

Status of Little Higgs Models in 2014

Jürgen Reuter

DESY

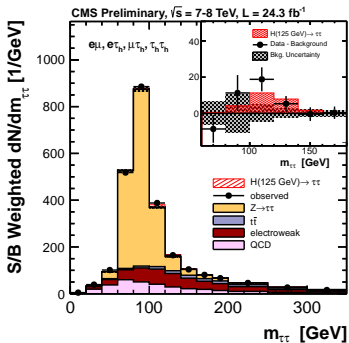
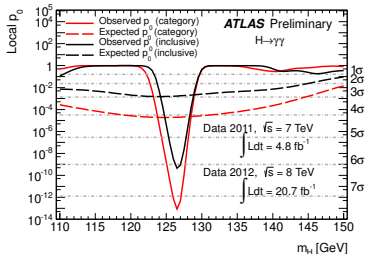
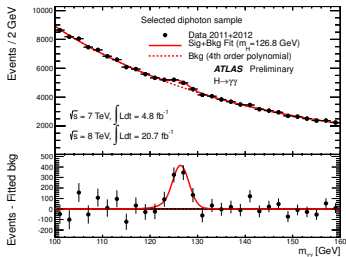


JRR/Tonini/de Vries, **JHEP 1402** (2014) 053; arXiv:1307.5010; JRR/Tonini, **JHEP 1302** (2013) 077; Kilian/JRR/Rainwater **PRD 74** (2006), 095003; **PRD 71** (2005), 015008;
Kilian/JRR **PRD 70** (2004), 015004

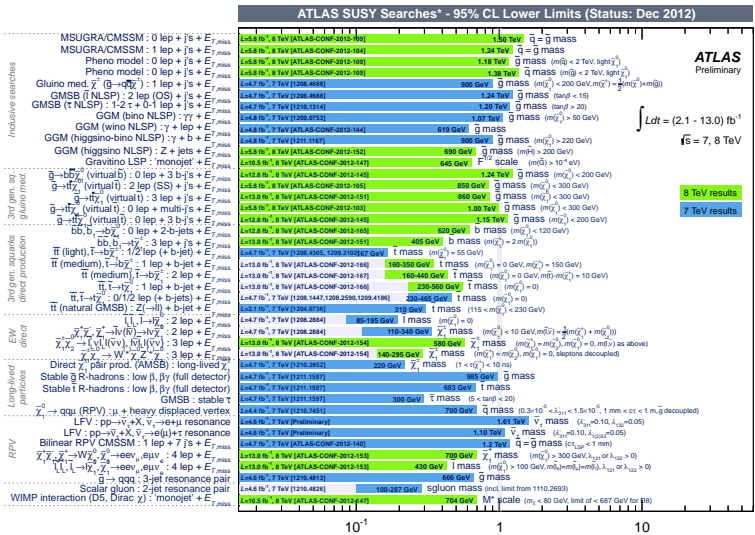
GK Seminar, KIT, Karlsruhe, 28.11.2014

Standard Model Triumph:

- ▶ 2012: Discovery of a Higgs boson

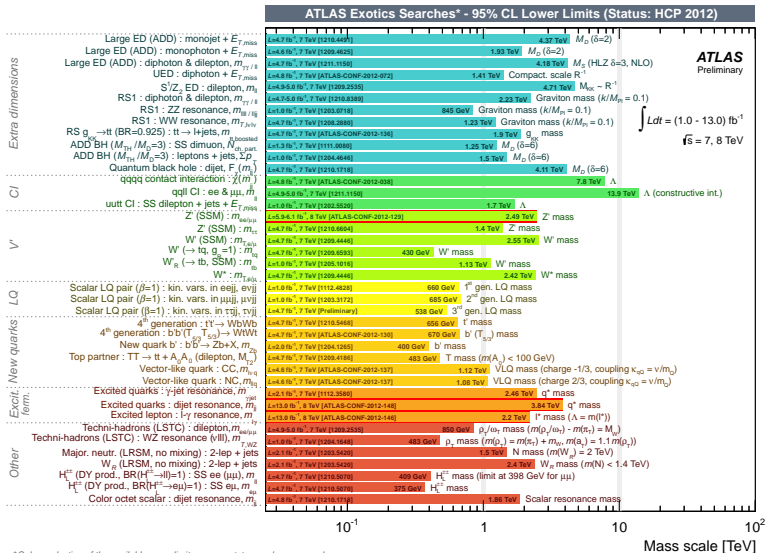


... and what now?



* Only a selection of the available mass limits on new states or phenomena shown.
 All limits quoted are observed minus 1σ theoretical signal cross section uncertainty.

... and what now?

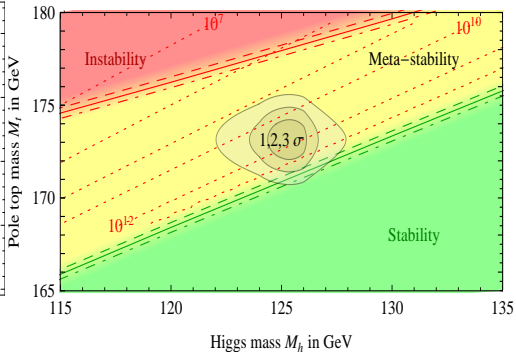
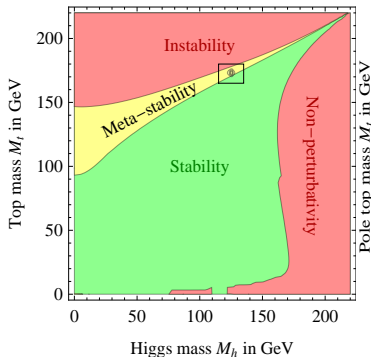


Electroweak vacuum stability

- ▶ Recent analysis: **Metastable vacuum with lifetime longer than the age of the universe** Degrassi et al., arXiv:1205.6497

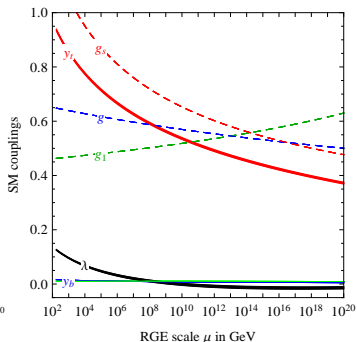
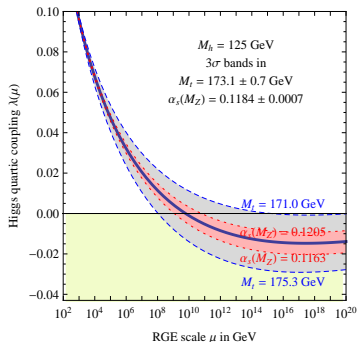
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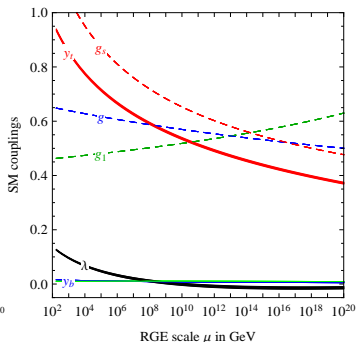
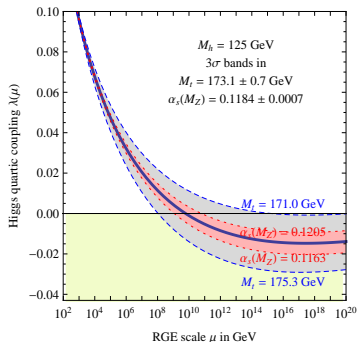
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- Could the Higgs field ever have fallen in the correct vacuum?** Hertzberg, arXiv:1210.3624

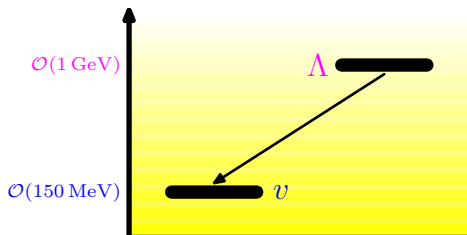
Higgs as Pseudo-Goldstone boson

Nambu-Goldstone Theorem: For each *spontaneously broken global symmetry generator* there is a **massless boson** in the spectrum.

Old idea:

Georgi/Pais, 1974; Georgi/Dimopoulos/Kaplan, 1984

Light Higgs as **(Pseudo)-Goldstone boson** of a spontaneously broken global symmetry



Analogous: QCD

Scale Λ : chiral symmetry breaking, quarks, $SU(3)_c$

Scale ν : pions, kaons, ...

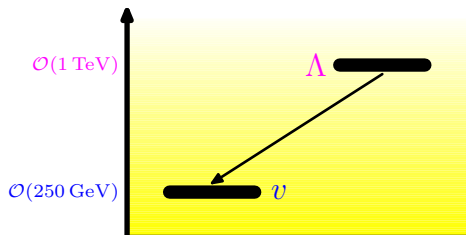
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Scale Λ : global symmetry breaking, new particles, new (gauge) IA

Scale v : Higgs, W/Z , ℓ^\pm , ...

Without Fine-Tuning: experimentally excluded

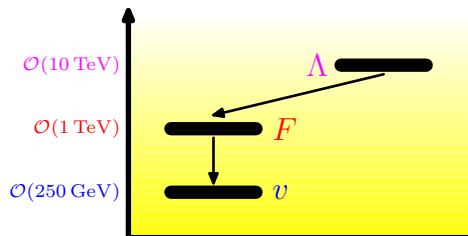
Collective symmetry breaking and 3-scale models

Collective symmetry breaking: Arkani-Hamed/Cohen/Georgi/Nelson/..., 2001

2 different global symmetries; one of them unbroken \Rightarrow Higgs
exact Goldstone boson

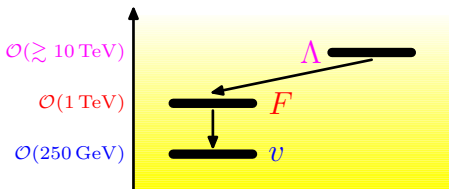
Coleman-Weinberg: boson masses by radiative corrections, but: m_H only at 2-loop level

$$m_H \sim \frac{g_1}{4\pi} \frac{g_2}{4\pi} \Lambda$$



Scale Λ : global SB, new IA
Scale F : Pseudo-Goldstone bosons, new vectors/fermions
Scale v : Higgs, W/Z , ℓ^\pm , ...

