

Theory Working Group Report.

Axions and Axion-Like Particles in Theory, Astrophysics and Cosmology

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Vistas in Axion Physics,
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Theory Working Group

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Towards a Roadmap for Theoretical and Experimental Axion Physics through 2025

- Experiments need guidance from theory, astrophysics, and cosmology
- What are the best motivated regions in the landscape of axion and axion-like particle (ALP) parameters?
 - THEORY:

What are the theoretically favored parameter ranges for the axion decay constant f_a and its couplings to standard model particles?

Are there good reasons to expect also ALPs? What are their decay constants and couplings?
 - ASTROPHYSICS:

Constraints from current and future observations in astrophysics which point to the existence of axions and ALPs? What are the accessible parameter ranges?
 - COSMOLOGY:

Are there favored regions in parameter space for axions and ALPs such that they may constitute CDM?
- There are some hot spots in this landscape: overlaps between theory, astrophysics and cosmology



THEORY

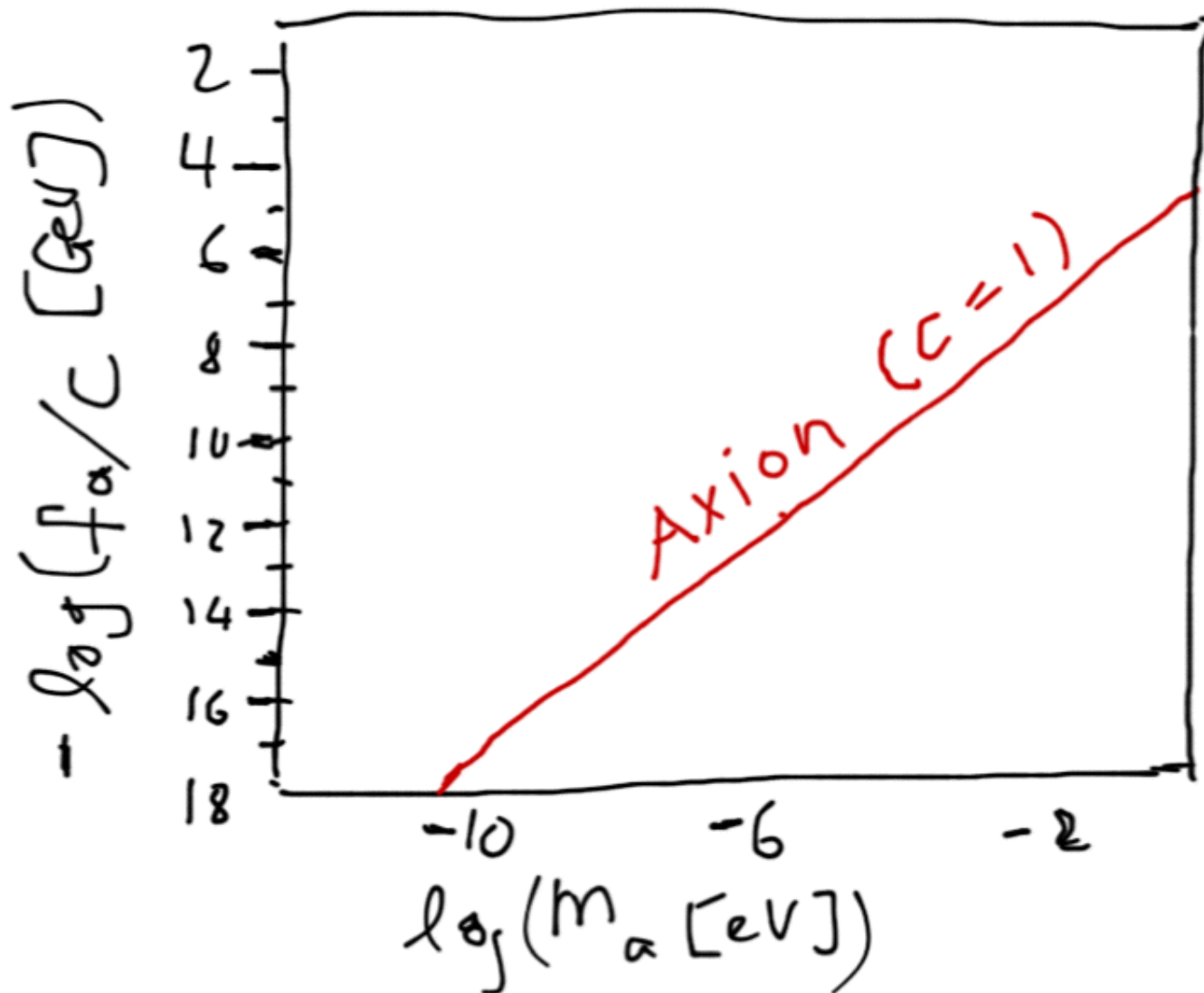
- > What are the theoretically favored parameter ranges for the axion decay constant f_a and its couplings to standard model particles?
- > Most general effective lagrangian for the axion at the weak scale in the standard model, cf. [Georgi, Kaplan, Randall '86](#)

$$\begin{aligned} \mathcal{L}^a = & \frac{1}{2}(\partial^\mu a)(\partial_\mu a) \\ & + (\partial^\mu a/f) \left(x_\varphi \varphi^\dagger i \overleftrightarrow{D}_\mu \varphi + \sum_\psi \bar{\psi}_L \gamma_\mu X_\psi \psi_L \right) \\ & - (a/f) [(g_3^2/32\pi^2) G\tilde{G} + C_{aWW}(g_2^2/32\pi^2) W\tilde{W} \\ & + C_{aYY}(g_1^2/32\pi^2) Y\tilde{Y}] . \end{aligned}$$

- > Scale effective field theory down below chiral symmetry breaking scale: model independent predictions, such as mass and coupling to photons and nucleons

$$m_a^2 = \frac{m_\pi^2}{(m_u + m_d) \text{tr}(M^{-1})} \frac{f_\pi^2}{f^2}$$





- > To be more predictive, need specific ultra-violet completion
 - KSVZ and DFSZ are particular examples with free parameter f_a
- > More predictive: completions which also solve other problems and link f_a to other scales such as the see-saw scale 10^{10} GeV, the GUT-scale 10^{16} GeV or the string scale (which can be both)
- > String compactifications generically contain pseudo-scalar fields with axionic coupling to gauge fields and anomalous global shift symmetry, cf. [Witten '84](#)

$$a_i F \tilde{F} \qquad a_i \longrightarrow a_i + \epsilon$$

- > These axion and axion-like particle (ALP) candidates arise in string compactifications as KK zero modes of antisymmetric tensor fields:

heterotic string : B_2

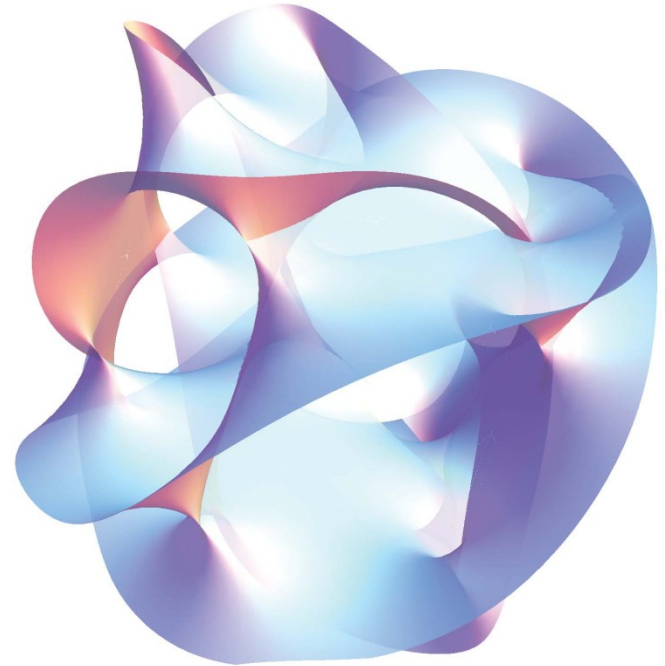
IIB string : C_2, C_4

- KK reduction (expansion in harmonic forms):

$$C_2 = c^a(x)\omega_a, \quad a = 1, \dots, h_-^{1,1}$$

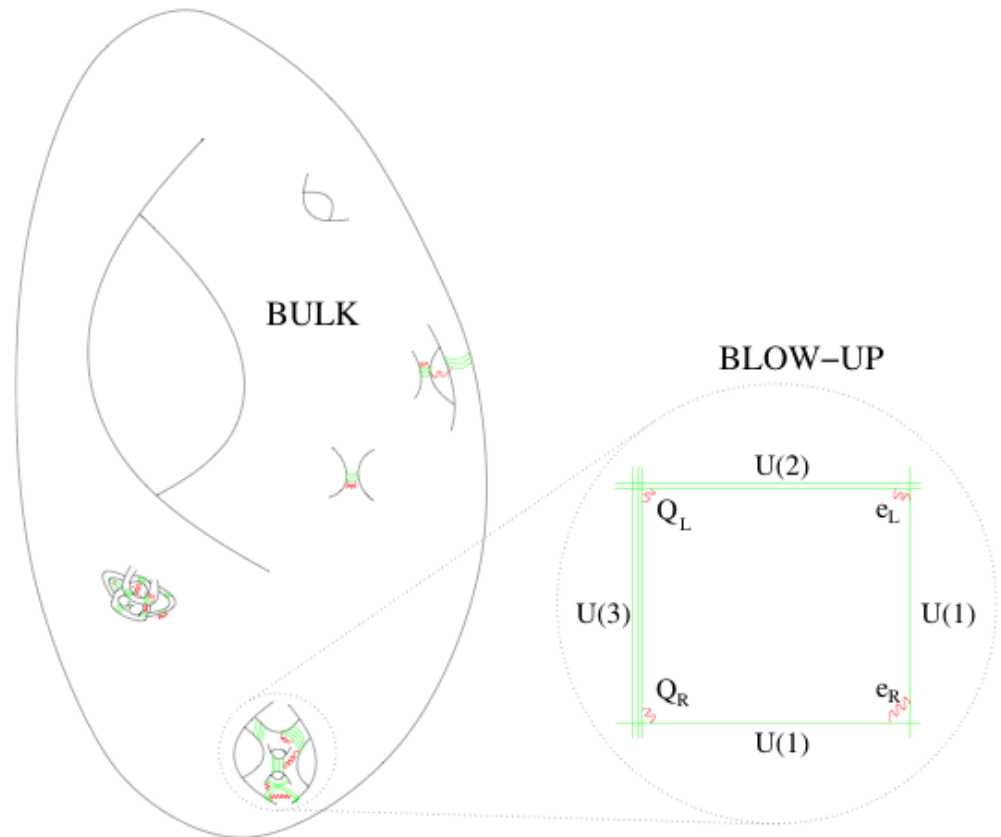
$$C_4 = c_\alpha(x)\tilde{\omega}^\alpha + \dots, \quad \alpha = 1, \dots, h_+^{1,1}$$

