

Excercise Sheet 5 to General Relativity

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Discussion on 23.11.2012 in the exercise classes

1. Geodesic Equation and Extremal Action

Show that the geodesic equation of motion in a gravitational field,

$$\frac{D}{d\tau} \frac{dx^\mu}{d\tau} = \frac{d^2 x^\mu}{d\tau^2} + \Gamma_{\rho\sigma}^\mu \frac{dx^\rho}{d\tau} \frac{dx^\sigma}{d\tau} = 0,$$

can also be derived from extremizing the action

$$S_{\text{particle}} = \frac{1}{c_0} \int ds = \int d\tau = \int d\lambda \mathcal{L}_{\text{particle}} = \int d\lambda \left(\frac{d\tau}{d\lambda} \right) = \int d\lambda \left(g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda} \right)^{1/2}. \quad (1)$$

under small variations $x^\mu(\lambda) \rightarrow x^\mu(\lambda) + \delta x^\mu(\lambda)$ of the path connecting two space-time points such that the variations δx^μ vanish at the end points. Hint: Transform to proper time as integration variable, using

$$d\lambda = \left(g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda} \right)^{-1/2} d\tau,$$

such that

$$g_{\mu\nu} \frac{dx^\mu}{d\tau} \frac{dx^\nu}{d\tau} = g_{\mu\nu} u^\mu u^\nu = 1,$$

and thus the variation of Eq. (1) has the form

$$\delta S_{\text{particle}} = \frac{1}{2} \int d\tau \delta \left(g_{\mu\nu} \frac{dx^\mu}{d\tau} \frac{dx^\nu}{d\tau} \right),$$

and leads to the equation of motion

$$\frac{d}{d\tau} \frac{\partial \mathcal{L}_{\text{particle}}}{\partial (dx^\mu/d\tau)} = 2 \frac{d}{d\tau} \left(g_{\mu\nu} \frac{dx^\nu}{d\tau} \right) = \frac{\partial \mathcal{L}_{\text{particle}}}{\partial x^\mu} = (\partial_\mu g_{\rho\sigma}) \frac{dx^\rho}{d\tau} \frac{dx^\sigma}{d\tau}.$$

please turn over

2. Torsion and Riemann Tensor as Multi-Linear Maps

The covariant derivative with respect to a vector field X is defined by

$$\nabla_X \equiv X^\mu \nabla_\mu$$

and defines a linear map from second vector field Y to a third vector field $\nabla_X Y$. Using this notation show the following:

(a) The torsion tensor with the components $T^\lambda_{\mu\nu} \equiv \Gamma^\lambda_{\mu\nu} - \Gamma^\lambda_{\nu\mu}$ for an arbitrary connection $\Gamma^\lambda_{\mu\nu}$ defines a bilinear map from two vector fields $X = X^\mu \partial_\mu$ and $Y = Y^\nu \partial_\nu$ to a third vector field whose components are given by

$$T^\lambda_{\mu\nu} X^\mu Y^\nu.$$

Show that this bilinear map can be written as

$$T(X, Y) = \nabla_X Y - \nabla_Y X - [X, Y].$$

(b) The Riemann tensor defines a trilinear map from three vector fields $X = X^\mu \partial_\mu$, $Y = Y^\nu \partial_\nu$, and $Z = Z^\sigma \partial_\sigma$ to a fourth vector field whose components are given by

$$R^\rho_{\sigma\mu\nu} X^\mu Y^\nu Z^\sigma.$$

Show that this trilinear map can be written as

$$R(X, Y)Z = \nabla_X \nabla_Y Z - \nabla_Y \nabla_X Z - \nabla_{[X, Y]} Z.$$

To demonstrate this express the components $R^\rho_{\sigma\mu\nu}$ of the Riemann tensor in terms of general connection coefficients $\Gamma^\lambda_{\mu\nu}$ as derived in the lecture.

3. Properties of the Riemann Tensor

(a) Show that the sum over cyclic permutations of the last three indices of the Riemann tensor vanishes,

$$R_{\rho\sigma\mu\nu} + R_{\rho\mu\nu\sigma} + R_{\rho\nu\sigma\mu} = 0.$$

(b) Show that this is equivalent to

$$R_{\rho[\sigma\mu\nu]} = 0.$$

Hint: Use the expression for the Riemann tensor in a locally inertial frame.