

S. Petersburg, DIS 2003

April 23rd, 2003

Tests of perturbative QCD with hadronic final states in hadron-induced reactions

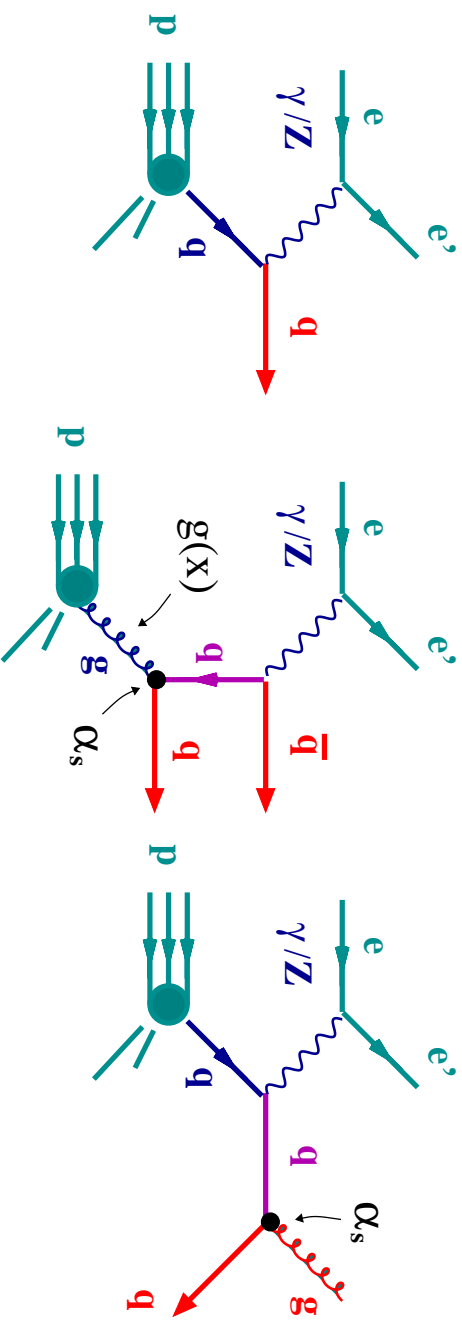
Juan Terrón (Universidad Autónoma de Madrid, Spain)



- **Outline**
 - **Deep inelastic scattering (HERA)**
 - **γp collisions (HERA)**
 - **$\gamma\gamma$ collisions (LEP)**
 - **$p\bar{p}$ collisions (TEVATRON)**

Jet Production in Neutral Current Deep Inelastic Scattering

- Jet production in neutral current deep inelastic scattering up to $\mathcal{O}(\alpha_s)$:



- Perturbative QCD calculations of jet cross sections:

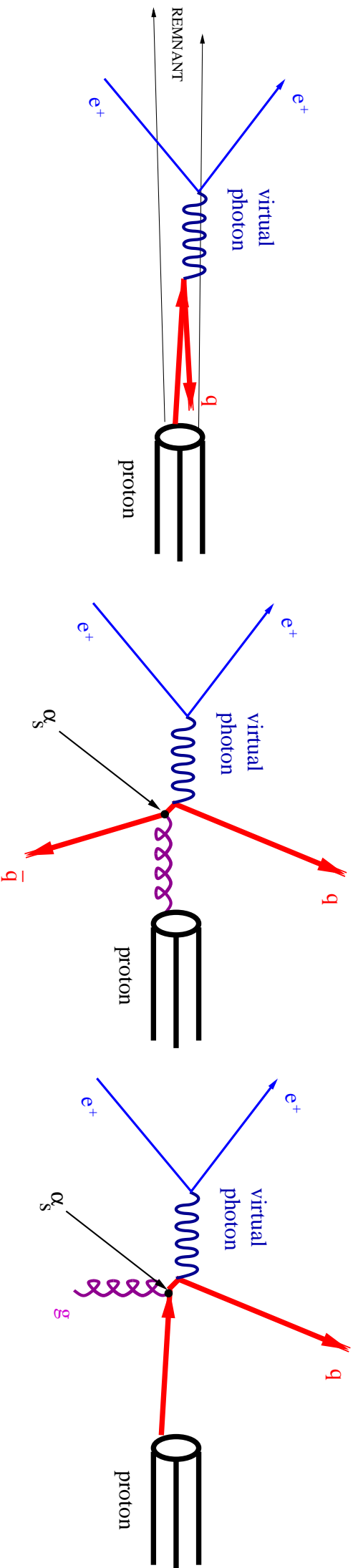
$$d\sigma_{jet} = \sum_{a=q, \bar{q}, g} \int dx f_a(x, \mu_F^2) d\hat{\sigma}_a(x, \alpha_s(\mu_R), \mu_R^2, \mu_F^2)$$

- f_a : parton a density in the proton, determined from experiment; **long-distance structure of the target**
- $\hat{\sigma}_a$: subprocess cross section, calculable in pQCD; **short-distance structure of the interaction**

Jet Production in Neutral Current Deep Inelastic Scattering

- In the region where the wealth of data from fixed-target and collider experiments has allowed an accurate determination of the proton PDFs, measurements of jet production in NC DIS provide
 - a sensitive test of the pQCD predictions of the short-distance structure
 - a determination of the strong coupling constant α_s
- To perform a stringent test of the pQCD predictions and a precise determination of α_s :
 - * Observables for which the predictions are directly proportional to α_s
 - Jet cross sections in the Breit frame
 - * Small experimental uncertainties → Jets with relatively high transverse energy
 - * Small theoretical uncertainties → NLO QCD calculations
 - Jet algorithm: longitudinally invariant k_T cluster algorithm (Catani et al)
 - (small parton-to-hadron effects, infrared safe, suppression of beam-remnant jet)
 - Jet selection criteria
- Exploration of the parton evolution at low x ⇒ footprints of BFKL effects?
- Exploration of the low Q^2 (transition) region ⇒ resolved virtual photons?

High- E_T Jet Production in the Breit Frame

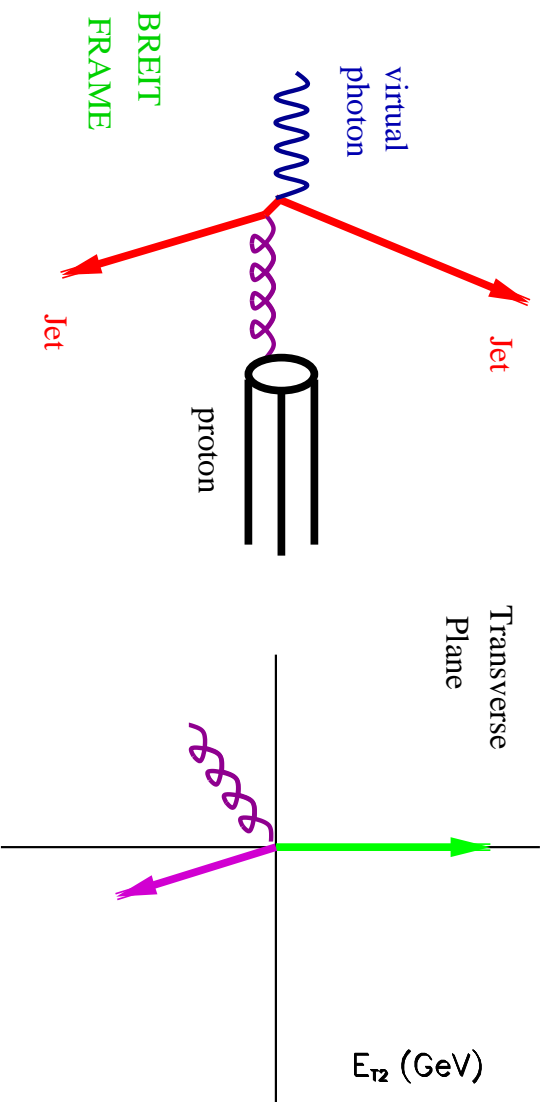


- In the Breit frame the virtual boson collides head-on with the proton
- High- E_T jet production in the Breit frame
 - suppression of the Born contribution (struck quark has zero E_T)
 - suppression of the beam-remnant jet (zero E_T)
 - lowest-order non-trivial contributions from $\gamma^* g \rightarrow q\bar{q}$ and $\gamma^* q \rightarrow qg$
 - ⇒ directly sensitive to hard QCD processes (α_s)

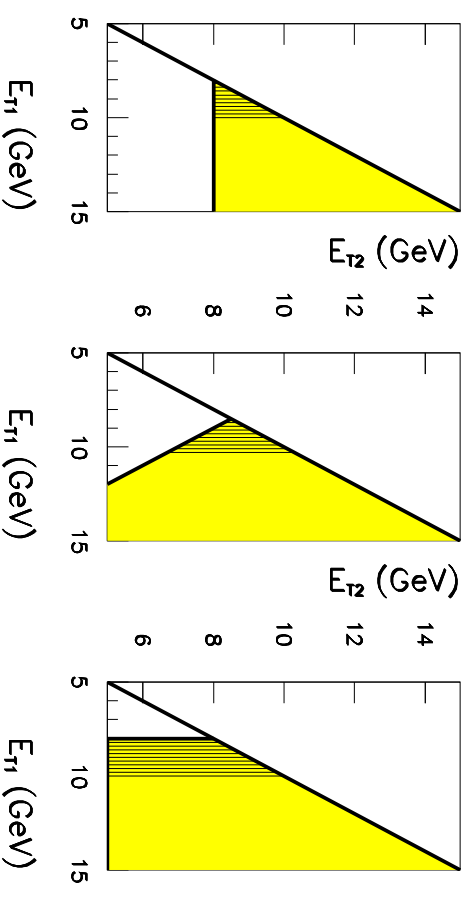
NLO QCD Calculations of Jet Cross Sections in DIS

- Several NLO QCD programs are available for performing jet cross section calculations → DISENT (Catani and Seymour), MEPPJET (Mirkes and Zeppenfeld), DISASTER++ (Graudenz), NLOJET (Nagy and Trocsanyi)
- NLO corrections → virtual corrections with internal particle loops
→ real corrections with a third parton in the final state
- Different methods to calculate real corrections:
→ phase space slicing method (M), subtraction method (D, D++, NJ)
- Since there are two hard scales in jet production, the renormalisation and factorisation scales can be chosen as one of the two, μ_R , $\mu_F = Q$ or E_T^{jet}
- The calculations are for jets of partons and the measurements are done at the hadron level → need to correct the calculations for hadronisation effects
- Theoretical uncertainties:
→ terms beyond NLO, which are usually estimated by varying μ_R by factor 2
→ uncertainties on $\alpha_s(M_Z)$ and the proton PDFs
→ uncertainty coming from the hadronisation corrections

Jet Finding and Selection Criteria for Dijet Events



NLO QCD Dijet Cross Section ($\mu_R=Q$) $Q^2 > 470 \text{ GeV}^2$

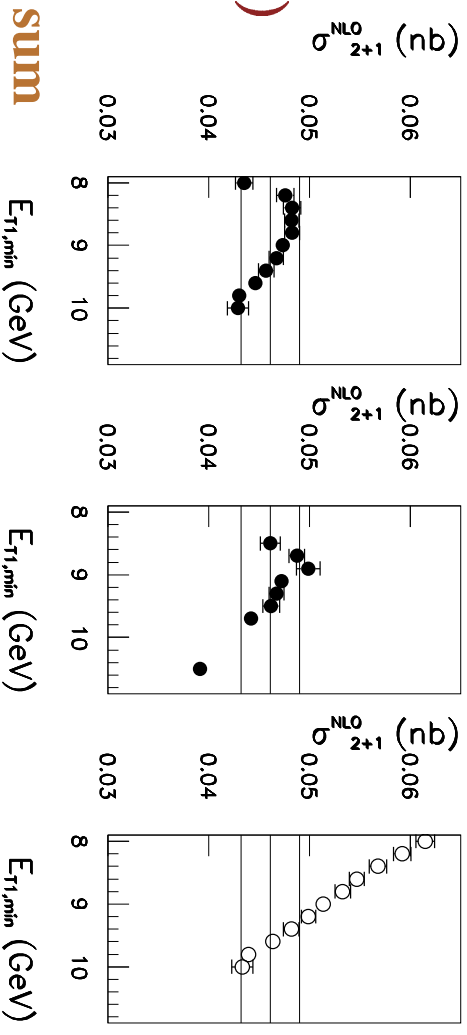


- Longitudinally invariant k_T -cluster algorithm in the η - ϕ plane of Breit frame

$$d_{ij} = \min(E_{T,i}^2, E_{T,j}^2) \cdot (\Delta\eta_{ij}^2 + \Delta\phi_{ij}^2)$$

- Dijet selection criteria:

- Symmetric cuts on $E_T^{jet,1(2)}$ ⇒ danger
- Symmetric cuts on $E_T^{jet,1(2)}$ and cut on sum
- Asymmetric cuts on $E_T^{jet,1(2)}$



⇒ NLO calculations for dijet cross sections can be (infrared) sensitive to the selection criteria

Dijet Cross Sections in NC DIS ($5 < Q^2 < 15000 \text{ GeV}^2$)

- Measurement of differential dijet cross sections

over a wide range in $Q^2 \rightarrow 5 < Q^2 < 15000 \text{ GeV}^2$

and $0.2 < y < 0.6$ for dijet production with

$$E_T^{jet,1(2)} (\text{Breit}) > 5 \text{ GeV}$$

$$E_T^{jet,1} (\text{Breit}) + E_T^{jet,2} (\text{Breit}) > 17 \text{ GeV}$$

$$-1 < \eta^{jet,1(2)} (\text{Lab}) < 2.5$$

- Detailed investigation of the jet algorithms:

→ Smallest parton-to-hadron effects: inclusive k_T

- Comparison with NLO QCD calculations:

$$\rightarrow \mu_R = \bar{E}_T, \mu_F = \sqrt{200} \text{ GeV}$$

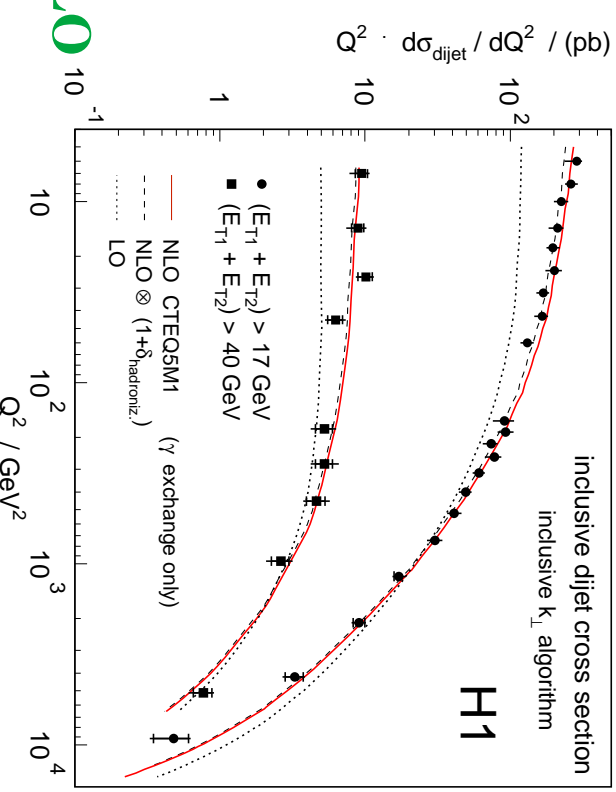
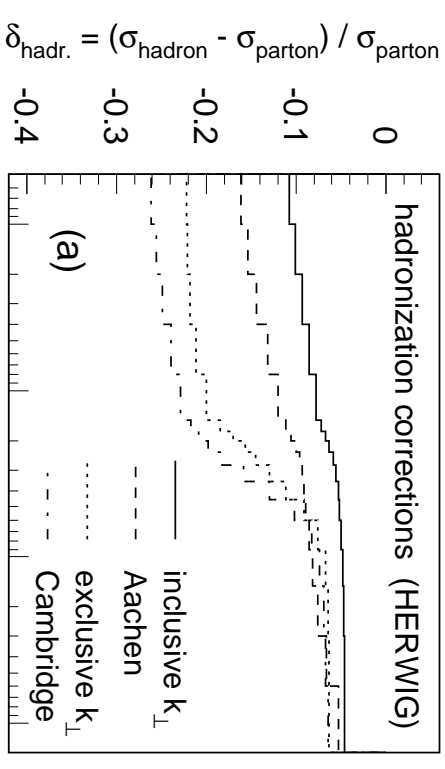
→ CTEQ5M1 parametrisations of proton PDFs

→ parton-to-hadron corrections applied

- NLO QCD gives a good description of the data over

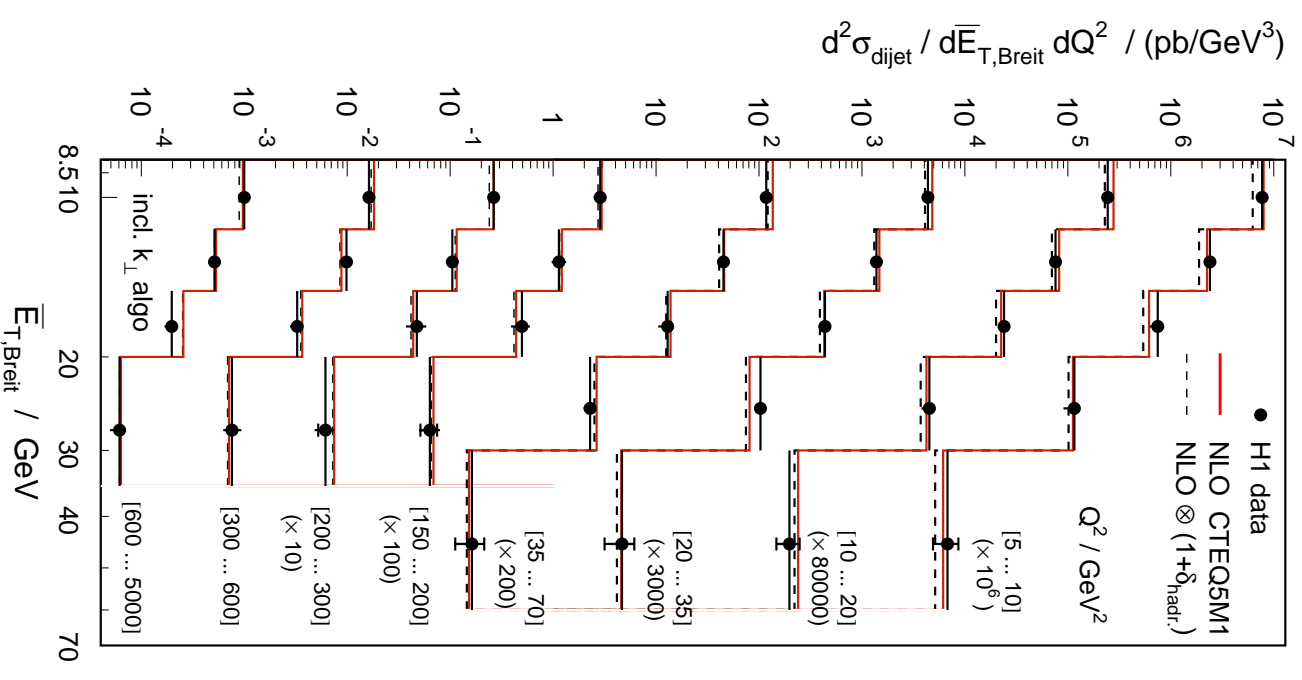
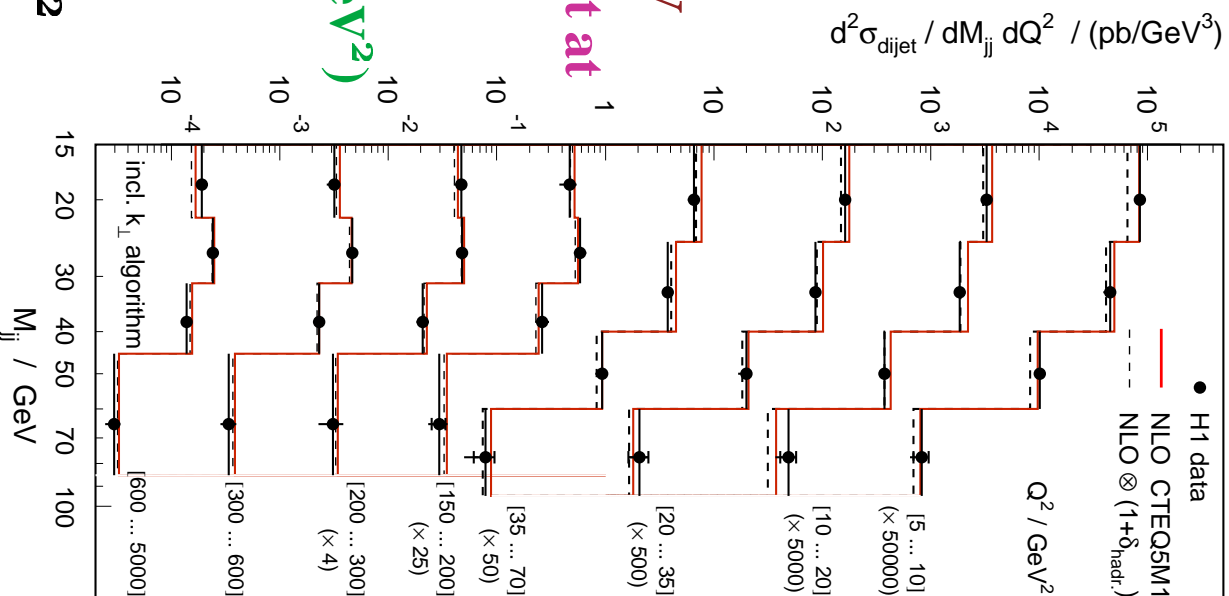
a wide range in Q^2 and E_T ; the Q^2 dependence is

observed to be reduced at high- E_T and described by NLO



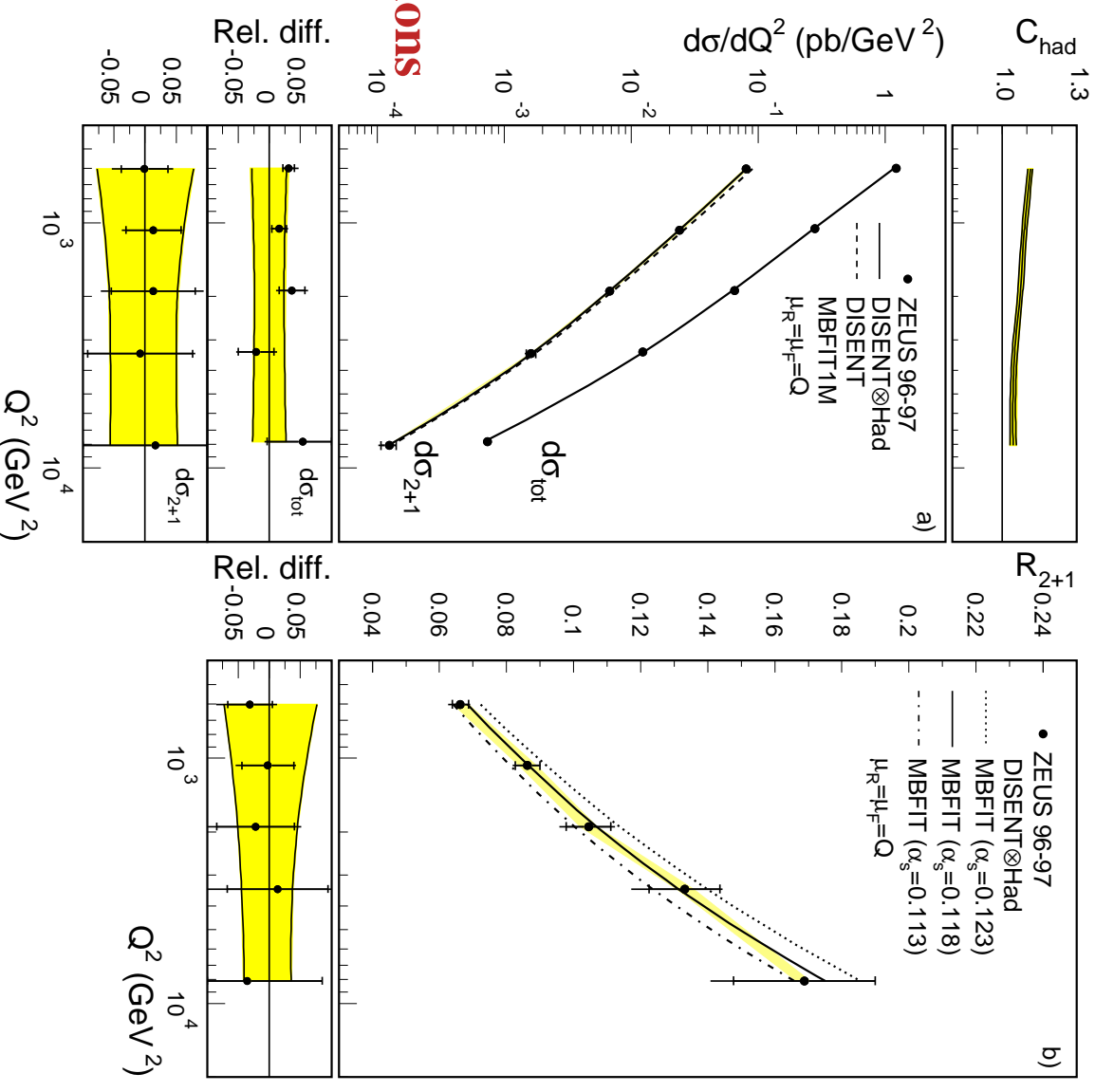
Dijet Cross Sections in NC DIS

- **Measurement of double differential cross sections**
 $d\sigma/dM_{JJ}dQ^2, d\sigma/d\bar{E}_T dQ^2$
over $5 < Q^2 < 5000 \text{ GeV}^2$
- It is observed that the spectra get harder as Q^2 increases
- **NLO QCD describes well the data over $15 < M_{JJ} < 95 \text{ GeV}$ and $8.5 < \bar{E}_T < 60 \text{ GeV}$ except at low Q^2 , where the shape is ok but not the normalisation**
- **Overview: at high Q^2 ($> 70 \text{ GeV}^2$) NLO describes the data well; as Q^2 decreases the theoretical uncertainties become large and NLO fails for $Q^2 < 10 \text{ GeV}^2$**