- Number of axionic fields, $h^{1,1}$, determined by topology of the 6D compactification manifold: number of topologically non-equivalent 2D or 4D submanifolds called cycles
- This number can be large



THEORY

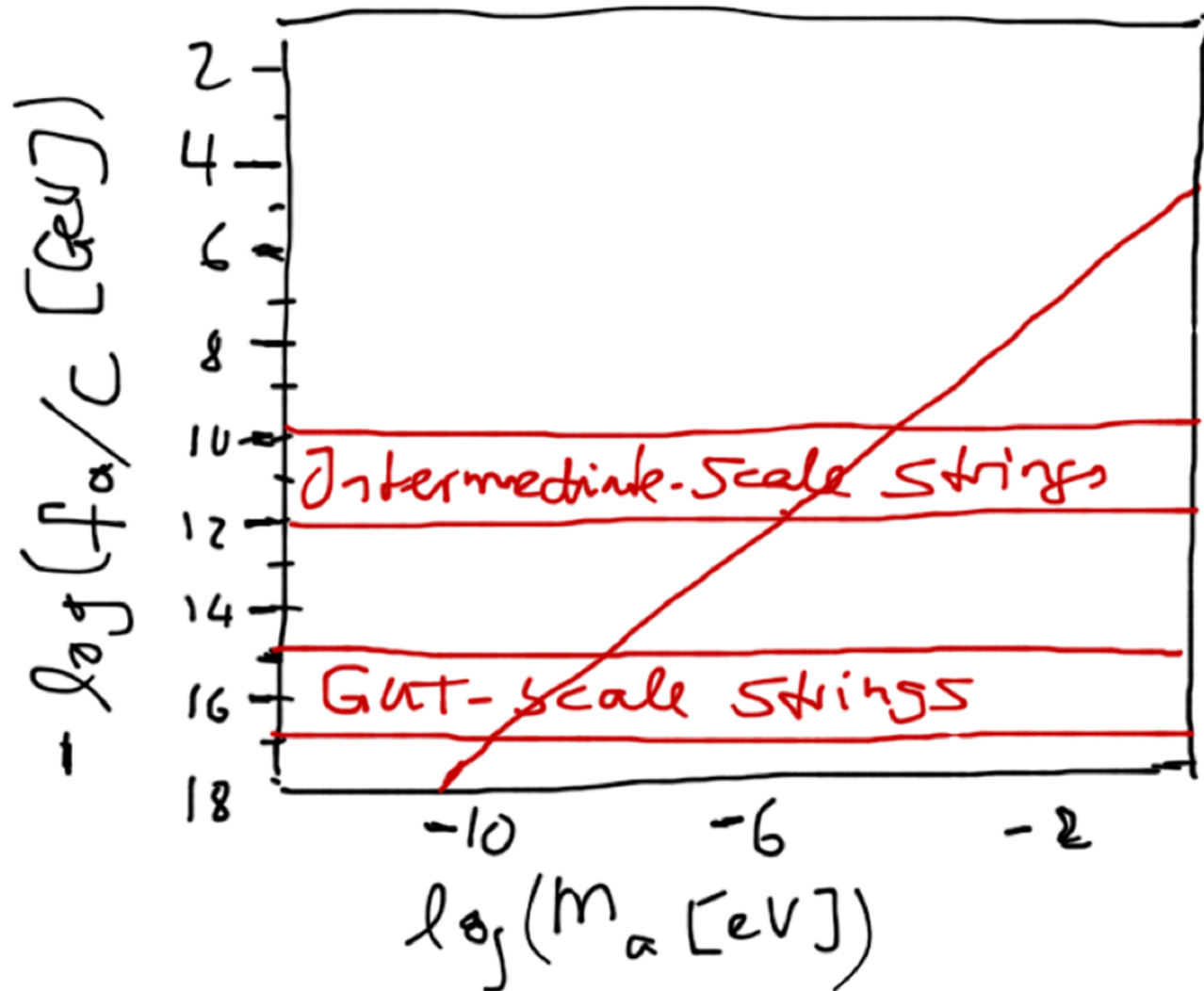
- Cycles can be wrapped by space-time filling branes
- Each of these branes gives rise to a gauge theory at low energy
- Axion and ALP decay constants as well as couplings to gauge bosons, e.g. photons, can be extracted after integrating out heavy KK fields



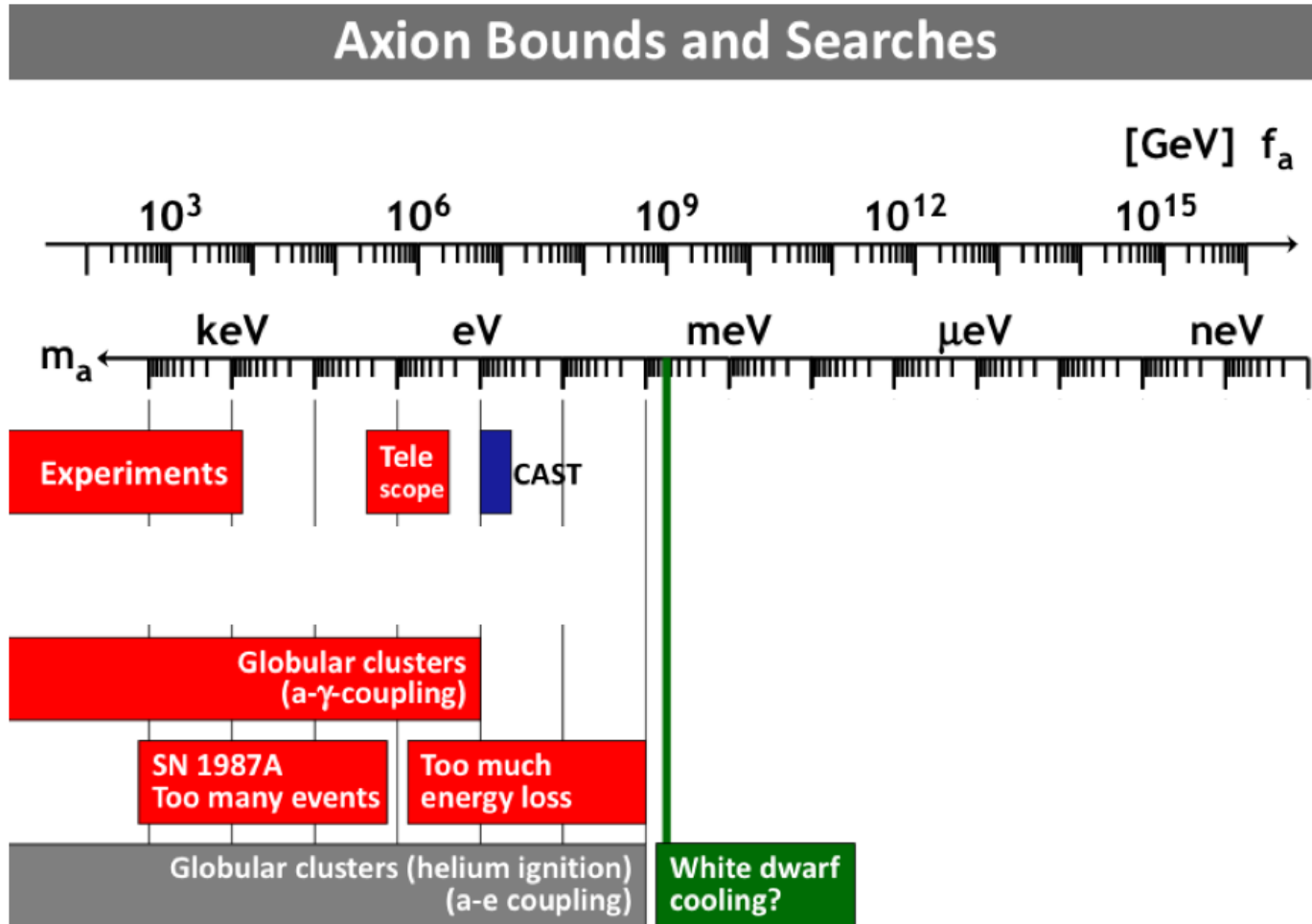
- > QCD axion plus many ALPs – an axiverse – guaranteed from string theory? cf. [Arvanitaki et al. '09](#)
- > [Acharya, Bobkov, Kumar '11](#): examples in M-theory with $f \sim 10^{16}$ GeV
- > [Cicoli, Goodsell, AR, DESY 12-058, 1205.nnnn](#): explicit globally consistent chiral model examples in IIB-theory (CYs with $h_{11}=4,5$, few, plus brane setups and fluxes), cf. [Cicoli, Mayrhofer, Valandro '11](#):
 - GUT-like model with 5 magnetised D7 branes: QCD axion with $f \sim 10^{16}$ GeV
 - MSSM-like model: QCD axion plus one light ALP, both have $f \sim 10^{16}$ GeV
 - Chiral model with QCD axion plus one light ALP, both have $f \sim M_s \sim 10^{11}$ GeV
- > In string compactifications with a QCD axion one expects that typically there are in addition further ALPs which have approximately the same decay constant and coupling to photons as the QCD axion

➤ Low energy effective lagrangian:

$$\begin{aligned}
 \mathcal{L} \supset & \frac{1}{2} \partial_\mu a_i \partial^\mu a_i - \frac{g_3^2}{32\pi^2} \left(\theta_0 + C_{i33} \frac{a_i}{f_{a_i}} \right) F_{3,\mu\nu} \tilde{F}_3^{b,\mu\nu} \\
 & - \frac{g_2^2}{32\pi^2} C_{iWW} \frac{a_i}{f_{a_i}} F_{W,\mu\nu} \tilde{F}_W^{b,\mu\nu} - \frac{g_Y^2}{32\pi^2} C_{iYY} \frac{a_i}{f_{a_i}} F_{Y,\mu\nu} \tilde{F}_Y^{\mu\nu} \\
 & - \sum_{\psi_L, \psi_R} (\psi_R M \psi_L + \bar{\psi}_L M^\dagger \bar{\psi}_R) + \sum_{\psi_L} \frac{X_{\psi_L}^i}{f_{a_i}} \psi_L \sigma^\mu \bar{\psi}_L \partial_\mu a_i + \sum_{\psi_R} \frac{X_{\psi_R}^i}{f_{a_i}} \psi_R \sigma^\mu \bar{\psi}_R \partial_\mu a_i \\
 & - \sum_{\psi_L, \psi_R} (i \psi_R g^i \psi_L \frac{a_i}{f_{a_i}} + h.c.) - V(a_i), \tag{A.1}
 \end{aligned}$$

