

solutions - exercises 2:

problem 1:

We can go to a frame where $U^\mu = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}$

$$U^\mu U_\mu = V^\mu V_\mu = -1$$

In this frame V^μ will be of the form

$$V^\mu = \gamma \begin{pmatrix} 1 \\ \vec{v} \end{pmatrix}$$

$$\Rightarrow U^\mu V_\mu = -\gamma = -\sqrt{\frac{1}{1-\vec{v}^2}}$$

$$\Rightarrow |\vec{v}| = \sqrt{1 - \frac{1}{(\gamma^2)^2}}$$

problem 2:

The product of the two boosts is given by

$$\Lambda^\mu{}_\nu(\varphi_1) \Lambda^\nu{}_\lambda(\varphi_2) = \begin{pmatrix} A & B & & \\ B & A & & \\ & & 1 & \\ & & & 1 \end{pmatrix}$$

$$A = \cosh \varphi_1 \cosh \varphi_2 + \sinh \varphi_1 \sinh \varphi_2 = \cosh(\varphi_1 + \varphi_2)$$

$$B = \cosh \varphi_1 \sinh \varphi_2 + \cosh \varphi_2 \sinh \varphi_1 = \sinh(\varphi_1 + \varphi_2)$$

$$\Lambda^{\mu}_{\nu}(\varphi_1) \Lambda^{\nu}_{\rho}(\varphi_2) = \Lambda^{\mu}_{\rho}(\varphi_1 + \varphi_2)$$

problem 3:

One can write the rotation as

$$\Lambda^{\mu}_{\nu} = \left(\begin{array}{c} 1 \\ R_{ij} \end{array} \right)^{\mu}_{\nu}$$

And the boost along x as

$$\Lambda^{\mu}_{\nu} = \begin{pmatrix} c & -s \\ -s & c \\ & & 1 & \\ & & & \ddots \end{pmatrix}$$

In order to generate a general boost, one can rotate + boost + rotate back

$$\left(\begin{array}{c} 1 \\ R \end{array} \right) \left(\Lambda^{\mu}_{\nu} \right) \left(\begin{array}{c} 1 \\ R^T \end{array} \right)$$

The general boost is given by

$$\Lambda = \begin{pmatrix} \gamma & -\gamma v_i \\ -\gamma v_i & \delta_{ij} + (\gamma-1) \frac{v_i v_j}{v^2} \end{pmatrix}$$

$$\left(\begin{array}{c} 1 \\ R \end{array} \right) \Lambda \left(\begin{array}{c} 1 \\ R^T \end{array} \right)$$

$$= \begin{pmatrix} \gamma & -\gamma (Rv)_i \\ -\gamma (Rv)_i & R \delta_{ij} R^T + (\gamma-1) \frac{(Rv)_i (Rv)_j}{v^2} \end{pmatrix}$$

The final result is that rotation+boost+rotation back is equivalent to a boost with a rotated velocity.

problem 4:

The proof follows the same proof for the energy momentum tensor in the lecture.

$$\begin{aligned}
 \dot{J}^{\mu} &= \sum q_n \delta^3(x-x_n) \frac{dx_n^{\mu}}{dt} \\
 \partial_i \dot{J}^i &= \sum q_n \partial_i \delta(x-x_n) \frac{dx_n^i}{dt} \\
 &= \sum q_n \left(\frac{\partial}{\partial x_n} \right) \delta(x-x_n) \frac{dx_n^i}{dt} \\
 &= \sum q_n (-\partial_i) \delta(x-x_n) \\
 &= -\partial_i \dot{J}^0 \Rightarrow \partial_{\mu} \dot{J}^{\mu} = 0
 \end{aligned}$$

problem 5:

$$\begin{aligned}
 \partial_{\mu} \dot{J}^{\mu} = 0 &\Rightarrow \\
 0 &= \int d^3x \partial_{\mu} \dot{J}^{\mu} = \int d^3x (\partial_t \dot{J}^0) \\
 &\quad + \underbrace{\int d^3x \vec{\nabla} \cdot \vec{\dot{J}}}_{\text{boundary term}}
 \end{aligned}$$

boundary term = 0
(by assumption)

$$\rightarrow 0 = \partial_t \int d^3x \gamma^0 = \partial_t Q$$

problem 6:

$$T_{em}^{\mu\nu} = F^\mu_\alpha F^{\nu\alpha} - \frac{1}{4} \eta^{\mu\nu} F_{\alpha\beta} F^{\alpha\beta}$$

$$\partial_\mu T_{em}^{\mu\nu} = (\partial_\mu F^\mu_\alpha) F^{\nu\alpha} + F^\mu_\alpha \partial_\mu F^{\nu\alpha} - \frac{1}{2} \eta^{\mu\nu} (\partial_\mu F_{\alpha\beta}) F^{\alpha\beta}$$

Bianchi: $\partial_\mu F_{\alpha\beta} + \partial_\beta F_{\mu\alpha} + \partial_\alpha F_{\beta\mu} = 0$

$$F^{\alpha\beta} (\partial_\mu F_{\alpha\beta} + \partial_\beta F_{\mu\alpha} + \partial_\alpha F_{\beta\mu})$$

$$= F^{\alpha\beta} (\partial_\mu F_{\alpha\beta} + \partial_\beta F_{\mu\alpha} - \partial_\beta F_{\alpha\mu})$$

$$= F^{\alpha\beta} (\partial_\mu F_{\alpha\beta} + 2\partial_\beta F_{\mu\alpha})$$

$$\Rightarrow \partial_\mu T_{em}^{\mu\nu} = -\partial_\alpha F^{\nu\alpha}$$

problem 7:

First, we show the following relation

$$R^{ia} R^{ib} R^{ic} \epsilon^{abc} = \det R \epsilon^{i'k}$$

First observe that the left-hand side is completely antisymmetric in the indices ijk .

So

$$R^i R^j R^k \epsilon_{ijk} \propto \epsilon_{ijk}$$

The proportionality is given by choosing $ijk = 123$:

$$R^{1a} R^{2b} R^{3c} \epsilon^{abc} = \det R$$

If R is a rotation matrix then

$$\det R = 1 \quad ; \quad R^{ai} R^{bi} = \delta^{ab}$$

$$R^{ia} R^{ib} \cancel{R^{jd} R^{kc}} \epsilon^{abc} = R^{kd} \epsilon^{ijk}$$

δ^{dc}

$$R^{ia} R^{ib} \epsilon^{abd} = R^{kd} \epsilon^{ijk}$$

The tensor components F^{ij} transform under the spatial rotations as

$$\begin{aligned} F^{ij} &\rightarrow R^{ia} R^{jb} F^{ab} = R^{ia} R^{jb} \epsilon^{ijk} B^k \\ &= R^{kd} \epsilon^{ijk} B^d \\ &= \epsilon^{ijk} (R^{kd} B^d) \\ B^d &\rightarrow R^{kd} B^d = R \cdot B \end{aligned}$$