Dijet Cross Sections at $Q^2 > 470 \text{ GeV}^2$ and extraction of α_s

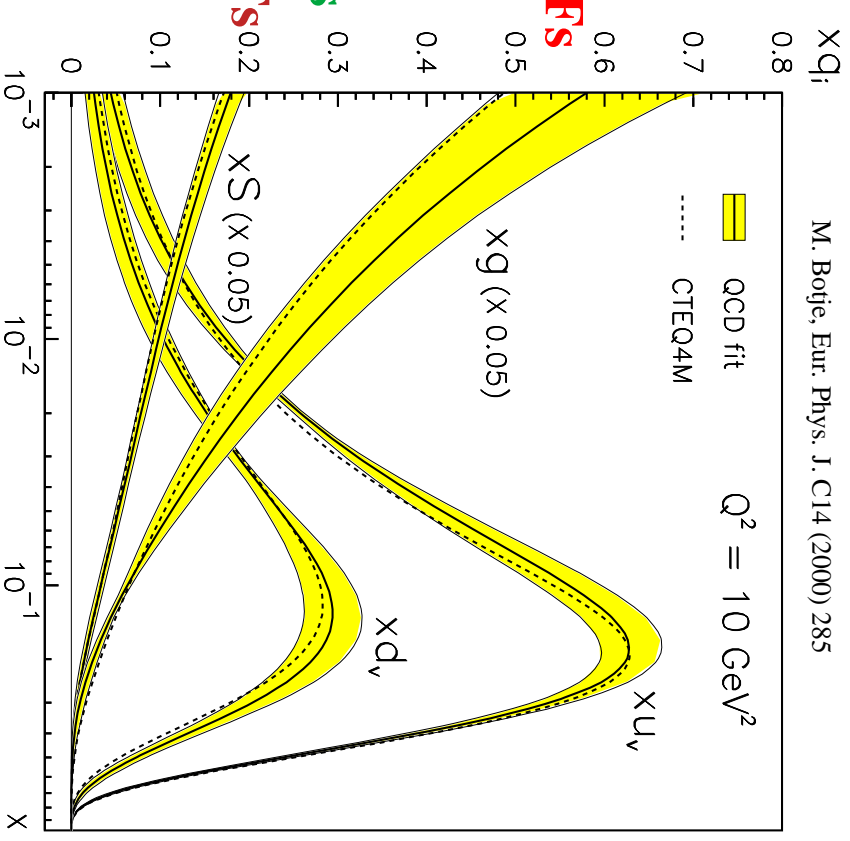
- **Dijet cross section $d\sigma_{2+1}/dQ^2$ for**
 - $470 < Q^2 < 20000 \text{ GeV}^2$
 - $E_T^{jet,1} \text{ (Breit)} > 8 \text{ GeV}$
 - $E_T^{jet,2} \text{ (Breit)} > 5 \text{ GeV}$
 - $-1 < \eta^{jet,1(2)} \text{ (Lab)} < 2$
- **Ratio $R_{2+1} \equiv \frac{d\sigma_{2+1}/dQ^2}{d\sigma_{tot}/dQ^2}$**
- **Small experimental uncertainties.**
- **Comparison with NLO QCD calculations**
- **Small theoretical uncertainties:**
 - **uncertainties on the proton PDFs**
 - **hadronisation corrections**
 - **higher-order terms ($> \text{NLO}$)**



Uncertainties of the Proton PDFs: effects on jet cross sections

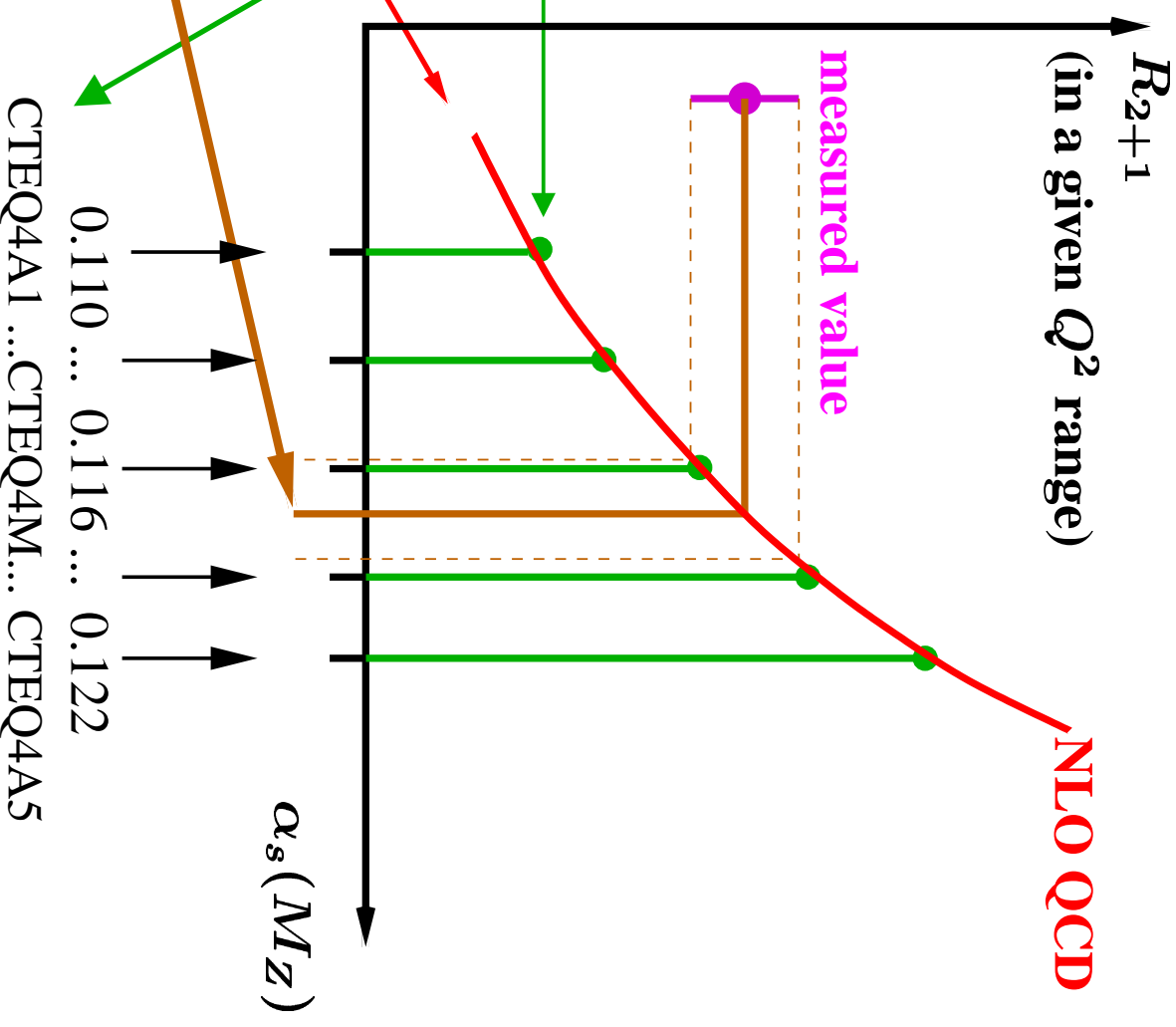
- Comparison of jet cross-section calculations using different parametrisations of the proton PDFs (e.g. MRST vs CTEQ) **DOES NOT give a reliable estimation of the uncertainties due to the proton PDFs**
- Several groups have developed methods to quantify these uncertainties by accounting (properly) for
 - the statistical and correlated systematic uncertainties of each data set used in the determination of the PDFs
 - the theoretical uncertainties affecting the extraction of the PDFs in the DGLAP fits
- Botje's analysis provides the covariance matrix (V_{p_μ, p_λ}) of the fitted PDF parameters ($\{p_\lambda\}$), the derivatives as functions of Q^2 and x (e.g. $\partial g(x, Q^2)/\partial p_\lambda$) and several sets obtained under different theoretical assumptions
 - ⇒ evaluation of the uncertainty on any function of the proton PDFs (e.g. σ_{jet})

$$(\Delta\sigma_{jet})^2 = \sum_{\lambda, \mu} \frac{\partial\sigma_{jet}}{\partial p_\mu} V_{p_\mu, p_\lambda} \frac{\partial\sigma_{jet}}{\partial p_\lambda}$$



Dijet Cross Sections at $Q^2 > 470 \text{ GeV}^2$ and extraction of $\alpha_s(M_Z)$

- NLO QCD calculations of $d\sigma_{2+1}/dQ^2$ depend on $\alpha_s(M_Z)$ through
 - Matrix Elements: $\hat{\sigma} \sim A \cdot \alpha_s + B \cdot \alpha_s^2$
 - proton PDFs: α_s assumed in evolution
- To take into account the correlation the NLO QCD calculations are performed using various sets of proton PDFs which assume different values of α_s
- The resulting NLO QCD calculations are parametrised as a function of $\alpha_s(M_Z)$ in each region of Q^2 of the measurements
- From the measured value of R_{2+1} in each region of Q^2 the value of $\alpha_s(M_Z)$ and its uncertainty are extracted



Dijet Cross Sections at $Q^2 > 470 \text{ GeV}^2$ and extraction of α_s

- Study of the **scale dependence of $\alpha_s(Q)$** :

from the measured $R_{2+1}(Q^2)$ in each Q^2 region

→ $\alpha_s (< Q >)$ is extracted

The measurements are consistent with

the running of α_s predicted by perturbative QCD

- A combined value of $\alpha_s(M_Z)$ has been extracted:

$$\alpha_s(M_Z) = 0.1166 \pm 0.0019 \text{ (stat.)}$$

$$+0.0024 \text{ (exp.) } +0.0057 \text{ (th.)}$$

$$-0.0033 \text{ (exp.) } -0.0044 \text{ (th.)}$$

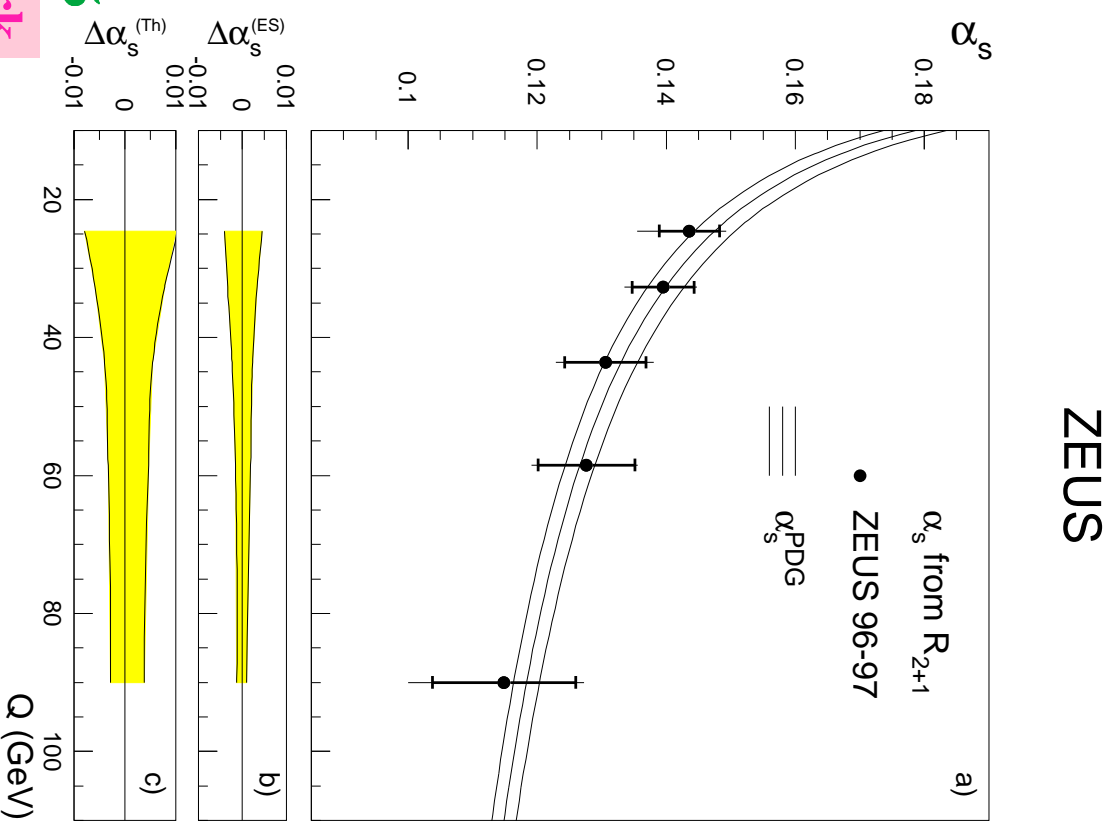
- The theoretical uncertainty dominates:

→ terms beyond NLO $\Delta\alpha_s(M_Z) = +0.0055$
 -0.0042

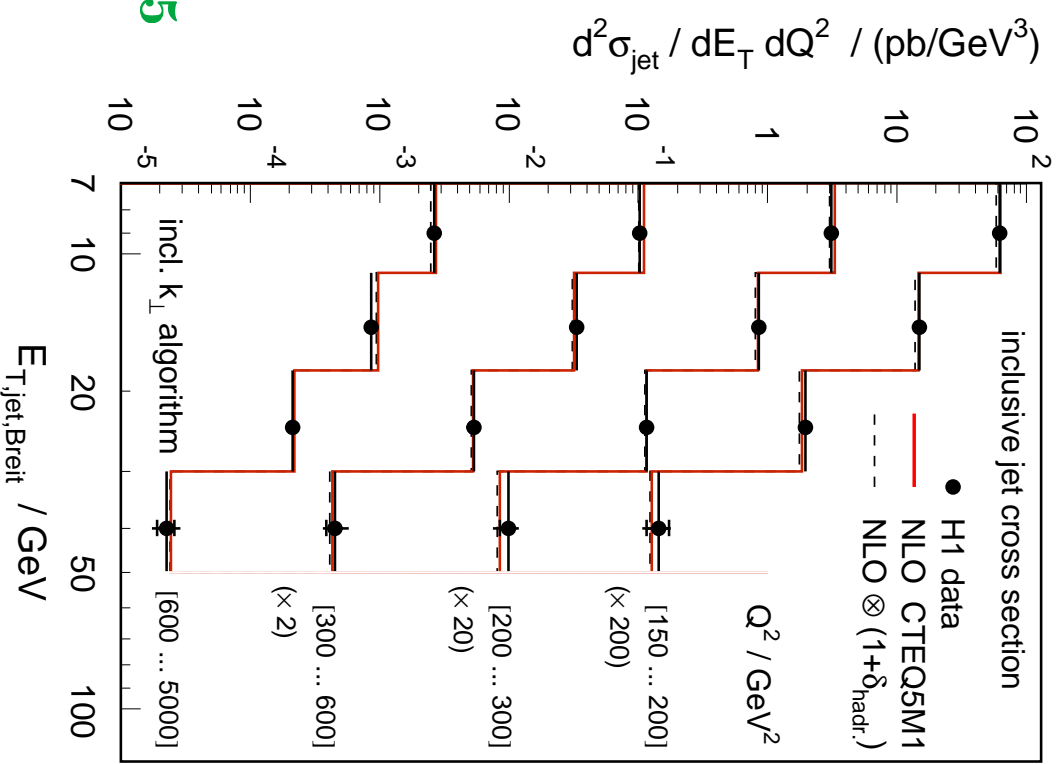
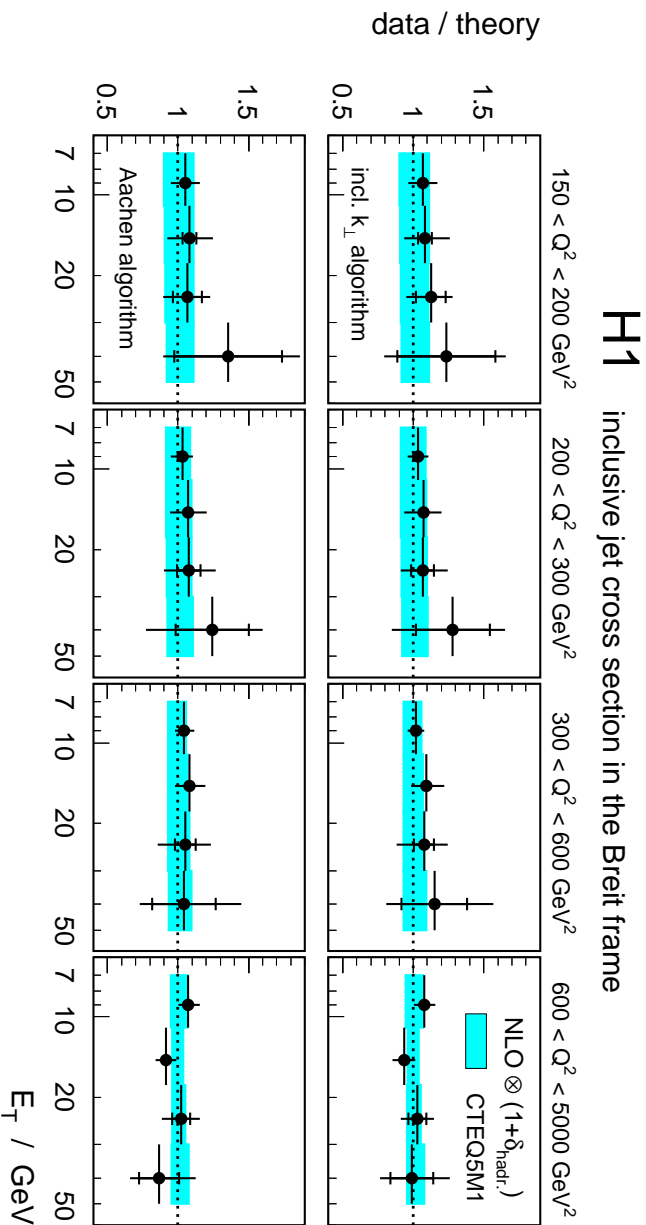
→ uncertainties proton PDFs $\Delta\alpha_s(M_Z) = +0.0012$
 -0.0011

→ hadronisation corrections $\Delta\alpha_s(M_Z) = \pm 0.0005$

Improvements depend upon further Theoretical Work



Inclusive jet Cross Sections in NC DIS and extraction of α_s



- Measurement of the differential cross section

$d\sigma / dE_T dQ^2$ for inclusive jet production with

E_T^{jet} (Breit) > 7 GeV and $-1 < \eta^{jet,1(2)}$ (Lab) < 2.5

over $150 < Q^2 < 5000$ GeV²

- The complication of selecting dijet events is absent!