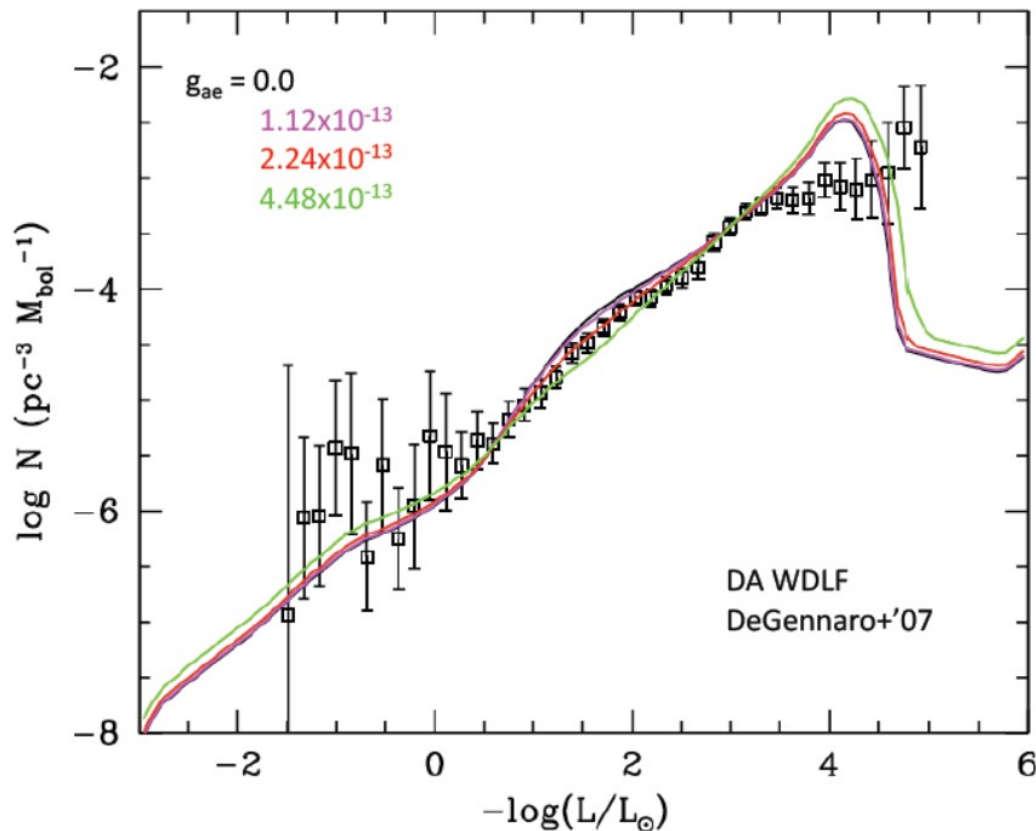


- Astrophysics constrains $f/C > 10^9$ GeV for C_γ and C_e , but at the border there are some intriguing observations



- Standard model does not fit to white dwarf luminosity function
- Extra energy loss compatible with axion or ALP bremsstrahlung

Isern et al. '08



- Required parameters:

$$m_{a_i} \lesssim \text{keV}, \quad g_{ie} \equiv \frac{C_{ie} m_e}{f_{a_i}} = (0.6 \div 1.7) \times 10^{-13}$$

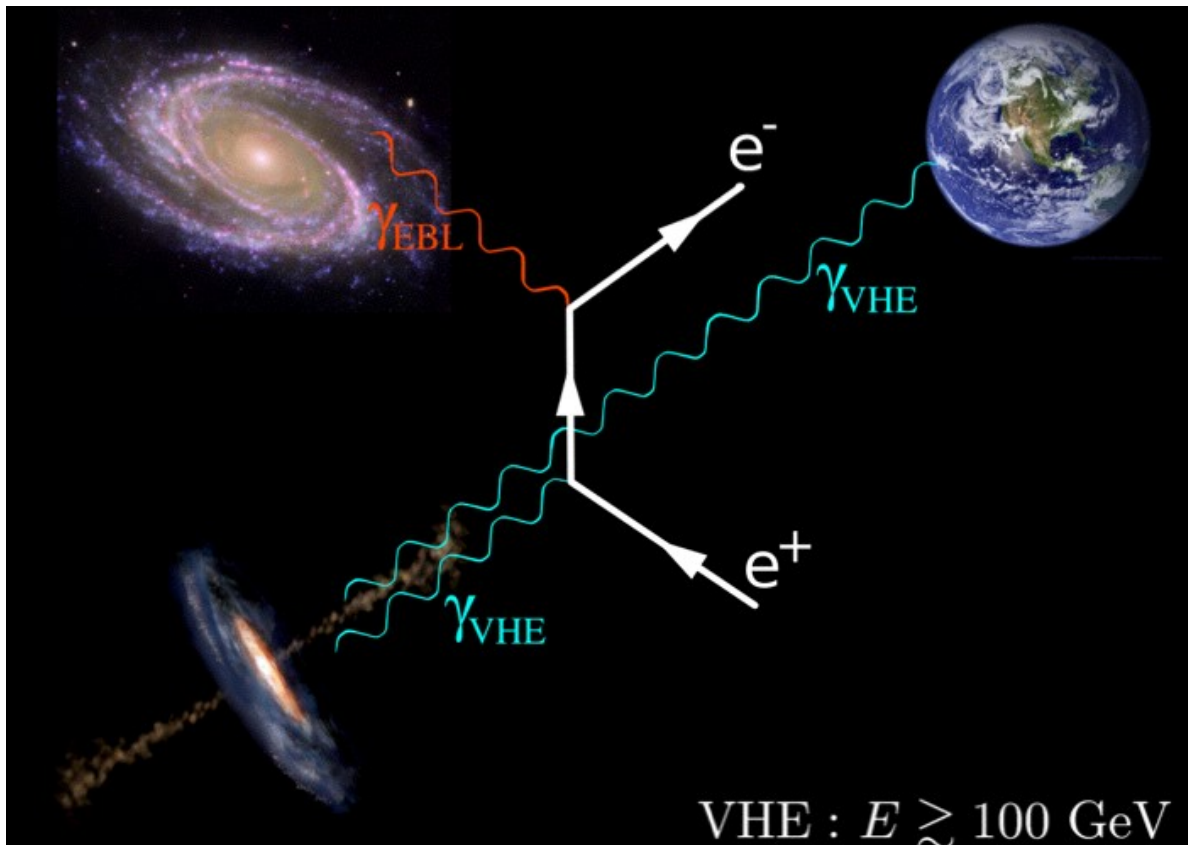
$$\Rightarrow f_{a_i} \sim 10^{10} C_{ie} \text{ GeV}$$

- Can be axion or ALP



ASTROPHYSICS

- TeV photon spectra of distant AGNs should show absorption feature due to electron-positron pair production on extra-galactic background light

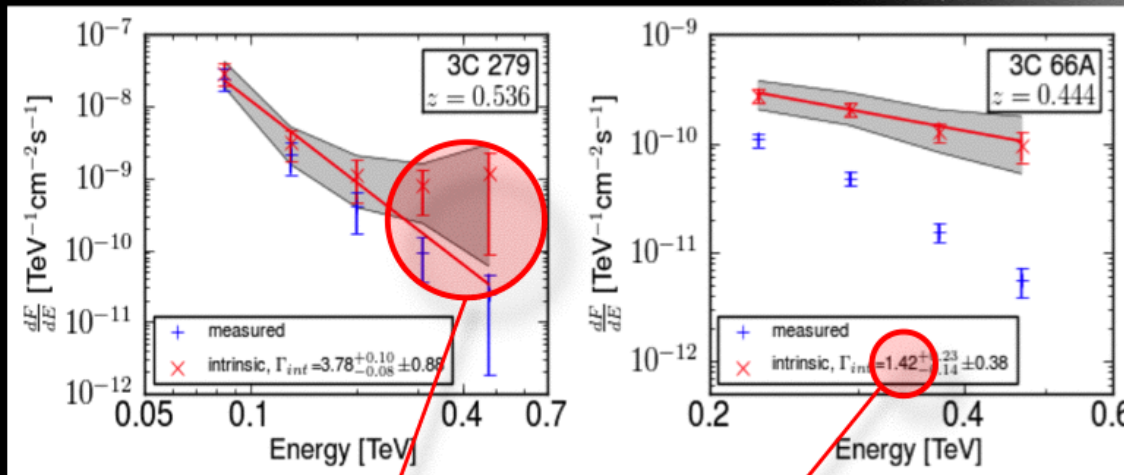




Hints for new physics?

- Some recent gamma-ray observations pose substantial challenges to the conventional models.
 - Intrinsic spectrum deviates from a power-law: pile-up problem (Dominguez+12).
 - Very hard intrinsic spectrum of FSRQs (e.g. Albert+08, Alecsik+11, Wagner+10)
 - Extremely rapid and intense flares (Tavecchio+12).
 - GeV spectral breaks!

[Domínguez et al. 2011]



See also Horns+12 !

PILE-UP!

