

# Charged multiplicity distributions in deep inelastic scattering at HERA

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## Abstract.

The hadronic final state has been investigated in inclusive neutral current deep inelastic  $ep$  scattering with the ZEUS detector at HERA, using an integrated luminosity of  $38.6 \text{ pb}^{-1}$ . The mean charged multiplicity has been measured for the hadrons belonging to the current region of the Breit frame, as well as for those belonging to the photon (current) fragmentation region of the hadronic centre-of-mass frame (HCM). The results are compared to leading-logarithm parton-shower Monte Carlo predictions as well as to results of  $e^+e^-$  and  $pp$  measurements.

**Keywords:** multiplicity distributions, hadronic final states

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## INTRODUCTION

The measurements of multiplicities of charged particles at colliders have yielded insights into hadronisation mechanisms. It has been found that charged particle multiplicities measured as a function of the centre-of-mass (cms) energy,  $\sqrt{s}$ , at  $e^+e^-$  colliders are the same as those measured at  $pp$  colliders as a function of  $\sqrt{q_{\text{tot}}^{\text{had}}} = \sqrt{[(q_1^{\text{inc}} - q_1^{\text{leading}}) + (q_2^{\text{inc}} - q_2^{\text{leading}})]^2}$ , where  $q_{1,2}^{\text{inc}}$  and  $q_{1,2}^{\text{leading}}$  are the four-momenta of the incoming protons and leading particles that escape down the beampipe, respectively [1]. It is therefore interesting to study charged particle multiplicities in  $ep$  collisions. At the HERA  $ep$  collider, because of the large asymmetry between the  $e$ -beam and  $p$ -beam energies, a large part of the final hadronic system produced near the proton falls outside the region of acceptance of the detectors. Therefore, to perform studies similar to  $e^+e^-$  and  $pp$  also in  $ep$  collisions at HERA, only hadrons belonging to the photon (current) fragmentation region of the hadronic centre-of-mass (HCM) frame and to the current region of the Breit frame were used.

## DATA SELECTION

The data were collected with the ZEUS detector during the 1996 and 1997 running periods, using protons of energy  $E_p = 820 \text{ GeV}$  and positrons of energy  $E_e = 27.5 \text{ GeV}$ , and correspond to an integrated luminosity of  $38.6 \pm 0.6 \text{ pb}^{-1}$ .

The ZEUS detector is described in detail elsewhere [2]. The most important components used in the current analysis are the central tracking detector (CTD) which was

used to measure the charged-particle multiplicity and the uranium-scintillator calorimeter (CAL) which measured the energy of the final-state hadrons, charged and neutral.

Neutral current deep inelastic scattering (DIS) events were selected in the range  $Q^2 > 25 \text{ GeV}^2$  and  $70 < W < 225 \text{ GeV}$ , where  $Q^2$  is the photon virtuality and  $W$  is the  $\gamma^*P$  centre-of-mass energy. Diffractive contributions were reduced by rejecting events where the pseudorapidity of the most forward energy deposit in the CAL is large. The reconstructed tracks were required to be associated with the primary vertex and have a sufficient  $p_T$  to ensure good reconstruction. A cut on the pseudorapidity of the measured tracks restricts the analysis to a region of high CTD acceptance where detector response and systematics are well understood.

In the Breit frame, the exchanged virtual boson is purely spacelike, with three-momentum  $\mathbf{q} = (0,0,-Q)$ . The particles produced in the interaction can be assigned to one of two regions: the current region if their longitudinal momentum in the Breit frame is negative, and the target region if their longitudinal momentum is positive. The hadronic system of the current region used in this analysis is almost fully (about 95%) contained within the acceptance of the CTD.

For the HCM frame the final particles are separated into the photon (current) and proton (target) fragmentation regions. The target region of the HCM can not be observed in the CTD and only 60-80% of the hadrons belonging to the photon (current) region are detected.