- NLO QCD gives a good description of the data over the whole range in E_T and Q^2

Determination of α_s

- The inclusive jet cross section $d\sigma/dE_T dQ^2$ over

$150 < Q^2 < 5000 \text{ GeV}^2$ has been fitted using NLO

QCD calculations and the CTEQ5M1 proton PDFs

to determine $\alpha_s(E_T)$ and a combined value of $\alpha_s(M_Z)$

$$\alpha_s(M_Z) = 0.1186 \pm 0.0030 \text{ (exp.)}$$

$$+0.0039 \text{ (th.)} +0.0033 \text{ (pdf.)}$$

$$-0.0045 \text{ (th.)} -0.0023 \text{ (pdf.)}$$

- No significant dependence on the α_s assumed in PDFs

- The theoretical uncertainty dominates:

→ higher orders and hadronisation $\Delta\alpha_s = +0.0039$
 -0.0045

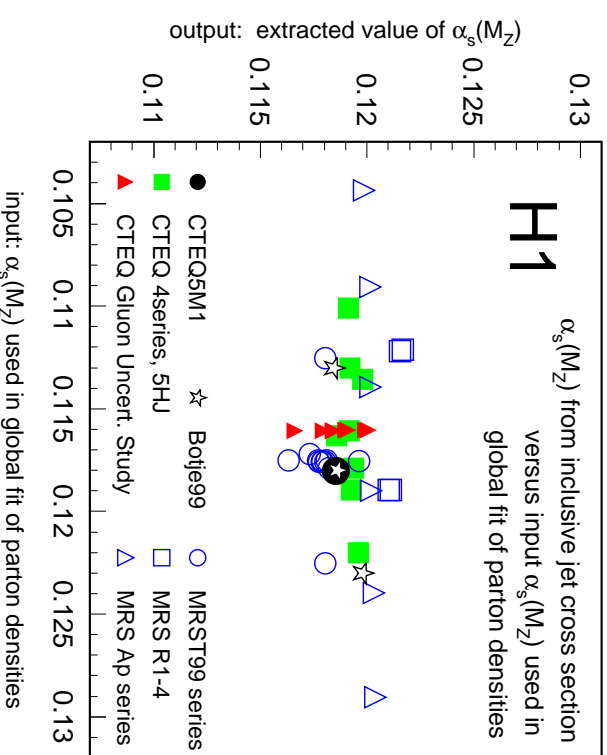
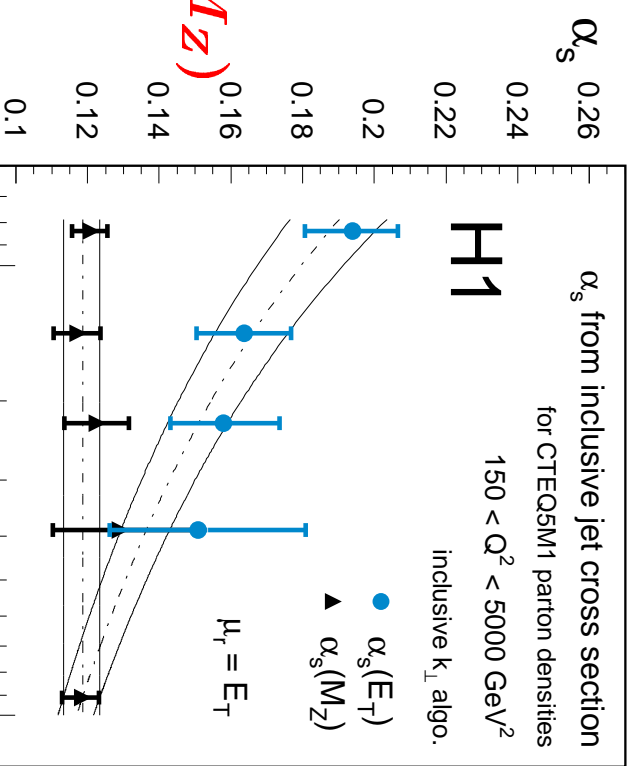
→ uncertainties on the proton PDFs $\Delta\alpha_s = +0.0033$
 -0.0023

- The observed E_T dependence of α_s is consistent

with the running predicted by QCD

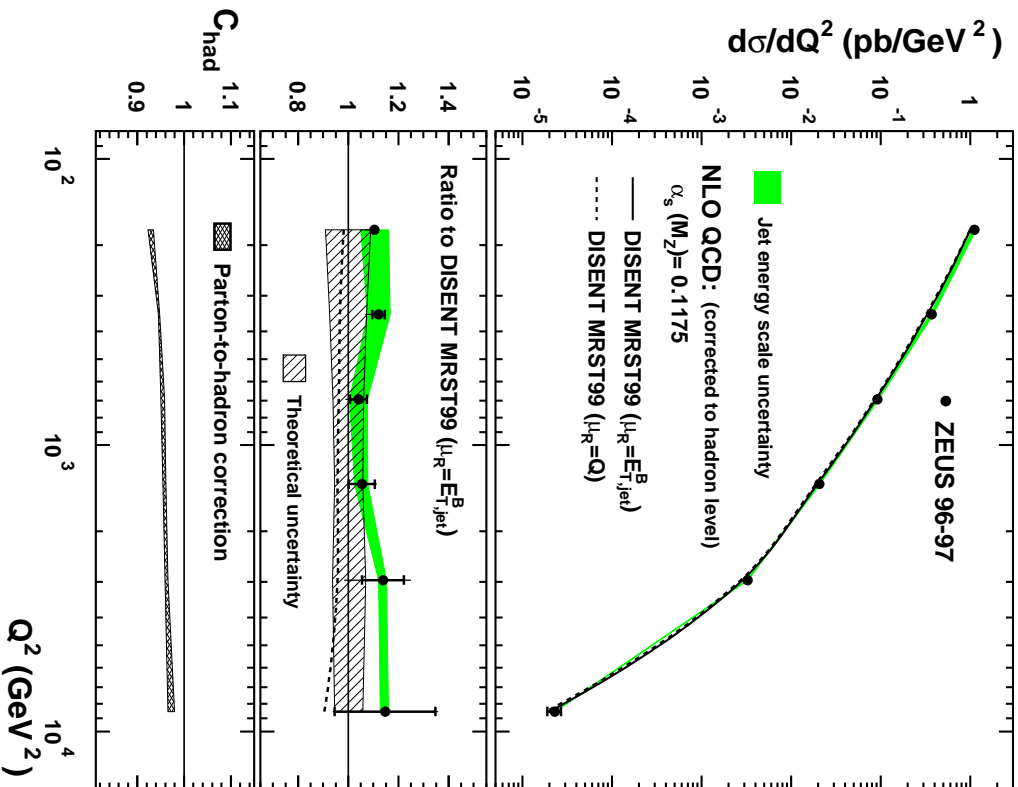
- Consistent with determination of α_s in e^+e^-

- Competitive accuracy in determination of α_s

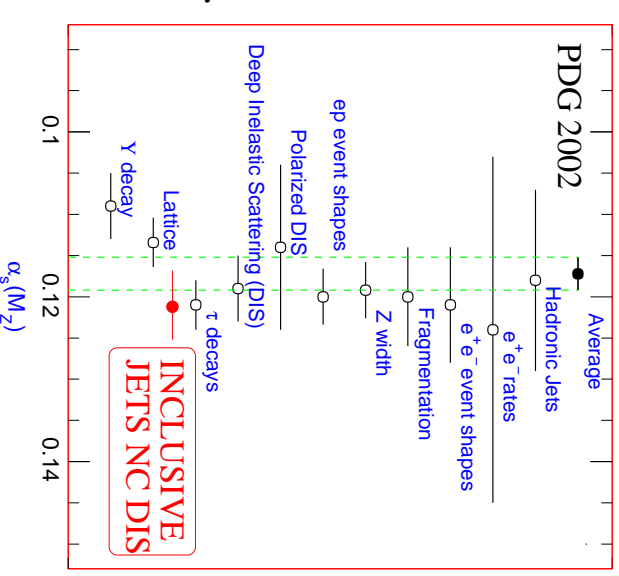


Inclusive Jet Production in NC DIS ($Q^2 > 125 \text{ GeV}^2$) and extraction of α_s

ZEUS



- Measurement (**NEW**) of inclusive jet cross sections in the Breit frame at $Q^2 > 125 \text{ GeV}^2$ for jets with $E_{T,jet}^B > 8 \text{ GeV}$ and $-2 < \eta_{jet}^B < 1.8$
 - no cut is applied in the laboratory frame
 - Extraction of α_s from $d\sigma/dQ^2$ at $Q^2 > 500 \text{ GeV}^2$
- $\alpha_s(M_Z) = 0.1212 \pm 0.0017 \text{ (stat.)}$
 $+0.0023 \text{ (exp.)} +0.0028 \text{ (th.)}$
 -0.0031
- Theoretical uncertainties:
 - terms beyond NLO (3%)
 - unc. proton PDFs (1%)
 - **Very precise determination of $\alpha_s(M_Z)$!**
 - Talk in WG-C by ZEUS speaker



Three-jet cross sections in NC DIS

- Three-jet cross sections test QCD beyond LO directly $\rightarrow \sigma_{3jet} \propto \alpha_s^2$

- At least three jets with

$$E_T^{jet} (\text{Breit}) > 5 \text{ GeV and}$$

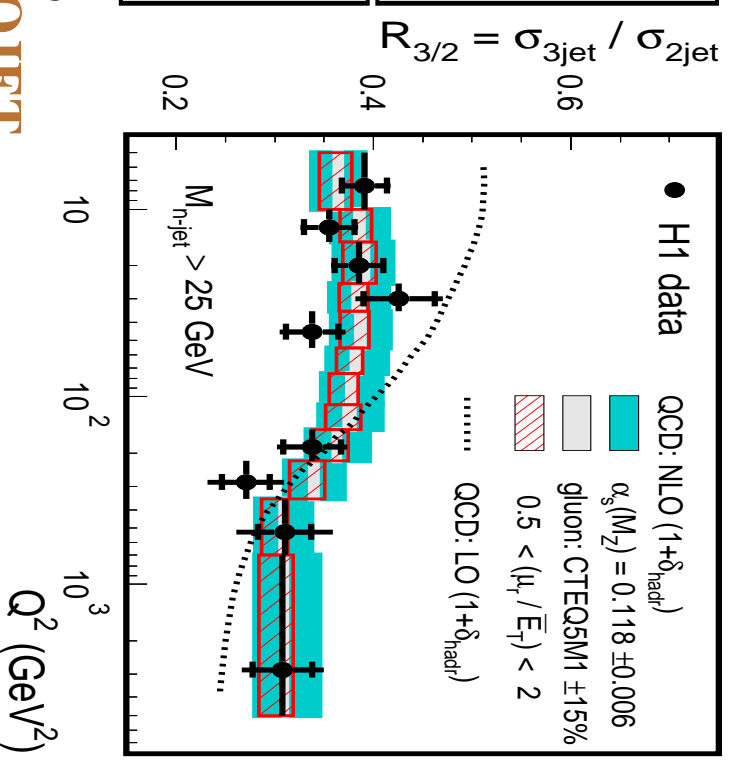
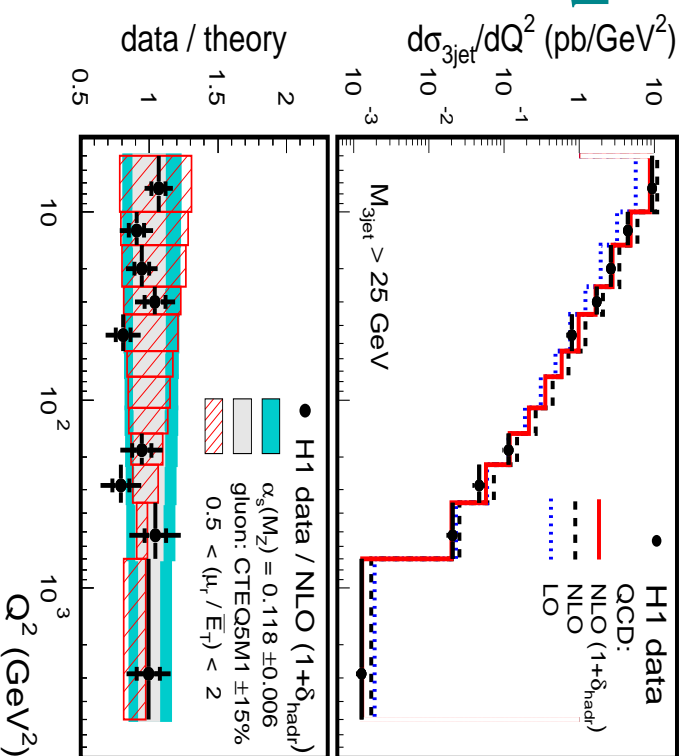
$$-1 < \eta^{jet} (\text{Lab}) < 2.5$$

- Kinematic region:

$$M^{3j} > 25 \text{ GeV,}$$

$$X_3 < 0.95 \text{ and}$$

$$|\cos \theta_3| < 0.8$$



- Comparison with NLO ($\mathcal{O}(\alpha_s^3)$) calculations using NLOJET

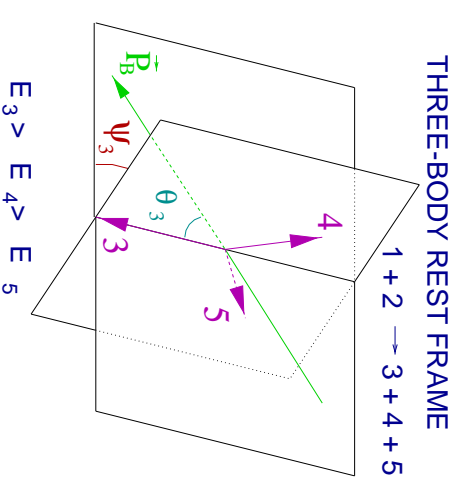
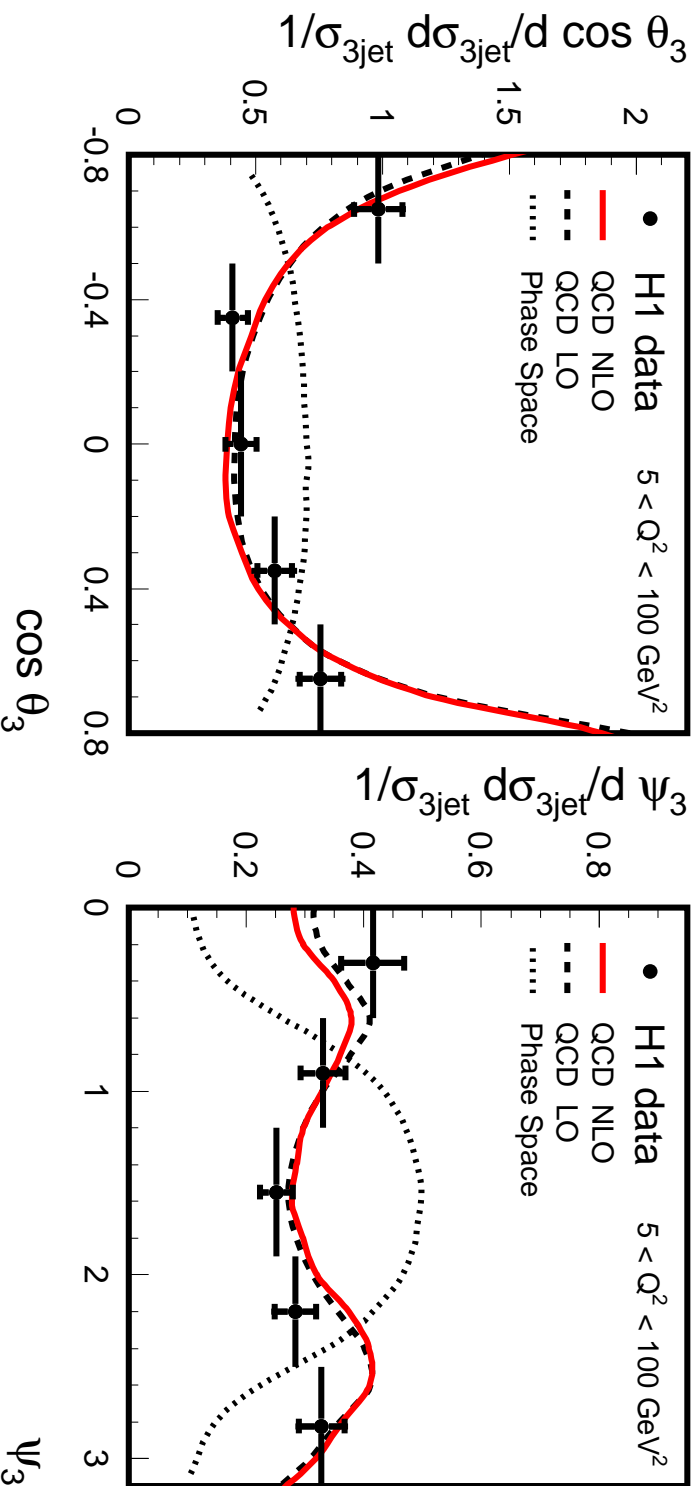
\rightarrow The NLO calculations give a good description of the data over the whole Q^2 range

- Theoretical uncertainties from PDFs and higher order terms reduced in the three-jet to dijet ratio but relatively large hadronisation corrections (10% – 18%)

\rightarrow Potentially useful observable to make an accurate determination of α_s

Three-jet cross sections in NC DIS

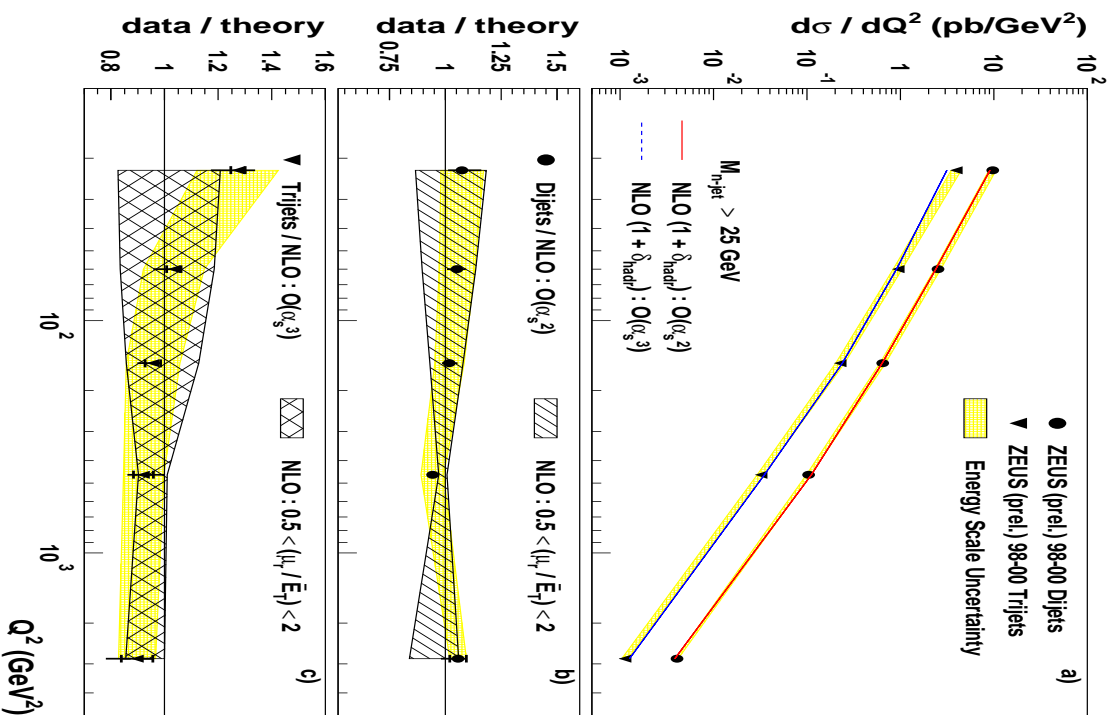
- θ_3 is sensitive to the spin of the exchanged particle
- ψ_3 reflects the orientation of the lowest energy jet



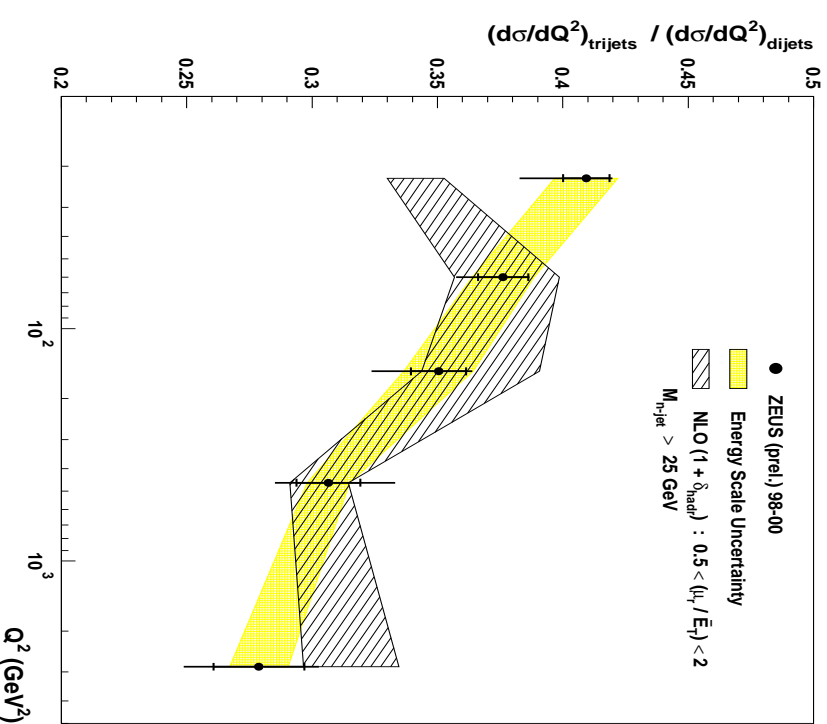
- The $\cos \theta_3$ distribution indicates that the highest energy jet tends to go either towards the proton or towards the photon directions
- The ψ_3 distribution probes the dynamics beyond LO directly
- The measured angular distributions are drastically different from phase space and in good agreement with QCD

Three-jet cross sections in NC DIS

ZEUS



ZEUS



- Measurements (**NEW**) of cross sections for three jet production in the Breit frame with $\mathcal{L} = 82 \text{ pb}^{-1}$
- Kinematic region: $10 < Q^2 < 5000 \text{ GeV}^2$
- $E_{T,jet}$ (Breit) $> 5 \text{ GeV}$, $-1 < \eta_{jet}(\text{Lab}) < 2.5$
- $M_{\text{jets}} > 25 \text{ GeV}$
- Talk in WG-C by N. Krumnack

Inclusive Jet Cross Sections in NC DIS at low Q^2 ($5 < Q^2 < 100 \text{ GeV}^2$)

- Measurement of the differential cross section

$d\sigma / dE_T^{jet}$ (Breit) for inclusive jet production with $E_T^{jet} > 5 \text{ GeV}$

- Comparison to NLO calculations

with $\mu_R = E_T^{jet}$ and $\mu_R = Q$

→ good description of the data

in the region $-1 < \eta_{jet} < 0.5$

→ failure to describe the data at low E_T^{jet}

in the region $1.5 < \eta_{jet} < 2.8$,

where the NLO corrections and

uncertainties (terms beyond NLO)

are largest, and largest differences

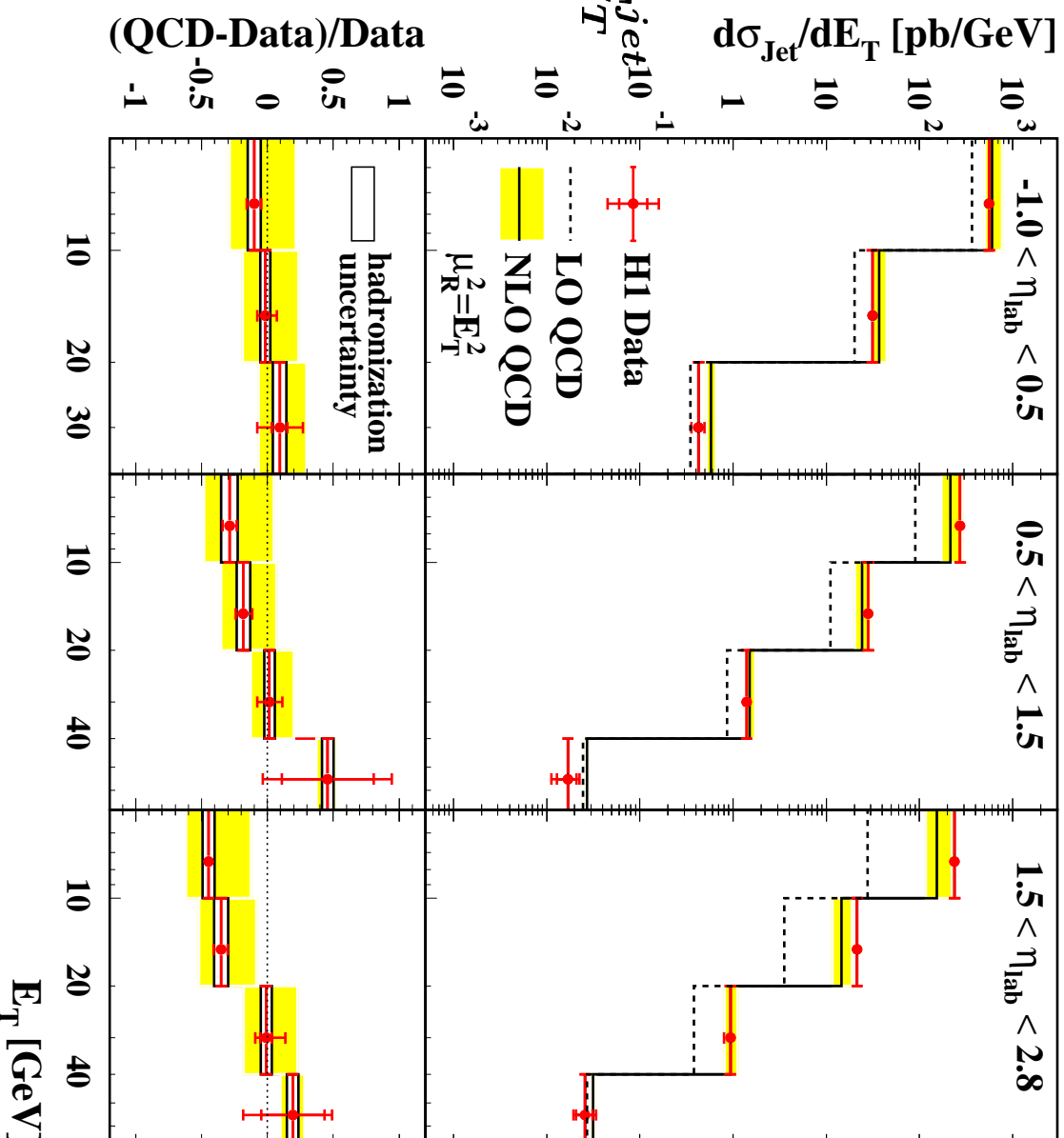
for $5 < Q^2 < 10 \text{ GeV}^2$

⇒ Improved calculations needed

to understand jet production at low Q^2

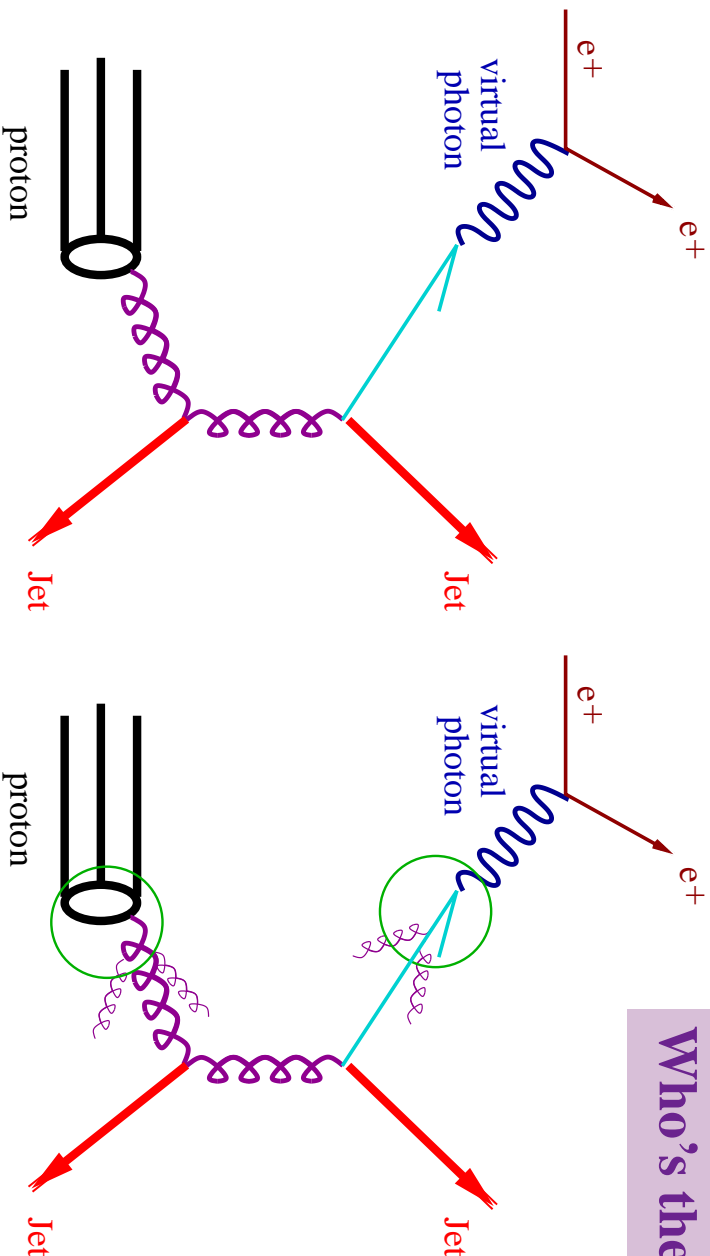
- Talk in WG-C by A. Specka

H1 Inclusive Jets



Jet Production in NC DIS at low Q^2

Who's the probe/target when $E_T^{jet} \gg Q$?



