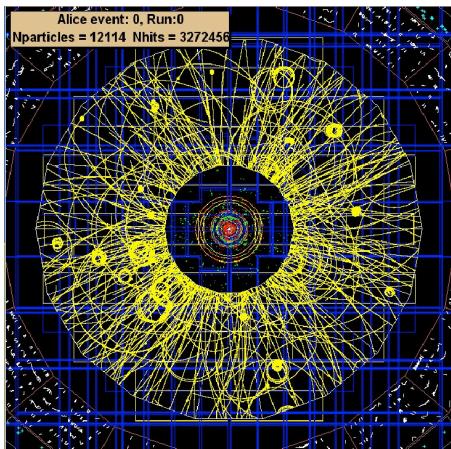
Experiment in preparation at Paul Scherrer Institut (PSI, Switzerland):

- search for lepton-number violating process: $\mu \rightarrow e \gamma$ sensitivity goal: 10^{-13} !
- needs excellent energy resolution, high event rate, but small track multiplicity per event
- start full data taking in 2007



ALICE @ LHC

Heavy Ion Physics: this simulation shows 1/10 of all 10000-20000 expected tracks in a typical event. The separation of all these tracks puts very high demands on the position resolution and double hit separation of the device.

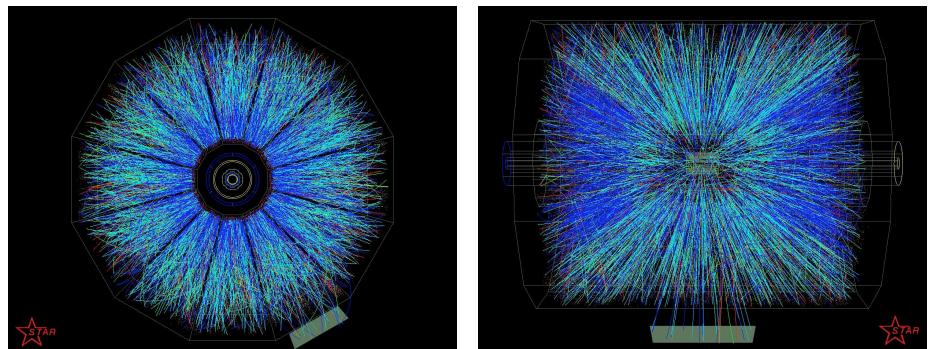


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Real Event in STAR at RHIC



≈ 2000 tracks per event

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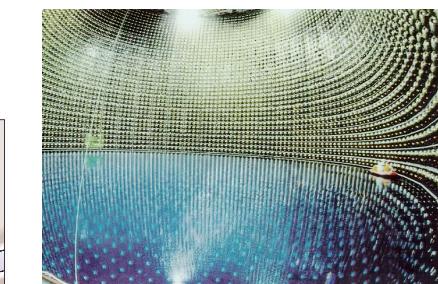
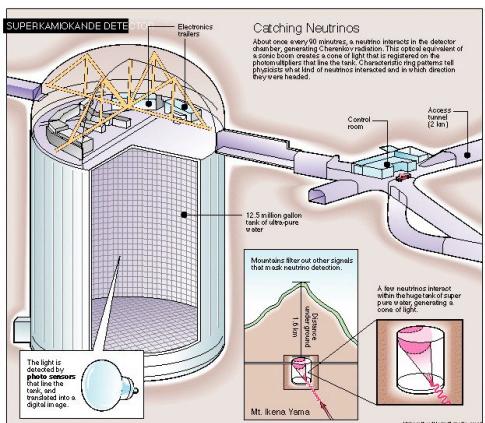
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Super Kamiokande (Japan)

Search for proton-decay and for neutrino oscillations

- 50000 tons of water
- 12000 photo tubes



Reactions :