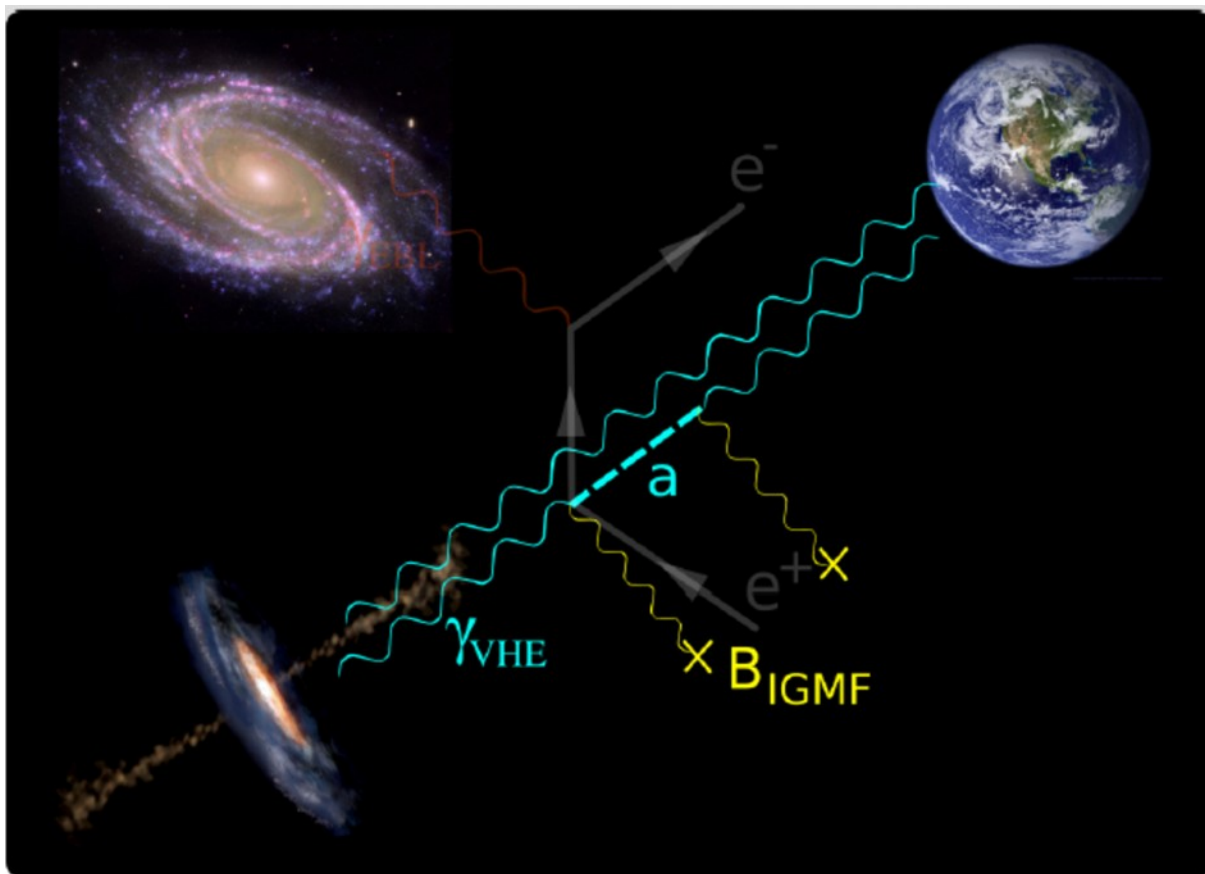
Modeling of AGN emission mechanisms typically assume spectral index > 1.5

cf. Sanchez-Conde



ASTROPHYSICS

- Possible explanation: photon $\leftrightarrow$ ALP oscillations in astrophysical magnetic fields



- Required parameters: cf.

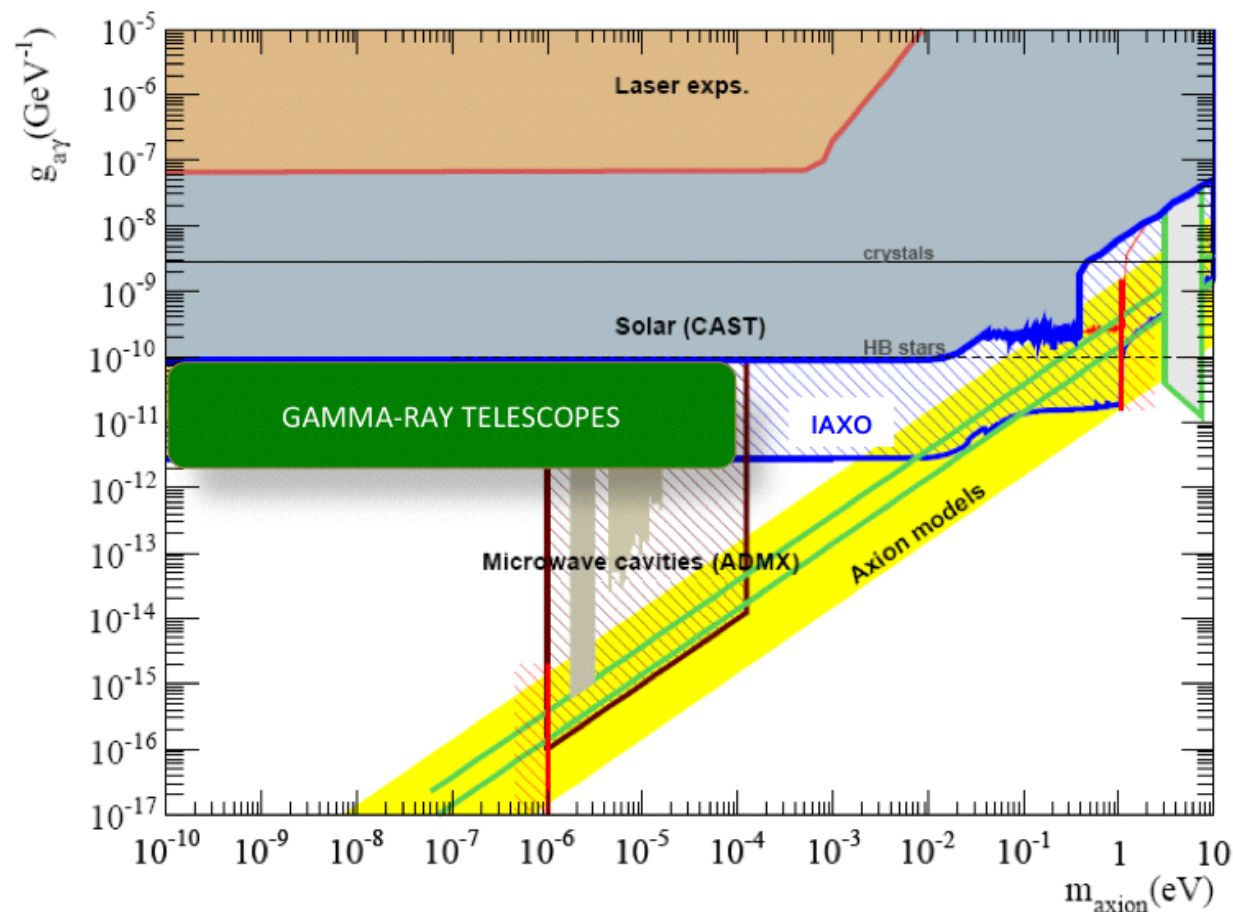
$$m_{a_i} \lesssim \text{neV}, \quad g_{i\gamma} \equiv \frac{\alpha}{2\pi} \frac{C_{i\gamma}}{f_{a_i}} \sim 10^{-12} \div 10^{-11} \text{ GeV}^{-1}$$

$$\Rightarrow f_{a_i} \sim 10^8 \div 10^9 C_{i\gamma} \text{ GeV}$$

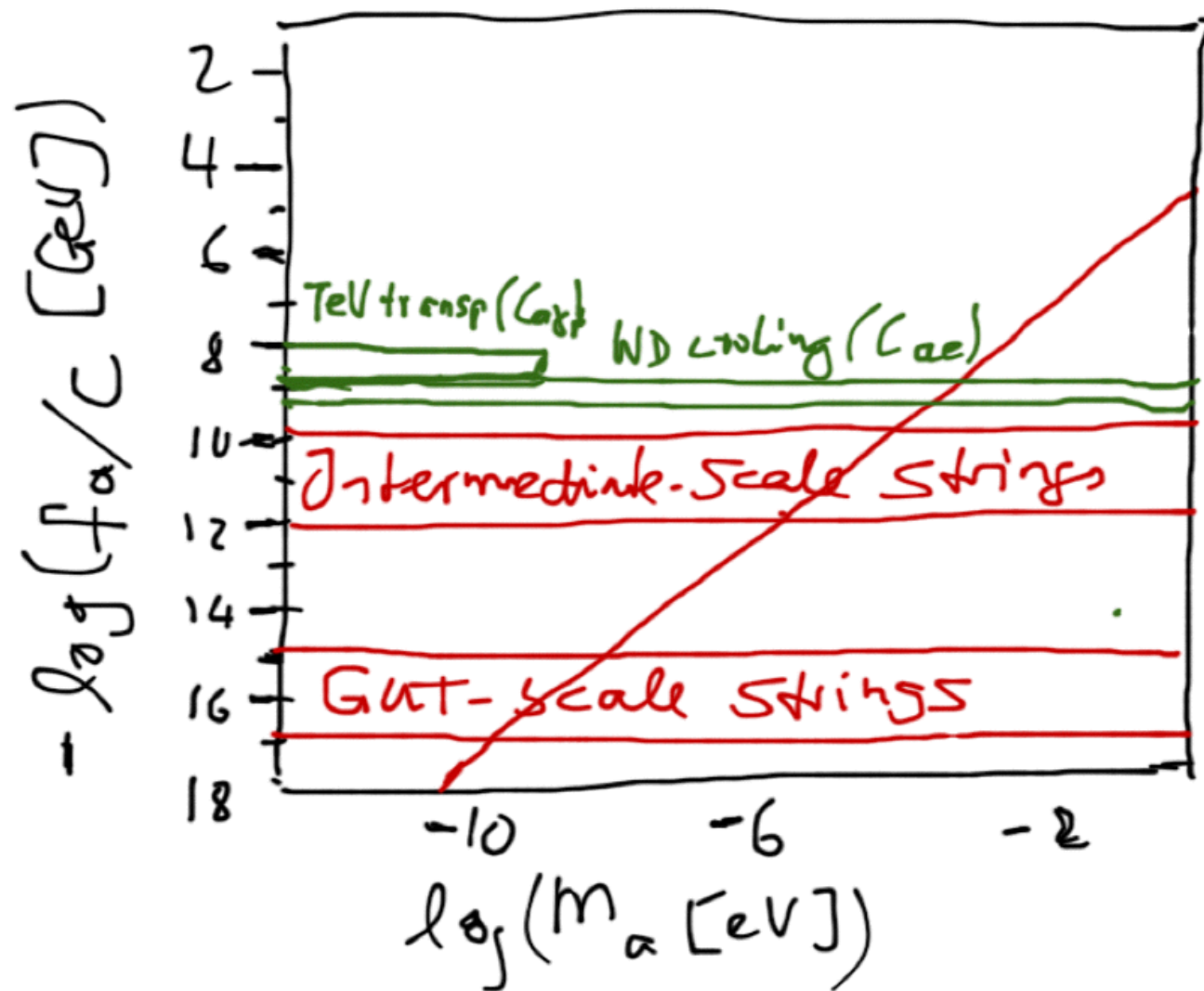
- Must be an ALP, too light for an axion with such a decay constant



Fermi + IACTs can explore a region of the ALP parameter space that is difficult to explore otherwise!