- Large logarithms $\log E_T / Q$ can be resummed by defining PDFs for a virtual photon in a similar fashion as for the proton
- Improving the calculations by including higher-order terms

- Several models for the virtual photon PDFs exist, which interpolate between the
 - leading-log part of the real photon PDFs $\Rightarrow \log E_T / \Lambda_{QCD}$
 - and the asymptotic behaviour predicted by pQCD $\Rightarrow \log E_T / Q$
- Parametrisations of the virtual photon PDFs:
 - Drees and Godbole, Schuler and Sjöstrand, Glück et al

Dijet Production in NC DIS at low Q^2 ($1.6 < Q^2 < 80 \text{ GeV}^2$)

- Measurement of dijet cross sections for jets defined in the γ^*p centre-of-mass system

using the inclusive k_T algorithm

- Phase-space region:

$$|\eta^{jet,1} - \eta^{jet,2}| < 1$$

$$-3 < \bar{\eta} < -0.5$$

$$\bar{E}_T^2 > 30 \text{ GeV}^2 \text{ and } |E_T^{jet,1} - E_T^{jet,2}| / (E_T^{jet,1} + E_T^{jet,2}) < 0.25$$

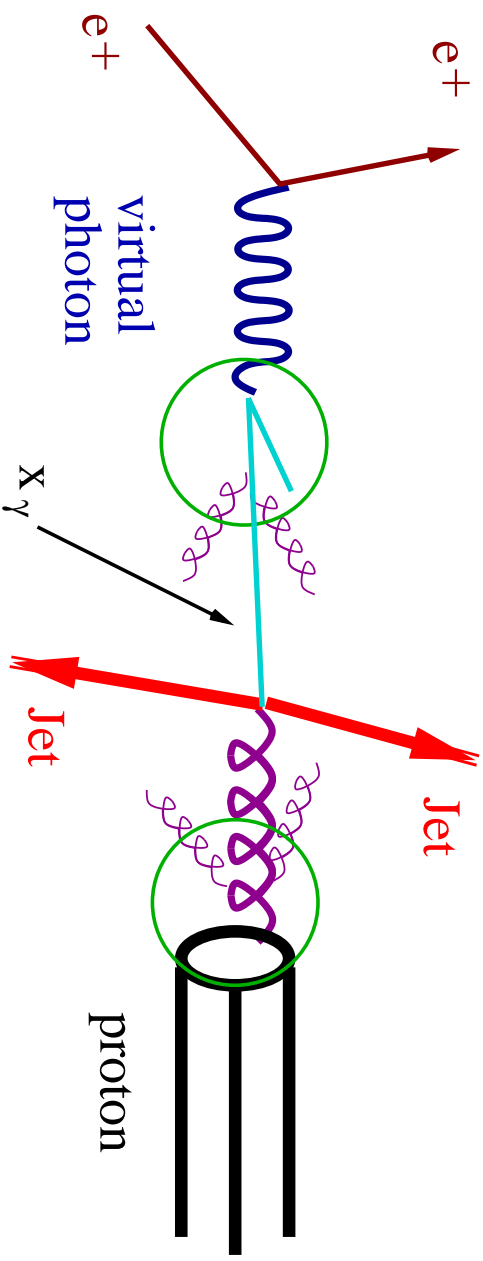
$$1.6 < Q^2 < 80 \text{ GeV}^2 \text{ and } 0.1 < y < 0.7$$

- Useful variable: the longitudinal momentum fraction of the photon's momentum carried by the interacting parton

$$x_\gamma^{jets} = \frac{\sum_{jets\ 1,2} (E^{jet} - p_z^{jet})}{\sum_{\text{all hadrons}} (E^h - p_z^h)} \quad (\text{sum over the two highest-} E_T^{jet} \text{ jets})$$

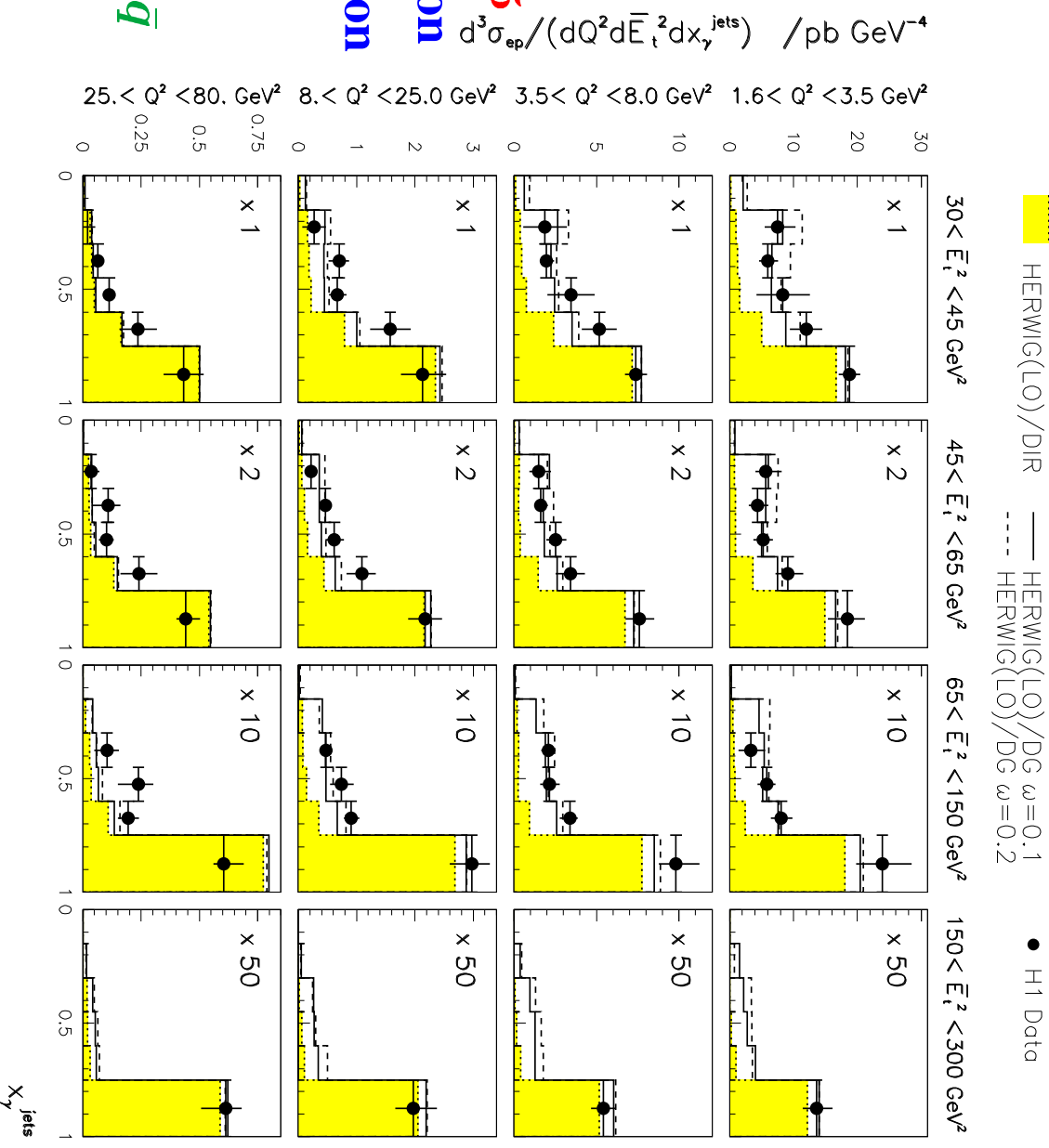
→ to separate “resolved” photons ($x_\gamma^{jets} < 1$) from “direct” photons ($x_\gamma^{jets} \approx 1$)

⇒ Measurement of the triple-differential cross section $d^3\sigma / dQ^2 d\bar{E}_T^2 dx_\gamma^{jets}$



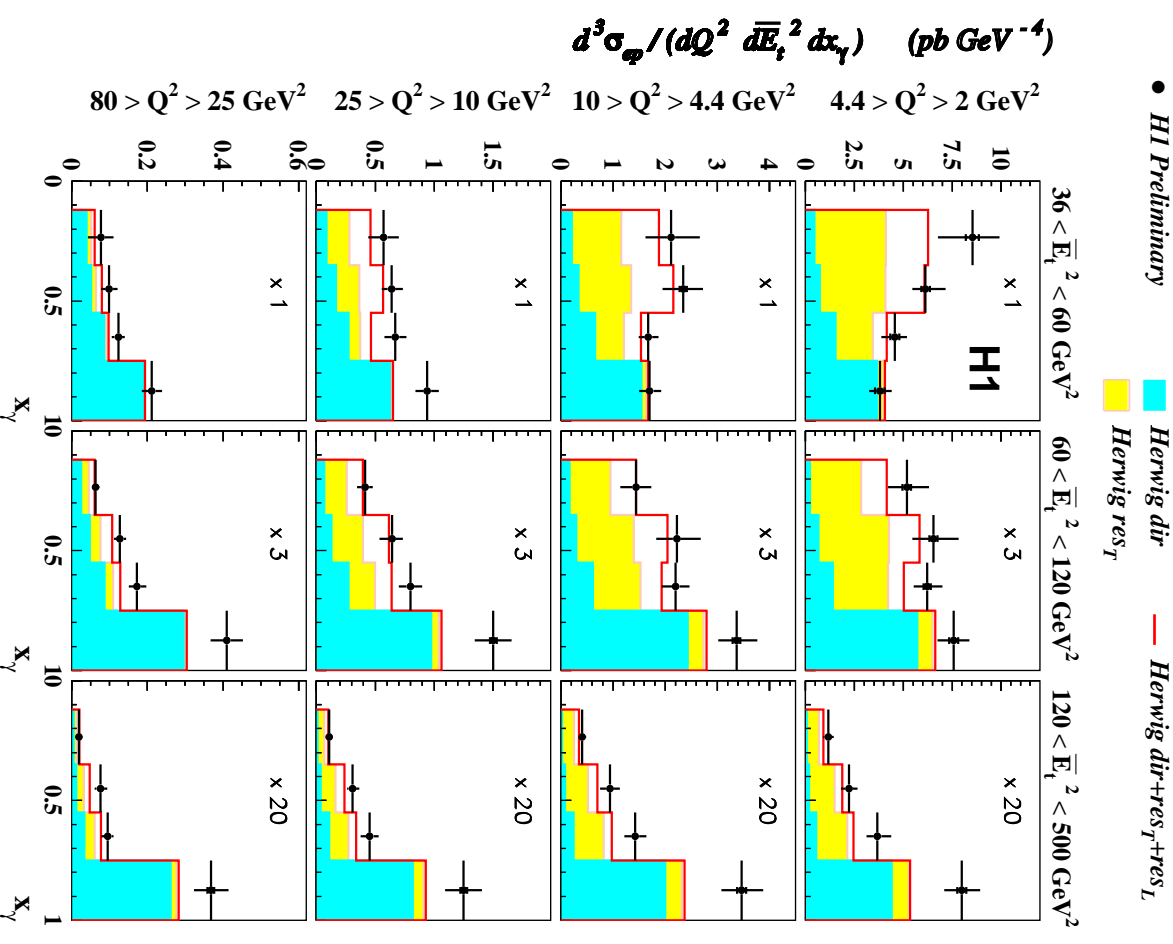
Dijet Production in NC DIS at low Q^2 ($1.6 < Q^2 < 80 \text{ GeV}^2$)

- The distributions peak towards $x_\gamma^{jets} = 1$
 - “direct” photon contribution
- As Q^2 increases at fixed \bar{E}_T the cross section in the region $x_\gamma^{jets} < 0.75$ (“resolved” photon) decreases
- Comparison with models:
 - “direct” photon only does not reproduce the data for $x_\gamma^{jets} < 0.75$
 - inclusion of the “resolved” photon contribution improves the description
 - ⇒ Consistent with the pQCD expectation that as $Q^2 \rightarrow \bar{E}_T^2$ the photon structure reduces to $\gamma \rightarrow q\bar{q}$



Dijet Production in NC DIS at low Q^2 ($2 < Q^2 < 80 \text{ GeV}^2$)

- **Measurements (NEW) of dijet cross sections in the kinematic region $2 < Q^2 < 80 \text{ GeV}^2$ and $0.1 < y < 0.85$ for jets with $E_T^{jet} > 5 \text{ GeV}$, $\bar{E}_T > 6 \text{ GeV}$ and $-2.5 < \eta_{jet}(\gamma^*p \text{ CMS}) < 0$**
- **Comparison with models which incorporate**
 - **only “direct” photon contribution**
 - **transversally-polarised “resolved” photons**
 - **longitudinally-polarised “resolved” photons**
- k_T -unordered initial-state parton cascades based on CCFM evolution equations and an unintegrated gluon density function
- **Improved description of the data**
- Talk in WG-C by H. Jung



Parton evolution at low x

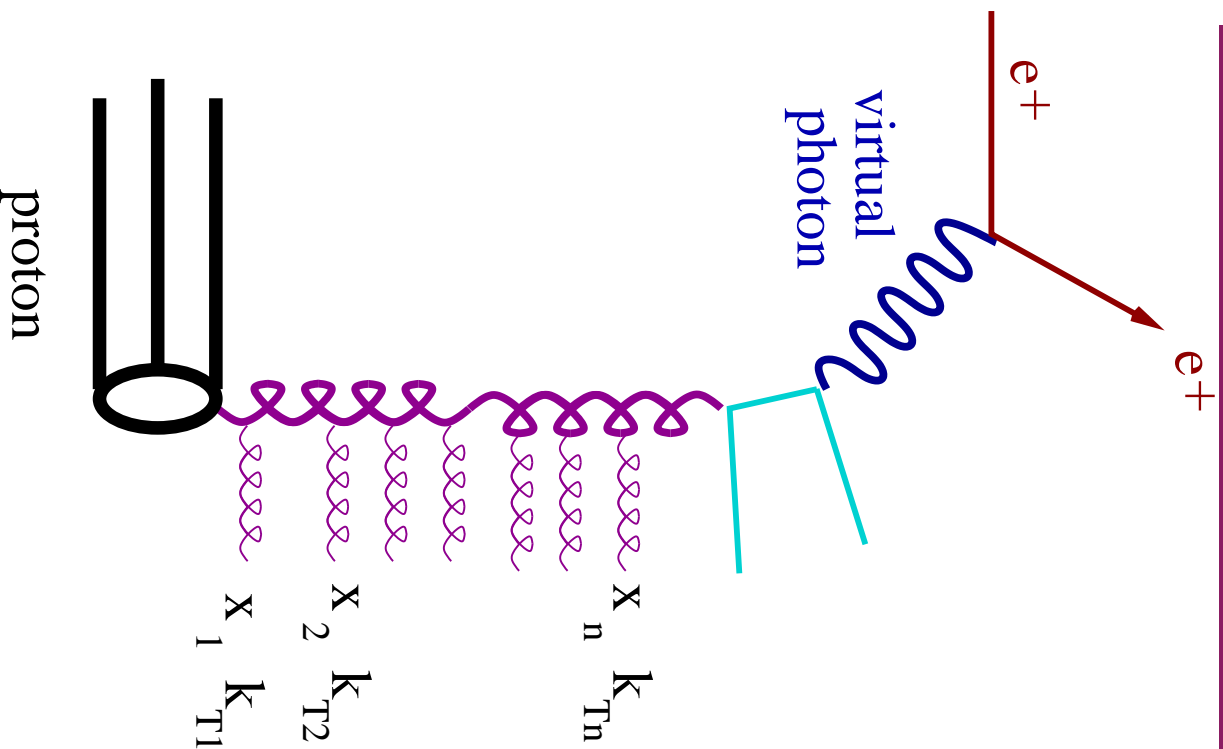
Searching for BFKL-induced effects

- DGLAP equations sum the leading powers of $\alpha_s \log Q^2$ in the region of strongly-ordered transverse momenta

$$Q^2 \gg k_{Tn}^2 \gg \dots \gg k_{T2}^2 \gg k_{T1}^2$$

- When $\log Q^2 \ll \log 1/x$ terms proportional to $\alpha_s \log 1/x$ become important and need to be summed
- the BFKL equation accomplishes that; the integration is taken over the full k_T phase space of the gluons
- ⇒ no k_T ordering

- Mueller and Navelet's proposal:
forward (proton's direction) jet production
with x_1/x as large as possible
and $k_{T1} \sim Q$



Measurement of Forward Jet Production at low x

- Measurement of the differential cross section $d\sigma/dx$ for jet production (cone algor.) in the kinematic region

$$10^{-4} < x < 4 \cdot 10^{-3}$$

$$x_{jet} = E_{jet}/E_p > 0.035$$

$$0.5 < p_{T,jet}^2/Q^2 < 2, 7^\circ < \theta_{jet} < 20^\circ$$

$$p_{T,jet} > 3.5 \text{ GeV (5 GeV)}$$

- Strong rise towards low x is observed

- Comparison to calculations

→ fixed-order ($\mathcal{O}(\alpha_s^2)$) lie well below

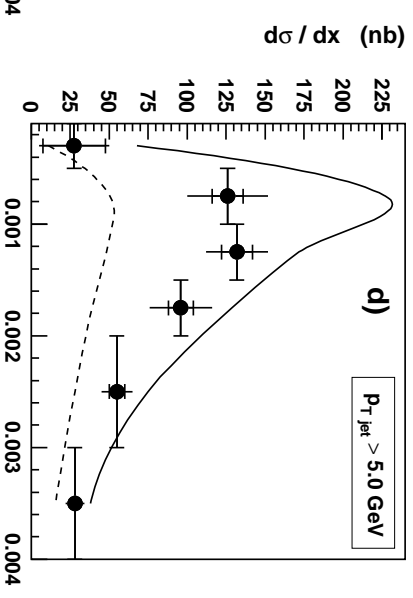
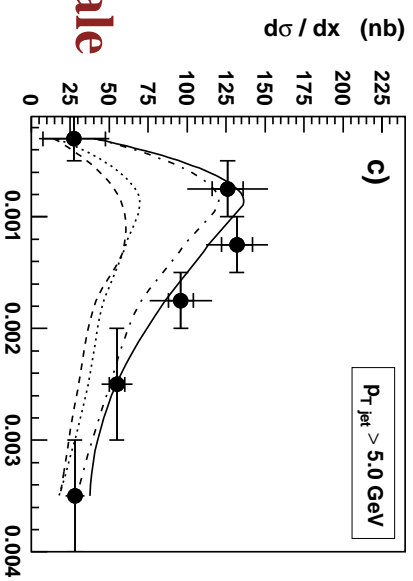
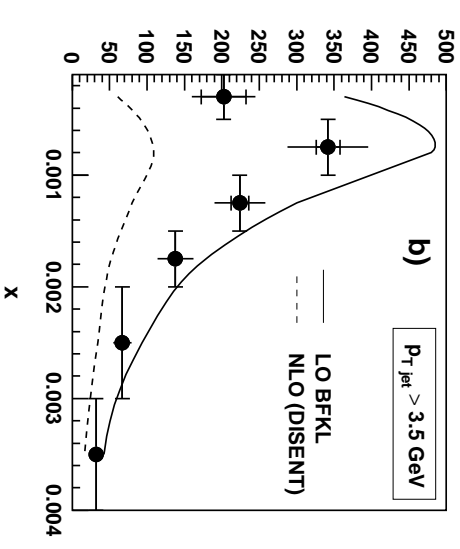
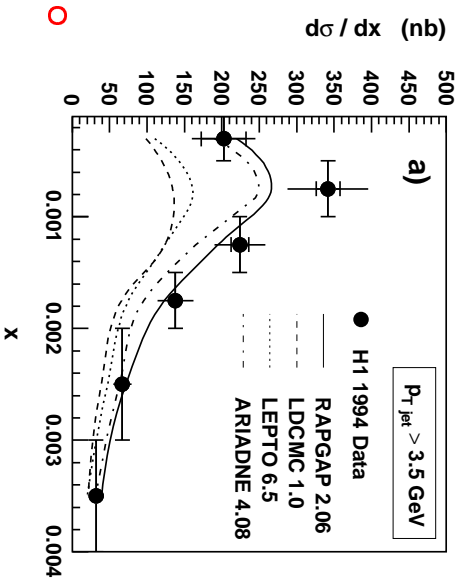
the data at low x ; variation of the scale

(Q^2) leads to 40% uncertainty

→ comparison to BFKL calculation hampered by lack of higher-order and hadronisation corrections, and implementation of jet algorithm