- $\nu_\mu N \rightarrow \mu N$ Cherenkov
- $\nu_e N \rightarrow e N$ shower

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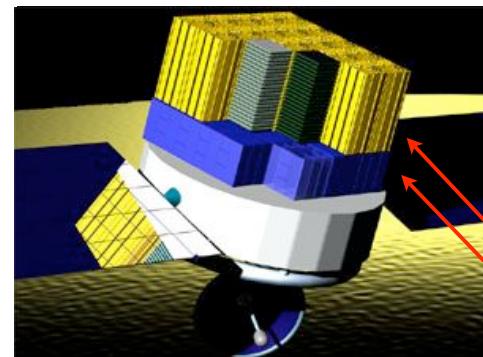
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Satellite based Detectors



Liftoff scheduled for August 2007



GLAST Gamma-Ray observatory for high energy photons in the range 20MeV to >300 GeV

Astro particle physics

- history of star formation
- acceleration mechanism of AGN's
- sources of gamma ray bursts
- nature of dark matter

Components (need highest reliability !)

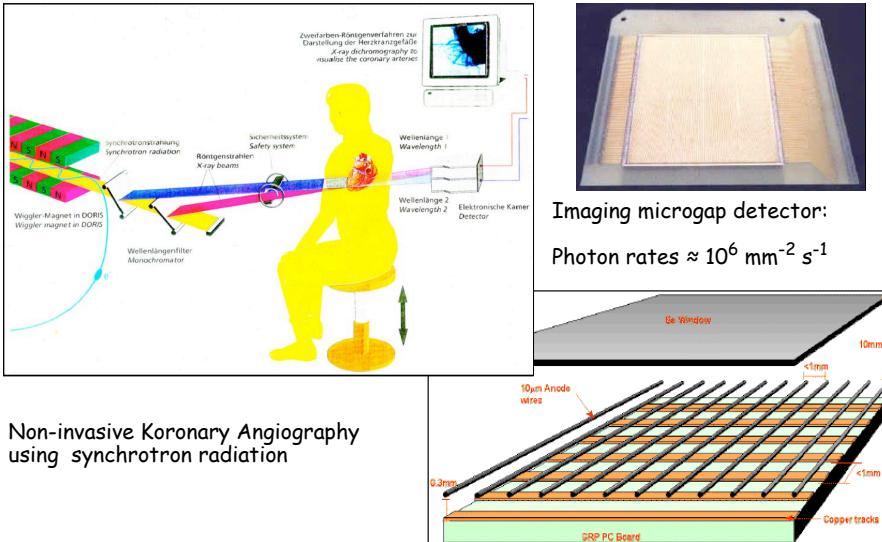
- precision tracker (Si-strips)
- calorimeter (CsI(Tl))
- data acquisition system
- anticoincidence detector

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Applications in Medicine

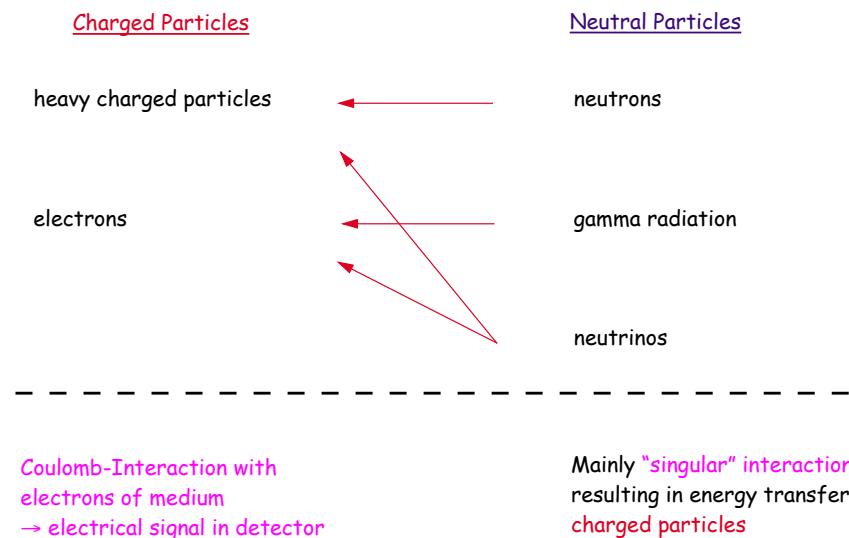


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Interaction of Radiation with Matter



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Interplay between Physics and Technology

Almost all effects used in particle detectors are based on the **electromagnetic interaction** only. Most modern detectors convert the absorbed energy into an electrical signal.