The boost to the corresponding frames was performed using the positron four-momentum as reconstructed using the scattered-positron angle,  $\theta_e$ , and the angle  $\gamma_H$ , where in the naive quark parton model,  $\gamma_H$  is the angle of the scattered massless quark in the laboratory frame.

## MONTE CARLO MODELS, ACCEPTANCE CORRECTIONS AND SYSTEMATIC ERRORS

Samples of neutral current DIS events were generated using the HERACLES 4.6.1 [3, 4] MC program with the DJANGO 1.1 [5, 6] interface to the hadronisation programs. The QCD cascade is simulated using the colour-dipole model as implemented in ARIADNE 4.08 [7] or with the MEPS model of LEPTO 6.5 [8]. Both ARIADNE and LEPTO use the Lund string model [9] for the hadronisation. All event samples were generated using the CTEQ4D parameterisation of the parton distribution functions in the proton. The generated hadron distributions do not include charged particles produced from weak decays with lifetimes below  $3 \cdot 10^{-10}$  seconds. The charged-particle decay products of  $K_S^0$  and  $\Lambda$  were excluded.

The MC event samples were passed through reconstruction and selection procedures identical to those for the data. Monte Carlo studies were used to determine the event and track acceptances in the selected kinematic region of  $Q^2$  and  $W^2$ .

The corrections applied to the data accounted for the effects of acceptance and resolution of the detector, event selection cuts, QED-radiative effects, track reconstruction, track selection cuts, the decay products of  $K_S^0$  and  $\Lambda$  which were assigned to the primary vertex, and energy losses in the inactive material in front of the calorimeter in the case of the energy measurement.

Two different correction procedures for the multiplicity distributions were used. One of them is based on the matrix unfolding method as used in earlier studies [10]. The second is a bin-by-bin method. Both methods give similar results. The difference between them is included as a systematic uncertainty.

The dominant sources of systematic uncertainties were investigated. The main uncertainty for the  $W$  measurement arises from different MC models used in the correction procedure and it amounts to up to 5%. To reduce the uncertainty the average correction factor between ARIADNE and LEPTO was taken, thus reducing the uncertainty by a factor of 2. For the Breit frame measurement this uncertainty is typically below 0.5%. Other sources of uncertainties are (typical values of the uncertainties are shown in brackets): uncertainty in the CAL energy scale (0.5 – 1.1%), event reconstruction and selection (< 0.5%), track reconstruction and selection (< 0.5%) and method of correction (matrix or modified bin-by-bin, 0.6 – 1.3%). The contaminations due to migrations from  $Q^2 < 25 \text{ GeV}^2$  and from diffractive events are negligible.

## RESULTS

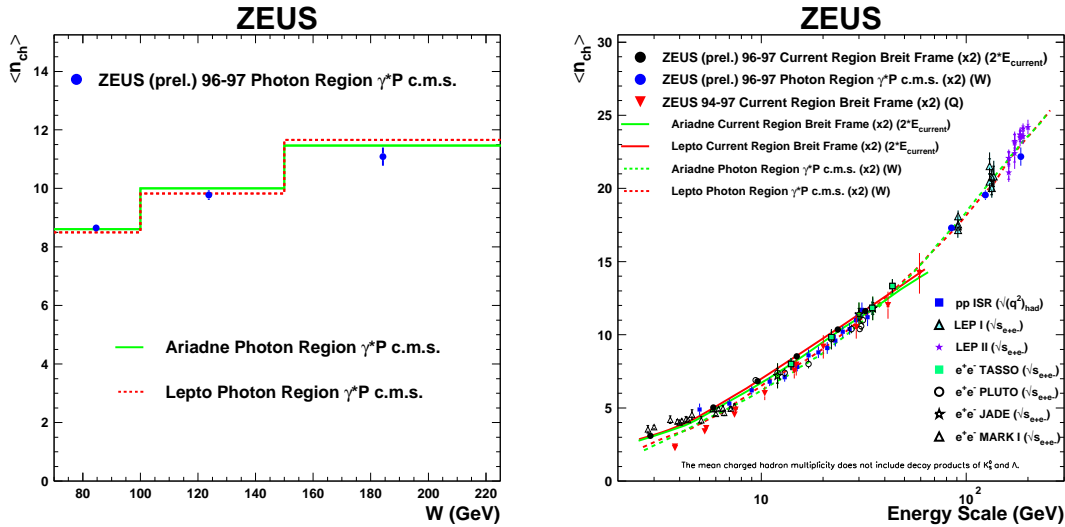
The current region of the Breit frame is analogous to a single hemisphere of  $e^+e^-$  annihilation. In  $e^+e^- \rightarrow q\bar{q}$  the two quarks are produced with equal and opposite momenta,  $\pm\sqrt{s_{ee}}/2$ . The fragmentation of these quarks can be compared to that of the quark struck from the proton in DIS. This quark has an outgoing momentum  $-Q/2$  in the Breit frame. The multiplicity of the current region of the Breit frame is expected to have a dependence on  $Q$  similar to that in  $e^+e^-$  annihilation versus the energy  $\sqrt{s_{ee}} = Q$ . To take into account contributions from soft and hard QCD processes that lead to a decrease of the energy and the number of particles in the current region of the Breit frame, in this analysis  $2E_{\text{current}}$  was used instead of  $Q$ , where  $E_{\text{current}}$  is the energy of all the particles in the current region of the Breit frame.