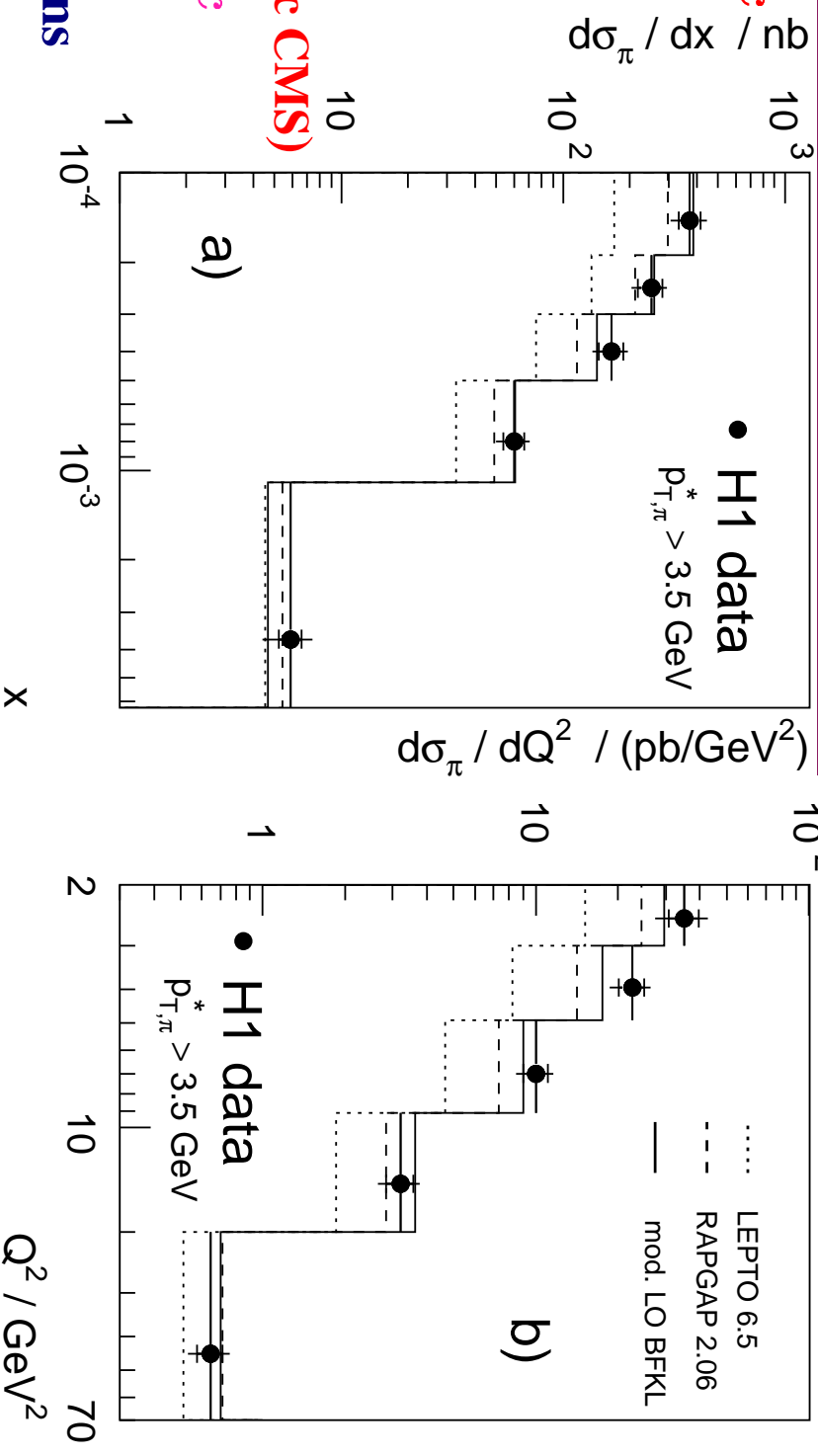
→ models with extra parton radiation reproduce the data (interpretation?)

- Talks in WG-C by L. Goerlich, S. Lammers and R. Poeschl



Measurement of Forward π^0 Production at low x

- Measurement of $d\sigma_\pi/dx$ for π^0 production in the kinematic region
 - $2 < Q^2 < 70 \text{ GeV}^2$
 - $x_\pi = E_\pi/E_p > 0.01$
 - $5^\circ < \theta_\pi < 25^\circ$
 - $p_{T,\pi}^* > 3.5 \text{ GeV}$ (hadronic CMS)
- Strong rise towards low x is observed



- Comparison to calculations

→ a model (LEPTO) based on parton emission according to the DGLAP splitting functions agrees with the data at high x , but falls below the data at low x

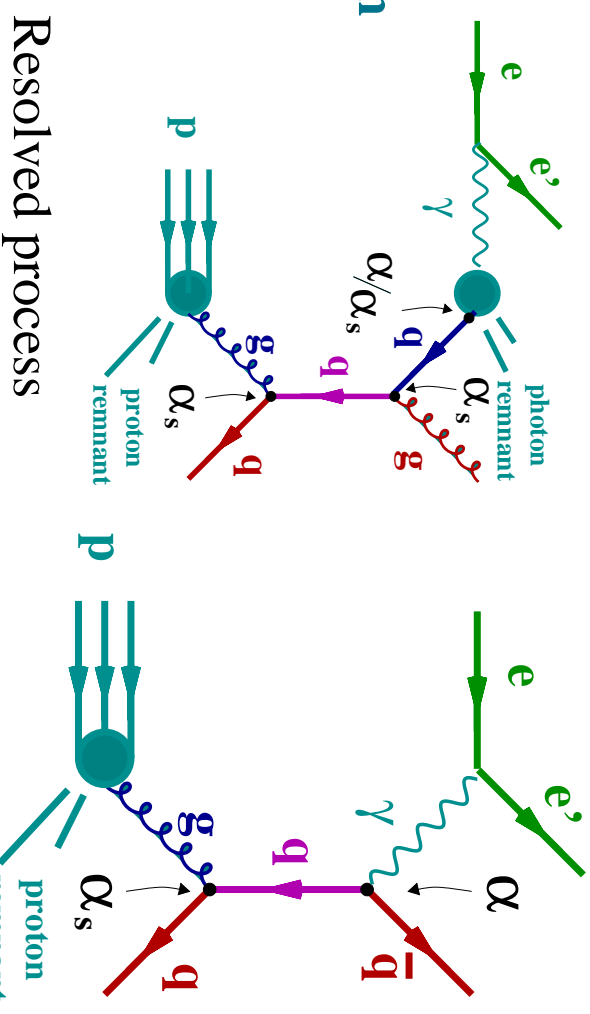
→ a modified LO BFKL evolution equation (Kwiecinski et al) convoluted with a π^0

fragmentation function gives a good description; 60% scale uncertainty on normalisation

- Talk in WG-C by L. Goerlich

Photoproduction of Jets

- Production of jets in γp collisions has been measured via ep scattering at $Q^2 \approx 0$
- At lowest order QCD, two hard scattering processes contribute to jet production \Rightarrow
- pQCD calculations of jet cross sections



$$d\sigma_{jet} = \sum_{a,b} \int_0^1 dy f_{\gamma/e}(y) \int_0^1 dx_\gamma f_{a/\gamma}(x_\gamma, \mu_{F\gamma}^2) \int_0^1 dx_p f_{b/p}(x_p, \mu_{Fp}^2) d\hat{\sigma}_{ab \rightarrow jj}$$

longitudinal momentum fraction of γ/e^+ (y), parton a/γ (x_γ), parton b /proton (x_p)

$\rightarrow f_{\gamma/e}(y)$ = flux of photons in the positron (WW approximation)

$\rightarrow f_{a/\gamma}(x_\gamma, \mu_{F\gamma}^2)$ = parton densities in the photon (for direct processes $\delta(1 - x_\gamma)$)

$\rightarrow f_{b/p}(x_p, \mu_{Fp}^2)$ = parton densities in the proton

$\rightarrow \sigma_{ab \rightarrow jj}$ subprocess cross section; short-distance structure of the interaction

Photoproduction of Jets

● Measurements of jet photoproduction provide

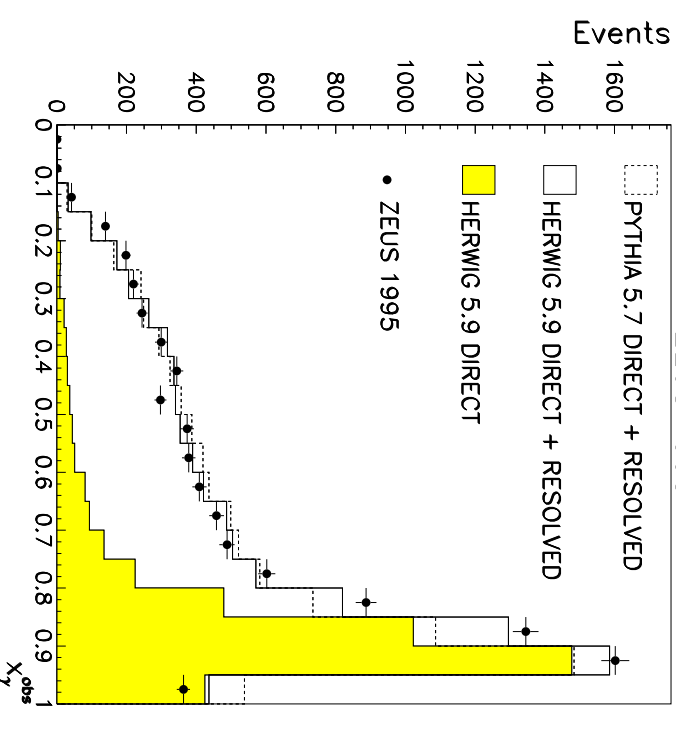
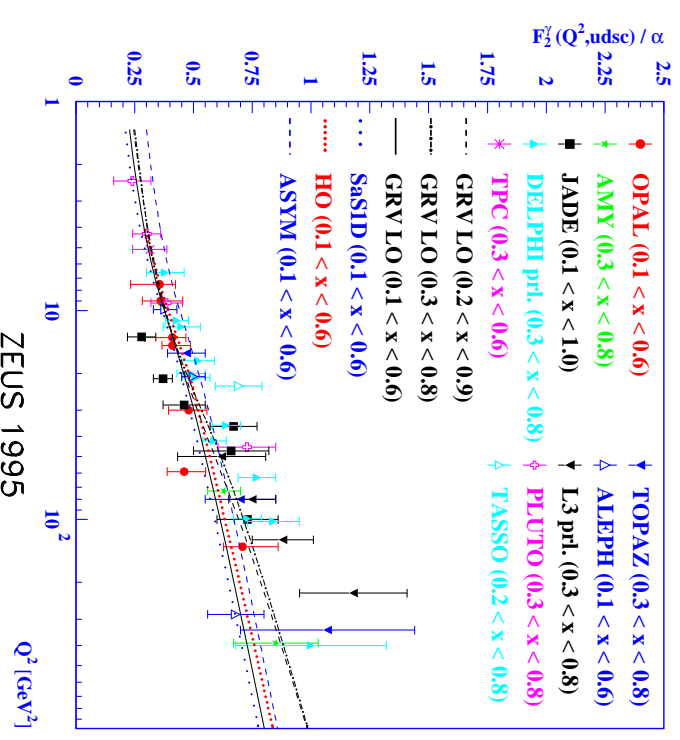
- Test of NLO QCD predictions based on current parametrisations of the proton and photon PDFs
- Dynamics of resolved and direct processes
- Photon structure: information on quark densities from F_2^γ in e^+e^- ; gluon density poorly constrained.

Jet cross sections in photoproduction are sensitive to both the quark and gluon densities in the photon at larger scales $\mu_{F_\gamma}^2 \sim E_{T,jet}^2$ ($200 - 10^4 \text{ GeV}^2$)

→ Proton structure: well constrained by DIS except for the gluon density at high x . Jet cross sections in γp are sensitive to parton densities at x_p up to ~ 0.6

● Observable to separate the contributions: the fraction of the photon's energy participating in the production of the dijet system

$$x_\gamma^{OBS} = \frac{1}{2E_\gamma} \sum_{i=1}^2 E_T^{jet_i} e^{-\eta^{jet_i}}$$

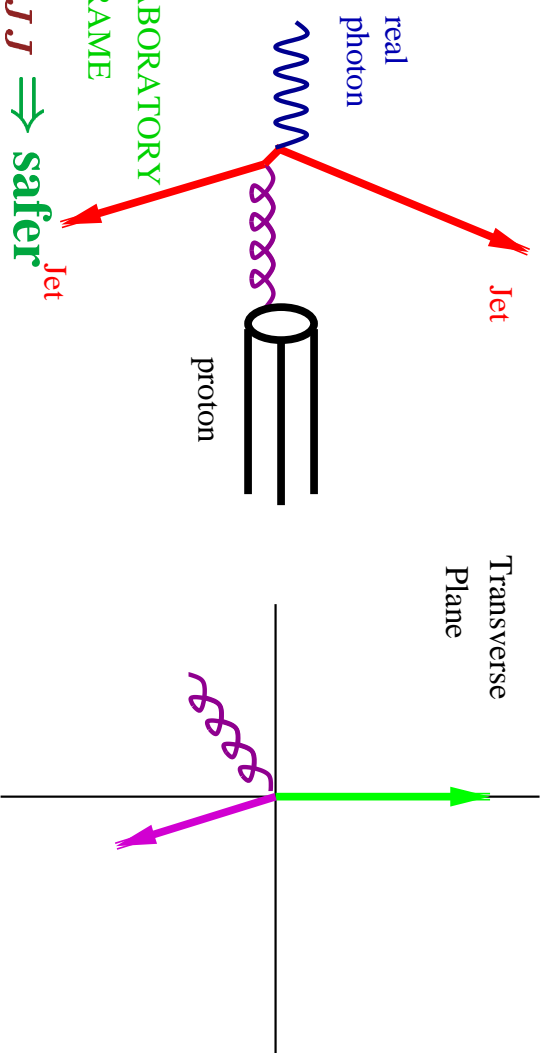


NLO QCD Calculations of Jet Cross Sections in Photoproduction

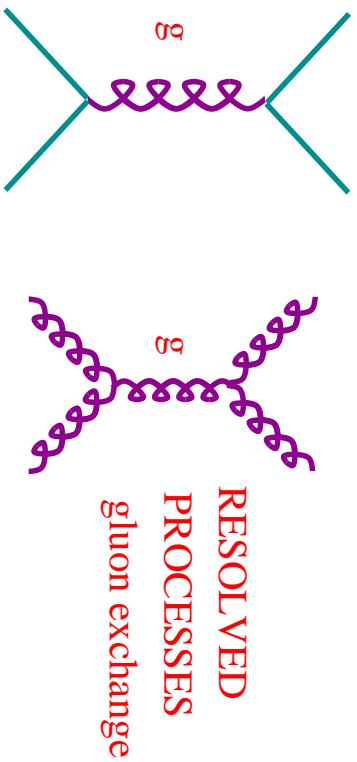
- Longitudinally invariant k_T -cluster algorithm in the η - ϕ plane of the laboratory frame
- Several NLO QCD calculations are available (phase-space slicing, subtraction method)
 - Klasen and Kramer, Harris and Owens, Aurenche et al, Frixione and Ridolfi
 - choice of scales: $\mu_R = \mu_F \sim E_T^{jet}$ (Reminder: $Q^2 \approx 0$!)
 - photon PDFs: NLO QCD parametrisations GRY-HO, AFG-HO
 - proton PDFs: NLO QCD parametrisations CTEQ5M, CTQ5HJ, MRST99
- Parton-to-hadron corrections estimated with HERWIG, PYTHIA, PHOJET
- Effects from the underlying event estimated with: multiparton interactions in PYTHIA, SUE model of HERWIG, PHOJET model; parameters tuned to reproduce the energy flow outside of jets.

- Dijet selection criteria:

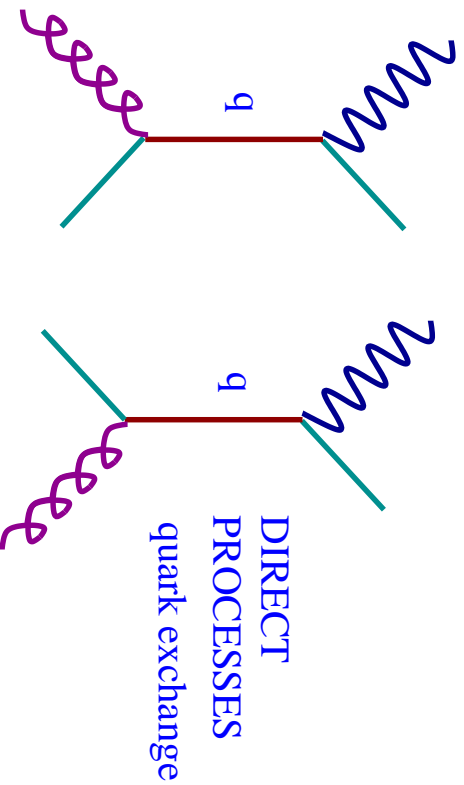
- Symmetric cuts on $E_T^{jet,1(2)} \Rightarrow$ danger
- Asymmetric cuts on $E_T^{jet,1(2)} \Rightarrow$ safer
- Symmetric cuts on $E_T^{jet,1(2)}$ and cut on $M_{JJ} \Rightarrow$ safer



Dijet Photoproduction: the dynamics of resolved and direct processes



RESOLVED
PROCESSES
gluon exchange



DIRECT
PROCESSES
quark exchange

- The dynamics of dijet production has been investigated by studying the variable:

$$\cos \theta^* \equiv \tanh\left(\frac{1}{2}(\eta^{jet,1} - \eta^{jet,2})\right)$$

- for two-to-two parton scattering θ^* coincides with the scattering angle in the dijet CMS

- QCD predicts different dijet angular distributions for resolved and direct:

- **Resolved (gluon-exchange dominated)**

$$d\sigma/d|\cos \theta^*| \sim \frac{1}{(1-|\cos \theta^*|)^2}$$

- **Direct (quark-exchange only)**

$$d\sigma/d|\cos \theta^*| \sim \frac{1}{(1-|\cos \theta^*|)^1}$$

- The dijet angular distribution $d\sigma/d|\cos \theta^*|$ for $x_\gamma^{OBS} < 0.75$ (“resolved”) should be steeper than that of $x_\gamma^{OBS} > 0.75$ (“direct”) as $|\cos \theta^*| \rightarrow 1$

Dijet Photoproduction: the dynamics of resolved and direct processes

- Measurement of the dijet differential cross

section $d\sigma/d|\cos\theta^*|$ for dijet events with

$$E_T^{jet,1} > 14 \text{ GeV}, E_T^{jet,2} > 11 \text{ GeV}$$

$$-1 < \eta^{jet} < 2.4 \text{ (both jets)}$$

in the kinematic region

$$Q^2 < 1 \text{ GeV}^2 \text{ and } 134 < W_{\gamma p} < 277 \text{ GeV}$$

- Phase-space region:

$$|\cos\theta^*| < 0.8, \quad M_{JJ} > 42 \text{ GeV}$$

$$0.1 < \frac{1}{2}(\eta^{jet,1} + \eta^{jet,2}) < 1.3$$

- Comparison with NLO QCD calculations:

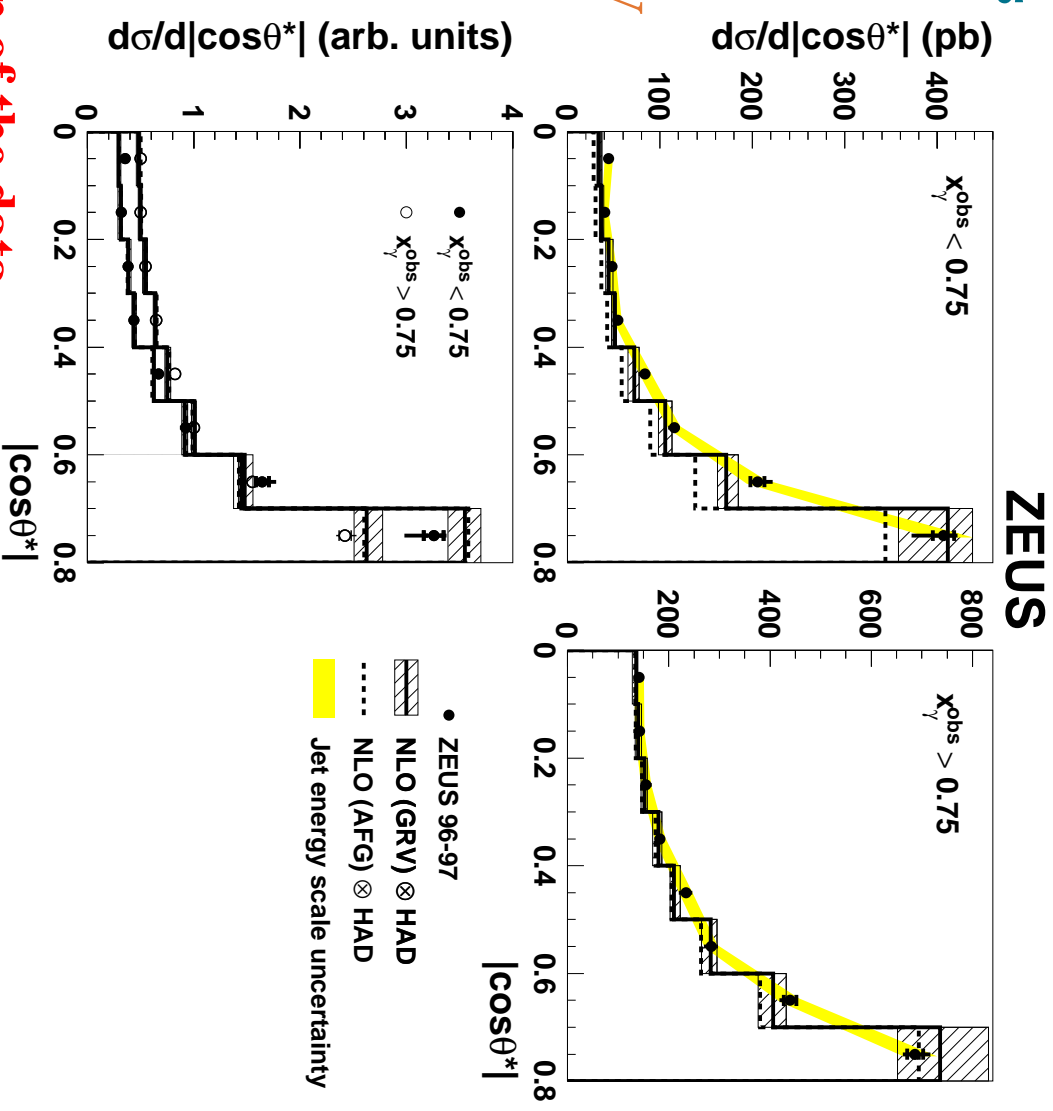
→ High- $x_\gamma^{OB\!S}$ (“direct”): NLO describes

the shape and normalisation of the data

→ Low- $x_\gamma^{OB\!S}$ (“resolved”): NLO describes

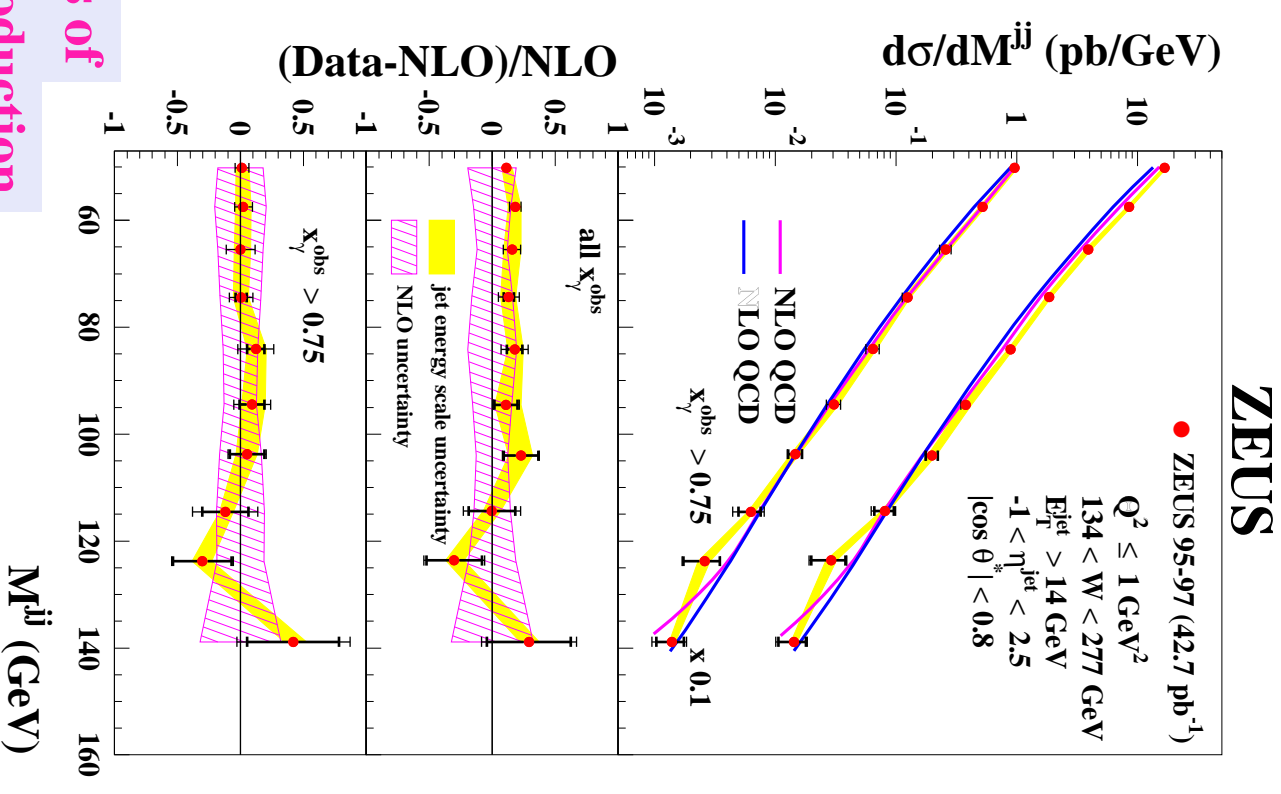
the shape and (reasonably) the normalisation of the data