The detection sensitivity and detector performance depends on

- statistical processes in the detector
- fluctuations in the electronics

To maximize detection sensitivity and resolution one must consider and optimize

- signal formation in the detector
- coupling of the detector to the readout electronics
- noise generated in the electronics

Understanding of e.g. a modern tracking detector in high-energy physics or a medical imaging system thus requires knowledge of

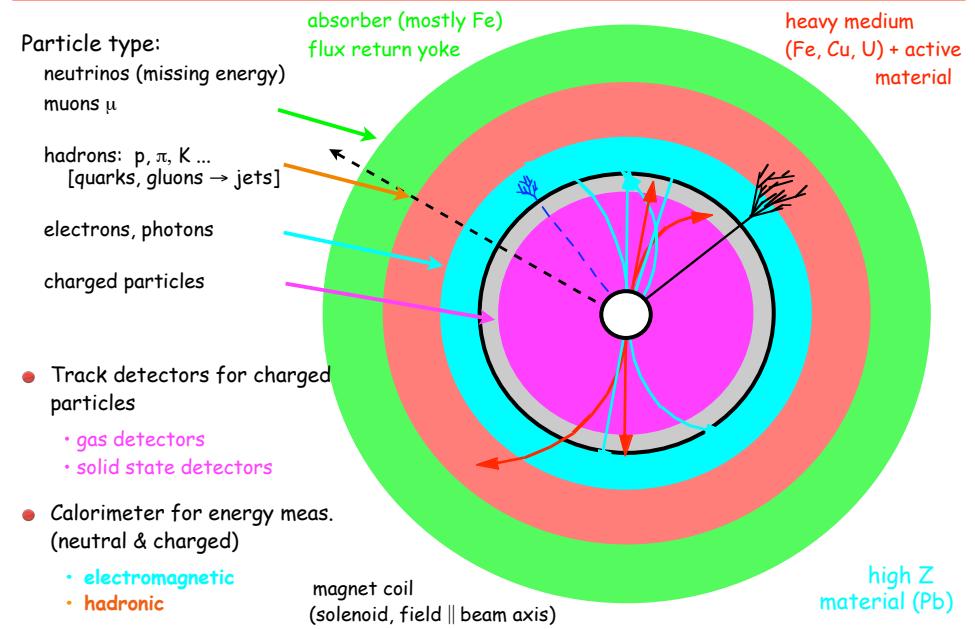
- solid state physics
- semiconductor device physics
- semiconductor fabrication technology
- low-noise electronics techniques
- analog and digital microelectronics
- high-speed data transmission
- computer-based data acquisition systems

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Cross Section of a typical Collider Detector

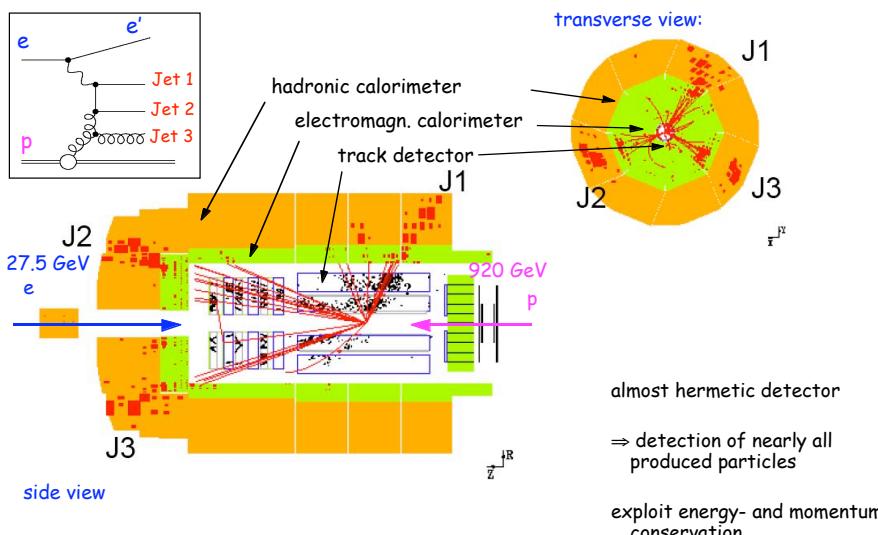


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Example 1: HERA ep Event with 3 Jets