> cf. Sanchez-Conde

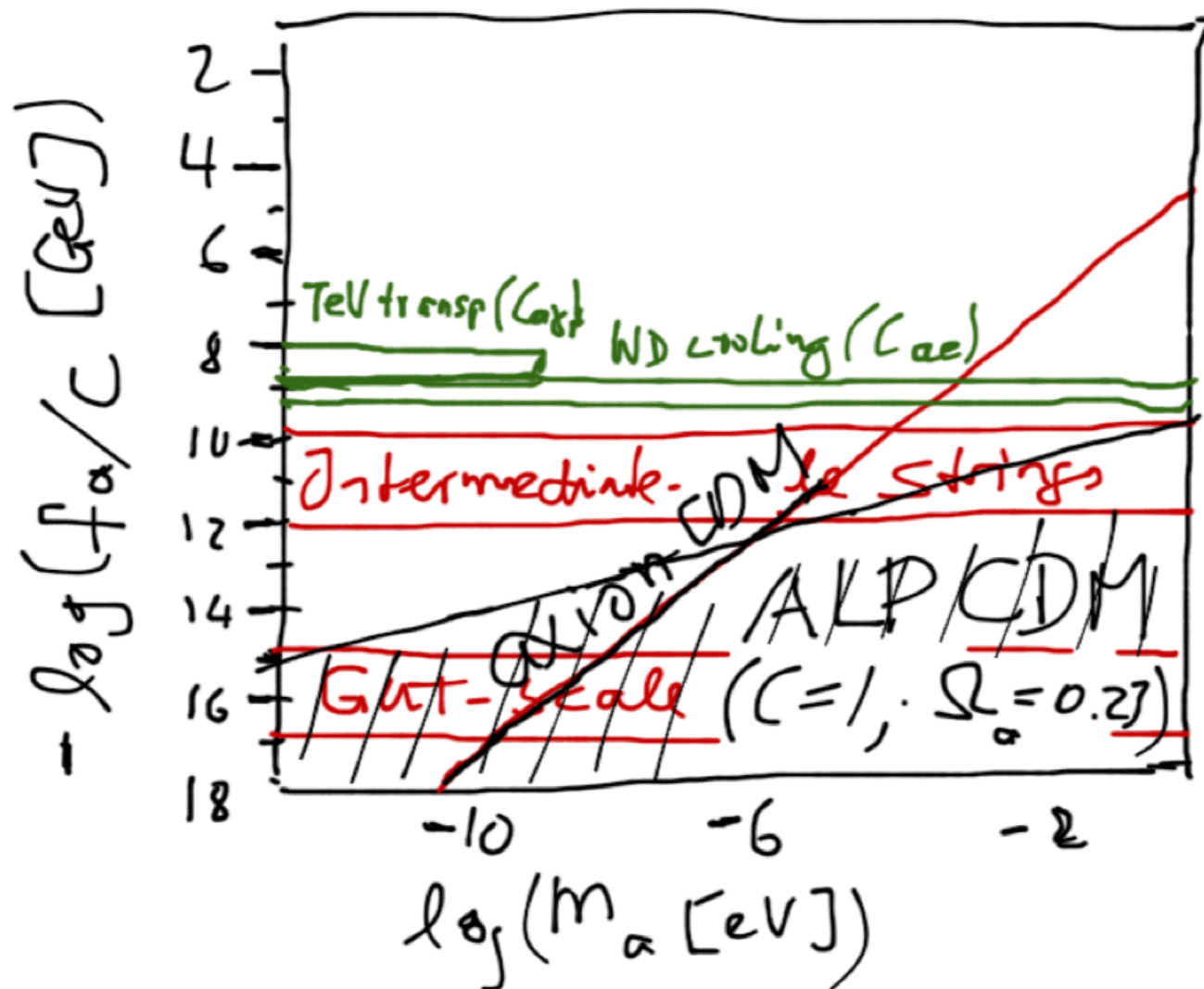


- > Are there favored regions in axion or ALP parameter space such that they constitute CDM?
- > Axion produced via misalignment mechanism,

$$\Omega_a h^2 \approx 0.7 \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{7/6} \left(\frac{\theta_i}{\pi} \right)^2$$

- > Generic ALP also produced via misalignment mechanism. Constitutes all CDM for (cf. [Redondo](#))

$$\Omega_{a_i} h^2 \approx 10^{-2} \times \left(\frac{m_{a_i}}{\mu\text{eV}} \right)^{1/2} \left(\frac{f_{a_i}}{10^{12} \text{ GeV}} \right)^2 \left(\frac{\theta_{a_i}}{\pi} \right)^2$$



- > Wide mass range where axion or ALP could be dark matter:

Conclusions

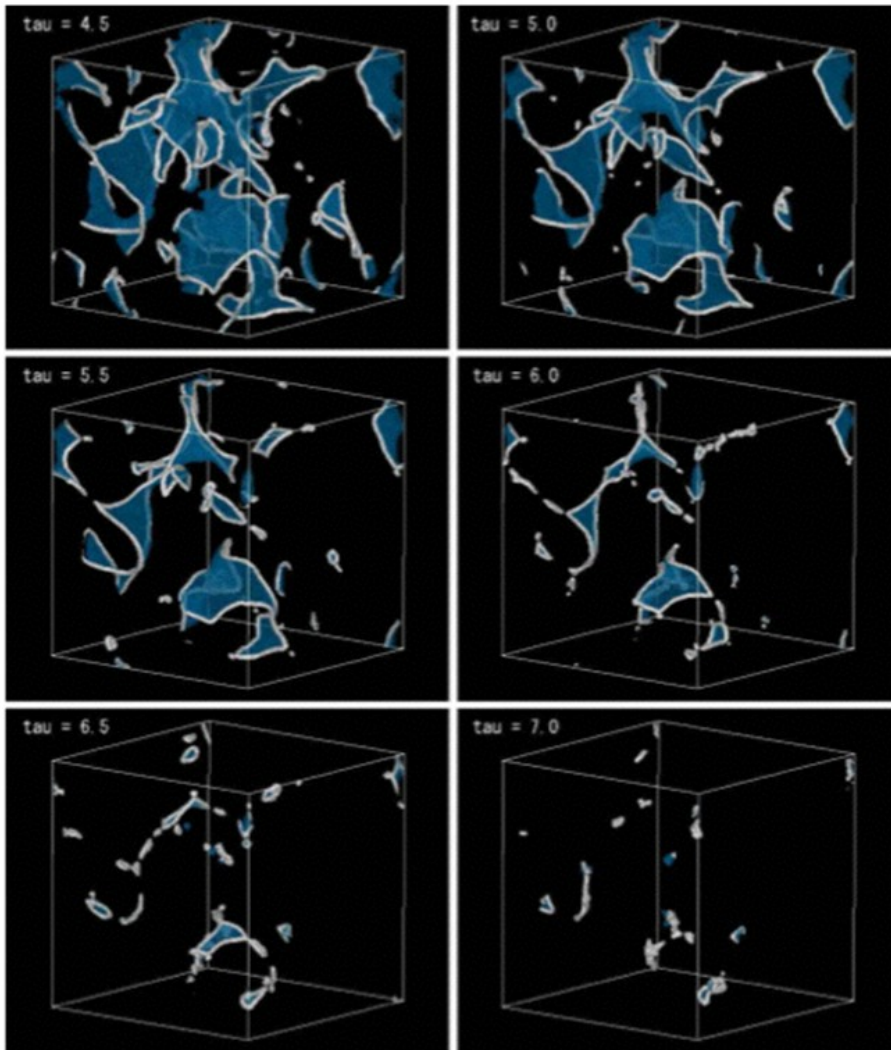
For axions to be 100% of cold dark matter....

- If the Peccei-Quinn symmetry breaks after inflation ends, the axion mass must be $m_a = 83 \mu\text{eV}$ in standard cosmology
 - much smaller m_a in LTR cosmology
 - much larger m_a in kination cosmology $0.1 \mu\text{eV} < m_a < 15 \text{ meV}$
- If the Peccei-Quinn symmetry breaks before inflation ends, an initial misalignment angle θ_i can be chosen for any $m_a < 15 \text{ meV}$
 - larger allowed region and larger θ_i in LTR cosmology
 - smaller allowed region and smaller θ_i in kination cosmology

- > cf. Gondolo



Axion Production by Domain Wall and String Decay



Recent numerical studies of collapse of string-domain wall system

$$\Omega_a h^2 = (16 \pm 6) \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{1.19} \times \left(\frac{g_{*,1}}{70} \right)^{-0.41} \left(\frac{\Lambda}{400 \text{ MeV}} \right)$$