- The dijet angular distribution of the “resolved” sample is steeper than that of “direct”



High- M_{JJ} Dijet Photoproduction

- Measurement of the dijet differential cross section $d\sigma/dM_{JJ}$ in the range $47 < M_{JJ} < 160$ GeV for dijet events with $E_T^{jet} > 14$ GeV, $-1 < \eta^{jet} < 2.5$ and $|\cos \theta^*| < 0.8$
- Small experimental uncertainties:
 - jet energy scale known to 1% \Rightarrow 5% on $d\sigma/dM_{JJ}$
- Small theoretical uncertainties:
 - higher-order terms (varying μ_R) below 15%
 - γ PDFs (GRV-HO, AFG-HO) below 10%
 - resolved processes suppressed at high M_{JJ}
 - small hadronisation corrections, below 5%
- NLO QCD calculations describe the shape and normalisation of the measurements well
 - Validity of the pQCD description of the dynamics of parton-parton and γ -parton interactions in photoproduction



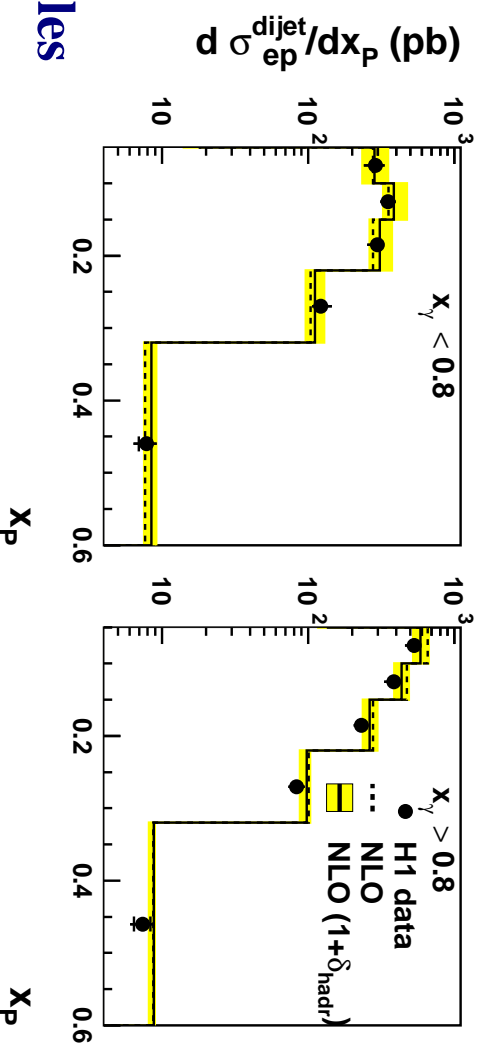
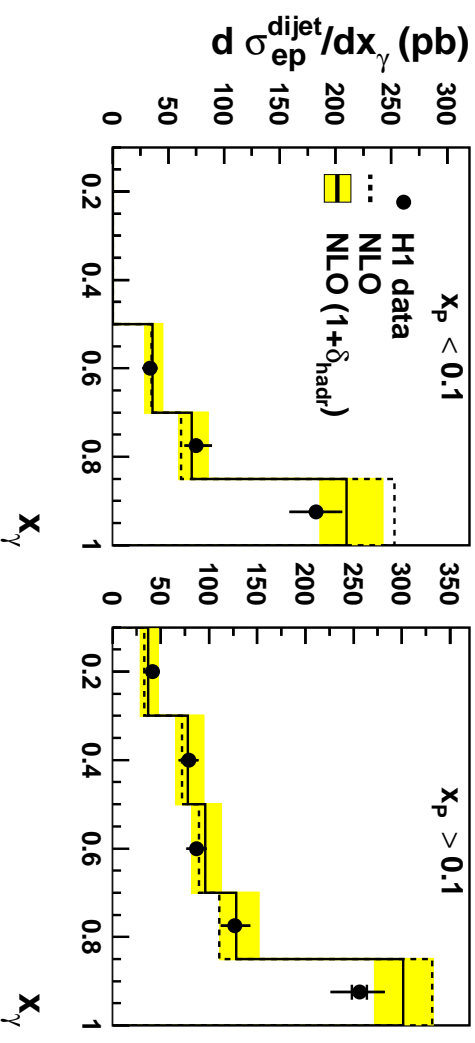
Dijet Photoproduction: photon and proton structure

- Measurement of the dijet cross sections $d\sigma/dx_\gamma$ and $d\sigma/dx_p$ for dijet events with

$$E_{T, \text{max}} > 25 \text{ GeV}, E_{T, \text{second}} > 15 \text{ GeV} \text{ and } -0.5 < \eta^{\text{j}et} < 2.5 \text{ (both jets)}$$

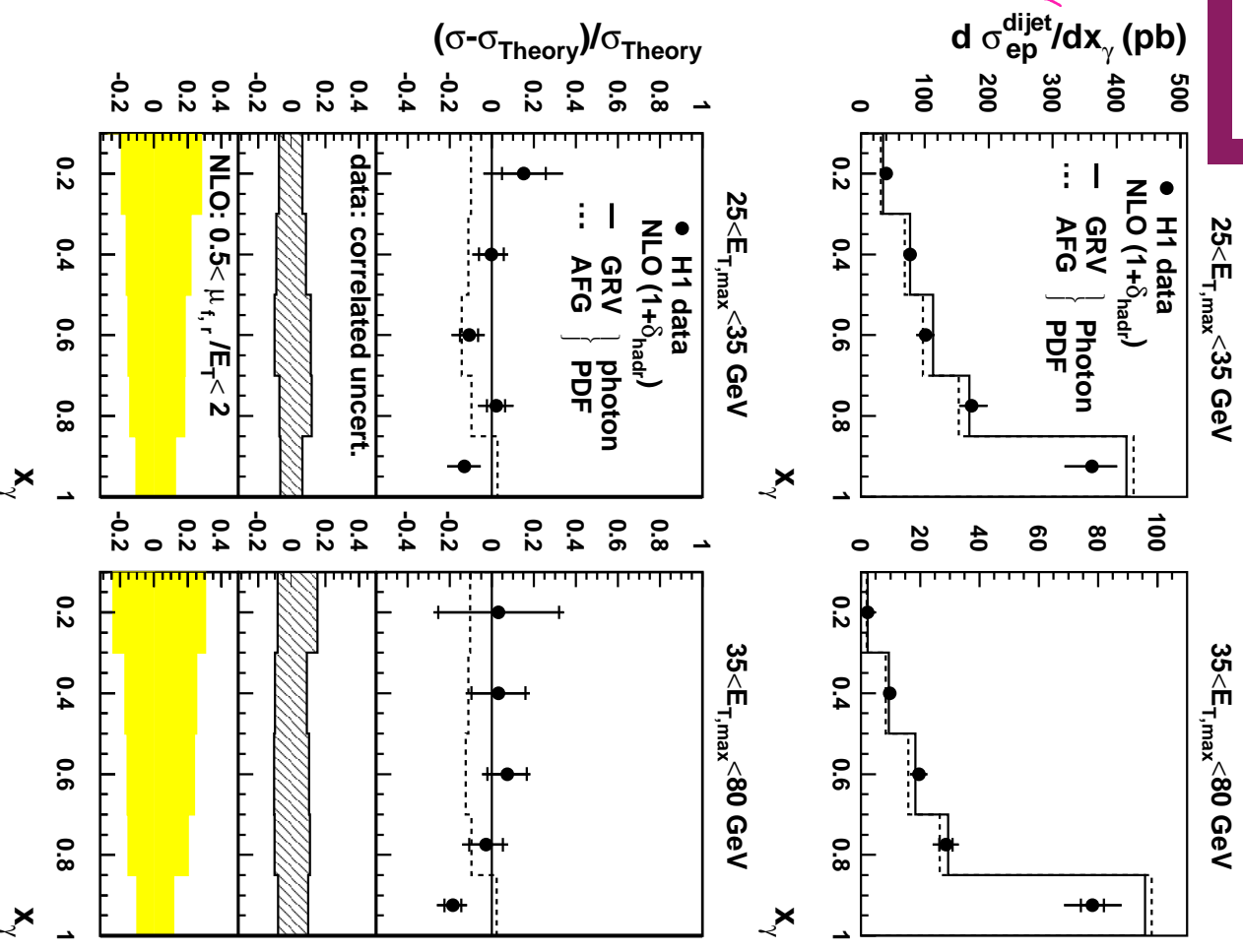
in the kinematic region $Q^2 < 1 \text{ GeV}^2$ and $95 < W_{\gamma p} < 285 \text{ GeV}$

- x_p variable: $x_p = \frac{1}{2E_p} \sum_{i=1}^2 E_T^{\text{j}et_i} e^{\eta^{\text{j}et_i}}$
- NLO calculations using CTEQ5M (proton) and GRV-HO (photon) describe the data
- Theoretical uncertainties:
 - terms beyond NLO \Rightarrow 10-20%
 - uncertainties of proton PDFs
 - < 5% (up to 15%) for $x_p < 0.1$ (> 0.1)
- Even up to the highest x_p , where 40% of $d\sigma/dx_p$ arises from gluon $_p$ -induced processes, the data is described by NLO
- Consistent with QCD-evolved photon PDFs determined from measurements at lower scales

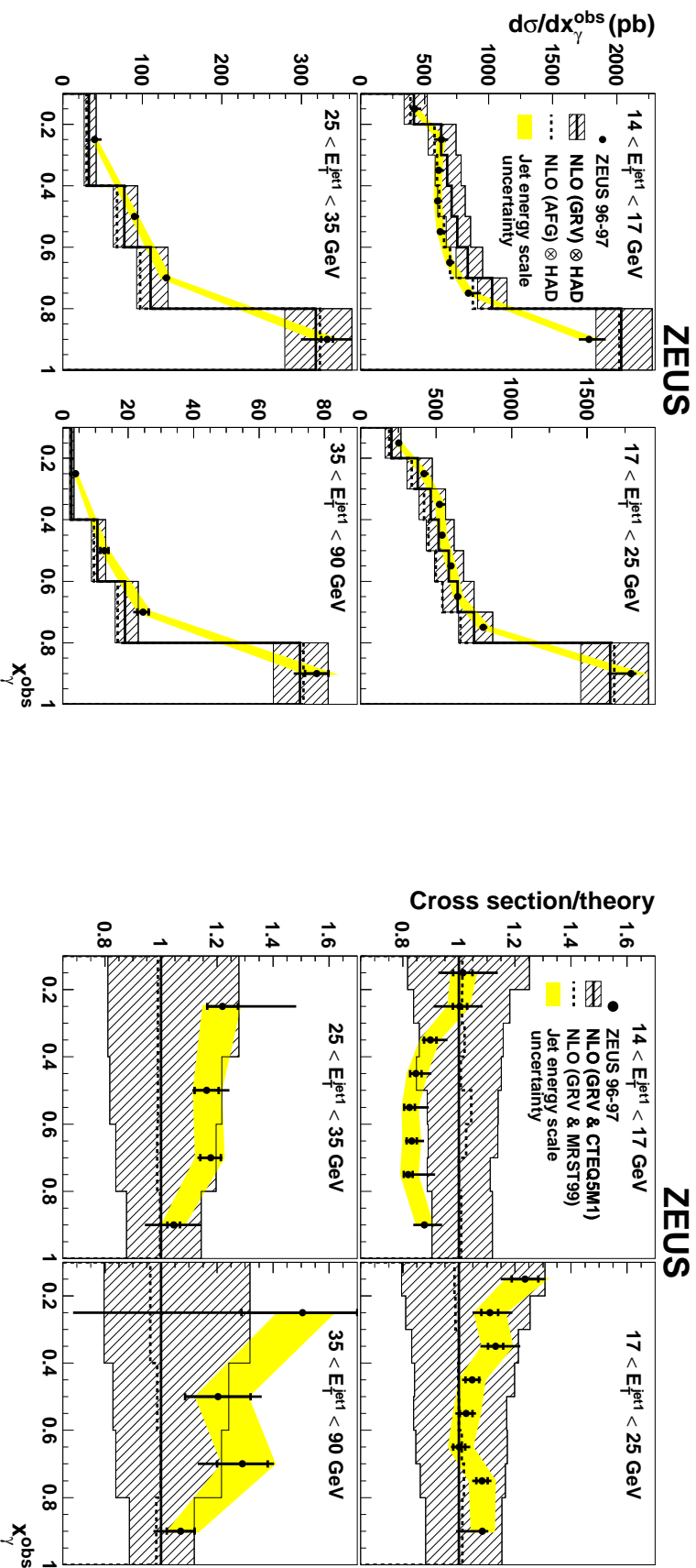


Dijet Photoproduction: photon structure

- **Measurement of the dijet cross section $d\sigma/dx_\gamma$ for two $E_{T,max}$ regions ($\mu_{F\gamma}$ where γ is probed)**
- **Comparison with NLO QCD using various sets of photon PDFs (GRV-HO, AFG-HO):**
 - NLO calculations describe the data well
 - the variation due to the choice of γ PDFs is smaller than the NLO uncertainty due to higher orders and the correlated experimental uncertainties (jet energy scale)
- **To constrain the photon PDFs further (rather than testing existing sets)**
 - reduction of experimental uncertainties and
 - improved calculations are needed



Dijet Photoproduction: photon structure



- Measurement of the dijet differential cross section $d\sigma/dx_\gamma$ for dijet events with $E_T^{jet,1} > 14$ GeV, $E_T^{jet,2} > 11$ GeV, $-1 < \eta^{jet} < 2.4$ (both jets)

- Small experimental uncertainties: jet energy scale known to 1%

- Comparison with NLO QCD calculations (CTEQ5M1 proton, GRV/AFG-HO photon):

the data at low x_γ^{OBS} fall less steeply with E_T^{jet} than NLO; sensitivity to γ -PDFs

\Rightarrow Useful constrain in a global determination of γ -PDFs; improved calculations would help

Dijet Photoproduction: H1 vs ZEUS

- Differences between (ZEUS vs NLO) and (H1 vs NLO)

→ The differences appear to be due to the different cuts on the second jet's E_T^{jet} ($E_{T,2}^{jetcut}$)!

→ 11 GeV in ZEUS; 15 GeV in H1

- $E_{T,2}^{jetcut}$ dependence of the NLO QCD calculations:

→ the measured dependence **not** reproduced by NLO

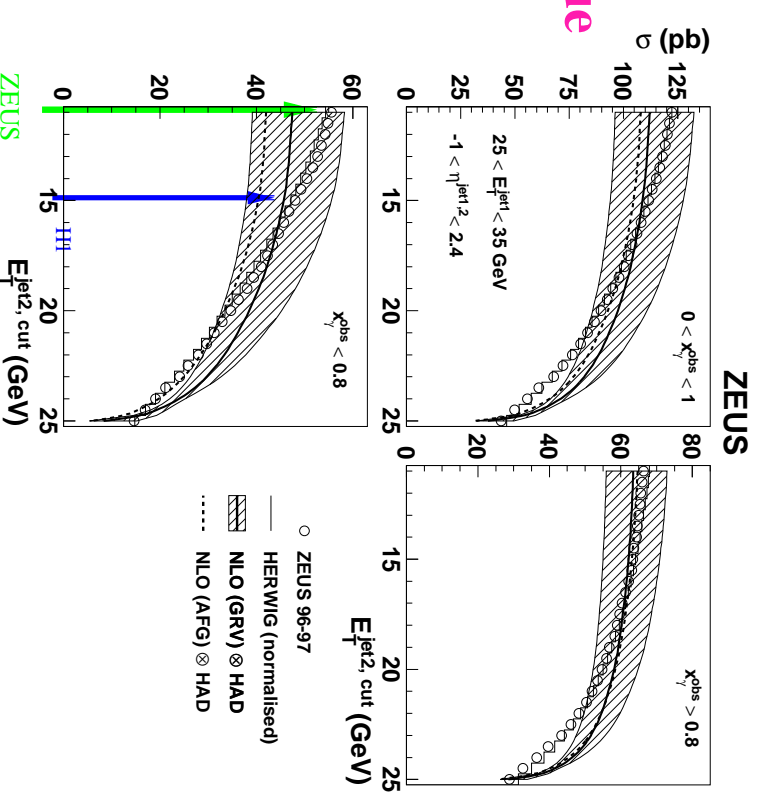
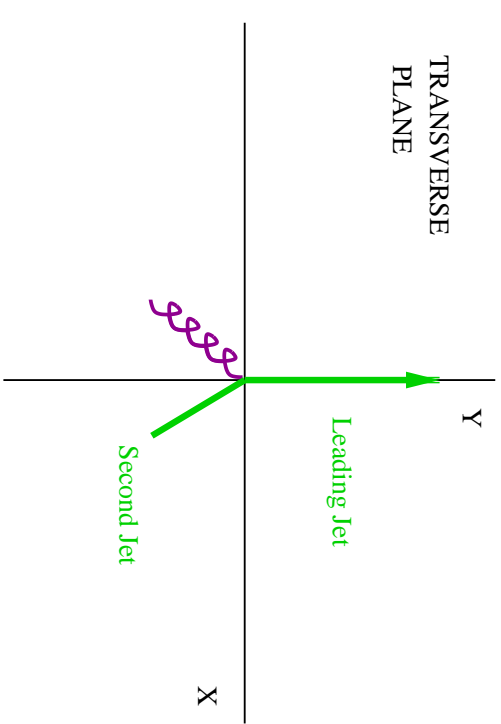
(HERWIG reproduces the measured dependence)

→ the comparison data vs NLO depends on the cut value

→ for $x_\gamma^{OBS} > 0.8$ the NLO converges to the data as the cut on $E_{T,2}^{jetcut}$ is decreased; there the cross is relatively insensitive to the cut value

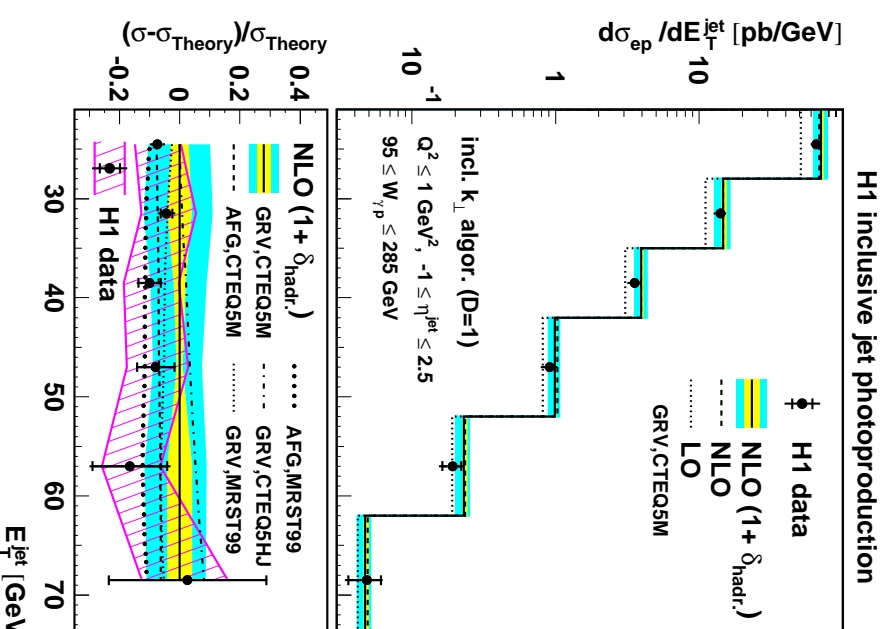
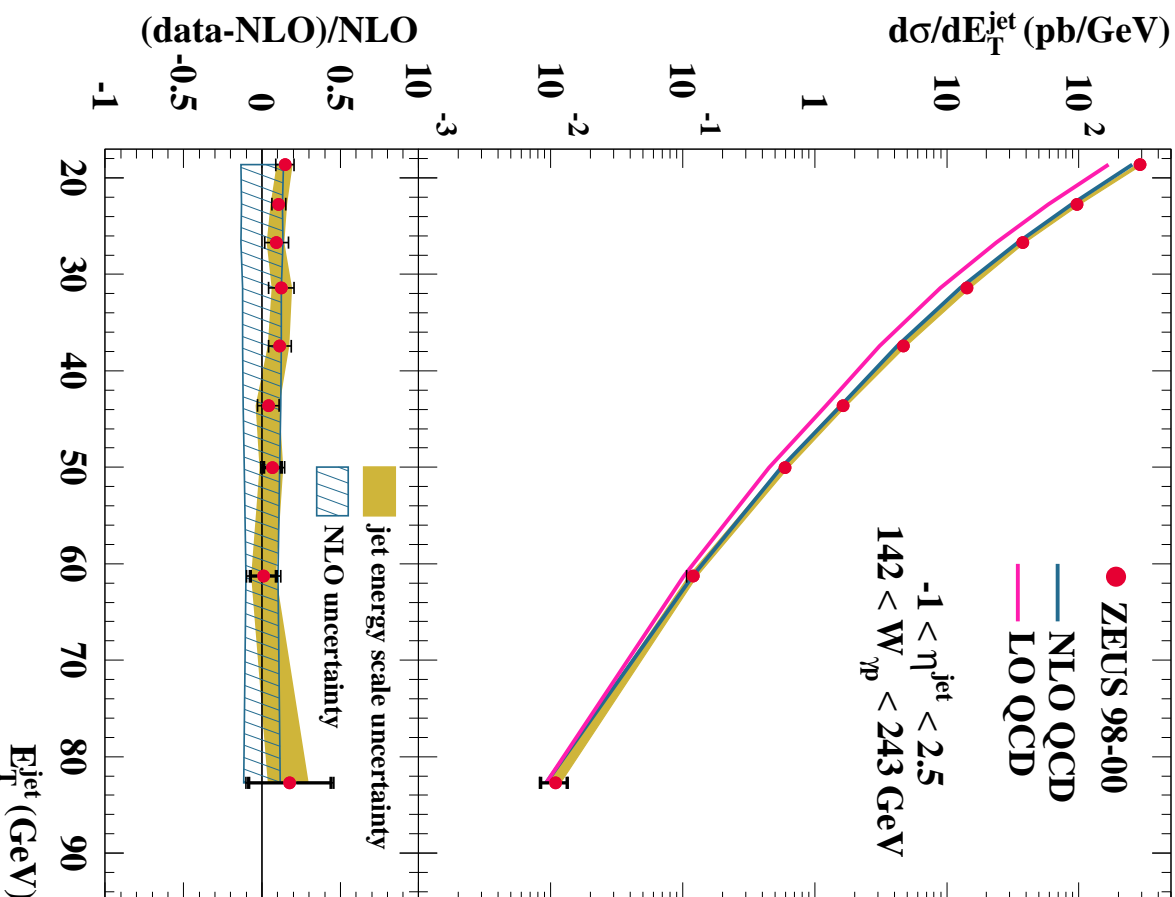
→ accounted for by the large μ_R uncertainty which is comparable to the differences between various γ PDFs

- Theoretical work on improving the dijet calculations (NNLO? selection criteria?) is needed!