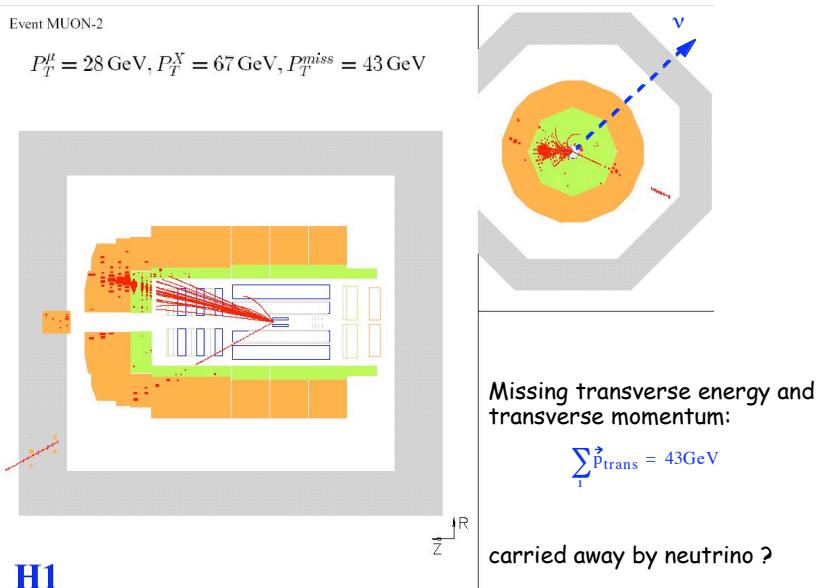


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Example 3: Neutrinos



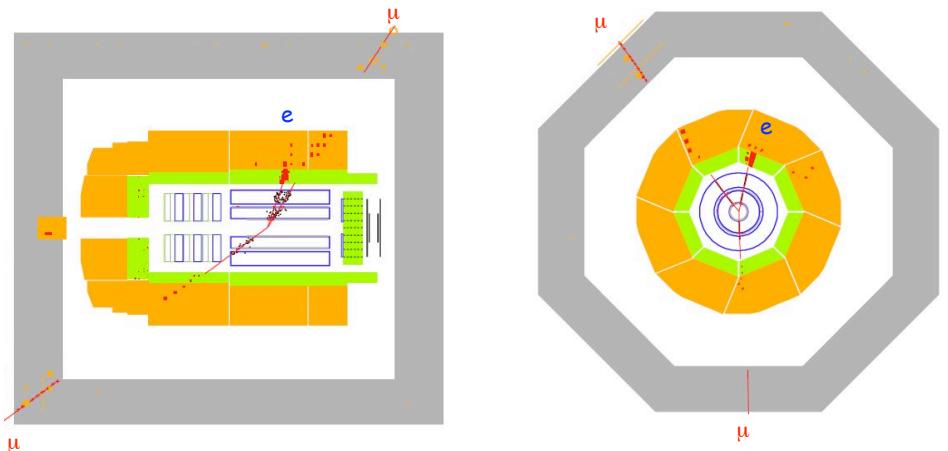
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Example 2: Muon Detection

Because muons do not interact strongly and because of their large mass (compared to electrons) they don't shower so easily and thus can penetrate calorimeter and iron yoke

