

# Polarised Gas Targets and Polarised Ion Sources for Accelerators

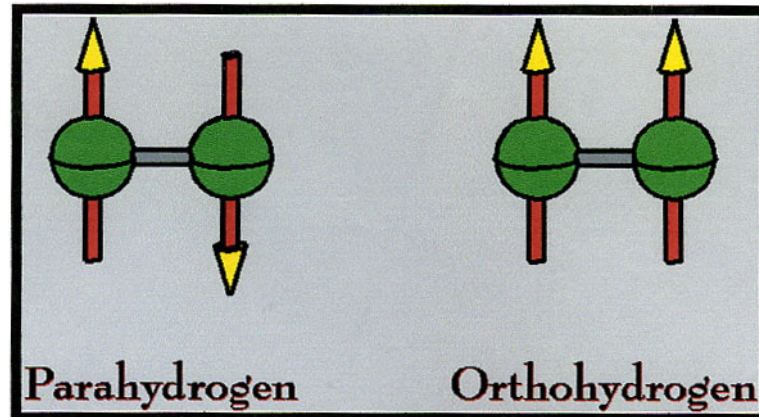
*Geoff Court – Physics Dept., Liverpool University*

Mainly principles – time limitation

Protons only – ideas apply for other nuclei (D,  $^3\text{He}$  .....

Start point – gaseous hydrogen molecules ( $\text{H}_2$ )

$I = 0$   
singlet state



$I = 1$   
triplet state

## **End point for polarised beam sources**

- a) positive atomic ion beam with single spin state
- b) negative atomic ion beam with single spin state

## **End point for gaseous polarised target**

- c) neutral atomic beam (gas flow) with single spin state
- d) neutral molecular ortho beam (gas) - single spin state

Applied magnetic field  $B$  - quantisation axis

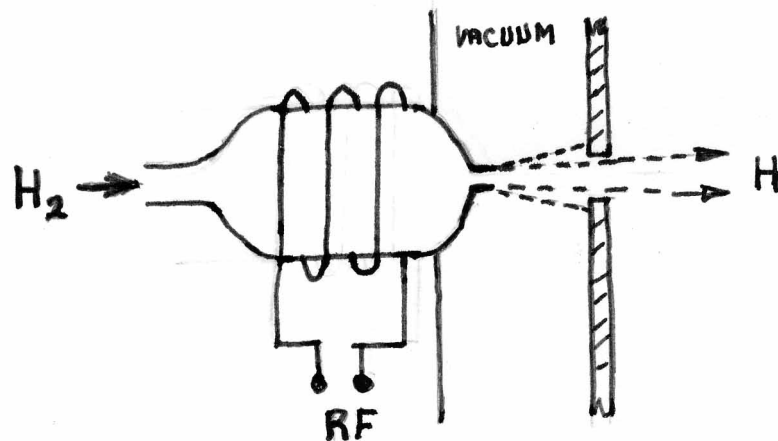
Polarisation  $P = (N_{\uparrow} - N_{\downarrow}) / (N_{\uparrow} + N_{\downarrow})$  with limits  $\pm 1$

+ 1 is parallel and -1 is anti-parallel to  $B$

Normal first stage for a), b) and c) is dissociation in a RF discharge



Second stage is to form an atomic beam (AB) (or use gas flow)