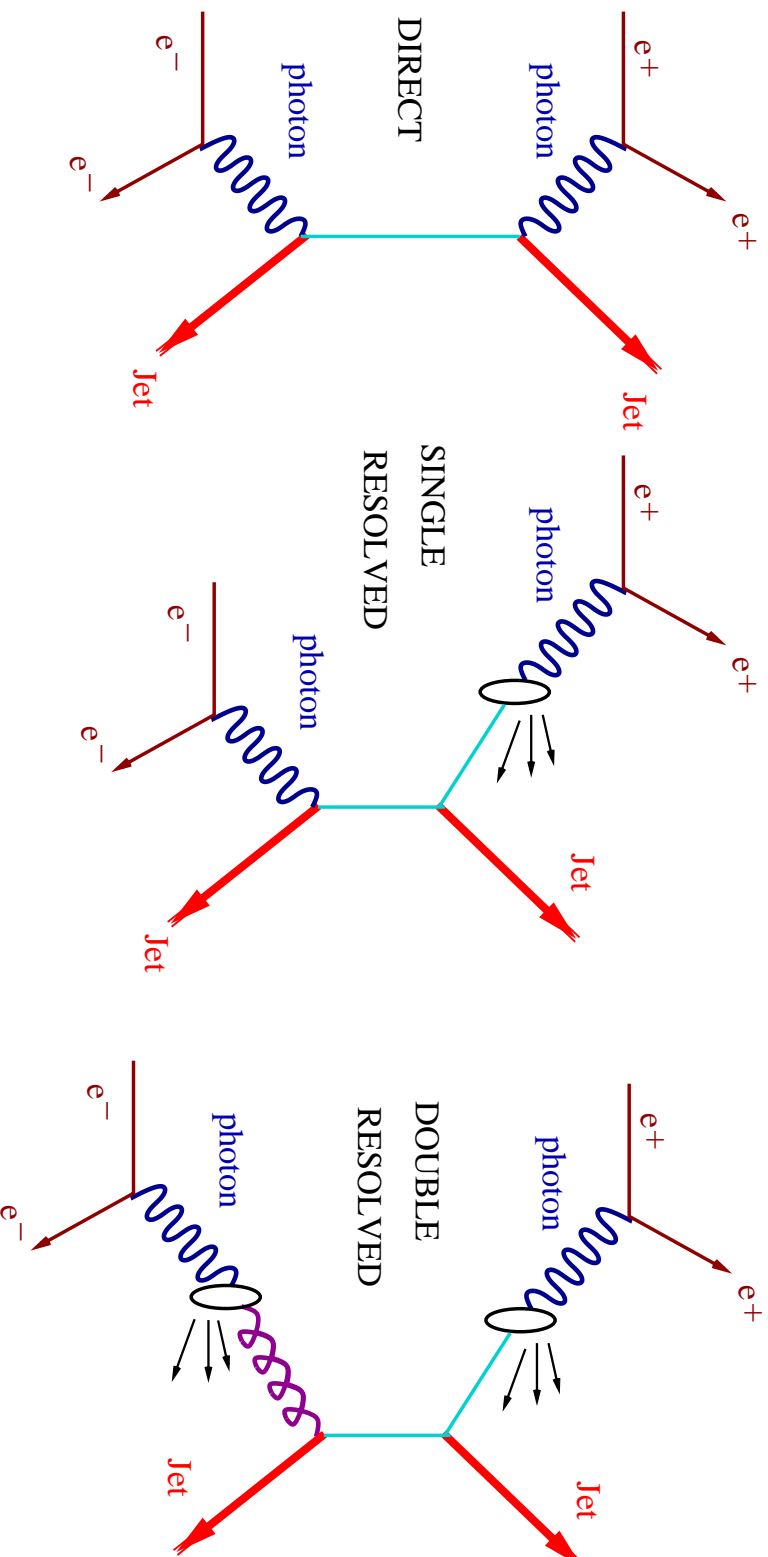
Inclusive Jet Photoproduction

ZEUS



- Measurements (**NEW**) of differential cross sections for **inclusive jet photoproduction**
 - **Smaller theoretical uncertainties**
 - ⇒ Precise test of NLO QCD calculations
 - ⇒ **Determination of α_s and its scale dependence!**
- Talks in WG-C by B. Andrieu and M. Sutton

Jet Production in $\gamma\gamma$ collisions



- Jet production in $\gamma\gamma$ collisions has been studied in e^+e^- interactions
- Three contributing processes: **direct**, **single-resolved** and **double-resolved**
- Measurements of jet production allow the study of
 - dynamics of $\gamma\gamma$, γ -parton and parton-parton interactions
 - photon structure: sensitivity to quark and gluon densities in the photon

Dijet Production in $\gamma\gamma$ collisions

OPAL

- Measurement of dijet production in

$\gamma\gamma$ collisions from e^+e^- data at

$\sqrt{s} = 161$ and 172 GeV; jets

defined using cone algorithm ($R = 1$)

with $E_T^{jet} > 3$ GeV and $|\eta^{jet}| < 2$

- Observables to separate the contributions:

$$x_\gamma^\pm = \frac{\sum_{\text{jets } 1,2} (E^{jet} \pm p_z^{jet})}{\sum_{\text{all hadrons}} (E^h \pm p_z^h)}$$

→ “direct-enhanced”: $x_\gamma^\pm > 0.8$

→ “single-resolved-enhanced”: $x_\gamma^+ > 0.8$

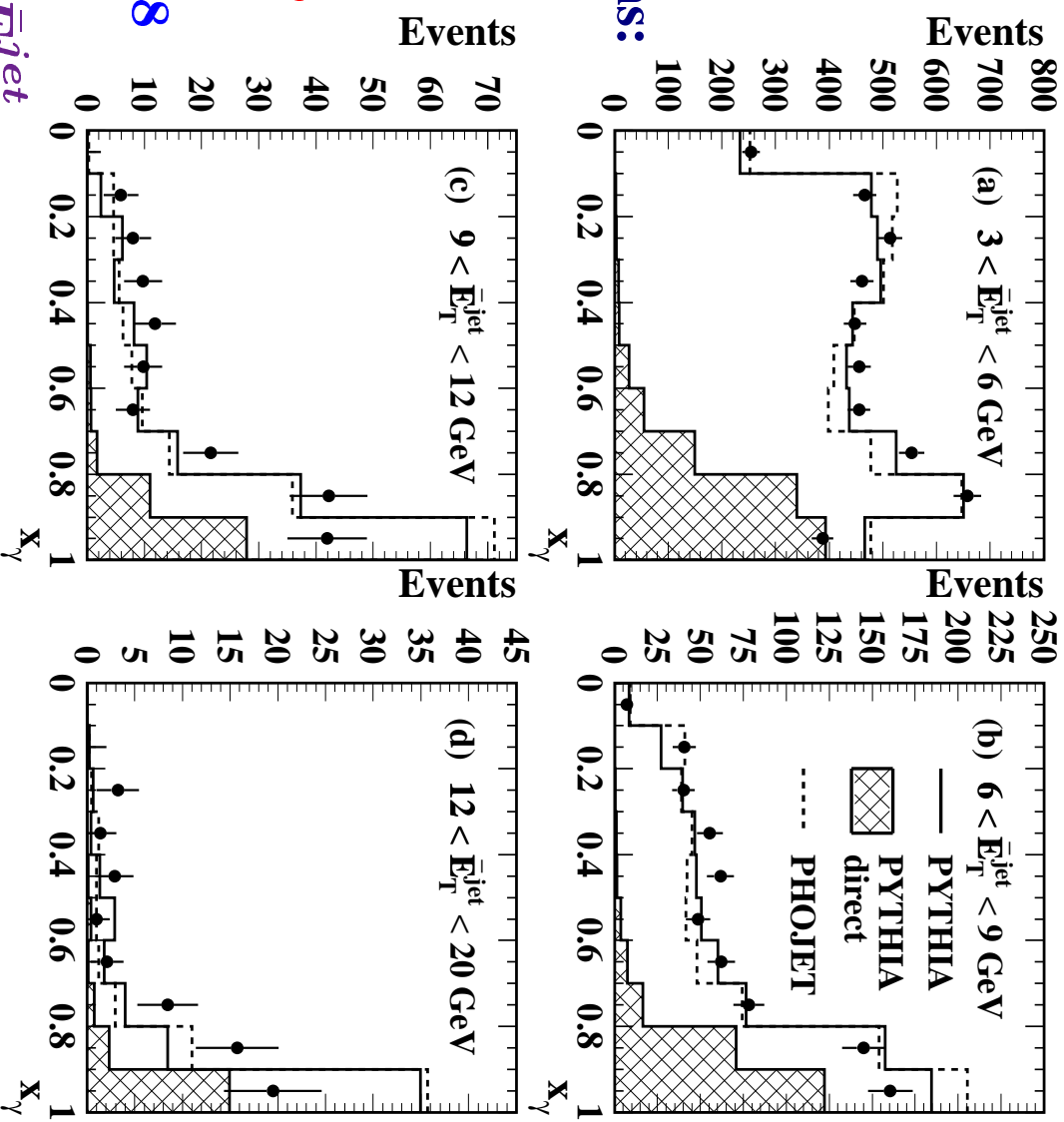
and $x_\gamma^- < 0.8$ (or viceversa)

→ “double-resolved-enhanced”: $x_\gamma^\pm < 0.8$

- Dependence of x_γ distribution on \bar{E}_T^{jet}

→ resolved contribution dominant at low \bar{E}_T^{jet}

→ increasing direct contribution as \bar{E}_T^{jet} increases



Dijet Production in $\gamma\gamma$ collisions: the dynamics

- Measurement of $d\sigma/d|\cos\theta^*|$ in the phase-space region defined by

$$M_{jj} > 12 \text{ GeV}, |\bar{\eta}^{jet}| < 1$$

- “Direct-enhanced sample”:

→ dominated by $\gamma\gamma \rightarrow q\bar{q}$ (q-exchange)

$$d\sigma/d|\cos\theta^*| \sim \frac{1}{(1-|\cos\theta^*|)^1}$$

- “Double-resolved-enhanced sample”:

→ dominated by parton-parton processes

(g-exchange)

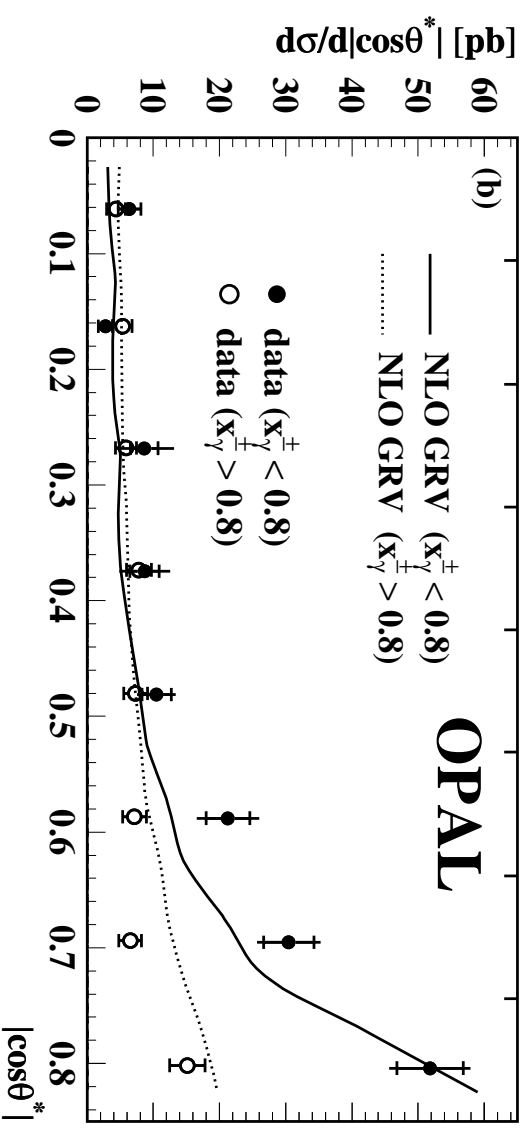
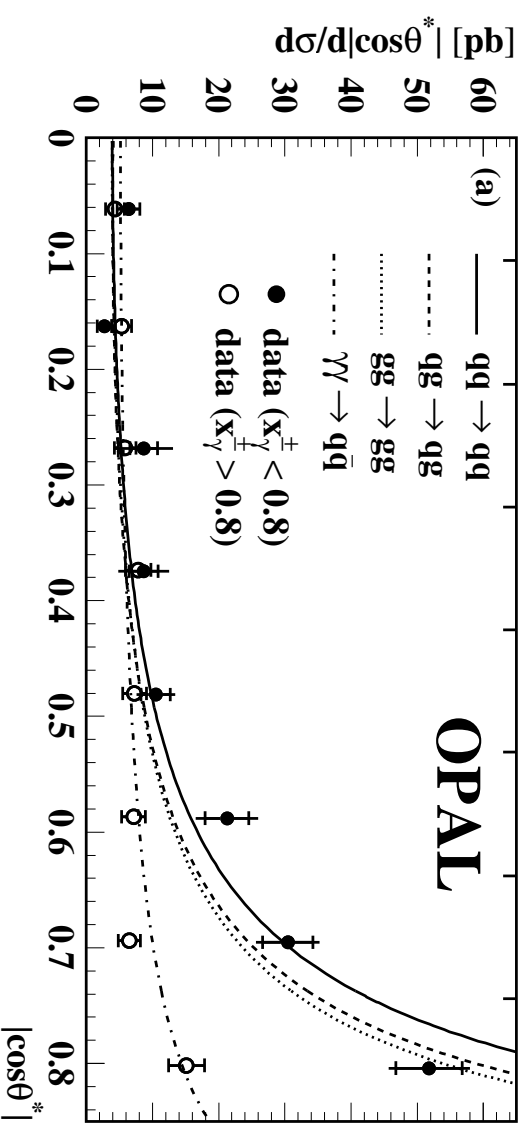
$$d\sigma/d|\cos\theta^*| \sim \frac{1}{(1-|\cos\theta^*|)^2}$$

- Stronger rise as $|\cos\theta^*| \rightarrow 1$ in the

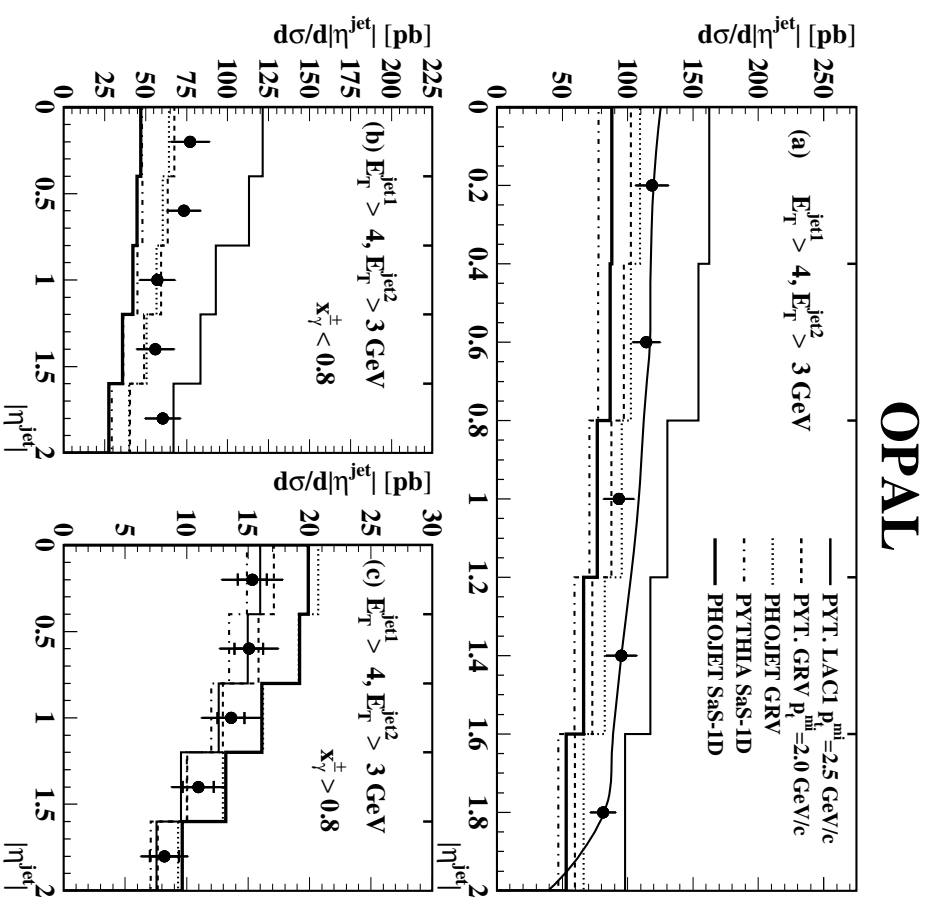
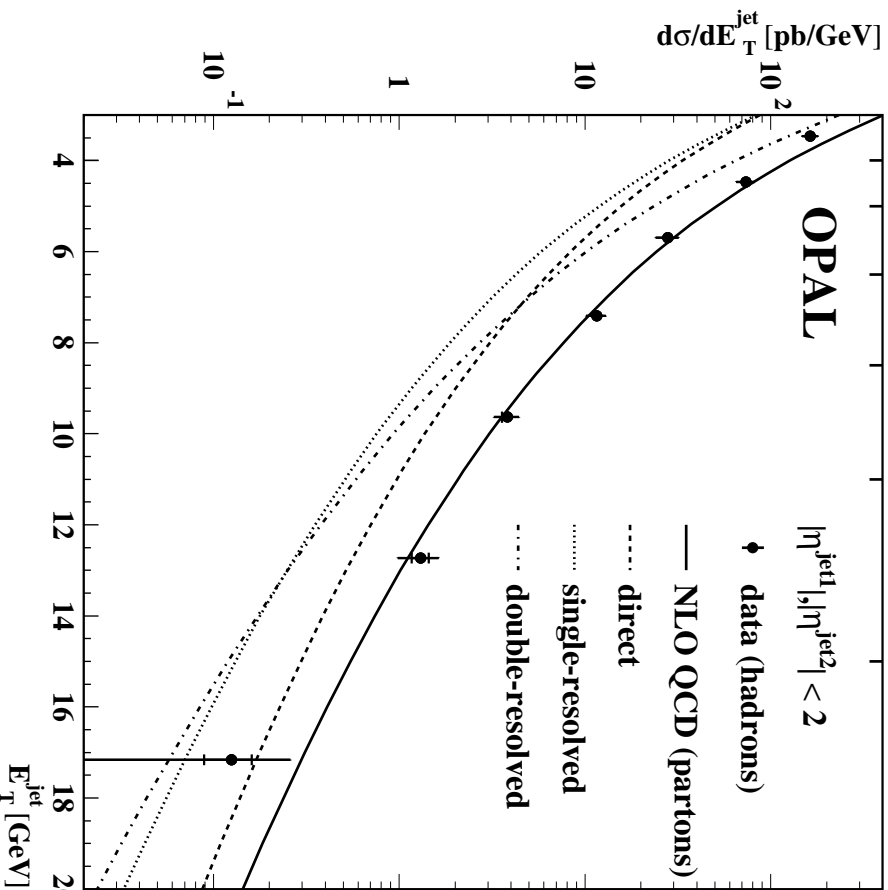
“double-resolved-enhanced sample”

than in the “direct-enhanced sample”

⇒ Shapes reproduced by LO-QCD models (PYTHIA, PHOJET) and NLO QCD

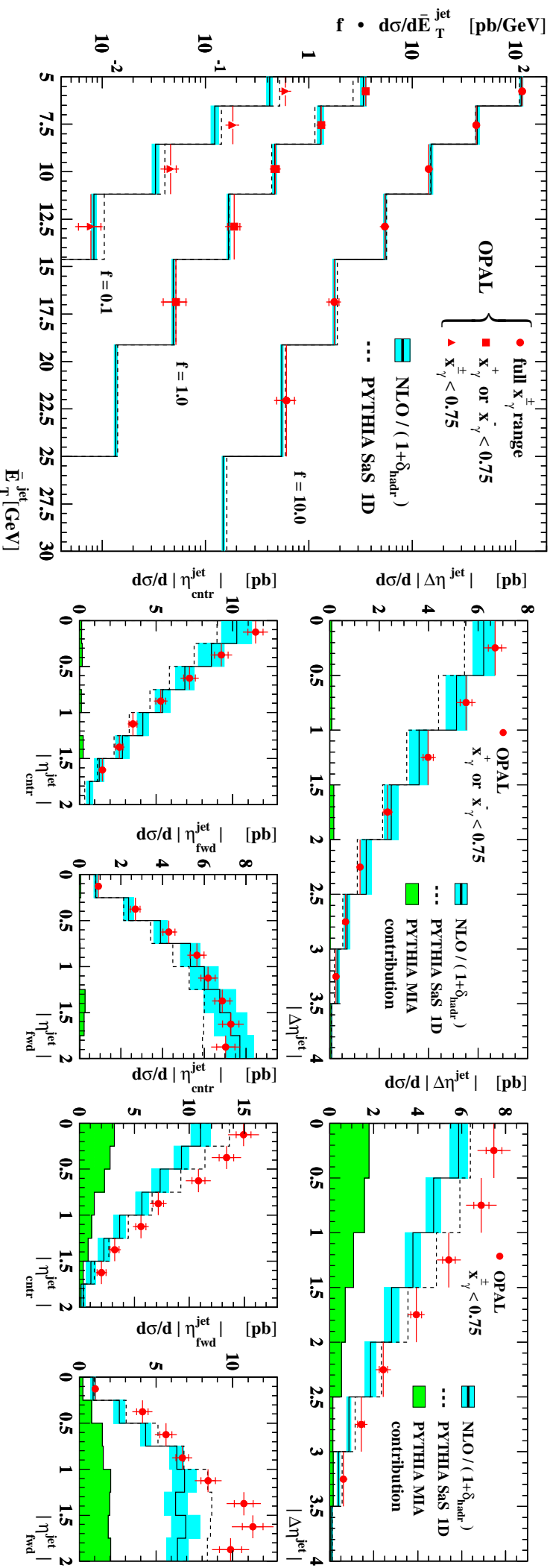


Dijet Production in $\gamma\gamma$ collisions: photon structure



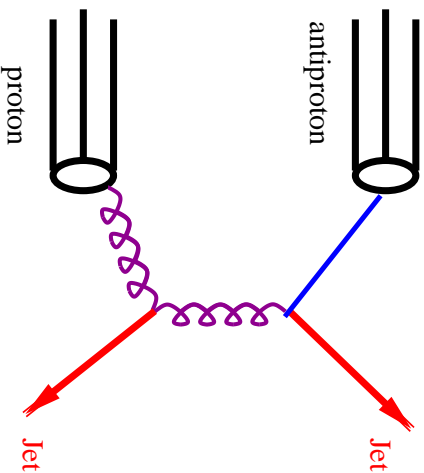
- **NLO QCD (without corr. for hadronis. or “underlying event?”) using GRV γ -PDFs** reproduces the shape and normalisation of the data; **NLO by Klasen et al**
- **Comparison to LO-QCD models (including hadronis. and “underlying event” effects)** → **sensitivity to γ -PDFs**