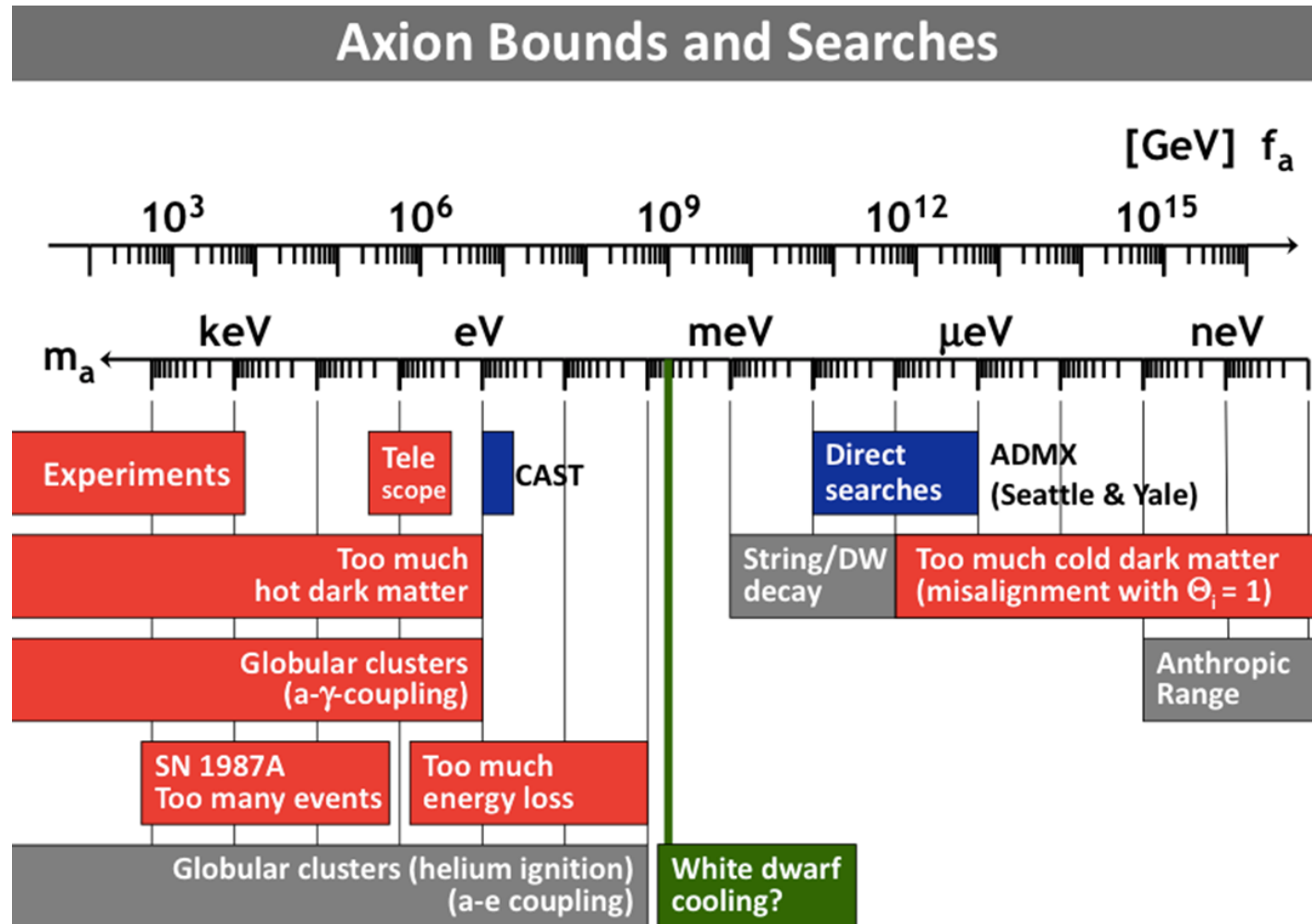
Implies a CDM axion mass of

$$m_a \sim 1 \text{ meV}$$

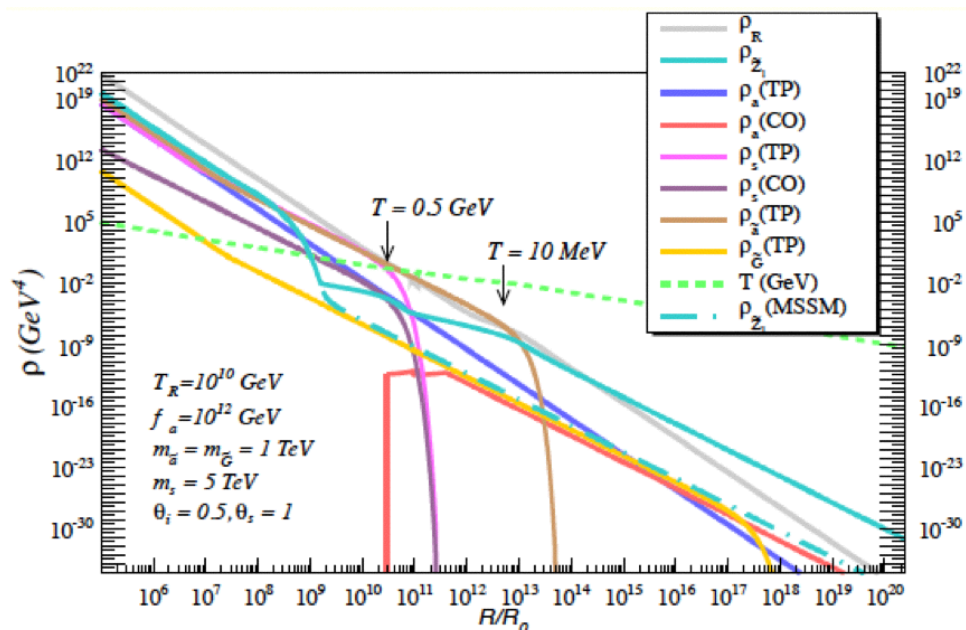
Hiramatsu, Kawasaki, Saikawa, Sekiguchi, arXiv:1202.5851 (2012)

Remains to be confirmed, interpretation of numerical studies not entirely straightforward

- Wide mass range where axion or ALP could be dark matter:



- The SUSY solution to gauge hierarchy problem and axion solution to strong CP problem are complementary: one needs both.
- Since lack of p-decay suggests R-parity conservation, then in this case one would get mixed axion-LSP dark matter.
- In this case, one must calculate the dark matter abundance including effects of SUSY LSP, axion, axino, saxion and gravitino. LSP might be neutralino, axino or gravitino, cf. [Baer](#)

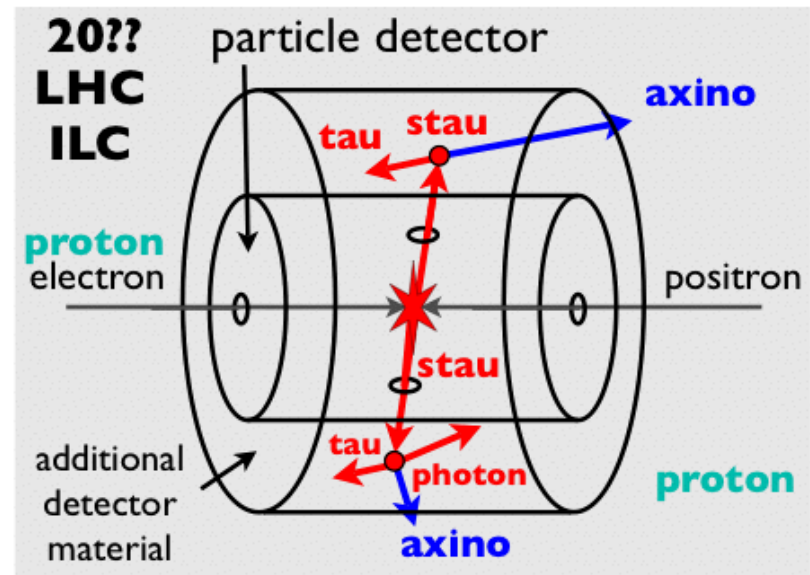
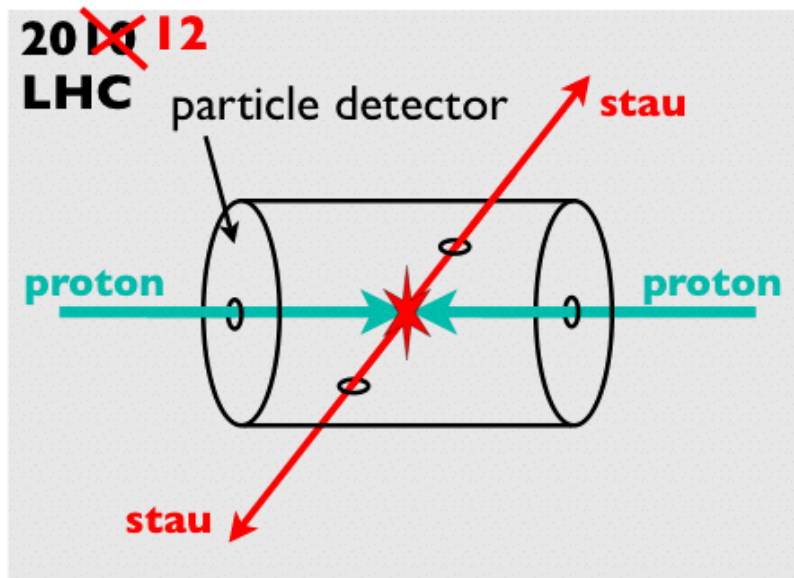


- > Much wider range of f_a is now allowed due to possibility of entropy dilution at temperature $T_{\text{fr}} > T > T_{\text{BBN}}$, and also variable θ_i . f_a can lie between 10^9 - 10^{16} GeV, cf. [Baer](#)
- > Interesting link to collider physics:
 - Depending on what LHC (and later ILC) might see, we may be able to distinguish nature of LMSSMP, whether it is bino-like, higgsino-like, stau, or other.
 - SUSY models with bino-like LMSSMP usually give standard overabundance which can be diminished by decay to light MeV scale axino.
 - SUSY models with higgsino or wino-like LMSSMP and hence standard underabundance of neutralinos favor a weak scale axino mass since thermal axino production in early universe followed by decay to neutralino builds up the neutralino abundance.



- Probing the axion decay constant at the LHC, cf. [Steffen](#)