Atoms effuse thru channel,  
beam formed if  $\text{MFP} \gg \text{diam.}$

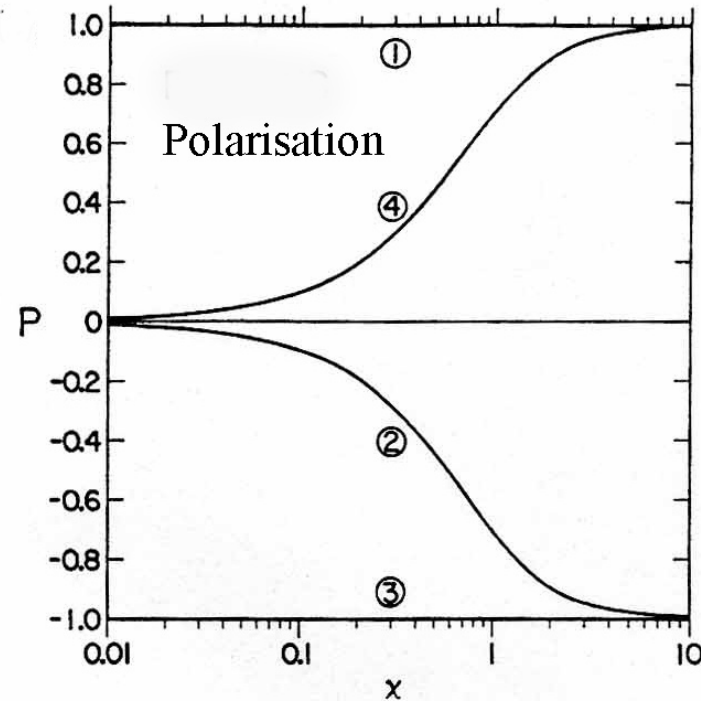
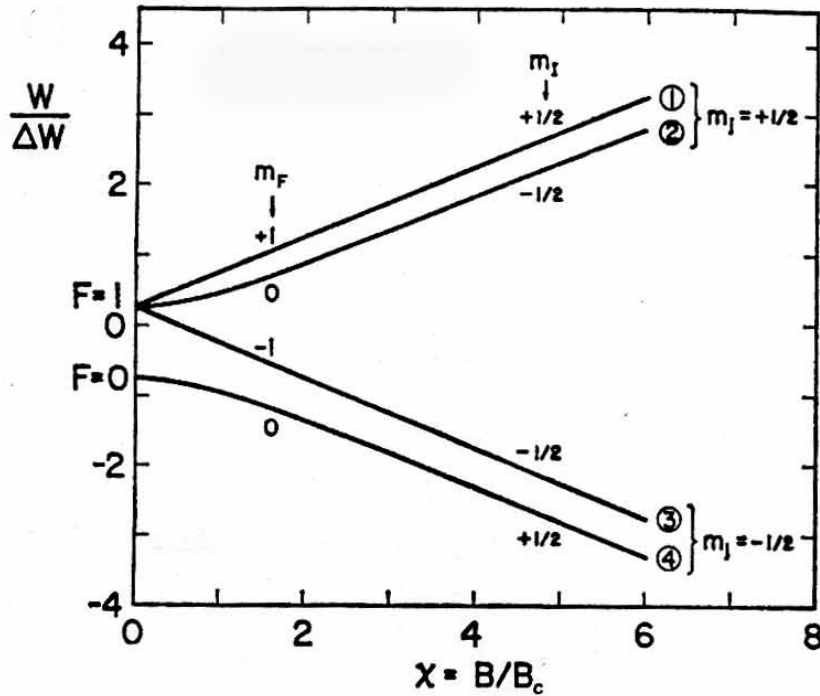
Thermal energy (E)– large  $\Delta E$