Characteristics and Spectra



Scale Λ : “hidden sector”, symmetry breaking

Scale F : new particles

Scale v : $h, W/Z, \ell^\pm, \dots$

Terascale: new particles to stabilize the hierarchy

SUSY

\tilde{q}_L ——— \tilde{t}_2
 \tilde{b}_2 ——— \tilde{q}_R
 \tilde{b}_1 ———
 ——— \tilde{t}_1

H, A ——— H^\pm

h ———

$\tilde{\tau}_2$ ——— $\tilde{\ell}_L$
 $\tilde{\nu}_\ell$ ——— $\tilde{\ell}_R$
 $\tilde{\tau}_1$ ——— $\tilde{\ell}_R$

$\tilde{\chi}_4^0, \tilde{\chi}_2^\pm$ ———
 $\tilde{\chi}_3^0$ ———

$\tilde{\chi}_2^0, \tilde{\chi}_1^\pm$ ———
 t ———
 $\tilde{\chi}_1^0$ ———

$M[\text{TeV}]$

Φ ———
 Φ^\pm ——— $\Phi^{\pm\pm}$
 Φ_P ———

h ———

η ———

Little Higgs

Z' ——— W'^\pm ——— T ———

γ' ———

U, C ———

W^\pm ——— Z ———

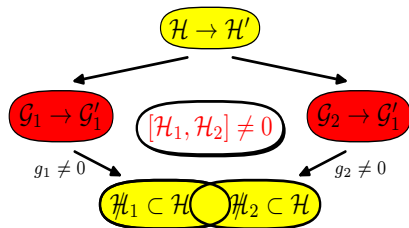
t ———

Generic properties of Little-Higgs models

- Extended global symmetry (extended scalar sector)
- **Specific functional form of the potential**
- Extended gauge symmetry: $\gamma' \equiv A_H, Z' \equiv Z_H, W'^{\pm} \equiv W_H$
- New heavy fermions: T , but also U, C, \dots

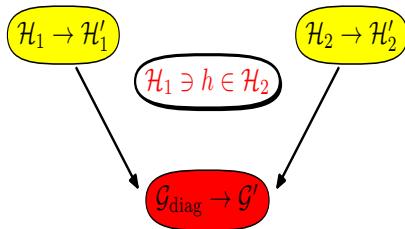
Product Group Models

(e.g. Littlest Higgs)



Simple Group Models

(e.g. Simplest Little Higgs)



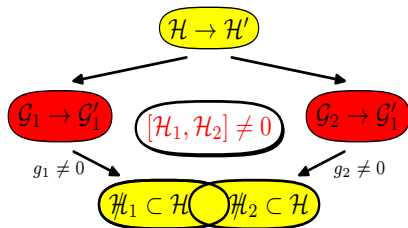
- discrete $T(\text{TeV})$ parity: pair production, cascades, DM

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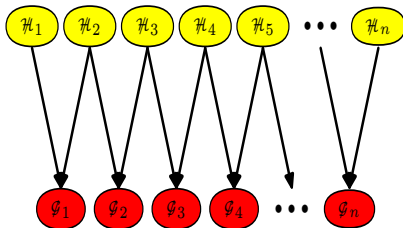
Product Group Models

(e.g. Littlest Higgs)



Moose Models

(e.g. Minimal Moose Model)



- **discrete $T(\text{TeV})$ parity: pair production, cascades, DM**

Prime Example: Simple Group Model

- ▶ enlarged gauge group: $SU(3) \times U(1)$; globally $U(3) \rightarrow U(2)$
- ▶ **Two** nonlinear Φ representations $\mathcal{L} = |D_\mu \Phi_1|^2 + |D_\mu \Phi_2|^2$

$$\Phi_{1/2} = \exp \left[\pm i \frac{f_{2/1}}{f_{1/2}} \Theta \right] \begin{pmatrix} 0 \\ 0 \\ f_{1/2} \end{pmatrix} \quad \Theta = \frac{1}{\sqrt{f_1^2 + f_2^2}} \begin{pmatrix} \eta & 0 & h^* \\ 0 & \eta & \\ h^T & & \eta \end{pmatrix}$$

Coleman-Weinberg mechanism: Radiative generation of potential

$$= \frac{g^2}{16\pi^2} \Lambda^2 (|\Phi_1|^2 + |\Phi_2|^2) \sim \frac{g^2}{16\pi^2} f^2$$

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but:

$$\Phi_1^\dagger \text{---} \text{---} \text{---} \Phi_2^\dagger = \frac{g^4}{16\pi^2} \log \left(\frac{\Lambda^2}{\mu^2} \right) |\Phi_1^\dagger \Phi_2|^2 \Rightarrow \frac{g^4}{16\pi^2} \log \left(\frac{\Lambda^2}{\mu^2} \right) f^2 (h^\dagger h)$$

Cancellations of Divergencies in Yukawa sector

$$\propto \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2(k^2 - m_T^2)} \left\{ \lambda_t^2 (k^2 - m_T^2) + k^2 \lambda_T^2 - \frac{m_T}{F} \lambda_T k^2 \right\}$$

Little Higgs global symmetry imposes relation

$$\frac{m_T}{F} = \frac{\lambda_t^2 + \lambda_T^2}{\lambda_T}$$

 \Rightarrow

Quadratic divergence cancels

Collective Symm. breaking: $\lambda_t \propto \lambda_1 \lambda_2$, $\lambda_1 = 0$
 or $\lambda_2 = 0 \Rightarrow SU(3) \rightarrow [SU(3)]^2$

$$\sim \frac{\lambda_1^2 \lambda_2^2}{16\pi^2} \log \left(\frac{\Lambda^2}{\mu^2} \right) |\Phi_1^\dagger \Phi_2|^2$$

Constraints from Oblique Corrections: S, T, U



$$\Delta T \sim \Delta \rho \sim \Delta M_Z^2 Z \cdot Z$$



$$\Delta S \sim W^0_{\mu\nu} B^{\mu\nu}, \Delta U \sim W^0_{\mu\nu} W^{0\mu\nu}$$

◇ All low-energy effects order v^2/F^2 (Wilson coefficients)

$\Delta S, \Delta T$ in the Littlest Higgs model, violation of **Custodial SU(2)**: Csáki et al., 2002; Hewett et al., 2002; Han et al., 2003; Chen/Dawson, 2003; Kilian/JRR, 2003

$$\frac{\Delta S}{8\pi} = - \left[\frac{c^2(c^2-s^2)}{g^2} + 5 \frac{c'^2(c'^2-s'^2)}{g'^2} \right] \frac{v^2}{F^2} \rightarrow 0 \quad \alpha \Delta T \rightarrow \frac{5}{4} \frac{v^2}{F^2} - \frac{2v^2 \lambda_{2\phi}^2}{M_\phi^4} \gtrsim \frac{v^2}{F^2}$$

Constraints from **contact IA**: $(f_{JJ}^{(3)}, f_{JJ}^{(1)})$ $4.5 \text{ TeV} \lesssim F/c^2$ $10 \text{ TeV} \lesssim F/c'^2$

◇ **Constraints evaded** $\iff c, c' \ll 1$
 B', Z', W'^{\pm} superheavy ($\mathcal{O}(\Lambda)$) decouple from fermions

Motivation

How to constrain a generic model in *HEP*?

- ▶ direct searches of resonances
- ▶ electroweak precision tests
- ▶ flavour constraints
- ▶ nowadays: Higgs sector

Higgs sector is the key to understand EW-scale physics (and beyond?)

Statistical analysis

We considered the three most popular Little Higgs models:

- ▶ Simplest Little Higgs (*SLH*) [Schmaltz]
- ▶ Littlest Higgs (*L²H*) [Arkani-Hamed et al.]
- ▶ Littlest Higgs with *T*-parity (*LHT*) [Low et al.]

and realized a χ^2 analysis on their parameter spaces, taking into account the whole set of 7+8 TeV Higgs searches by *ATLAS* and *CMS*, and by fitting 21 different *EW* Precision Observables:

$$\chi^2 = \sum_i \frac{(\mathcal{O}_i - \mathcal{O}_i^{\text{exp}})^2}{\sigma_i^2}$$

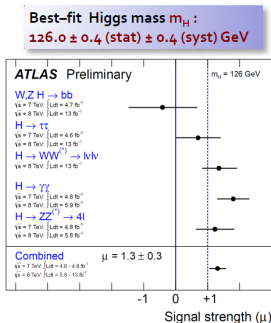
where \mathcal{O}_i depends on the free parameters of the model considered.

Data used: Higgs sector

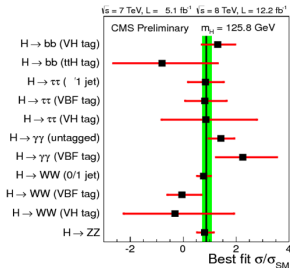
the Higgs results are expressed in terms of a *signal strength modifier*

$$\mu_i = \frac{\sum_p \epsilon^p_i \sigma_p}{\sum_p \epsilon^p_i \sigma_p^{SM}} \cdot \frac{BR(h \rightarrow X_i X_i)}{BR(h \rightarrow X_i X_i)_{SM}}$$

we included in our χ^2 analysis the best-fit values of μ_i reported by the Collaborations for all the different 7+8 TeV channels i :



- $M = 125.8 \pm 0.4$ (stat) ± 0.4 (syst) GeV



- $\sigma/\sigma_{SM} = 0.88 \pm 0.21$

Data used: *EWPD*

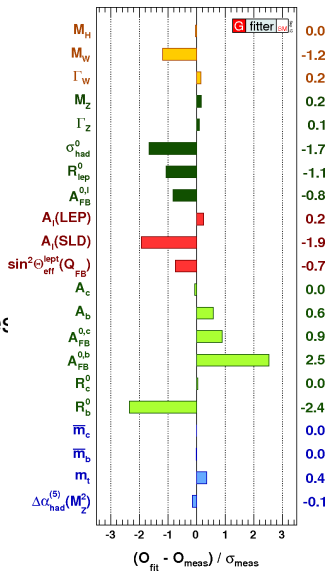
every extension of *SM* has to satisfy at least the precision constraints of the electroweak sector:

- ▶ low-energy observables

e.g. ν -scattering, parity violation observables

- ▶ Z -pole observables

e.g. m_Z , Γ_Z , Z -pole asymmetries...



LH Smoking guns

Where do the *LH* corrections to the *SM* quantities come from?

