

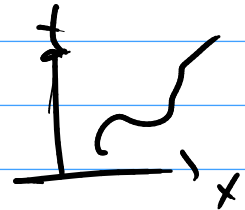
Energy and momentum

Let's say you have a particle moving and the path is given by

$$\vec{x}(t)$$

$$x^\mu = \begin{pmatrix} t(\lambda) \\ \vec{x}(\lambda) \end{pmatrix}$$

$$\lambda = t$$



proper time:

$$ds^2 = dx^\mu \eta_{\mu\nu} dx^\nu = -dt^2 + dx^2$$

$$d\tau^2 = -ds^2 \leq 0$$

$$\underline{\lambda = \tau} :$$

$$\frac{dx^\mu}{d\tau} \frac{dx_\mu}{d\tau} = \frac{ds^2}{d\tau^2} = -1$$

Energy, momentum:

$$\vec{p} \approx m \vec{v} \quad ; \quad \frac{dx^\mu}{d\tau} = \gamma \begin{pmatrix} 1 \\ \vec{v} \end{pmatrix}$$

$$p^\mu = m \frac{dx^\mu}{d\tau}$$

rest energy
↓
kinetic energy

$$p^0 = \gamma m = m \sqrt{1/\gamma^2 - v^2} \approx m + \frac{1}{2} m v^2 + \dots$$

Evidently:

$$p_\mu p^\mu = -m^2$$

[onshell condition, dispersion relation]

forces:

If we believe that everything is Lorentz covariant, we can start in the restframe of the particle and then boost the result

$$\vec{f} = m\vec{a} = m \frac{d\vec{x}}{dt}$$

The acceleration has a strange transformation property. It coincides with

$$\frac{d^2\vec{x}}{d\tau^2} \quad \text{when} \quad \frac{d\vec{x}}{d\tau} \rightarrow 0 \quad (\tau \approx t)$$

$$\Rightarrow f^\mu = m \frac{dx^\mu}{d\tau^2}$$

where

$$\begin{aligned} f^i &= \vec{f} \\ f^0 &= 0 \end{aligned} \quad \text{in the restframe}$$

Notice that

$$\frac{dx^\mu}{d\tau} \frac{dx_\mu}{d\tau} = -1$$

$$\hookrightarrow \left(\frac{d}{d\tau} \frac{dx^\mu}{d\tau} \right) \eta_{\mu\nu} \frac{dx^\nu}{d\tau} = 0$$

$$\rightarrow f^\mu \eta_{\mu\nu} \frac{dx^\nu}{d\tau} = 0$$

$$\Leftrightarrow f^0 = 0 \quad \text{in the restframe}$$

So any consistent force has to fulfill the constraint

$$\int^{\tau} \frac{dx^{\nu}}{d\tau} \eta_{\mu\nu} = 0$$

Example: particle moving in an electric field:

$$\frac{d\vec{p}}{dt} = q \vec{E} + q \vec{v} \times \vec{B}$$

The generalization reads

$$f^{\mu} = \frac{dp^{\mu}}{d\tau} = m \frac{dx^{\mu}}{d\tau}$$

$$= q \cdot \eta_{\beta\gamma} F^{\mu\beta} \frac{dx^{\gamma}}{d\tau}$$

$$\eta_{\mu} \int^{\tau} \frac{dx^{\nu}}{d\tau} \stackrel{!}{=} 0 \quad [\text{check consistency}]$$

$$= q F^{\alpha\beta} \frac{dx^{\alpha}}{d\tau} \frac{dx^{\nu}}{d\tau} \eta_{\mu\alpha} \eta_{\nu\beta}$$

$$= 0 \quad (\neq \text{antisym.})$$

Currents and densities

The current and density of a set of moving particles is given by

$$\rho(\vec{x}, t) = \sum_n \overset{\text{charges}}{e_n} \delta^3(\vec{x} - \vec{x}_n(t))$$

$$\vec{j}(\vec{x}, t) = \sum_n e_n \delta^3(\vec{x} - \vec{x}_n(t)) \frac{d\vec{x}_n}{dt}$$

This can be cast into a four-vector

$$j^\mu = \sum_n e_n \delta^3(\vec{x} - \vec{x}_n) \frac{dx_n^\mu}{dt}$$

In order to show that this is indeed a four-vector, we write

$$j^\mu(\vec{x}, t) = \sum_n \int dt' e_n \delta^4(x^\mu - x_n^\mu(t')) \frac{dx_n^\mu}{dt}$$

$$j^\mu(\vec{x}, t) = \sum_n \int dx e_n \overset{x_n^0 = t'}{\delta^4(x^\mu - x_n^\mu(\lambda))} \frac{dx_n^\mu}{d\lambda}$$

Also notice (see exercises)

$$\partial_\mu j^\mu = 0$$