Large angular spread

Limitation - recombination ( $E = 3.4 \text{ eV}$ ), but not in gas phase

Conservation of Energy and momentum – only on surfaces

# Atomic hydrogen – spin energy levels in mag field



Hyperfine levels

$\Delta W(\nu)$  for  $F = 0$  to  $F = 1$  is 1420 MHz

States 1 and 3 are QM pure spin states, 2 and 4 are mixed states

## State selection type 1.

Stern – Gerlach spatial separation technique

Spins experience force in inhomogenous B-field

$F = \mu(dB_z/dz)$  with z axis is in direction of gradient

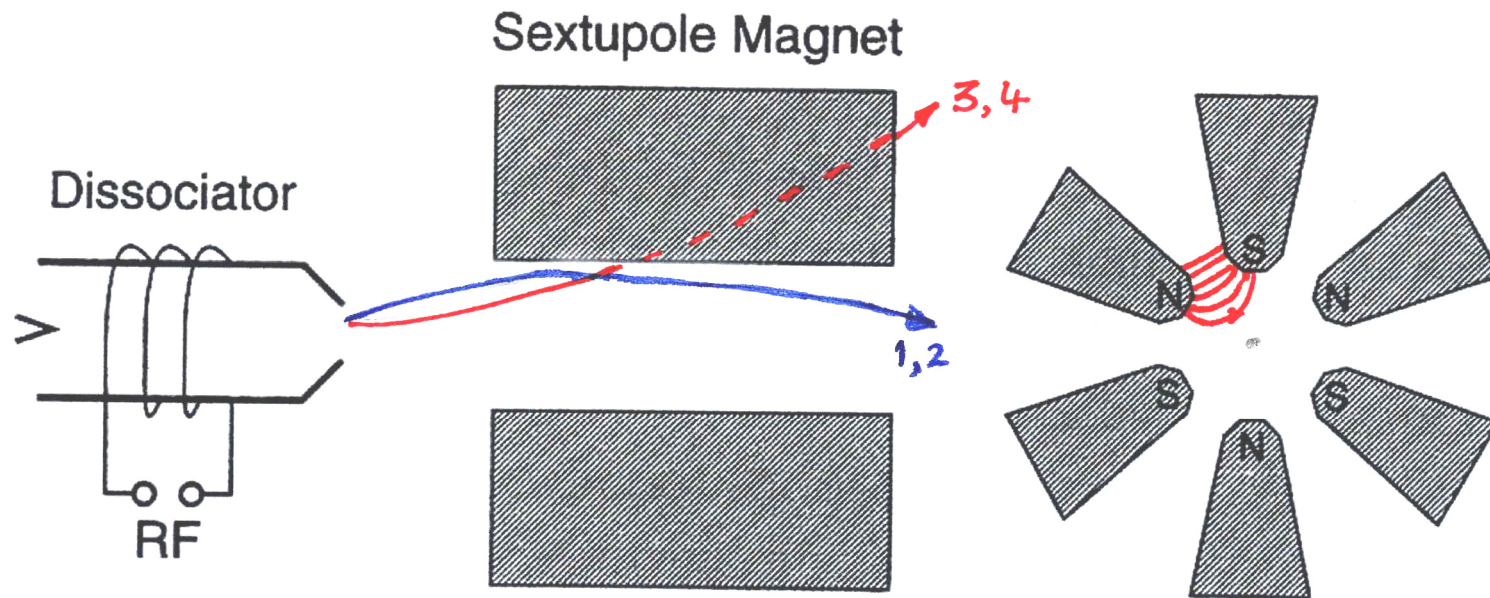
$\mu$  is effective magnetic moment

Hence in field B the force on the **electron spin** makes

States 1 and 2 ( $m_j = + 1/2$ ) **low** field seekers (LFS)