- ▶ new decay channels of the Higgs, e.g. $h \rightarrow A_H A_H$ in *LHT*
- ▶ modified Higgs couplings with *SM* fermions and vector bosons

$$\text{e.g. } 2 \frac{m_W^2}{v} y_W h W^+ W^-, \quad y_W = \begin{cases} 1 & \text{SM} \\ 1 + \mathcal{O}(v^2/f^2) & \text{LH} \end{cases}$$

- ▶ interaction terms of Higgs with new fermions/vector bosons

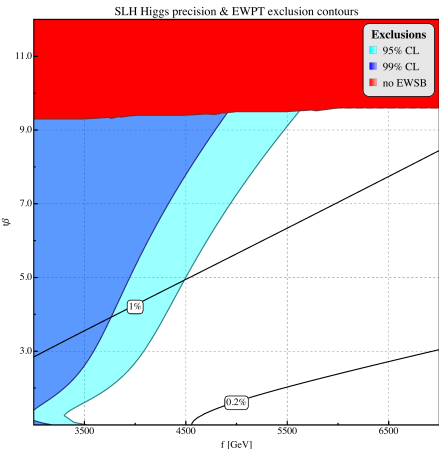
$$\text{e.g. } \frac{m_T}{v} y_T h \bar{T} T \quad m_T \sim f, y_T \sim \mathcal{O}(v^2/f^2)$$

- ▶ modified neutral- and charged-currents

$$\text{e.g. } \frac{g}{c_W} \sum_f \bar{f} \gamma^\mu \left((g_L^{SM} + \delta g_L) P_L + (g_R^{SM} + \delta g_R) P_R \right) f Z_\mu$$

SLH results

JRR/Tonini, JHEP 1302 (2013) 077; JRR/Tonini/de Vries, JHEP 1402 (2014) 053



$$\chi_{\min}^2/\text{d.o.f.} = 1.043$$

$$\chi_{\text{SM}}^2/\text{d.o.f.} = 1.048$$

- ▶ free parameters: f SSB scale, t_β ratio of vevs of scalar fields $\phi_{1,2}$
- ▶ $f_{\min}^{99\%} = 2.88$ TeV, translates into lower bounds on new states' masses, e.g.

$$m_{W'} \gtrsim 1.35 \text{ TeV}$$

$$m_T \gtrsim 2.81 \text{ TeV}$$

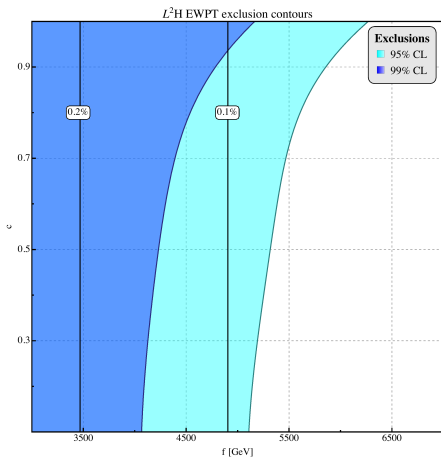
- ▶ min. required fine tuning: $\sim 1\%$, defined as

$$\Delta = \frac{|\delta\mu^2|}{\mu_{\text{obs}}^2}$$

- ▶ results mainly driven by *EWPD*
- ▶ includes data from Moriond 2013

L^2H results

JRR/Tonini, JHEP 1302 (2013) 077; JRR/Tonini/de Vries, JHEP 1402 (2014) 053



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- ▶ free parameters: f SSB scale, c mixing angle in gauge sector
- ▶ $f_{\min}^{99\%} = 3.20$ TeV, translates into lower bounds on new states' masses, e.g.

$$m_{W'} \gtrsim 2.13 \text{ TeV}$$

$$m_T \gtrsim 4.50 \text{ TeV}$$

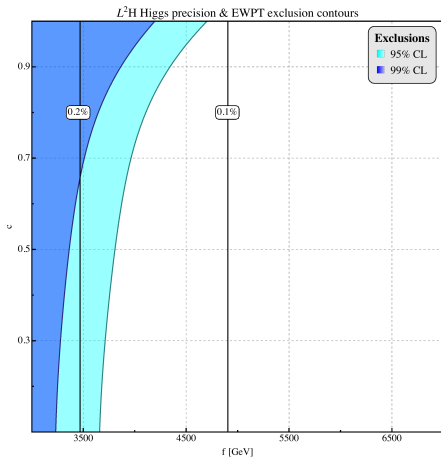
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- ▶ Exclusion gets weaker by Higgs data (d.o.f.)!

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Partial decay widths in LH

- ▶ 1-loop decays

$$\Gamma(h \rightarrow gg)_{LH} \sim \frac{\alpha_s^2 m_h^3}{32\pi^3 v^2} \left| \sum_{f, \text{col}} -\frac{1}{2} F_{\frac{1}{2}}(x_f) y_f \right|^2$$

$$\Gamma(h \rightarrow \gamma\gamma)_{LH} \sim \frac{\alpha^2 m_h^2}{256\pi^3 v^2} \left| \sum_{f, \text{ch}} \frac{4}{2} F_{\frac{1}{2}}(x_f) y_f + \sum_{v, \text{ch}} F_1(x_v) y_v + \sum_{s, \text{ch}} F_0(x_s) y_s \right|^2$$

where $x_i = \frac{4m_i^2}{m_h^2}$, $F_i(x_i)$ are loop functions, y_i the modified Yuk. coupl.

$$\Rightarrow \text{ narrow-width approximation: } \frac{\sigma_{LH}}{\sigma_{SM}}(gg \rightarrow h) = \frac{\Gamma(h \rightarrow gg)_{LH}}{\Gamma(h \rightarrow gg)_{SM}}$$

- ▶ tree-level decays

$$\Gamma(h \rightarrow VV)_{LH} \sim \Gamma(h \rightarrow VV)_{SM} \left(\frac{g_{hVV}}{g_{hVV}^{SM}} \right)^2$$

$$\Gamma(h \rightarrow f\bar{f})_{LH} \sim \Gamma(h \rightarrow f\bar{f})_{SM} \left(\frac{g_{hff}}{g_{hff}^{SM}} \right)^2$$

where $g_{hVV} = \frac{m_V^2}{v} y_V$ and $g_{hff} = \frac{m_f}{v} y_f$

LHT: Littlest Higgs with T parity

- ▶ Goldstone boson matrix:

$$\Sigma = e^{2i\Pi/f} \quad \Pi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & H & \sqrt{2}\Phi \\ H^\dagger & 0 & H^t \\ \sqrt{2}\Phi^\dagger & H^* & 0 \end{pmatrix} \quad \Phi \propto \begin{pmatrix} \sqrt{2}\phi^{++} & \phi^+ \\ \phi^+ & \phi^0 + i\phi^P \end{pmatrix}$$

- ▶ Discrete T parity:

$$T: \quad \Pi \rightarrow -\Omega \Pi \Omega \quad \Omega = \text{diag}(1, 1, -1, 1, 1)$$

$$V_{CW} = \lambda_{\phi^2} f^2 \text{Tr}(\phi^\dagger \phi) + i\lambda_{h\phi h} f (H\phi^\dagger H^t - H^* \phi H^\dagger) - \mu^2 H H^\dagger + \lambda_{h^4} (H H^\dagger)^2 + \\ + \lambda_{h\phi\phi h} H\phi^\dagger \phi H^\dagger + \lambda_{h^2\phi^2} H H^\dagger \text{Tr}(\phi^\dagger \phi) + \lambda_{\phi^2\phi^2} [\text{Tr}(\phi^\dagger \phi)]^2 + \lambda_{\phi^4} \text{Tr}(\phi^\dagger \phi \phi^\dagger \phi).$$

$$\lambda_{\phi^2} = 2(g^2 + g'^2) + 8\lambda_1^2 \quad \lambda_{h^4} = \frac{1}{4}\lambda_{\phi^2} \quad \lambda_{h\phi\phi h} = -\frac{4}{3}\lambda_{\phi^2} \\ \lambda_{h^2\phi^2} = -16\lambda_1^2 \quad \lambda_{\phi^4} = -\frac{8}{3}(g^2 + g'^2) + \frac{16}{3}\lambda_1^2$$

- ▶ Yukawa couplings $k, R \equiv \lambda_1/\lambda_2$

$$\mathcal{L}_k = -kf \left(\bar{\Psi}_2 \xi \Psi_c + \bar{\Psi}_1 \langle \Sigma \rangle \Omega \xi^\dagger \Omega \Psi_c \right) - m_q \bar{u}'_c u_c - m_q \bar{d}'_c d_c - m_\chi \bar{\chi}'_c \chi_c + \text{h.c.}$$

$$\mathcal{L}_t = -\frac{\lambda_1 f}{2\sqrt{2}} \epsilon_{ijk} \epsilon_{xy} \left[(\bar{\Psi}_{1,t})_i \Sigma_{jx} \Sigma_{ky} - (\bar{\Psi}_{2,t} \langle \Sigma \rangle)_i \Sigma'_{jx} \Sigma'_{ky} \right] t'_R - \lambda_2 f (\bar{T}_{L1} T_{R1} + \bar{T}_{L2} T_{R2})$$

T parity and Dark Matter

Cheng/Low, 2003; Hubisz/Meade, 2005; Wang/Yang/Zhu,

2013

- ▶ **T parity**: $T^a \rightarrow T^a$, $X^a \rightarrow -X^a$, automorphism of coset space analogous to R parity in SUSY, KK parity in extra dimensions
- ▶ Bounds on F MUCH relaxed, $F \sim 0.5 - 1 \text{ TeV}$
but: Pair production!, typical **cascade decays**
- ▶ Lightest T -odd particle (LTP) \Rightarrow **Candidate for Cold Dark Matter**

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Littlest Higgs: A' LTP

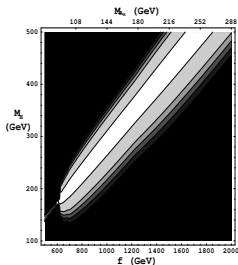
$W', Z' \sim 650 \text{ GeV}$, $\Phi \sim 1 \text{ TeV}$

$T, T' \sim 0.7-1 \text{ TeV}$

Annihilation: $A'A' \rightarrow h \rightarrow WW, ZZ, hh$

Hubisz/Meade, 2005

0/10/50/70/100



T parity and Dark Matter

Cheng/Low, 2003; Hubisz/Meade, 2005; Wang/Yang/Zhu,

2013

- ▶ T parity: $T^a \rightarrow T^a$, $X^a \rightarrow -X^a$, automorphism of coset space analogous to R parity in SUSY, KK parity in extra dimensions
- ▶ Bounds on F MUCH relaxed, $F \sim 0.5 - 1$ TeV
but: Pair production!, typical **cascade decays**
- ▶ Lightest T -odd particle (LTP) \Rightarrow **Candidate for Cold Dark Matter**

Littlest Higgs: A' LTP

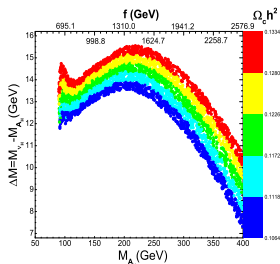
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Wang/Yang/Zhu, 2013

Relic density/SI cross section



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Wang/Yang/Zhu, 2013

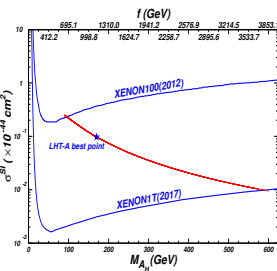
Relic density/SI cross section

- ▶ T parity Simplest LH: **Pseudo-Axion η LTP**

Z' remains odd: good or bad (?)