Dijet Production in $\gamma\gamma$ collisions

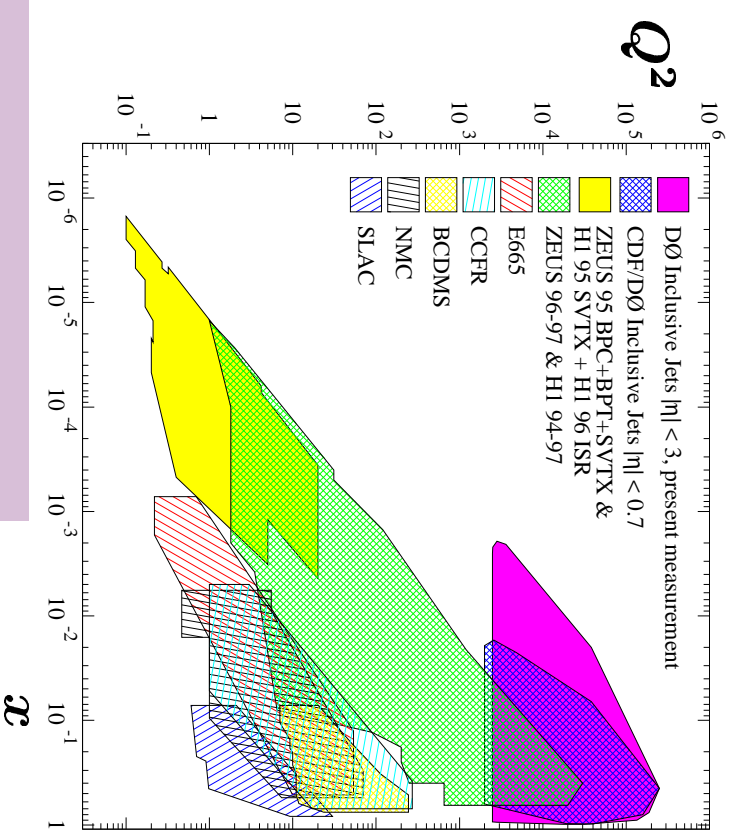


- Measurements (**NEW**) of dijet production in $\gamma\gamma$ from e^+e^- at $\sqrt{s} = 189\text{-}209$ GeV with larger luminosity and using the inclusive k_T jet algorithm
- Measurements of differential cross sections in different regions of the $x_\gamma^+ - x_\gamma^-$ plane:
 - either x_γ^+ or $x_\gamma^- < 0.75$: resolved-photon processes dominate; less sensitive to “underlying event” effects; NLO QCD (GRV-HO γ -PDFs) agree well with the data
- Talks in WG-C by K. Krueger and P. Acharid

Jet Production in $p\bar{p}$ collisions



- pQCD calculations of jet cross sections



$$d\sigma_{jet} = \sum_{a,b} \int_0^1 dx_{\bar{p}} f_{a/\bar{p}}(x_{\bar{p}}, \mu_{F\bar{p}}^2) \int_0^1 dx_p f_{b/p}(x_p, \mu_{Fp}^2) d\hat{\sigma}_{ab \rightarrow jj}$$

longitudinal momentum fraction of parton a /antiproton ($x_{\bar{p}}$), parton b /proton (x_p)

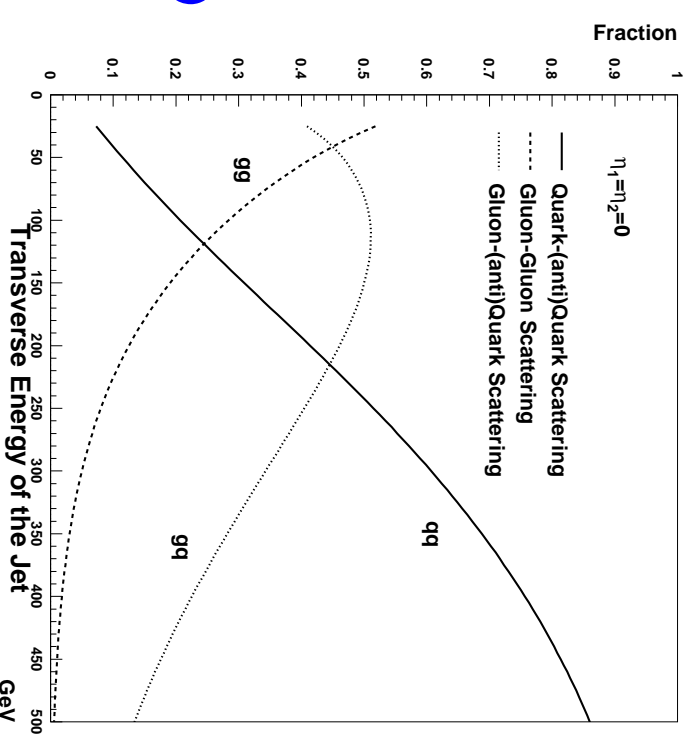
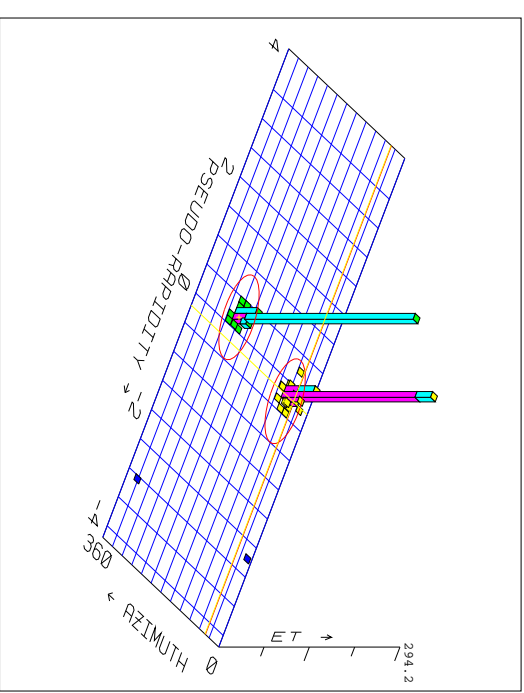
→ $f_{b/p}(x_p, \mu_{Fp}^2)$ = parton densities in the proton

→ $\sigma_{ab \rightarrow jj}$ subprocess cross section; short-distance structure of the interaction

- Dynamics of parton-parton interactions at the smallest distances!
- Proton structure: sensitivity to quark and gluon densities at large x and large Q^2

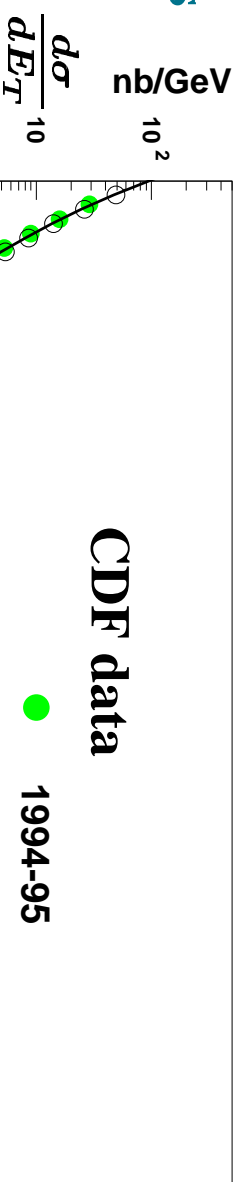
NLO QCD Calculations of Jet Cross Sections

- **Iterative cone algorithm in the η - ϕ plane (laboratory)**
 - **new measurements using inclusive k_T jet algorithm**
- **Several NLO QCD calculations ($\mathcal{O}(\alpha_s^3)$) available**
 - **Ellis et al, Aversa et al, Giele et al**
 - **choice of scales: $\mu_R = \mu_F \sim E_T^{jet}$**
 - **proton PDFs: current NLO QCD parametrisations**
- **Cone algorithm: introduction of R_{sep} parameter to approximate experimental effects on clustering**
- **Effects from the underlying event estimated from minimum bias data** → Talk in WG-C by S. Lami
- **At low E_T^{jet} jet production is dominated by (incoming) gg and gq scattering**
- **At high E_T^{jet} jet production is dominated by (incoming) qq scattering; gq about 30% at $E_T^{jet} = 350$ GeV**



Measurement of Inclusive Jet Production in $p\bar{p}$ collisions

- Measurement of the differential cross section $d\sigma/dE_T^{jet}$ for inclusive jet production using the cone algorithm



($R = 0.7$) in the region

$$0.1 < |\eta^{jet}| < 0.7$$

- Comparison to NLO QCD calculations:

$$\rightarrow \mu_R = \mu_F = E_T^{jet}/2, R_{sep} = 1.3$$

\rightarrow proton PDFs CTEQ4M

- Theoretical uncertainties:

\rightarrow higher-order terms (varying μ_R): 20%

\rightarrow variation of R_{sep} : 5%

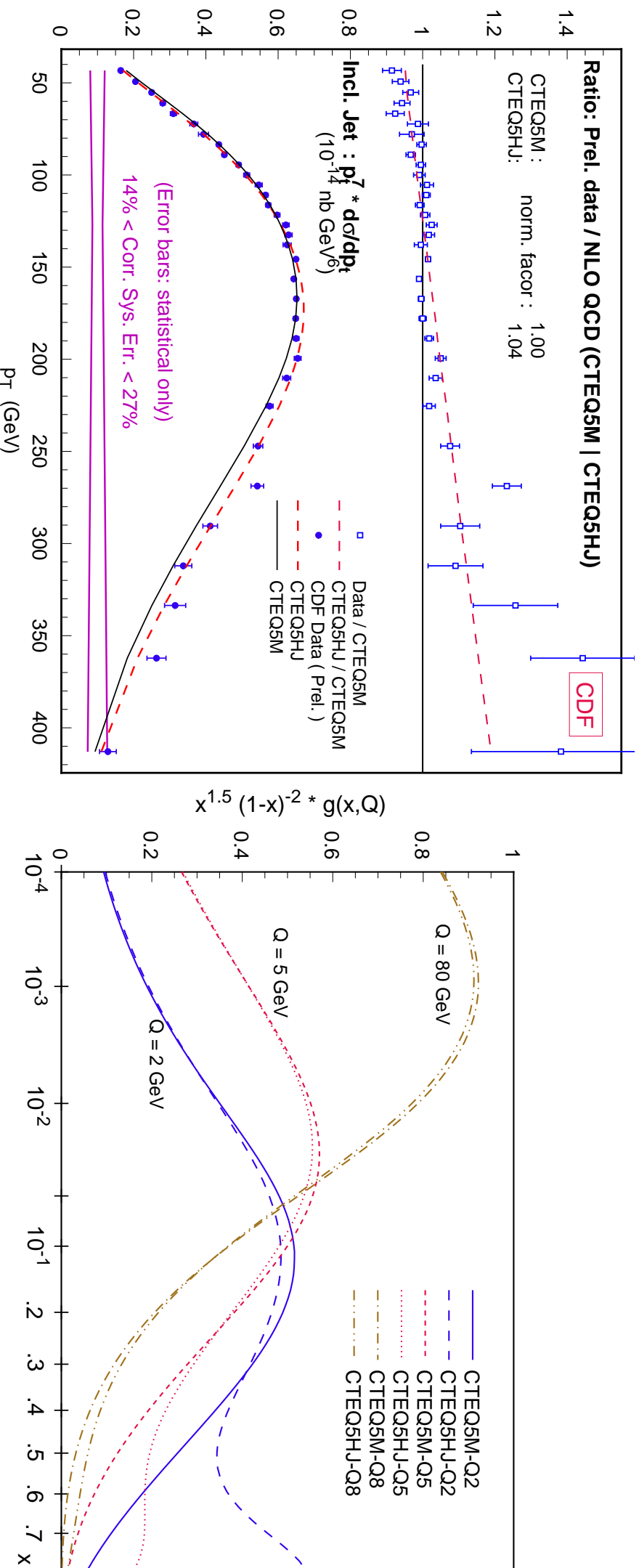
- NLO QCD calculations based on proton

PDFs CTEQ4M describe the shape and

normalisation of the data well for $E_T^{jet} < 250$ GeV

\Rightarrow Excess at high E_T^{jet} ! New physics? Parametrisations of the proton PDFs?

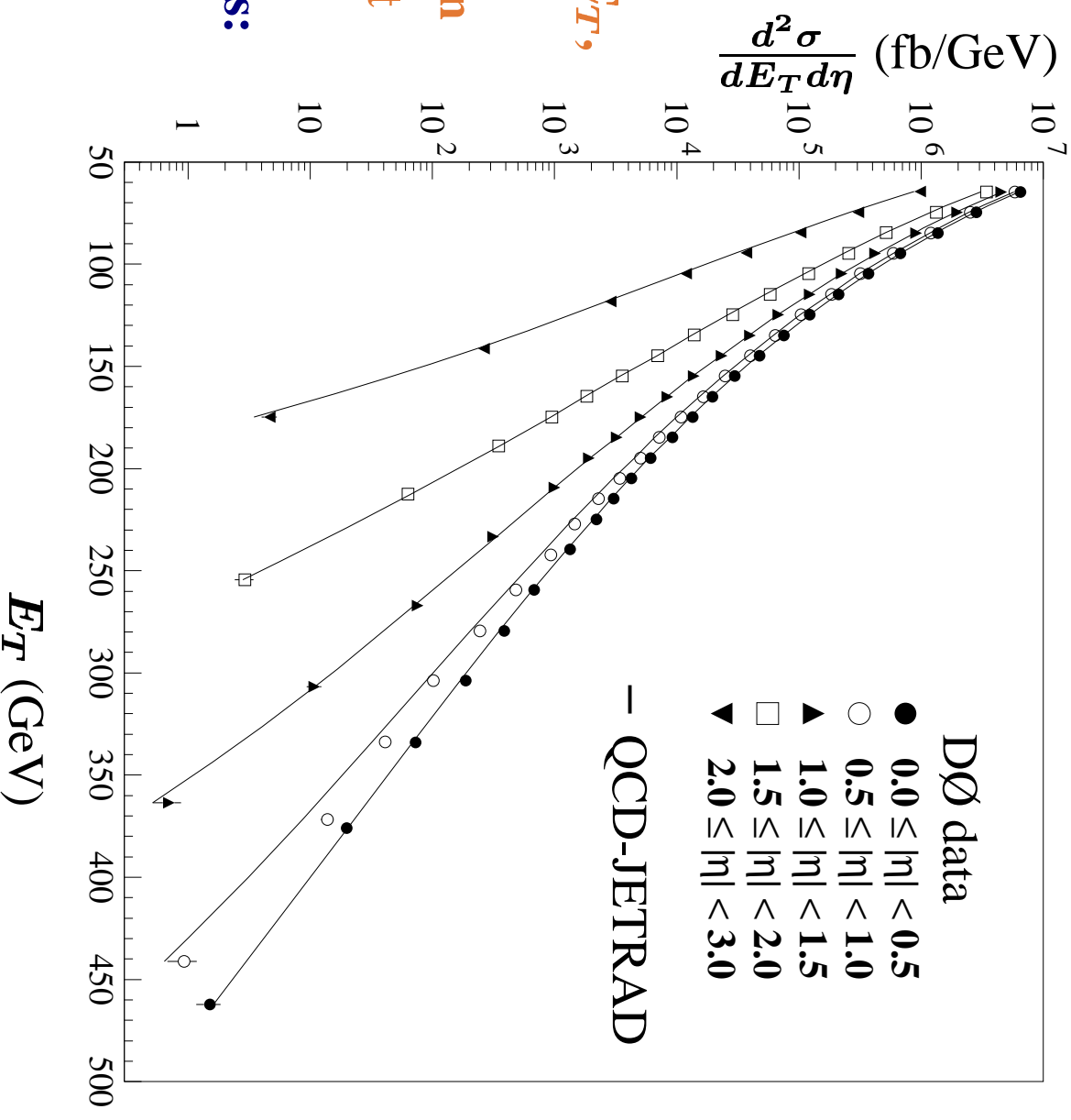
Revisiting the gluon density at high x



- It is possible to describe the data by increasing the gluon density at high x while maintaining a good description of other data sets used in global analyses of proton PDFs
- The CTEQ5HJ parametrisations were obtained putting extra weight to the CDF data; calculations based on the central fit, CTEQ5M, lie below the data at high E_T^{jet}

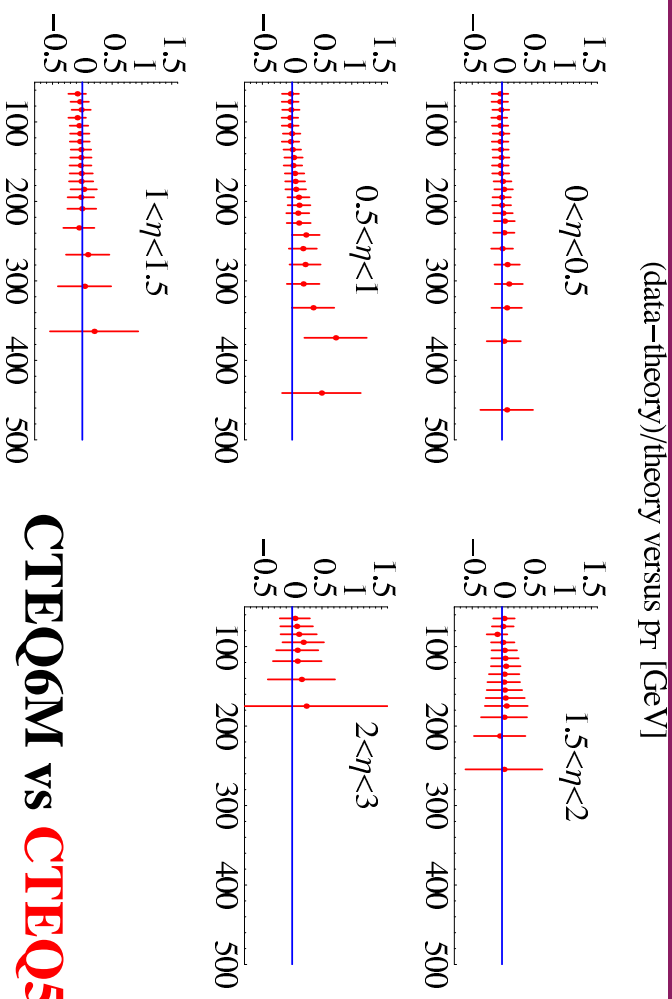
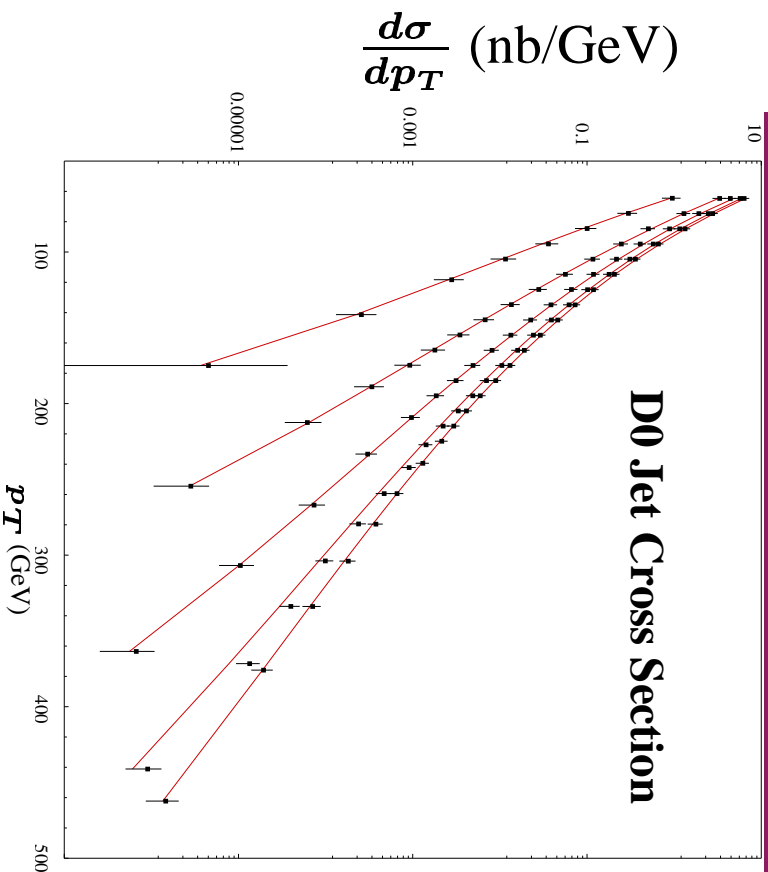
Measurement of Inclusive Jet Production in $p\bar{p}$ collisions

- Measurement of the differential cross section $d^2\sigma/dE_T^{jet}d\eta^{jet}$ for inclusive jet production using the cone algorithm ($R = 0.7$) in the region $|\eta^{jet}| < 3$
- **Extended range in the (x, Q^2) plane!**
- **Experimental uncertainties:**
 - jet energy scale: 12 – 20% at low E_T , 35 – 80% at high E_T
 - jet-energy-resolution parametrisation and unfolding procedure: 3 – 5% at low E_T , 10 – 20% at high E_T
- **Comparison to NLO QCD calculations:**
 - ⇒ preference for proton PDFs with larger gluon content at high x



- Talk in WG-C by N. Skatchkov

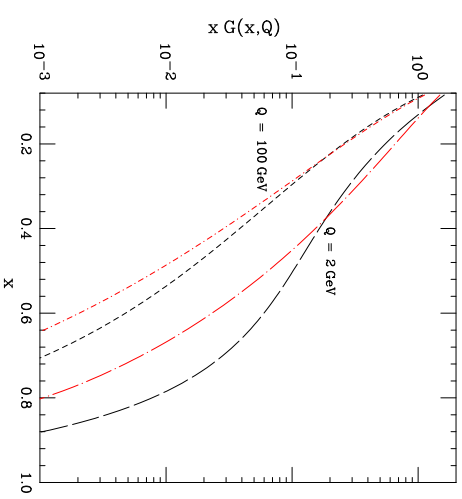
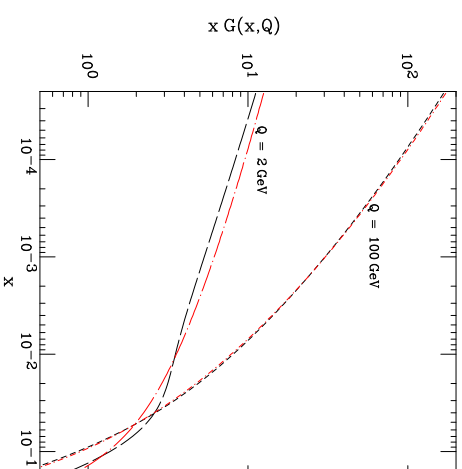
Inclusive Jet Production in $p\bar{p}$ collisions: the gluon density at high x



CTEQ6M vs CTEQ5M1

- DØ and CDF data used in the global determination of proton PDFs (CTEQ6M) without “extra” weight \Rightarrow constraining the gluon density at high x
- \Rightarrow **good description of the data**

- Run II measurements will help to confirm or rule out the “high- x gluon solution” \rightarrow Talk in WG-C by V. O’Dell



Last Remarks

- **In many areas further progress is limited by theoretical uncertainties!**
- Measurements of jet production in $p\bar{p}$ included in global determination of the proton PDFs
 - **Precise measurements of jet production in ep , γp and $\gamma\gamma$ demand a similar effort!**
- Apologies for many omissions due to lack of time (and expertise!)