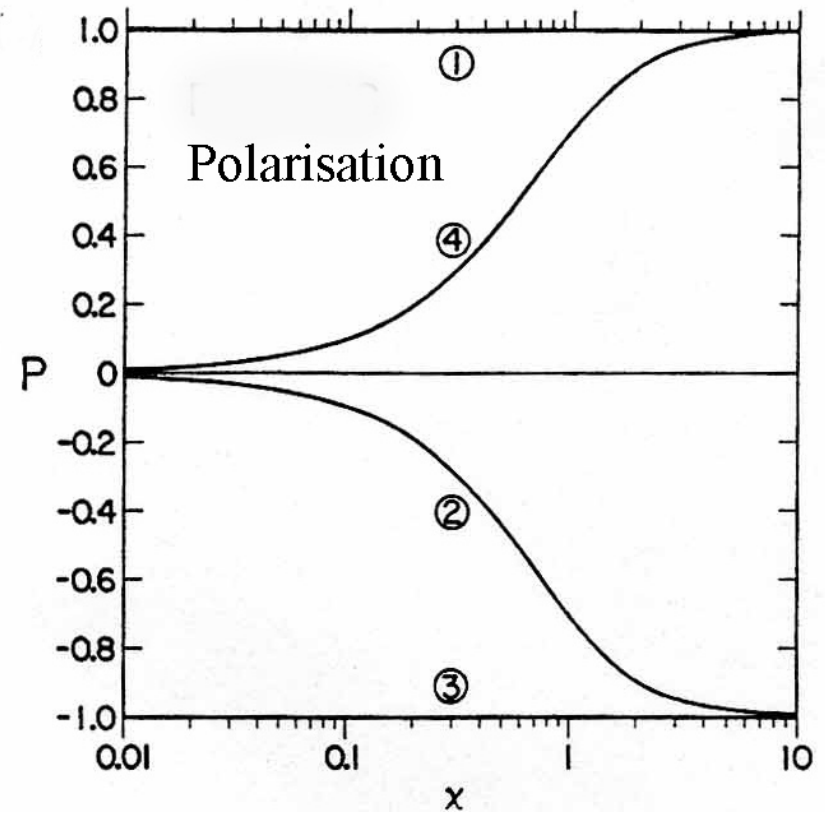
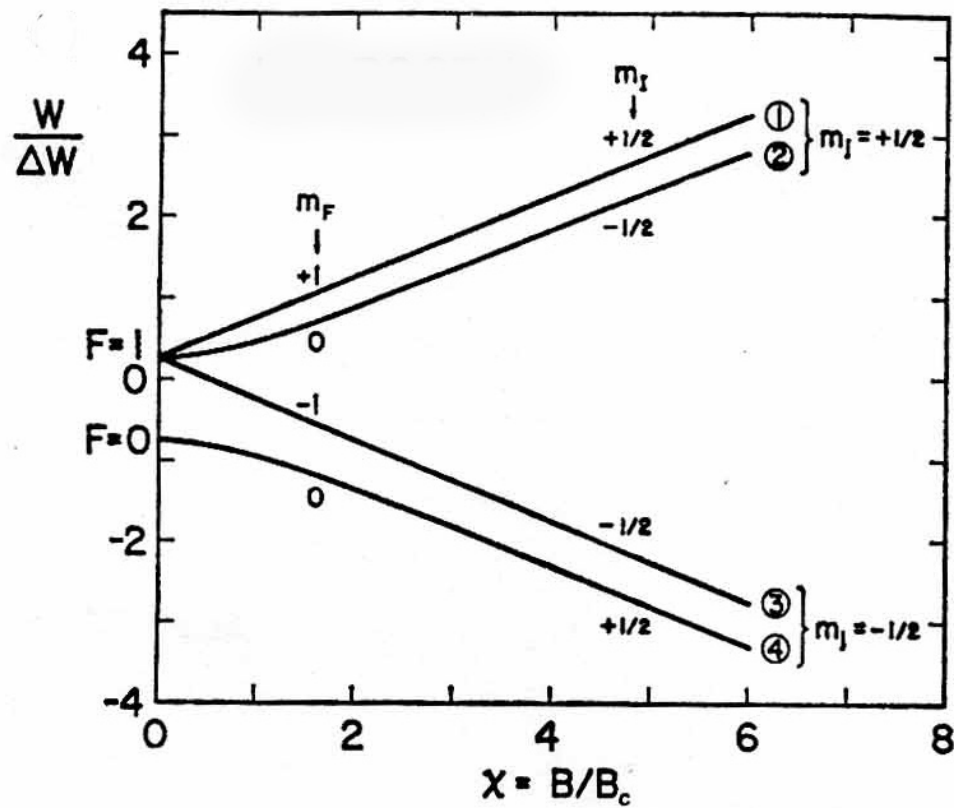
States 3 and 4 ( $m_j = - 1/2$ ) **high** field seekers.(HFS)

Remember  $\mu_p \ll \mu_e$



LFS focussed and HFS defocussed,  
 $P_e = 1$  and  $P_n = 0.5$  in low B and  $P_n = 0$  in high B