Martin, 2006; JRR/Tonini, in prep.

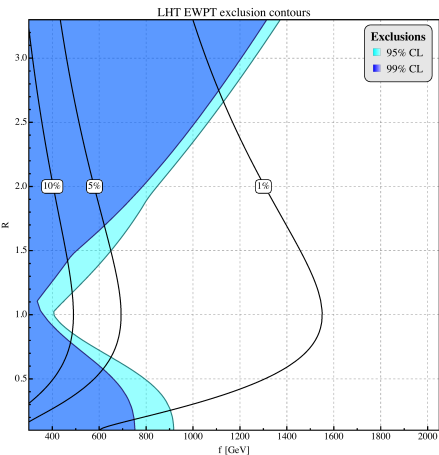
- ▶ T parity might be anomalous (???)



Hill/Hill, 2007

LHT results

JRR/Tonini, JHEP 1302 (2013) 077; JRR/Tonini/de Vries, JHEP 1402 (2014) 053



$$\chi_{\min}^2/\text{d.o.f.} = 1.048$$

$$\chi_{\text{SM}}^2/\text{d.o.f.} = 1.053$$

- ▶ free parameters: f SSB scale, R ratio of Yukawa couplings in top sector
- ▶ $f_{\min}^{99\%} = 405.9$ GeV, translates into lower bounds on new states' masses, e.g.

$$m_{W'} \gtrsim 269.6 \text{ GeV}$$

$$m_T \gtrsim 553.6 \text{ GeV}$$

- ▶ min. required fine tuning: $\sim 10\%$, defined as

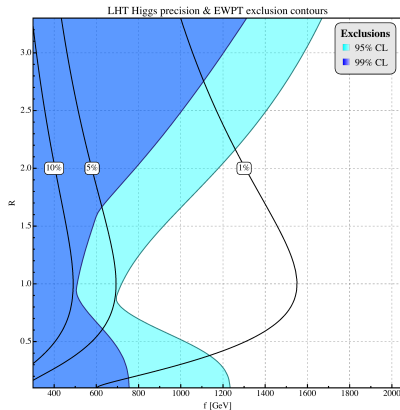
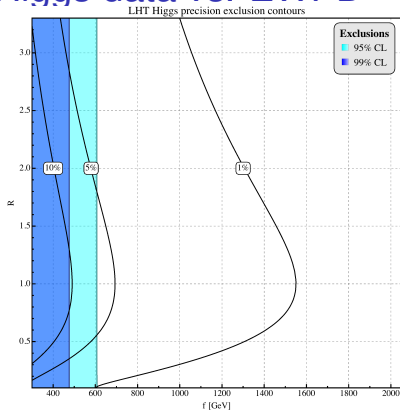
$$\Delta = \frac{|\delta\mu^2|}{\mu_{\text{obs}}^2}$$

- ▶ results mainly driven by EWPD (see next slide)

EWPT \Rightarrow

$$f \gtrsim 405 \text{ GeV}$$

Higgs data vs. *EWPD*

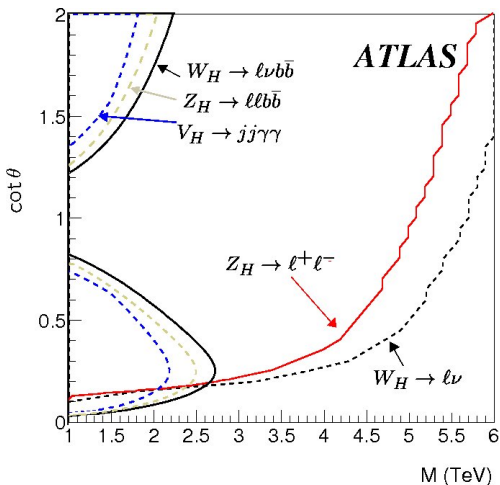


- ▶ the shape of result driven by *EW* constraints (much smaller uncertainties)
- ▶ Higgs data only: for $v/f \gtrsim 0.6$ decay $h \rightarrow A_H A_H$ open and dominant
- ▶ Higgs data only: subdominant dependence on R w.r.t. f is a consequence of the Collective Symmetry Breaking mechanism

EWPT and Higgs data \Rightarrow

$$f \gtrsim 694 \text{ GeV}$$

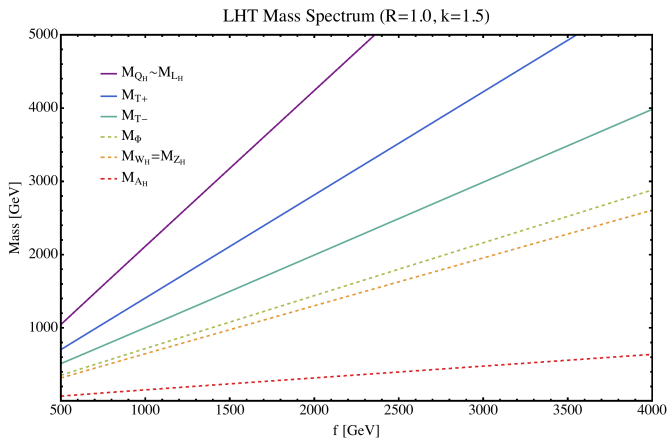
Direct searches: Drell-Yan mainly



Reach in the gauge boson sector: depends on mixing angle

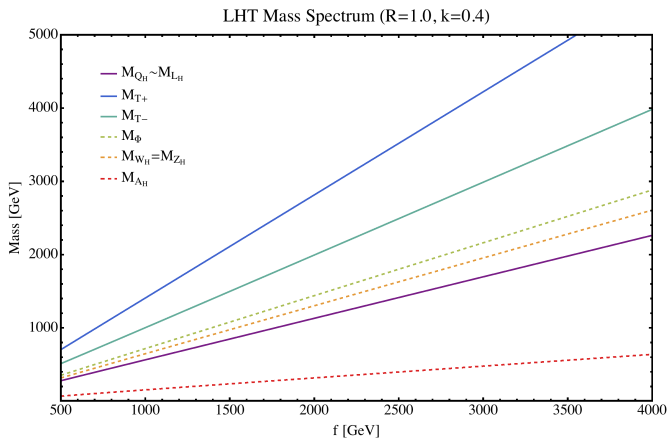
Direct Searches: Focus on LHT

- ▶ Defining two benchmark scenarios: 1. heavy quarks



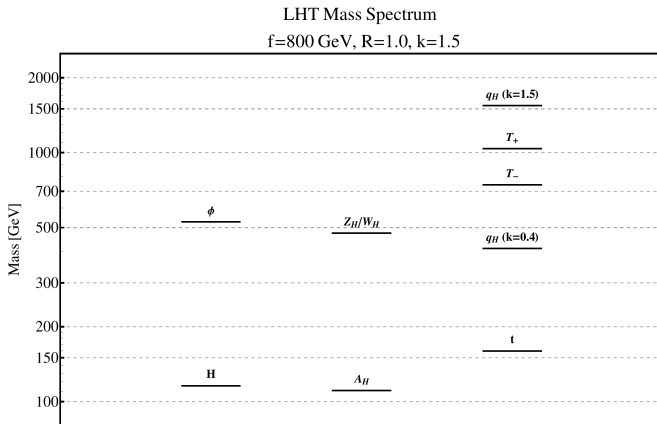
Direct Searches: Focus on LHT

- ▶ Defining two benchmark scenarios: 2. heavy top/vectors



Direct Searches: Focus on LHT

- Defining two benchmark scenarios: 1. $k = 1.5$, 2. $k = 0.4$



Branching Ratios

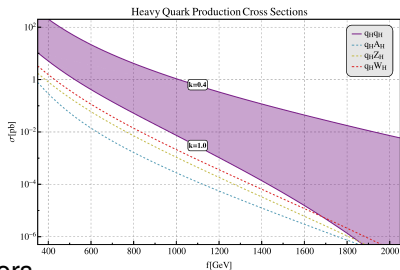
► Decay patterns:

Particle	Decay	BR _{k=1.0}	BR _{k=0.4}
l_H^\pm	$W_H^\pm \nu$	62%	0%
	$Z_H l^\pm$	31%	0%
	$A_H l^\pm$	6%	100%
ν_H^\pm	$W_H^\pm l^\mp$	61%	0%
	$Z_H \nu$	30%	0%
	$A_H \nu$	9%	100%
T_H^+	$W^+ b$	46%	45%
	$Z t$	22%	22%
	$H t$	21%	21%
	$T_H^- A_H$	11%	11%
A_H	stable		
Z_H	$A_H H$	100%	2%
	$d_H d$	0%	41%
	$u_H u$	0%	30%
	$l_H^\pm l^\mp$	0%	14%
	$\nu_H \nu$	0%	14%

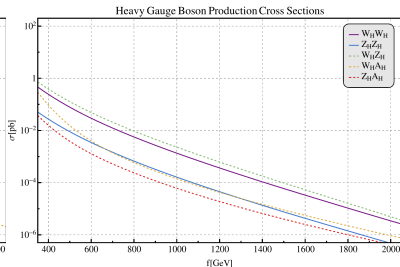
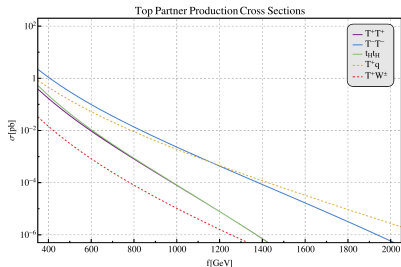
Particle	Decay	BR _{k=1.0}	BR _{k=0.4}
d_H	$W_H^- u$	62%	0%
	$Z_H d$	30%	0%
	$A_H d$	6%	100%
u_H	$W_H^+ d$	58%	0%
	$Z_H u$	30%	0%
	$A_H u$	9%	100%
T_H^-	$A_H t$	100%	100%
	$Z_H t$	0%	0%
$\Phi^{0/P}$	$A_H H$	100%	100%
Φ^\pm	$A_H W^\pm$	100%	100%
$\Phi^{\pm\pm}$	$A_H (W^\pm)^2$	100%	96%
W_H^\pm	$A_H W^\pm$	100%	2%
	$u_H d$	0%	44%
	$d_H u$	0%	27%
	$l_H^\pm \nu$	0%	16.5%
	$\nu_H l^\pm$	0%	16.5%