For the photon (current) fragmentation region of the HCM frame the total number of charged particles can be studied as a function of  $W$ . Additional studies of the energy in this region demonstrated that there are no significant migrations from the photon (current) to the target region and the variable  $W$  can be safely used.

Figure 1 shows the measured mean charged multiplicity,  $\langle n_{\text{ch}} \rangle$ , in the photon (current) region of the HCM frame as a function of  $W$ . The measurement as a function of  $W$  can be only performed at high  $W$ , because of the acceptance of the ZEUS detector. The predictions of ARIADNE and LEPTO are also shown.

To compare the results of measurements in the Breit and HCM frame with results of the  $e^+e^-$  and  $pp$  experiments the mean charge multiplicity was multiplied by 2. Figure 2 shows twice the measured mean charged multiplicity,  $2 \cdot \langle n_{\text{ch}} \rangle$ , in the current region of the Breit frame plotted versus  $2 \cdot E_{\text{current}}$  and twice the measured mean charged multiplicity in the photon (current) region of the HCM frame plotted versus  $W$ . Also shown are the predictions of ARIADNE, LEPTO and the measurements from  $e^+e^-$  and  $pp$

experiments together with a previous ZEUS measurement in the current region of the Breit frame [11]. For the previous ZEUS measurement,  $2 \cdot \langle n_{\text{ch}} \rangle$  is measured as a function of  $Q$ .



**FIGURE 1.** Mean charged multiplicity,  $\langle n_{ch} \rangle$ , in the photon (current) fragmentation region of the HCM as a function of  $W$ . Also shown are the predictions from ARIADNE and LEPTO.

**FIGURE 2.** Mean charged multiplicity,  $\langle n_{ch} \rangle$ , in the current region of the Breit frame multiplied by 2 plotted versus  $2E_{current}$ , where  $E_{current}$  is the sum of the energies of the particles (charged and neutral) in the current region. The blue dots represent the  $\langle n_{ch} \rangle$  in the photon (current) region of the HCM frame, multiplied by 2, plotted versus  $W$ . The downward triangles show the ZEUS measurement of ref. [11]. The prediction of the ARIADNE and the LEPTO, as well as the measurements of  $pp$  and  $e^+e^-$  interactions are also shown.

The ZEUS measurements presented in this paper agree with the  $e^+e^-$  and  $pp$  measurements. At low values of energy, the measurement as a function of  $2 \cdot E_{current}$  agrees better with  $e^+e^-$  than the measurement as a function of  $Q$ , since the migrations of final state particles out of the current region are larger at the low values of energy and they are properly taken into account in  $E_{current}$ .

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