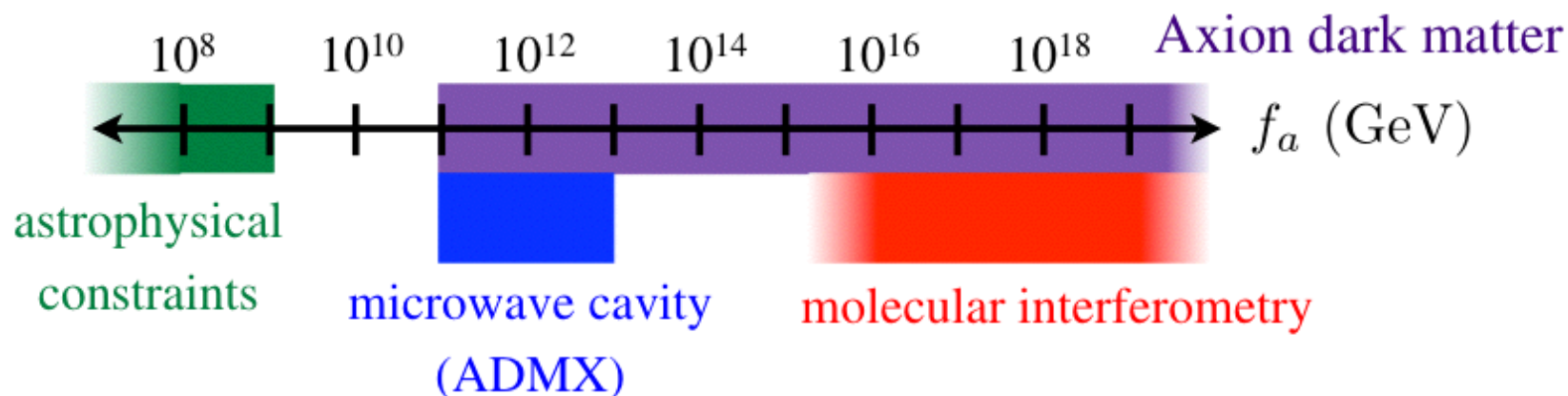
Stopping of long-lived staus



Probing f_a @ Colliders

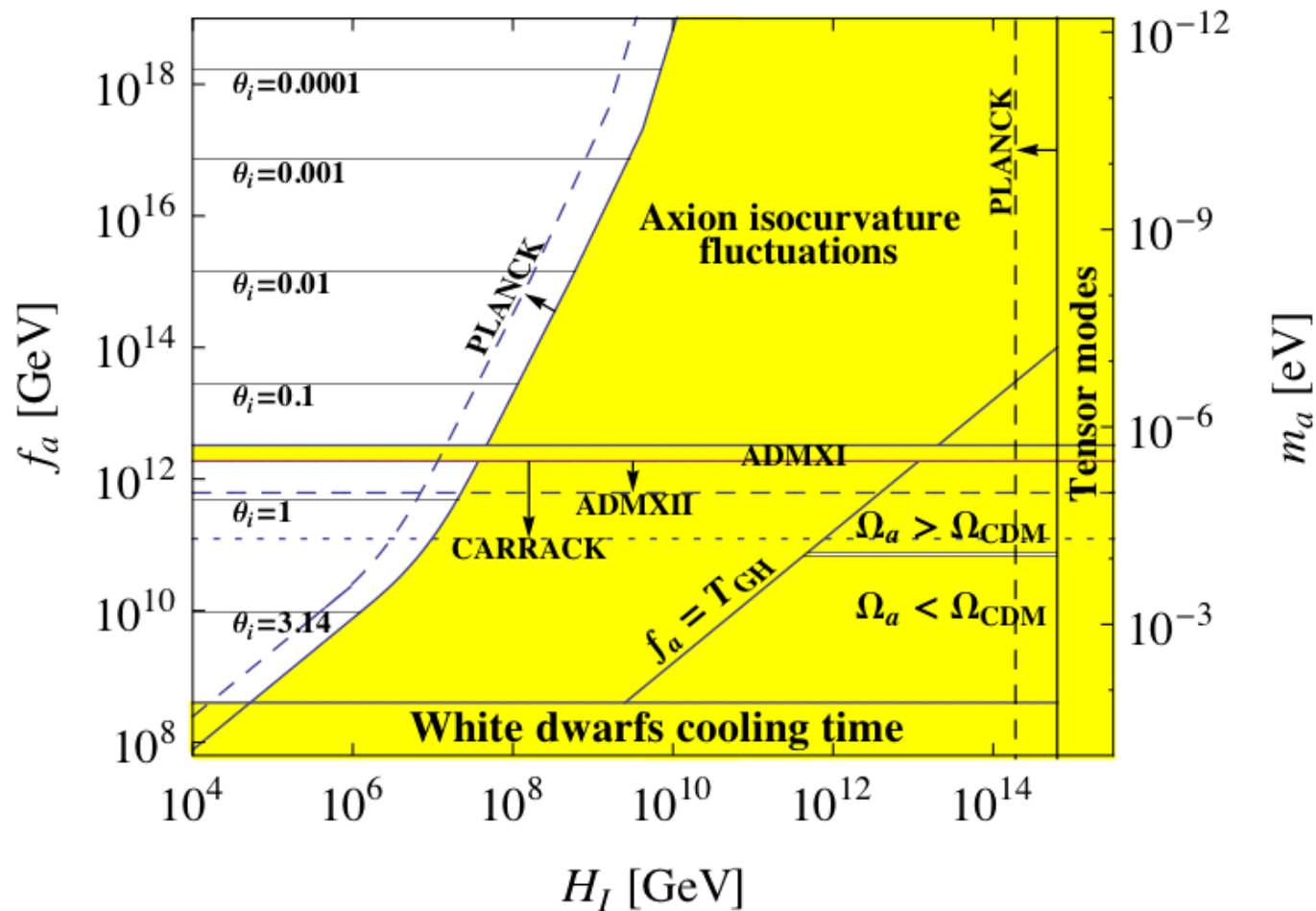
- > Direct axion and ALP CDM searches should be extended towards smaller and larger masses
- > $f_a \sim 10^{12}$ GeV, $m_a \sim \text{micro-eV}$ axion or ALP CDM doable
- > $f_a \sim 10^{16}$ GeV, $m_a \sim 10^{-10}$ eV axion CDM possibility big challenge
 - Direct detection of time varying shifts of atomic energy levels due to the coupling between internal atomic fields and time varying CP-odd nuclear moments? cf.

Graham, Rajendran '11



> CMB probes:

- Isocurvature fluctuations as a probe of large f_a : require low H_I , cf. [Gondolo](#)

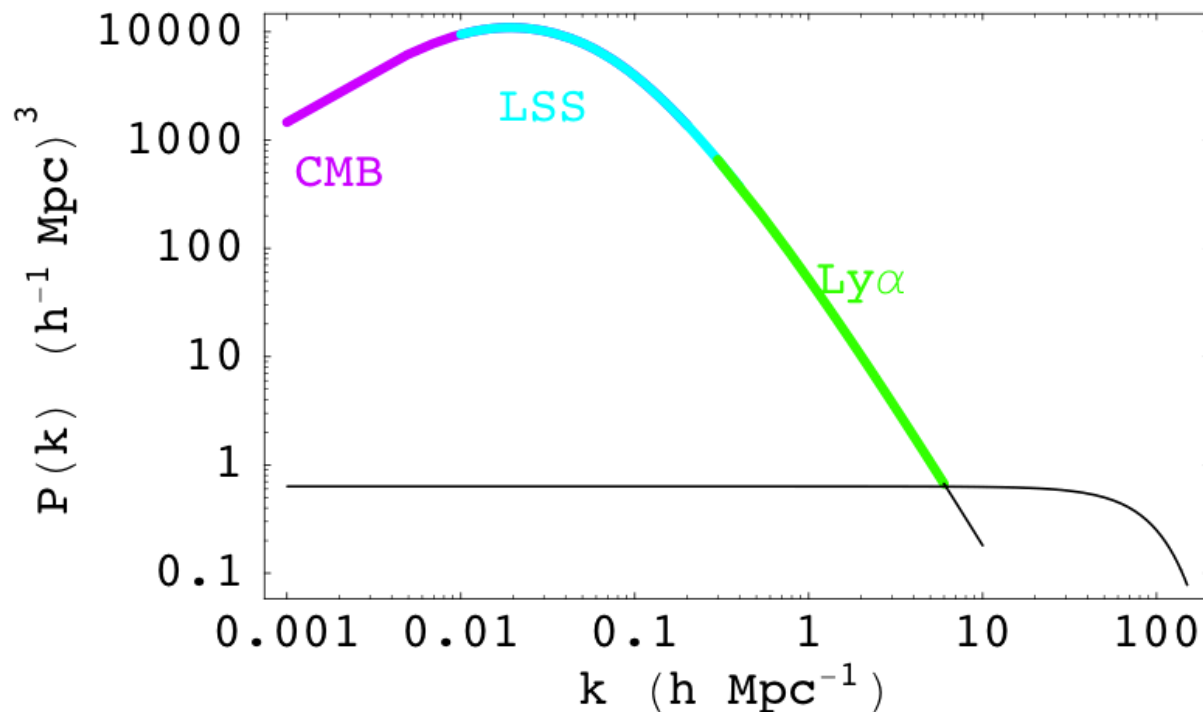


> CMB probes:

- Isocurvature fluctuations as a probe of large f_a

> Large scale and CDM structure probes:

- Miniclusters: enhancements in power spectrum, cf. [Zurek, Hogan, Quinn '07](#)

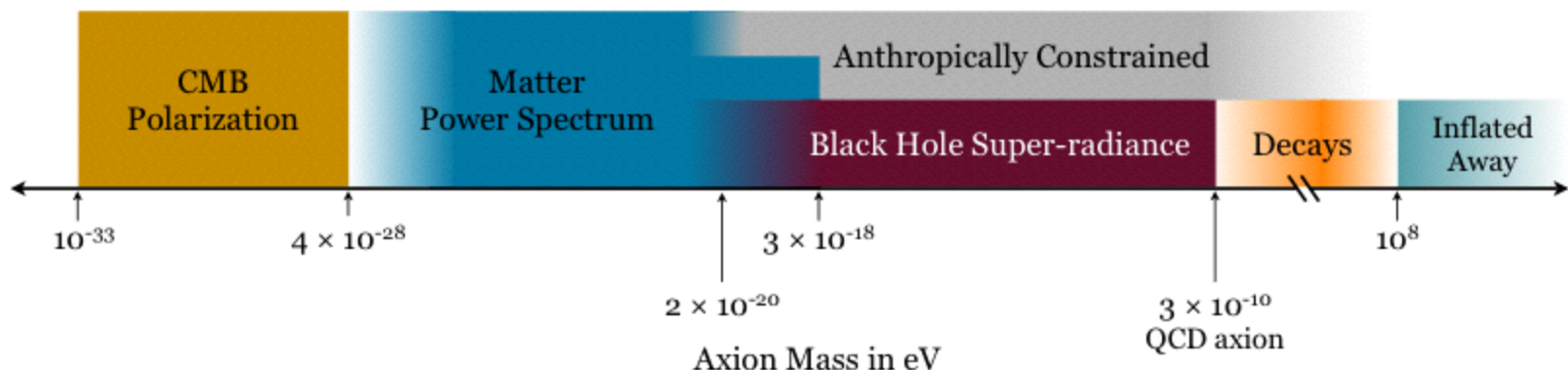


> CMB probes:

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> Large scale and CDM structure probes:

- Miniclusters: enhancements in power spectrum
- If axions or ALPs are Bose Einstein Condensate: possible solutions to CDM structure problems, cf. [Salucci, Sikivie, Quinn](#)
- Ultra-light ALPs: effects on large scale structure and CMB, cf. [Marsh, Macaulay, Trebitsch, Ferreira '11](#)



cf. [Arvanitaki et al. '10](#)