Cross Sections (I)

► Heavy Quarks



► Heavy Top and Vectors



Channels and signatures: Parameters

final state			modes	params	final state			modes	params
leptons	# jets	\cancel{E}_T			leptons	# jets	\cancel{E}_T		
0	1	✓	$q_H A_H$	f, k	l^\pm	2	✓	$W_H^\pm W_H^\mp$ $W_H^\pm Z_H$ $q_H q_H$	f, k f, k f, k
0	2	✓	$q_H q_H$	f, k	l^\pm	3	✓	$q_H W_H^\pm$ $T^+ q$	f, k f, k, R
0	3	✓	$q_H W_H^\pm$	f, k	l^\pm	4	✓	$q_H q_H$ $T^- T^-$	f, k f, k, R
0	4	✓	$q_H q_H$ $W_H^\pm W_H^\mp$ $W_H^\pm Z_H$ $Z_H Z_H$	f, k f, k f, k f, k	$l^+ l^-$	0	✓	$W_H^\pm W_H^\mp$	f, k
0	4	✗	$T^+ q$	f, k, R	$l^+ l^-$	1	✓	$q_H W_H^\pm$	f, k
0	5	✓	$q_H W_H^\pm$	f, k	$l^+ l^-$	2	✓	$q_H q_H$ $T^- T^-$	f, k f, k, R
0	6	✓	$q_H q_H$ $T^- T^-$	f, k f, k, R	$l^\pm l^\pm$	2	✓	$q_H q_H$	f, k

Channels and signatures (I)

final state			production modes	$\sigma_{8\text{ TeV}} \times \text{Br (fb)}$		$\sigma_{14\text{ TeV}} \times \text{Br (fb)}$	
$\# l^\pm$	$\# \text{ jets}$	\cancel{E}_T		$k = 1.0$	$k = 0.4$	$k = 1.0$	$k = 0.4$
0	1	✓	$q_H A_H$	0.24	1.1×10^2	2.1	4.5×10^2
0	2	✓	$q_H q_H$	0.56	5.6×10^3	5.2	3.2×10^4
0	3	✓	$q_H W_H^\pm$	0.73	14	8.0	77
			$q_H Z_H$	0.76	8.6	8.0	49
0	4	✓	$q_H q_H$	4.0	9.1×10^2	35	5.6×10^3
			$W_H^\pm W_H^\mp$	1.9	low	9.1	low
			$W_H^\pm Z_H$	4.8	low	23	low
			$Z_H Z_H$	0.56	low	3.0	low
0	4	✗	$T^+ q$	2.0	2.0	17	17
0	5	✓	$q_H W_H^\pm$	5.1	✗	54	✗
			$q_H Z_H$	4.1	✗	44	✗
0	6	✓	$q_H q_H$	1.6	9.7×10^2	1.7×10^2	6.0×10^3
			$T^- T^-$	2.5	2.5	25	25

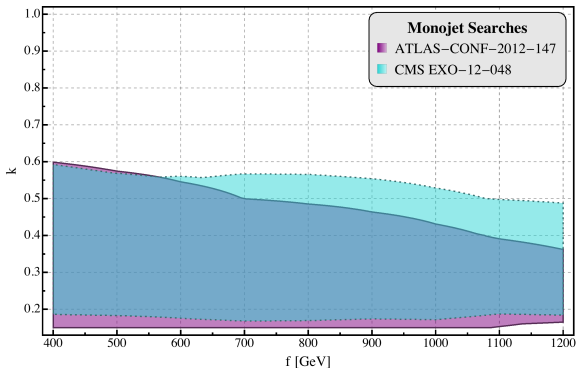
Channels and signatures (II)

final state			production modes	$\sigma_{8 \text{ TeV}} \times \text{Br (fb)}$		$\sigma_{14 \text{ TeV}} \times \text{Br (fb)}$	
# l^\pm	# jets	\cancel{E}_T		$k = 1.0$	$k = 0.4$	$k = 1.0$	$k = 0.4$
l^\pm	2	✓	$q_H q_H$	0.058	9.0×10^2	1.1	5.6×10^3
			$W_H^\pm W_H^\mp$	0.77	low	3.9	low
			$W_H^\pm Z_H$	2.1	low	10	low
			$T^+ q$	1.3	1.2	10	10
l^\pm	3	✓	$q_H W_H^\pm$	3.5	✗	37	✗
			$q_H Z_H$	0.99	✗	11	✗
l^\pm	4	✓	$q_H q_H$	7.4	9.7×10^2	82	6.0×10^3
			$T^- T^-$	2.2	2.2	21	21
$l^+ l^-$	0	✓	$W_H^\pm W_H^\mp$	0.32	low	1.7	low
$l^+ l^-$	1	✓	$q_H W_H^\pm$	0.54	✗	5.8	✗
$l^+ l^-$	2	✓	$q_H q_H$	1.1	✗	11	✗
			$T^- T^-$	0.47	0.47	4.6	4.6
$l^\pm l^\pm$	2	✓	$q_H q_H$	0.37	✗	2.7	✗

Recasting results

JRR/Tonini/deVries,2013

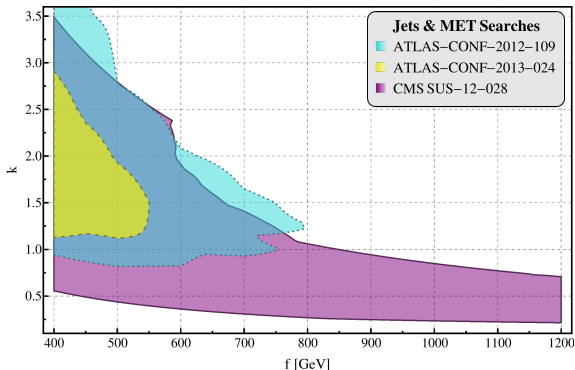
- **95% CL from Monojets + \cancel{E}_T from LHC8**
- 1 hard jet, \cancel{E}_T , no leptons, 2nd jet w. $p_T > 30$ GeV
signal regions: ATLAS (p_T, \cancel{E}_T) > 120/220/350/500 GeV, CMS: $\cancel{E}_T > 250/300/350/400/450/500/550$ GeV
- Dijet suppression: ATLAS $\Delta\phi(\cancel{E}_T, j_2) > 0.5$, CMS $\Delta\phi(j_1, j_2) < 2.5$
- $pp \rightarrow q_H q_H, pp \rightarrow q_H A_H$



Recasting results

JRR/Tonini/deVries,2013

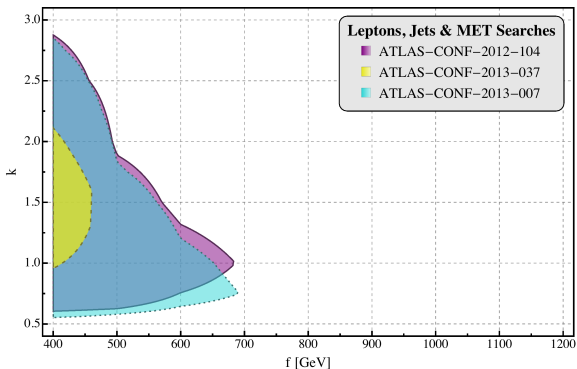
- **95% CL from Jets + \cancel{E}_T from LHC8**
- ≥ 2 hard jets, \cancel{E}_T , no leptons
- signal regions: ATLAS $\cancel{E}_T > 200/300/350$ GeV, CMS:
 $(N_j, N_b) = (2 - 3, 0); (2 - 3, 1 - 2); (\geq 4, 1 - 2); (\geq 4, 0); (\geq 4, \geq 2)$
- QCD suppression: ATLAS $\Delta\phi(\cancel{E}_T, j_2) > 0.5$, \cancel{E}_T/m_{eff} , CMS $\Delta\phi(j_1, j_2) < 2.5$
- $pp \rightarrow q_H q_H \rightarrow (j A_H)(j A_H)$



Recasting results

JRR/Tonini/deVries, 2013

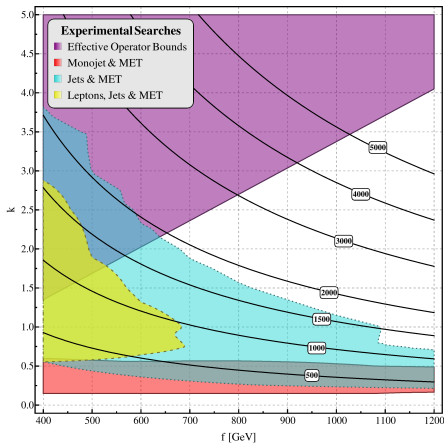
- 95% CL from Leptons + Jets + \cancel{E}_T from LHC8
- single isolated lepton, ≥ 2 hard jets, \cancel{E}_T ,
- signal regions: ATLAS $\cancel{E}_T > 200/300/350$ GeV
- Cuts: $\cancel{E}_T > 250$ GeV, $m_T(l, \cancel{E}_T) > 250$ GeV, $\cancel{E}_T/m_{\text{eff}} > 0.2$, $m_{\text{eff}}^{\text{inc}} > 800$ GeV
- $pp \rightarrow q_H q_H$ with $q_H \rightarrow W_H q, Z_H q, t_H \rightarrow t A_H, Z_H \rightarrow H A_H$



Combined analysis

JRR/Tonini/deVries, 2013

- ▶ Operator bounds: $\mathcal{O}_{4f} = -\frac{k^2}{128\pi^2 f^2} \bar{\psi}_L \gamma^\mu \psi_L \bar{\psi}'_L \gamma_\mu \psi'_L + \mathcal{O}\left(\frac{g}{k}\right)$
Hubisz/Meade/Noble/Perelstein, 2005



- ▶ Bound from combined analysis: $f \gtrsim 638\text{GeV}$

Conclusions

- ▶ *Little Higgs* models are an appealing solution to the hierarchy problem, alternative to weakly coupled solutions like *SUSY*
- ▶ most of the parameter space of three popular *Little Higgs* models is still compatible at $\sim 99\%$ *CL* with the early results of the 7+8 TeV Higgs searches
- ▶ electroweak precision data represent still the most severe constraints
- ▶ fine-tuning as a guideline to understand the naturalness of a model: *Little Higgs* models require a minimum level of $\sim 10\%$ of fine tuning
- ▶ Limits on the LHT:

1. EWPO:

$$f \gtrsim 405 \text{ GeV@95\% CL}$$

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- ▶ We need more data!