Modify state populations with adiabatic RF transitions  
 e.g. exchange population from 1 to 3 then  $P_e = 0$ ,  $P_n = -1$

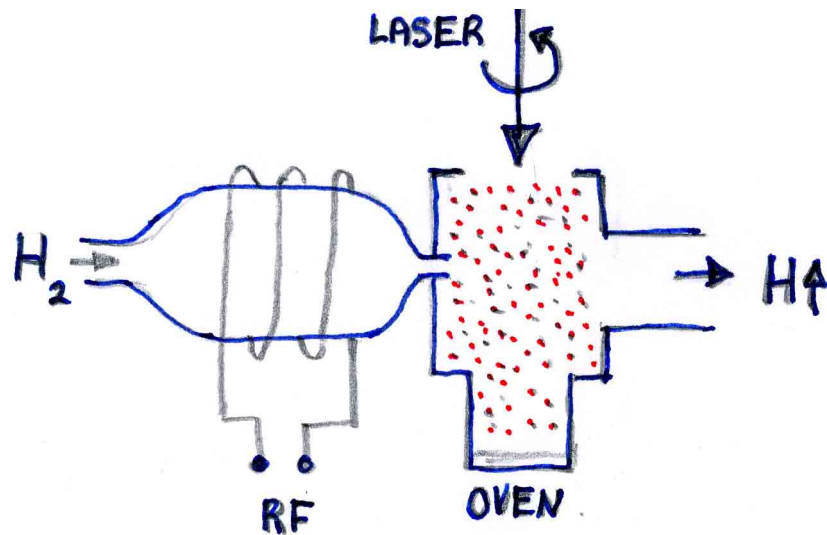


## State selection 2.

Optical pumping technique – select spin state using **polarised** laser source at transition frequency (LDS)

Problem with H – atomic transition levels are in far UV

Practical solution - pass beam through optically electron polarised alkali metal vapour (visible or IR)



Spin exchange interaction gives initially – high  $P_e$  with high  $P_n$  after large number of collisions

Uses Na, K, Rb or Cs

Result from 1 or 2 is a neutral atomic beam or gas flow

In principle with  $P_e$  or  $P_n = \pm 1$  (actually 0.8 to 0.9)

### Use as a gaseous target

#### Polarised jet target

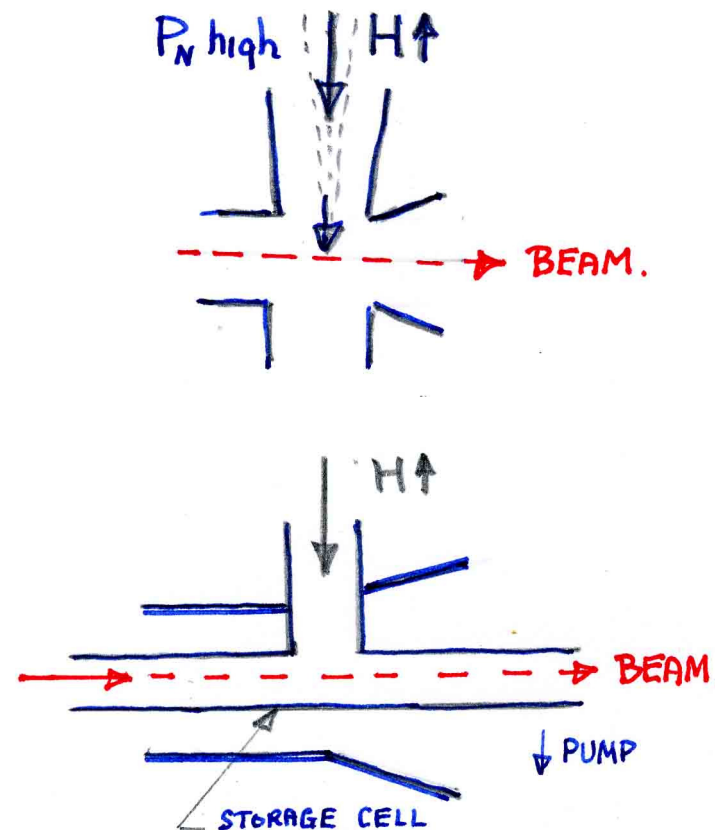
Needs cool  $H\uparrow$  beam

Very low luminosity

#### Storage cell target

Needs cooled cell

Higher lumi ( $\approx \times 100$ )



