

Anti-de Sitter space (AdS_d)

like de Sitter space hyperboloid,
defined by embedding into $(d+1)$ -dim

flat space:

$$-X_0^2 + \sum_{i=1}^{d-1} X_i^2 - X_d^2 = -\ell^2$$

negative
curvature

$$d\Delta^2 = -dX_0^2 - dX_d^2 + \sum_{i=1}^{d-1} dX_i^2$$

by construction: isometry $SO(2, d-1)$,
 $\left(G_{\mu\nu} = \begin{pmatrix} -1 & & & \\ & -1 & & \\ & & \dots & \\ & & & 1 \end{pmatrix} \right)$ maximal symmetry,
homogeneous + isotropic

coordinates:

$$X_0 = \ell \cdot \cosh\left(\frac{\rho}{\ell}\right) \cdot \cos\left(\frac{\tau}{\ell}\right)$$

$$X_i = \ell \cdot \sinh\left(\frac{\rho}{\ell}\right) \cdot \omega_i, \quad \sum_{i=1}^{d-1} \omega_i^2 = 1$$

$$X_d = \ell \cdot \cosh\left(\frac{\rho}{\ell}\right) \cdot \sin\left(\frac{\tau}{\ell}\right)$$

$$\rho \geq 0, \quad 0 \leq \frac{\tau}{\ell} < 2\pi$$

metric on AdS_d :

$$d\Delta^2 = -\cosh^2\left(\frac{\rho}{\ell}\right) d\tau^2 + d\rho^2 + \ell^2 \sinh^2\left(\frac{\rho}{\ell}\right) d\Omega_{d-2}^2$$

coordinates periodic in τ ; closed time-like
curves?

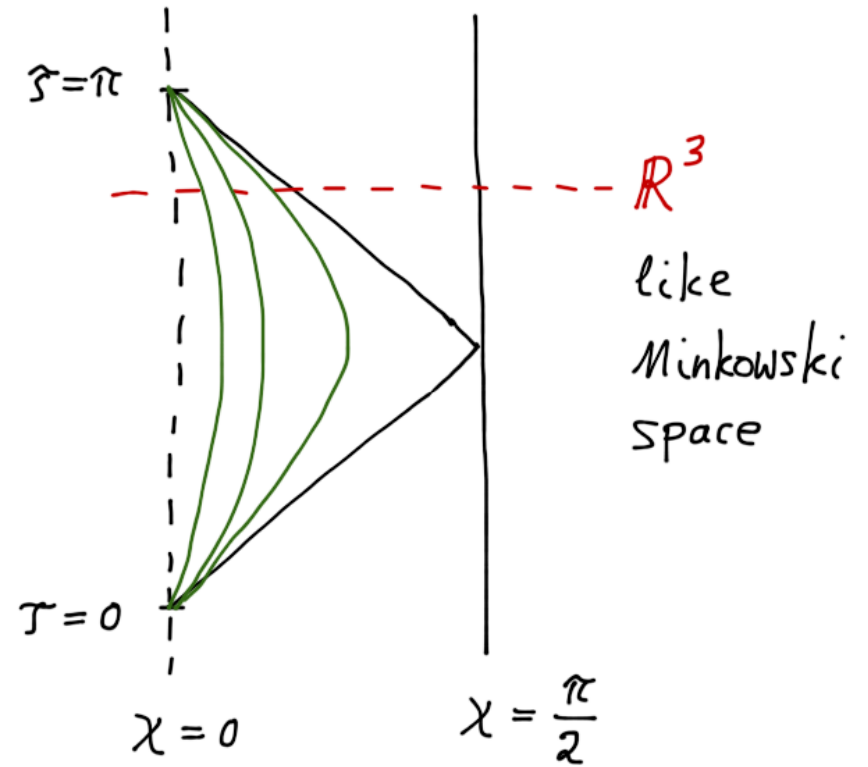
→ causal space-time by changing
 $S^1 \rightarrow \mathbb{R}$ (covering space)
 \leadsto same metric!

conformal coordinates (cf. de Sitter):

$$\cosh\left(\frac{\rho}{l}\right) = \frac{1}{\cos \chi}, \quad 0 \leq \chi \leq \frac{\pi}{2}$$

$$\leadsto d\rho^2 = \frac{1}{\cos^2 \chi} \left[-d\tau^2 + l^2 (d\chi^2 + \sin^2 \chi \cdot d\Omega_{d-2}^2) \right]$$

conformal diagram:



II.4 Horizons & Holography

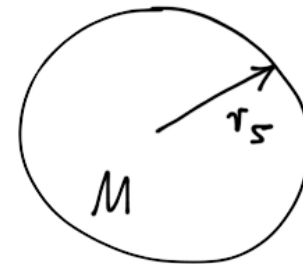
some literature:

- Susskind & Lindesay '05:
" [...] The holographic universe "
- Susskind & Witten '98:
"The holographic bound in AdS space"
hep-th/9805114
- Kaloper & Linde '99:
"Cosmology vs. Holography"
hep-th/9904120

black hole entropy

for outside observer BH of mass M has temperature:

$$T = \frac{1}{4\pi \cdot r_S}, \quad r_S = 2GM$$



1st law of thermodynamics implies entropy:

$$dE = T \cdot dS \quad \left(\begin{array}{l} \text{energy of infalling matter,} \\ \underline{\text{no work done on BH}} \end{array} \right)$$

$$\leadsto dE = dM = \frac{1}{8\pi G M} dS, \quad S|_{M=0} = 0$$

$$\Leftrightarrow S_{\text{BH}} = 4\pi G M^2 = \frac{4\pi \cdot r_s^2}{4G} = \frac{A_{\text{horizon}}}{4G}$$

← horizon area

↑
Bekenstein-Hawking entropy of BH

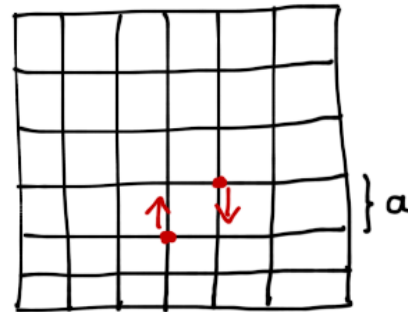
↓
Bekenstein-Hawking entropy
= $\frac{1}{4}$ area of event horizon in Planck units
Planck length: $l_p = \sqrt{\frac{G\hbar}{c^3}}$, $G = \frac{l_p^2 c^3}{\hbar}$

$$\leadsto S_{\text{BH}} = \frac{\hbar}{4c^3} \cdot \frac{A_{\text{horizon}}}{l_p^2}$$

$l_p \approx 10^{-33}$ cm ($k=1$ Boltzmann const.)
quantum effect ($\mathcal{O}(\hbar)$)!!

Maximum entropy

system of spins in volume \mathcal{V} ;



\mathcal{V}
a: lattice spacing

number of states:

$$N_{\text{tot}} = 2 \frac{V}{a^3}$$

entropy of macroscopic state with statistical weight ΔP

$$\left(\text{e.g. } \frac{\Delta P \Delta q}{2\pi \hbar} \right)$$

(# of corresponding quantum states):

$$S = \ln \Delta P$$

no knowledge about system:

$$\Delta P = N_{\text{tot}}$$

maximum entropy of spin system:

$$S = \frac{V}{a^3} \cdot \ln 2$$

proportional to volume, extensive quantity.

effect of gravity