Lessons from Lepton Photon 2013 ...

There are either colored exotics ...



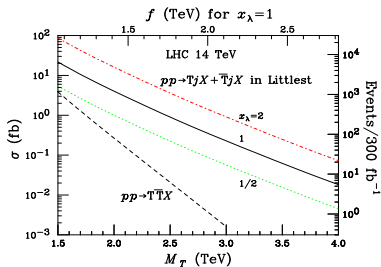
Lessons from Lepton Photon 2013 ...

... or the world is fine tuned



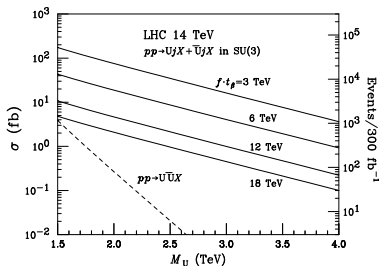
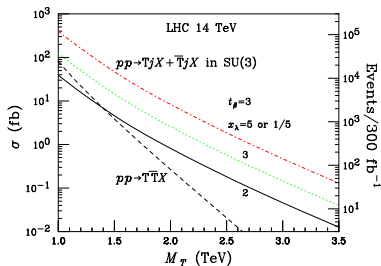
Direct Searches – Heavy Quark States

- EW single dominates QCD pair production: Perelstein/Peskin/Pierce, '03



Direct Searches – Heavy Quark States

- EW single dominates QCD pair production: Perelstein/Peskin/Pierce, '03



- Characteristic branching ratios :

$$\Gamma(T \rightarrow th) \approx \Gamma(T \rightarrow tZ) \approx \frac{1}{2} \Gamma(T \rightarrow bW^+) \approx \frac{M_T \lambda_T^2}{64\pi}, \quad \Gamma_T \sim 10-50 \text{ GeV}$$

- Proof of T as EW singlet; but: $T \rightarrow Z'T, W'b, t\eta$!

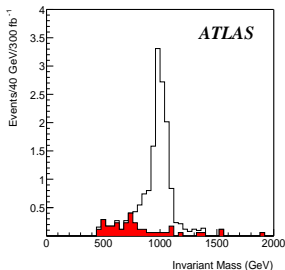
AIM: Determination of $M_T, \lambda_T, \lambda_{T'}$

$\lambda_{T'}$ indirect ($T\bar{T}h$ impossible)

$$T \rightarrow Zt \rightarrow \ell^+ \ell^- \ell \nu b$$

SN-ATLAS-2004-038

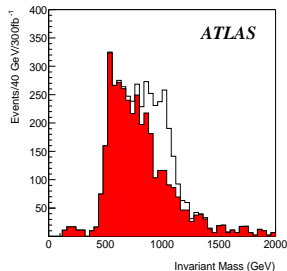
-
- ▶ $\cancel{E}_T > 100 \text{ GeV}$, $l\ell l$, $p_T > 100/30 \text{ GeV}$,
 b , $p_T > 30 \text{ GeV}$
 - ▶ Bkgd.: WZ , ZZ , btZ
 - ▶ Observation for $M_T \lesssim 1.4 \text{ TeV}$



$$T \rightarrow Wb \rightarrow \ell\nu b$$

 SN-ATLAS-2004-038

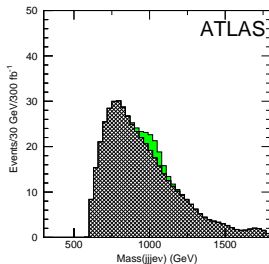
- ▶ $\cancel{E}_T > 100 \text{ GeV}$, $\ell, p_T > 100 \text{ GeV}$,
 $b, p_T > 200 \text{ GeV}$, max. $jj, p_T > 30 \text{ GeV}$
- ▶ Bkgd.: $t\bar{t}$, $Wb\bar{b}$, single t
- ▶ Observation for $M_T \lesssim 2.5 \text{ TeV}$



$T \rightarrow th \rightarrow \ell\nu bbb$

SN-ATLAS-2004-038

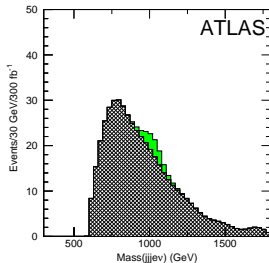
-
- ▶ $\ell, p_T > 100 \text{ GeV}$, $jjj, p_T > 130 \text{ GeV}$,
at least 1 b -tag
 - ▶ Bkgd.: $t\bar{t}$, $Wb\bar{b}$, single t
 - ▶ Observation for $M_T \lesssim 2.5 \text{ TeV}$



$$T \rightarrow th \rightarrow \ell\nu bbb$$

SN-ATLAS-2004-038

- ▶ $\ell, p_T > 100 \text{ GeV}$, $jjj, p_T > 130 \text{ GeV}$, at least 1 b -tag
- ▶ Bkgd.: $t\bar{t}$, $Wb\bar{b}$, single t
- ▶ Observation for $M_T \lesssim 2.5 \text{ TeV}$

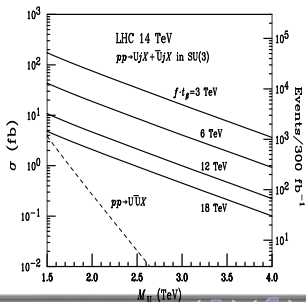


Additional heavy quarks (Simple Group Models): U, C or D, S

Han et al.,

05

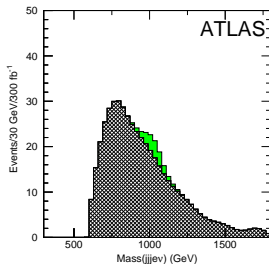
- ▶ Large cross section: u or d PDF
- ▶ Huge final state ℓ charge asymmetry
- ▶ Good mass reconstruction



$T \rightarrow th \rightarrow l\nu bbb$

SN-ATLAS-2004-038

- ▶ $\ell, p_T > 100 \text{ GeV}$, $jjj, p_T > 130 \text{ GeV}$,
at least 1 b -tag
- ▶ Bkgd.: $t\bar{t}$, $Wb\bar{b}$, single t
- ▶ Observation for $M_T \lesssim 2.5 \text{ TeV}$

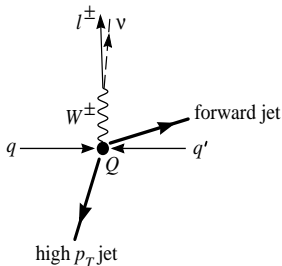


Additional heavy quarks (Simple Group Models): U, C or D, S

Han et al.,

05

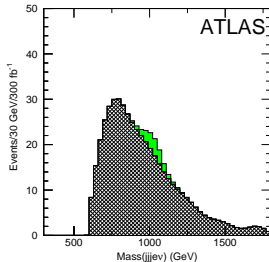
- ▶ Large cross section: u or d PDF
- ▶ Huge final state ℓ charge asymmetry
- ▶ Good mass reconstruction



$T \rightarrow th \rightarrow \ell\nu bbb$

SN-ATLAS-2004-038

- ▶ $\ell, p_T > 100 \text{ GeV}, jjj, p_T > 130 \text{ GeV}$, at least 1 b -tag
- ▶ Bkgd.: $t\bar{t}, Wb\bar{b}$, single t
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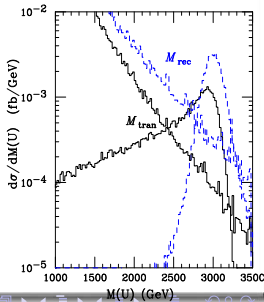
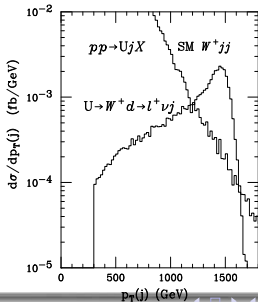


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Direct Searches – Heavy Vectors

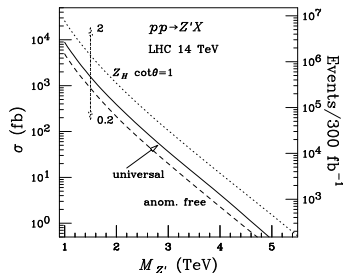
Drell-Yan Production: Tevatron Limits $\sim 500 - 600$ GeV

► **Dominant decays:**

Product group: $Z' \rightarrow Zh, WW,$

$W' \rightarrow Wh, WZ$

Simple group: $Z' \rightarrow qq, \quad X \rightarrow fF$



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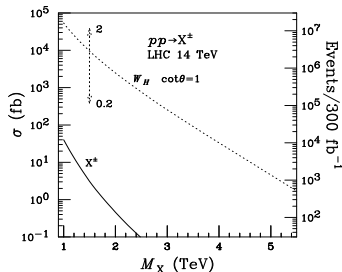
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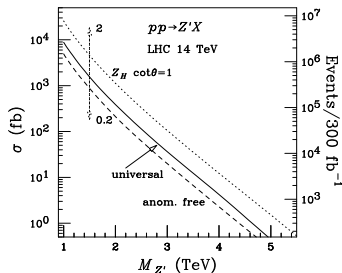
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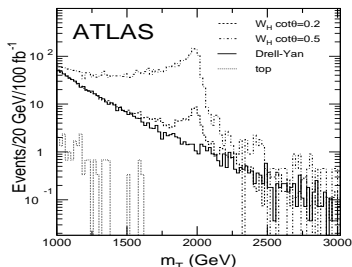
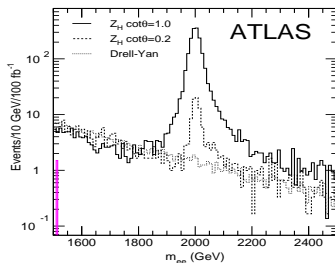
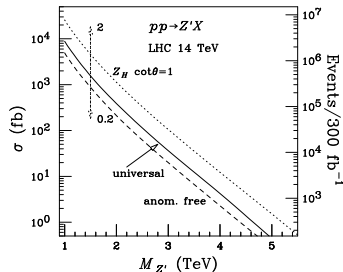
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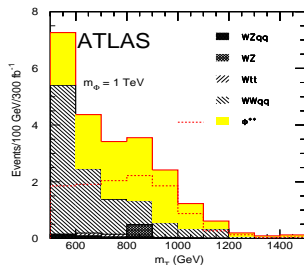