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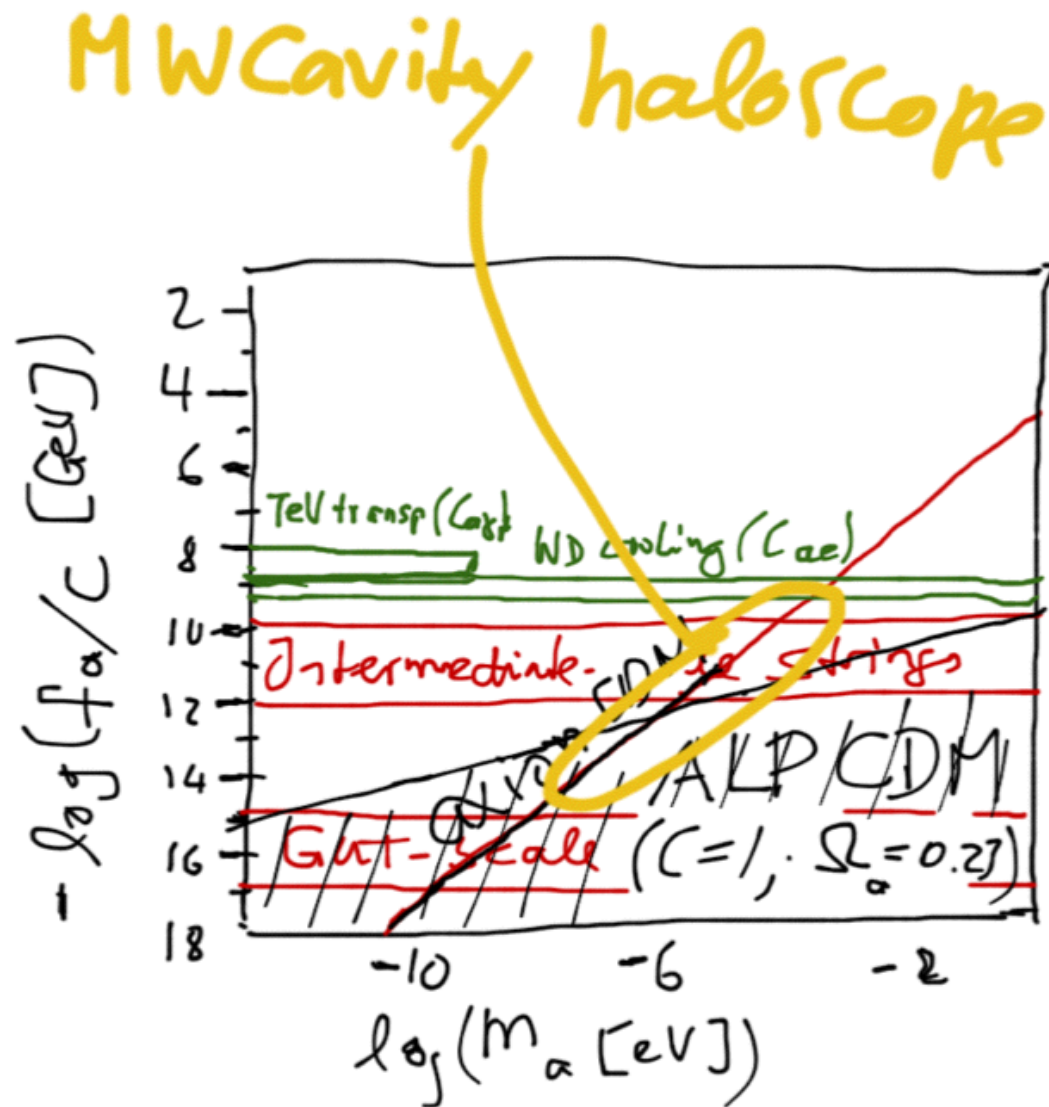
> Combination of both:

- Tilted universe (cf. [Turner](#)) where the photon rest frame differs from the matter rest frame, may be a probe for axions with GUT scale f_a , cf. [Kaplan](#)



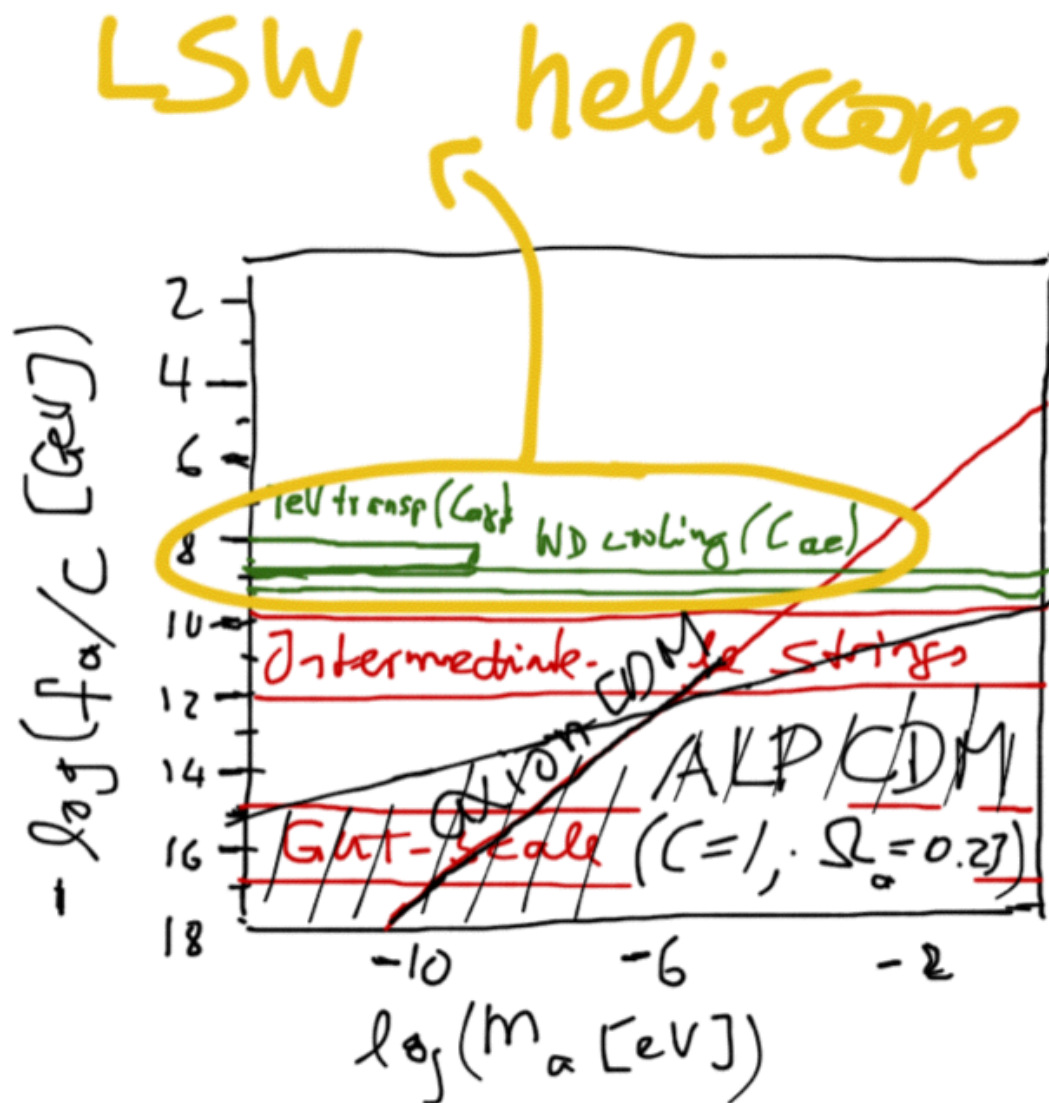
CONCLUSIONS

- Hot spots in the landscape of axions and ALPs:



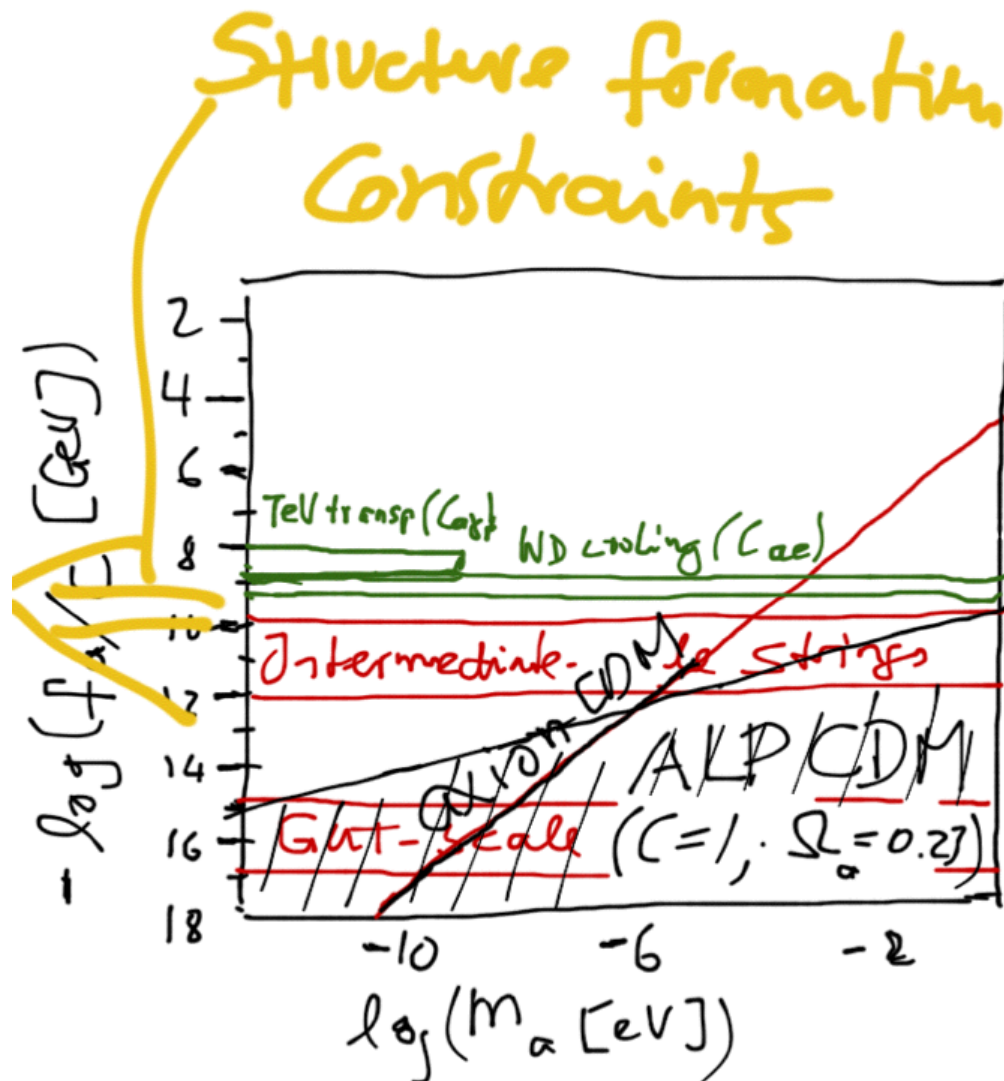
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