## Polarised ion sources

1. Using an ABS or LDS - thermal atomic beam(gas flow)

a) Convert  $H^0$  beam to  $H^+$  beam with an ioniser

RF discharge - electron cyclotron resonance (ECR)

High static field, electrons circulate in resonance with RF

High efficiency in small volume, field inhibits  $P_p$  loss

Ion extraction, beam formation via DC potential ( $\sim 3KV$ )

b) Convert  $H^0$  to  $H^-$

Pass  $H^0$  beam through Na vapour cell, picks up  $e^-$

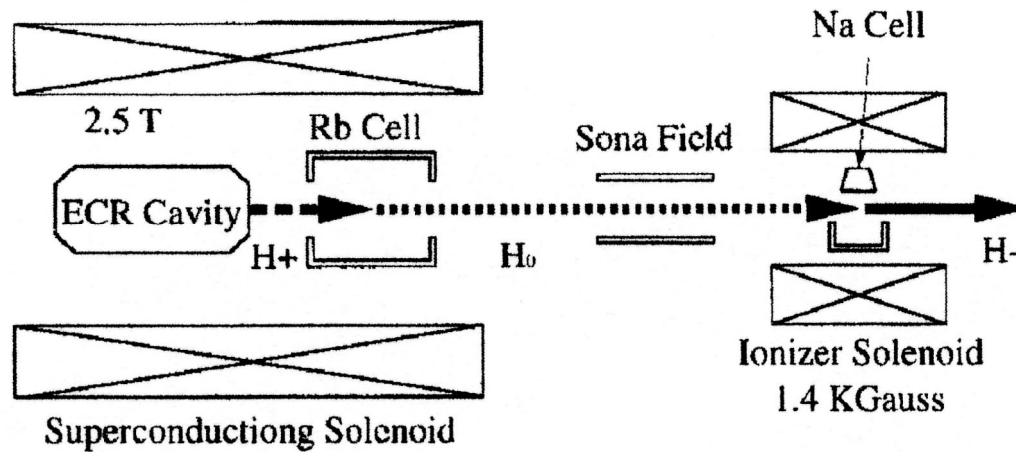
Works best with higher energy  $H^0$  beam

## 2. Using a high energy $H^0$ beam (OPPIS)

- a) Generate  $\sim 3.5$  KeV  $H^+$  beam with ECR
- b) Capture polarised electron from optically pumped Rb atoms giving  $H^0 (e\uparrow)$ .
- c) Transfer electron spin to proton via Sona \* transition giving  $H^0 (p\uparrow)$ .
- d) Ionise to  $H^-$  with Na cell

\* Sona transition is spin state population redistribution when beam crosses a field zero