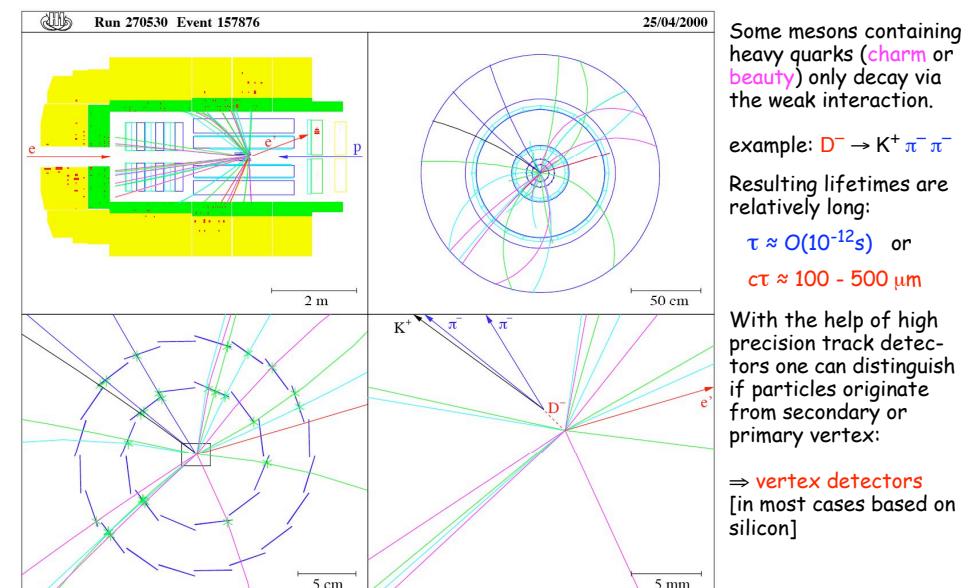


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Example 4: Secondary Vertices



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Interaction of Charged Particles

There are two principal effects which characterize the passage of charged particles through matter:

- energy loss
- change of direction

both effects result from the following electromagnetic processes:

- inelastic collisions with shell electrons of medium
- elastic scattering off nuclei

relevant is the statistical sum of many such interactions.

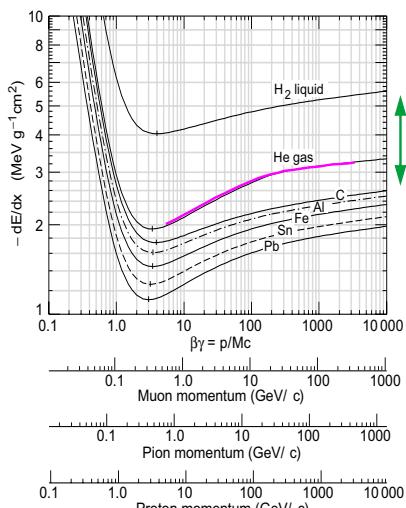
In addition there are the following processes:

- bremsstrahlung
- emission of Cherenkov radiation
- nuclear reactions
- emission of transition radiation

which however in general occur much less frequent than atomic collisions.

For charged particles one must distinguish **light particles** (i.e. e^+ , e^-) and **heavy particles** (i.e. all the rest: μ , π , p , α , light nuclei, ...)

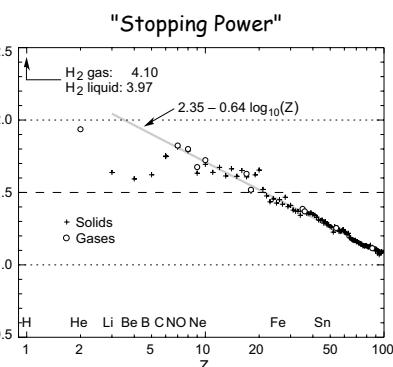
Material Dependence



Energy loss is a function of Z/A , I , and the density ρ . Hence use instead :

$$X = x \cdot \rho \Rightarrow \frac{dE}{dX} = \frac{1}{\rho} \cdot \frac{dE}{dx} \quad [\text{MeV g}^{-1} \text{cm}^2]$$

