Exercises, week 2

1 Relative velocity

Since velocities do not behave additively in special relativity, there is the question how to determine the relative velocities between two objects in a frameindependent way. Given two velocities U^{μ} and V^{μ} , one way of defining this is using

$$\Delta v = \sqrt{1 - \frac{1}{(U_{\mu}V^{\mu})^2}}.$$
(1)

Derive this expression by boosting into the frame where one of the objects is at rest.

2 Rapidities

We have seen that the Lorentz-transformation for a boost in the x-direction is of the form $\begin{pmatrix} x & y \\ y & y \end{pmatrix} = \begin{pmatrix} x & y \\ y & y \end{pmatrix}$

$$\Lambda^{\mu}_{\ \nu}(\phi) = \begin{pmatrix} \cosh(\phi) & -\sinh(\phi) & 0 & 0\\ -\sinh(\phi) & \cosh(\phi) & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & 1 \end{pmatrix}.$$
 (2)

Show that the rapidity ϕ indeed behaves additively by calculating $\Lambda^{\mu}_{\ \nu}(\phi_1)\Lambda^{\nu}_{\ \kappa}(\phi_2)$.

3 General boosts

Show that boosts in a general direction are given by

$$\Lambda^{\mu}_{\ \nu}(\phi) = \begin{pmatrix} \gamma & -\gamma \vec{v}_i \\ -\gamma \vec{v}_j & \delta_{ij} + (\gamma - 1) \vec{v}_i \vec{v}_j / v^2 \end{pmatrix}$$
(3)

by rotating a boost in the x-direction into other directions.

4 Current conservation

Show that the following current is conserved $(q_n \text{ is the charge for the particle } n)$

$$J^{\mu} = \sum_{n} q_n \delta(\vec{x} - \vec{x}_n) \frac{dx_n^{\mu}}{dt} \,. \tag{4}$$

5 Charge conservation

Show that current conservation $\partial_{\mu}J^{\mu}$ implies that the following charge is timeindependent (assume that the current falls off quickly enough at infinity)

$$Q = \int d^3x \, J^0 \,. \tag{5}$$

6 Electromagnetic energy-momentum tensor

Show that the following tensor involving the field strength $F^{\mu\nu}$

$$T_{em}^{\mu\nu} = F^{\mu}_{\ \alpha} F^{\nu\alpha} - \frac{1}{4} \eta^{\mu\nu} F_{\alpha\beta} F^{\alpha\beta} \tag{6}$$

has as the divergence

$$\partial_{\mu}T^{\mu\nu}_{em} = -F^{\nu}_{\ \gamma}J^{\gamma}\,,\tag{7}$$

by writing the field strength tensor in terms of the potential A^{μ} and using Maxwell's equations (you also can use the Bianchi identity from last week to simplify this derivation).

7 Magnetic field

In special relativity, the spatial coordinates of the field strength tensor are given by the magnetic field

$$F^{ij} = \epsilon^{ijk} B_k \tag{8}$$

Show that B_k transforms as 3-vector under normal spatial rotations even though F^{ij} transform as an antisymmetric tensor.