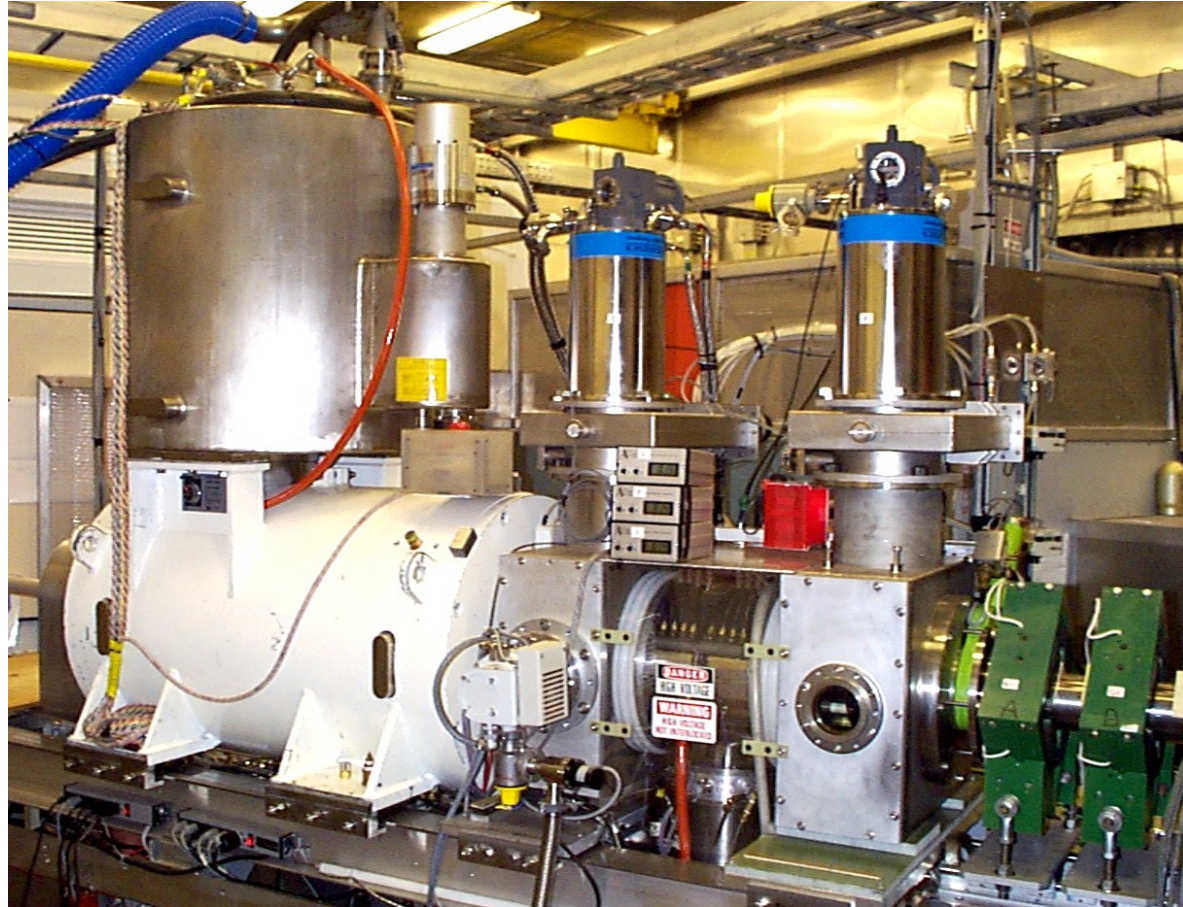
# Typical OPPIS schematic layout



Polarisation reversal - optical technique

Change sense of circular polarisation of laser light

# A High Intensity Polarised Ion Source



KEK OPPIS  
upgraded at TRIUMF

75 - 80 % Polarization

$15 \times 10^{11}$  protons/pulse  
at source

$6 \times 10^{11}$  protons/pulse  
at end of LINAC

For use in RHIC at Brookhaven

## General issues

### 1. Polarisation reversal

Needs to be fast and systematic effect free

- With ABS change RF transition frequency
- With LDS or OPPIS change laser polarisation

### 2. Which technique?

ABS (LDS?) for polarised targets

OPPIS for polarised ion source

### 3. Other Nuclei

a) Deuteron - spin 1

Techniques mainly as for proton

b)  $^3\text{He}$  - spin  $\frac{1}{2}$ ,

Optical pumping only as it is a two electron atom  
(no electron polarisation possible)

## **Information and references**

Polarised Gas Targets - Steffens and Haberli  
Rep. Prog. Phys. **66**(2003) 1187-1935

Polarised Ion Sources for High Energy  
Accelerators – Levy and Zelenski  
Rev. Sci. Inst. **69** - 2 (1998) 732-736

International Symposium on High Energy Spin  
Physics - Biennial Conference series with 17<sup>th</sup> in  
Japan (SPIN06) this year.