$\Rightarrow dE/dX$ essentially a function of Z/A



Bethe-Bloch Formula

Energy loss by ionisation:

$$-\frac{dE}{dx} = 4\pi N_A r_e^2 m_e c^2 \frac{z^2}{A} \cdot \frac{1}{\beta^2} \cdot \left[\ln \frac{2m_e c^2 \beta^2 \gamma^2}{I} - \beta^2 - \frac{\delta}{2} \right] \quad [\text{MeV/cm}]$$

incident particle

- $\beta = v/c$ velocity of particle
- z charge of particle

medium

- Z, A of medium

- $I \approx 16 \cdot Z^{0.9}$ average ionisation potential in [eV]
- δ describes density effect due to polarisation of medium (\Rightarrow saturation of relativistic rise)

other constants

- N_A Avogadro's number
- $r_e = 2.8 \text{ fm}$ classical electron radius
- m_e electron mass

Properties:

- dependence only on particle **velocity** **not on particle mass**
- at small energies
- drops like $\propto 1/\beta^2$
- at relativistic energies i.e. $\beta\gamma \gg 1$
- logarithmic rise $\propto \ln(\beta\gamma) = \ln \frac{p}{m}$



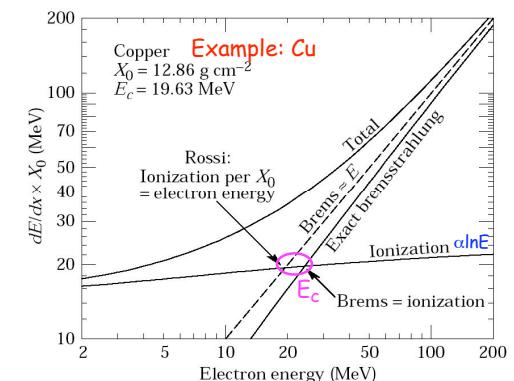
- approximate range of validity:
 $10 \text{ MeV}/c \leq p \leq 50 \text{ GeV}/c$

Energy Loss of Electrons

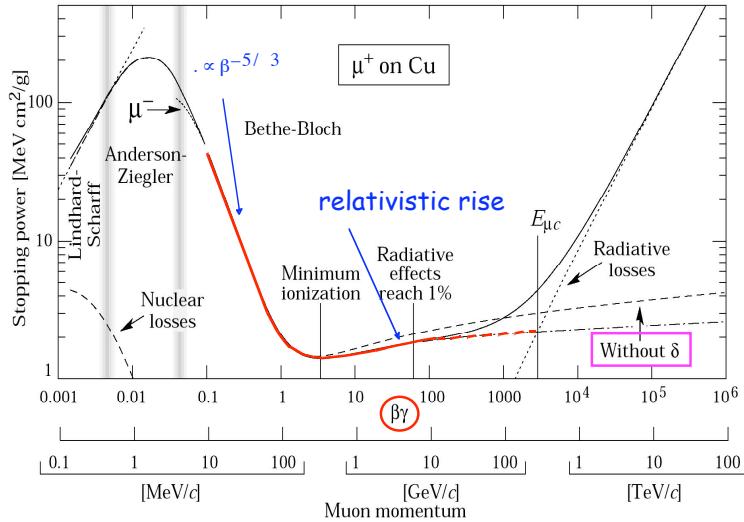
In addition to energy loss by ionisation high energy particles also loose energy due to interaction with the Coulomb field of the **nuclei**: **Bremsstrahlung**

Due to their small mass this effect is especially prominent for electrons (positrons):

- $dE/dx \propto Z^2 \cdot E/m^2$
- it is useful to introduce **radiation length** X_0
- energy attenuation: $E = E_0 \exp\left(-\frac{x}{X_0}\right)$
- approx.: $X_0 = \frac{716 \text{ cm}^{-2} A}{Z(Z+1) \ln(287/\sqrt{Z})}$
- critical energy E_c :
$$\frac{dE_{\text{Brems}}}{dx} = \frac{dE_{\text{collision}}}{dx}$$
- approximately: $E_c = \frac{610 \text{ MeV}}{Z+1.24}$