Direct Searches – Heavy Scalars

Generally: **Large model dependence**

no states complex singlet **complex triplet**

- ▶ **Littlest Higgs**, complex triplet:
 $\Phi^0, \Phi_P, \Phi^\pm, \Phi^{\pm\pm}$
- ▶ Cleanest channel: $q\bar{q} \rightarrow \Phi^{++}\Phi^{--} \rightarrow llll$:
 Killer: PS
- ▶ WW -Fusion: $dd \rightarrow uu\Phi^{++} \rightarrow uuW^+W^+$
- ▶ 2 hard forward jets, hard close l^+l^+
 p_T -unbalanced



Alternative: Model-Independent search in WW fusion:

ILC: Beyer/Kilian/Krstonosic/Mönig/JRR/Schmidt/Schröder, 2006

LHC: Alboteanu/Kilian/JRR, 2008; Kilian/JRR/Sekulla, 2013

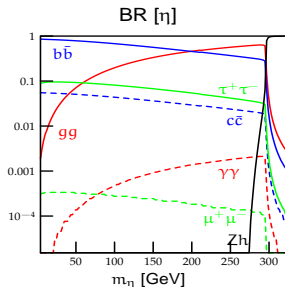
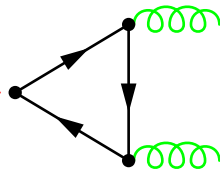
Pseudo-Axions in Little Higgs

Kilian/Rainwater/JRR, 2004, 2006; JRR,

2007

- gauged $U(1)$ group: Z' \longleftrightarrow ungauged: η
- couples to fermions like a pseudoscalar
- $m_\eta \lesssim 400$ GeV
- SM singlet, couplings to SM particles v/F suppressed
- η axion-like particle:

Anomalous $U(1)$: $---$



$$\longrightarrow \frac{1}{F} \frac{\alpha_s}{8\pi^2} \eta F_{\mu\nu} F_{\rho\sigma} \epsilon^{\mu\nu\rho\sigma}$$

- $U(1)$ explicitly broken \Rightarrow Axion limits from astroparticle physics not applicable

Classification of Axions in Little Higgs Models

Number of Pseudo-Axions: $n = g - l$

Mismatch between global (g) and local rank reduction (l)

Product Group Models

Arkani-Hamed, ...

- ▶ Doubling of electroweak gauge group: $SU(2) \times SU(2) \rightarrow SU(2)_L$,
 $U(1) \times U(1) \rightarrow U(1)_Y$ (latter not necessary) $\Rightarrow l = 1$
 - ▶ Littlest Higgs, $g: SU(5) \rightarrow SO(5) \Rightarrow n = (4 - 2) - 1 = 1$
 - ▶ antisymmetric, $g: Sp(6)/SO(6)$, $n = (3 - 2) - 1 = 0$

Simple Group Models

Kaplan, Schmaltz, ...

- ▶ Simple gauge group: $SU(N) \times U(1) \rightarrow SU(2) \times U(1) \Rightarrow l = N - 2$
- ▶ Higgs is distributed over several global symmetry multiplets
- ▶ Simplest Little Higgs, $g: [SU(3)]^2/[SU(2)]^2$ $n = g - l = 2 - 1 = 1$
- ▶ Original Simple Group Model, $g: [SU(4)]^3/[SU(3)^3 \times SU(2)]$,
 $l: SU(4) \rightarrow SU(2)$ $n = g - l = 4 - 2 = 2$

Moose Models

Arkani-Hamed, ...

- ▶ "Minimal" Moose: $g [SU(3)]^4 \rightarrow SU(3)$, $l [SU(3) \times SU(2)]/SU(2)$
 $n = g - l = 6 - 2 = 4$
- ▶ 3-site model: $g [SU(2)]^4/[SU(2)]^2$, $l [SU(2)]^2 \rightarrow SU(2)$, $n = 2 - 1 = 1$

$ZH\eta$ coupling as a discriminator

Kilian/Rainwater/JRR, 2006

- ▶ pseudo-axion: $\xi = \exp[i\eta/F]$, $\Sigma = \exp[i\Pi/F]$ non-linear representation of the remaining Goldstone multiplet Π

$$\mathcal{L}_{\text{kin.}} \sim F^2 \text{Tr} [(D^\mu(\xi\Sigma))^\dagger (D_\mu(\xi\Sigma))] = \dots - 2F(\partial_\mu\eta) \text{Im Tr} [(D^\mu\Sigma)^\dagger \Sigma] + O(\eta^2)$$

- ▶ Use special structure of covariant derivatives:

$$D_\mu \Sigma = \partial_\mu \Sigma + A_{1,\mu}^a (T_1^a \Sigma + \Sigma (T_1^a)^T) + A_{2,\mu}^a (T_2^a \Sigma + \Sigma (T_2^a)^T),$$

$$\text{Tr} [(D^\mu \Sigma)^\dagger \Sigma] \sim W_\mu^a \text{Tr} [\Sigma^\dagger (T_1^a + T_2^a) \Sigma + (T_1^a + T_2^a)^*] = 0.$$

- ▶ Little Higgs mechanism cancels this coupling
- ▶ Simple Group Models: $\Phi = \exp[i\Sigma/F]$, $\zeta = (0, \dots, 0, F)^T$ VEV directing in the N direction

$$\begin{aligned}\mathcal{L}_{\text{kin.}} &\sim F^2 D^\mu (\zeta^\dagger \Phi^\dagger) D_\mu (\Phi \zeta) = \dots + \frac{i}{F} (\partial_\mu \eta) \zeta^\dagger (\Phi^\dagger (D_\mu \Phi) - (D_\mu \Phi^\dagger) \Phi) \zeta \\ &= \dots + iF (\partial_\mu \eta) (\Phi^\dagger (D_\mu \Phi) - (D_\mu \Phi^\dagger) \Phi)_{N,N} .\end{aligned}$$

$$\Sigma = \begin{pmatrix} 0 & h \\ h^\dagger & 0 \end{pmatrix}, \quad \mathbb{V}_\mu = \begin{pmatrix} \mathbb{W}_\mu & 0 \\ 0 & 0 \end{pmatrix} + \text{heavy vector fields}$$

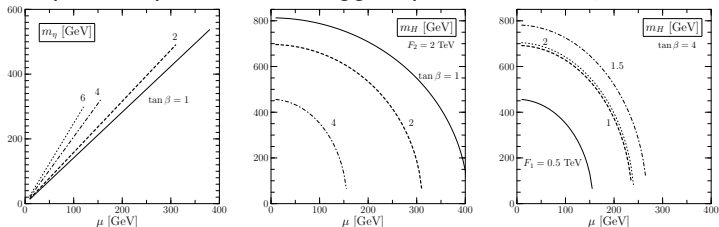
$$\begin{aligned}&\mathbb{V}_\mu + \frac{i}{F} [\Sigma, \mathbb{V}_\mu] - \frac{1}{2F^2} [\Sigma, [\Sigma, \mathbb{V}_\mu]] + \dots \\ &= \begin{pmatrix} \mathbb{W}_\mu & 0 \\ 0 & 0 \end{pmatrix} + \frac{i}{F} \begin{pmatrix} 0 & -\mathbb{W}_\mu h \\ h^\dagger \mathbb{W}_\mu & 0 \end{pmatrix} - \frac{1}{2F^2} \begin{pmatrix} hh^\dagger \mathbb{W} + \mathbb{W} h h^\dagger & 0 \\ 0 & -2h^\dagger \mathbb{W} h \end{pmatrix} + \dots\end{aligned}$$

- ▶ 1st term cancels by multiple Goldstone multiplets
- ▶ 2nd term cancels by EW symmetry
- ▶ 3rd term

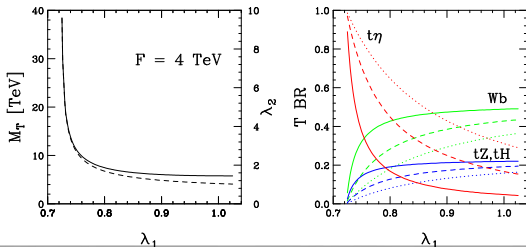
$$(\partial^\mu \eta) h^\dagger \mathbb{W}_\mu h \sim v H Z_\mu \partial^\mu \eta .$$

More properties of Pseudo-Axions

- ▶ Take e.g. one specific model: Simplest Little Higgs Schmaltz, 2004
- ▶ Simple Group Model, two Higgs-triplets with a $\tan\beta$ -like mixing angle



- ▶ $\tan\beta \sim 1$: heavy Higgs, (very) light pseudoscalar
- ▶ Heavy top decays: Kilian/Rainwater/JRR, 2006



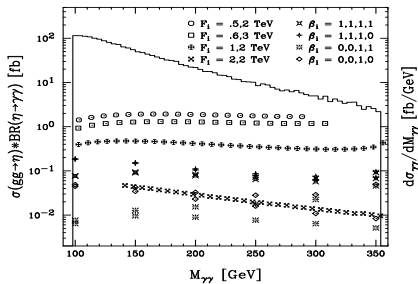
Discovery of Pseudo-axions

Kilian/Rainwater/JRR, 2004, 2006

LHC: Gluon fusion, diphoton
signal for $m_\eta \gtrsim 200$ GeV, 7σ
possible

LHC: $T \rightarrow t\eta$

ILC: $e^+e^- \rightarrow t\bar{t}\eta$



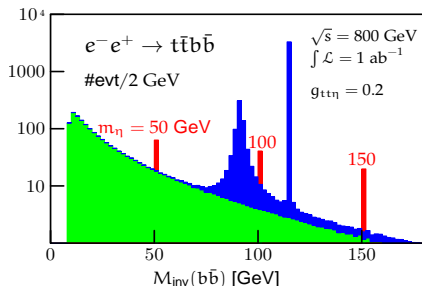
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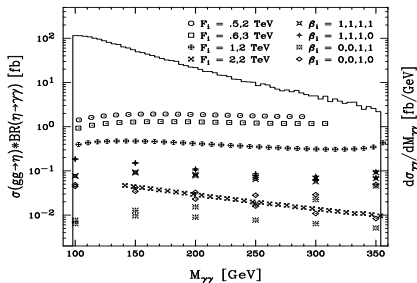
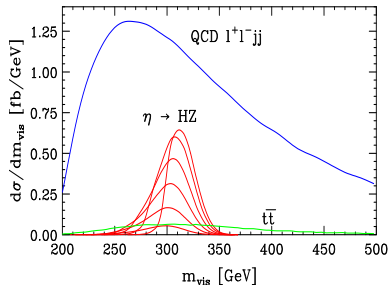
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$ZH\eta$ coupling

forbidden in Product Group Models

Discriminator of diff. model classes

$$gg \rightarrow \left\{ \begin{array}{ll} H \rightarrow Z\eta & \rightarrow llbb \\ \eta \rightarrow ZH & \rightarrow llbb, llljj \end{array} \right\}$$