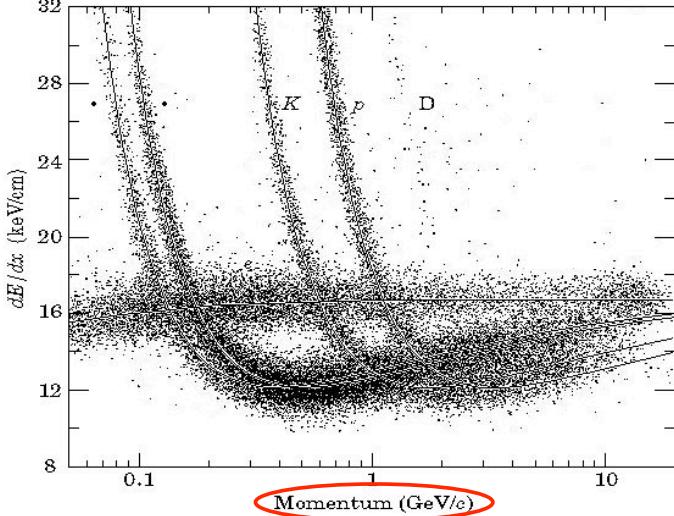
Energy Loss of Muons



MIP = minimum-ionising particles

dE/dx in a TPC

Measurements in PEP4/9-TPC (Ar-CH₄ = 80:20 @ 8.5atm)



If dE/dx is plotted versus **momentum** of particle the curves are shifted horizontally for different masses

Application: if also the momentum of the particle is known the measurement of the **specific ionisation** can be used for **particle identification**

In this example each dot represents ≈ 185 single measurements in a drift chamber

dE/dx Applications for different Detector Types

Gas

- ionisation

⇒ proportional or drift chamber

Liquid

- local heating

⇒ bubble chamber

- ionisation

⇒ calorimeter (Liquid Argon or Krypton)

Solid

- excitation of electrons

→ conversion into light

⇒ scintillators

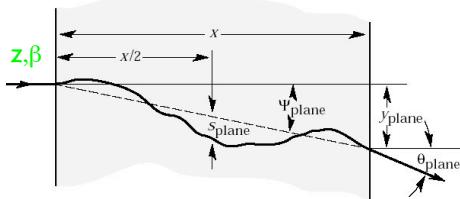
- creation of electron-hole pairs

⇒ solid state detectors

Properties of some Materials (PDG)

Material	Z	A	$\langle Z/A \rangle$	Nuclear ^a interaction length λ_T (g/cm^2)	dE/dx _{min} (MeV/ g/cm^2)	Radiation length ^c X_0 (g/cm^2)	Density (g/cm^3) (g/ℓ) for gas	Liquid boiling point at atm(K) atm(K) for gas	Refractive index n ($((n-1) \times 10^6$)
H ₂ gas	1	1.00794	0.99212	43.3	50.8	(4.103)	61.28 ^d (731000)	0.0838 [0.0899]	[139.2]
H ₂ liquid	1	1.00794	0.99212	43.3	50.8	(4.034)	61.28 ^d (2.052)	0.0708 [0.179]	20.39 [1.112]
D ₂	1	2.0140	0.49652	45.7	54.7	(2.052)	22.4	0.1690 [0.179]	23.65 [1.128] [138]
He	2	4.002602	0.49968	49.9	65.1	(1.937)	94.32	756	0.1249 [0.1786]
Li	3	6.941	0.43221	54.6	73.4	1.639	82.76	155	0.534
Be	4	9.012182	0.44384	55.8	75.2	1.594	65.19	35.28	1.848
C	6	12.011	0.49954	60.2	86.3	1.745	42.70	18.8	2.265 ^e
N ₂	7	14.00674	0.49976	61.4	87.8	(1.825)	37.99	47.1	0.8073 [1.250]
O ₂	8	15.9994	0.50002	63.2	91.0	(1.801)	34.24	30.0	1.141 [1.428]
F ₂	9	18.9984032	0.47372	65.5	95.3	(1.675)	32.93	21.85	1.507 [1.696]
Ne	10	20.1797	0.49555	66.1	96.6	(1.724)	28.94	24.0	1.204 [0.9005]
Al	13	26.981539	0.48181	70.6	106.4	1.615	24.01	8.9	2.70
Si	14	28.0855	0.49848	70.6	106.0	1.664	21.82	9.36	2.33
Ar	18	39.948	0.45059	76.4	117.2	(1.519)	19.55	14.0	1.396 [1.782]
Ti	22	47.867	0.45948	79.9	124.9	1.476	16.17	3.56	4.54
Fe	26	55.845	0.46556	82.8	131.9	1.451	13.84	1.76	7.87
Cu	29	63.546	0.45636	85.6	134.9	1.403	12.86	1.43	8.96
Ge	32	72.61	0.44071	88.3	140.5	1.371	12.25	2.30	5.323
Sn	50	118.710	0.42120	100.2	163	1.264	8.82	1.21	7.31
Xe	54	131.29	0.41130	102.8	169	(1.255)	8.48	2.87	2.953 [5.858]
W	74	183.84	0.40250	110.3	185	1.145	6.76	0.35	19.3
Pt	78	195.08	0.39984	113.3	189.7	1.129	6.54	0.305	21.45
Pb	82	207.2	0.39575	116.2	194	1.123	6.37	0.35	—
U	92	238.0289	0.38651	117.0	199	1.082	6.00	≈0.32	≈18.95
Air, (20°C, 1 atm), [STP]					(1.815)	36.66	[30420]	(1.205)[1.2931]	78.8 [273] [293]
H ₂ O					(1.5509)	60.1	83.6	1.991 [36.08]	36.1 [1.00]
CO ₂ gas					(0.49989)	62.4	89.7	(1.819)[36.2]	373.15 [410]
								[1.977]	

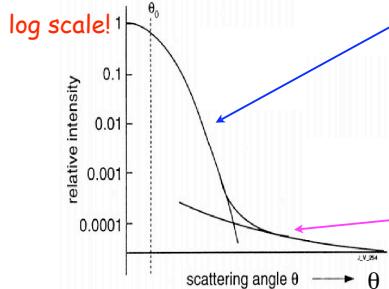
Multiple Scattering



Scattering of charged particles off the atoms in the medium causes a change of direction:

the statistical sum of many such small angle scatterings results in a gaussian angular distribution with a width given by:

$$\theta_0 = \frac{13.6 \text{ MeV}}{\beta p c} z \sqrt{\frac{x}{X_0}} \left[1 + 0.038 \ln\left(\frac{x}{X_0}\right) \right]$$



example: $p = 1 \text{ GeV}/c$, $d = 300 \mu\text{m Si}$, $X_0 = 9.4 \text{ cm}$
 $\Rightarrow \theta_0 \approx 0.8 \text{ mrad}$. For 10cm distance this corresponds to $80 \mu\text{m}$, which is significantly larger than typical resolution of Si-strip detector

the less likely scattering off the atomic nuclei causes large scattering angles resulting in a deviation from a gaussian distribution at large angles

Summary Part I

- properties of different particles require many different types of detectors
- rough classification
 - track/position detectors (**non destructive**) → Part II
 - calorimeters (**destructive**) → Part III
- basically all detectors based on electromagnetic interaction
- detectors also more and more used for other applications (e.g. medical appl.)
- energy loss of charged particles:** Bethe-Bloch describes loss due to ionisation
 - $z^2/\beta^2 \rightarrow \ln E \rightarrow$ Fermiplateau
 - minimum at $\beta\gamma \approx 3$: 1-2 MeV per gcm⁻²
 - however light particles (electrons, muons) at high energies lose energy predominantly by bremsstrahlung: $-dE/dx \propto Z^2 \cdot E/m^2$
- multiple scattering**
 - gaussian core of distribution with angular spread: $\theta_0 \sim \frac{1}{p} \sqrt{\frac{x}{X_0}}$
 - deviation from gaussian distribution at **large angles** due to Rutherford scattering (nuclei)

Part II

• Use of Track Detectors for Momentum Measurement

• Gas Detectors

- Proportional Chamber
- Drift Chamber
- TPC
- MSGC, RPC, GEM

• Solid State Detectors

- Strip Detectors
- Pixel Detectors

Common lecture on Fr, 10.8.
by Georg Steinbrück