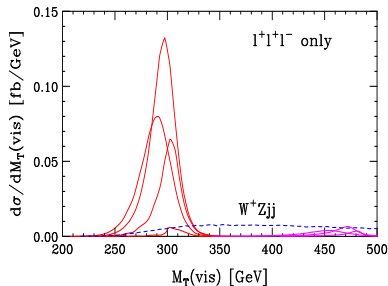
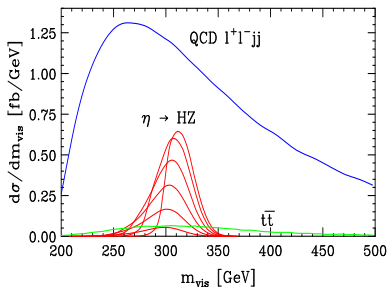
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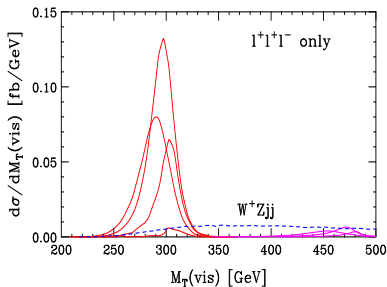
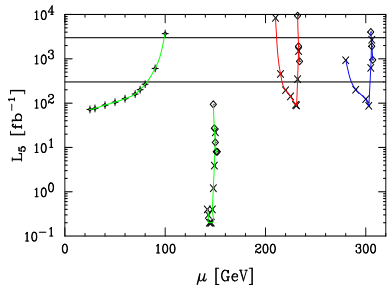
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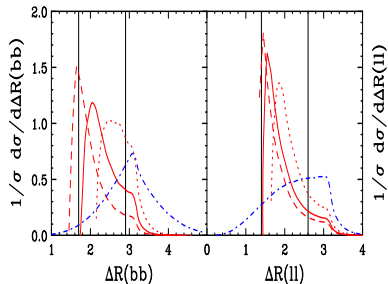
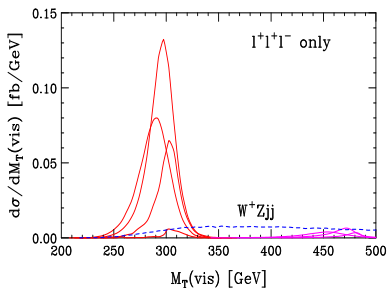
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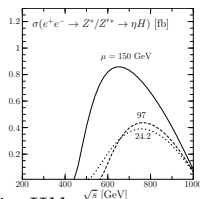
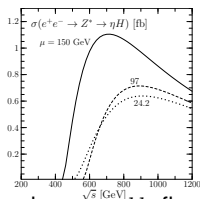
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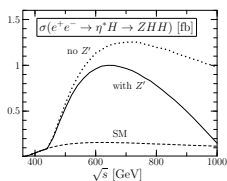
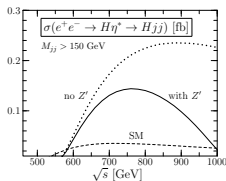
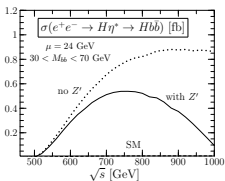
η pheno at ILC

Kilian/Rainwater/JRR, 2006

If $ZH\eta$ coupling present: $H\eta$ production in analogy to HA :



- Light pseudoaxion, $\eta \rightarrow b\bar{b}$, final state $Hb\bar{b}$
- Intermediate range, $\eta \rightarrow g\bar{g}$, final state $Hj\bar{j}$
- $\eta \rightarrow ZH$: ZHH final state

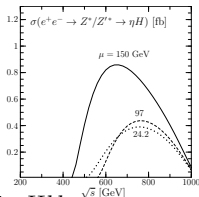
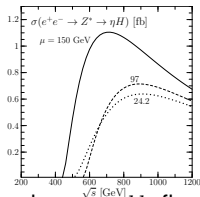


More detailed insights from photon collider option

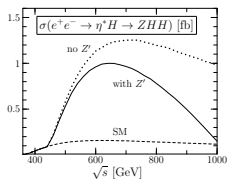
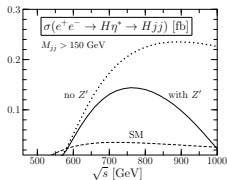
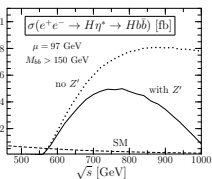
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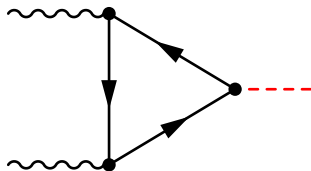
- Light pseudoaxion, $\eta \rightarrow bb$, final state Hbb
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- $\eta \rightarrow ZH$: ZHH final state



More detailed insights from photon collider option

Pseudo Axions at the Photon Collider

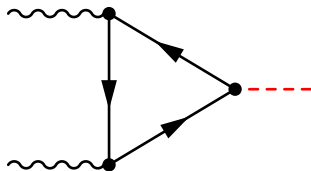
- ▶ **Photon Collider** as precision machine for Higgs physics (s channel resonance, anomaly coupling)



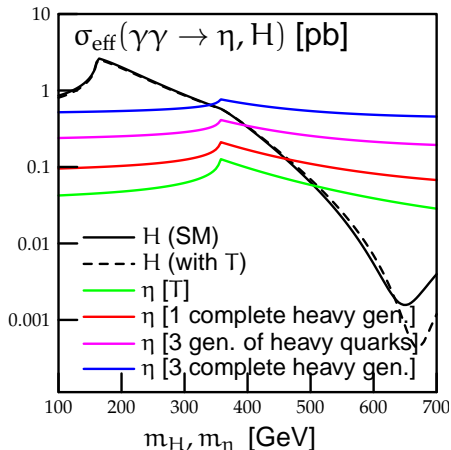
- ▶ S/B analogous to LC
- ▶ η in the μ model with (almost) identical parameters as A in MSSM
(\leftrightarrow Mühlleitner et al. (2001))

Pseudo Axions at the Photon Collider

- ▶ **Photon Collider** as precision machine for Higgs physics (s channel resonance, anomaly coupling)

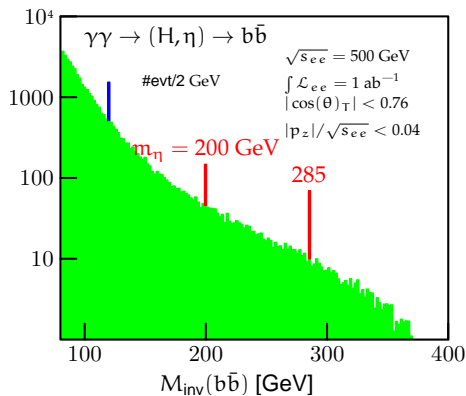
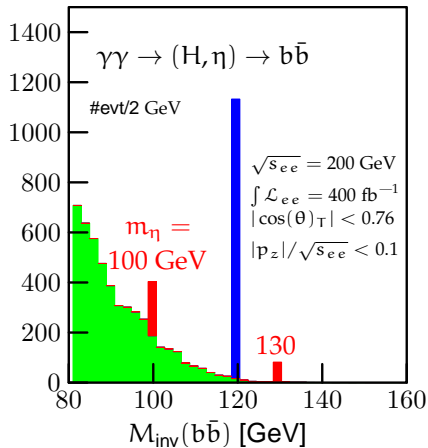


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$$g_{bb\eta} = 0.4 \cdot g_{bbh}$$

m_η	100	130	200	285
$\Gamma_{\gamma\gamma}$ [keV]	0.15	0.27	1.1	3.6



Simplest Little Higgs (“ μ Model”)

Schmaltz '04, Kilian/Rainwater/JRR '04

Field content ($SU(3)_c \times SU(3)_w \times U(1)_X$ quantum numbers)

$$\begin{array}{lll}
 \Phi_{1,2} & : & (1, 3)_{-\frac{1}{3}} \\
 \Psi_Q & : & (3, 3)_{\frac{1}{3}} \\
 \Psi_\ell & : & (1, 3)_{-\frac{1}{3}} \\
 d^c & : & (\bar{3}, 1)_{\frac{1}{3}} \\
 u_{1,2}^c & : & (\bar{3}, 1)_{-\frac{2}{3}} \\
 e^c, n^c & : & (1, 1)_{1,0}
 \end{array}$$

Lagrangian $\mathcal{L} = \mathcal{L}_{\text{kin.}} + \mathcal{L}_{\text{Yuk.}} + \mathcal{L}_{\text{pot.}}$ $\Psi_{Q,L} = (u, d, U)_L, \Psi_\ell = (\nu, \ell, N)_L$:

$$\begin{aligned}
 \mathcal{L}_{\text{Yuk.}} = & -\lambda_1^u \bar{u}_{1,R} \Phi_1^\dagger \Psi_{T,L} - \lambda_2^u \bar{u}_{2,R} \Phi_2^\dagger \Psi_{T,L} - \frac{\lambda^d}{\Lambda} \epsilon^{ijk} \bar{d}_R^b \Phi_1^i \Phi_2^j \Psi_{T,L}^k \\
 & - \lambda^n \bar{n}_{1,R} \Phi_1^\dagger \Psi_{Q,L} - \frac{\lambda^e}{\Lambda} \epsilon^{ijk} \bar{e}_R \Phi_1^i \Phi_2^j \Psi_{Q,L}^k + \text{h.c.},
 \end{aligned}$$

$$\mathcal{L}_{\text{pot.}} = \mu^2 \Phi_1^\dagger \Phi_2 + \text{h.c.}$$

Hypercharge embedding ($\text{diag}(1, 1, -2)/(2\sqrt{3})$):

$$Y = X - T^8/\sqrt{3} \qquad D_\mu \Phi = (\partial_\mu - \frac{1}{3} g_X B_\mu^X \Phi + ig W_\mu^w) \Phi$$